

High-Definition X-Ray Polarimetry

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Workshop - EMMI

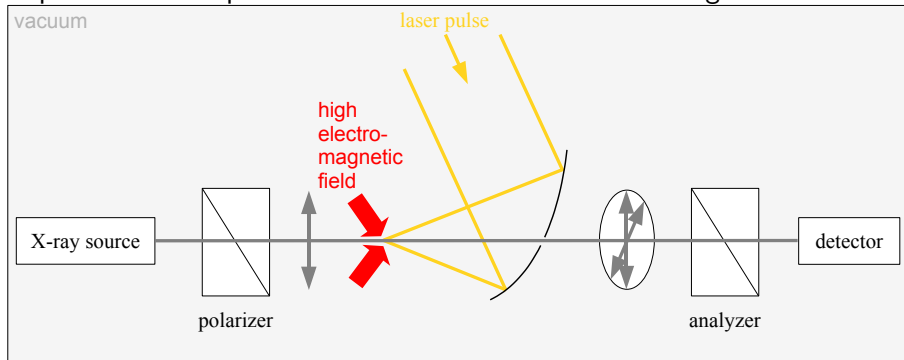
I.Uschmann, S. Höfer, R. Loetzsch, O. Wehrhan, E. Förster, M. Kaluza, G.G. Paulus, C. Detlefs,
T. Roth, J. Härtwig

Outline

- 1 Motivation
- 2 Principles
 - Kinematical theory of diffraction
 - Dynamical theory of diffraction
 - Polarimetry
- 3 Measurements
 - Polarization measurements at the ESRF Grenoble
- 4 Summary
- 5 Outlook
- 6 Acknowledgement

Vacuum birefringence induced by intense laser fields

Experimental setup for the verification of vacuum birefringence:



Vacuum birefringence induced by intense laser fields

resulting ellipticity δ

$$\delta^2 \cong \frac{2\alpha}{15} \frac{z_0}{\lambda} \frac{I_0}{I_C} \kappa$$

T. Heinzl et al. 2006

I_0 -peak intensity

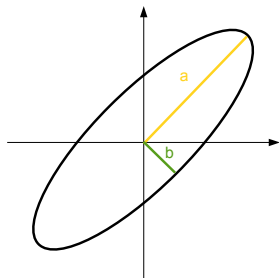
λ -probe pulse wavelength

z_0 -Rayleigh length

α -fine-structure constant

κ -correction factor

I_C -critical intensity



Ellipticity, which would be possible with the petawatt laser system at Jena:

$$I_0 = 10^{22} \frac{W}{cm^2}$$

$$\lambda = 0.08 nm$$

$$z_0 = 25 \mu m$$

$$\Rightarrow \delta^2 = 4.8 \cdot 10^{-10}$$

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Kinematical theory

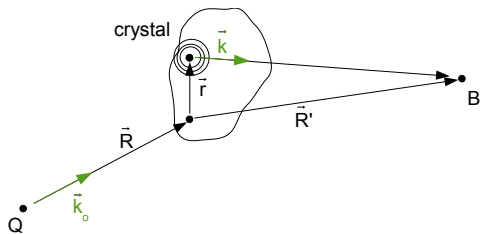


fig.: Diffraction

- no change of wavenumber in the crystal
- decrease of intensity of the primary beam is neglected
- no secondary scattering
 \Rightarrow only valid for mosaic crystals and thin crystals

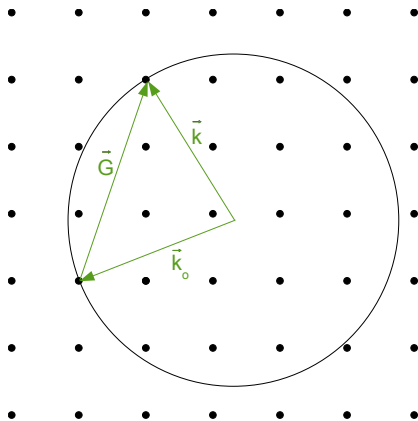
Intensity of the diffracted beam at the observation place

$$I(\vec{K}) \cong \frac{|A_0|^2}{R'^2} \left| \sum_{\vec{G}} \rho_{\vec{G}} \int e^{i(\vec{G}-\vec{K})\vec{r}} d\vec{r} \right|^2 \quad \text{with} \quad \vec{K} = \vec{k} - \vec{k}_0$$

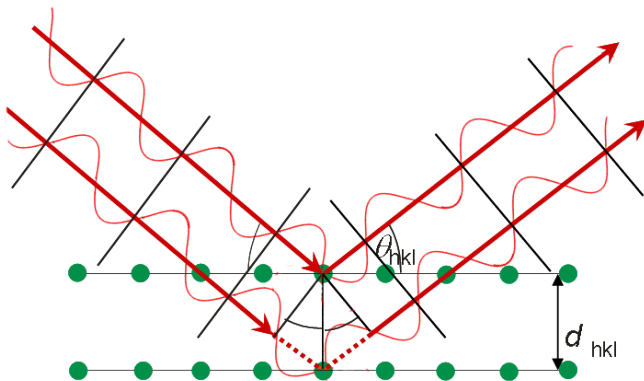
Diffraction condition

Laue condition

$$\vec{G} = \vec{k} - \vec{k}_0$$



Diffraction condition



Bragg's law

$$2d_{hkl} \sin(\theta_{B_{hkl}}) = n\lambda$$

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Dynamical theory of diffraction

- solves Maxwell's equations for the propagation of electromagnetic waves in the crystal
⇒ is valid for perfect crystals

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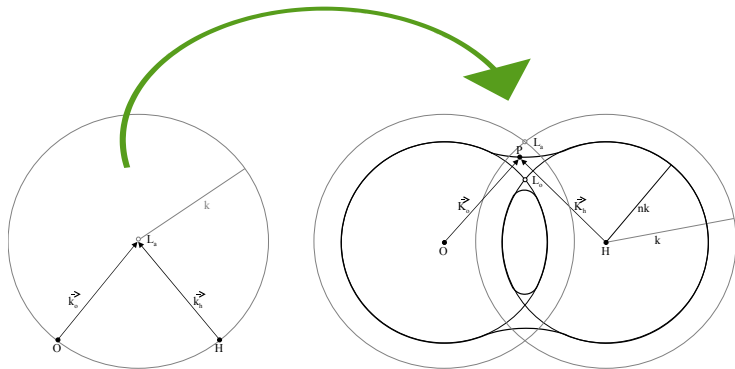
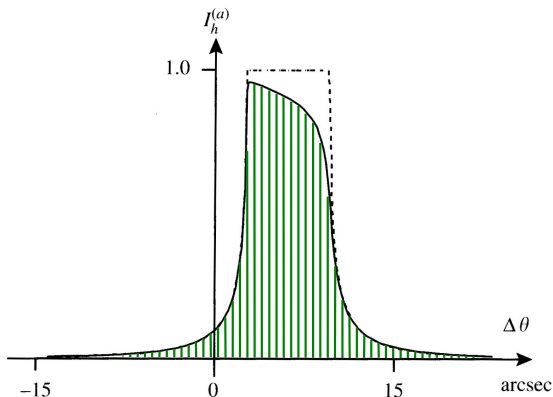


fig.: —Dispersion surface

Integrated intensity of a perfect crystal



Integrated intensity

$$R_{\sigma/\pi} = \frac{8R\lambda^2 |C| \sqrt{|\gamma|}}{\pi V \sin(2\theta_B)} \sqrt{F_h F_{\bar{h}}}$$

fig.: Comparison of the rocking curves for
 ··· non-absorbing
 — absorbing crystals

Integrated intensity of a perfect crystal

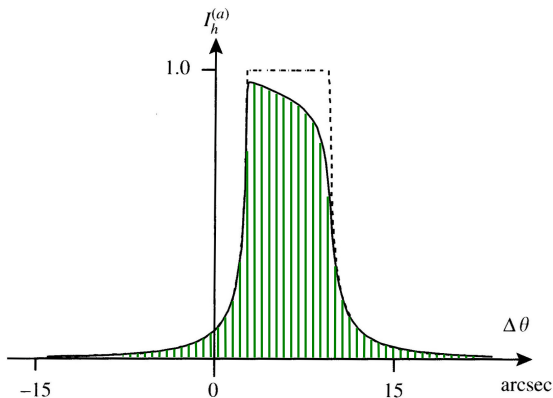


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Integrated intensity

$$R_{\sigma/\pi} = \frac{8R\lambda^2 |C| \sqrt{|\gamma|}}{\pi V \sin(2\theta_B)} \sqrt{F_h F_{\bar{h}}}$$

Polarization factor

$$C = \begin{cases} 1 & \text{for } \sigma \\ \cos 2\theta_B & \text{for } \pi \end{cases}$$

⇒ Suppression of the
 π -polarization at $\theta_B = 45^\circ$

Integrated intensity of a perfect crystal

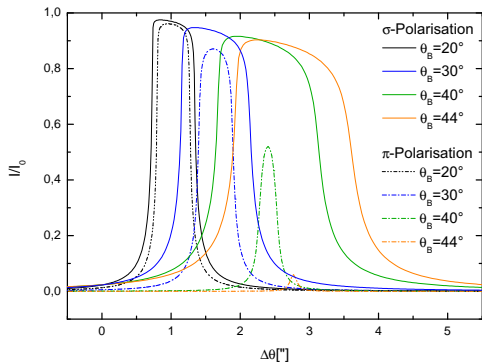


fig.: Comparison of rocking curves for different Bragg angles

Integrated intensity

$$R_{\sigma/\pi} = \frac{8R\lambda^2 |C| \sqrt{|\gamma|}}{\pi V \sin(2\theta_B)} \sqrt{F_h F_{-h}}$$

Polarization factor

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Polarization purity

Transmitted flux

$$T_{(\pi,\sigma)} = \int \int I_{(\pi,\sigma)}^n(\theta, E_0 + \epsilon) B(\theta, E_0 + \epsilon) d\theta d\epsilon$$

- $B(\theta, E_0 + \epsilon)$ - spectral brilliance
- E_0 - Bragg energy
- θ - angle around θ_B (Bragg angle)
- $I_{(\pi,\sigma)}(\theta, E_0 + \epsilon)$ - single-crystal reflectivity
- n - number of reflections inside the channel-cut

Polarization purity

$$\delta_0 = \frac{T_\pi}{T_\sigma}$$

Suppression of π -polarization

Si 333, Cu $K\alpha$, $\theta_B = 47.5^\circ$

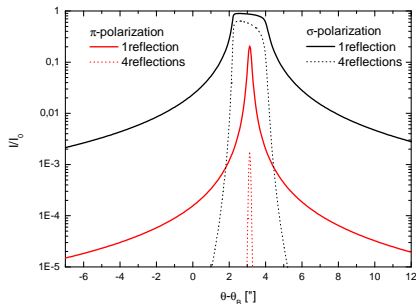
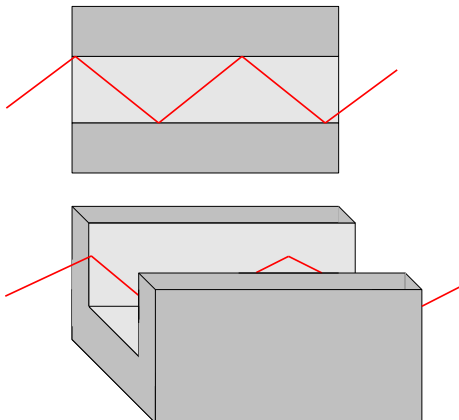


fig.: multiple reflections

$$\delta_1 = 2,91 \cdot 10^{-2} \leftrightarrow \delta_4 = 2,20 \cdot 10^{-4}$$

Multiple scattering

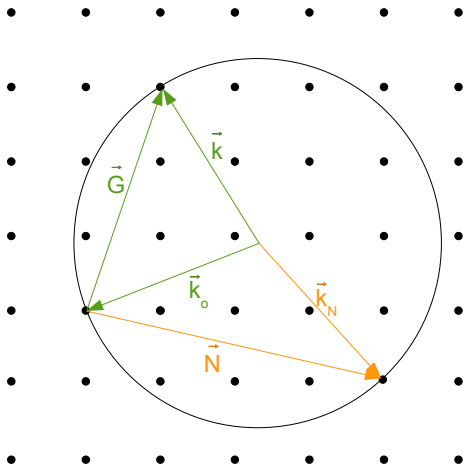


fig.: Appearance of multiple scattering demonstrated at the Ewald sphere

Umweganregungen

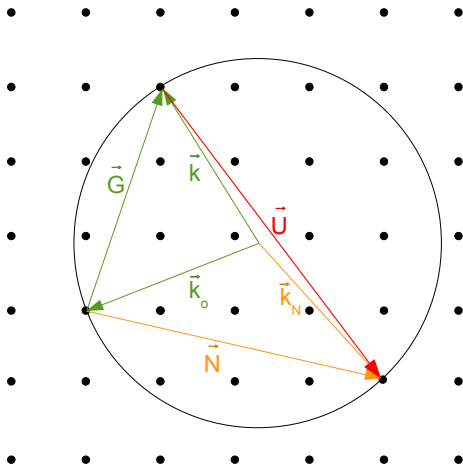


fig.: Appearance of indirect reflections demonstrated at the Ewald sphere

Multiple scattering

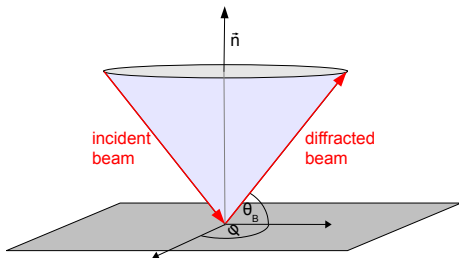


fig.: Kossel cone

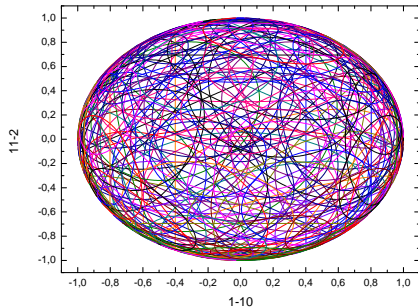


fig.: Kossel pattern of a (333) silicon crystal, $\text{CuK}\alpha$, $E=8.05\text{keV}$

Multiple scattering

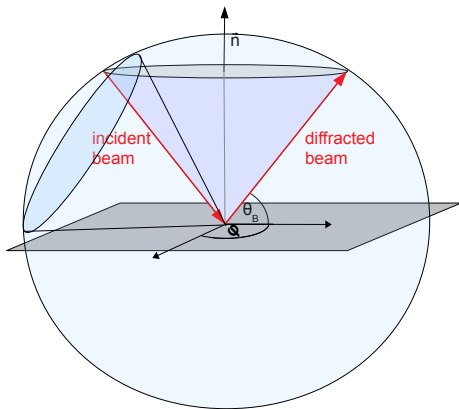


fig.: Kossel cone

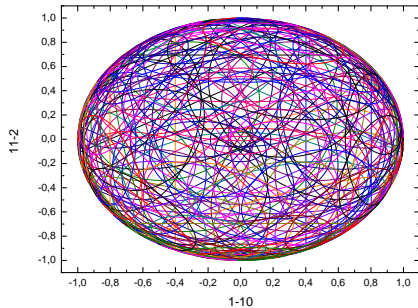


fig.: Kossel pattern of a (333) silicon crystal, $\text{CuK}\alpha$, $E=8.05\text{keV}$

Umweganregungen

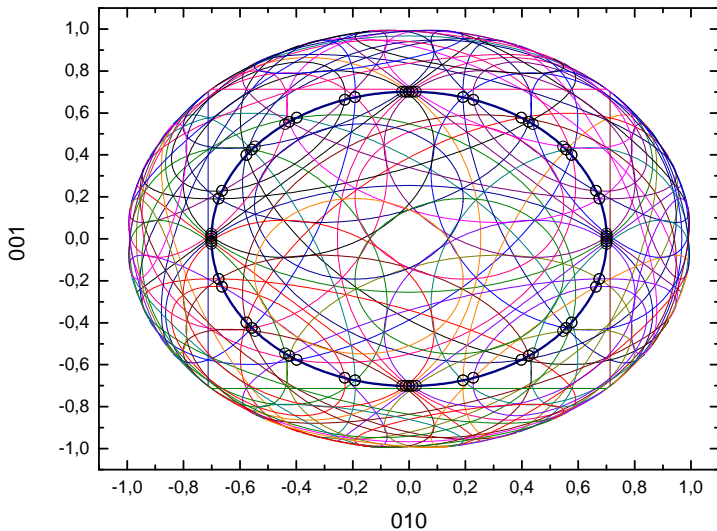


fig.: Kossel pattern of a (400) silicon crystal, $\text{FeK}\alpha$, $E=6.4\text{keV}$, $\theta_B = 45,48^\circ$

Umweganregungen

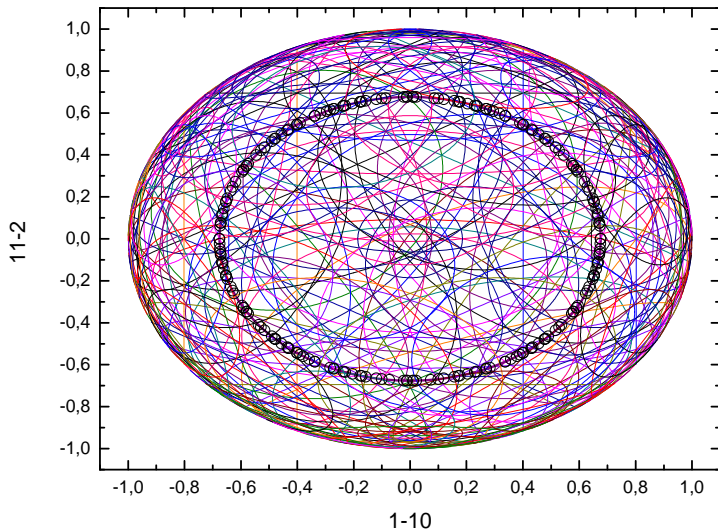


fig.: Kossel pattern of a (333) silicon crystal, $\text{CuK}\alpha$, $E=8.05\text{keV}$, $\theta_B = 47,5^\circ$

Umweganregungen

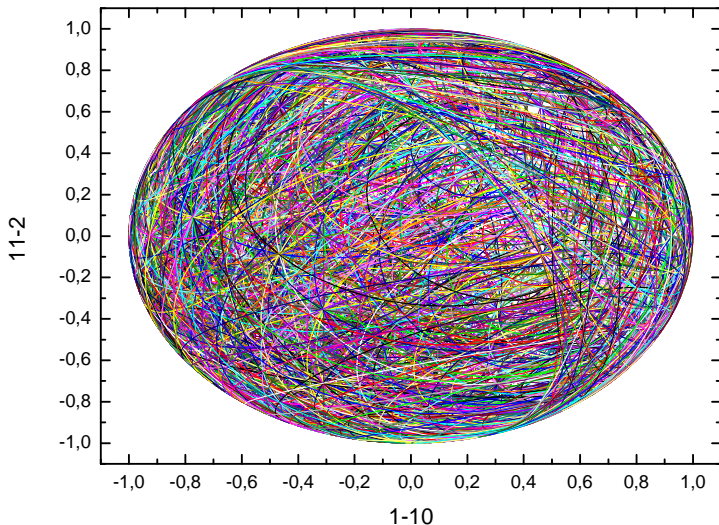
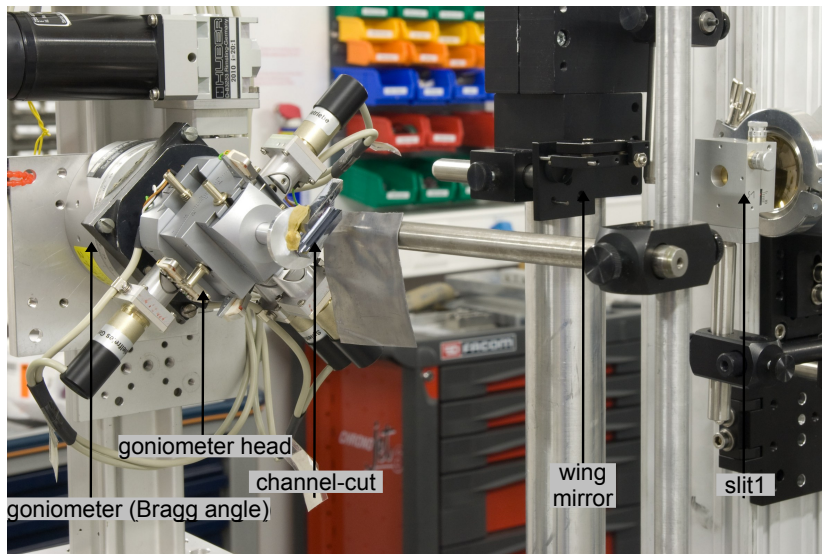
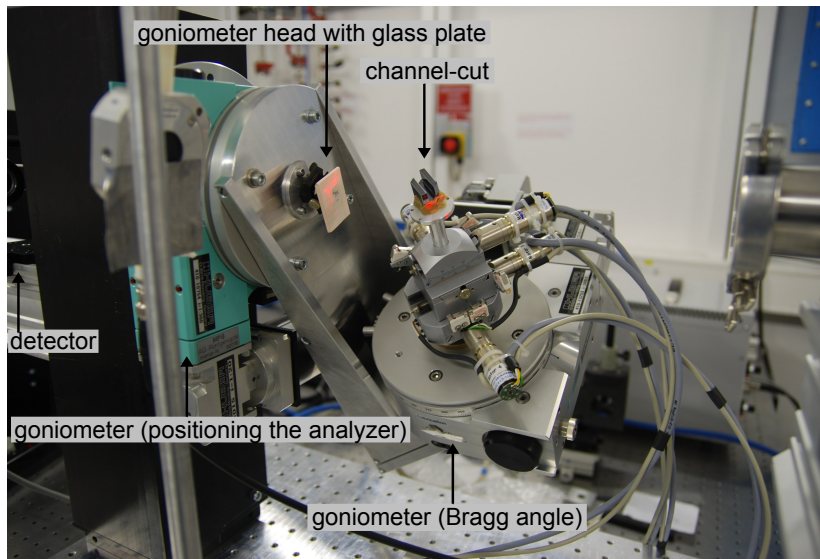


fig.: Kossel pattern of a (888) silicon crystal, $\text{AgK}\alpha$, $E=22.2\text{keV}$, $\theta_B = 45, 53^\circ$

Polarimeter setup: polarizer



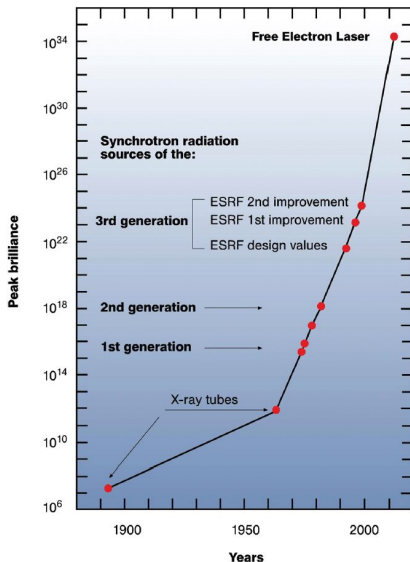
Polarimeter setup: analyzer



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ESRF



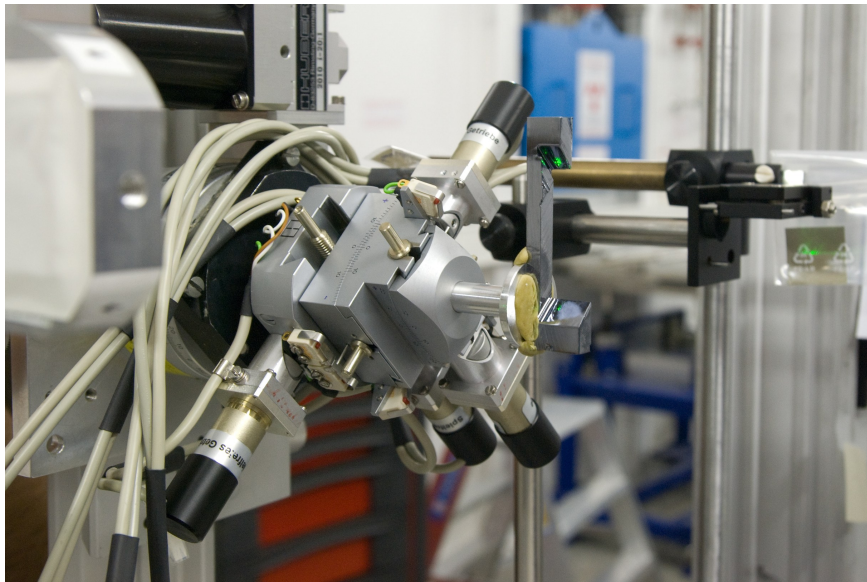
<http://www.saxier.org/>

<http://www.xfel.eu/>

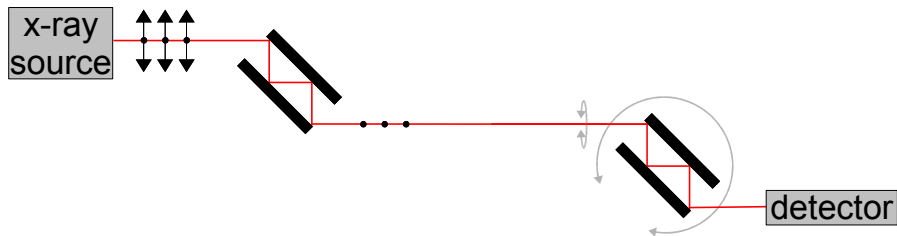
Source parameters

- 16 bunch filling mode, $I < 90mA$
- beam is premonochromized (Si 111 2 reflections)
- vertical divergence $4.35\mu rad$
- horizontal divergence $15.45\mu rad$
- 2 undulators ($< 10keV$, $> 10keV$)
- Number of photons: $10^{10} \frac{ph}{s} - 10^{12} \frac{ph}{s}$

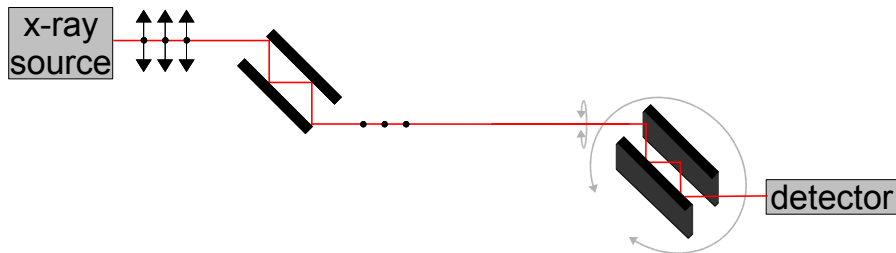
Setting up the wavelength



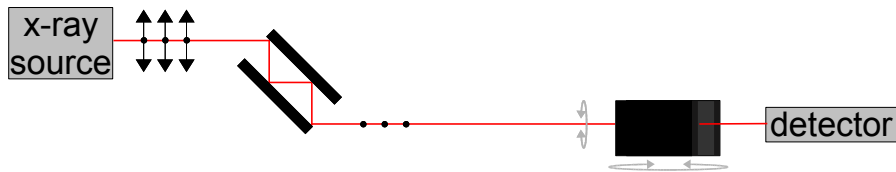
Principle of measurement



Principle of measurement



Principle of measurement



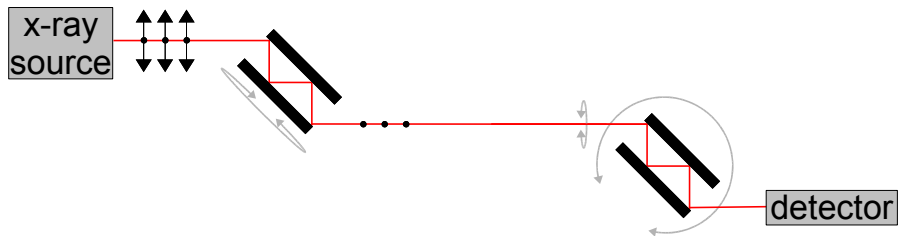
Results polarization purity

Intensity suppression for Si(400) with $E=6.457\text{keV}$ ($\theta_B = 45.000^\circ$).

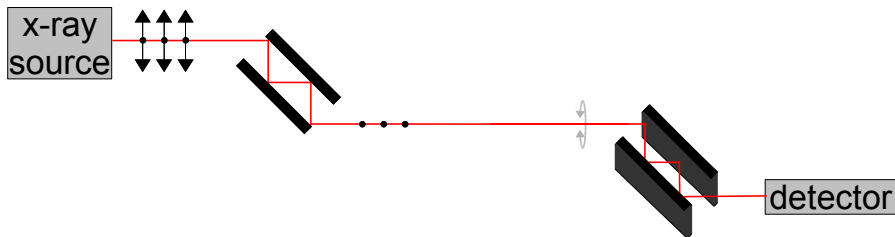
Results polarization purity

Intensity suppression for Si(800) with $E=12.914\text{keV}$ ($\theta_B = 45.000^\circ$).

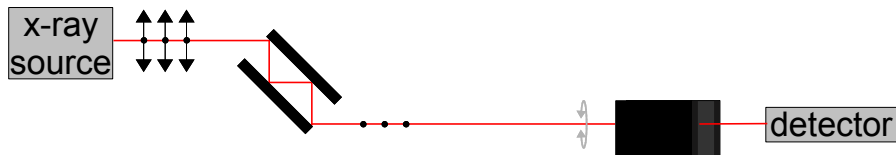
Varying the polarizer azimuth



Varying the polarizer azimuth



Varying the polarizer azimuth



Rocking curves depending on the analyzer position

Intensity suppression for Si(444) by varying the polarizer azimuth at $E=11.1838\text{keV}$.

Summary

Outlook

- polarization measurement at the ESRF using
 - more reflections (6,8)
 - other materials
 - crystal cooling
- full theoretical description of the polarimeter

Acknowledgement

- 
- I. Üschmann¹, S. Höfer, R. Loetzsch, O. Wehrhan, E. Förster¹,
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