

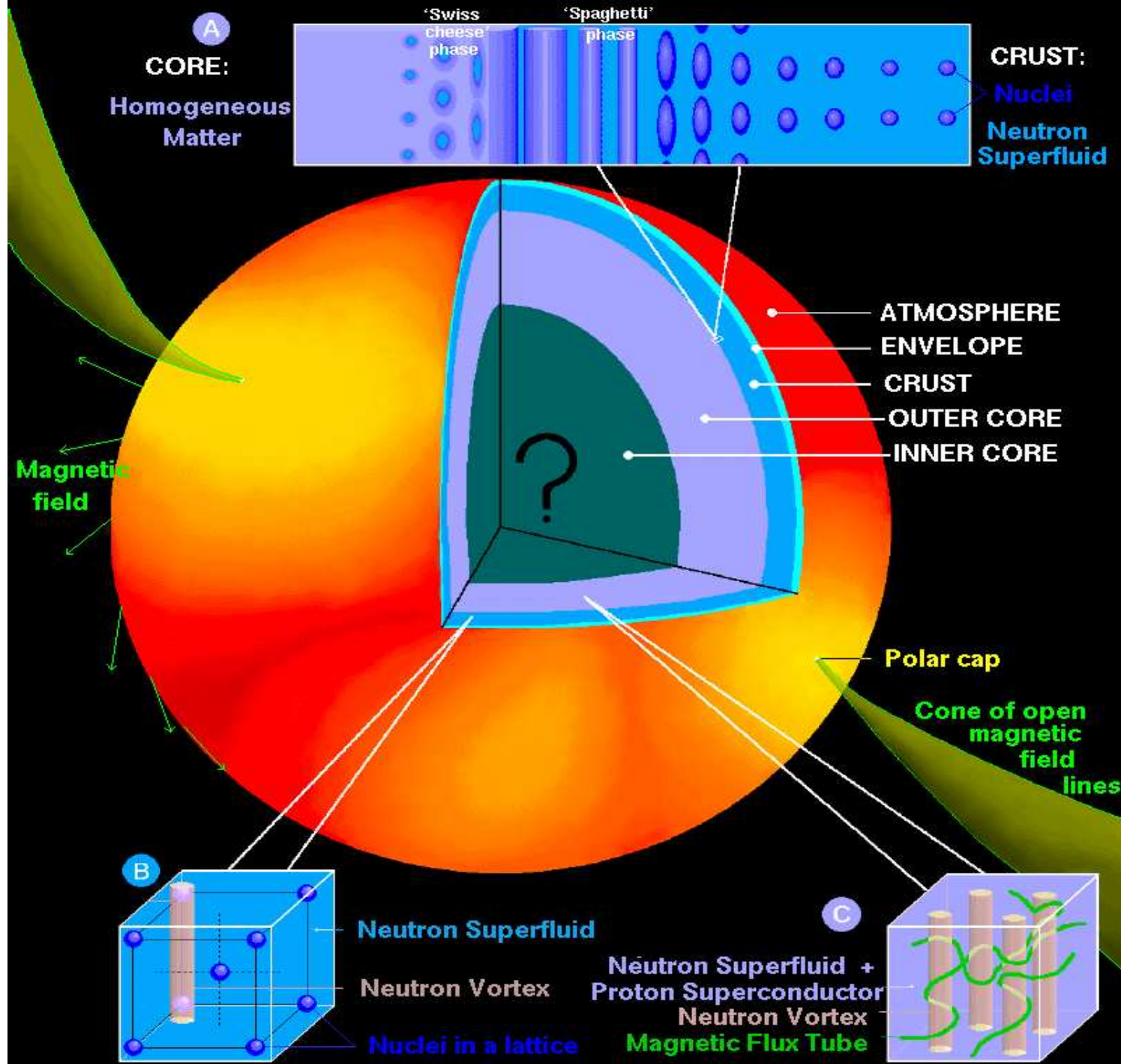
Neutron Star Masses and Radii

James M. Lattimer

`lattimer@astro.sunysb.edu`

Department of Physics & Astronomy
Stony Brook University

A NEUTRON STAR: SURFACE and INTERIOR

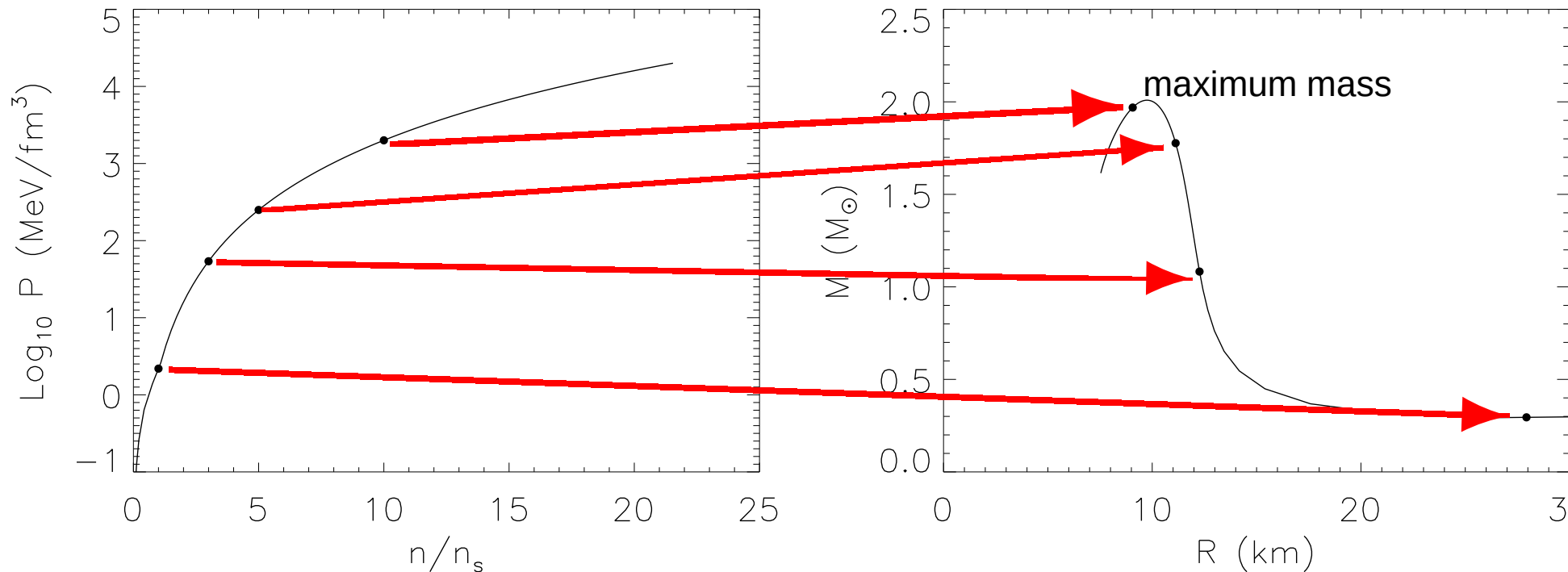


Credit: Dany Page, UNAM

Neutron Star Structure

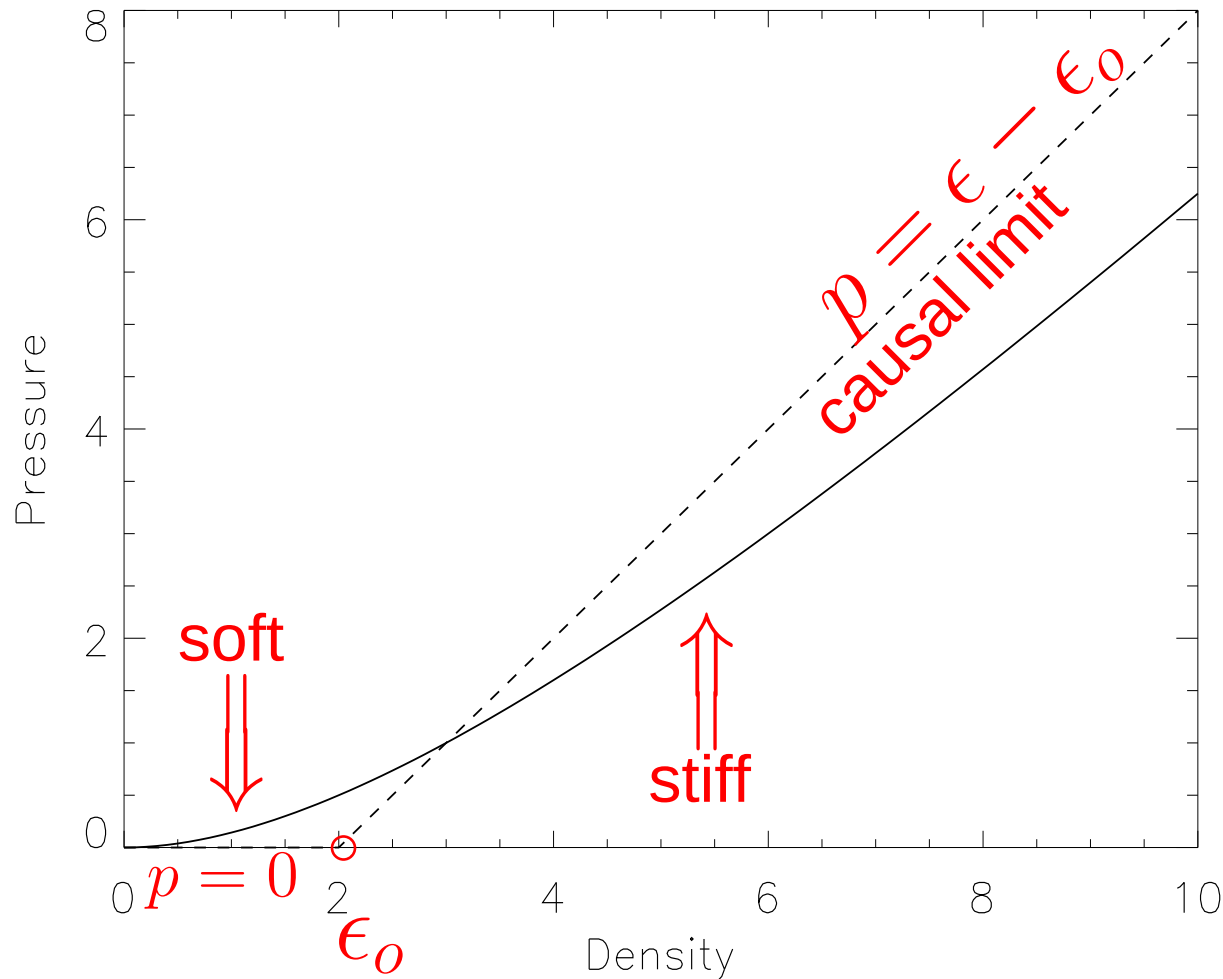
Tolman-Oppenheimer-Volkov equations

$$\frac{dp}{dr} = -\frac{G}{c^2} \frac{(m + 4\pi pr^3)(\epsilon + p)}{r(r - 2Gm/c^2)}$$
$$\frac{dm}{dr} = 4\pi \frac{\epsilon}{c^2} r^2$$



Extreme Properties of Neutron Stars

- The most compact configurations occur when the low-density equation of state is "soft" and the high-density equation of state is "stiff".



ϵ_0 is the only EOS parameter

The TOV solutions scale with ϵ_0

Maximum Mass, Minimum Period

Theoretical limits from GR and causality

- $M_{max} = 4.2(\epsilon_s/\epsilon_f)^{1/2} M_\odot$ Rhoades & Ruffini (1974), Hartle (1978)

- $R_{min} = 2.9GM/c^2 = 4.3(M/M_\odot) \text{ km}$

Lindblom (1984), Glendenning (1992), Koranda, Stergioulas & Friedman (1997)

- $\epsilon_c < 4.5 \times 10^{15} (M_\odot/M_{largest})^2 \text{ g cm}^{-3}$ Lattimer & Prakash (2005)

- $P_{min} \simeq 0.74(M_\odot/M_{sph})^{1/2} (R_{sph}/10 \text{ km})^{3/2} \text{ ms}$

Koranda, Stergioulas & Friedman (1997)

- $P_{min} \simeq 0.96 \pm 0.03 (M_\odot/M_{sph})^{1/2} (R_{sph}/10 \text{ km})^{3/2} \text{ ms}$

(empirical)

Lattimer & Prakash (2004)

- $\epsilon_c > 0.91 \times 10^{15} (1 \text{ ms}/P_{min})^2 \text{ g cm}^{-3}$ (empirical)

- $cJ/GM^2 \lesssim 0.5$ (empirical, neutron star)

Mass-Radius Diagram and Theoretical Constraints

GR:

$$R > 2GM/c^2$$

$P < \infty$:

$$R > (9/4)GM/c^2$$

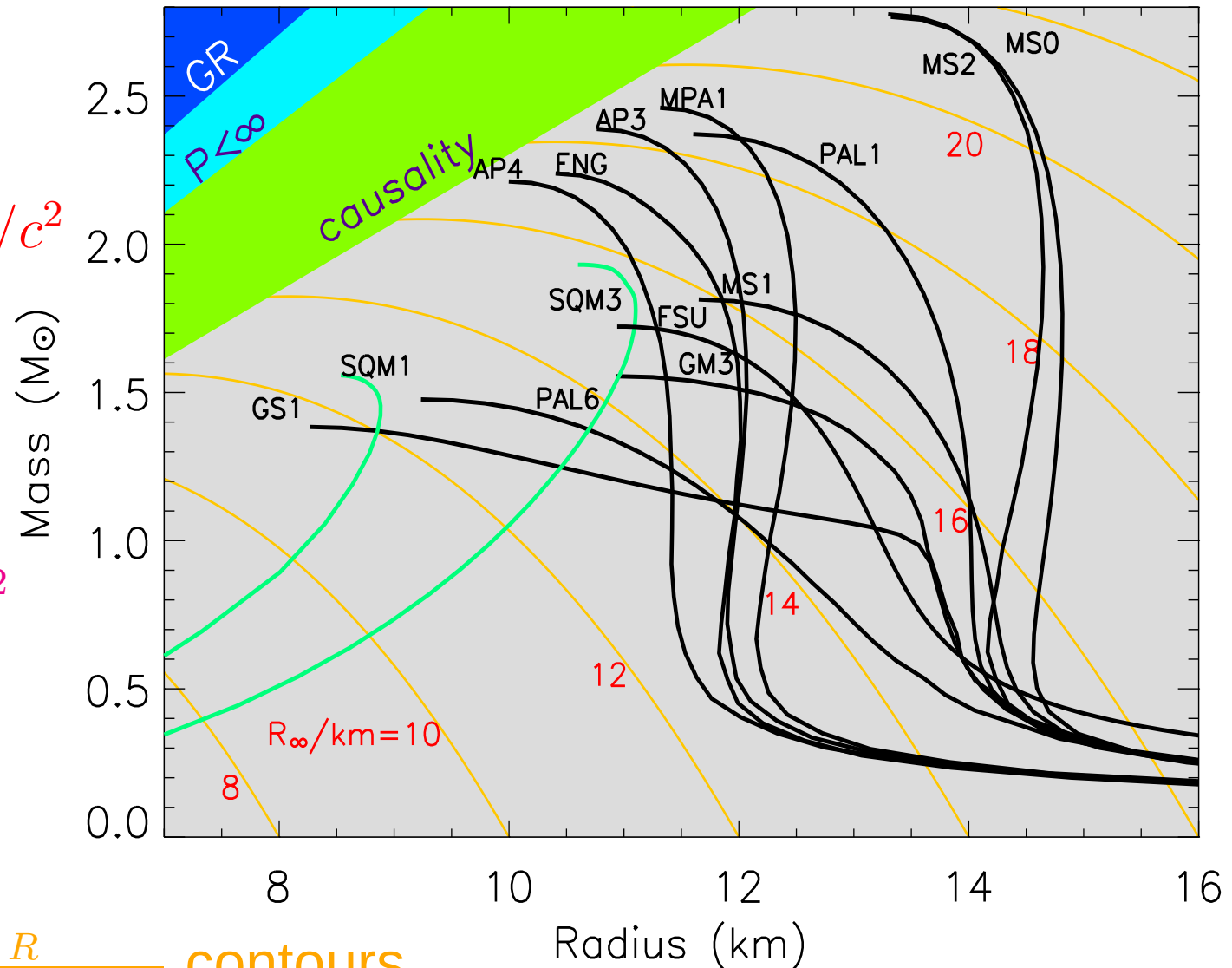
causality:

$$R \gtrsim 2.9GM/c^2$$

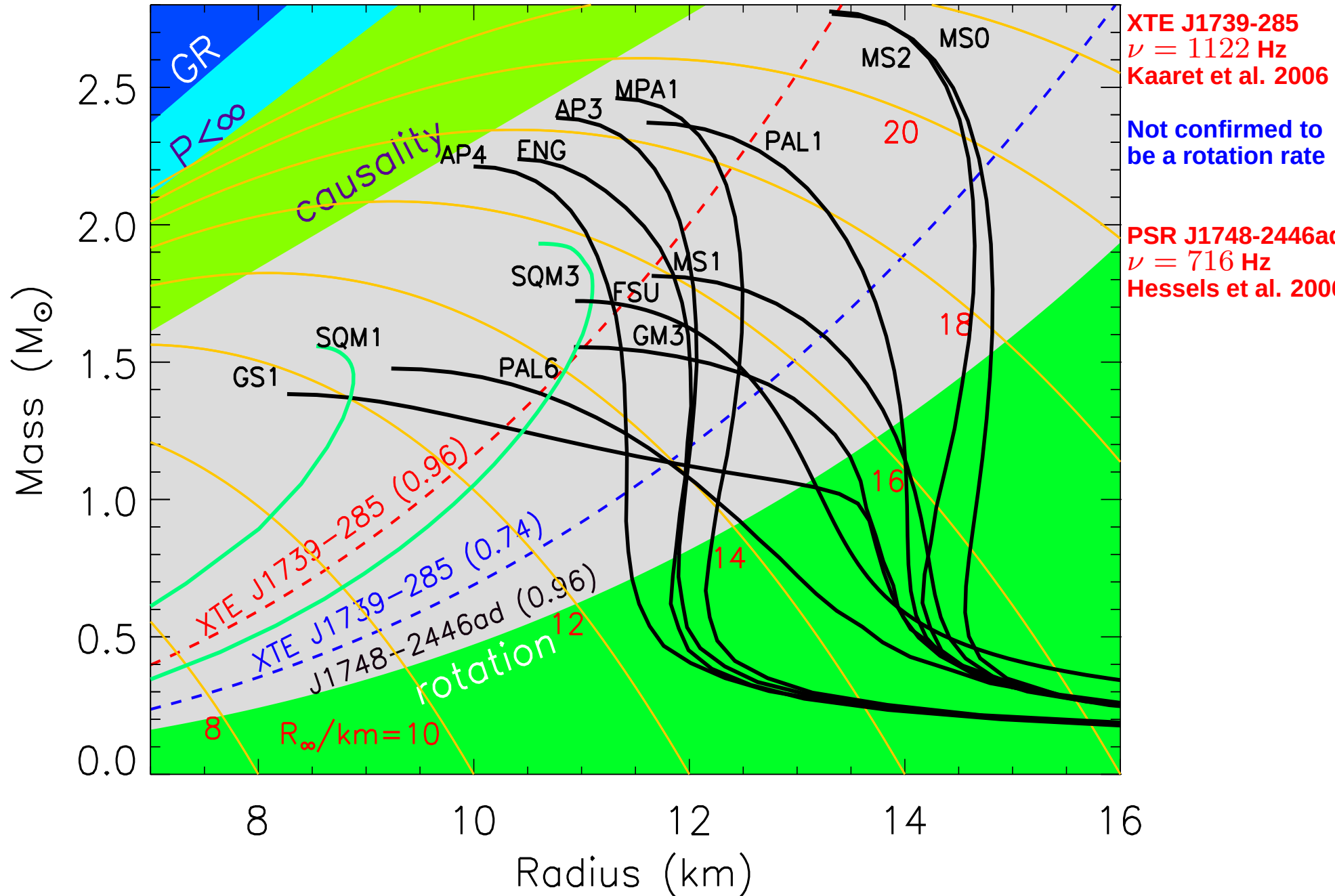
— normal NS

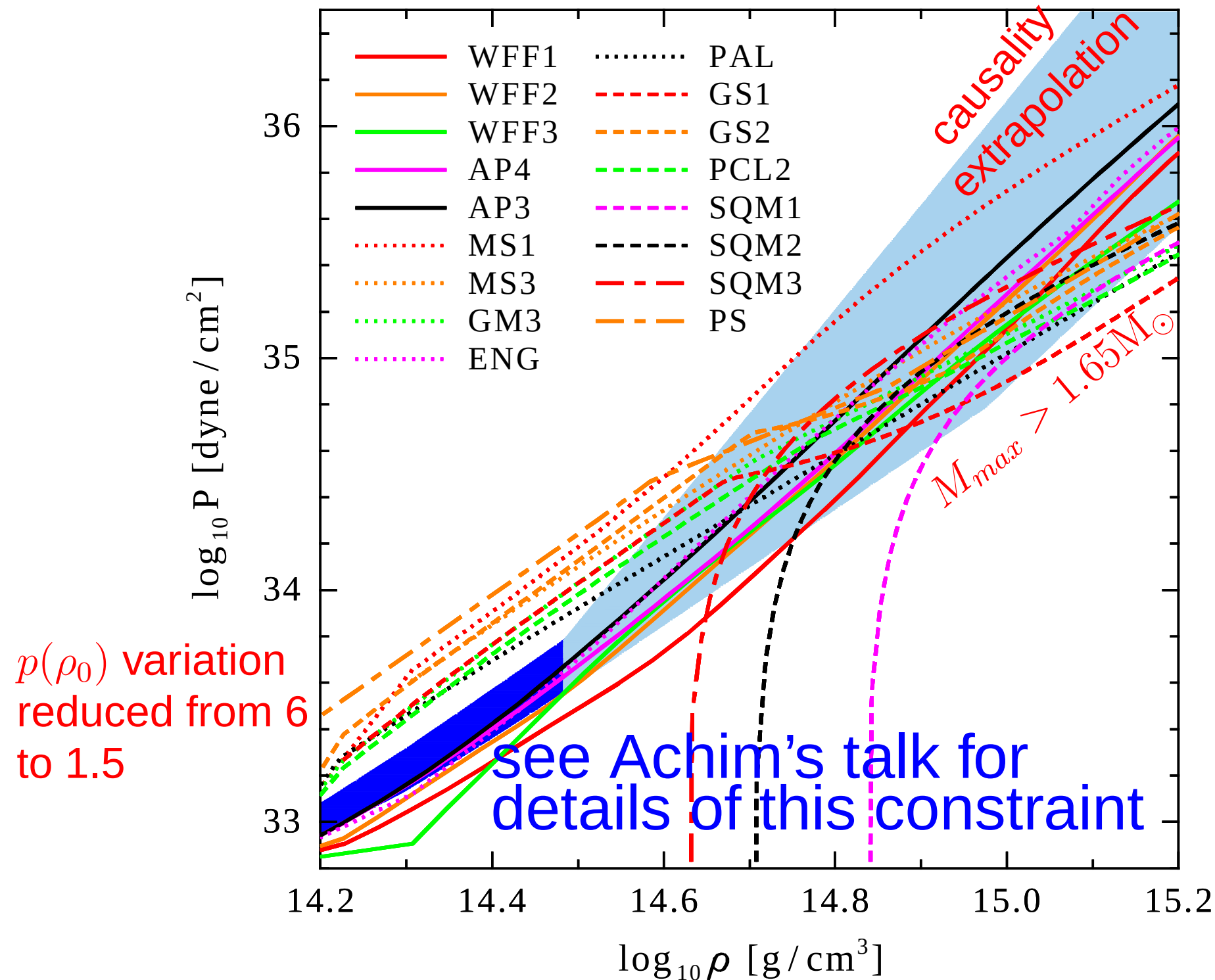
— SQS

$$— R_\infty = \frac{R}{\sqrt{1-2GM/Rc^2}} \text{ contours}$$



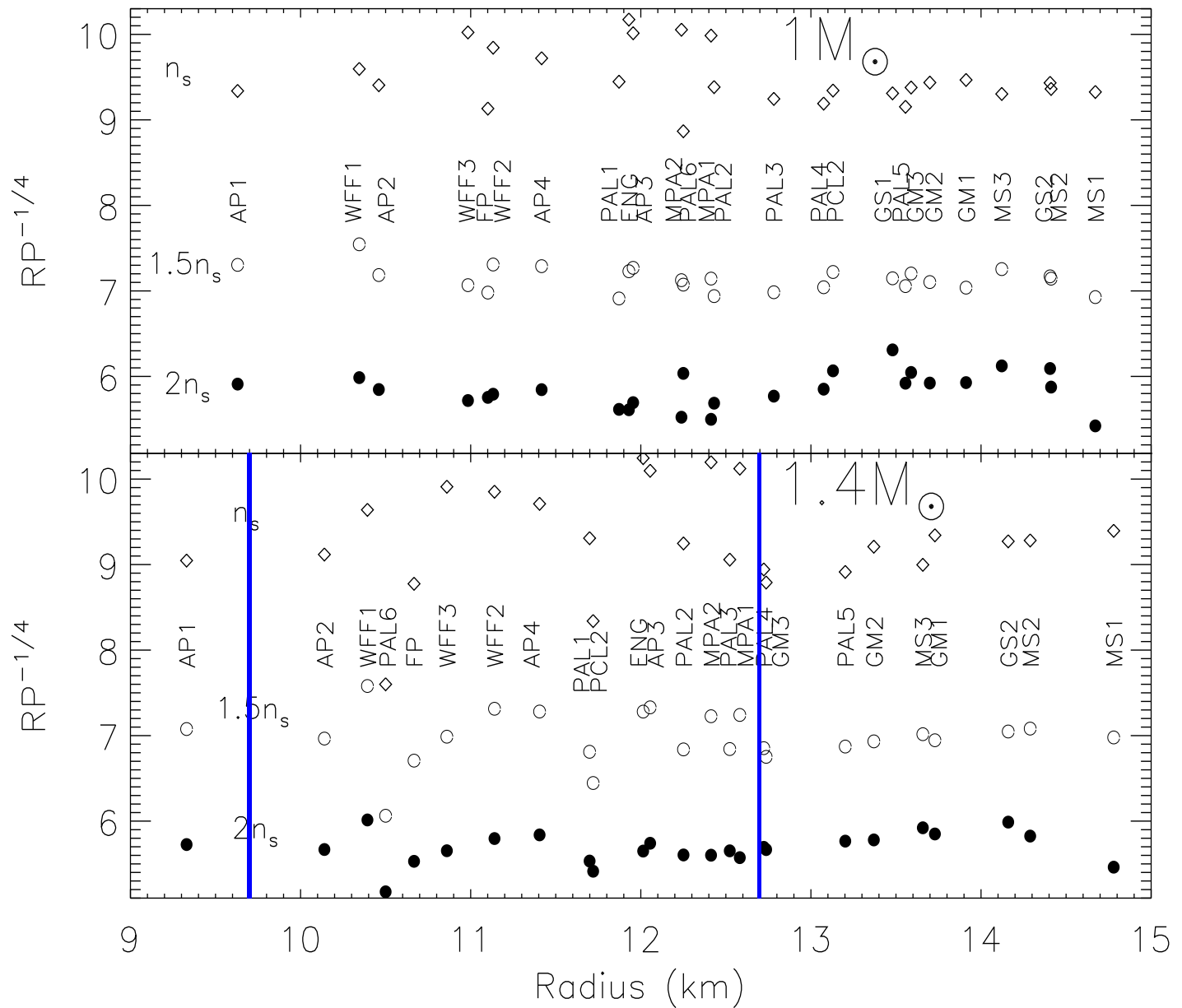
Constraints from Pulsar Spins





The Radius – Pressure Correlation

$$R \propto p^{1/4}$$



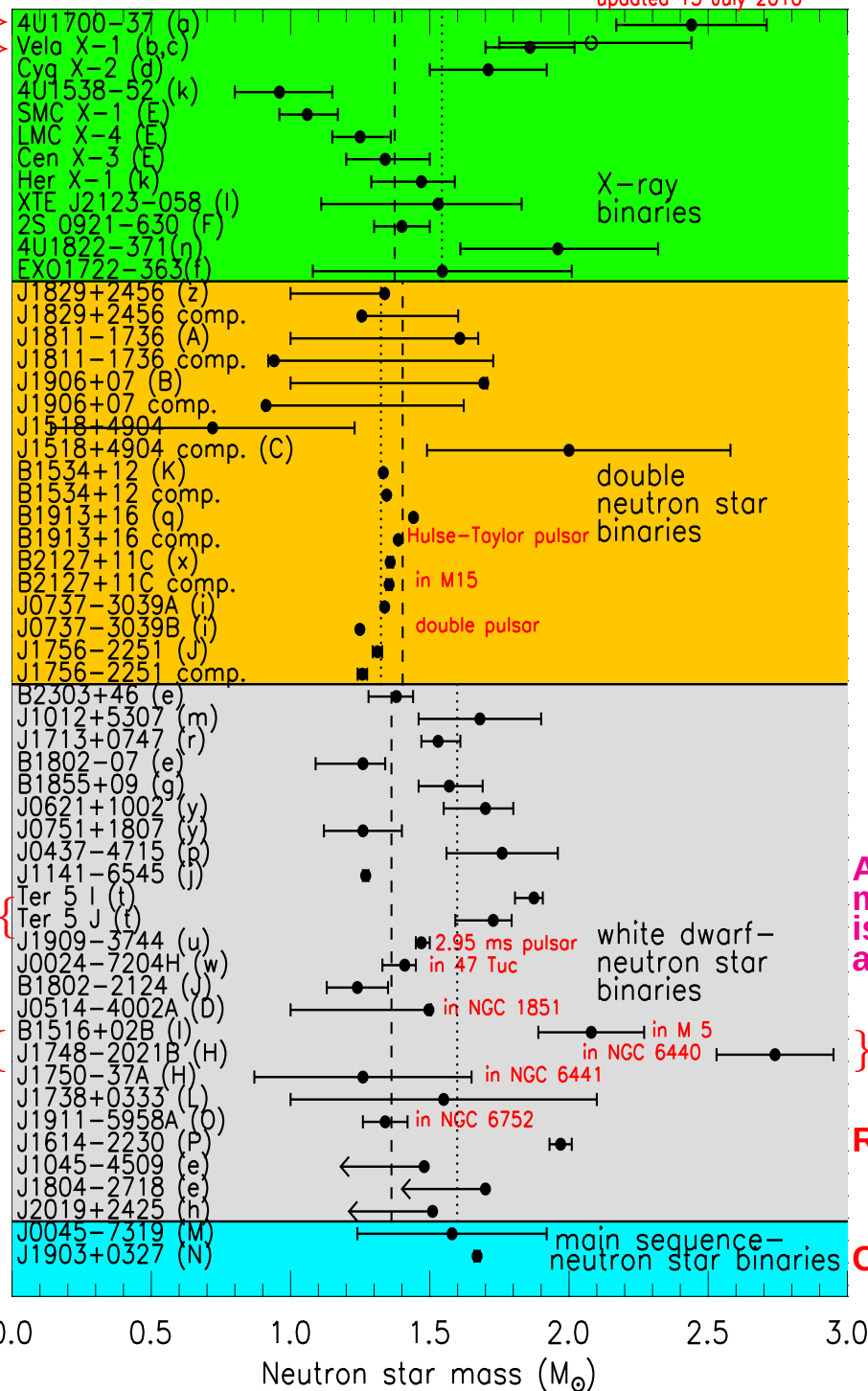
Lattimer & Prakash (2001)

Black hole? $\Rightarrow$
Firm lower mass limit? $\Rightarrow$

updated 13 July 2010

$M > 1.68 M_{\odot}$, 95% confidence

Freire et al. 2007



Although simple average mass of w.d. companions is $0.27 M_{\odot}$ larger, weighted average is $0.08 M_{\odot}$ smaller

w.d. companion? statistics?

Ransom et al. 2010

Champion et al. 2008

Possible Kinds of Observations

- Maximum and Minimum Mass (binary pulsars)
- Minimum Rotational Period*
- Radiation Radii or Redshifts from X-ray Thermal Emission*
- Crustal Cooling Timescale from X-ray Transients*
- X-ray Bursts from Accreting Neutron Stars*
- Seismology from Giant Flares in SGR's*
- Neutron Star Thermal Evolution (URCA or not)*
- Moments of Inertia from Spin-Orbit Coupling*
- Neutrinos from Proto-Neutron Stars (Binding Energies, Neutrino Opacities, Radii)*
- Pulse Shape Modulations*
- Gravitational Radiation from Neutron Star Mergers* (Masses, Radii from tidal Love numbers)

* Significant dependence on symmetry energy

Potentially Observable Quantities

- Apparent angular diameter from flux and temperature measurements

$$\beta \equiv GM/Rc^2$$

$$\frac{R_\infty}{D} = \frac{R}{D} \frac{1}{\sqrt{1-2\beta}} = \sqrt{\frac{F_\infty}{\sigma}} \frac{1}{f_\infty^2 T_\infty^2}$$

- Redshift

$$z = (1 - 2\beta)^{-1/2} - 1$$

- Eddington flux

$$F_{Edd} = \frac{GMc}{\kappa c^2 D^2} (1 - 2\beta)^{1/2}$$

- Crust thickness

$$\frac{m_b c^2}{2} \ln \mathcal{H} \equiv \int_0^{p_t} \frac{dp}{n} = \mu_{n,t} - \mu_{n,t}(p=0)$$

$$\frac{\Delta}{R} \equiv \frac{R - R_t}{R} \simeq \frac{(\mathcal{H} - 1)(1 - 2\beta)}{\mathcal{H} + 2\beta - 1} \simeq (\mathcal{H} - 1) \left(\frac{1}{2\beta} - 1 \right).$$

- Moment of Inertia

$$I \simeq (0.237 \pm 0.008) M R^2 (1 + 2.84\beta + 18.9\beta^4) \text{ M}_\odot \text{ km}^2$$

- Crustal Moment of Inertia

$$\frac{\Delta I}{I} \simeq \frac{8\pi}{3} \frac{R^6 p_t}{I M c^2}$$

- Binding Energy

$$\text{B.E.} \simeq (0.60 \pm 0.05) \frac{\beta}{1 - \beta/2}$$

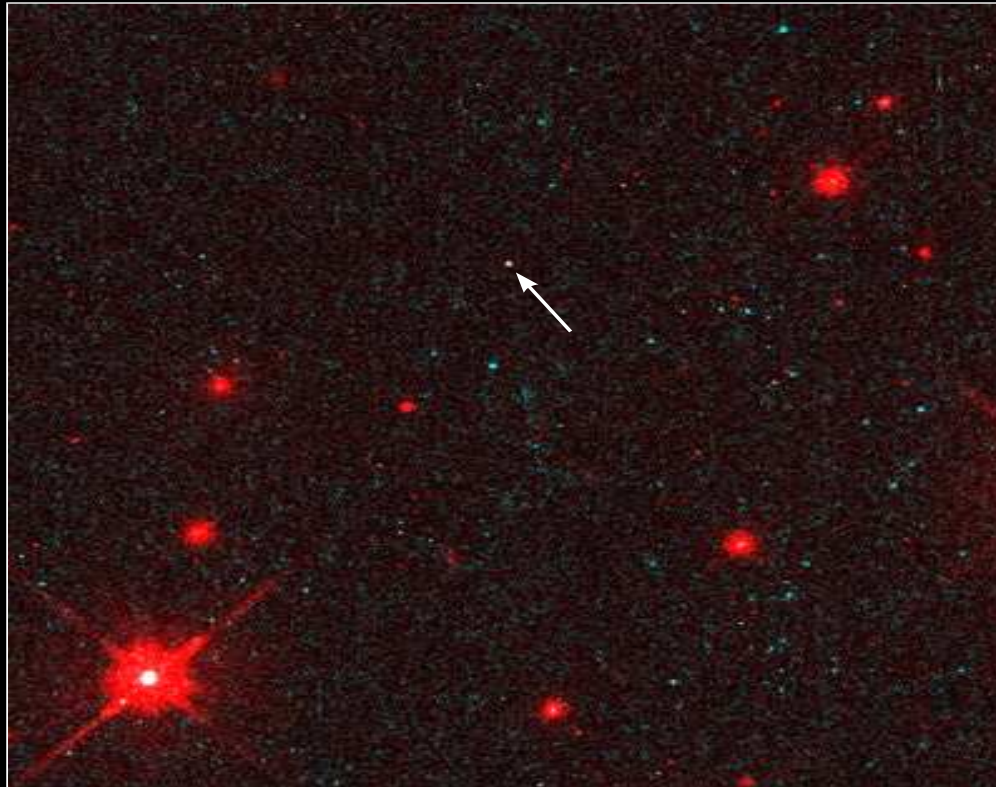
Radiation Radius

- Combination of flux and temperature measurements yields apparent angular diameter (pseudo-BB):

$$\frac{R_\infty}{d} = \frac{R}{d} \frac{1}{\sqrt{1 - 2GM/Rc^2}}$$

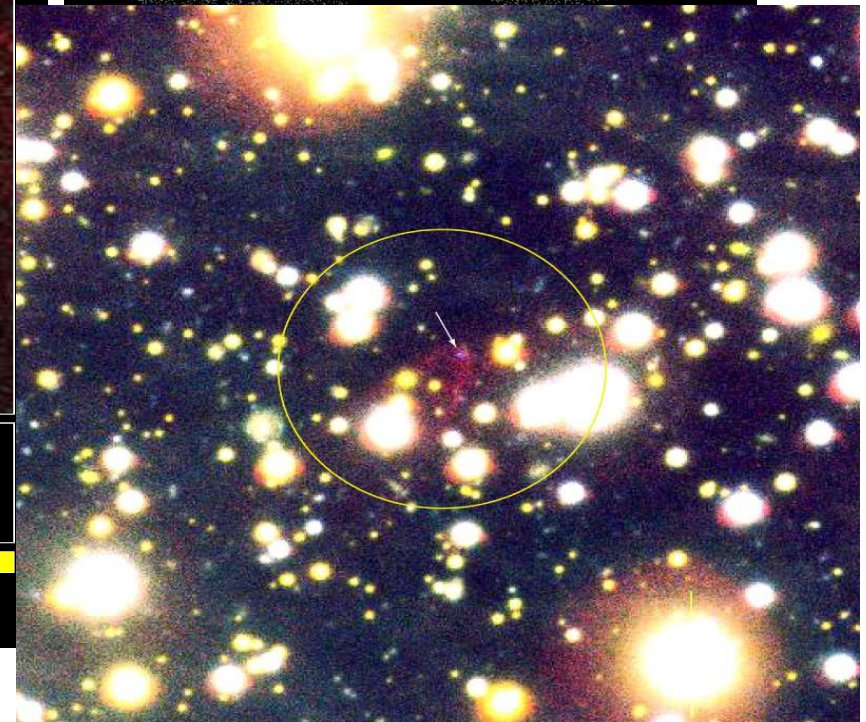
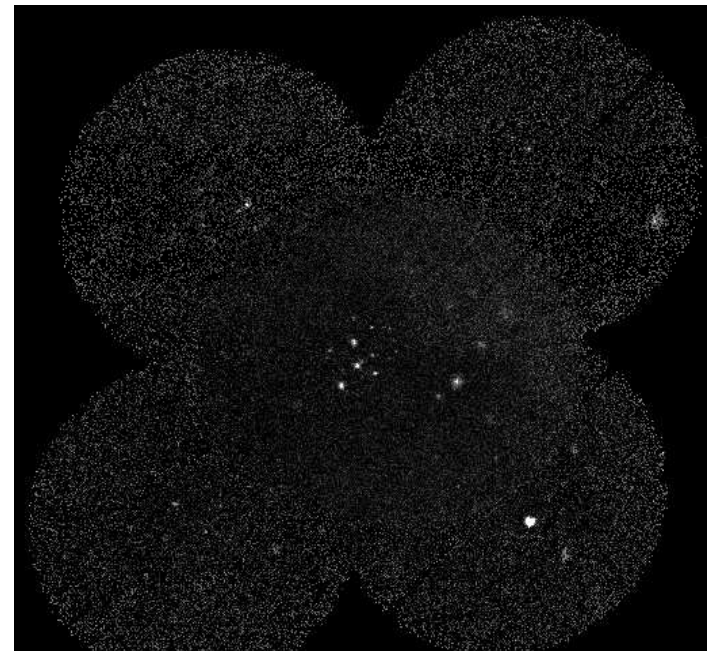
- Observational uncertainties include distance, interstellar H absorption (hard UV and X-rays), atmospheric composition
- Best chances for accurate radii are from
 - Nearby isolated neutron stars (parallax measurable)
 - Quiescent X-ray binaries in globular clusters (reliable distances, low *B* H-atmospheres)

RX J1856-3754



Isolated Neutron Star RX J185635-3754
Hubble Space Telescope • WFPC2

PRC97-32 • ST ScI OPO • September 25, 1997
F. Walter (State University of New York at Stony Brook) and NASA



A Bowshock Nebula Near the Neutron Star RX J1856.5-3754 (Detail)
(VLT KUEYEN + FORS2)

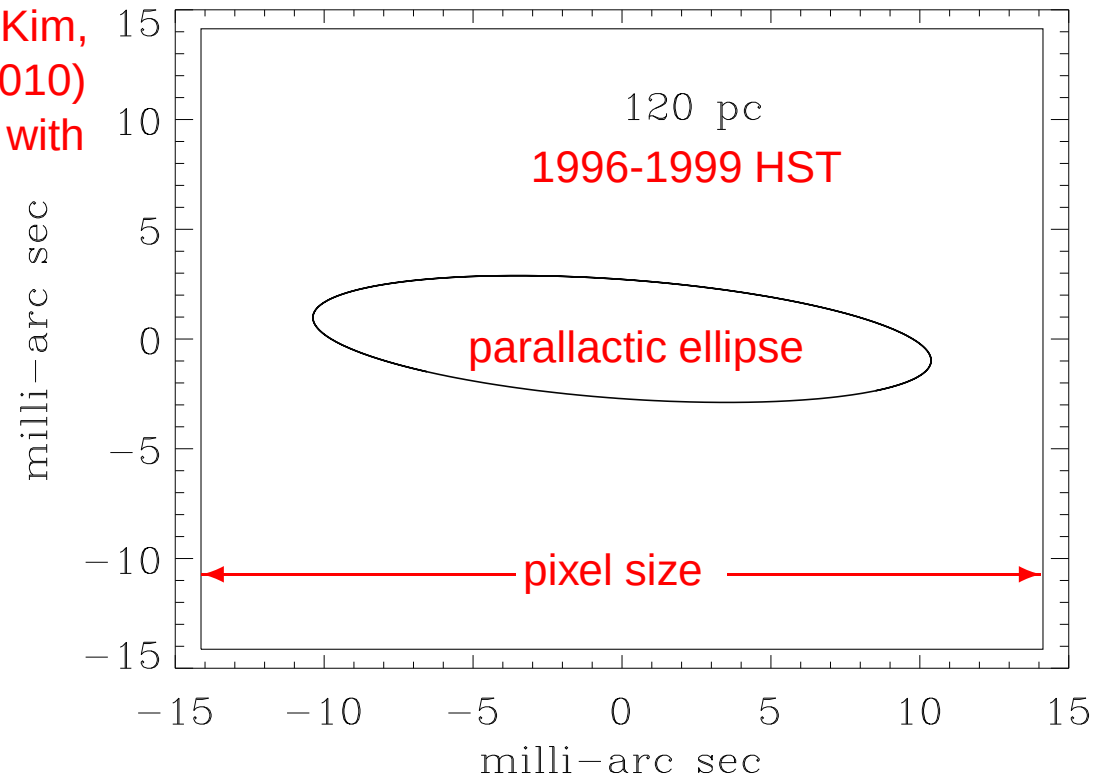
ESO PR Photo 23b/00 (11 September 2000)

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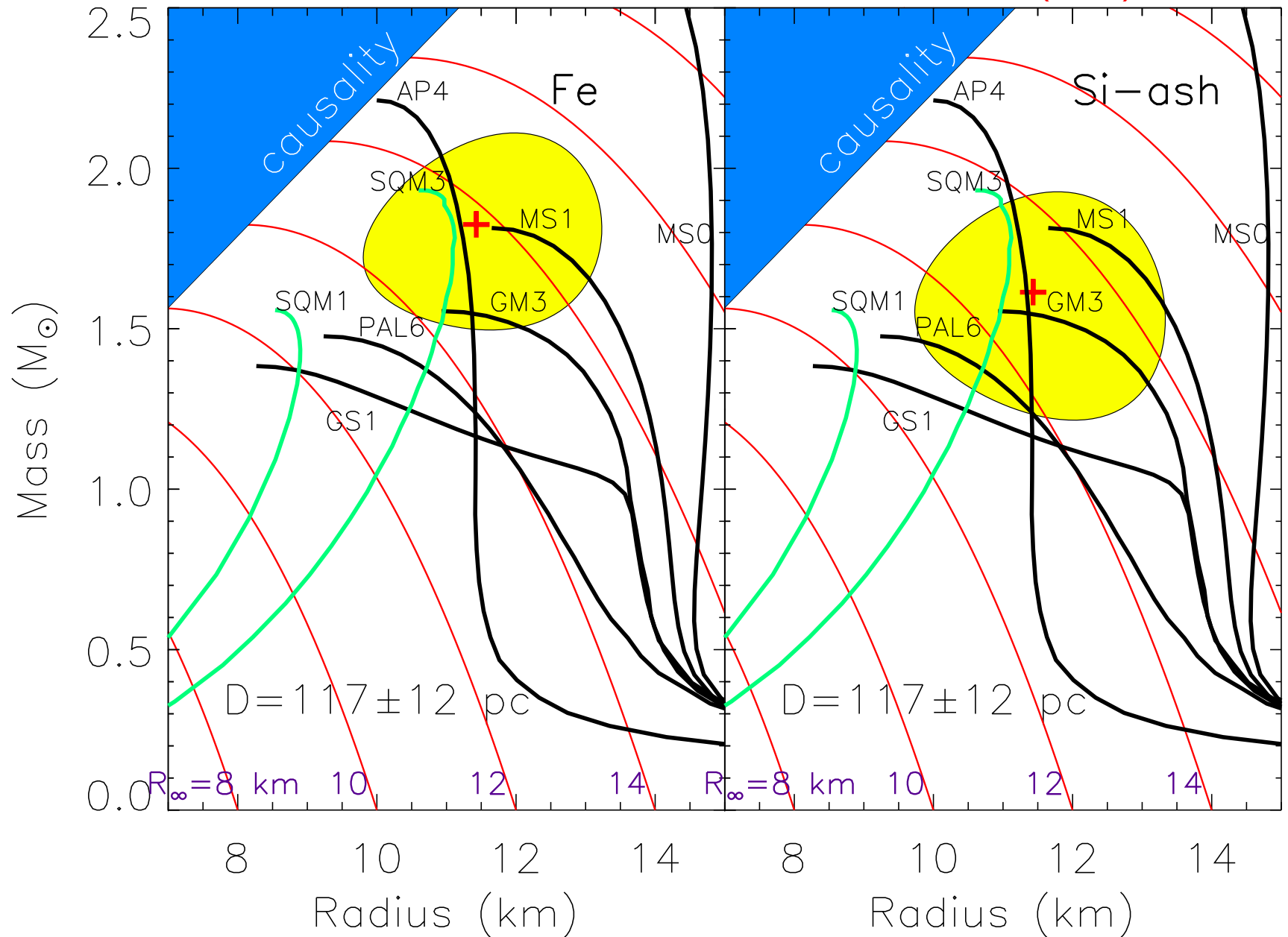
Astrometry of RXJ 1856-3754

- Walter & Lattimer (2002) determined $D = 117 \pm 12$ pc and $v \simeq 190$ km/s from 1996-1999 HST Planetary Camera observations
- Star's age is probably 0.5 million years
- Kaplan, van Kerkwijk & Anderson (2002): $D = 140 \pm 40$ pc using same data
- van Kerkwijk & Kaplan (2007, conference proceeding) revised this to $D = 161 \pm 16$ pc based on 2002-2004 High-Resolution Camera of the Advanced Camera for Surveys HST observations (double the resolution)
- Walter, Eisenbeiß, Lattimer, Kim, Hambaryan & Neuhäuser (2010) determined $D \simeq 115 \pm 8$ pc with 2002-2004 HST data

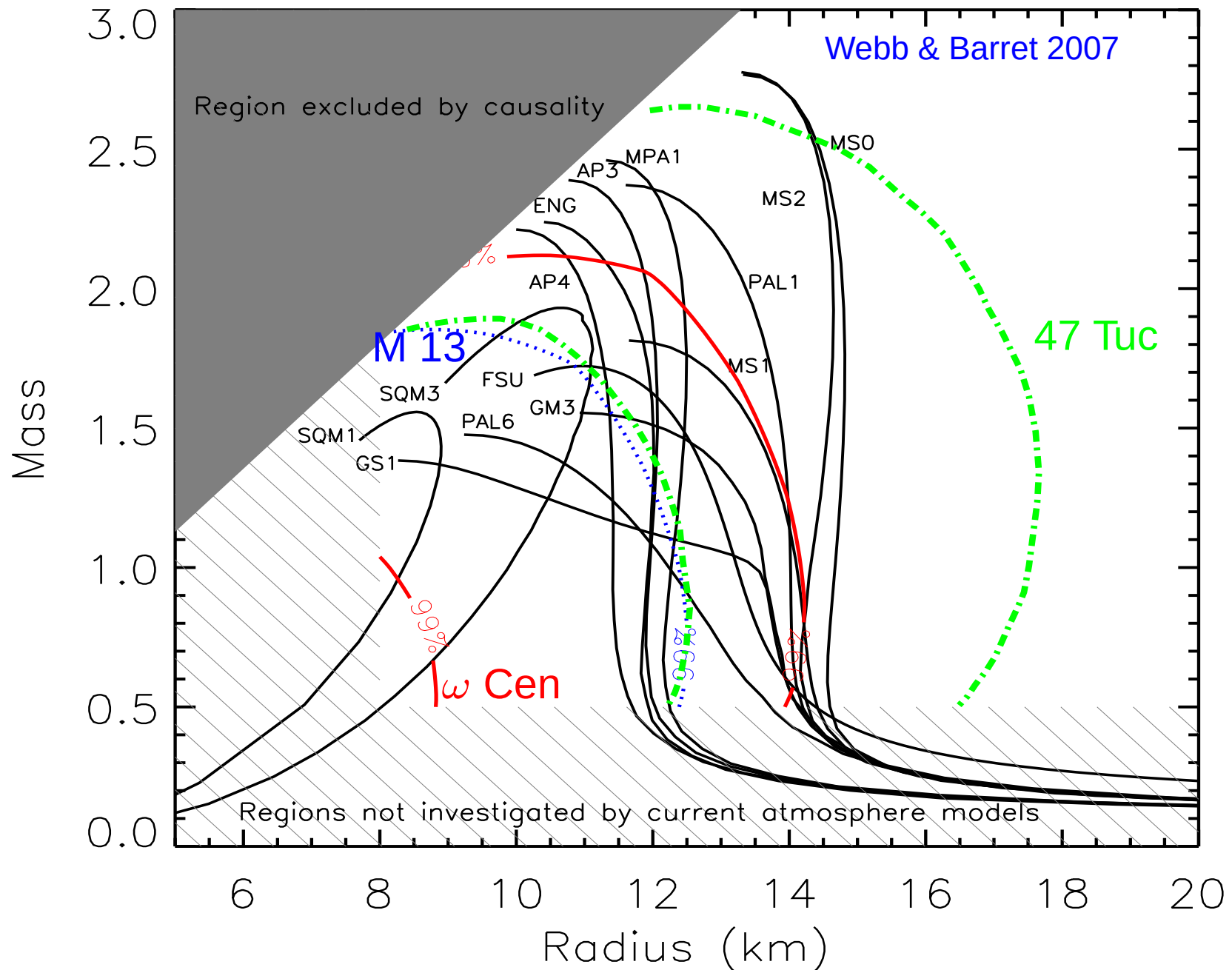


RX J1856-3754

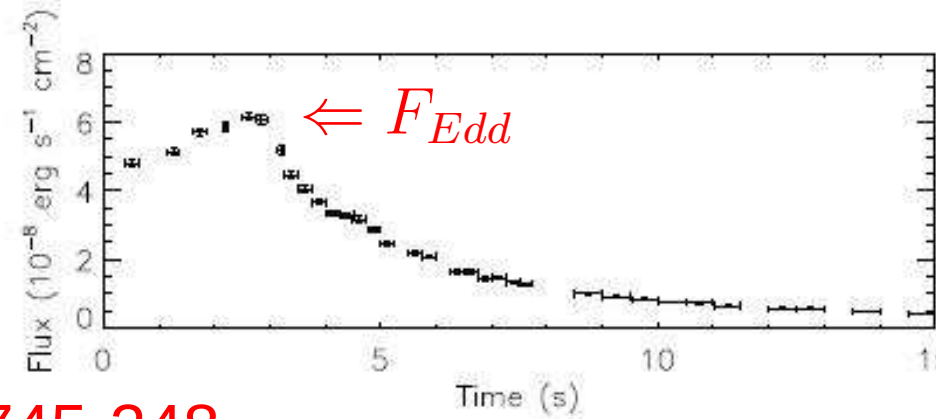
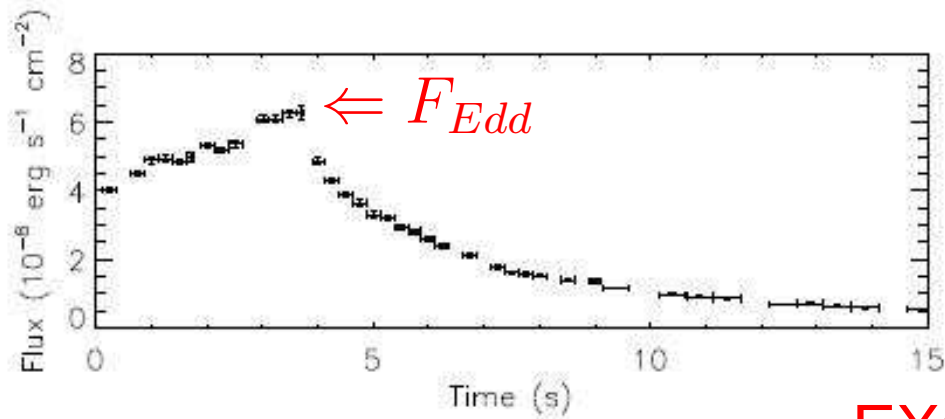
Walter & Lattimer (2002)



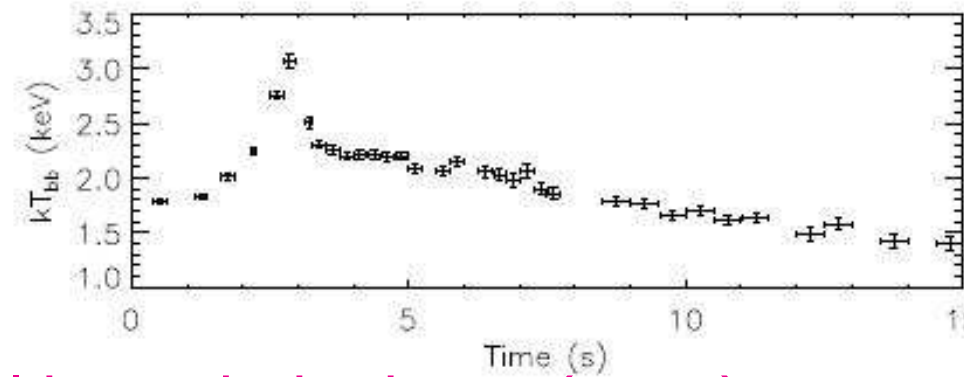
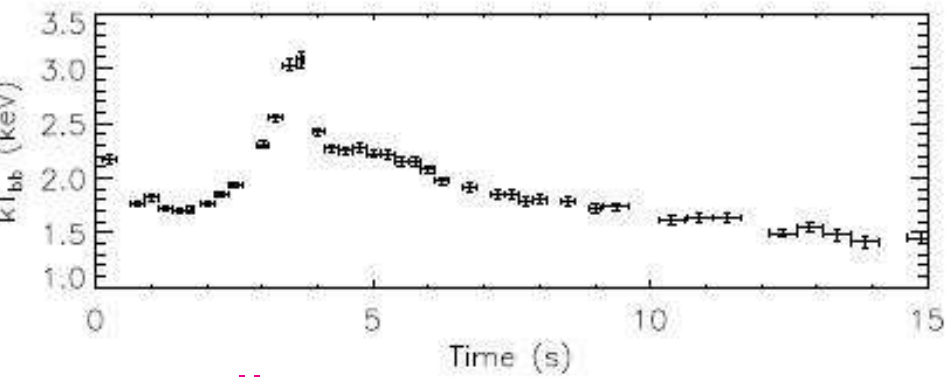
Radiation Radius: Globular Cluster Sources



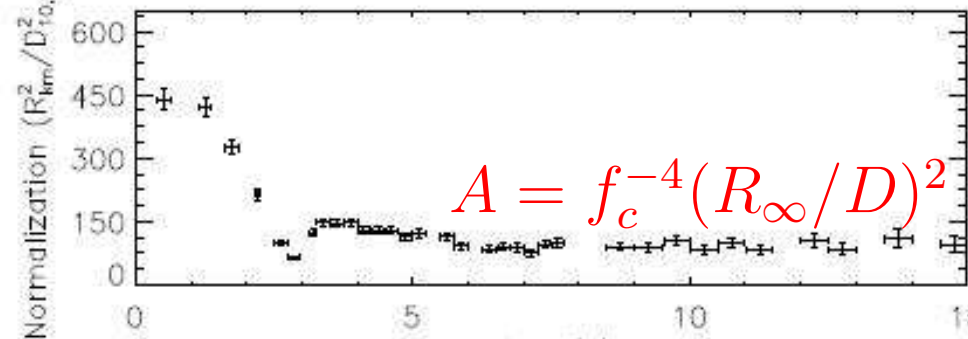
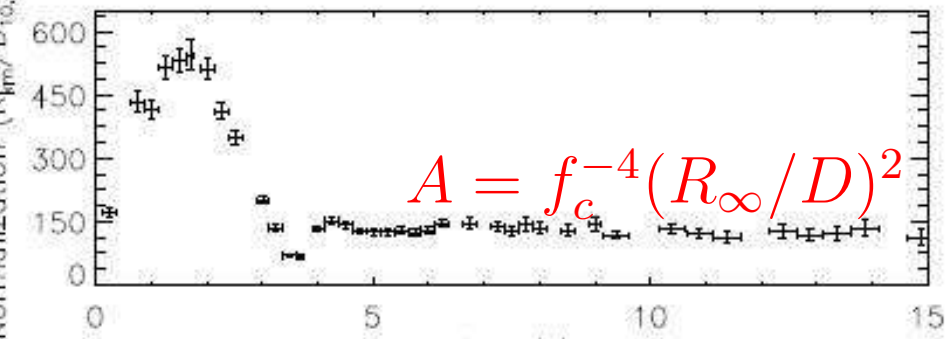
Photospheric Radius Expansion X-Ray Bursts



EXO 1745-248



Galloway, Muno, Hartman, Psaltis & Chakrabarty (2006)



Systematics with $R_{ph} = R$

$$F_{Edd} = \frac{GMc}{\kappa D^2} \sqrt{1 - 2\frac{GM}{R_{ph}c^2}} = \frac{GMc}{\kappa D^2} \sqrt{1 - 2\beta}$$

$$\kappa \simeq 0.2(1 + X) \text{ cm}^2 \text{g}^{-1}$$

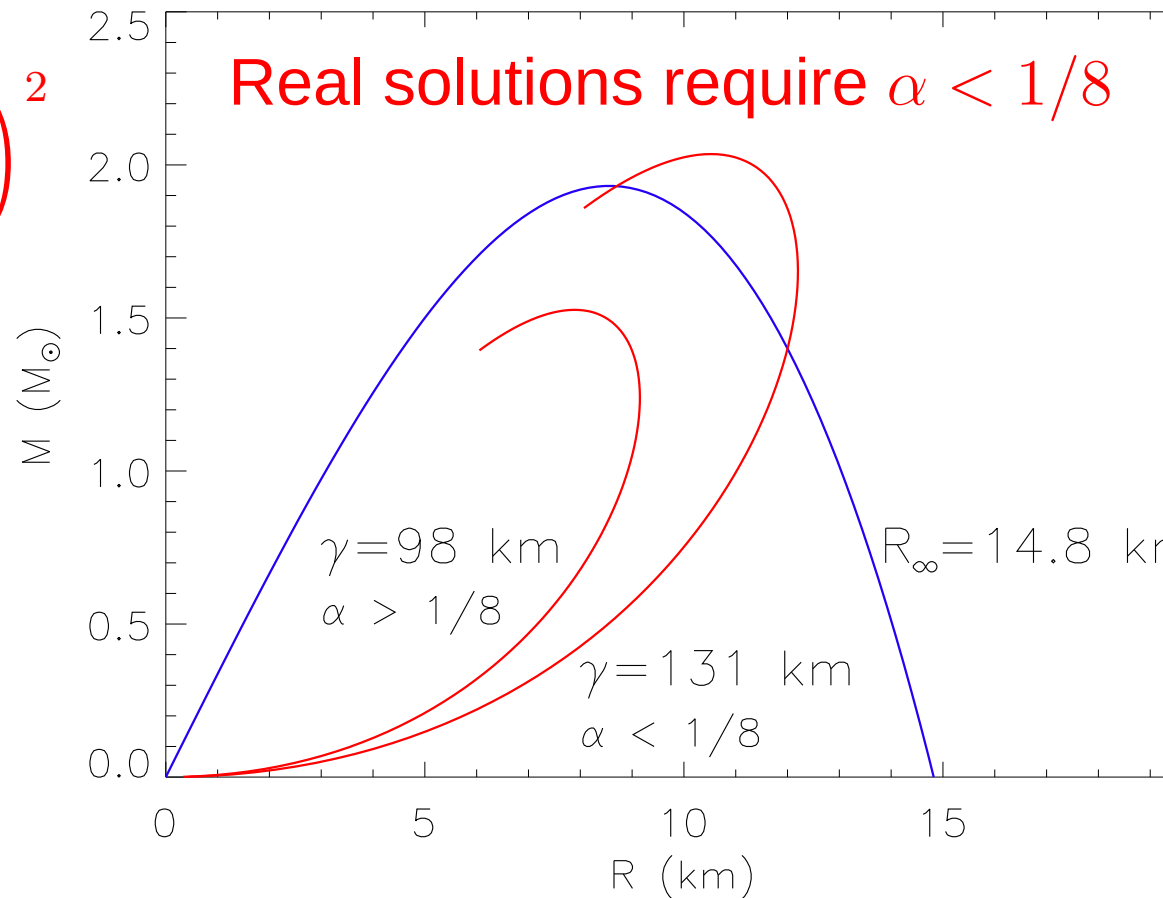
$$A = \frac{F_{\infty}}{\sigma T_{\infty}^4} = f_c^{-4} \left(\frac{R_{\infty}}{D} \right)^2$$

$$\alpha = \frac{F_{Edd}}{\sqrt{A}} \frac{\kappa D}{c^3 f_c^2} = \beta(1 - 2\beta)$$

$$\gamma = \frac{Ac^3 f_c^4}{F_{Edd} \kappa} = \frac{R}{\beta(1 - 2\beta)^{3/2}}$$

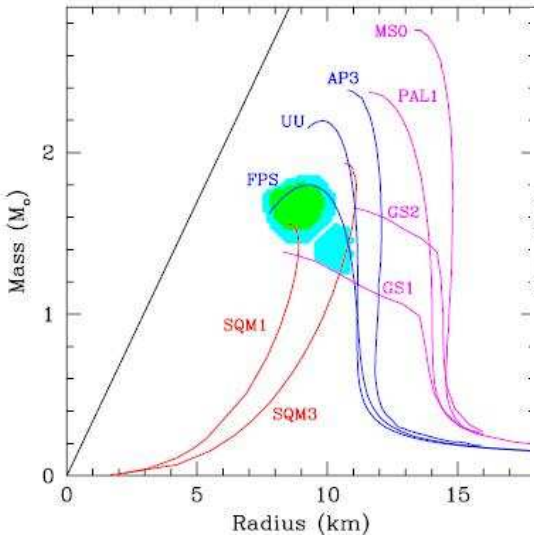
$$\beta = \frac{1}{4} \pm \frac{1}{4} \sqrt{1 - 8\alpha}$$

$$R_{\infty} = \alpha \gamma$$

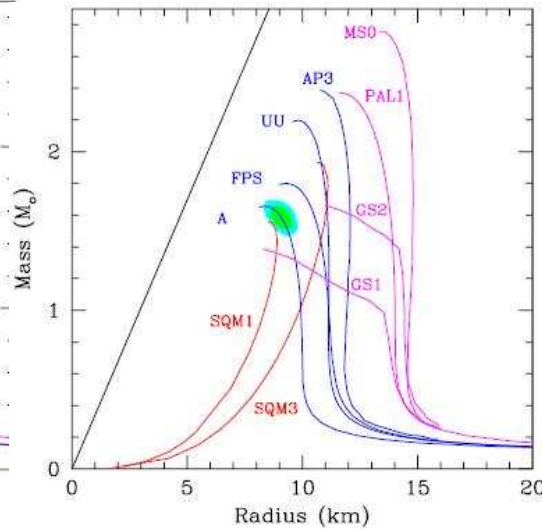


$$R_{ph} = R$$

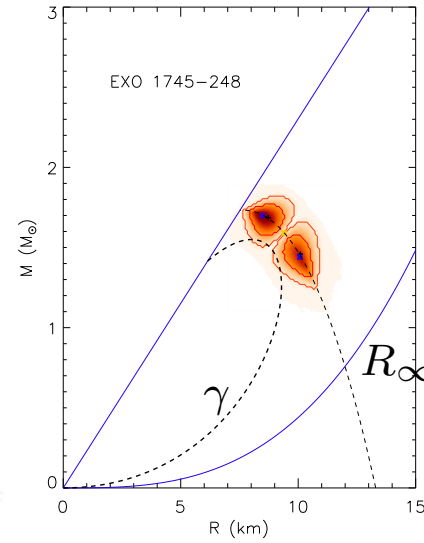
EXO 1745-248
 $\alpha = 0.14 \pm 0.02$



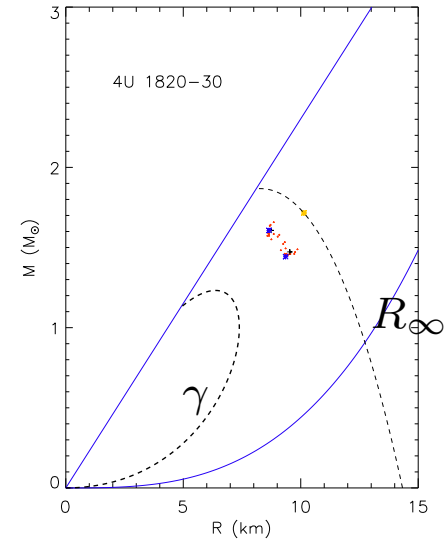
4U 1820-30
 $\alpha = 0.18 \pm 0.02$



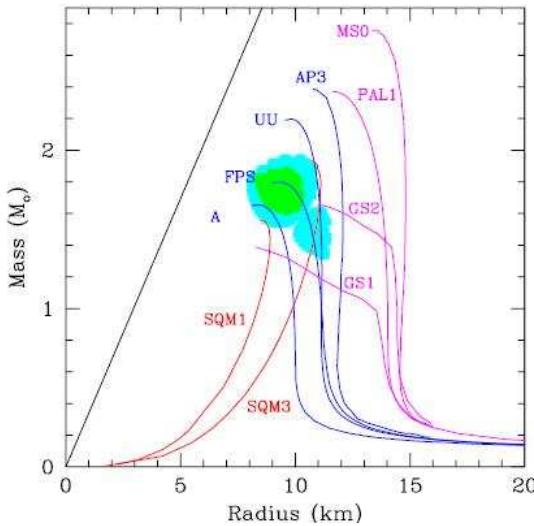
EXO 1745-248
 $\alpha = 0.14 \pm 0.02$



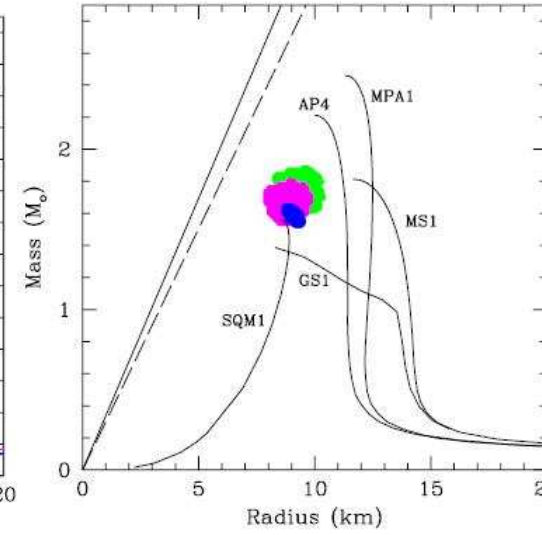
4U 1820-30
 $\alpha = 0.18 \pm 0.02$



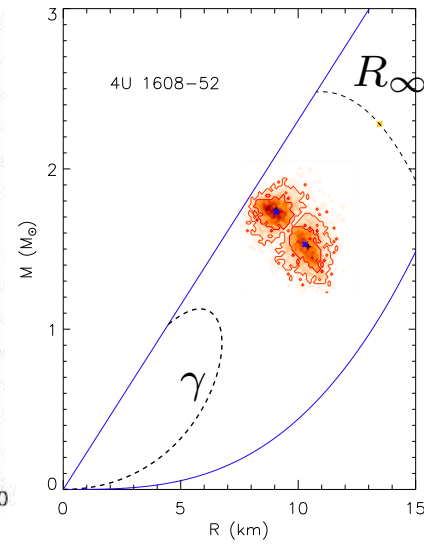
4U 1608-52
 $\alpha = 0.26 \pm 0.11$



Ozel, Baym & Güver 2010



4U 1608-52
 $\alpha = 0.26 \pm 0.11$



Systematics with $R_{ph} \gg R$

$$F_{Edd} = \frac{GMc}{\kappa D^2} \sqrt{1 - 2 \frac{GM}{R_{ph} c^2}} \simeq \frac{GMc}{\kappa D^2}$$

$$\kappa \simeq 0.2(1 + X) \text{ cm}^2 \text{g}^{-1}$$

$$A = \frac{F_{\infty}}{\sigma T_{\infty}^4} = f_c^{-4} \left(\frac{R_{\infty}}{D} \right)^2$$

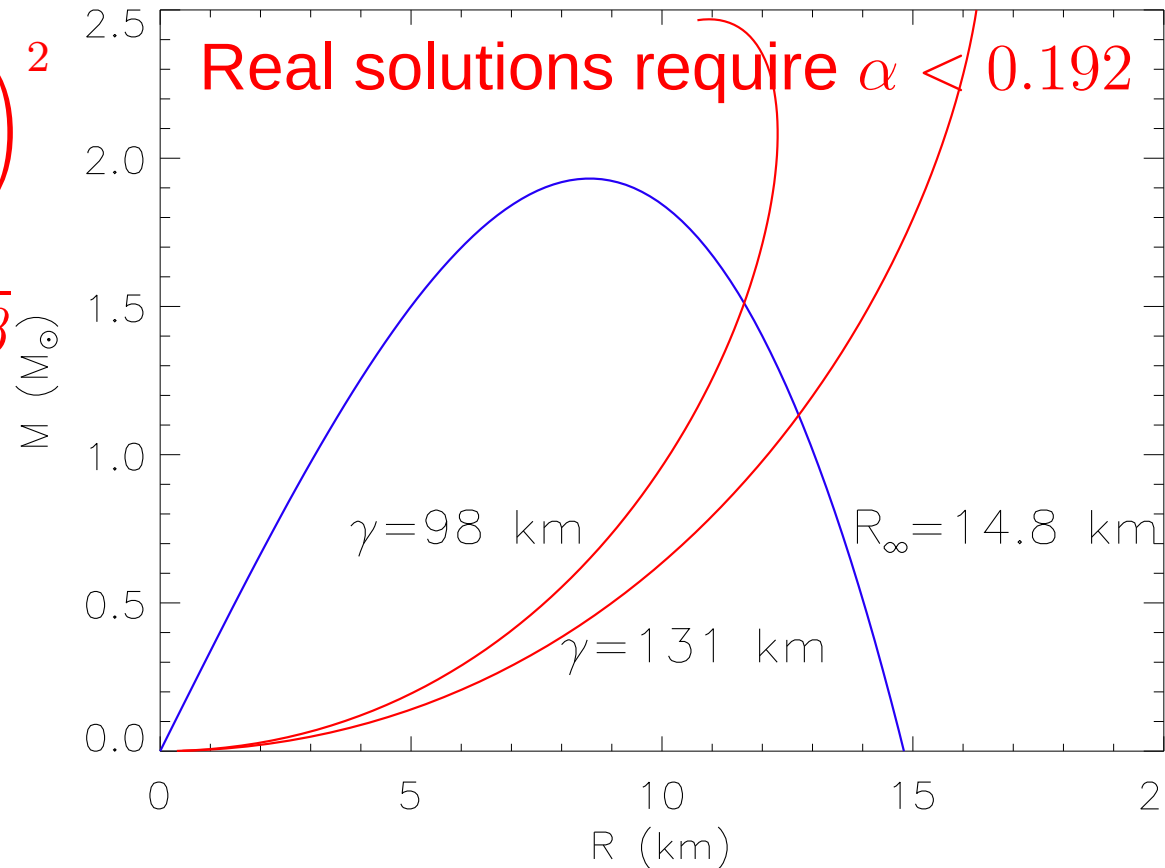
$$\alpha = \frac{F_{Edd}}{\sqrt{A}} \frac{\kappa D}{c^3 f_c^2} = \beta \sqrt{1 - 2\beta}$$

$$\gamma = \frac{Ac^3 f_c^4}{F_{Edd} \kappa} = \frac{R}{\beta(1 - 2\beta)}$$

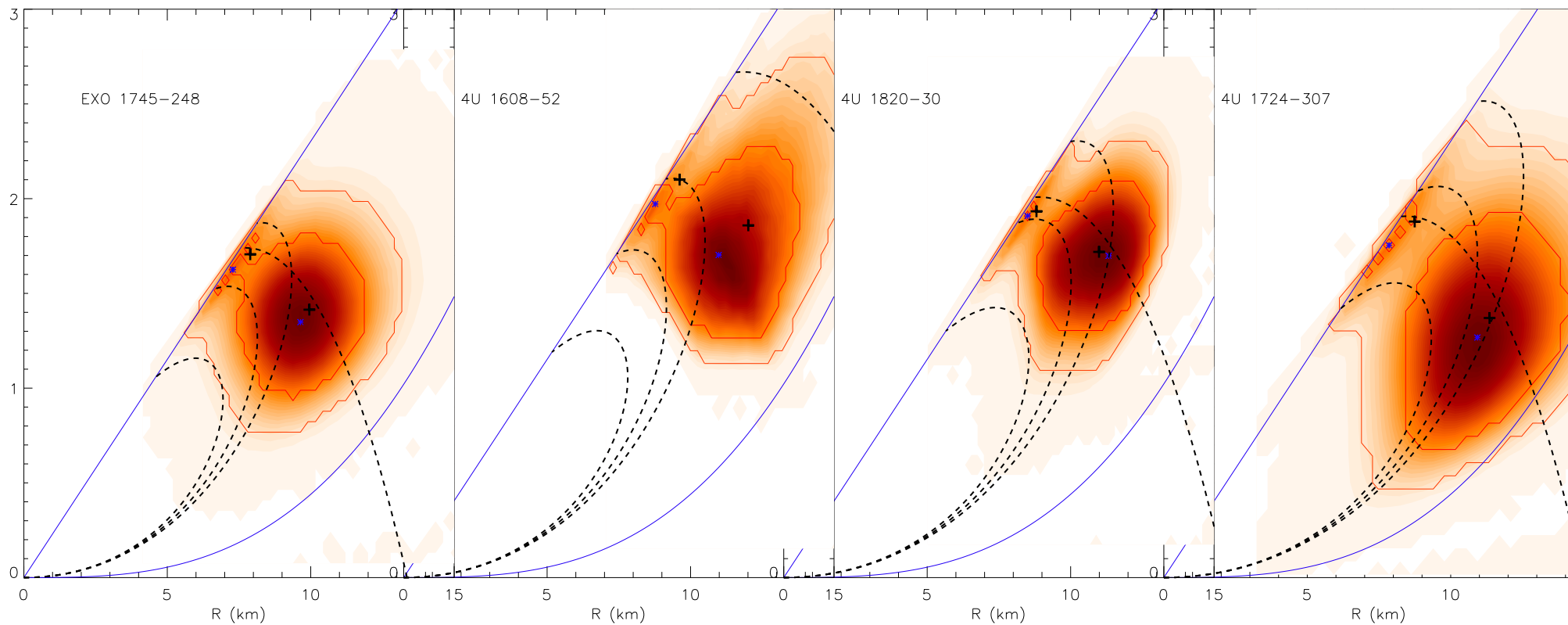
$$\beta = \frac{1}{6} \left[1 + \sqrt{3} \sin \theta - \cos \theta \right]$$

$$\theta = \frac{1}{3} \cos^{-1}(1 - 54\alpha^2)$$

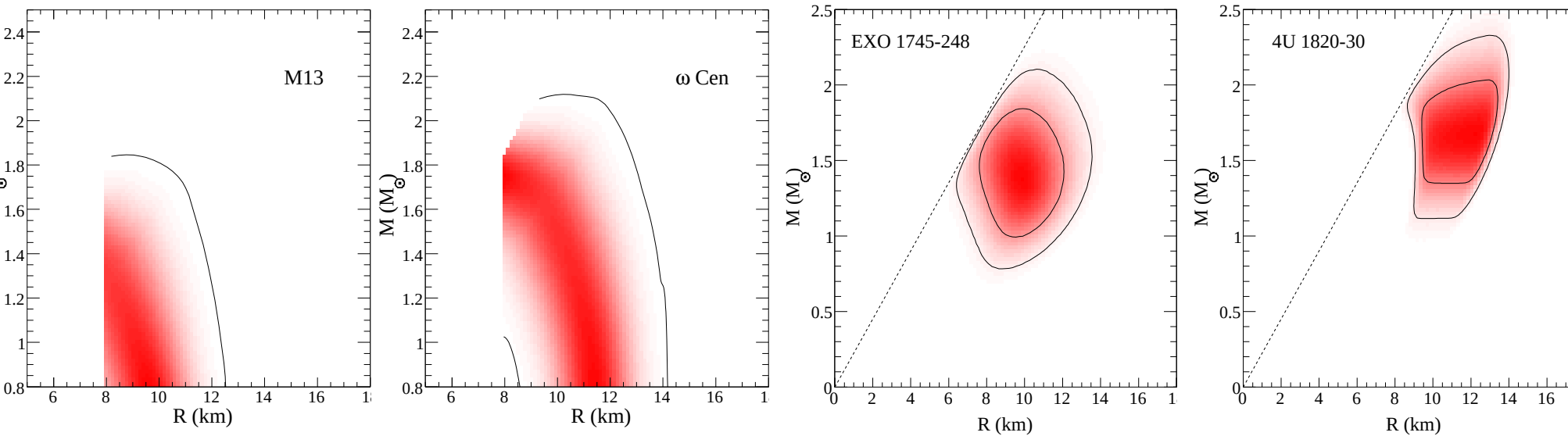
$$R_{\infty} = \alpha \gamma$$



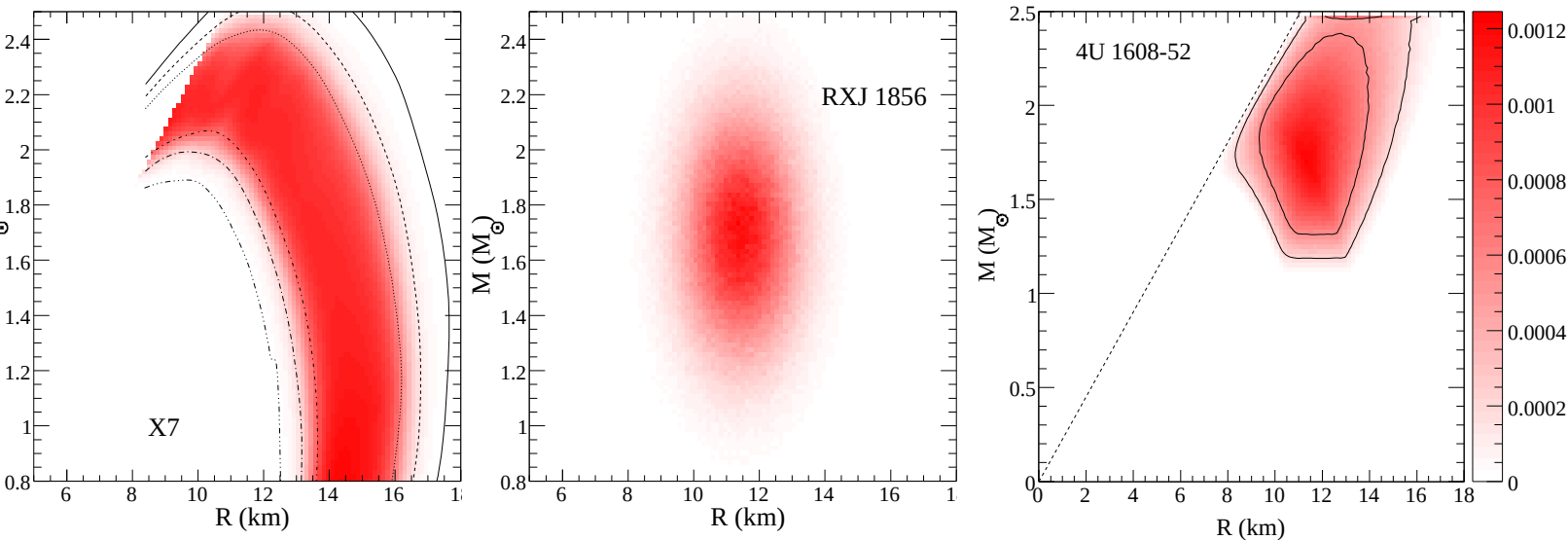
Arbitrary r_{ph}/R



M-R Probability Distributions



Steiner, Lattimer & Brown 2010



Bayesian Analysis

$P(\mathcal{M}_i|D)$: conditional probability of the model $\mathcal{M}_i$ given the data D (what we want).

$$P(\mathcal{M}_i|D) = \frac{P(D|\mathcal{M}_i)P(\mathcal{M}_i)}{\sum_k P(D|\mathcal{M}_k)P(\mathcal{M}_k)} \propto P(D|\mathcal{M}_i).$$

$P(\mathcal{M}_i)$: prior probability of the model $\mathcal{M}_i$ without any information from the data D (assumed uniform).

$P(D)$: prior probability of the data D ($M - R$ probabilities).

$P(D|\mathcal{M}_i)$: conditional probability of the data D given the model $\mathcal{M}_i$. i is the EOS parameter set and the set of specific masses j .

$$P(D|\mathcal{M}_i) \propto \prod_{k,j} \mathcal{D}_{k,j}|_{M=M_j, R=R_k(M_j)}$$

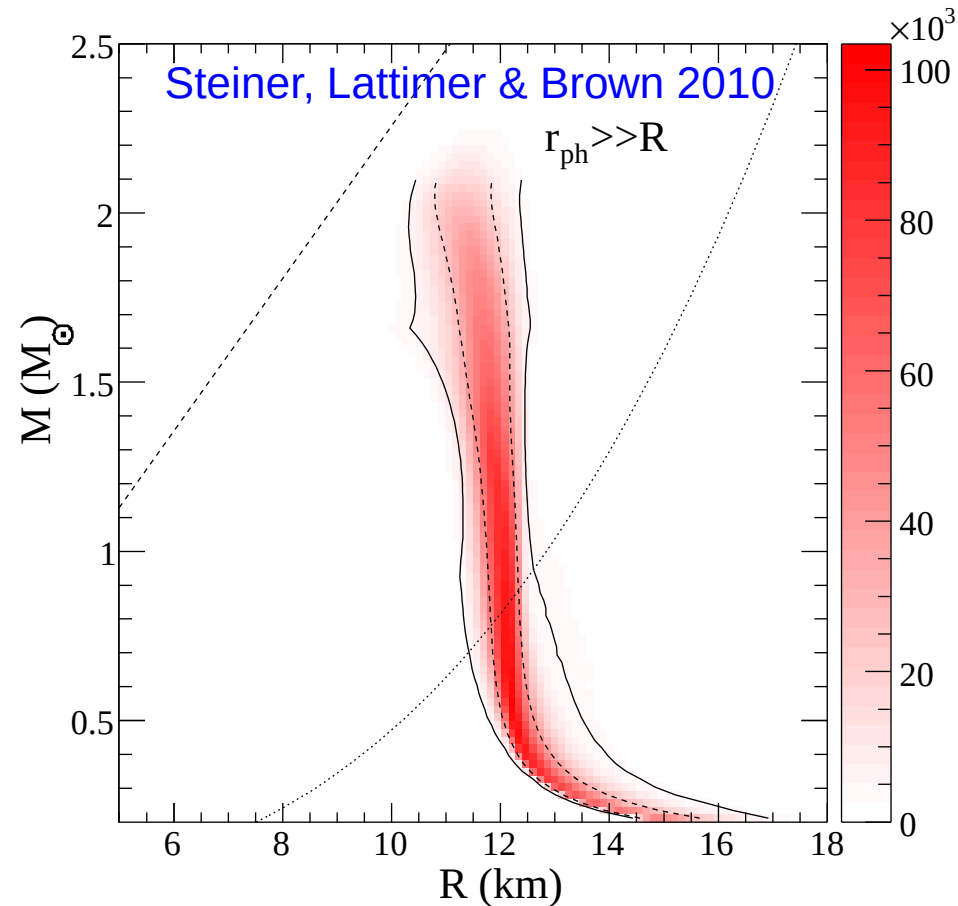
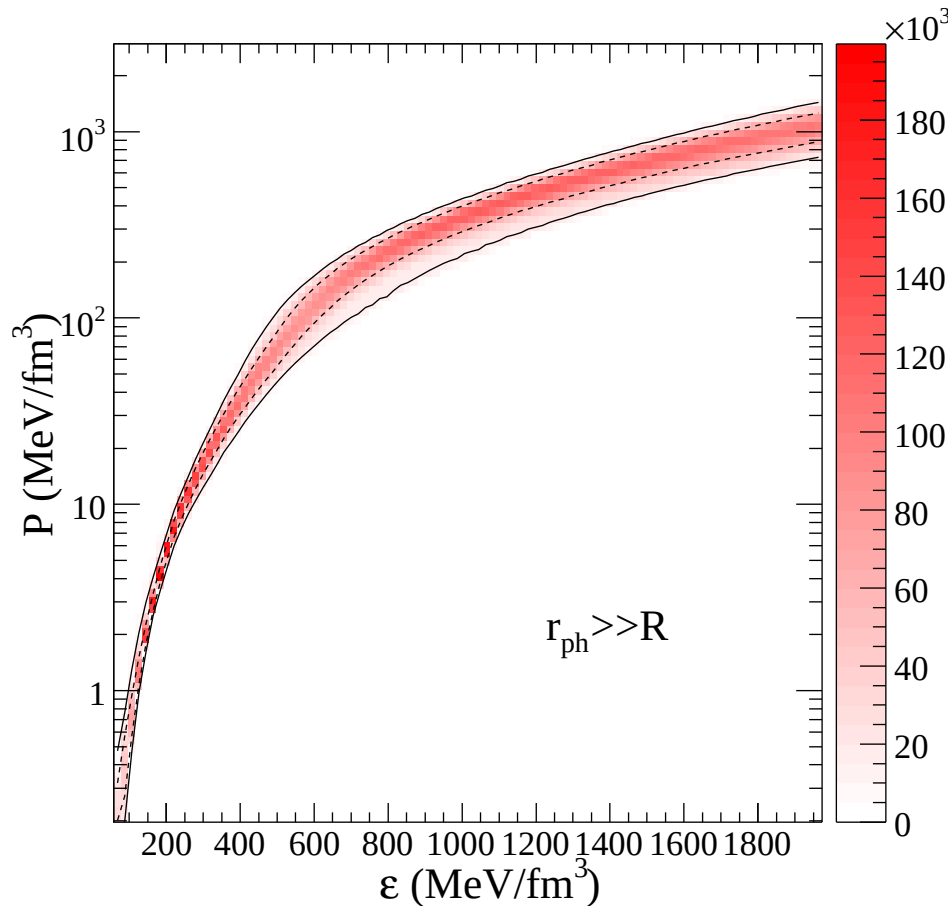
$\mathcal{D}_{k,j}$: probability of observing the radius R_j given a mass M_j in observation k .

The posterior probability for a model parameter p_i is

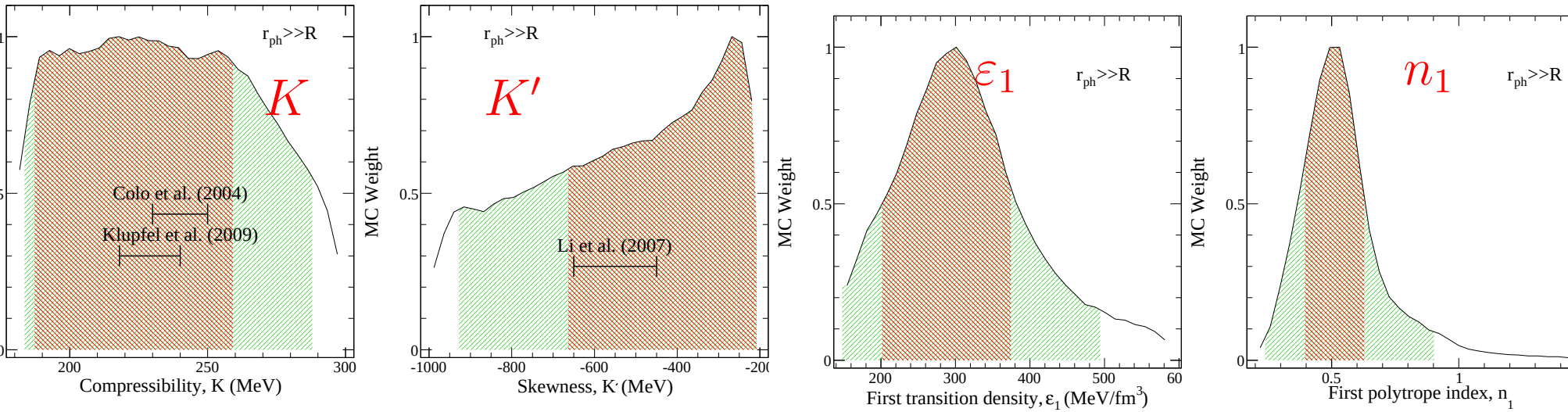
$$P(p_i|D) \propto \int P(D|\mathcal{M}_i) dp_1 dp_2 \dots dp_{i-1} dp_{i+1} \dots dp_{N_p} dM_1 dM_2 \dots dM_{N_M}$$

Bayesian TOV Inversion

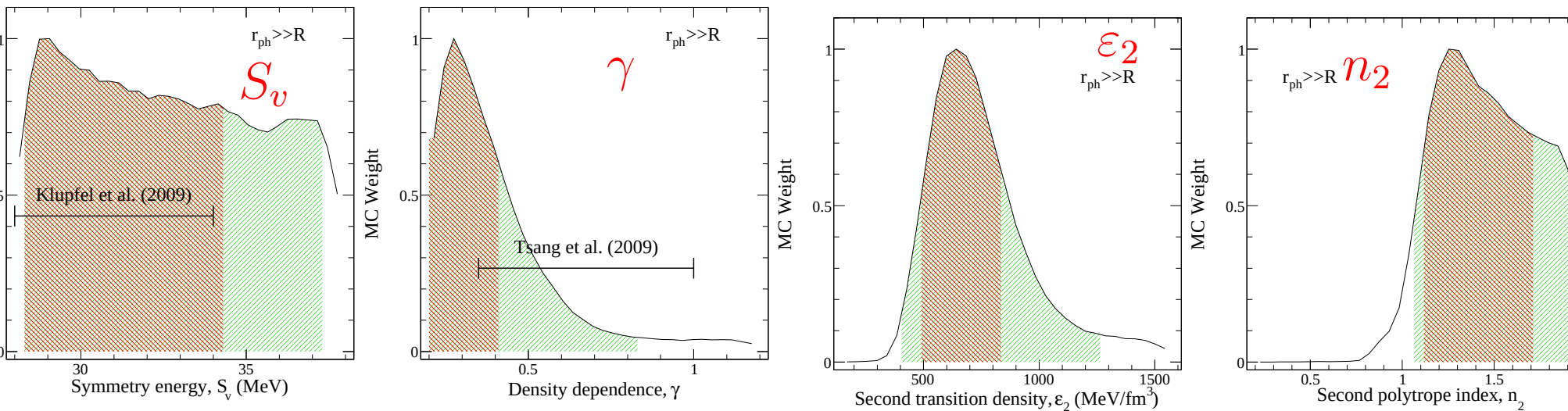
- $\varepsilon < 0.5\varepsilon_0$: EOS from BBP and NV
- $0.5\varepsilon_0 < \varepsilon < \varepsilon_1$: EOS parametrized by K, K', S_v, γ
- $\varepsilon_1 < \varepsilon < \varepsilon_2$: EOS is polytrope with n_1 ; $\varepsilon > \varepsilon_2$: EOS is polytrope with n_2
- A-priori EOS parameters ($K, K', S_v, \gamma, \varepsilon_1, n_1, \varepsilon_2, n_2$) uniformly distributed
- M and R probability distributions for 7 neutron stars treated equally ($0.8 M_\odot < M < 2.5 M_\odot$; $5 \text{ km} < R < 18 \text{ km}$)



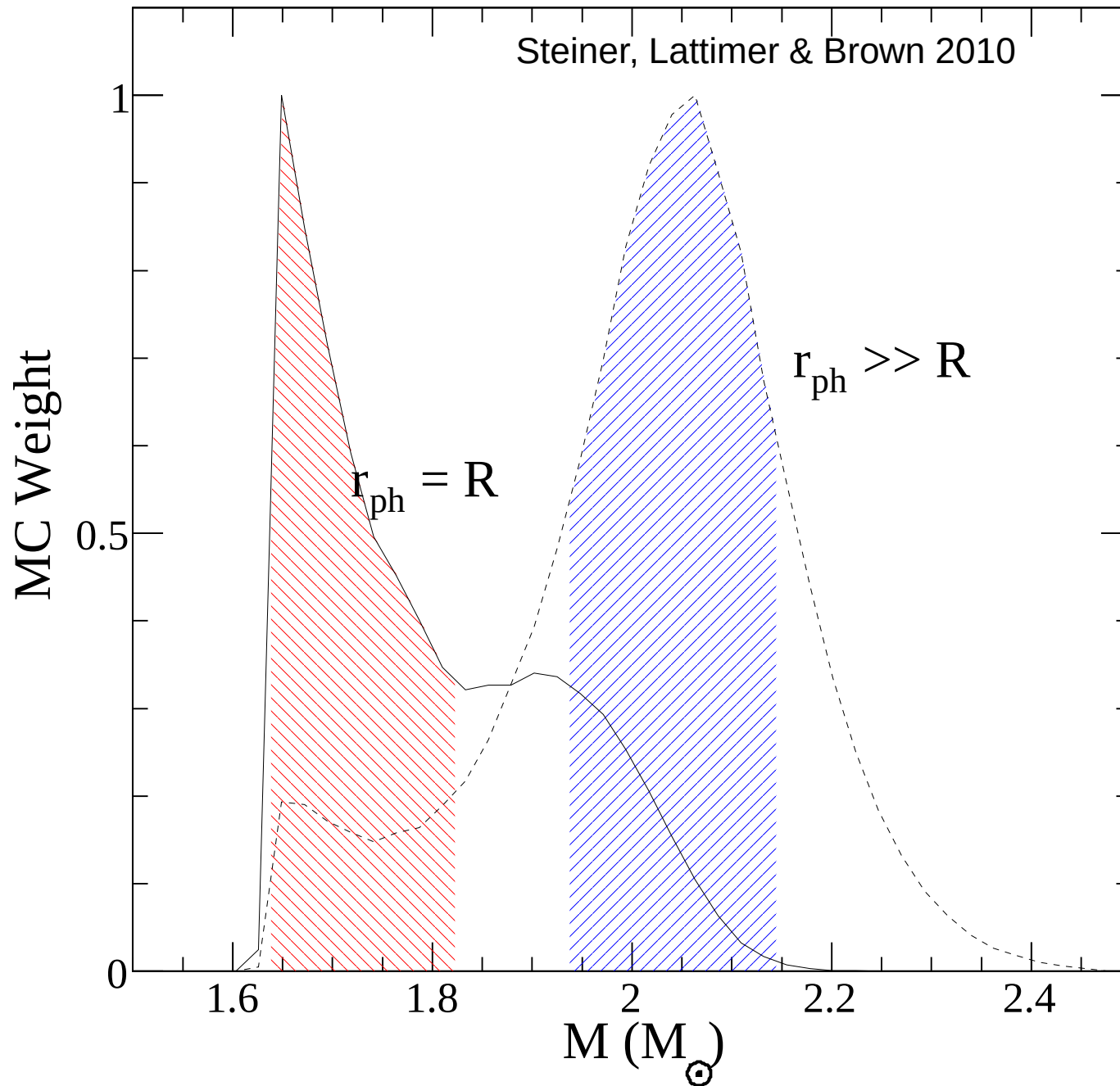
Inferred Model EOS Parameters



Steiner, Lattimer & Brown 2010



Maximum Mass Probability Distributions



Conclusions

- In spite of large estimated errors for individual stars, taken as an ensemble, they predict a remarkably tight pressure-density relation.
- Nuclear symmetry energy is predicted to be very soft, $\gamma = 0.3 \pm 0.1$
- Maximum Neutron Star Mass may be above $2 M_{\odot}$
- Radii of $1.4 M_{\odot}$ stars are apparently on the small side, $11.3 \pm 0.3 \text{ km}$ ($1 - \sigma$)