

Generation of high energy electrons and x-rays under the action of short intense laser pulses: theory and experiment

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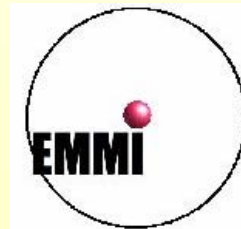


ExtreMe Matter Institute EMMI

Workshop on Plasma Physics
with Intense Laser and Heavy Ion Beams

May 20 - 21, 2010 - Moscow, Russia

Organized by JIHT RAS and EMMI



Outline

- **Introduction**

- *high energy electrons – the origin of secondary sources*

- **X-ray radiation using solid and nano-structured targets**

- *mechanisms of hot electrons production*
- *last experiments at JIHT laser complex*
- *optimization of K_{α} x-ray yield*

- **Laser wakefield electron acceleration**

- *predictions and perspectives*
- *last experimental and theoretical results*

➤ *What can we do with PW lasers...?*

secondary sources ***of high energy particles and radiation***

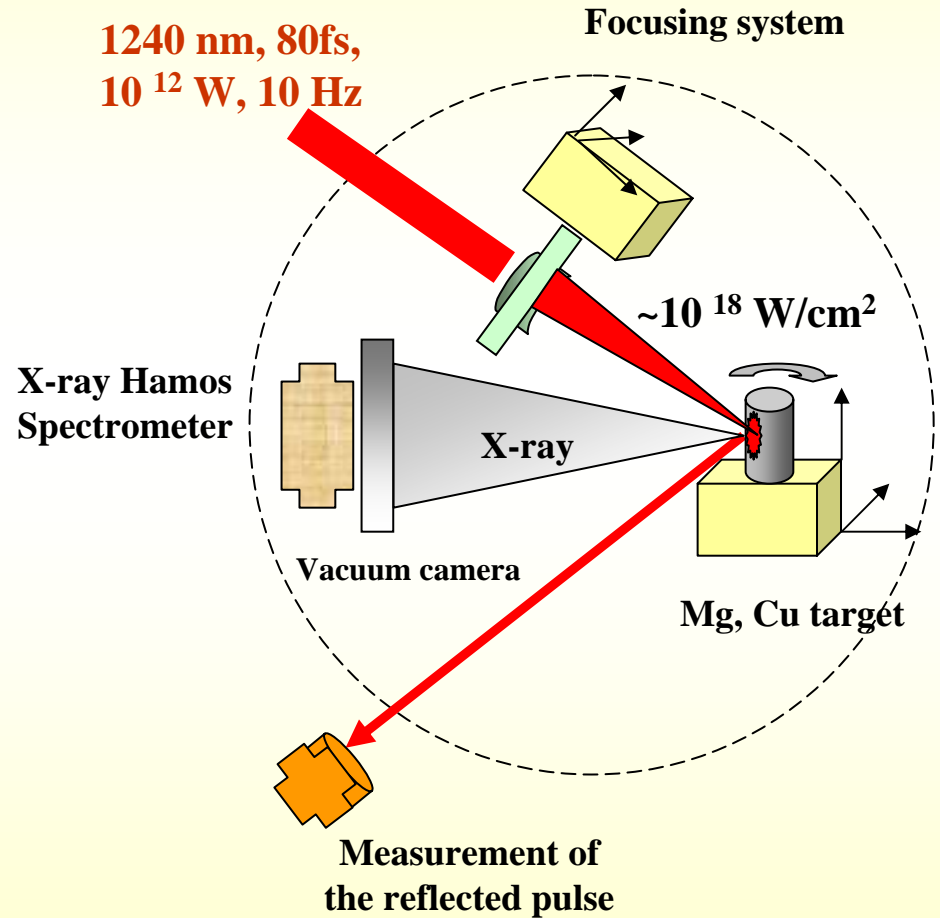
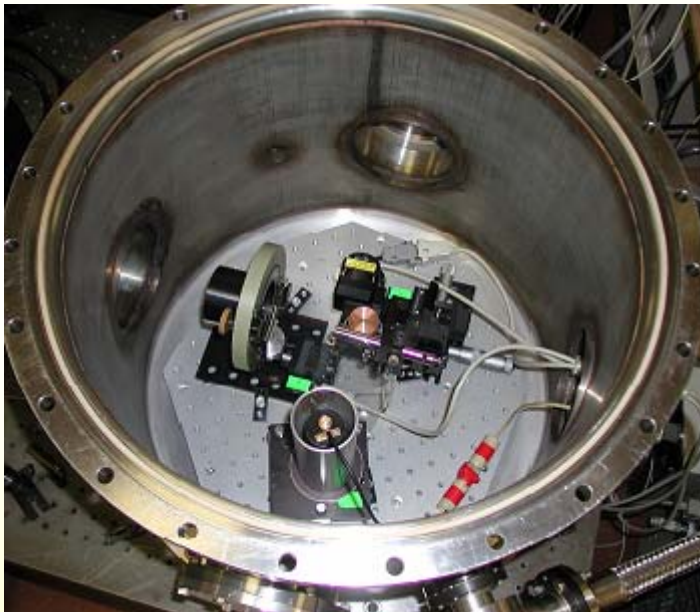
The particular subjects could be:

- production of quasi-monoenergetic and stable source of protons**
(*e.g., with the energy ~ 200 MeV for medical applications*),
- new sources of X-ray radiation using solid and nano-structured targets,**
- laser wake field electron acceleration,**
- powerful sources of THz radiation,**
- laser-driven gamma-ray sources**

➤ new sources of X-ray radiation using solid and nano-structured targets

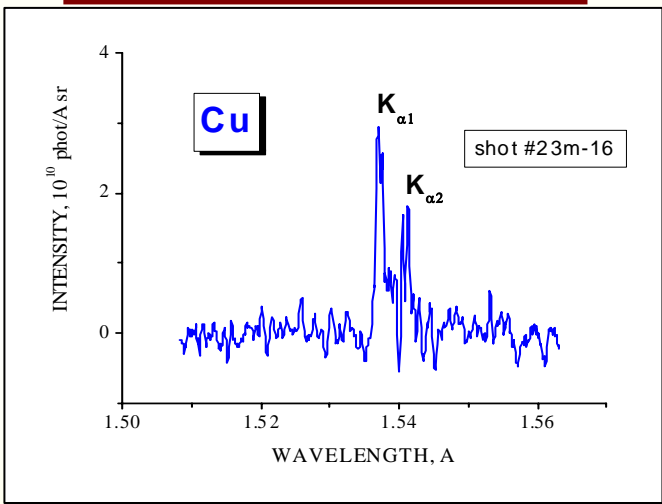


X-ray radiation of laser plasma

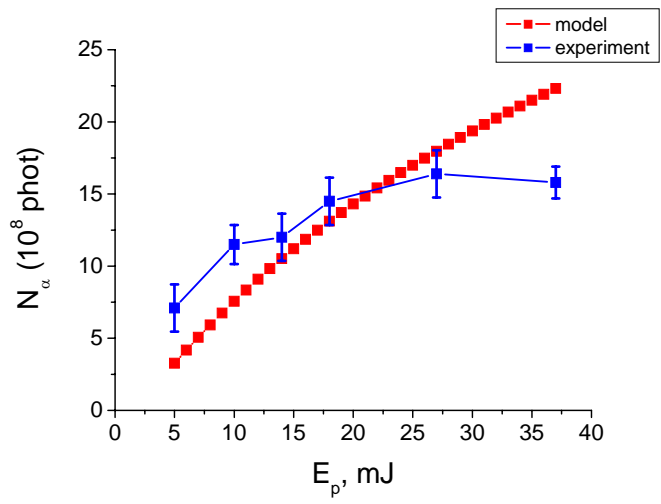


K_α radiation from Cu laser-produced plasmas
 Cr:Forsterite laser system, ω_0 , p -polarization,
 $E_L=30$ mJ, $\tau_L = 80$ fs, $10^{16} \div 10^{17}$ W/cm²

Single laser short experiment



Optimization of K_α x-ray yield



The model calculation of K_α yield against pulse energy in comparison with experimental data received at IHED laser facility:
 $\lambda = 1.24 \mu\text{m}$, $\tau_p = 80$ fs,
 $d_f = 10 \mu\text{m}$, $\theta = 45^\circ$, Fe target

$$E_0 = E_L [1 + (1 - f)^{1/2}] \sin \theta$$

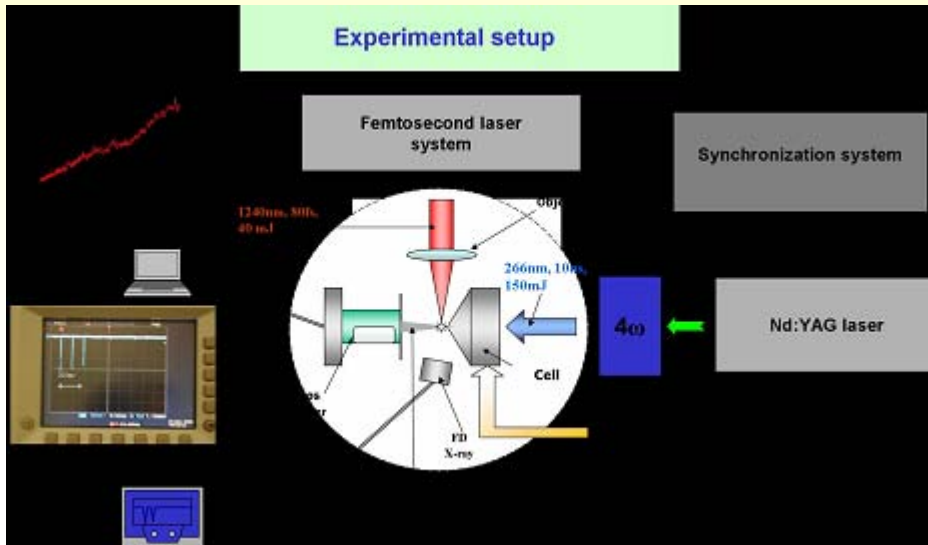
$$T_{h,keV} = 7.6 I_{L,16} \lambda_{\mu\text{m}}^2 \alpha^2 \sin^2 \theta$$

E_0 – driving, E_L – laser field, f – absorption, $\eta = 1.57$, $I_{L,16} = I_L / 10^{16}$ W/cm² – laser intensity

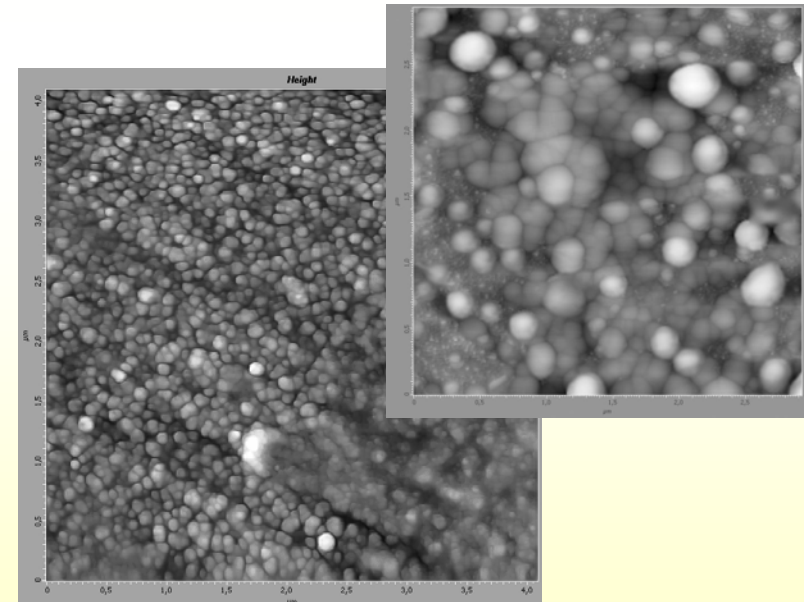
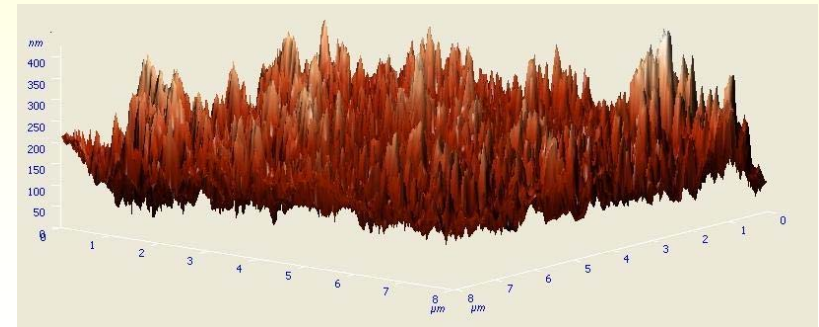
There is a range of laser pulse parameters in which K_α yield may be described by vacuum heating mechanism

O. Kostenko, Friday, May 21st, 15:00

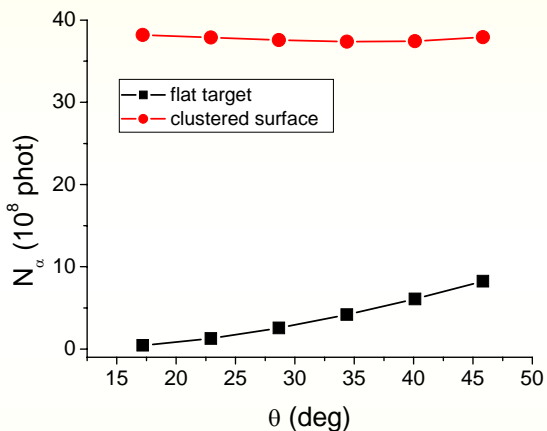
nano-structured targets



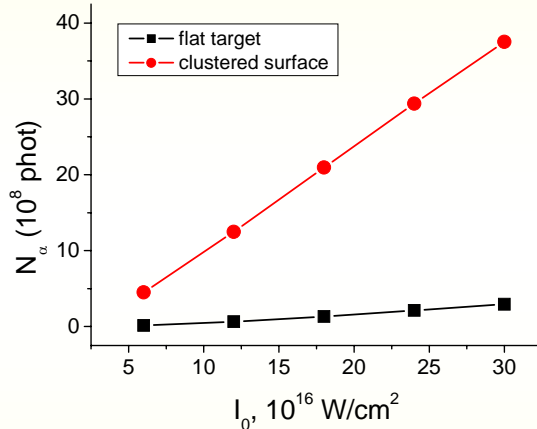
Fe clusters on Cu target



Fe clusters on Cu target



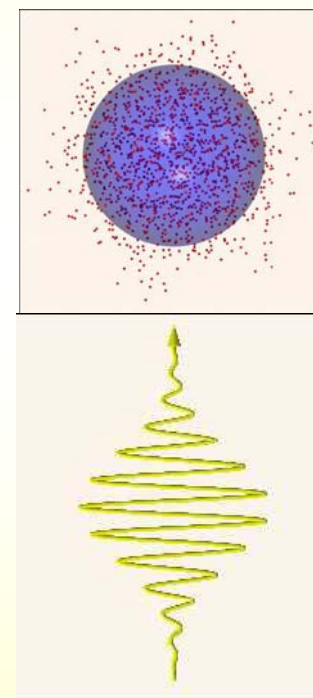
Dependence of K_α yield on angle of incidence for peak intensity $3 \times 10^{17} \text{ W/cm}^2$



Dependence of K_α yield on peak intensity for angle of incidence 30°

$$T_{h,keV} = 68.41 \times I_{L,16} \lambda_{\mu m}^2 \cos^2 \theta_1$$

$$E_0 = 3E_L \cos \theta_1$$

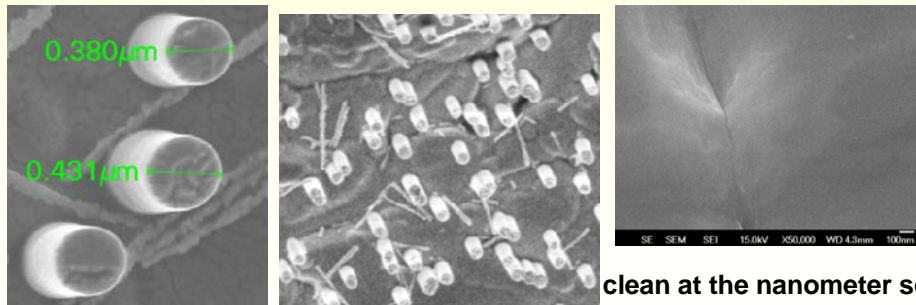


Enhancement of driving electric field at a cluster surface and favorable conditions for K_α photons to escape from the wafer lead to considerable enhancement of K_α yield

First Experiments with high Z- Nanostructures at

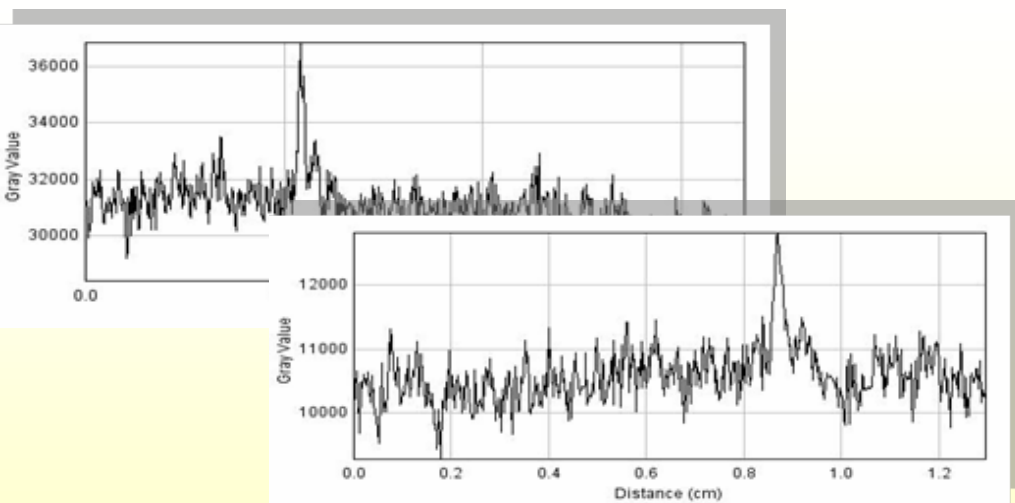
Cu-nano-hairs of 100-500nm, 1-5 μm high on the 8 μm Cu layer

GSI, Material Research; IMP, Lanzhou



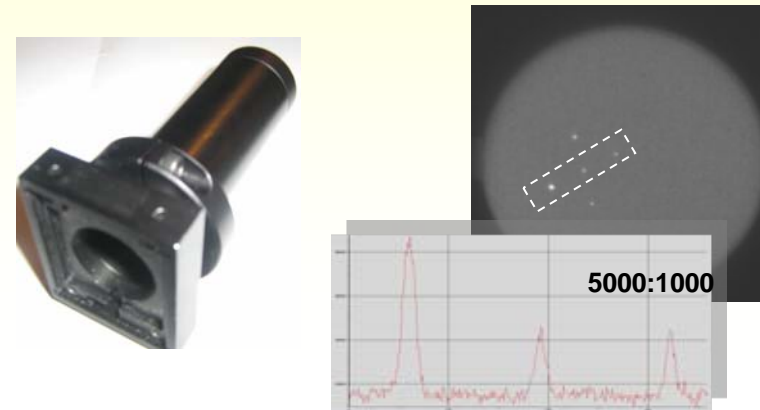
Cu-K α obtained by laser irradiation of the Cu-foil

18.12.08, 20 J, 5°, Cu-foil 33 μm

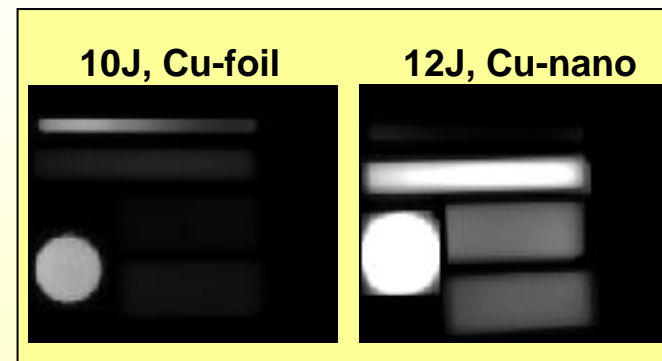


X-ray pin-hole camera images

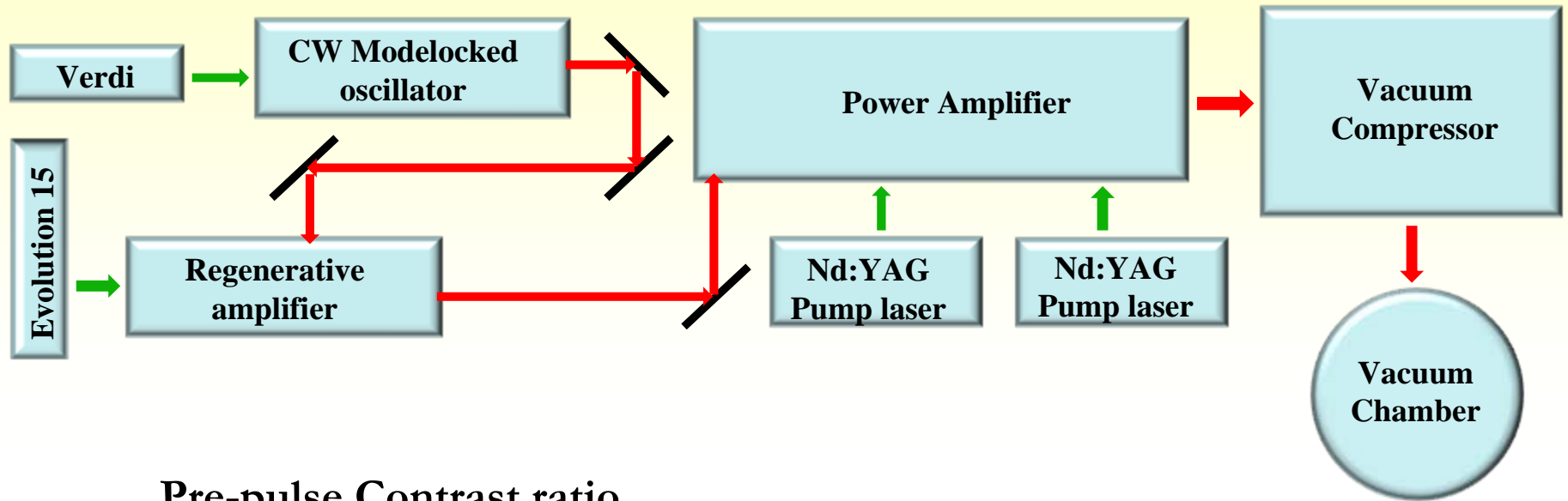
16.12.08 2J, 5°, Cu-foil



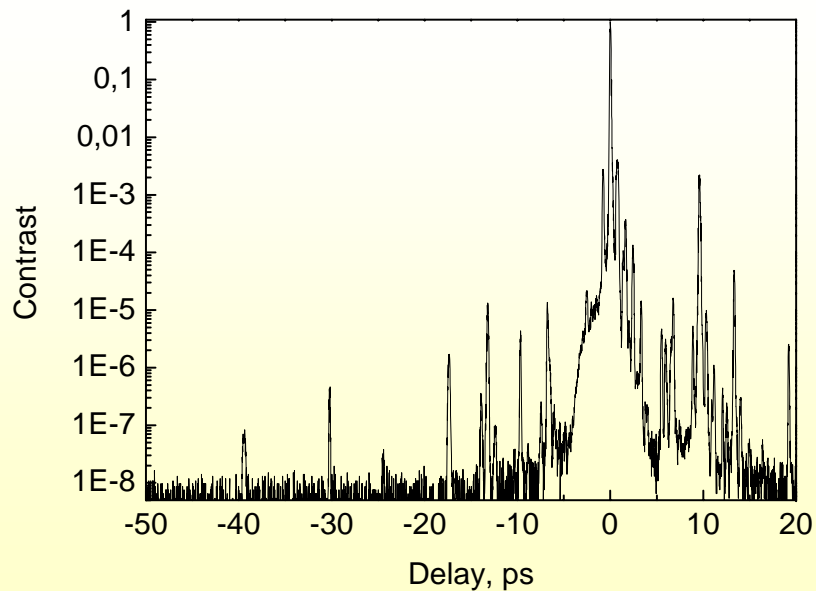
Increased hard x-ray yield by irradiation of nanostructures



Femtosecond Laser System at JIHT



Pre-pulse Contrast ratio



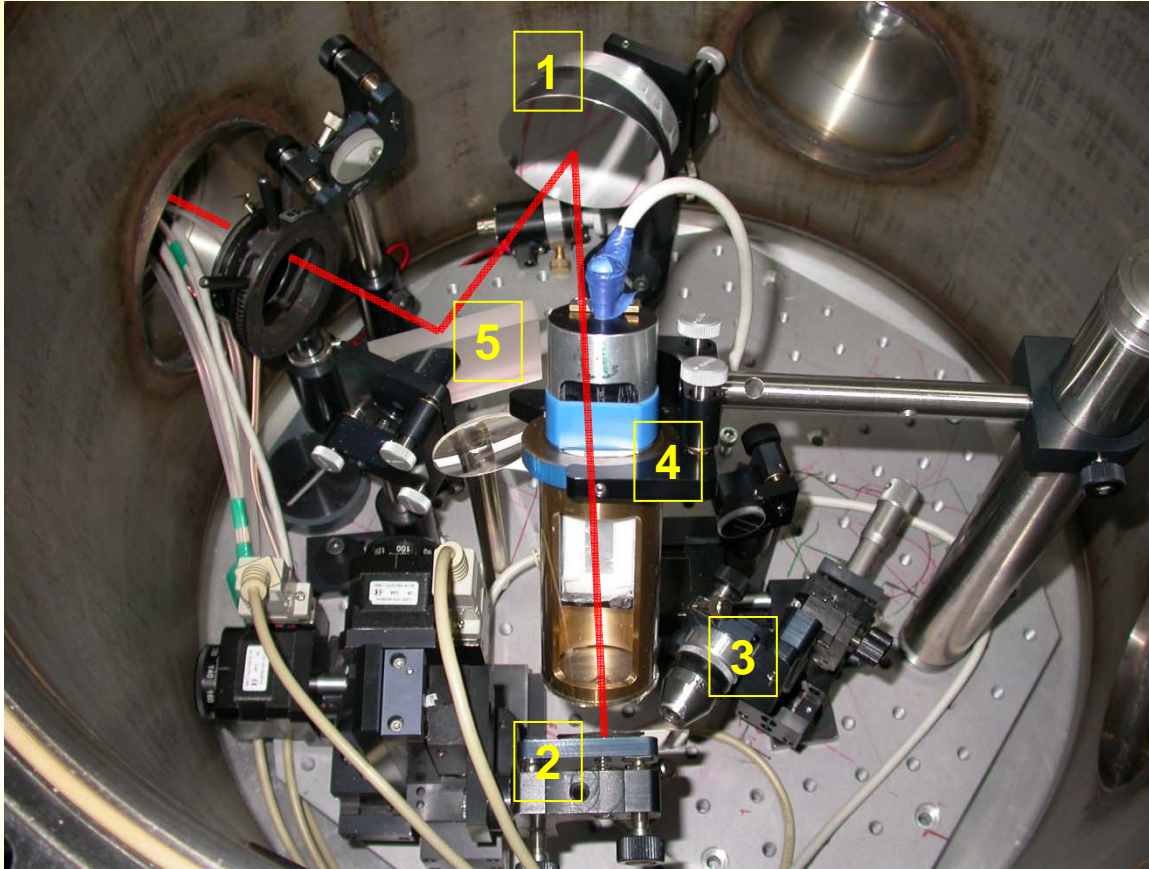
Laser parameters

<i>Center wavelength, nm</i>	~ 800
<i>Pulse duration (FWHM), fs</i>	37 ± 5
<i>Output energy</i>	> 250 mJ
<i>M² Figures</i>	$M_x^2 = 1.5$ $M_y^2 = 1.35$
<i>Bandwidth (nm)</i>	~ 28 nm
<i>Beam diameter (1/e²)</i>	30 mm

10 TW JIHT Femtosecond Laser System



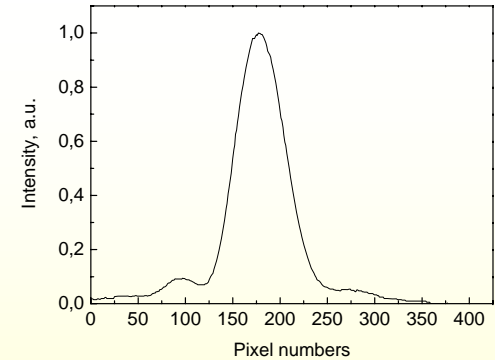
Experimental setup



Focal spot

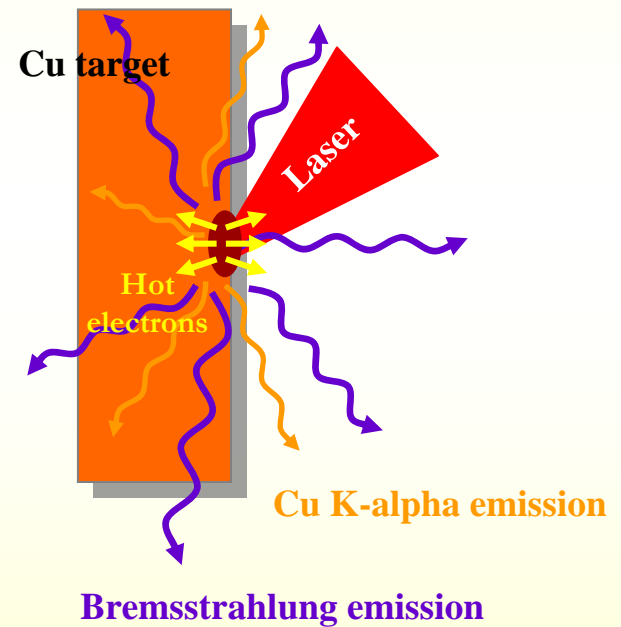
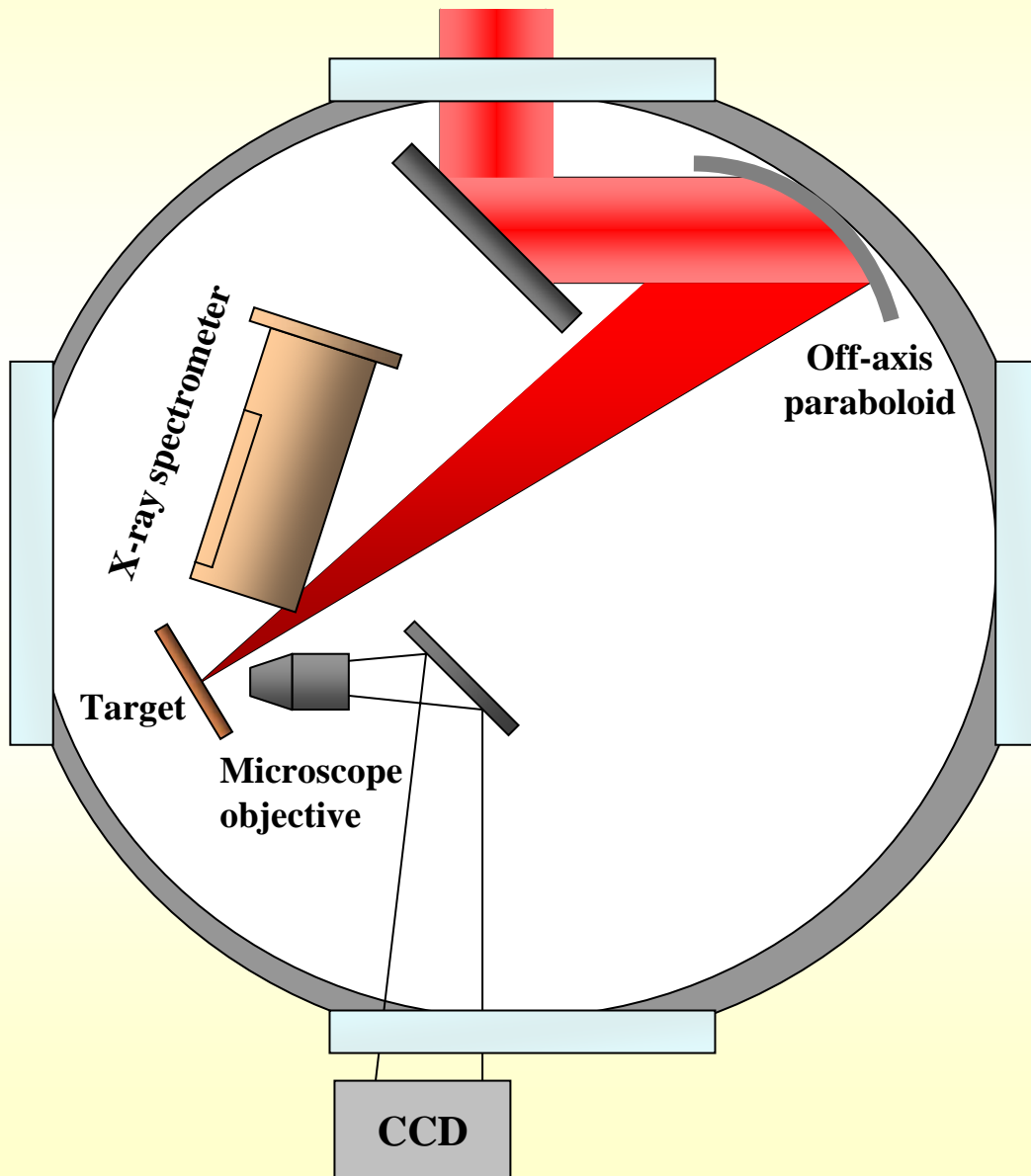


$D_{1/e} = 14\mu\text{m}$



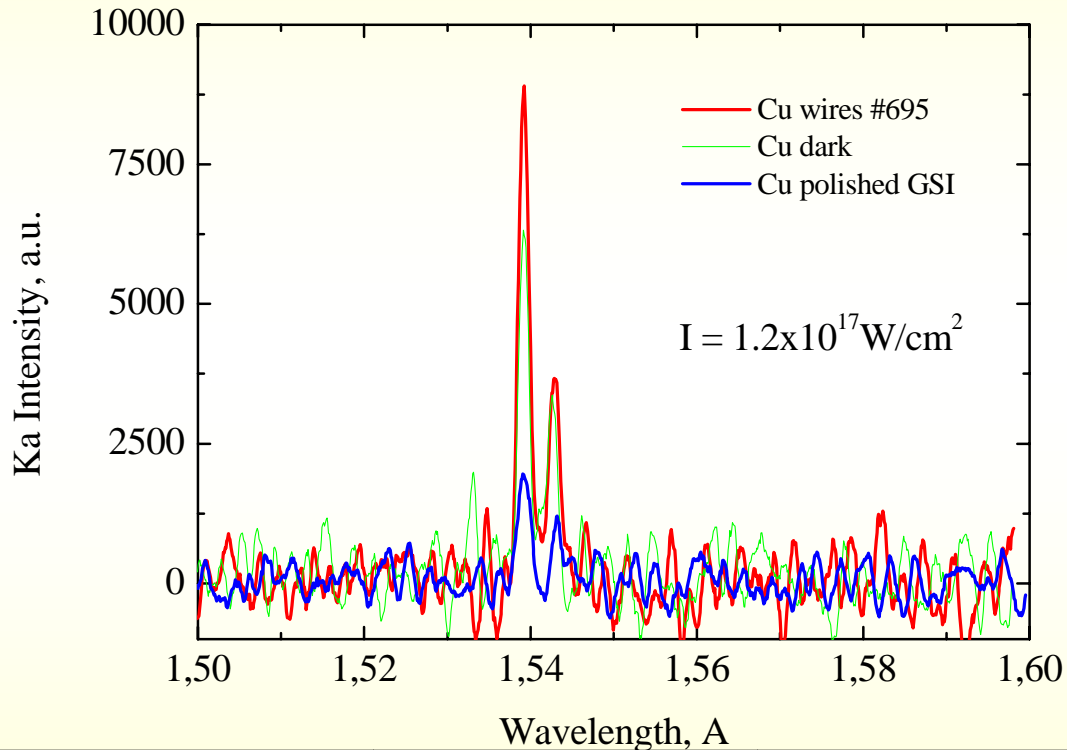
- 1 - Off-axis paraboloid (focal length 254mm).
- 2 - Motorized target unit with target holder.
- 3 - System for interactive control of focus spot.
- 4 - Von Hamos spectrometer
- 5 - Mirror (R=100%)

Experimental scheme



X-ray radiation using solid and nano-structured targets

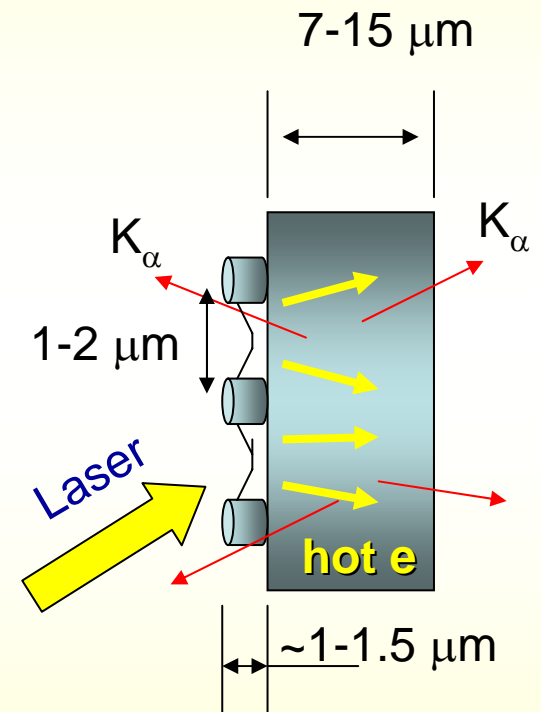
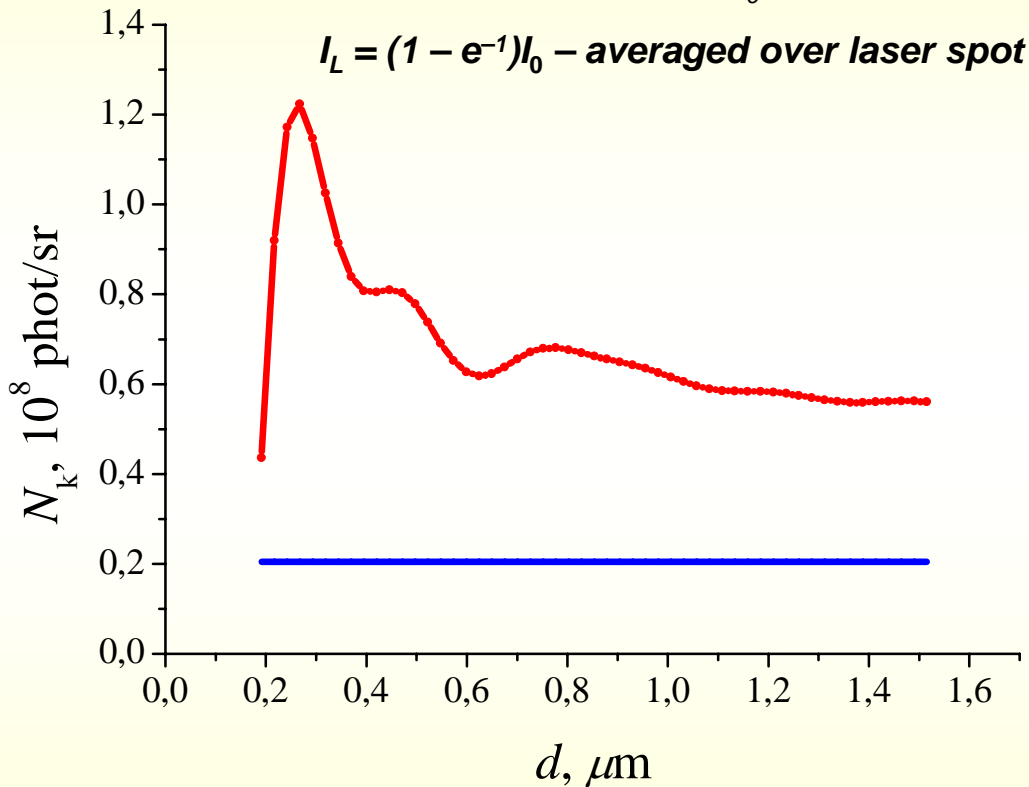
$I_0 = 1.2 \times 10^{17} \text{ W cm}^{-2}$ (peak intensity), angle of incidence $\theta = 30^\circ$,
 laser wavelength $0.8 \mu\text{m}$, pulse duration 37 fs, P-polarization
 focal spot diameter $2r_0 = 14 \mu\text{m}$



Target	X-Ray integral, a.u.	Abs yield of K_{α} , ph/(ster pulse)
Cu polished foil (GSI), 8um	5.98×10^4	8.7×10^7
Cu dark foil (GSI), 12.5 um	2.12×10^5	3.1×10^8
Cu foil with wires №695 (GSI)	2.26×10^5	3.3×10^8

K_α radiation using solid and nano-structured targets

$I_0 = 1.2 \times 10^{17} \text{ W cm}^{-2}$ (peak intensity), angle of incidence $\theta = 30^\circ$,
 laser wavelength $0.8 \mu\text{m}$, pulse duration 37 fs, P-polarization
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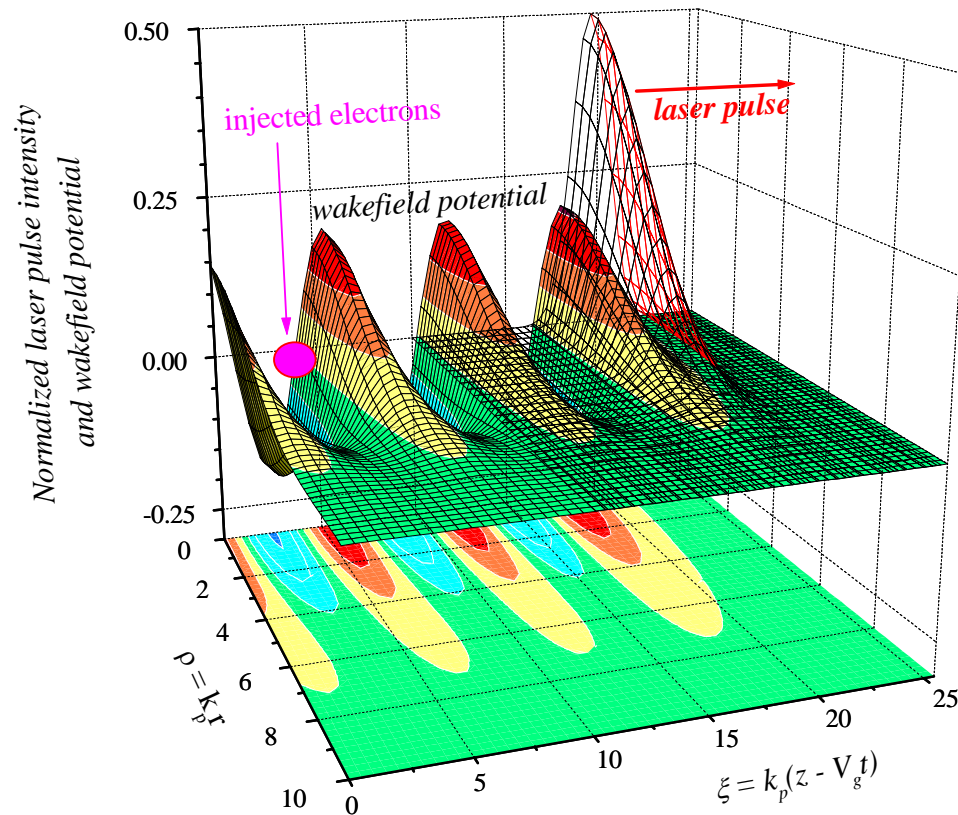
@ $d = 0.5 \mu\text{m}$ *theoretical ratio nano wires / polished* = 3.8

experimental ratio = $3.3 \times 10^8 / 8.7 \times 10^7 = 3.8$

➤ laser wake field electron acceleration



Laser-Plasma Acceleration of electrons with record gradients ~ 10 GV/m



The result of optimization of laser and plasma parameters
for the monoenergetic acceleration of a short electron bunch in the wakefield
generated by the 50 TW, 100 fs laser pulse in the plasma channel



Expected potential of laser – plasma acceleration of electrons

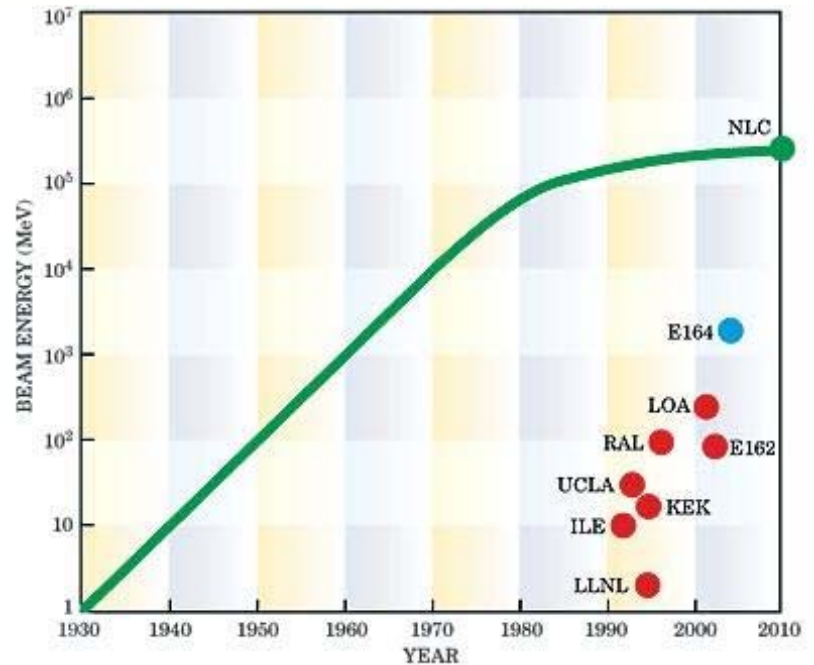
Electric field of plasma wave (with phase velocity $\sim c$, $\lambda_p = 2\pi c/\omega_p$):

$$E_p \text{ [V/m]} \approx 10^2 \alpha (n_e \text{ [cm}^3\text{]})^{1/2} \propto \gamma_g^{-1} = \omega_p / \omega_0$$

$\alpha = \delta n / n_0$ – plasma wave amplitude; at $\alpha = 0.3 \div 1.0$, $n_e = 10^{17} \div 10^{18} \text{ cm}^{-3}$:
 $E_p = 10 \div 100 \text{ GV/m}$

maximum of accelerating gradient
in traditional accelerators (RF linac):
 $E_{RF} \sim 10 - 100 \text{ MV/m}$

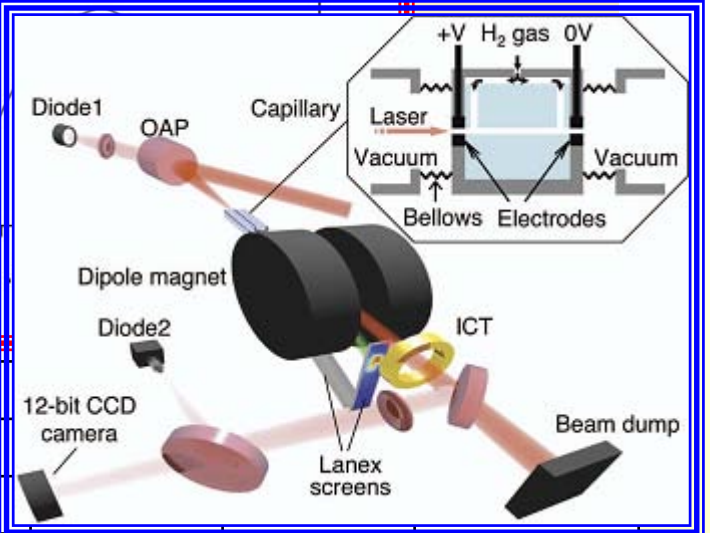
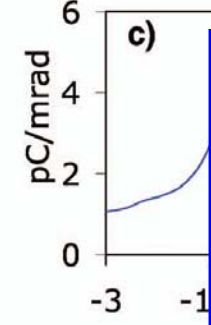
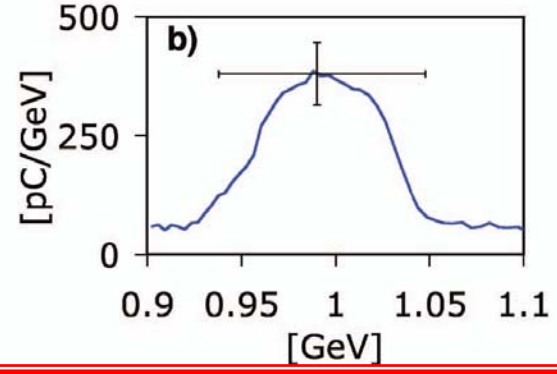
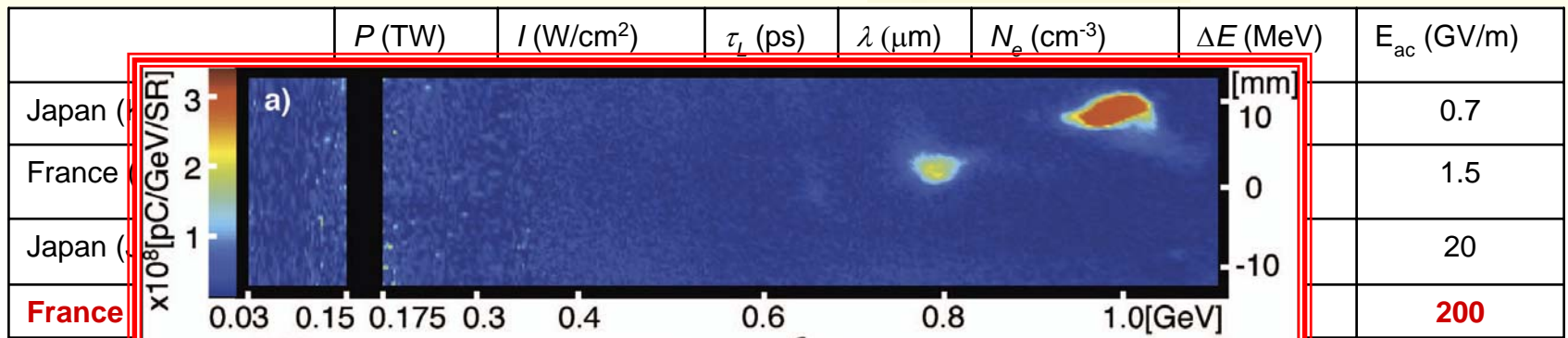
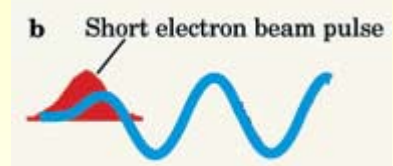
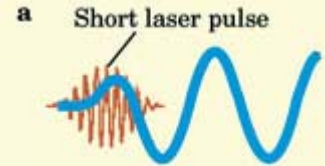
Exponential growth of “the Livingston curve” began tapering off around 1980



Parameters and results of some experiments

for standard LWFA scheme

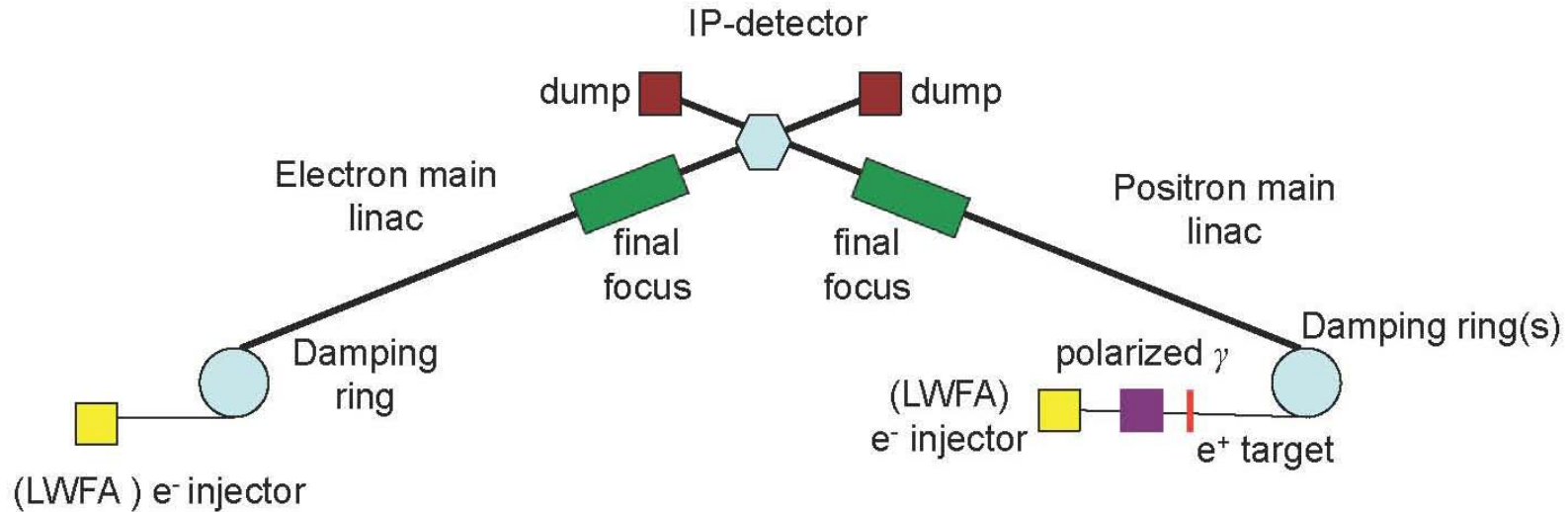
$$c\tau_L \approx \lambda_p/2$$



for Self-N

	P (TW)	I (W/cm ²)	τ_L (ps)	λ (μ m)	N_e (cm ⁻³)	ΔE (MeV)	E_{ac} (GV/m)
Japan (KEK)	3	10^{17}	1.0	1.05	10^{19}	17	30
USA (NRL)	2.5	4×10^{18}	0.4	1.05			
USA (CUOS)	4.5	10^{18}	0.4	1.05			
France (CEA)	5	10^{19}	0.3	1.05			
UK (RAL)							
France							200
Japan (KEK)							20
France							1.5
Japan (KEK)							0.7

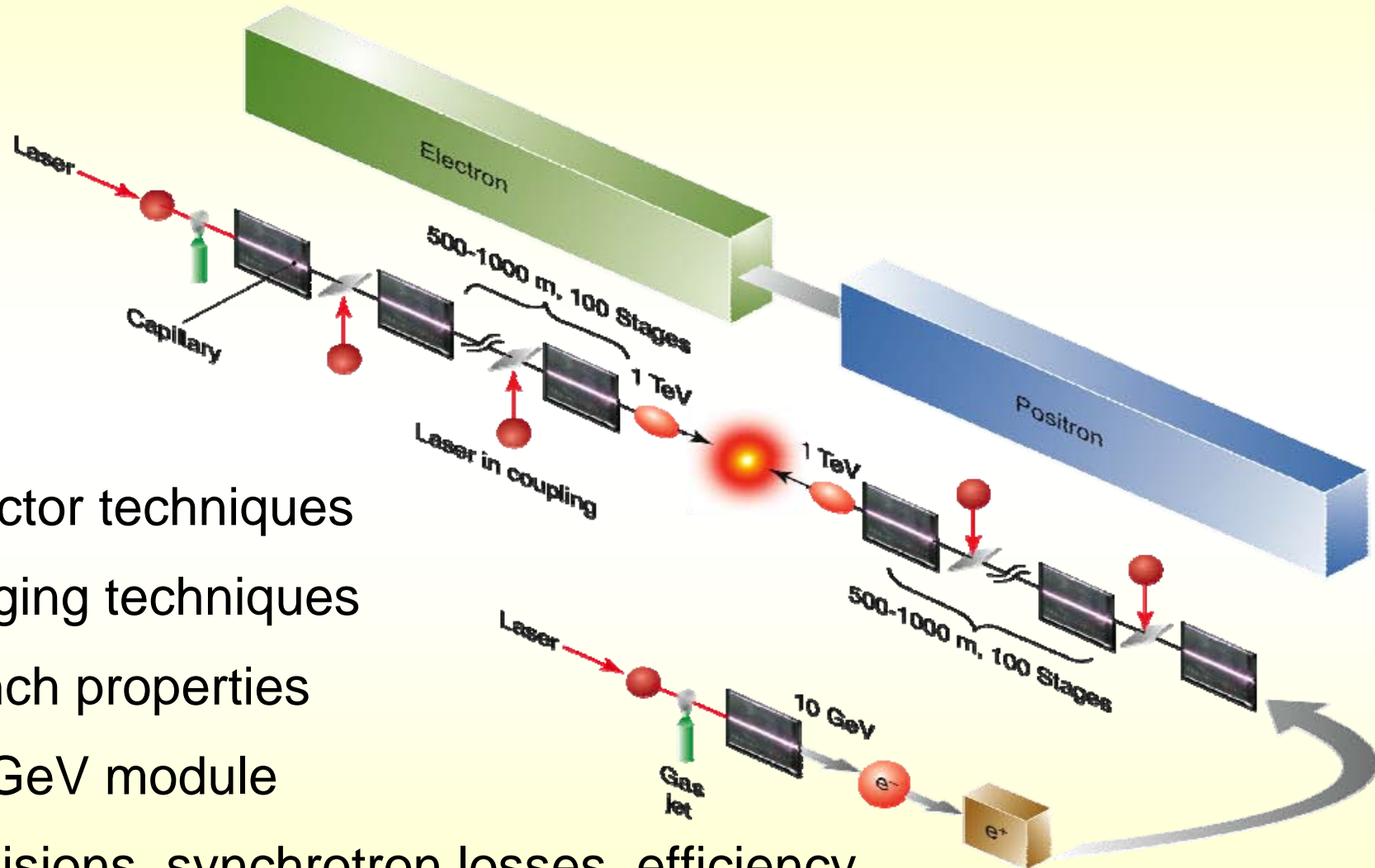
Electron-Positron Linear Collider



Conventional technology:

- Current generation of future linear collider designs based on existing technology (e.g., ILC): $E_{cm} \sim 0.5$ TeV; gradient ~ 0.03 GV/m; ~ 30 km (\sim multi- $\$B$).
- Higher energy collider with existing technology: 5 TeV \rightarrow >100 km, $>$ tens of $\$B$

Laser plasma accelerator based concept for a Laser Plasma Linear Collider



- Injector techniques
- Staging techniques
- Bunch properties
- 10 GeV module
- Collisions, synchrotron losses, efficiency

Main Processes Accompanying the Propagation of Short Intense Laser Pulses in Gases

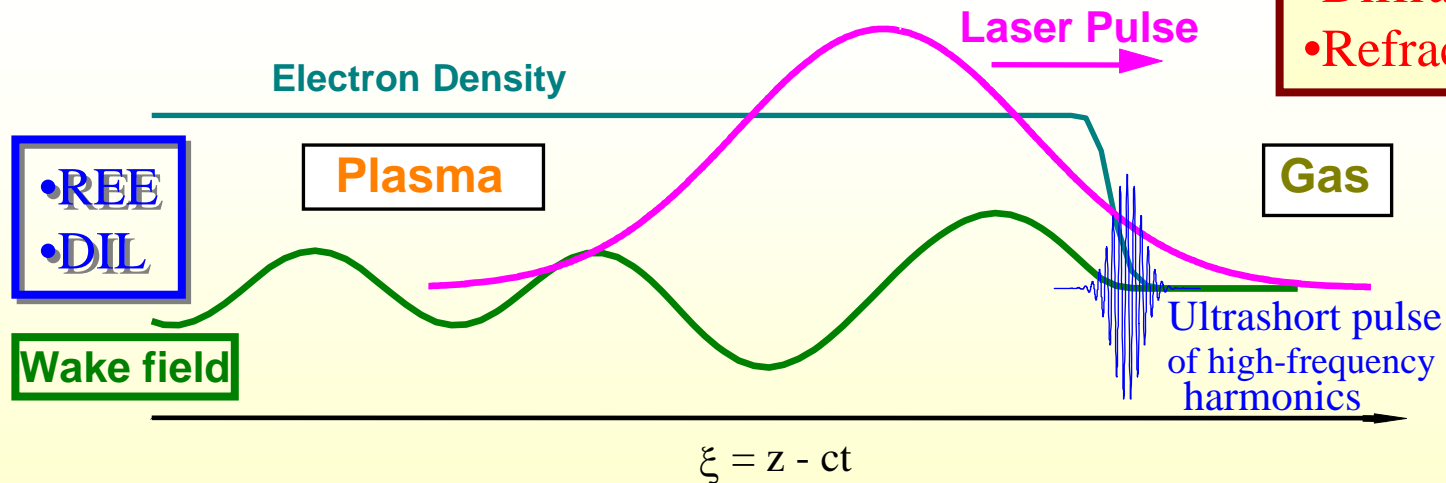
Typical parameters

Laser pulse:
 $I_{\max} > 10^{17} \text{ W/cm}^2$, FWHM < 1 ps, $\lambda \approx 1 \text{ mkm}$

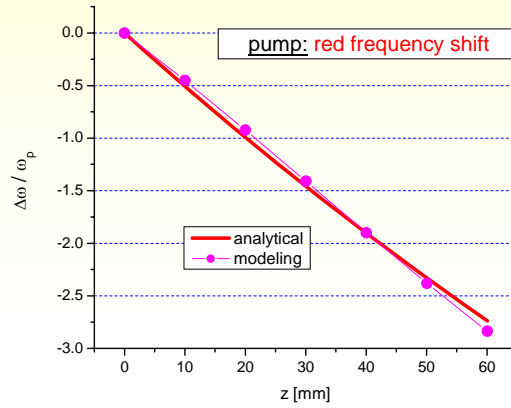
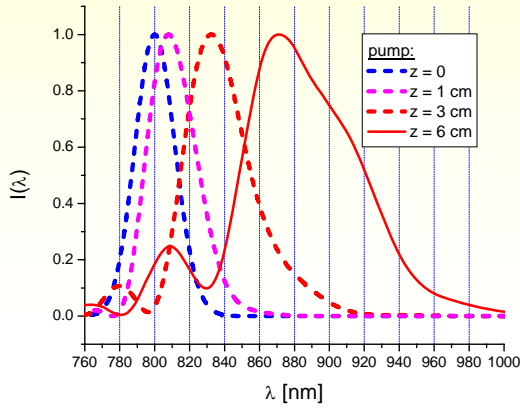
Gas:
 comparatively light atoms $Z \approx 1$, $n_{\text{at}} < 10^{19} \text{ cm}^{-3}$

Laser Pulse Distortions:

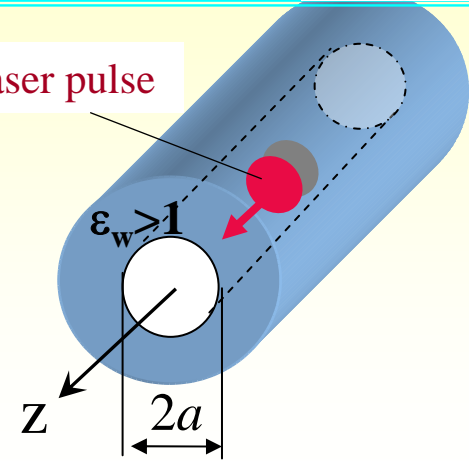
- LASRS
- RS-F & S-M
- Filamentation
- Diffraction
- Refraction



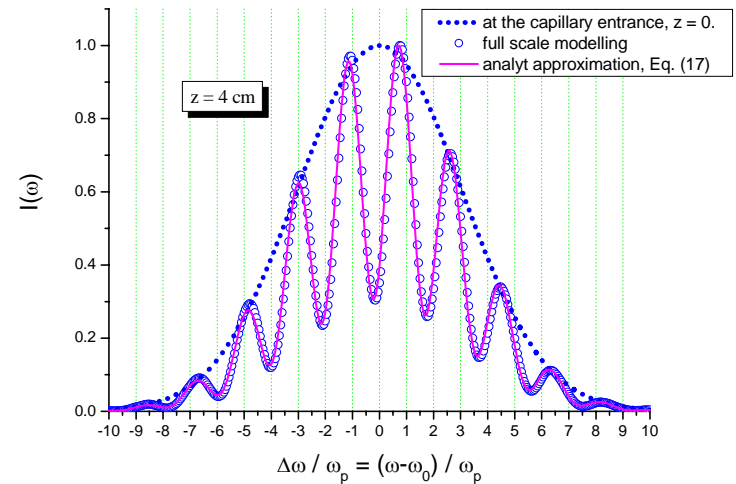
Frequency Shift of the Pump Pulse



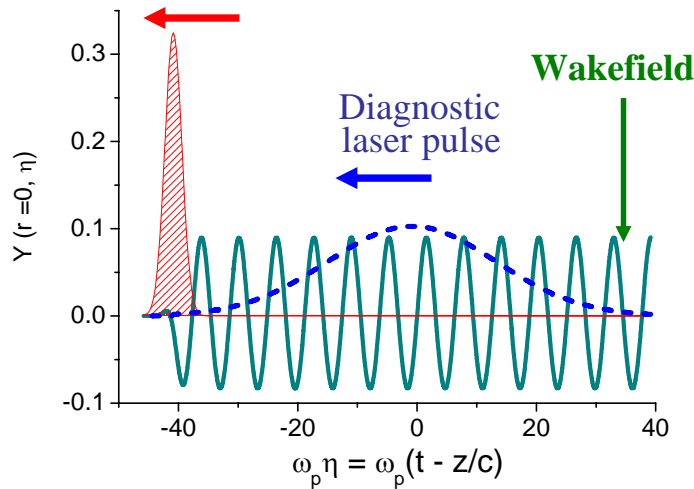
Laser pulse



Spectrum Modulation of the Probe Pulse



Short intense laser pulse



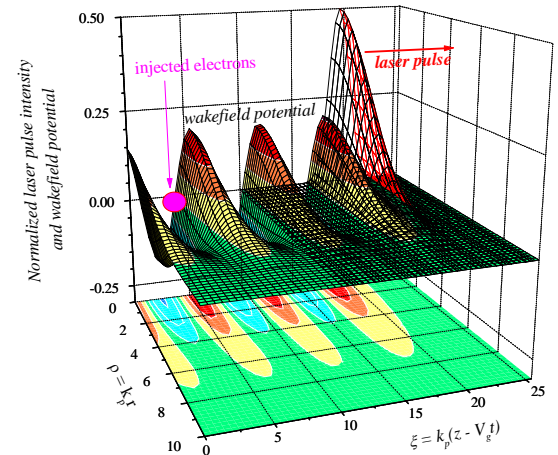
$$\bar{A}(z) = C_r \frac{k_0}{k_p^2 z} \left[\sin \left(\frac{1}{4} \frac{\Omega^2 \beta}{1 + \beta^2} \right) \right]^{-1} [\delta \bar{I}(z) - \delta \bar{I}(z = 0)]$$

Spectroscopic Diagnostics of the Plasma Wakefield

$$\langle \omega^2 \rangle_{out} - \omega_0^2 = \frac{1}{\epsilon_{out}} \left\{ \int_V d^3\mathbf{r} \int_{-\infty}^{+\infty} dt \left[\omega_0^2 \mathbf{E} \cdot \mathbf{j} - \frac{\partial \mathbf{E}}{\partial t} \cdot \frac{\partial \mathbf{j}}{\partial t} \right] + \frac{1}{8\pi} \int_V d^3\mathbf{r} \left[\omega_0^2 (\mathbf{E}^2 + \mathbf{B}^2) - \left(\left(\frac{\partial \mathbf{E}}{\partial t} \right)^2 + \left(\frac{\partial \mathbf{B}}{\partial t} \right)^2 \right) \right] \right\}_{t=-\infty}^{t=+\infty}$$

$$\langle \omega^2 \rangle_{\alpha} = \oint_S \int_0^{+\infty} \omega^2 I_{\alpha}(\omega, \mathbf{R}) d\omega n d\mathbf{s} \left[\oint_S \int_0^{+\infty} I_{\alpha}(\omega, \mathbf{R}) d\omega n d\mathbf{s} \right]^{-1}$$

$\delta \bar{n}_e$ – Wakefield Amplitude
 $I(\mathbf{r}, t)$ – Laser Pulse Intensity



$$\langle \omega^2 \rangle_{out} - \langle \omega^2 \rangle_{in} = -\frac{6\pi e^2}{m_e c \epsilon_{out}} \int_V \int_{-\infty}^{+\infty} \delta \bar{n}_e \frac{\partial I(\mathbf{r}, t)}{\partial t} dt d^3\mathbf{r} + \frac{\omega_0^2}{8\pi \epsilon_{out}} \int_V \left(1 - \frac{n_0(\mathbf{r})}{n_c} \right) \mathbf{E}_{p, \max}^2(\mathbf{r}) d\mathbf{r}$$

$$\langle \omega^2 \rangle_{out} - \omega_0^2 = -\frac{\omega_0^2}{4\pi \epsilon_{out}} \int_V \mathbf{E}_{p, \max}^2 d^3\mathbf{r}$$

$$\mathbf{E}_{p, \max}^2 = \frac{m^2 c^4}{e^2} \frac{\omega_p^2}{16} \left\{ k_p^2 \left| \int_{-\infty}^{\infty} dt e^{i\omega_p t} |\vec{a}|^2 \right|^2 + \left| \frac{\partial}{\partial r} \int_{-\infty}^{\infty} dt e^{i\omega_p t} |\vec{a}|^2 \right|^2 \right\}$$

Spectroscopic Diagnostics of the Plasma Wakefield ...

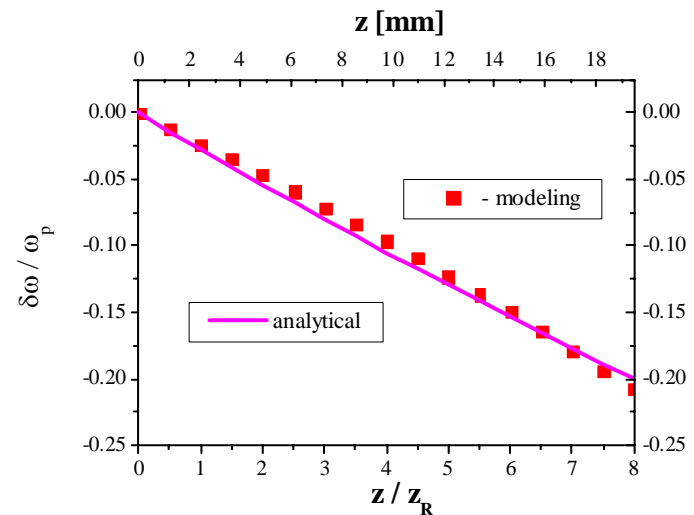
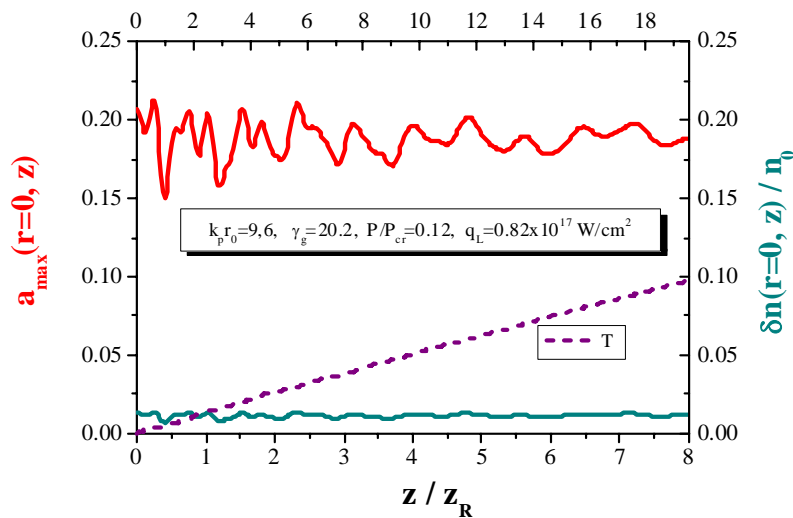
Nonlinear laser pulse propagation in a gas-filled capillary

$$I(t, z, \mathbf{r}_\perp) = I_{\max} \exp\left(-2\frac{(t - z/c)^2}{\tau^2}\right) f(r)$$

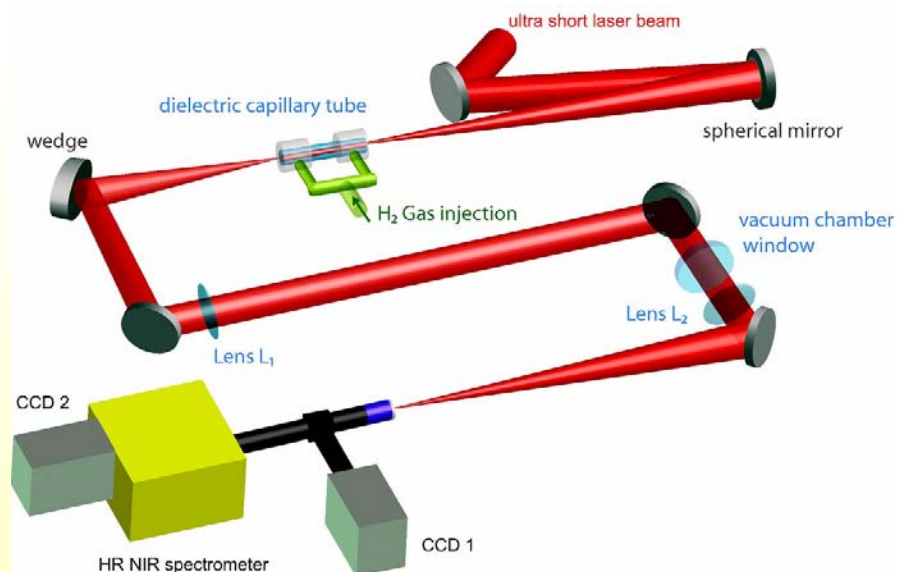
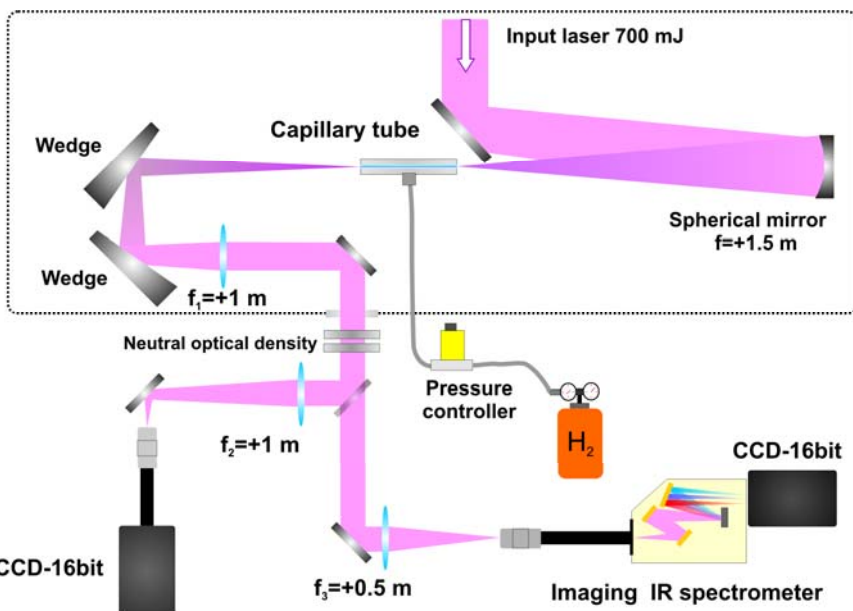
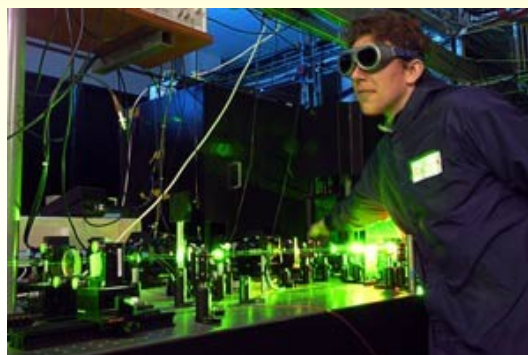
$$f(r) \cong J_0^2\left(\frac{u_1 r}{R_{\text{cap}}}\right), \quad J_0(u_1) = 0$$

$$\frac{\langle \omega^2 \rangle_{\text{out}} - \omega_0^2}{2\omega_0\omega_p} = \frac{\delta\omega}{\omega_p} = -\frac{1}{64} \left(\frac{\pi}{2}\right)^{1/2} (1.13k_p^2 r_0^2 + 8.72) a_0^2 \Omega \exp\left(-\frac{\Omega^2}{4}\right) z/z_R, \quad \Omega = \omega_p \tau$$

Hydrogen: $N_a = 4.1 \times 10^{18} \text{cm}^{-3}$, $P_L = 0.8 \text{ TW}$, $k_p r_0 = 9.5$, $D_{\text{cap}} = 75 \text{ mkm}$, $P_L/P_{cr} = 0.12$

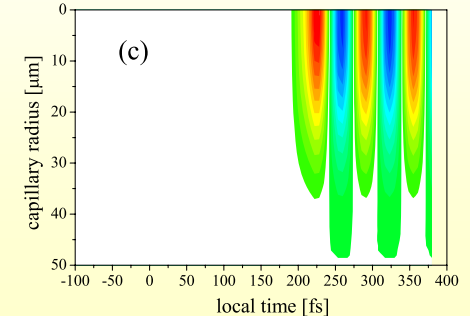
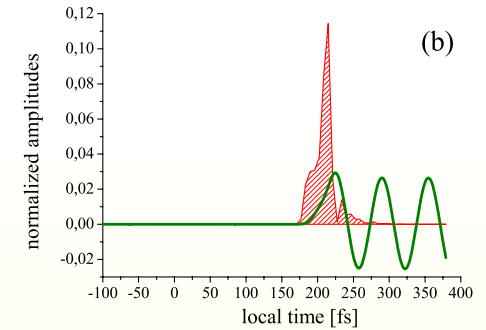
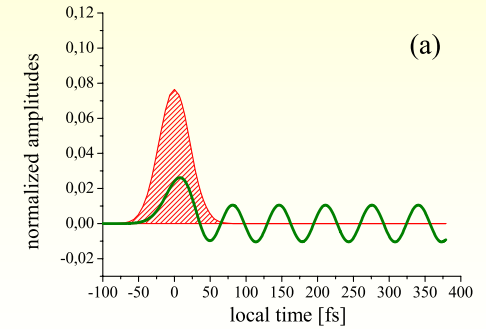
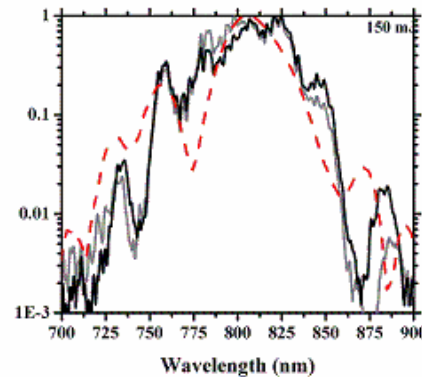
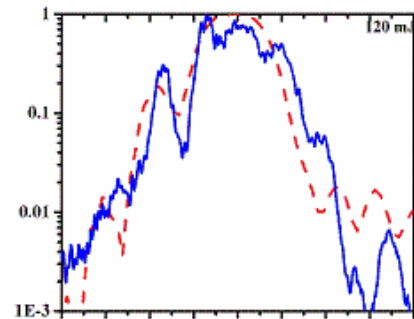
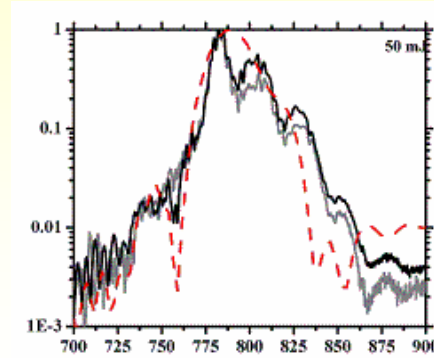
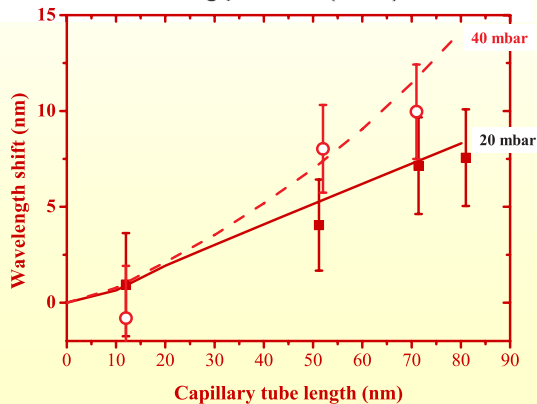
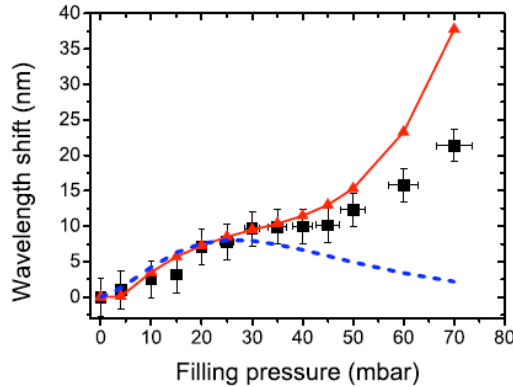
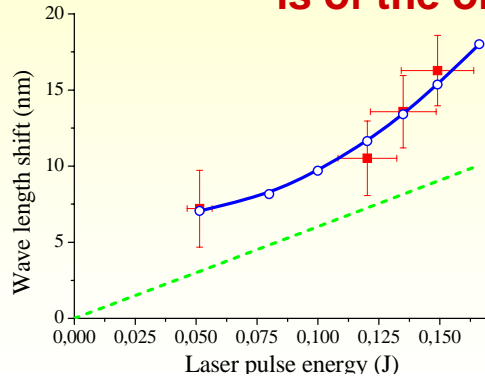


Laser plasma electron acceleration experiments - 2009

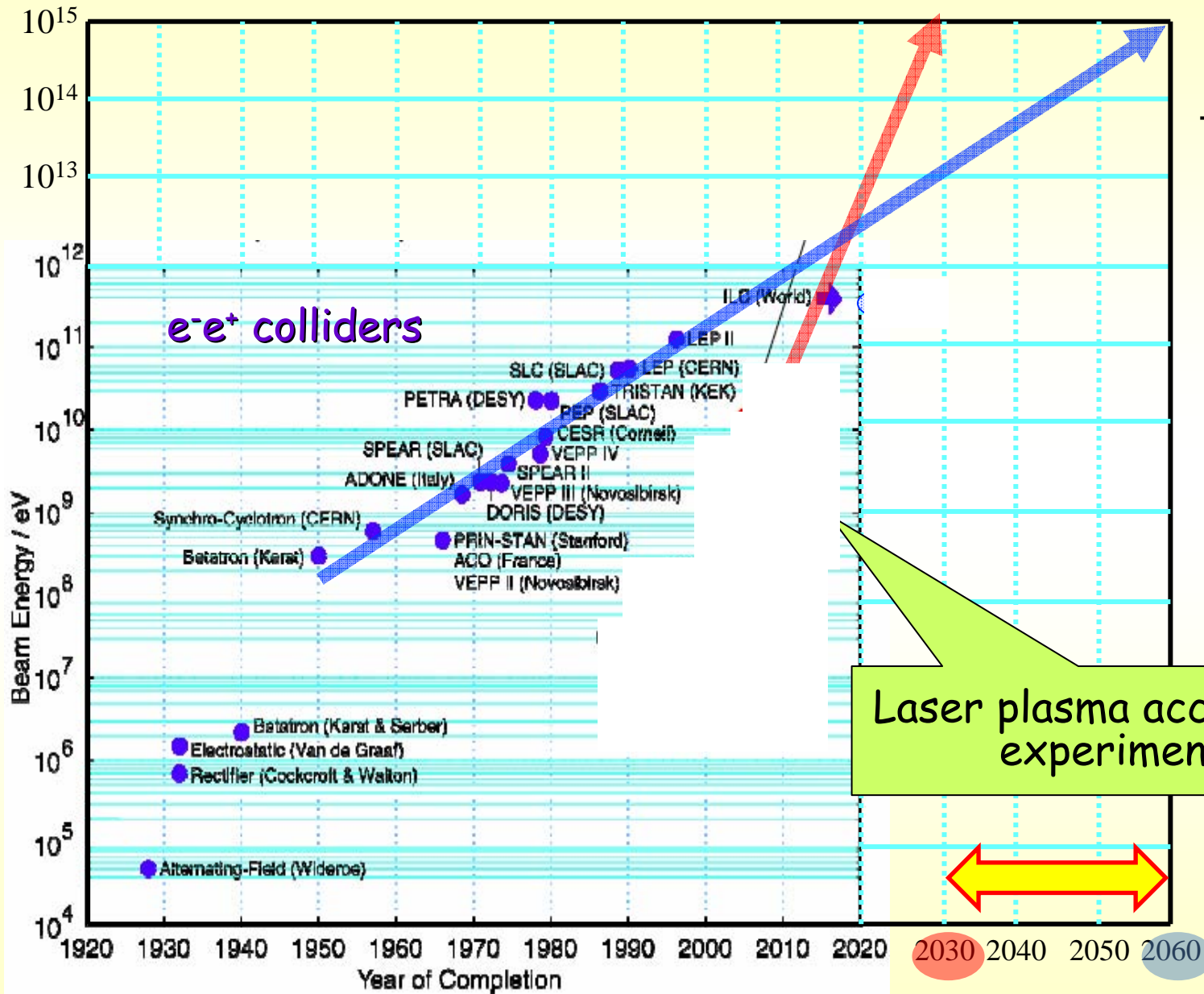


Spectral diagnostics of the laser wake fields in capillary tubes

The average product of gradient and length achieved in this experiment is of the order of 0.4 GV at a pressure of 50 mbar



When can we reach 1 PeV?: Suzuki Challenge(2)



(Suzuki,2009)

Toshi Tajima



ELI Grand Challenges

Paris, April
27th-28th 2009

ExtreMe Matter Institute EMMI

Workshop on Plasma Physics
with Intense Laser and Heavy Ion Beams

May 20 - 21, 2010 - Moscow, Russia



Organized by Joint Institute for High Temperatures RAS, Russia, and
ExtreMe Matter Institute EMMI, Germany

Conference Chairmen

Vladimir E. Fortov, JIHT RAS, Moscow
Thomas Stöhlker, GSI, Darmstadt

Main Topics

Physics of laser-matter interaction, experiments, and theory
Interaction of heavy ions with matter: experiments and theory
Scientific activities and perspectives at VNIIEAN laser facilities LULSE
Scientific activities and perspectives at TWAC/LS18 and perspectives for FAIR
Target development for laser and heavy ion experiments

Thank

You for attention!

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