

# 3rd EMMI workshop: anti-matter, hypermatter and exotica production at the LHC Wrap-up day 2

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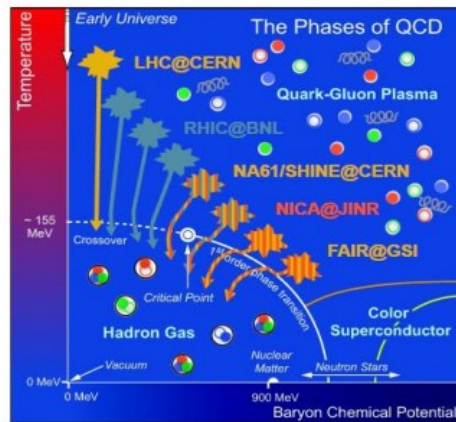
Ramona Lea

University and INFN (Trieste)

# So... how are produced light nuclei?

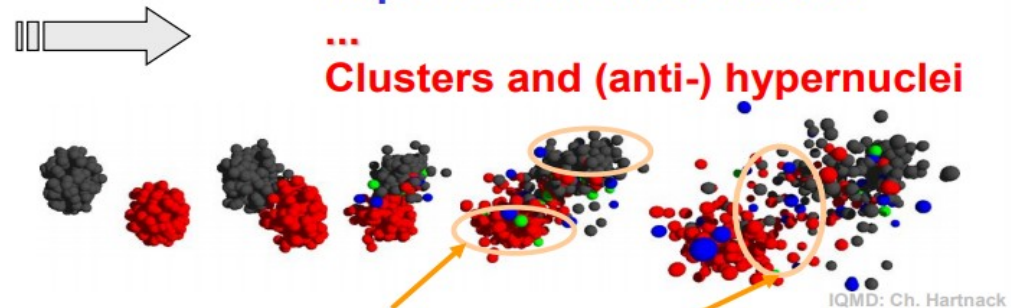
## The ,holy grail' of heavy-ion physics:

### The phase diagram of QCD



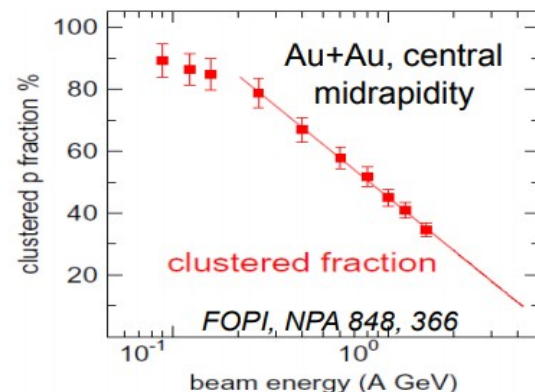
### Experimental observables:

#### ... Clusters and (anti-) hypernuclei



- projectile/target spectators → heavy cluster formation
- midrapidity → light clusters

- Clusters are very abundant at low energy

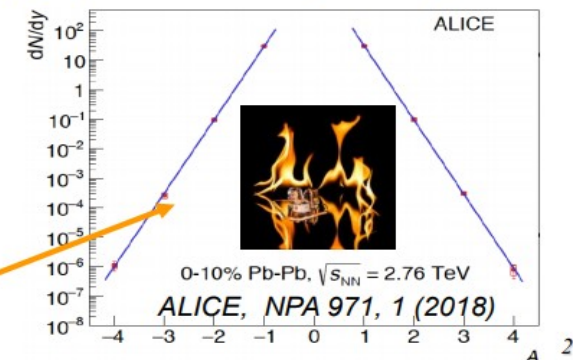


! Hyperons are created in participant zone

#### (Anti-) hypernuclei production:

- at mid-rapidity by  $\Lambda$  coalescence during expansion
- at projectile/target rapidity by rescattering/absorption of  $\Lambda$  by spectators

High energy HIC:  
,Ice in a fire' puzzle:  
how the weakly bound  
objects can be formed  
in a hot environment ?!

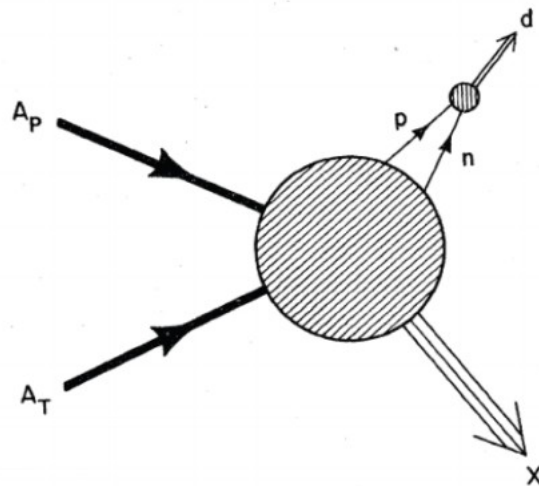


# Thermal vs coalescence

## The sudden formation of clusters

**Statistical model:** describes very well the multiplicities in central collisions but not the spectra (yield  $V, T, \mu$ )  
difficult to imagine how the cluster production takes place  
d:  $E_b = 2.2$  MeV , rms radius = 1.7 fm  
does not survive in heat bath of  $T > 100$  MeV  
“ice in fire”, “snowball in hell”

**Coalescence:** goes back to Butler and Pearson PR129,836 (p+A)



d-production is a 3-body process  
momentum has to be transferred to the third body

QM: in a static potential  $\sim 1/p^2$

$$n_d(p) \sim \frac{1}{p^2} n_n^2\left(\frac{p}{2}\right)$$

# Development of different models

## Why does one need a new model ?

### Present microscopic approaches:

- ❑ VUU(1985), BUU(1985), HSD(1996), PHSD(2008), SMASH(2016) solve the time evolution of the one-body phase-space density in a mean field  
→ **no dynamical fragments**
- ❑ UrQMD is a n-body model but makes clusterization via coalescence and a statistical fragmentation model
- ❑ QMD is a n-body model but is limited to energies  $< 1.5$  AGeV  
→ describes fragments at SIS energies,  
but conceptually not adapted for NICA/FAIR energies and higher

In order to understand the **microscopic origin of cluster formation** one needs:

- a realistic model for the dynamical time evolution of HICs
- **dynamical modelling of cluster formation** based on interactions

**Dynamical modelling of cluster formation** is a complex task which involves:  
the fundamental nuclear properties, quantum effects, variable timescales

# Development of different models

## PHQMD

**The goal:** to develop a **unified n-body microscopic transport approach** for the description of heavy-ion dynamics and dynamical cluster formation from low to ultra-relativistic energies

**Realization:** combined model **PHQMD = (PHSD & QMD) & SACA**

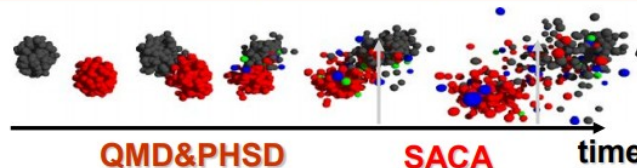
### Parton-Hadron-Quantum-Molecular Dynamics

Initialization → propagation of baryons:  
**QMD (Quantum-Molecular Dynamics)**

Propagation of partons (quarks, gluons) and mesons  
+ **collision integral** = interactions of hadrons and partons (QGP)  
from **PHSD (Parton-Hadron-String Dynamics)**

Clusters recognition:  
**SACA (Simulated Annealing Clusterization Algorithm)**  
vs. **MST (Minimum Spanning Tree)**

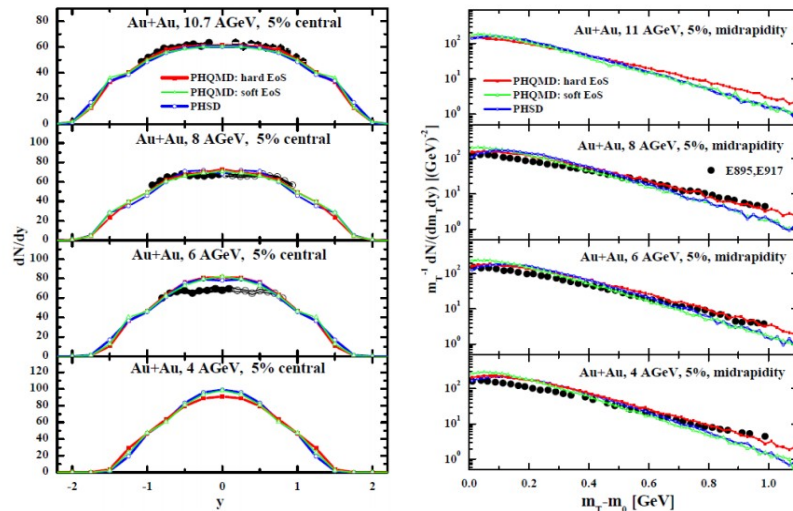
(Can start at any time)



# Does microscopic approach work?

## PHQMD: ,bulk' dynamics at AGS/FAIR/NICA energies

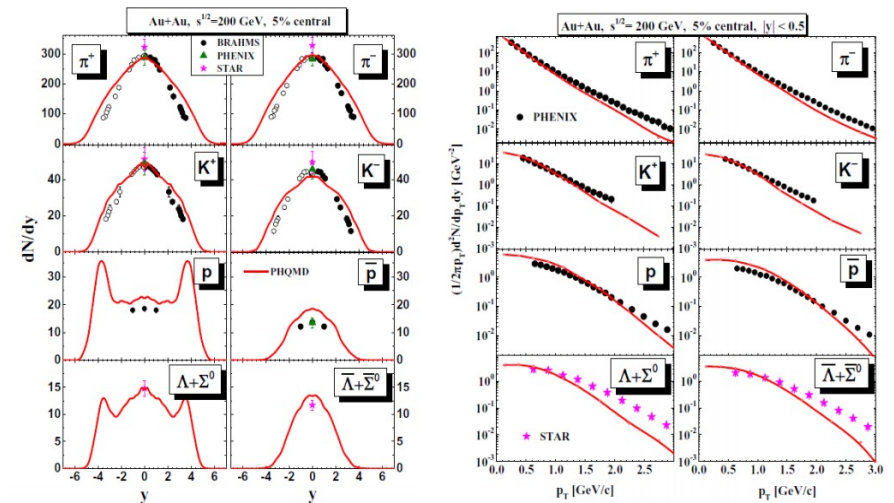
The rapidity and  $m_T$  distributions for protons from 5% central Au+Au collisions at 4, 6, 8, 10.7 A GeV



- the influence of EoS is slightly visible in rapidity spectra of protons
- $m_T$  spectra of protons from PHQMD with a 'hard' EoS are harder than with 'soft' EoS
- PHQMD results for the  $m_T$  spectra with 'soft' EoS are in a good agreement with the PHSD spectra (using 'soft' EoS in default PHSD4.0 version)
- QMD and MF dynamics gives similar results with similar EoS

## PHQMD: ,bulk' dynamics at RHIC

The rapidity and  $p_T$  distributions for p, anti-p,  $\pi^+$ ,  $\pi^-$ ,  $K^+$ ,  $K^-$ ,  $\Lambda + \Sigma^0$ ,  $\bar{\Lambda} + \bar{\Sigma}^0$  from 5% central Au+Au collisions at  $s^{1/2} = 200$  GeV

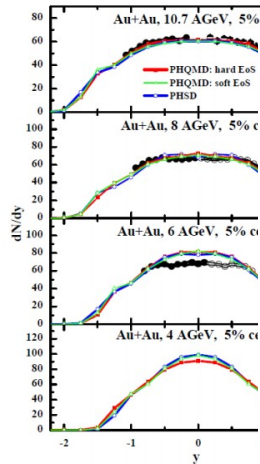


PHQMD: results are similar to PHSD - since at RHIC energies the dynamics is dominated by collisions of partons/hadrons rather than nuclear potential interactions!

# Does microscopic approach work?

## PHQMD: 'bulk' dynamics at AGS/FAIR/NICA energies

The rapidity and  $m_T$  distributions for protons from 5% central Au+Au collisions at 4, 6, 8, 10.7 A GeV



- the influence of EoS is slightly
- $m_T$  spectra of protons from PHQMD
- PHQMD results for the  $m_T$  spectra (using 'soft' EoS in de → QMD and MF dynamics give



## Summary

The **PHQMD** is a **microscopic n-body transport approach** for the description of heavy-ion dynamics and cluster formation

combined model **PHQMD = (PHSD & QMD) & SACA**

### PHQMD

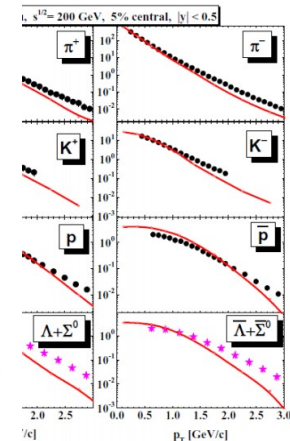
- provides the good description of **hadronic 'bulk' observables** from SIS to RHIC energies
- shows **sensitivity to EoS**:  $m_T$  spectra of baryons
- predicts the **dynamical formation of clusters** from low to ultra-relativistic energies
- allows to understand the proton spectra and the **properties of clusters** ( $dn/dp_T dy$ ,  $v_1, v_2$ , fluctuations)
- allows to understand clusters formation in the **participant and spectator region** (consistent with available cluster data by ALADIN collaboration)
- allows to understand the **formation of hypernuclei**

❖ **PHQMD : under development**



## at RHIC

$\Lambda(\Lambda+\Sigma^0)$  from 5% central Au+Au

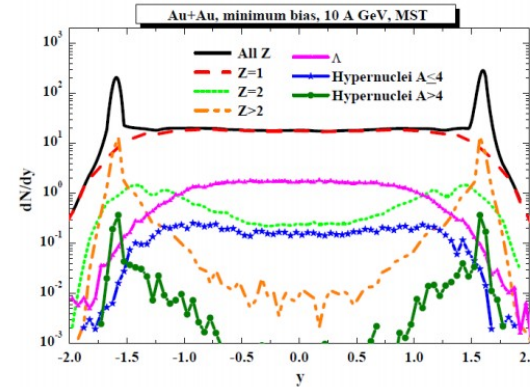
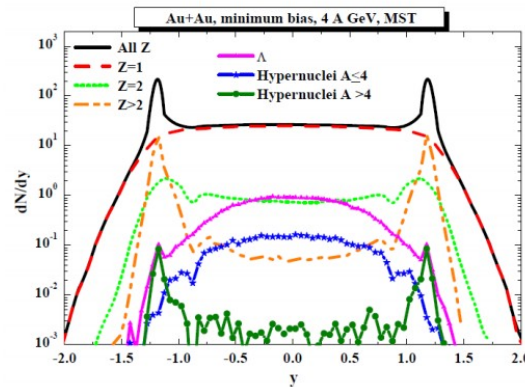


energies the dynamics is 1 nuclear potential interactions!

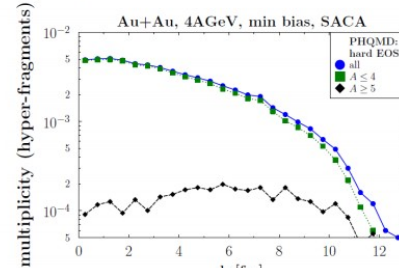
# Hypernuclei production

## PHQMD: hypernuclei

PHQMD results (with a **hard EoS** and **MST algorithm**) for the rapidity distributions of all charges,  $Z = 1$  particles,  $Z=2$ ,  $Z>2$ , as well as  $\Lambda$ 's, hypernuclei  $A \leq 4$  and  $A > 4$  for Au+Au at 4 and 10 AGeV



The multiplicity of light hypercluster vs. impact parameter  $b$  for Au+Au, 4 AGeV



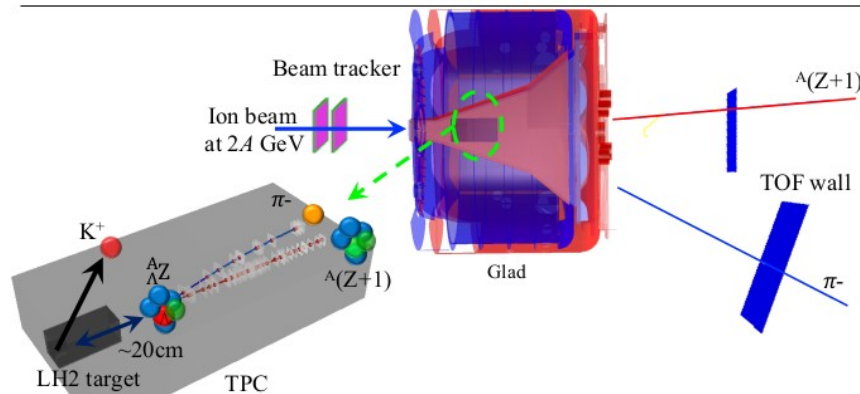
- Central collisions  $\rightarrow$  light hypernuclei
- Peripheral collisions  $\rightarrow$  heavy hypernuclei

Penetration of  $\Lambda$ 's, produced at midrapidity, to target/projectile region due to rescattering

$\rightarrow$  Possibility to study  $\Lambda N$  interaction

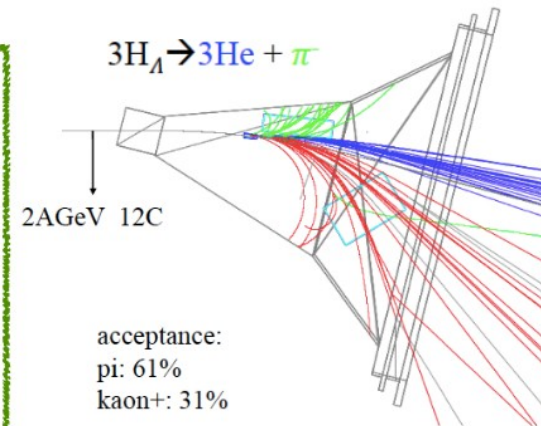
# Future plans for hypernuclei

## Future plans at R3B, GSI/FAIR



### Hypernuclei production in inverse kinematics

- Inverse kinematics with **hydrogen target**
  - increase luminosity
  - reaction mechanism
  - decrease background
- Strangeness production by kaon tagging
- **TPC** inside GLAD dipole (invariant mass)
- Starting project (prototype under design)

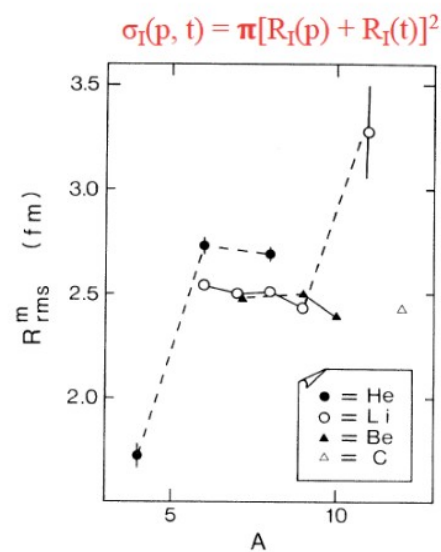


# Future plans for hypernuclei

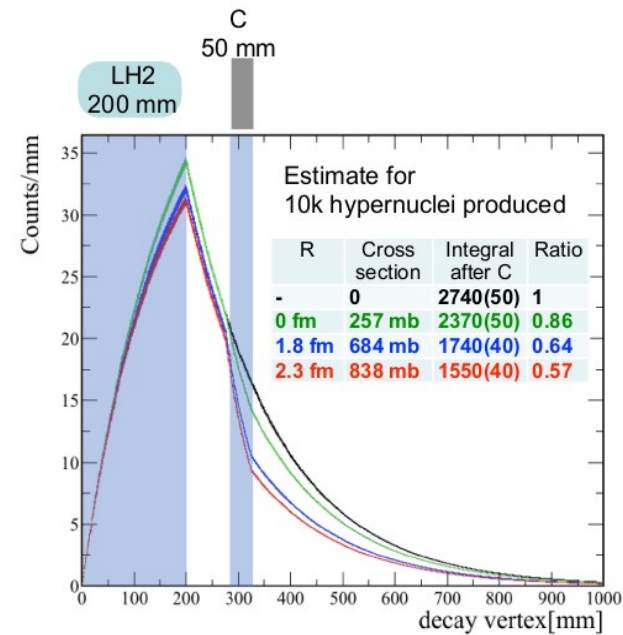
## Future plans at R3B, GSI/FAIR



- “Are  $^3\text{H}$ -lambda and  $^7\text{Be}$ -lambda halos?”
- Hyper-version of the discovery of halo nuclei by Tanihata (1985)



I. Tanihata et al., PLB 160 (1985)  
I. Tanihata et al., PRL 55 (1985)



# Coalescence vs thermal

## Can light nuclei exist in a fireball?

- ▶ Interparticle spacing in a hadron gas is about 1.5 fm at  $T = 156$  MeV.
- ▶ Root mean square radius of a deuteron is 2.0 fm.
- ▶ Binding energy of a deuteron is  $\epsilon_B = 2.2$  MeV.
- ▶ A characteristic time of deuteron formation is  $1/\epsilon_B = 100$  fm/c.
- ▶ A hadron gas at  $T = 156$  MeV is essentially a classical system.

- *Snowflakes in hell ?*  
- *No, snowflakes from hell.*



# Coalescence from CF

Coalescence from correlation functions:

$$B_2(p) \approx \frac{3}{2m} \int d^3q \mathcal{D}(\vec{q}) C_2^{\text{PRF}}(\vec{p}, \vec{q})$$

R. Scheibl, U. Heinz, Phys.Rev. C59 (1999) 1585-1602



$$B_2 = \frac{3 \pi^{3/2} \langle C_d \rangle}{2m_t \mathcal{R}_\perp^2(m_t) \mathcal{R}_\parallel(m_t)}. \quad (6.3)$$

$$\langle C_d \rangle \approx \frac{1}{\left(1 + \left(\frac{d}{2\mathcal{R}_\perp(m)}\right)^2\right) \sqrt{1 + \left(\frac{d}{2\mathcal{R}_\parallel(m)}\right)^2}}. \quad (4.12)$$

Using this approach to coalescence, the same results obtained by Scheibl and Heinz (but here no hydro is needed) and by Mrowczynski using WF

S. Mrowczynski, Acta Phys.Polon. B48 (2017) 707



assumed: 1D GSM

$$D(\mathbf{r}) = \frac{e^{-\frac{\mathbf{r}^2}{4R_{\text{kin}}^2}}}{(4\pi R_{\text{kin}}^2)^{3/2}} \Rightarrow C_2^{\text{PRF}} = e^{-R_{\text{kin}}^2 q^2}$$

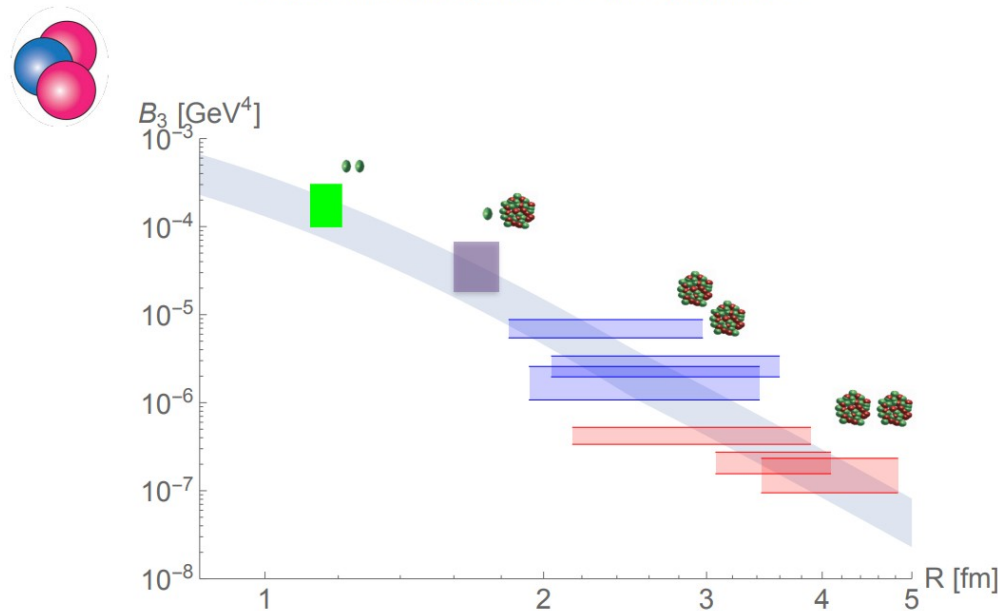
$\gamma_d \rightarrow 1$

$$\text{obtained: } \frac{dN_d}{d^3\mathbf{p}} = \mathcal{A} \frac{dN_p}{d^3(\frac{1}{2}\mathbf{p})} \frac{dN_n}{d^3(\frac{1}{2}\mathbf{p})}, \quad \mathcal{A} = \frac{3}{4} \frac{\pi^{3/2}}{(R_{\text{kin}}^2 + R_d^2)^{3/2}} = \boxed{\gamma_d} \frac{m}{2} B_2$$

# Coalescence from CF

Nuclear coalescence from correlation functions

KB, M. Takimoto, Phys.Rev. C99 (2019) no.4, 044913



Kfir BLUM

The theory seems to work, but one can see if it can be falsified.

How:

Reproduce “theoretical” results by using the same event,  $p_T$  and centrality intervals → bring together HBT and nuclei analyses

Measure accurately of HBT radii and multiplicity (→ calibrate the x-axis) : different model would have different dependencies.

# Is it possible to falsify thermal model?

## ${}^4\text{He}$ vs. ${}^4\text{Li}$

- ▶ The thermal and coalescence models give different predictions on the ratio of yields of  ${}^4\text{Li}$  to  ${}^4\text{He}$ .
- ▶ In the thermal model the ratio of yields is independent of collision centrality.
- ▶ In the coalescence model the ratio is maximal for central collisions and rapidly decreases when one goes to peripheral collisions.
- ▶ Since  ${}^4\text{Li}$  can be observed through the correlation function of  ${}^3\text{He}$ - $p$ , the correlation needs to be measured.

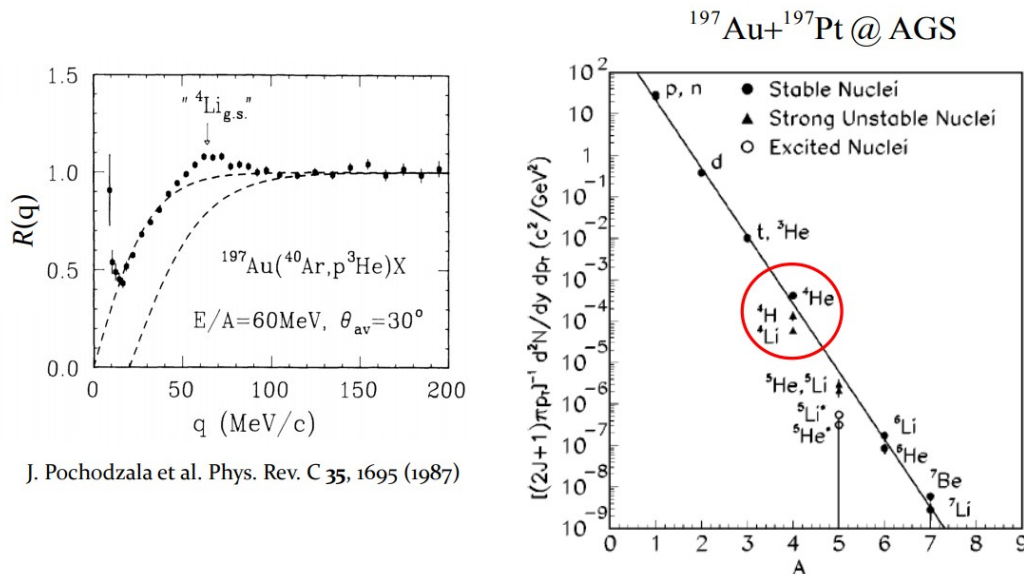
## $h$ - $D$ correlations

- ▶ Hadron-deuteron correlations carry information about source of deuterons.
- ▶ Measurement of  $h$ - $D$  &  $h$ - $p$  correlation function can tell us whether deuterons are directly emitted from a fireball like protons or deuterons are formed due to final state interactions.
- ▶  $p$ - $D$  correlation functions show a sufficient sensitivity to a size of particle source to falsify the thermal or coalescence model.

# ${}^4\text{He}$ vs ${}^4\text{Li}$

## How to observe ${}^4\text{Li}$ ?

Measurement of the correlation function of  ${}^3\text{He}$ - $p$  is needed



J. Pochodzala et al. Phys. Rev. C 35, 1695 (1987)

T. A. Armstrong et al. Phys. Rev. C 65, 014906 (2001)

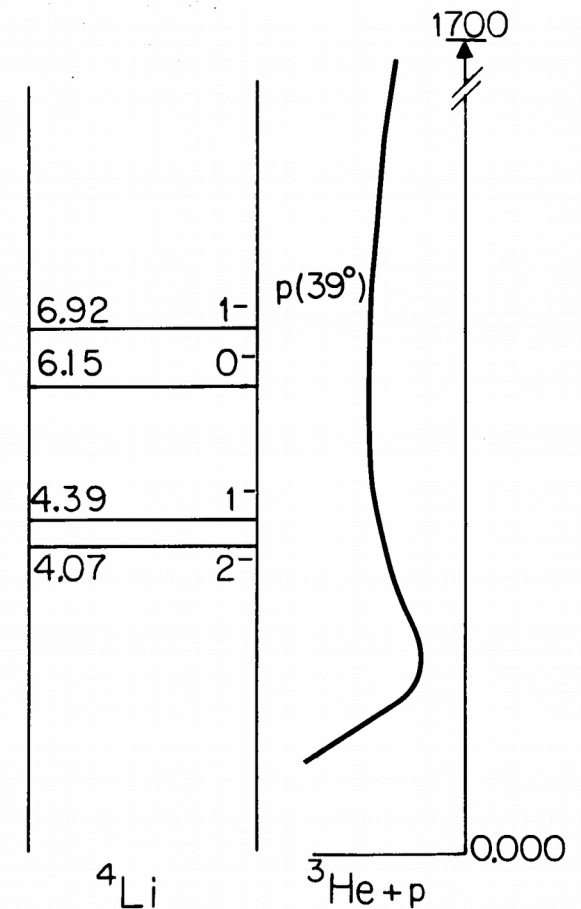
A difference between  ${}^4\text{He}$  and  ${}^4\text{Li}$  should be seen in data.

Experimentally this is very difficult because of the production yield is small and  ${}^4\text{Li}$  is a wide and “complex” resonance.

→ One should maybe repeat the theoretical exercise to look at  ${}^3_{\Lambda}\text{H}$  and  ${}^3\text{He}$

$$\frac{7.718}{n+3p}$$

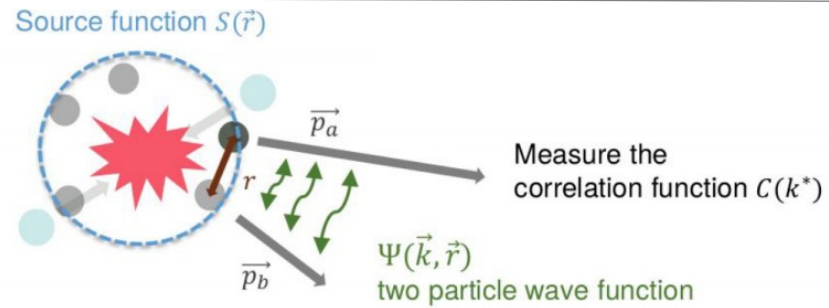
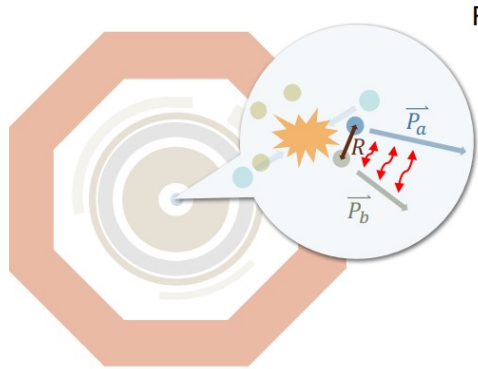
$$\frac{5.494}{d+2p}$$



19

Stanislaw MROWCZYNSKI

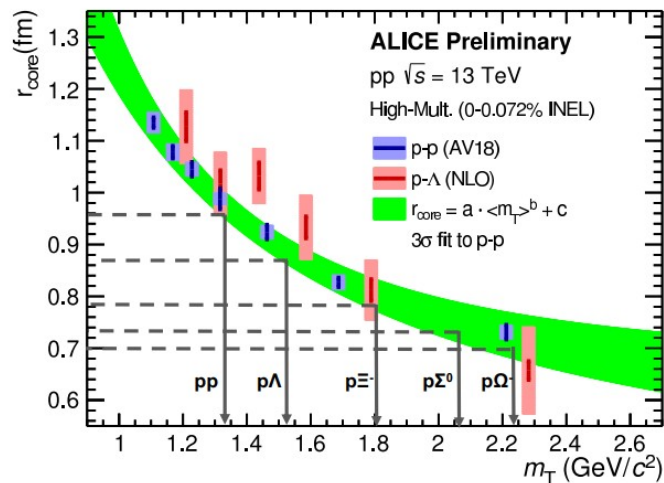
# Two particle correlation



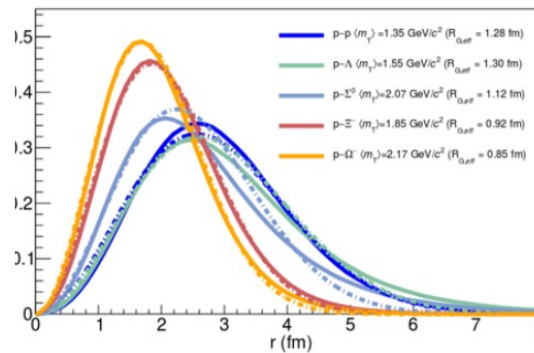
$$C(k^*) = \mathcal{N} \frac{N_{\text{Same}}(k^*)}{N_{\text{Mixed}}(k^*)} = \int S(\vec{r}) |\Psi(\vec{k}^*, \vec{r})|^2 d^3\vec{r} \xrightarrow{k^* \rightarrow \infty} 1$$

$>1$  : **Attractive** Interaction  
 $<1$  : **Repulsive** Interaction

Relative distance / reduced momentum in the rest frame of the pair



Global Source for each Pair

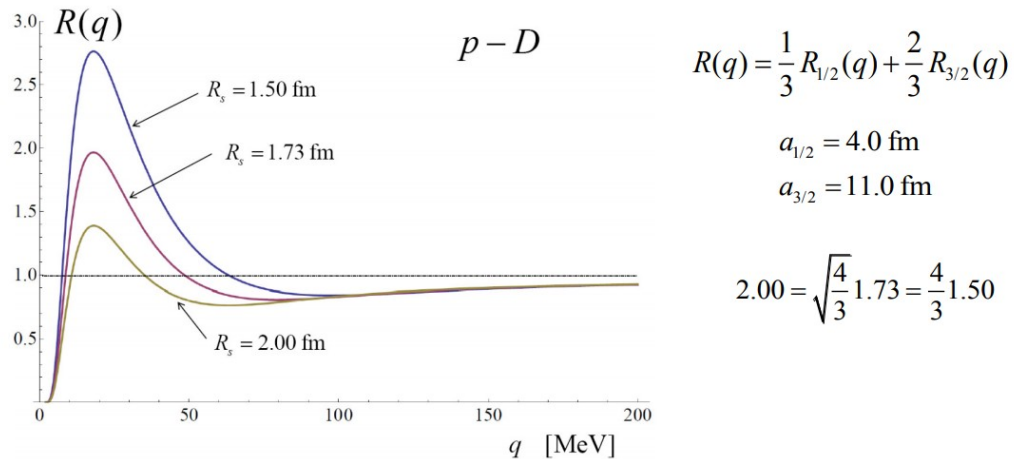


The source size is extracted from pp pairs and a common “core-baryon” source is obtained

Laura FABBIETTI

# h-D correlation

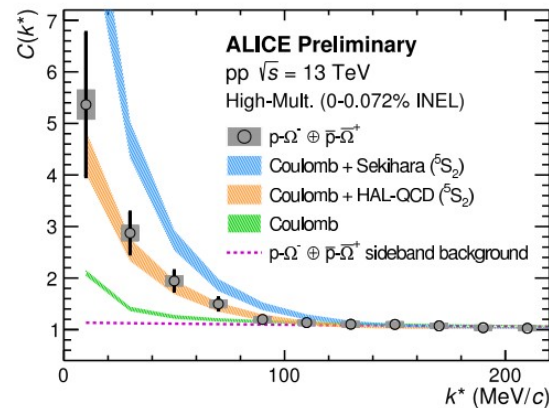
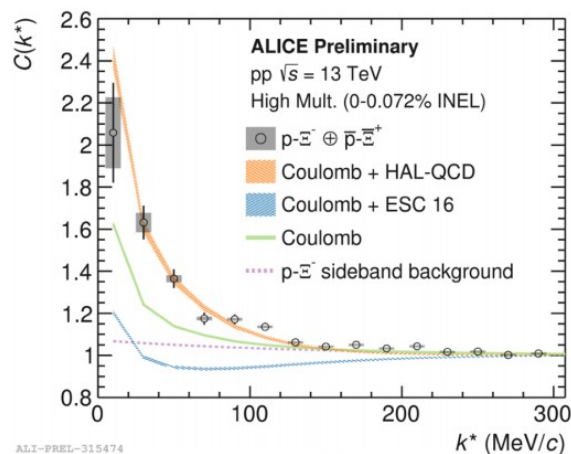
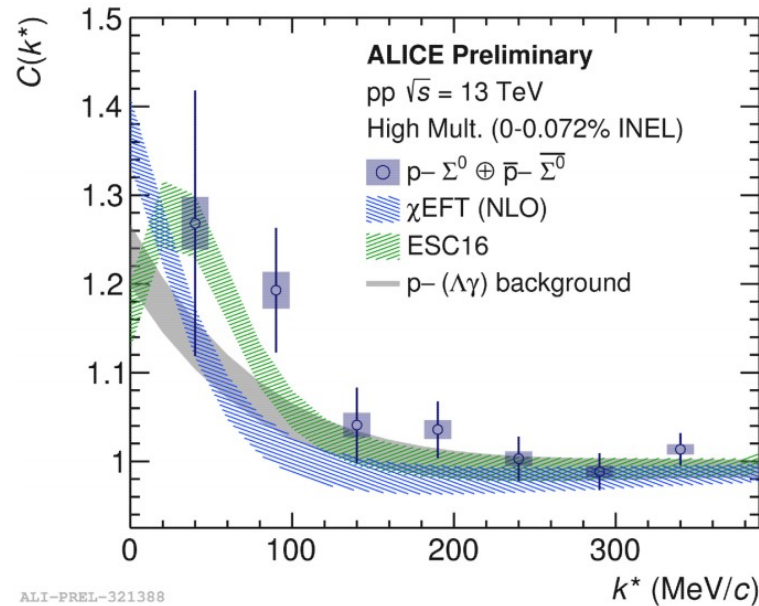
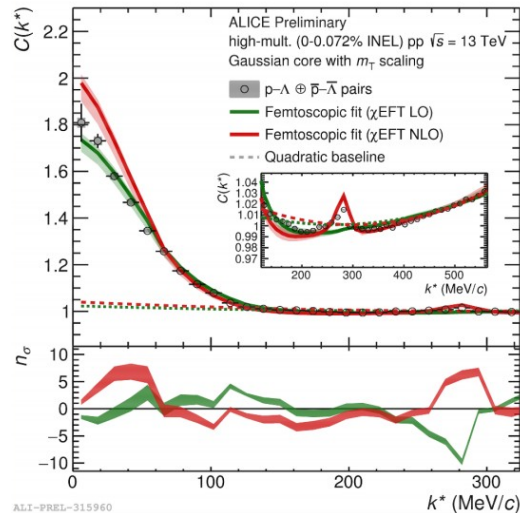
## *p-D* correlation functions



$R_s$  from *p-D* correlation function vs.  $R_s$  from *p-p* correlation function

CF are sensitive to source radius  
so looking at p-p and p-d CF it  
would be possible to measure if  
the source size changes for the  
different particle pairs

# Study baryon-baryon interaction



Plenty of new results that can be used to as alternative (or complementary) to scattering experiment to study very in detail the NY and YY interactions

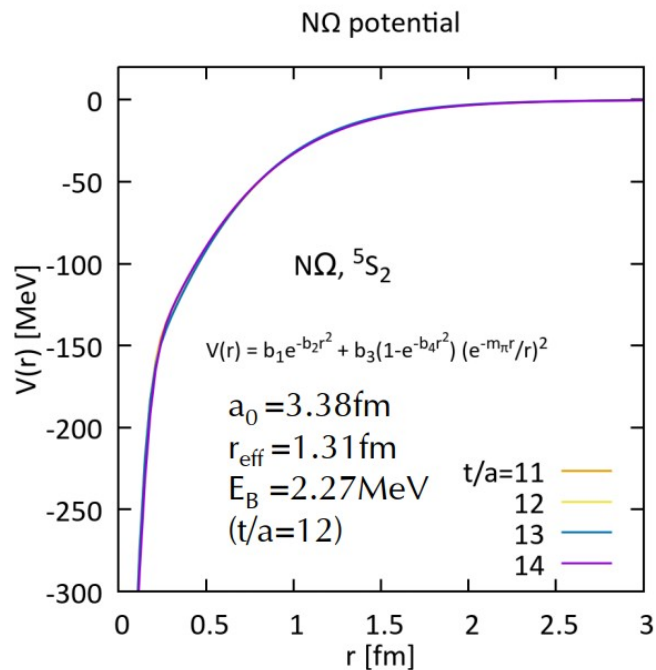
# CF from theory

Kenji Morita (QST/Riken)

3 Dec. 2019

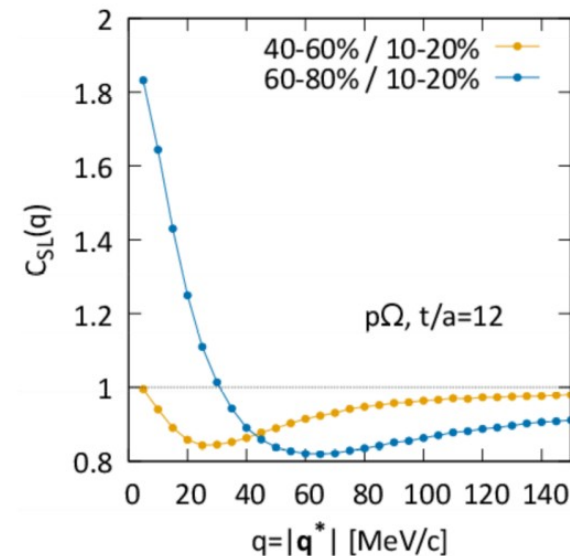
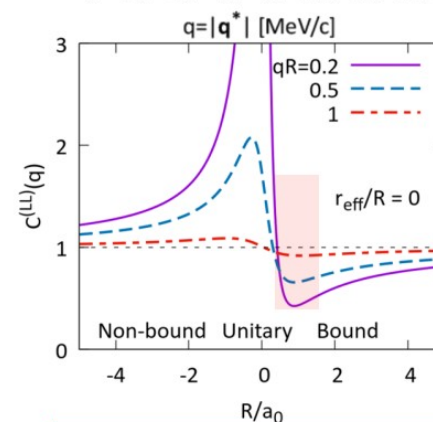
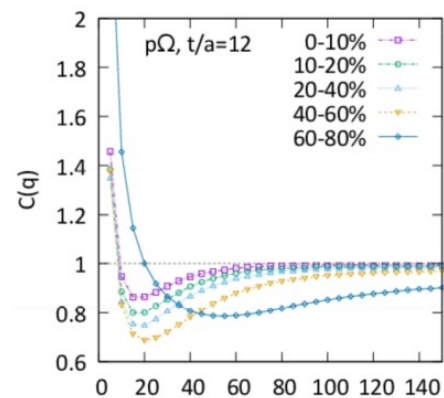
## Example : pΩ Correlation

■ Theory : HAL QCD potential + Expanding Source model



## Example: pΩ Correlation

KM+ '19 (PRC in press)



40-60% Bound regime  
 60-80% Unitary regime

pp unitary measured: Need HI to confirm  
 bound regime (STAR, but more statistics!)

Kenji MORITA

3rd EMMI Workshop on anti-matter, hyper-matter and exotica production at the LHC

03/12/2019

# QCD exotica

## Search for $X(5568)$ in $B_s^0 \pi^+$ state

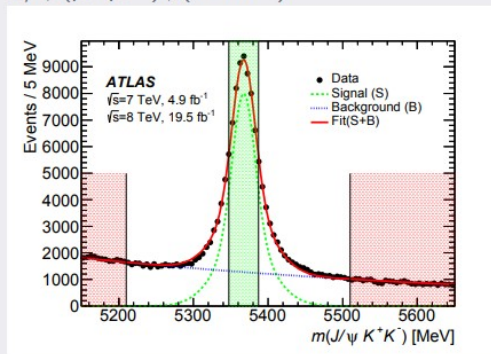
Phys. Rev. Lett. 120, 202007 (2018)



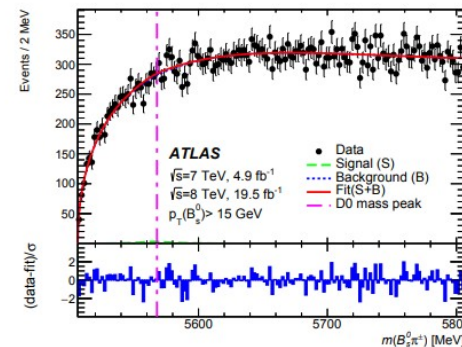
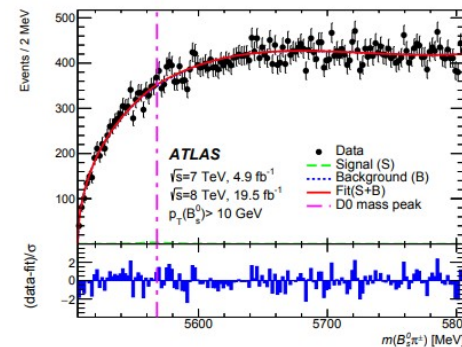
### Analysis overview

D0 reported evidence of a narrow structure,  $X(5568)$ , which was interpreted as a tetraquark with four different quark flavors:  $b, s, u, d$ . [ PRL 117, 022003 (2016) ]

- ▶  $4.9 \text{ fb}^{-1}$  of  $pp$  collision data at 7 TeV and  $19.5 \text{ fb}^{-1}$  at 8 TeV
- ▶ Reconstruct  $B_s^0$  mesons in their decays to  $J/\psi(\mu^+ \mu^-) \phi(K^+ K^-)$



- ▶ Combine each of the tracks forming the selected PV with the selected  $B_s^0$  candidate



Y. Fang | Overview on exotica production at ATLAS

Yi FANG

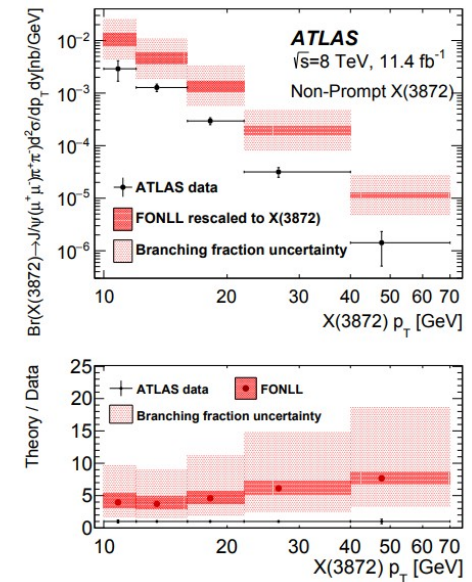
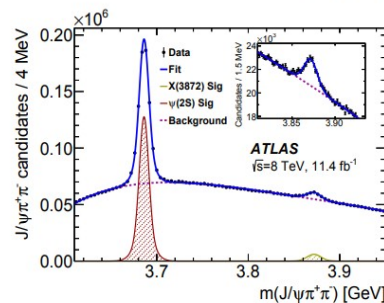
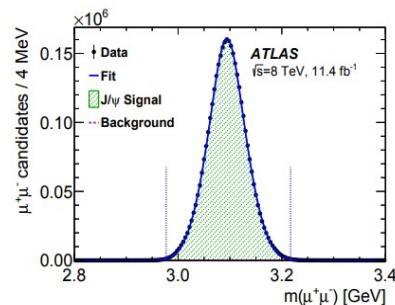
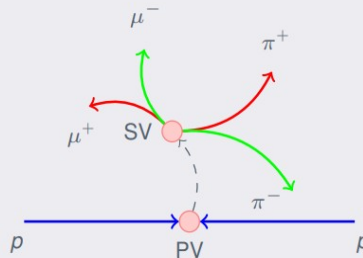
# QCD exotica

## $X(3872)$ production at 8 TeV JHEP 01 (2017) 117



### Analysis overview

- ▶  $11.4 \text{ fb}^{-1}$  of  $pp$  collision data at 8 TeV
- ▶  $J/\psi$  candidates formed by fitting a muon pair to a common vertex
- ▶ Combine a  $J/\psi$  candidate with another 2 tracks assumed as pions
- ▶ Di-muon mass is constrained to the  $J/\psi$  mass in a 4-prong vertex fit



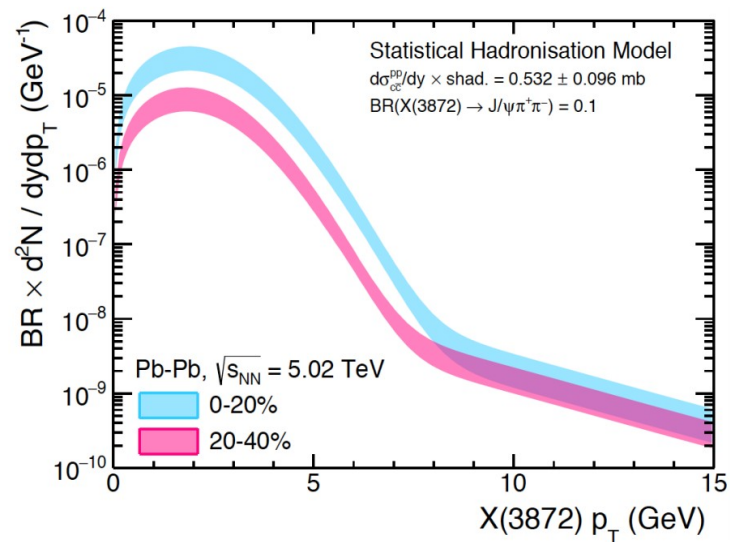
Y. Fang | Overview on exotica production at ATLAS

Y. Fang | Overview on exotica production at ATLAS

Yi FANG

# X(3872)

transverse momentum spectrum for X(3872) in the  
statistical hadronization model  
Pb-Pb collisions at 5 TeV/u



Measurement of X(3872)  
→ relevant role in Pb-Pb collisions

# Back to particle production...

## (Sub-Threshold) Strangeness Production

Unique observable:

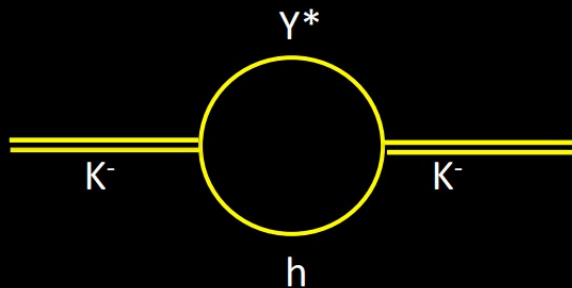
Not produced in binary NN collisions at  $\sqrt{s_{NN}} = 2.4$  GeV (no obvious elementary reference)

$NN \rightarrow NYK^+$ :  $\sqrt{s_{NN}} = 2.55$  GeV,

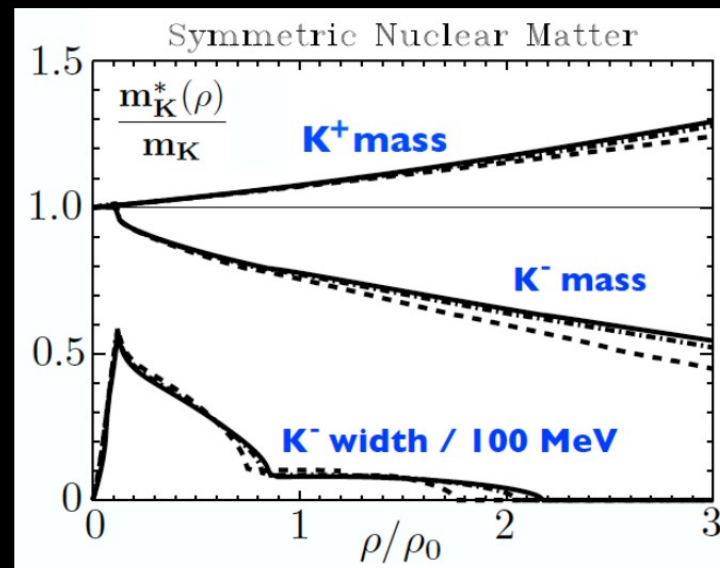
$NN \rightarrow NNK^+K^-$ :  $\sqrt{s_{NN}} = 2.86$  GeV

(strong  $K^-$  suppression).

Energy must be provided from the system.



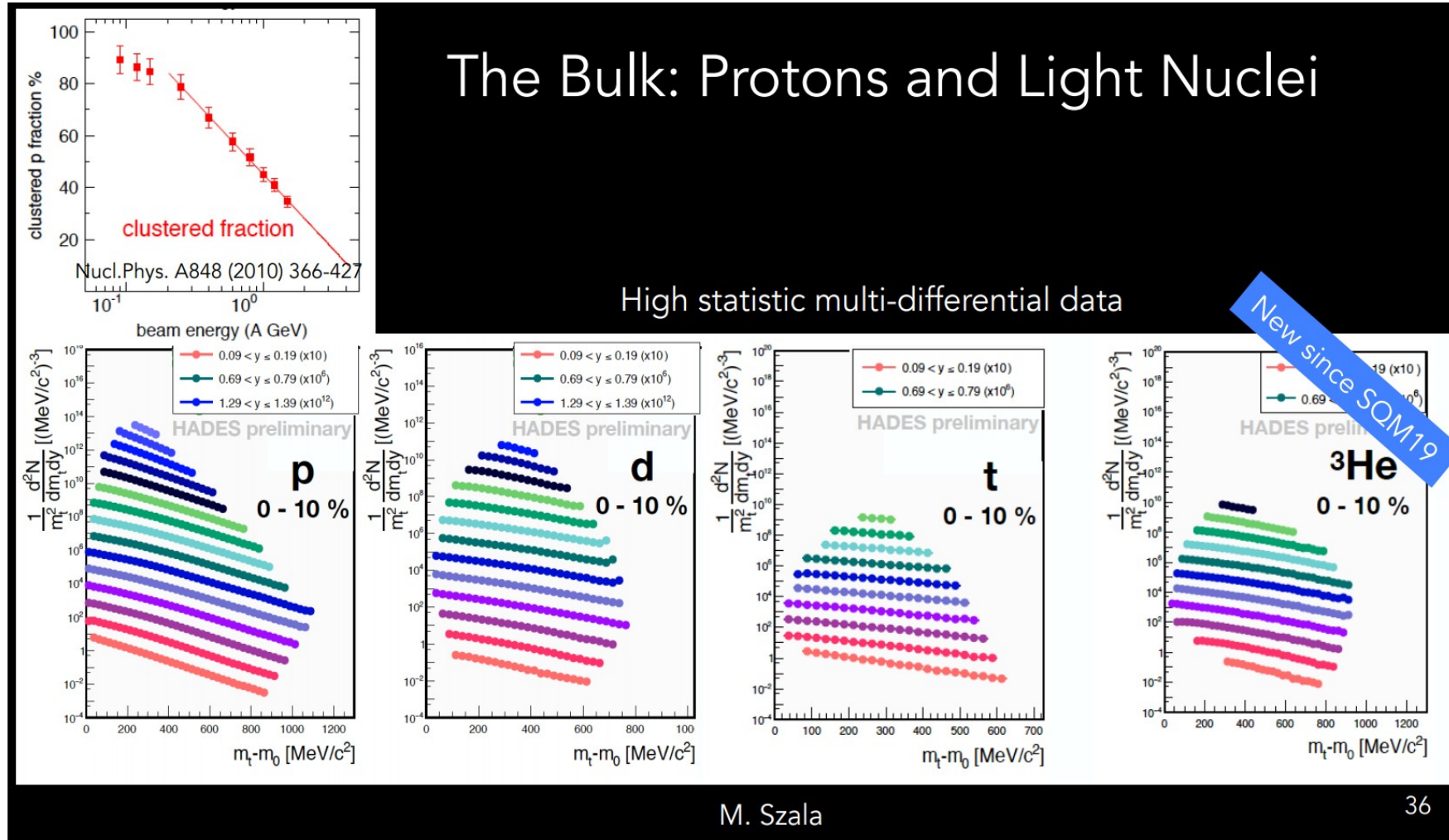
Coupling of  $K^-$  to baryons  
and strangeness exchange reactions  
e.g.  $\pi Y \rightarrow NK^-$ .



Repulsive  $K^+N$  potential  
Attractive  $K^-N$  potential  $\rightarrow K^-$  condensation?  
(complicated form due to resonances)

T. Waas, M. Rho and W. Weise Nucl. Phys. A617 (1997) 449-463  
N. Kaiser et al. Nucl. Phys. A594 (1995) 325-345

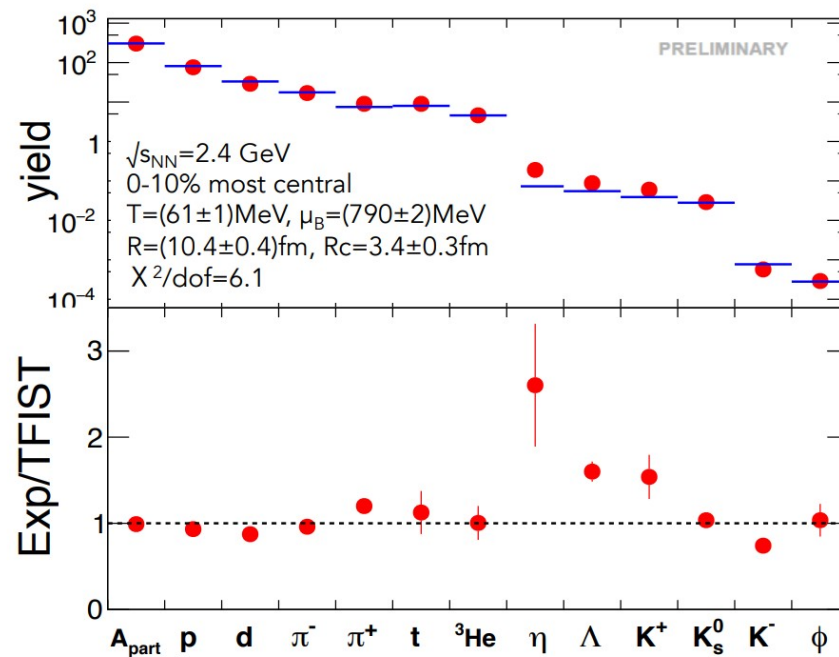
# Back to particle production...



# Thermal model applied to HADES data

## Macroscopic description of yields

Thermal FIST: V. Vovchenko H. Stoecker, Comput. Phys. Commun. 244 (2019) 295.



Momentum distribution not isotropic even in most central events.

Fit to HADES data consistent with previous works when same selections of hadron species are used (p, d,  $\pi$ ,  $K^+$ )

J. Cleymans, H. Oeschler, K. Redlich, Phys.Rev. C59 (1999)

R. Auerbeck, R. Holzmann, V. Metag, R.S. Simon. Phys.Rev. C67 (2003)

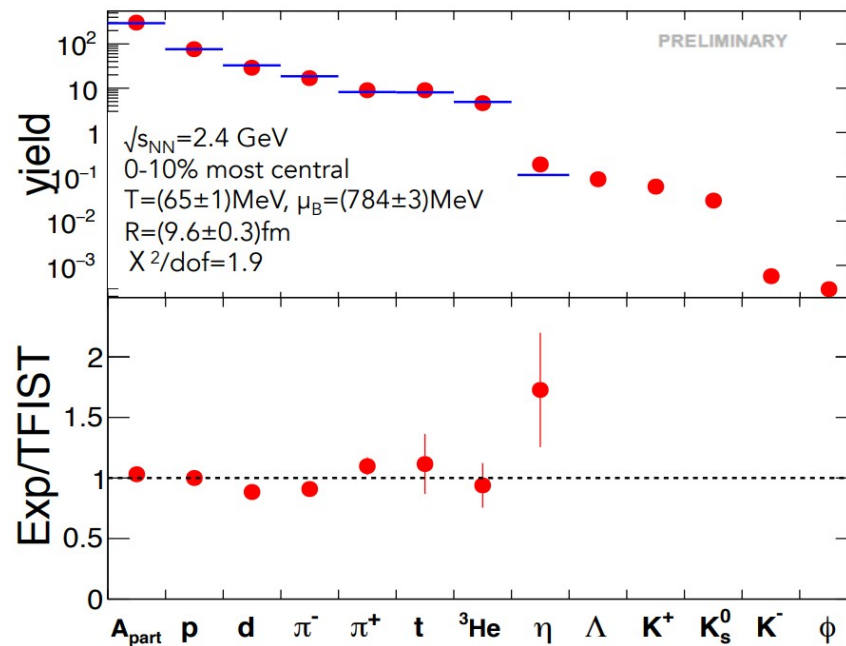
Fit to full hadron spectrum results in large  $\chi^2$ !

Inclusion of excited nuclei:  
As proposed by e.g. E. Shuryak.  
→ Small improvement in  $\chi^2$ .

# Thermal model applied to HADES data

## Macroscopic description of yields

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Fit excluding strangeness and but including excited nuclei results in small  $\chi^2$ !