

# Bottomonium thermal limit in the KSU approach

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Suppression and (re)generation of quarkonium in heavy-ion collisions at the LHC

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**ENERGY**

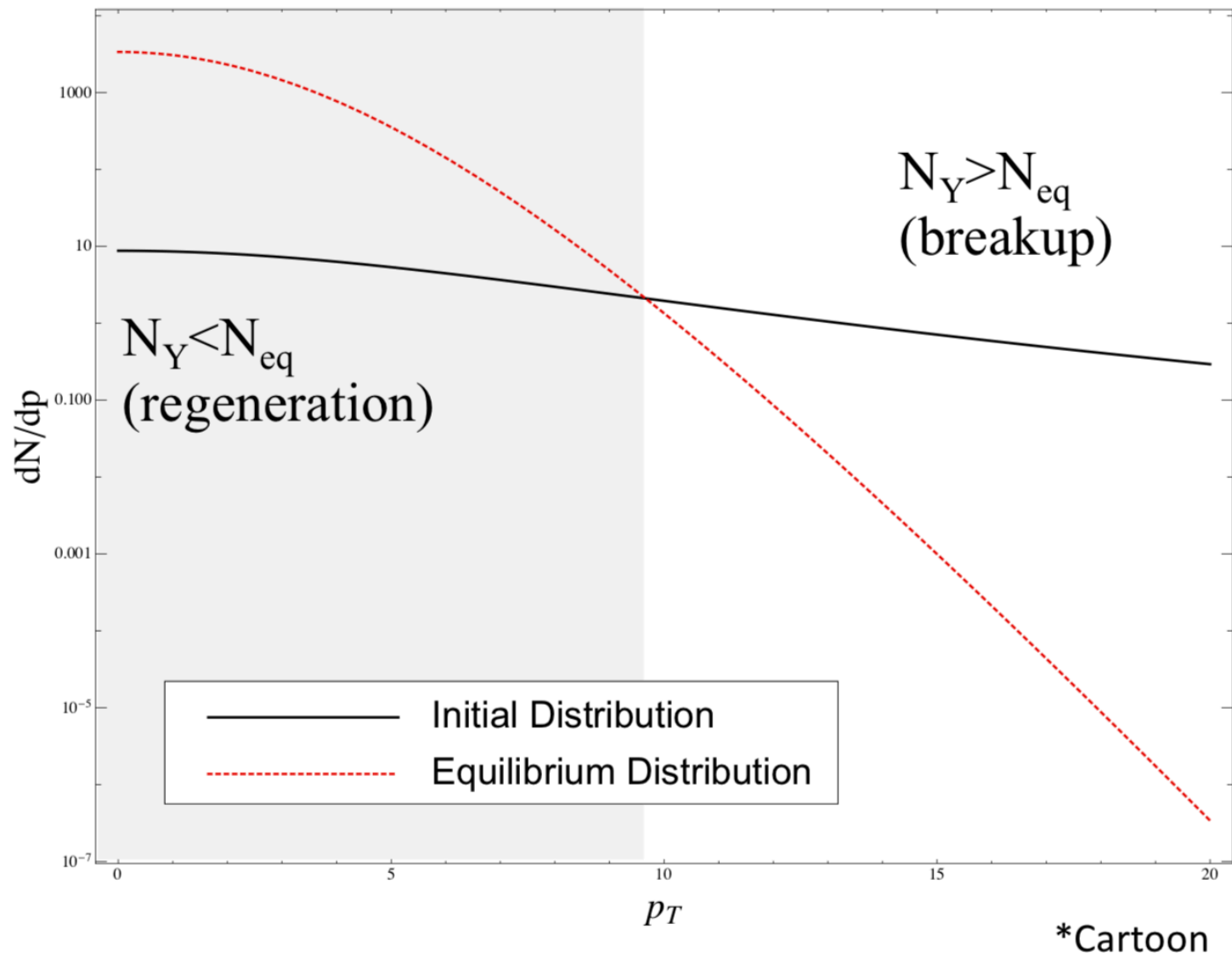
# Including regeneration

- In my former student Brandon's dissertation he included regeneration by generalizing the TAMU prescription X. Du, R. Rapp, and M. He, (2017), [arXiv:1706.08670](https://arxiv.org/abs/1706.08670) [hep-ph].

$$\frac{dn(\tau, \mathbf{x})}{d\tau} = -\Gamma(\tau, \mathbf{x}) \left[ n(\tau, \mathbf{x}) - n_{\text{eq}}(T(\tau, \mathbf{x})) \right],$$

$$f_{\text{eq}} = \frac{3}{(2\pi)^3} \gamma^2 \text{Exp} \left[ -\cosh(y) \sqrt{p_T^2 + m^2} / T(\xi, \Lambda) \right],$$

- Once  $\sigma_{q\bar{q}}$  is scaled to heavy-ion collisions with  $N_{\text{coll}}$ , we divide by the total inelastic cross section to determine the number of  $q\bar{q}$  pairs one expects in a heavy-ion collision.



# Fugacity computation

- Within  $f_{\text{eq}}$  is a fugacity factor  $\gamma_q$  which dictates the amount of local regeneration in the plasma
- The first task is to derive  $\gamma_q$  as a function of the local effective temperature,  $T$  [Braun-Munzinger, Stachel, arXiv: 0007059 \[nucl-th\]](#)

$$N_{q\bar{q}} = \frac{1}{2} N_{\text{op}} \frac{I_1(N_{\text{op}})}{I_0(N_{\text{op}})} + N_{\text{hid}},$$

$$N_{\text{op}} = \gamma_q V_{\text{coll}} d_b \int \frac{d^3p}{(2\pi)^3} f^q(p; T),$$

$$N_{\text{hid}} = \gamma_q^2 V_{\text{coll}} \sum_{\text{states}} d_{\text{states}} \int \frac{d^3p}{(2\pi)^3} f^{\text{states}}(p; T),$$

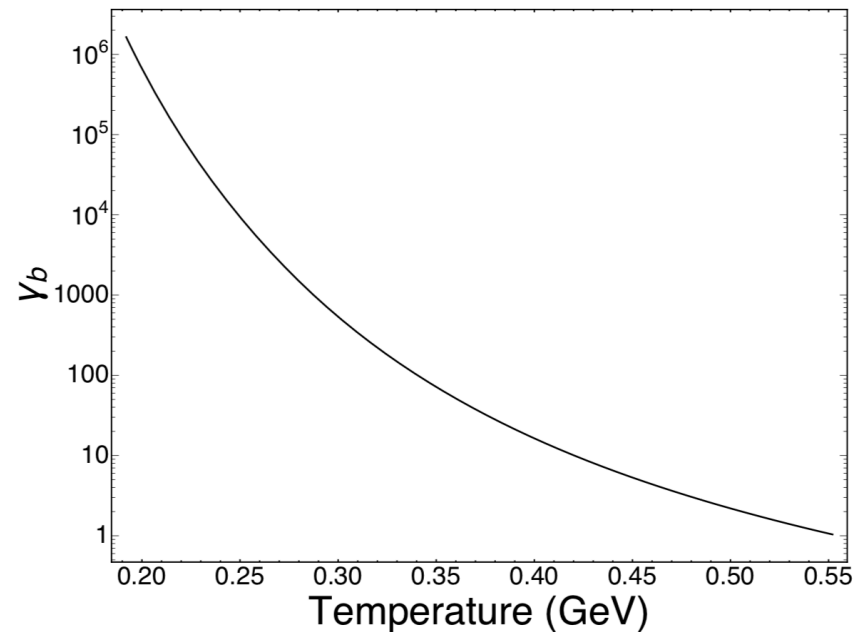
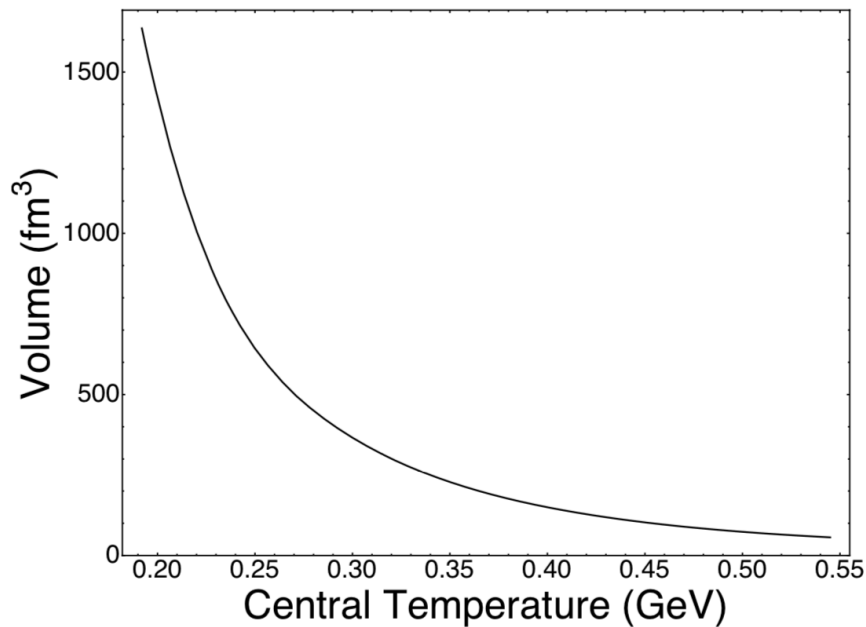
Full Cross Section $\sigma_{b\bar{b}}$	0.2 TeV	2.76 TeV	5.02 TeV
pp $\rightarrow b\bar{b} + X$ [ $\mu\text{b}$ ] ( $ y  < 0.5$ )	0.92	-	-
pp $\rightarrow b\bar{b} + X$ [ $\mu\text{b}$ ] ( $ y  < 2.4$ )	-	82.56	156.96
pp $\rightarrow b\bar{b} + X$ [ $\mu\text{b}$ ] ( $2.5 < y < 4.0$ )	-	12.90	24.60

TABLE I.  $b\bar{b}$  production cross sections for various experimental windows in rapidity.

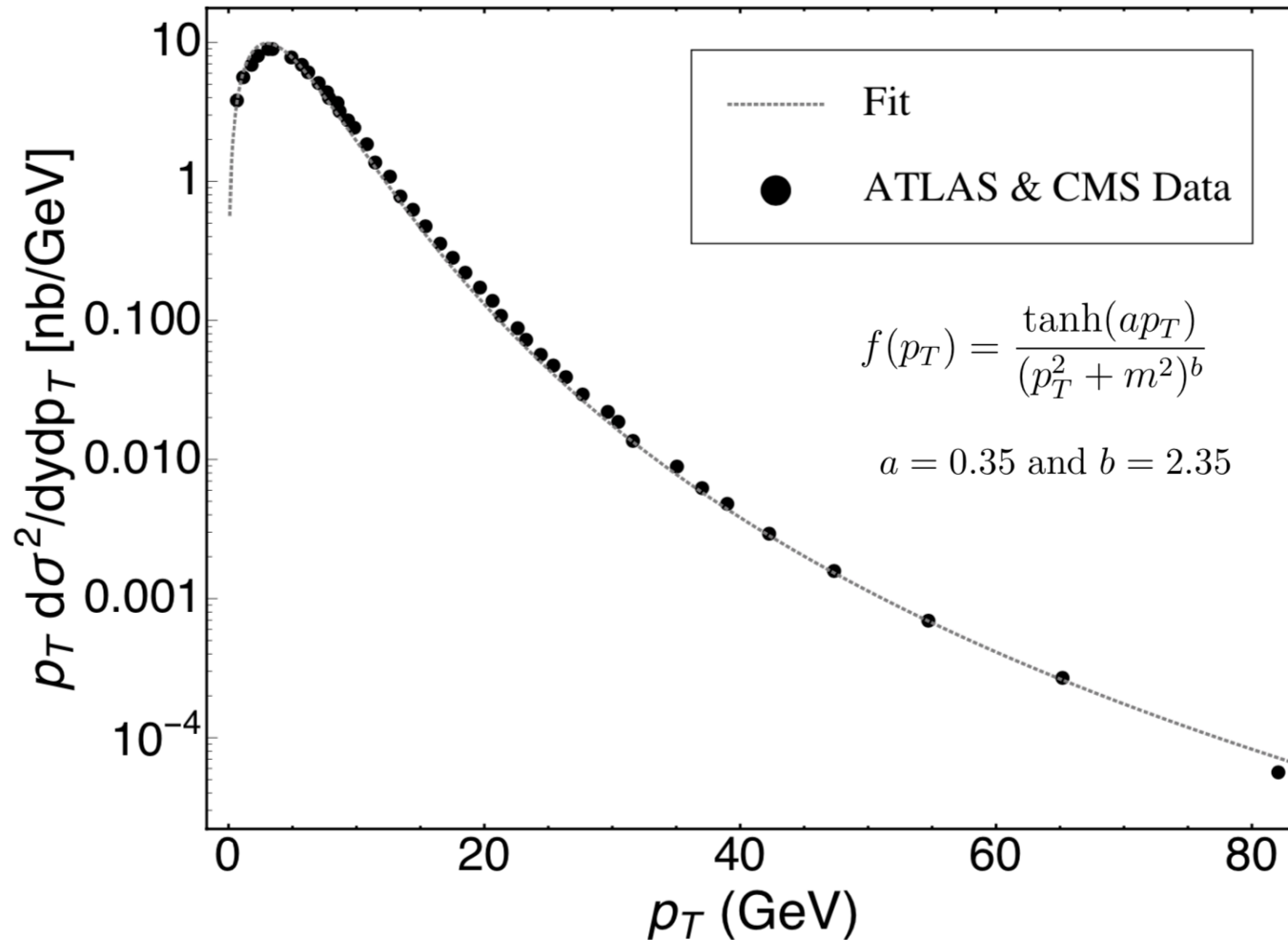
# Fugacity computation example

$$N_{q\bar{q}} = \frac{1}{2} N_{\text{op}} \frac{I_1(N_{\text{op}})}{I_0(N_{\text{op}})} + N_{\text{hid}},$$

$$N_{\text{op}} = \gamma_q V_{\text{coll}} d_b \int \frac{d^3 p}{(2\pi)^3} f^q(p; T), \quad N_{\text{hid}} = \gamma_q^2 V_{\text{coll}} \sum_{\text{states}} d_{\text{states}} \int \frac{d^3 p}{(2\pi)^3} f^{\text{states}}(p; T),$$



# Initial conditions 1

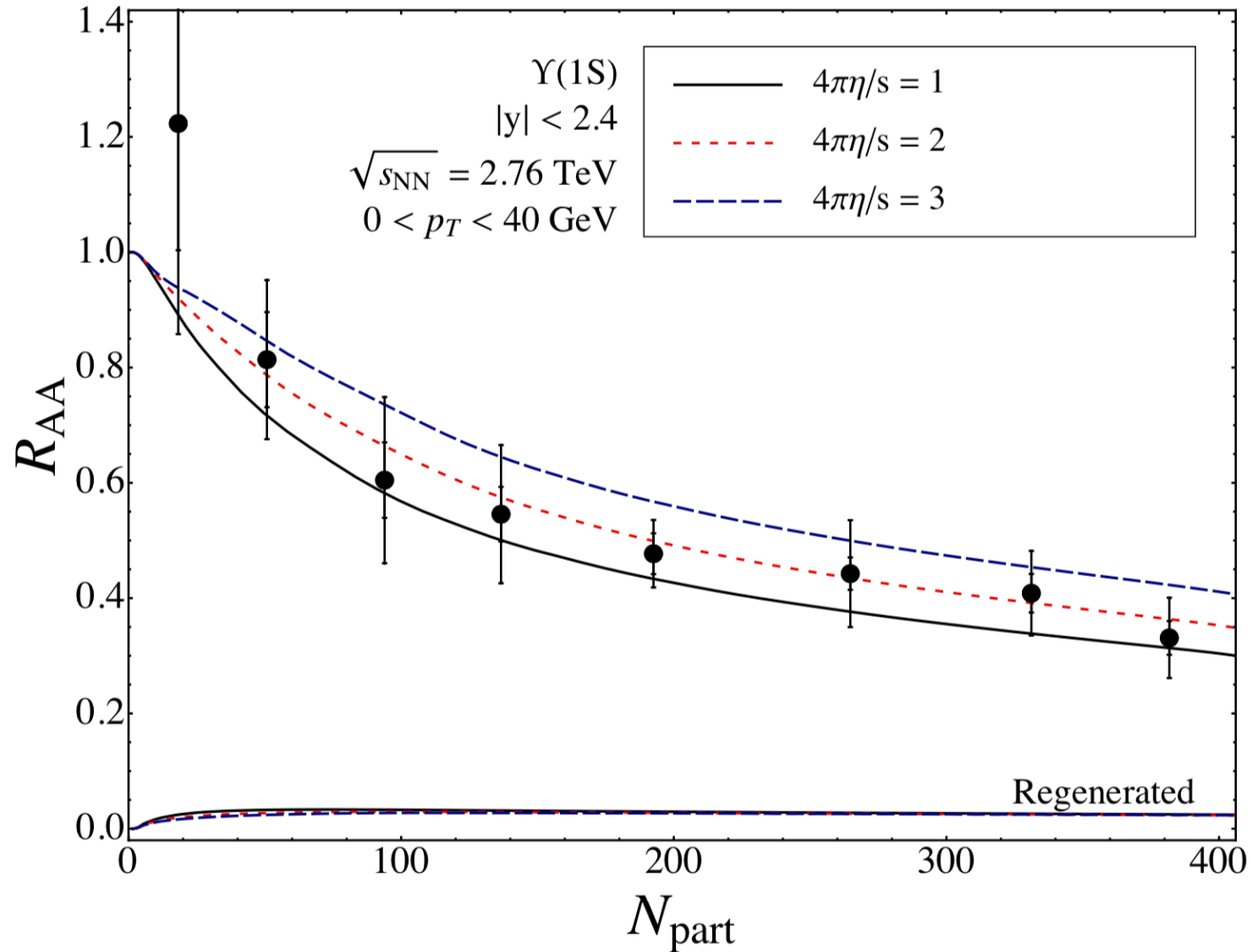


# Initial conditions 2

Full Cross Section	0.2 TeV	2.76 TeV	5.02 TeV
$pp \rightarrow \Upsilon(1S) ( y  < 0.5) \text{ [nb]}$	2.35	-	-
$pp \rightarrow \Upsilon(1S) ( y  < 2.4) \text{ [nb]}$	-	145.44	276.48
$pp \rightarrow \Upsilon(1S) (2.5 < y < 4.0) \text{ [nb]}$	-	22.65	43.20
$pp \rightarrow \Upsilon(2S) ( y  < 0.5) \text{ [nb]}$	0.77	-	-
$pp \rightarrow \Upsilon(2S) ( y  < 2.4) \text{ [nb]}$	-	48.00	91.20
$pp \rightarrow \Upsilon(2S) (2.5 < y < 4.0) \text{ [nb]}$	-	7.50	14.25
$pp \rightarrow \Upsilon(3S) ( y  < 0.5) \text{ [nb]}$	0.13	-	-
$pp \rightarrow \Upsilon(3S) ( y  < 2.4) \text{ [nb]}$	-	0.01	0.02
$pp \rightarrow \Upsilon(3S) (2.5 < y < 4.0) \text{ [nb]}$	-	0.002	0.003
$pp \rightarrow \chi_b(1P) ( y  < 0.5) \text{ [nb]}$	0.40	-	-
$pp \rightarrow \chi_b(1P) ( y  < 2.4) \text{ [nb]}$	-	24.72	47.00
$pp \rightarrow \chi_b(1P) (2.5 < y < 4.0) \text{ [nb]}$	-	3.85	7.34
$pp \rightarrow \chi_b(2P) ( y  < 0.5) \text{ [nb]}$	0.12	-	-
$pp \rightarrow \chi_b(2P) ( y  < 2.4) \text{ [nb]}$	-	7.27	13.82
$pp \rightarrow \chi_b(2P) (2.5 < y < 4.0) \text{ [nb]}$	-	1.13	2.16
$pp \rightarrow \chi_b(3P) ( y  < 0.5) \text{ [nb]}$	0.02	-	-
$pp \rightarrow \chi_b(3P) ( y  < 2.4) \text{ [nb]}$	-	1.09	2.07
$pp \rightarrow \chi_b(3P) (2.5 < y < 4.0) \text{ [nb]}$	-	0.16	0.32

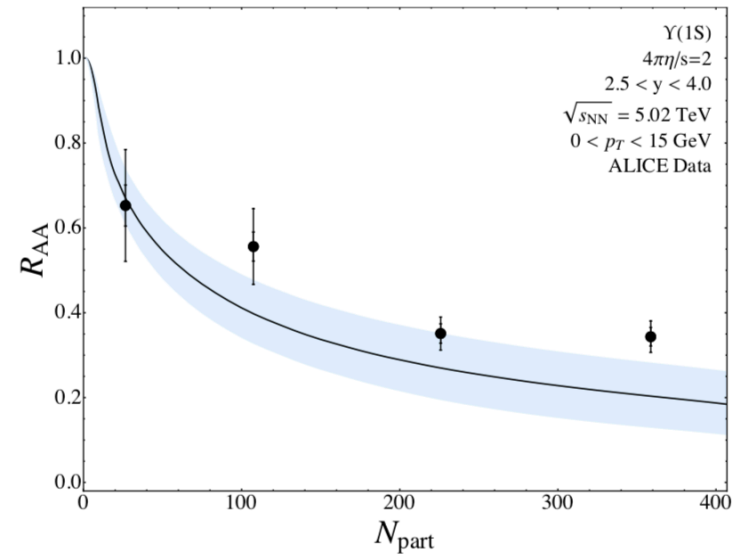
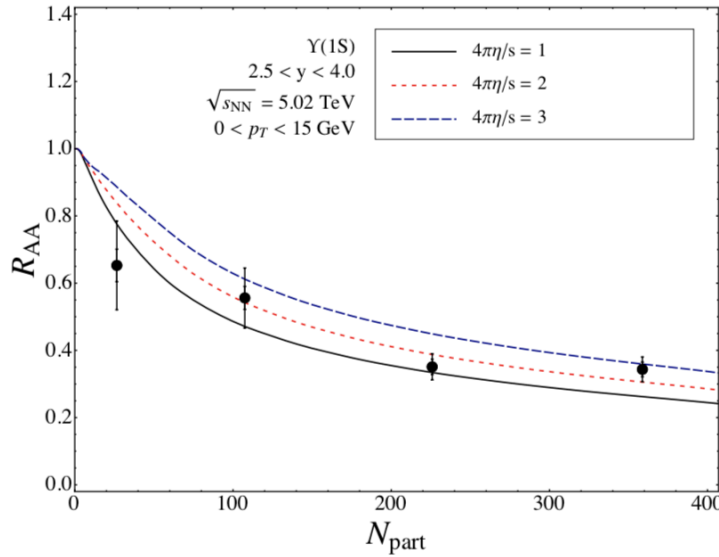
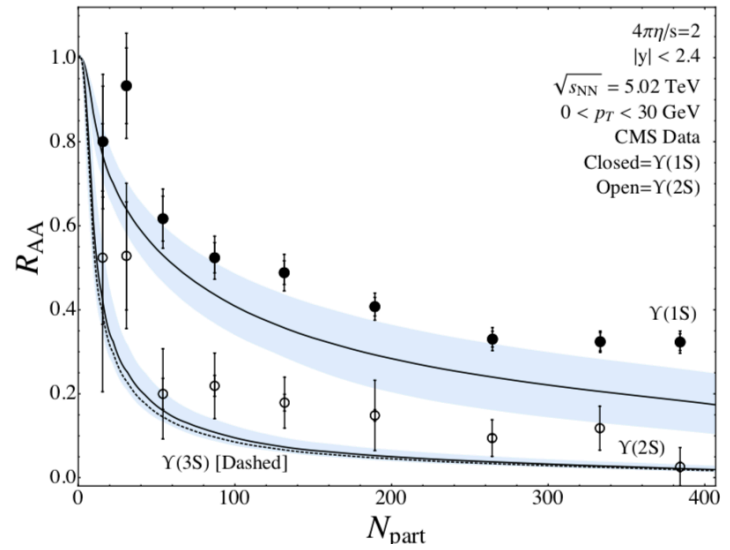
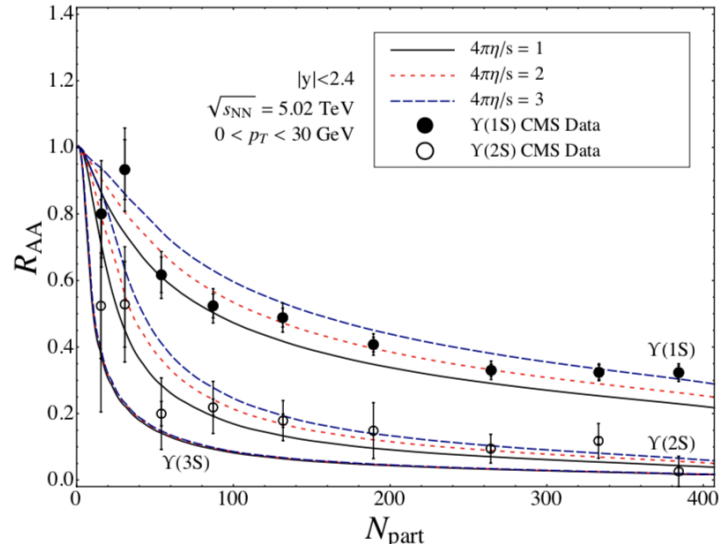
TABLE II. Bottomonium state cross sections for  $\Upsilon(nS)$  and  $\chi_b(mP)$  states in various experimental windows in rapidity.

# Including regeneration

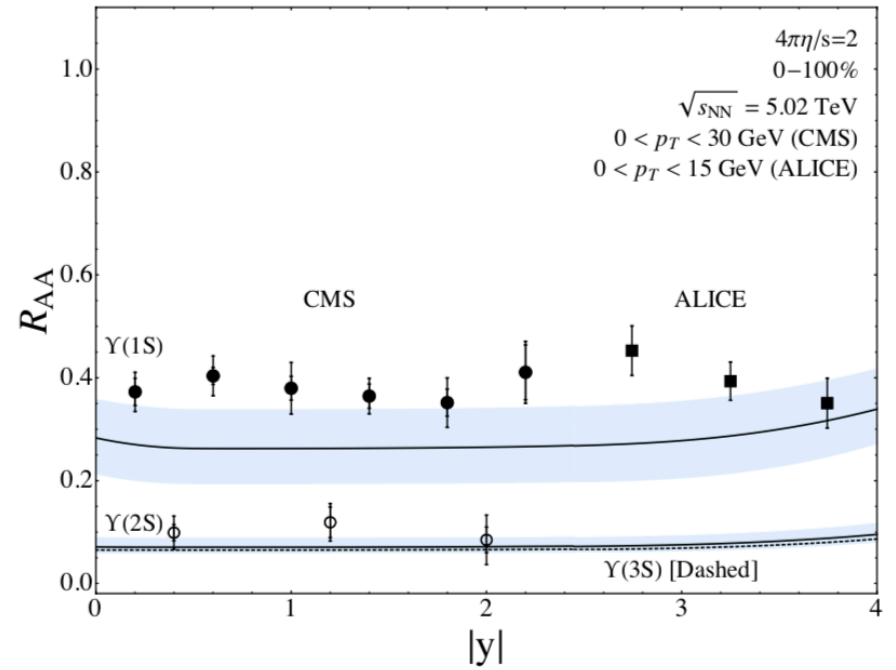
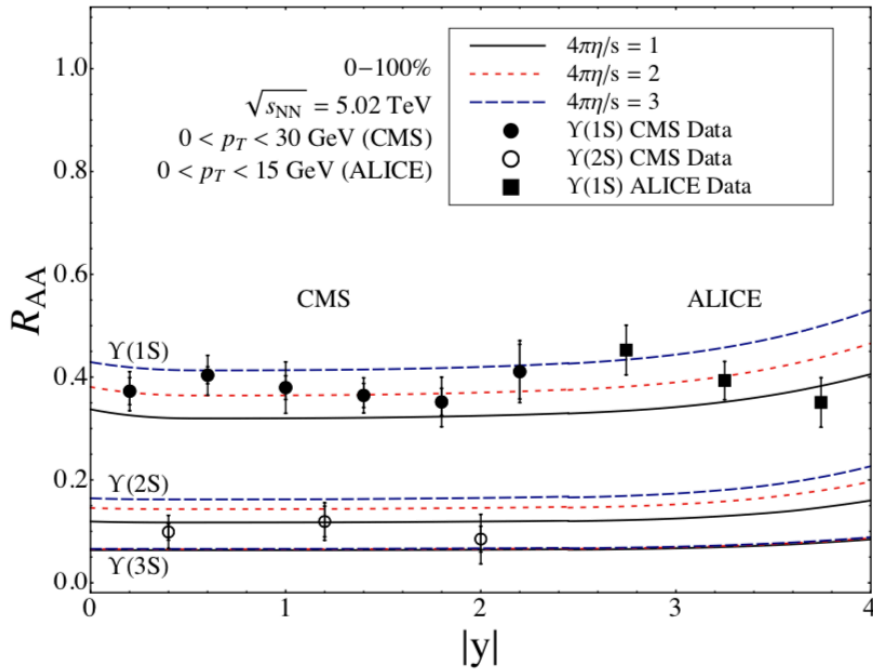




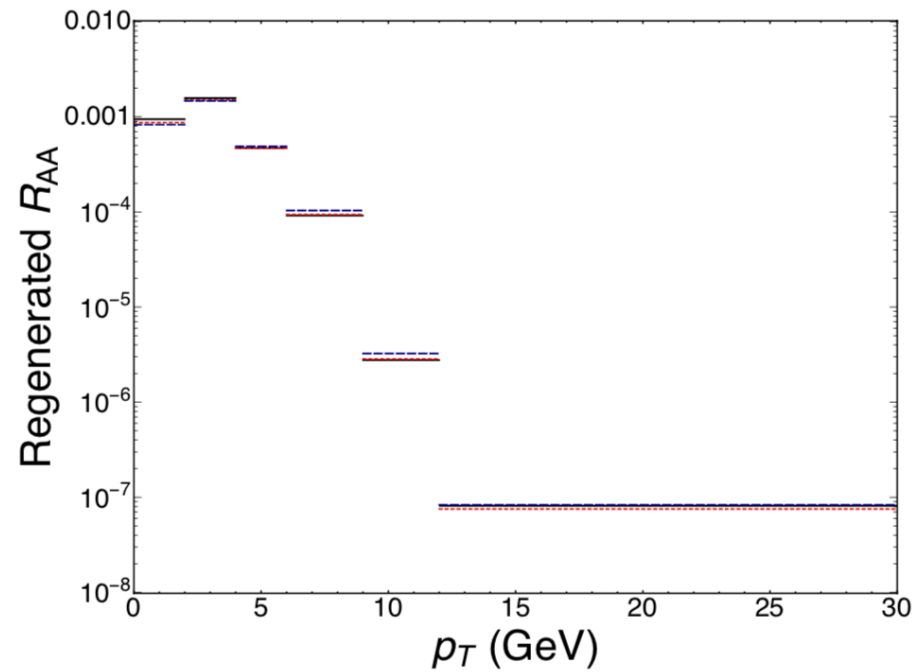
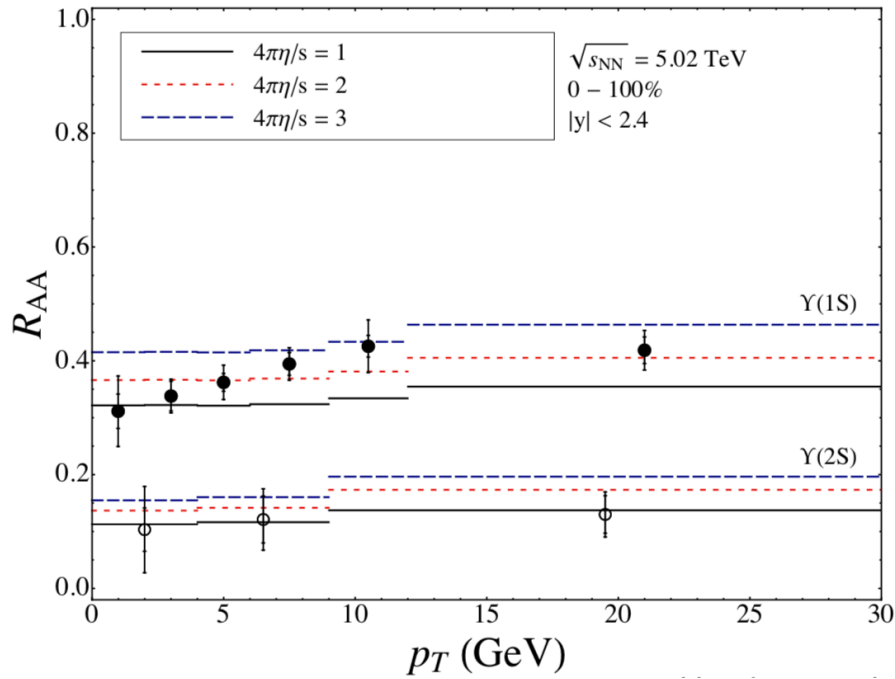
# Results 1



# Results 2



# Results 2 Continued



# Conclusions

- Bottomonium regeneration is small

# Charmonium results – 2.76 TeV Pb-Pb

- Raw suppression yields for three charmonium states
- Shielding of suppression due to large  $\tau_{\text{form}}$

$$\tau_{\text{form}} = (E_T/M) * \tau_{\text{form}}^0$$

$$p_{T,\text{min}} = 6.5 \text{ GeV}/c$$

$$\tau_{\text{form}}^0 = 0.8 \text{ fm}/c$$

QGP present after 1.7 fm/c

$$\tau_{\text{form}}^0 = 3.9 \text{ fm}/c$$

QGP not present after 7.9 fm/c

