

Beta-decay of very neutron-rich nuclei in $N=126$ region

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**Continuum QRPA framework
based
on the self-consistent g.s. from EDF theory.**

DF+CQRPA

A quest for spin-current part of (universal) NDF

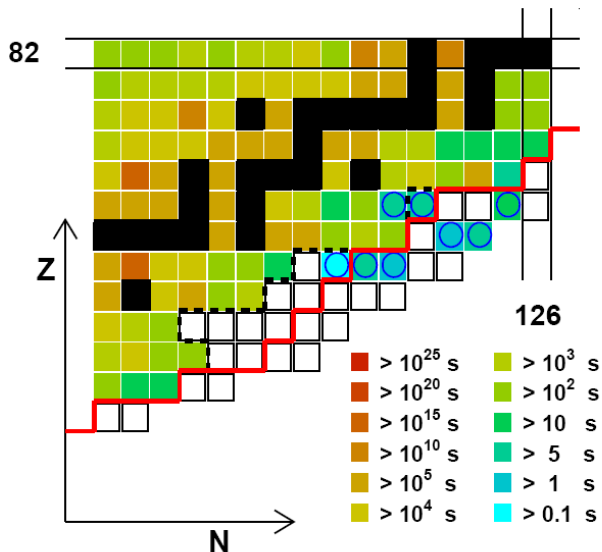
Nuclear structure at the limit of a high (N-Z).

Important applications:

- ❑ RNB experiments on beta-decay of short-lived neutron-rich nuclei.
- ❑ Astrophysical R-process modeling (masses, beta-rates, (n, γ)- rates...).
- ❑ Accelerator Driven Systems (1GeV p + 208Pb and 238U targets).

GSI experiments near N=126

The experimental technique used for the production of heavy neutron-rich nuclei was the in-flight fragmentation of relativistic heavy projectiles. In these collisions, the fragmentation process can populate a reaction channel referred to as 'cold fragmentation', where only protons are removed from the projectile and the excitation energy of the resulting projectile pre-fragment is so low that no neutron evaporation follows fragmentation. When using a ^{208}Pb beam, this reaction mechanism allows to produce heavy neutron-rich nuclei along the closed shell N=126.



GSI 2007

8 nuclides: 75 Re, 76 Os, 77 Ir ...

Colour scale: the known half-lives (NuBase).

Circles: T. Kurtukian-Nieto , J.Benlliure et al. Nucl-exp.0711.0101v1.

Empty boxes: identified with no half-life information.

GSI 2010

213Tl , 221,222 Po, 224At, 236Ac

L. Chen et al. Nucl-exp.0110.

Identified with the half-life information

Large-scale explorations of the nuclear mass-surface and beta-decay “ landscape”.

Gamow-Teller vs first-forbidden decays near closed shell N=126

EDF. Self-Consistent Ground State

$$E_{\text{Kohn-Sham}}[\rho, \nu] = \text{Tr} \left(\frac{p^2}{2M} \rho \right) + E_{\text{int}}[\rho, \nu]$$

$$E_{\text{int}} = \sum_{\text{main, Coul, sl}} \varepsilon_n[\rho] + \frac{1}{2} \nu^* F^\xi[\rho] \nu$$

"main" \propto volume + fin.range surface

$F^\xi \propto$ volume + surface

$$H = \begin{pmatrix} h - \mu & -\Delta \\ -\Delta & \mu - h \end{pmatrix}$$

$$h = \frac{p^2}{2m} + \frac{\delta E}{\delta \rho} \sim \rho$$

$$\Delta = \frac{\delta E_{\text{int}}}{\delta \nu}$$

$$\rho_0, \nu_0 \Rightarrow h_0, \Delta_0 \Rightarrow \rho_1, \nu_1 \Rightarrow h_1, \Delta_1$$

□ Energy Density functional

S.A. Fayans, S.V. Tolokonnikov, E. Trykov, D. Zawischa, Nucl. Phys. A676 (2000) 49.

I.N. Borzov, S.A. Fayans, E. Kromer, D. Zawischa Z. Phys. A335(1996) 117

□ DF3

Phenomenological (local) energy-density functional by **S.A. Fayans** et al.,
 δ -function + density dependent (volume+surface) pairing.

Fitted to the g.s. properties of **very neutron-rich nuclei** near "magic cross" $Z=50/N=82$ at ^{132}Sn .

$$F^{\omega}_{\tau\tau} = \frac{\delta^2 E}{\delta\rho^{\tau} \delta\rho^{\tau}}$$

$$F^{\xi}_{\tau\tau} = \frac{\delta^2 E}{\delta\nu^{\tau} \delta\nu^{\tau}}$$

*The ground state properties are less sensitive to spin-isospin components of EDF.
The spin-isospin (time odd) parts of the effective NN-interaction can be defined independently of the scalar (time-even) parts.*

Continuum QRPA based on the self-consistent ground state.

*Universal effective NN-interaction (A-independent).
Universal quenching $e_q^2=(0.9)^2$ (A, E-independent).*

T=1, ph

$$F^{\omega}_{\tau\tau} \propto g'_{LM} \delta_{r-r'} + \pi + \rho$$

T=0, pp

$$F^{\xi}_{\tau\tau} \propto g'_{pp} \delta_{r-r'}$$

Q_β - values: maximal beta-decay energy releases

Accurate description of the Q_β -values is crucial for beta-decay studies.

Q_β is correlated with the qp-energies, as both are obtained from the same DF framework.

DFs with $m^*/M \sim 1$ reproduce Q_β -values well enough :

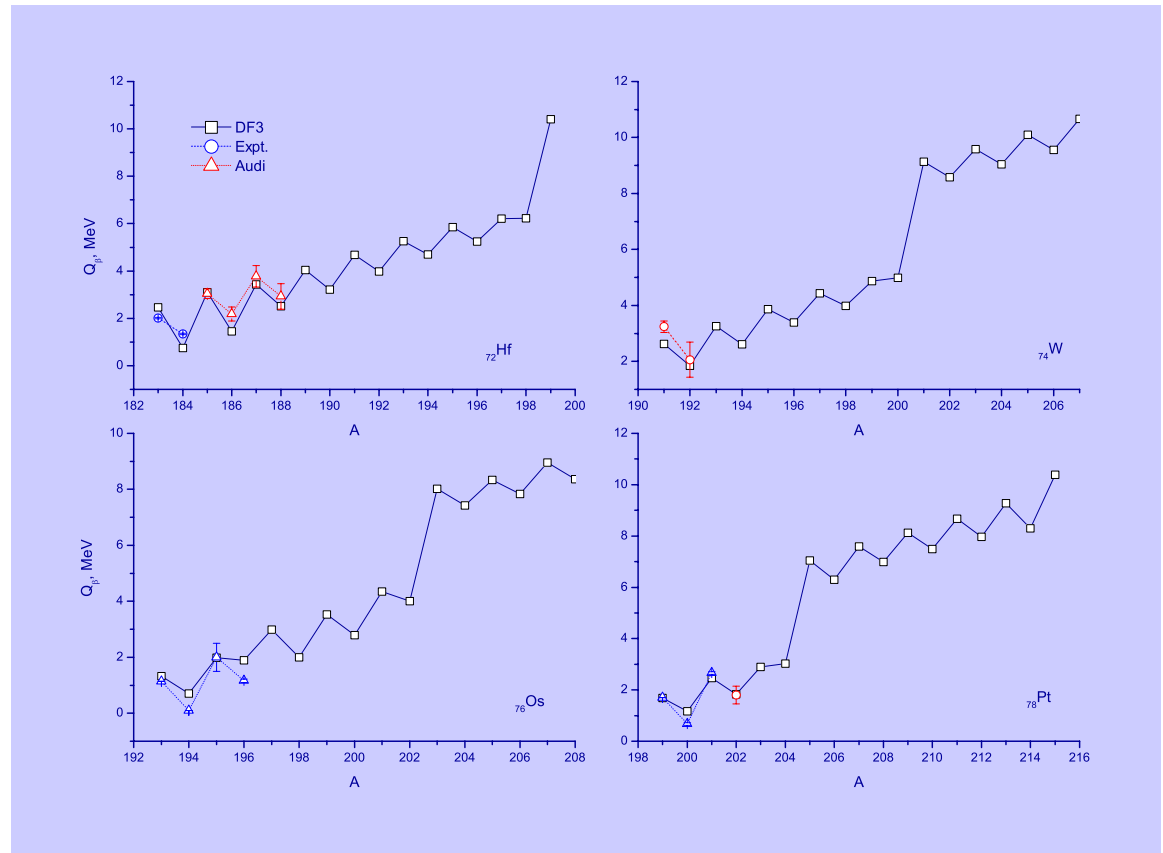
For N=82

Typical deviation from the data is 0.5 - 1.5 MeV

For N=126

Deviation is 0.5 - 0.6 MeV

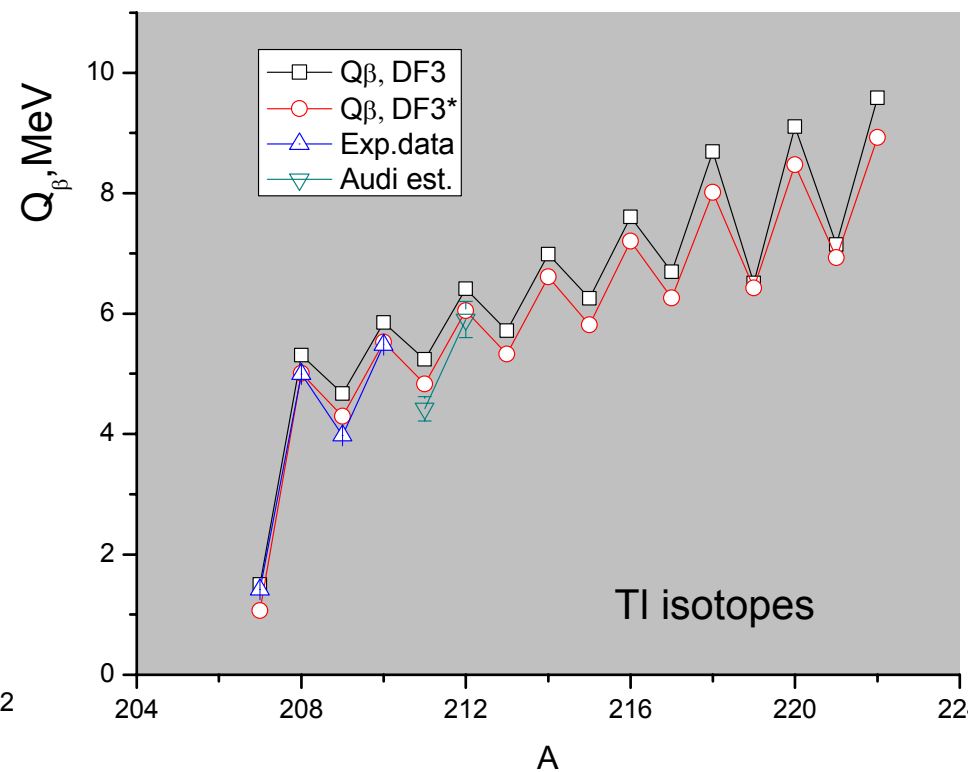
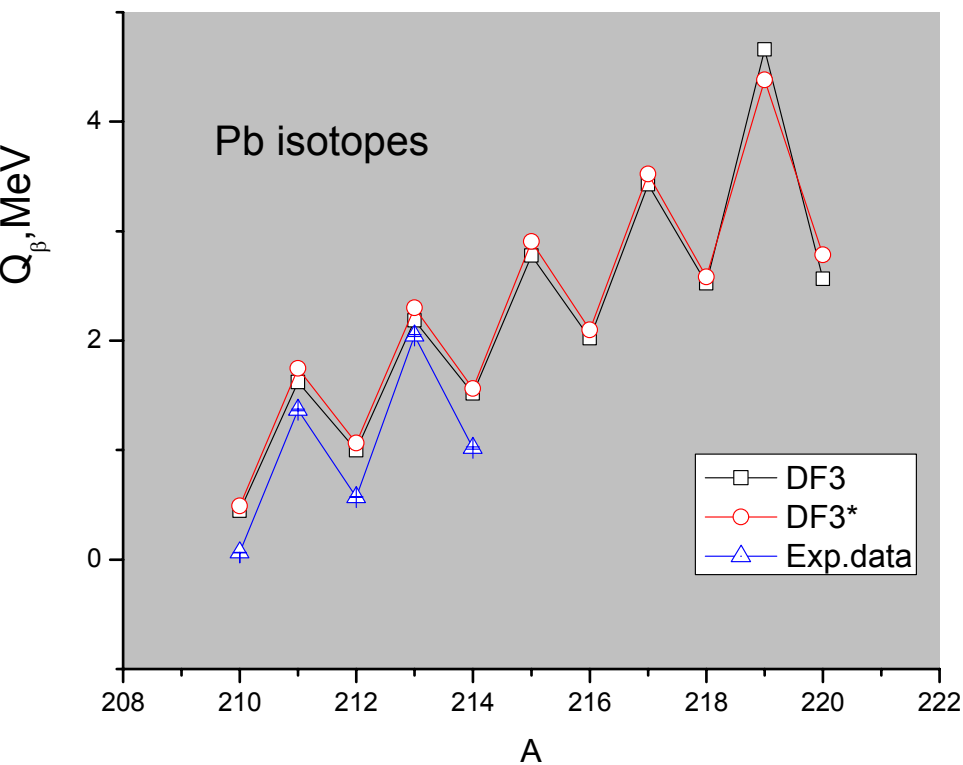
(Recent RDFs underestimate the phase-space, overshooting $T_{1/2}$)



Extended mass measurements at N=126 are of high value for improving the DF

Possible improvements: *Spin-orbital part of the Df3*

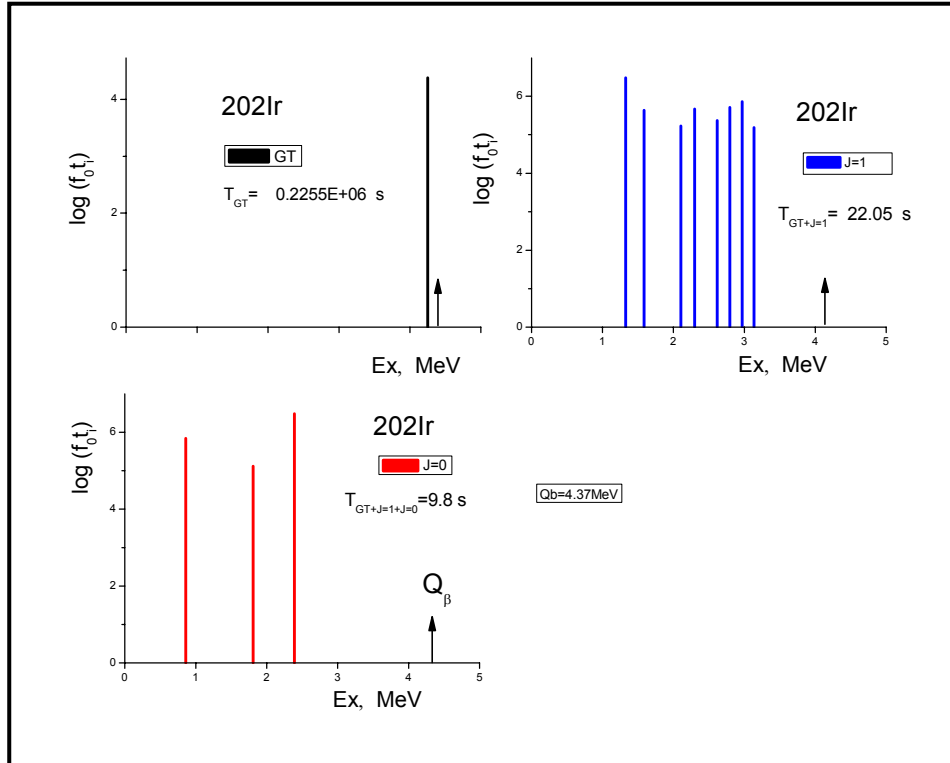
Example: Qbeta-values sensitivity to the velocity dependent interaction ($g_{1,np}$)



Spin-orbital splitting needs re-adjustment in a wide A-region (especially for the nuclei with high spin-orbital density, Z=80-90 region)

The GT and FF strength-functions are calculated within the single DF3+CQRPA framework

GT
 $n1h9/2 \rightarrow p1h11/2$



FF
 $n1i13/2 \rightarrow p1h11/2$

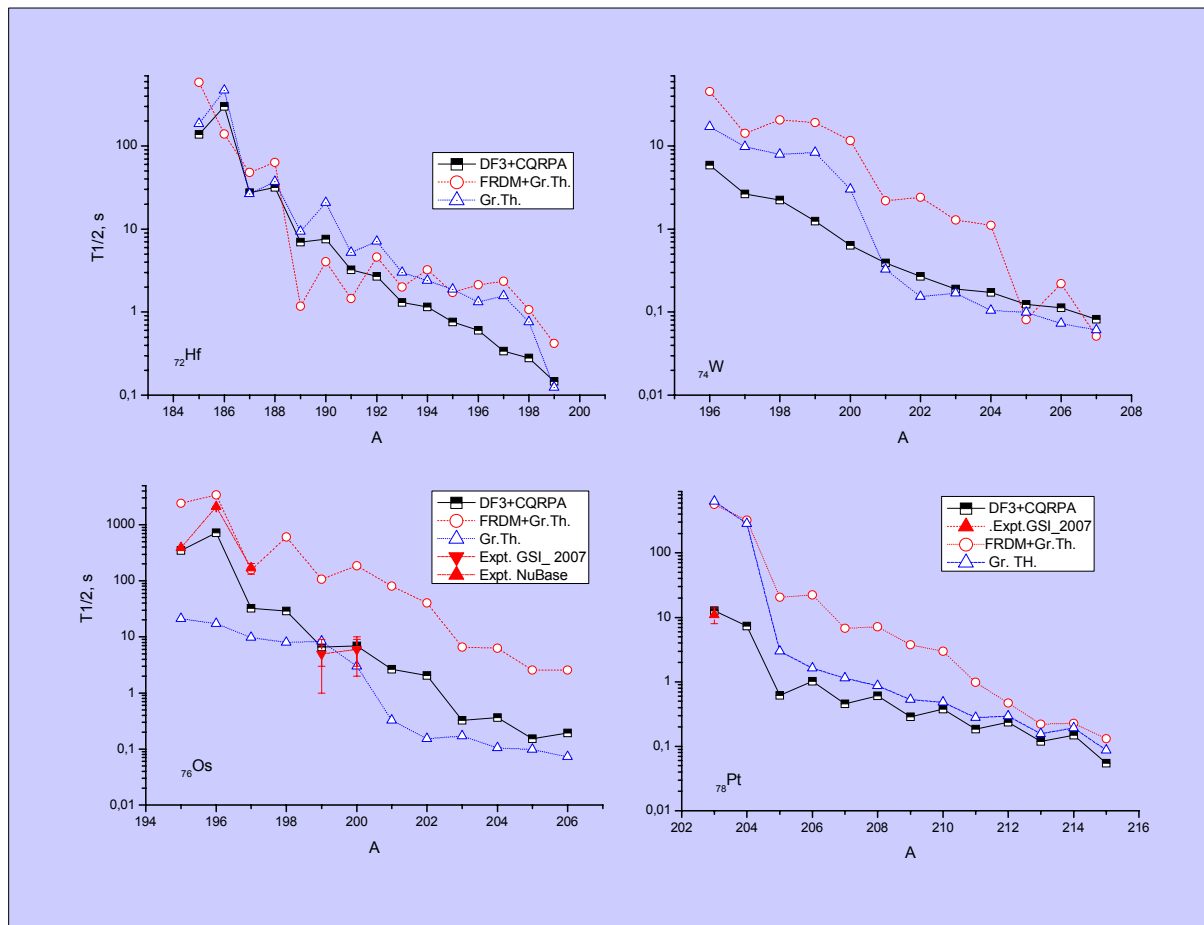
$n3p3/2 \rightarrow p3d3/2$
 $n3p1/2 \rightarrow p3s1/2$

The GT decays are retarded by

- the low “phase-space” factors $f(Q_\beta - Ex)$
- high occupancy factor of the $1p11/2$ orbital.

The high-energy FF decays compete with the low transition energy GT decays.

72Hf - 78Pt.



Half-lives approaching N=126

Typical deviation from the exp. data

FRDM = 30

DF3 = 1 - 2

A,Z	T1/2 (s)	DF3	FRDM
199Os	5 +4-2	6.6	106.8
200Os	6+4-3	6.9	187.1

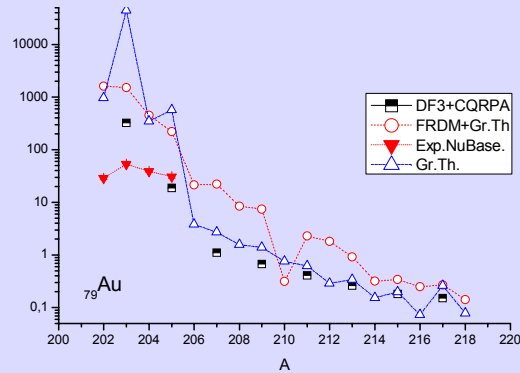
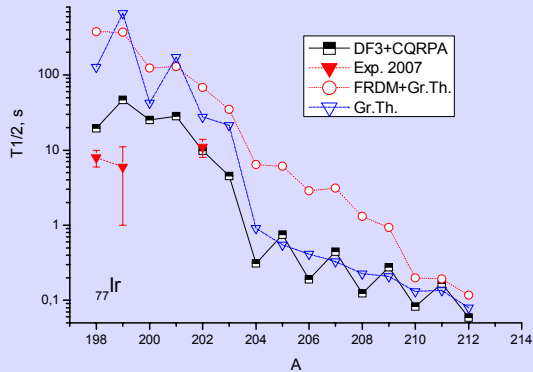
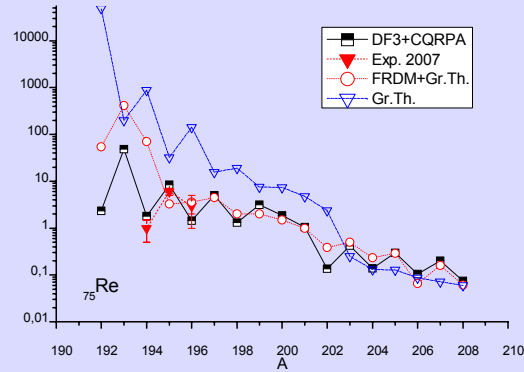
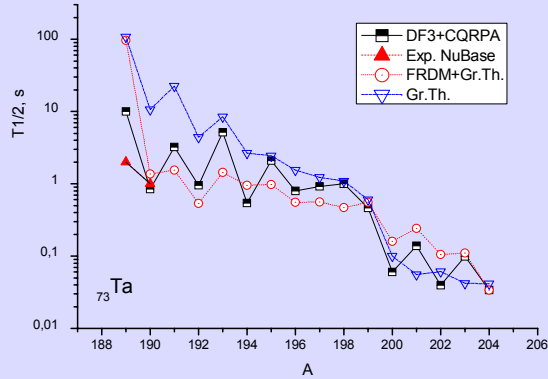
New GSI FRS-ESR campaigns

TRIAC, KEK Japan, Proposal 09

These Proposals are aimed to pin down the r-process waiting point nuclei

198Hf, 200W, 202Os, 204Pt

73Ta - 79Au.



Half-lives of odd-Z nuclides approaching N=126

Typical deviation from the exp. data

FRDM = up to 70

DF3 = up to 8

A,Z	T1/2 (s)	DF3	FRDM
194Re	1 +/- 0.5	2.1	70.8
195Re	6 +/- 1	8.5	3.3
196Re	3 +/- 1-2	1.4	3.6

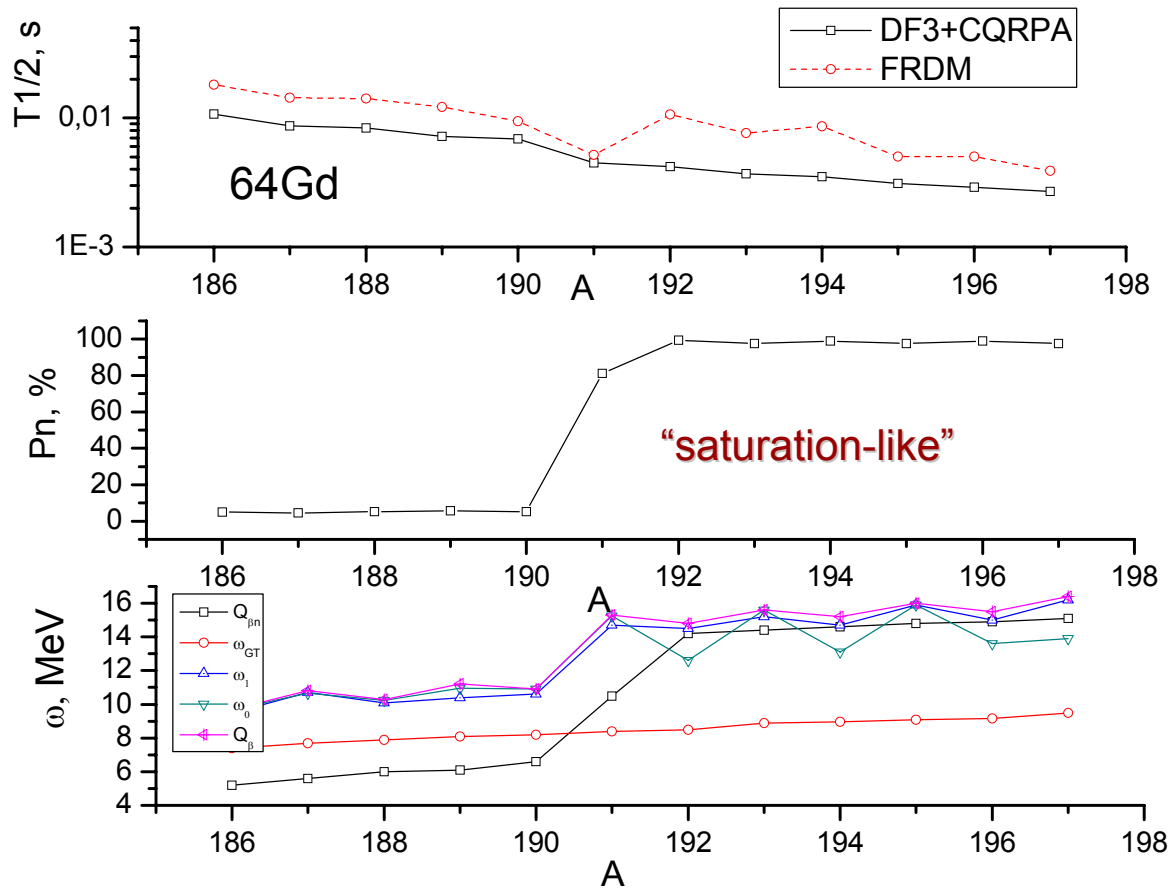
205Au

Exp. 31 +/- 2 (s)

FRDM = 221,0 s (7)

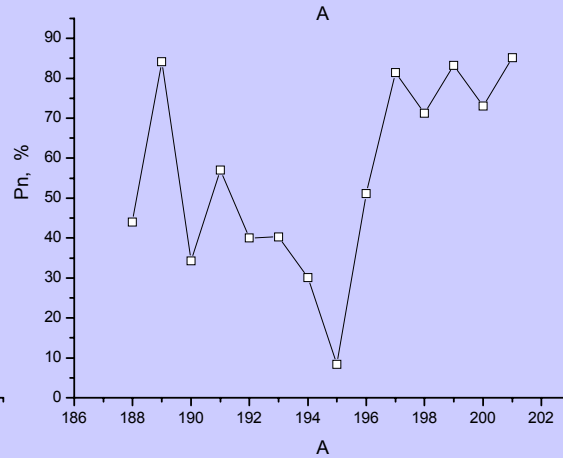
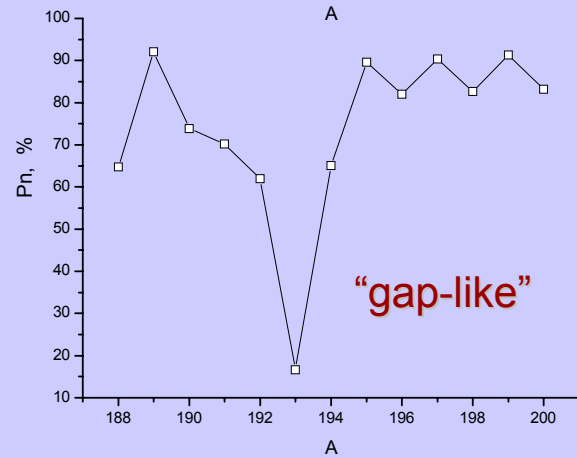
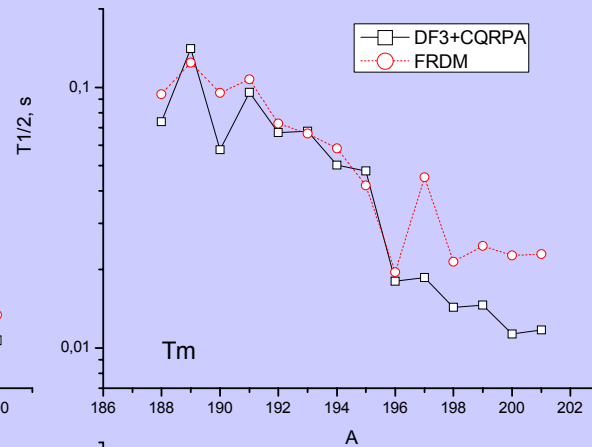
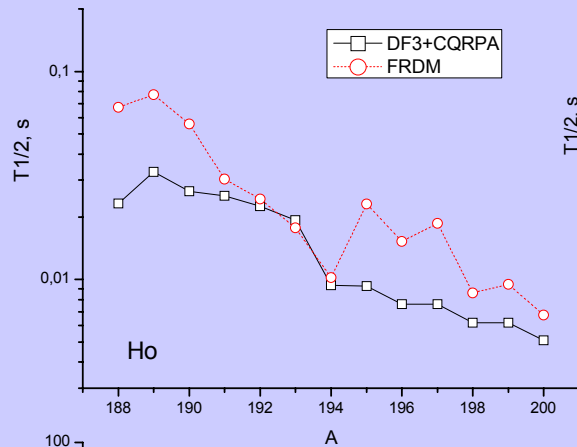
DF3 = 18.7 s (1.6)

^{64}Gd isotopes near $N=126$.

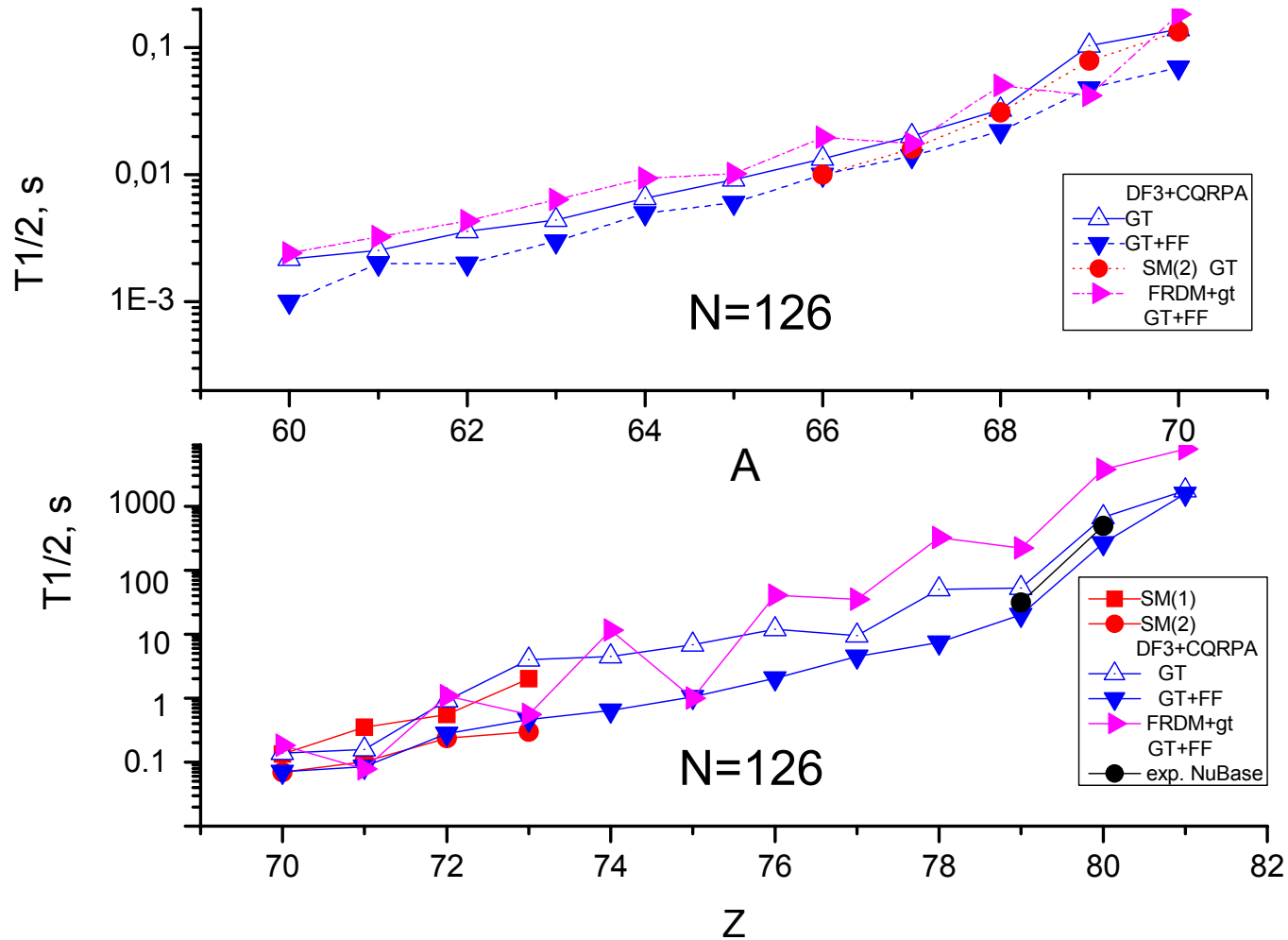


$^{67}\text{Ho}, ^{69}\text{Tm}$ isotopes near $N=126$.

$Z=67, 69, N \sim 126$



N=126 waiting point nuclei

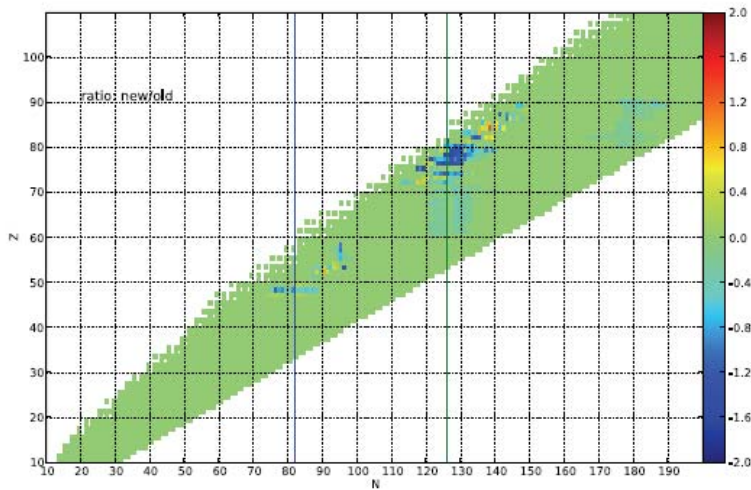


Impact of the half-lives on R-process abundances at N=126

(A.Arcones e.a. 2010)

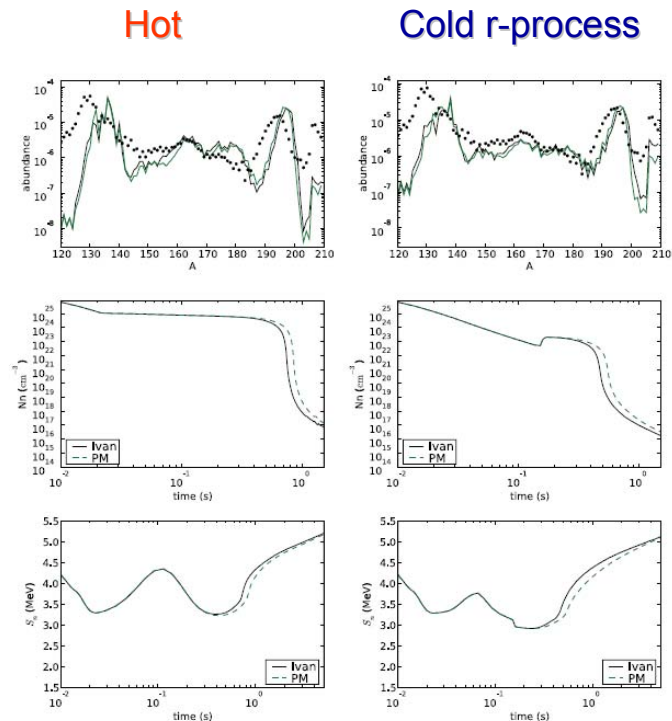
FRDM half-lives were replaced by the DF3 ones
for the nuclides near N=126 only.

Changes the abundances near the third peak (A~195)
and for fissioning nuclei



The colors represent $\log(t_{1/2}(I)/t_{1/2}(PM))$.
One can easily see which rates are updated.

The main effect occurs at N=126, where the new half-lives are smaller. This produces a faster evolution after reaching this region and it is visible in the figure of the average neutron separation. We have identified the minima of S_n with the moments when the path reaches N=82 and 126. Since the evolution is faster neutrons are consumed more rapidly and neutron density starts to drop early. Therefore, this could have some consequence for the freeze-out.



I. Summary (Beta-decay)

- Global calculations of the β -decay half-lives.
Systematic studies in a wide ($Z=15 - 92$) region:
Table I. Spherical nuclei. Used for new experiments
at GSI, KEK, HRIBF
- GT decay dominates for $Z \sim 28, 50$ for $N < N_{\text{mag}} = 50, 82$
Competition of GT and high-energy FF decays
for heavier $Z = 60 - 82$, N approaching $N_{\text{mag}} = 126$.
- Signatures of the GT/FF competition: “saturation-like” and
“saw-like” patterns in $P_n(A)$ behavior.
- The shorter half-lives make evolution faster, break through
the $N=126$ waiting points faster. The matter flow to heavier
(fissioning) nuclei should be also faster in DF3+CQPRA.

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I.N.B, K. Langanke, G. Martinez-Pinedo, F. Montes Nucl.Phys. A814,159,2008

J.A.Winger ... INB et al., HRIBF Collaboration Phys.Rev.Lett 102 (2009) 142502

T.Kurtukian-Nietto, J.Benlliure, .INB...et al.GSI Collaboration. nucl-ex 0711.0101v1, (→Phys.Lett B, 2010)