

Hadronic Correlation and Spectral Functions at Finite Temperature

Olaf Kaczmarek



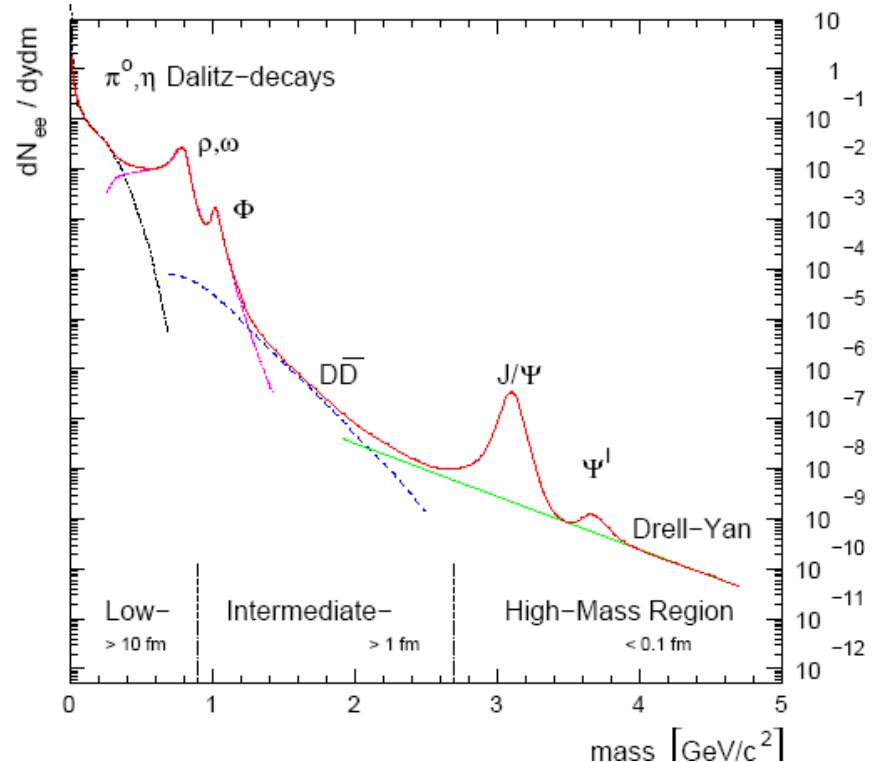
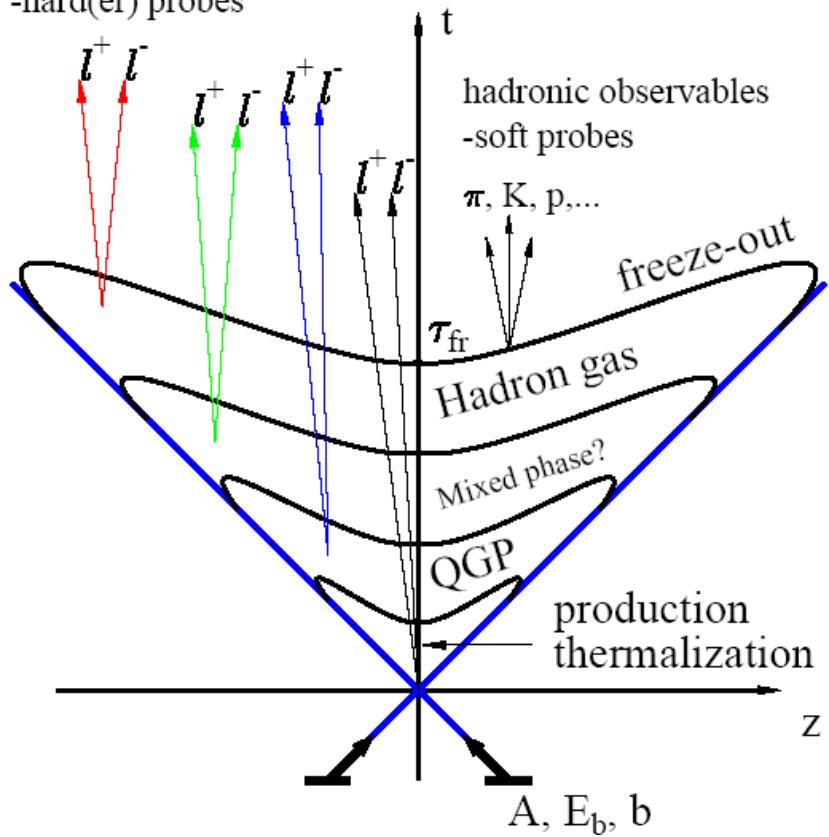
Universität Bielefeld

FAIR Lattice QCD Days
GSI Darmstadt, November 23, 2009

Hard Probes in Heavy Ion Collisions - Dileptons

electromagnetic observables

-hard(er) probes



Dileptonrate directly related to vector spectral function:

$$\frac{dW}{d\omega d^3p} = \frac{5\alpha^2}{27\pi^2} \frac{1}{\omega^2(e^{\omega/T} - 1)} \sigma_V(\omega, \vec{p}, T)$$

Hard Probes in Heavy Ion Collisions – RHIC results

pp-data well understood by hadronic cocktail

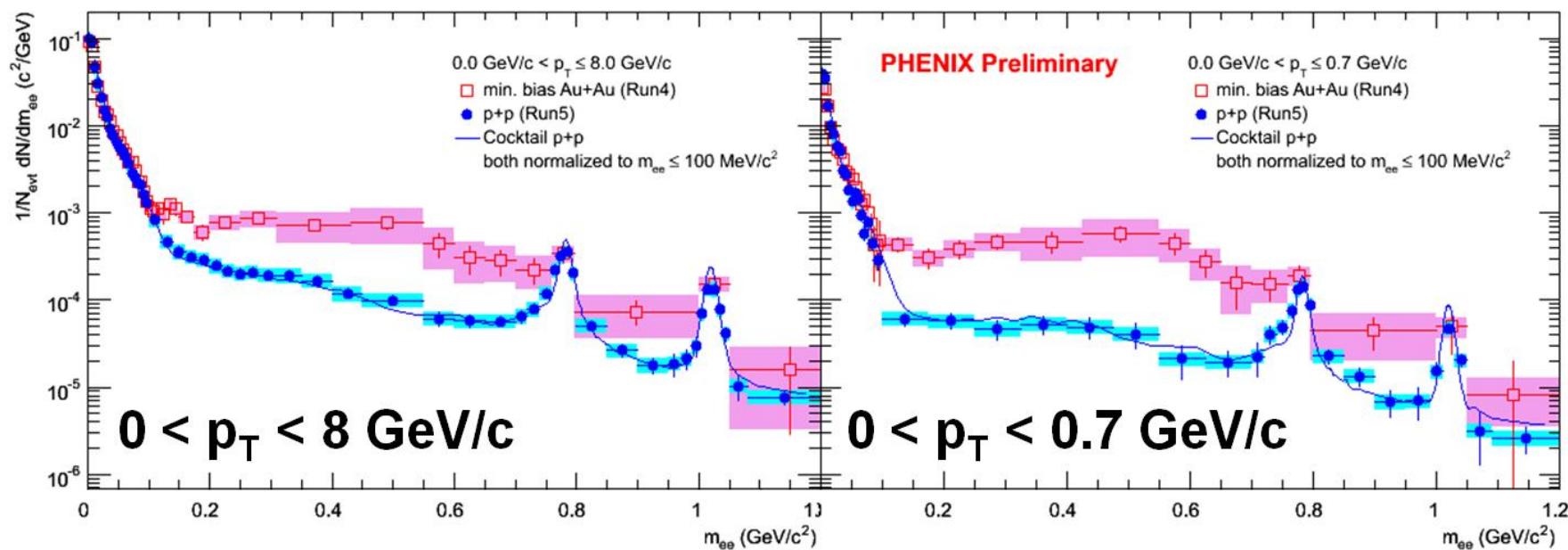
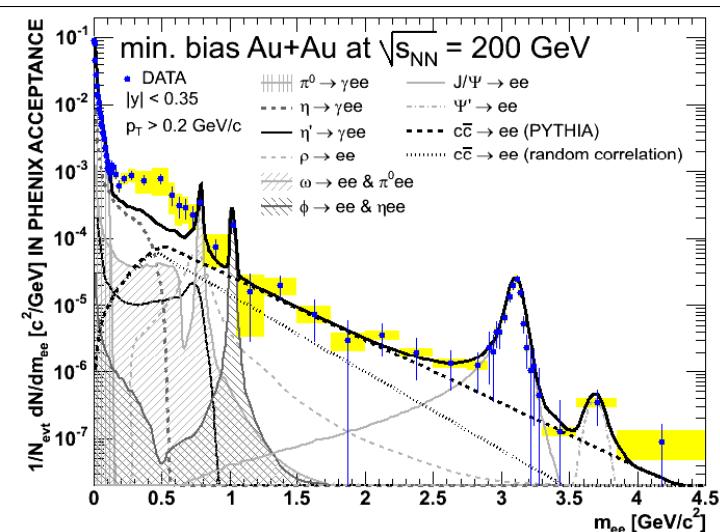
low invariant mass region <150 MeV similar in Au-Au

large enhancement between 150-750 MeV

indications for thermal effects!? Also at higher m_{ee} ?

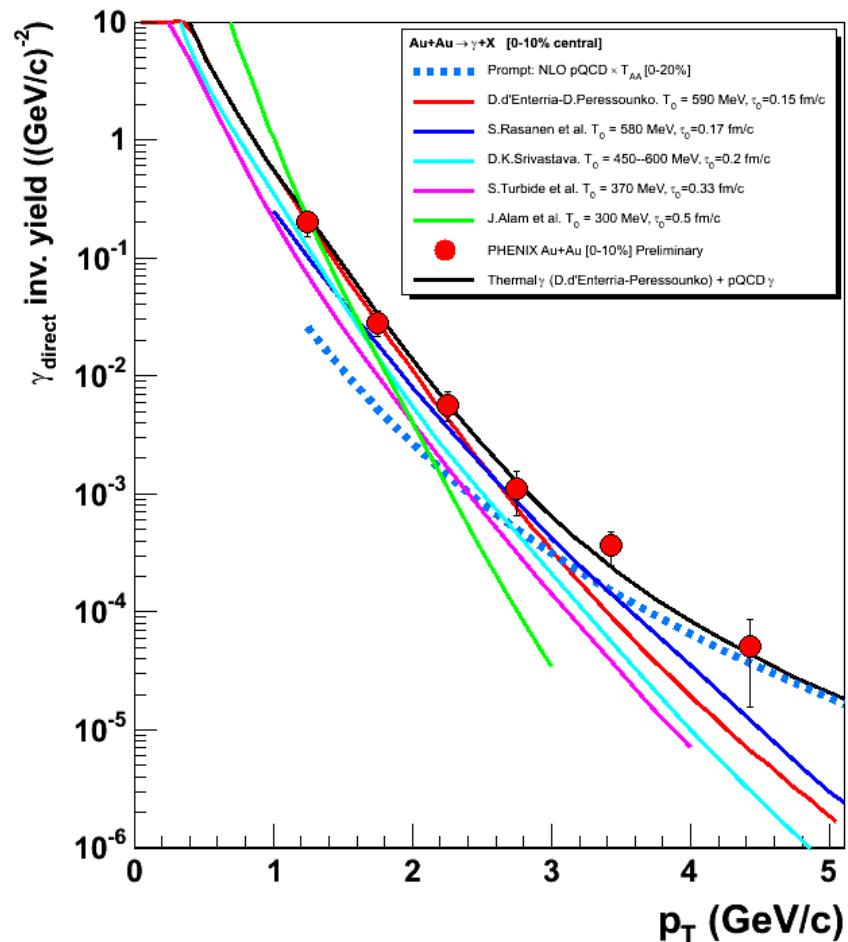
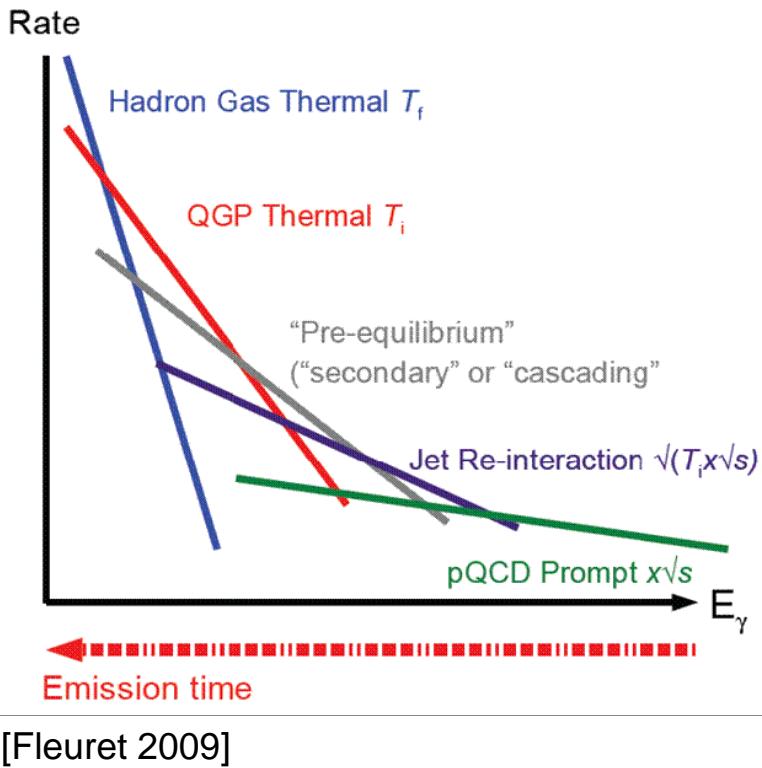
Need to understand the contribution from QGP!

→ spectral functions from lattice QCD



Hard Probes in Heavy Ion Collisions - Photons

Direct and fragmentation photon relative contribution



Photonrate directly related to vector spectral function (at finite momentum):

$$\omega \frac{dN_\gamma}{d^4x d^3q} = \frac{5\alpha}{6\pi^2} \frac{1}{e^{\omega/T} - 1} \sigma_V(\omega = |\vec{p}|, T)$$

Transport Coefficients

The small ω limit of σ_V is related to transport coefficients (Kubo-Formulas)

Light quark sector → **electrical conductivity**:

$$\sigma_{el} = \lim_{\omega \rightarrow 0} \frac{\sigma_V^{ii}(\omega, \vec{p} = 0, T)}{6\omega}$$

Heavy quark sector → **heavy quark diffusion constant**:

$$D = \frac{1}{3\chi^{00}} \lim_{\omega \rightarrow 0} \frac{\sigma_V^{ii}(\omega, \vec{p} = 0, T)}{\omega}$$

Free (non-interacting) spectral function [Karsch et al. 03, Aarts et al. 05]

$$\sigma_H = \frac{N_c}{8\pi^2} \Theta(\omega^2 - 4m^2) \textcolor{blue}{\omega^2} \tanh\left(\frac{\omega}{4T}\right)$$

$$\times \sqrt{1 - \left(\frac{2m}{\omega}\right)^2} \left[a_H + \left(\frac{2m}{\omega}\right)^2 b_H \right]$$

$$+ \textcolor{magenta}{\frac{N_c}{3} \frac{T^2}{2} f_H \omega \delta(\omega)}$$

with interactions:

$$\delta(\omega) \rightarrow \frac{1}{\pi} \frac{\eta}{\omega^2 + \eta^2} \quad \begin{array}{l} \text{[Petreczky+Teaney 06]} \\ \text{Aarts et al. 05] } \end{array}$$

Introduction

Hadronic correlators for light quarks ($m_q=0$)

Screening masses in the thermodynamic and continuum limit

Temporal correlators vs. free correlators

Spectral functions and Dilepton rates

Charmonium hadronic correlators ($m_q=m_c$)

Screening Masses below and above T_c

Temporal correlators vs. free correlators

Spectral functions below and above T_c

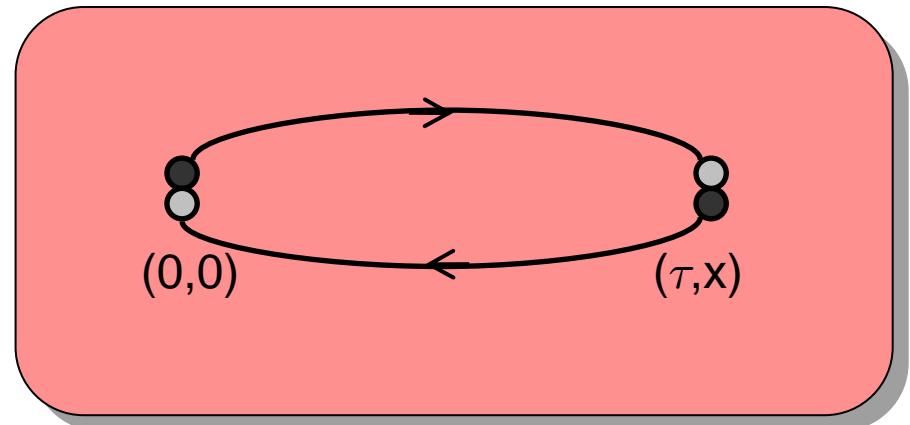
Temporal correlators vs. reconstructed correlators

Zero mode contributions

Hadronic correlators – Lattice setup

Thermal hadronic correlation functions

$$J_H = \bar{q}(\tau, \vec{x}) \Gamma_H q(\tau, \vec{x})$$



$$G_H(\tau, T, \vec{p}) = \sum_{\vec{x}} e^{-i\vec{p}\vec{x}} \langle J_H(0, 0) J_H^\dagger(\tau, \vec{x}) \rangle$$

O(a)-improved Clover improved fermionic action

on large quenched lattice configurations up to $128^3 \times 16/32/48$

includes all the relevant physics (in the quenched limit)

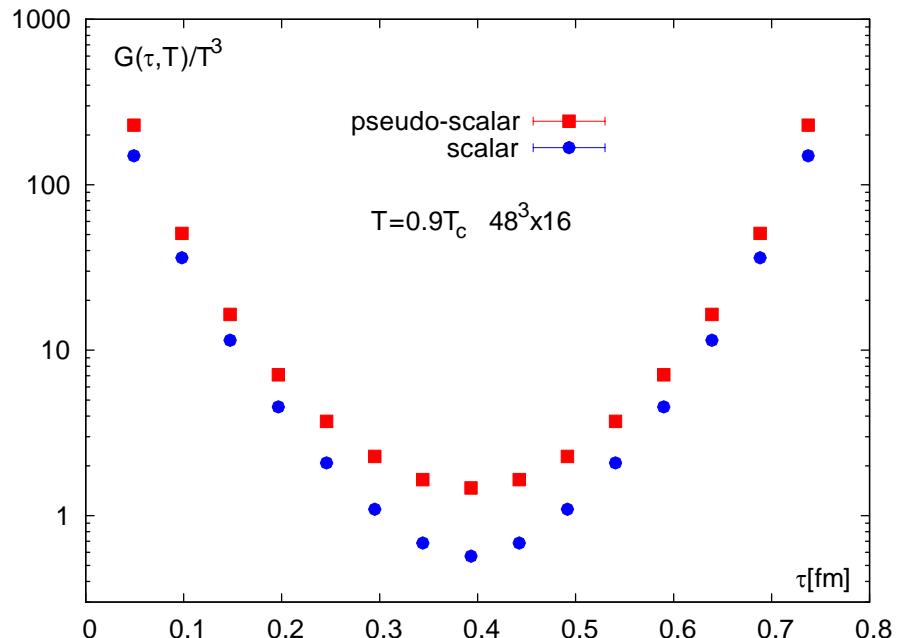
how to extract it?

directly from the correlators?

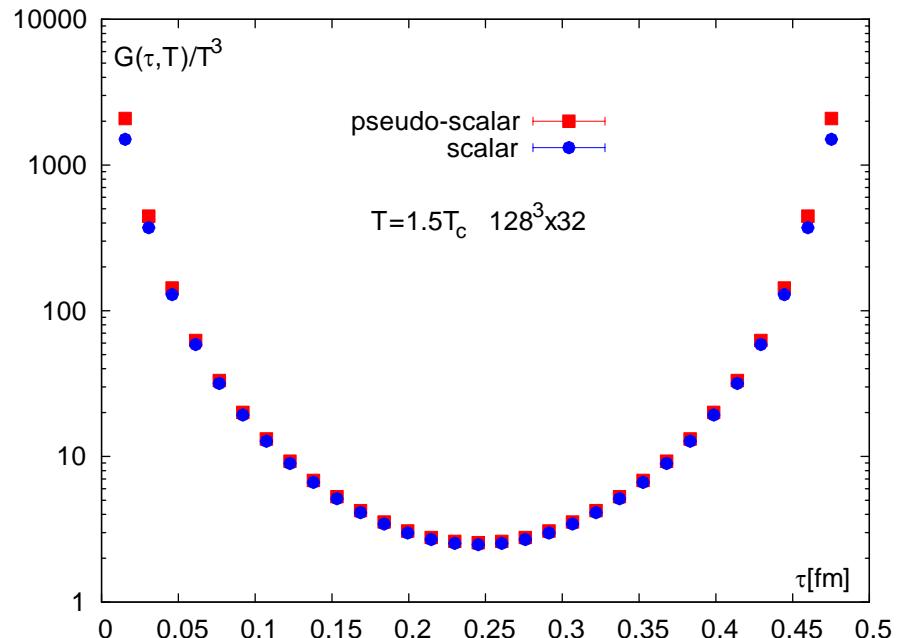
spectral functions using MEM?

Light Quark Correlators – Chiral Symmetry

Temporal Correlators:



splitting below T_c due to
chiral and axial U(1) symmetry breaking

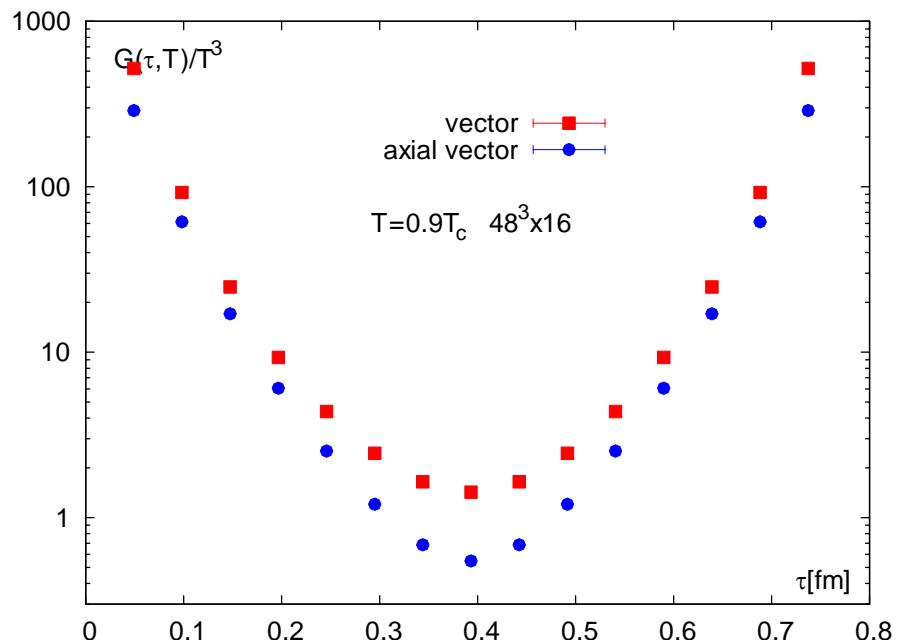


degenerate states at $1.5 T_c$
symmetry restauration above T_c

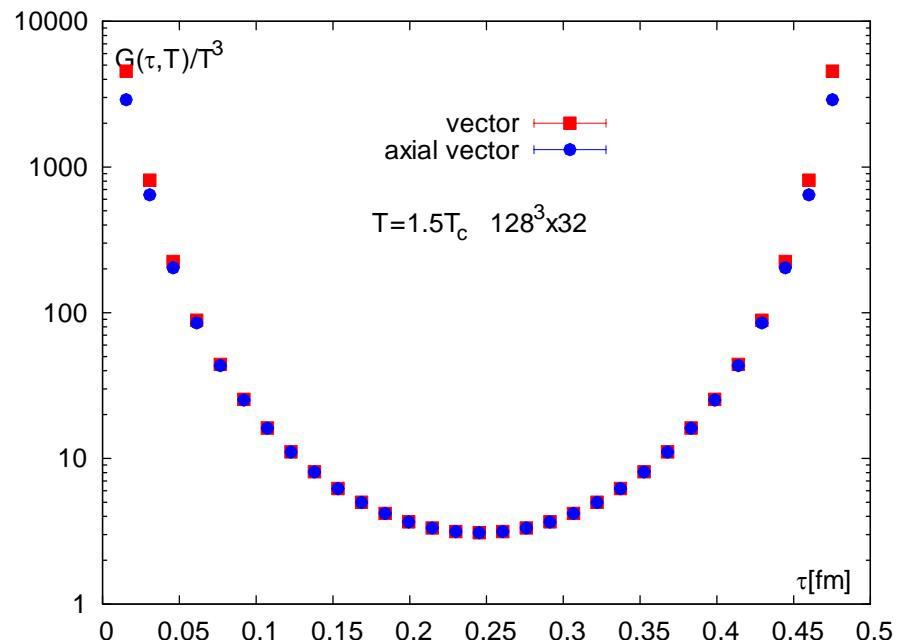
deviations at small τ due to
different cut-off effects

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Light Quark Screening Masses – Definition

use **Spatial Correlators**

$$G_H(z, T, \vec{p}_\perp) = \sum_{\tau, \vec{x}_\perp} e^{-i\vec{p}_\perp \cdot \vec{x}_\perp} \langle J_H(0, 0) J_H^\dagger(\tau, \vec{x}) \rangle$$

correlation function depends on the same spectral density,
but the relation is more involved

$$G_H(z) = \int_{-\infty}^{+\infty} \frac{dp_z}{2\pi} e^{ip_z z} \int_{-\infty}^{+\infty} dp_0 \frac{\sigma_H(p_0, \vec{0}_\perp, p_z)}{p_0} \xrightarrow[z \rightarrow \infty]{} \text{Ampl.} \times \exp(-m_{\text{screen}} z)$$

however, $m_{\text{screen}}(T) \neq m_{\text{pole}}(T)$ in general :

look for zeros of $G^{-1}(p) = p_0^2 + \vec{p}^2 + m_0^2 + \Pi(p_0, \vec{p}, T)$

$$\vec{p} = 0 : \quad -p_0^2 = m_0^2 + \Pi(p_0, \vec{0}, T) = (m_{\text{pole}}(T))^2$$

$$p_0 = 0 : \quad -\vec{p}^2 = m_0^2 + \Pi(0, \vec{p}, T) = (m_{\text{screen}}(T))^2$$

$$\Rightarrow m_{\text{screen}}(T) = \frac{m_{\text{pole}}(T)}{A(T)}$$

Light Quark Screening Masses – Thermodynamic Limit, $V \rightarrow \infty$

large collection of lattices ranging from $16^3 \times 8$ to $128^3 \times 16$

allowing for thermodynamic limit $V \rightarrow \infty$ at $N_t=8, 12$ and 16

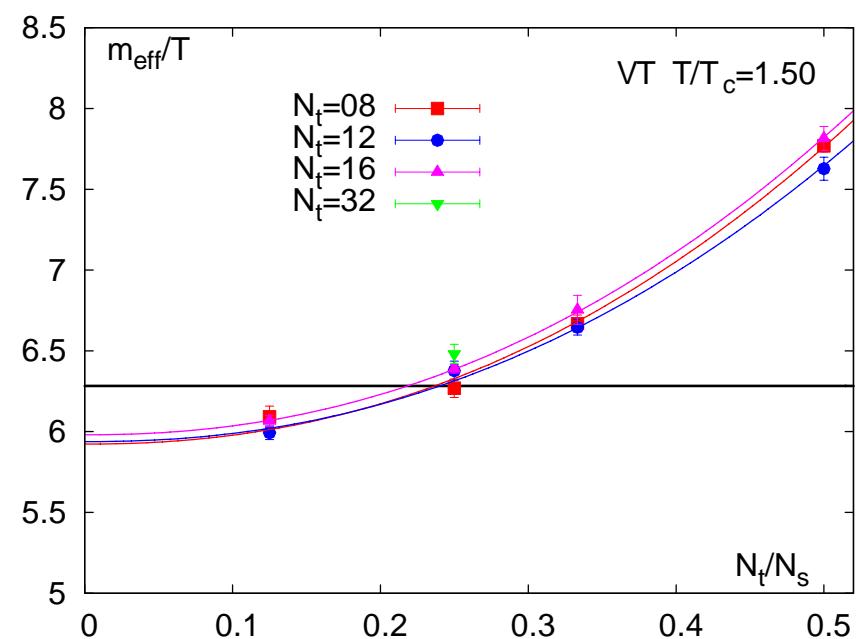
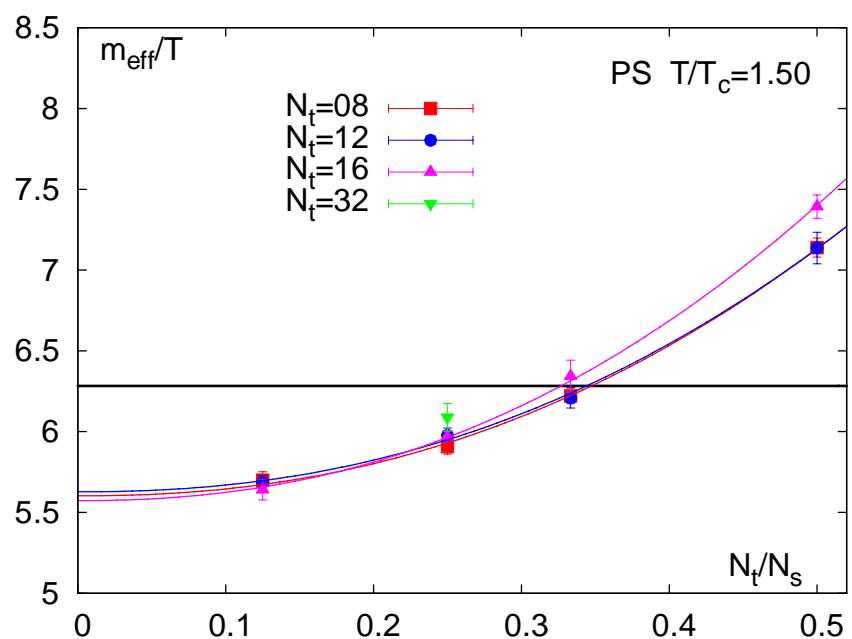
$$T \leq T_c : m_{\text{screen}}(L, a) = m_{\text{screen}}(a) \left[1 + \gamma_V \left(\frac{N_\tau}{N_\sigma} \right)^3 \right]$$

$$T > T_c : m_{\text{screen}}(L, a) = m_{\text{screen}}(a) \left[1 + \gamma_V \left(\frac{N_\tau}{N_\sigma} \right)^p \right]$$

$$T = \infty : m_{\text{screen}}(L, a) = m_{\text{screen}}(a) \left[1 + \gamma_V \left(\frac{N_\tau}{N_\sigma} \right)^1 \right]$$

combined fit:

p	PS	V
$1.5 T_c$	$2.22(10)$	$2.18(13)$
$3.0 T_c$	$2.06(7)$	$2.05(11)$



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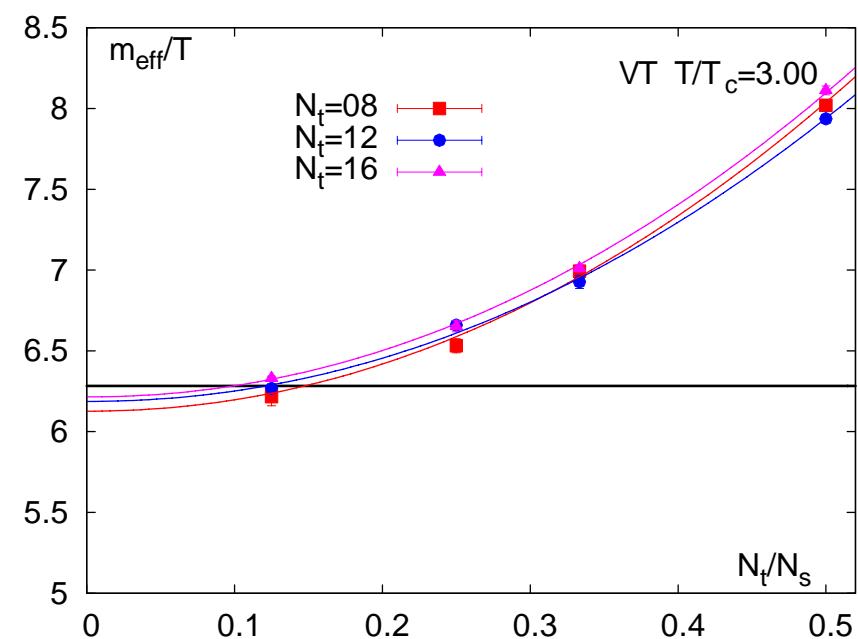
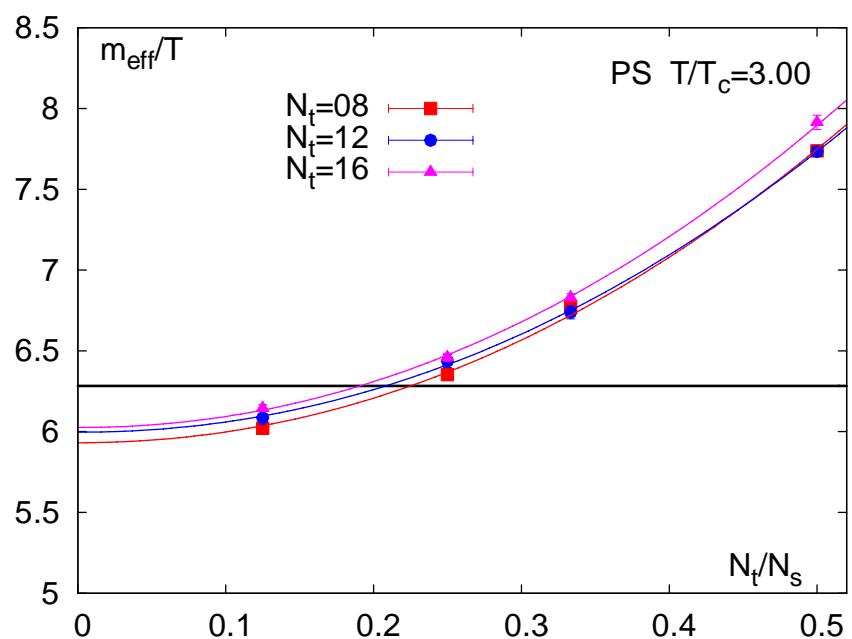
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Light Quark Screening Masses – Continuum Limit

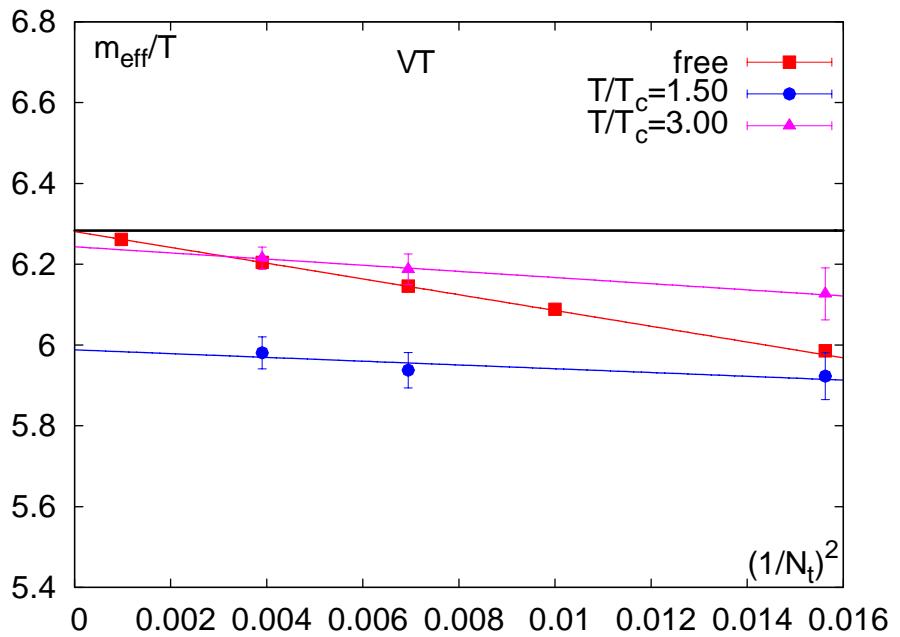
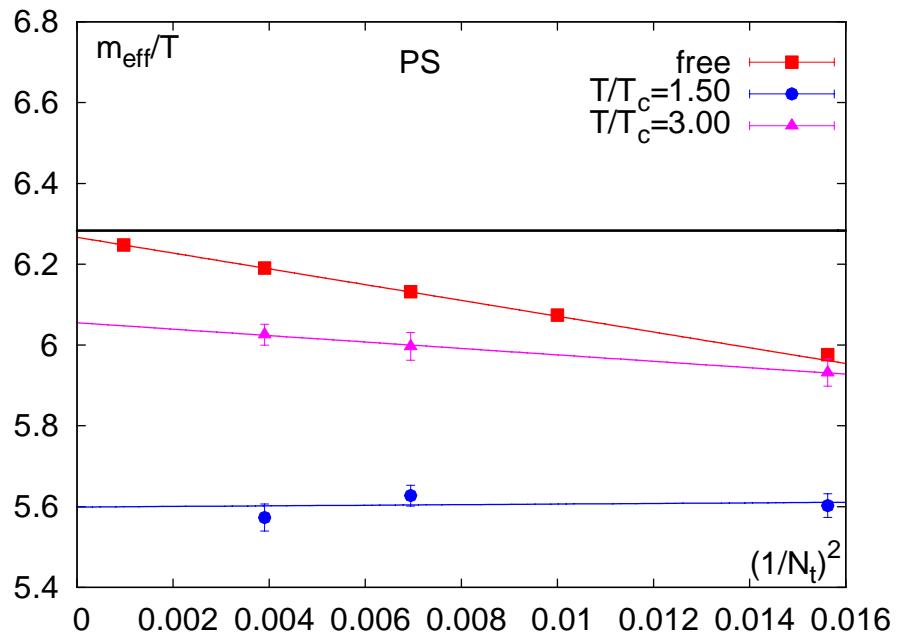
Continuum Limit

lattice spacing $a \rightarrow 0$

Non-perturbatively improved action

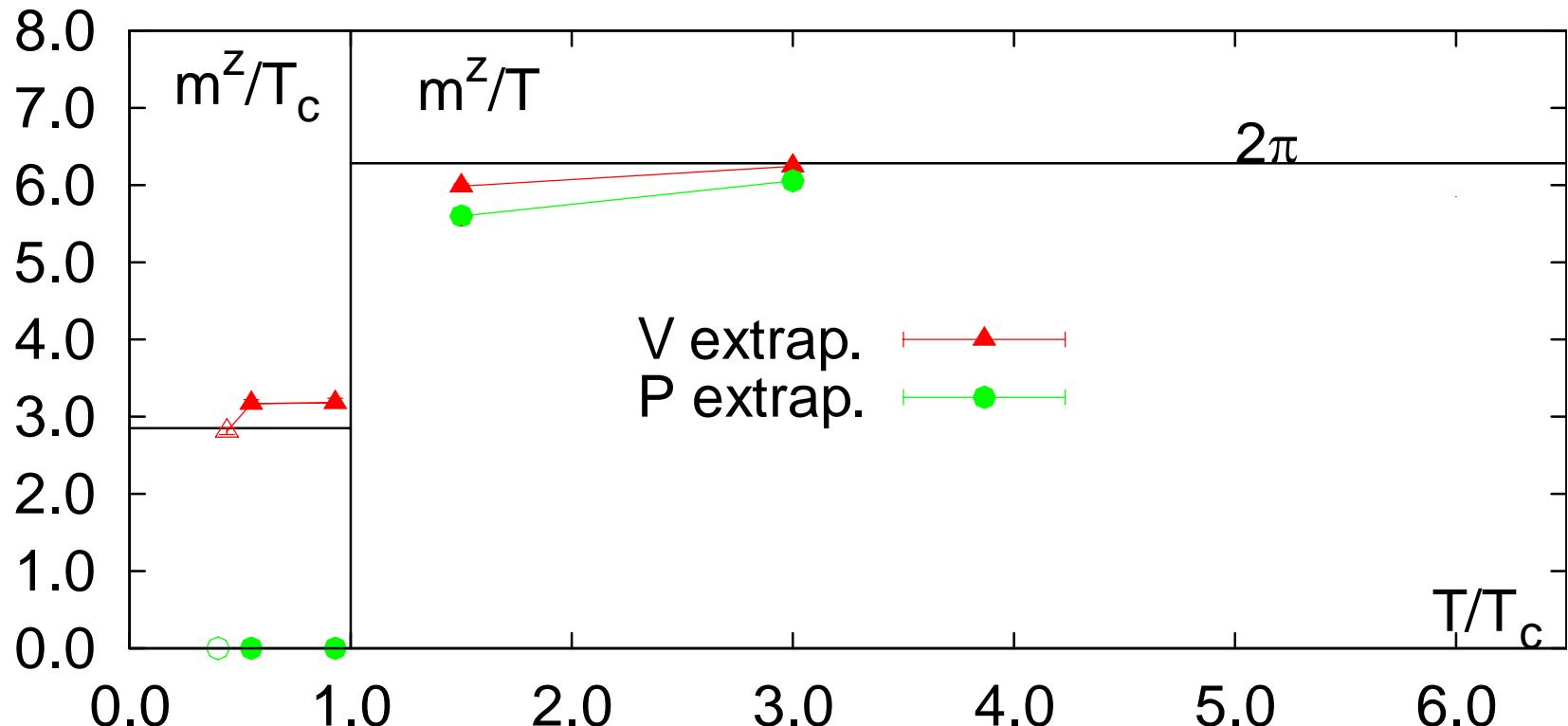
\Rightarrow discretization errors $O(a^2)$

$$T = \infty : \quad \frac{m^z(a)}{T} = \frac{m^z}{T} - \lambda \left(\frac{1}{N_\tau} \right)^2$$



Thermodynamic and Continuum Limit:

$V \rightarrow \infty$ and $a \rightarrow 0$



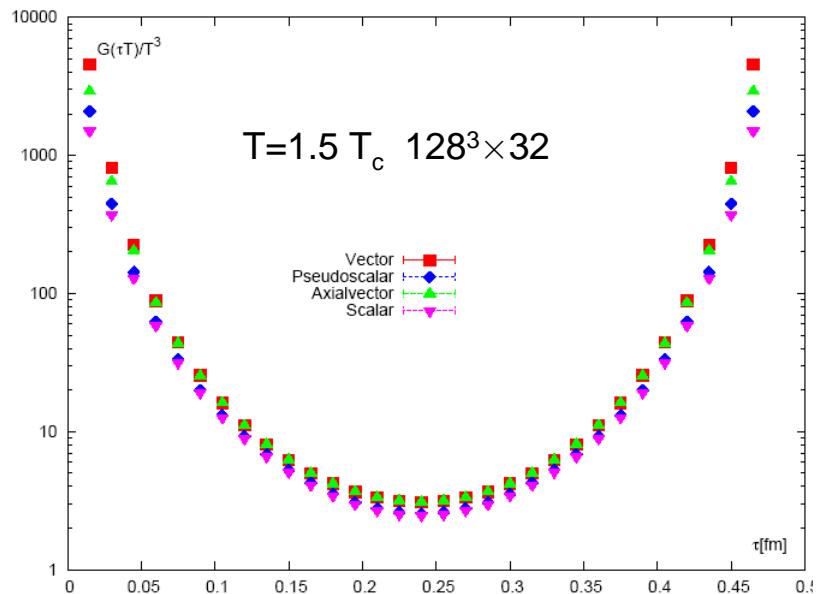
weak temperature dependence below T_c

data still below free limit ($2\pi T$) at $3 T_c$, vector closer to free case

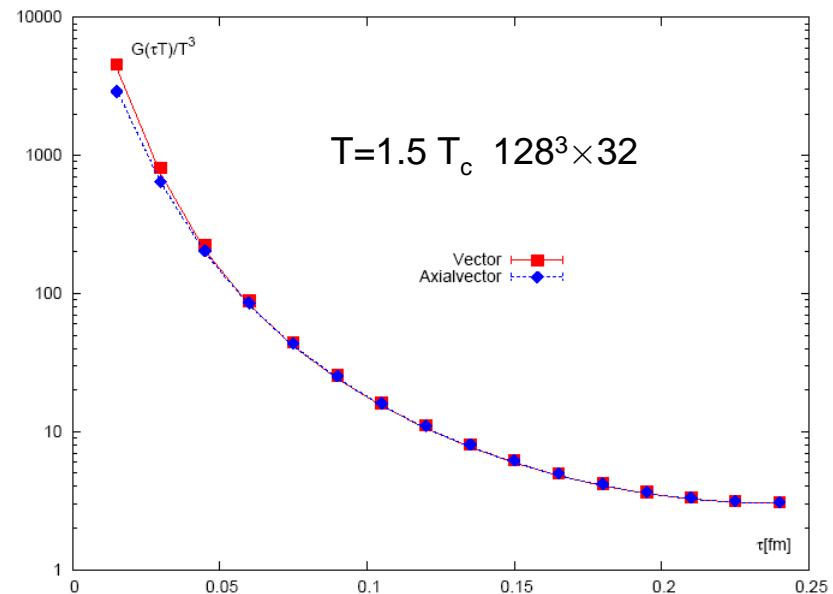
perturbative limit reached from above [Laine, Vepsäläinen]

→ need higher temperatures to verify this

Light Quark Correlators – Temporal direction



$T=1.5 T_c$ $128^3 \times 32$



$T=1.5 T_c$ $128^3 \times 32$

comparison with free (non-interacting) high temperature lattice correlator

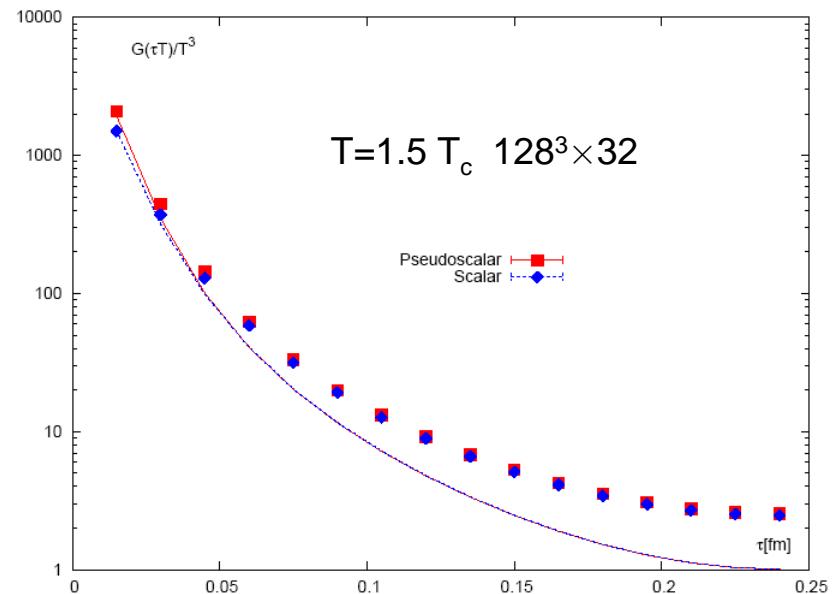
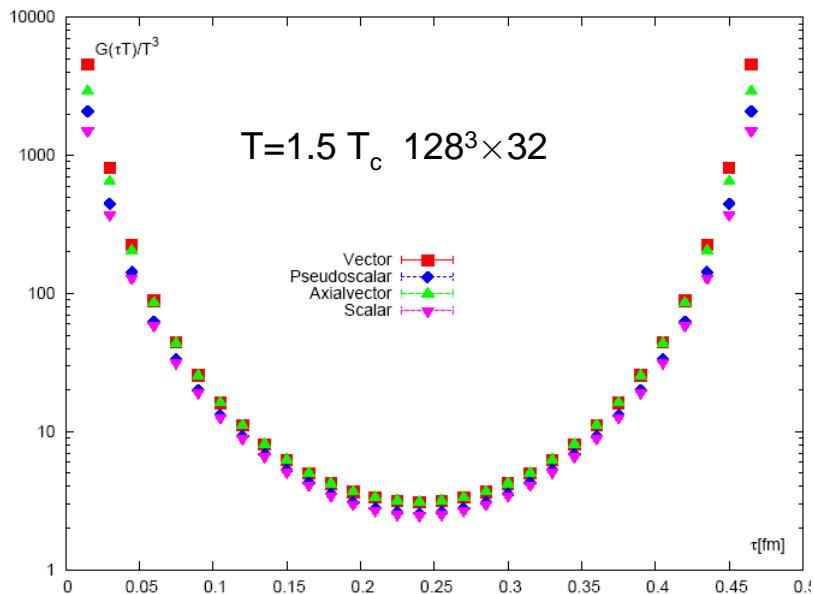
vector and axial-vector close to free case

cut-off effects are well described by free lattice correlator

this explains the difference at small τ

in the following only vector and pseudo-scalar are discussed

Light Quark Correlators – Temporal direction



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vector and axial-vector close to free case

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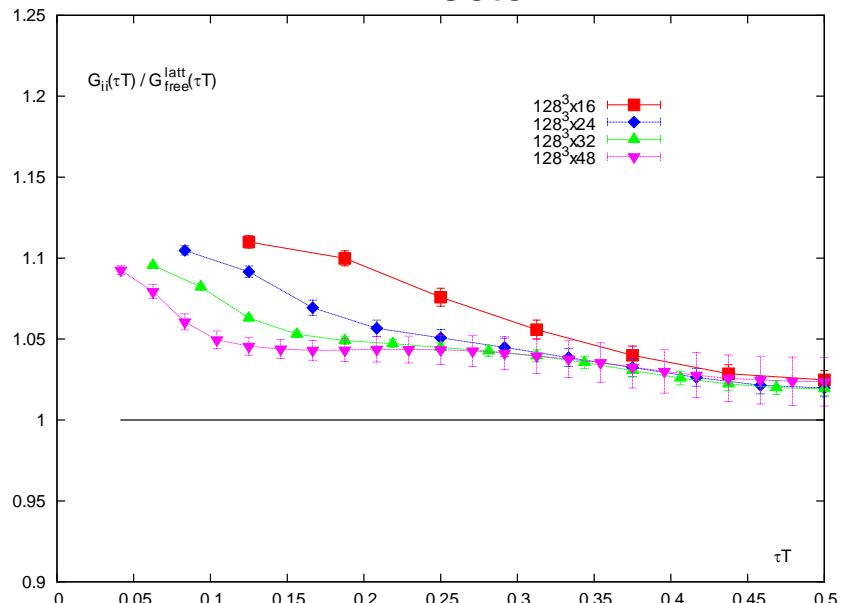
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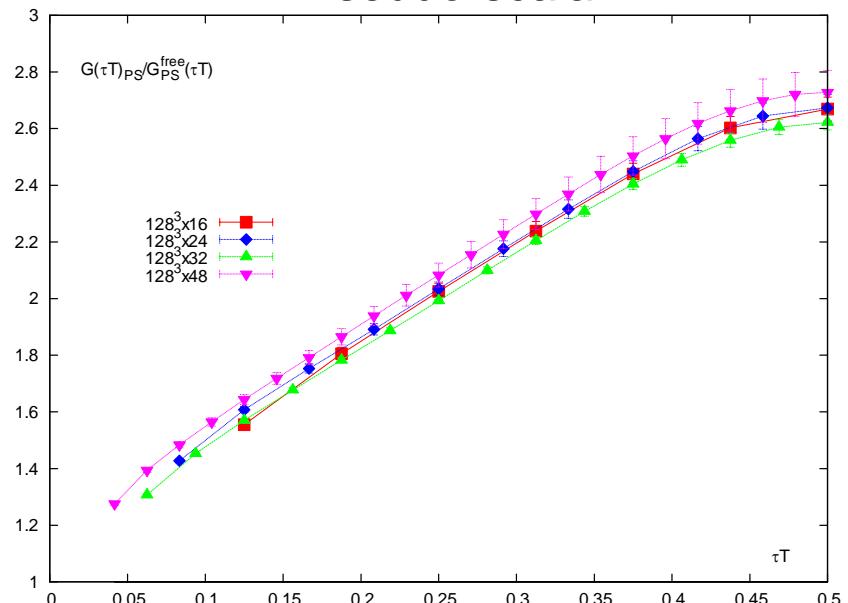
still strong correlations in scalar and pseudo-scalar channel!

Light Quark Correlators vs Free Correlators

Vector



Pseudo-scalar



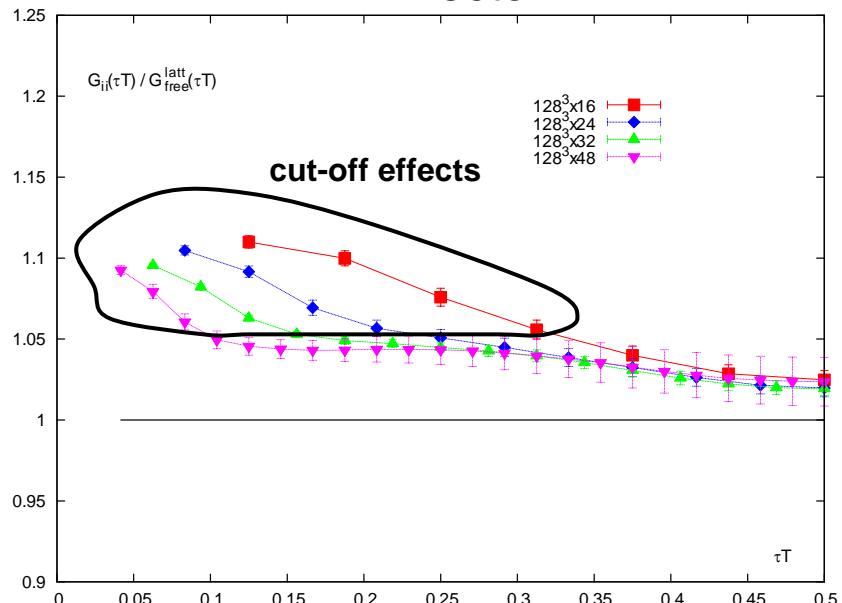
comparison with free (non-interacting) high temperature lattice correlator

vector and axial-vector close to free case

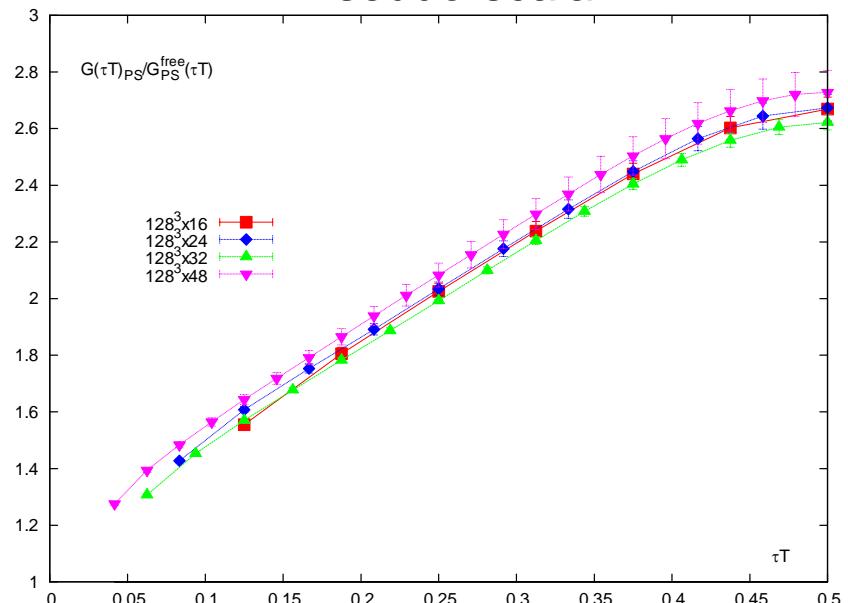
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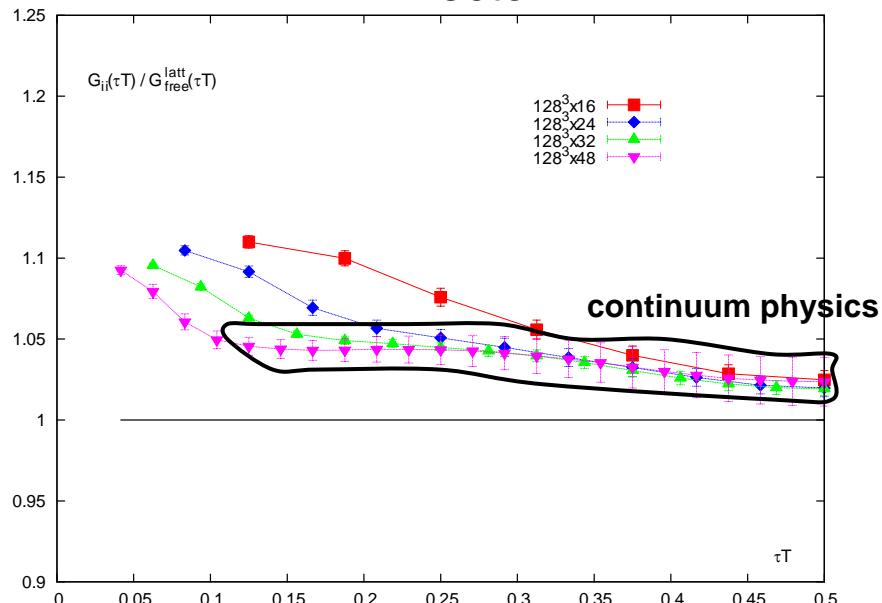
vector and axial-vector close to free case

still strong correlations in scalar and pseudo-scalar channel

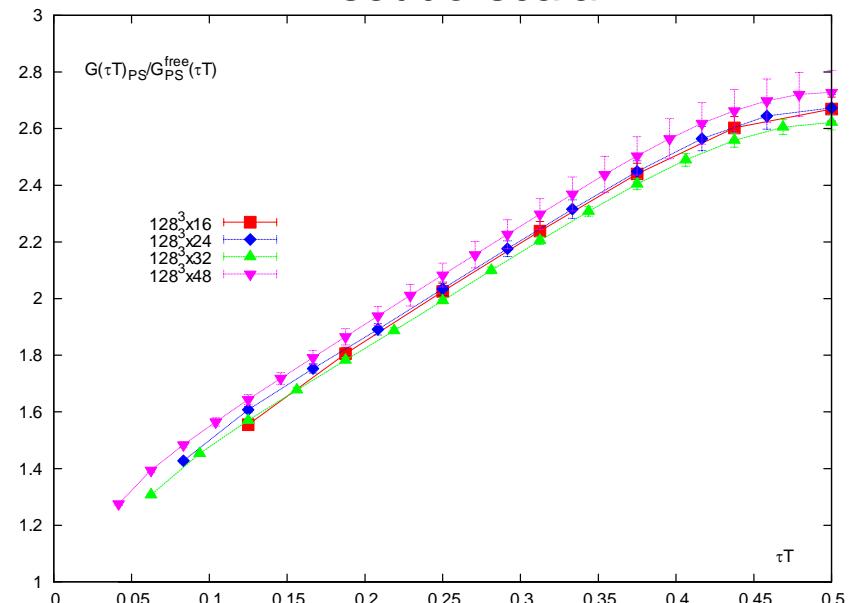
first 4-5 distances still dominated by cut-off effects

Light Quark Correlators vs Free Correlators

Vector



Pseudo-scalar



comparison with free (non-interacting) high temperature lattice correlator

vector and axial-vector close to free case

still strong correlations in scalar and pseudo-scalar channel

first 5 points still dominated by cut-off effects!

Nt=32 and 48 needed to extract continuum physics!

Spectral Functions – Maximum Entropy Method

How to obtain continuous spectral function $\sigma(\omega, T)$

from discrete (and small) number of correlators?

$$G(\tau, T) = \int_0^\infty d\omega K(\tau, \omega, T) \sigma(\omega, T)$$

$$K(\tau, \omega, T) = \frac{\cosh(\omega(\tau - \frac{1}{2T}))}{\sinh(\frac{\omega}{2T})}$$

Best method on the market: Maximum Entropy Method (MEM)

based on Bayesian theorem [Asakawa et al. 01] → most probable spectral function
properly renormalized correlators as input

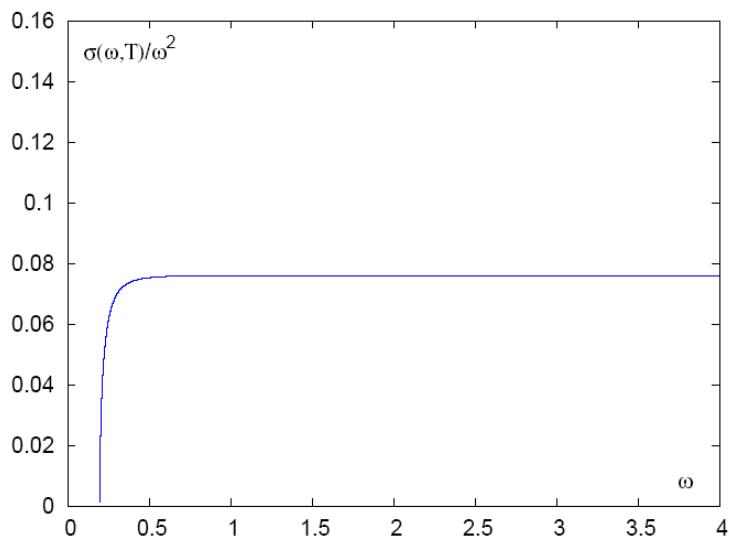
non-perturbative renormalization constants for vector [Lüscher et al. 1997]

TI perturbative renormalization constants for pseudo-scalar
prior knowledge needed as input → default model $m(\omega)$
result should be independent of default model ← usually not the case

MEM – Free spectral function

Free (non-interacting) spectral function [Karsch et al. 03, Aarts et al. 05]

$$\begin{aligned}\sigma_H = & \frac{N_c}{8\pi^2} \Theta(\omega^2 - 4m^2) \omega^2 \tanh\left(\frac{\omega}{4T}\right) \\ & \times \sqrt{1 - \left(\frac{2m}{\omega}\right)^2} \left[a_H + \left(\frac{2m}{\omega}\right)^2 b_H \right] \\ & + \frac{N_c}{3} \frac{T^2}{2} f_H \omega \delta(\omega)\end{aligned}$$



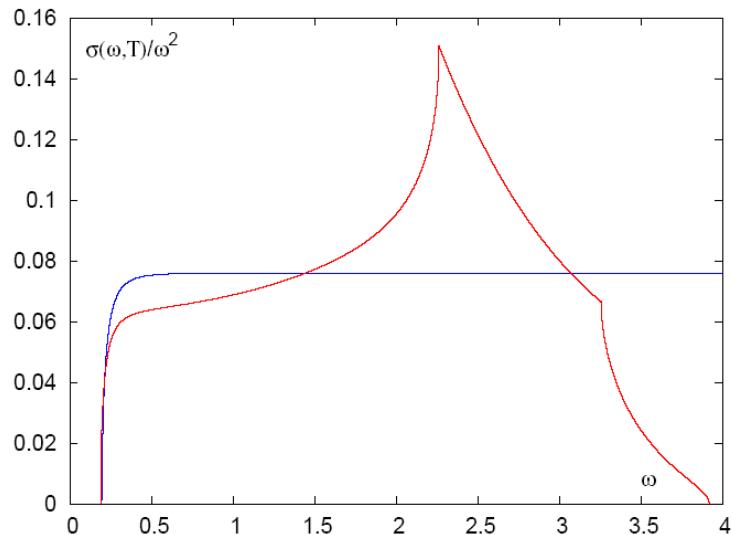
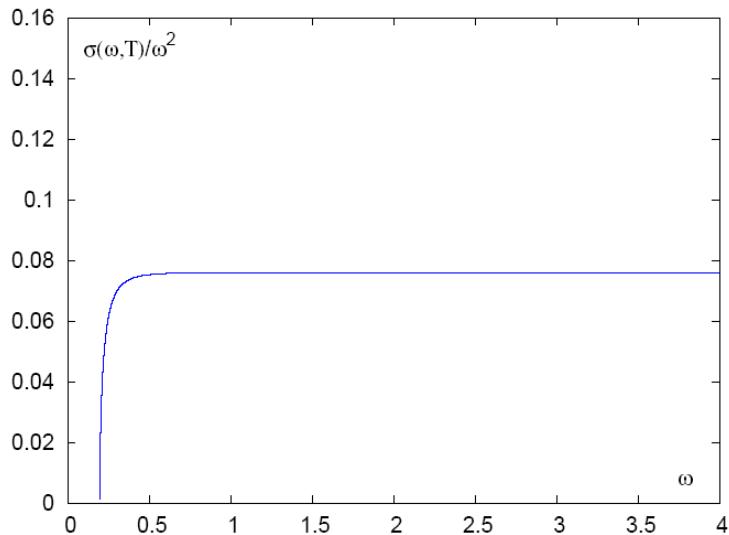
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Lattice cut-off effects

$$\omega_{max} = 2 \log(7 + ma)$$



MEM – Free spectral function

Free (non-interacting) spectral function [Karsch et al. 03, Aarts et al. 05]

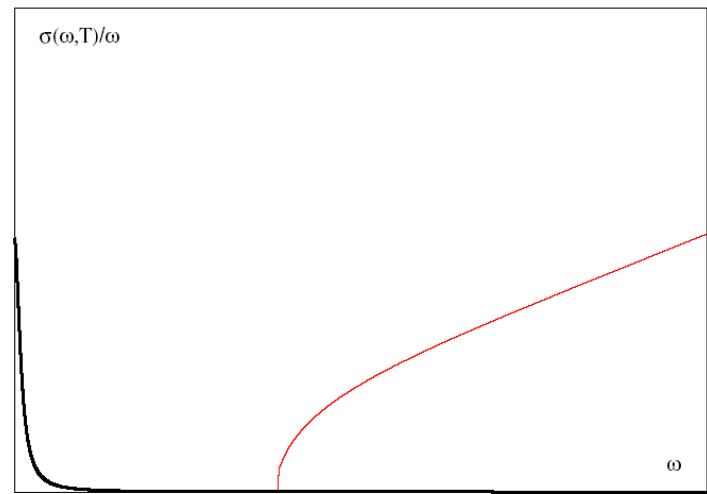
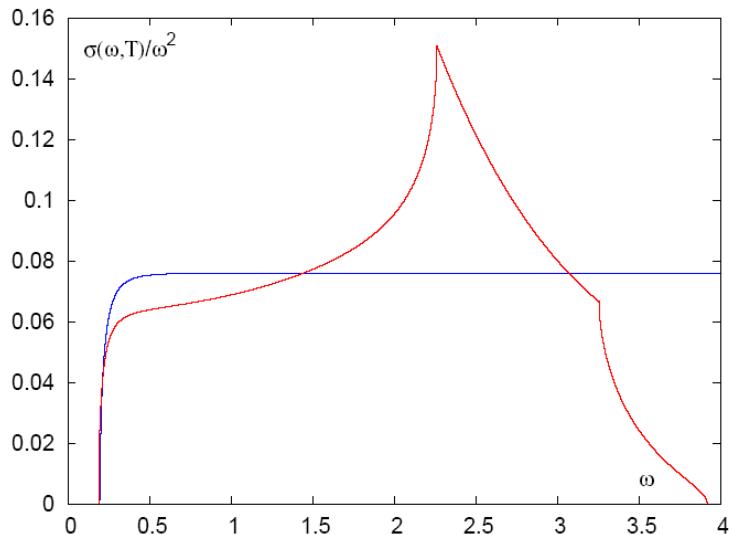
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zero mode contribution at $\omega \simeq 0$ [Umeda 07]

with interactions:

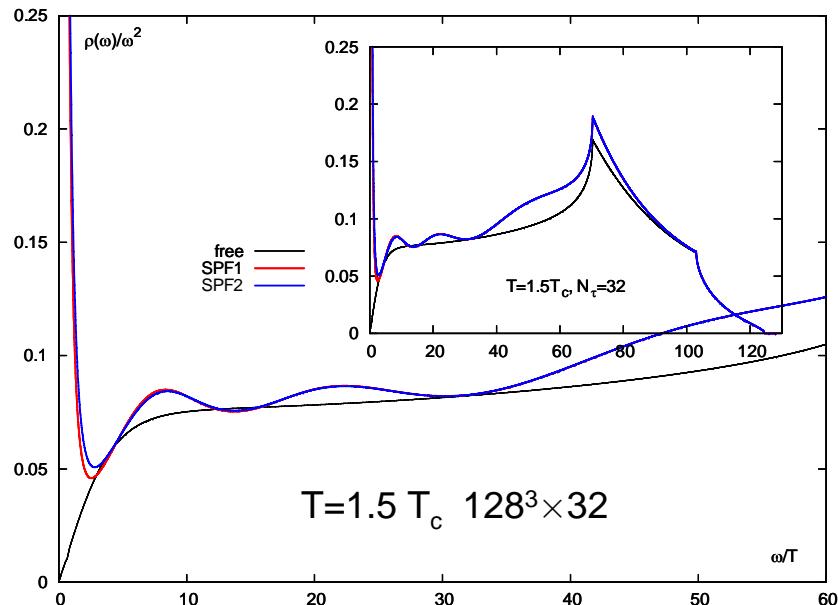
$$\delta(\omega) \rightarrow \frac{1}{\pi} \frac{\eta}{\omega^2 + \eta^2}$$

[Petreczky+Teaney 06
Aarts et al. 05]

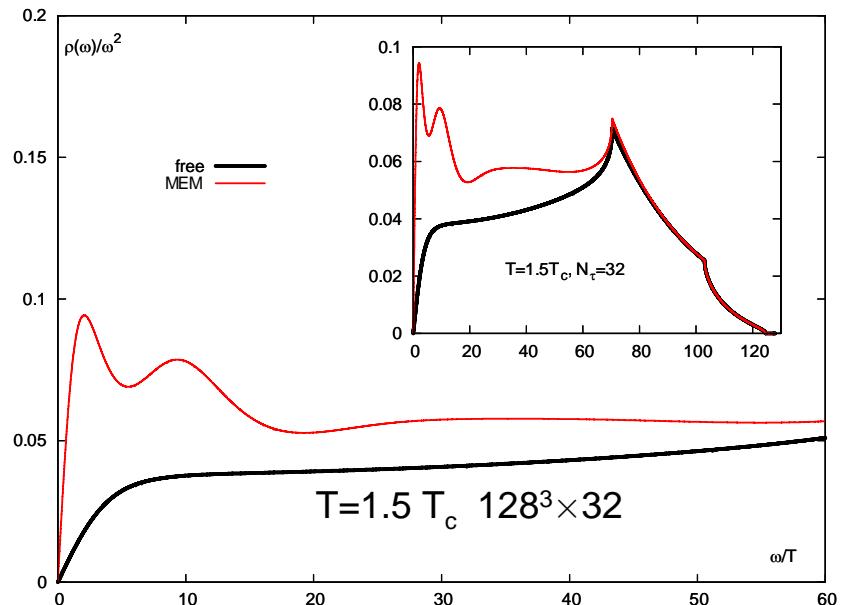


Spectral Function – Light Quark Sector

Vector



Pseudo-scalar



large ω behaviour well described by free lattice SPF

cut-Off effects are under control and well separated from physical interesting region

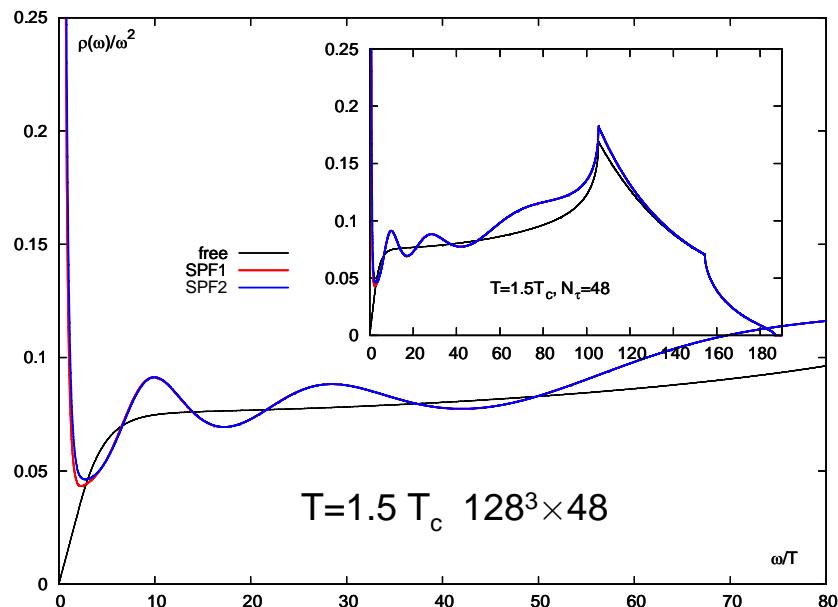
Vector SPF close to free case except at small ω

still large correlations in the Pseudoscalar sector

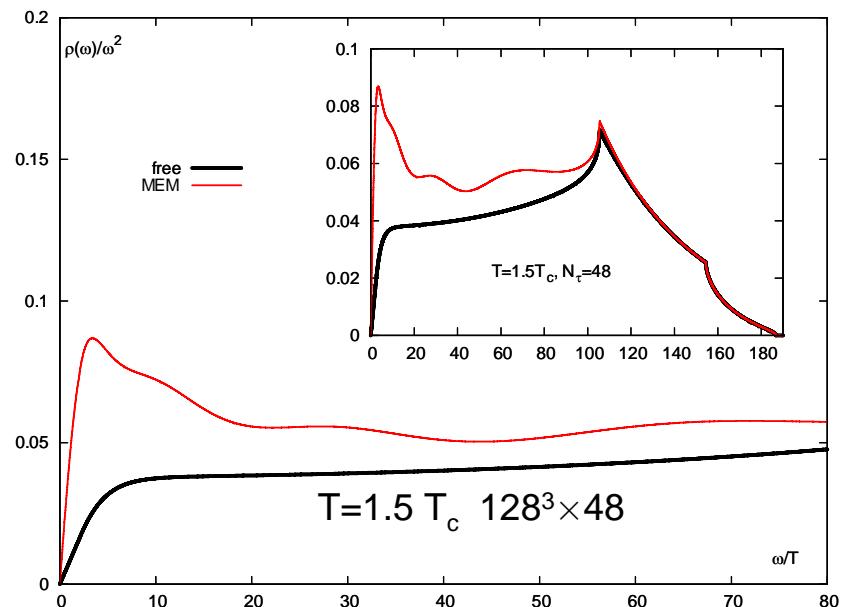
small ω region accessible → hope to extract transport properties
 → higher statistics needed

Spectral Function – Light Quark Sector

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Pseudo-scalar



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Vector SPF close to free case except at small ω

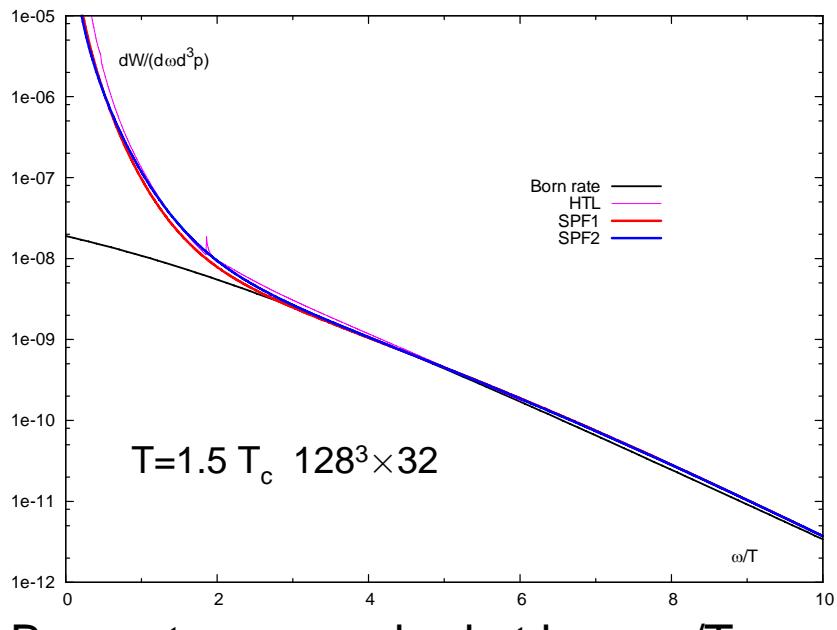
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Dilepton Rates

Dileptonrate directly related to vector spectral function:

$$\frac{dW}{d\omega d^3p} = \frac{5\alpha^2}{27\pi^2} \frac{1}{\omega^2(e^{\omega/T} - 1)} \sigma_V(\omega, \vec{p}, T)$$



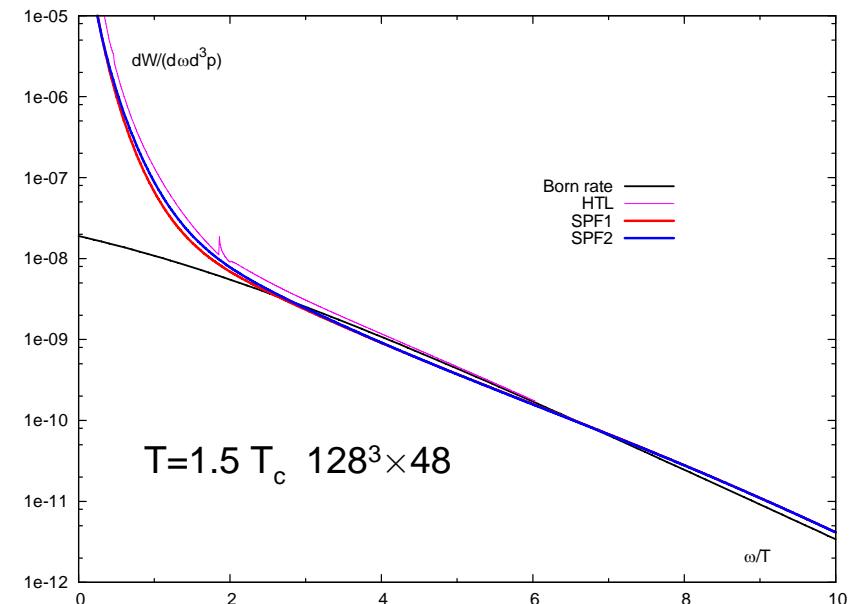
Born rate approached at large ω/T

consistent with HTL calculations at intermediate ω/T

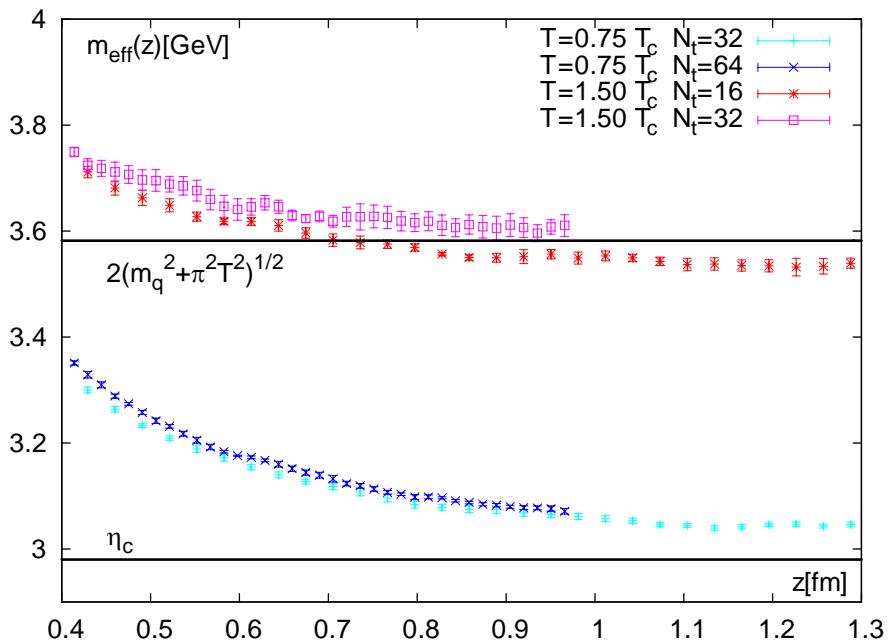
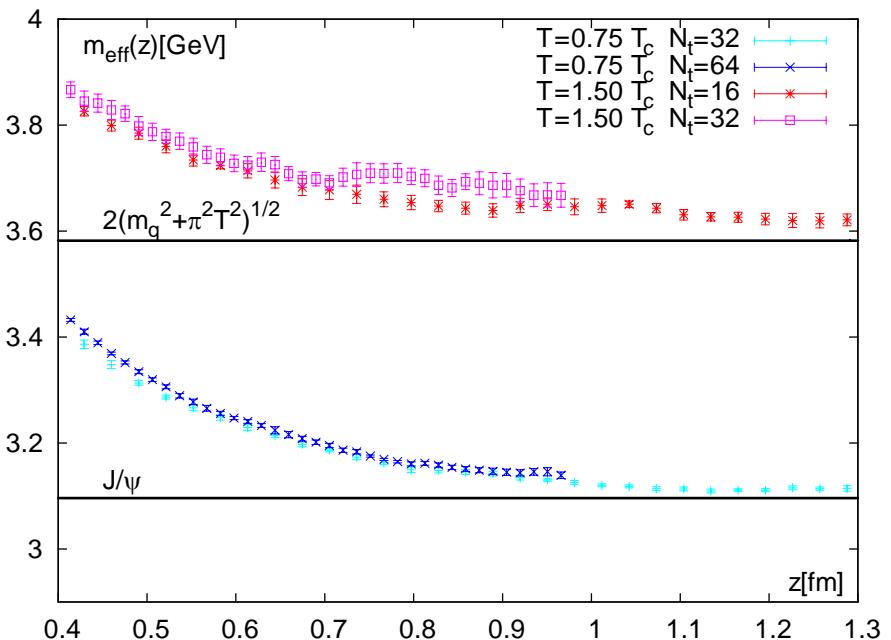
better behaved $\sim 1/\omega^2$ at small ω/T [Moore et al., Teaney, ...]

higher statistics needed to resolve details at small ω/T

consistent behavior at both lattice spacings \rightarrow continuum physics



Charmonium Correlators – Screening Masses



screening masses at $1.50 T_c$ already close to the free case

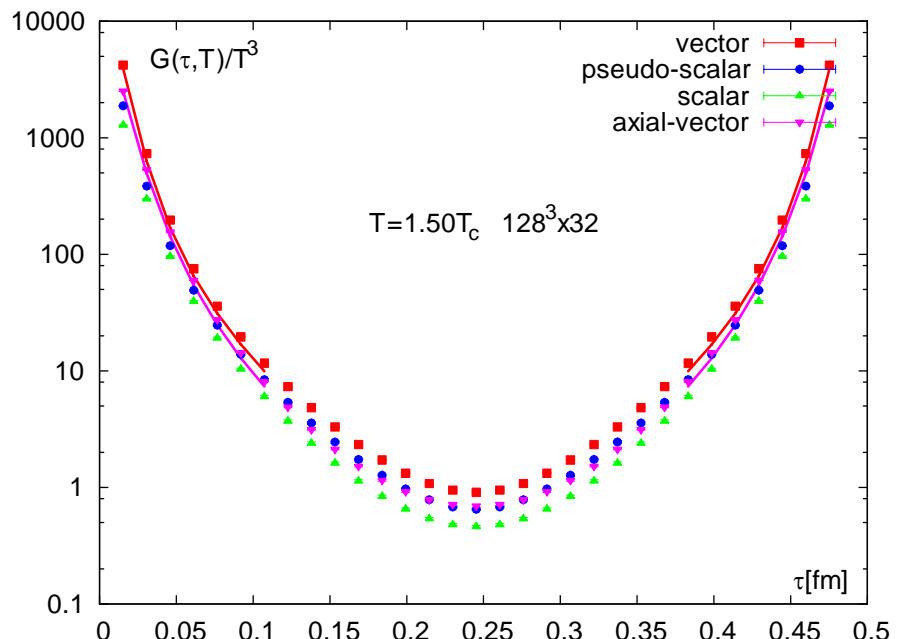
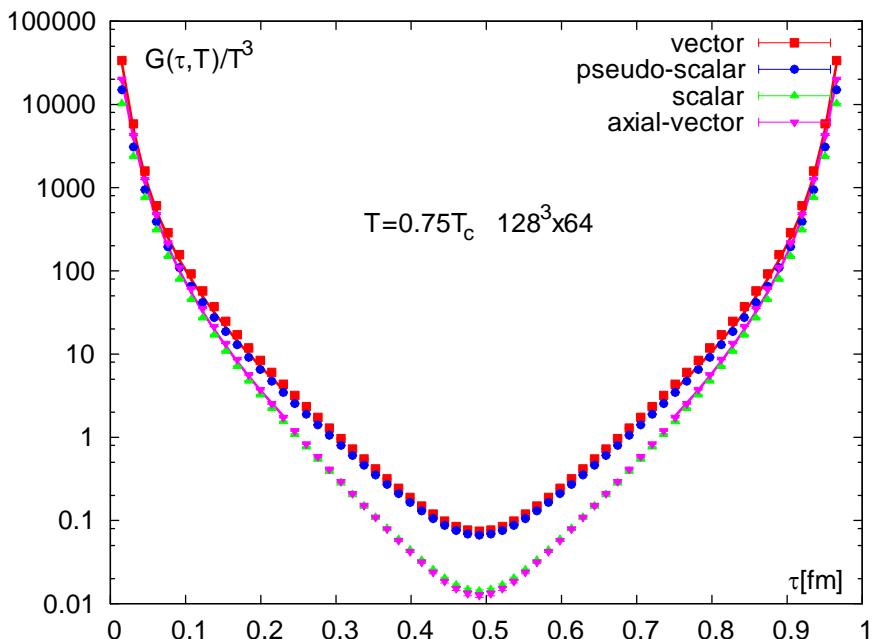
$$m_{\text{free}}^{\text{scr}}(T) = 2\sqrt{(\pi T)^2 + m_c^2}.$$

does this tell us anything about dissociation?

need to understand the momentum dependence of $G_H(\tau, T, p)$ and $\sigma(\omega, T, p)$ in detail

thermodynamic and continuum limit not performed yet

Charmonium Correlators – Temporal Correlators



non-degenerate states still at $1.50 T_c$

(almost) close to free correlators at (very) small separations

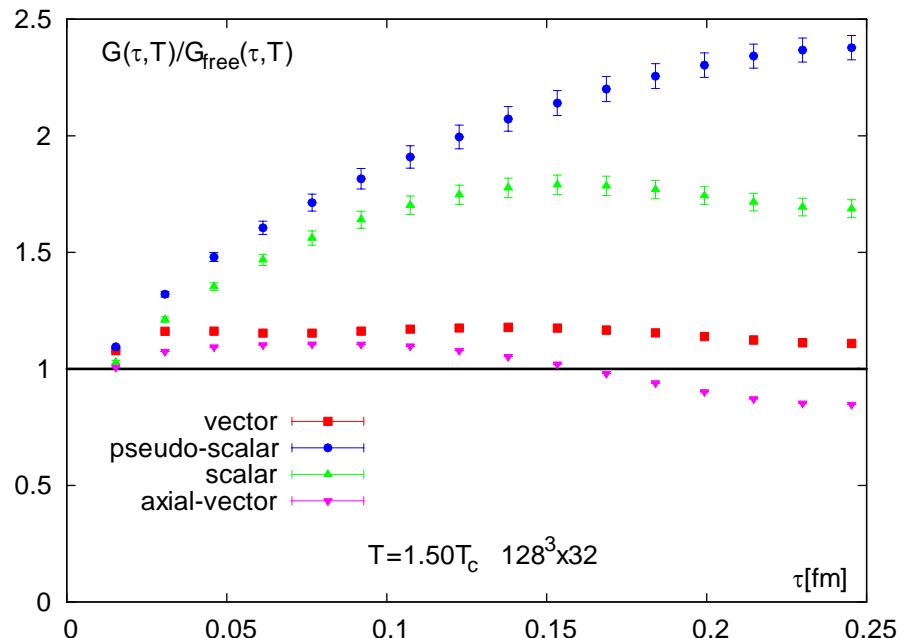
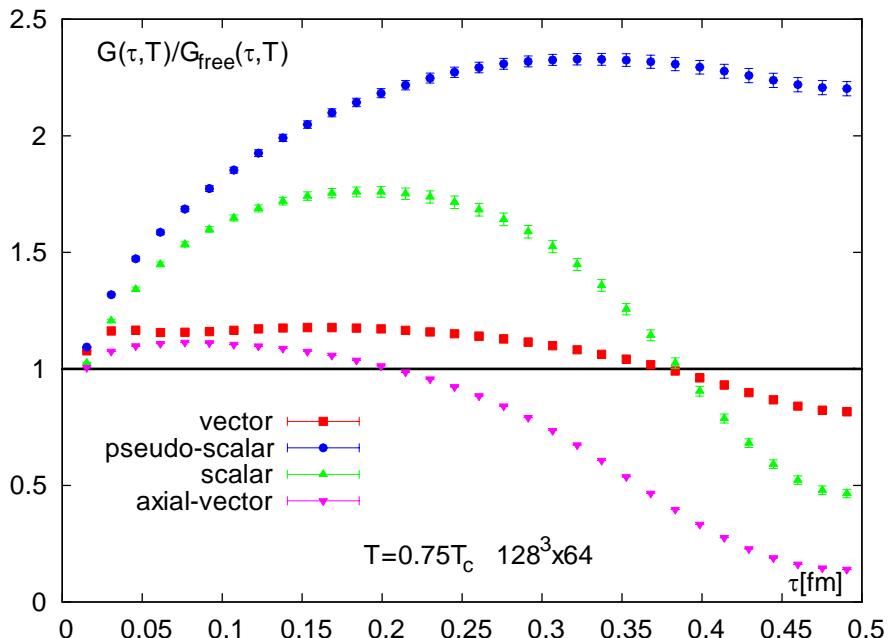
largest distance 0.25 fm due to compact temporal direction

only small distance regime (0.1-0.25 fm) relevant

for thermal effects

for bound state effects

Charmonium Correlators vs Free Correlators



non-degenerate states still at $1.50 T_c$

(almost) close to free correlators at (very) small separations

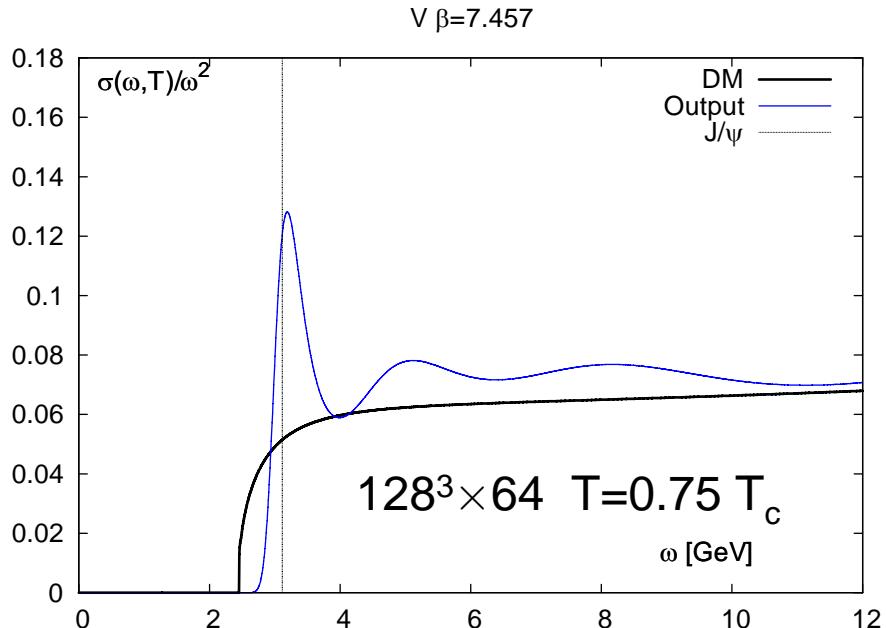
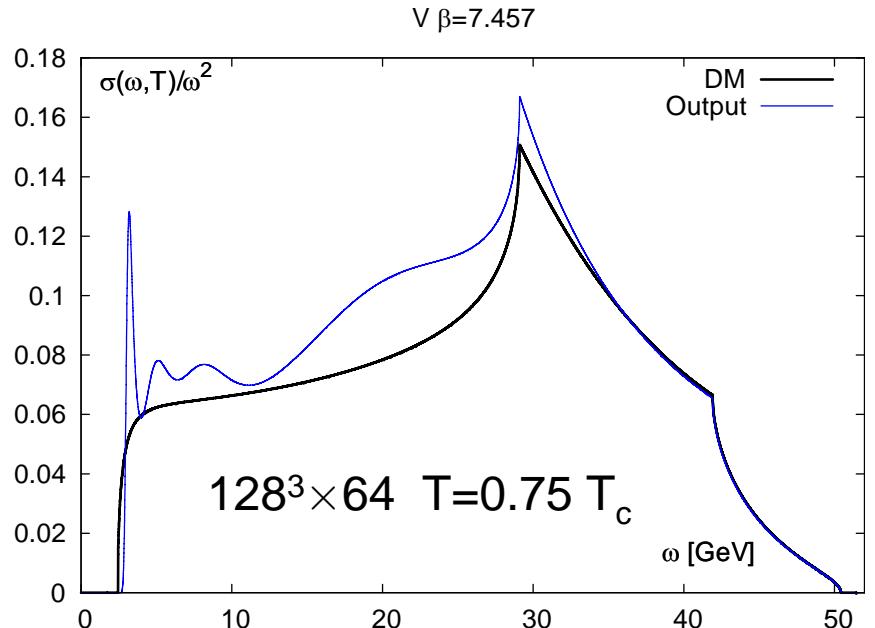
largest distance 0.25 fm due to compact temporal direction

only small distance regime relevant

for thermal effects

for bound state effects

MEM – Spectral function below T_c

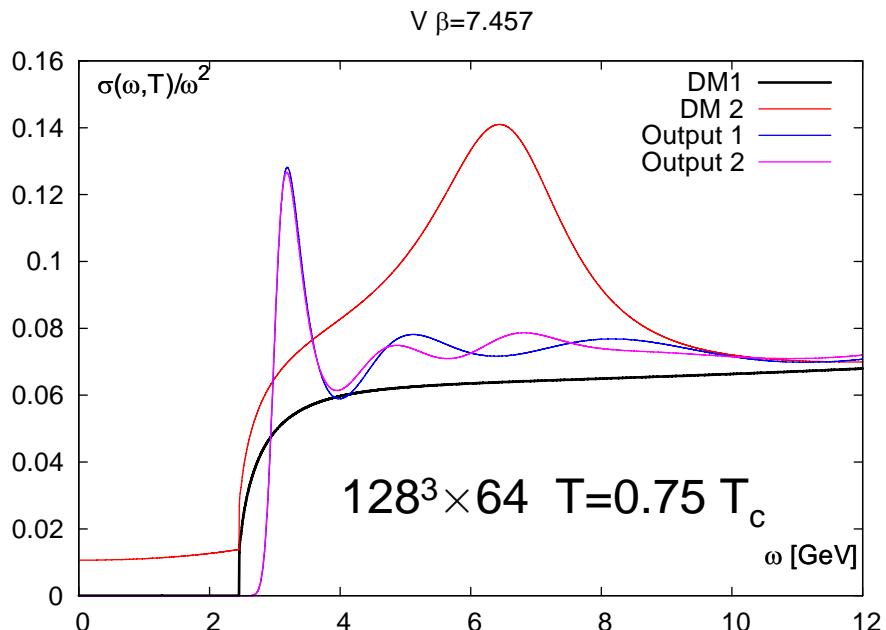
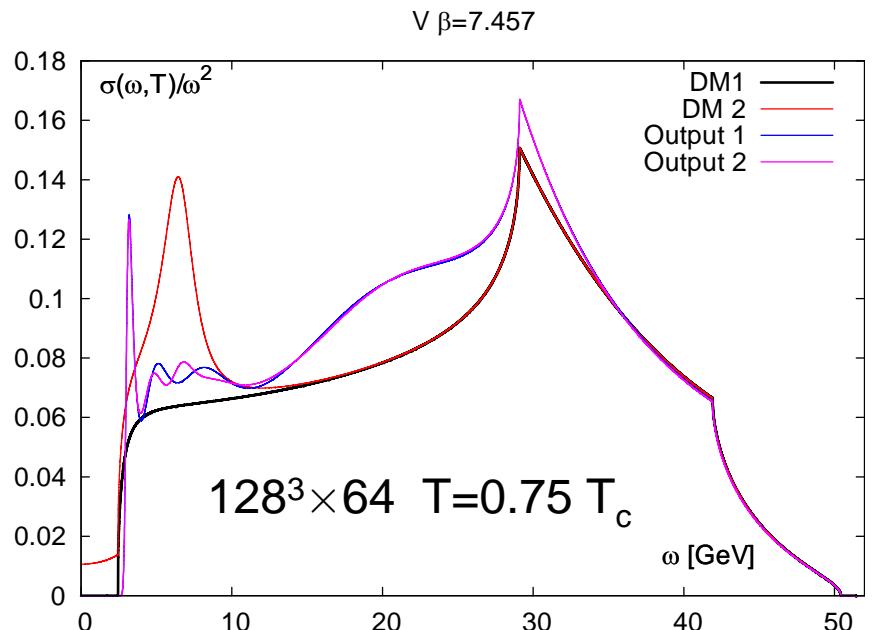


$N_\sigma=128$ and $N_\tau=64$ ($a^{-1} \approx 13 \text{ GeV}$) \rightarrow cut-off effects well separated

Pronounced ground state peak close to J/ψ mass

no zero mode contributions observed below T_c in all channels

MEM – Spectral function below T_c



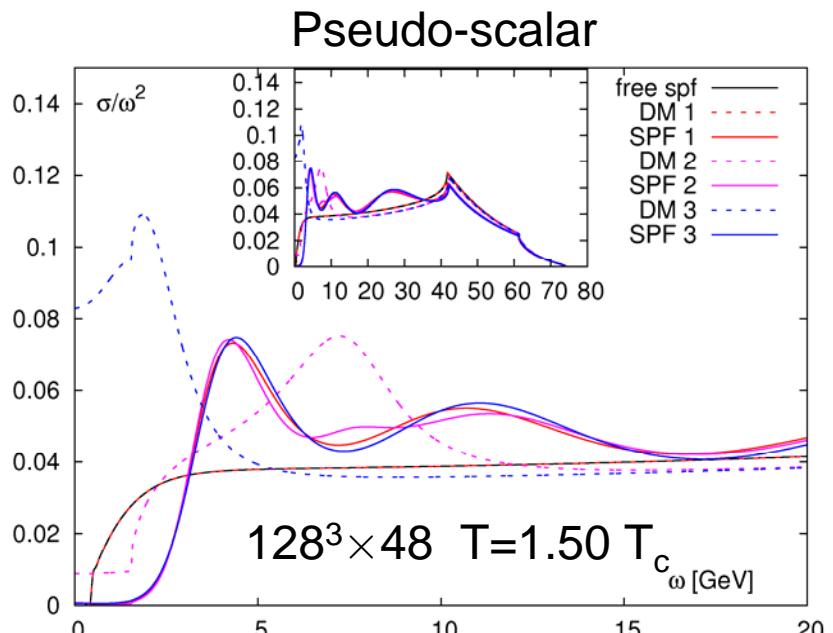
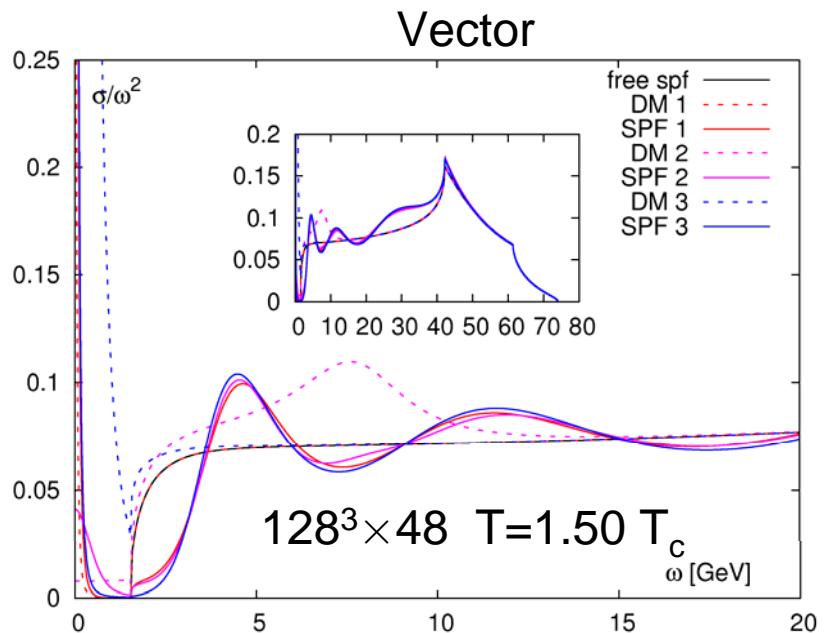
$N_\sigma=128$ and $N_\tau=64$ ($a^{-1} \approx 13$ GeV) \rightarrow cut-off effects well separated

pronounced ground state peak close to J/ψ mass

no zero mode contributions observed below T_c in all channels

first peak independent of default model

MEM – Spectral function above T_c



$N_\sigma = 128$ and $N_\tau = 48$ ($a^{-1} \approx 19$ GeV) → cut-off effects well separated

small ω region accessible → hope to extract transport properties

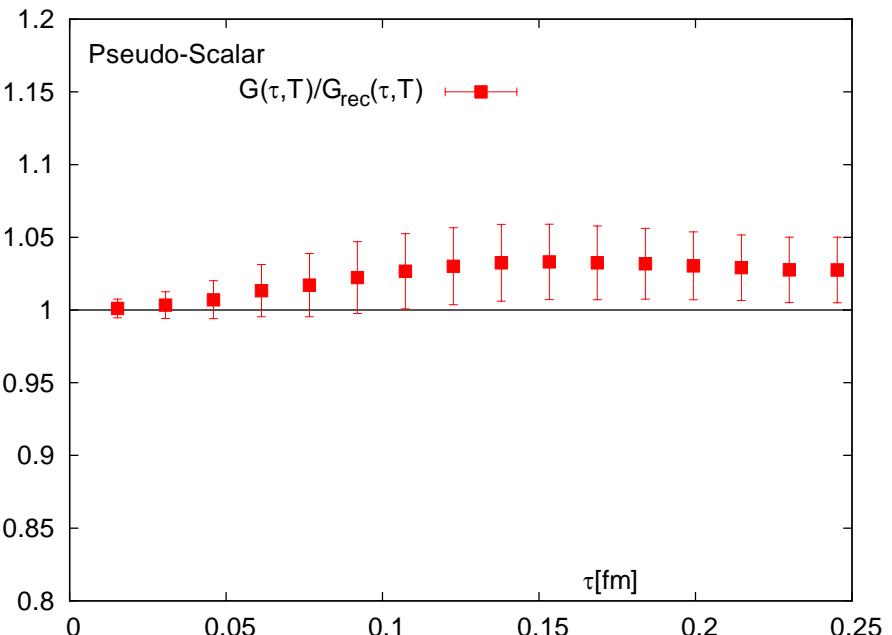
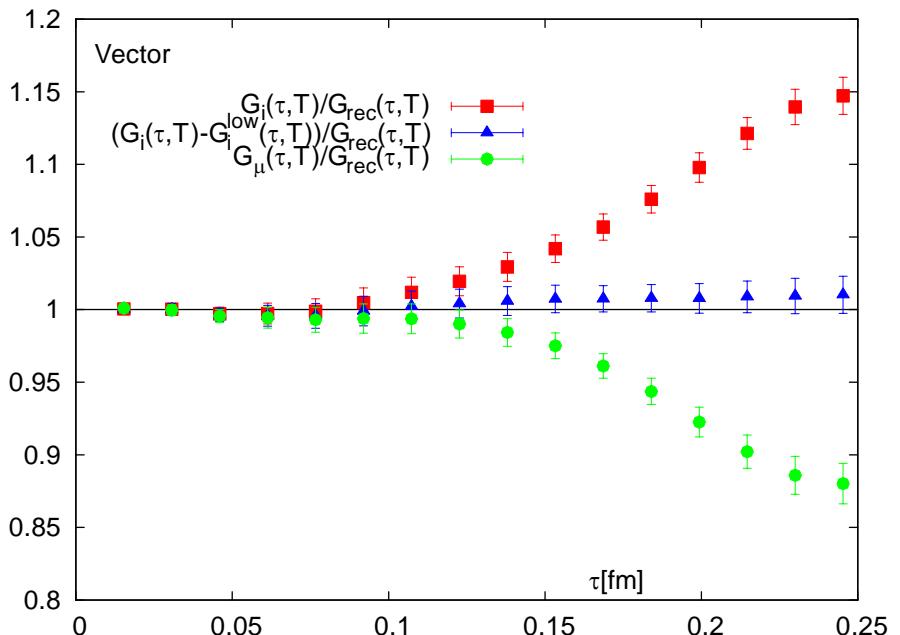
default model dependence only in small ω region → higher statistics needed

no default model dependence in the intermediate ω region

no zero-mode contribution in pseudo-scalar channel observed

no pronounced peak → bound states melted? / threshold enhancement?

Charmonium Correlators – Zero Mode Contributions



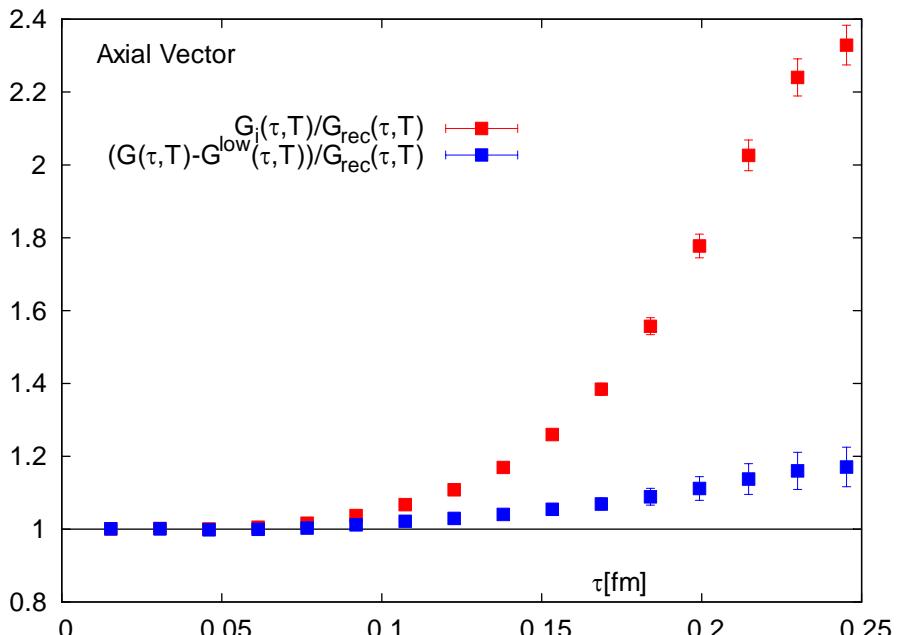
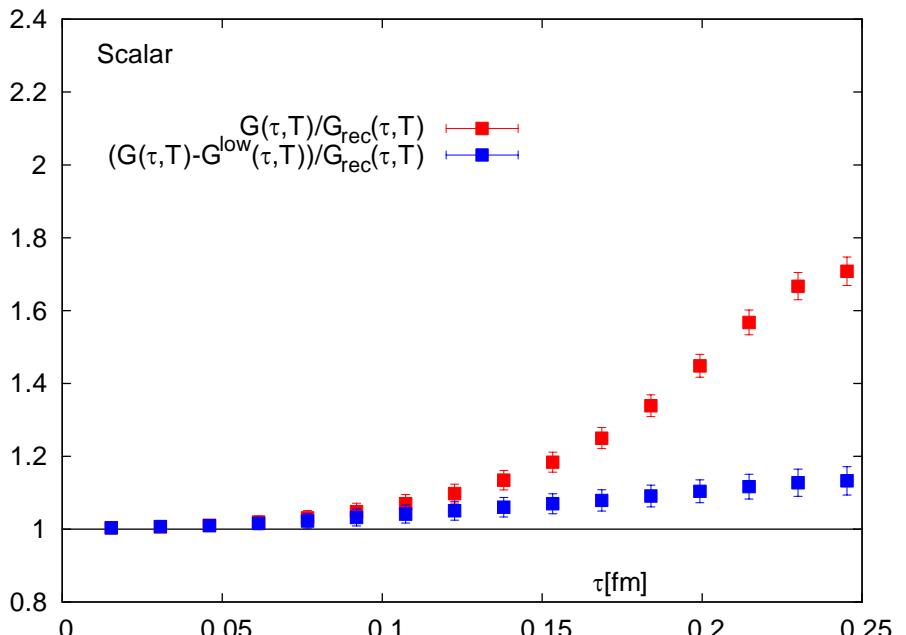
$$G_{rec}(\tau, T) = \int \sigma_0(\omega, 0.75T_c) K(\omega, \tau, T)$$

$$G_{rec}^{low}(\tau, T) = \int_0^{2m_c} \sigma_T(\omega, T) K(\omega, \tau, T)$$

$$G_{\mu\mu}(\tau, T) = G_{ii}(\tau, T) + G_{44}(\tau, T)$$

- main T-effect due to zero-mode contribution
 - well described by small ω -part of $\sigma_T(\omega, T)$
 - smaller than $G_{00}(\tau, T) = \chi(T)\tau$
 - no zero-mode contribution in PS-channel
- (similar to discussions by Umeda, Petreczky)

Charmonium Correlators – Zero Mode Contributions



$$G_{rec}(\tau, T) = \int \sigma_0(\omega, 0.75T_c) K(\omega, \tau, T)$$

$$G_{rec}^{low}(\tau, T) = \int_0^{2m_c} \sigma_T(\omega, T) K(\omega, \tau, T)$$

- larger zero-mode contribution in S-wave
- larger T-effect in the S-wave states

systematic uncertainties in reconstruction and low- ω part of spectral function
 high quality data and small lattice spacing + large momenta (volume) needed

Conclusions and Outlook

Hadronic correlators for light quarks ($m_q=0$)

Thermodynamic and Continuum Limit of screening masses!

Spectral functions → Dilepton rates, Transport coefficients?

Momentum dependence needs to be analyzed → Photonrates

Comparison with HTL calculations and experiment

Charmonium hadronic correlators ($m_q=m_c$)

What can we learn from Hadronic correlators on Dissociation?

Momentum dependence vs. Spatial correlators/screening masses?

Spectral functions → Dilepton rates, Transport coefficients?

Momentum dependence needs to be analyzed → Photonrates

Comparison with HTL/NRQCD calculations and experiment

Many Thanks to

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