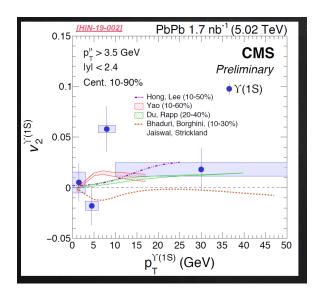
## Bottomonium elliptic flow from anisotropic escape

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#### Recent results from CMS



#### The model

- aHydro output for 5.02 TeV Pb+Pb collisions from KSU.
- We use Optical Glauber model for spatial initialization of the fireball energy density profile.
- The spatial distribution of the bottomonium production points in the transverse plane is assumed to follow the number of binary collisions.
- We consider transverse momentum  $(p_T)$  distribution of the  $\Upsilon$ 's obtained from PYTHIA simulations for p+p collisions, scaled by the mass number of the colliding nuclei [K. Zhou, N. Xu and P. Zhuang, Nucl. Phys. A 931, 654.].
- We adopt the isotropic thermal decay widths Γ(T) of the bottomonium states from: [M. Strickland and D. Bazow, Nucl. Phys. A 879, 25]
- Then we allow the bottomonium states to propagate through the medium and decay/survive according to their widths.

#### Formation time of bottomonium states

- The formation of each bound bottomonium state requires a finite formation time  $\tau_{\rm form}$ .
- $\bullet$  The value  $\tau_{\rm form}^0$  of the latter in the bottomonium rest frame is assumed to be proportional to the inverse of the vacuum binding energy for each state.
- For the  $\Upsilon(1S)$ ,  $\Upsilon(2S)$ ,  $\Upsilon(3S)$ ,  $\chi_b(1P)$  and  $\chi_b(2P)$  states we use  $\tau_{\rm form}^0 = 0.2, 0.4, 0.6, 0.4, 0.6$  fm/c, respectively [B. Krouppa and M. Strickland, arXiv:1605.03561].
- In the laboratory frame, relative to which a bottomonium state with mass M has transverse momentum  $p_T$ , the formation time becomes  $\tau_{\rm form} = E_T \tau_{\rm form}^0/M$  with  $E_T = \sqrt{p_T^2 + M^2}$ .

Strickland, PRC 100, 051901 (2019)].

- We assume that the bottomonium states propagate quasi freely following nearly straight-line trajectories.
- They are either dissociated (absorbed) or survive, no elastic scattering.
- If  $(x_0, y_0)$  denotes the position in the transverse plane where a bottomonium with momentum  $p_T$  is at time  $\tau_i$ , it will at a later time  $\tau = \tau_i + \tau'$  be at

$$x' = x_0 + v_T \tau' \cos \phi \quad , \quad y' = y_0 + v_T \tau' \sin \phi,$$

where  $v_T = p_T/E_T$  is the bottomonium transverse velocity and  $\phi$  the azimuthal angle of its transverse momentum.

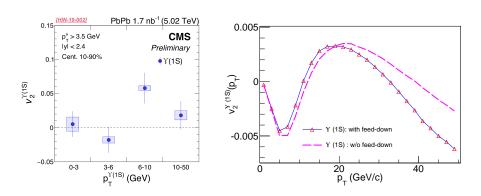
### Dissociation and width of bottomonium

- From hydro output we can get the temperature experienced by the  $b\bar{b}$  pair at (x',y') at  $\tau'$ .
- For this temperature, we obtain the thermal decay widths  $\Gamma(T(x',y',\tau'))$  of the bottomonium states, adopting the recent state-of-the-art results of in-medium dissociation of different bound  $b\bar{b}$  states [M. Strickland and D. Bazow, Nucl. Phys. A 879, 25].
- ullet The final transmittance for a  $bar{b}$  bound state labelled by j is given by

$$\mathcal{T}_{j}(x, y, p_{T}, \phi) = \exp \left[ -\Theta(\tau_{f} - \tau_{j}^{\text{form}}) \int_{\max(\tau_{j}^{\text{form}}, \tau_{i})}^{\tau_{f}} d\tau' \, \Gamma_{j} \left( T(x', y'; \, \tau') \right) \right],$$

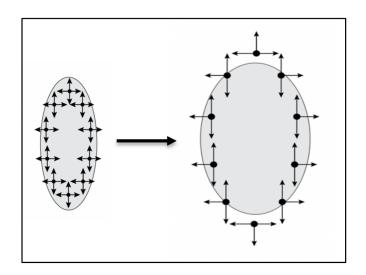
where  $\Theta$  is the usual step function.

# Estimation of $R_{AA}(p_T)$ and $v_2(p_T)$ for $\Upsilon(1S)$ states

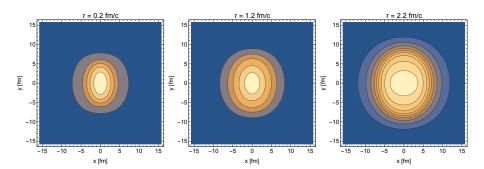


- Feed down from excited states is included.
- Excited states are produced late in the plasma frame.

## Expansion geometry



## Hydrodynamic evolution



### Effect of time dilation

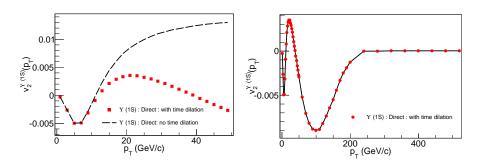
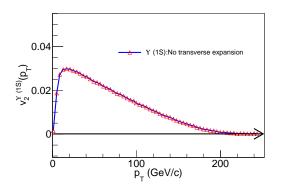


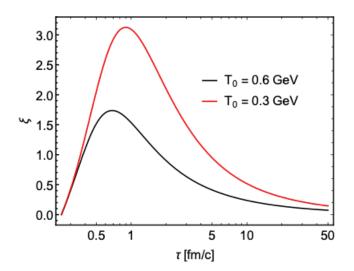
Figure:  $p_T$  dependence of  $v_2$  for directly produced  $\Upsilon(1S)$  states with and w/o time dilation effect. Feed down correction not included.

## Effect of transverse expansion



- Transverse expansion is turned off; initial profile from aHydro.
- No negative  $v_2$  observed. At very high  $p_T$  bottomonia are formed late in the plasma frame, and eventually  $v_2 \to 0$ .
- Temperature is scaled by time dependent central temperature:  $T(\tau, x, y, z = 0) = T(\tau_0, x, y) \times \frac{T(\tau, 0, 0, 0)}{T_0}$

## Backup 1



## Backup 2

