Dimensional crossover in ultracold Fermi gases

from Functional Renormalisation

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EMMI Workshop on Functional Methods in Strongly Correlated Systems Hirschegg, Austria

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Physics of ultracold atoms

Scale hierarchy & Hamiltonian





Effective Hamiltonian valid on scales $\gg \ell_{vdW}$:

$$\hat{H} = \int_{\vec{x}} \left[\hat{a}^{\dagger}(\vec{x}) \left(-\frac{\hbar \, \Delta}{2 \, M} + V_{\rm ext}(\vec{x}) \right) \hat{a}(\vec{x}) + g_{\Lambda} \, \hat{n}(\vec{x})^2 \right]$$

with $\hat{n}=\hat{a}^{\dagger}\,\hat{a}$ and $g_{\Lambda}=\frac{4\,\pi\,\hbar^{2}}{M}\,a$

Feshbach resonances

adapted from Boettcher et al. Nuclear Physics B (2012)



The BCS-BEC crossover in 3D



 $\begin{array}{l} k_F \,{=} \\ (3\,\pi^2\,n)^{1/3} \end{array}$

Randeria Nature (2010) BCS-BEC physics from Functional Renormalisation

Microscopic action

$$S = \int_X \left\{ \psi^* \left(\partial_\tau - \Delta - \mu \right) \psi + \phi^* \left(\partial_\tau - \frac{\Delta}{2} + \nu - 2\mu \right) \phi - h \left(\phi^* \, \psi_1 \, \psi_2 - \phi \, \psi_1^* \, \psi_2^* \right) \right\}$$

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- Truncation: derivative expansion
- Litim-type regulator (fermions: regularise around Fermi surface)

$$U(\rho) = \sum_{n=1}^{2} \frac{u_n}{n!} \left(\rho - \rho_0\right)^n - n_k \left(\mu - \mu_0\right) + \alpha \left(\mu - \mu_0\right) \left(\rho - \rho_0\right)$$

3D BCS-BEC crossover



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T = 0: (with $c = k_F a$)



T > 0:



3D BCS-BEC crossover

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: (with $c = k_F a$)



 $\xi_{\text{exp}} = 0.376(5), \quad \xi_{\text{calc}} \simeq 0.55$

T > 0:



Dimensional crossover

Why quasi-2D systems?

- promising materials: graphene, high $T_c\mbox{-superconductors},$ layered semiconductors
- pronounced influence of quantum fluctuations
- · experimental accessibility via highly anisotropic trapping potentials
- \cdot for insufficient anisotropy \longrightarrow dimensional crossover

\longrightarrow disentangle dimensionality from many-body physics



Zürn et al. PRL (2015)

Boundary conditions

- delimit z-direction by potential well of length L

$$V_{\rm box}(z) = \begin{cases} 0 & 0 \leq z \leq L \\ \infty & {\rm else} \end{cases}$$

- impose **periodic** boundary conditions: $\Psi = \{\psi, \phi\}$

$$\Psi(\tau,x,y,z=0)=\Psi(\tau,x,y,z=L)$$

 \rightarrow quantisation of momentum in *z*-direction:

$$q_z \to k_n = \frac{2 \, \pi \, n}{L}, \qquad n \in \mathbb{N}$$

• spatial Matsubara sum

$$\int \frac{d^d q}{(2\pi)^d} = \frac{1}{L} \sum_{k_n} \int \frac{d^{d-1}q}{(2\pi)^{d-1}}$$

Litim-type regulator (as before) with $\vec{q}=\hat{\vec{q}}+q_z\rightarrow\hat{\vec{q}}+k_n$

Litim-type regulator (as before) with $\vec{q} = \hat{\vec{q}} + q_z \rightarrow \hat{\vec{q}} + k_n$

Idea:

- + initialise RG flow at UV scale $k=\Lambda$ where $\Gamma_{\Lambda}=S_{3D}$
- *L* introduces new length scale to 3D system
- following RG flow successively integrates out 3rd dimension

• correct 3D-limit



• correct 3D-limit



 $a_{\rm 2D}^{\rm (pbc)} = L\,\exp\left(-\frac{1}{2}\,\frac{L}{a_{\rm 3D}}\right)$

 qualitatively correct behaviour of equation of state (except very small confinements)



Finite temperature & phase diagram



[cf. A. Fischer and M. Parish (2014) for mean-field study (BCS-limit)]

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Finite temperature & phase diagram



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- \cdot 3D phase diagram reproduced for large L
- increased T_c/T_F around $\ln(k_F a_{\rm 2D}) \sim 1$

Comparison to experiment



experimental data from Ries et al. PRL (2015)

Determination of the density

$\mu\text{-dependence}$ of initial conditions



μ -dependence of initial conditions

Free energy density, f, normalised in the vacuum,

$$\partial_t f(T,\mu) := \left(\frac{\partial_t \Gamma_k[\phi_{\mathrm{EoS},k};T,\mu]}{\mathcal{V}_T} - \frac{\partial_t \Gamma_k[\phi_{\mathrm{EoS},k};0,0]}{\mathcal{V}_0} \right) \,,$$

has polynomial growth with k,

$$\partial_t f(T,\mu) \to c_{d_f-2} k^{d_f} \hat{k}^{-2} + c_{d_f-4} k^{d_f} \hat{k}^{-4} + k^{d_f} O(\hat{k}^{-6}) \,.$$

with $\hat{k} = k/\sqrt{\mu}$ (non-rel.), $\hat{k} = k/\mu$ (rel).

 \longrightarrow initial conditions for couplings g_i with scaling dimension $d_{g_i} \geq 2$ $\mu\text{-dependent},$

i.e. fermionic mass parameter (including the density itself)

• For non-relativistic case, equation of density (from its flow)

$$\partial_t n = \frac{1}{\operatorname{Vol}} \frac{d\,\partial_t \Gamma_k}{d\mu} \to c_{n,3} k^3 + c_{n,1} \mu\, k + O(\hat{k}^{-1})\,,$$

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- while flow of second μ -derivative tends to zero

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• Write flow of density as

$$\partial_t n_k = \frac{d\partial_t \Gamma_k}{d\mu} = \left. \partial_\mu \right|_{\vec{g}} \partial_t \Gamma_k + \frac{dg_i}{d\mu} \partial_{g_i} \partial_t \Gamma_k \,.$$

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- Each partial $\mu\text{-}$ and $dg_i/d\mu\,\partial_{g_i}\text{-}\mathrm{derivative}$ lowers effective $k\text{-}\mathrm{dimension}$ by two.
- The coefficients $g_i^{(1)}=dg_i/d\mu$ follow from their flow
- · Can be iteratively extended to higher derivatives



Summary and Outlook

Conclusion

- dimensional crossover in Fermi gas with FRG
- qualitatively comparable to experiments
- preliminary results for density

Outlook

- iterative calculation of density for whole crossover (3d, 2d-3d)
- employ harmonic trapping potential
- explore scale anomaly in dimensional crossover

Thank you for your attention!

Phase diagram from experiment



Ries et al. PRL (2015)