

Simulation of Fragment Separators

Helmut Weick, GSI Helmholtzzentrum Accelerator Seminar, 21st Feb 2019

- Production Reactions 3-10
- Fragment Separation 11-13
- Simulation Techniques 14-18
- Resolution 19-22
- Examples 23-24
- New Separators 25-36 (BigRIPS, Super-FRS, FRIB)





Fragment Separators Worldwide = devices to separate nuclides produced in-flight



adapted from T. Kubo, NIM B 376,102 (2016)

How to produce Rare Isotopes ?



The Abrasion Model for Fragmentation

Mass loss is given by geometrical cut



Assumption: Trajectories of nucleons are straight lines. Masses of projectile spectator and target spectator are given by the geometrical overlap as a function of impact parameter.

Velocity of projectile fragment ≅ velocity of projectile

friction is small

Distribution of Fragments



Mass Distribution

Mass distribution ¹³⁶Xe + Pb, 1 A GeV



thesis Antoine Bacquias, U. Strasbourg (2008)

Production Cross sections







Momentum Distribution in Fragmentation Reaction

$$\sigma_{p_{\parallel}}^2 = \sigma_{p_{\parallel Fermi}}^2 + \sigma_{p_{\parallel recoil}}^2 + \sigma_{p_{\parallel Coul}}^2$$

- Fermi momentum of bound nucleons
- Mom. transfer by evaporated nucleons
- Coulomb expansion in multi fragment.



Βρ-Δ**Ε**-**Β**ρ Separation Method

scheme of FRS @ GSI, L=72m

Production Target



Βρ-Δ**Ε**-**Β**ρ Separation Method

scheme of FRS @ GSI, L=72m



Βρ-Δ**Ε**-**Β**ρ Separation Method



 $\Delta p/p_1 = \Delta p/p_2$ with $p_1 > p_2$

Non-Liouvillean system.

Coupling of longitudinal and transverse emittance, $\Delta p \rightarrow \Delta p / f$ (f>1), but $\varepsilon_x \rightarrow \varepsilon_x * f$ In addition angular and energy-loss straggling.

Input for Simulation

nuclear physics:



Monte Carlo MOCADI, LISE⁺⁺ MC

Optimize for beamlines with matter:

Ions always fly in forward direction, no multiplicity, physics routines adjusted for beamline needs, GEANT would be orders of magnitude slower.

Parametrized, pre-calculated physics:

Do not evaluate physics for each ion, optics by transfer matrices, no magnetic fields, use parametrization in target (EPAX, Goldhaber, Morrissey) pre-calculated results for energy loss (ATIMA spline tables),

Biasing:

10¹⁰ ions/s like in reality are impossible for Monte Carlo, Do not create fragments with probability like in reality, calculate certain number for one nuclide, biasing is done by the very different production cross sections.

Convolution Technique

Each target, each piece of matter in beamline, each collimator reshapes the distribution in position, angle, energy of an ion species





- + Fast, only a low number of points have to be calculated. Calculation of all fragments possible while watching.
- + Parametrization on log scale, even small tails are still visible
- Becomes very difficult with many cross correlations.
- Usually limited to linear transformations (only 1st order optics).

Simplification to Blocks

Example FRS:

Star shaped vacuum chamber is difficult to describe in convolution technique. Only use independent cuts in x or y distribution.



Details cannot be taken into account \rightarrow do not try. Replace single aperture cuts by effective cut for whole section.

One cut in x, y after each separator stage,

one angular acceptance a_{max} , b_{max} for each stage (S0-S1, S1-S2, S2-S3, S3-S4).



Values in LISE adjusted to results of MOCADI simulation. In normal FRS operation agreement of transmission within 20%

Transfer Matrix Description

Ion-optical coordinates

x, $a=p_x/p_0$, y, $b=p_y/p_0$, $\delta = \Delta B\rho/B\rho_0$ $B\rho = c_0\beta\gamma m/q = \nu m/q$ $\delta_m = \Delta m/m$, $\delta_q = \Delta q/q$, $\delta_\nu = \Delta \nu/\nu_0$

Separator stage, D

Degrader wedge W

($x|\delta$) -($x|\delta$) ($x|\delta$) x ((X|X))х (a|x) (a|a) (a| δ) -(a| δ) (a| δ) a a $\left|\begin{array}{cccccc} 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 \end{array}\right|$ δ_m = δ_m δ_q δ_q $\delta_v f$ $\left(\delta_{v}\right)_{i}$ 0 0 х х 0 0 а a 0 δ_m = 0 δ_m δ_q δ_q 1 0 δ_v) $(\delta_{\mathcal{V}}|\mathbf{x}) \quad (\delta_{\mathcal{V}}|\mathbf{a}) \quad (\delta_{\mathcal{V}}|\delta_{\mathcal{M}}) \quad (\delta_{\mathcal{V}}|\delta_{\mathcal{Q}}) \quad (\delta_{\mathcal{V}}|\delta_{\mathcal{V}})$ δ_v

Full system

 $= D_2 \cdot W \cdot D_1$





Resolving power R

$$R = \frac{1}{x_0 a_0} \int \frac{\overrightarrow{B(s)}}{B\rho} d\vec{f}$$

For given emittance $x_0 a_0$ the B-field area covered by the beam defines R.

R is limited by momentum spread

 $\begin{array}{ll} 1^{st} half: & \Delta p/p_{nucl.\ reaction} \sim 0.4 - 8 \ \% \\ 2^{nd} half: & \Delta p/p_{E-loss\ fluct.} \sim 0.05 - 0.3 \ \% \\ & statistical\ energy-loss\ fluctuation \end{array}$

Transfer Matrix Description

Geissel, Weick, Winkler, Münzenberg, Yavor, NIM B 247 (2006) 368.

Momentum resolving power of one half

 $R = \frac{(\mathbf{x}|\delta)}{\Delta \mathbf{x}_{f}} \sim 1000-2000$

Energy-loss straggling

$$\delta_{v,str} \sim 1/1000$$

 $\frac{\text{Resolution in mass, charge}}{\Delta X_{f}} \quad R_{q} = \frac{(X|\delta_{q})}{\Delta X_{f}} \quad R_{q} = \frac{(X|\delta_{q})}{\Delta X_{f}}$

for achromatic separator $(x|\delta_v)_{tot}=0$

$$\begin{split} \mathbf{R}_{\mathbf{q}} &= \begin{array}{c} \frac{\left[\left(\delta_{\mathcal{V}}|\delta_{q}\right) - 1 + \left(\delta_{\mathcal{V}}|\delta_{\mathcal{V}}\right)\right] / \left(\delta_{\mathcal{V}}|\delta_{\mathcal{V}}\right)}{\sqrt{1/\mathbf{R}^{2}} + \left(\delta_{\mathbf{v},\text{str}}/\left(\delta_{\mathcal{V}}|\delta_{\mathcal{V}}\right)\right)^{2}} \end{array} = \mathbf{R} \begin{array}{c} \frac{\left[\left(\delta_{\mathcal{V}}|\delta_{q}\right) - 1 + \left(\delta_{\mathcal{V}}|\delta_{\mathcal{V}}\right)\right]}{\left(\delta_{\mathcal{V}}|\delta_{\mathcal{V}}\right)}\right]}{\sqrt{1/\mathbf{R}^{2}} + \left(\delta_{\mathbf{v},\text{str}}/\left(\delta_{\mathcal{V}}|\delta_{\mathcal{V}}\right)\right)} \end{array} = \mathbf{R} \begin{array}{c} \frac{\left[\left(\delta_{\mathcal{V}}|\delta_{m}\right) + 1 - \left(\delta_{\mathcal{V}}|\delta_{\mathcal{V}}\right)\right]}{\left(\delta_{\mathcal{V}}|\delta_{\mathcal{V}}\right)}\right]}{\left(\delta_{\mathcal{V}}|\delta_{\mathcal{V}}\right)} = \mathbf{R} \begin{array}{c} \frac{\left[\left(\delta_{\mathcal{V}}|\delta_{m}\right) + 1 - \left(\delta_{\mathcal{V}}|\delta_{\mathcal{V}}\right)\right]}{\left(\delta_{\mathcal{V}}|\delta_{\mathcal{V}}\right)}\right]}{\left(\delta_{\mathcal{V}}|\delta_{\mathcal{V}}\right)} \end{split}$$

Effect of degrader at different velocity



Separator Setting for ²¹²Pb (exp. S350 in 2009)



Details with Monte-Carlo

primary U beam at S1







MOCADI

Super-FRS



²³⁸U²⁸⁺ 0.4 – 2.7 GeV/u



All fission fragments after target, 1.5 GeV/u ²³⁸U --> 4 g/cm² ¹²C





Separation only by $B\rho$





Separation after pre-separator ($B\rho$ - ΔE - $B\rho$)



Separation

Separation after main separator ($B\rho$ - ΔE - $B\rho$) x ($B\rho$ - ΔE - $B\rho$)





For fission fragments separation is difficult, other beams more pure.

Q-State Effect on Separator

Separation of ²¹³Fr

optimal combination of q-states and stripper materials at different energy, FRS with achromatic S2 degrader in LISE⁺⁺



Change in q in degrader fools the Bp- Δ E-Bp separation

Fragments from Degrader BigRIPS Comissioning at RIKEN



BigRIPS Commissioning in 2006 Setting ⁸⁶Kr -> ⁷⁶Ni



Without degrader ⁷⁶Ni region not visible, with degrader still lots of lighter fragments.

BigRIPS ⁸⁶Kr -> ⁷⁶Ni Simulation with LISE⁺⁺



Ratio: total rate / 76 Ni rate (=S/N).

FRIB Separator



- Asymmetric pre-separator to compress B ρ spread by factor 3
- Pre and main separator in different planes

M. Hausmann, M. Portillo, C. Wilson et al.

Separator Comparison

			in LAB		normalized to Bp =18 Tm for m/q = 2.56			
		[Tm]	[mrad]	[mrad]	[%]	[mrad]	[mrad]	[%]
	stages	Brho	Α	В	dp/p	Α	В	dp/p
FRS	1	18	11	17	1.2	11.0	17.0	1.2
A1900-MSU	1	6	30	50	2.9	10.0	16.7	1.7
BigRIPS	2	9	55	40	3.0	27.5	20.0	2.2
Super-FRS	2	20	40	20	2.5	44.4	22.2	2.2
ARIS-FRIB	3	7	40	40	4.5	15.6	15.6	2.9

Forward focusing (symmetric in projectile system, Lorentz transformation) transverse: $A = p'_x / p_0$ longitudinal: $\Delta p_z / p = \gamma p'_z / p_0$

	[mm] x 0	$R=\frac{(x \delta)}{2(x x)x_0}$	[m ²] dipole area*
FRS	1.2	3385	2.36
A1900-MSU	1.0	1480	0.68
BigRIPS	0.5	2850	1.28
Super-FRS	1.0	2900	4.08
ARIS-FRIB	0.5	1300	0.96 * for 2 dipoles

Resolution limit by energy-loss straggling in half range thickness degrader, $R \sim 800$.

→ Use large area Super-FRS dipoles for high-resolution physics experiments

Summary

Most production reactions require large acceptance. Fragments more forward focused at higher energies. Hunt for very rare cases covered by much more intense species.

Simulation must include ion optics, atomic, nuclear physics. Many orders of magnitude rate difference, many paths to follow.

Simplified convolution method for fast rate calculation and optimisation, Monte Carlo for precision and details. \rightarrow LISE⁺⁺, MOCADI.

B ρ - Δ **E**-**B** ρ method works well, but is limited in resolution, has problems with many atomic charge states.

Problem of too high rates for particle tracking requires many stage separators like Super-FRS, BigRIPS, ARIS.