

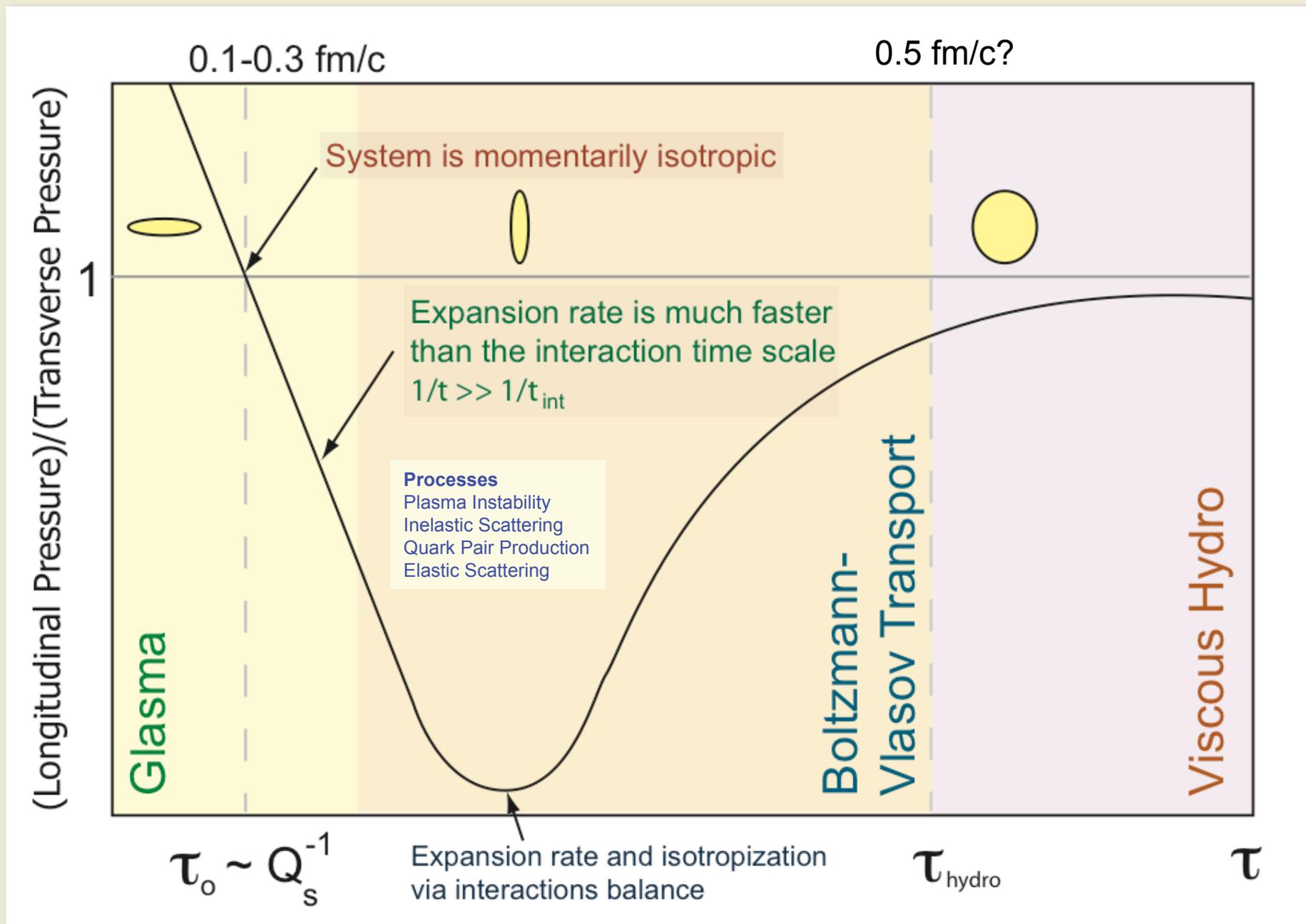
# Is early isotropization in heavy ion collisions a necessary condition to describe HIC data?

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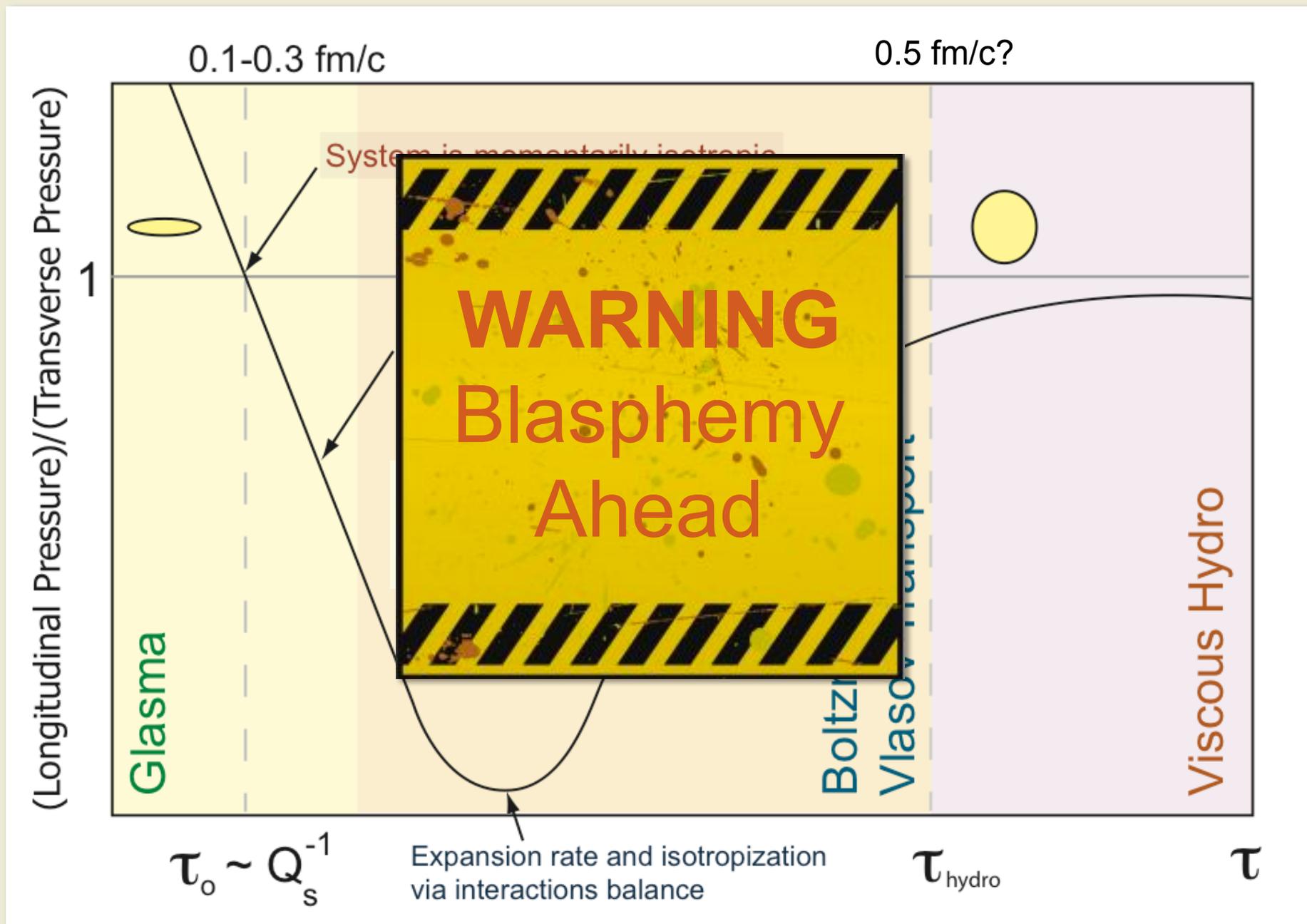
Heidelberg Thermalization Workshop  
Dec 12, 2011



# QGP momentum anisotropy



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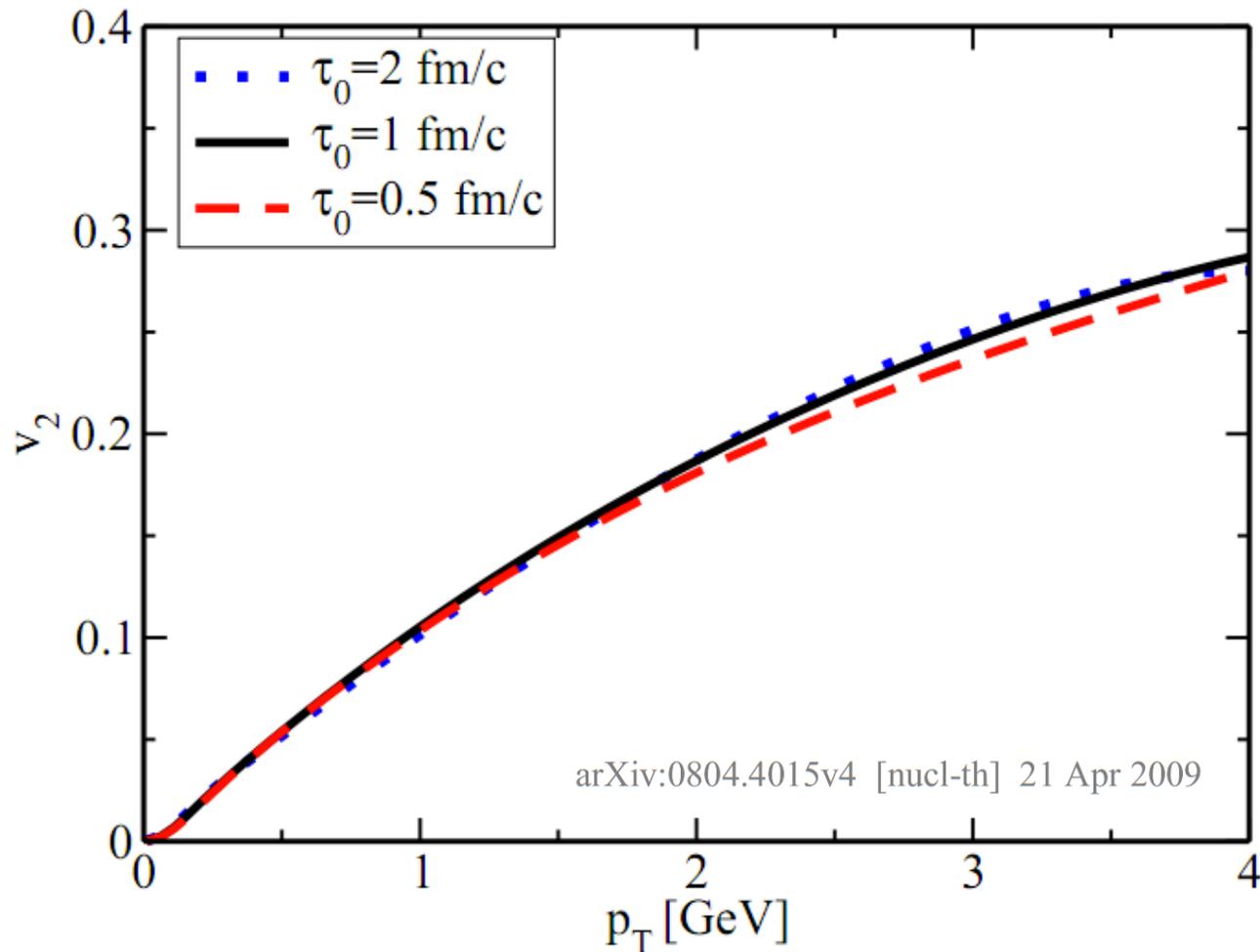
# Come ye of little faith ...



- It has been taken as gospel that agreement with experimental data for elliptic flow requires early isotropization/thermalization at times on the order of  $0.5 \text{ fm}/c$ .
- Is that true within viscous hydro?
- Let's ask some experts...

# Conformal Relativistic Viscous Hydrodynamics: Applications to RHIC results at $\sqrt{s_{NN}} = 200$ GeV

Matthew Luzum<sup>1</sup> and Paul Romatschke<sup>2</sup>

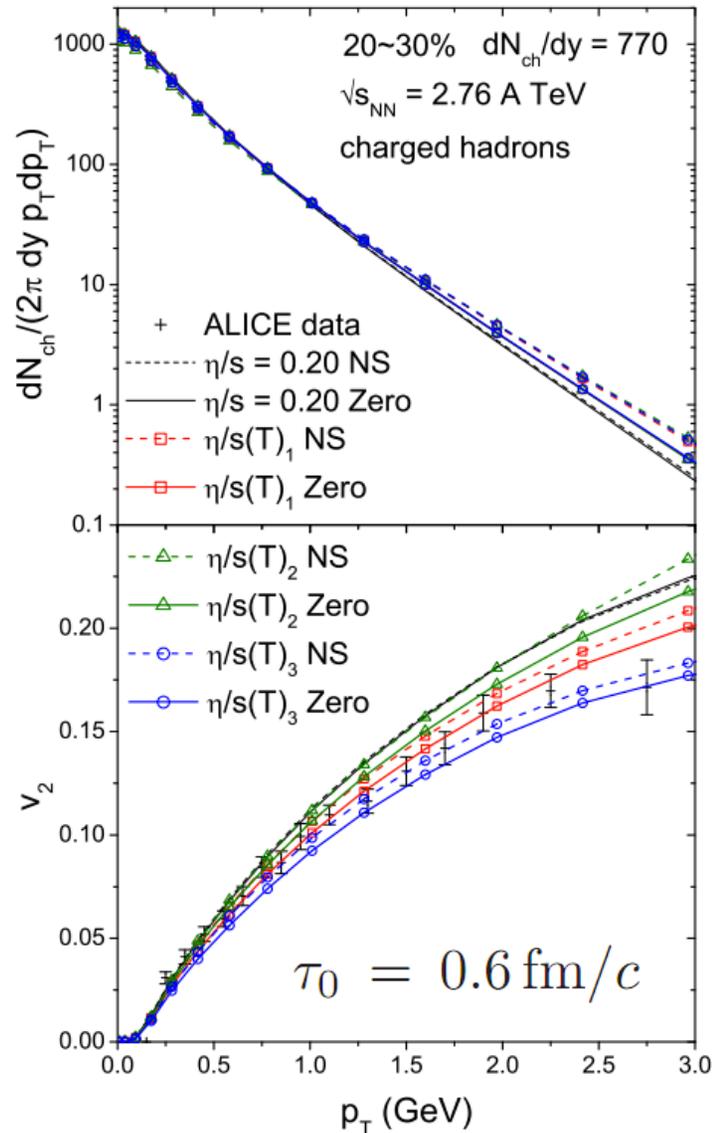


Simulation parameters were  $T_i = 0.29$  GeV,  $T_f = 0.14$  GeV for  $\tau_0 = 2$  fm/c,  $T_i = 0.36$  GeV,  $T_f = 0.15$  GeV for  $\tau_0 = 1$  fm/c, and  $T_i = 0.43$  GeV,  $T_f = 0.16$  GeV for  $\tau_0 = 0.5$  fm/c.

# Radial and elliptic flow in Pb+Pb collisions at the Large Hadron Collider from viscous hydrodynamics

Chun Shen,<sup>1,\*</sup> Ulrich Heinz,<sup>1,†</sup> Pasi Huovinen,<sup>2,‡</sup> and Huichao Song<sup>3,§</sup>

arXiv:1105.3226v2 [nucl-th] 9 Sep 2011



$\eta/s$ model	$\pi_0^{\mu\nu}$	$s_0$ (fm <sup>-3</sup> )	$T_0$ (MeV)
$\eta/s = 0.2$	0	191.6	427.9
	NS	172.4	413.9
$(\eta/s)_1(T)$	0	179.6	419.2
	NS	119.3	368.7
$(\eta/s)_2(T)$	0	179.6	419.2
	NS	115.6	365.1
$(\eta/s)_3(T)$	0	175.2	416.0
	NS	116.6	366.1

NS = Navier Stokes

$$\xi = \frac{\langle p_T^2 \rangle}{2\langle p_L^2 \rangle} - 1 \quad \xi_{NS} = \frac{10}{T\tau} \frac{\eta}{S}$$

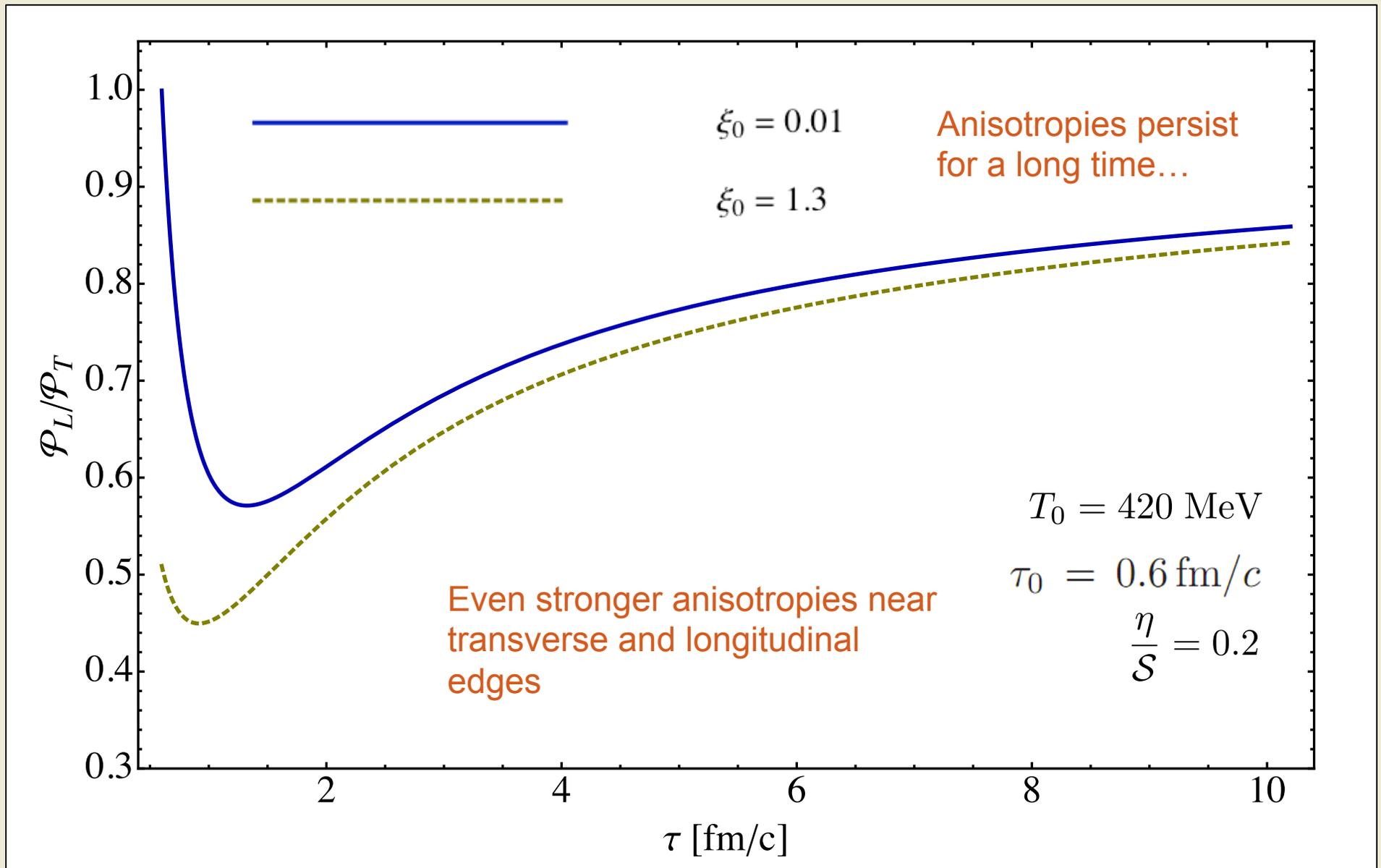
Using values from the paper I obtain

$$\xi_{0,NS} \simeq 1.3$$

Which corresponds to

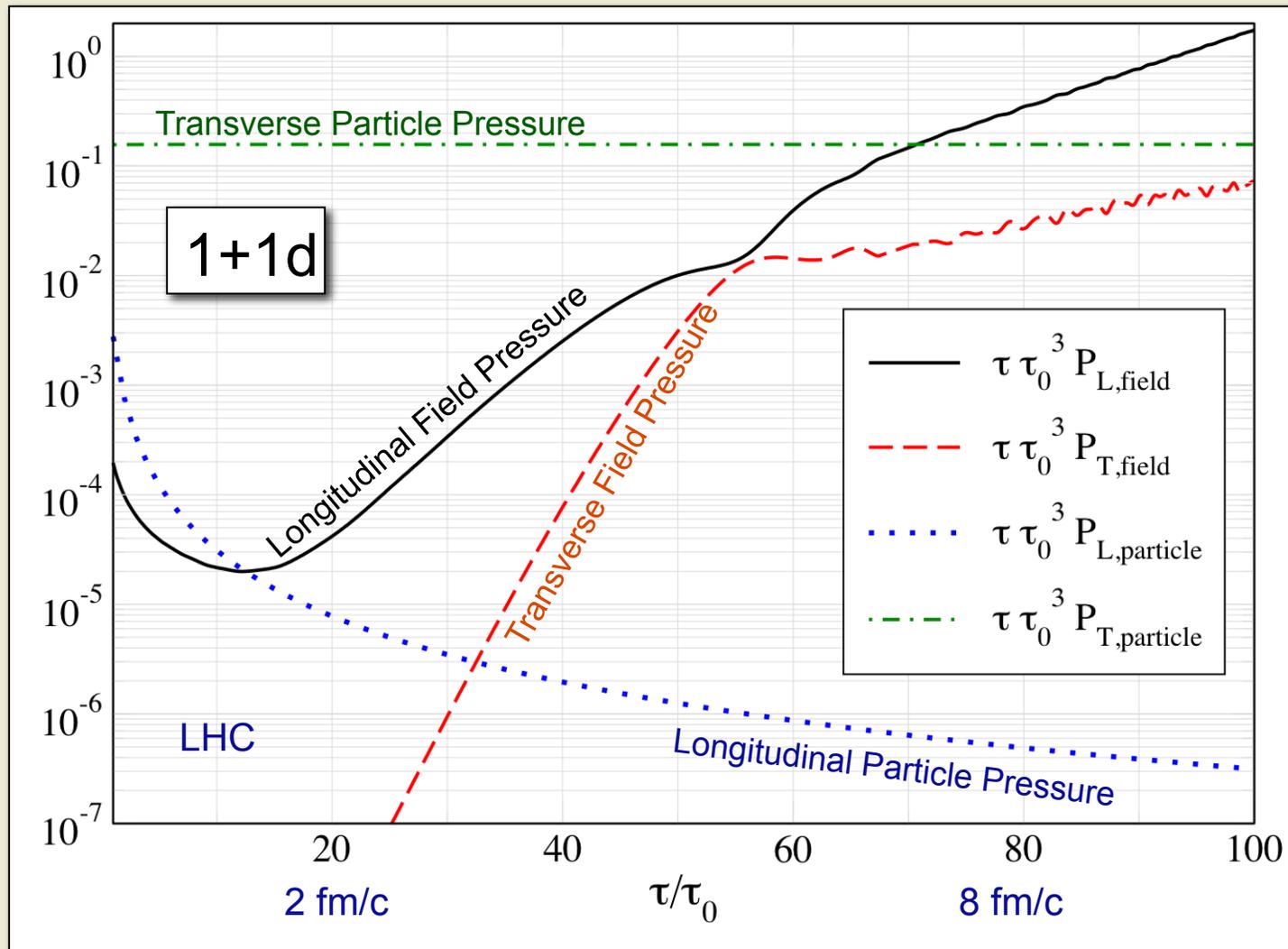
$$\mathcal{P}_L/\mathcal{P}_T \simeq 0.51$$

# Pressure Anisotropy as a Function of Time



# Hard loops in a free streaming background

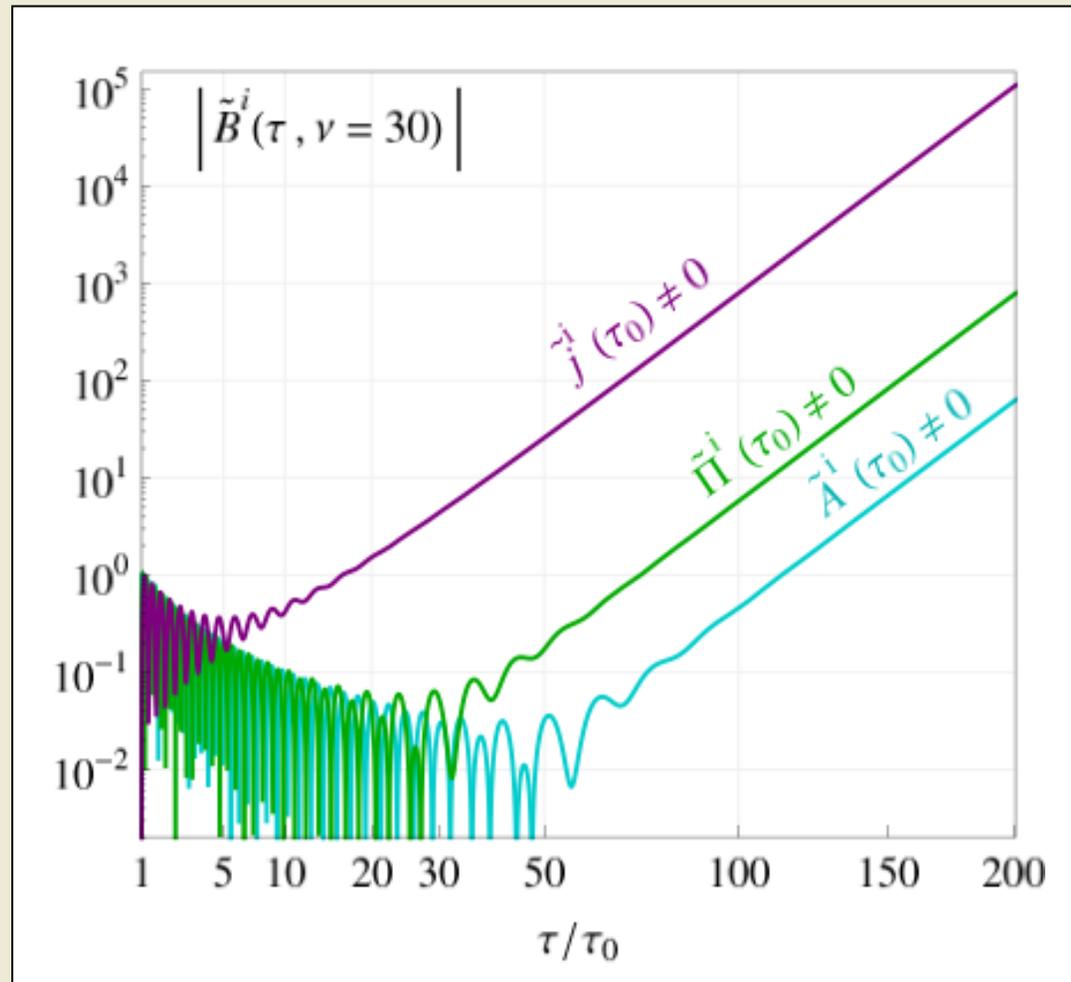
Particle and field pressures  $\rightarrow$  Isotropization [Rebhan, Strickland, and Attems, 2008]



3+1d results forthcoming! [Attems, Rebhan, Strickland 2012] Requires 5d Lattice Calc  $\sim 96^5$ !

# HEL - Different Initial Conditions

Rebhan and Steineder (2009) have shown the onset of unstable growth is accelerated dramatically if one has current fluctuations in the initial condition.



# Start over from scratch

## Viscous Hydrodynamics Expansion

$$f(\mathbf{x}, \mathbf{p}, \tau) = \underline{f_{\text{eq}}(|\mathbf{p}|, T(\tau))} + \delta f_1 + \delta f_2 + \dots$$

↑ Isotropic in momentum space

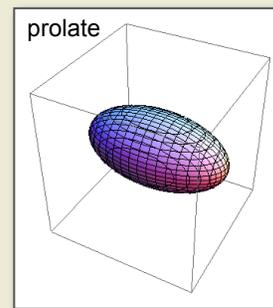
## Anisotropic Hydrodynamics (AHYDRO) Expansion

$$f(\mathbf{x}, \mathbf{p}, \tau) = f_{\text{aniso}}(\mathbf{p}, p_{\text{hard}}(\tau), \xi(\tau)) + \delta f'_1 + \delta f'_2 + \dots$$

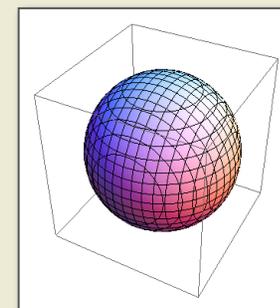
M. Martinez and MS, Nuclear Physics A 848, 183 (2010).

$$\xi = \frac{\langle p_T^2 \rangle}{2\langle p_L^2 \rangle} - 1$$

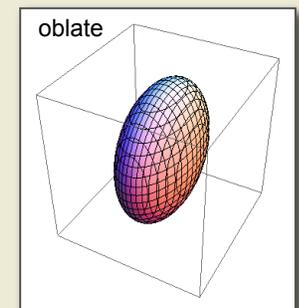
$$\begin{aligned} f(\tau, \mathbf{x}, \mathbf{p}) &= f_{RS}(\mathbf{p}, \xi(\tau), p_{\text{hard}}(\tau)) \\ &= f_{\text{iso}}([\mathbf{p}^2 + \xi(\tau)p_z^2]/p_{\text{hard}}^2(\tau)) \end{aligned}$$



$$-1 < \xi < 0$$



$$\xi = 0$$



$$\xi > 0$$

Using relaxation-time approximation scattering kernel gives in 1d

## 0<sup>th</sup> Moment of Boltzmann EQ

$$\partial_\alpha N^\alpha \neq 0$$

$$\frac{1}{1+\xi} \partial_\tau \xi - \frac{2}{\tau} - \frac{6}{p_{\text{hard}}} \partial_\tau p_{\text{hard}} = 2\Gamma \left[ 1 - \mathcal{R}^{3/4}(\xi) \sqrt{1+\xi} \right]$$

## 1<sup>st</sup> Moment of Boltzmann EQ

$$\partial_\alpha T^{\alpha\beta} = 0$$

$$\frac{\mathcal{R}'(\xi)}{\mathcal{R}(\xi)} \partial_\tau \xi + \frac{4}{p_{\text{hard}}} \partial_\tau p_{\text{hard}} = \frac{1}{\tau} \left[ \frac{1}{\xi(1+\xi)\mathcal{R}(\xi)} - \frac{1}{\xi} - 1 \right]$$

M. Martinez and MS, Nuclear Physics A 848, 183 (2010).

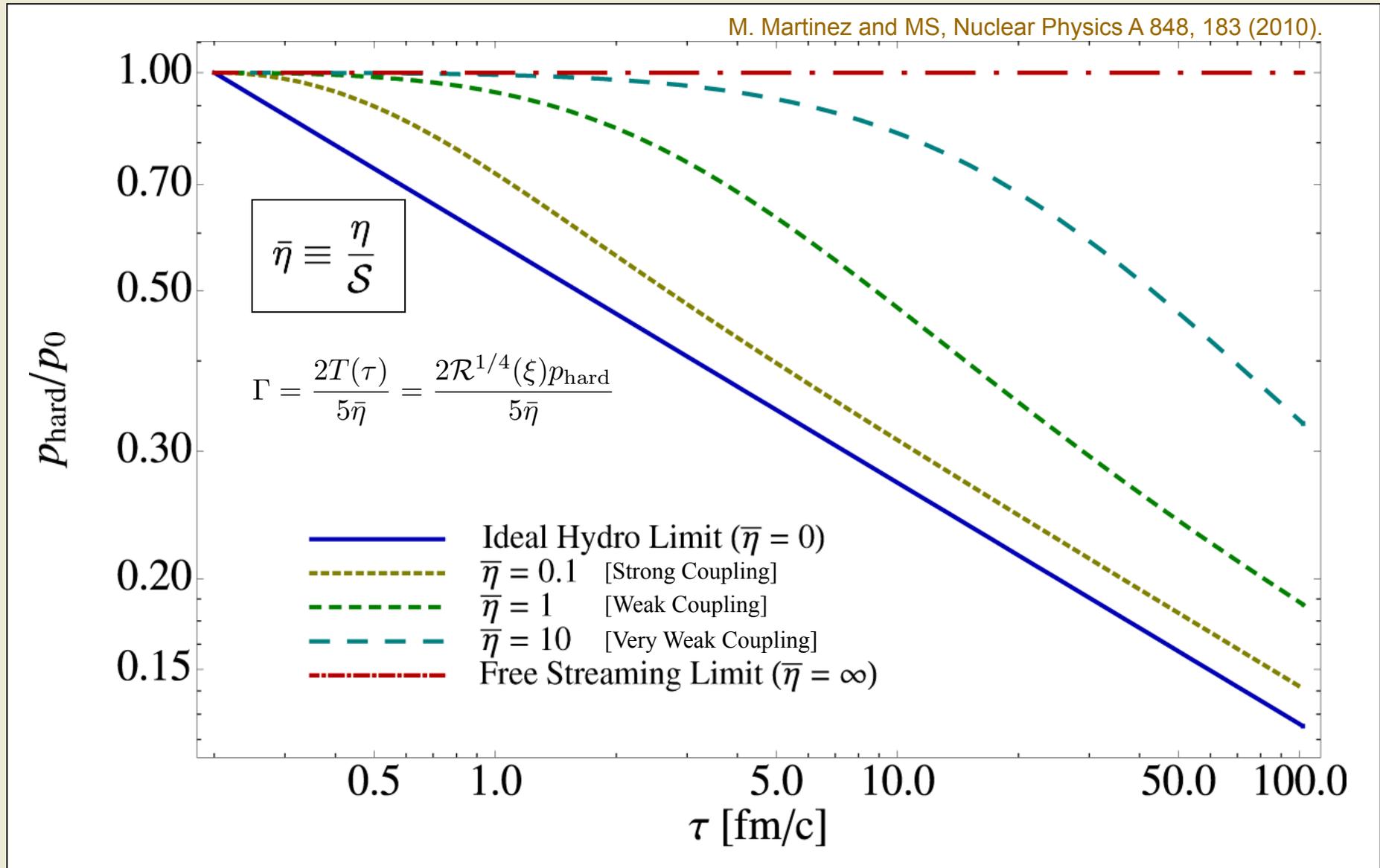
where

$$\mathcal{R}(\xi) \equiv \frac{1}{2} \left( \frac{1}{1+\xi} + \frac{\arctan \sqrt{\xi}}{\sqrt{\xi}} \right) \quad \Gamma = \frac{2T(\tau)}{5\bar{\eta}} = \frac{2\mathcal{R}^{1/4}(\xi)p_{\text{hard}}}{5\bar{\eta}}$$

# Hard Momentum vs Time

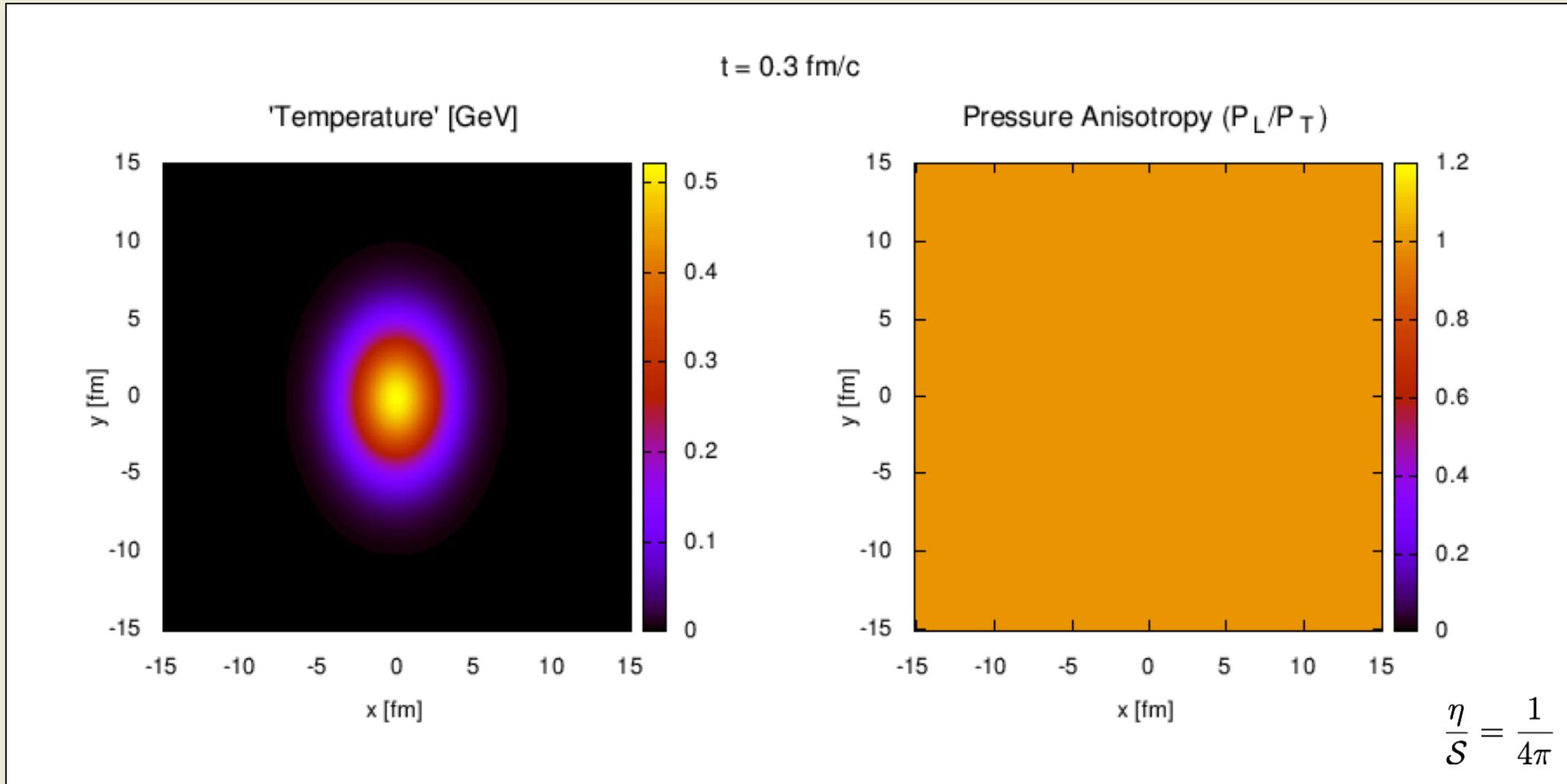
$$\frac{1}{1+\xi} \partial_\tau \xi - \frac{2}{\tau} - \frac{6}{p_{\text{hard}}} \partial_\tau p_{\text{hard}} = 2\Gamma \left[ 1 - \mathcal{R}^{3/4}(\xi) \sqrt{1+\xi} \right]$$

$$\frac{\mathcal{R}'(\xi)}{\mathcal{R}(\xi)} \partial_\tau \xi + \frac{4}{p_{\text{hard}}} \partial_\tau p_{\text{hard}} = \frac{1}{\tau} \left[ \frac{1}{\xi(1+\xi)\mathcal{R}(\xi)} - \frac{1}{\xi} - 1 \right]$$



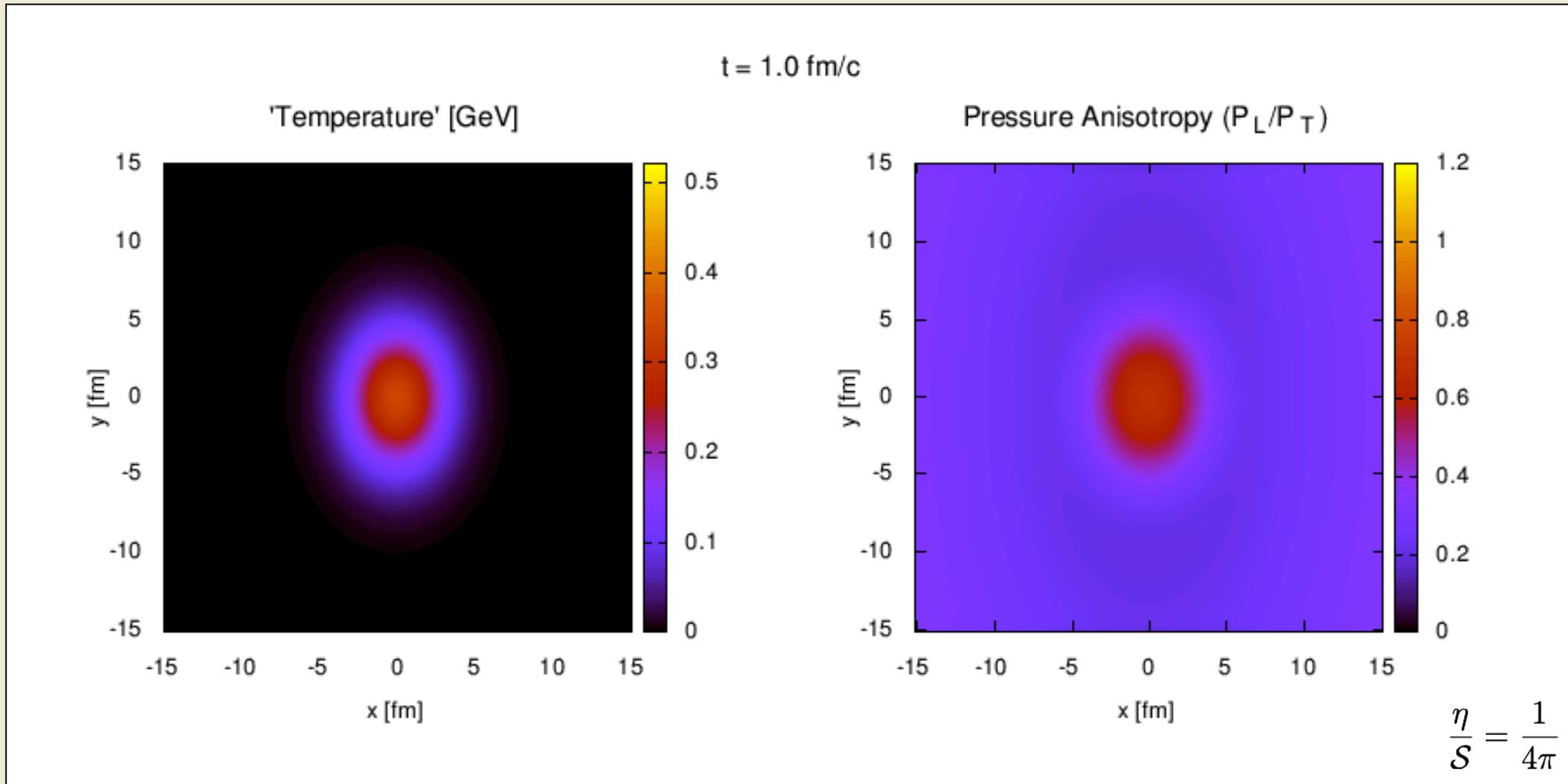
# Including Transverse Dynamics

W. Florkowski, M. Martinez, R. Ryblewski, and MS, forthcoming.



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# Conclusions

- Perhaps fast isotropization isn't required after all?
- Even in 2<sup>nd</sup> order viscous hydro local momentum-space anisotropies can be relatively large and persist for a long time
- New general class of codes which can be applied to highly anisotropic systems (trans/long edges, large shear viscosity) → AHYDRO  
Work in progress ... stay tuned