

A Laser-Accelerated Th Beam is Used to Produce Neutron-Rich Nuclei Around the $N = 126$ Waiting Point of the r-Process Via the Fission-Fusion Reaction Mechanism

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A laser-accelerated Th beam is used to produce neutron-rich nuclei around the waiting point $N = 126$ of the r-process via the fission-fusion mechanism

- Radiation Pressure Acceleration (RPA)
- Atomic stopping of dense ion bunches
- The astrophysical r-process and the $N = 126$ waiting point
- Experimental setup

D. Habs et al., “Introducing the Fission-Fusion Reaction Process: Using a Laser-Accelerated Th Beam to produce Neutron-Rich Nuclei towards the $N=126$ Waiting Point of the r Process”, arXiv:1007.1251v1 [nucl-ex], submitted to Appl. Phys. B

APOLLON laser stand alone

- 3·10 PW
- 15 fs
- $\sim 1/\text{min}$
- 10^{24} W/cm^2
- $\sim 10^{15} \text{ V/m}$

γ beam stand-alone

- $E_{\text{max}} = 13 \text{ MeV}$ (19 MeV)
- 12 kHz
- ring-down cavity for photons
- warm electron linac, 600 MeV
- high brilliance ($\Delta E_{\gamma}/E_{\gamma} \geq 10^{-3}$)
- high flux ($I = 10^{13} \text{ s}^{-1}$)

APOLLON + γ beam

- $E_{\text{max}} \approx 100 \text{ MeV}$
- $\sim 1/\text{min}$
- flux: $I = 10^8 \text{ s}^{-1}$
- higher harmonics + e bunches
- timing with fibre optics (fs)



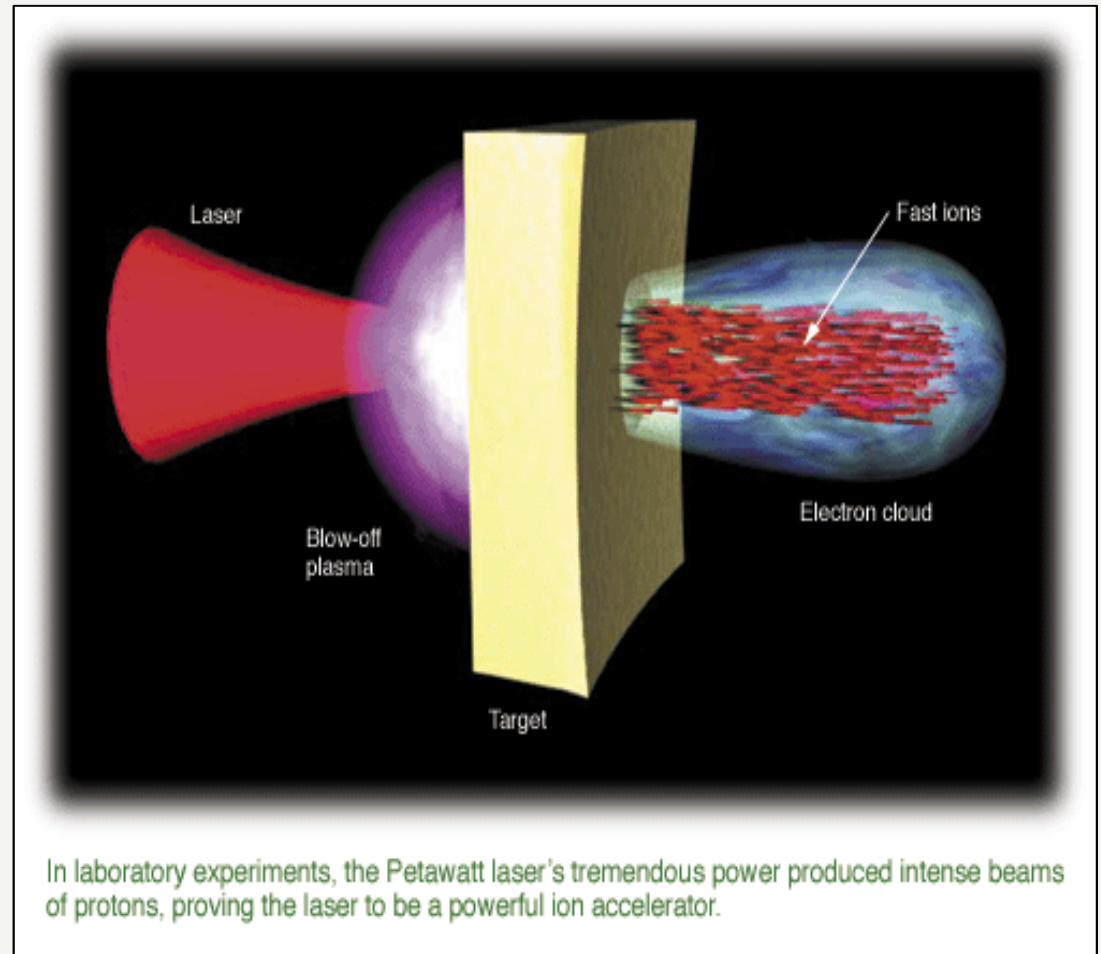
Ion acceleration

TNSA

(target-normal sheath acceleration)

- Low conversion efficiency
- Huge lasers are required

$$E_{\text{ion}} \propto \sqrt{I_{\text{Laser}}}$$

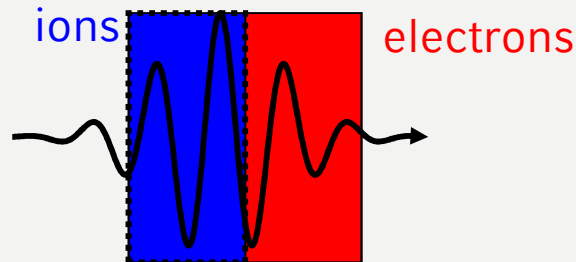


S.C. Wilks et al., Phys. Plasmas **8**, 542 (2001).



Optimum ion acceleration

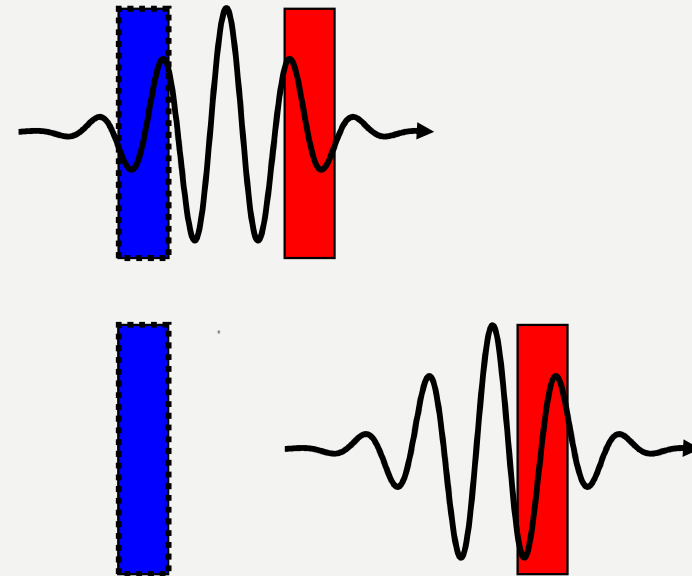
$$a_L \approx \sigma$$



$$D \approx 4nm \text{ for } a_L = 5$$

Optimum electron acceleration

$$a_L \approx \sigma \cdot 2\pi$$



$$D \approx 0.65nm \text{ for } a_L = 5$$

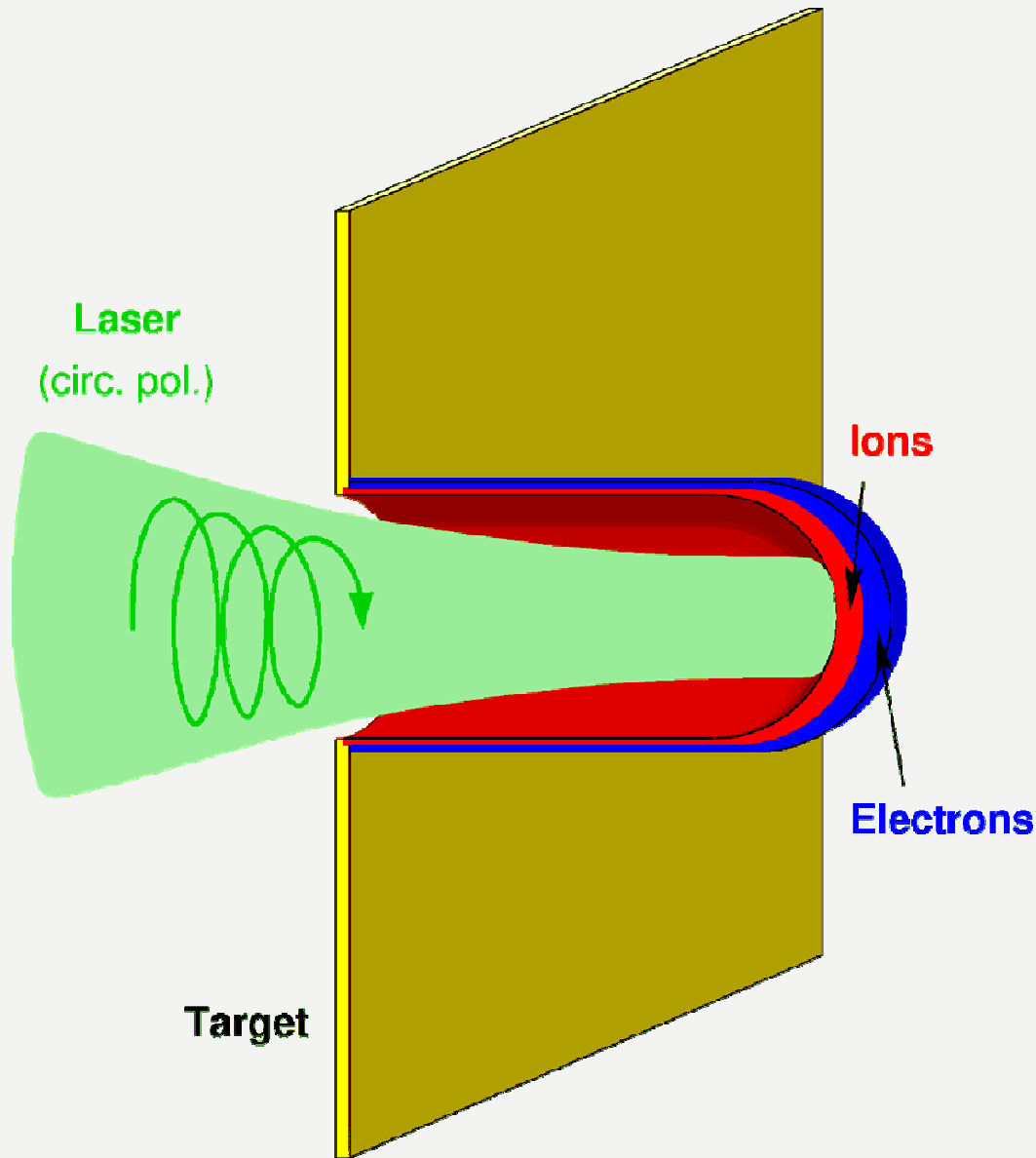
Normalized areal electron density: $\sigma = (n_e \cdot D) / (n_c \cdot \lambda) = \left(\frac{n_e}{n_c} \right) \left(\frac{D}{\lambda} \right) = \text{dimensionless}$

Normalized vector potential: $a_L^2 = I \lambda^2 / 1.38 \times 10^{18} W cm^{-2} \mu m^2 = \text{dimensionless}$

S.G. Rykovanov et al.,
New J. Phys. **10**, 113005 (2008).

O. Klimo et al.,
Phys. Rev. ST AB **11**, 031301 (2008).

Radiation pressure acceleration (RPA)



Cold compression of electron sheet.

Rectified dipole field between electrons and ions.

Neutral bunch of ions + electrons accelerated.

Solid-state density: $10^{24} \text{ e cm}^{-3}$

Classical bunches: 10^8 e cm^{-3}

$$E_{\text{ion}} \propto I_{\text{Laser}}$$

Very efficient!

Radiation pressure acceleration (RPA)



radiation pressure = intensity / c

high power (10^{15} W)

high intensities ($>10^{21}$ W/cm²)

short acceleration times (10's, 100's fs)

Target 5 nm, 2 μ m focus: many J rest energy

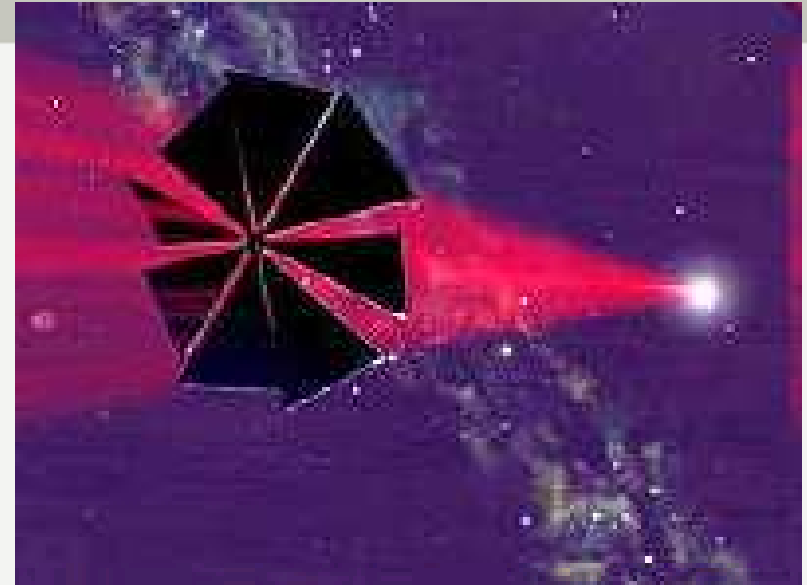


Photo courtesy [The Planetary Society](#)

small intensities (0.14 W/cm²)

“infinite” acceleration times (years)

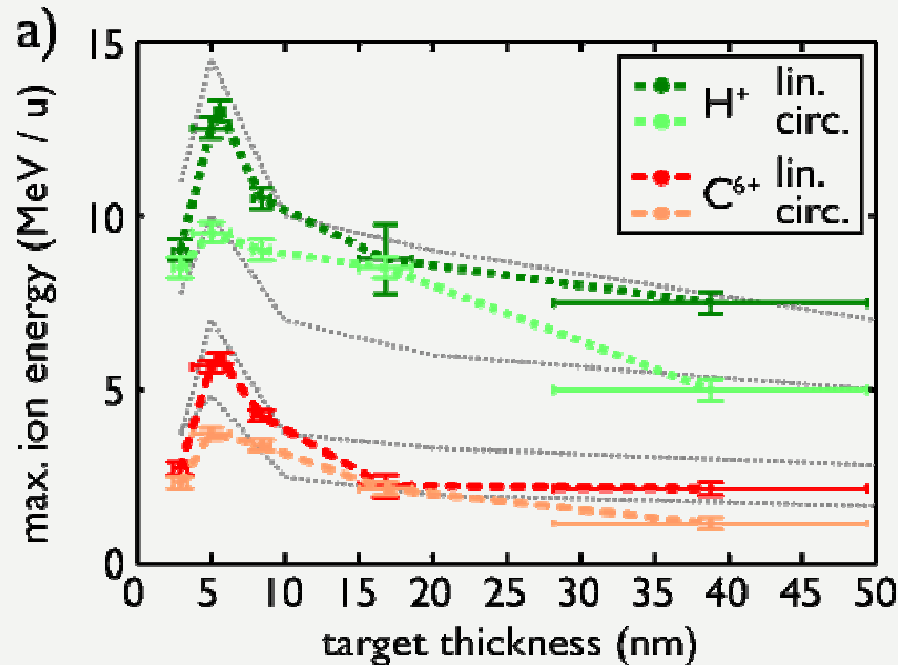
JFL Simmons et al, Am J Phys 1993 (Marx Nat 1966)

fuel is avoided but at a cost. A laser of 10 GW powered over 10 years or so would provide an energy equivalent to about the rest mass of a vehicle of 30 kg, and so would be sufficient to accelerate it to relativistic speeds. In fact, the

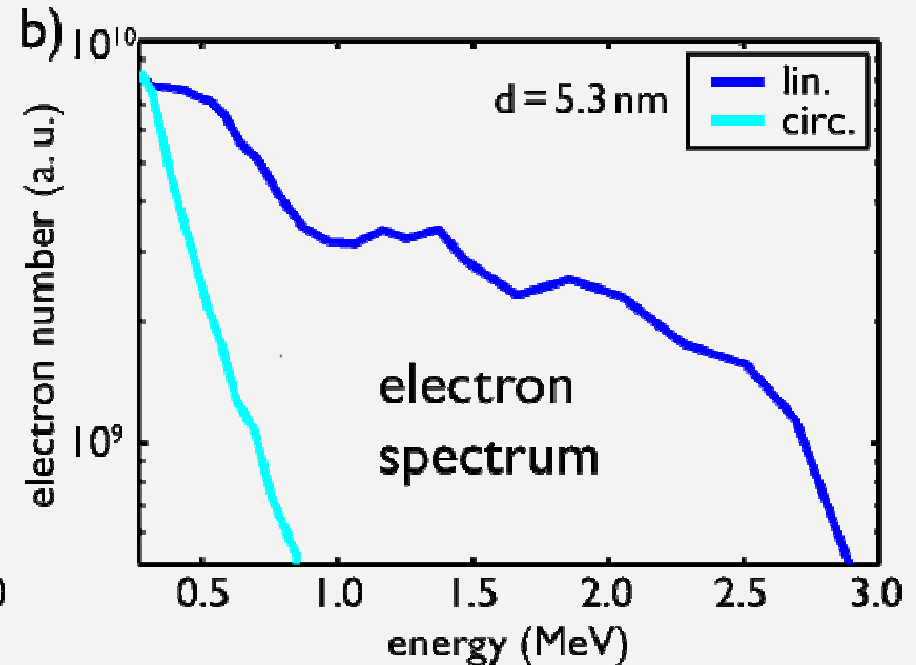


Max-Born Institute (MBI),
Berlin

Laser Power: 15 TW (700 mJ in 45 fs)
Focused Intensity: $a_L = 5$, Contrast: $> 10^{11}$



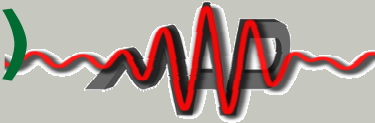
Peak at very low target thickness of 5.6 nm



Cold target for circular polarization

$$a_L \approx \sigma, \quad \sigma = \left(\frac{n_e}{n_c} \right) \cdot \left(\frac{D}{\lambda} \right), \quad a_L = \sqrt{I_L \cdot \lambda^2}$$

A. Henig et al., "Radiation pressure acceleration of ion beams driven by circularly polarized laser pulses",
Phys. Rev. Lett. **103**, 245009 (2009).



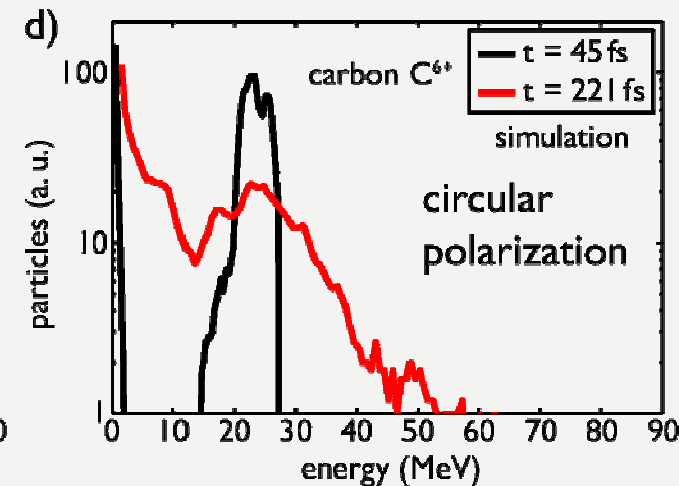
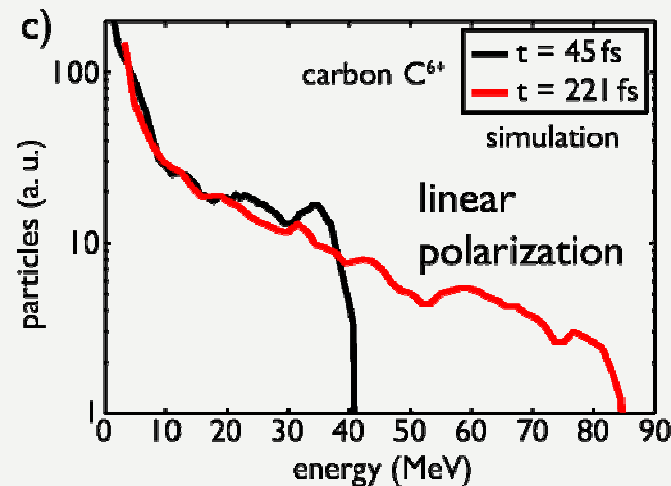
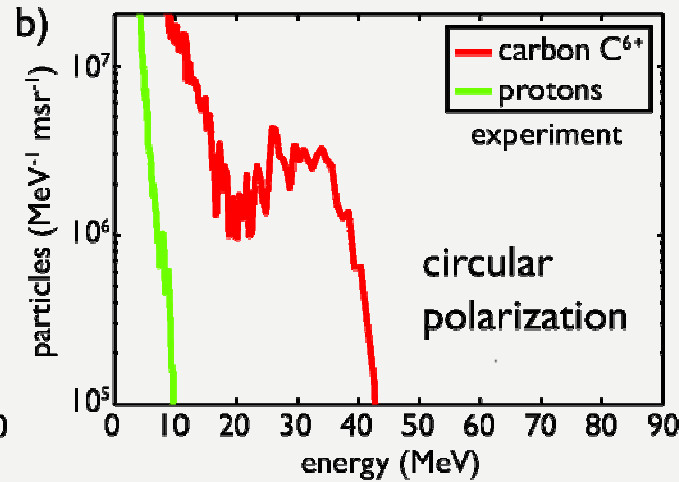
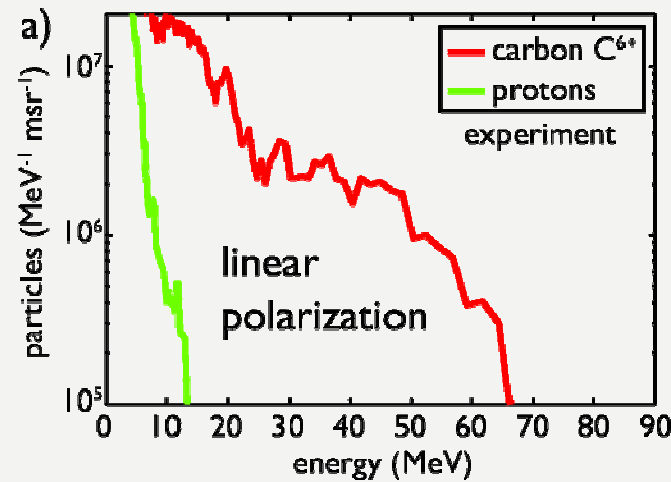
Hot electrons

cold electrons

Experiment

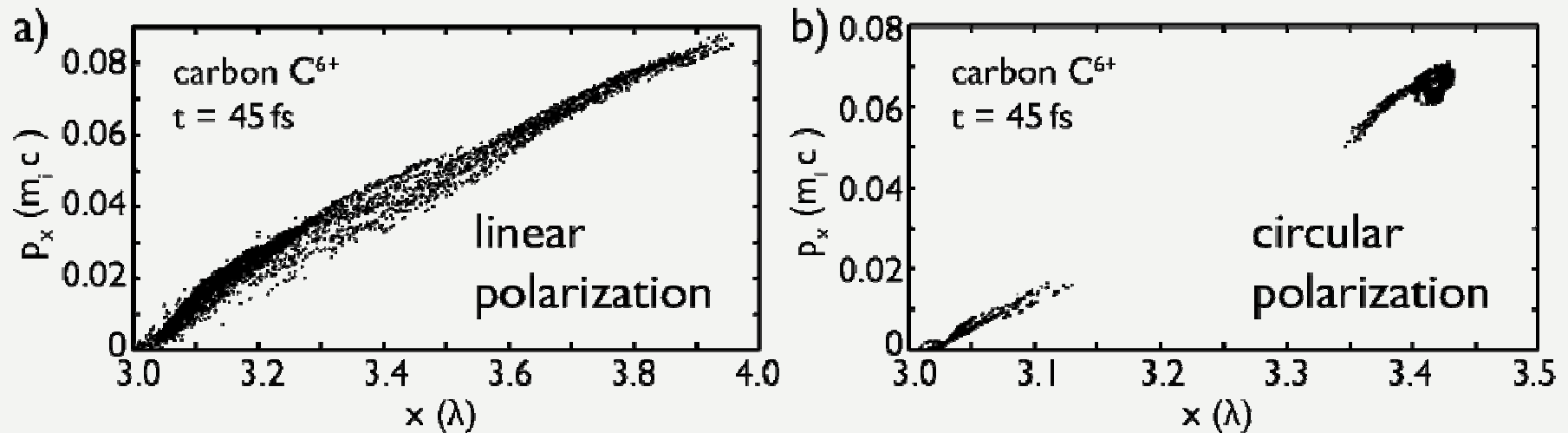
Theory

2D PIC simulations

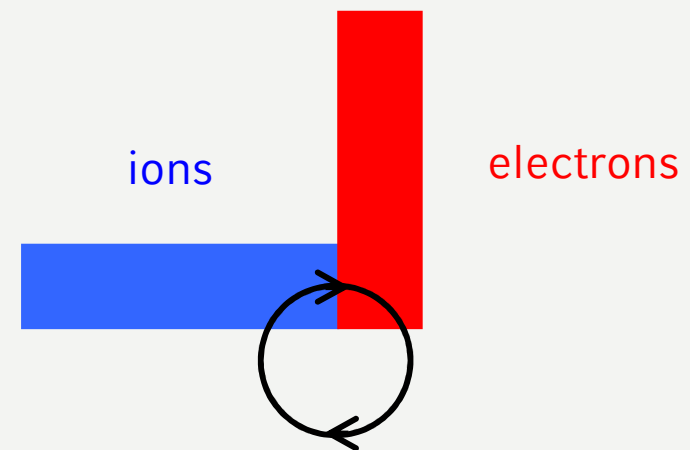




2D PIC simulations



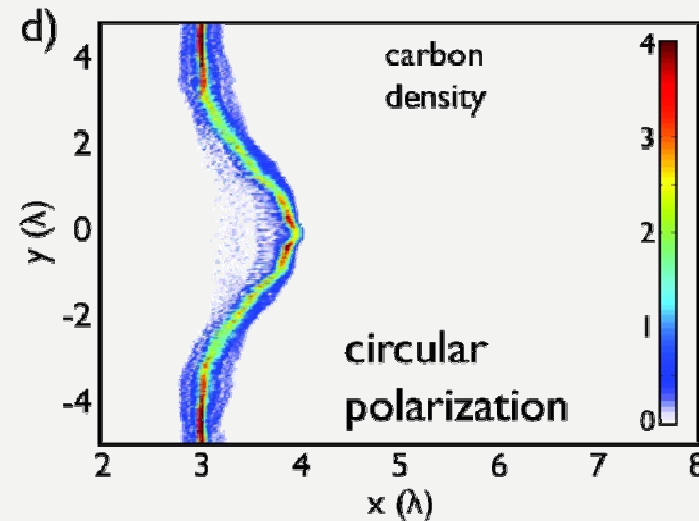
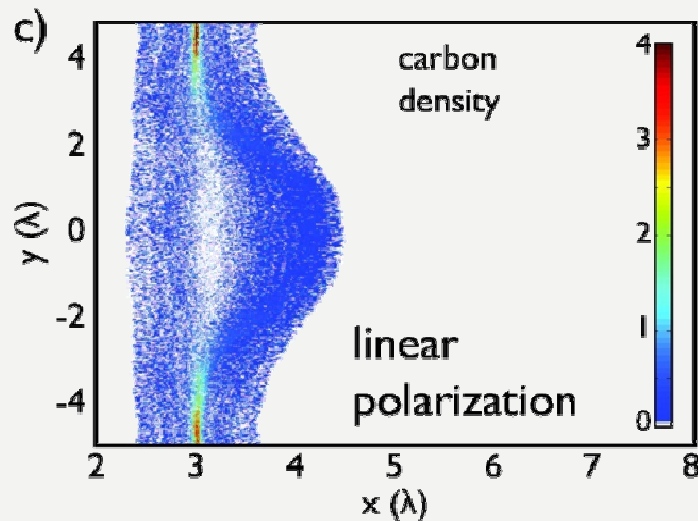
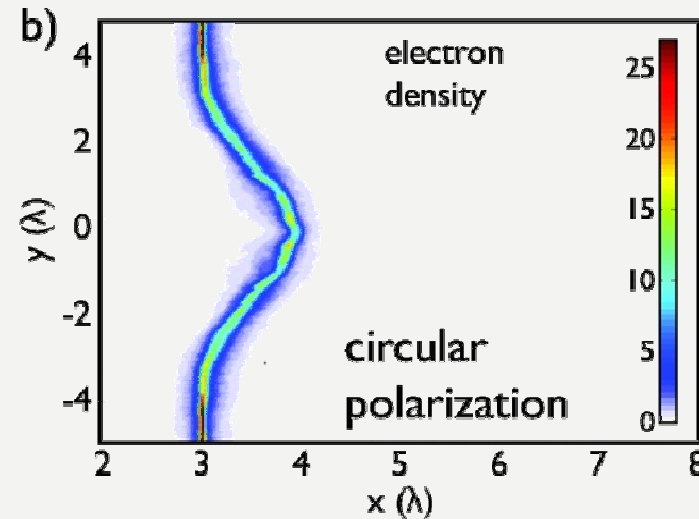
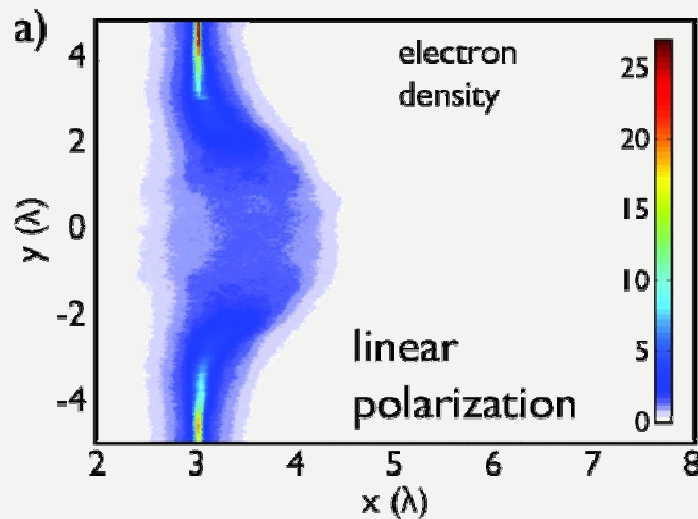
Phase space rotation
Peak in ion spectrum

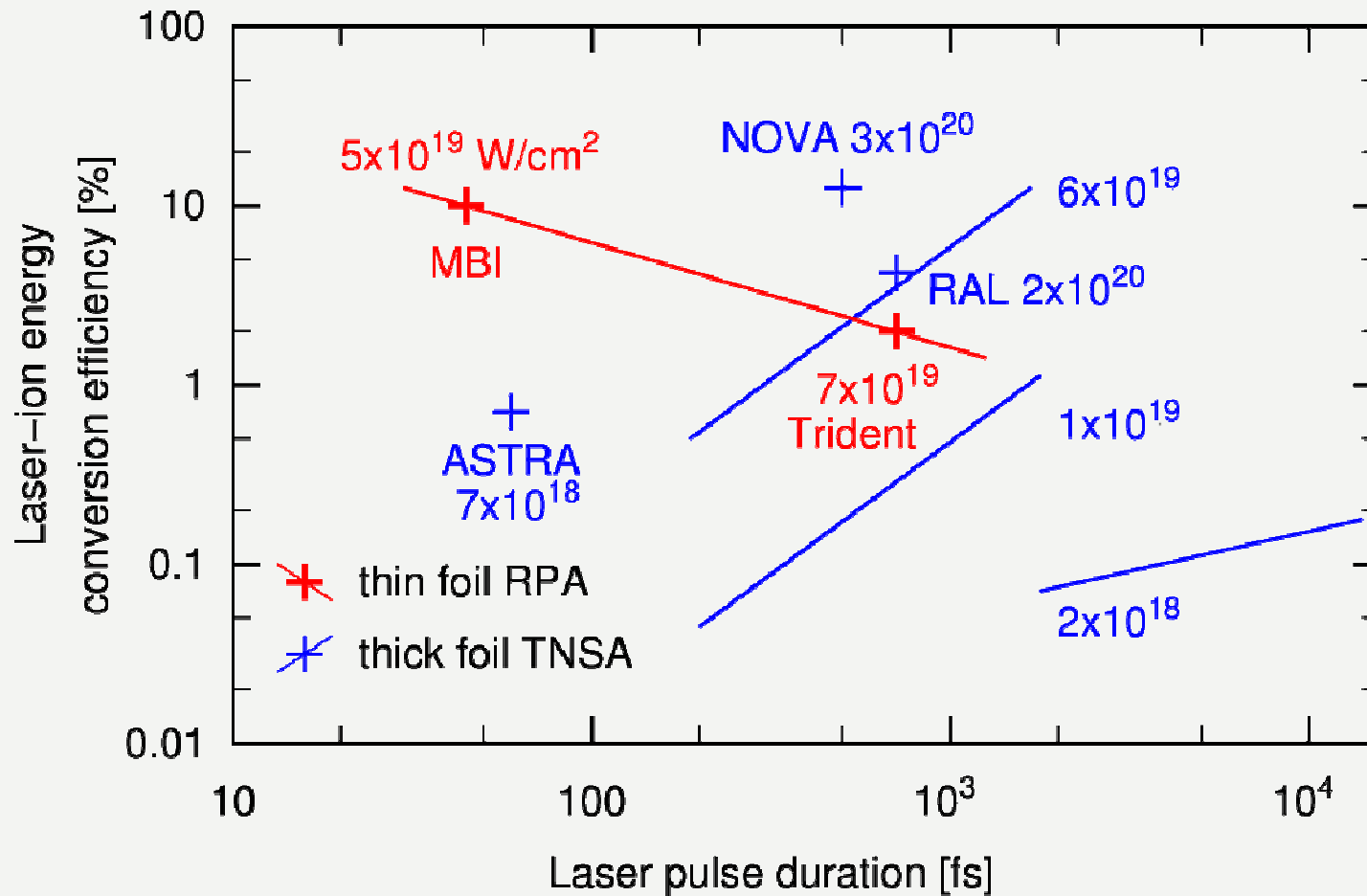




Hot electrons exploding foil

Cold electrons

 $\tau = 61$ fs

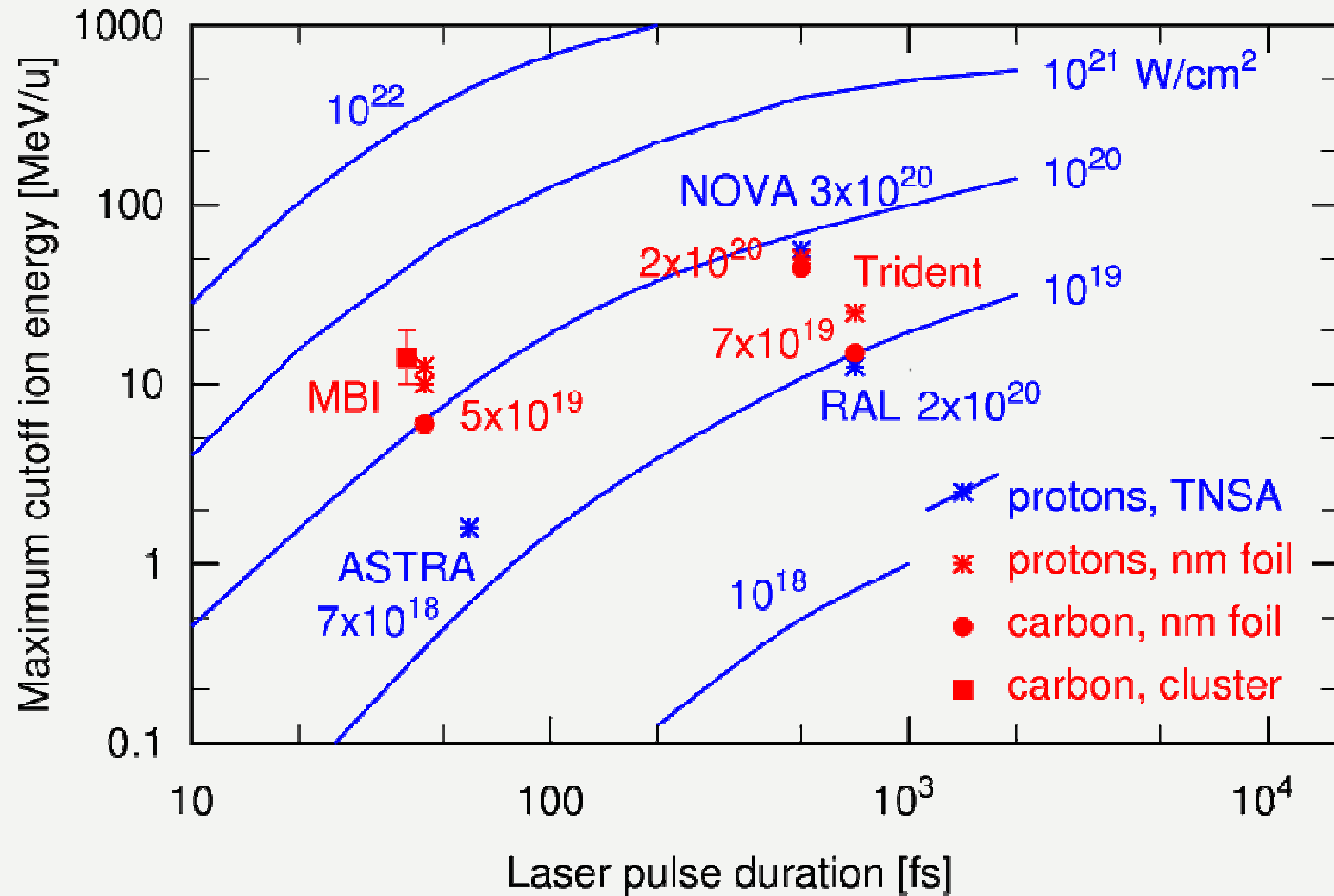


Ultra-thin foils:

MBI-laser (Berlin)
45 fs

Trident-laser (LANL)
500 fs

RPA: much higher conversion efficiency at short pulse durations.



Stopping power

of ion bunches with solid state density



Bethe-Bloch formula for individual ion:

$$-\frac{dE}{dx} = 4\pi n_e \frac{Z_{\text{eff}}^2 e^4}{m_e v^2} \left(\ln \left(\frac{m_e v^2}{e^2 k_D} \right) + \ln \left(\frac{k_D v}{\omega_p} \right) \right)$$

binary collisions
 k_D = Debye wave number

long-range
 collective interaction
 ω_p = plasma frequency

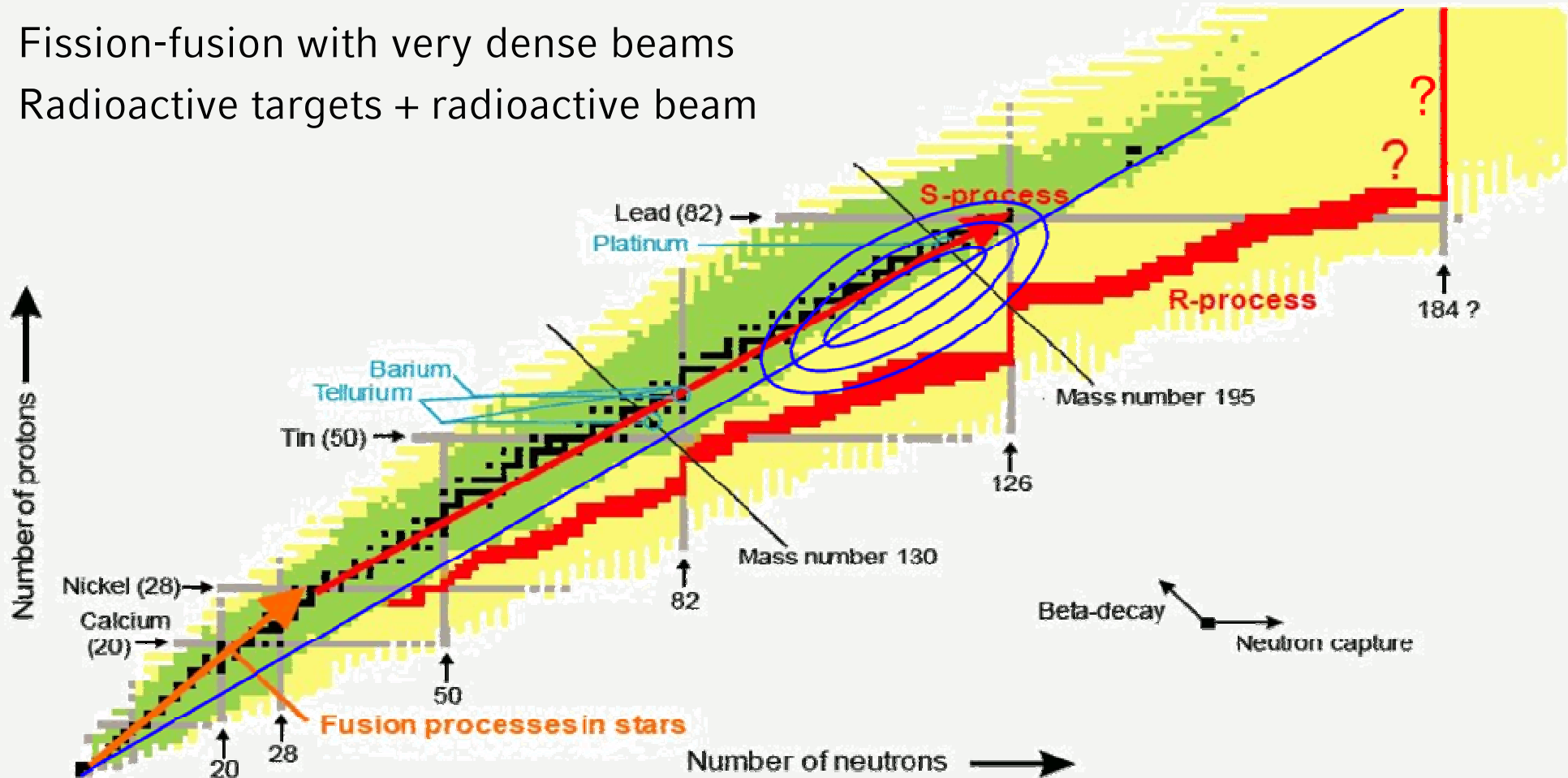
- a) enhanced stopping ($10^5 \times$) in low-density targets
 dense bunch interacts with collective wake
 → **reduced fraction of nuclear reaction**
- b) reduced stopping in solid target
 first electrons of bunch kick out electrons of foil like a snow plow.
 → **enhanced fraction of nuclear reactions.**

Chart of the Nuclides

r-process and waiting points



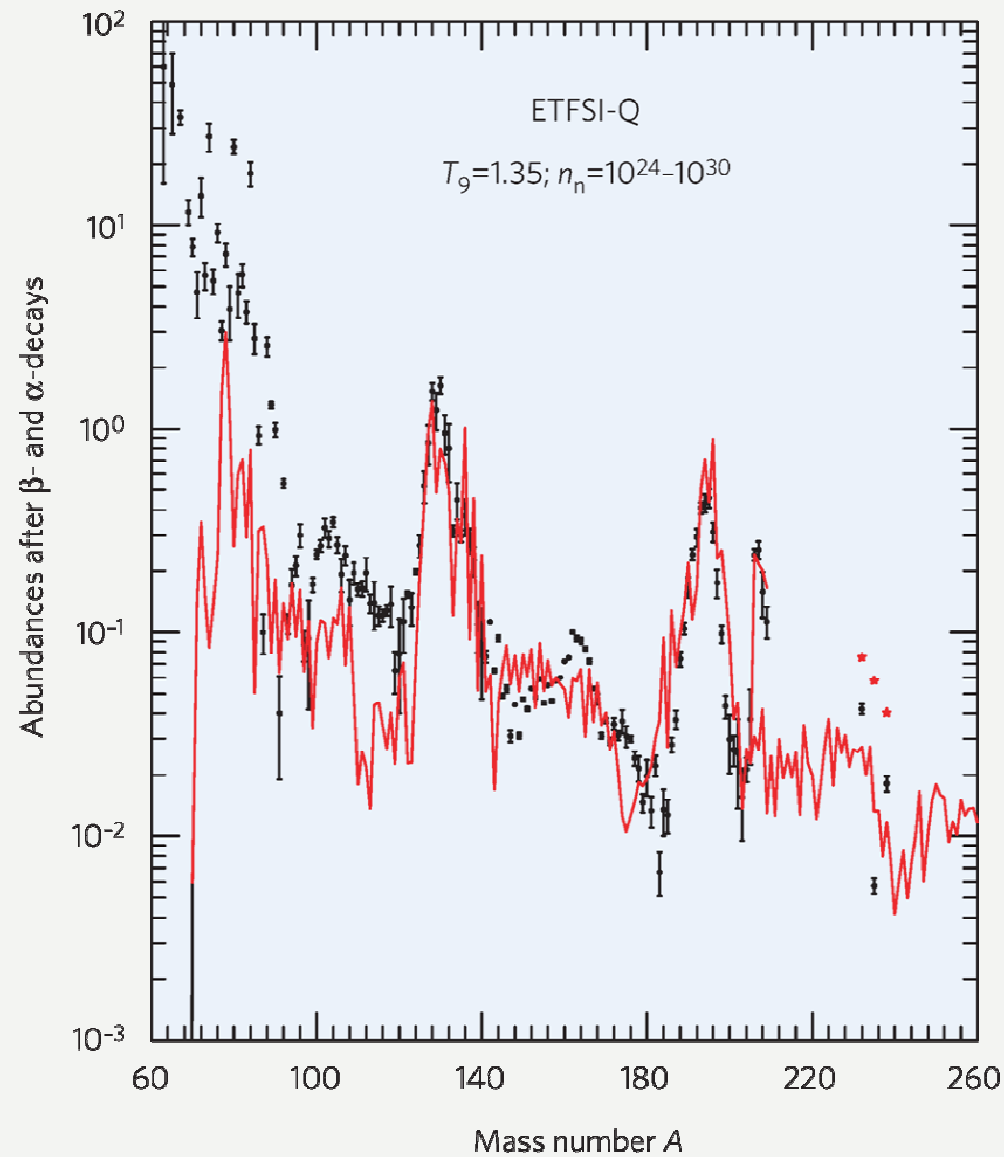
Fission-fusion with very dense beams
Radioactive targets + radioactive beam



- Superheavies: $Z = 110$, $T_{1/2} = 10^9$ a ?
- recycling of fission fragments ?

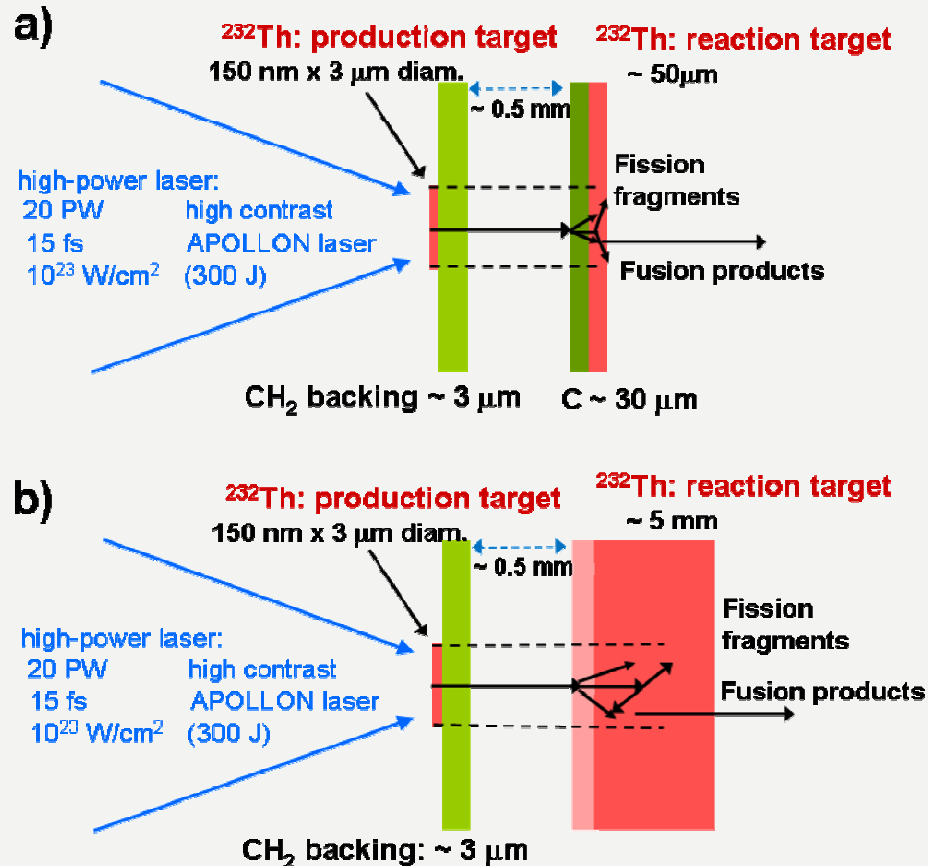
r-process

Solar abundances



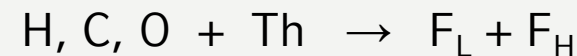
Fission-fusion reaction

very neutron-rich nuclei



7 MeV/u H, C, O, ^{232}Th , beam + ^{232}Th target

1) Fission



fission fragments in target

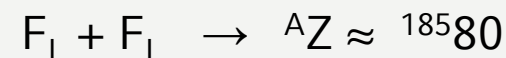


Reaction of radioactive short-lived light

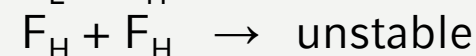
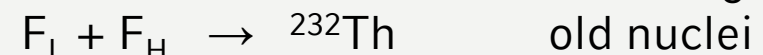
fission fragments of beam +

Radioactive short-lived light fission
fragments of the target

2) Fusion:



nuclei close to N=126 waiting point



a) Normal Bethe-Bloch stopping

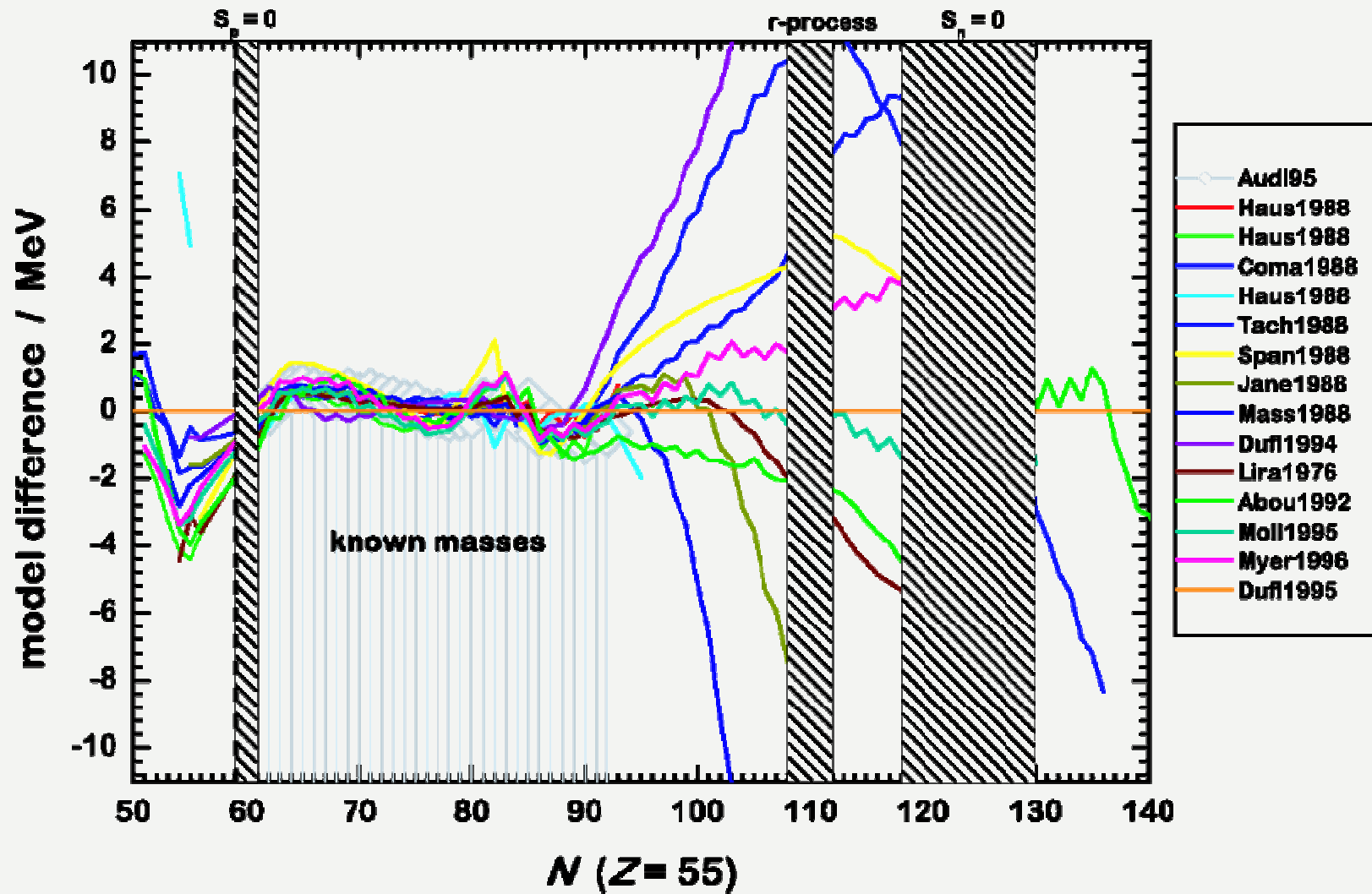
$$N_{\text{fissfus}} = 5 \cdot 10^2$$

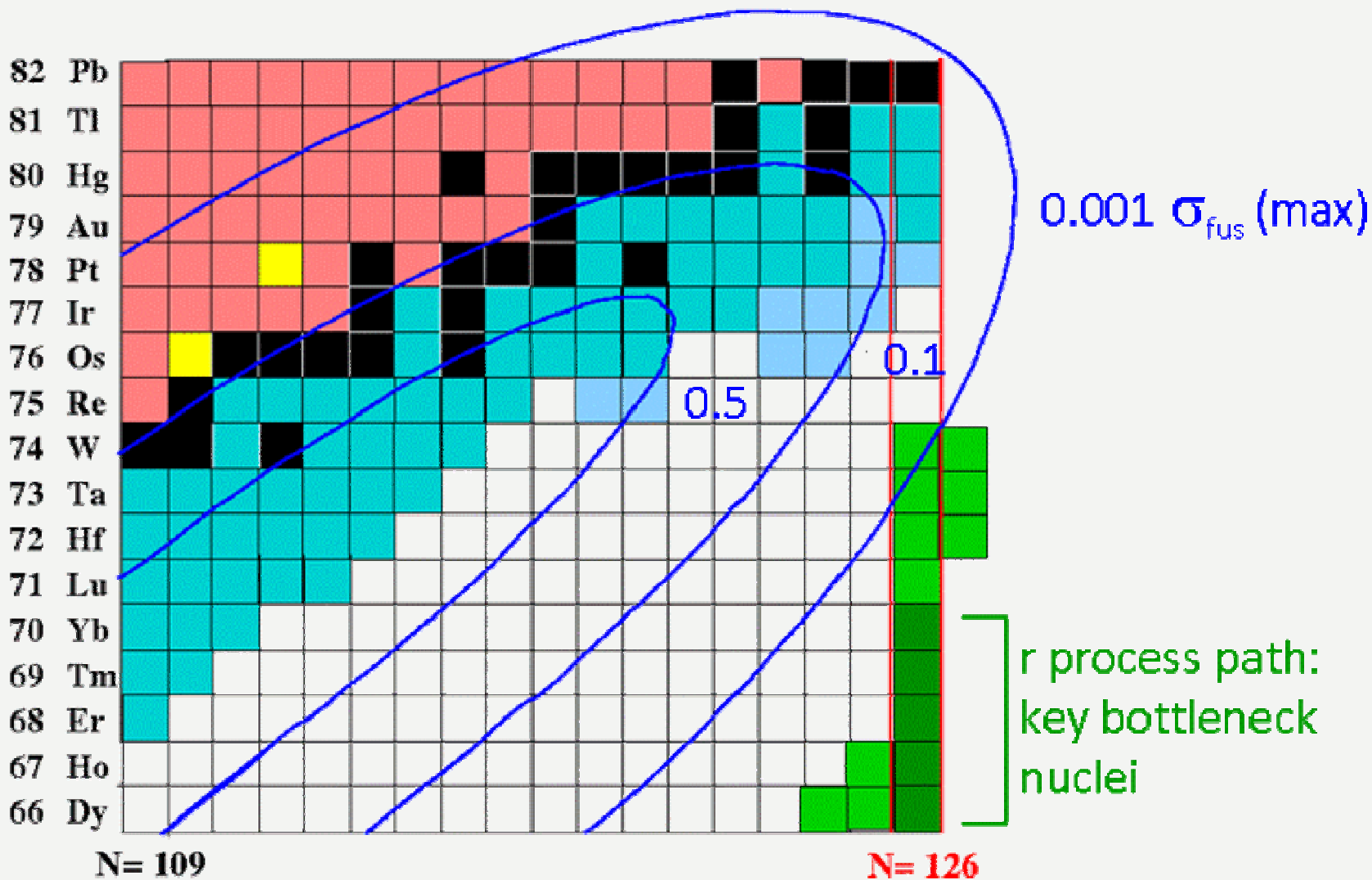
b) Reduced Bethe-Bloch stopping

$$N_{\text{fissfus}} = 10^6$$

Nuclear masses

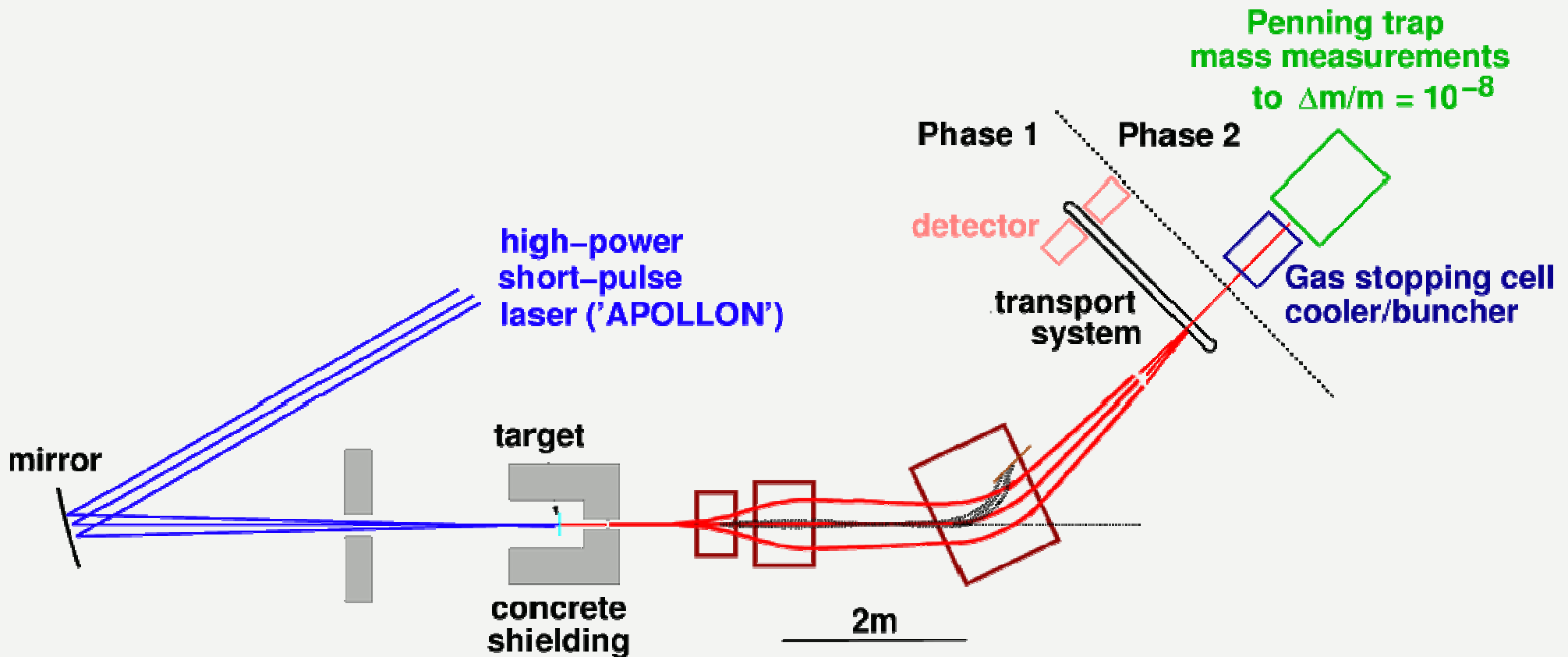
Theories and experiment





Experimental setup

neutron-rich nuclei in fission-fusion



- New scheme: Radioactive beam + radioactive target produced at the same time
- Perspective: With diode pumping of fiber lasers much higher repetition rates are possible:
from 1/min $\rightarrow$ 1 kHz
from PW peak power $\rightarrow$ kW to MW average power
factories for HiPER, etc. will make pumping devices much cheaper