Nuclear Energy Density Functionals



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Energy Density Functionals



Nuclear Energy Density Functionals: the many-body problem is mapped onto a one body problem without explicitly involving inter-nucleon interactions!

Self-consistent Kohn-Sham DFT: includes correlations and therefore goes beyond the Hartree-Fock. It has the advantage of being a **local scheme**.

$$v_s[\rho(\mathbf{r})] = v(\mathbf{r}) + U[\rho(\mathbf{r})] + v_{xc}[\rho(\mathbf{r})]$$

external potentialHartree termexchange-correlation $v_{xc}[\rho(\mathbf{r})] = \frac{\delta E_{xc}[\rho(\mathbf{r})]}{\delta \rho(\mathbf{r})}$

The practical usefulness of the Kohn-Sham scheme depends entirely on whether accurate approximations for E_{xc} can be found!

... accurate and controlled approximations for the nuclear exchange-correlation energy functional

... microscopic foundation for a universal EDF framework, related to and constrained by low-energy QCD

... extensions that include non-nucleonic degrees of freedom

... correlations related to restoration of broken symmetries and fluctuations of collective coordinates

Exchange-correlation functional:

The true E_{xc} is a universal functional of the density: it has the same functional form for all systems.





...vary smoothly with nucleon number! Can be included implicitly in an effective Energy Density Functional. A **microscopic nuclear energy density functional** must start from the relevant active degrees of freedom at low energy:

PIONS & NUCLEONS



Effective Field Theory of low-energy in-medium NN interactions \Rightarrow approximations to **the exact exchange-correlation functional.**



relevant scale: Fermi momentum

$$k_f \approx 2m_\pi << 4\pi f_\pi$$



The density functional involves an expansion of nucleon self-energies in **powers of the Fermi momentum.**



Inclusion of the $\Delta(1232)$ degree of freedom:

 $M_{\Delta} - M_N \approx k_f \approx 2m_{\pi}$

Two-pion exchange diagrams with single and double $\Delta(1232)$ excitations:



Model for Finite Nuclei

P. Finelli, N. Kaiser, D. Vretenar, W. Weise, Nucl. Phys. A 735 (2004) 449, A 770 (2006) 1.

... universal exchange-correlation functional $E_{xc}[\rho]$

Ist step: Local Density Approximation

$$E_{xc}^{LDA} \equiv \int \varepsilon^{ChPT} [\rho(\mathbf{r})] \rho(\mathbf{r}) d^3 r$$

2nd step: second-order gradient correction to the LDA

ChPT calculations for inhomogeneous nuclear matter:

$$\mathcal{E}(\rho, \nabla \rho) = \rho \overline{E}(k_f) + (\nabla \rho)^2 F_{\nabla}(k_f) + \dots$$



Charge form factors of ⁴⁸Ca, ⁹⁰Zr and ²⁰⁸Pb calculated with the FKVW functional, in comparison with the experimental form factors.

CHIRAL EFT provides a consistent microscopic framework for the isoscalar and isovector channels of a **universal nuclear energy density functional**.

Hypernuclear single particle spectra based on in-medium chiral SU(3) dynamics

P. Finelli, N. Kaiser, D. Vretenar, W. Weise, Nucl. Phys. A 831 (2009) 163.

Hypernuclear energy density functional:

$$E[\rho] = E^{N}[\rho] + E^{\Lambda}_{\text{free}}[\rho] + E^{\Lambda}_{\text{int}}[\rho]$$

hypernuclear ground state

$$E_{\rm free}^{\Lambda} = \int d^3 r \langle \Phi_0 | \bar{\psi}_{\Lambda} [-i \boldsymbol{\gamma} \cdot \boldsymbol{\nabla} + M_{\Lambda}] \psi_{\Lambda} | \Phi_0 \rangle$$

coupling to the scalar nucleon density

$$E_{\rm int}^{\Lambda} = \int d^3r \left\{ \underbrace{\Phi_0 | G_S^{\Lambda}(\rho) \left(\bar{\psi} \psi \right) \left(\bar{\psi}_{\Lambda} \psi_{\Lambda} \right) | \Phi_0 \right)}_{+ \langle \Phi_0 | G_V^{\Lambda}(\rho) \left(\bar{\psi} \gamma_{\mu} \psi \right) \left(\bar{\psi}_{\Lambda} \gamma^{\mu} \psi_{\Lambda} \right) | \Phi_0 \rangle} \right. \\ \left. + \langle \Phi_0 | D_S^{\Lambda} \partial_{\mu} (\bar{\psi} \psi) \partial^{\mu} (\bar{\psi}_{\Lambda} \psi_{\Lambda}) | \Phi_0 \rangle \right\}$$
coupling to the vector nucleon density

derivative term - from a gradient expansion of the EDF

Binding energies of the Λ in different s, p,... orbitals of six hypernuclei:



Collective correlations: restoration of broken symmetries and fluctuations of collective variables

Nikšić, Vretenar, Ring Phys. Rev. C **73**, 034308 (2006) Phys. Rev. C **74**, 064309 (2006)

- 1. Mean-field calculations, with a constraint on the quadrupole moment.
- 2. Angular-momentum and particle-number projection.
- 3. Generator Coordinate Method ⇒ configuration mixing



... larger variational space for projected GCM calculations!



Five-dimensional collective Hamiltonian

Nikšić, Li, Vretenar, Prochniak, Meng, Ring, Phys. Rev. C 79, 034303 (2009)

... nuclear excitations determined by quadrupole vibrational and rotational degrees of freedom

$$\begin{aligned} H_{\rm coll} &= \mathcal{T}_{\rm vib}(\beta,\gamma) + \mathcal{T}_{\rm rot}(\beta,\gamma,\Omega) + \mathcal{V}_{\rm coll}(\beta,\gamma) \\ \mathcal{T}_{\rm vib} &= \frac{1}{2} B_{\beta\beta} \dot{\beta}^2 + \beta B_{\beta\gamma} \dot{\beta} \dot{\gamma} + \frac{1}{2} \beta^2 B_{\gamma\gamma} \dot{\gamma}^2 \\ \mathcal{T}_{\rm rot} &= \frac{1}{2} \sum_{k=1}^3 \mathcal{I}_k \omega_k^2 \end{aligned}$$



Li, Nikšić, Vretenar, Meng, Lalazissis, Ring, Phys. Rev. C 79, 054301 (2009)

EDF description of nuclear Quantum Phase Transitions

Nikšić, Vretenar, Lalazissis, Ring, Phys. Rev. Lett. **99**, 092502 (2007)



... detailed spectroscopy in the EDF framework!

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