WHAT CAN WE LEARN ABOUT R-PROCESS AND NUCLEAR EOS FROM NS-MERGER OBSERVATIONS?

OLIVER JUST ASTROPHYSICAL BIG BANG LABORATORY RIKEN

> **6TH FAIRNESS WORKSHOP** GENOVA, MAY 20TH,

WITH: H.-TH.-JANKA, A. BAUSWEIN, S. GORIELY, R. ARDEVOL, M. OBERGAULINGER, S. NAGATAKI, M. WU, I. TAMBORRA M., AND OTHERS







Max Planck Institute for Astrophysics

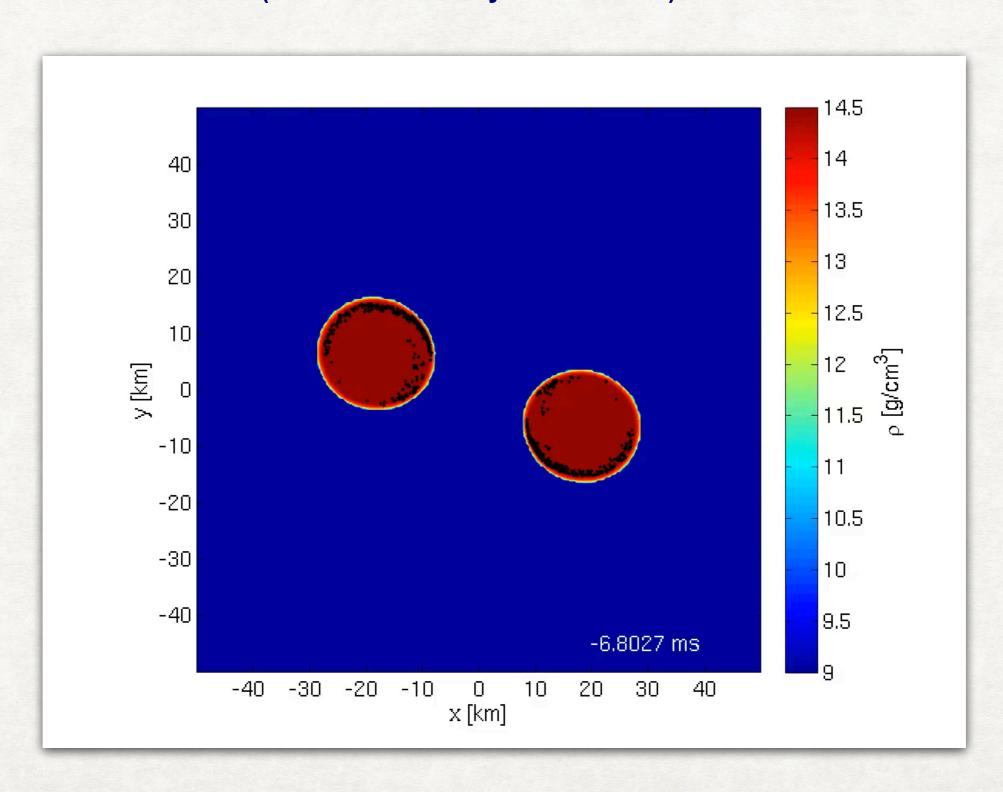






Established by the European Commission

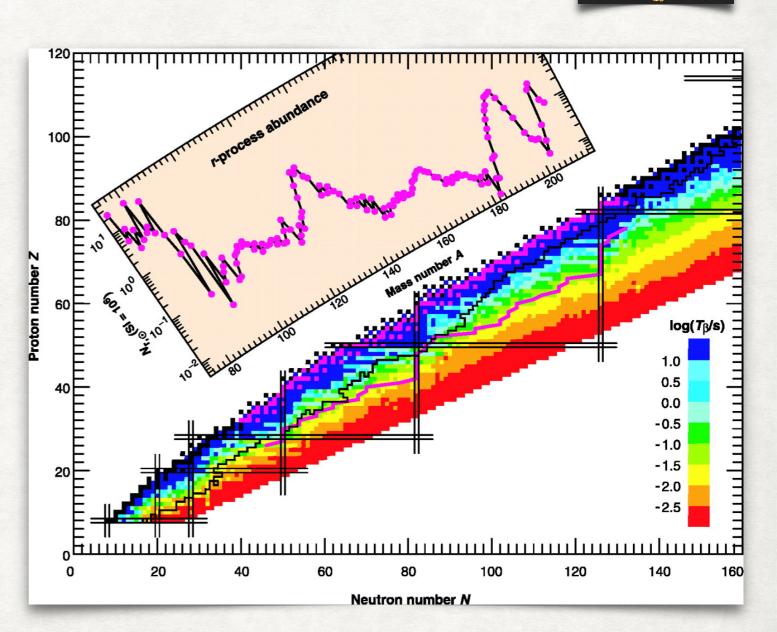
Movie: NS-NS Merger (SPH simulation by A. Bauswein)



(...SUCH AS GOLD)

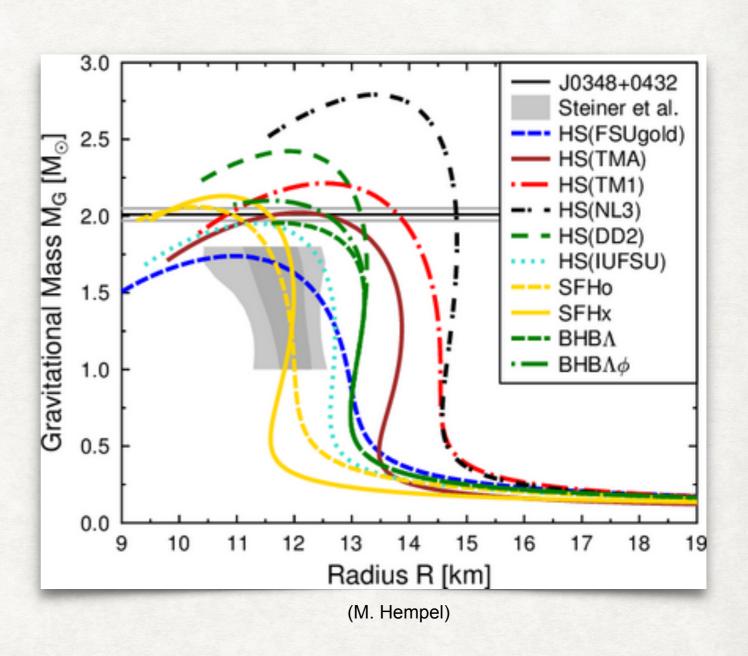
NEUTRON-STAR MERGERS AS SOURCES OF R-PROCESS ELEMENTS

- Are neutron-star mergers significant/dominant sources?
- need explosive conditions with very high neutron densities
- neutron density depends sensitively on neutrino interactions
- suggested alternatives:
 - core-collapse supernovae
 - jets of magneto-rotational supernovae
 - accretion disks in collapsars



NS MERGERS AS KEY TOWARDS UNDERSTANDING THE NUCLEAR EOS?

- EOS determines Mass-Radius relationship for neutron stars
- each EOS characterized by typical NS radius and maximum mass



ejecta mass+velocity

> ejection mechanism

nucleosynthesis yields



collapse time of HMNS

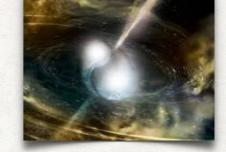
nuclear EOS constraints

jet launching mechanism gravitational waves

radio

optical

UV



X-ray

gamma-ray

ejecta mass+velocity

ejection mechanism



nucleosynthesis yields



collapse time of HMNS



nuclear EOS constraints

modeling "black box"

- ★ hydrodynamic models from numerical simulations including GR, neutrino transport, MHD
- nucleosynthesis yields from nuclear physics calculations
- ★ atomic shell calculations and radiative transfer for EM light curve

gravitational waves

radio

optical





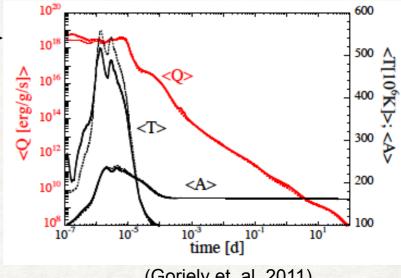
X-ray

gamma-ray

jet launching mechanism

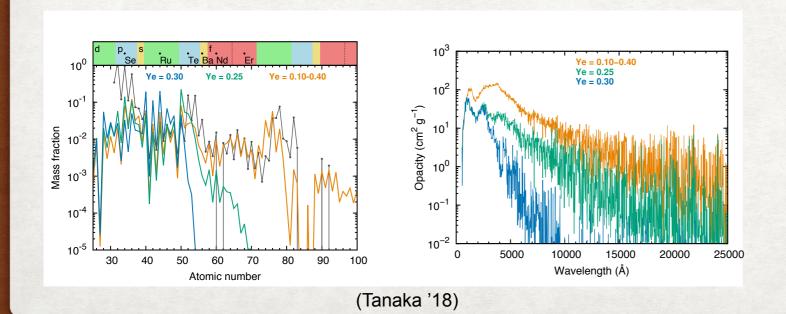
Kilo-/Macronovae

- radioactive decay heats newly synthesized material
- light curves contain valuable information about mass, velocity, and composition



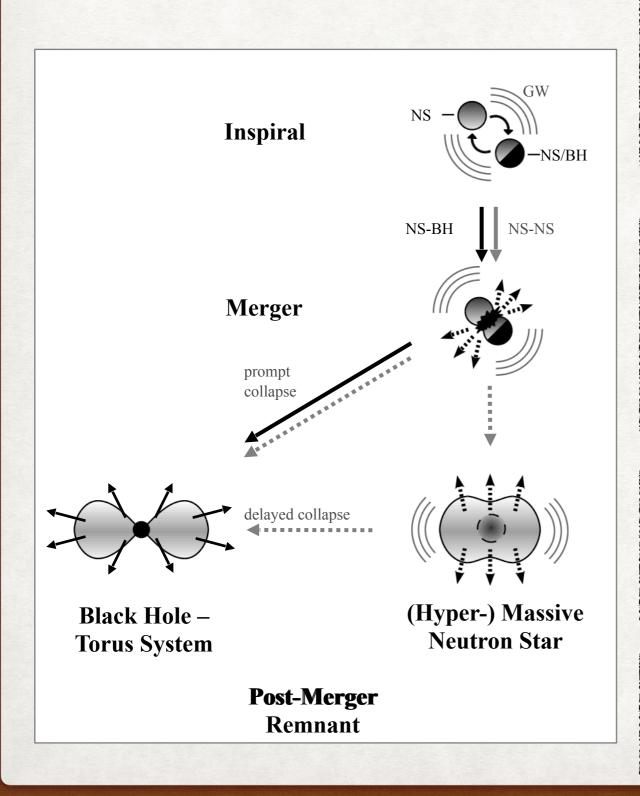
(Goriely et. al. 2011)

$$\begin{split} t_{\rm peak} &\approx \sqrt{\frac{\kappa M_{\rm ej}}{4\pi c V_{\rm ej}}} \xi \approx 1.5 \, {\rm days} \, \left(\frac{V_{\rm ej}}{0.1 \, c}\right)^{-1/2} \left(\frac{M_{\rm ej}}{0.03 \, M_{\odot}}\right)^{1/2} \left(\frac{\kappa}{0.3 \, {\rm cm}^2 \, {\rm g}^{-1}}\right)^{1/2} \xi^{1/2}, \\ L_{\rm peak} &\approx \frac{f M_{\rm ej} c^2}{t_{\rm peak}} \approx 4.3 \times 10^{41} {\rm erg \, s}^{-1} \, \left(\frac{f}{10^{-6}}\right) \left(\frac{V_{\rm ej}}{0.1 \, c}\right)^{1/2} \left(\frac{M_{\rm ej}}{0.03 \, M_{\odot}}\right)^{1/2} \left(\frac{\kappa}{0.3 \, {\rm cm}^2 \, {\rm g}^{-1}}\right)^{-1/2} \xi^{-1/2}, \\ T_{\rm eff,peak} &\approx \left[\frac{L_{\rm peak}}{4\pi (V_{\rm ej} t_{\rm peak})^2 \sigma_{\rm SB}}\right]^{1/4} \approx 8 \times 10^3 \, {\rm K} \, \left(\frac{f}{10^{-6}}\right)^{1/4} \left(\frac{V_{\rm ej}}{0.1 \, c}\right)^{1/8} \left(\frac{M_{\rm ej}}{0.03 \, M_{\odot}}\right)^{-1/8} \left(\frac{\kappa}{0.3 \, {\rm cm}^2 \, {\rm g}^{-1}}\right)^{-3/8} \xi^{-3/8} \end{split}$$



(Barnes & Kasen 2013)

Modeling challenges (apart from nuclear physics challenges)



→ Neutrino transport:

- ★ full 6D Boltzmann eq. extremely expensive =>approximations inevitable
- ★ neutrino leakage schemes
- ★ one- or two-moment transport schemes
- ★ Monte-Carlo or ray-tracing schemes
- ★ neutrino oscillations???

→ Magnetic fields and turbulence

- ★ extremely fine resolution needed to resolve all relevant length scales (e.g. magneto-rotational instability), particularly in the HMNS
- many approaches employ effective (alpha-) viscosity

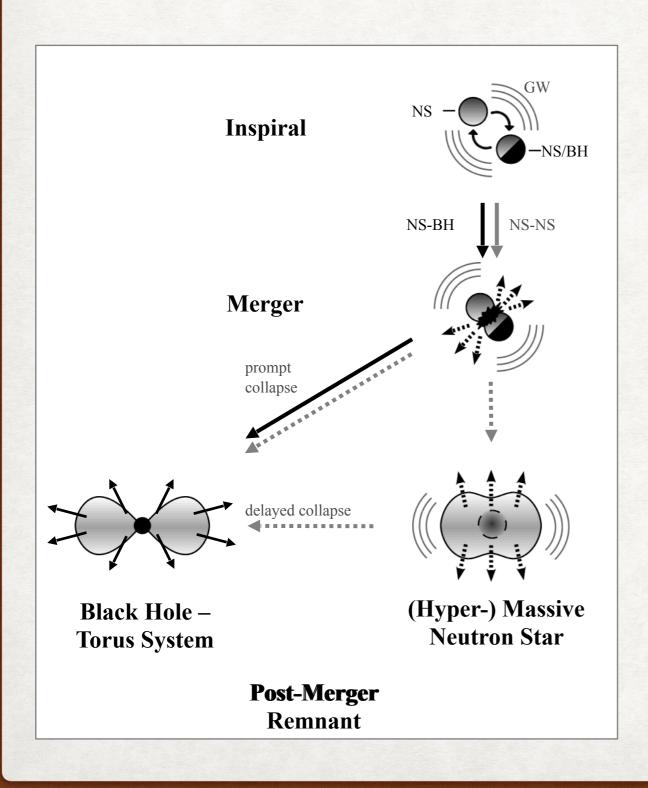
→ General/special relativistic effects

- ★ Newtonian, post-Newtonian methods
- ★ Conformally flat general relativity
- ★ full general relativity

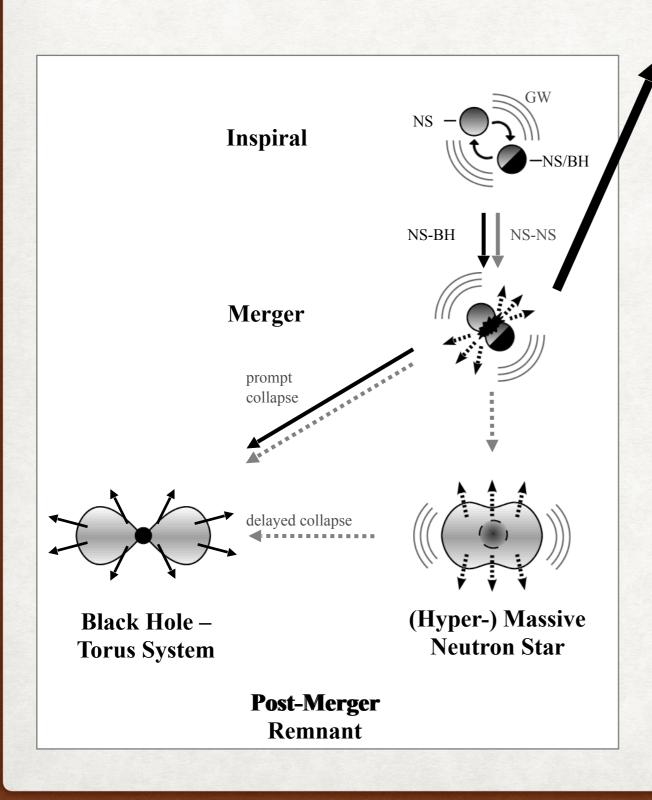
→ Photon transport for electromagnetic signal

- ★ opacities of heavy elements poorly known
- ★ computationally expensive

Neutron-star mergers: ejecta components



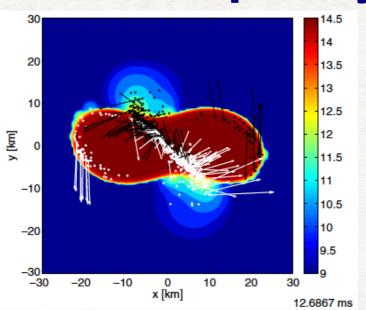
Neutron-star mergers: ejecta components

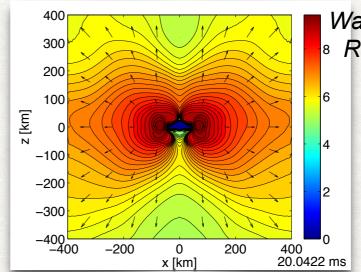


dynamical/prompt ejecta

- → tidal tails
- → shock-heated

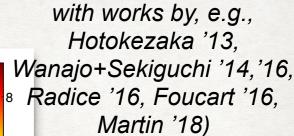
Prompt / dynamical ejecta



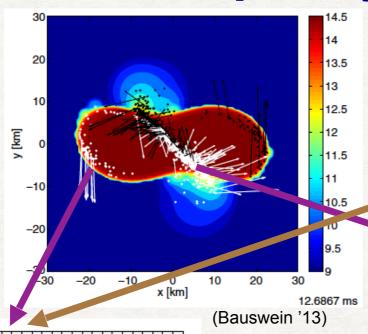


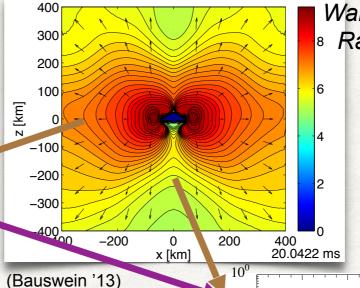
(qualitatively consistent with works by, e.g., Hotokezaka '13, Wanajo+Sekiguchi '14,'16, Radice '16, Foucart '16, Martin '18)

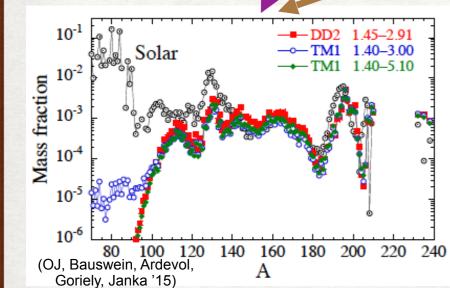
Prompt / dynamical ejecta

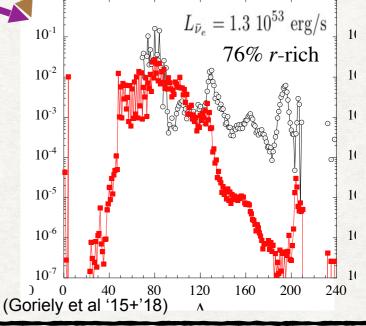


(qualitatively consistent









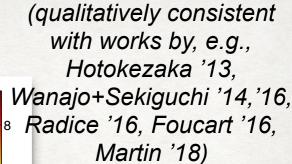
from tidal tails

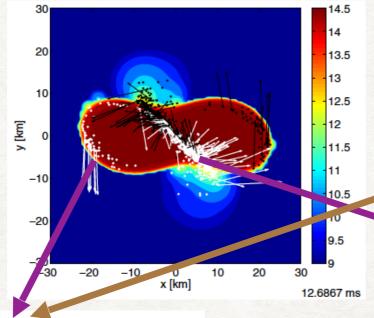
- -> low Ye
- -> more lanthanides
- -> higher opacity
- -> red Kilonova (if observed independently)

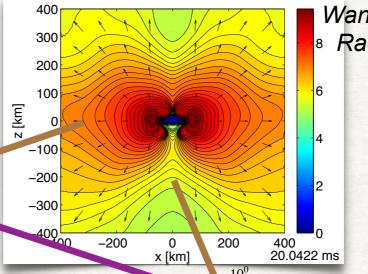
from collision shock

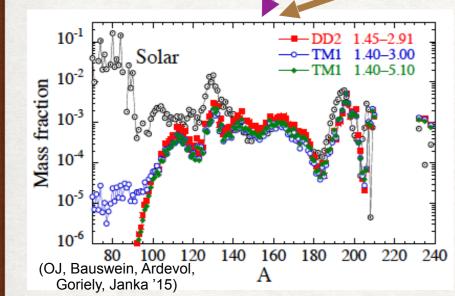
- -> high Ye
- -> less lanthanides
- -> lower opacity
- -> blue Kilonova (if observed independently)

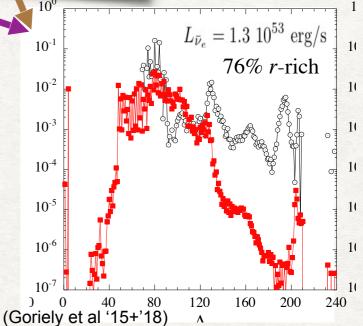
Prompt / dynamical ejecta







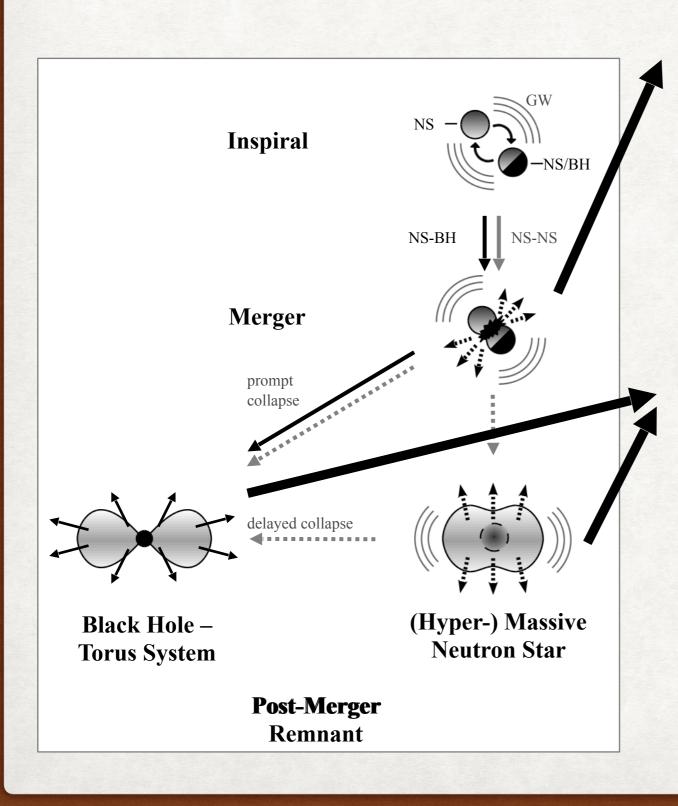




Further properties:

- ★ v ~ 0.1-0.3 c
- ★ M ~ 1E-4 1E-2 Msun
- ★ typically larger mass for softer EOS
- ★ very low mass for prompt collapse

Neutron-star mergers: ejecta components



dynamical/prompt ejecta

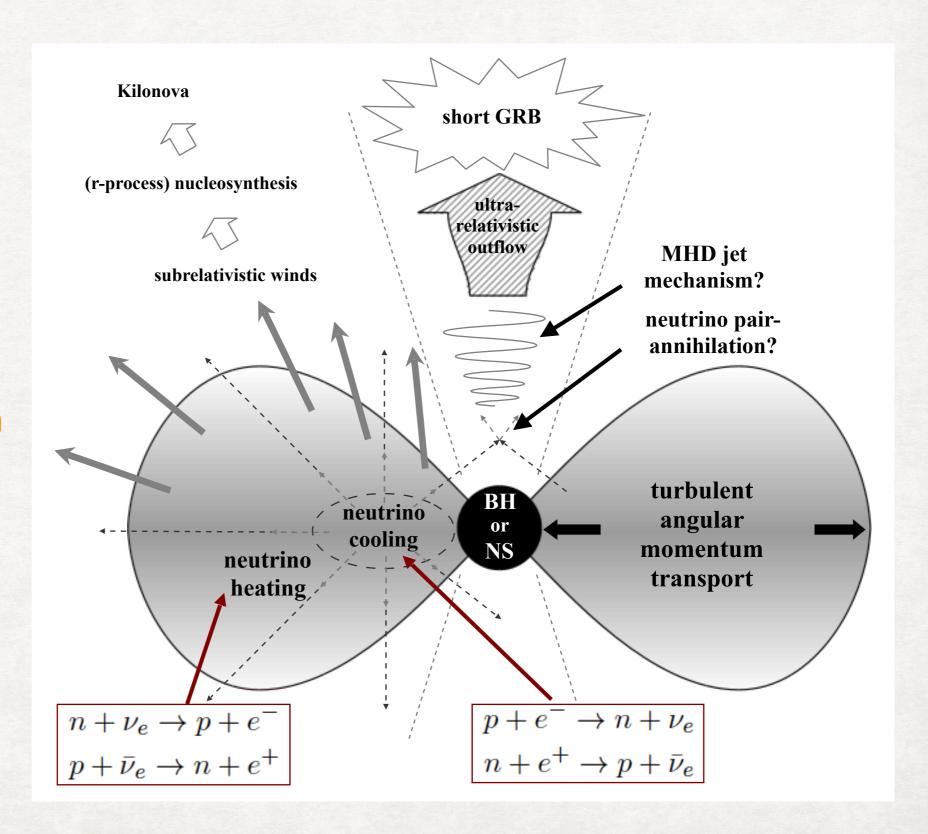
- → tidal tails
- → shock-heated

post-merger ejecta

- → neutrino-driven
- → viscous/MHD driven expansion
- → MHD turbulence

PHYSICS OF POST-MERGER CONFIGURATION

Mass accretion, wind generation, and jet launching highly sensitive to angular momentum transport and neutrino cooling and heating



Post-merger BH-torus remnant

Typical ejecta properties:

- outflow masses:
 - ~ 5-20% of torus mass
 - ~ 1E-4 1E-2 Msun
- electron fraction:

$$<$$
Ye> $\sim 0.1-0.3$

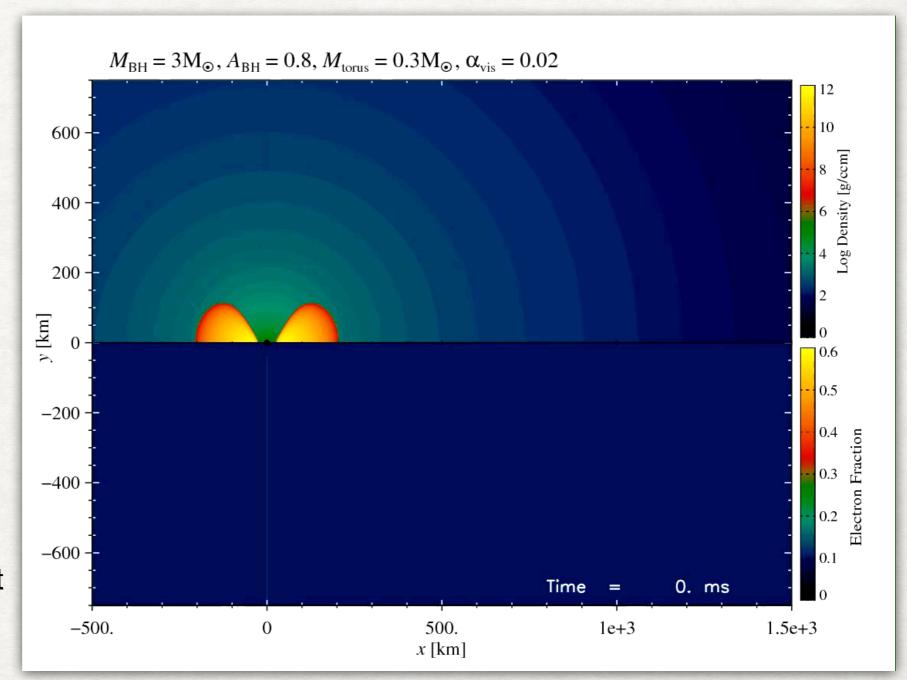
entropy per baryon:

$$< s > ~ 10 - 30 \text{ kB}$$

velocity:

$$<$$
v $> ~ 0.05-0.1 c$

- small neutrino-driven component
- large viscous component



(OJ, Bauswein, Ardevol, Goriely, Janka '15)

"ALCAR" NEUTRINO TRANSPORT MODULE

(OJ, OBERGAULINGER, JANKA '15, MNRAS, 453, 3386)

TWO-MOMENT TRANSPORT WITH ALGEBRAIC EDDINGTON FACTOR (AEF OR M1 SCHEME)

$$E = \int \mathrm{d}\Omega \, \mathcal{I}(\boldsymbol{x},\boldsymbol{n},\epsilon,t) \qquad \leftarrow \text{energy density,} \qquad \text{Oth-angular moment}$$

$$F^i = \int \mathrm{d}\Omega \, \mathcal{I}(\boldsymbol{x},\boldsymbol{n},\epsilon,t) \, n^i \qquad \leftarrow \text{momentum density,} \qquad \text{1st-angular moment}$$

$$P^{ij} = \int \mathrm{d}\Omega \, \mathcal{I}(\boldsymbol{x},\boldsymbol{n},\epsilon,t) \, n^i n^j \qquad \leftarrow \text{pressure,} \qquad \qquad \text{2nd-angular moment}$$

$$Q^{ijk} = \int \mathrm{d}\Omega \, \mathcal{I}(\boldsymbol{x},\boldsymbol{n},\epsilon,t) \, n^i n^j n^k \qquad \left(\partial_t E + \nabla_j \left(\alpha F^j + v^j E \right) + P^{ij} \nabla_i v_j + F^i \nabla_i \alpha \right)$$

evolution equations

$$\partial_{t}E + \nabla_{j}\left(\alpha F^{j} + v^{j}E\right) + P^{ij}\nabla_{i}v_{j} + F^{i}\nabla_{i}\alpha$$

$$-\partial_{\epsilon}\left[\epsilon\left(P^{ij}\nabla_{i}v_{j} + F^{i}\nabla_{i}\alpha\right)\right] = \alpha S^{(0)},$$

$$\partial_{t}F^{i} + \nabla_{j}\left(\alpha c^{2}P^{ij} + v^{j}F^{i}\right) + F^{j}\nabla_{j}v^{i} + c^{2}E\nabla^{i}\alpha$$

$$-\partial_{\epsilon}\left[\epsilon\left(Q^{ijk}\nabla_{j}v_{k} + c^{2}P^{ij}\nabla_{j}\alpha\right)\right] = \alpha S^{(1),i}$$

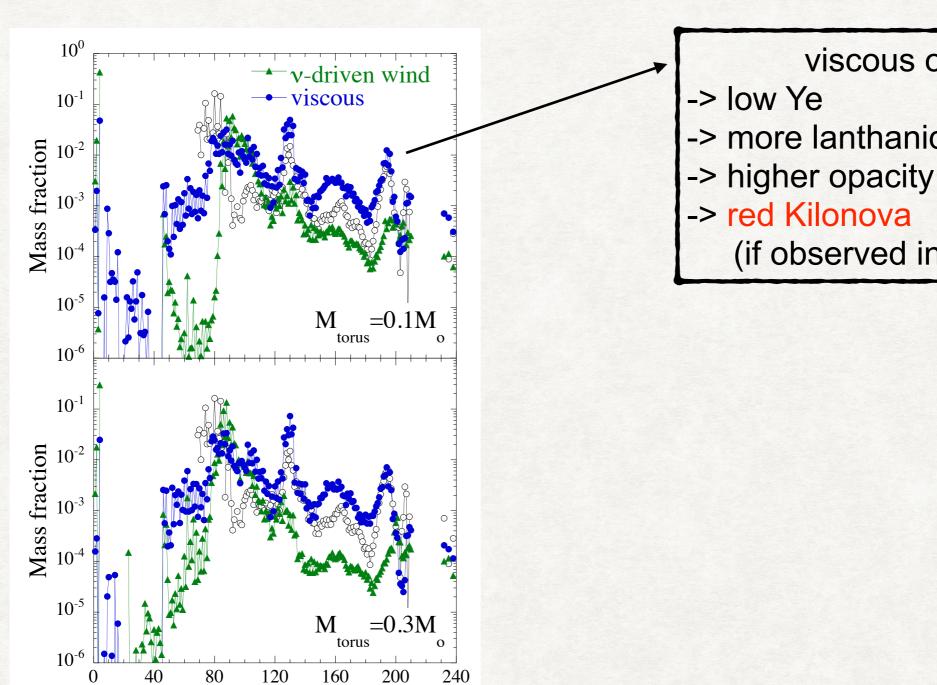
$$P^{ij} = P^{ij}(E, F^i)$$

$$Q^{ijk} = Q^{ijk}(E, F^i)$$

 $P^{ij} = P^{ij}(E, F^i)$ central approximation of M1: $Q^{ijk} = Q^{ijk}(E, F^i)$ local closure relation for higher moments (e.g. "M1 closure")

- => removes two degrees of freedom of nu-phase space large gain of computational efficiency
- => trade-off: potential loss accuracy (at least in optically thin regions)

Nucleosynthesis yields of BH-torus ejecta



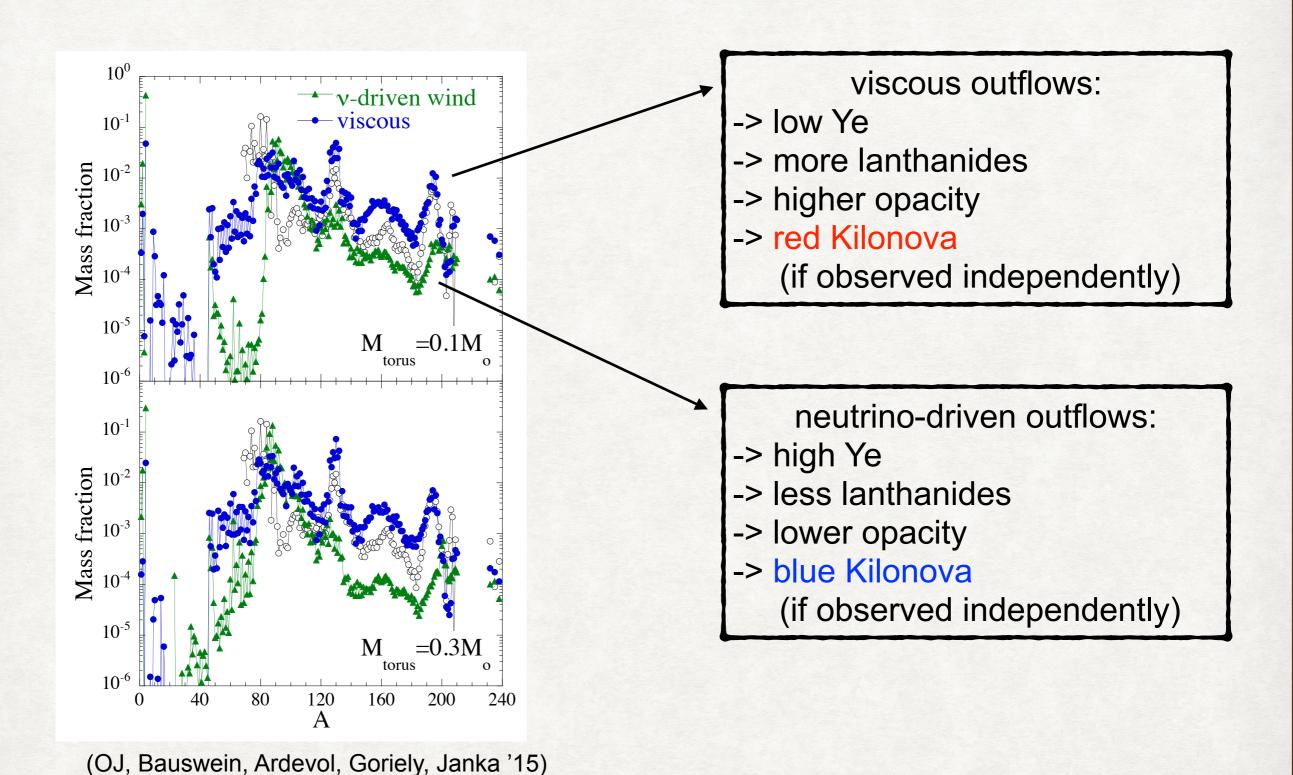
(OJ, Bauswein, Ardevol, Goriely, Janka '15)

A

viscous outflows:

- -> more lanthanides
- (if observed independently)

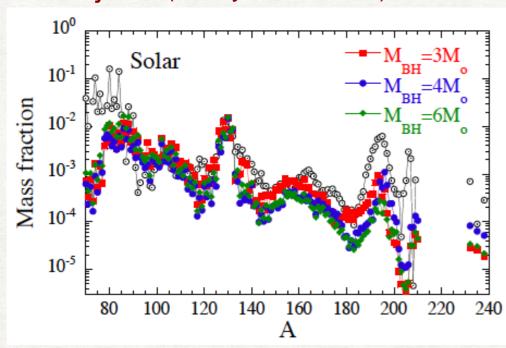
Nucleosynthesis yields of BH-torus ejecta



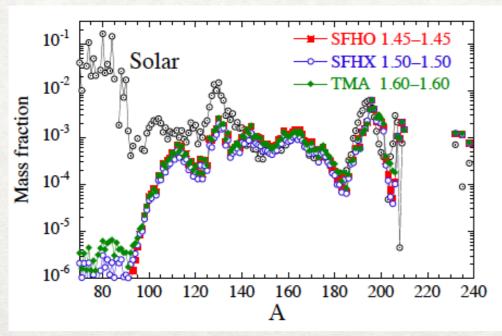
COMBINED NUCLEOSYNTHESIS YIELDS OF PROMPT AND TORUS EJECTA

(OJ, Bauswein, Ardevol, Goriely, Janka '15)

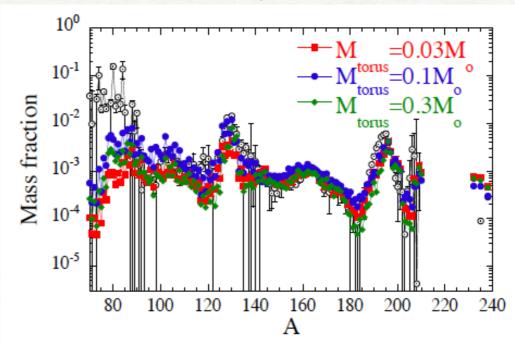
→ DISK ejecta (mainly A ~ 90 - 140)



→ PROMPT ejecta (mainly A ~ 140 - 210)



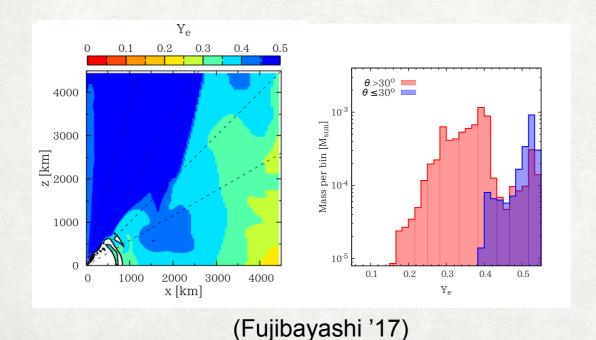
→ DISK + PROMPT ejecta

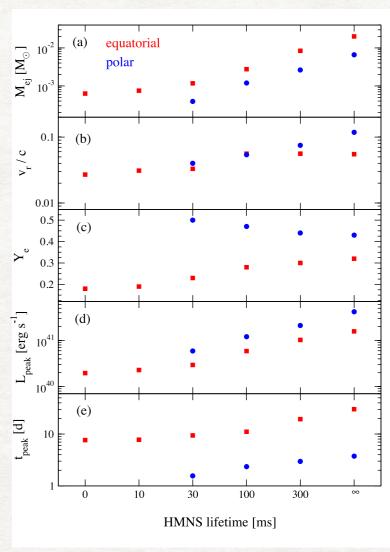


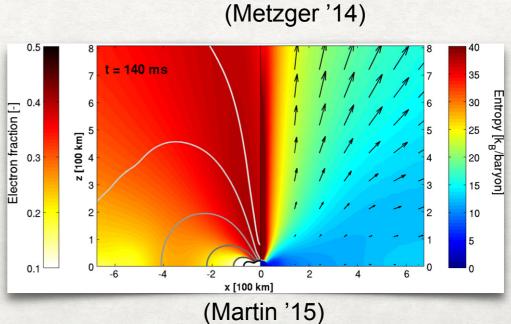
- → nicely recovers the full mass range A > 90
- → strong support for idea that multiple ejecta components of NS mergers are significant sources of r-process elements

NS-torus remnant

- very challenging to model, because of extremely small MRI length scales and complicated high-density neutrino physics
- NS represents much larger energy reservoir than BHtorus
- hence, ejecta mass likely grows with HMNS lifetime
- relative amount of (blue) neutrino-driven to (red)
 viscous ejecta higher



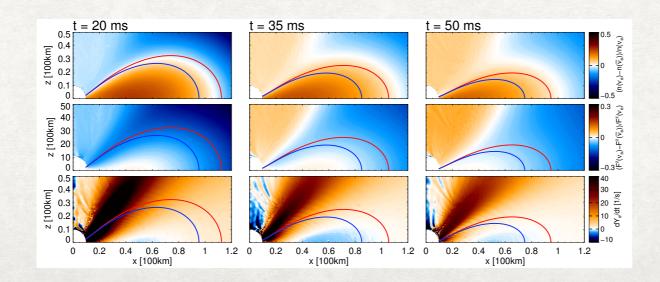


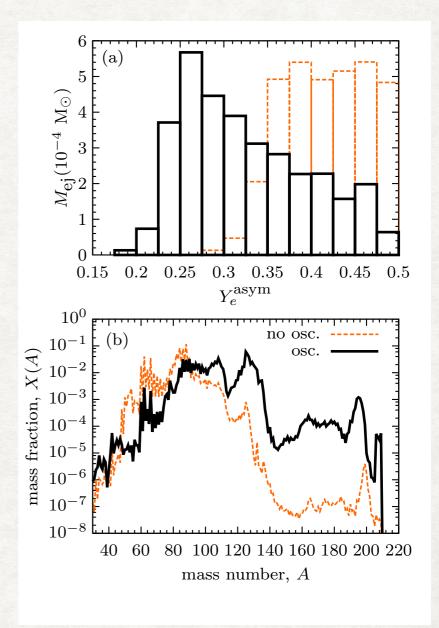


IMPACT OF NU-NU OSCILLATIONS ON THE NEUTRINO-DRIVEN WIND COMPONENT

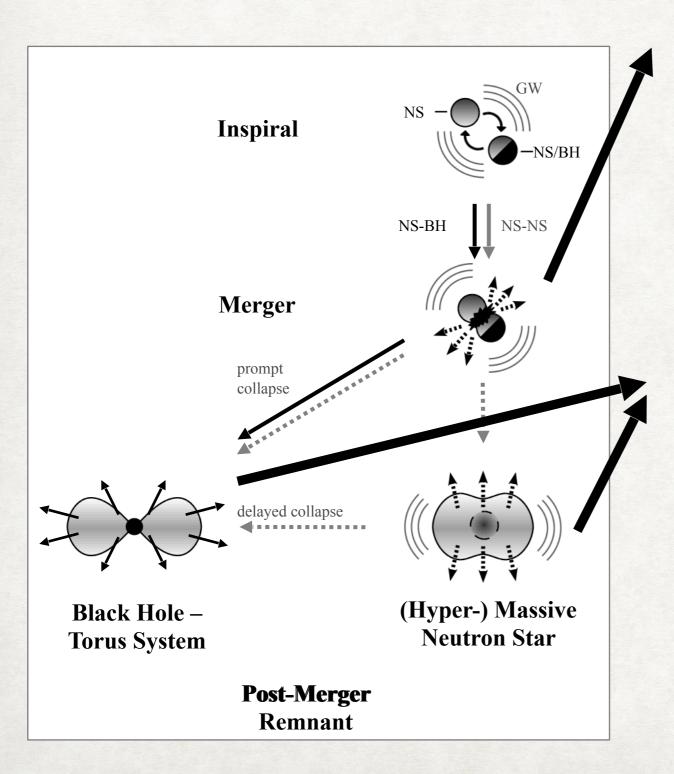
(Wu, Tamborra, OJ, Janka, 2017PhRvD, 96l3015W)

- recently (re-)discovered "fast pairwise flavor conversions" may lead to flavor equilibration on length scales of O(10cm) (e.g. Sawyer+ 05, 09, 16)
- take place whenever $n(\nu_e) n(\bar{\nu_e})$ changes sign in angular space
- our simplified, exploratory study indicates that neutrino-driven ejecta may remain more neutron rich
- neutrino oscillations remain major open question for neutrino effects in NS mergers





Neutron-star mergers: ejecta components



dynamical/prompt ejecta

- → tidal tails
- → shock-heated

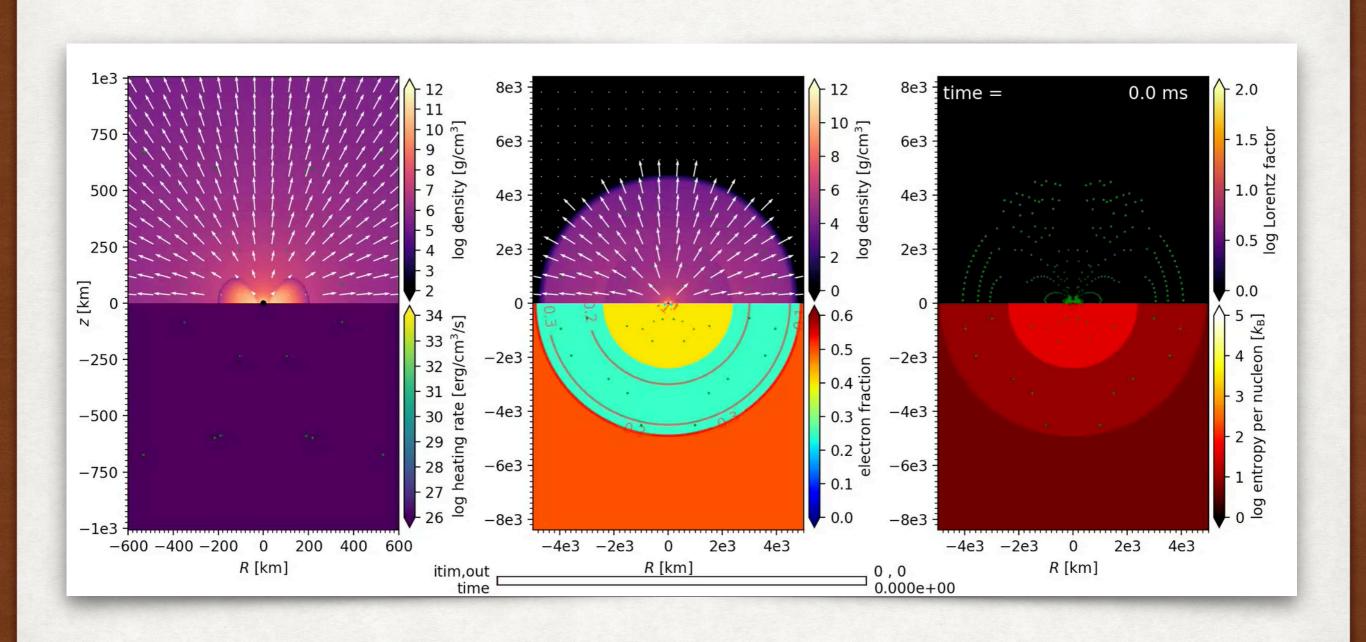
post-merger ejecta

- → neutrino-driven
- → viscous/MHD driven expansion
- → MHD turbulence

(short gamma-ray burst) jet

- → from BH by what mechanism?
- → Magnetar spindown?
- → choked or successful?

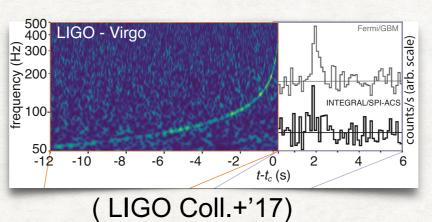
ONGOING WORK: IMPACT OF JET ON NUCLEOSYNTHESIS?

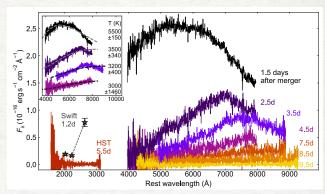


GW170817/AT2017GFO

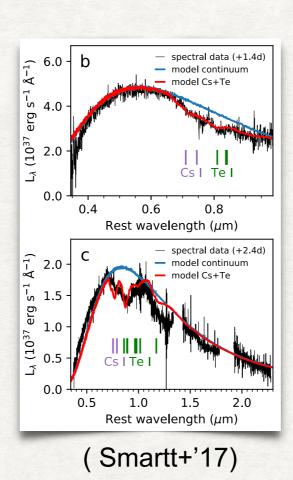
GW170817 + EM COUNTERPARTS

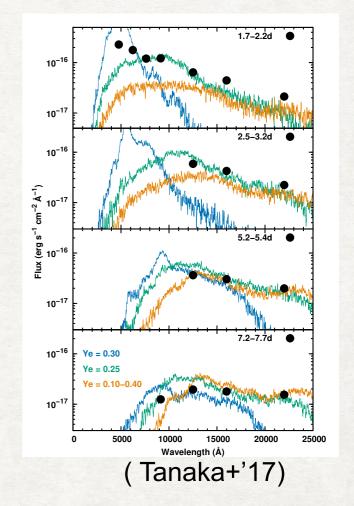
- → Mtot = M1 + M2 ~ 2.74 Msun
- → M1/M2 ~ 0.7 1
- → blue ejecta component with <Y_e> > 0.25 M ~ 0.01-0.03 Msun
- → red ejecta component with <Y_e> < 0.25 M ~ 0.01-0.03 Msun
- → low-luminosity gamma-ray burst with E_{peak} ~ 100keV





(Nicholl+'17)





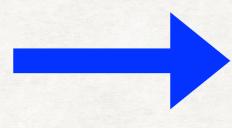
+ many other works by, e.g., Berger, Kasliwal, Kasen, Metzger, Mooley, Piran, Rosswog, Tanvir, Nakar, Gottlieb, MacFadyen, ...

GW170817 + EM COUNTERPARTS

→ Mtot = M1 + M2 ~ 2.74 Msun

→ many studies on interpretation of EM signals, e.g., Kasen '18, Shibata '18, Metzger '18, Mooley '18, Gottlieb '18, Bromberg '18, MacFadyen '18, ...

- \rightarrow M1/M2 ~ 0.7 1
- → blue ejecta component with <Y_e> > 0.25 M ~ 0.01-0.03 Msun



shock-heated dynamical ejecta and/or neutrino-processed ejecta launched from a HMNS remnant? High mass and velocity still enigmatic...

→ red ejecta component with <Y_e> < 0.25 M ~ 0.01-0.03 Msun



dynamical ejecta launched during merger or viscous ejecta from the remnant?

→ low-luminosity gamma-ray burst with E_{peak} ~ 100keV

shock breakout emission from choked jet or cocoon emission from structured jet. Recent observations favor successful jet (Mooley+18). High E_{peak} still puzzling...

→ major open questions remain...better models needed

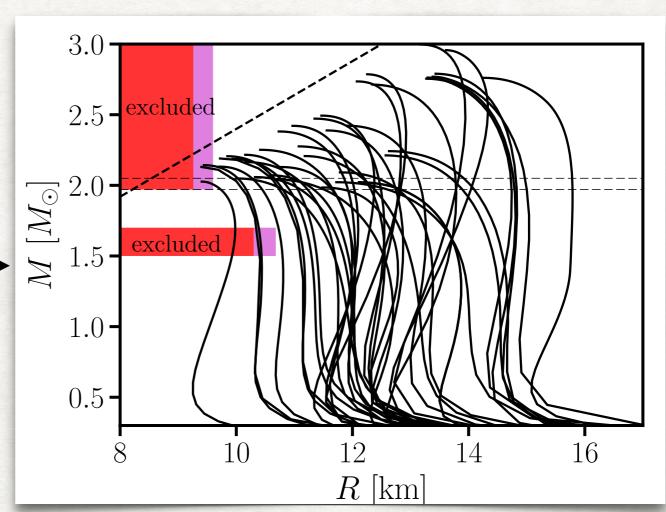
DELAYED COLLAPSE IS VERY LIKELY FOR GW170817: IMPLICATION FOR NUCLEAR EOS

(Bauswein, OJ, Janka, Stergioulas, 2017ApJ, 850L, 34B)

empirical relation deduced from NS merger simulations (Bauswein+'13, PRL 111, 131101)

$$M_{\text{thres}} = \left(-3.606 \frac{GM_{\text{max}}}{c^2 R_{1.6}} + 2.38\right) \cdot M_{\text{max}}$$

$$M_{\text{thres}} = \left(-3.38 \frac{GM_{\text{max}}}{c^2 R_{\text{max}}} + 2.43\right) \cdot M_{\text{max}}$$



conservative constraint assuming HMNS lifetime > 10 ms



 $R_{16} > 10.7 \text{ km}$

 $R_{max} > 9.6 \text{ km}$

(also several other EOS studies, e.g., Margalit+ '17, Rezzolla+ '17, Ruiz+ '17, Radice+ '17)

Take Home Messages

- NS mergers produce a variety of different outflow components, each with individual nucleosynthesis signature as well es EM signal
- ejecta are less neutron-rich (i.e. blue) if illuminated by neutrinos, e.g. for shock-heated dynamical ejecta or neutrino-driven winds from remnant
- ejecta more neutron-rich (i.e. red) if neutrino illumination is low, e.g. for tidal dynamical ejecta or viscous remnant ejecta
- ejecta more massive for longer lifetime of HMNS
- GW170817 suggests large amount of both red and blue ejecta, hence likely a delayed collapse
- likely mixture of dynamical, neutrino-driven, viscously driven, MHD-driven ejecta => however, safe identification not yet possible => better models needed
- new radius constraint: delayed collapse suggests R_NS > 10.7 km
- neutrino oscillations might have strong impact on Y_e, however, still extremely uncertain