

# Decoding the phase structure of QCD via particle production at high energy

the QCD phase diagram

EOS, the chiral and the deconfinement phase transitions in lattice QCD  
experimental access to the QCD phase diagram

- delineating the phase boundary from hadron yields
- access to chiral criticality via fluctuations of conserved charges
- deconfinement and quarkonia

work done over the past 18 years in collaboration  
with Peter Braun-Munzinger, Anton Andronic,  
Krzysztof Redlich  
see article in Nature 561 (2018) 321



Johanna Stachel – Universität Heidelberg  
EMMI Physics Day 2018  
GSI, November 20, 2018

# Phase diagram of strongly interacting matter

## at low temperature and normal density

colored quarks and gluons are bound in colorless hadrons - **confinement**

chiral symmetry is spontaneously broken (generating 99% of proton mass e.g.)

1972 QCD (Gross, Politzer, Wilczek)

asymptotic freedom at small distances

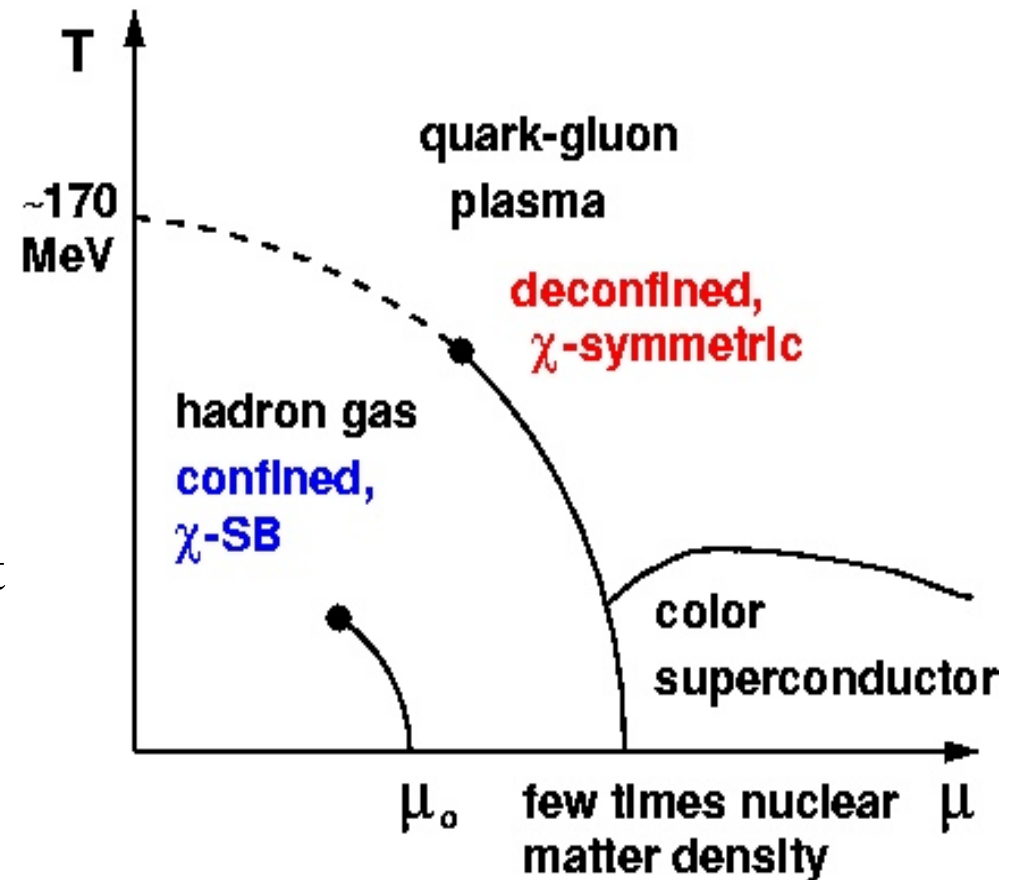
## at high temperature and/or high density

quarks and gluons freed from confinement

-> new state of strongly interacting matter

1975 (Collins/Perry and Cabibbo/Parisi)

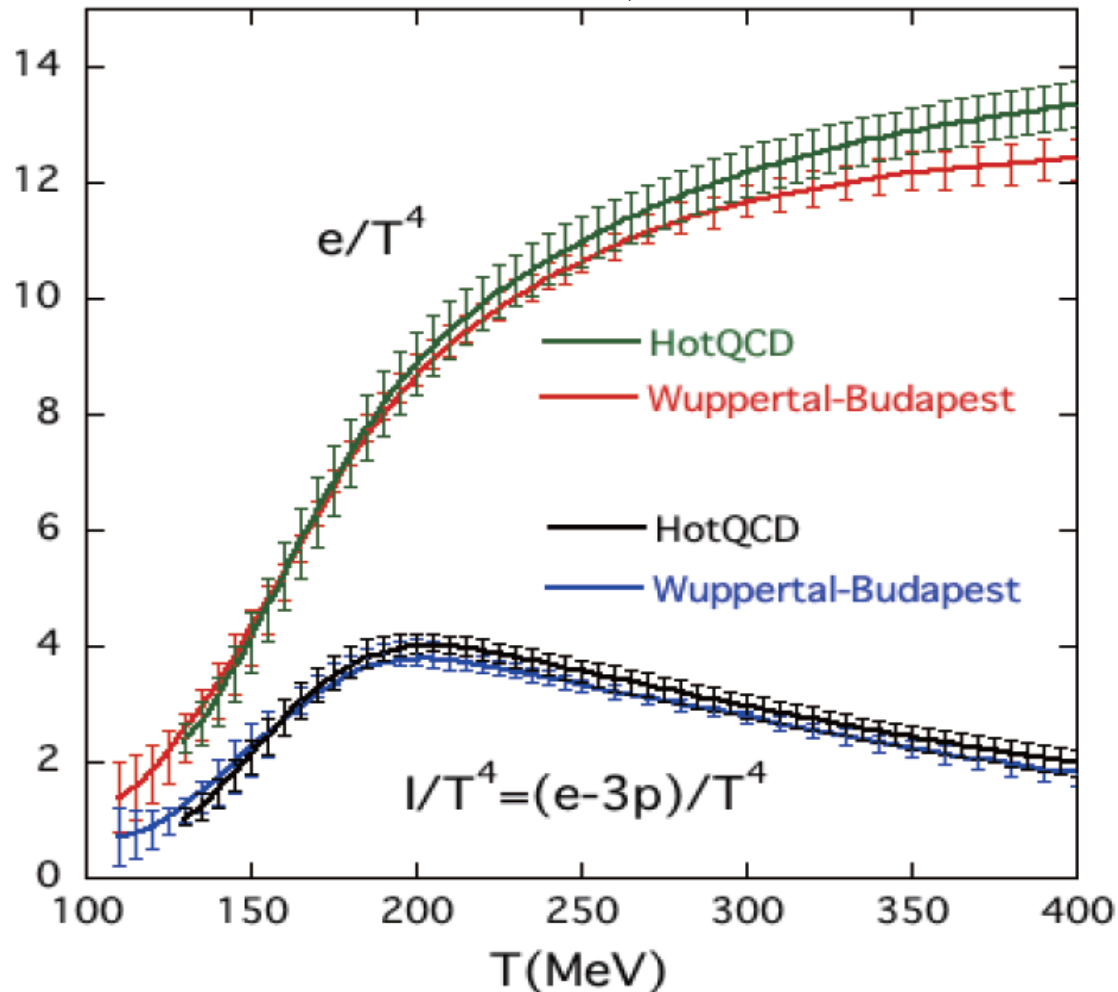
called **Quark-Gluon Plasma (QGP)**



# Equation of state of hot QCD matter in lattice QCD

computation of QCD EoS one of the major goals in lQCD community since 1980

A.Ukawa, arXiv:1501.04215



consolidated results on EoS from different groups, extrapolated to continuum and chiral limit

rapid rise of energy density (normalized to  $T^4$  rise for relativistic gas)

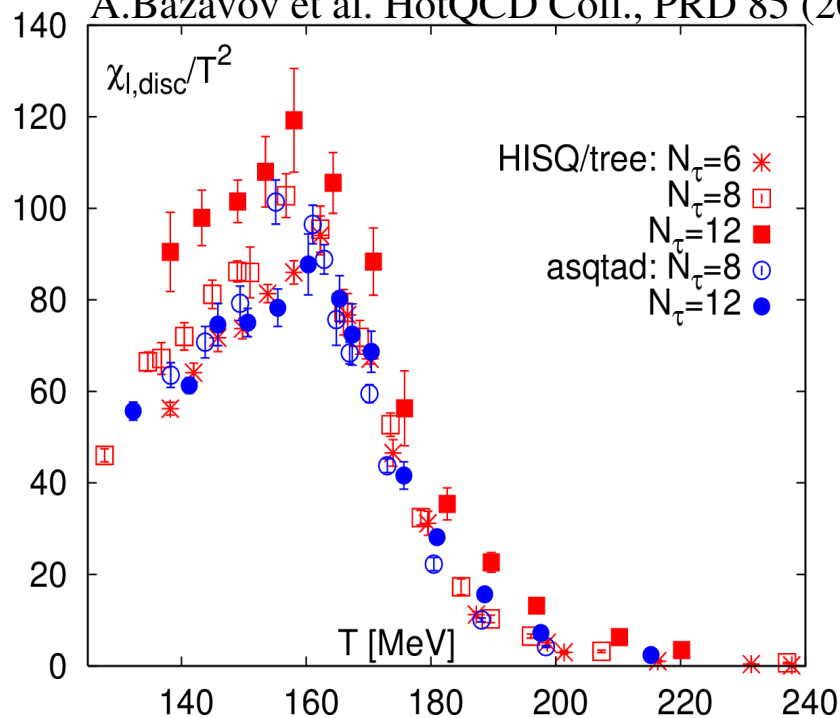
- signals rapid increase in degrees of freedom due to transition from hadrons to quarks and gluons
- lQCD points to continuous cross over transition

# Measure for chiral symmetry restoration in IQCD

order parameter: chiral condensate, its susceptibility peaks at  $T_c$

S.Borsanyi et al. Wuppertal-Budapest Coll., JHEP 1009 (2010) 073

A.Bazavov et al. HotQCD Coll., PRD 85 (2012) 054503



$$\langle \bar{\Psi} \Psi \rangle = \frac{T}{V} \frac{\partial \ln Z}{\partial m}$$

$$\chi_{\bar{\Psi} \Psi} = \frac{T}{V} \frac{\partial^2 \ln Z}{\partial m^2}$$

comparing different measures and different fermion actions, consensus:  
pseudocritical temperature  $T_c = 154 \pm 9$  MeV for chiral restoration

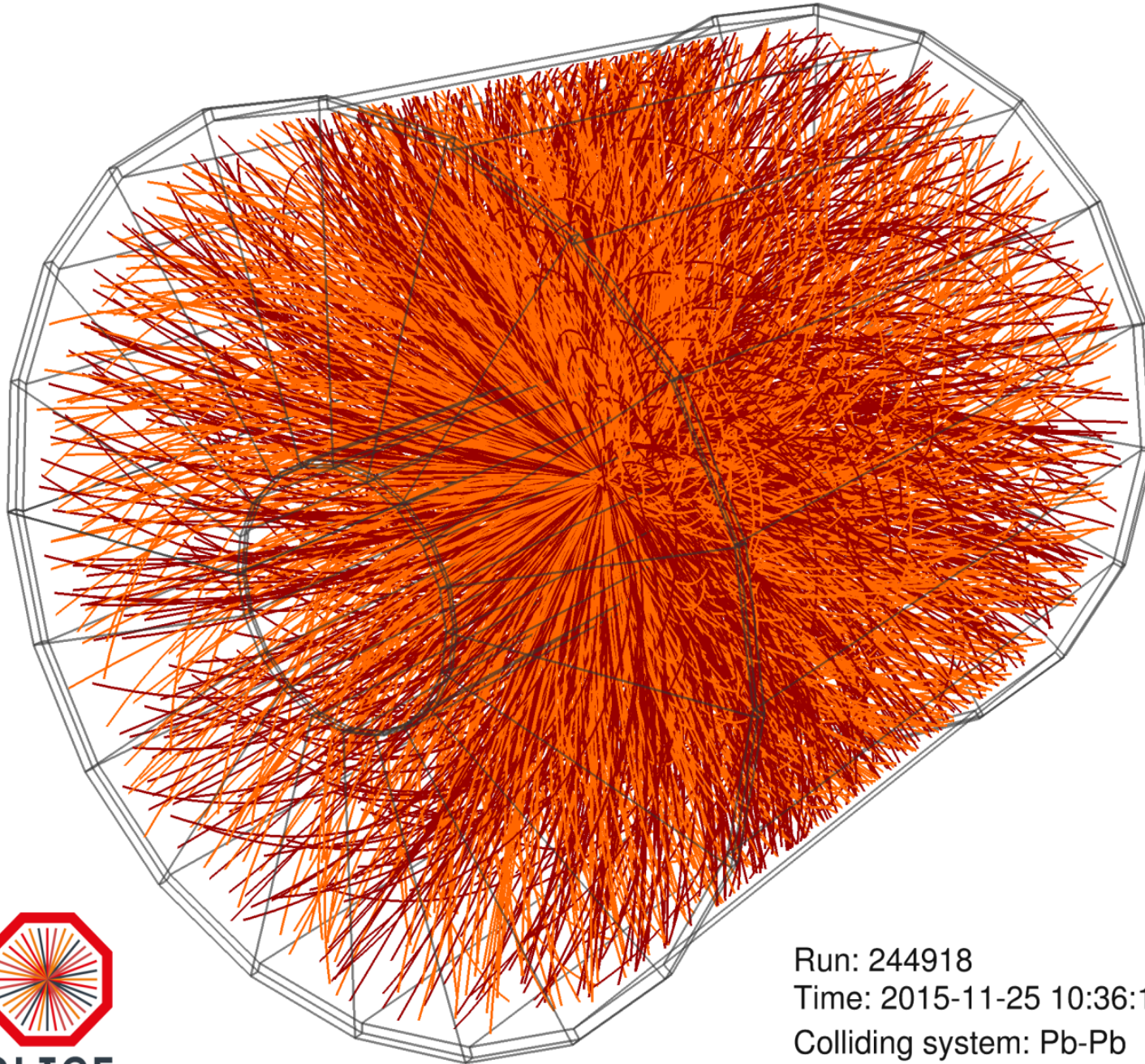
# Experiment

QGP and phase diagram studied in high energy collisions of nuclei  
since 1987 at AGS/SPS,  
since 2000 at RHIC,  
since 2010 at the LHC at  $\sqrt{s_{NN}} = 2.76$  TeV, now 5.02 TeV



nuclear collision rates	2018: 8 kHz, from 2021: 50 kHz
read-out rate	500 Hz                      50 kHz

# first PbPb collisions at LHC at $\sqrt{s} = 5.02$ A TeV



about 3750 charged particles in 1.8 units of pseudorapidity



Run: 244918  
Time: 2015-11-25 10:36:18  
Colliding system: Pb-Pb  
Collision energy: 5.02 TeV

# Experimental Observables

... and link to QCD phase diagram

# Hadro-chemical composition of the fireball

what are the 25 800 hadrons observed in the final state at LHC?  
(32 300 at full LHC energy)



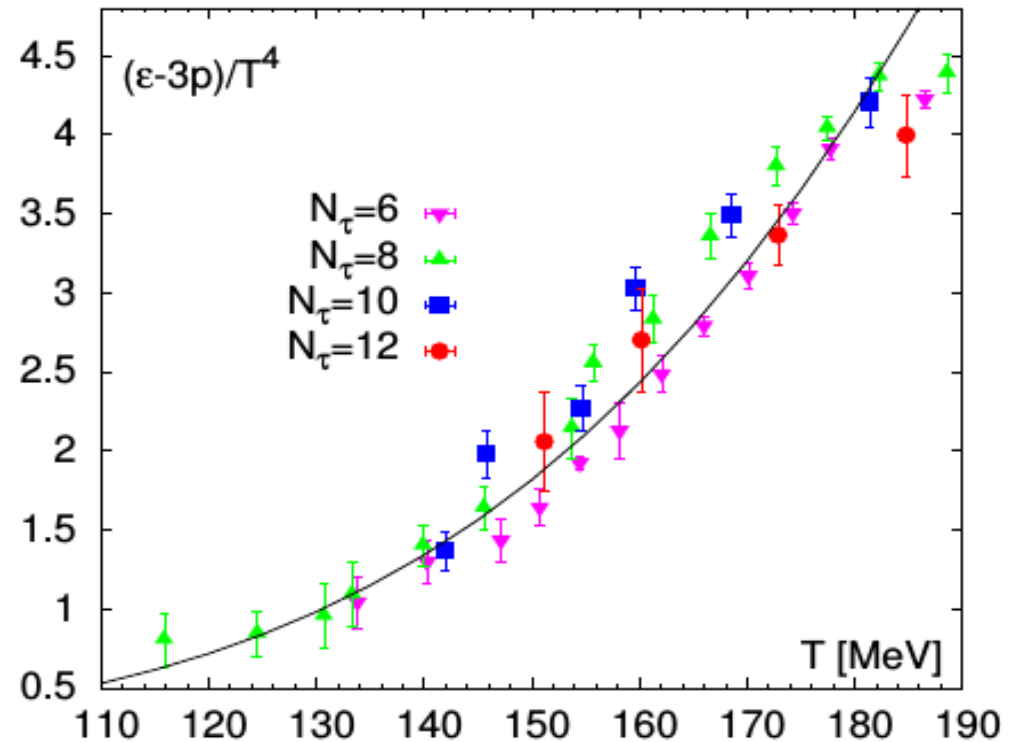
# Duality between hadrons and quarks/gluons (I)

all thermodynamic quantities derived from full QCD partition function  $Z$

e.g. the pressure  $\frac{p}{T^4} = \frac{1}{T^3} \frac{\partial \ln Z(V, T, \mu)}{\partial V}$

trace anomaly from lQCD  
full dynamical quarks with realistic pion  
mass (HotQCD coll. PRD 90 (2014) 094503)  
perfectly matched by  
hadron resonance gas prediction  
(solid line)

similar agreement seen for many other  
observables



# Duality between hadrons and quarks/gluons (II)

in the dilute limit  $T < 165$  MeV:

$$\ln Z(T, V, \mu) \approx \sum_{i \in \text{mesons}} \ln \mathcal{Z}_{M_i}^M(T, V, \mu_Q, \mu_S) + \sum_{i \in \text{baryons}} \ln \mathcal{Z}_{M_i}^B(T, V, \mu_b, \mu_Q, \mu_S)$$

- partition function of hadron resonance model expressed in mesonic and baryonic components.
- chemical potentials reflect the baryon number, charge and strangeness

# Thermal model of particle production and QCD

partition function  $Z(T,V)$  contains sum over the full hadronic mass spectrum and is fully calculable in QCD

for each hadron species  $i$ , the grand canonical statistical operator is:

$$\ln Z_i = \frac{V g_i}{2\pi^2} \int_0^\infty \pm p^2 dp \ln(1 \pm \exp(-(E_i - \mu_i)/T))$$

leading to particle densities:

$$n_i = N_i/V = -\frac{T}{V} \frac{\partial \ln Z_i}{\partial \mu} = \frac{g_i}{2\pi^2} \int_0^\infty \frac{p^2 dp}{\exp((E_i - \mu_i)/T) \pm 1}$$

for every conserved quantum number there is a chemical potential:

$$\mu_i = \mu_B B_i + \mu_S S_i + \mu_{I_3} I_i^3$$

but can use conservation laws to constrain  $V, \mu_S, \mu_{I_3}$

use full hadronic mass spectrum from the PDG to compute 'primordial yields' and feeding from strong decays



**Fit at each energy provides values for T and  $\mu_b$**

# Hadro-chemistry at the LHC

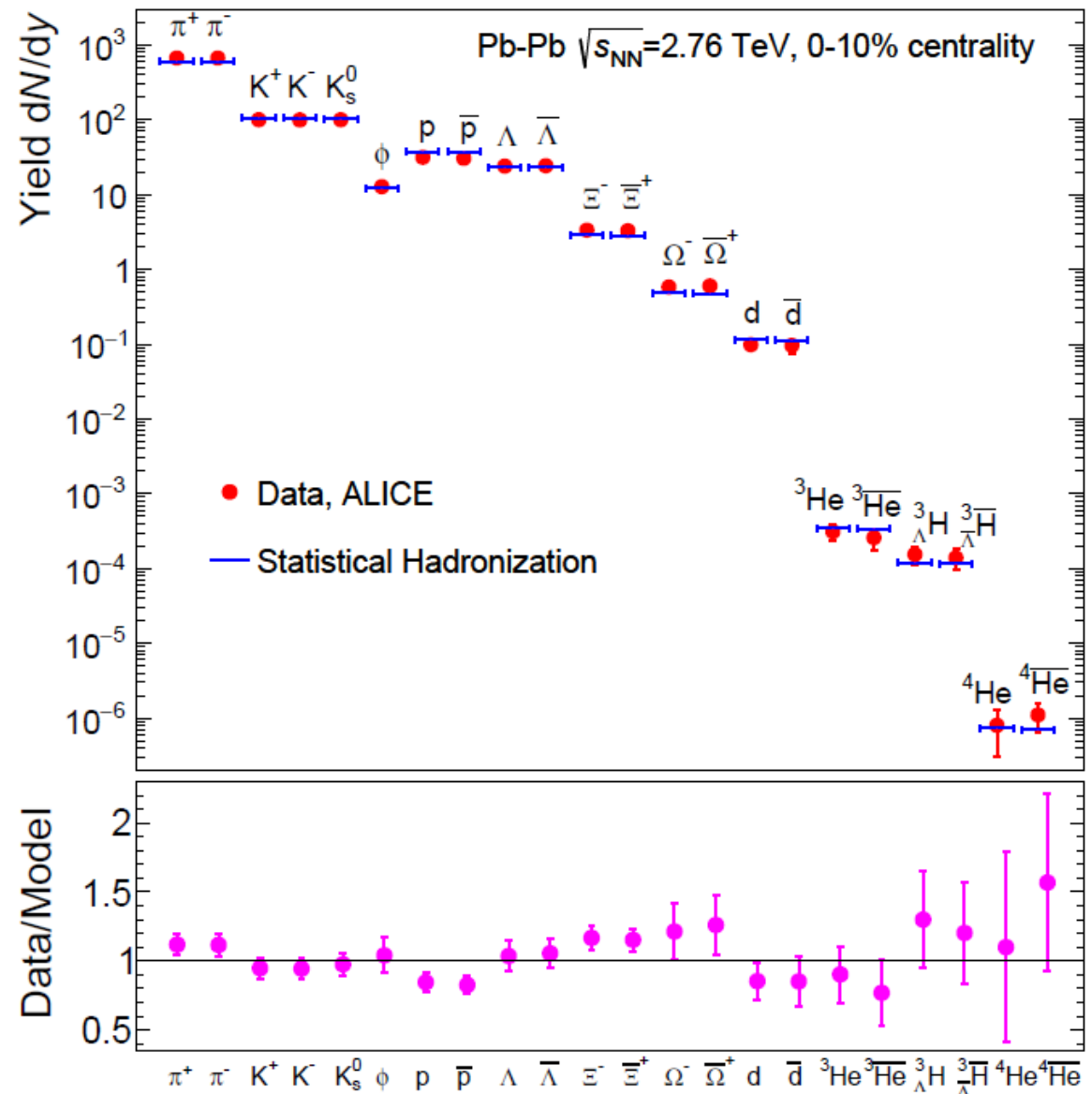
excellent description of  
ALICE@LHC data with  
grand canonical (GC)  
statistical ensemble  
 $T = 156.5 \pm 1.5 \text{ MeV}$

fit includes nuclei

2.7 sigma deviation for protons  
solved in the mean time

$\chi^2/\text{dof} = 19.7/19$

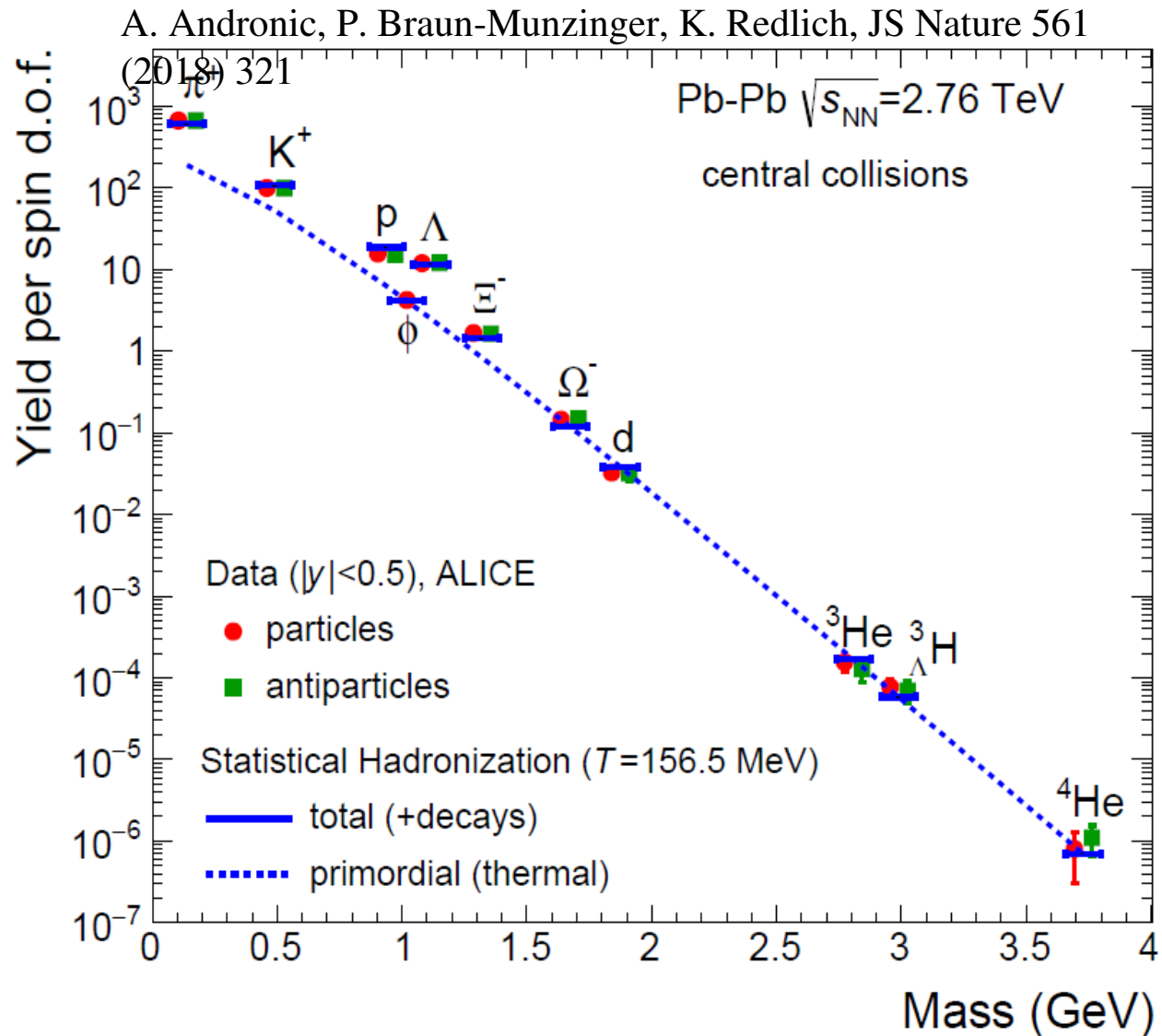
A. Andronic, P. Braun-Munzinger, K. Redlich, JS  
Nature (in print) arXiv: 1710.09425



# Statistical model (grand canonical) describes production of hadrons and (anti-)nuclei at LHC

1 free parameter: temperature  $T$   
 $T = 156.5 \pm 1.5 \text{ MeV}$

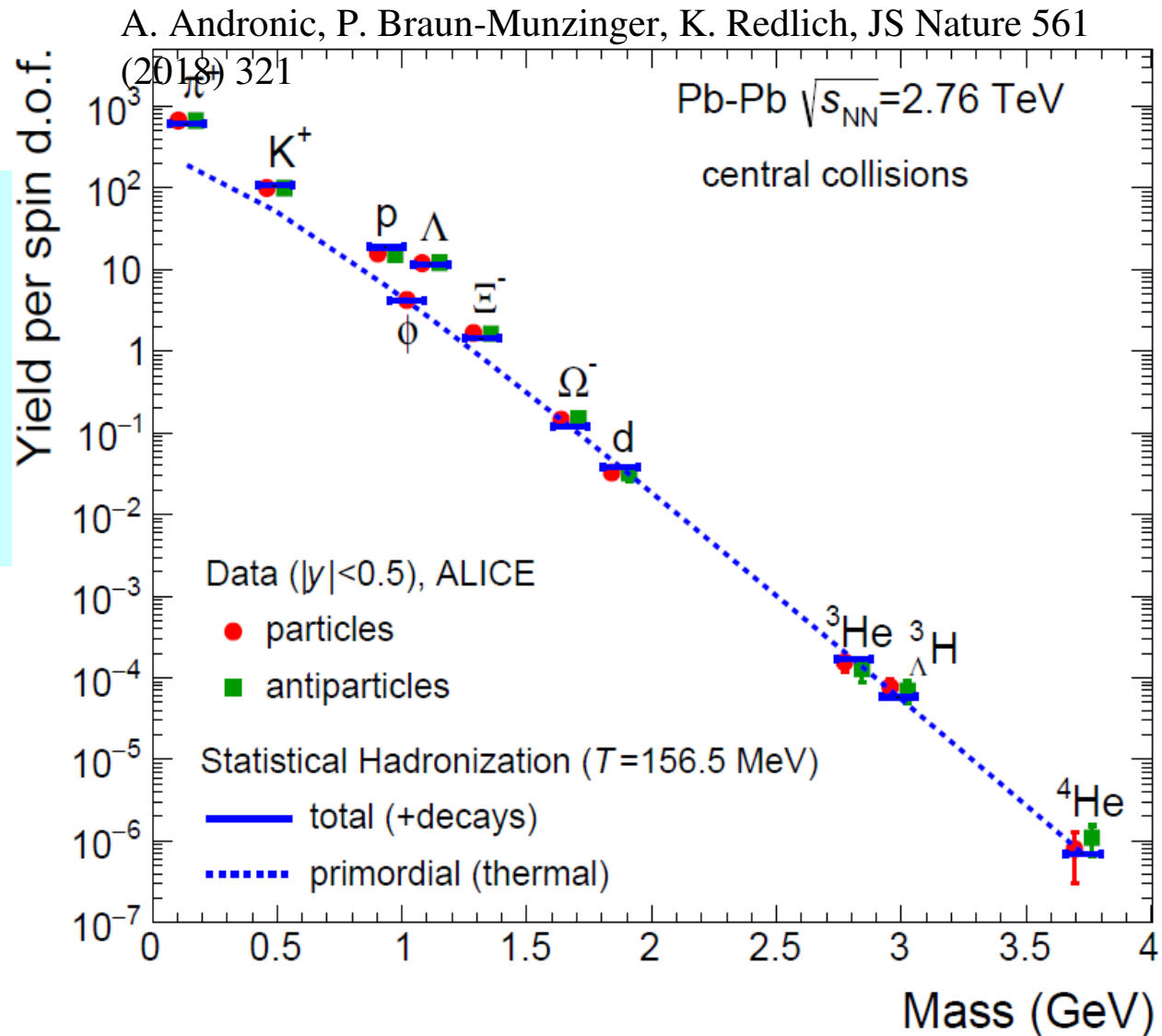
agreement over 9 orders of magnitude with QCD statistical operator prediction  
 (- strong decays need to be added)



# Statistical model (grand canonical) describes production of hadrons and (anti-)nuclei at LHC

1 free parameter: temperature  $T$   
 $T = 156.5 \pm 1.5 \text{ MeV}$

PbPb central collisions: even loosely bound nuclei produced with yields fixed at the phase boundary  
 how can it be?



# The Hypertriton

mass = 2990 MeV, binding energy = 2.3 MeV

$\Lambda$  separation energy = 0.13 MeV

molecular structure: (p+n) +  $\Lambda$

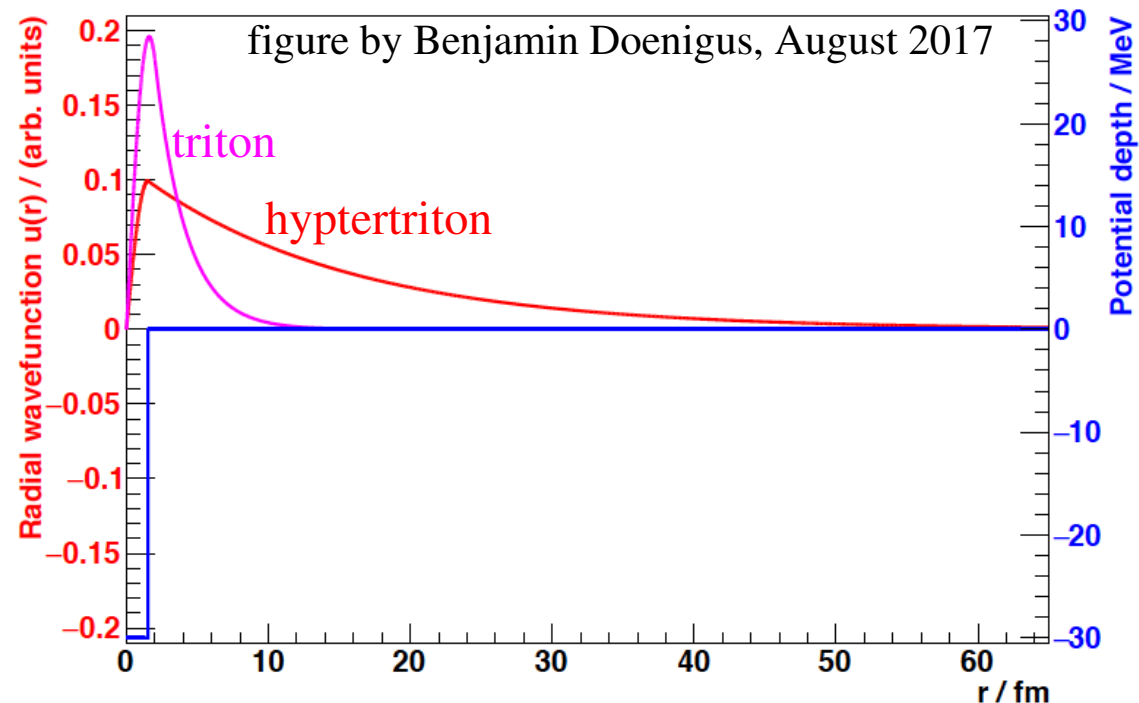
rms radius = rms separation between d and  $\Lambda$  =  $(4 \text{ B.E. } M_{\text{red}})^{-1/2} = 10.6 \text{ fm}$

in that sense: hypertriton = (pn $\Lambda$ ) = (d $\Lambda$ ) is the ultimate halo state

yet production yield is fixed at  
156 MeV temperature  
(about 1000 x separation energy.)

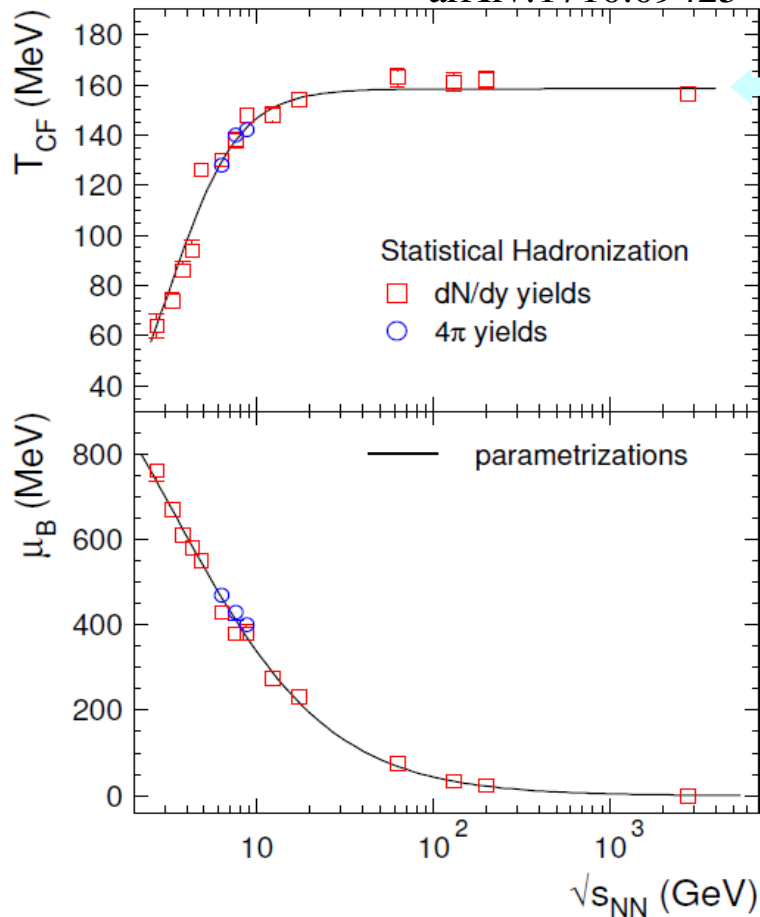
**hypothesis: all nuclei and hyper-nuclei are formed as compact multi-quark states at the phase boundary. Then slow time evolution into hadronic representation.**

Andronic, Braun-Munzinger, Redlich, JS,  
arXiv :1710.09425

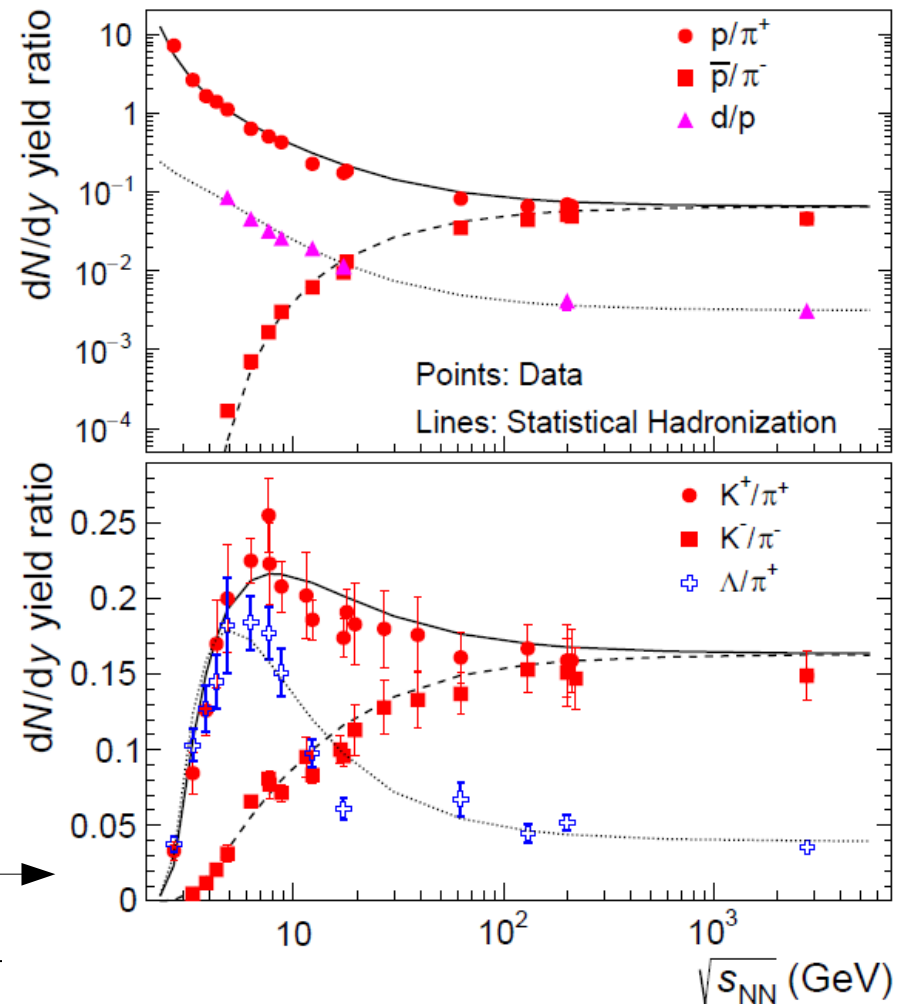


# Statistical analysis for lower collision energy data

arXiv:1710.09425



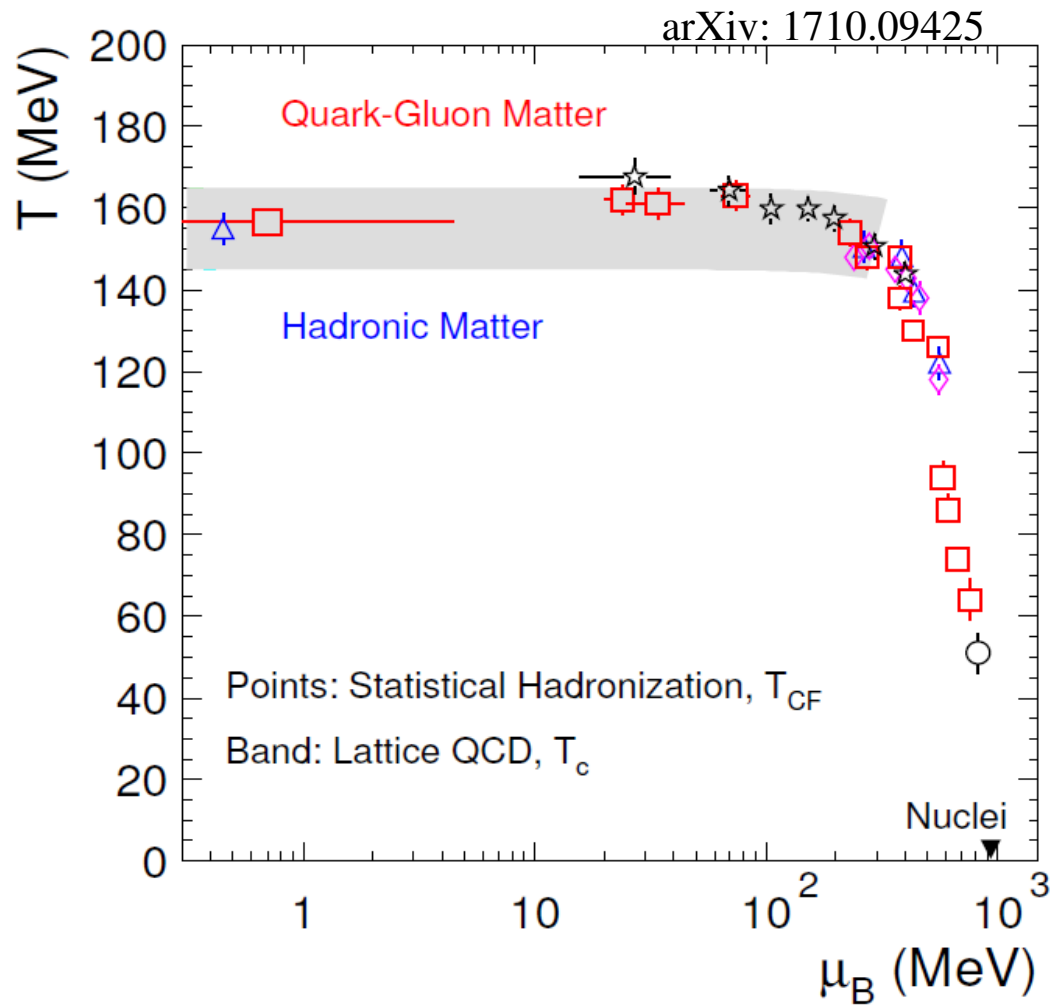
systematic evolution of chemical freeze-out parameters as function of cm energy



hadron yields for Pb-Pb central collisions from LHC down to RHIC, SPS, AGS well described by a statistical ensemble

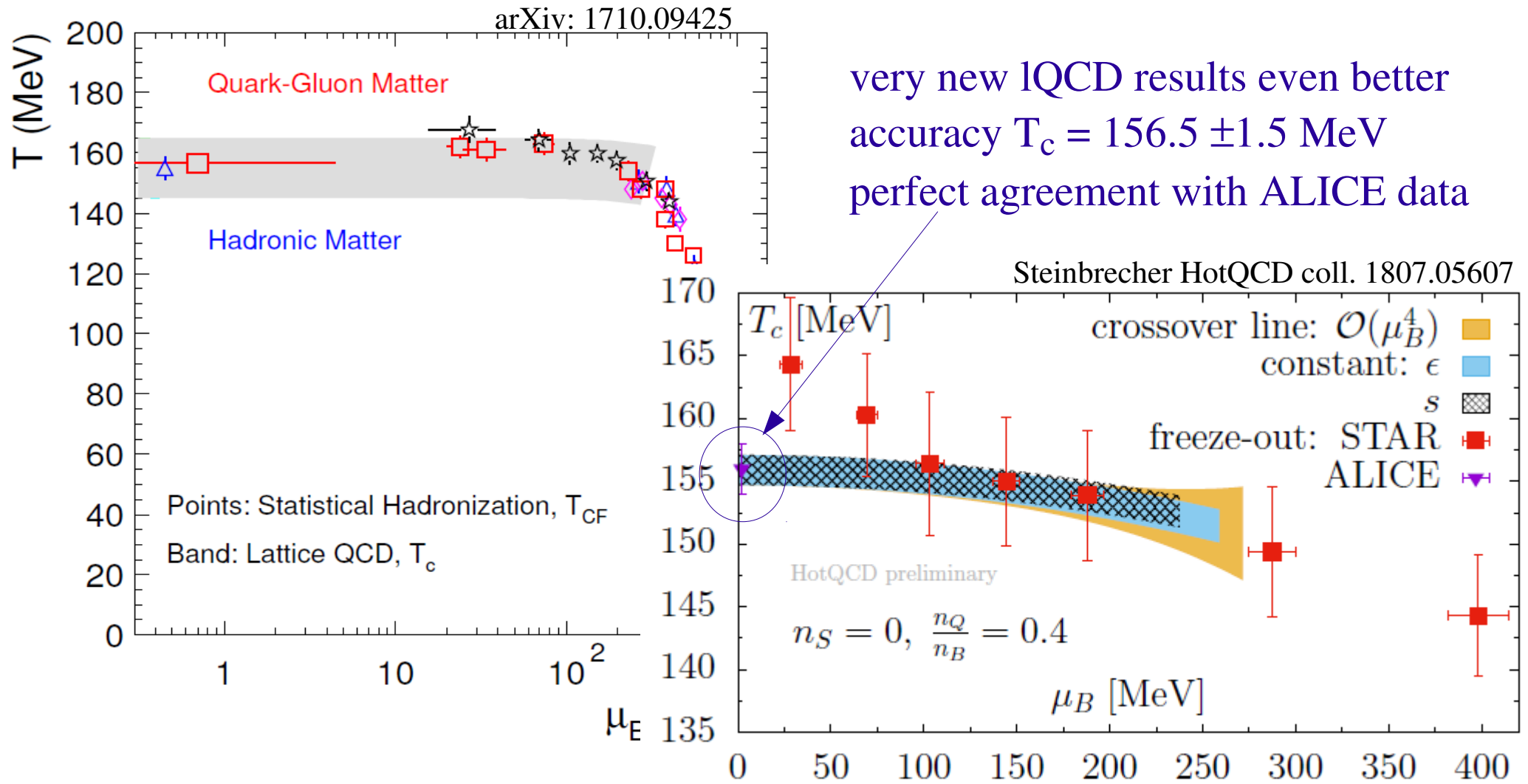


# Energy dependence of temperature and baryochemical potential



quantitative agreement of chemical freeze-out parameters with lQCD predictions of  $\mu_B < 300$  MeV or  $\sqrt{s_{NN}} \geq 10$  GeV

# Energy dependence of temperature and baryochemical potential

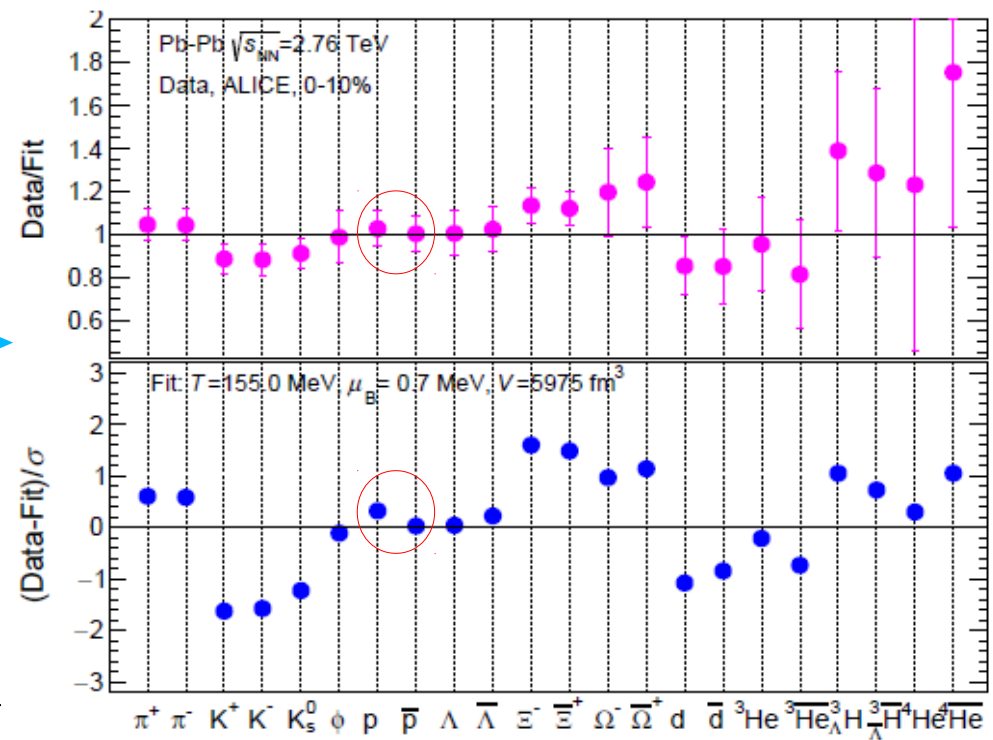
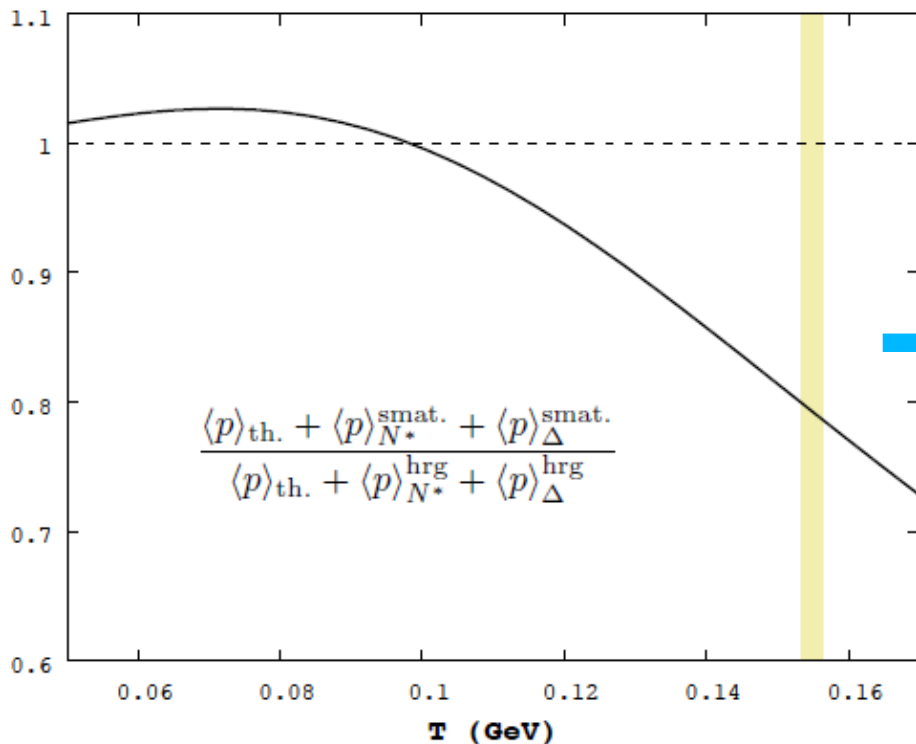


# Solution of the proton puzzle

The thermal proton yield anomaly in Pb-Pb collisions at the LHC and its resolution

Anton Andronic,<sup>1</sup> Peter Braun-Munzinger,<sup>2,3,4</sup> Bengt Friman,<sup>5</sup>  
 Pok Man Lo,<sup>6</sup> Krzysztof Redlich,<sup>6,7,2</sup> and Johanna Stachel<sup>3,2</sup>

use S-matrix formalism to include  $\pi N$  interaction in hadron resonance gas (analysis of measured phase shifts)



# Charmonia as probe of deconfinement

the original idea (Matsui and Satz 1986): implant charmonia into the QGP and observe their modification (Debye screening of QCD), in terms of suppressed production in nucleus-nucleus collisions with plasma formation

→ notion of charmonia as thermometer – sequential melting signature of deconfinement, but no direct link to phase boundary

new insight (Braun-Munzinger, J.S. 2000):

QGP screens all charmonia (as proposed by Matsui and Satz), but charmonium production takes place at the phase boundary,

enhanced production at colliders – signal for deconfinement  
production probability from thermalized charm quarks scales with  $N_{c\bar{c}}^2$

→ yields of charmonia (and open charm hadrons) directly linked to phase boundary and hadronization temperature  
still probe of deconfinement

# Extension of statistical model to include charmed hadrons

- assume: **all charm quarks are produced in initial hard scattering**; number not changed in QGP

$N_{c\bar{c}}^{direct}$  from data (total charm cross section) or from pQCD

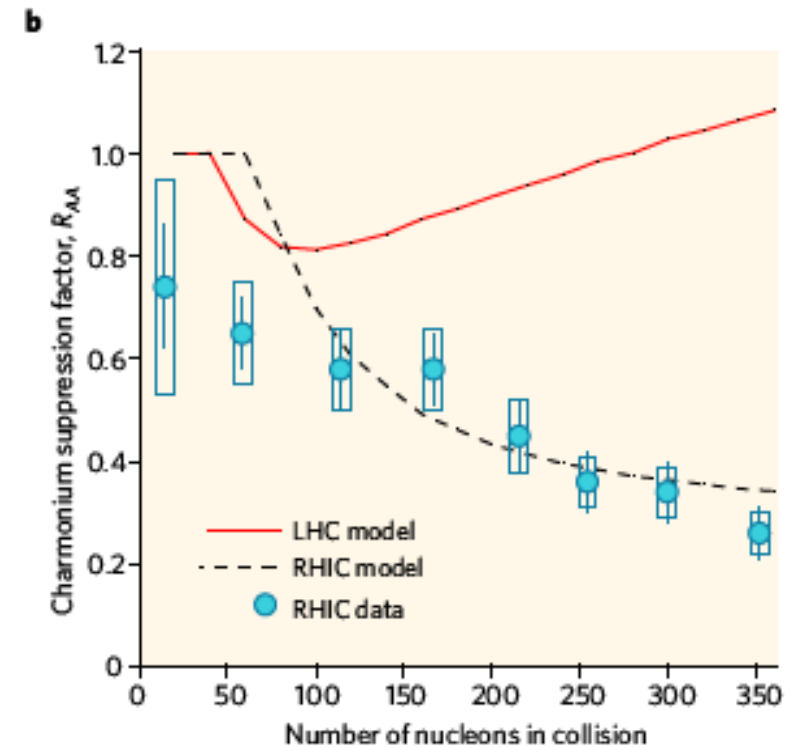
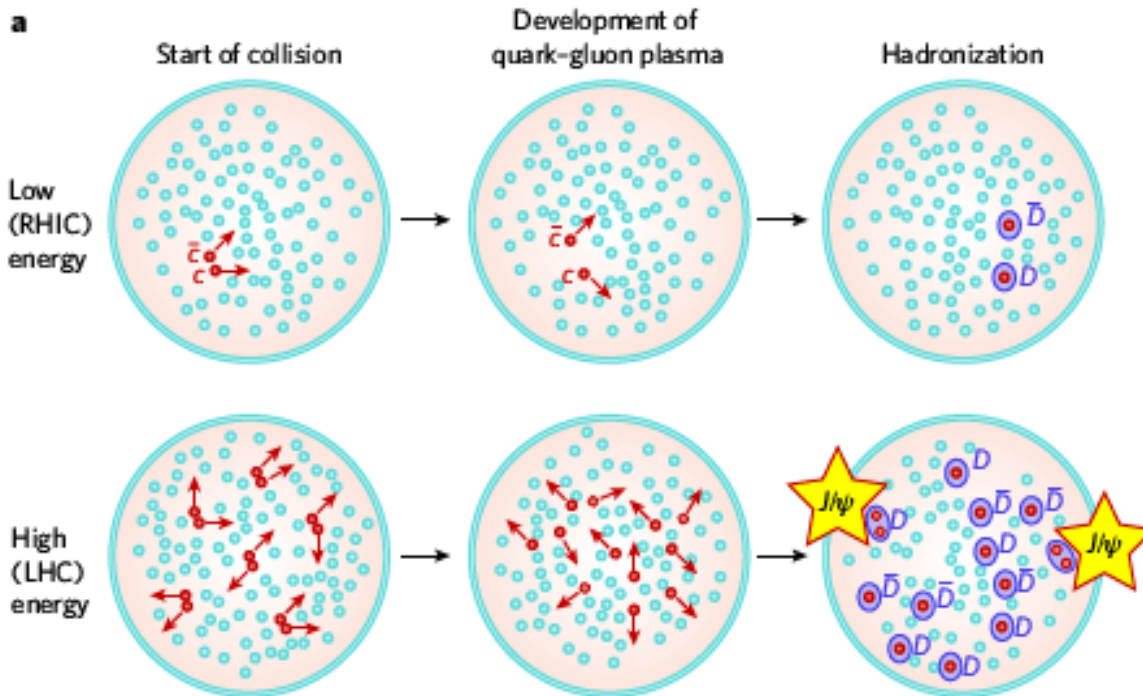
- **hadronization at  $T_c$  following grand canonical statistical model** used for hadrons with light valence quarks (canonical corr. if needed)  
technically number of charm quarks fixed by a charm-balance equation containing fugacity  $g_c$

$$N_{c\bar{c}}^{direct} = \frac{1}{2} g_c V \left( \sum_i n_{D_i}^{therm} + n_{\Lambda_i}^{therm} \right) + g_c^2 V \left( \sum_i n_{\psi_i}^{therm} \right) + \dots$$



the only additional free parameter

# Quarkonium as a probe for deconfinement at the LHC the statistical hadronization picture



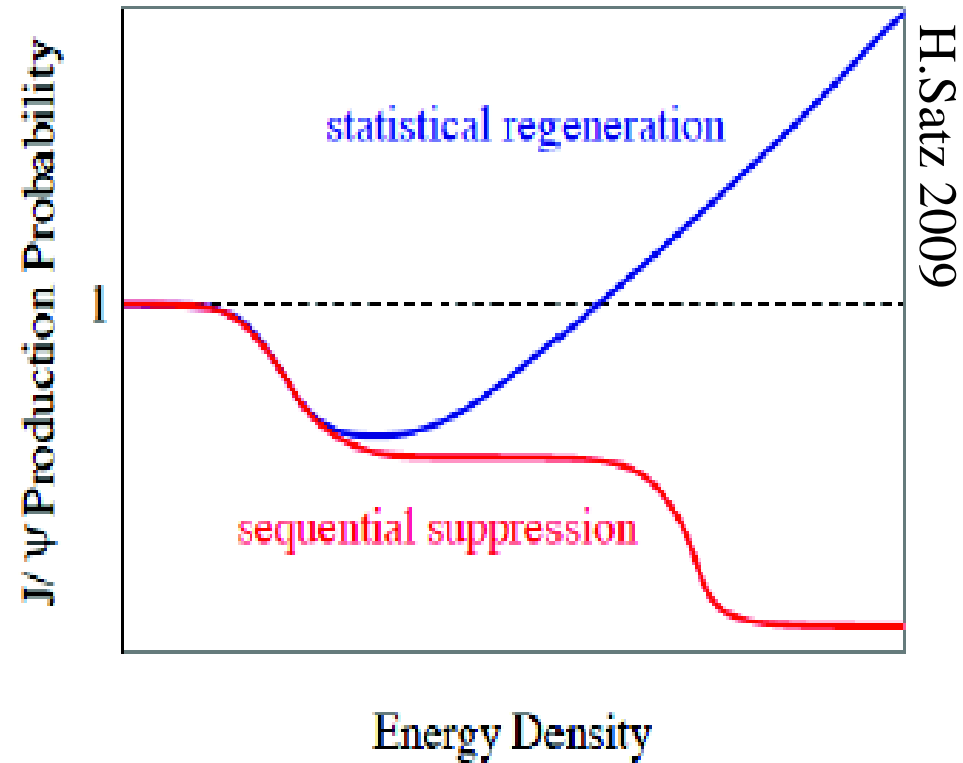
charmonium enhancement as fingerprint of deconfinement at LHC energy  
- a prediction!

Braun-Munzinger, J.S. Phys. Lett. B490 (2000) 196

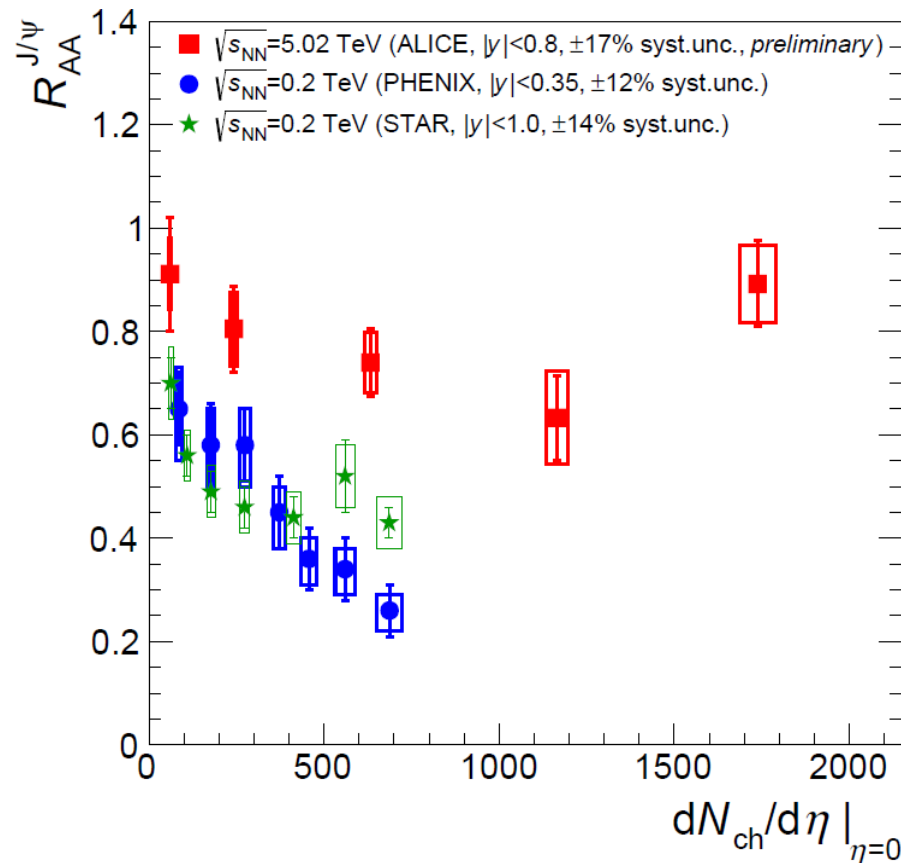
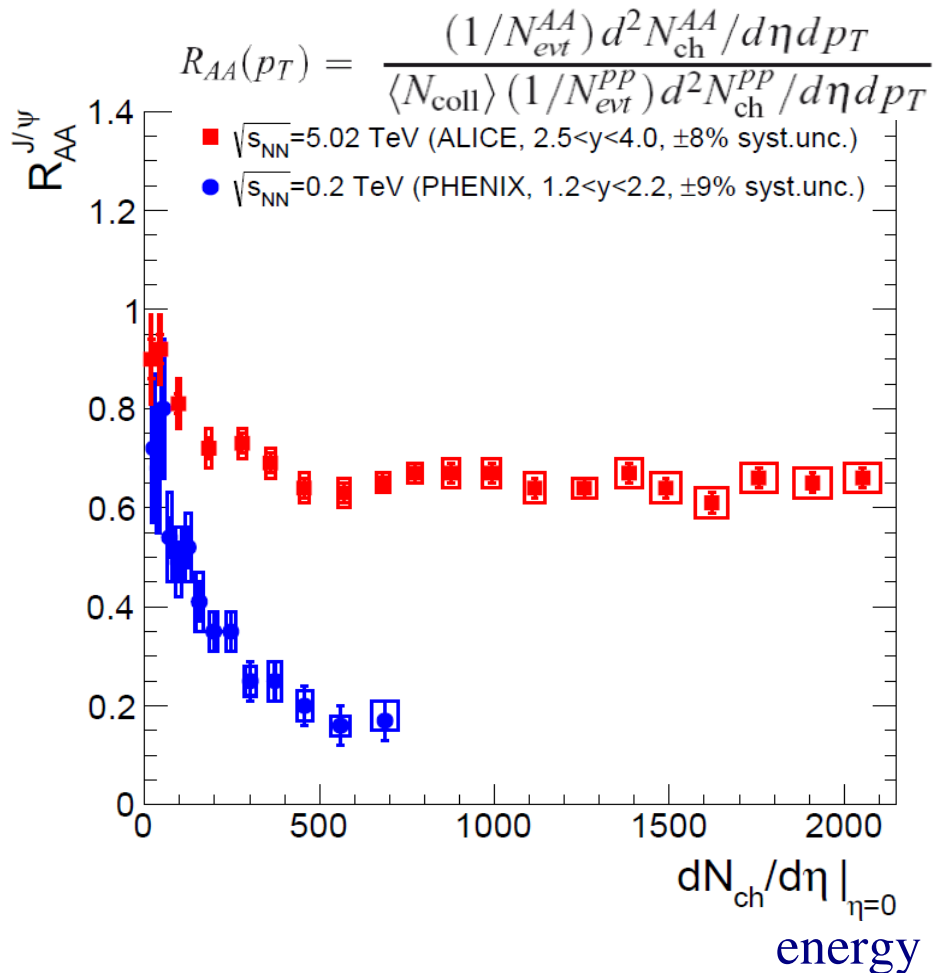
Andronic, Braun-Munzinger, Redlich, J.S., Phys. Lett. B652 (2007) 659

# Expectations for LHC

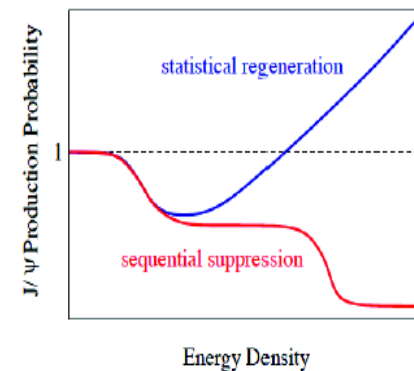
2 possibilities:



# J/ψ production in PbPb collisions: LHC rel. to RHIC

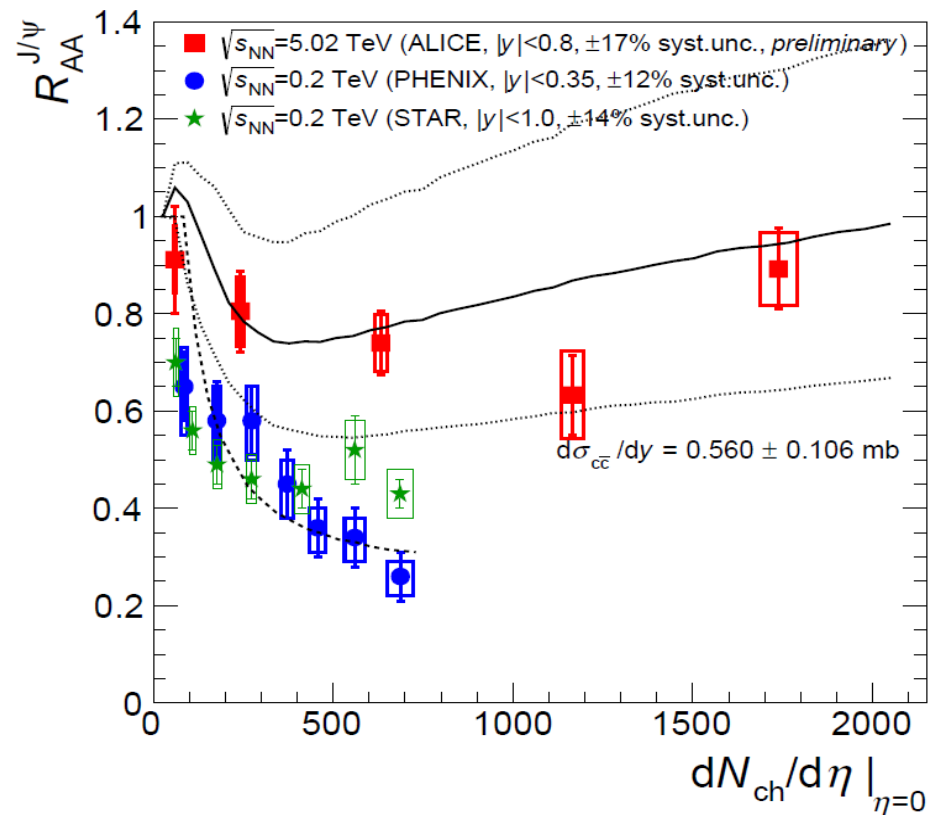
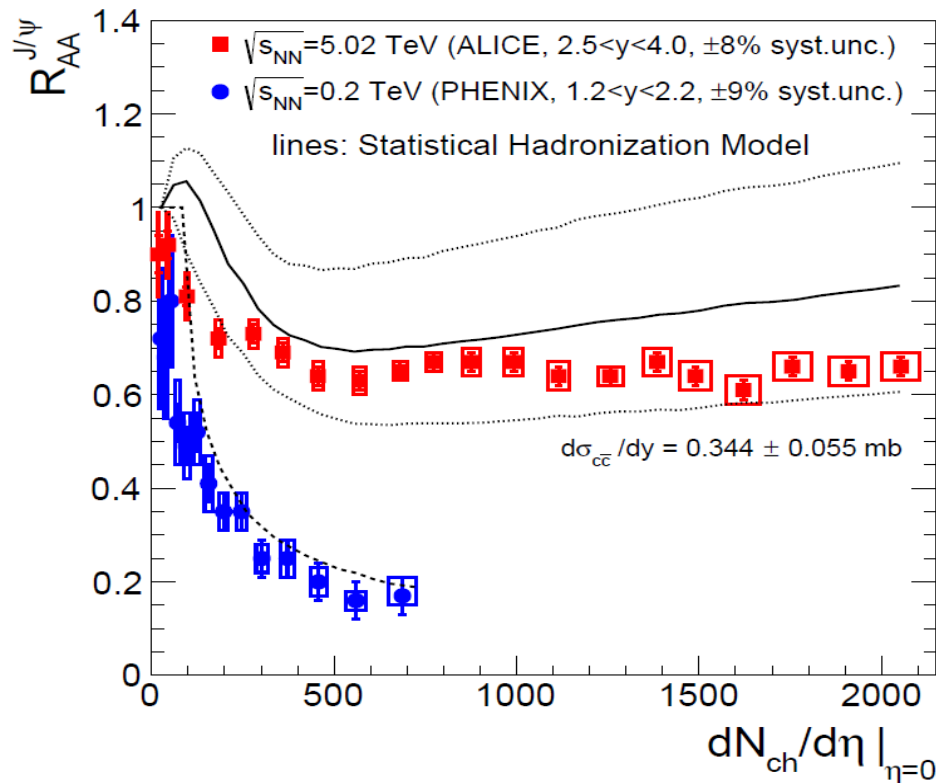


sequential melting scenario not observed  
 rather: **enhancement with increasing energy density!**  
 (from RHIC to LHC and from forward to mid-rapidity)



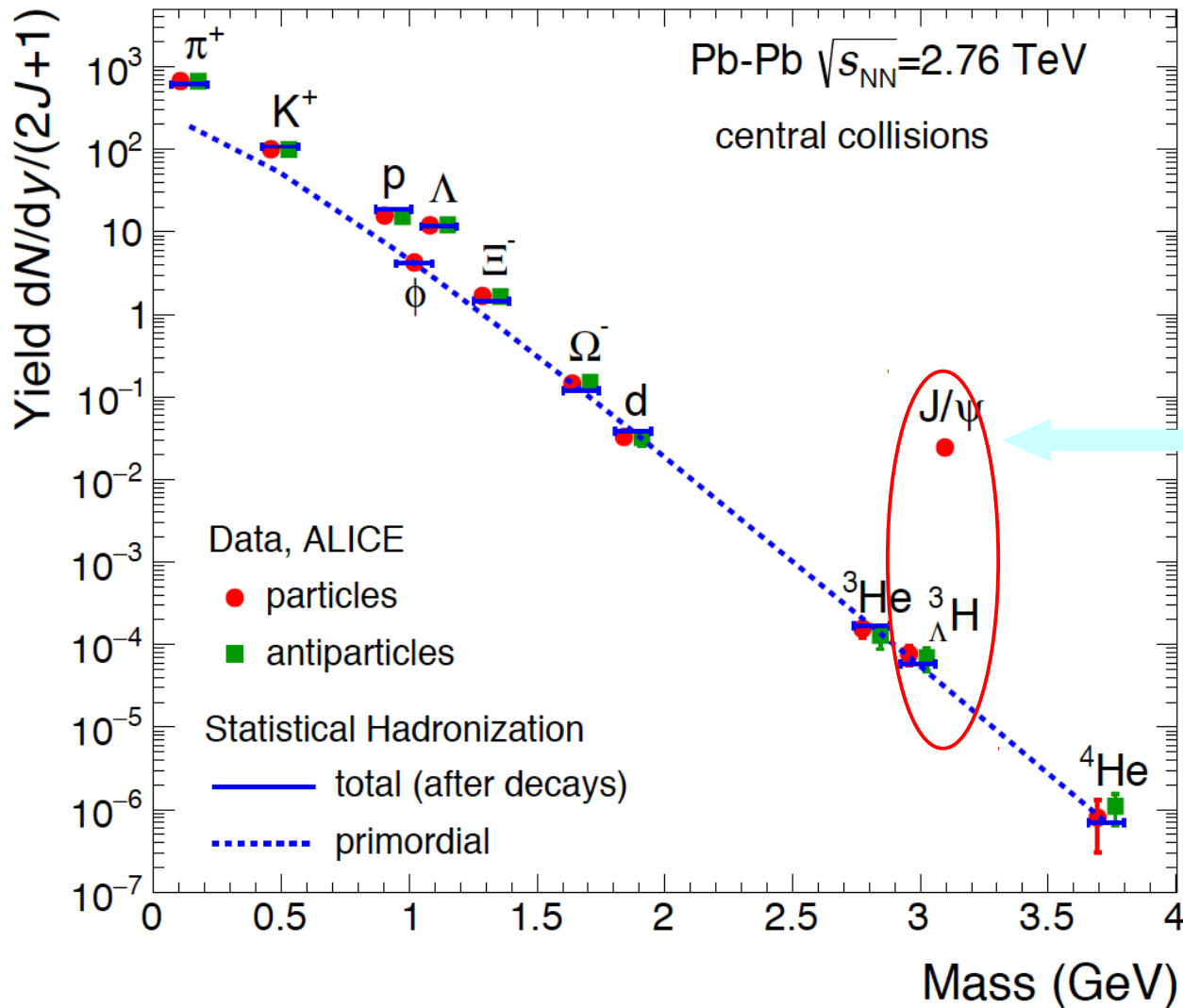


# J/ψ and statistical hadronization



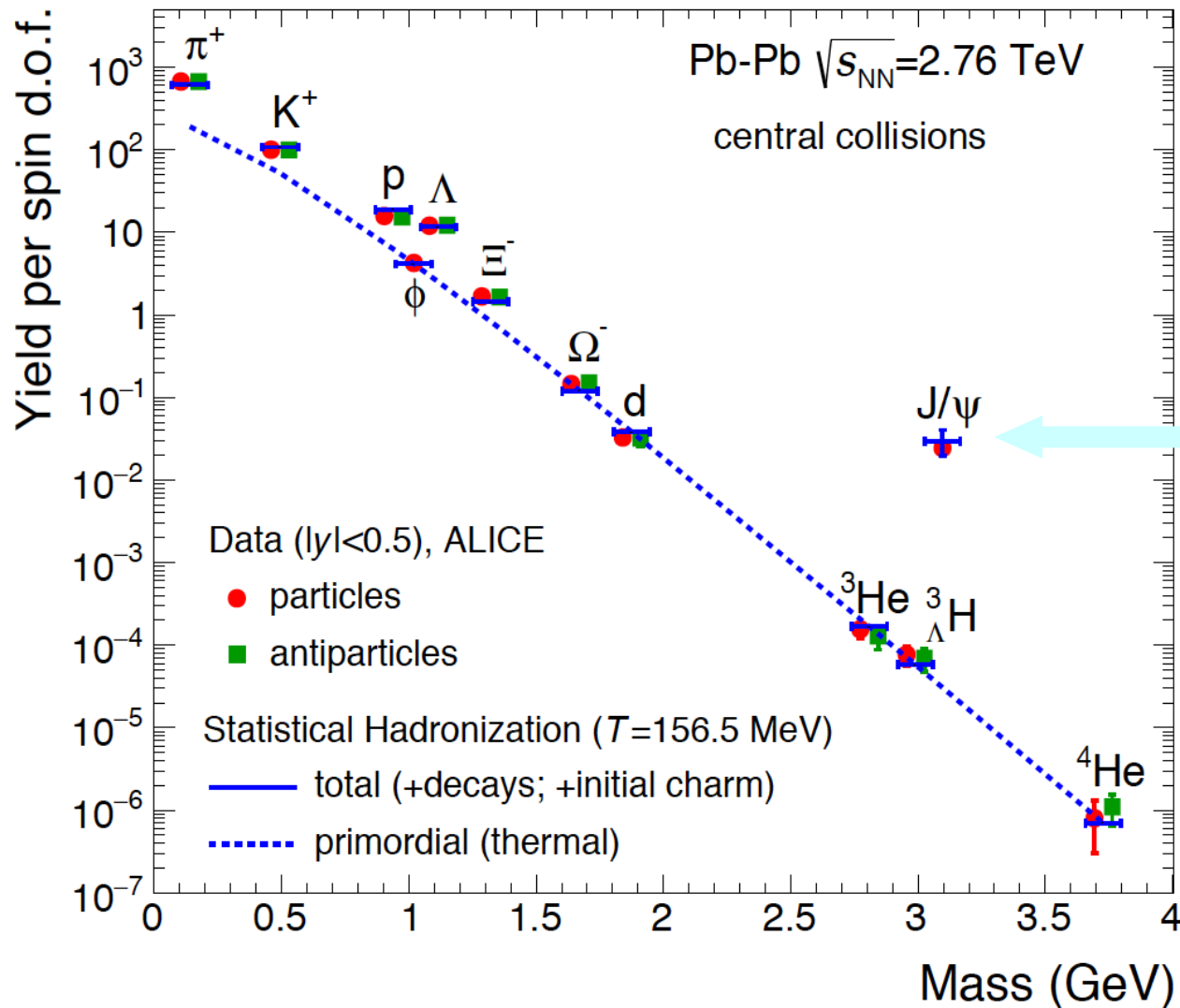
production in PbPb collisions at LHC consistent with deconfinement and subsequent statistical hadronization within present uncertainties  
 main uncertainties for models: open charm cross section, shadowing in Pb

# Systematics of hadron production in SHM



J/psi mass close to hypertriton where does 3 oom enhancement come from?

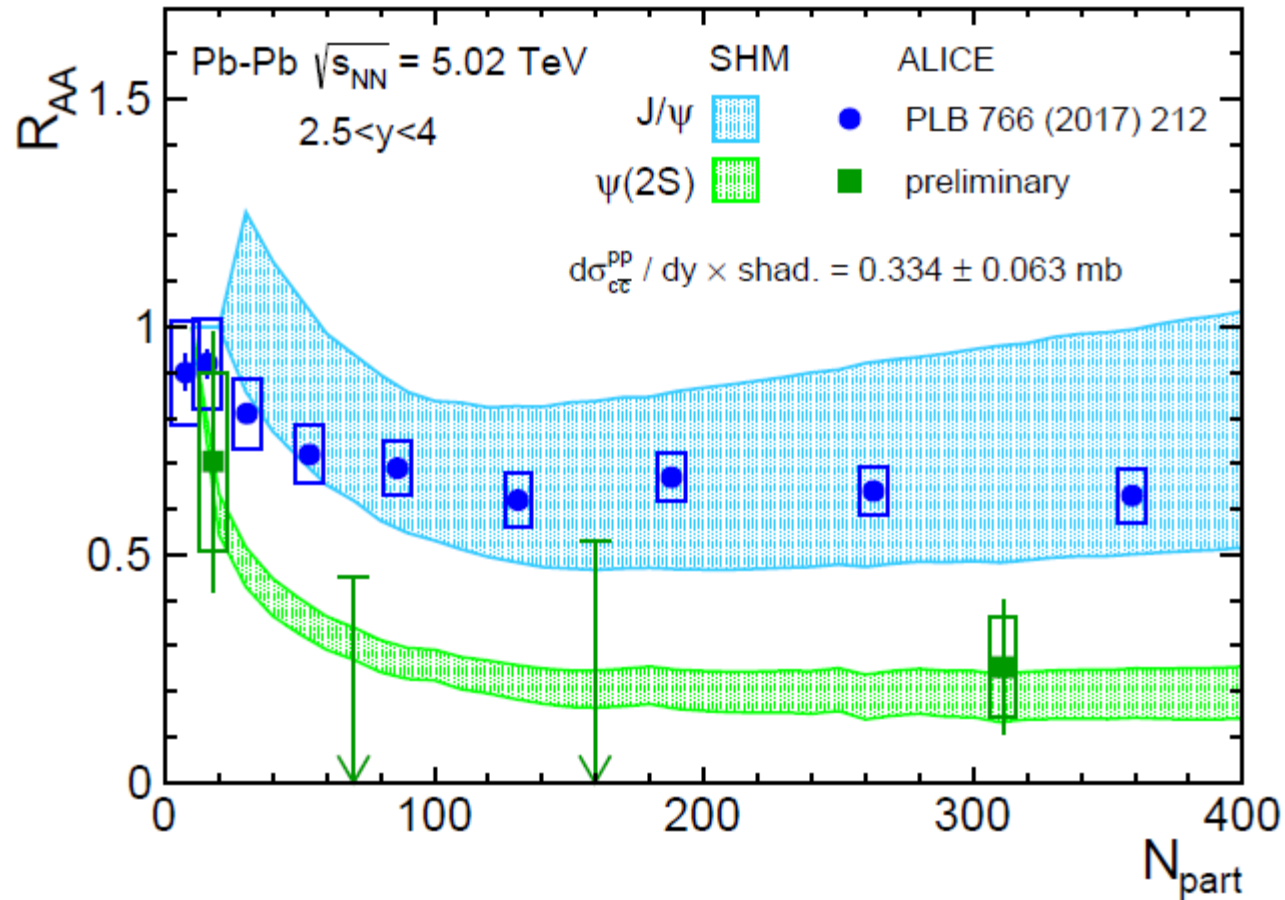
# Systematics of hadron production in SHM



yield exactly reproduced with  
stat hadr. of deconfined and  
thermalized c-quarks from  
initial hard scattering (fugacity)

# What about $\psi(2S)$ ?

M. Köhler, A. Andronic, P. Braun-Munzinger, JS, arXiv:1807.01236



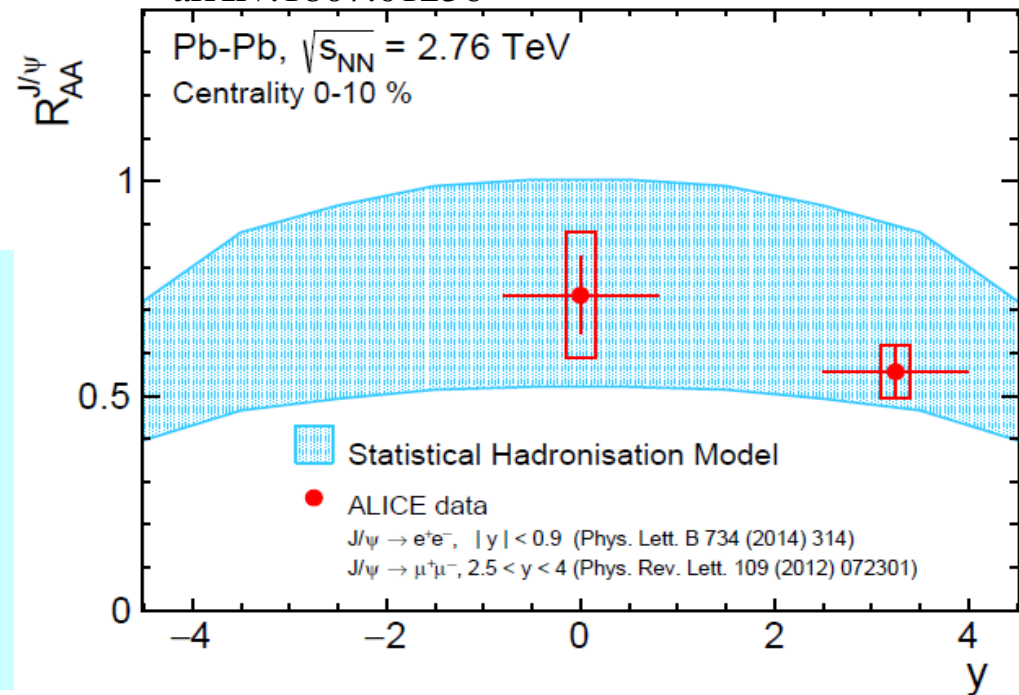
also excited state completely in line, suppressed by Boltzmann factor  
errors will decrease with more data

# Rapidity dependence of $R_{AA}^{J/\psi}$

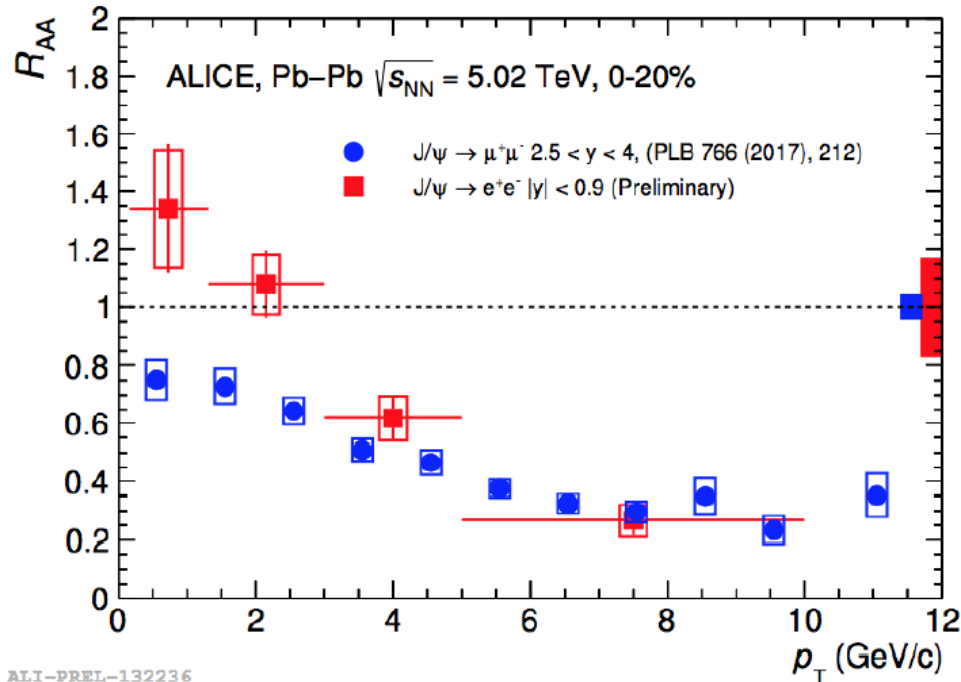
yield in PbPb peaks at mid- $y$   
where energy density is largest  
?

for statistical hadronization  $J/\psi$  yield  
proportional to  $N_c^2$  - higher yield at  
mid-rapidity predicted in line with  
observation  
(at RHIC and LHC)

M. Köhler, A. Andronic, P. Braun-Munzinger, JS  
arXiv:1807.01236



# Transverse momentum dependence



**compared to pp collisions**  
**enhancement at small  $p_t$ !**

– was predicted for statistical hadronization component

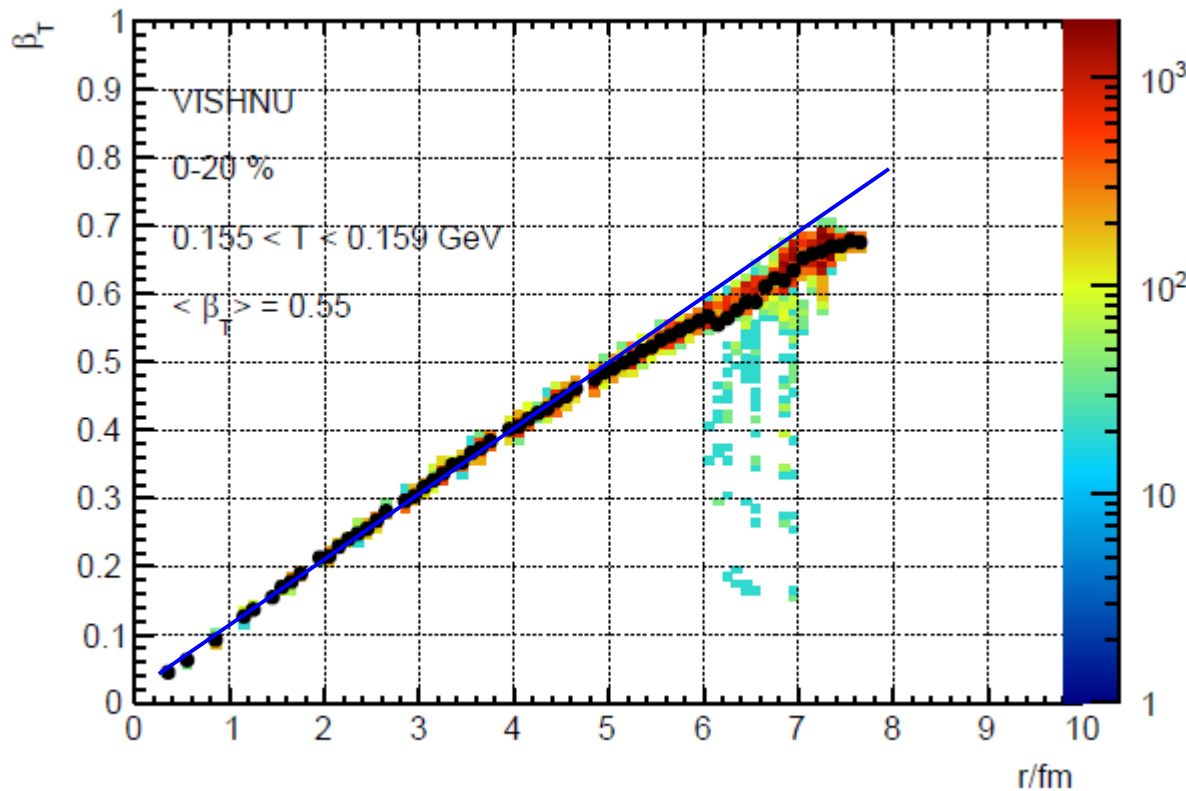
what does statistical hadronization have to say about  $p_t$  spectrum?

the physical picture: charmonia are formed at hadronization from charm quarks in the medium

implies: they should exhibit – as other hadrons – a spectrum characterized by the temperature and the flow of the surrounding medium

recipe: take flow characteristics at  $T_c$  from a good hydro describing the other light flavor observables, normalization given by  $c\bar{c}$  cross section

# Transverse hydro velocity profile at $T_c$

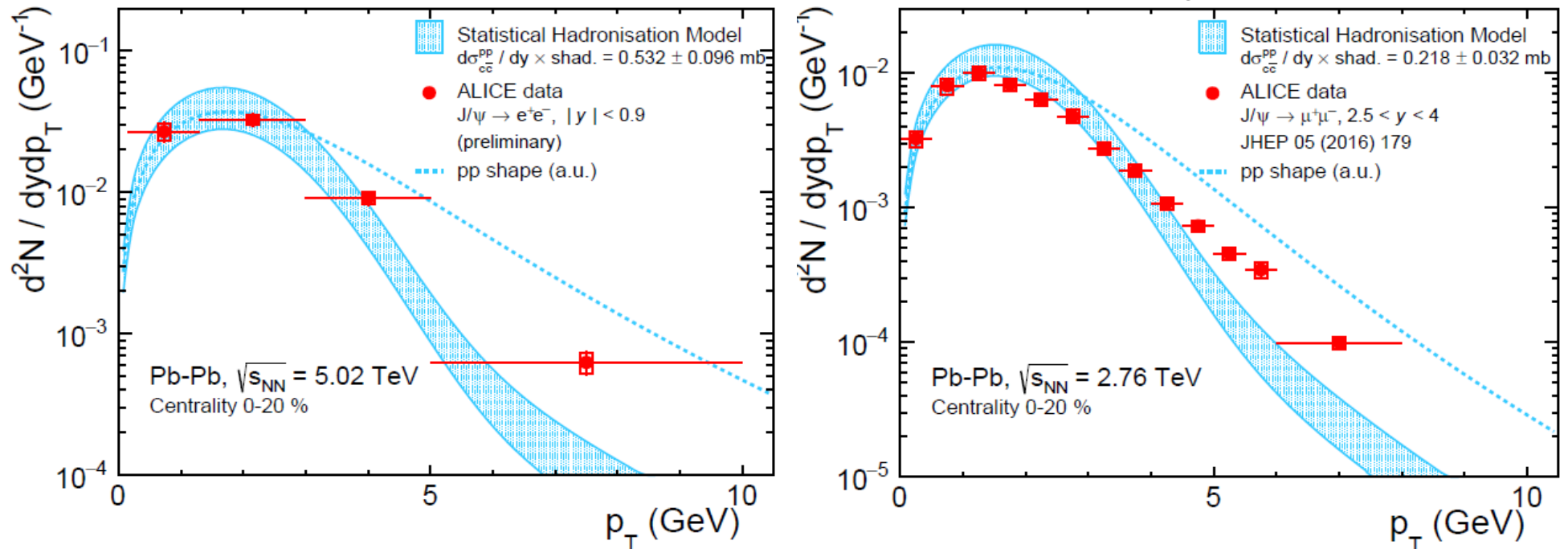


- velocity profile linear in  $r$
- average transverse velocity:  
 $0.55 c$

first approach: use blast wave parameterization with hydro input, i.e. linear velocity profile and correct mean velocity and  $T=T_c$  and  $m=m(J/\psi)$  for core and  $pp$  spectrum for corona

# J/ψ transverse momentum spectra from stat. hadr.

M. Köhler, A. Andronic, P. Braun-Munzinger, JS, arXiv:1807.01236



quite reasonable agreement without any free parameters

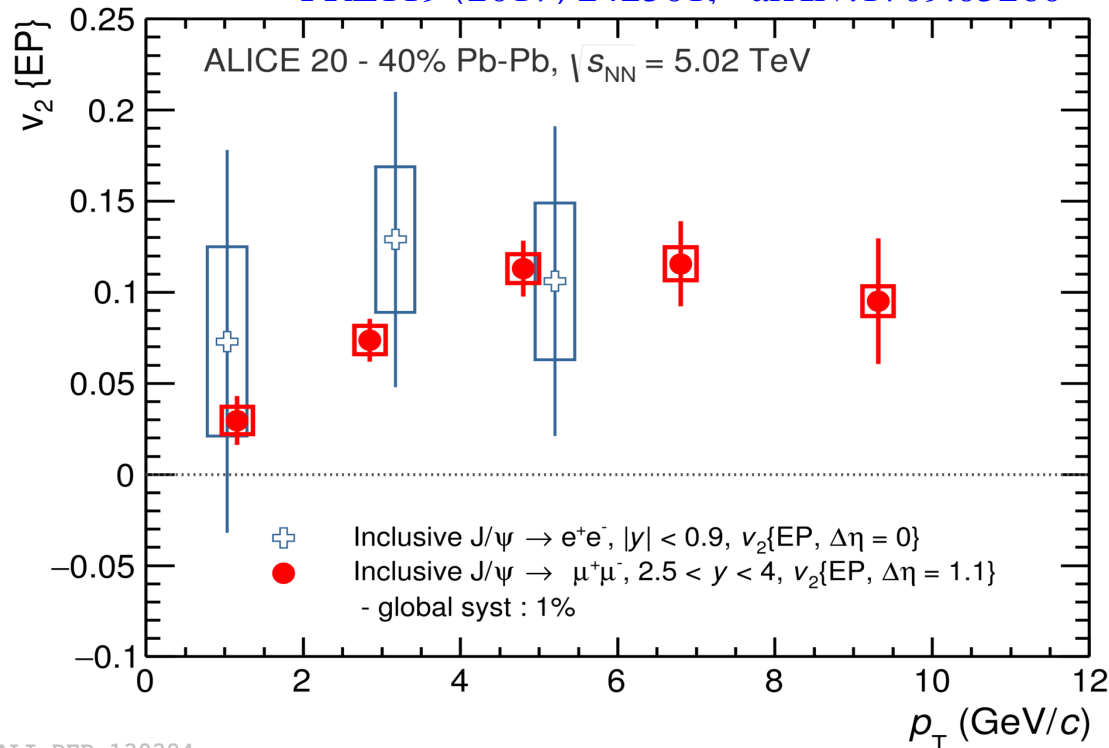
J/ψ formed at hadronization at  $T_c$  from thermalized charm quarks flowing with the rest of the medium



# Elliptic flow of $J/\psi$ vs $p_t$

semi-central collisions: asymmetric overlap region  $\rightarrow$  asym expansion velocity profile  
charm quarks thermalized in the QGP should exhibit the elliptic flow generated in this phase

PRL119 (2017) 242301, arXiv:1709.05260



- expect build-up with  $p_t$  as observed for  $\pi$ , p, K,  $\Lambda$ , ... and vanishing signal for high  $p_t$  region not dominated by flow

first observation of significant  $J/\psi$   $v_2$  in line with expectation from statistical hadronization

# Summary

Hadronization of the QGP delineates the phase boundary as computed with lattice QCD

Even yields of fragile nuclei determined by this temperature

Fluctuations of conserved charges developed as tool to access chiral pseudo-criticality, measurements of higher moments start appearing

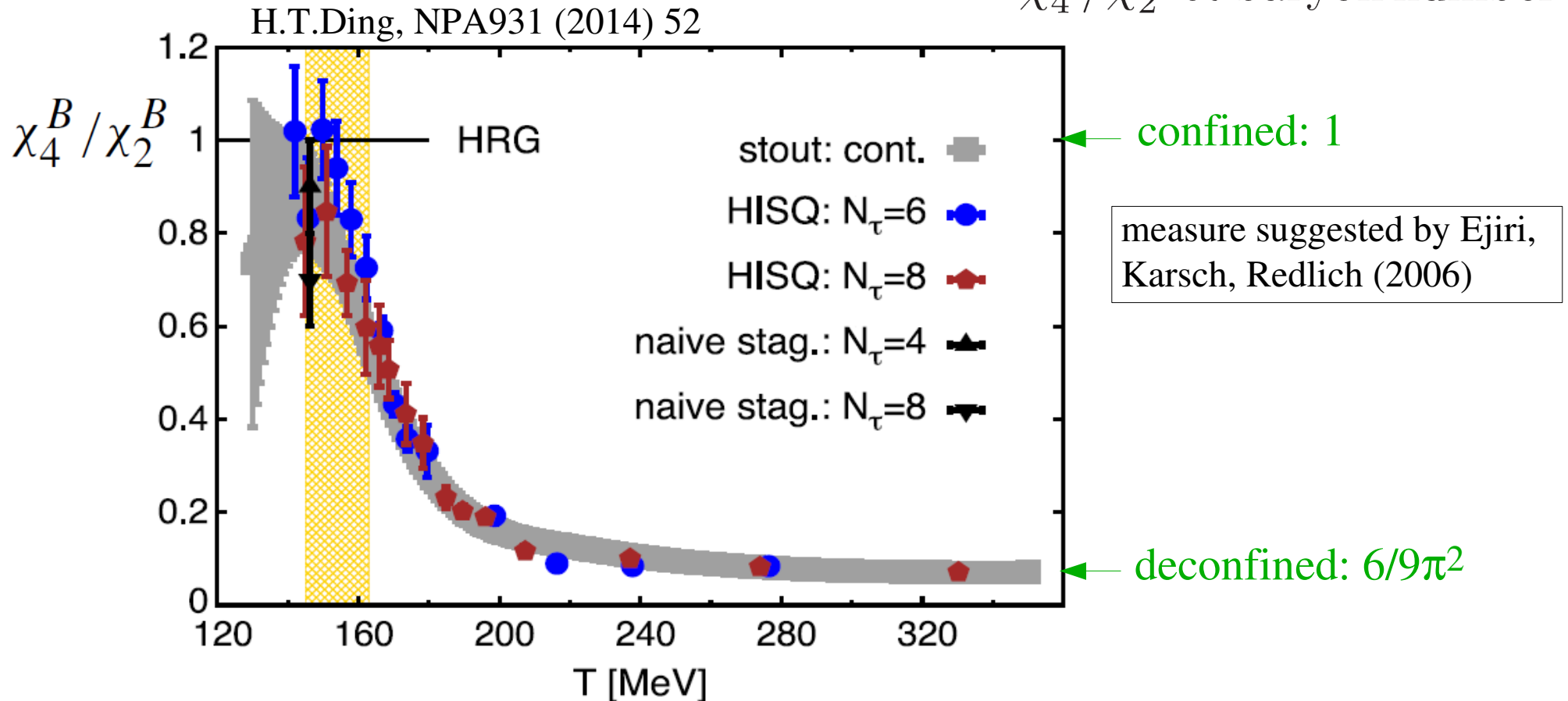
Global, Hubble-like expansion of the nuclear fireball

Charmonia give evidence for deconfinement, formation at hadronization of the fireball together with the rest

backup

# Measure of deconfinement in IQCD

$$\chi_4^B / \chi_2^B \propto \text{baryon number}^2$$



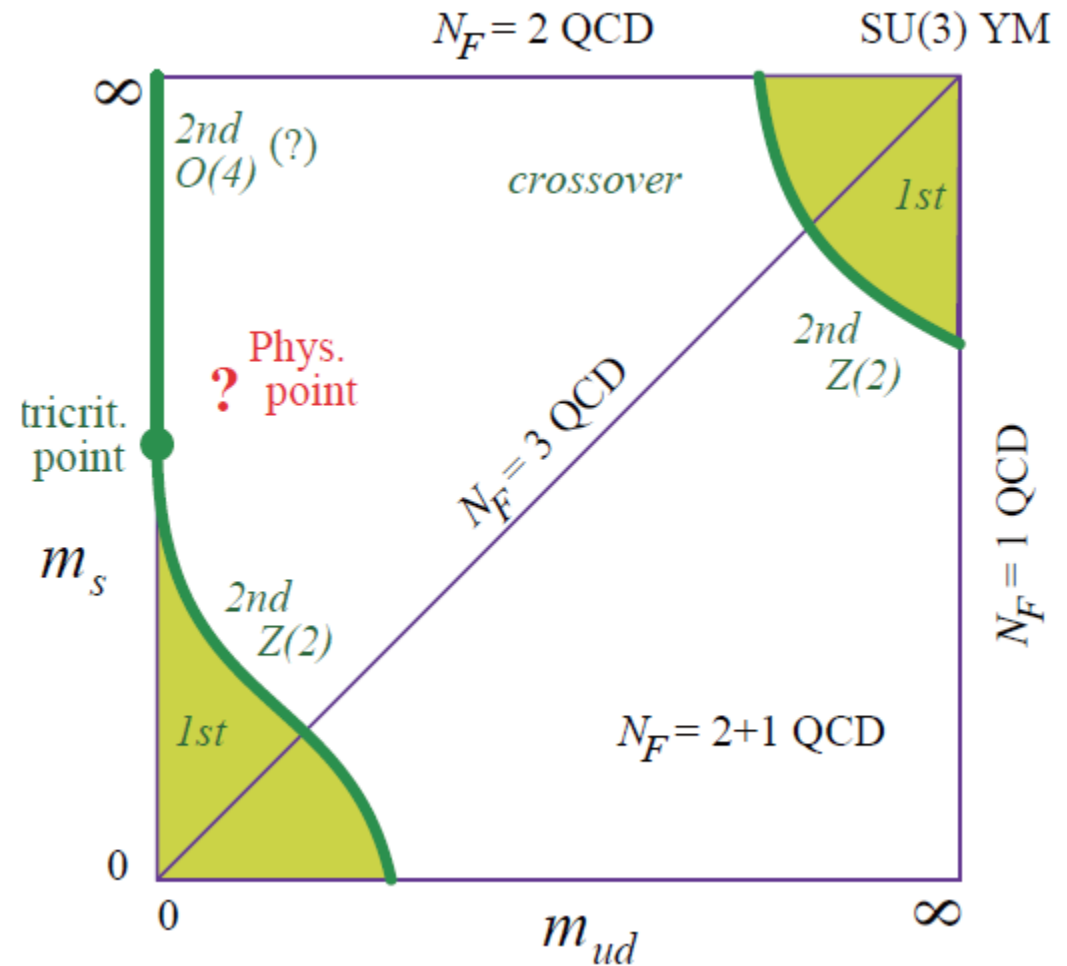
rapid drop suggests: chiral cross over and deconfinement appear in the same narrow temperature range

# Phase diagram of 2+1 flavor QCD from lattice

lQCD finds continuous analytic cross over for physical quark masses

meaning for deconfinement?

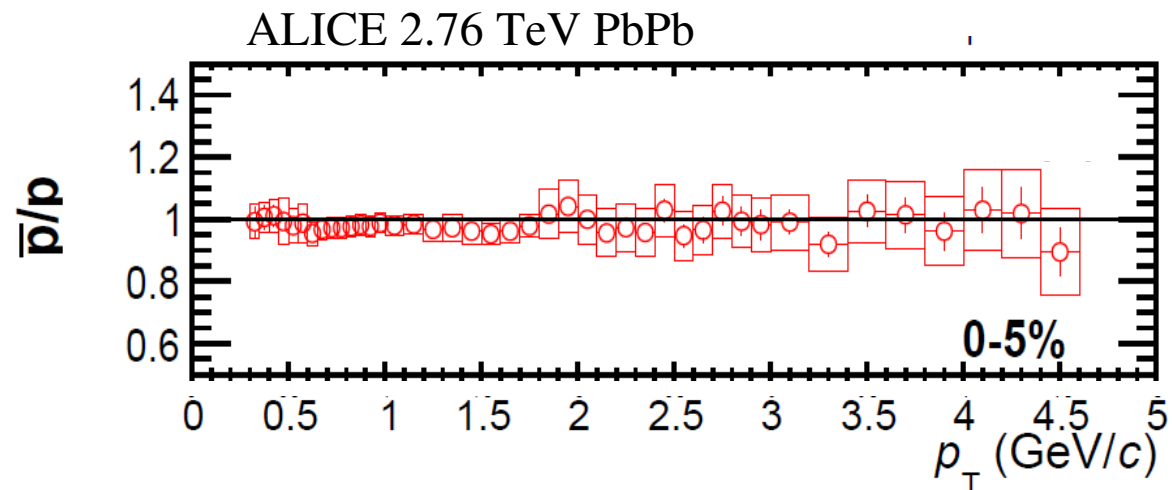
closeness of physical point to 2nd order  $O(4)$  transition could have observable consequences in fluctuations



'Columbia plot' Kanaya, Lattice 2010

# Biggest difference LHC compared to lower energies

- matter and anti-matter produced in equal proportions at LHC
- consistent with net-baryon free central region, ( $\mu_b = 0.7 \pm 3.8$  MeV)  
similar to early universe



- even 10 anti- $^4\text{He}$  nuclei observed!

# a direct comparison of LHC data and lattice QCD

fluctuations of conserved charges (baryon number, strangeness, charge) sensitive to criticality related to spontaneous breaking of chiral symmetry.

- in lQCD susceptibilities exhibit characteristic properties governed by universal part of free energy in vicinity of O(4) critical region of chiral transition

$$\chi_{ijk}^{BQS}(T) = \left. \frac{\partial P(T, \hat{\mu})/T^4}{\partial \hat{\mu}_B^i \partial \hat{\mu}_Q^j \partial \hat{\mu}_S^k} \right|_{\hat{\mu}=0} \quad \text{with } \hat{\mu}_X \equiv \mu_X/T$$

can we see signs of this criticality in experimental data?

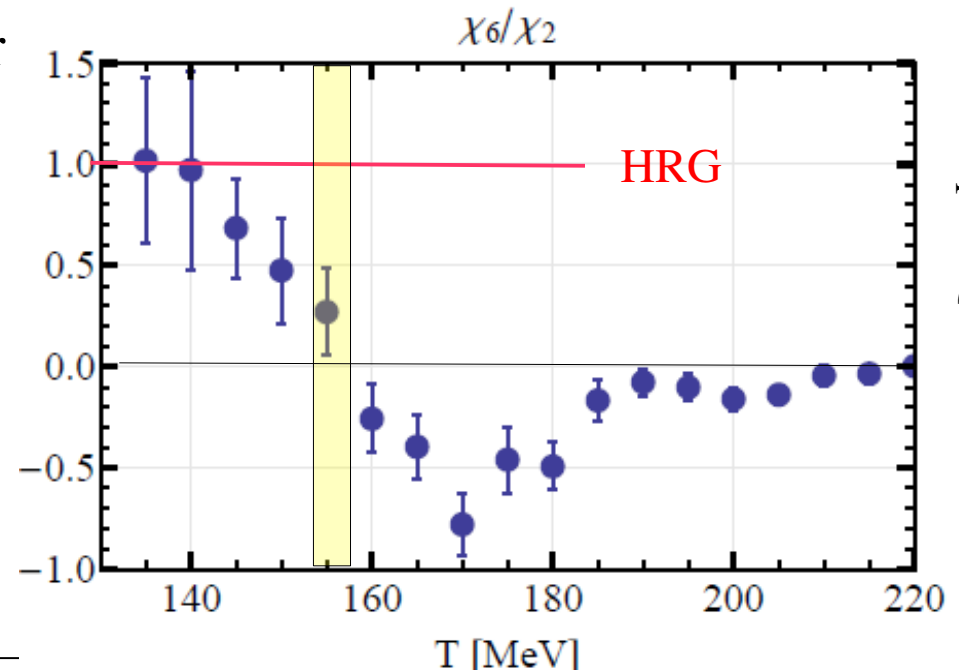
- look at moments of e.g. net baryon number  
 $\Delta N_B = N_B - \bar{N}_B$ ,  $\mu_i = \langle (\Delta N_B - \langle \Delta N_B \rangle)^i \rangle$   
 cumulants of this distribution are directly linked to lQCD susceptibilities

$$\kappa_2 = \mu_2 = VT^3 \chi_2^B$$

$$\kappa_3 = \mu_3 = VT^3 \chi_3^B$$

$$\kappa_4 = \mu_4 - 3\mu_2^2 = VT^3 \chi_4^B$$

....



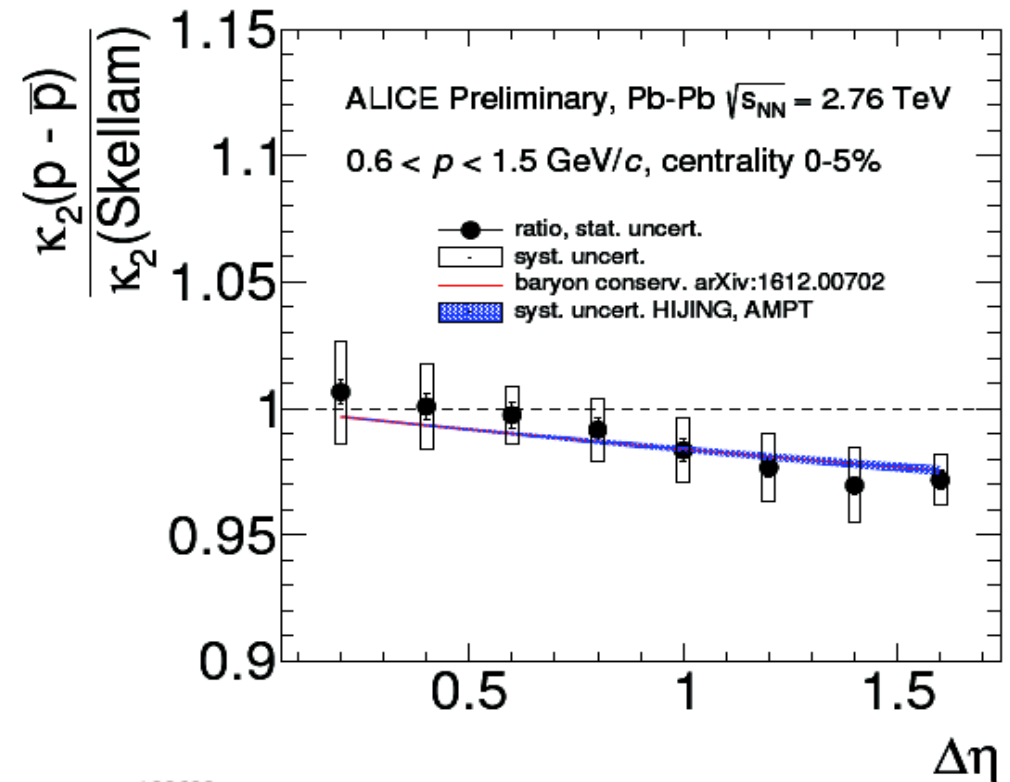
# Moments of net proton distribution

take net proton distribution as a proxy for net baryons  
need a number of corrections before comparing to IQCD

- correct for volume fluctuations
- correct for baryon number conservation
- ...

second moment of ALICE net proton distribution completely understood and comparison to IQCD baseline fulfilled

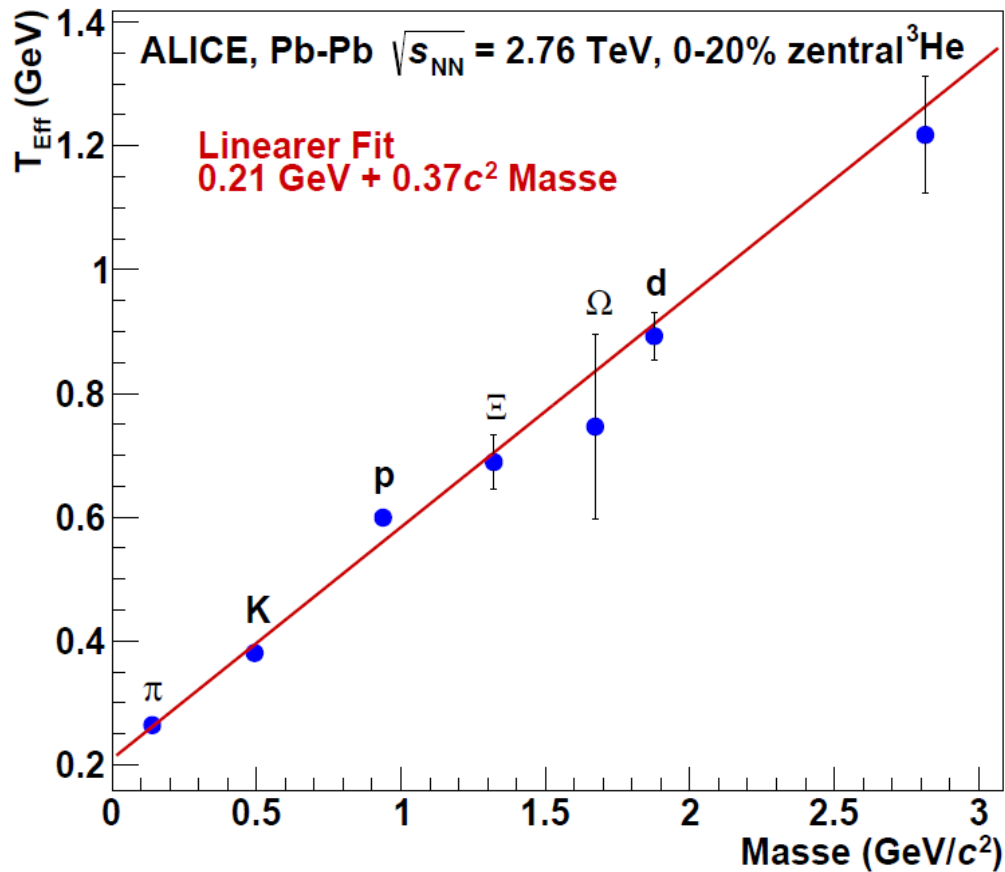
higher moments very statistics hungry and need very good understanding of all experimental fluctuation (efficiency 3rd and 4th moments from 2018 data up to 6th moment LHC run3)



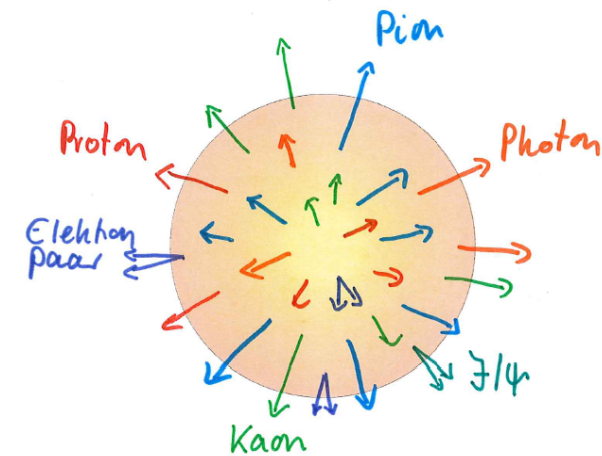
ALI-PREL-122602



# Rapid radial expansion of nuclear fireball

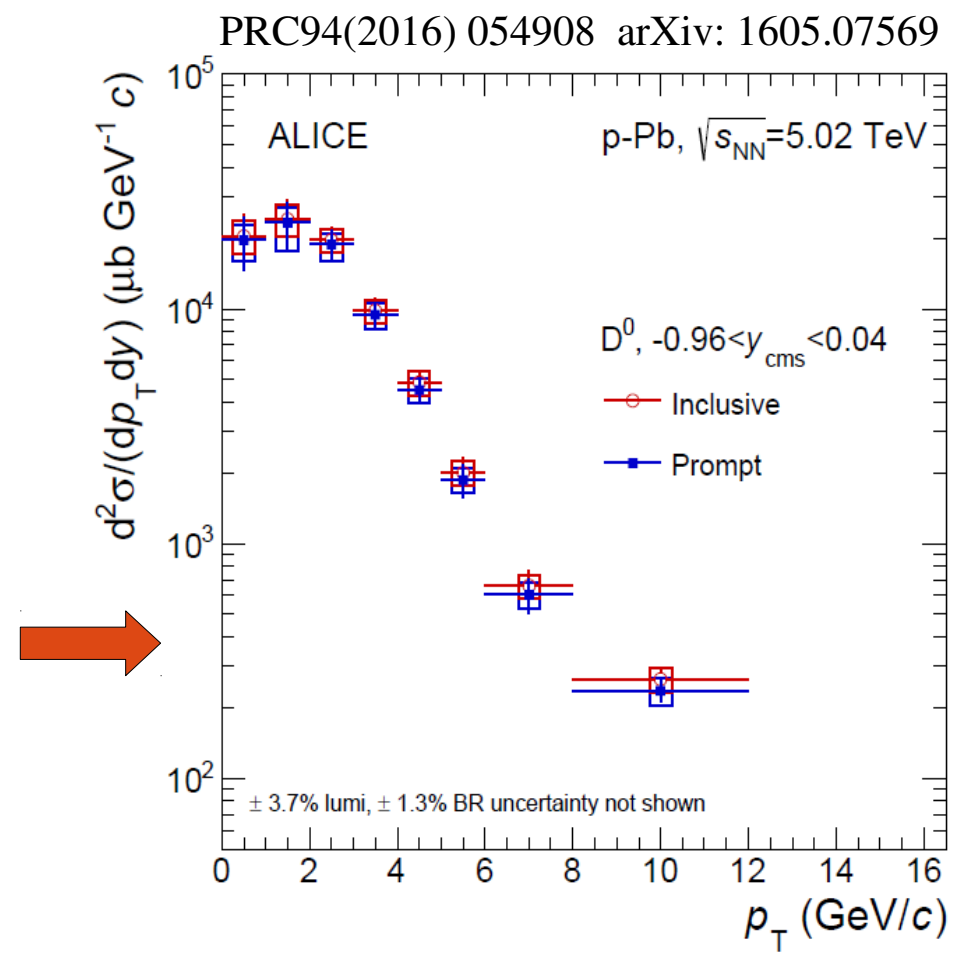
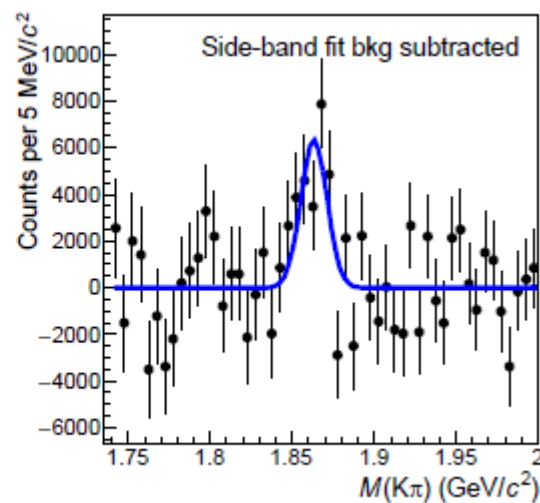
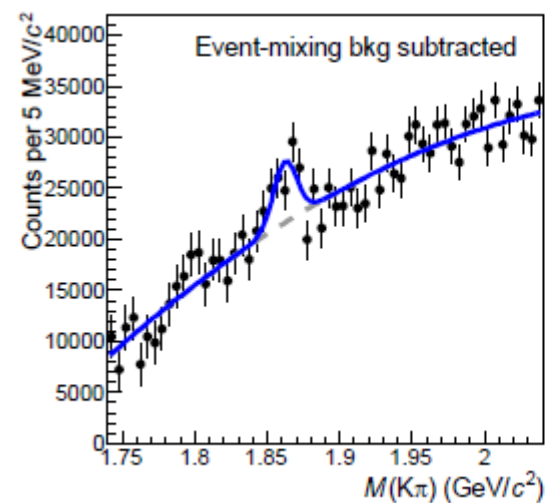
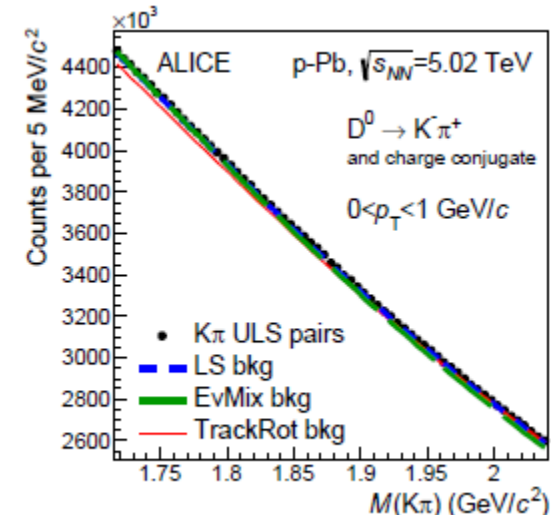


slope constant of spectra  $T_{\text{eff}} \propto m$   
 reflects superposition of random thermal motion and collective expansion  
 at surface velocity  $\frac{3}{4}$  speed of light  
 even fragile objects as deuteron follow radial flow



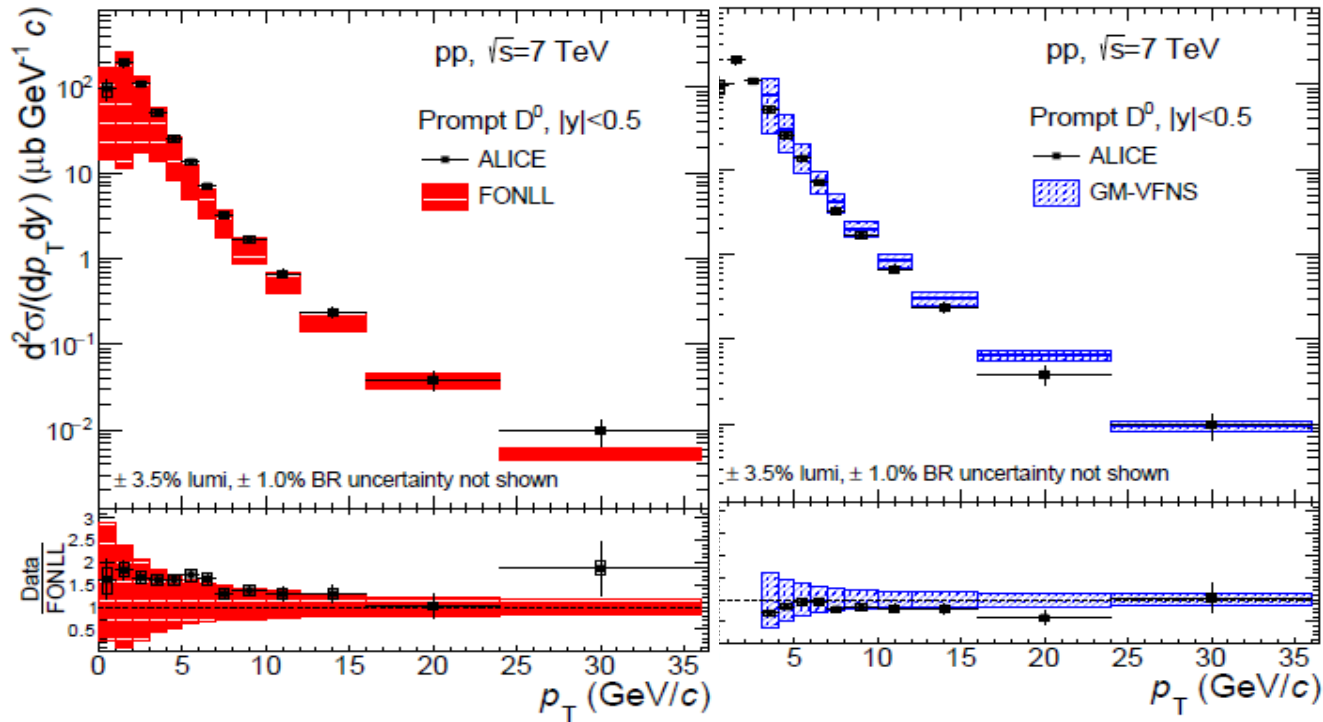
# Production of $c\bar{c}$ - open charm

# first measurements of open charm down to $p_t = 0$ at $y=0$



very hard struggle to deal with (irreducible) combinatorial background,  
 very recently successful for  $D^0$  in pp and pPb

# measurements in pp at 7 TeV agree well with state of the art pQCD calculations



ALICE: 1702.00766  
 FONLL: Cacciari et al., arXiv:1205.6344  
 GM-VFNS: Kniesl et al., arXiv:1202.0439

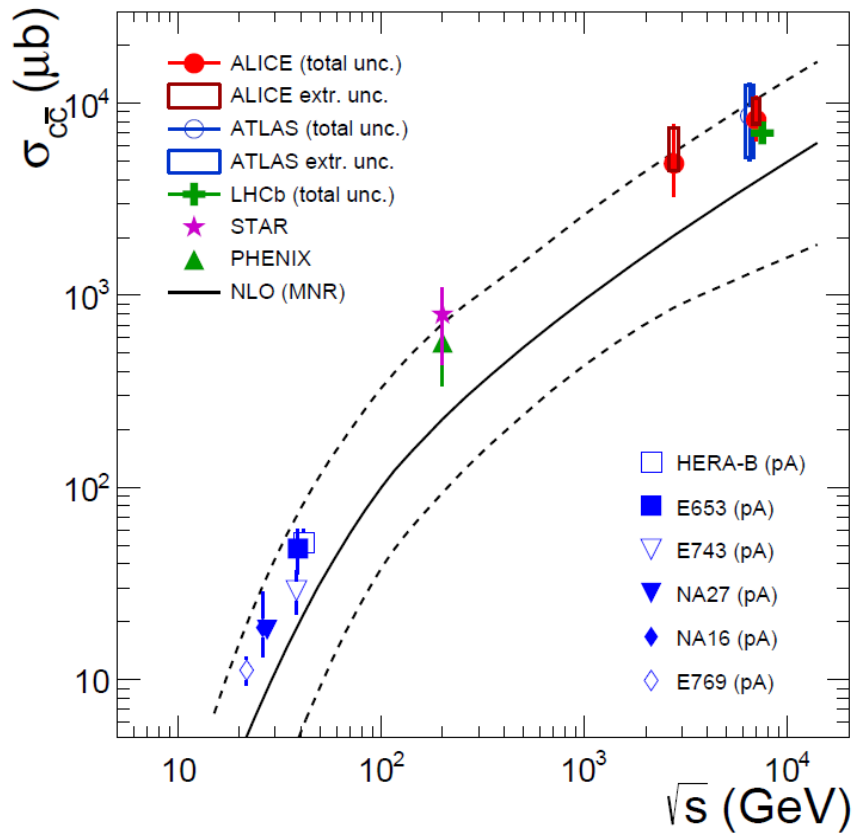
data are compared to perturbative QCD calculations  
 reasonable agreement  
 - at upper end of FONLL and at lower end of GM-VFNS

mid-y cross sections

	Extr. factor to $p_T > 0$	$d\sigma/dy  _{ y <0.5}$ ( $\mu\text{b}$ )
$D^0$	$1.0002^{+0.0004}_{-0.0002}$	$512 \pm 37(\text{stat}) \pm 39(\text{syst}) \pm 18(\text{lumi}) \pm 5(\text{BR})$
$D^+$	$1.25^{+0.29}_{-0.09}$	$235 \pm 19(\text{stat}) \pm 26(\text{syst}) \pm 8(\text{lumi}) \pm 6(\text{BR})^{+54}_{-16}(\text{extrap})$
$D^{*+}$	$1.21^{+0.28}_{-0.08}$	$251 \pm 29(\text{stat}) \pm 24(\text{syst}) \pm 9(\text{lumi}) \pm 3(\text{BR})^{+58}_{-16}(\text{extrap})$
$D_s^+$	$2.23^{+0.71}_{-0.65}$	$89 \pm 18(\text{stat}) \pm 11(\text{syst}) \pm 3(\text{lumi}) \pm 3(\text{BR})^{+28}_{-26}(\text{extrap})$

# currently best measurement of the total $c\bar{c}$ cross section in pp at LHC

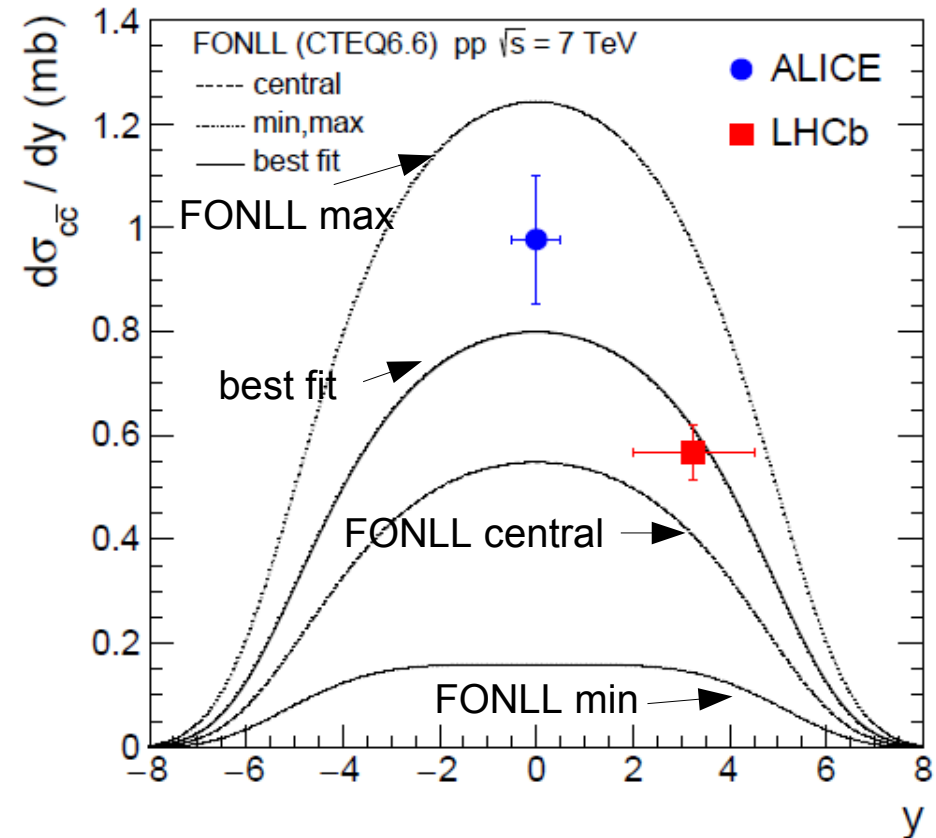
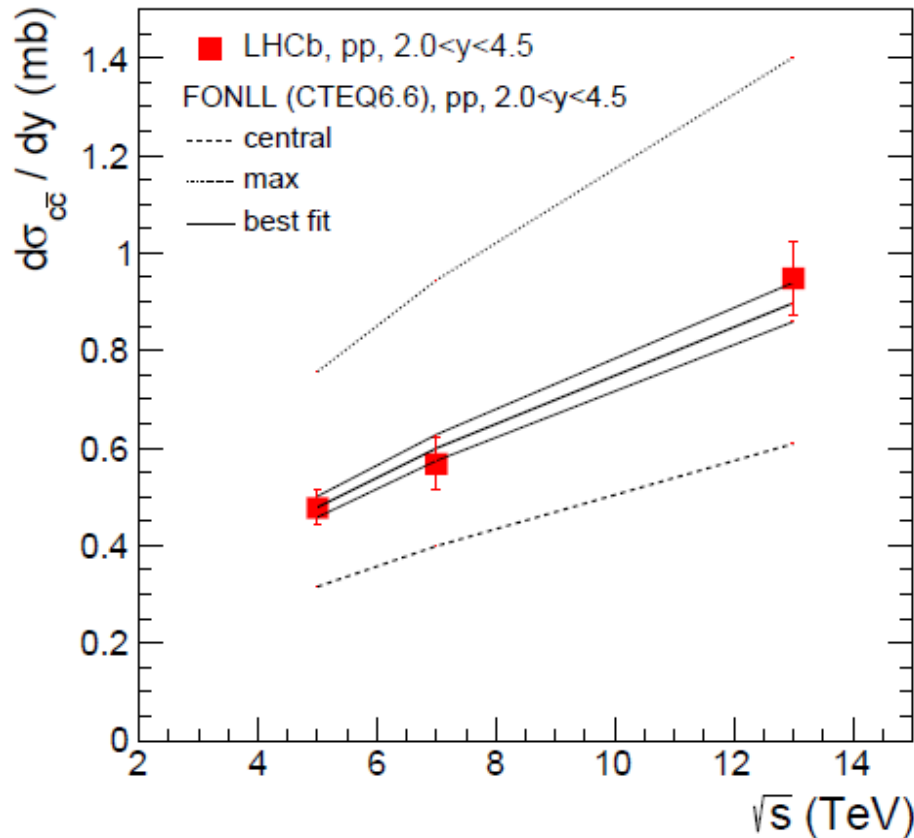
PRC94(2016) 054908 arXiv: 1605.07569



- cross sections in good agreement with NLO pQCD (at upper end of band but well within uncertainty)
- beam energy dependence follows well NLO pQCD

# the baseline for the interpretation of PbPb data

use shape of FONLL to interpolate to proper  $\sqrt{s}$  and  $y$ -interval

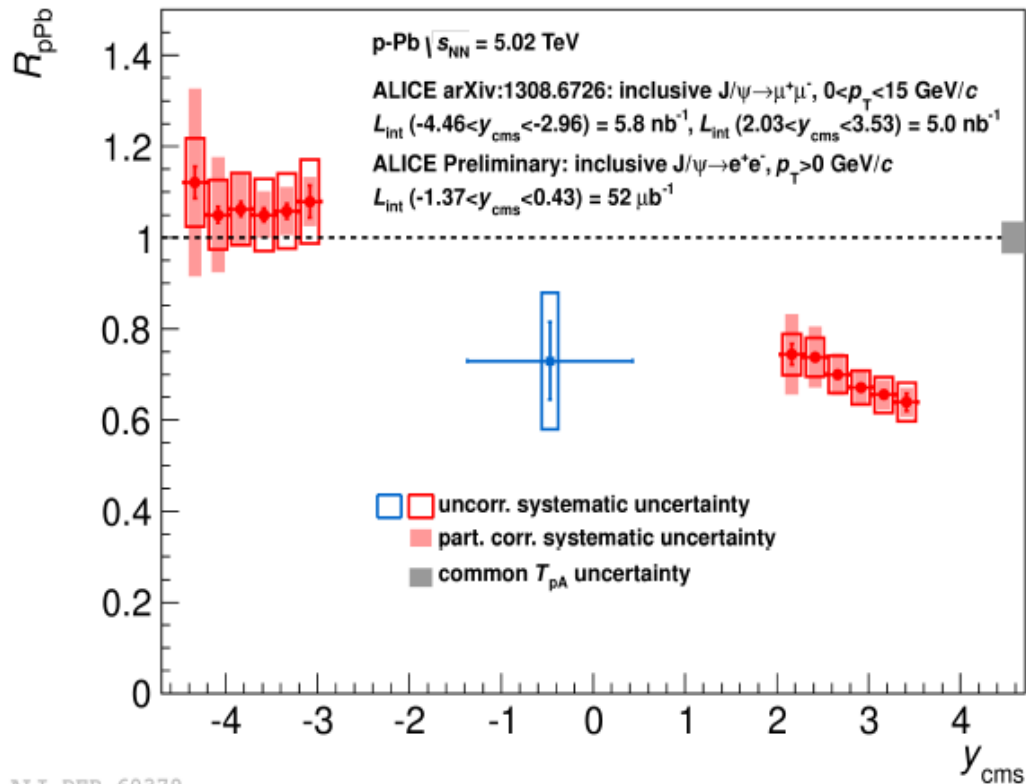


A. Andronic priv. Comm.

LHCb: 5 TeV arXiv:1610.02230  
 7 TeV NPB 871 (2013) 1  
 13 TeV JHEP 03 (2016) 159  
 plus erratum

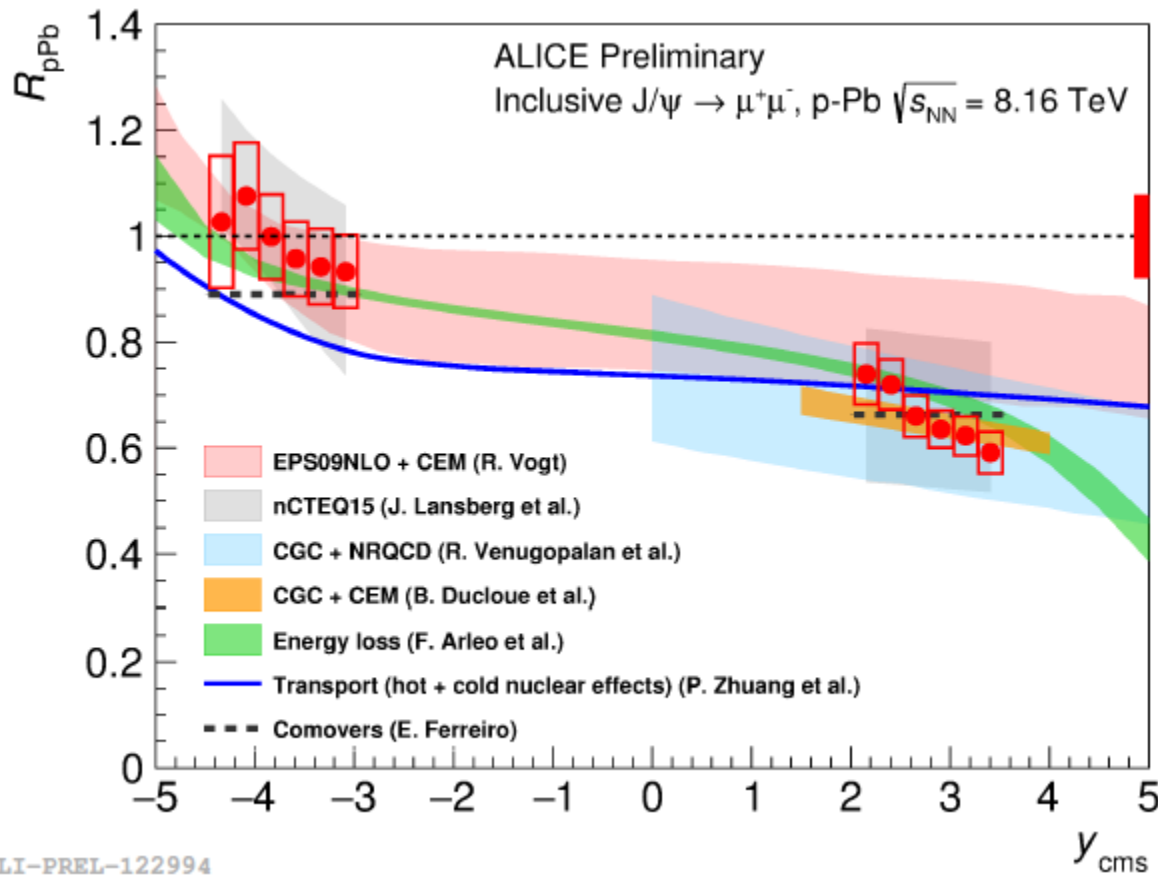
ALICE: 7 TeV PRC94(2016) 054908  
 and 1702.00766

# J/psi rapidity distribution in pPb compared to pp



ALICE forward/backward arXiv:1308.6726  
good agreement with LHCb arXiv:1308.6729  
ALICE mid-y hard probes 2013

# J/psi rapidity distribution in pPb compared to pp



ALICE new 8.16 TeV data

good agreement with shadowing calculations  
also with energy loss models wo shadowing  
and CGC calculation

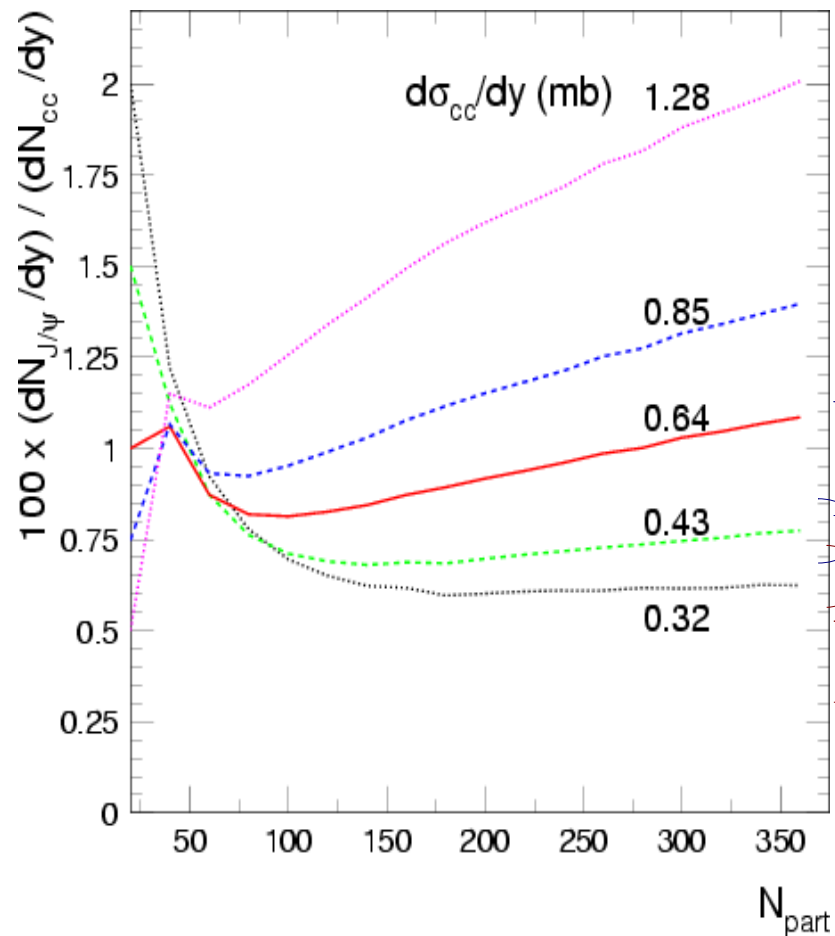
pp open charm  $d\sigma/dy$  plus  
nuclear effects from J/ψ in pPb  
form current baseline for  
charmonia in PbPb



# Expectations for LHC from measured $c\bar{c}$ cross section in pp collisions

A. Andronic, P. Braun-Munzinger, K. Redlich, J. Stachel Phys. Lett. B652 (2007) 259

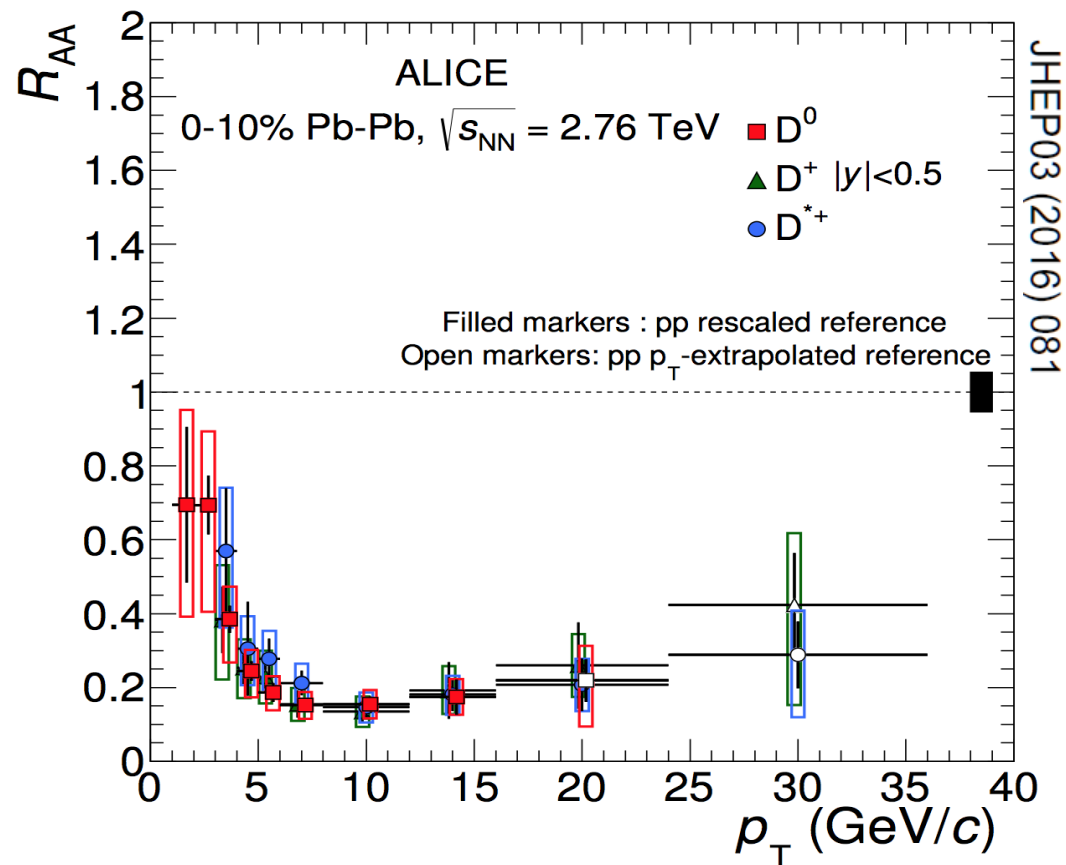
measured  $c\bar{c}$  cross sections at appropriate rapidity by ALICE and LHCb and shadowing from measured J/psi production in pPb collisions compared to pQCD



mid-y LHC 2.76  
and 5.02 TeV  
including shadowing

forward-y LHC 2.76  
and 5.02 TeV  
including shadowing

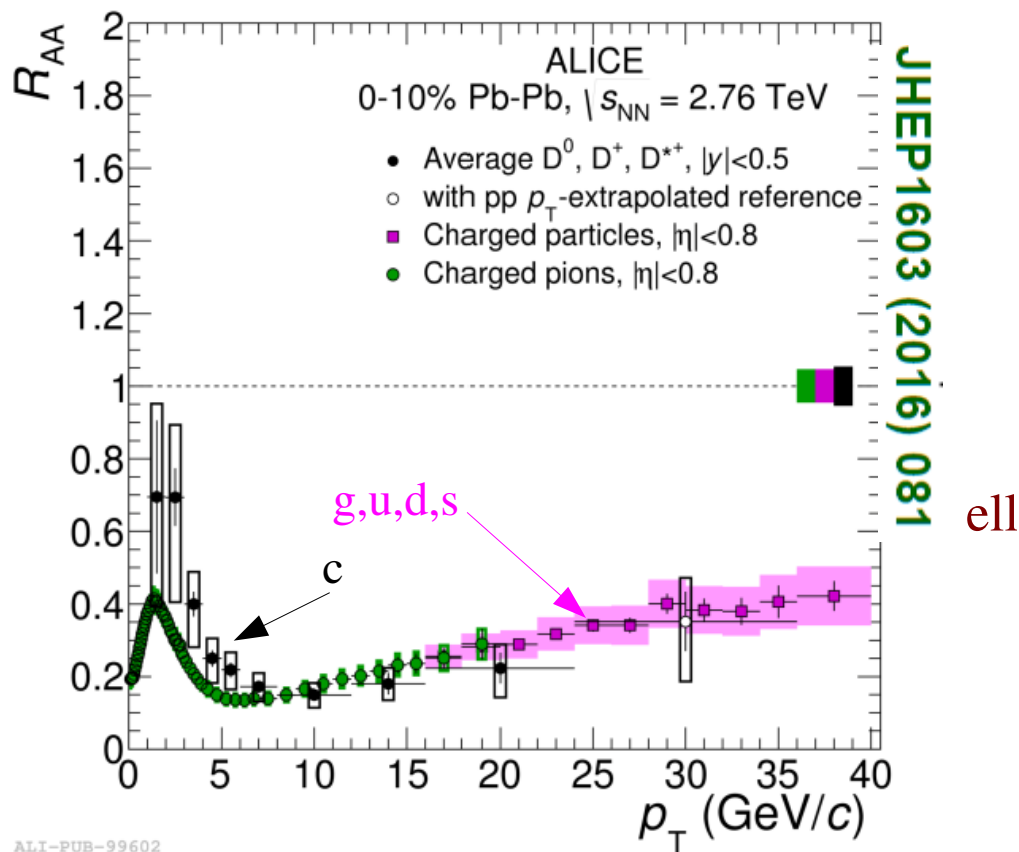
# suppression of charm at LHC energy



energy loss for all species of D-mesons within errors equal - not trivial  
energy loss of central collisions very significant - suppr. factor 5 for 5-15 GeV/c

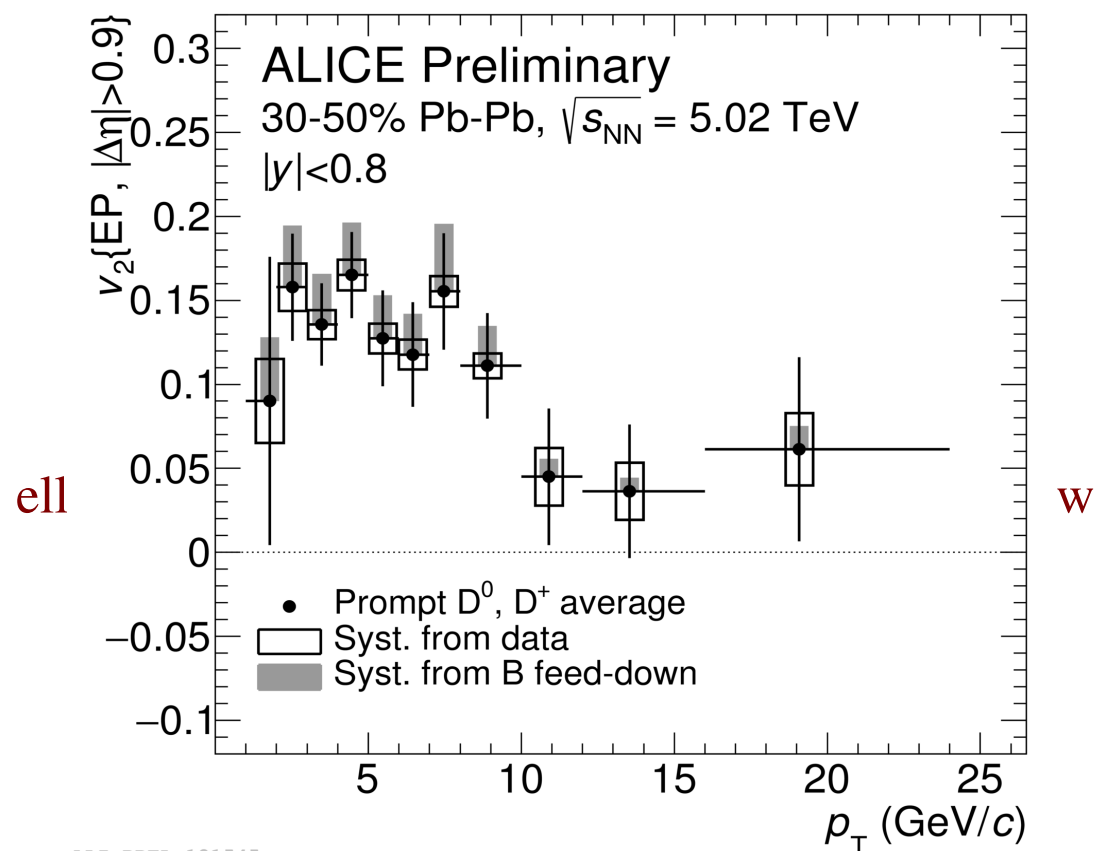
# charm quarks thermalize to large degree in QGP

strong energy loss of charm quarks



ALI-PUB-99602

elliptic flow for charm – participation in coll. flow



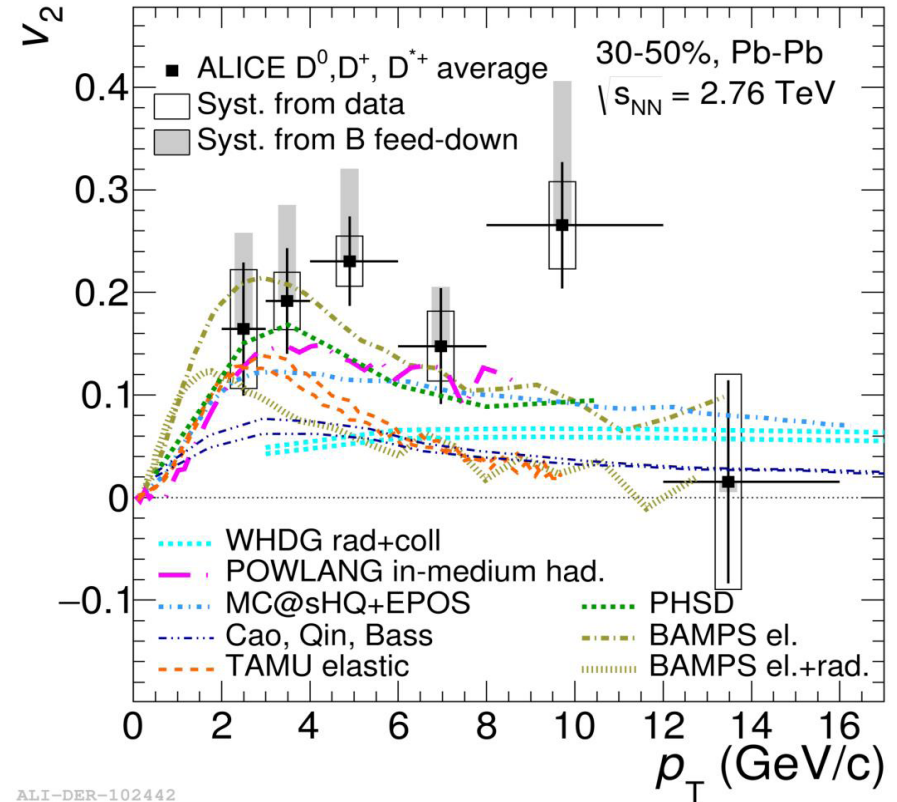
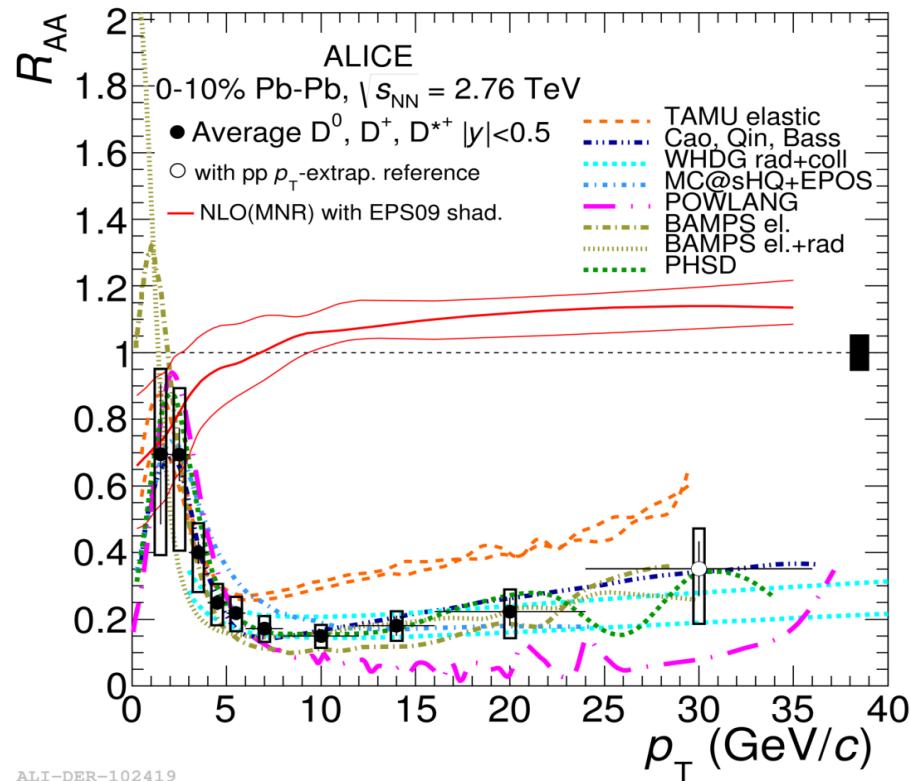
ALI-PREL-121545

M.Djordjevic, arXiv:1307.4098:  
equal  $R_{AA}$  is a conspiracy of different  
fragmentation functions of light quarks,  
gluons, charm and different color factors in  
energy loss



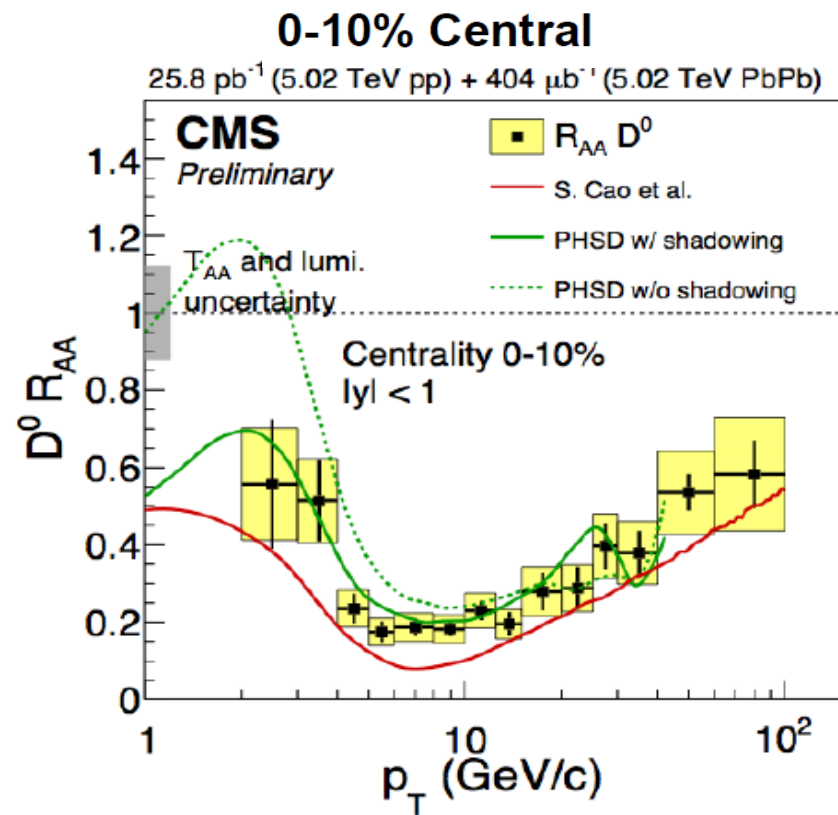
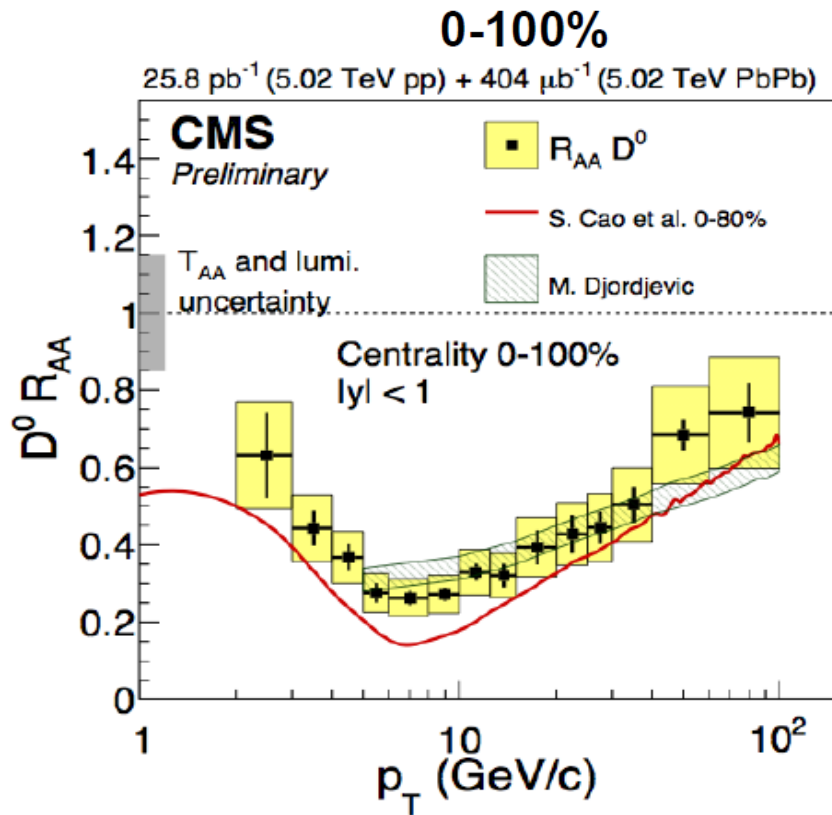
# models constrained by simultaneous fit of $R_{AA}$ and $v_2$

PRL 111 (2013) 102301, PRC 90 (2014) 034904



- models capture various relevant aspects leading to thermalization of charm
- serious need to put together a coherent picture
  - a difficult theoretical challenge, that is being addressed
  - recently an EMMI rapid reaction task force took up the issue (Andronic, Averbeck, Gossiaux, Masciocchi, Rapp)

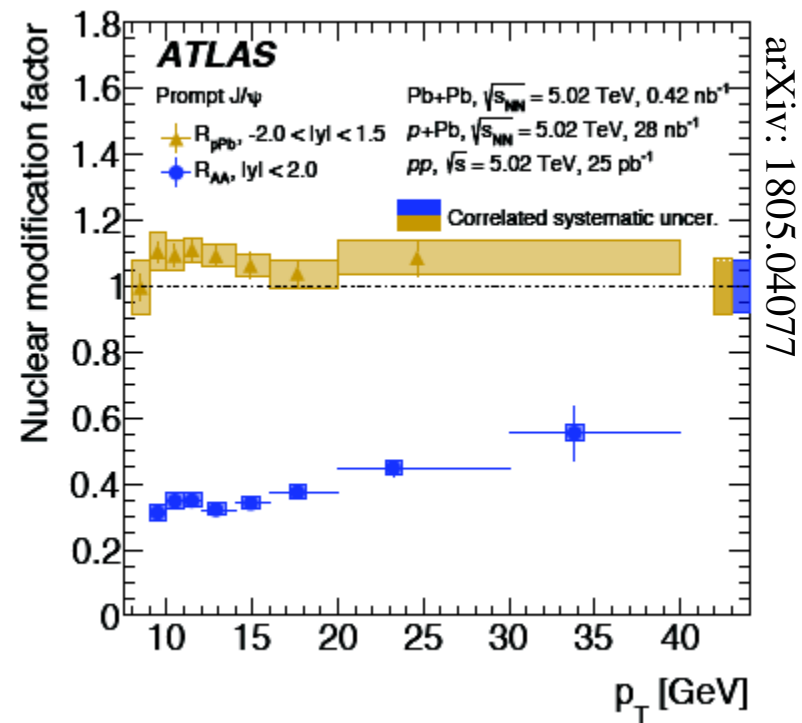
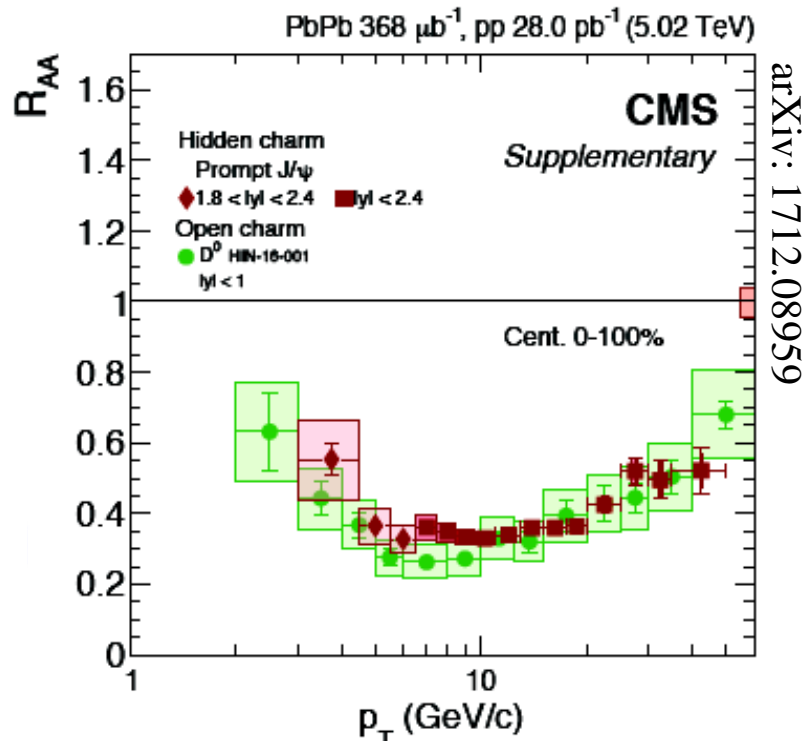
# D<sup>0</sup> R<sub>AA</sub> compared to models



models: predictions before run2 data

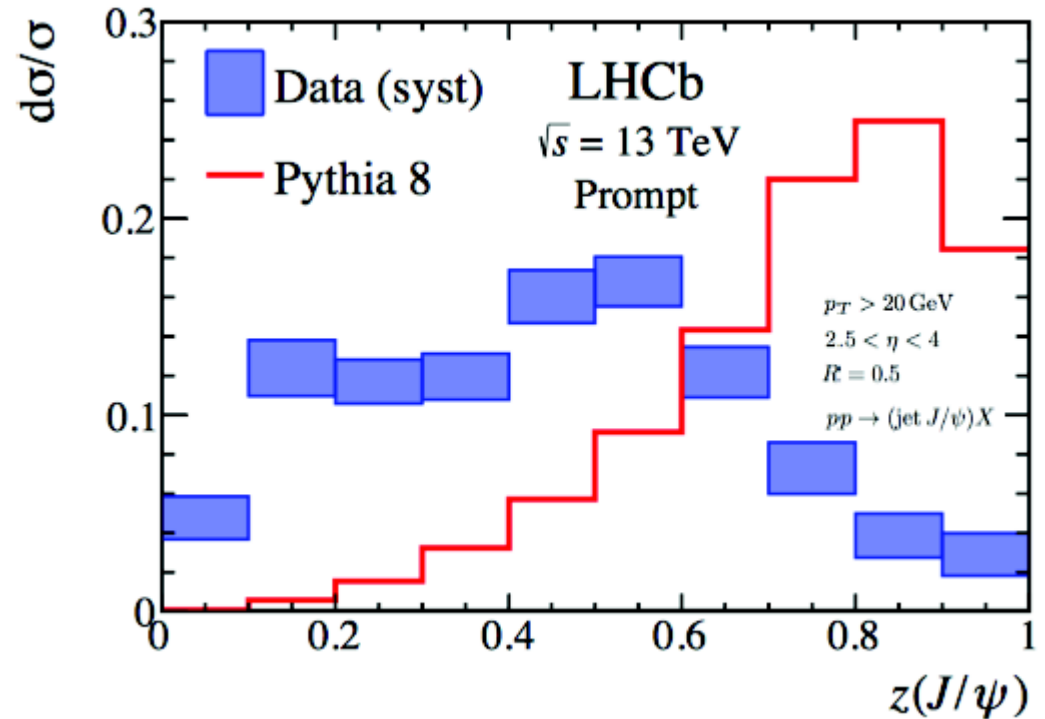
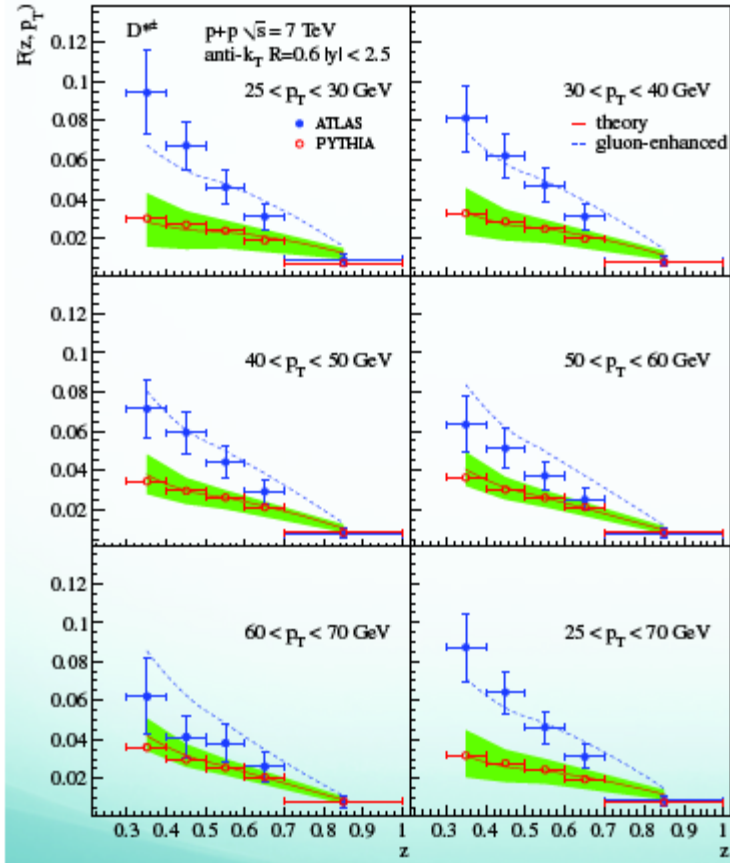
- PHSD (Parton-Hadron-String Dynamics model[2])
- S.Cao et al. ( Linearized Boltzmann transport model + hydro ) arXiv:1605.06447v1
- M. Djordjevic ( QCD medium of finite size with dynamical scattering centers with collisional and radiative energy loss ) Phys. Rev. C 92 (Aug, 2015) 024918

# J/psi $R_{AA}$ in central PbPb like open heavy and light flavor hadrons



prompt J/ψ suppression in PbPb collision –  $R_{AA}$  rising at high  $p_t$   
 same shape and magnitude as D-mesons charged particles  
 J/ψ from gluon fragmentation?

# D meson and J/ψ fragmentation functions surprising



H.Xing (Wuhan 10/2018) : data prefers that jet was initiated by a single parton fragmentation, while PYTHIA starts from a  $c\bar{c}$ bar

Using ZM-VFNS scheme:  
 Chien, Kang, Ringer, Vitev, Xing,  
 1512.06851, JHEP 16

$$- - - D_g^D(z, \mu) \rightarrow 2D_g^D(z, \mu)$$

Gluon fragmentation into J/ψ could well be the mechanism explaining the high  $p_t$   $R_{AA}$

charmonia



# Charmonia as probe of deconfinement

the original idea (Matsui and Satz 1986): implant charmonia into the QGP and observe their modification (Debye screening of QCD), in terms of suppressed production in nucleus-nucleus collisions with plasma formation

table from H. Satz, J. Phys. G32 (2006) 25

state	$J/\psi$	$\chi_c$	$\psi'$	$\Upsilon$	$\chi_b$	$\Upsilon'$	$\chi'_b$	$\Upsilon''$
mass [GeV]	3.10	3.53	3.68	9.46	9.99	10.02	10.26	10.36
$\Delta E$ [GeV]	0.64	0.20	0.05	1.10	0.67	0.54	0.31	0.20
$\Delta M$ [GeV]	0.02	-0.03	0.03	0.06	-0.06	-0.06	-0.08	-0.07
$r_0$ [fm]	0.50	0.72	0.90	0.28	0.44	0.56	0.68	0.78

in the QGP, the screening length  $\lambda_{\text{Debye}}(T)$  decreases with increasing  $T$ . If  $\lambda_{\text{Debye}}(T) < r_{\text{charmonium}}$  the system becomes unbound



→ notion of charmonia as thermometer – sequential melting signature of deconfinement, but no direct link to phase boundary

# Charmonia and statistical hadronization

new insight (Braun-Munzinger, J.S. 2000):

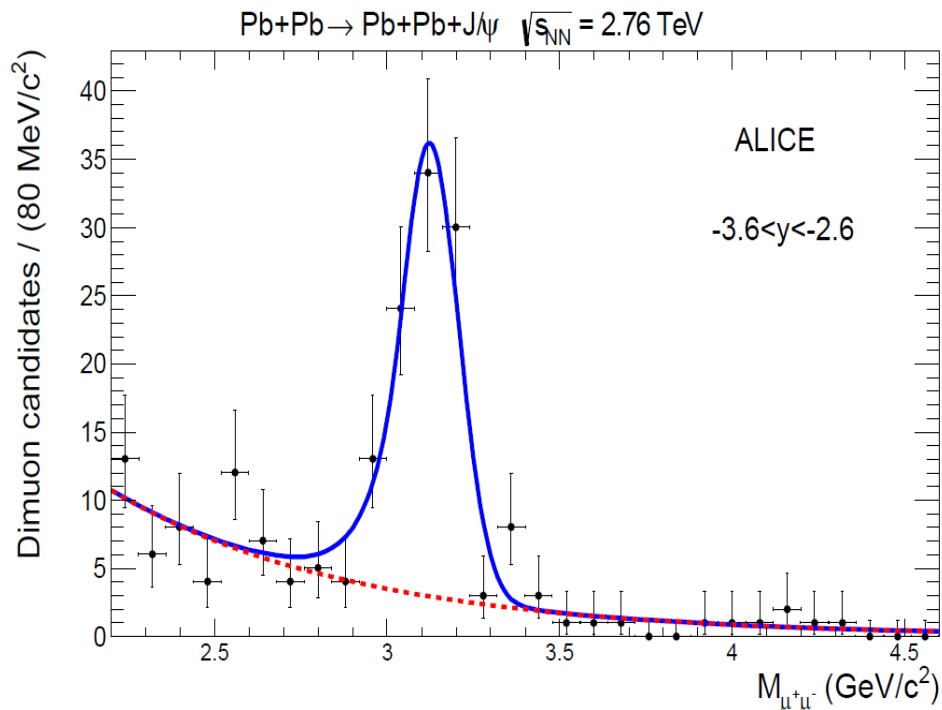
QGP screens all charmonia (as proposed by Matsui and Satz), but charmonium production takes place at the phase boundary,

→ enhanced production at colliders – signal for deconfinement  
production probability from thermalized charm quarks scales with  $N_{c\bar{c}}^2$

yields of charmonia (and open charm hadrons) directly linked to phase boundary and hadronization temperature  
still probe of deconfinement

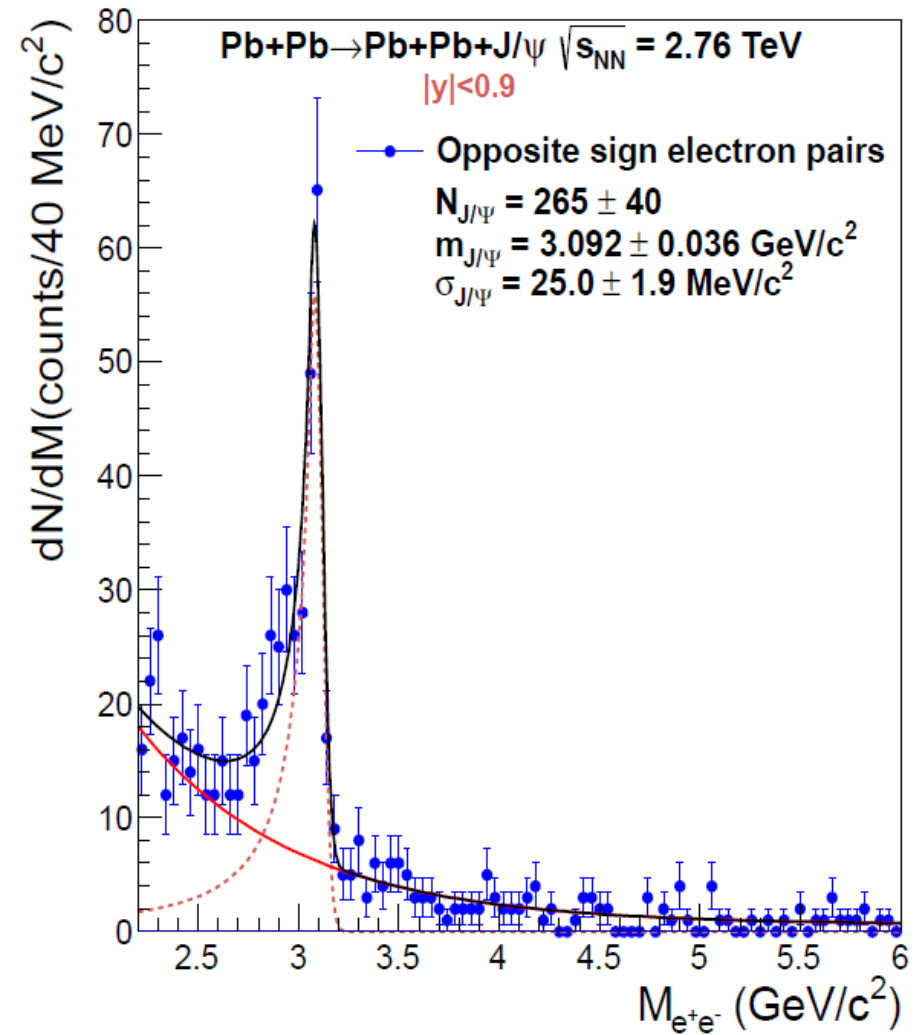
# reconstruction of $J/\psi$ via $\mu^+\mu^-$ and $e^+e^-$ decay

PLB 718 arXiv:1209.3715



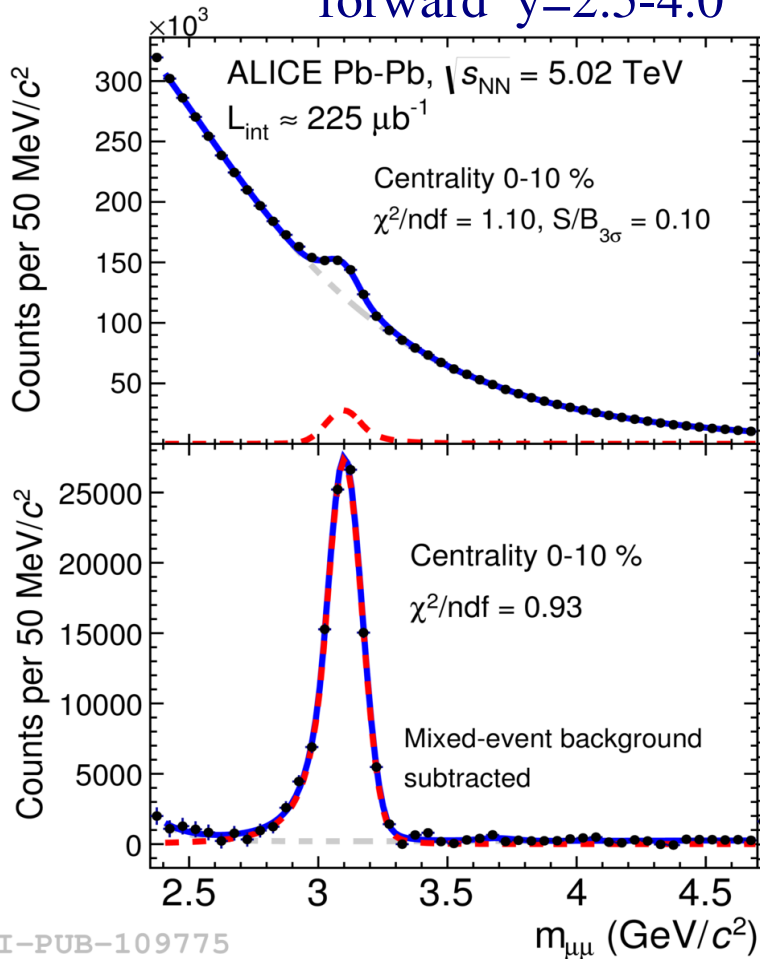
photoproduction in ultra-peripheral PbPb collisions – excellent signal to background  
very good understanding of line shape  
(probes nuclear gluon shadowing, not discussed here)

ALICE EPJ C73 arXiv:1305.1467

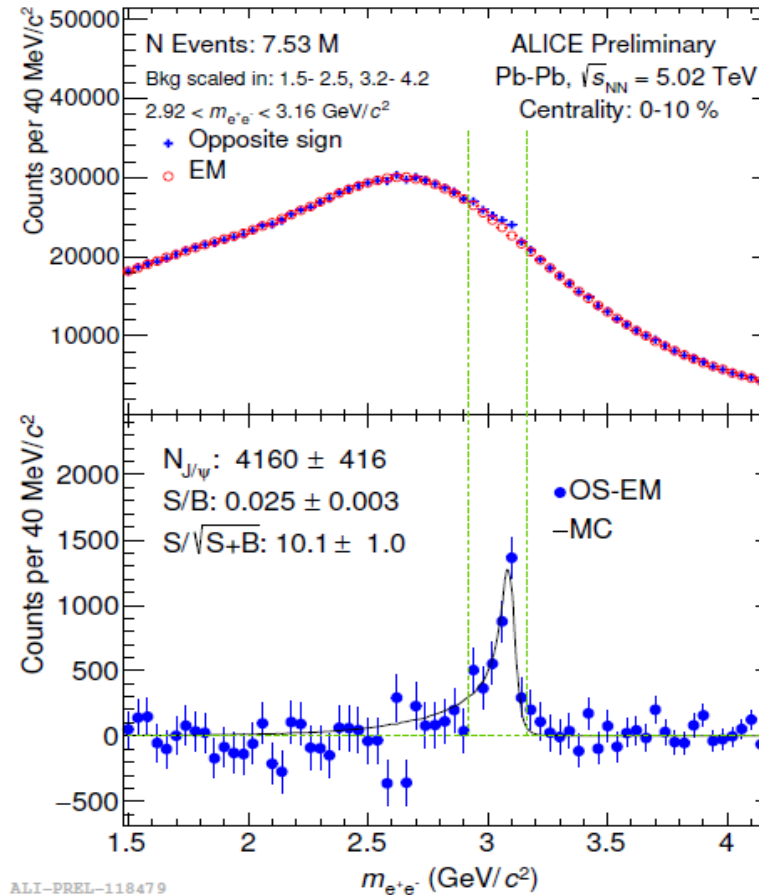


# Reconstruction of $J/\psi$ via $\mu^+\mu^-$ and $e^+e^-$ decays

forward  $y=2.5-4.0$



mid  $|y| < 0.8$

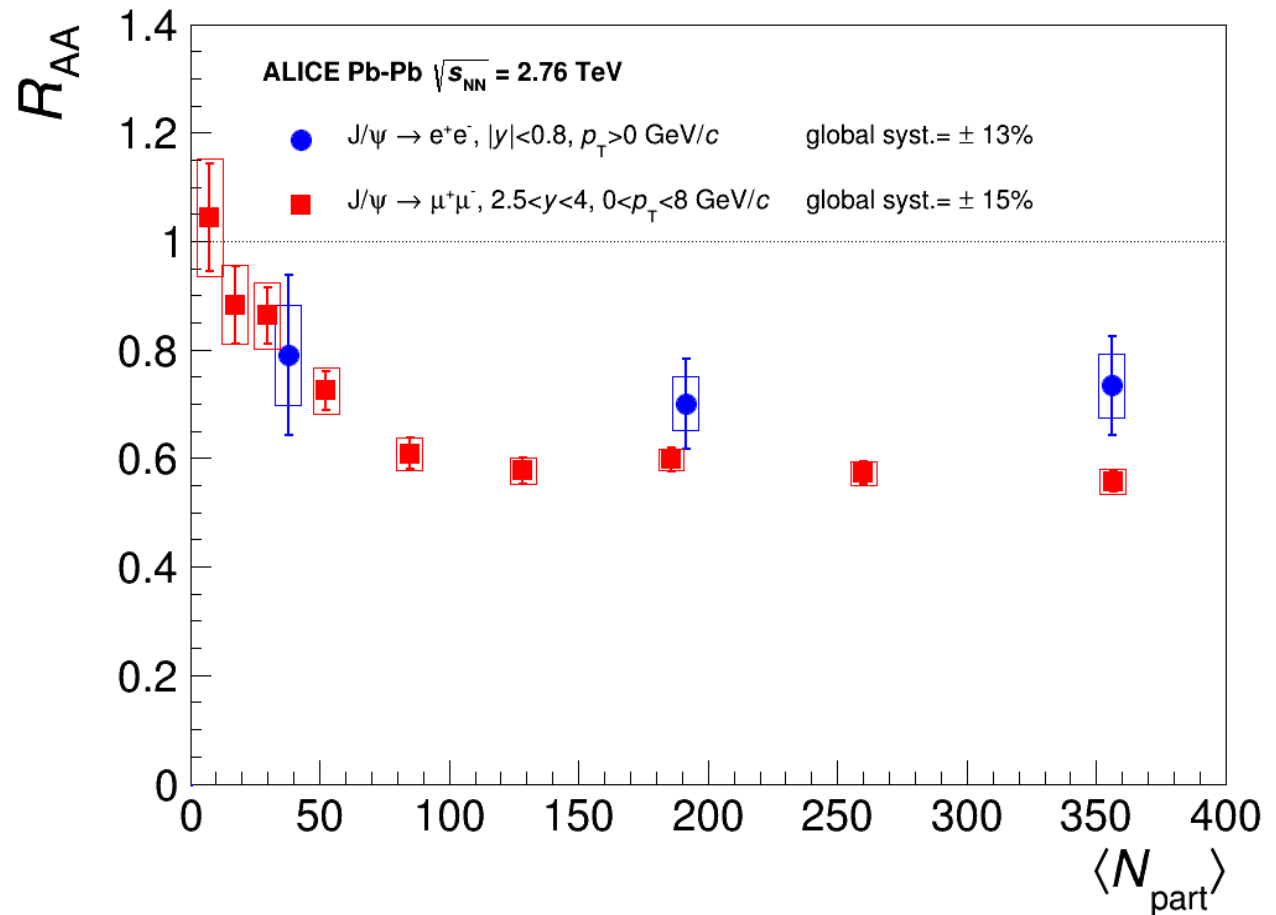


most challenging: central PbPb collisions

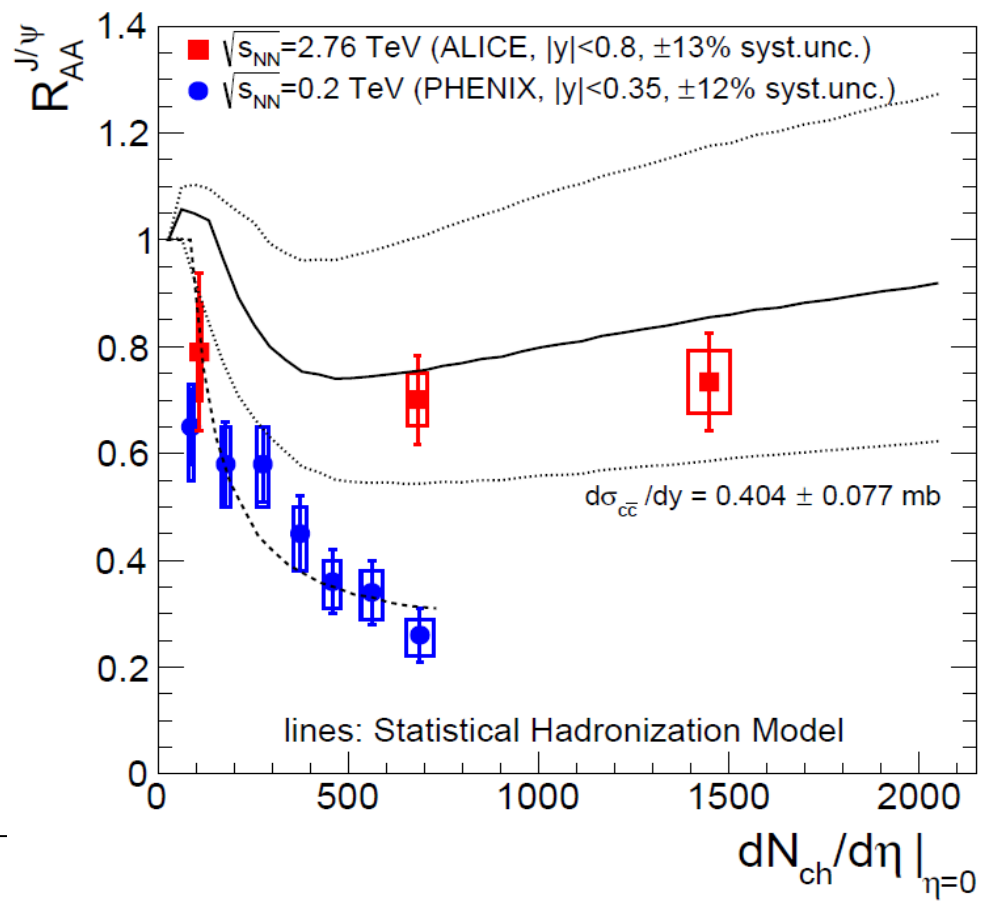
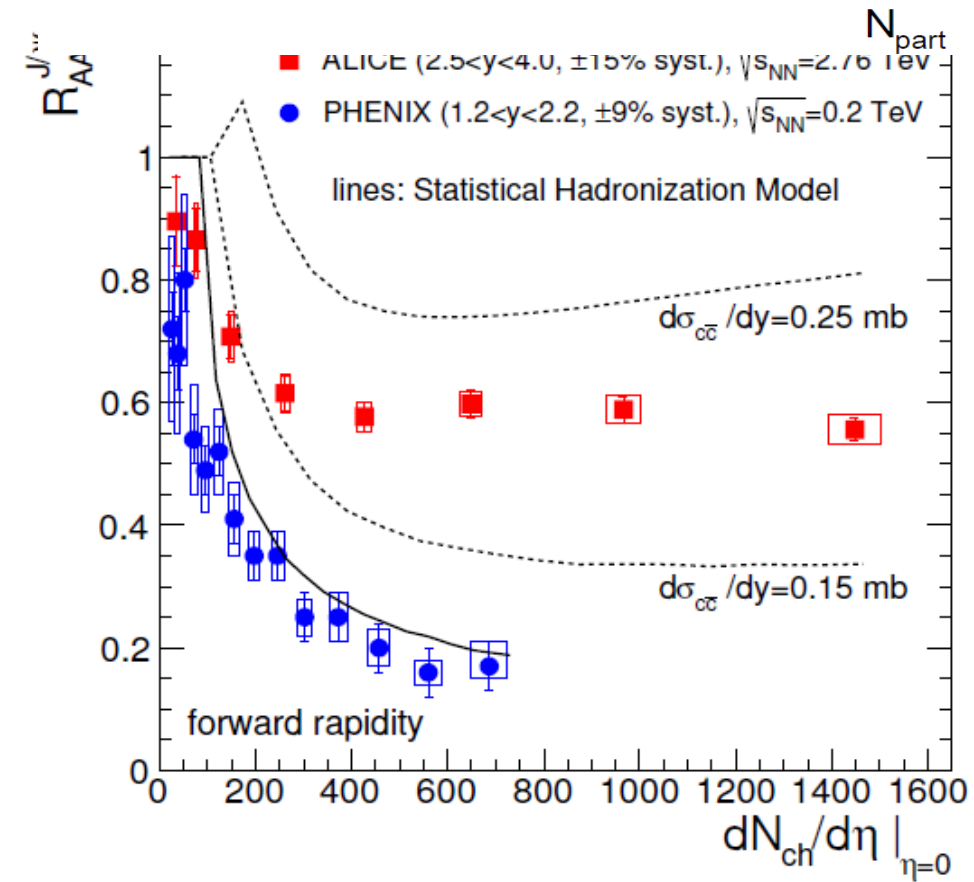
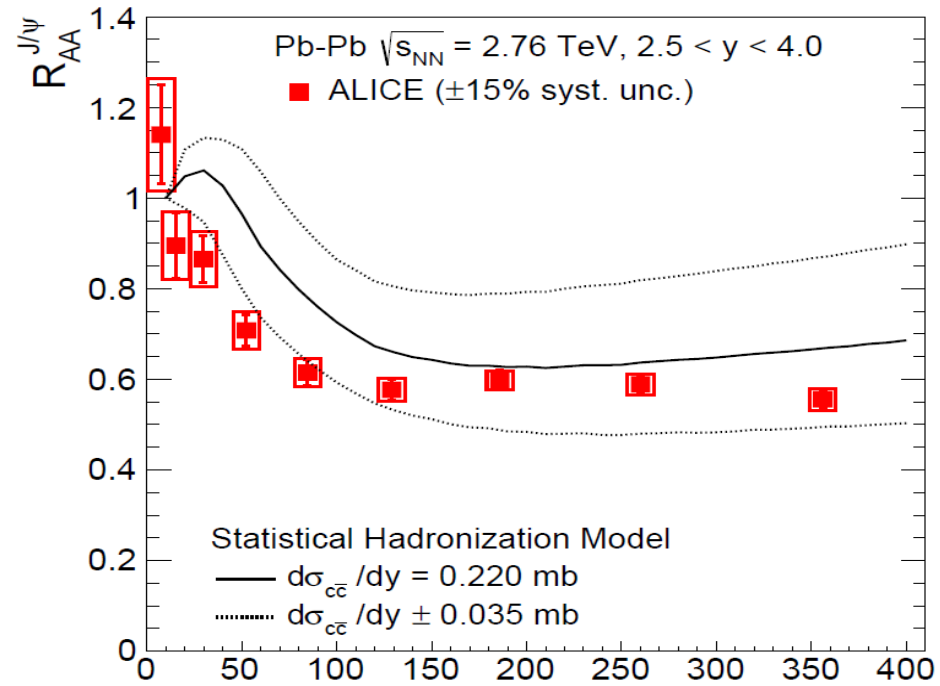
in spite of formidable combinatorial background

(true electrons, not from  $J/\psi$  decay but e.g. D- or B-mesons) resonance well visible

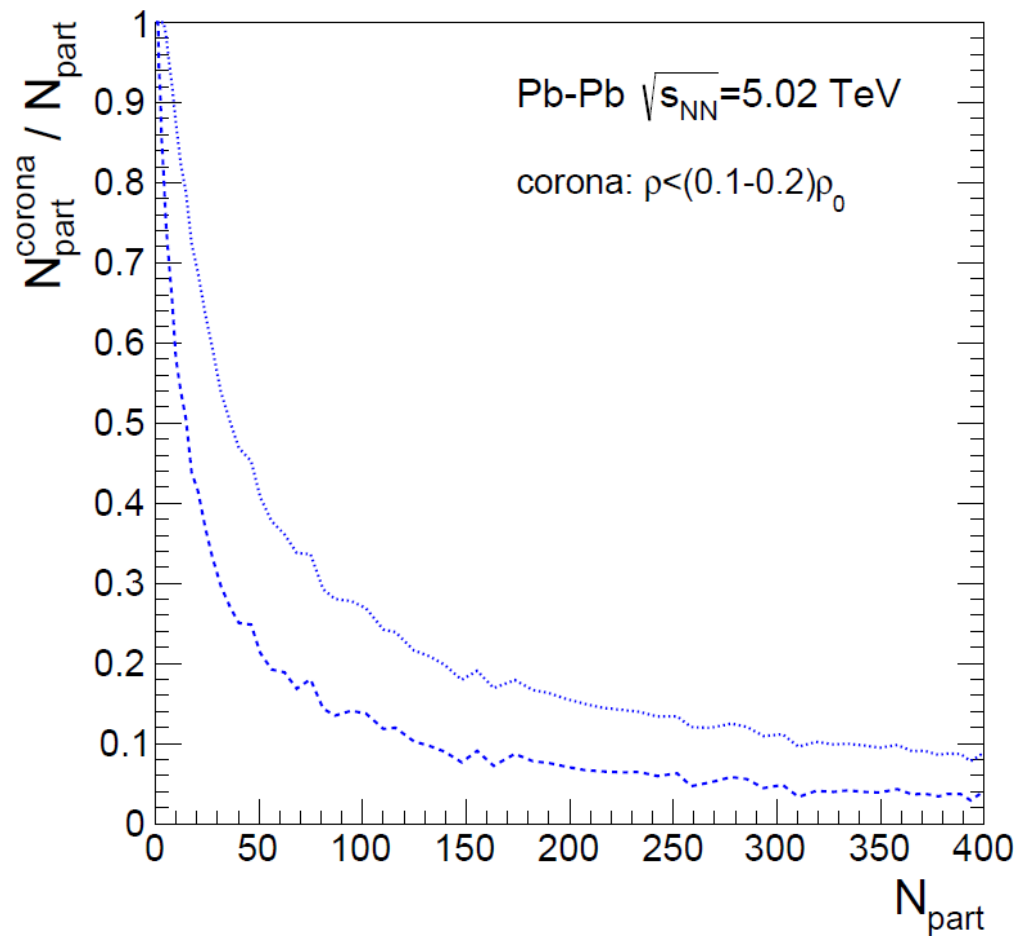
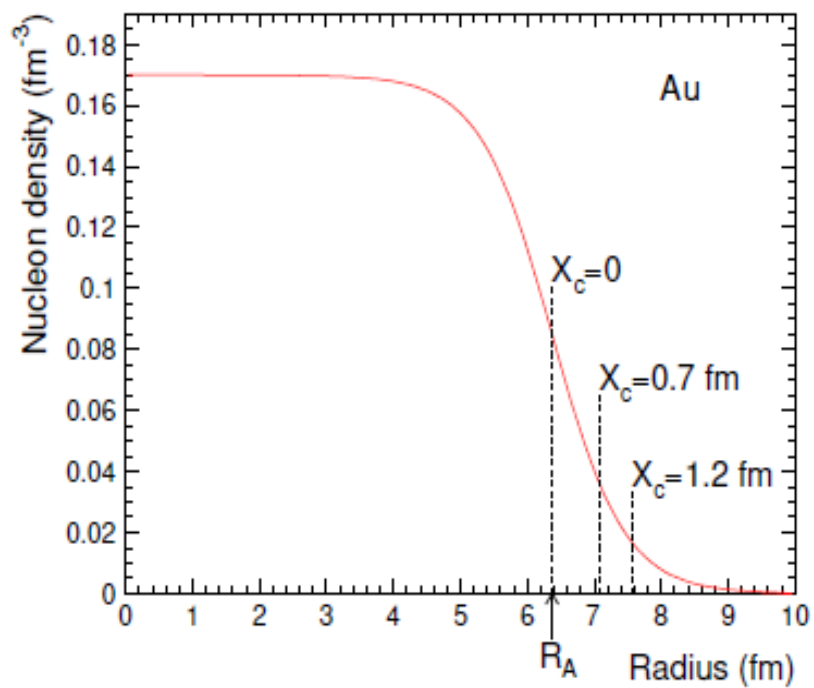
# J/psi in PbPb collisions relative to pp



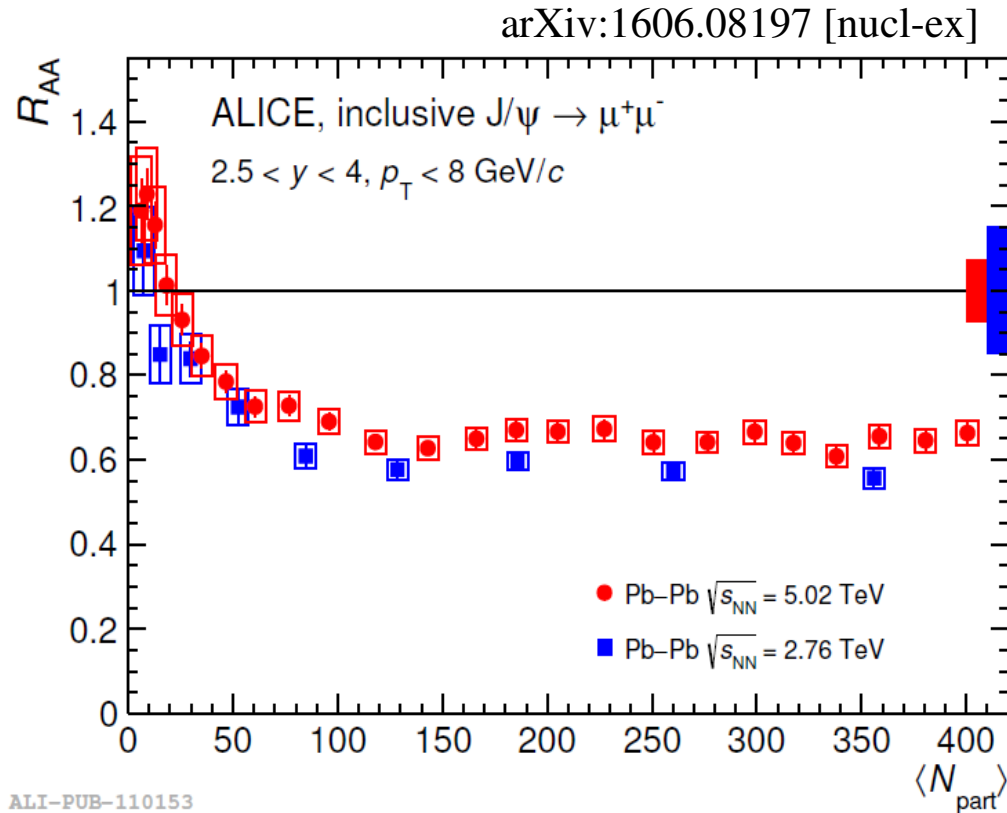
- nearly flat over large centrality range
- indication of rise for most central and mid-rapidity



# Corona fraction in PbPb collisions



# J/ψ in PbPb at $\sqrt{s_{NN}} = 5.02$ TeV

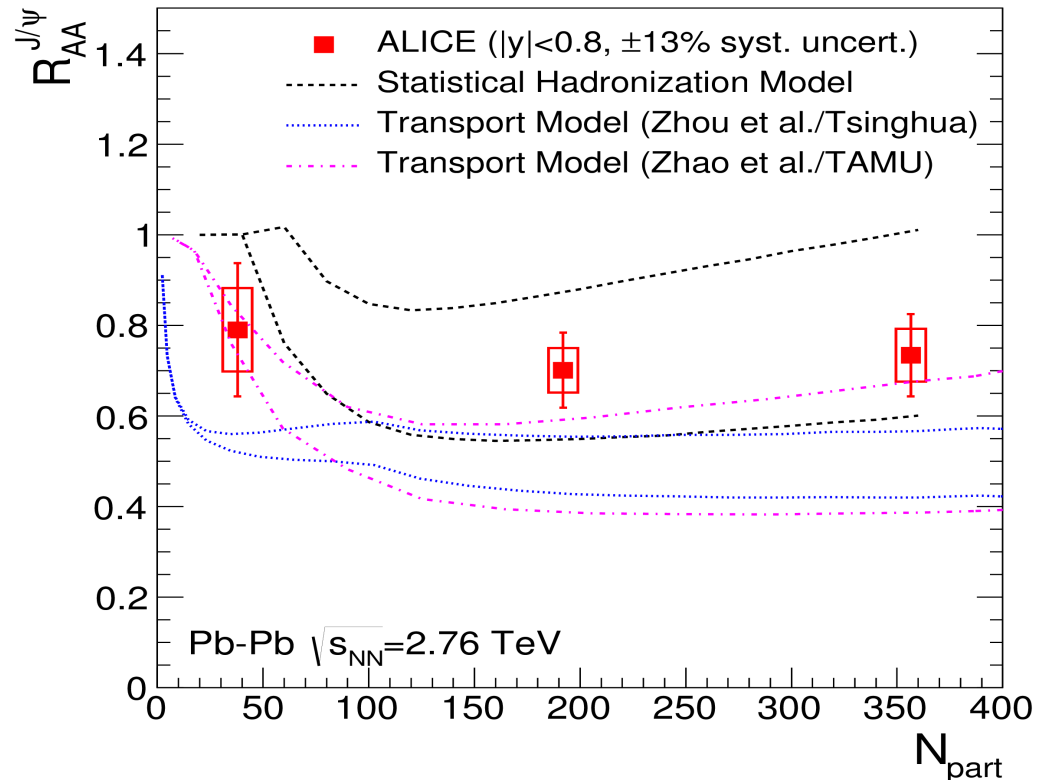
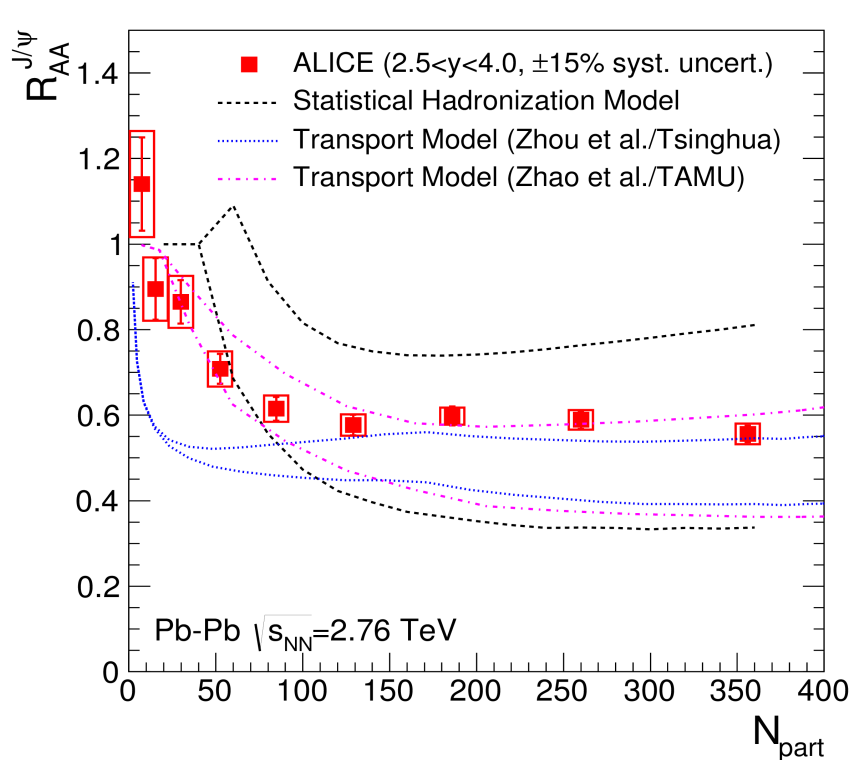


$$R_{AA}^{0-90\%}(5.02 \text{ TeV}) / R_{AA}^{0-90\%}(2.76 \text{ TeV}) = 1.13 \pm 0.02(\text{stat}) \pm 0.18(\text{syst})$$

increase of J/ψ  $R_{AA}$  for all centralities and over large range of  $p_t$  (but within  $1 \sigma$ )



# J/psi and transport models (and stat hadronization)



in transport models (Rapp et al. & P.Zhuang, N.Xu et al.) J/psi generated both in QGP and at hadronization

- transport models also in line with  $R_{AA}$

part of J/psi from direct hard production, part dynamically generated in QGP, part at hadronization, **but different open charm cross section used**

(0.5-0.75mb TAMU and 0.65-0.8 mb Tsinghua vs. 0.3-0.4 mb SHM)

# Attempt to determine Debye mass from data

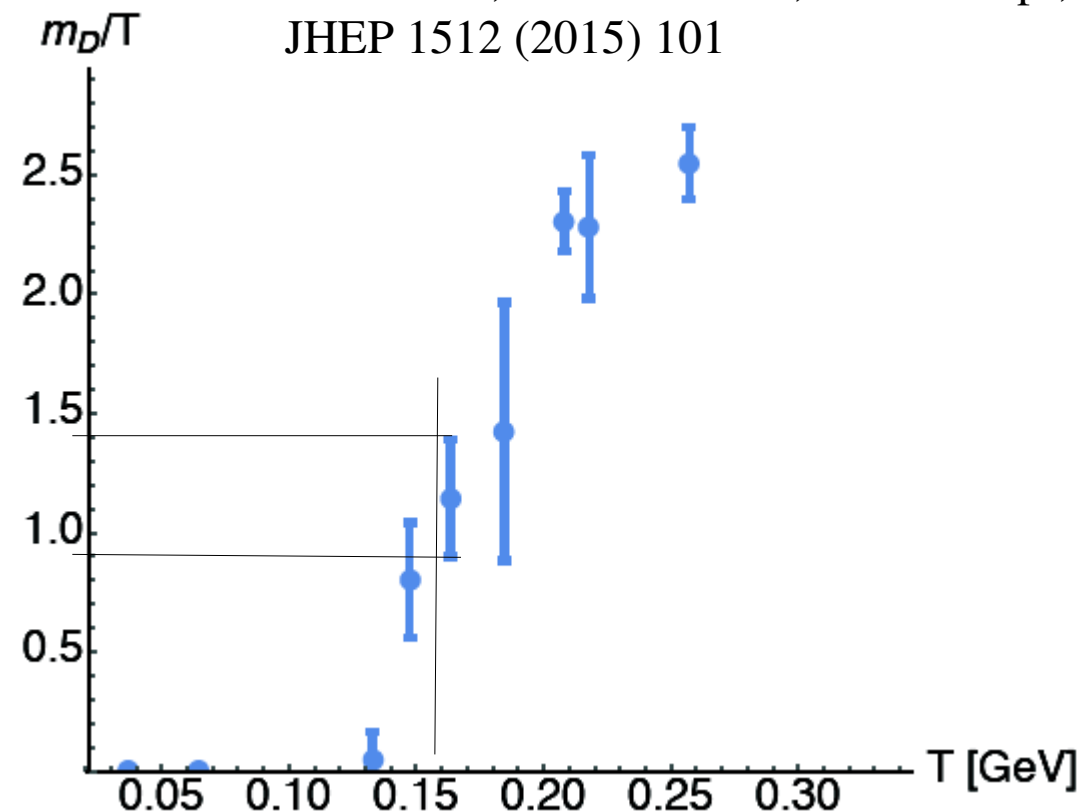
$J/\psi$  formation via statistical hadronization at  $T_c$  implies in classical picture:

$$\lambda_D < r_{J/\psi} \simeq 0.5 \text{ fm at } T = 156 \text{ MeV} \quad \text{or} \quad \omega_D/T > 2.5$$

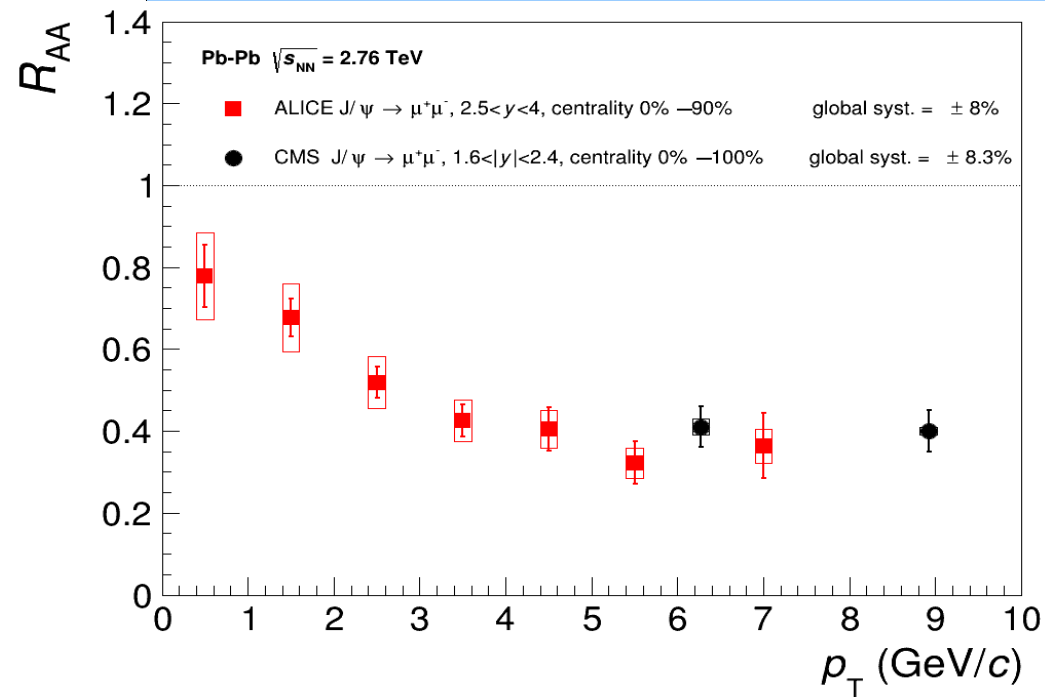
compare to recent finite temperature lQCD potential result:

- value at  $T_c$  lower
  - systematics?
  - and: lattice potential has real and imaginary part, both contribute to screening
- other observable to determine  $\lambda_D$  ?  
e.g.  $\psi'$  vs  $J/\psi$

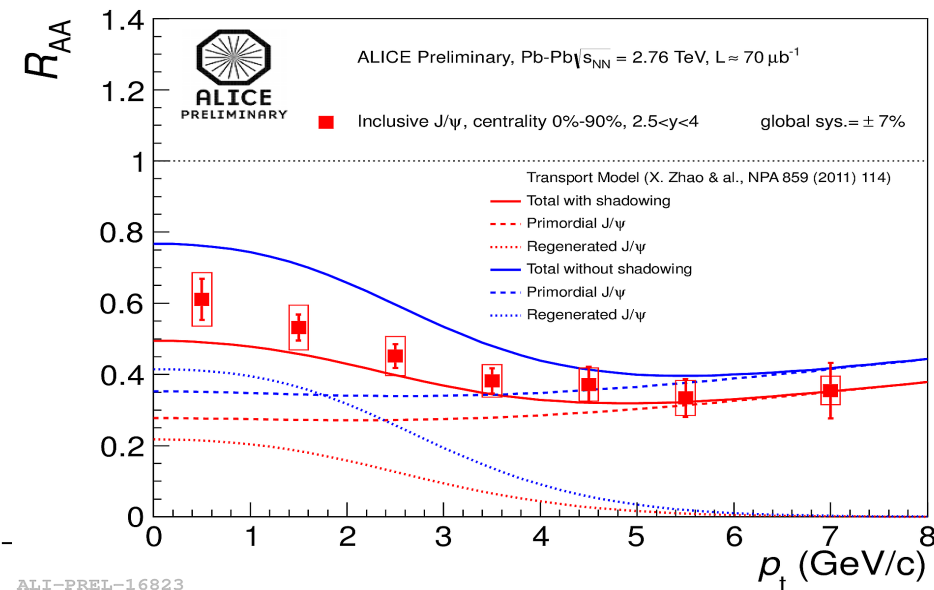
Y. Burnier, O. Kaczmarek, A. Rothkopf,  
JHEP 1512 (2015) 101



# $p_t$ Dependence of $J/\psi R_{AA}$

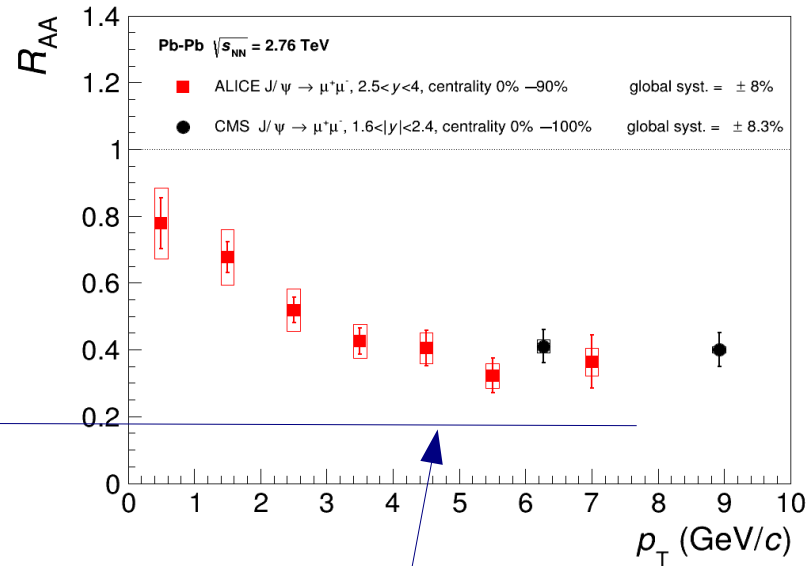
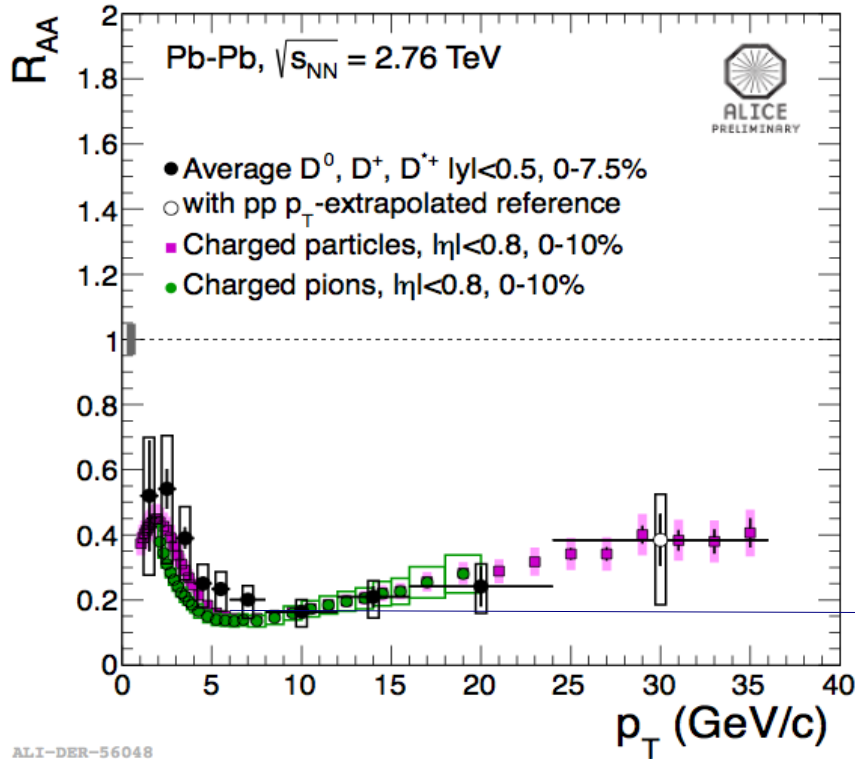


relative yield larger at low  $p_t$  in nuclear collisions  
 good agreement with CMS at high  $p_t$



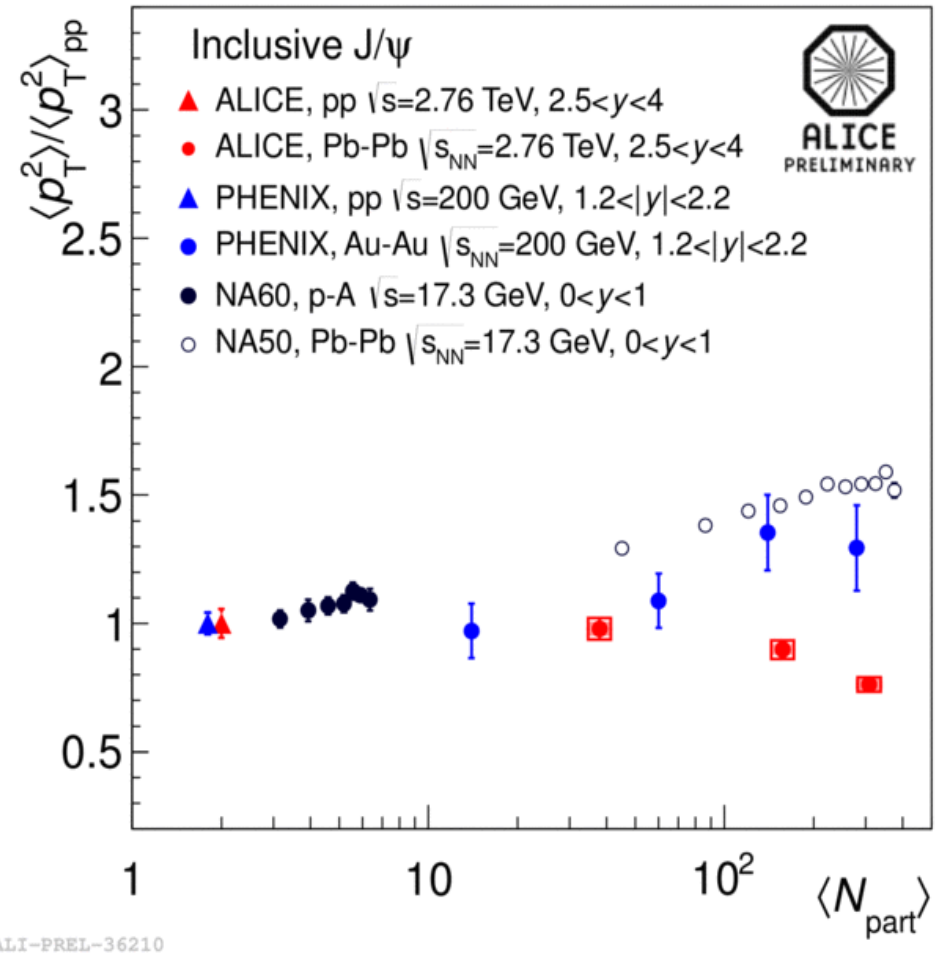
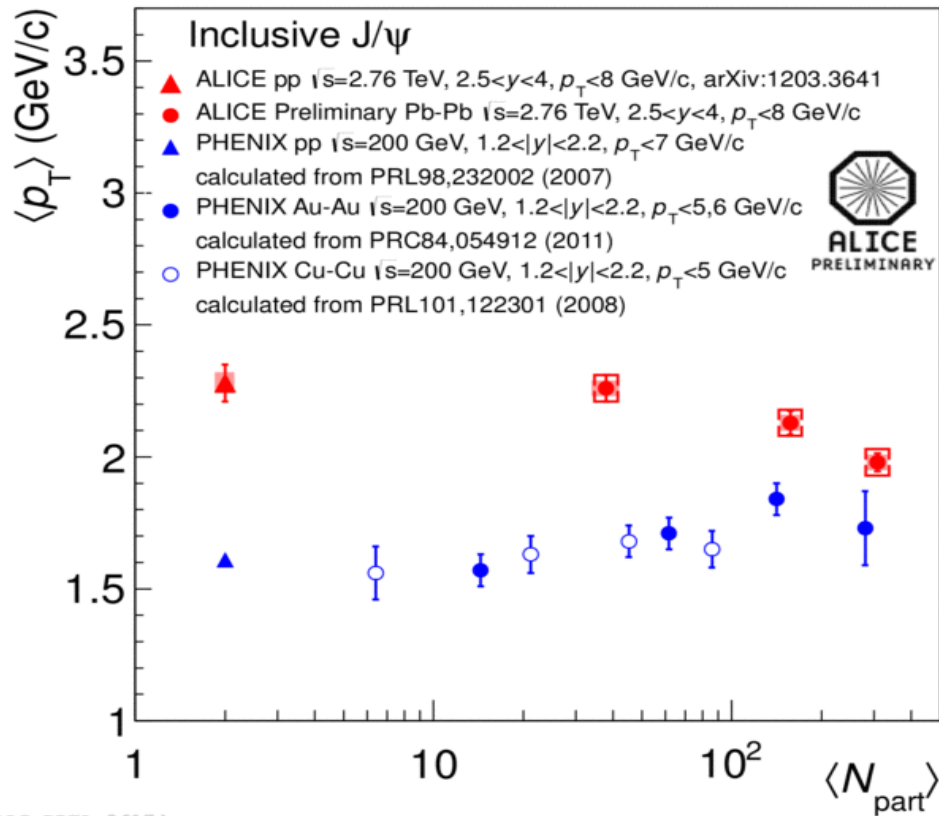
← statistical hadronization only expected for charm quarks thermalized in the QGP  
 $p_t$  dependence in line with this prediction  
 — in CMS only suppression

# $p_t$ dependence of $R_{AA}$



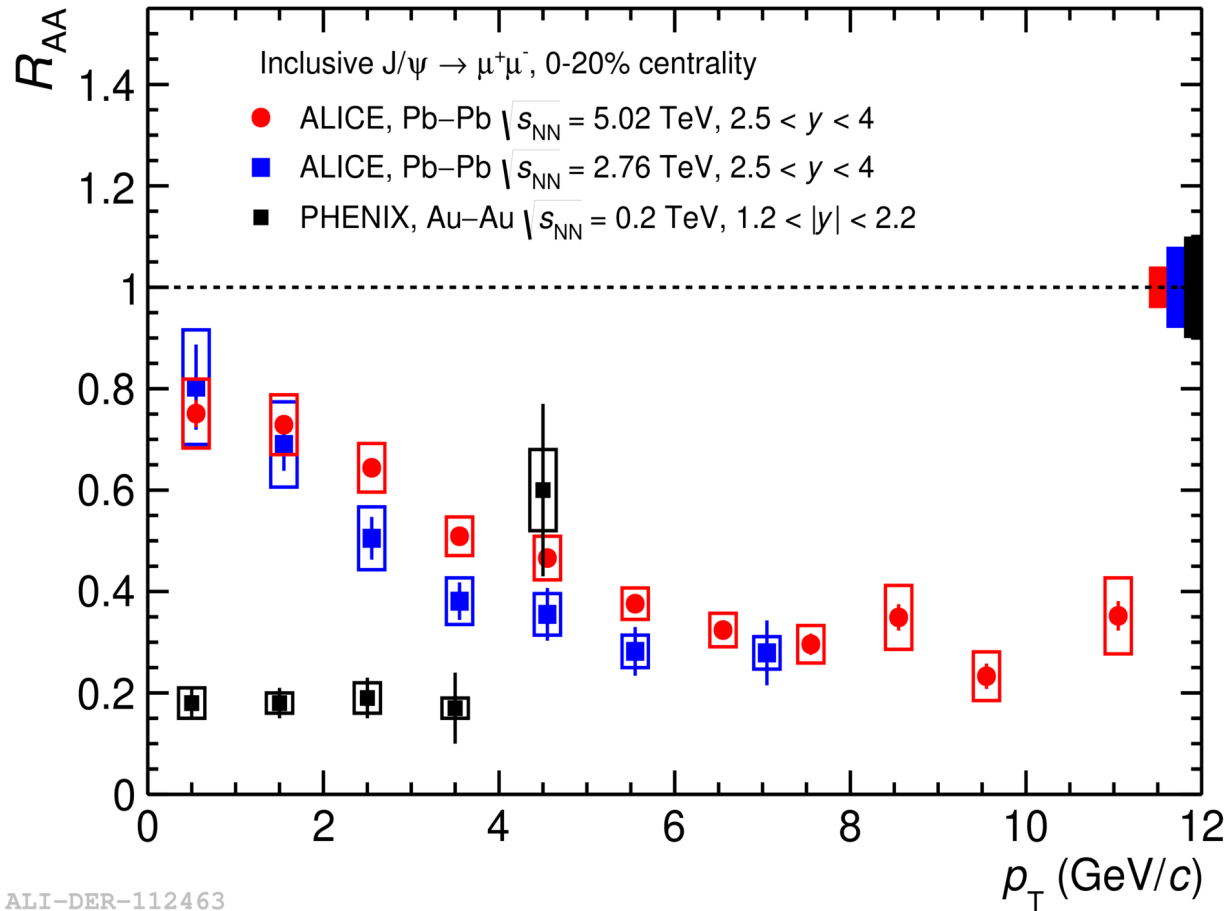
is high  $p_t$  part indicative of the same charm quark energy loss seen for D's  
 out to what  $p_t$  is statistical hadronization/regeneration relevant?

# Softening of J/psi $p_T$ distributions for central PbPb collisions



At LHC for central collisions softening relative to peripheral collisions and relative to pp (opposite trend to RHIC) - consistent with formation of J/psi from thermalized c-quarks

# Transverse momentum spectrum



softer in PbPb as compared to pp

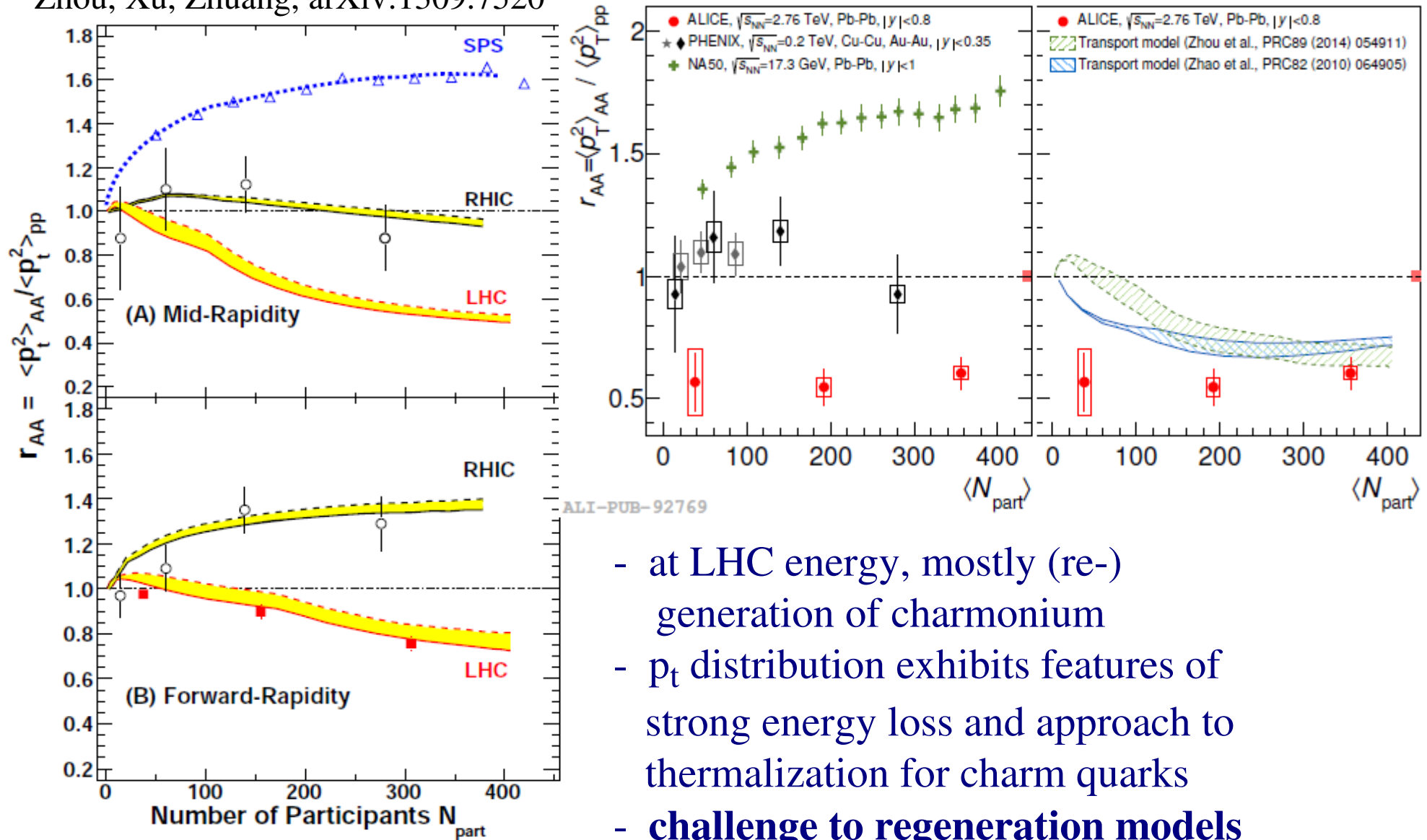
a qualitatively new feature as compared to RHIC where the trend is opposite

in line with thermalized charm in QGP at LHC, forming charmonia

ALI-DER-112463

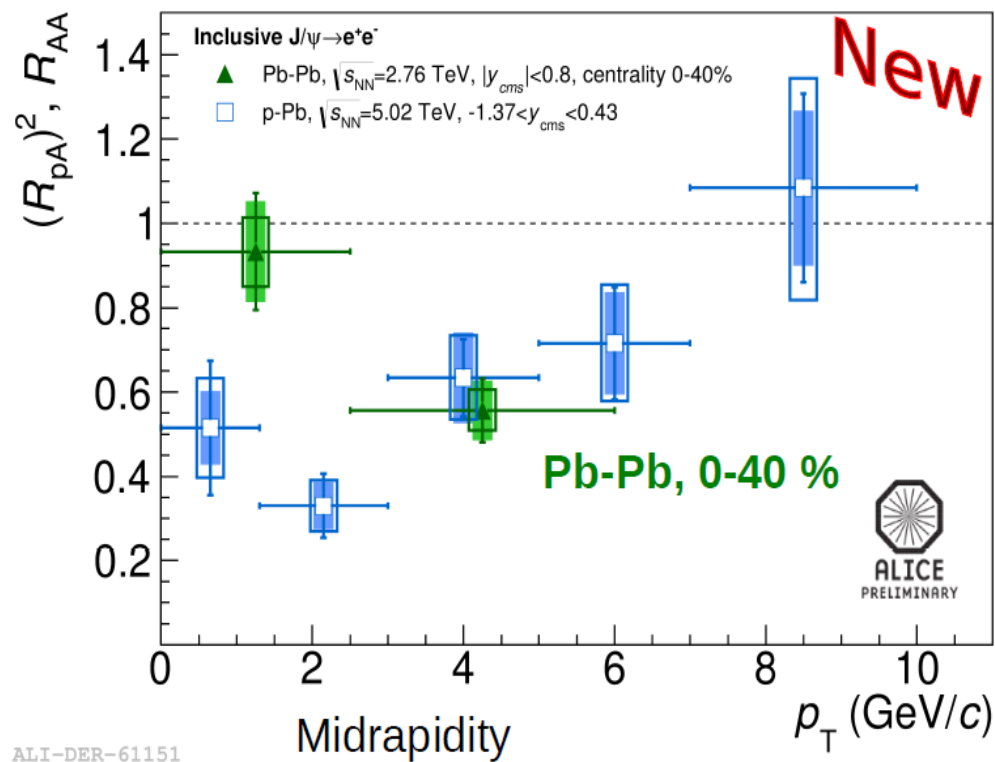
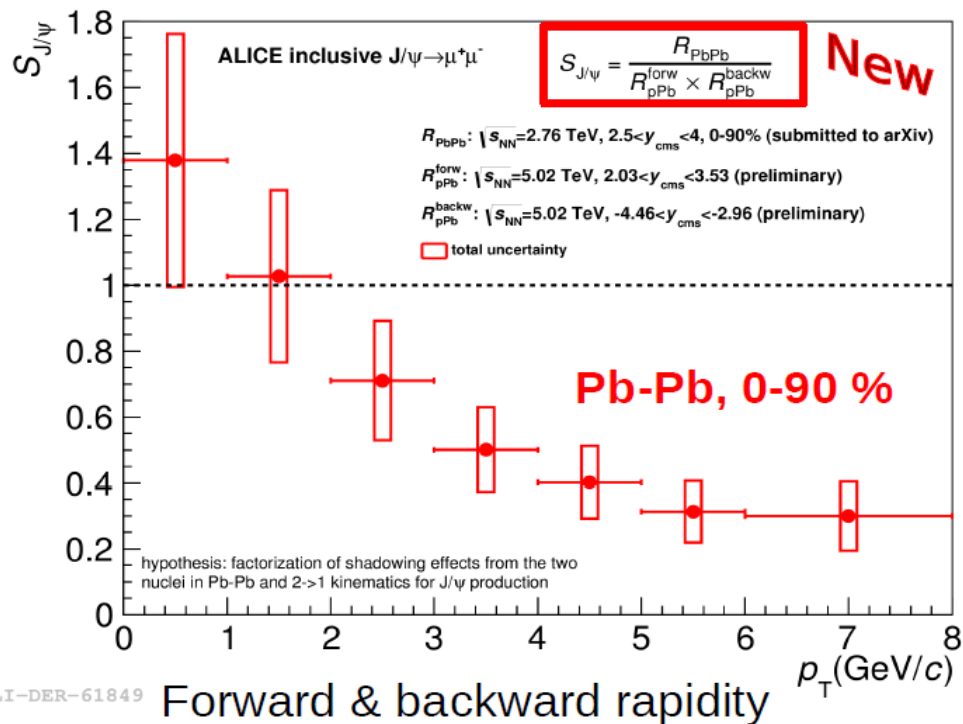
# Analysis of transverse momentum spectra

Zhou, Xu, Zhuang, arXiv:1309.7520



- at LHC energy, mostly (re-) generation of charmonium
- $p_t$  distribution exhibits features of strong energy loss and approach to thermalization for charm quarks
- **challenge to regeneration models**

# J/psi vs pt in PbPb collisions relative to pPb collisions



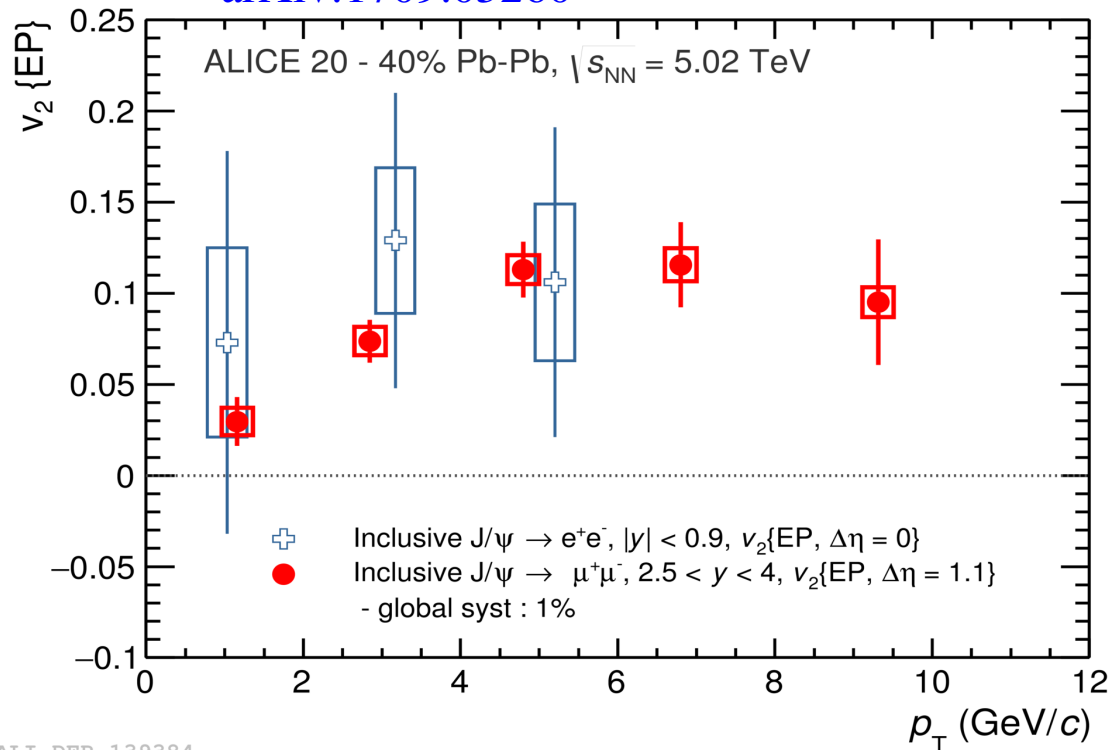
at low  $p_T$  yield in nuclear collisions above pPb collisions

$J/\psi$  production enhanced in nuclear collisions over mere shadowing effect



# Elliptic flow of J/ψ vs p<sub>t</sub>

arXiv:1709.05260



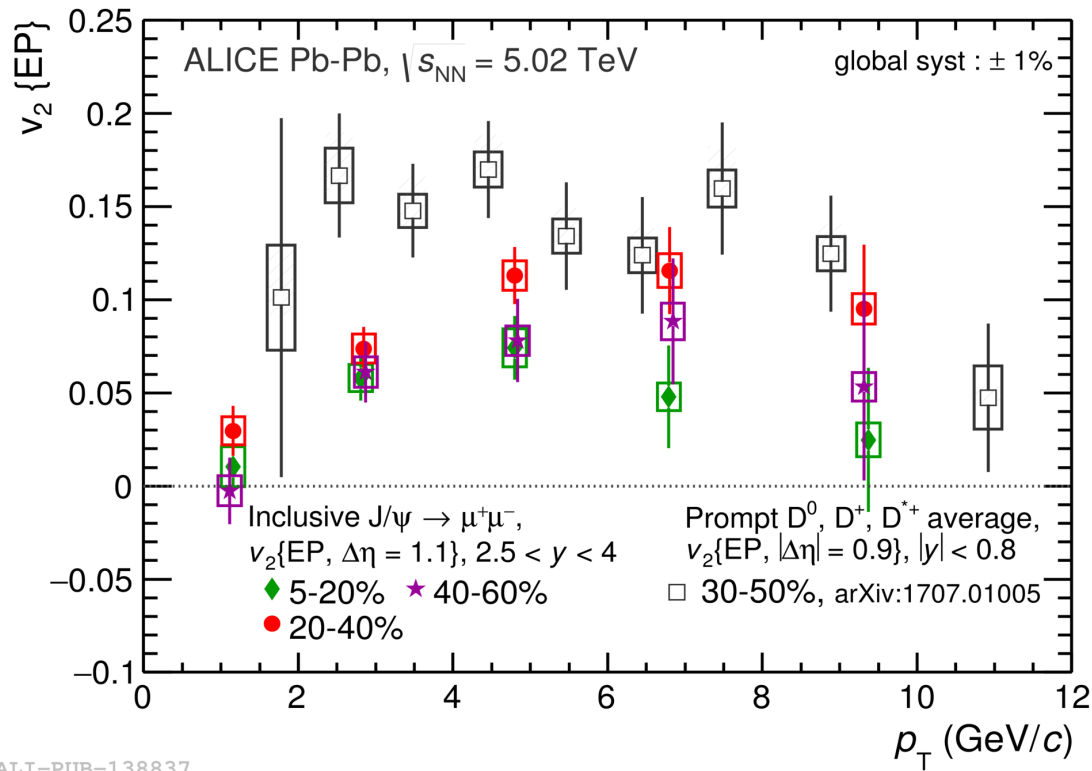
charm quarks thermalized in the QGP should exhibit the elliptic flow generated in this phase

- expect build-up with p<sub>t</sub> as observed for π, p, K, Λ, ... and vanishing signal for high p<sub>t</sub> region not dominated by flow

first observation of significant J/ψ v<sub>2</sub> in line with expectation from statistical hadronization can be computed following approach above with hydro velocity profile

# Elliptic flow of $J/\psi$ vs $p_t$

arXiv:1709.05260



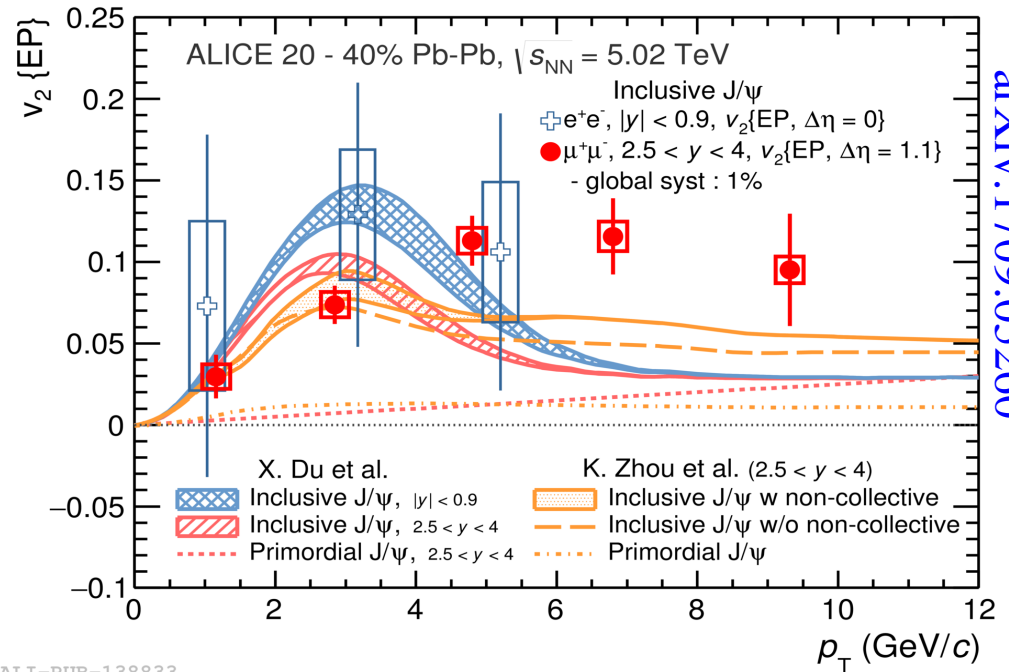
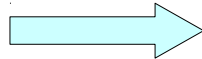
ALI-PUB-138837

Strength of  $J/\psi$   $v_2$  similar to D-mesons

# Elliptic flow of $J/\psi$

charm quarks thermalized in the QGP should exhibit the elliptic flow generated in this phase

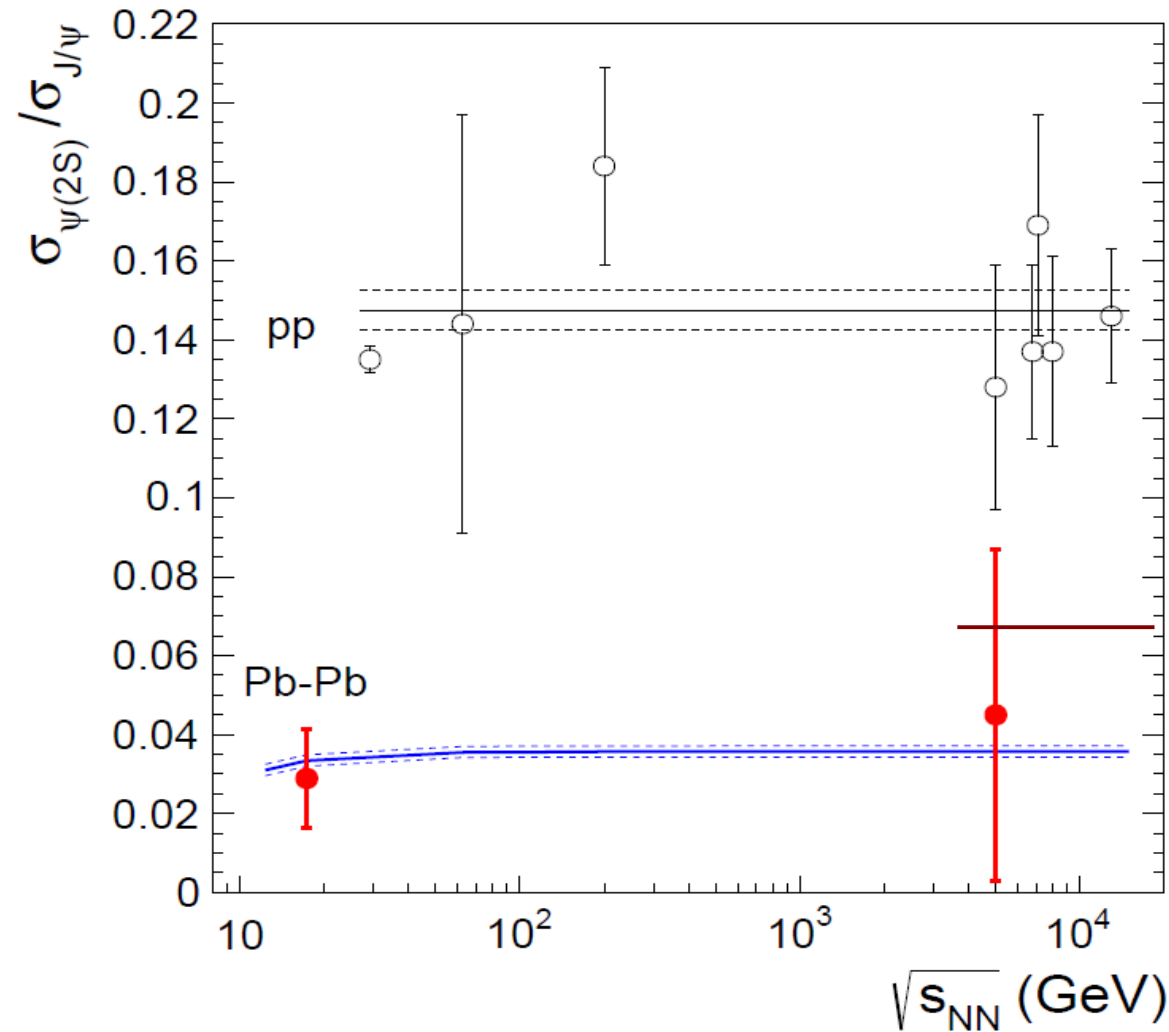
first observation of significant  $J/\psi$   $v_2$  both at forward and mid rapidity



arXiv:1709.05260

$J/\psi$  elliptic flow in line with expectation from statistical hadronization

# $\psi(2S)$



first data in line with expectation from statistical hadronization at phase boundary but transport model prediction also inside 1 sigma error

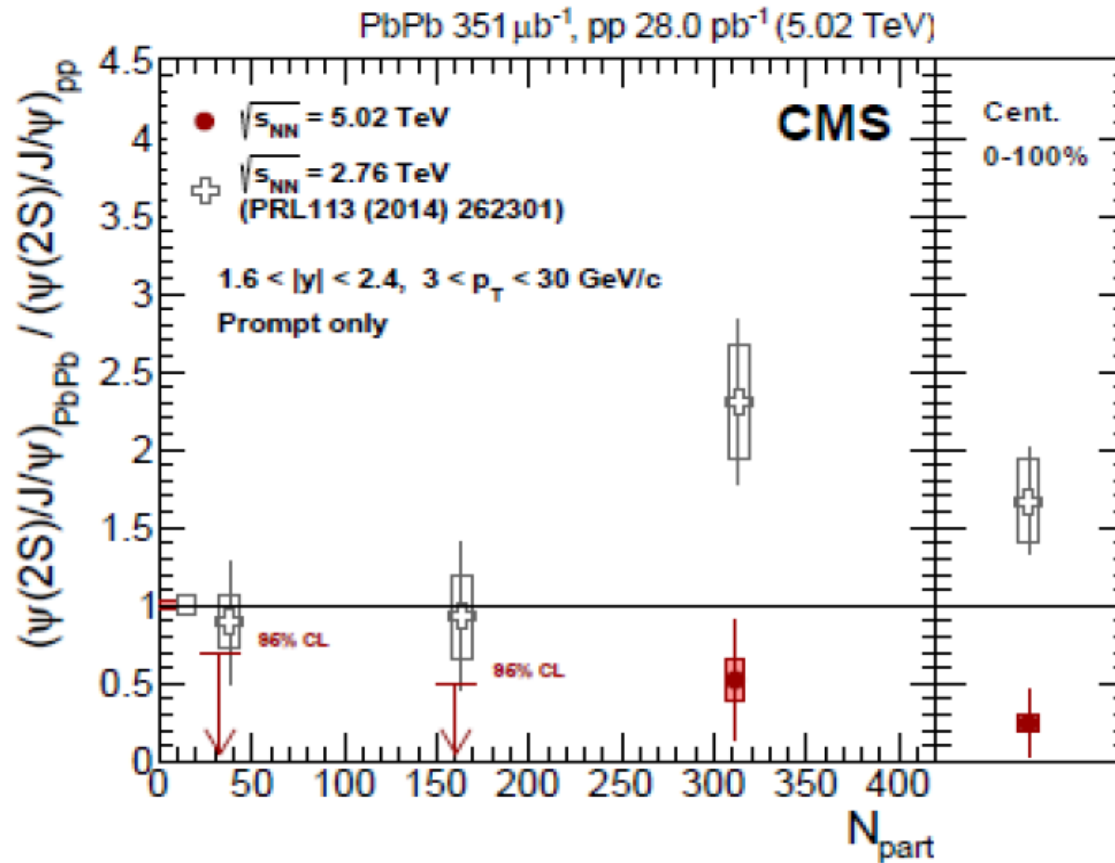
transport, Rapp

SH

# $\psi(2S)$



in picture where psi is created from deconfined quarks in QGP or at hadronization,  $\psi(2S)$  is suppressed more than  $J/\psi$  – run1 CMS results indicate the opposite!

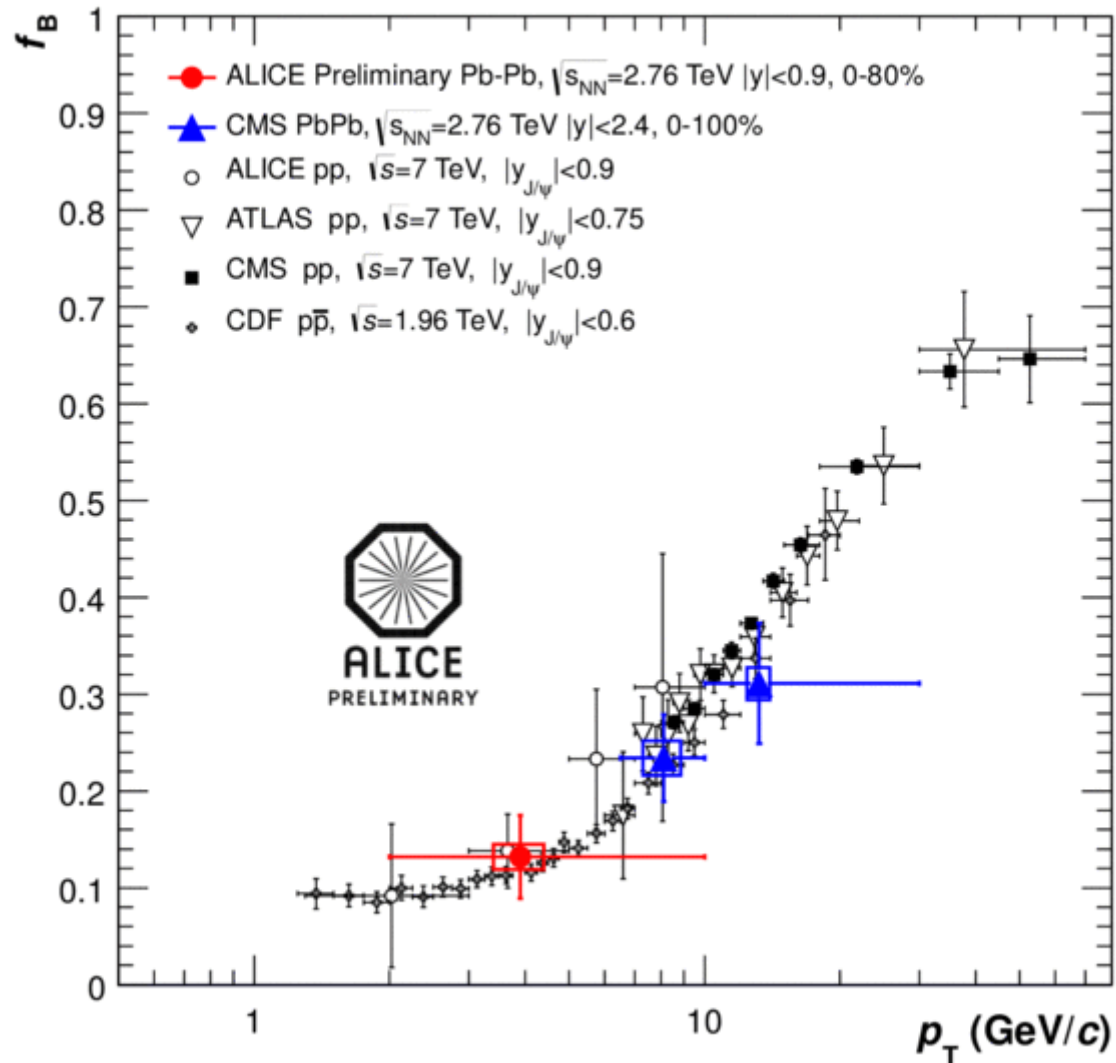


expect value of 1/3 for inclusive  $p_T$  and central collisions

the anomaly (enhancement relative to pp) from 2.76 TeV is not there at 5.02 TeV - very nice ALICE data from  $p_T=0$  to be approved this week



# Fraction of J/psi from B-decays



$p_T$  integrated non-prompt B-fraction of small  
within current errors no significant difference in pp and PbPb collisions

# Outlook – what ALICE can do in the future

## LHC run1:

2 PbPb runs

- 2010  $O(10 \mu\text{b}^{-1})$

- 2011  $O(150 \mu\text{b}^{-1})$

luminosity reached  $\mathcal{L}=2 \cdot 10^{26} \text{ cm}^{-2} \text{ s}^{-1}$  twice design lumi at this energy

1 pPb run

- 2012/2013  $O(30 \text{ nb}^{-1})$

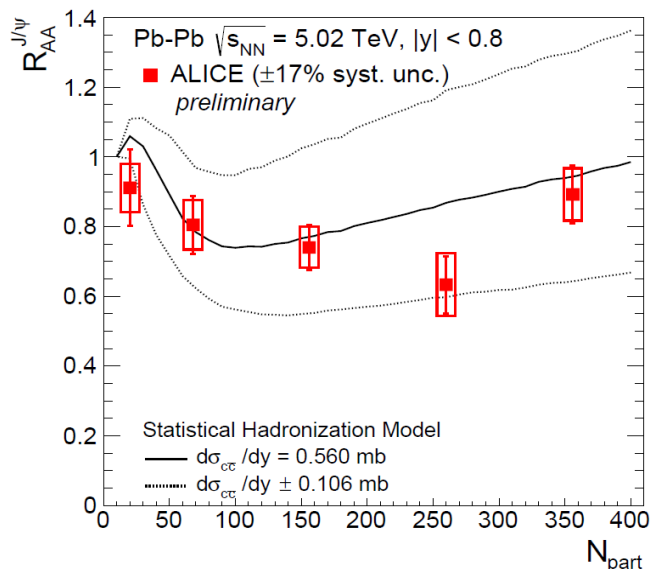
from 2/2013 until end of 2014 **LS1**: consolidation of LHC to allow full energy

**LHC run2:** 2015-2018 PbPb running at  $\sqrt{s_{\text{NN}}} = 5.5 \text{ TeV}$   
to achieve approved initial goal of  $1 \text{ nb}^{-1}$

2019 start **LS2** – increase of LHC luminosity und experiment upgrade, LHCb will join PbPb!

**LHC run3:** 2021 onwards - expect  $\mathcal{L}=6 \cdot 10^{27} \text{ cm}^{-2} \text{ s}^{-1}$  or PbPb interactions at 50 kHz  
achieve for PbPb  $10 \text{ nb}^{-1}$  corresponding to  $8 \cdot 10^{10}$  collisions sampled  
plus a low field run of  $3 \text{ nb}^{-1}$  + pp reference running + pPb - a program for about 6 years

# J/psi as probe of deconfinement



well on the way towards goal for run2  
 expect still a huge jump in performance  
 for runs3/4

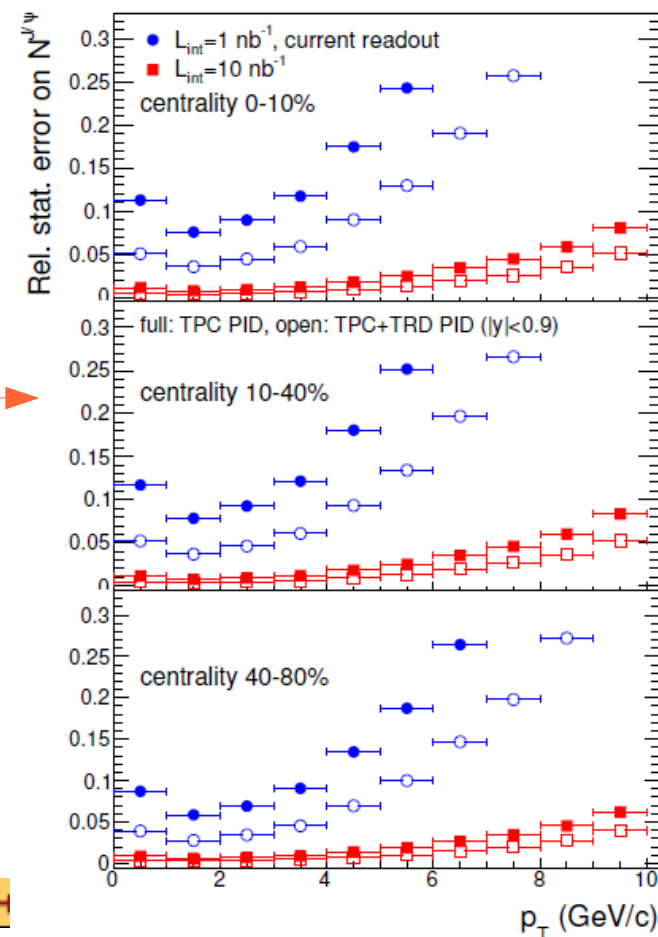
di-electrons statistics limited,  $10 \text{ nb}^{-1}$  will have huge effect

but also syst uncertainties will decrease with upgrade:

will also add TRD for electron id - reduced comb background

thinner ITS reduced radiation tail

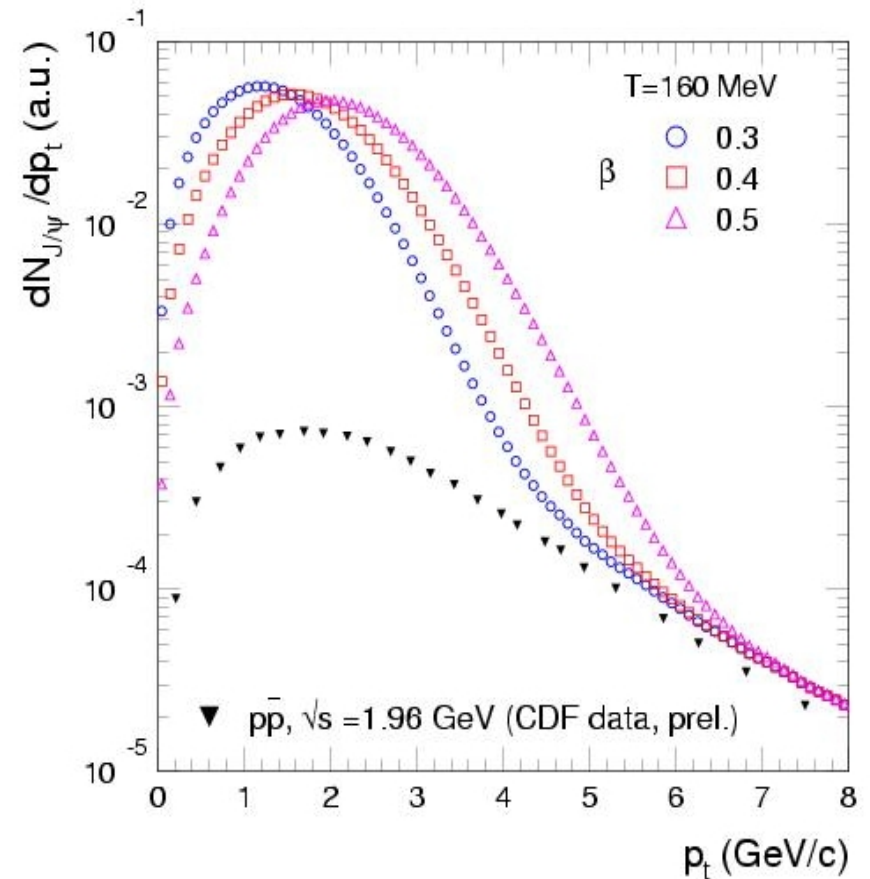
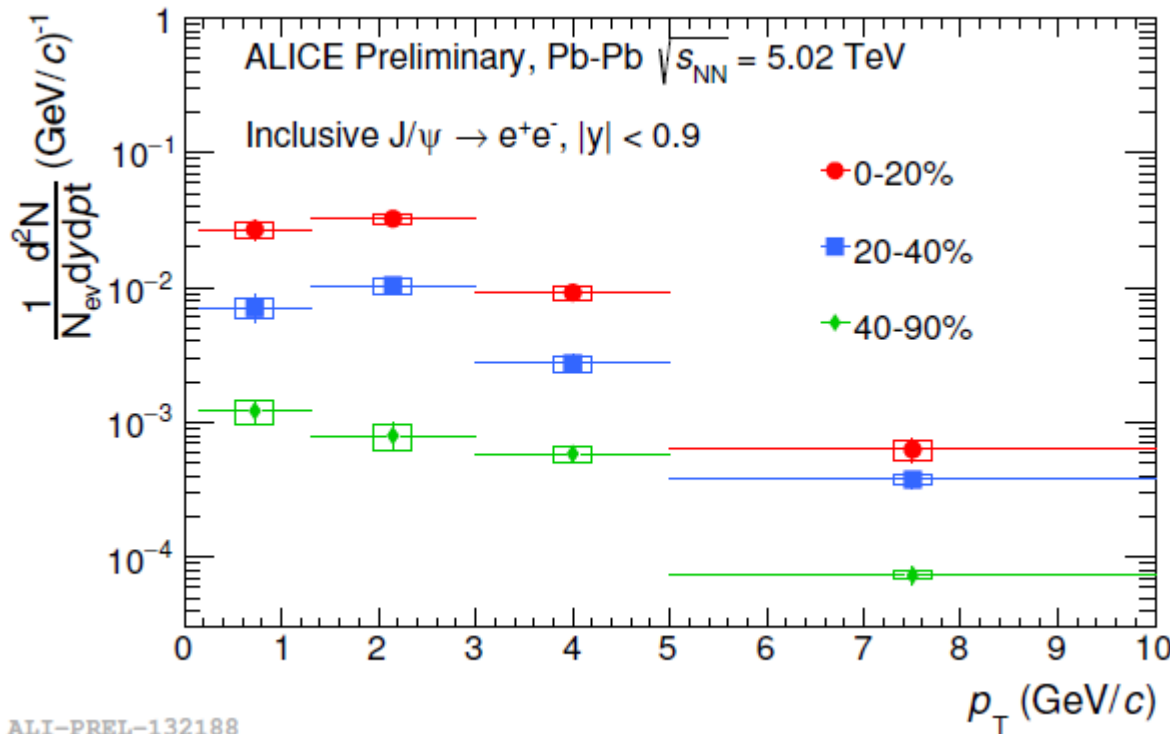
both affect signal extraction





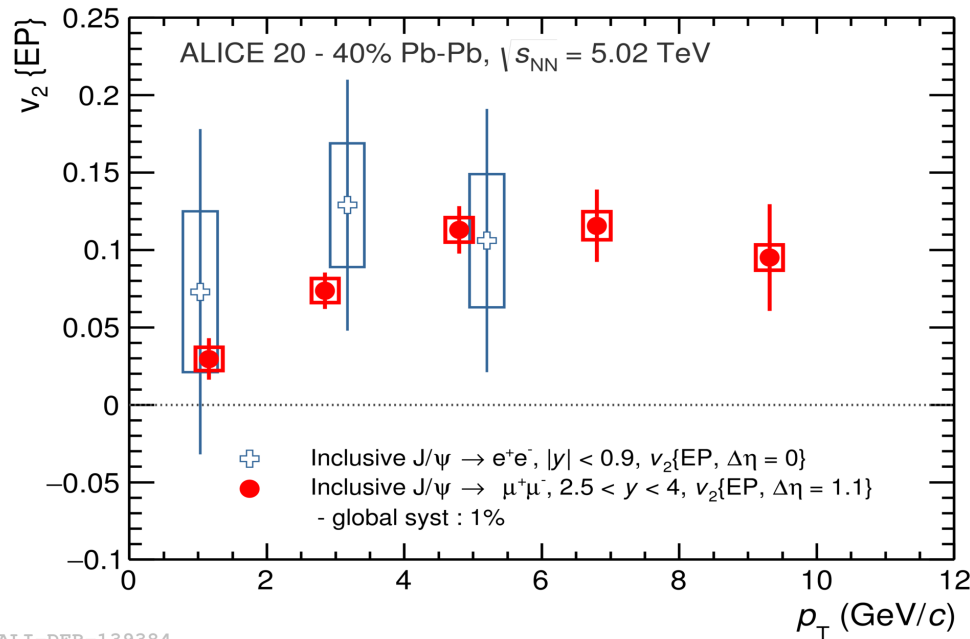
# spectral distribution is key to thermalization

if charm quark thermalize, their spectral distributions should also reflect collective flow of liquid



first spectra a mid-y appearing  
 much more to come  
 we are computing spectra

# J/psi elliptic flow



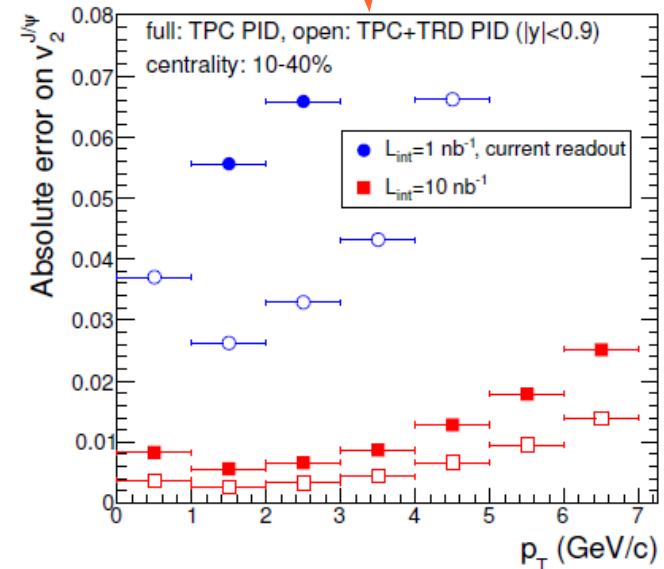
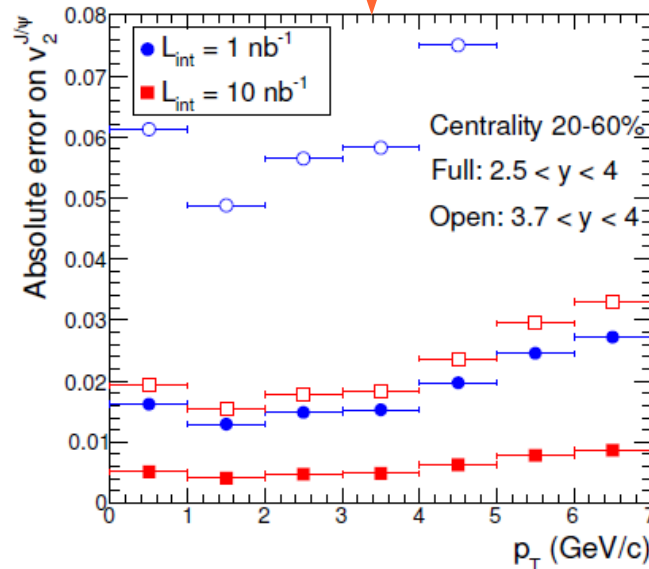
goal for run2 in muon arm already achieved  
for e+e- at mid-y getting there

future statistical errors

muon arm

central barrel

ALI-DER-139384

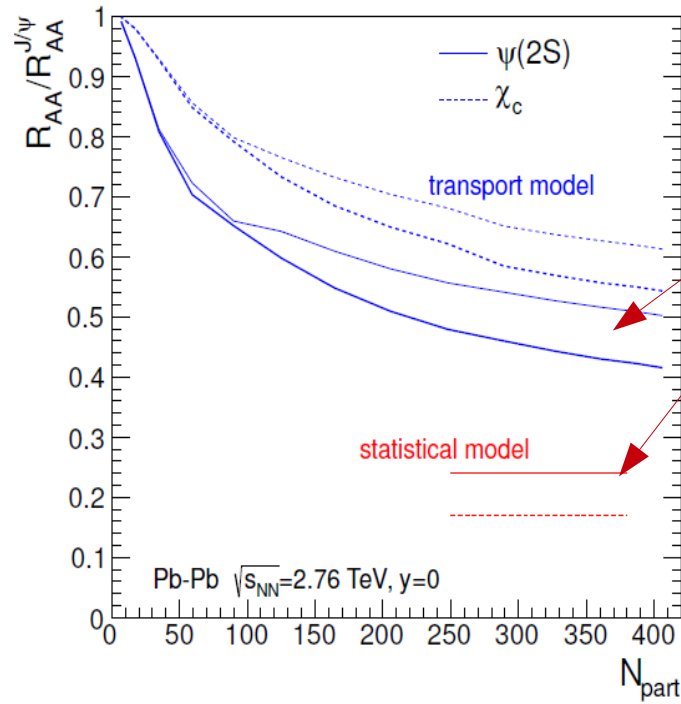
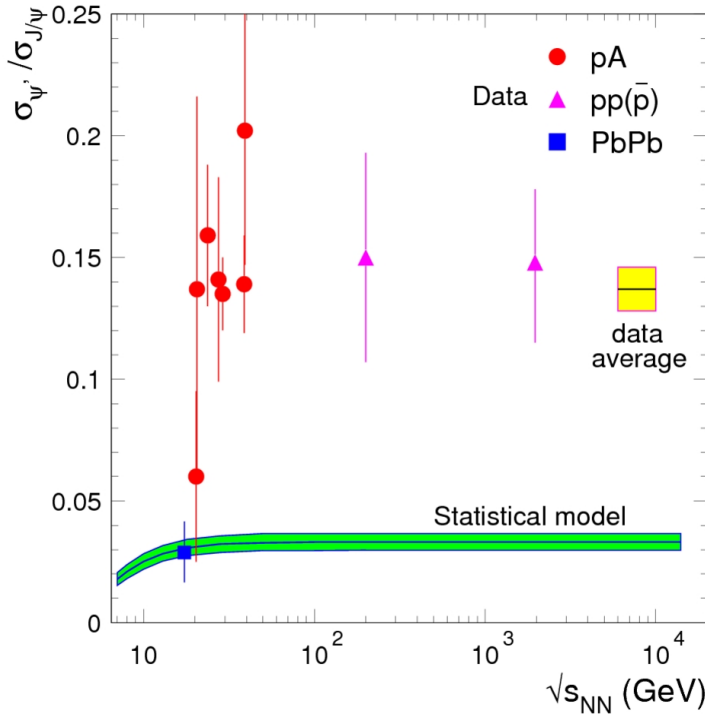


# How to distinguish between statistical hadronization and transport models with $J/\psi$ beyond $T_c$ ?

not a detail, which model is right, but fundamental question  
link to phase boundary and existence of bound states beyond  $T_c$  at stake

- $R_{AA}$  can be reproduced by both, albeit with different charm cross sections  
go away from  $R_{AA}$ , normalize to open charm cross
- spectra: transport models start to be challenged, need more precise data  
and more refined hydro based computation
- similar:  $v_2$  of  $J/\psi$
- maybe decisive: excited state population

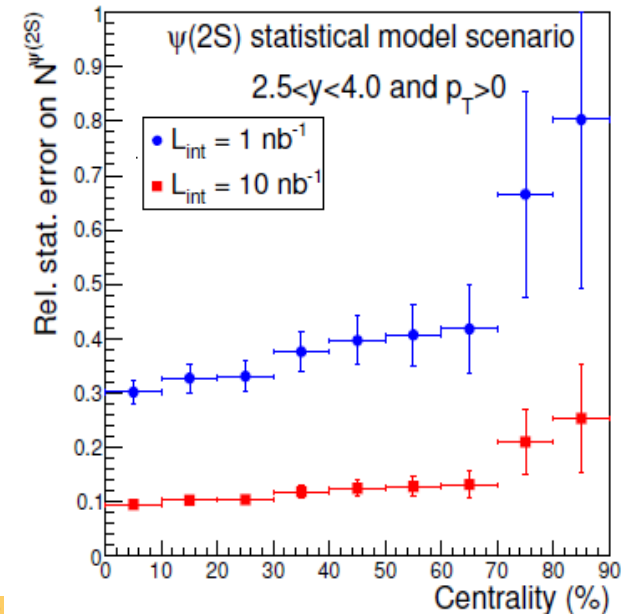
# excited charmonia crucial to distinguish between models



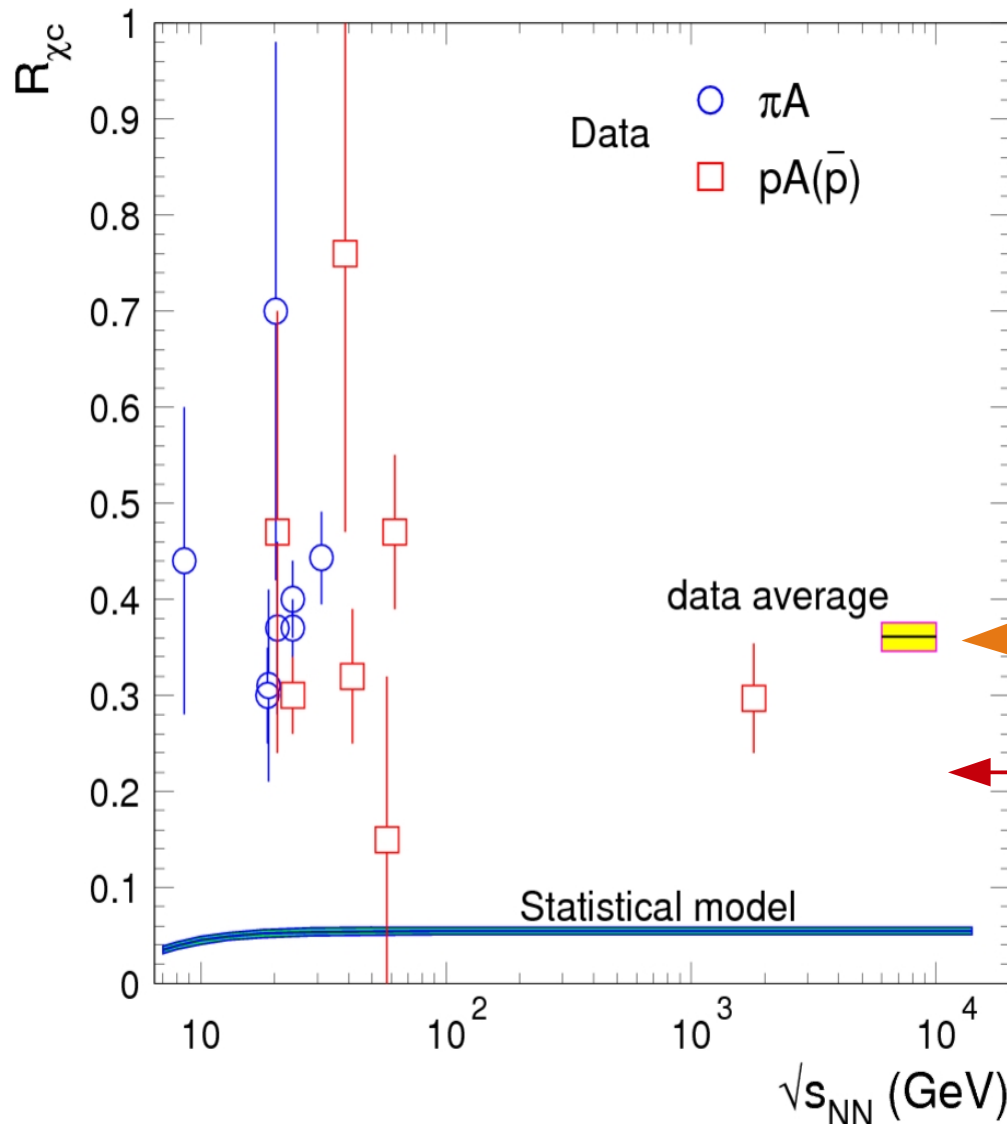
in fact **here** one can distinguish between the transport models that form charmonia already in QGP and statistical hadronization at phase boundary!

for statistical hadronization need to see suppression by Boltzmann factor  $\chi_c$  even bigger difference

expected ALICE performance  $\longrightarrow$   
muon arm run2 and run3



# Situation even more dramatic for P-states

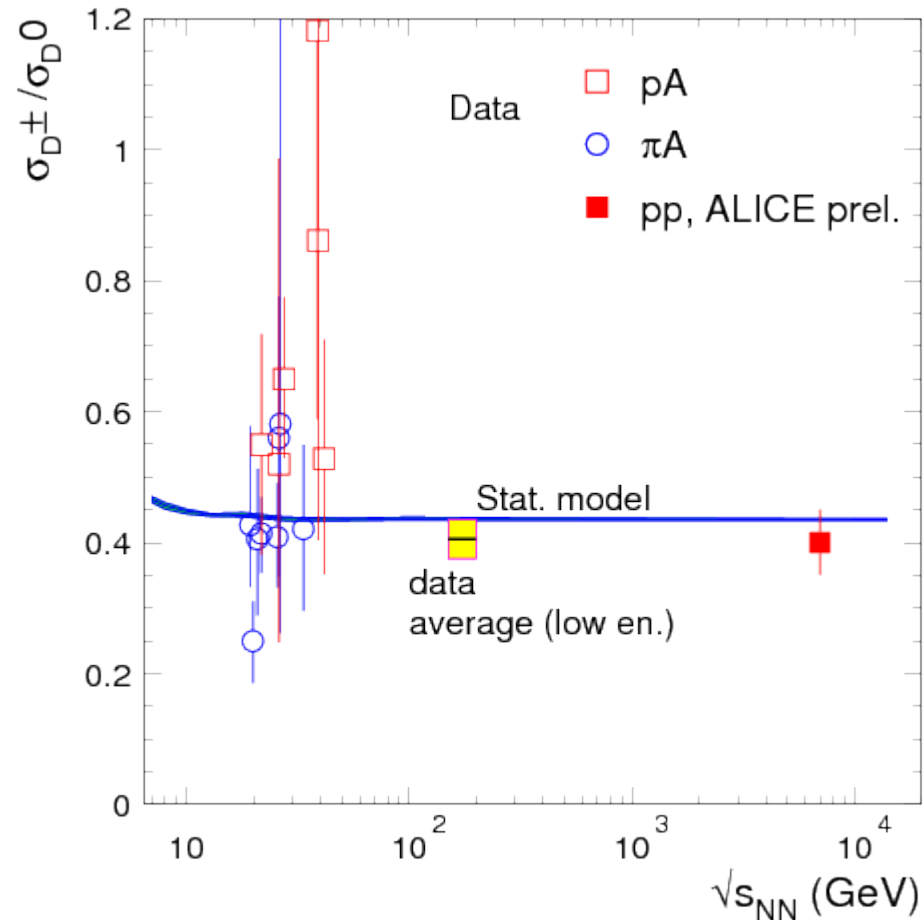


$p A$  and  $\pi A$  data on average factor 7 above statistical model prediction

Transport model (Rapp)

A. Andronic, F. Beutler, P. Braun-Munzinger, K. Redlich,  
J. Stachel Phys. Lett. B678 (2009) 350

# Charged to neutral D-mesons

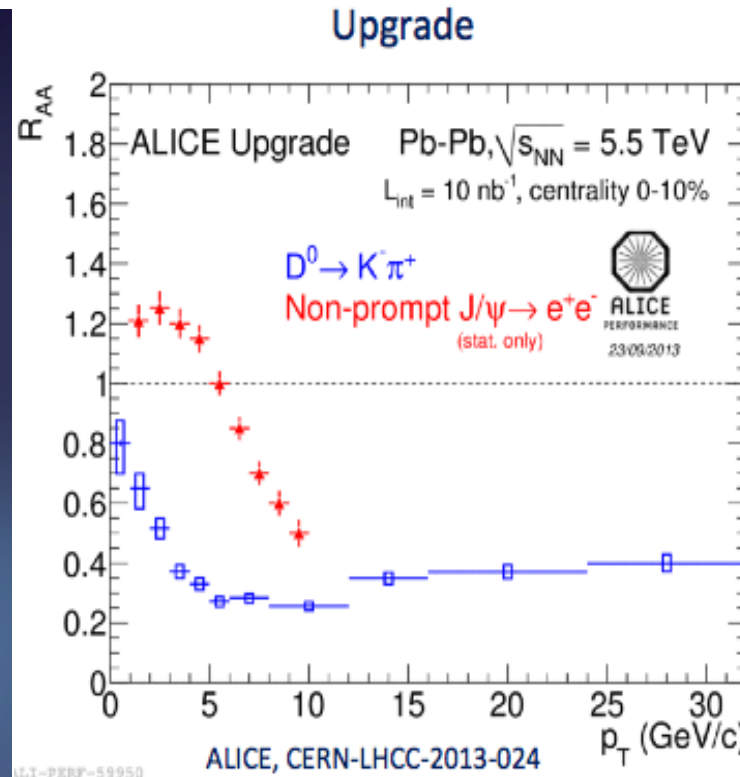
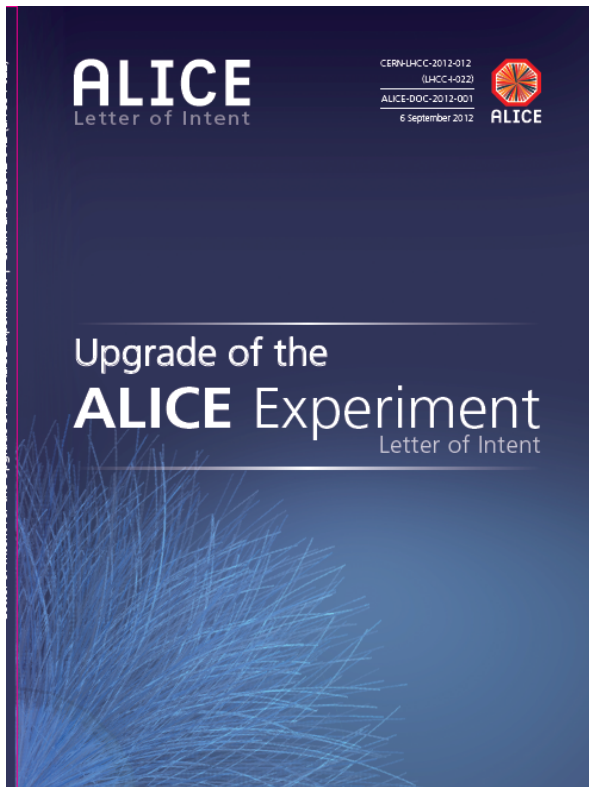


open charm hadrons in pp collisions consistent with quarks hadronizing at about  $T = 165$  MeV

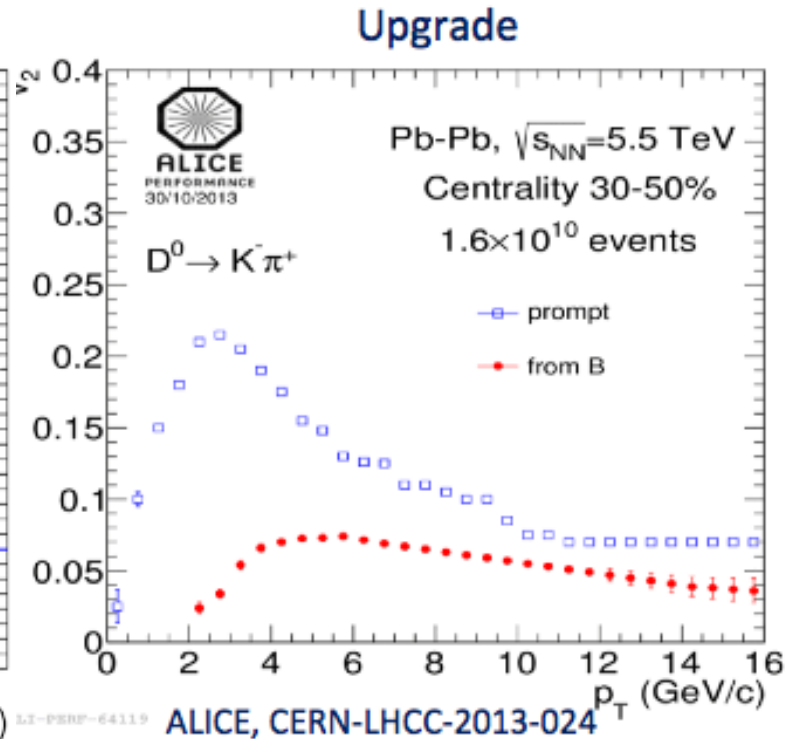
what about PbPb collisions? all D and  $\Lambda_c$  states predicted. Data to come soon!

# outlook open heavy flavor – LHC run3

new high performance ITS plus rate increase (TPC upgrade)



Charm and beauty  $R_{AA}$  down to  $p_T \sim 0$  using  $D^0$  and B-decay  $J/\psi$



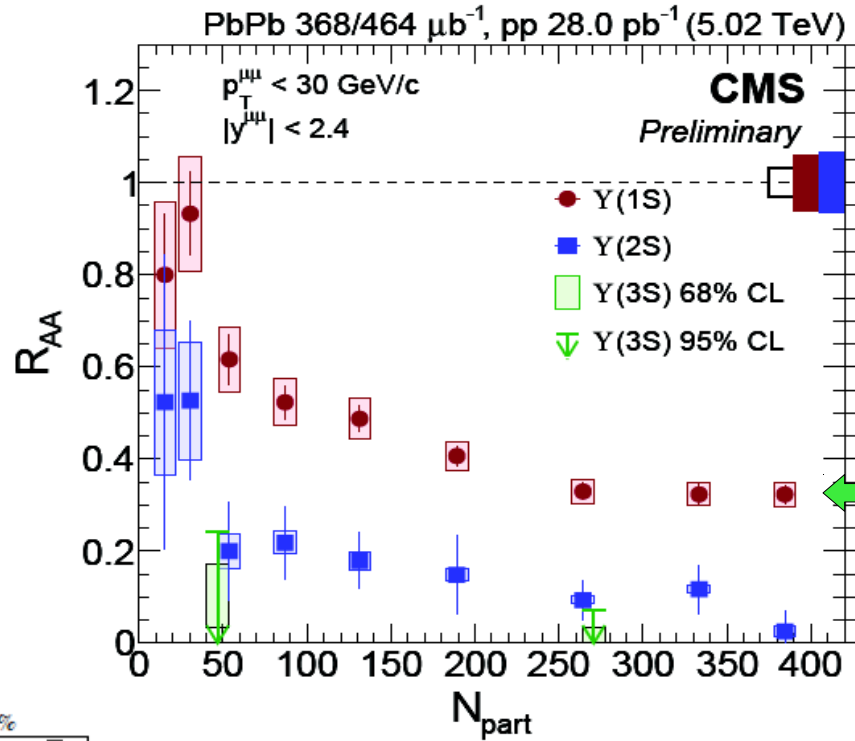
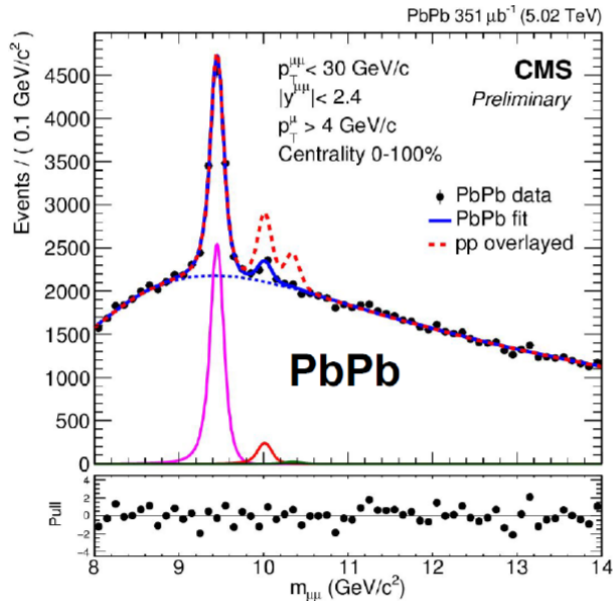
ALICE, CERN-LHCC-2013-024  
Input values from BAMPS model: C. Greiner et al. arXiv:1205.4945

Charm  $v_2$  down to  $p_T \sim 0$  using prompt and beauty  $v_2$  down to B  $p_T \sim 0$  using B-decay  $D^0$

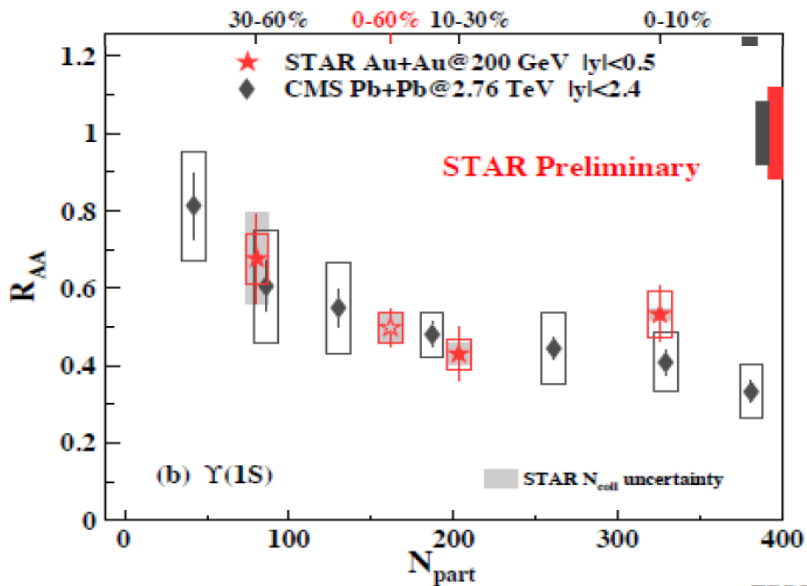
# bottomonia



# Suppression of Upsilon states



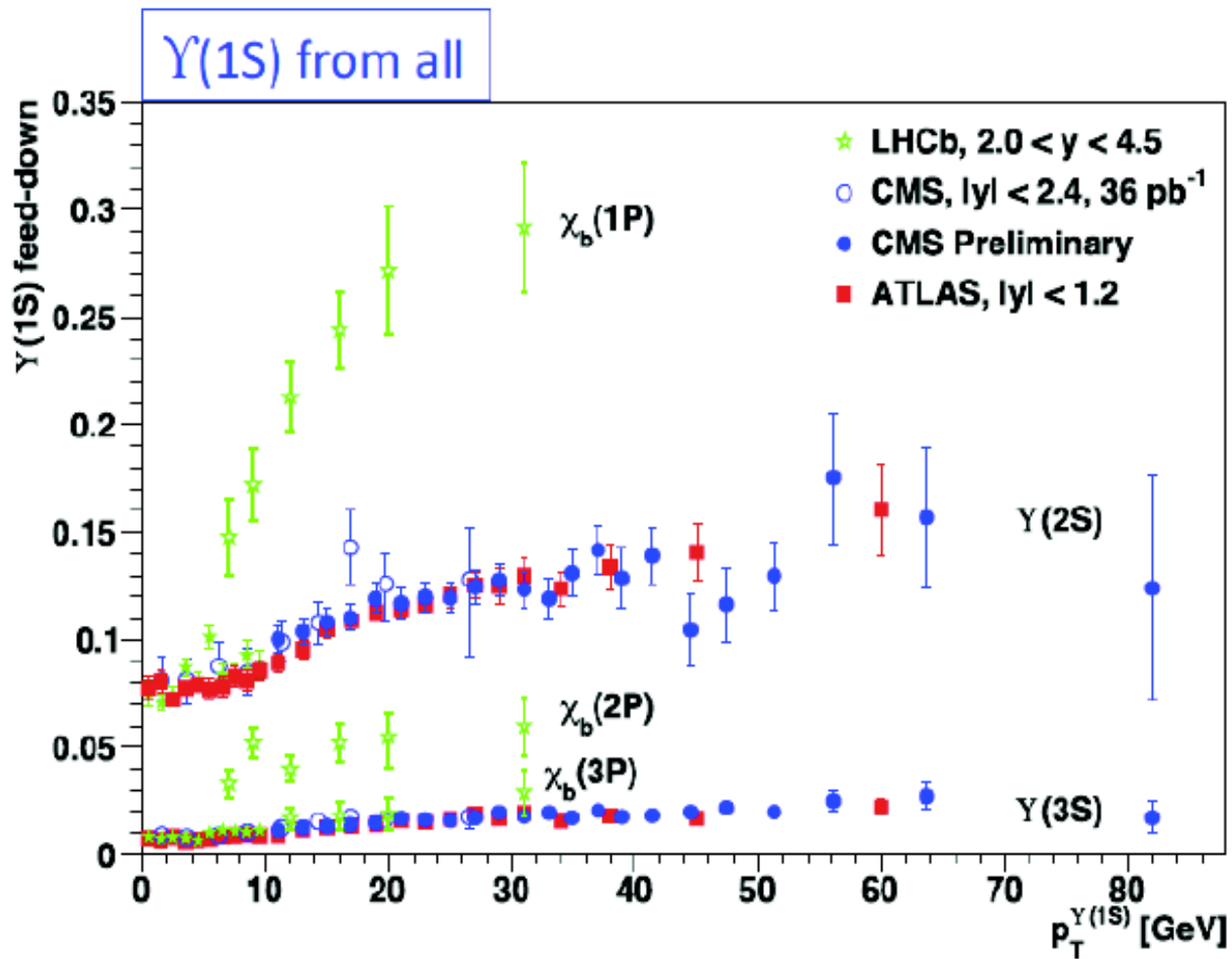
not consistent with just excited state suppression (LHCb data: only 25 % feed-down in pp at LHC)



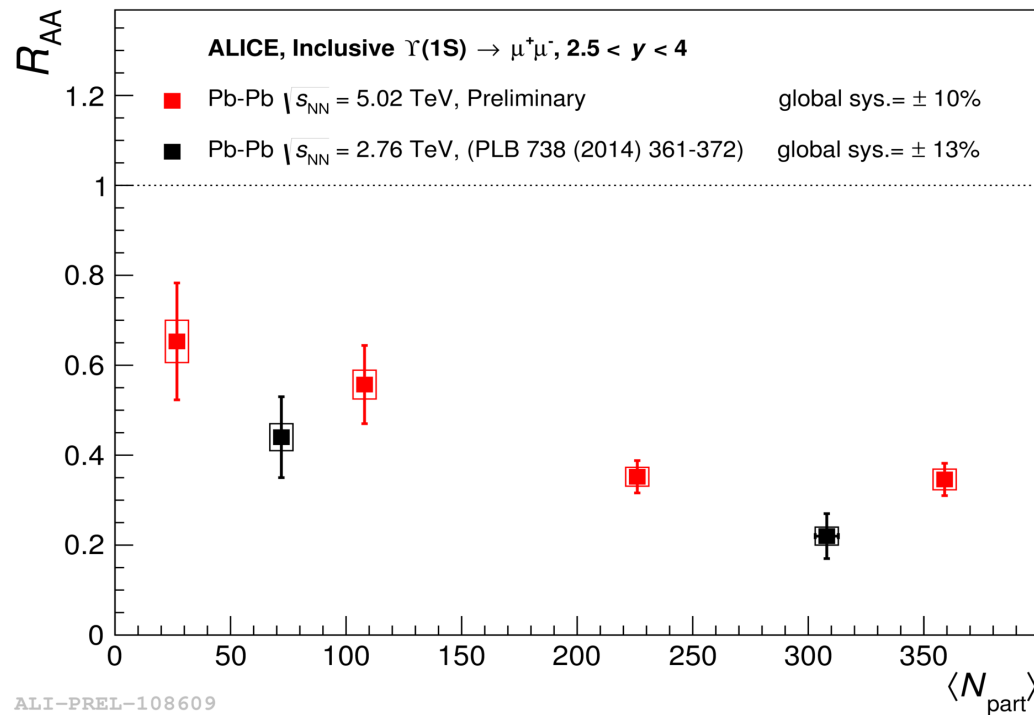
genuine Upsilon suppression

- real and imaginary part of potential at finite temperature play a role
- similarity of RHIC and LHC suppression reminiscent of SPS and RHIC for  $J/\psi$
- possibility of statistical hadronization?

# Feeding into Upsilon (1S)



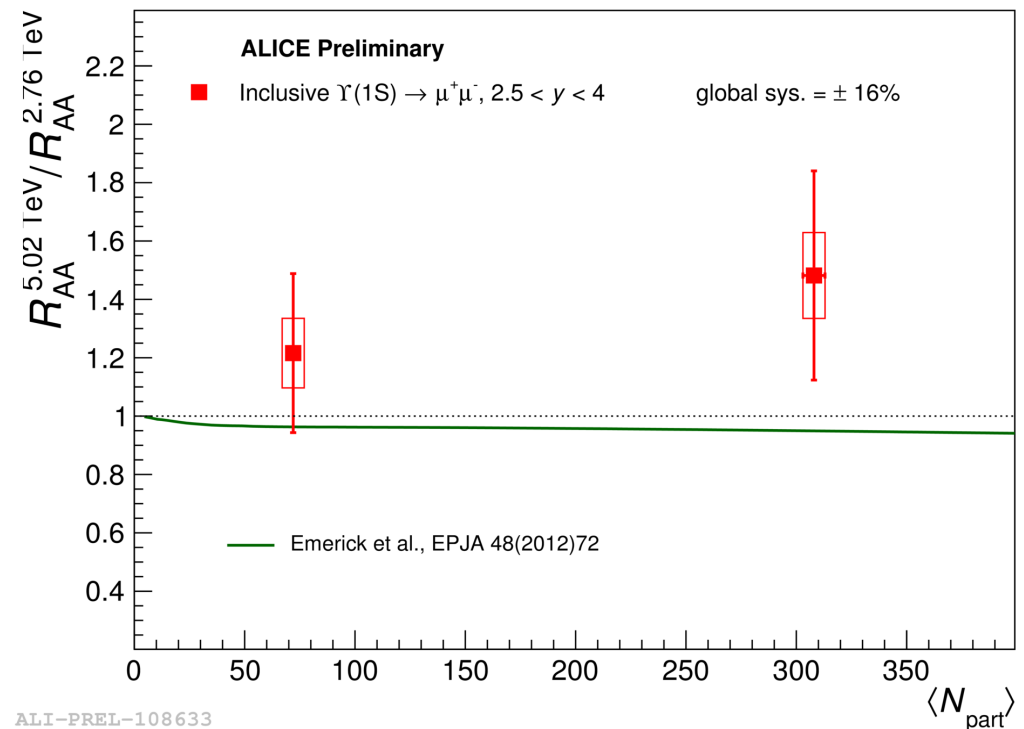
# Upsilon in PbPb at 5 TeV compared to 2.76 TeV



ALI-PREL-108609

dissociation of Upsilon in a hydrodynamically medium will not produce an increase with increasing energy density

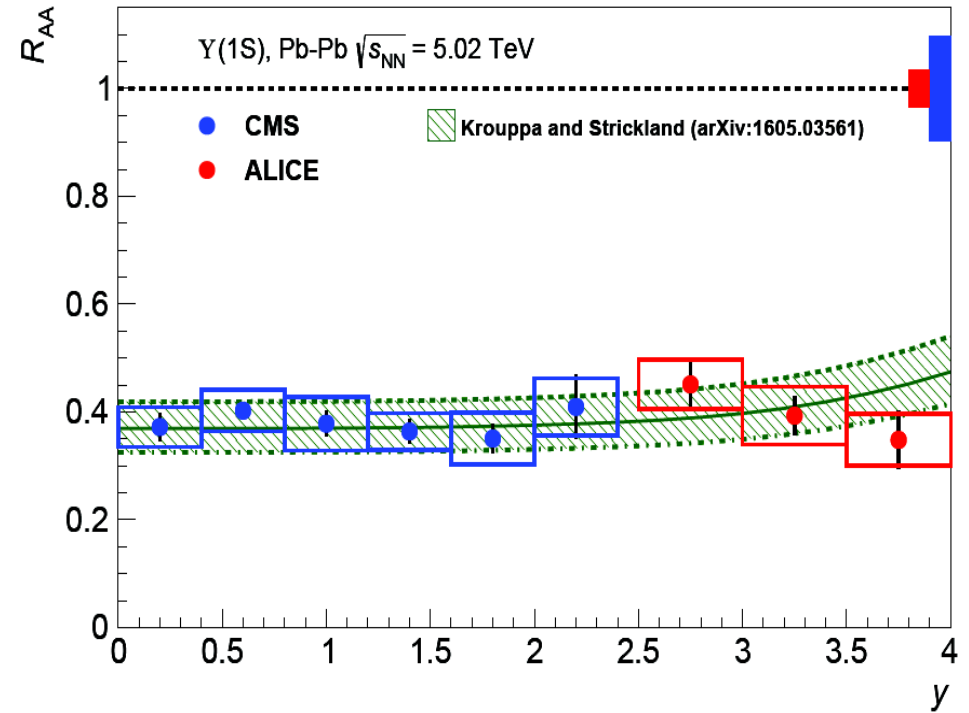
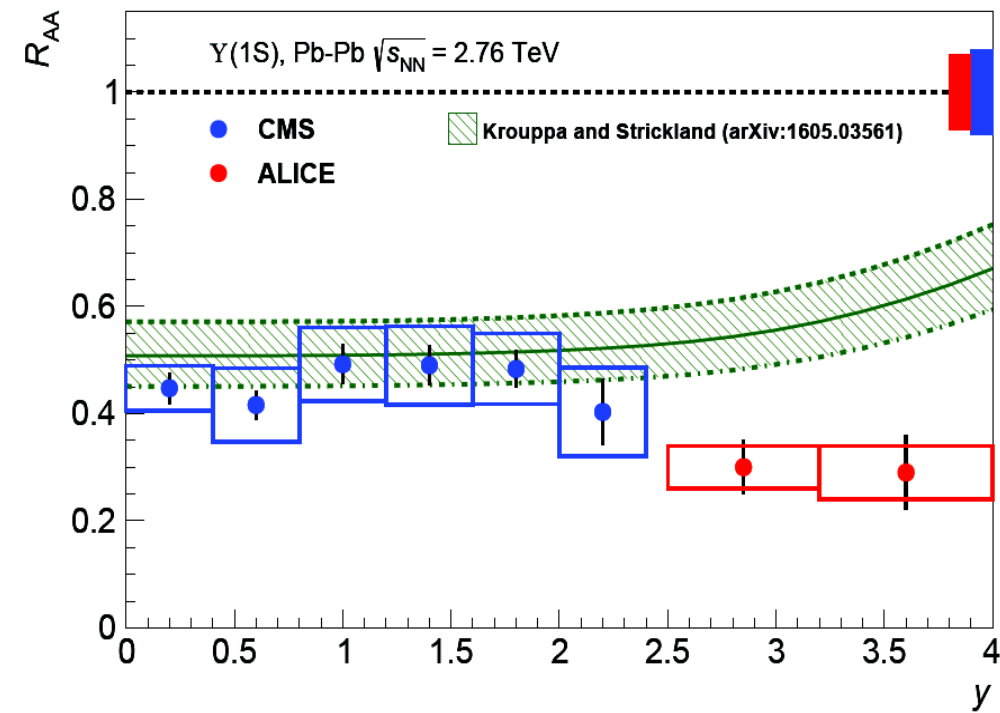
yield of Upsilon(1S) increases with beam energy



ALI-PREL-108633

$$R_{AA}^{0-90\%}(5.02 \text{ TeV}) / R_{AA}^{0-90\%}(2.76 \text{ TeV}) = 1.3 \pm 0.2(\text{stat}) \pm 0.2(\text{syst})$$

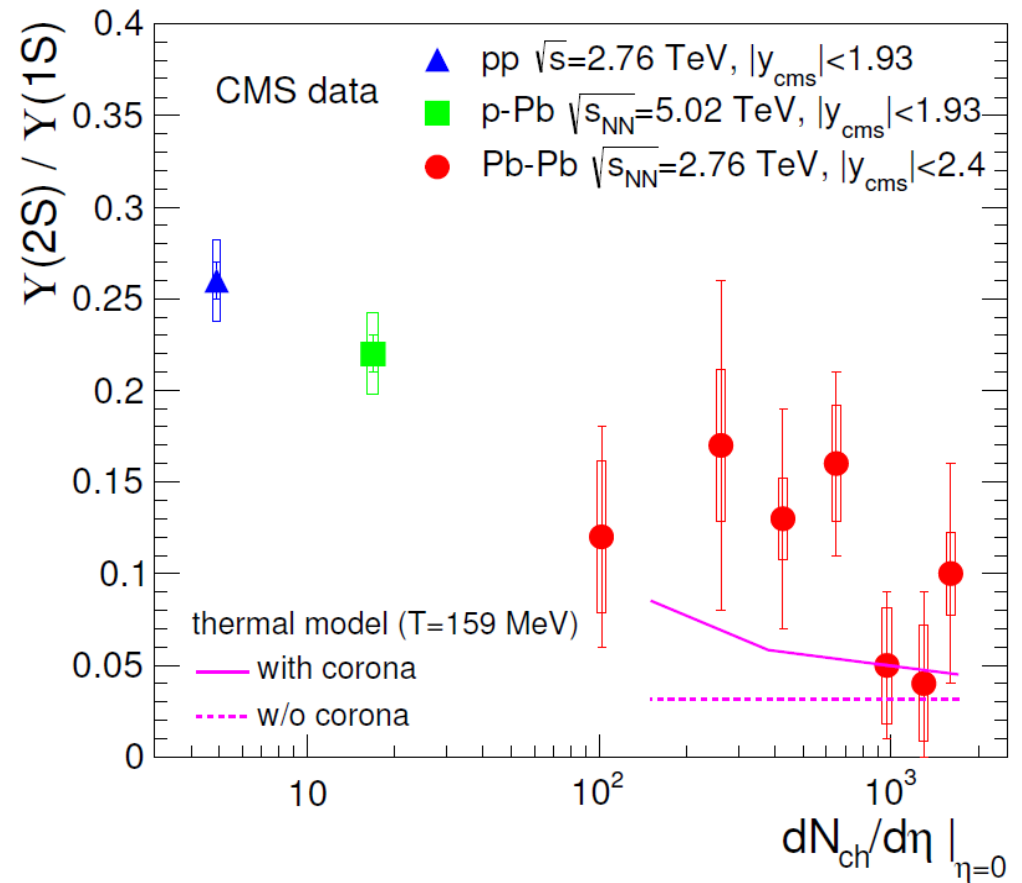
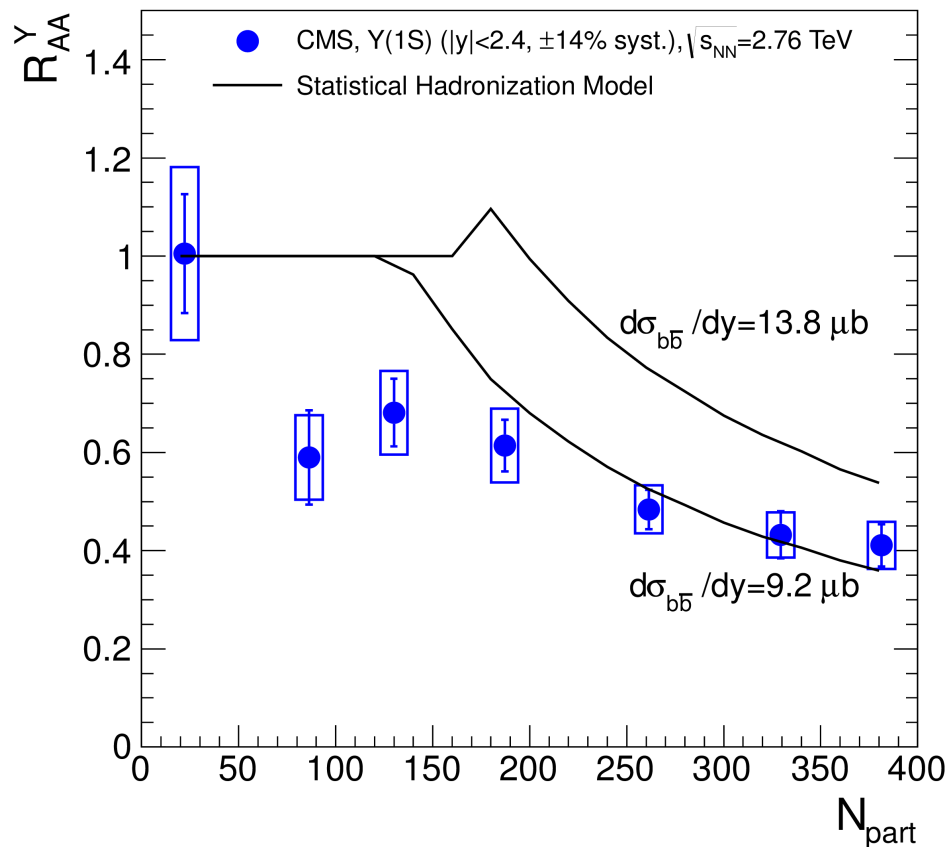
# Upsilon $R_{AA}$ rapidity dependence



Indication:  $R_{AA}$  peaked at mid- $y$  like for  $J/\psi$   
not in line with collisional damping in expanding medium

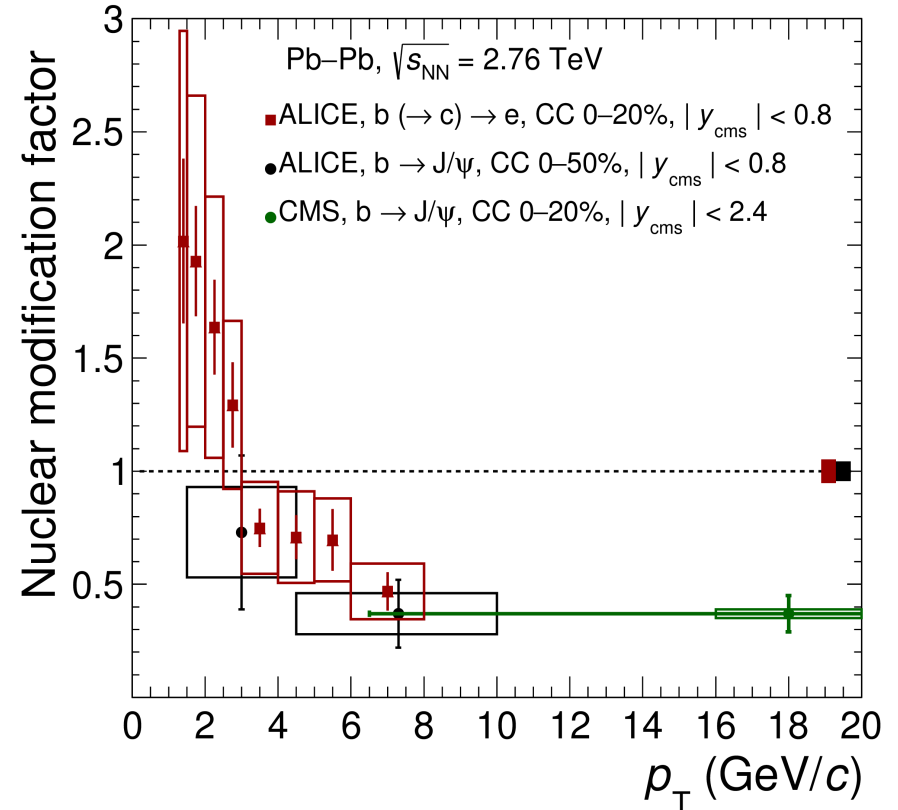
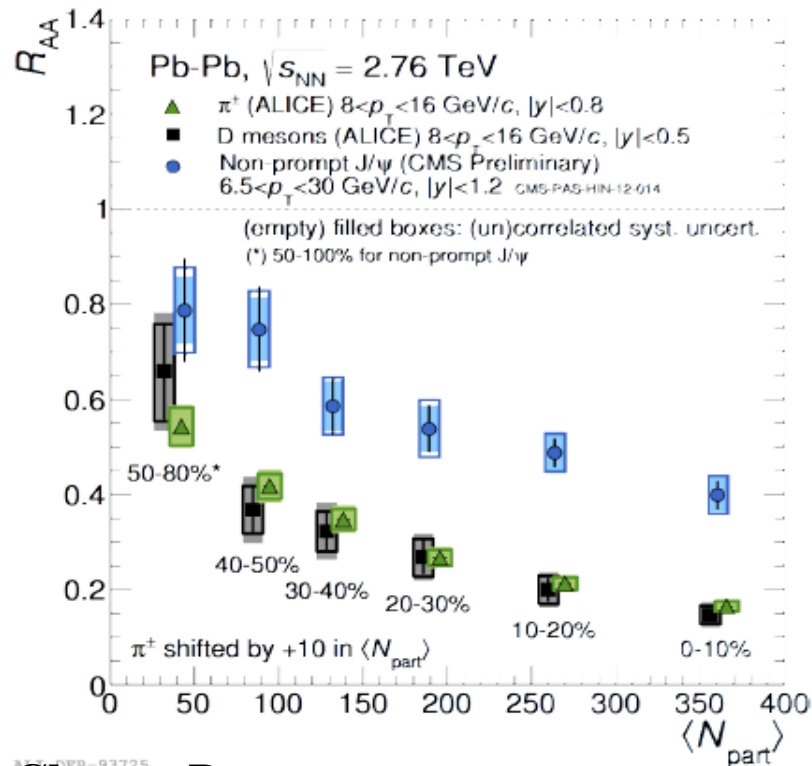
# the Upsilon could also come from statistical hadronization

SHM/thermal model: Andronic et al.



in this picture, the entire Upsilon family is formed at hadronization  
 but: need to know first – do b-quark thermalize at all? spectra of B  
 - total b-cross section in PbPb

# what about b-quark energy loss and thermalization?



Charm: D mesons

JHEP 11 (2015) 205

Beauty: non-prompt J/ $\psi$

arXiv:1610.00613

$b \rightarrow c \rightarrow e$  arXiv:1609.03898

separation via impact parameter distribution

- mass ordering between charm and beauty observed

- for more central collisions, electrons from b-decay show suppression for  $p_T > 3$  GeV/c