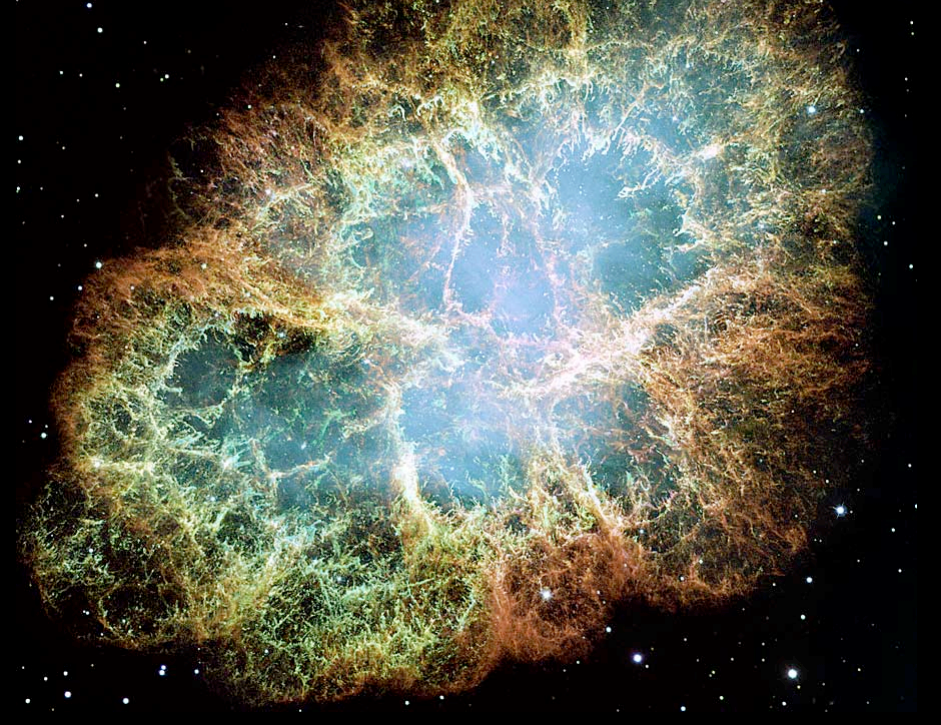


# Origin of the LEPP nuclei in supernovae

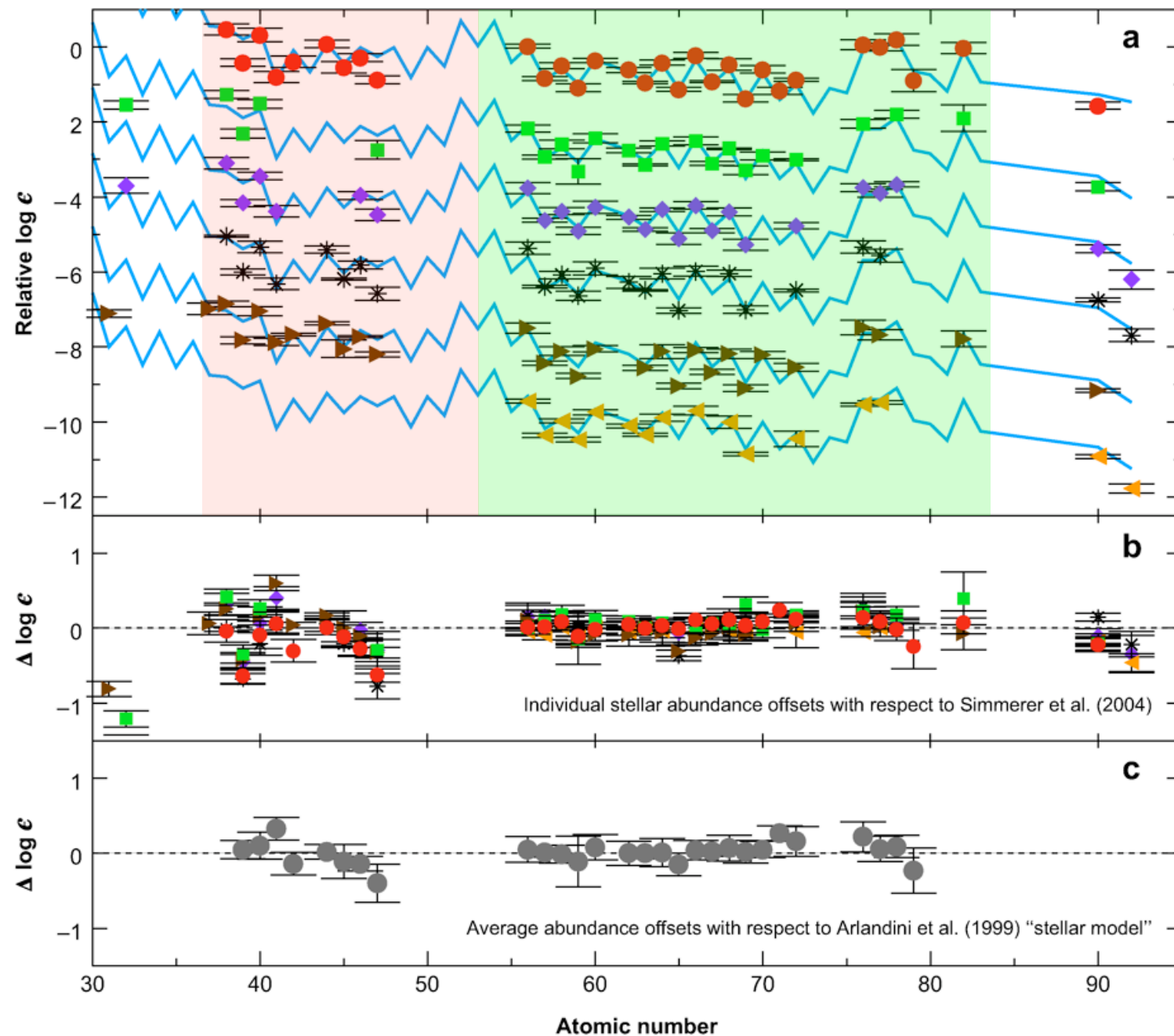


Almudena Arcones (Uni Basel) & Fernando Montes (NSCL/JINA)

arXiv:1007.1275v1

EMMI workshop:  
Neutron Matter in Astrophysics:  
From Neutron Stars to the r-Process

# Observations



Abundances of “r-process” elements:  
r-process-rich galactic halo (old) stars  
vs. [Solar system](#) abundances  
(r-process only)

Only few nucleosynthesis events have  
contributed to the abundances present  
in old stars.

Robust r-process for  $56 < Z < 83$   
but some scatter for  $Z < 47$

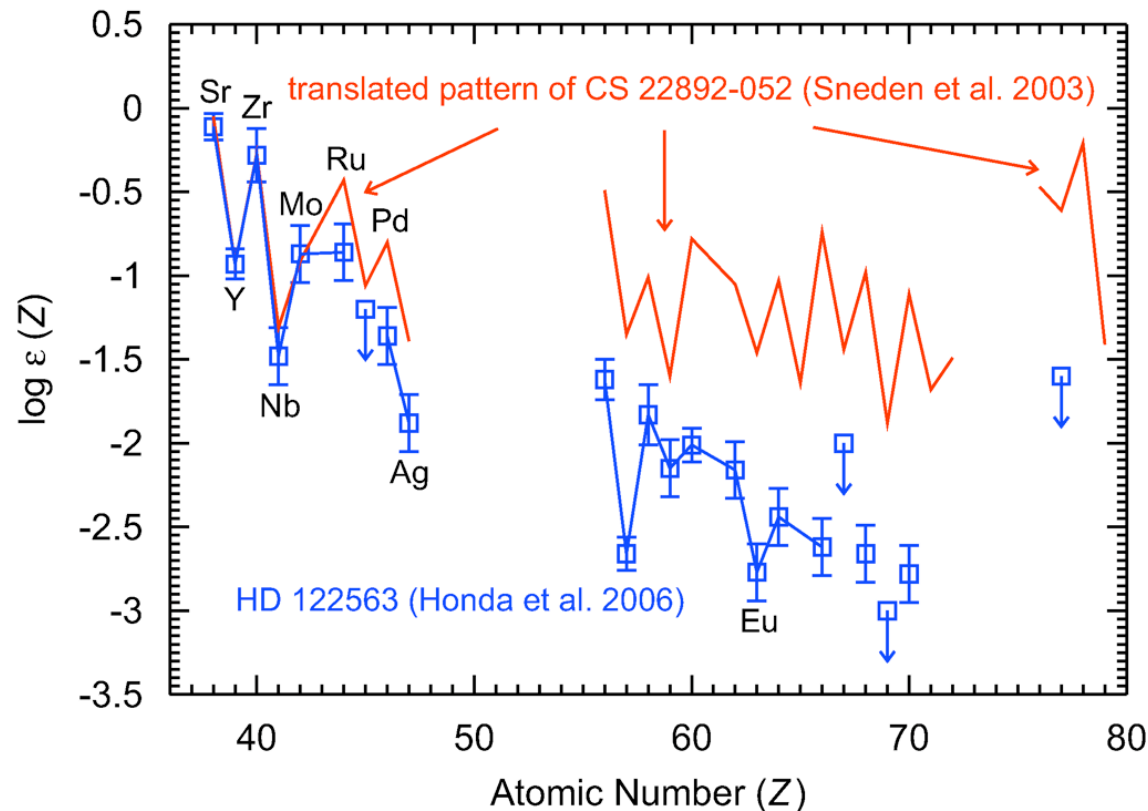
Suggestive of two components or sites:  
Qian & Wasserburg 2001...,  
Truran et al. 2002, Travaglio et al. 2004,  
Aoki et al. 2005, Otsuki et al. 2006.

- CS 22892-052: Sneden et al. (2003)
- HD 115444: Westin et al. (2000)
- ◆ BD+17°324817: Cowan et al. (2002)
- \* CS 31082-001: Hill et al. (2002)
- ▶ HD 221170: Ivans et al. (2006)
- ◀ HE 1523-0901: Frebel et al. (2007)

from Sneden, Cowan, Gallino 2008

# Two components of heavy element nucleosynthesis

Qian & Wasserburg: developed a model based on stars with high and low enrichment of heavy r-nuclei.



- In neutrino-driven winds when a neutron star forms, charged-particle reactions (CPR) produce nuclei with  $A \sim 90-110$  ( $Z < 47$ ).
- Observations of low-metallicity stars show that sites producing heavy r-nuclei do not produce Fe or any other elements between N and Ge. This suggests that heavy r-nuclei with  $A > 130$  ( $56 < Z < 83$ ) cannot be produced in every neutrino-driven wind.

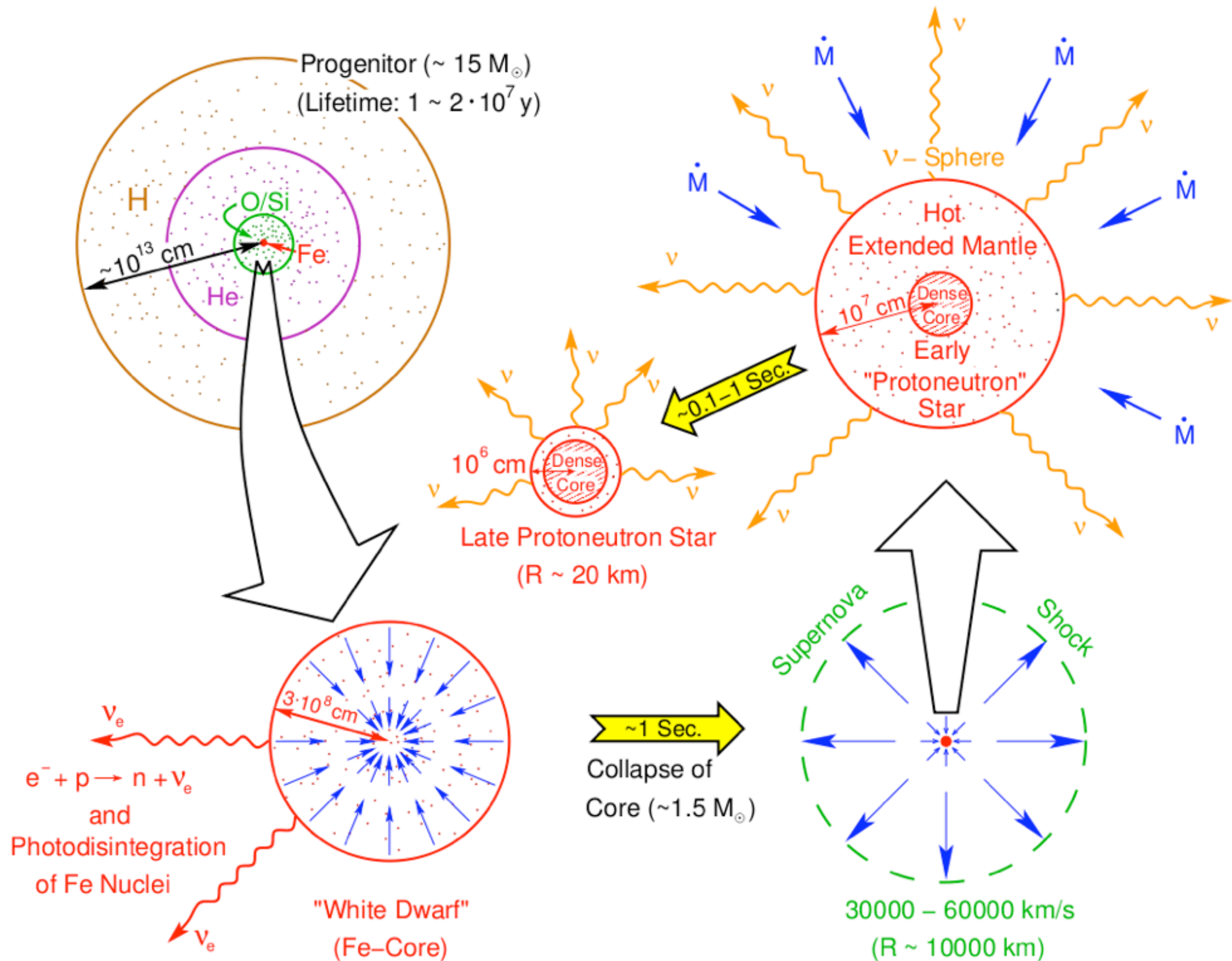
Travaglio et al 2004: Light Element Primary Process: LEPP = solar – r-process – s-process

Montes et al. 2007: LEPP creates a uniform and unique pattern

Can this be confirmed by state-of-the-art neutrino-driven wind simulations?

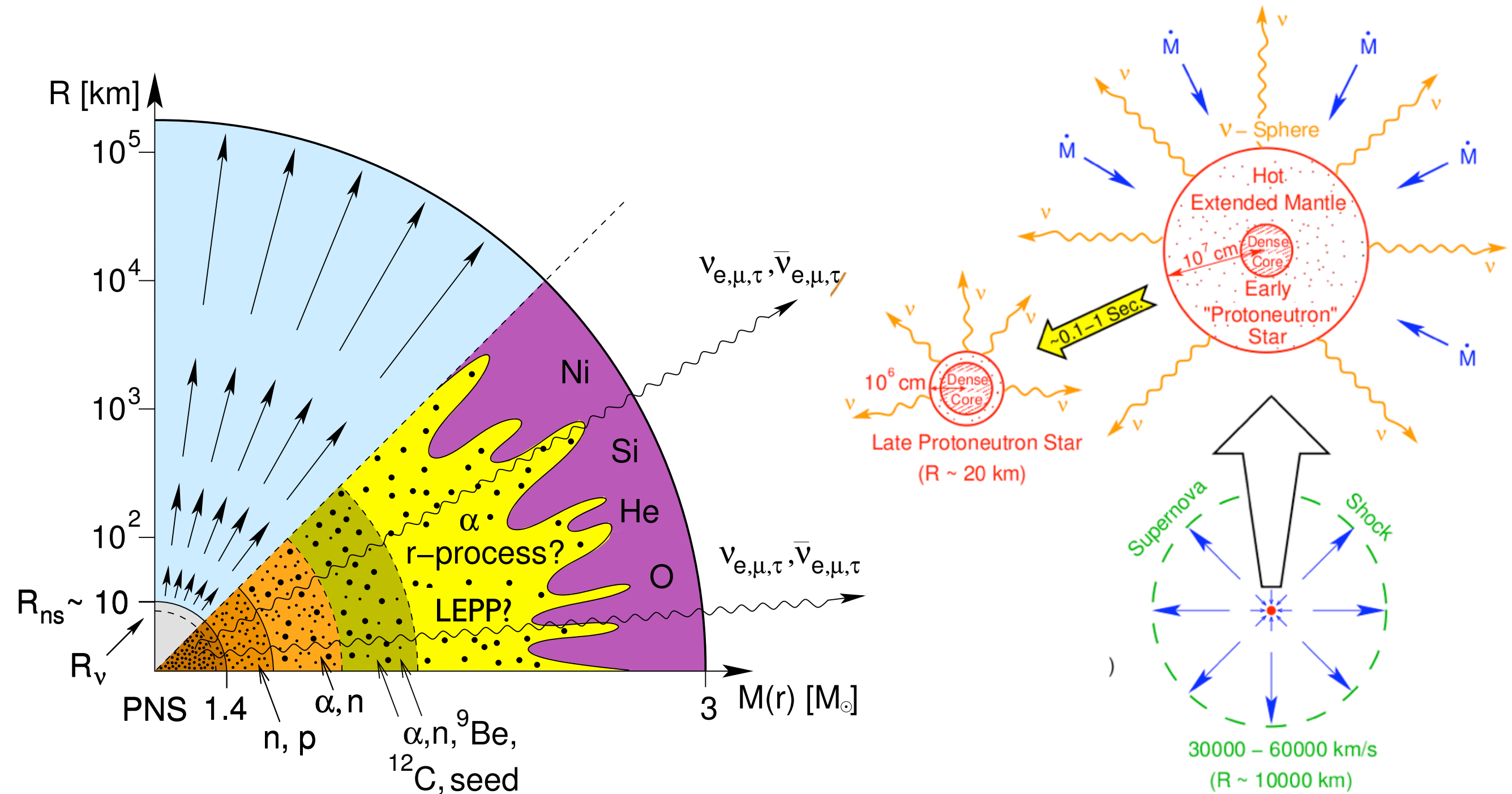
Do supernovae produce the LEPP pattern?

# Core-collapse supernovae





# Neutrino-driven wind



# Neutrino-driven wind simulations and nucleosynthesis networks

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## Simulations of core-collapse supernovae and the subsequent neutrino-driven winds

Problems: - explosion mechanism

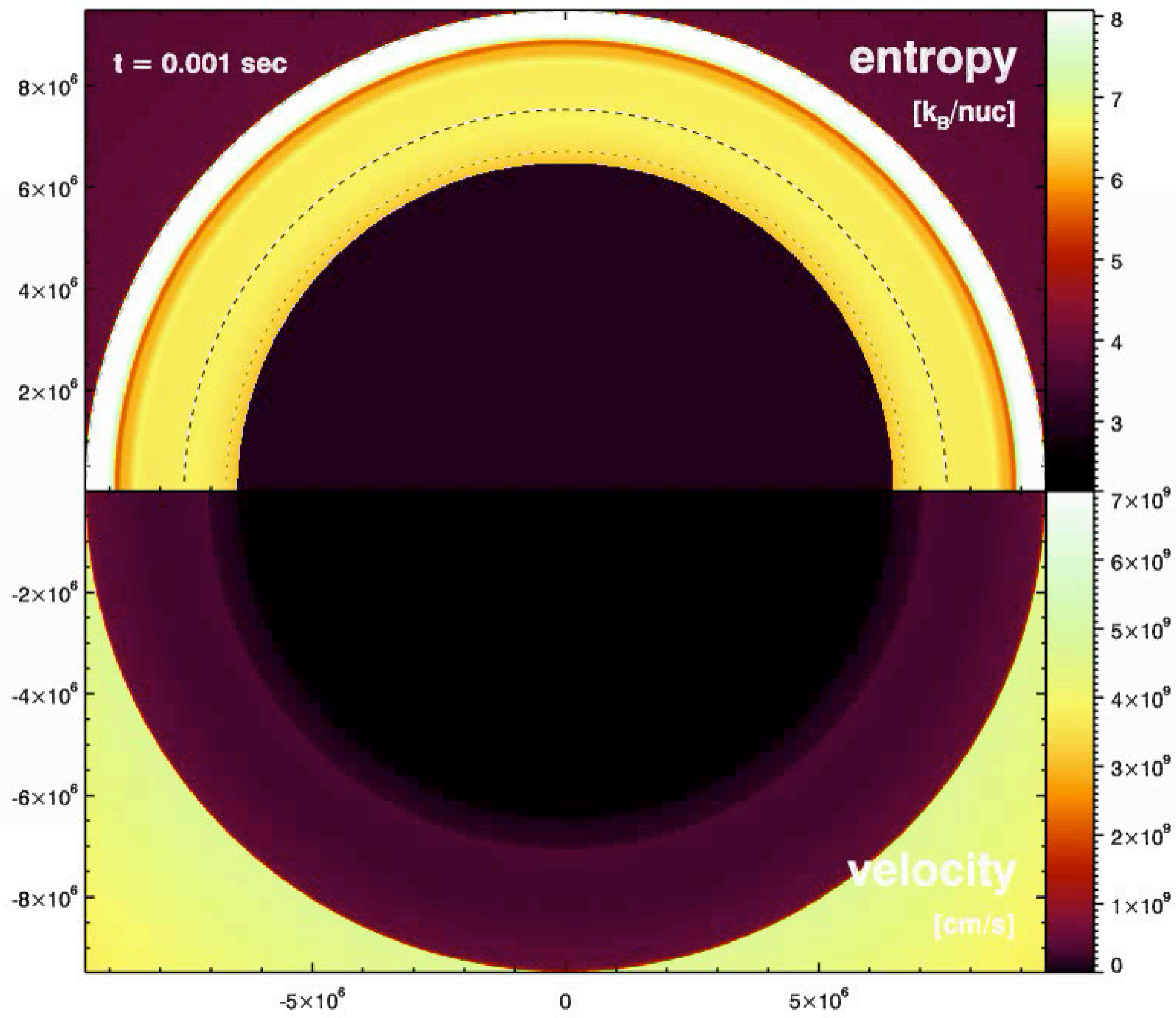
- simulations are computationally very expensive to follow the wind phase

Solutions: - steady-state wind models (Otsuki et al. 2000, Thompson et al. 2001, Wanajo)

- one-dimensional simulations with an artificial explosion  
(Arcones et al. 2007 (also 2d), Fischer et al. 2009)

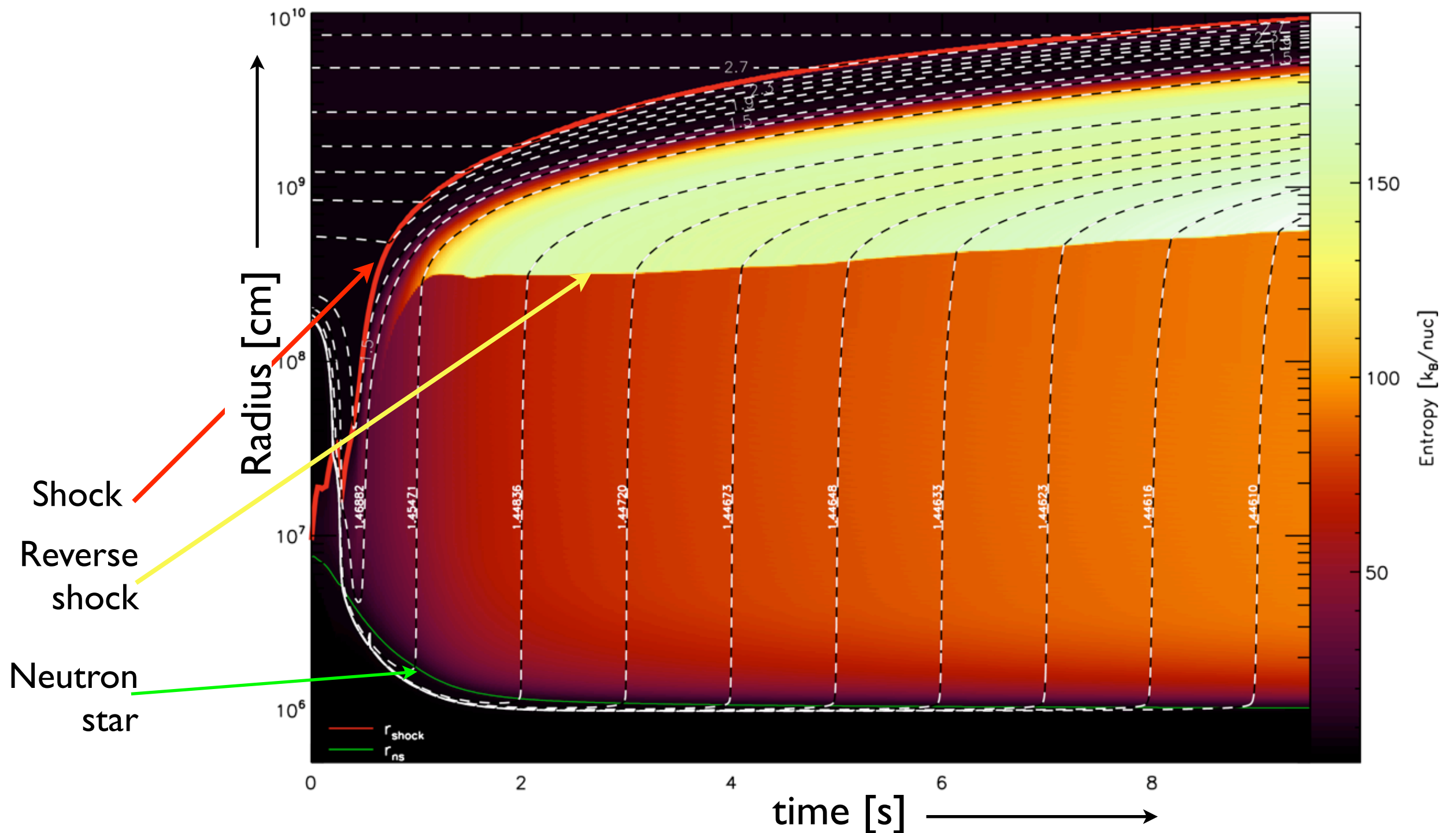
## Nucleosynthesis network including over 5000 nuclei from stability to drip lines

- Network input: trajectories ( $\rho, T$ ) from hydrodynamical simulations + initial  $Y_e$ .
- Starting composition at 10GK is given by nuclear statistical equilibrium (NSE).
- Extended nuclear reaction network including neutral and charged particle reactions from REACLIB (Fröhlich et al. 2006), and weak-reaction rates (Fuller et al. 1999, Langanke & Martinez-Pinedo 2000).



# Neutrino-driven wind results

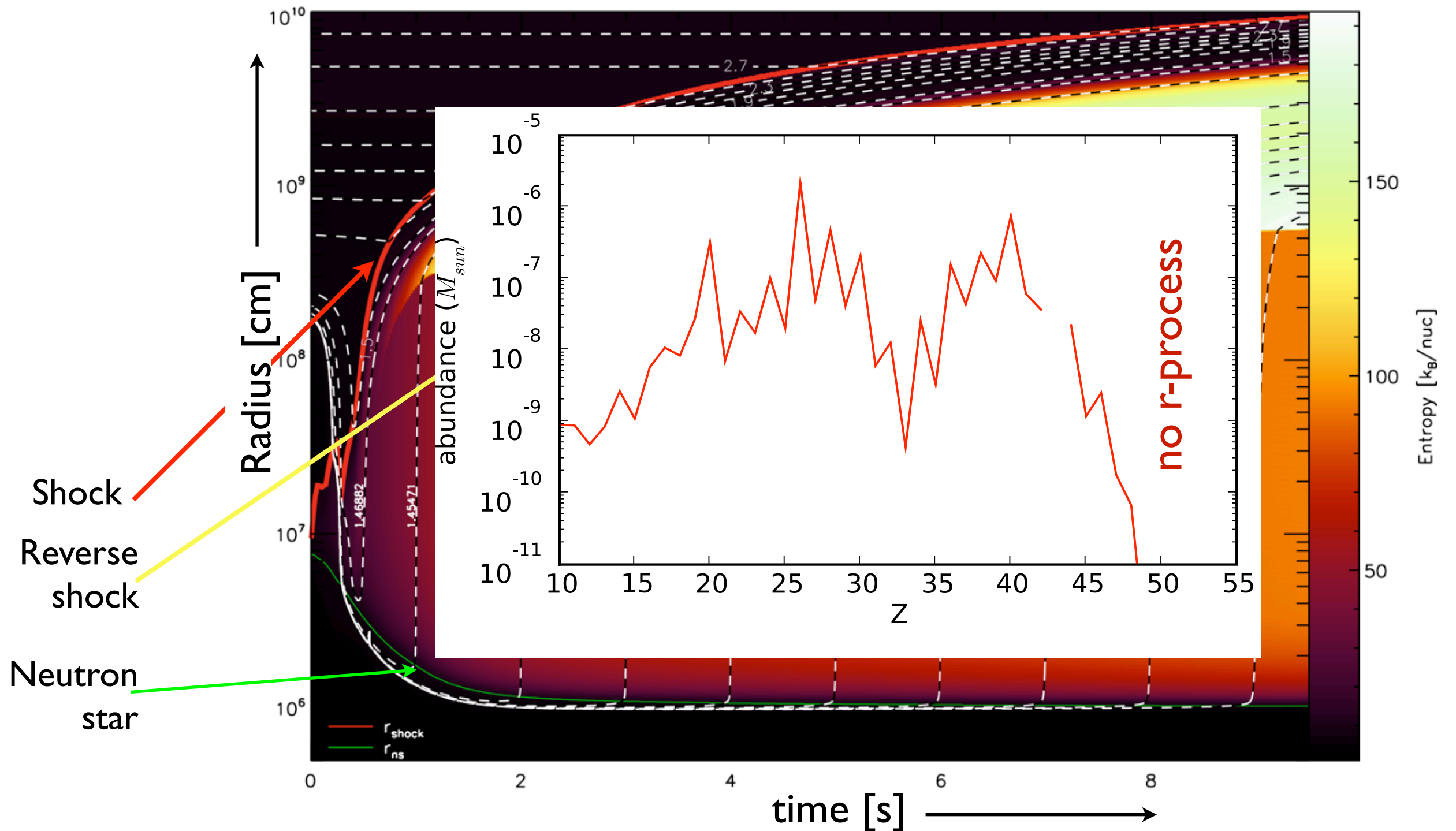
Arcones et al 2007





# Neutrino-driven wind results

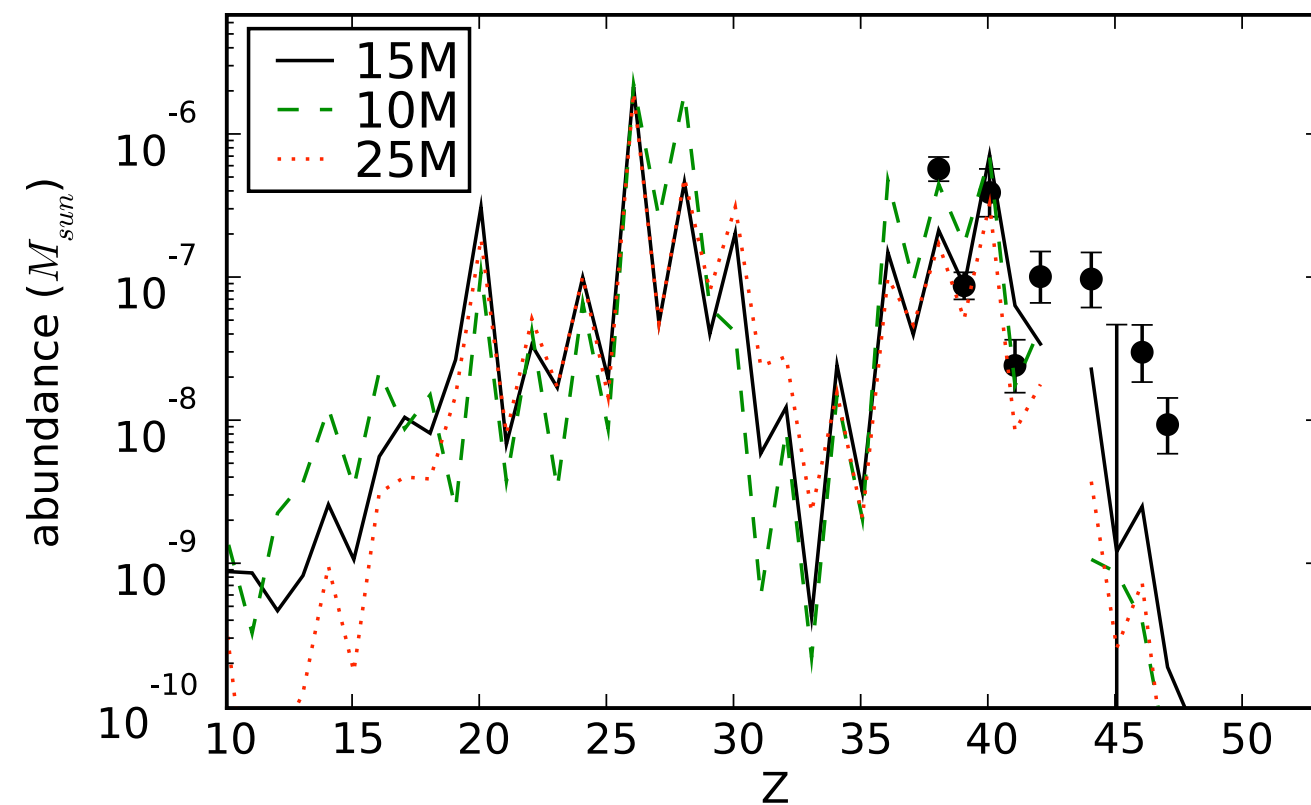
Arcones et al 2007



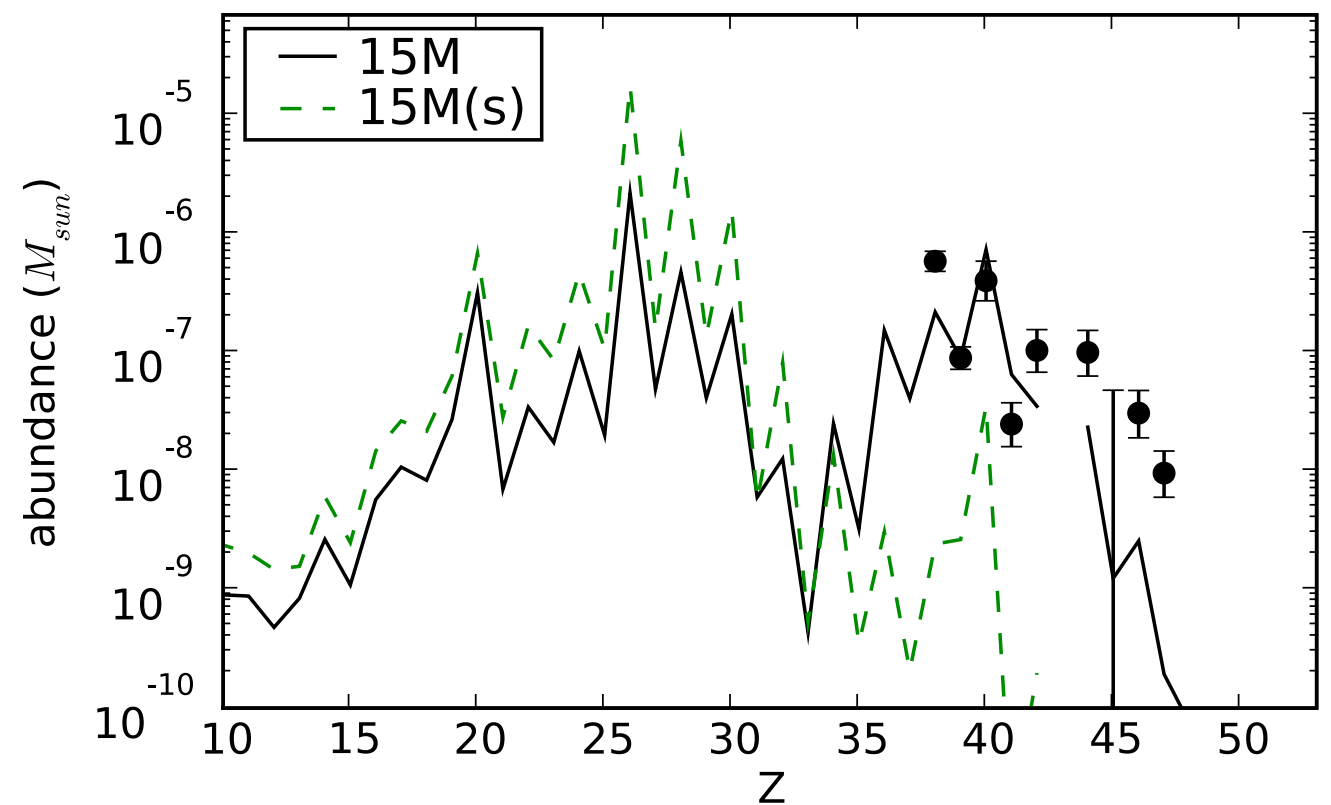
# Nucleosynthesis results

Integrated abundances based on the neutrino-driven wind trajectories compared to LEPP pattern (Montes et al. 2007)

LEPP elements are produced, but no heavy r-nuclei (Arcones & Montes, submitted).



Progenitor variation with same proto-neutron star evolution (cooling and contraction).



Variation of proto-neutron star evolution:

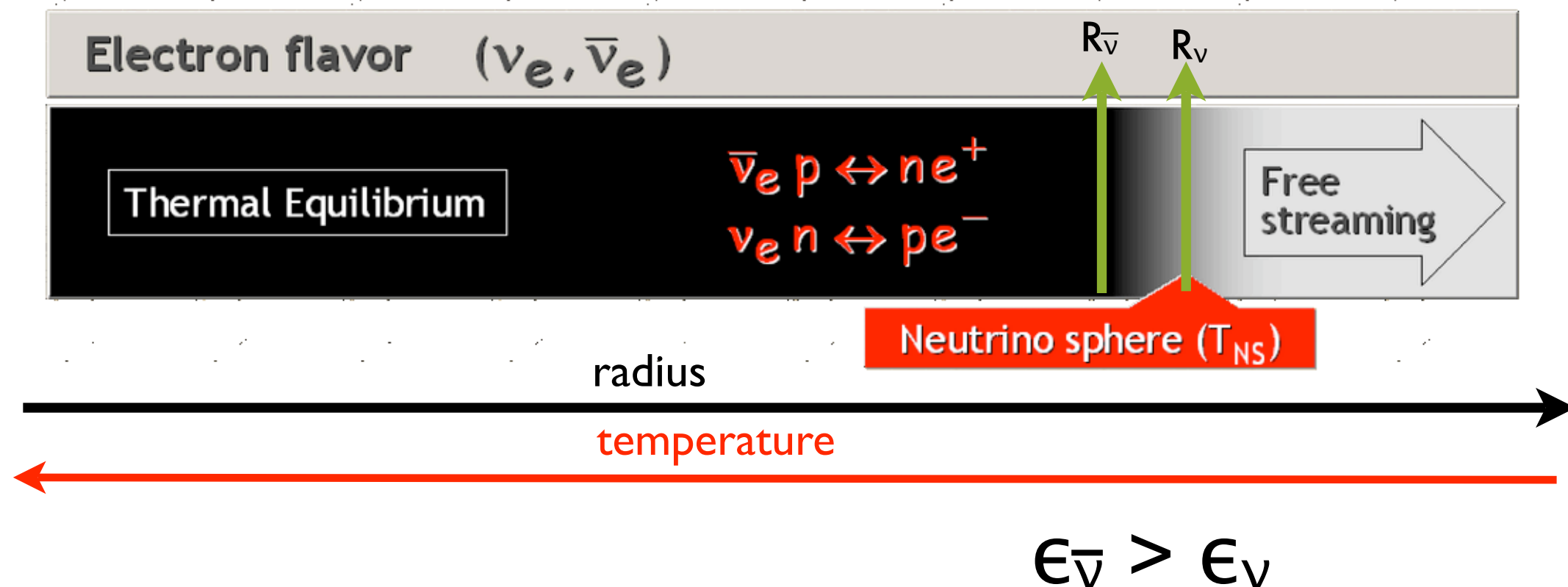
- 15M(s): slow contraction to 15km,  $S \approx 50 k_B/\text{nuc}$
- 15M: fast contraction to 11km,  $S \approx 80 k_B/\text{nuc}$

# Electron fraction and uncertainties

Electron fraction depends on accuracy of the supernova neutrino transport and on details of neutrino interactions in the outer layers of the neutron star.

$$Y_e = \frac{\lambda_{\nu_e, n}}{\lambda_{\nu_e, n} + \lambda_{\bar{\nu}_e, p}} = \left[ 1 + \frac{L_{\bar{\nu}_e} \varepsilon_{\bar{\nu}_e} - 2\Delta + 1.2\Delta^2/\varepsilon_{\bar{\nu}_e}}{L_{\nu_e} \varepsilon_{\nu_e} + 2\Delta + 1.2\Delta^2/\varepsilon_{\nu_e}} \right]^{-1} \quad (\Delta = m_n - m_p)$$

The neutrino energies are determined by the position (temperature) where neutrinos decouple from matter: neutrinosphere



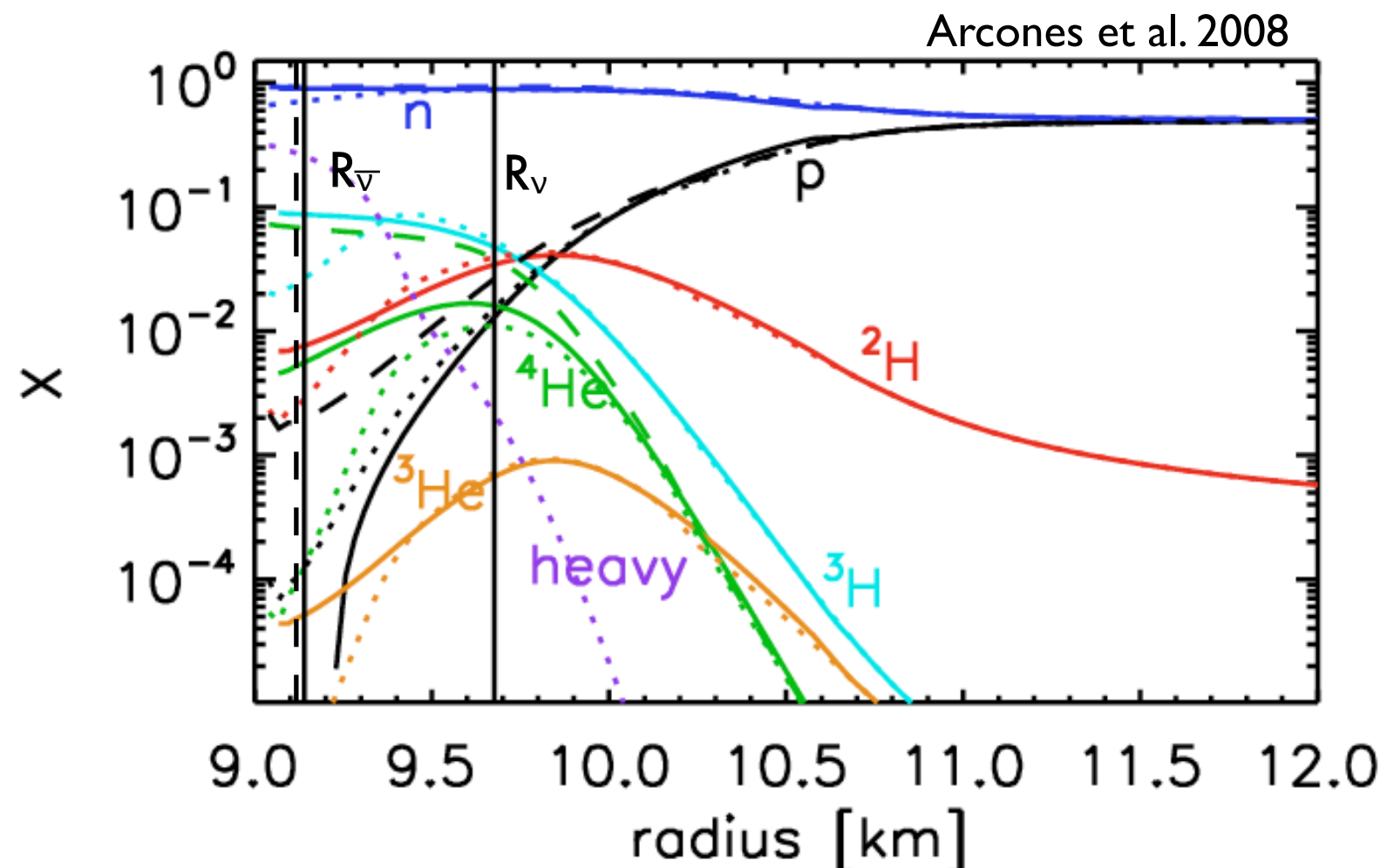
Raffelt 2001, ApJ, 561, 890

# Light nuclei

- The EOS used in supernova simulations include neutron, protons, alpha-particles, and a representative heavy nucleus (Lattimer & Swesty 1991, Shen et al. 1998, or EOS based on NSE). This is o.k. for low densities and high temperatures.
- However, they don't allow for the presence of light nuclei ( $A \leq 4$ ) which are present for densities  $\rho \approx 10^{12} \text{g/cm}^3$  (O'Connor et al. 2007, Sumiyoshi & Ropke 2008, Typel et al. 2009).

We have compared 3 EOS:

- non-interaction ideal gas of n,p, $\alpha$  (dashed lines)
- Virial (Horowitz & Schwenk 2006): nuclei with  $A \leq 4$  and interactions (solid lines)
- NSE: several thousands nuclei (dotted lines)



# Light nuclei

• The EOS use representative NSE). This is o.

• However, the densities  $\rho \approx 10$

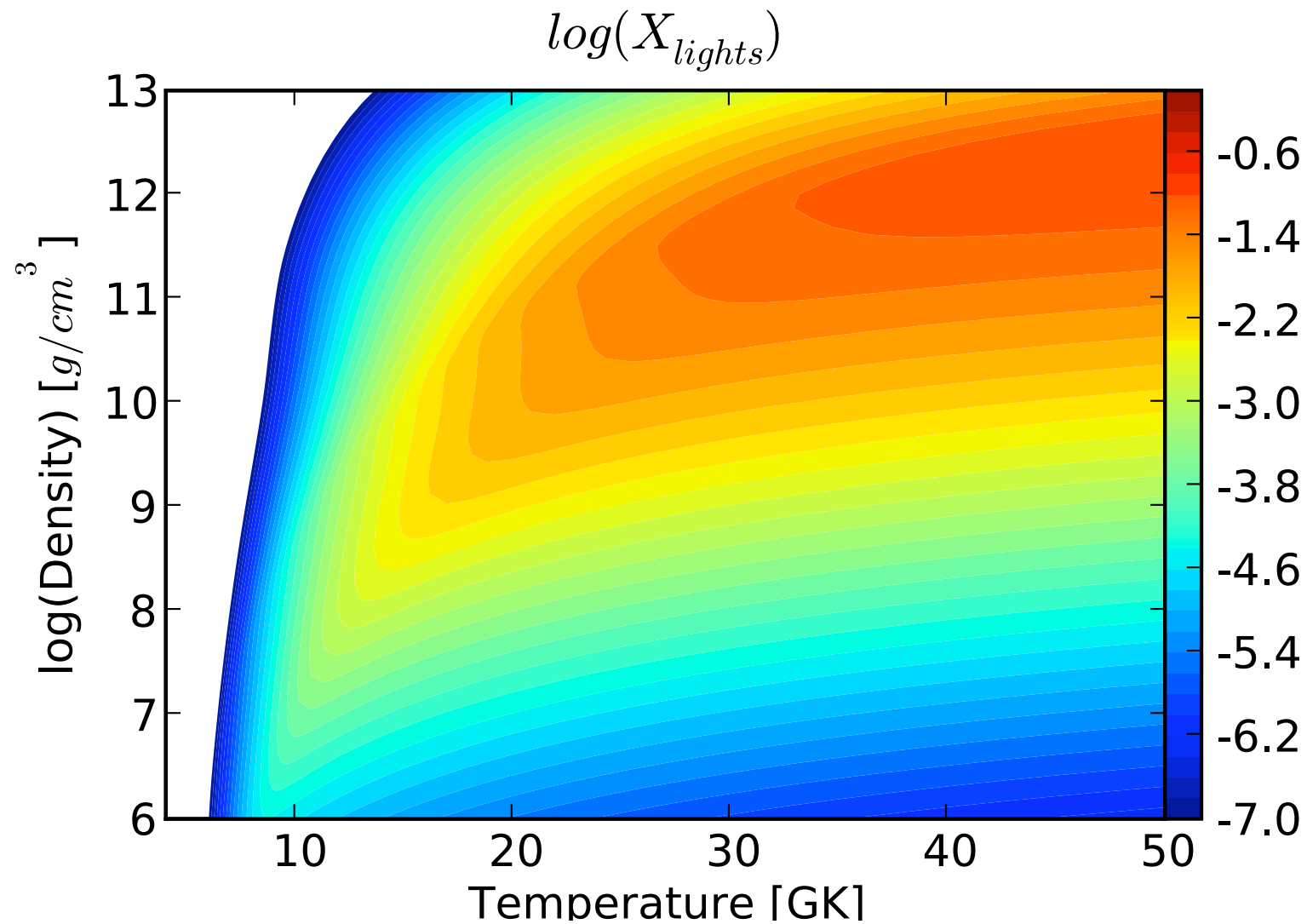
We have comp

• non-interacti  
(dashed lines)

• Virial (Horov  
nuclei with  $A \leq 4$   
(solid lines)

• NSE: several  
(dotted lines)

Under general assumptions (NSE+beta equilibrium)  $1 < A < 4$

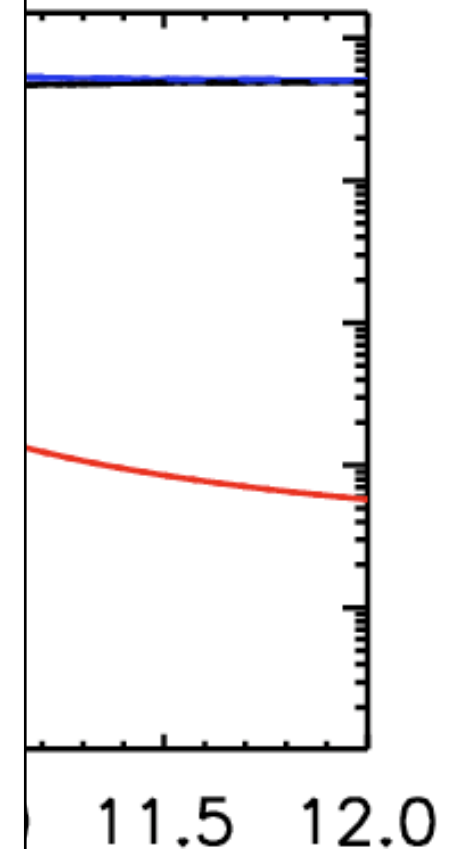


Arcones, Martinez-Pinedo, Roberts, Woosley 2010

and a  
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t for  
(2009).

Arcones et al. 2008



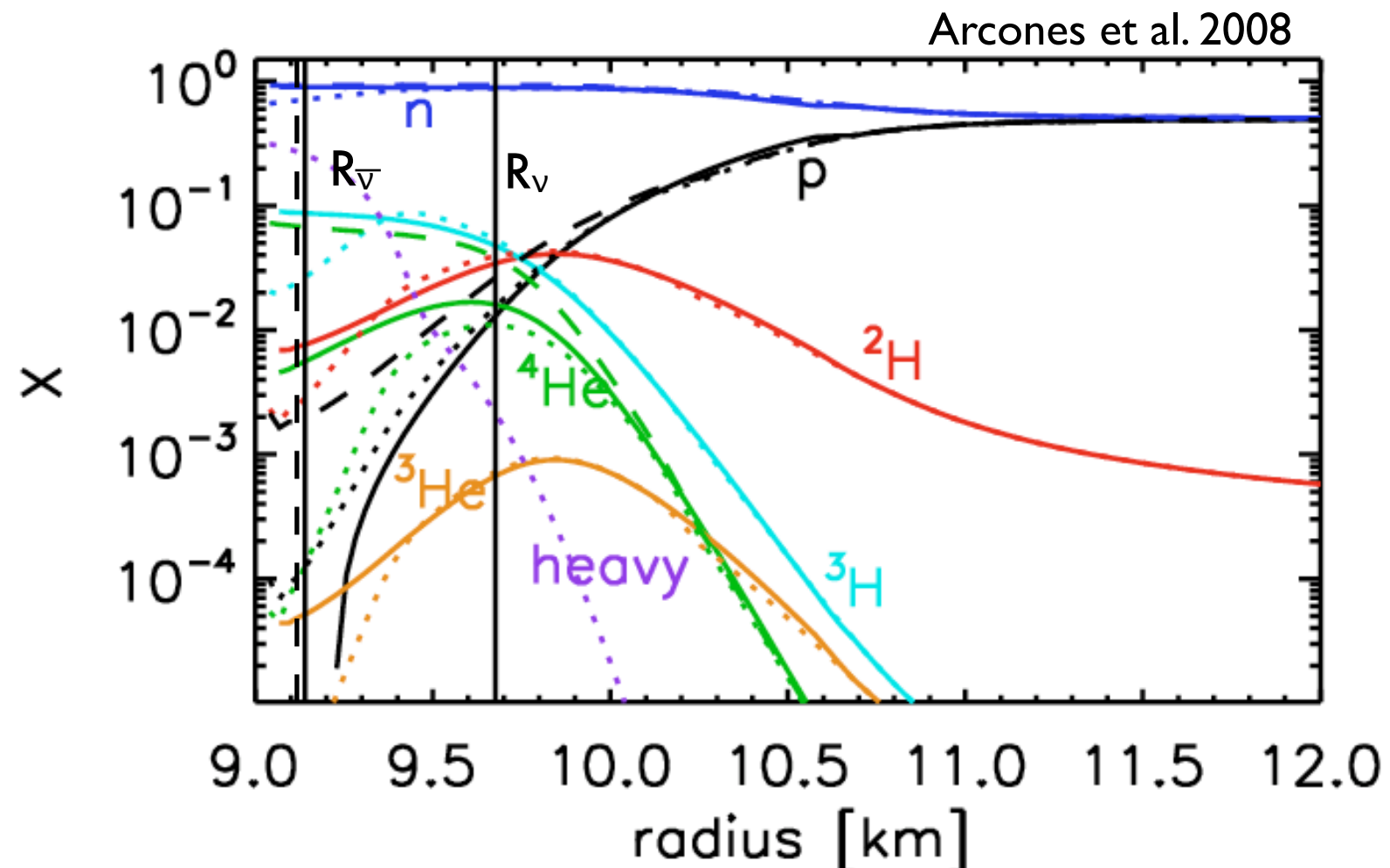


# Light nuclei

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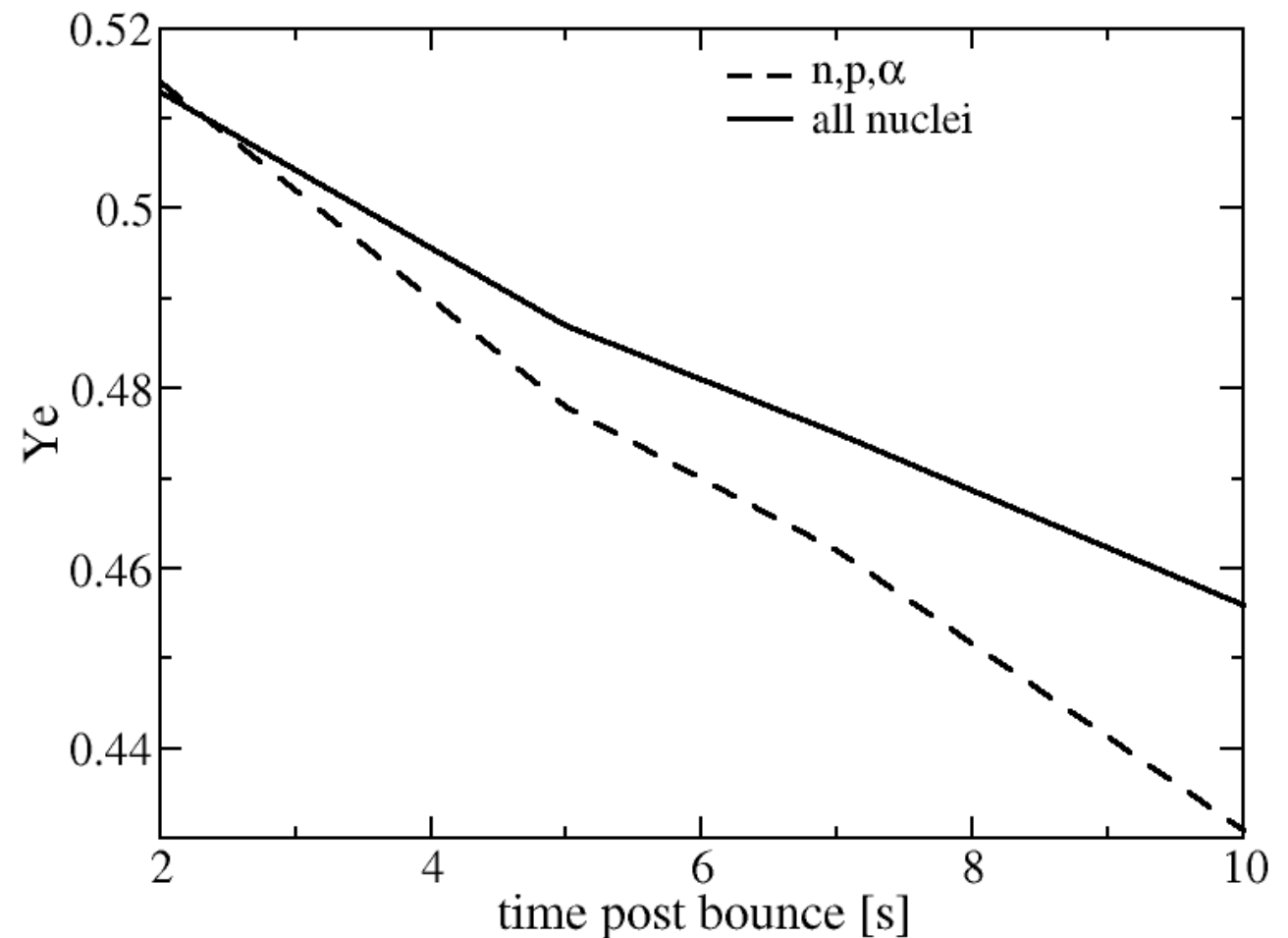
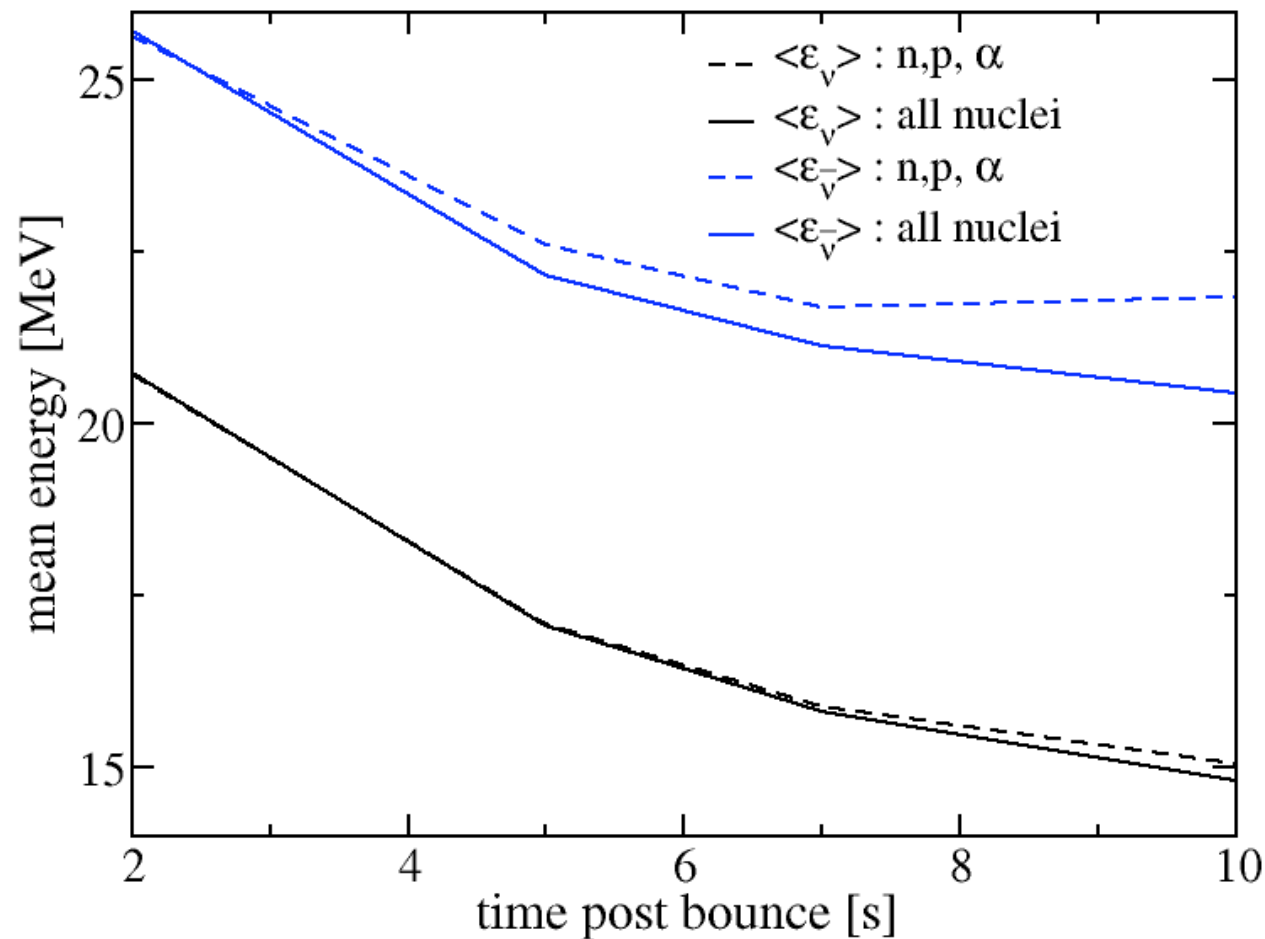
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# Light nuclei: antineutrino energy

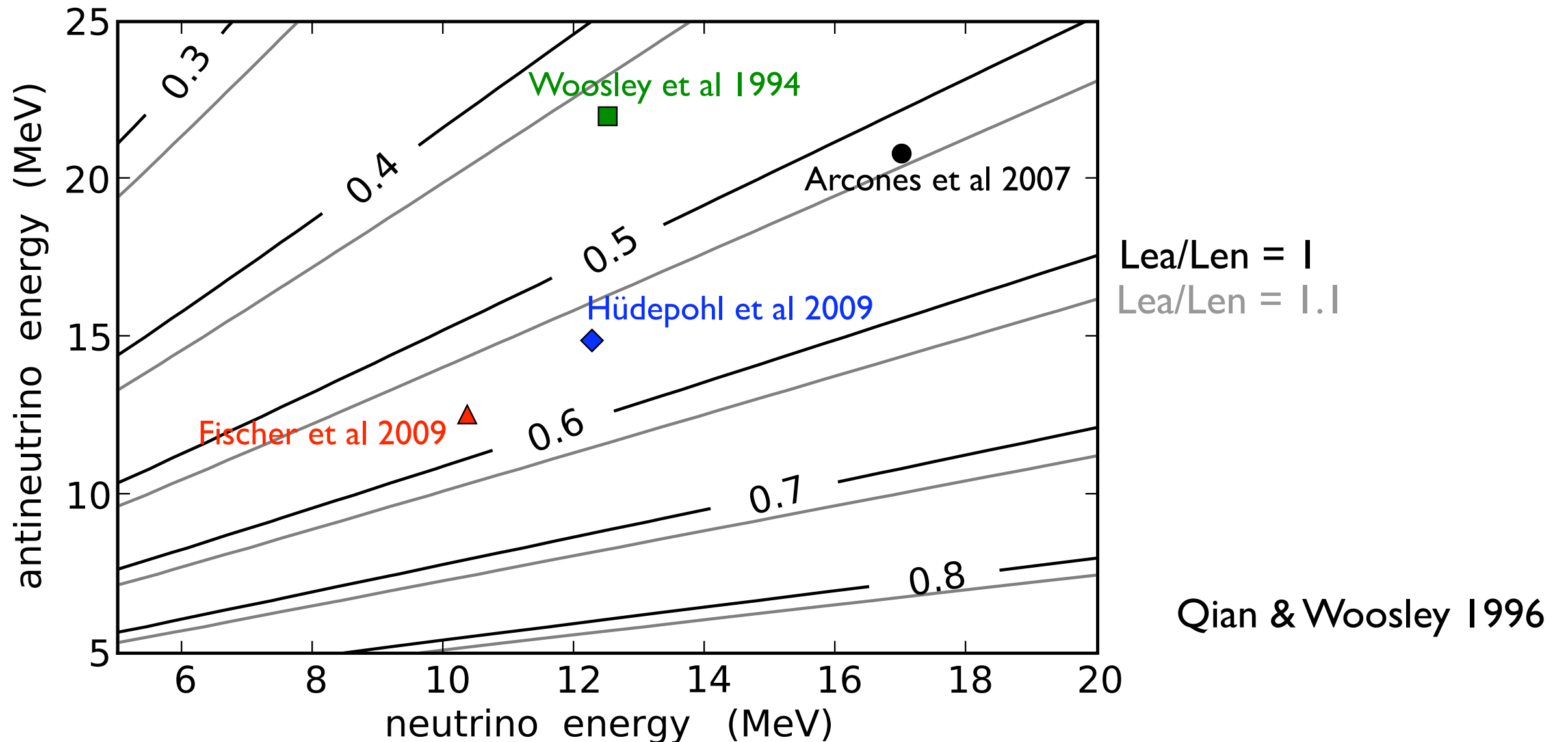
We have followed the evolution of the neutrinosphere taking into account the presence of light clusters in the outer layers of the neutron star (Arcones et al. 2008) and the neutrino interaction with them



# Wind models and electron fraction

Neutrino energies change with more realistic neutrino physics input

More recent simulations obtain lower antineutrino energies and therefore proton-rich conditions

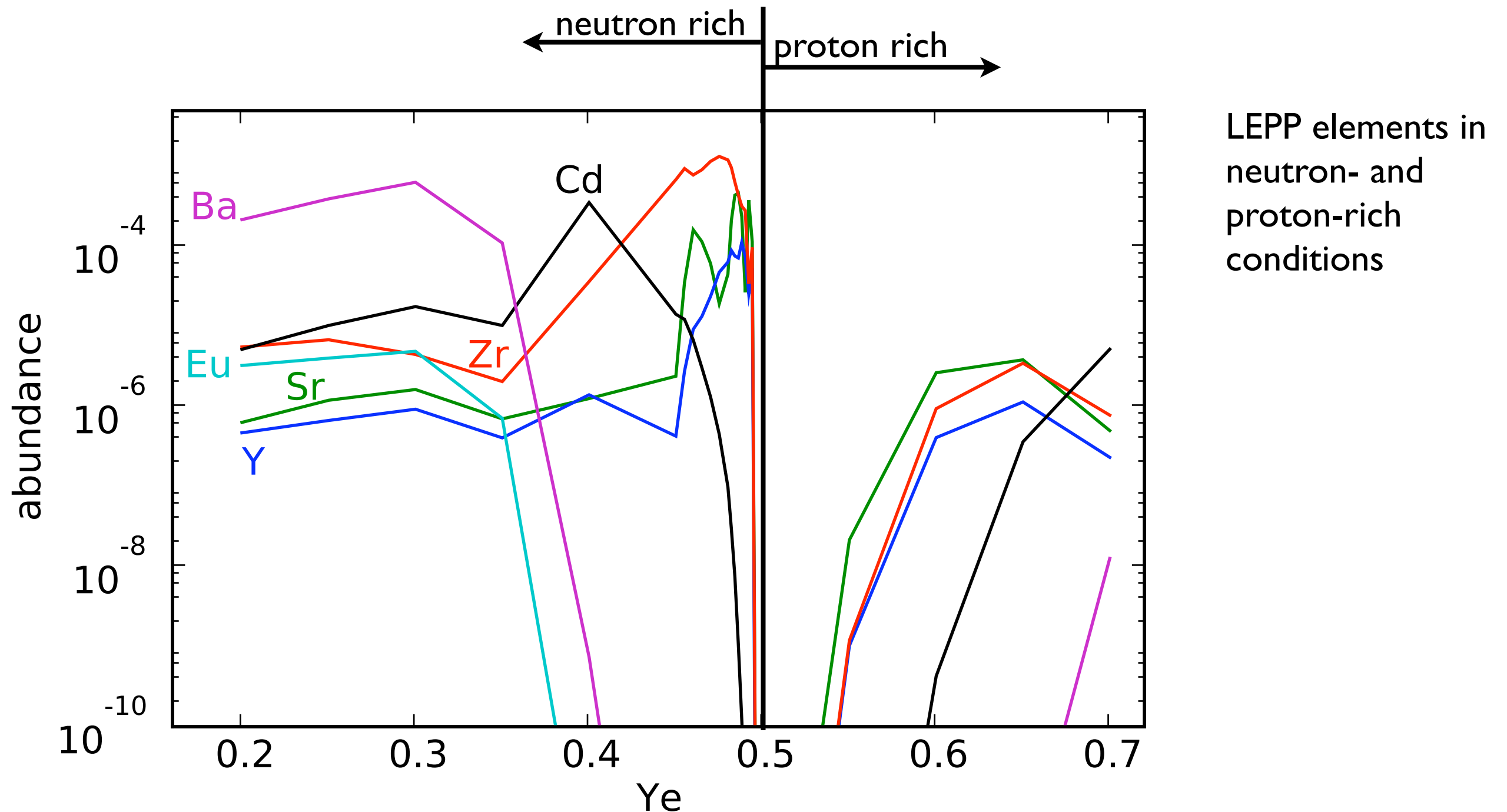


$$Y_e = \frac{\lambda_{\nu_e, n}}{\lambda_{\nu_e, n} + \lambda_{\bar{\nu}_e, p}} = \left[ 1 + \frac{L_{\bar{\nu}_e} \epsilon_{\bar{\nu}_e} - 2\Delta + 1.2\Delta^2/\epsilon_{\bar{\nu}_e}}{L_{\nu_e} \epsilon_{\nu_e} + 2\Delta + 1.2\Delta^2/\epsilon_{\nu_e}} \right]^{-1}$$

$$Y_e > 0.5: \epsilon_{\bar{\nu}} - \epsilon_{\nu} < 4\Delta$$

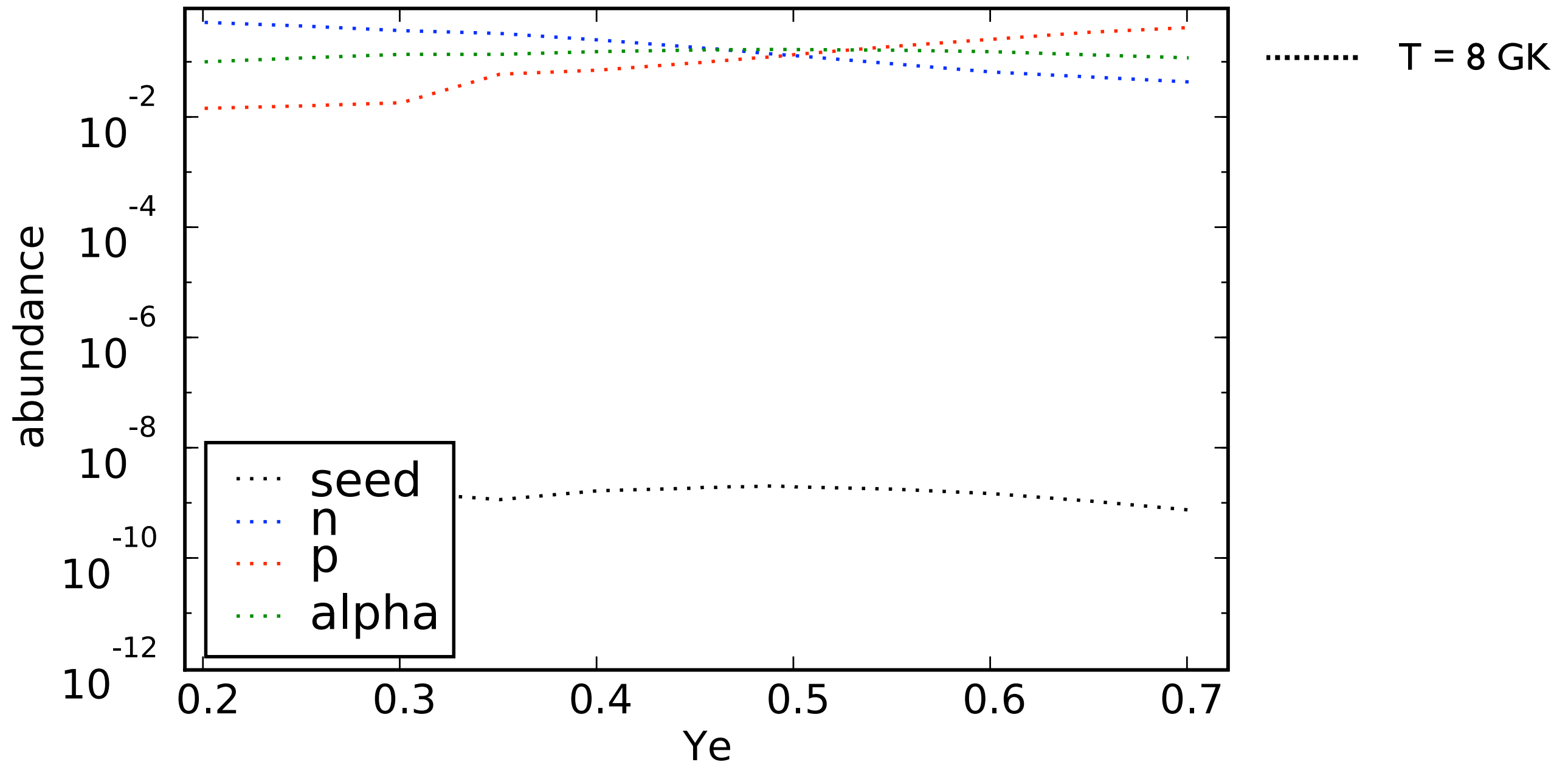
# Nucleosynthesis and electron fraction

Study the impact of the electron fraction on the production of LEPP elements (Sr, Y, Zr)  
(Arcones & Montes)



# Nucleosynthesis and electron fraction

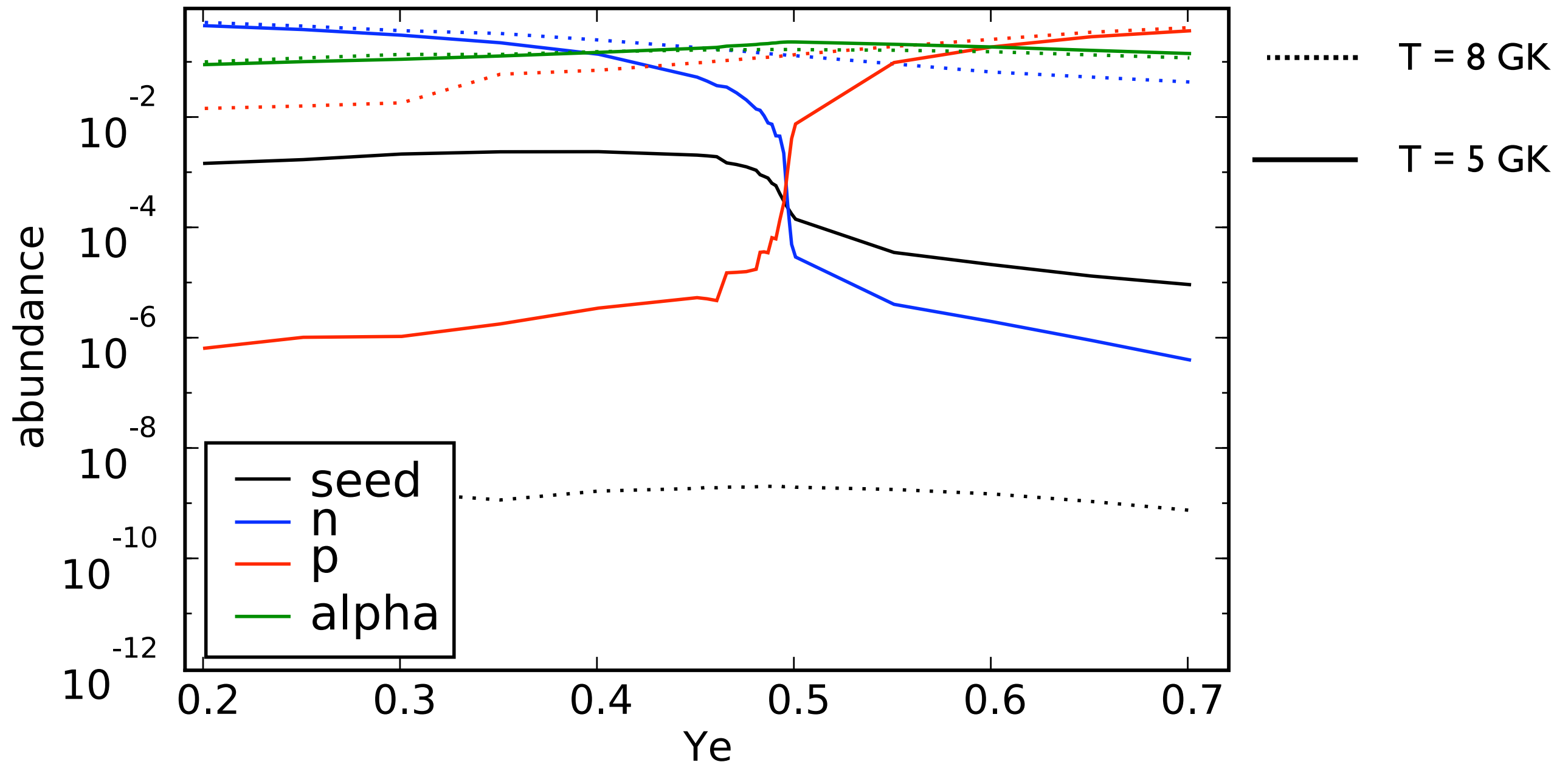
Initial composition is given by NSE, at high temperatures only n, p and alphas.





# Nucleosynthesis and electron fraction

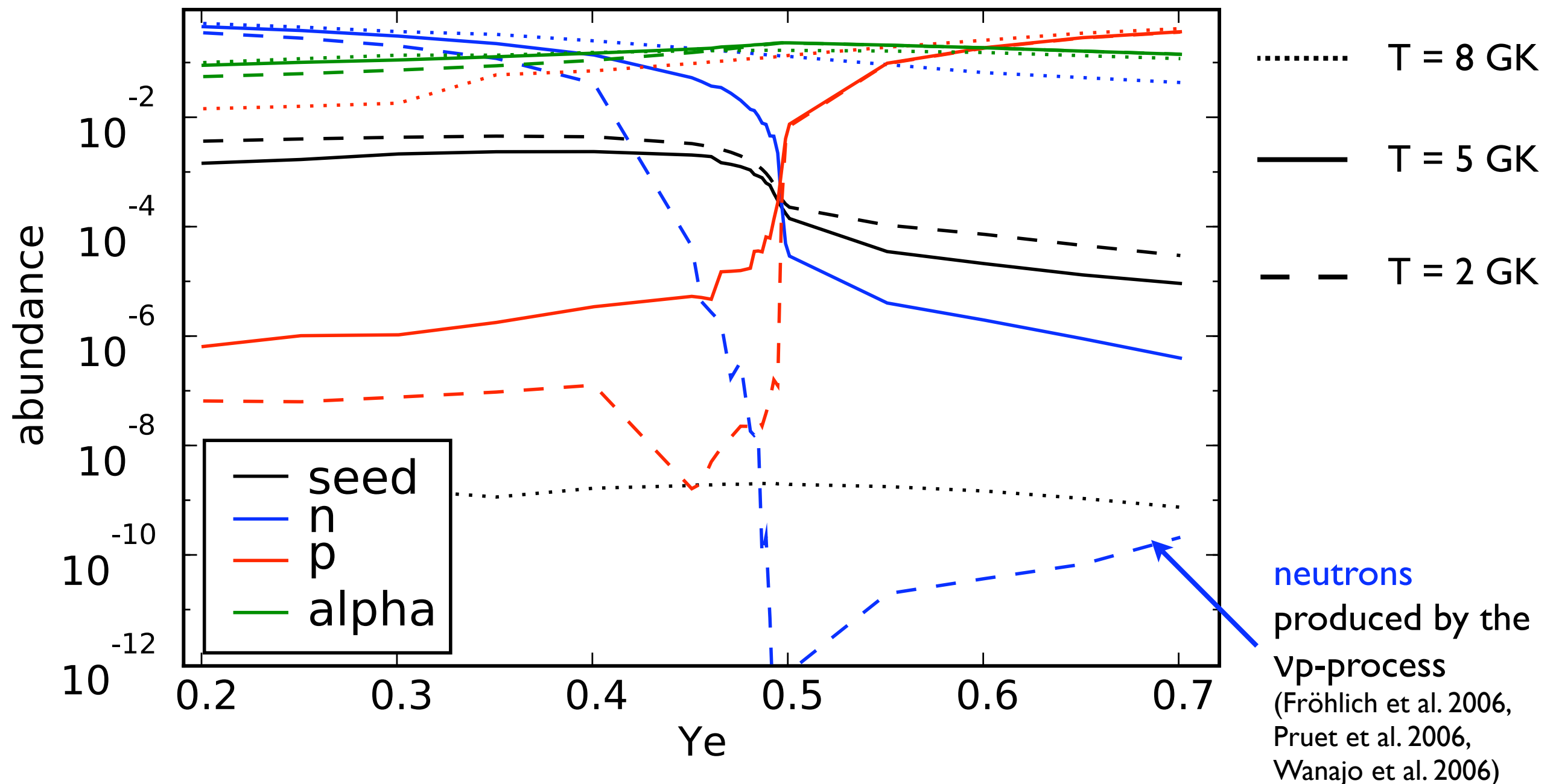
Initial composition is given by NSE, at high temperatures only n, p and alphas.  
Alpha particles recombine forming seed nuclei.

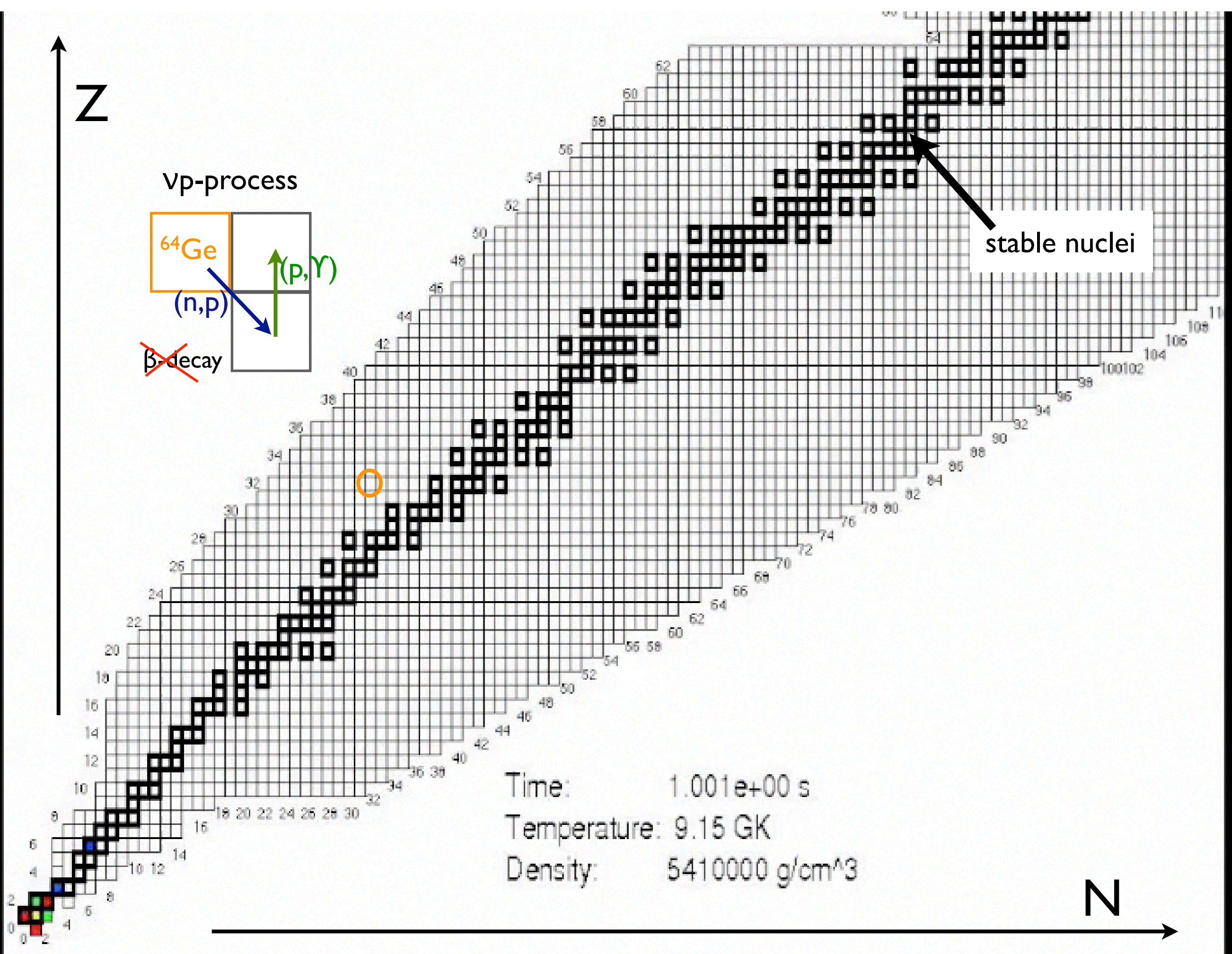


# Nucleosynthesis and electron fraction

Initial composition is given by NSE, at high temperatures only n, p and alphas.  
Alpha particles recombine forming seed nuclei.

At freeze-out neutron- and proton-to-seed ratio determine production of heavy elements.



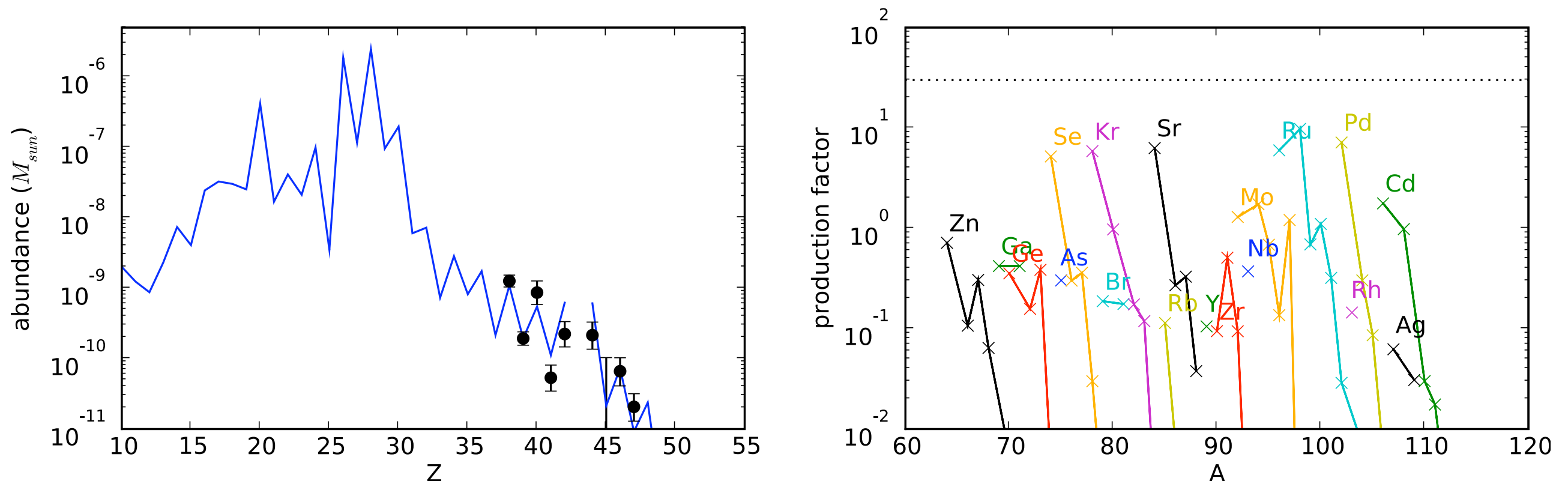


# LEPP in proton-rich ejecta

Exploration of the time dependence of the electron fraction.

Superposition of trajectories with  $Y_e > 0.5$  following most recent simulations (Basel and Garching 2009). Compare to LEPP pattern (rescaled to  $Z=39$ ).

Our results can explain the LEPP abundances in old halo stars and the origin of p-nuclei.



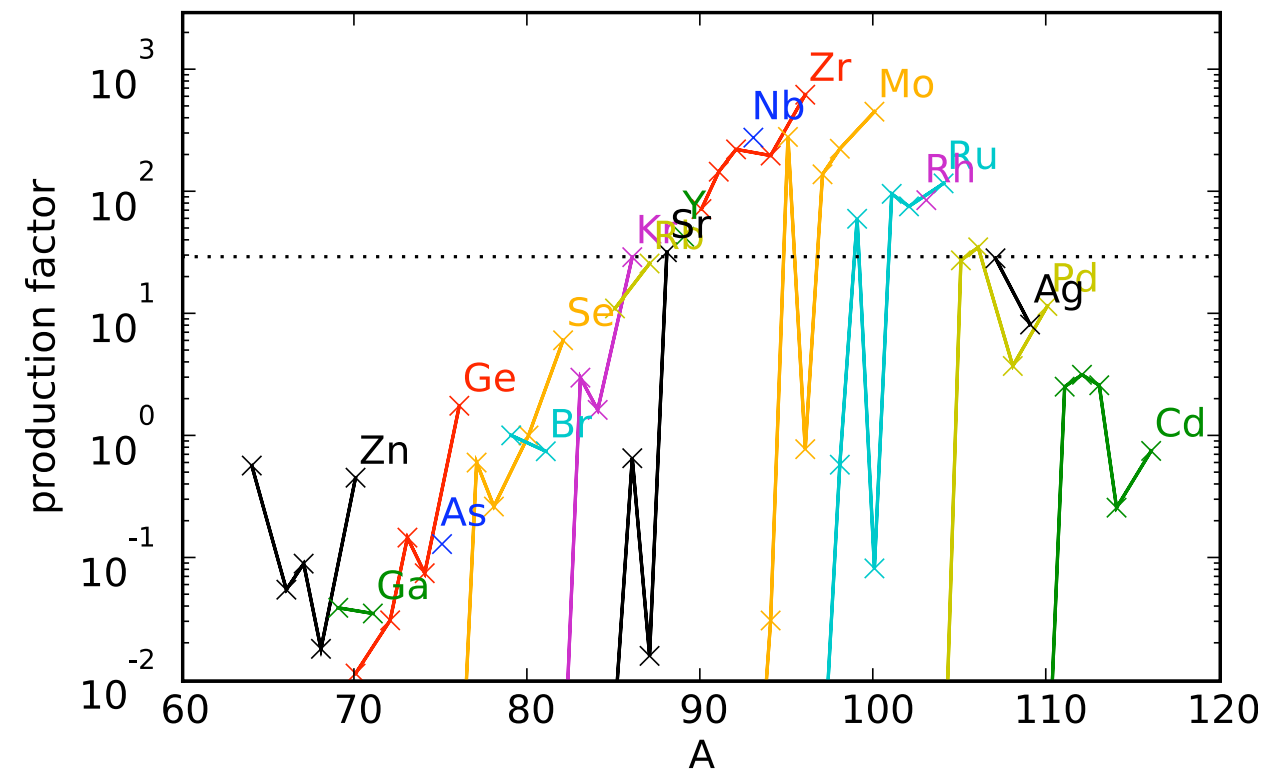
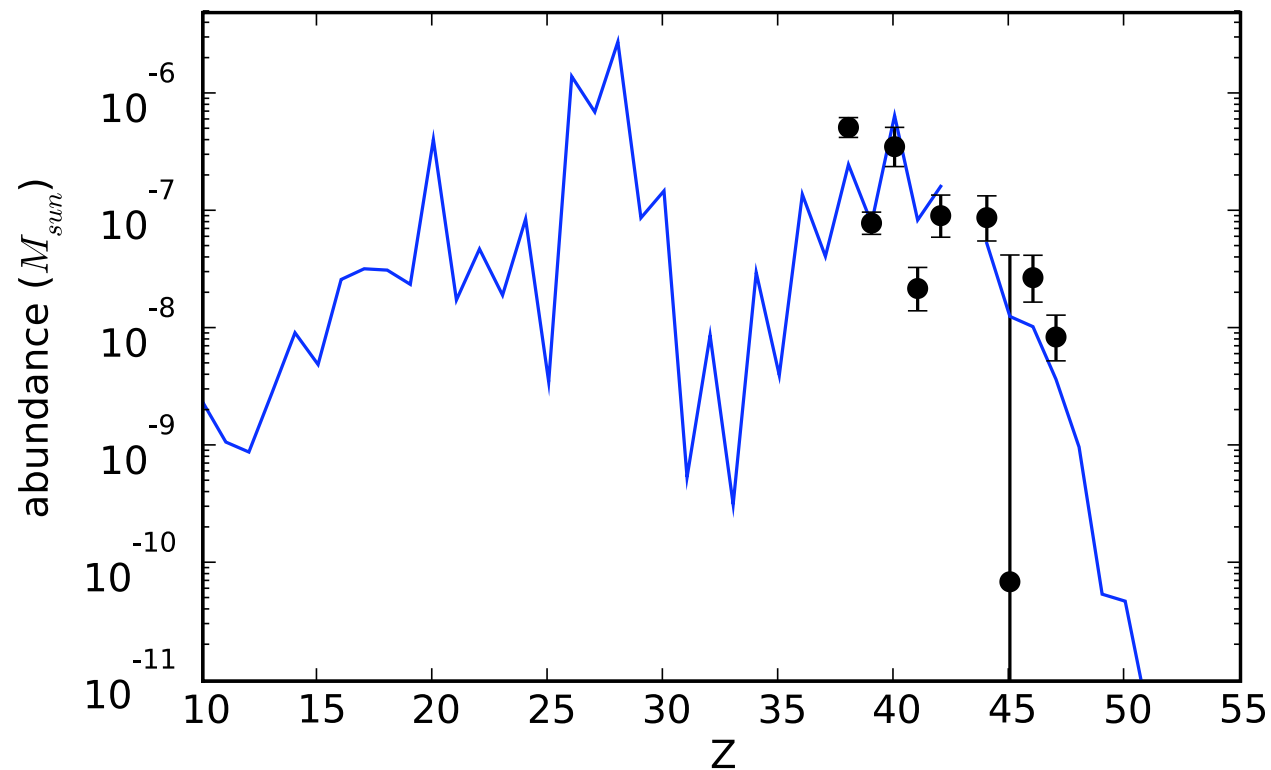
Problem: in the LEPP component of the solar system there are also neutron-rich nuclei.

Isotopic abundances from UMP stars will give rise new insights.

# LEPP in neutron-rich ejecta

Superposition of trajectories with neutron-rich conditions:  $0.5 > Y_e > 0.45$ .

LEPP elements are produced and also neutron-rich isotopes.



Problem: overproduction at  $A=90$  for magic neutron number  $N=50$  (Hoffman et al. 1996).

Suggest that if only a fraction of the supernovae eject neutron-rich material → can explain LEPP in the solar system.



# LEPP: neutron vs. proton-rich winds

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## Proton-rich wind:

- reproduces LEPP pattern found in UMP stars
- robust under small variations of wind parameters
- no overproduction but mainly p-nuclei produced, no neutron-rich isotopes as in solar system
- results depend strongly on neutrino properties:  $\nu p$ -process

## Neutron-rich wind:

- produces LEPP elements with neutron-rich isotopes, it could explain LEPP component in solar system
- no robustness under small variations of wind parameters
- overproduction at  $A \sim 90$ , only subset of supernovae or small amount of their ejecta can be neutron rich

# Conclusions

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- First comparison of the LEPP pattern and nucleosynthesis calculations based on hydrodynamical wind simulations.
- Electron fraction is key for abundances and depends on details of the composition and neutrino interactions in the outer layers of the proto-neutron star.
- Light elements are present in the outer layer of the proto-neutron star.
- LEPP pattern is reproduced in neutron- and proton-rich ejecta.

## and outlook

- Multi-dimensional simulations with detailed neutrino transport.
- Include light elements and their corresponding neutrino reactions.
- Observations of isotopic abundances in old stars can discriminate.
- Use abundances to constrain  $Y_e$  evolution.
- Galactic chemical evolution models.

# Thank you!

H.Th. Janka, K. Langanke, G. Martinez-Pinedo,  
H. Schatz, F. K. Thielemann

