

Neutron-rich nuclei, neutron matter and constraints on neutron star structure from chiral effective field theory interactions

Achim Schwenk



TECHNISCHE
UNIVERSITÄT
DARMSTADT



EMMI Workshop, GSI, July 15, 2010



OAK RIDGE NATIONAL LABORATORY



東京大学
THE UNIVERSITY OF TOKYO



NORDITA

Niels Bohr Institutet



Thanks to collaborators and main message

J. Menendez

K. Hebeler

S.K. Bogner

R.J. Furnstahl

A. Nogga

J.D. Holt

T. Otsuka

T. Suzuki

Y. Akaishi

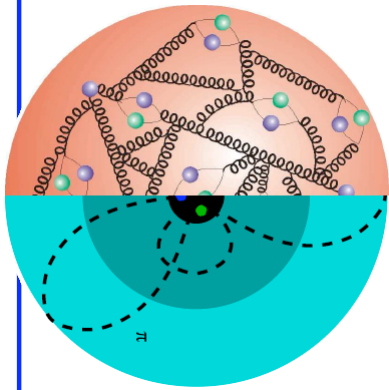
C.J. Pethick

J.M. Lattimer

3N forces are a frontier:
impact the structure and existence
of neutron-rich nuclei and
neutron-rich matter in astrophysics

Λ / Resolution dependence of nuclear forces

with high-energy probes:
quarks+gluons



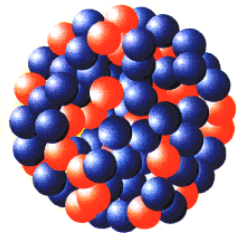
Effective theory for NN, 3N, many-N interactions and electroweak operators: resolution scale/ Λ -dependent

$$H(\Lambda) = T + V_{\text{NN}}(\Lambda) + V_{\text{3N}}(\Lambda) + V_{\text{4N}}(\Lambda) + \dots$$

Λ_{chiral}

momenta $Q \sim \lambda^{-1} \sim m_\pi$: chiral effective field theory

neutrons and protons interacting via pion exchanges
and shorter-range contact interactions



typical momenta in nuclei $\sim m_\pi$

$\Lambda_{\text{pionless}}$

$Q \ll m_\pi$

Chiral Effective Field Theory for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale ~ 500 MeV

NN

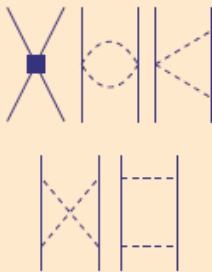
3N

4N

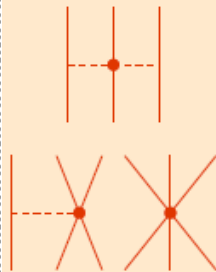
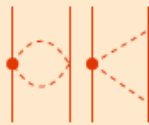
LO $\mathcal{O}\left(\frac{Q^0}{\Lambda^0}\right)$



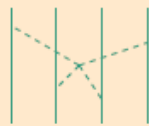
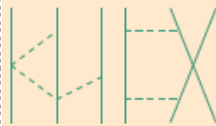
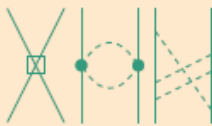
NLO $\mathcal{O}\left(\frac{Q^2}{\Lambda^2}\right)$



N²LO $\mathcal{O}\left(\frac{Q^3}{\Lambda^3}\right)$



N³LO $\mathcal{O}\left(\frac{Q^4}{\Lambda^4}\right)$



+ ...

+ ...

+ ...

limited resolution at low energies,
can expand in powers $(Q/\Lambda_b)^n$

expansion parameter $\sim 1/3$

Chiral Effective Field Theory for nuclear forces

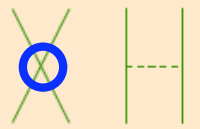
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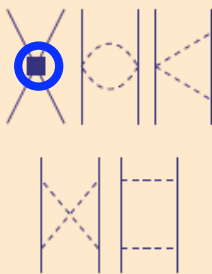
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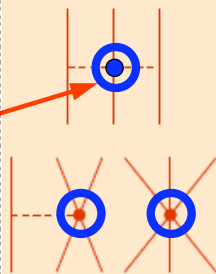
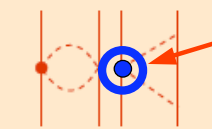
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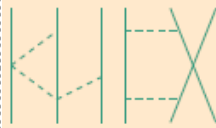
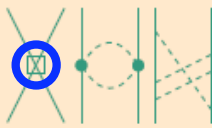
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+ ...

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include long-range pion physics

details at short distance not resolved

capture in few **short-range couplings**,
fit to experiment once, **Λ** -dependent

Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Meissner,...

Chiral Effective Field Theory for nuclear forces

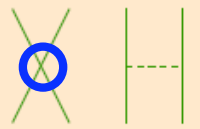
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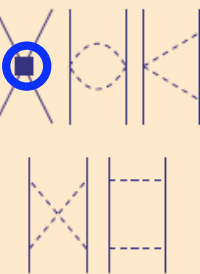
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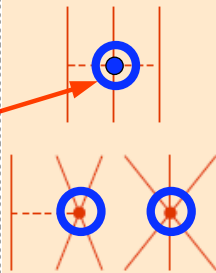
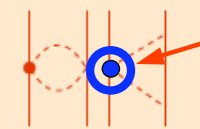
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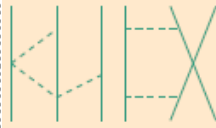
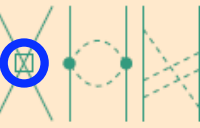
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include long-range pion physics

details at short distance not resolved

capture in few **short-range couplings**,
fit to experiment once, **Λ** -dependent

systematic: can work to desired
accuracy and obtain error estimates
from truncation order and Λ variation

can connect to lattice QCD

several open problems regarding
renormalization and power counting

Chiral Effective Field Theory for nuclear forces

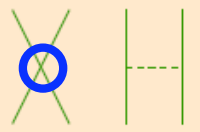
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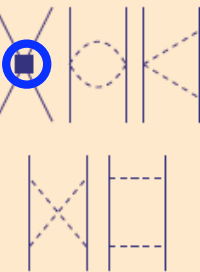
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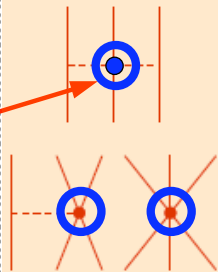
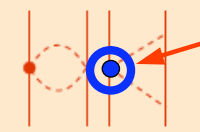
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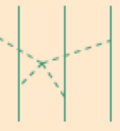
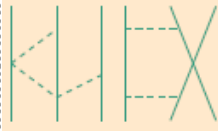
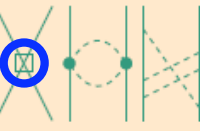
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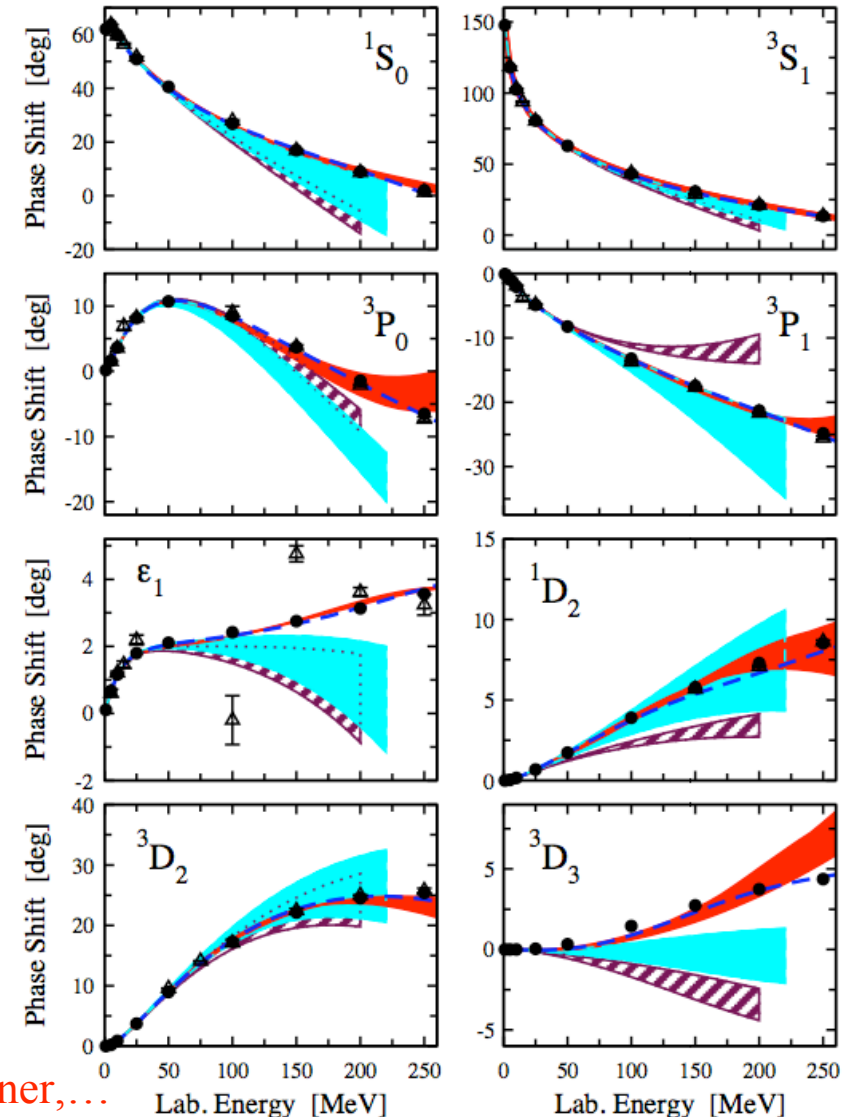
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accurate reproduction of
low-energy NN scattering at N³LO



Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Meissner,...

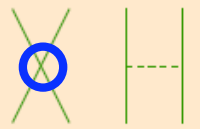
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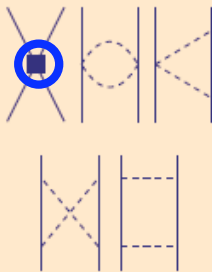
NN

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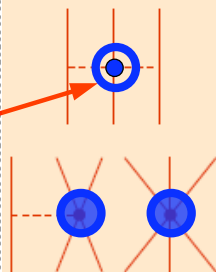
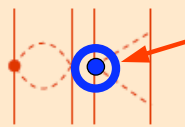
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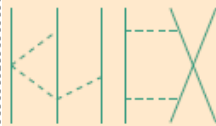
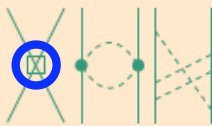
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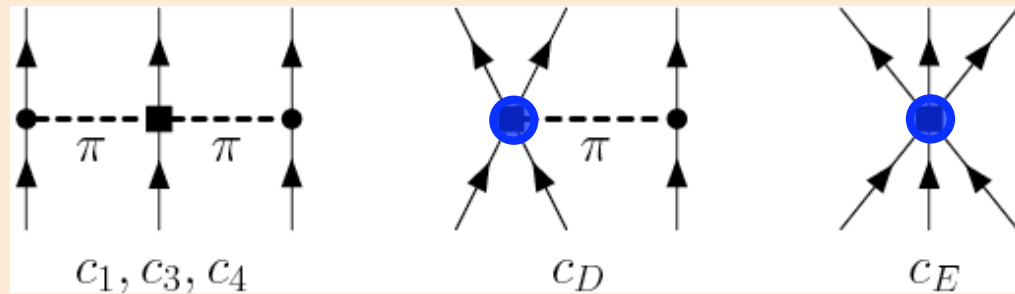


+ ...

+ ...

consistent NN-3N interactions

3N,4N: only 2 new couplings to N³LO



c_i from πN and NN Meissner et al. (2007)

$$c_1 = -0.9^{+0.2}_{-0.5}, \quad c_3 = -4.7^{+1.2}_{-1.0}, \quad c_4 = 3.5^{+0.5}_{-0.2}$$

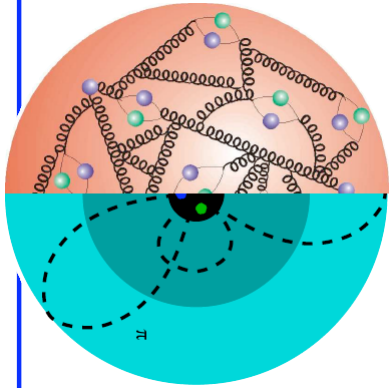
single- Δ = particular c_i values

c_D, c_E fit to ${}^3\text{H}$ binding energy and ${}^4\text{He}$ radius (or ${}^3\text{H}$ beta decay half-life)

Nuclear forces and the Renormalization Group (RG)

RG evolution to lower resolution/cutoffs

$$H(\Lambda) = T + V_{\text{NN}}(\Lambda) + V_{\text{3N}}(\Lambda) + V_{\text{4N}}(\Lambda) + \dots$$



Λ_{chiral}



Nuclear forces and the Renormalization Group (RG)

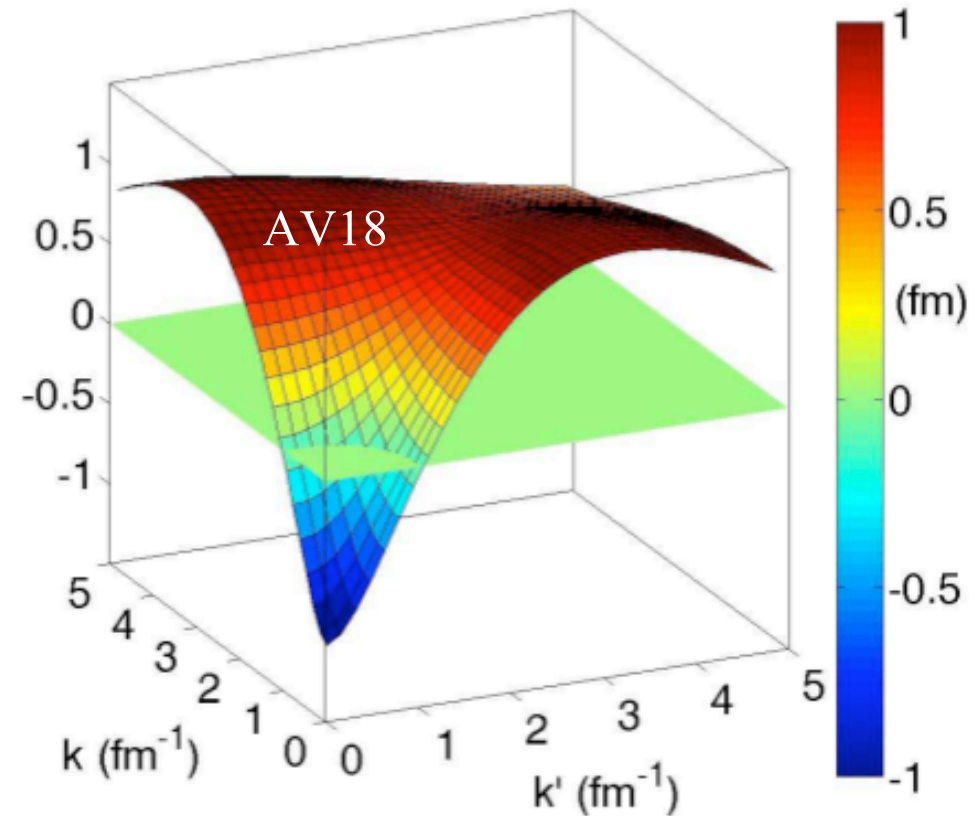
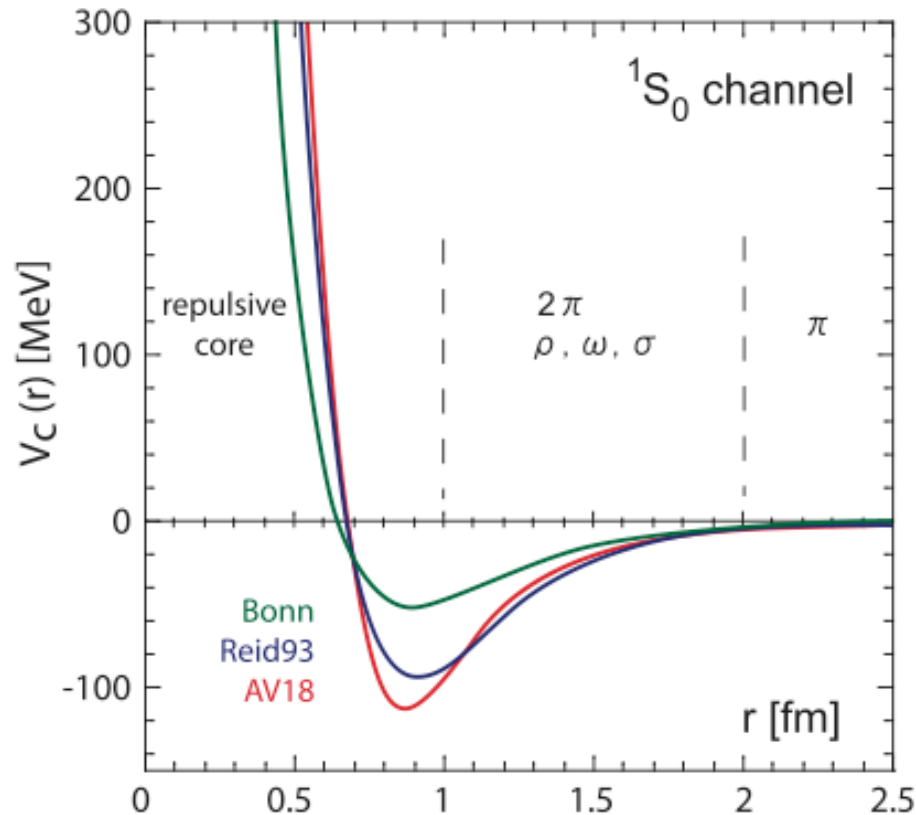
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for NN interactions
(preserves NN observables)

$$\frac{d}{d\Lambda} V_{\text{low } k}(k', k) = \frac{2}{\pi} \frac{V_{\text{low } k}(k', \Lambda) T_{\text{low } k}(\Lambda, k; \Lambda^2)}{1 - (k/\Lambda)^2}$$

Bogner, Kuo, AS, Furnstahl,...



red = short-range repulsion

Nuclear forces and the Renormalization Group (RG)

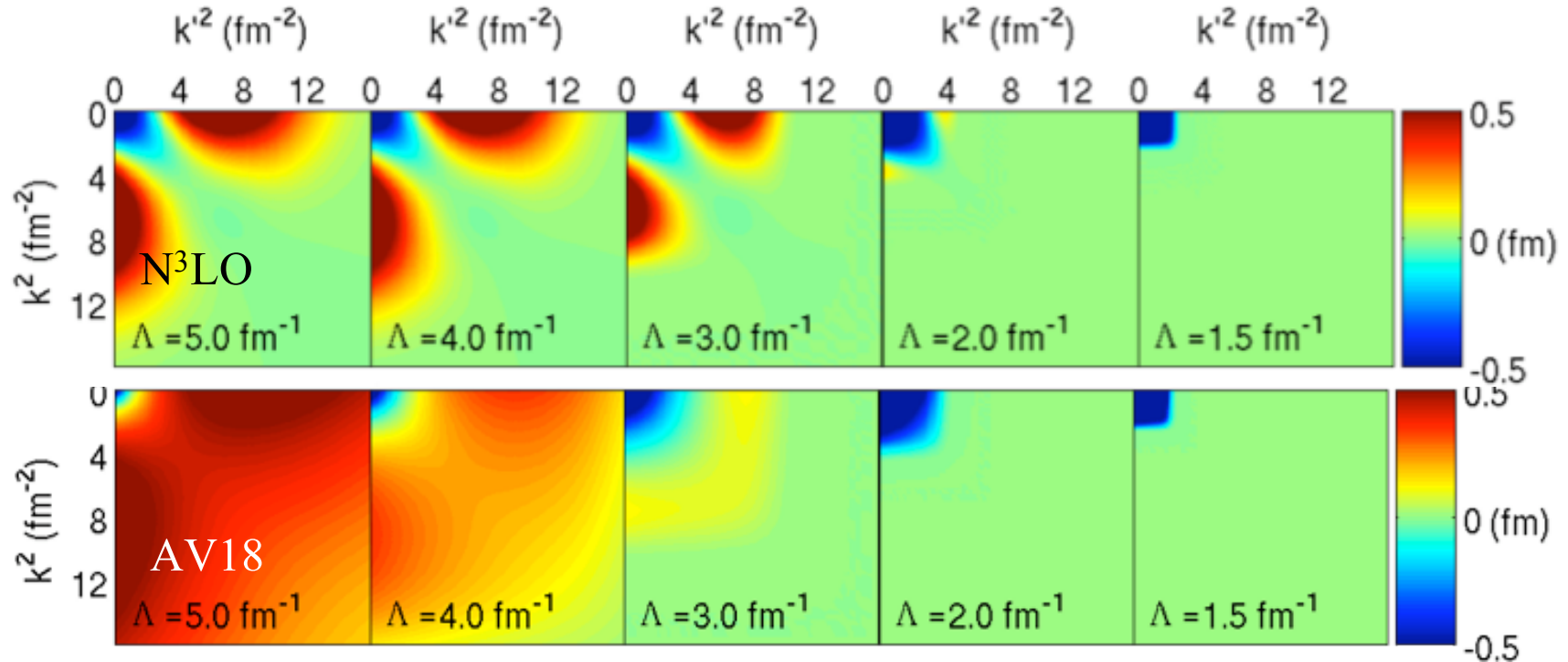
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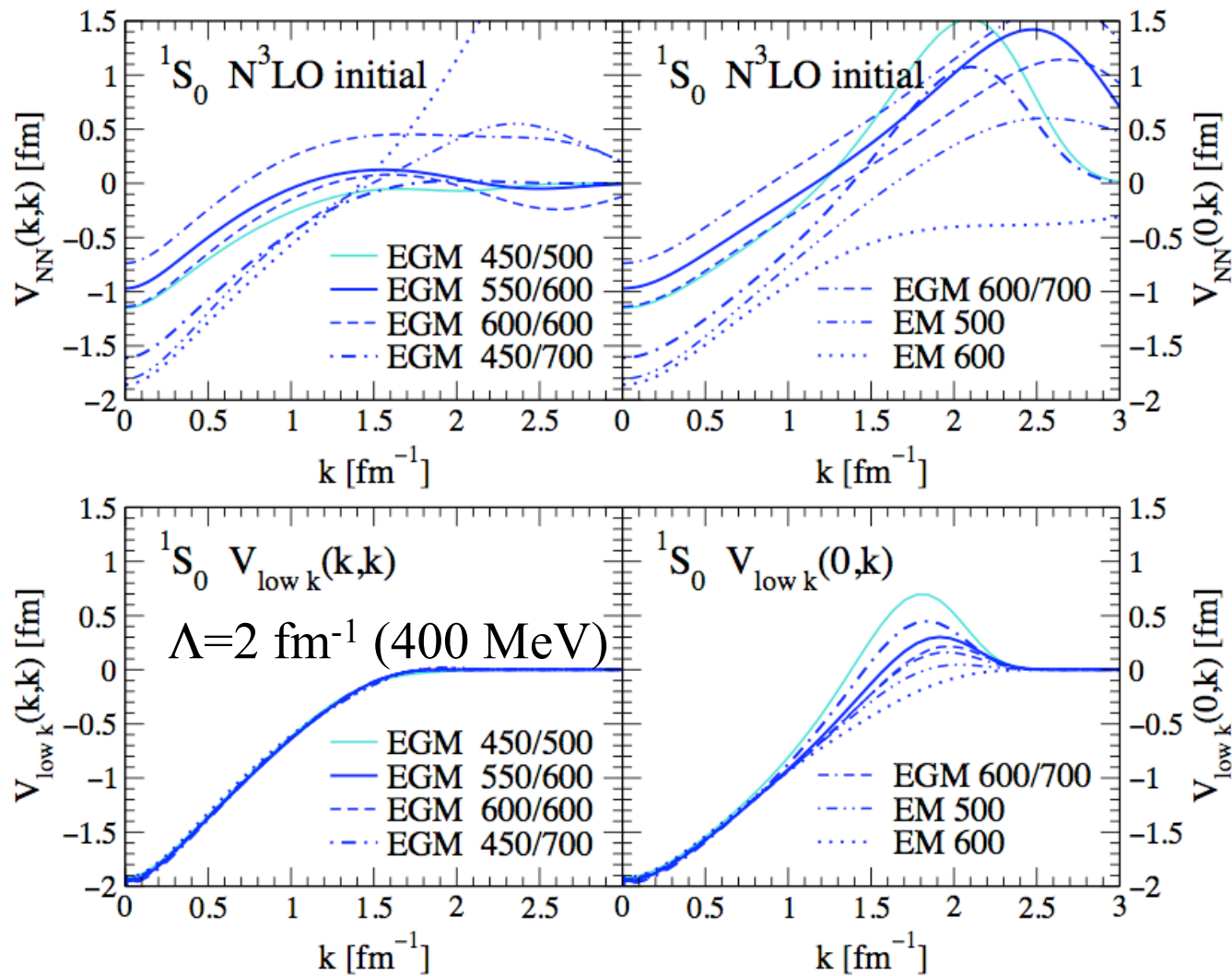
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Bogner, Kuo, AS, Furnstahl,...



low-momentum interactions $V_{\text{low } k}(\Lambda)$ with sharp or smooth regulators
decouples low-momentum physics from high momenta

Low-momentum universality



\approx **universality** from different chiral N^3LO potentials

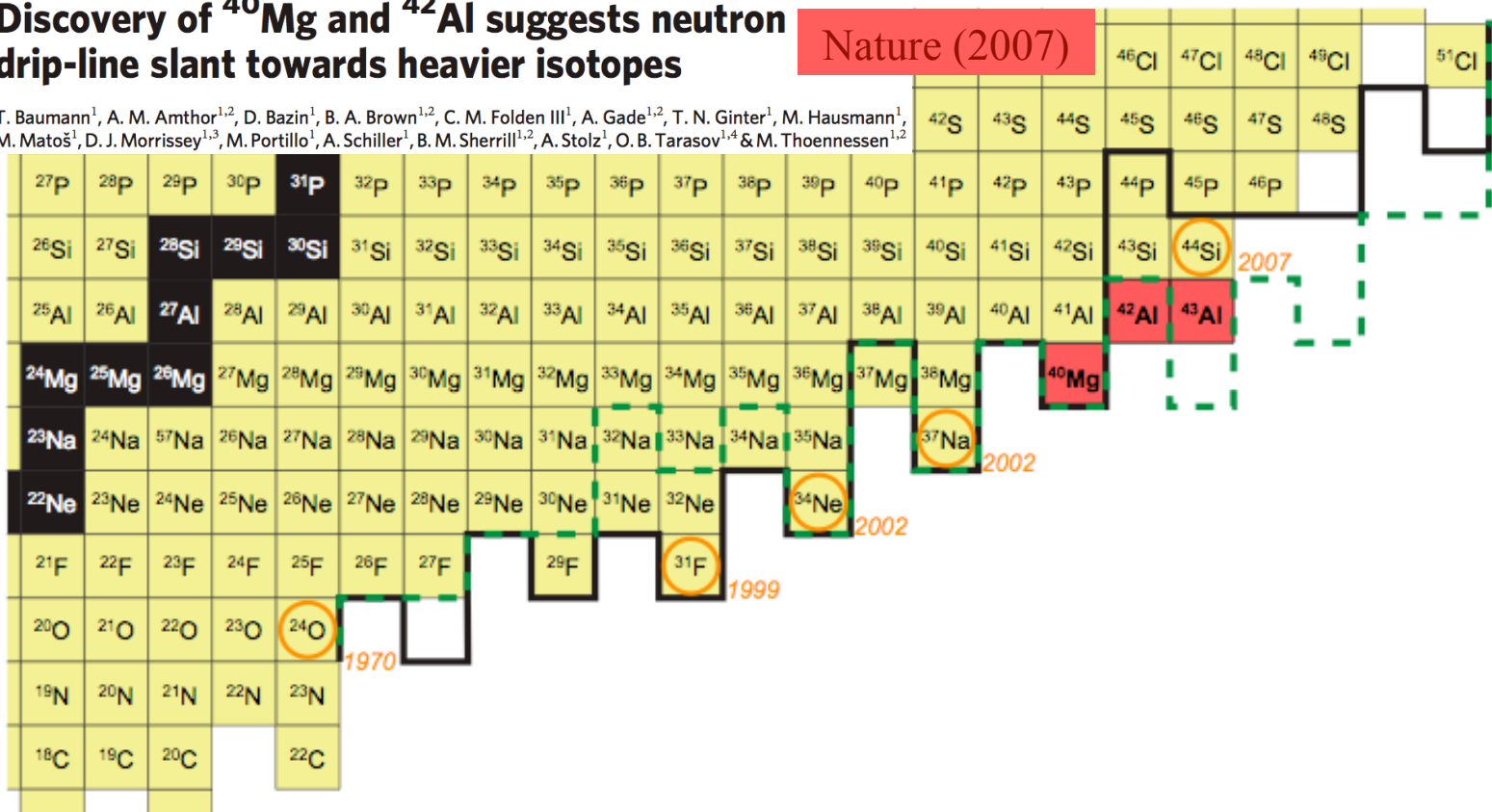
RG preserves NN observables and long-range parts
decouples low-momentum physics from high momenta

Towards the limits of existence - the neutron drip-line

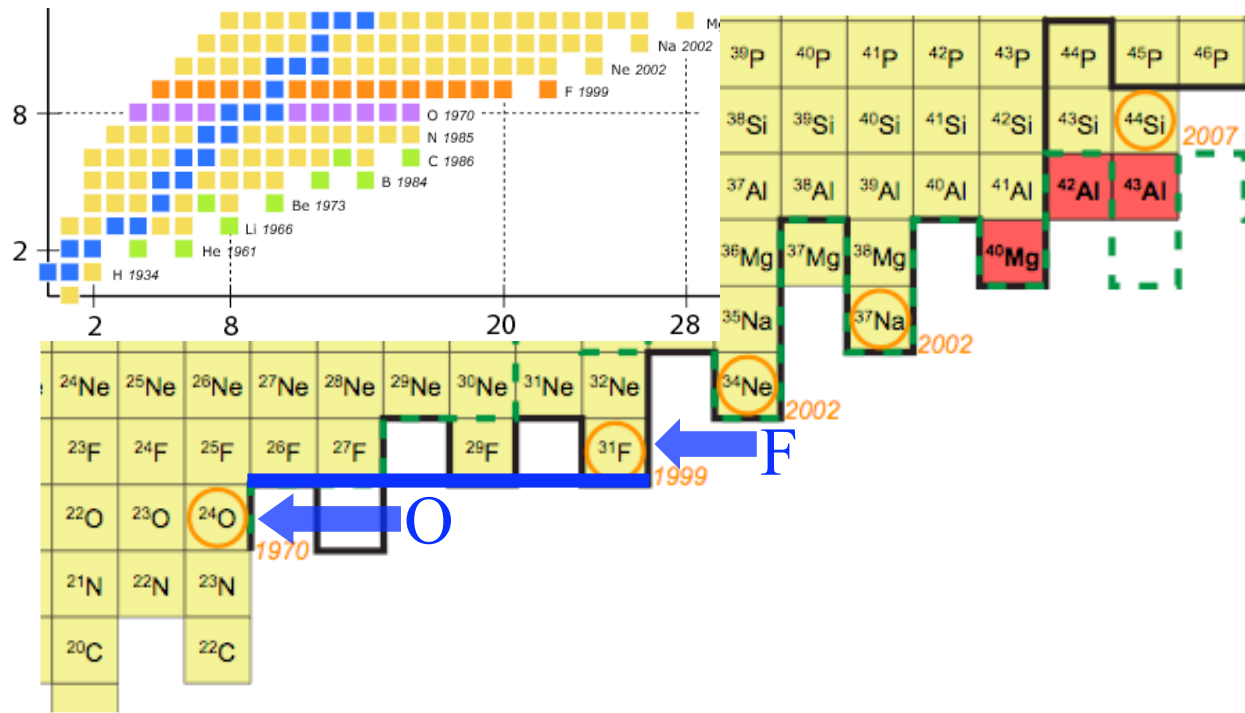
Discovery of ^{40}Mg and ^{42}Al suggests neutron drip-line slant towards heavier isotopes

Nature (2007)

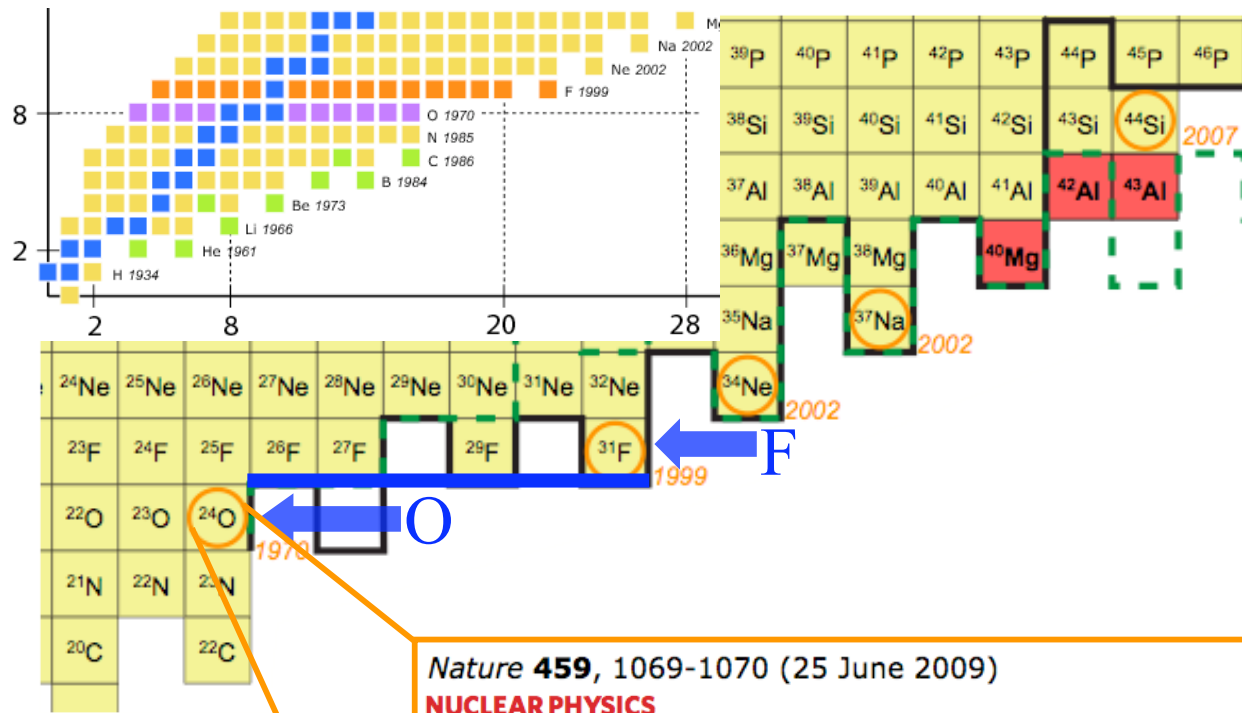
T. Baumann¹, A. M. Amthor^{1,2}, D. Bazin¹, B. A. Brown^{1,2}, C. M. Folden III¹, A. Gade^{1,2}, T. N. Ginter¹, M. Hausmann¹, M. Matoš¹, D. J. Morrissey^{1,3}, M. Portillo¹, A. Schiller¹, B. M. Sherrill^{1,2}, A. Stolz¹, O. B. Tarasov^{1,4} & M. Thoennessen^{1,2}



The oxygen anomaly



The oxygen anomaly



Nature **459**, 1069-1070 (25 June 2009)

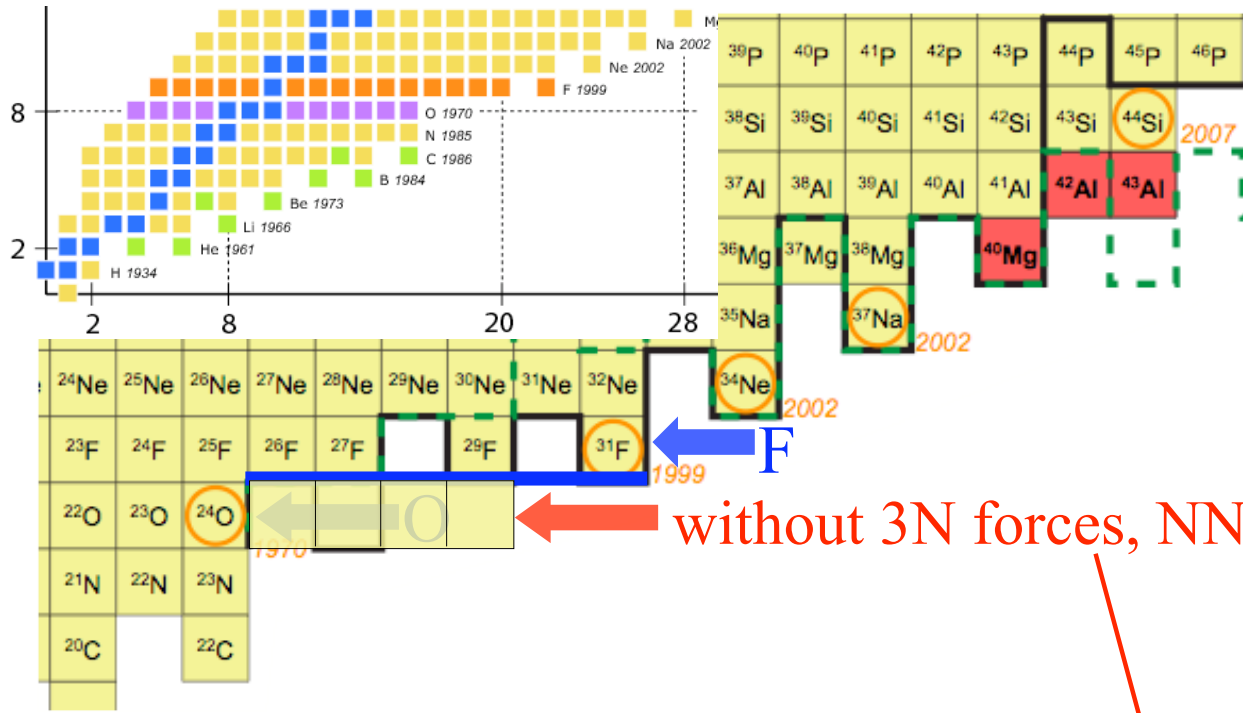
NUCLEAR PHYSICS

Unexpected doubly magic nucleus

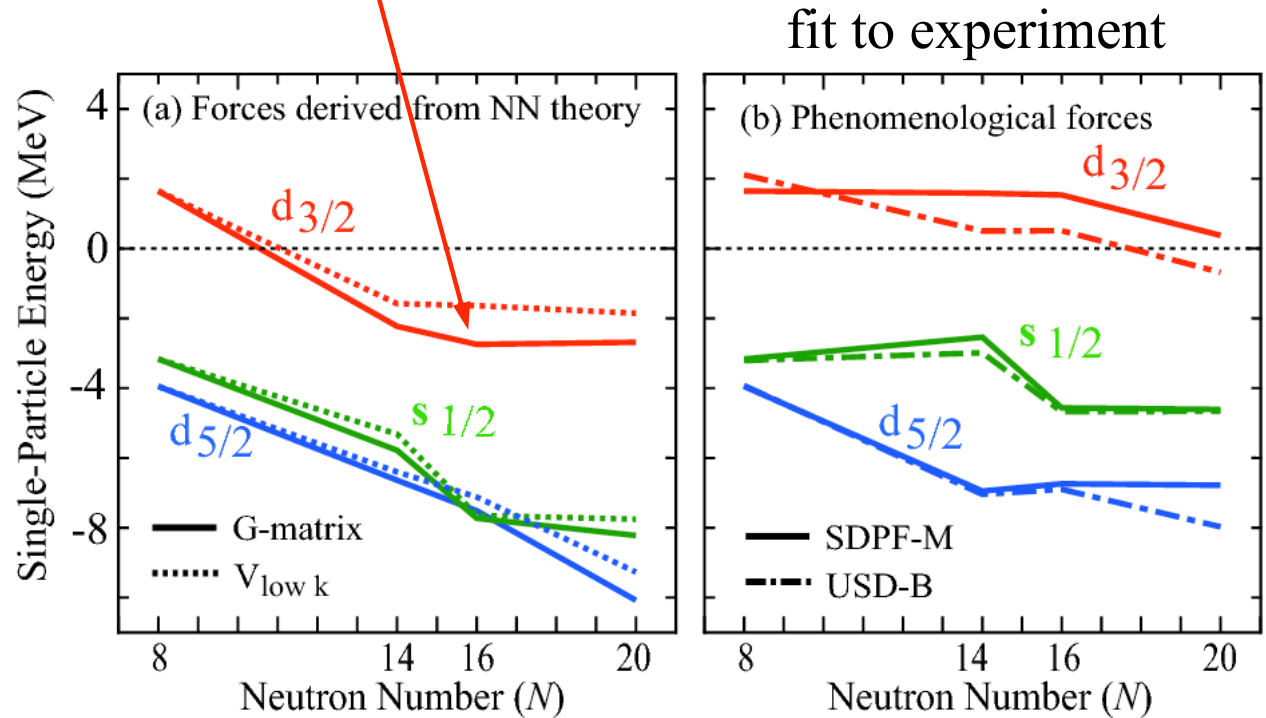
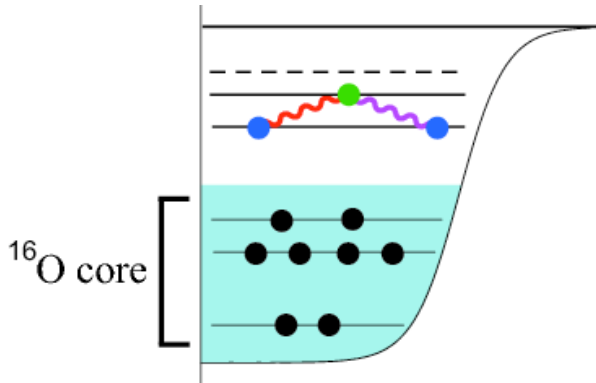
Robert V. F. Janssens

Nuclei with a 'magic' number of both protons and neutrons, dubbed doubly magic, are particularly stable. The oxygen isotope ^{24}O has been found to be one such nucleus — yet it lies just at the limit of stability.

The oxygen anomaly - not reproduced without 3N forces



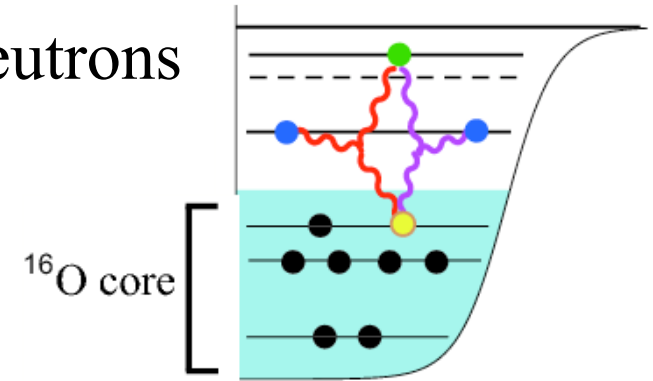
many-body theory based
on two-nucleon forces:
drip-line incorrect at ^{28}O



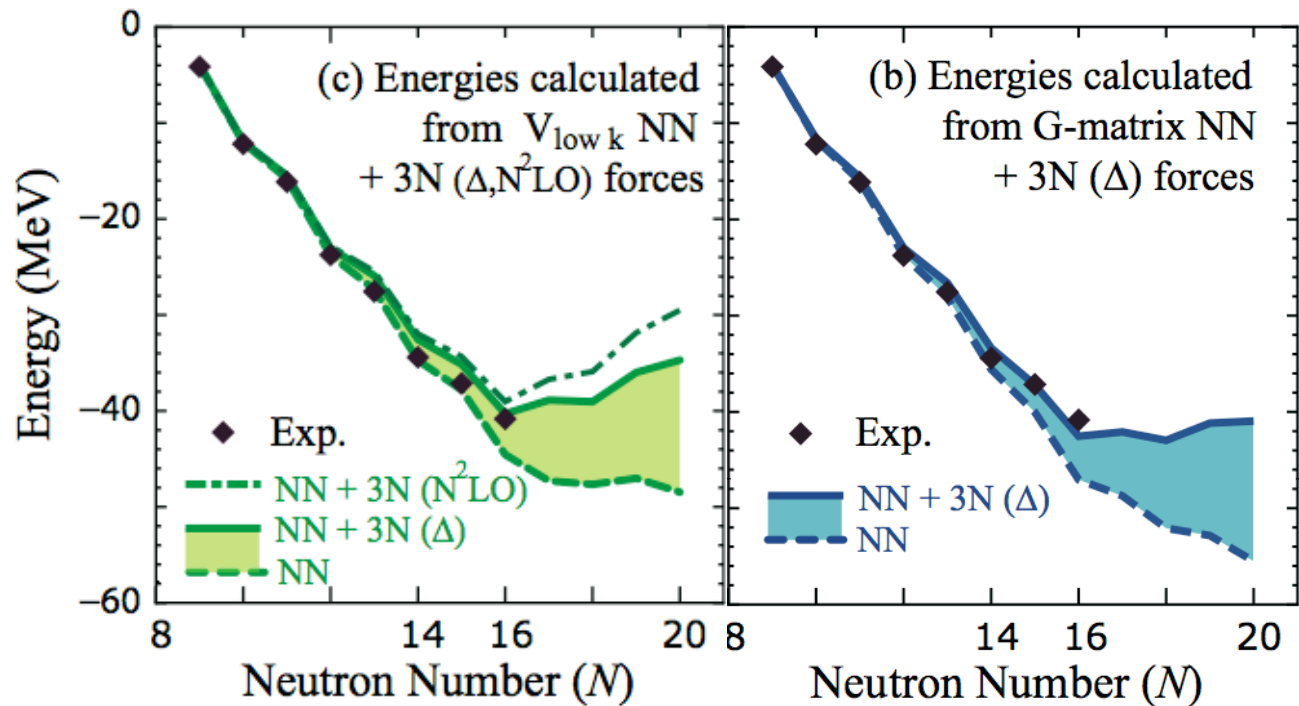
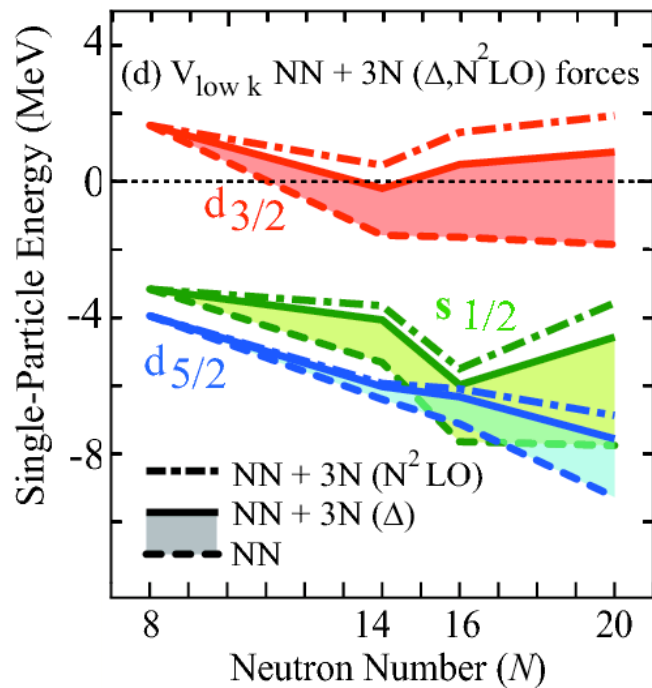
The oxygen anomaly - impact of 3N forces

include “normal-ordered” 2-body part of 3N forces (enhanced by core A)

leads to repulsive interactions between valence neutrons
(can understand partly based on Pauli principle)



$d_{3/2}$ orbital remains unbound from ^{16}O to ^{28}O



first microscopic explanation of the oxygen anomaly

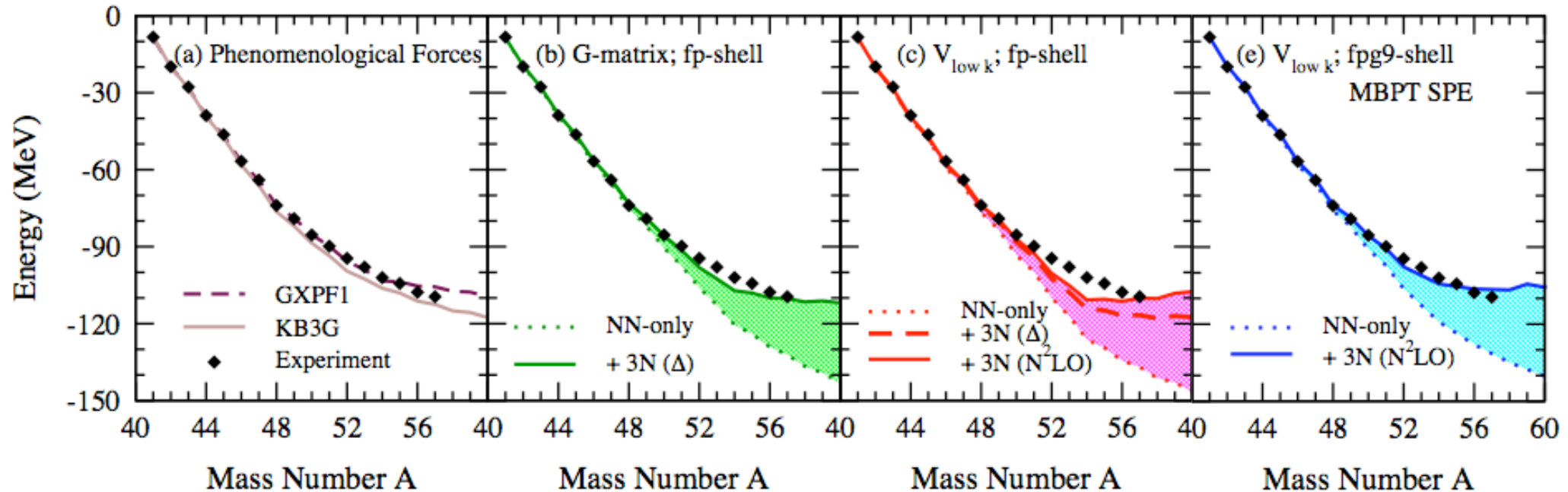
Otsuka, Suzuki, Holt, AS, Akaishi (2010)

Evolution to neutron-rich calcium isotopes

repulsive 3N contributions also key for calcium ground-state energies

Holt, Otsuka, AS, Suzuki, in prep.

fit to experiment



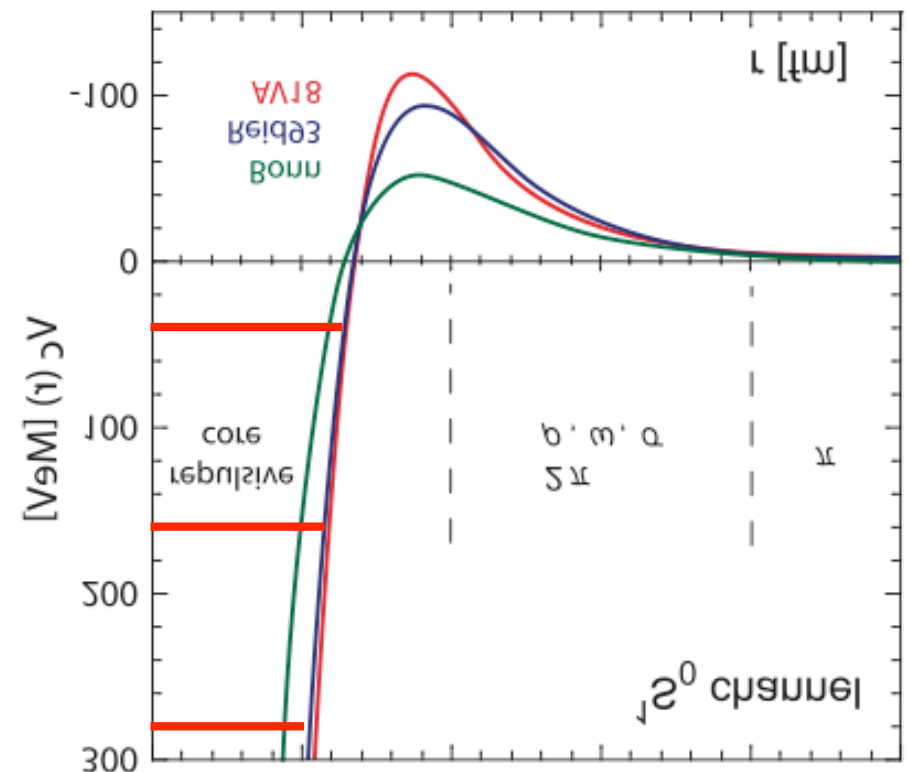
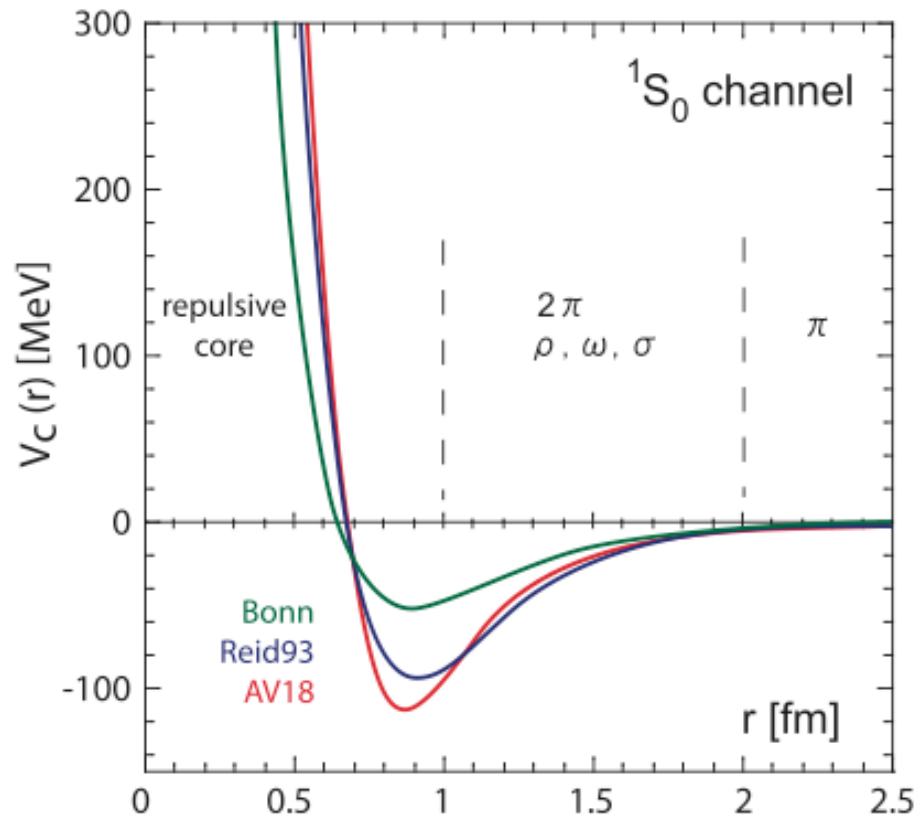
3N mechanism important for shell structure (e.g., $N=28$ and ^{48}Ca)

changes due to 3N forces amplified and testable in neutron-rich nuclei

impact on energy-density-functional based ground-state properties?

Convergence with low-momentum interactions

large cutoffs lead to **flipped-potential bound states**, even for small $-\lambda V$
requires nonperturbative expansion, leads to slow convergence for nuclei

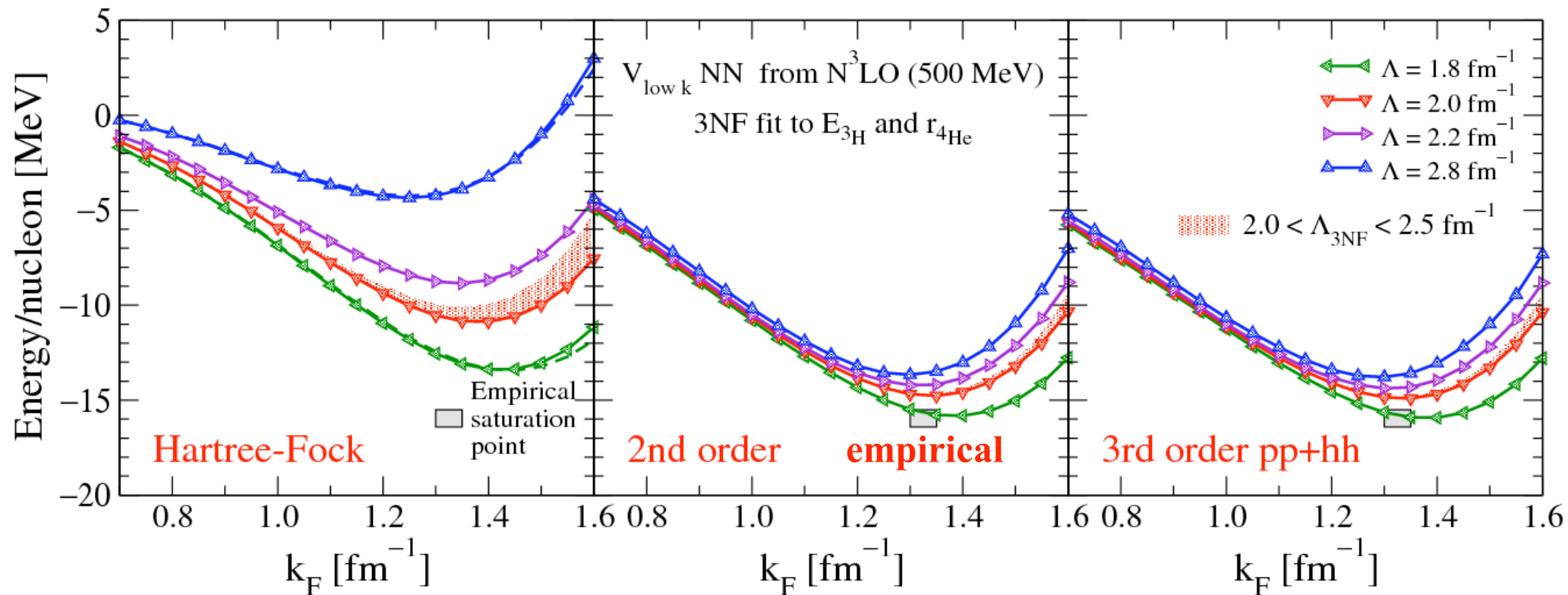


Weinberg eigenvalues show: two-body scattering becomes perturbative after RG evolution, except in channels with bound states

RG leads to improved convergence for nuclei and nuclear matter

Advances in nuclear matter theory

Is nuclear matter perturbative with chiral EFT and RG evolution?



Bogner, Furnstahl, Hebeler, Nogga, AS, in prep. and (2009)

exciting: empirical saturation with theoretical uncertainties

improved 3N improvements

input to develop a universal energy density functional for all nuclei

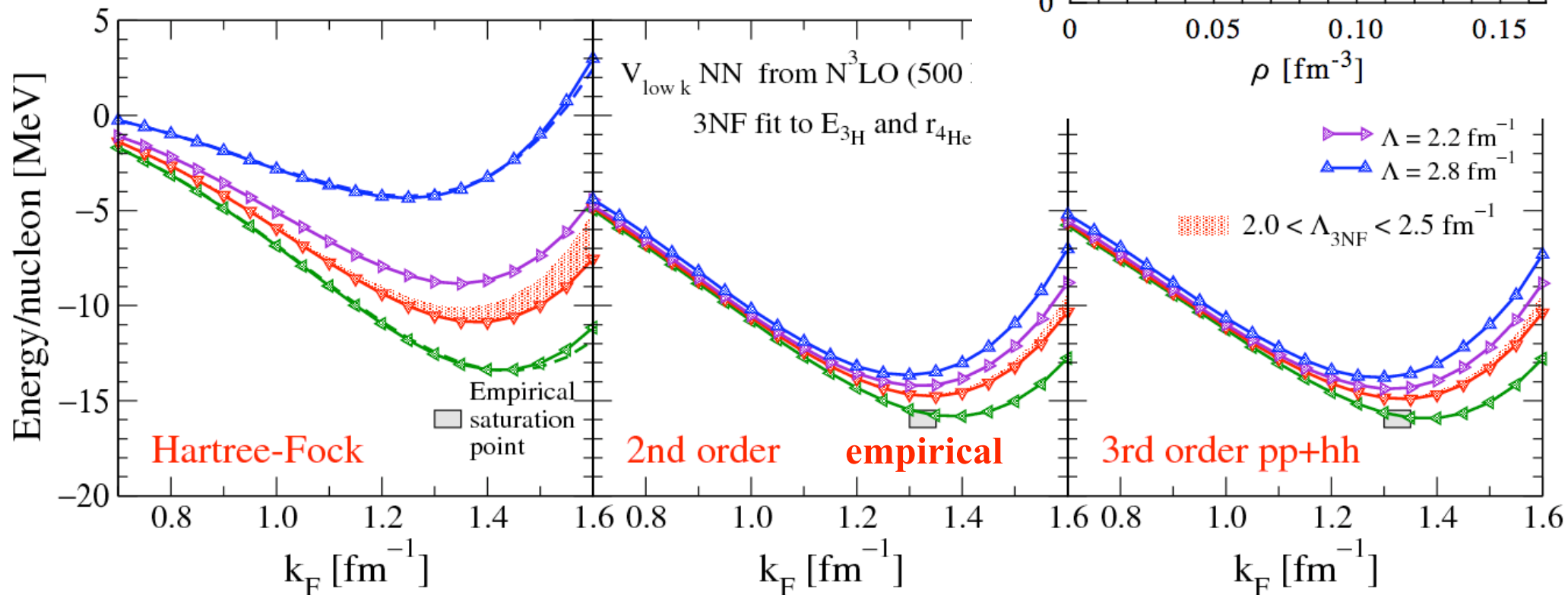
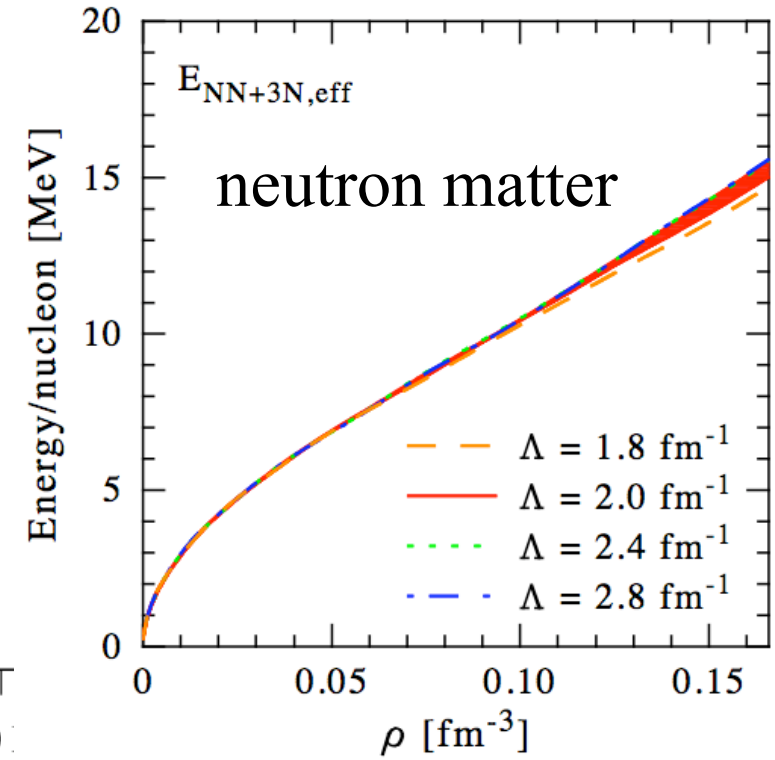
UNEDF SciDAC Collaboration

Universal Nuclear Energy Density Functional

Impact of 3N forces on neutron matter

Hebeler, AS (2009); Tolos, Friman, AS (2007)

only long-range parts of 3N forces
contribute to neutron matter (c_1 and c_3)



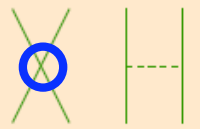
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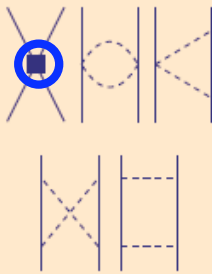
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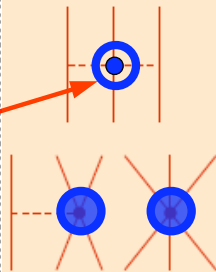
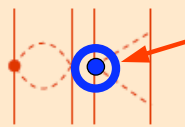
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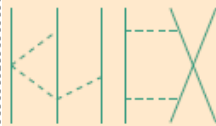
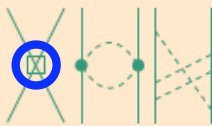
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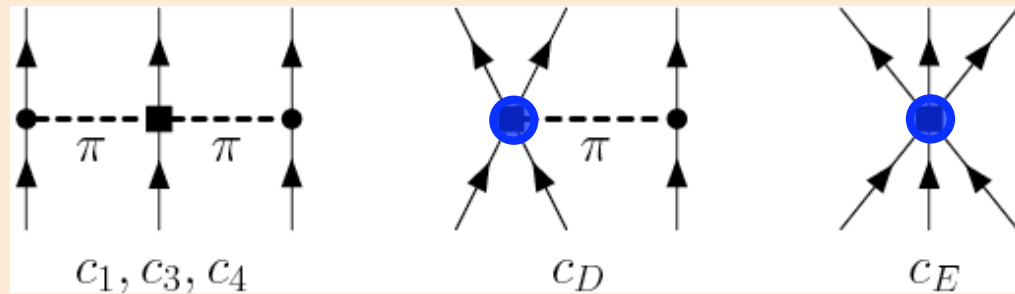


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c_1, c_3, c_4

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c_E

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single- Δ = particular c_i values

c_D, c_E fit to ${}^3\text{H}$ binding energy and ${}^4\text{He}$ radius (or ${}^3\text{H}$ beta decay half-life)

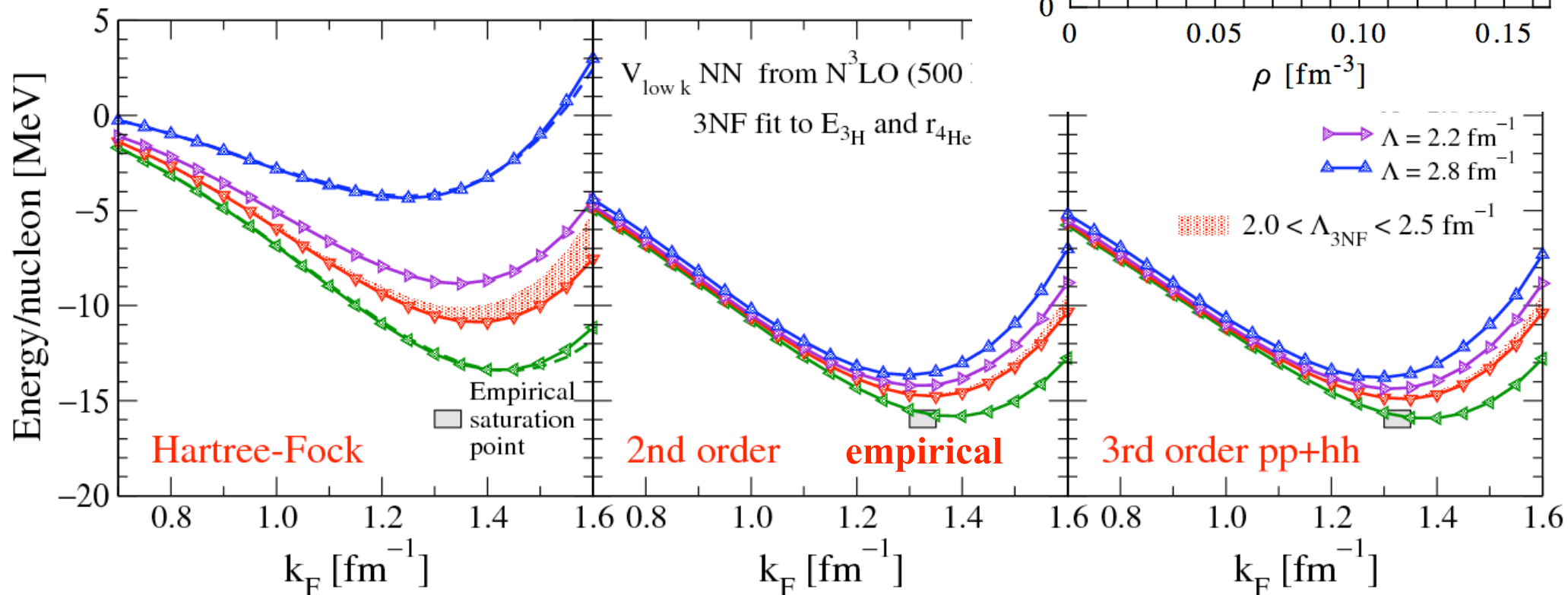
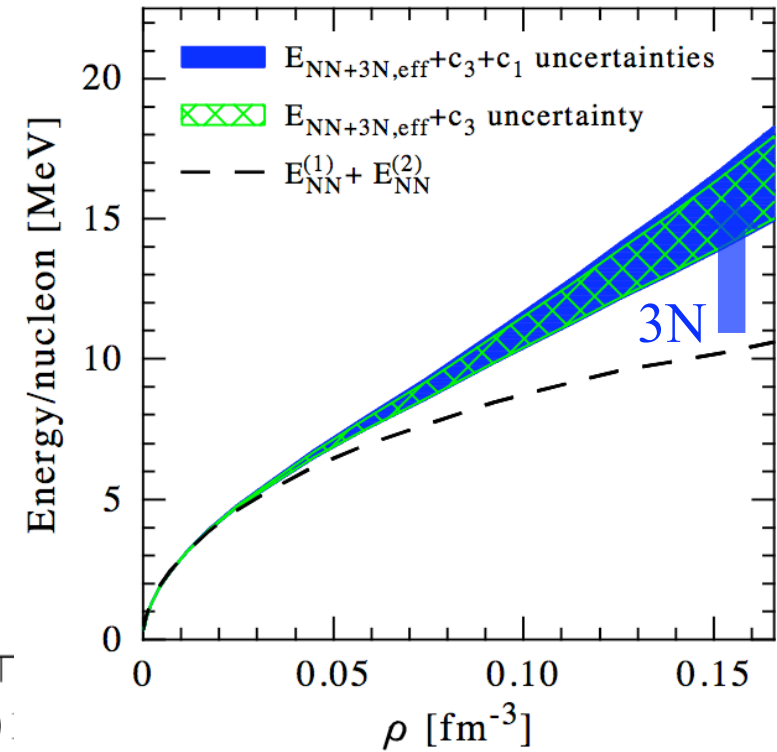
Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Meissner,...

Impact of 3N forces on neutron matter

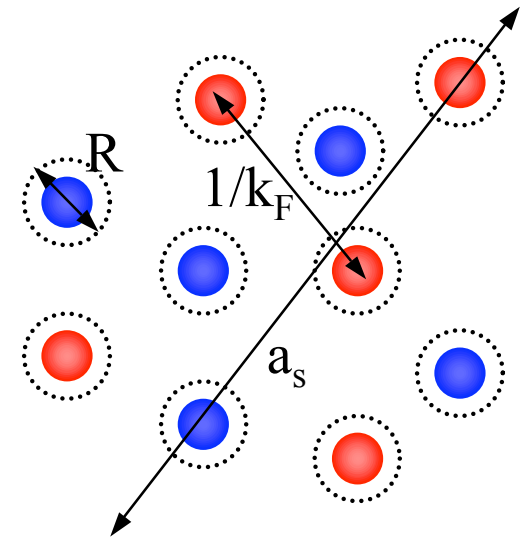
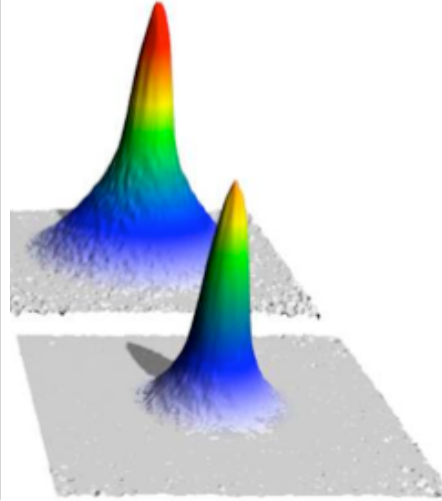
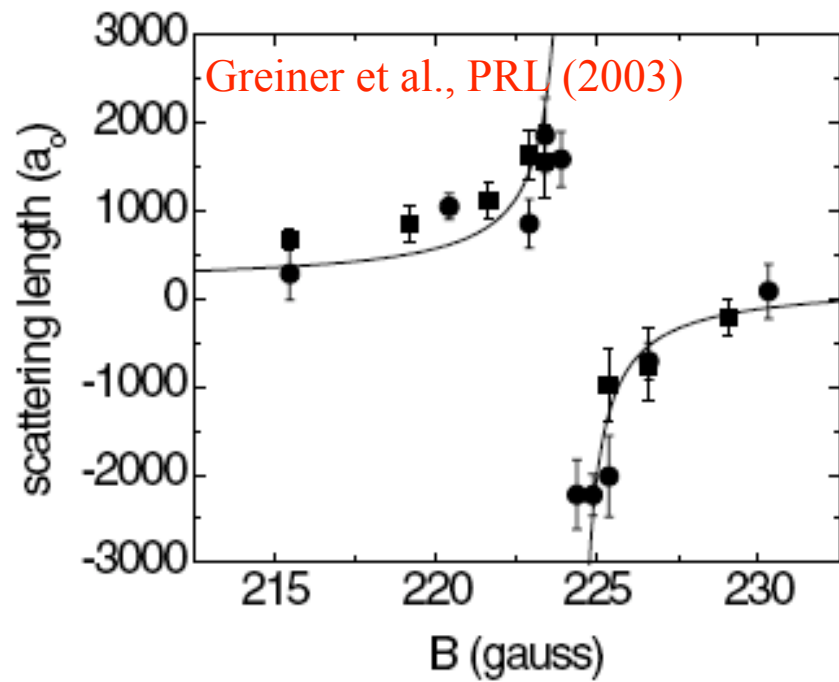
Hebeler, AS (2009); Tolos, Friman, AS (2007)

only long-range parts of 3N forces
contribute to neutron matter (c_1 and c_3)

uncertainties dominated by c_3 coupling



Large scattering lengths: Universal properties at low densities

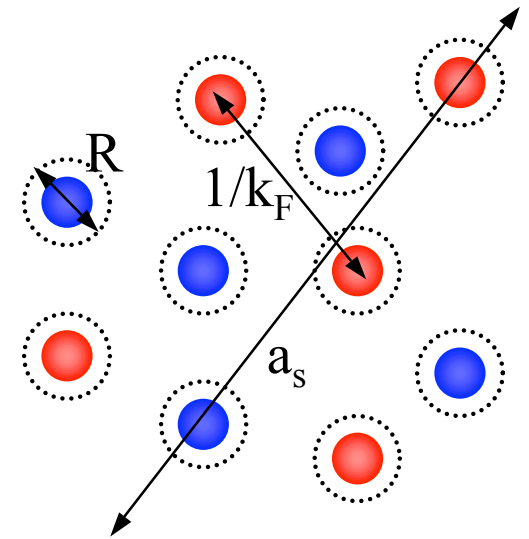
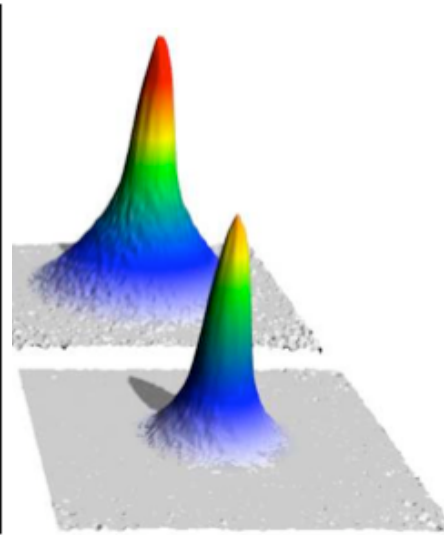
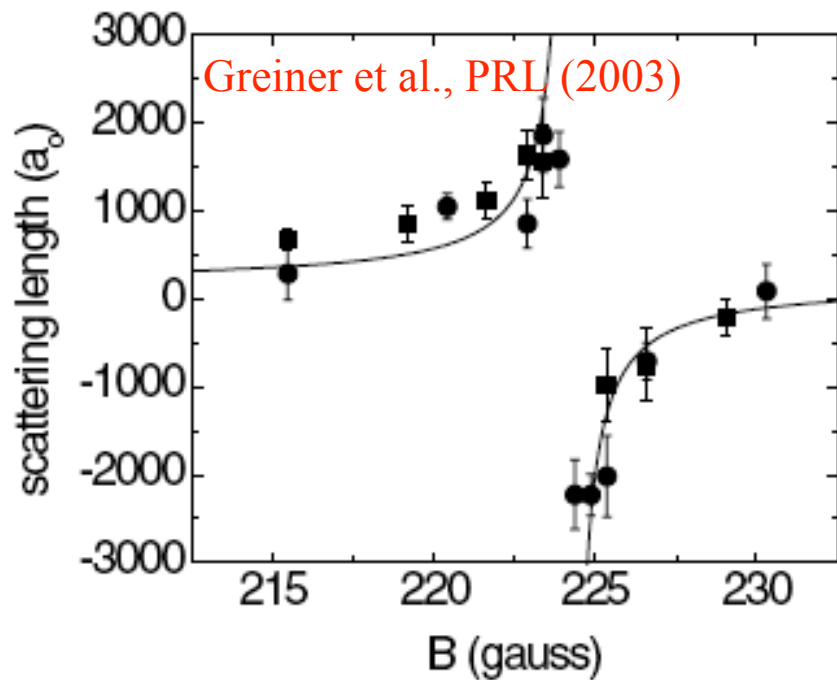


$$a_{nn} = -18.5 \pm 0.3 \text{ fm}$$

large for neutrons

strong interactions via Feshbach resonances

Large scattering lengths: Universal properties at low densities



$$a_{nn} = -18.5 \pm 0.3 \text{ fm}$$

strong interactions via Feshbach resonances

large for neutrons

dilute Fermi system with large scattering length has **universal properties**

$$0 \leftarrow 1/a_s \ll k_F \ll 1/r_e, 1/R, \dots \rightarrow \infty$$

strongly-interacting **dilute**

only Fermi momentum or density sets scale

physics is independent of interaction/system details:

from dilute neutron matter to resonant ${}^6\text{Li}$ or ${}^{40}\text{K}$ atoms in traps

Large scattering lengths: Universal thermodynamics

energy per particle $\frac{E}{N} = \xi \left(\frac{E}{N} \right)_{\text{free}} = \xi \frac{3k_F^2}{10m}$

with universal Bertsch parameter ξ

Quantum Monte Carlo: $\xi=0.40(1)$

Gezerlis, Carlson (2009),...

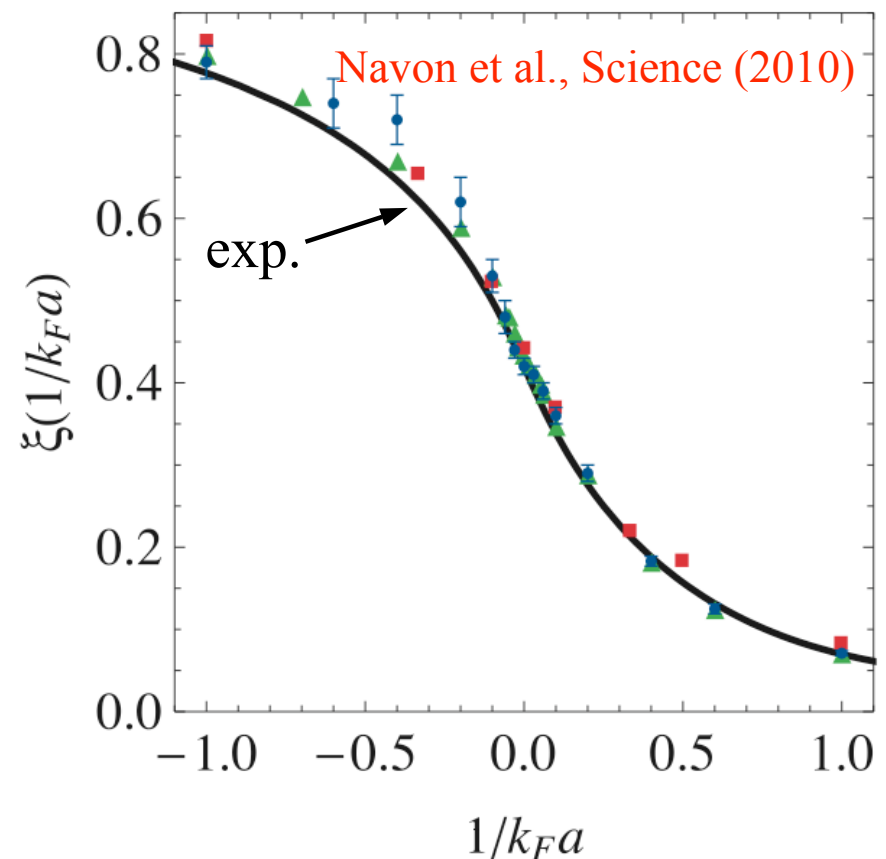
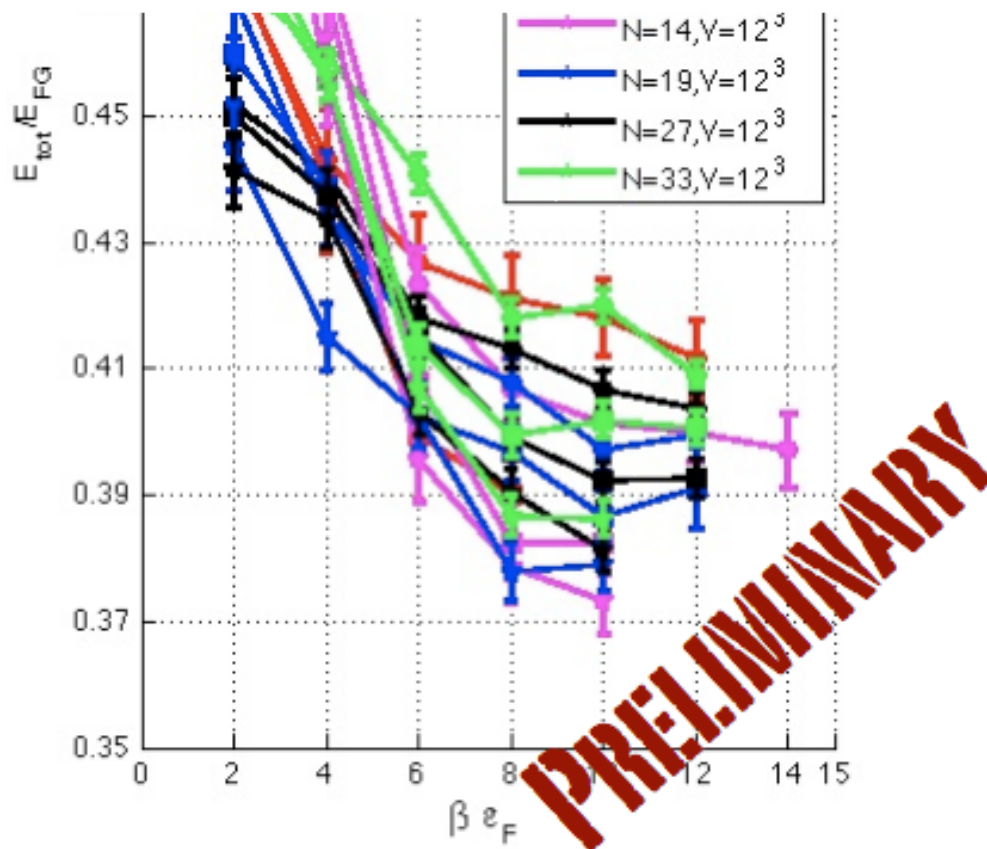
lattice: Auxiliary-Field Hybrid MC

Drut, Gezerlis, Laehde @ 2010 UNEDF meeting

most precise results with ^6Li

$\xi=0.39(2)$ and $0.41(2)$

cloud size and $E(S)$ Zuo, Thomas (2009)

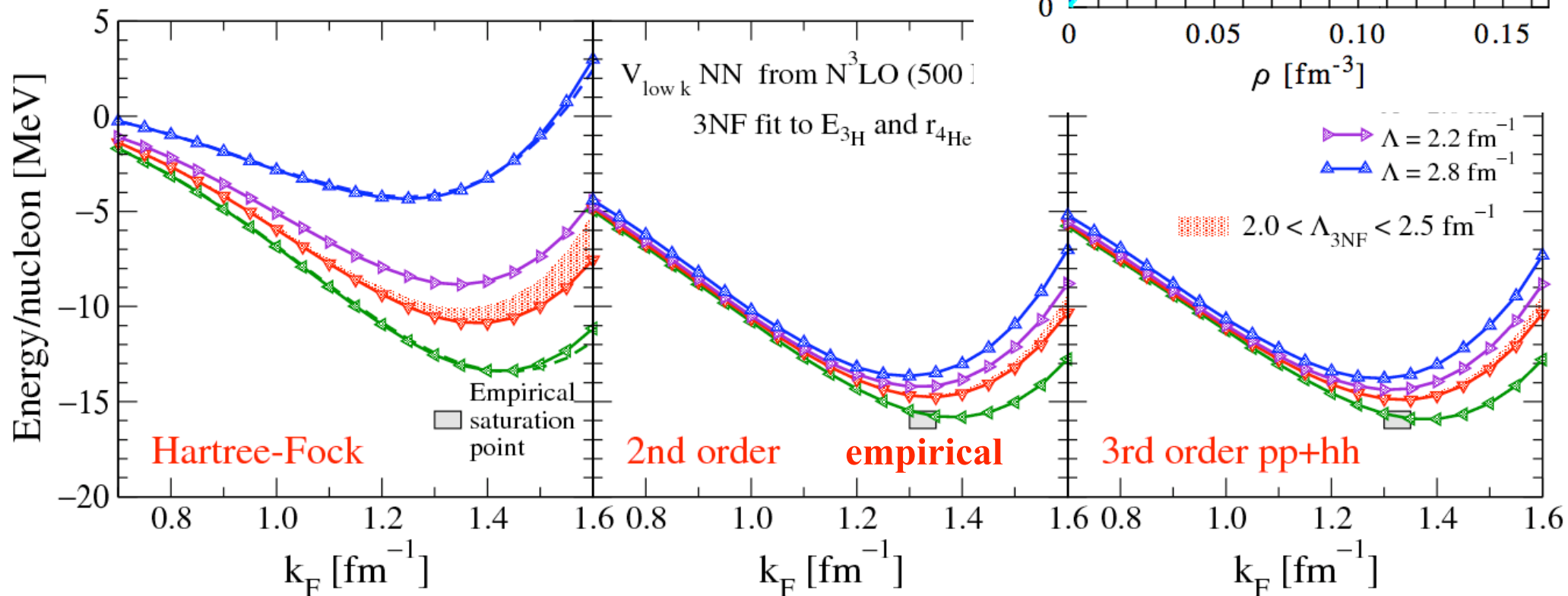
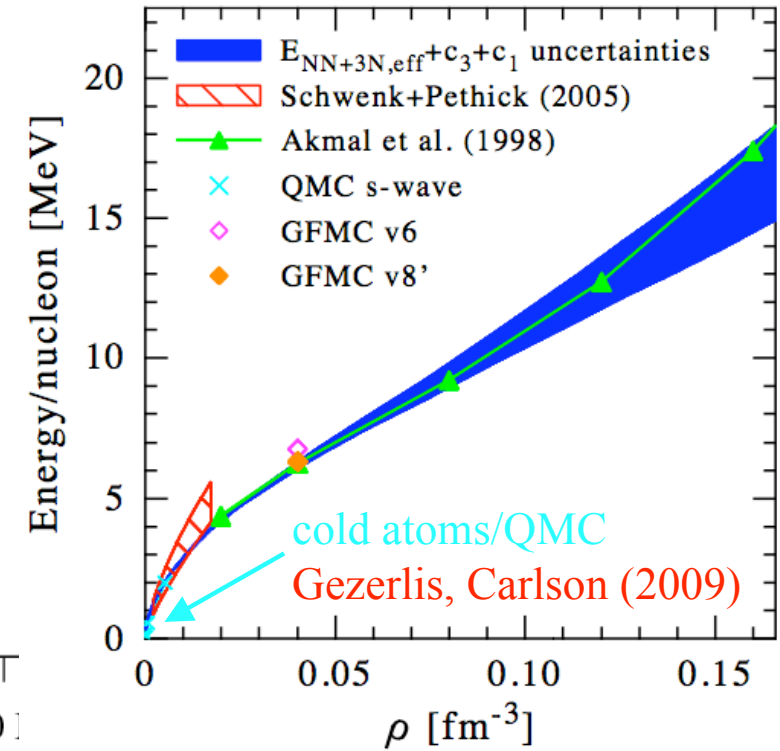


Impact of 3N forces on neutron matter

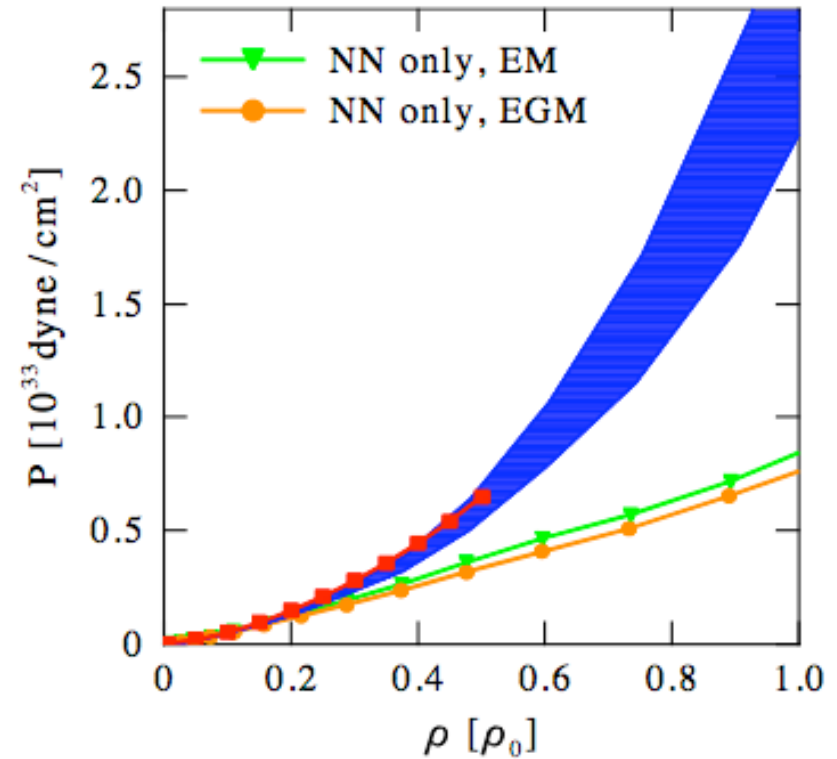
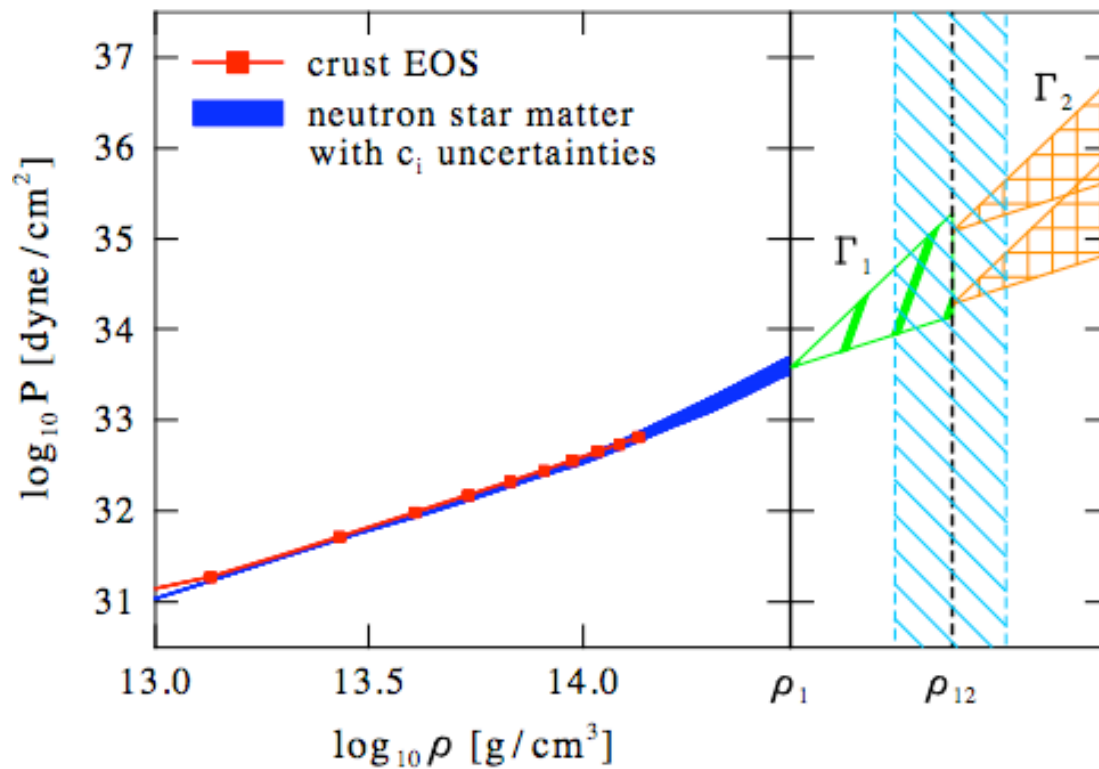
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Impact on neutron stars Hebeler, Lattimer, Pethick, AS (2010)

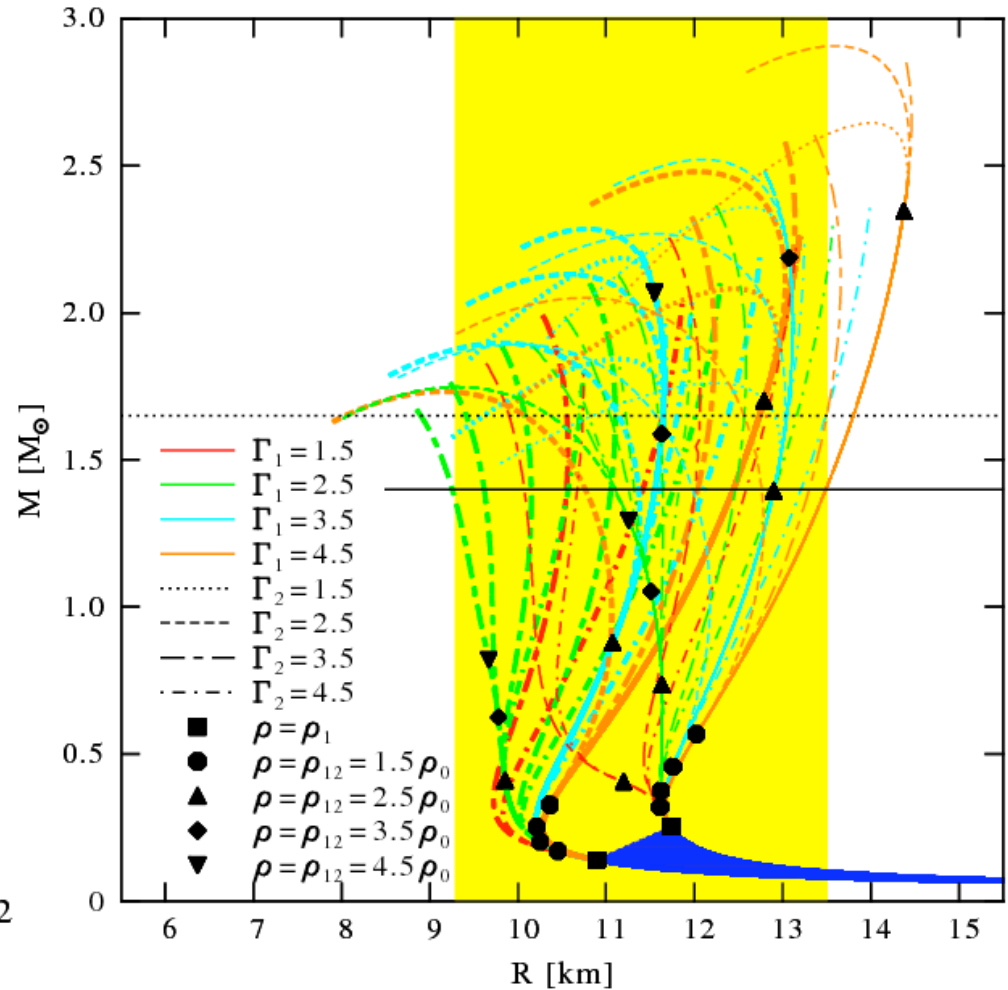
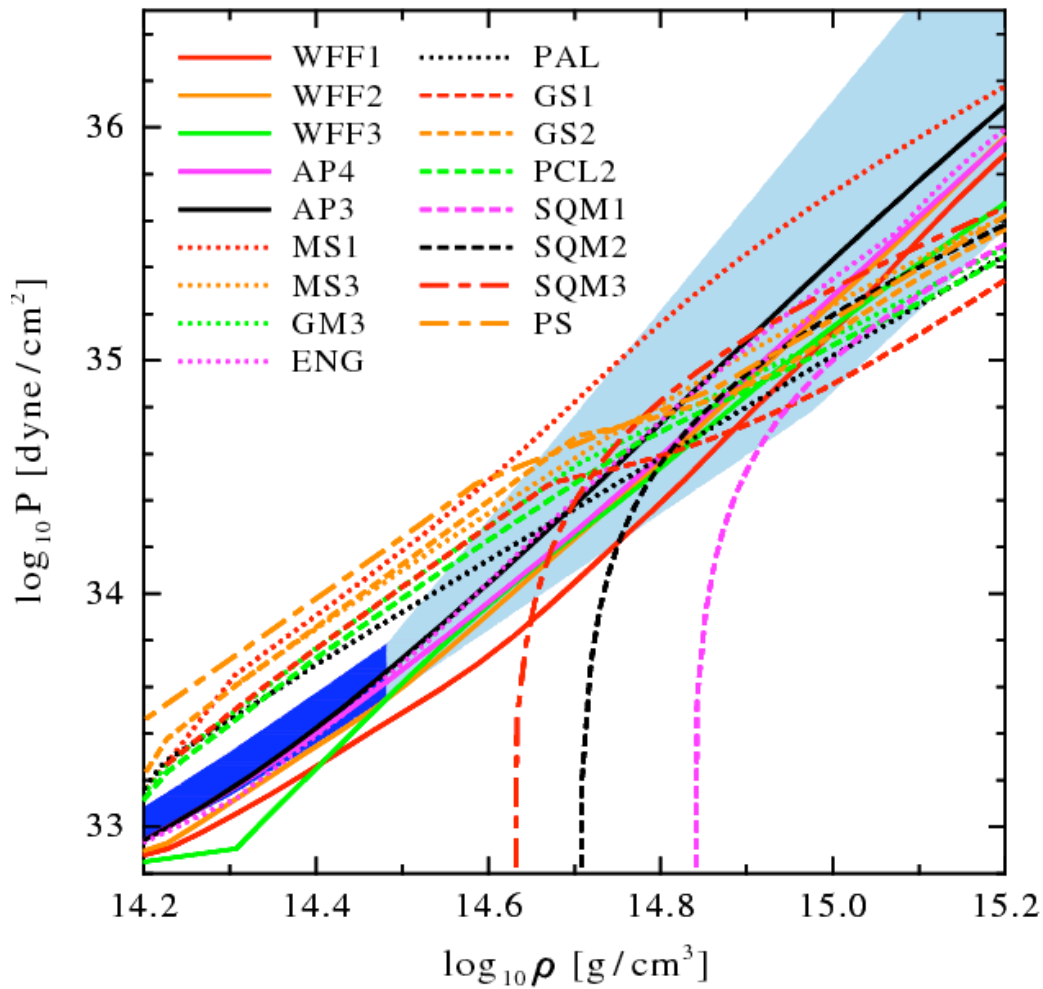


pressure below nuclear densities agrees with standard crust EOS
only after 3N forces are included

extend uncertainty band to higher densities using piecewise polytropes

constrain polytropes by causality and require to support $1.65 M_{\text{sun}}$ star

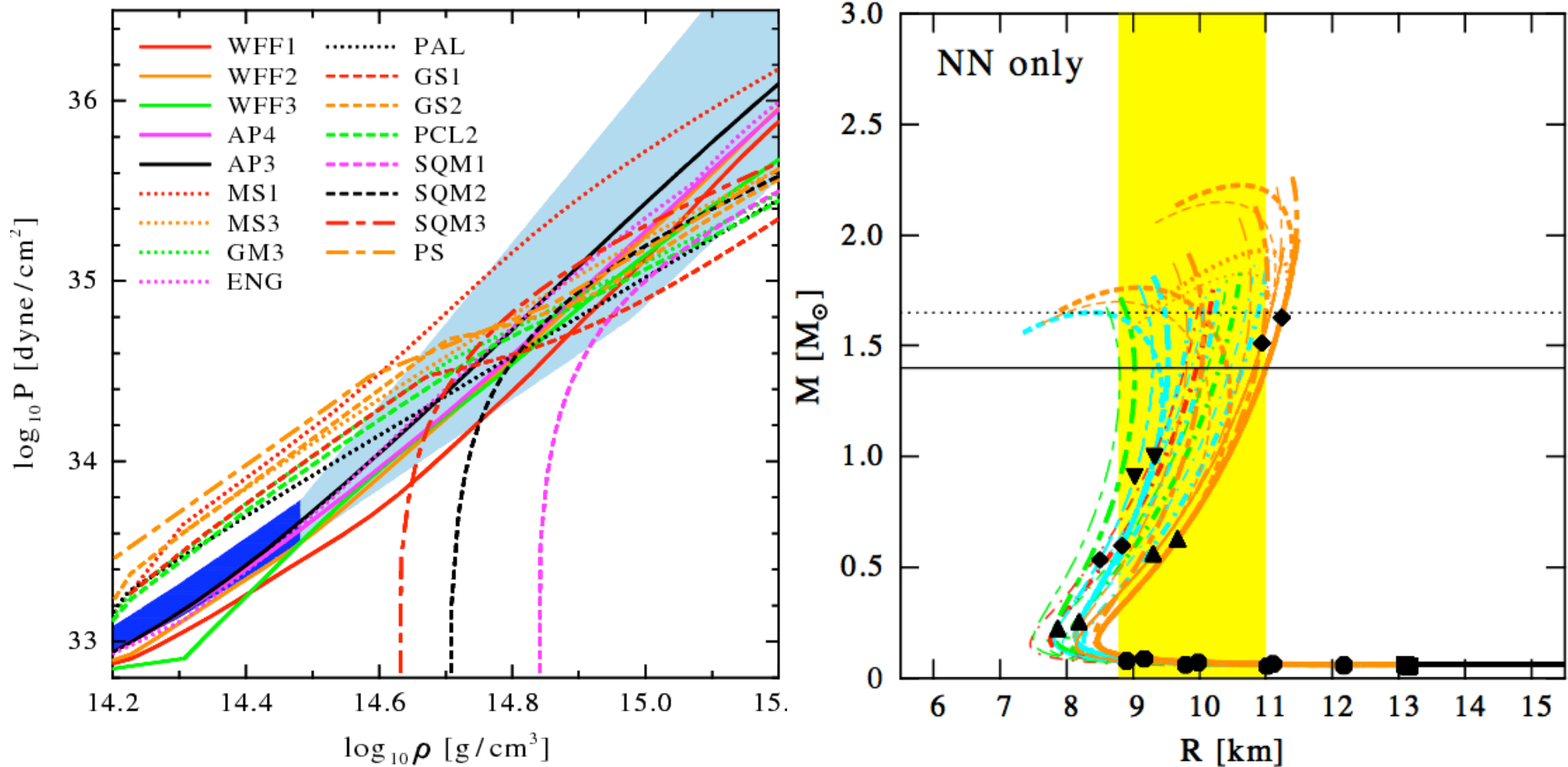
Impact on neutron stars Hebeler, Lattimer, Pethick, AS (2010)



low-density pressure sets scale, our results reduce spread at nuclear densities in current neutron star modeling from factor 6 to $\pm 25\%$

constrains neutron star radius to 11.8 ± 2.1 km for $M = 1.4 M_{\text{sun}}$

Impact on neutron stars Hebeler, Lattimer, Pethick, AS (2010)



low-density pressure sets scale, our results reduce spread at nuclear densities in current neutron star modeling from factor 6 to $\pm 25\%$

constrains neutron star radius to 11.8 ± 2.1 km for $M = 1.4 M_{\text{sun}}$
 9.9 ± 1.1 km without 3N forces

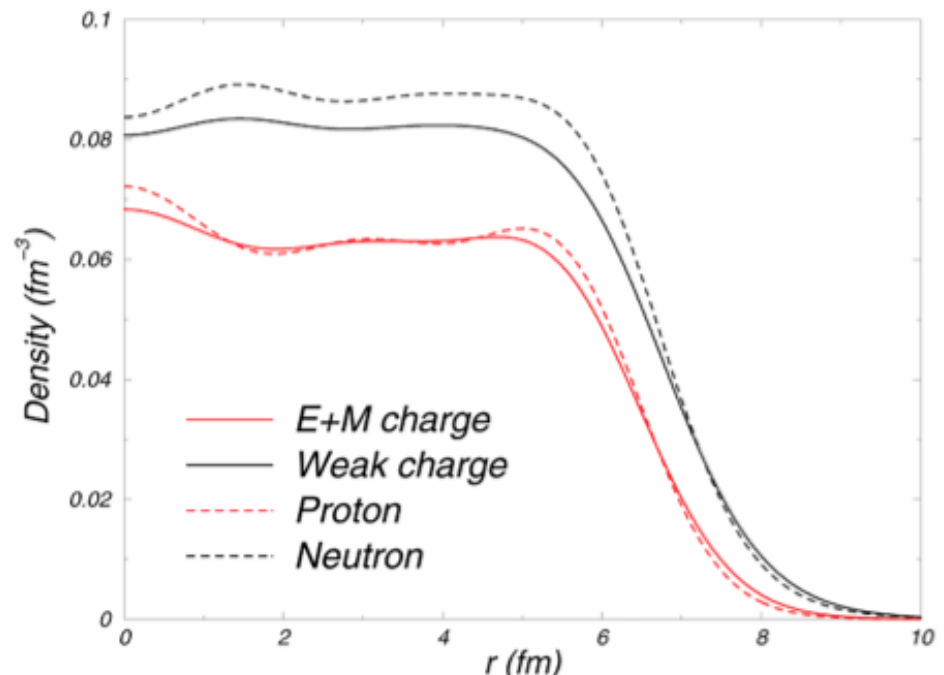
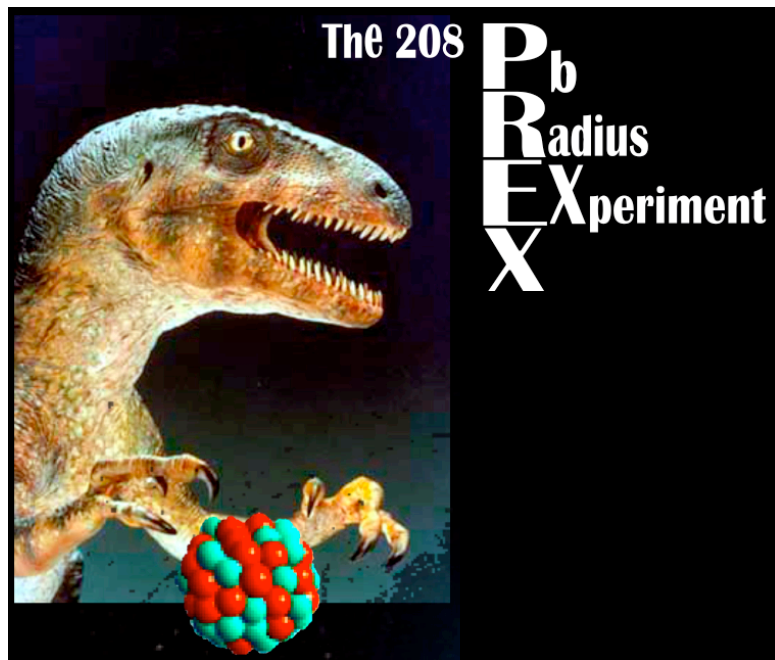
PREX - Neutron radius of lead P.A. Souder, G.M. Urciuoli, R. Michaels

parity-violating electron-scattering at JLAB, scheduled for spring 2010

Z-boson couples preferentially to neutrons

will measure neutron radius to 1% or +/- 0.05 fm

^{208}Pb



	up-quark	down-quark	proton	neutron
γ -coupling	+2/3	-1/3	+1	0
Z_0 -coupling	$\approx +1/3$	$\approx -2/3$	≈ 0	-1

$$g_v = 2t_z - 4Q \sin^2 \theta_W \approx 2t_z - Q$$

Neutron matter band predicts the neutron skin of lead to 0.17 ± 0.03 fm

Summary

Exciting era with advances on many fronts:
development of effective field theory and the renormalization group

enables a unified description from nuclei to matter in astrophysics

3N forces are a frontier:

in chiral EFT for nuclear forces (and in RG evolution)

key to explain why ^{24}O is the heaviest oxygen isotope

expect to be important for extrapolations towards neutron drip line

dominant uncertainty for pressure of neutron star matter below nuclear densities, important for radii and for comparison to standard crust EOS