

Exploring QCD Matter

at high baryon density



Daniel Brandenburg,
Brookhaven National Laboratory

FAIRNESS : FAIR Next Generation Scientists
20-24 May 2019 Grand Hotel Arenzano,
Genova, Italy

Talk Outline

2

Introduction

- Big Questions: Why Study QCD Matter?
- The QCD Phase Diagram
- What Theory can tell us (experimentalists perspective)

Studying QCD Matter with Heavy Ion Collisions

- Basics of Heavy Ion Collisions
- Scanning the phase diagram
- Search for 1st Order phase transition & Critical point
- Search for Chiral Symmetry Restoration

Summary & Outlook

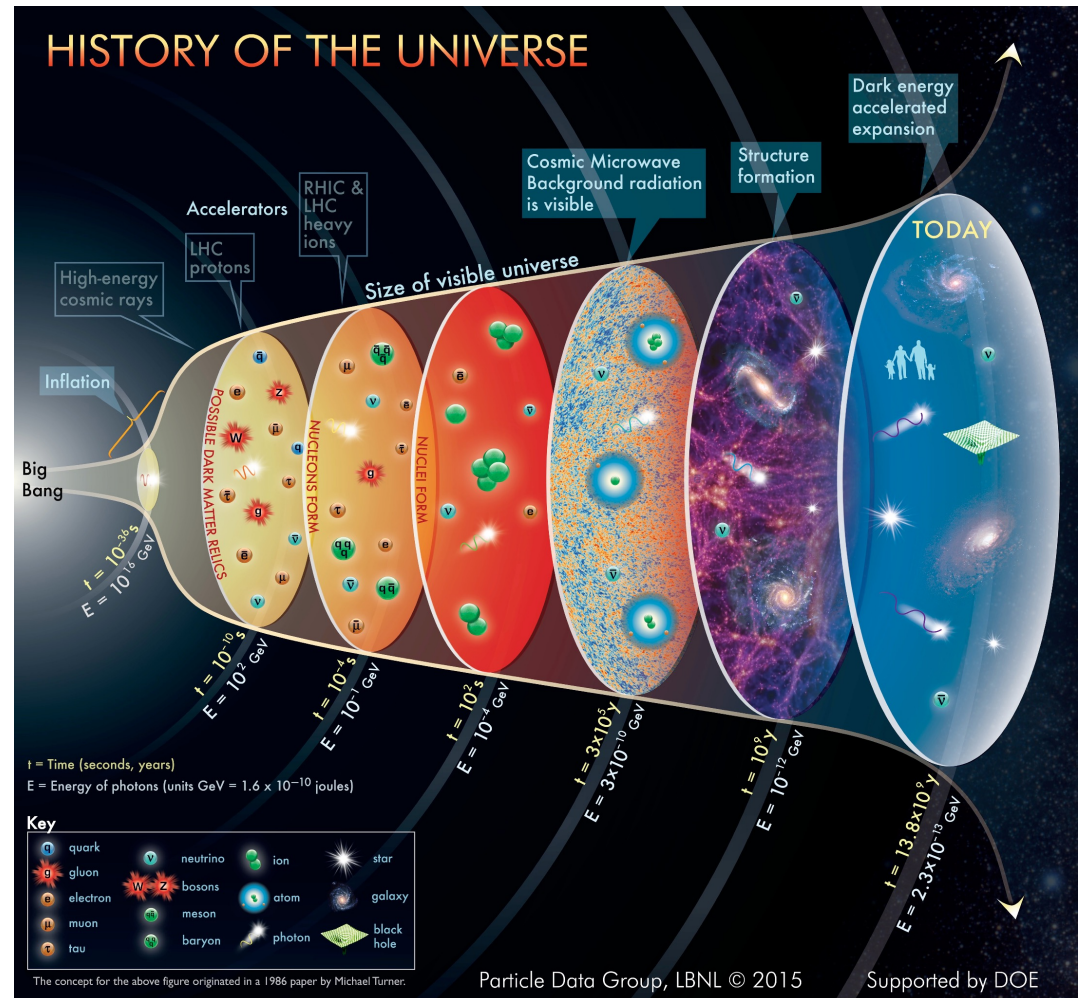
Why study QCD Matter?

3

The history of the Universe (as we currently understand it)

Early Universe :

- High temperature / energy density
- low density
- Too hot for “normal” matter to form
- How did the earliest forms of normal matter form?

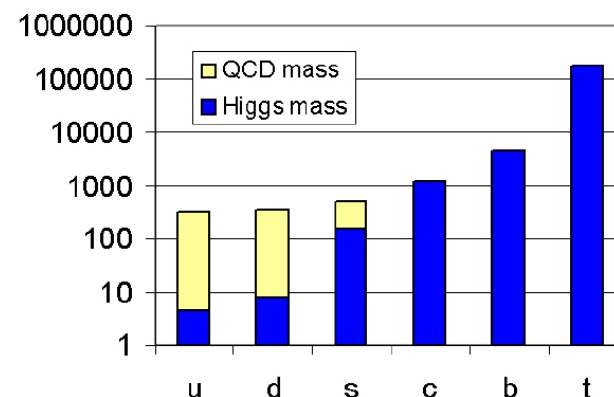
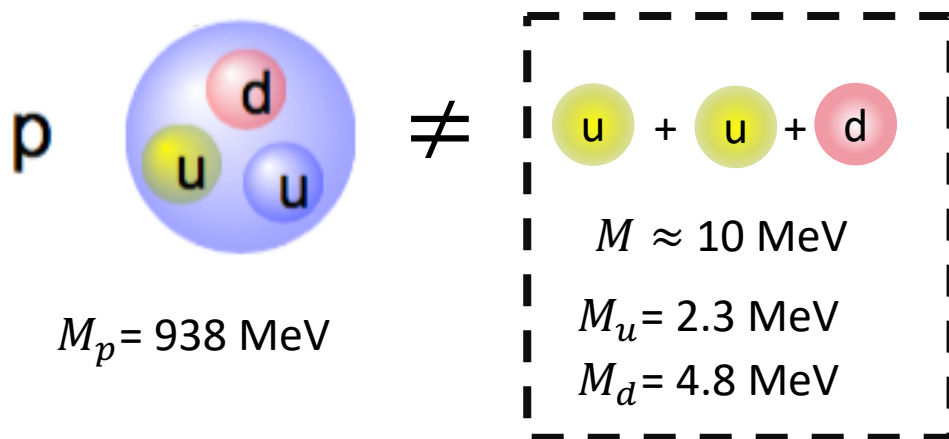


We study QCD matter to understand the physical evolution of the early Universe and emergence of structure

Chiral Symmetry and the Proton

4

- The proton (and neutron) make up almost all visible matter
- The proton is (much) more than the sum of its parts!



- Higgs Boson only explains the mass of fundamental particles!
- The remaining 99% of proton mass:
 - Kinetic energy
 - Gluon self interaction
 - Quark Interactions (chiral symmetry breaking)

We study QCD matter to understand the mass of normal matter

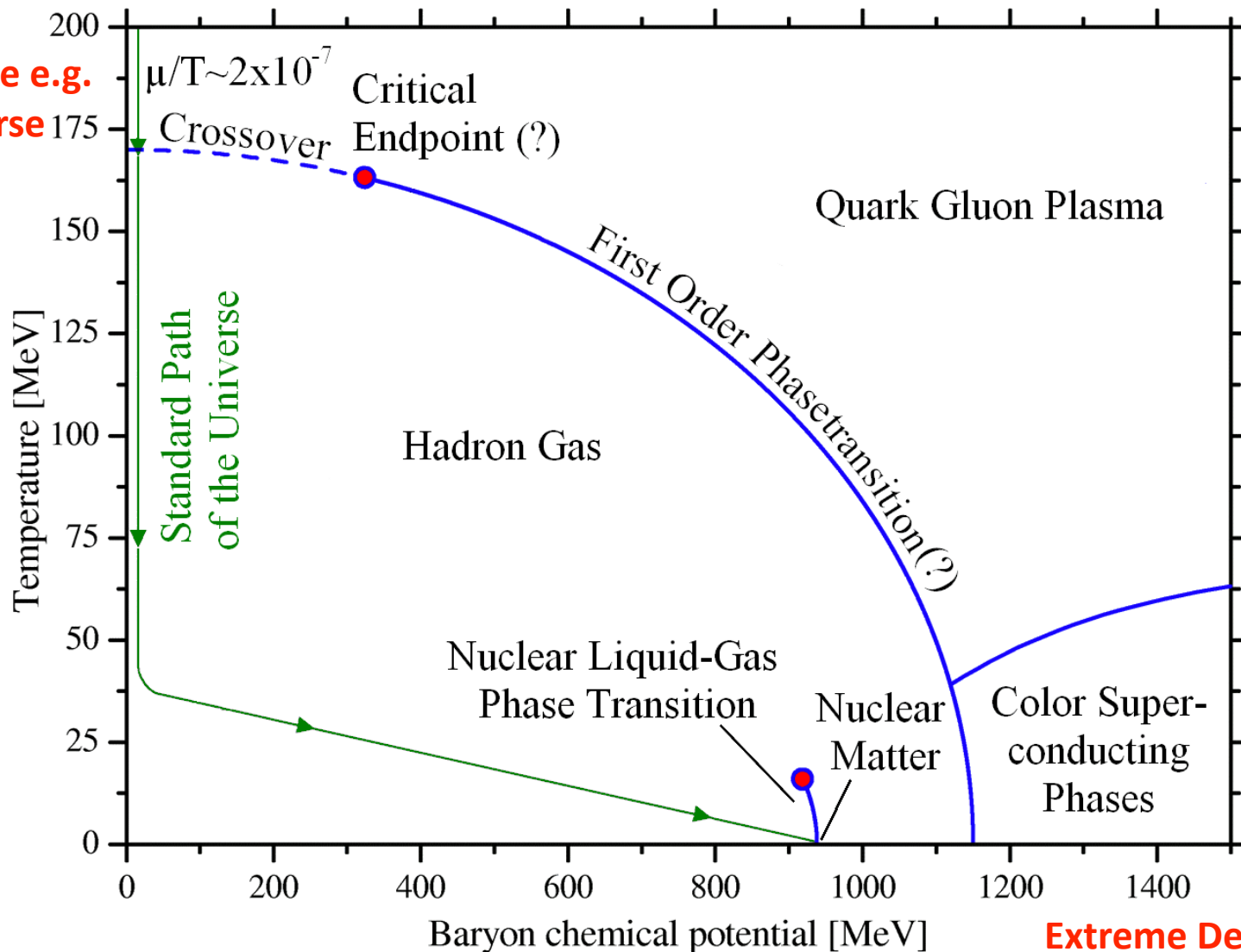
The QCD Phase Diagram

5

Extreme

Temperature e.g.

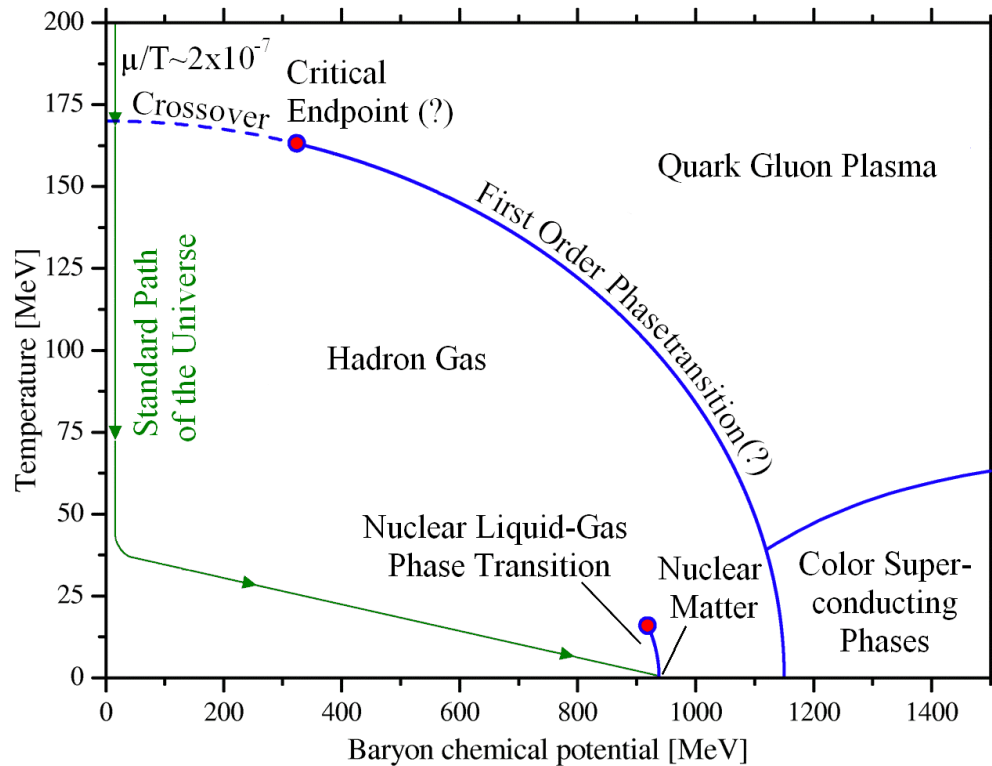
Early Universe



Extreme Density
e.g. Compact Stars

The QCD Phase Diagram

6



Big questions = questions about the QCD Phase Diagram

Is there a 1st Order Phase transition between the hadron gas and the Quark Gluon Plasma phases?

Is there a critical endpoint between phase transition and crossover region?

Does the transition to a chirally symmetric phase coincide with the deconfinement transition?

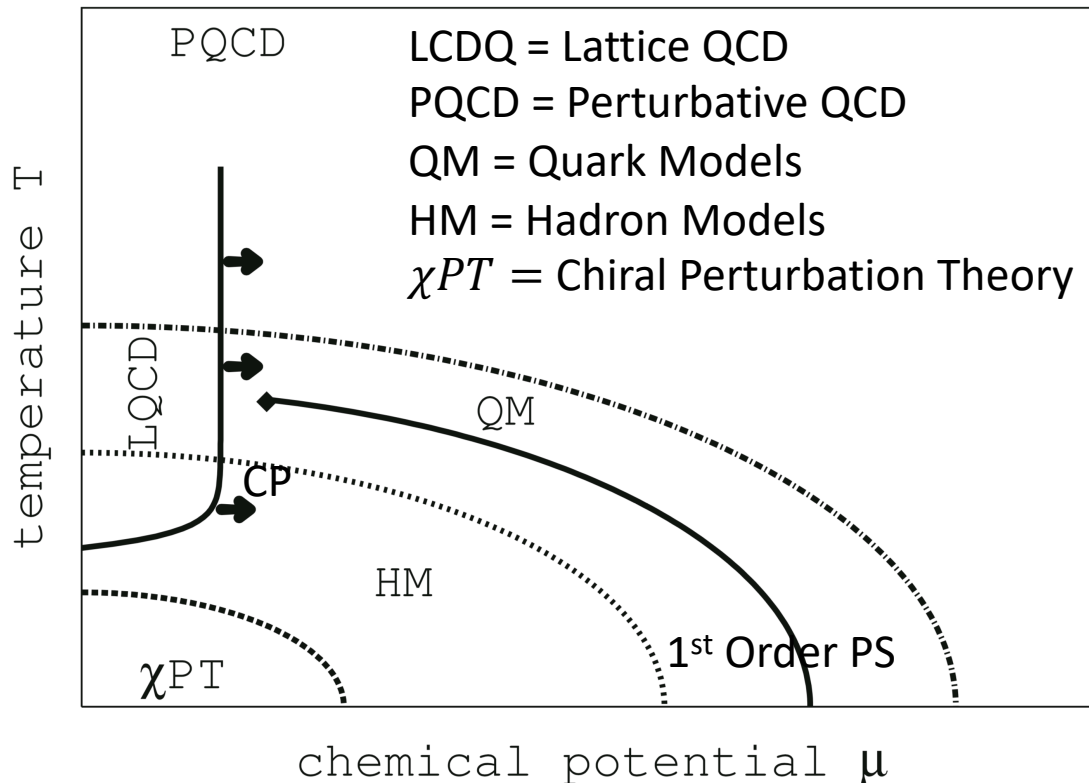
What is the equation of state for these types of matter?

$$\mathcal{L}_{\text{QCD}} = \bar{q}(i\gamma^\mu D_\mu - m)q - \frac{1}{4}F_{\mu\nu}^a F_a^{\mu\nu}$$

But we know QCD, why not just calculate it!

QCD Calculations

7

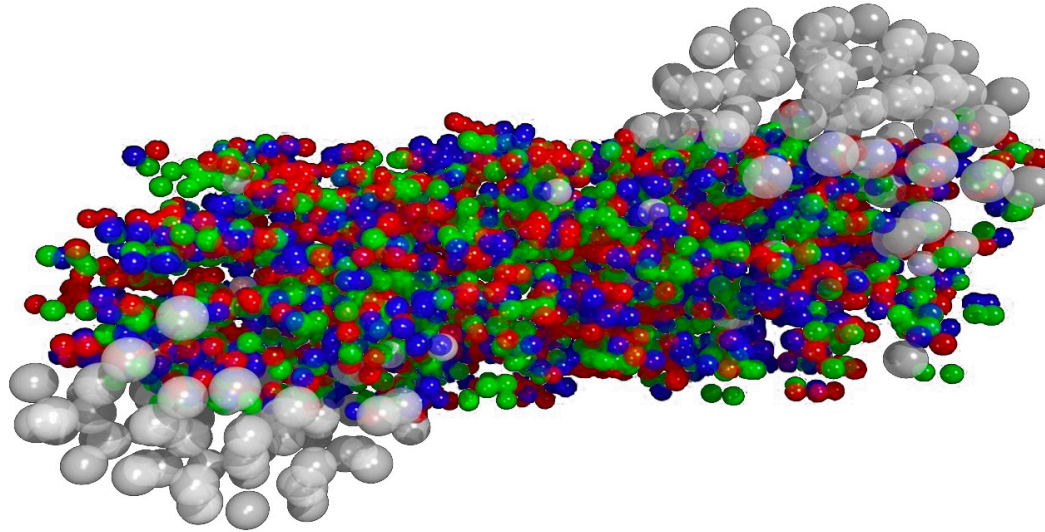


Theoretical Approaches:

- Lattice QCD : “first principles”
- Chiral Perturbation Theory : Systematic approach
- Hadron Gas Models & Quark Models : Can be verified in overlap region with LQCD and χ^{PT}
- Many other models – Sorry, I cannot possible list them all

- Only Lattice QCD is really QCD from first principles – though some caveats
- Sign problem prevents direct calculations at large baryon chemical potential
 - Progress is being made
 - Regions of overlap with LQCD can help verify other models

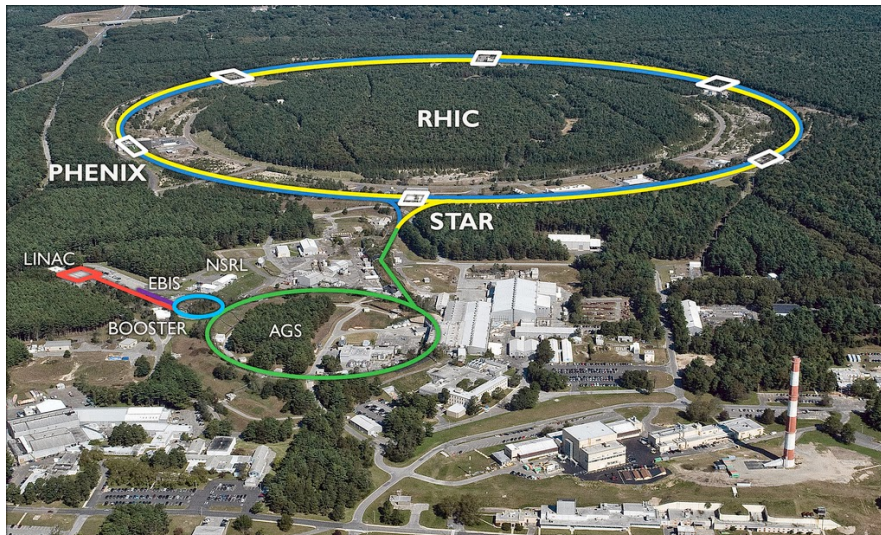
Experimental measurements are needed to determine key features of phase diagram

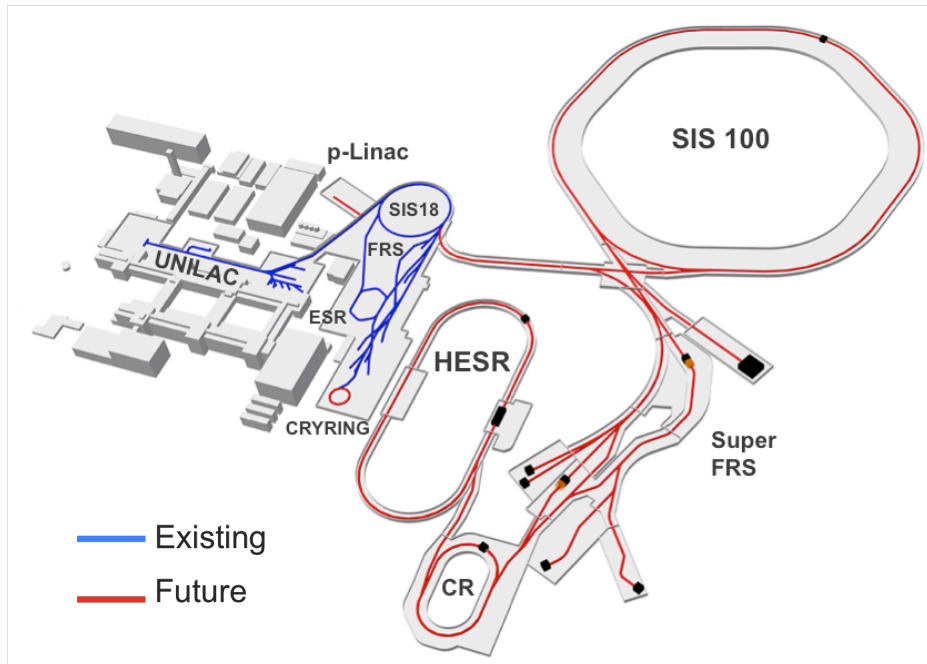


Studying QCD Matter with Heavy Ion Collisions

Facilities for Heavy Ion Collisions

9





Currently Running SIS18

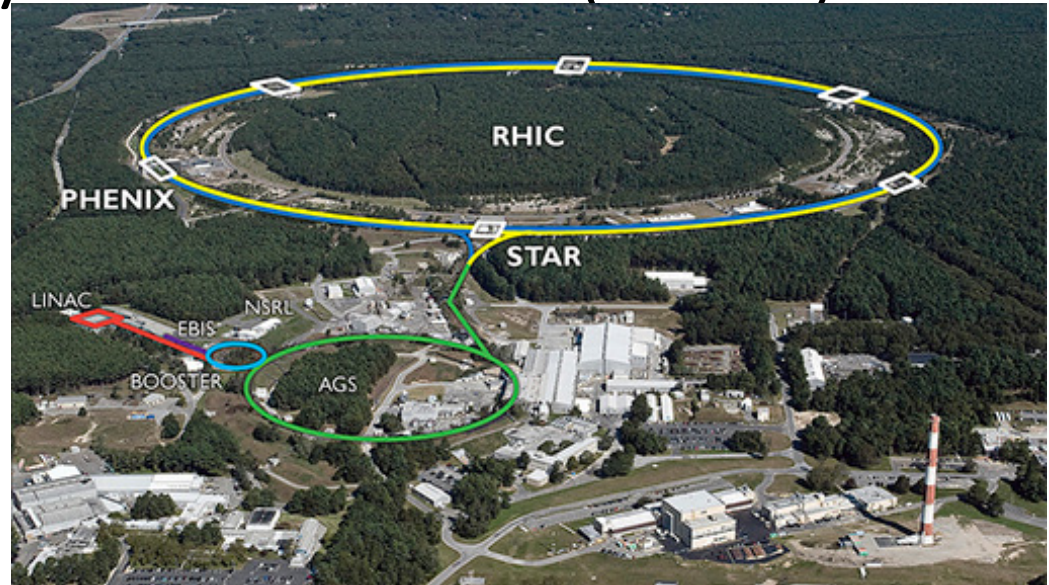
- Fixed target with ions, protons and π beams
- $\sqrt{s_{NN}} \approx 2.4 - 2.6$ GeV
- μ_B range : $\sim 880 - 670$
- Active experiments:
 - HADES (start 2012)
 - miniCBM (start 2018)

More about the Future
near the end of the talk

Relativistic Heavy Ion Collider (RHIC)

11

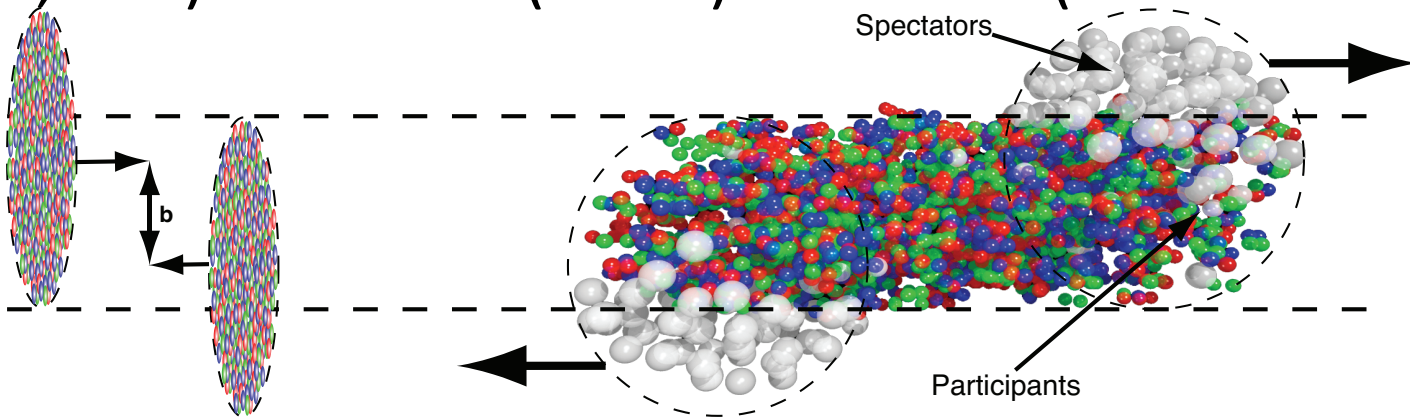
- Located at Brookhaven National Laboratory, Long Island, NY, USA
- Ions accelerated with \vec{E} , steered with \vec{B}
- Uses >1,700 dipole magnets in two 3.8 km rings



- Since beginning operation in 2000:
 - p+p, p+Au, p+Al, d+Au, $^3\text{He}+\text{Au}$, Cu+Cu, Cu+Au, Au+Au and U+U at $\sqrt{s_{NN}}=7.7 - 200$ GeV, p+p at $\sqrt{s} = 510$ GeV
- Four experiments:
 - STAR, PHENIX('16), BRAHMS('06), PHOBOS ('05) – STAR is only active experiment now

Studying the QCD Phase with Heavy Ion Collisions

Use SIS18 / RHIC / LHC to accelerate heavy ions (gold, lead, etc.) to $>0.90c$ (SIS18) or $>0.99c$ (RHIC & LHC)



$$\gamma^{RHIC} \approx 108$$

$$\gamma^{LHC} \approx 7000$$

At large γ - ions are compressed into thin “pancakes”

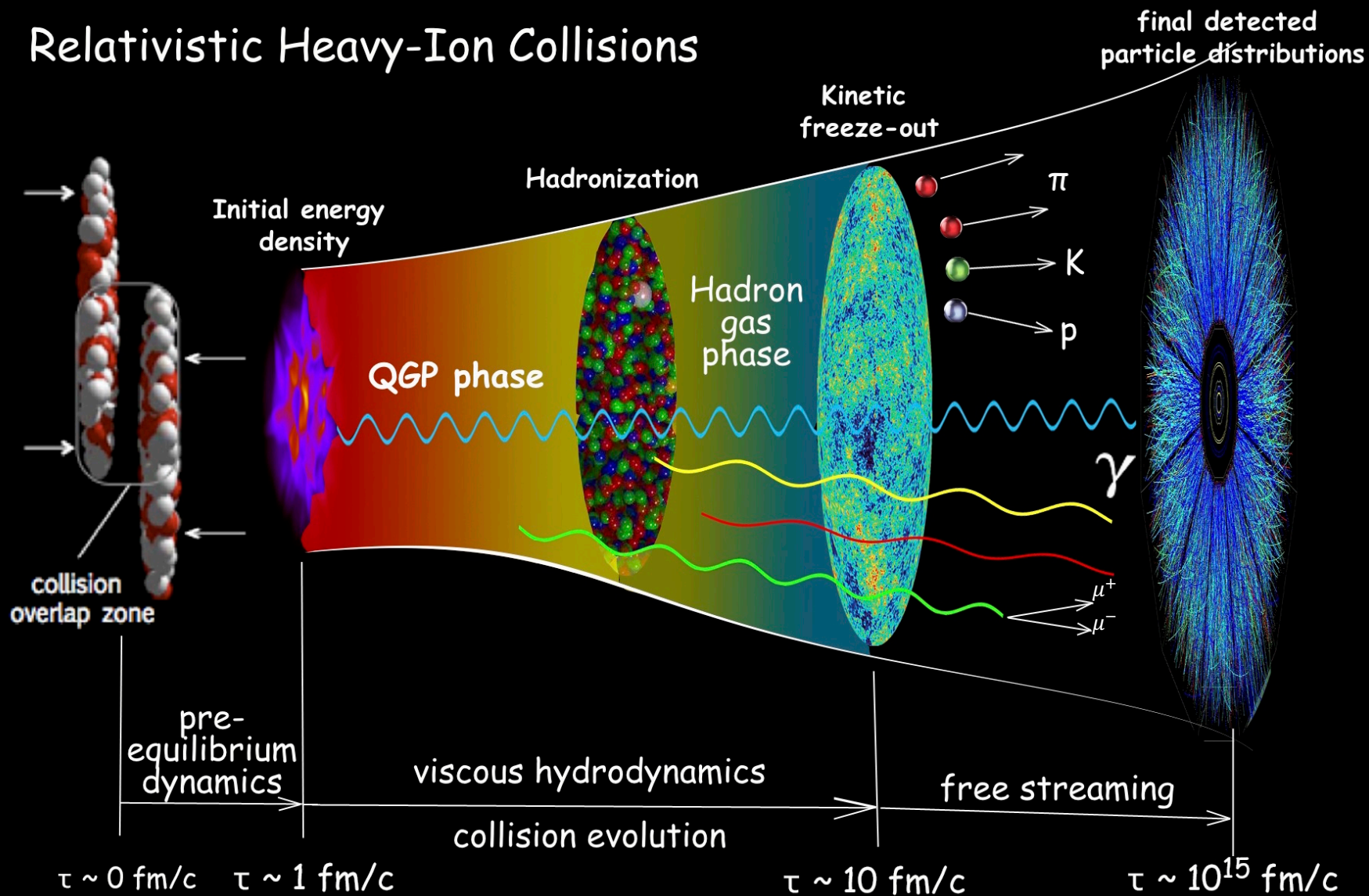
$$\gamma = \frac{1}{\sqrt{1-\beta^2}}, \quad L' = L/\gamma, \quad \beta = v/c$$

- Each HIC follows a specific trajectory through the QCD phase diagram
- Varying collision energy and impact parameter we can control:
 - Initial temperature, initial μ_B , system size, lifetime

Life of a Heavy Ion Collision

13

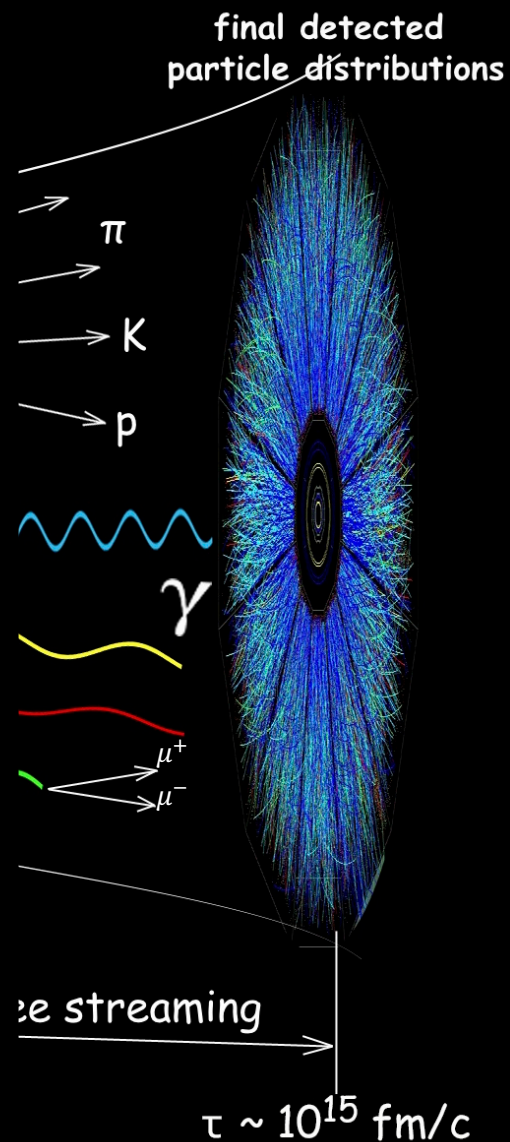
Relativistic Heavy-Ion Collisions



Life of a Heavy Ion Collision

14

Relativistic Heavy-Ion Collisions
As we actually observe them



The HADES Detector

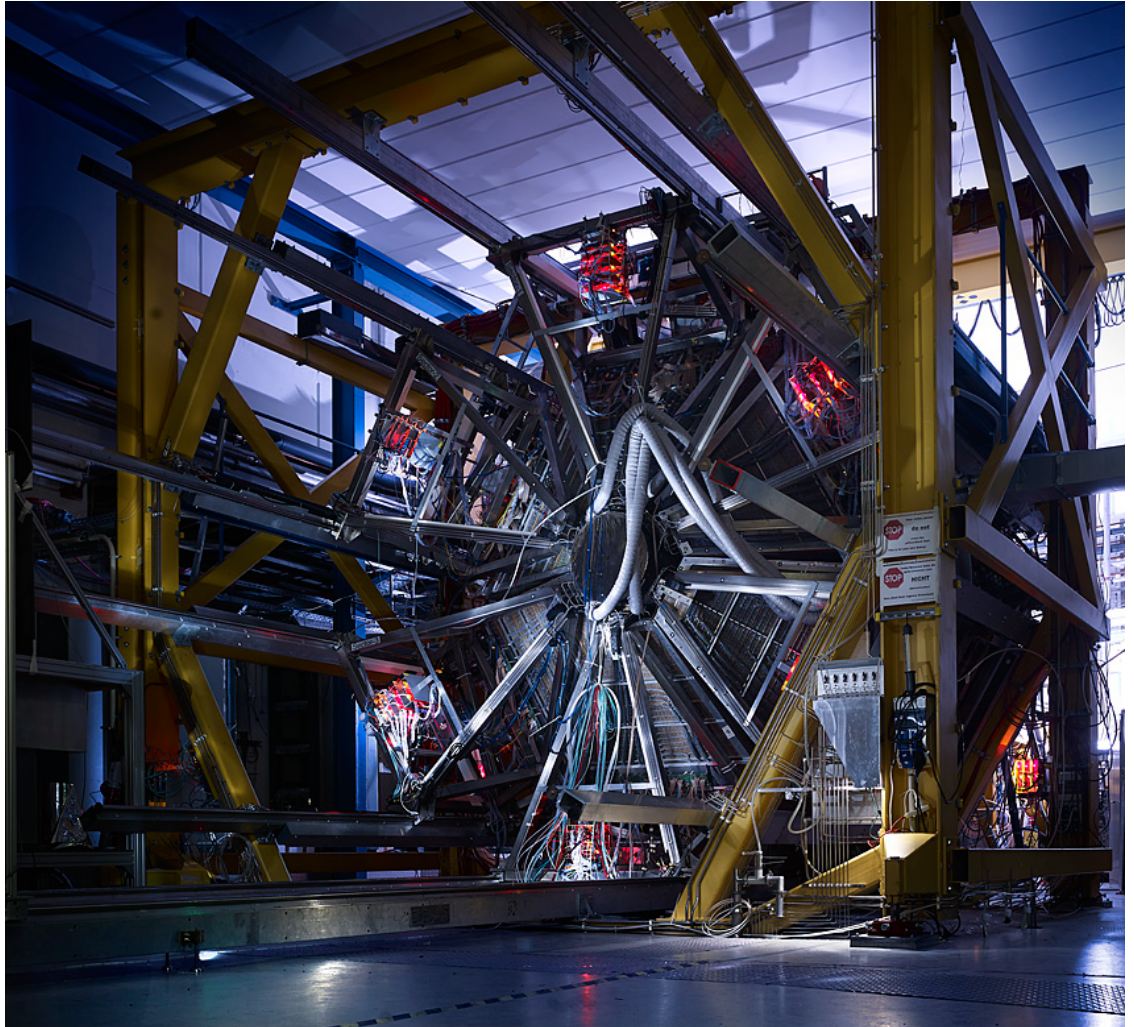
15

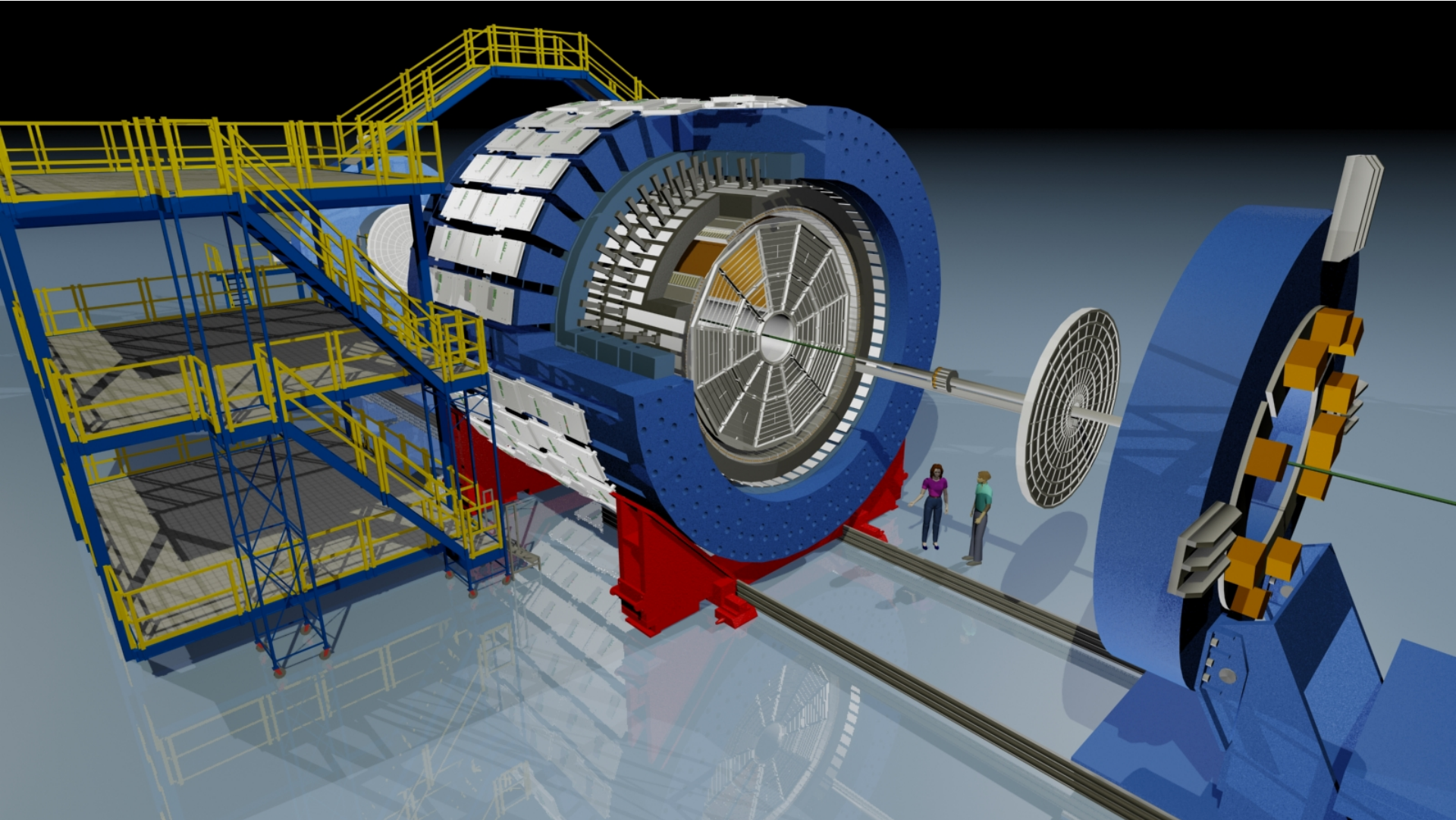
HADES = High Acceptance Dielectron Spectrometer

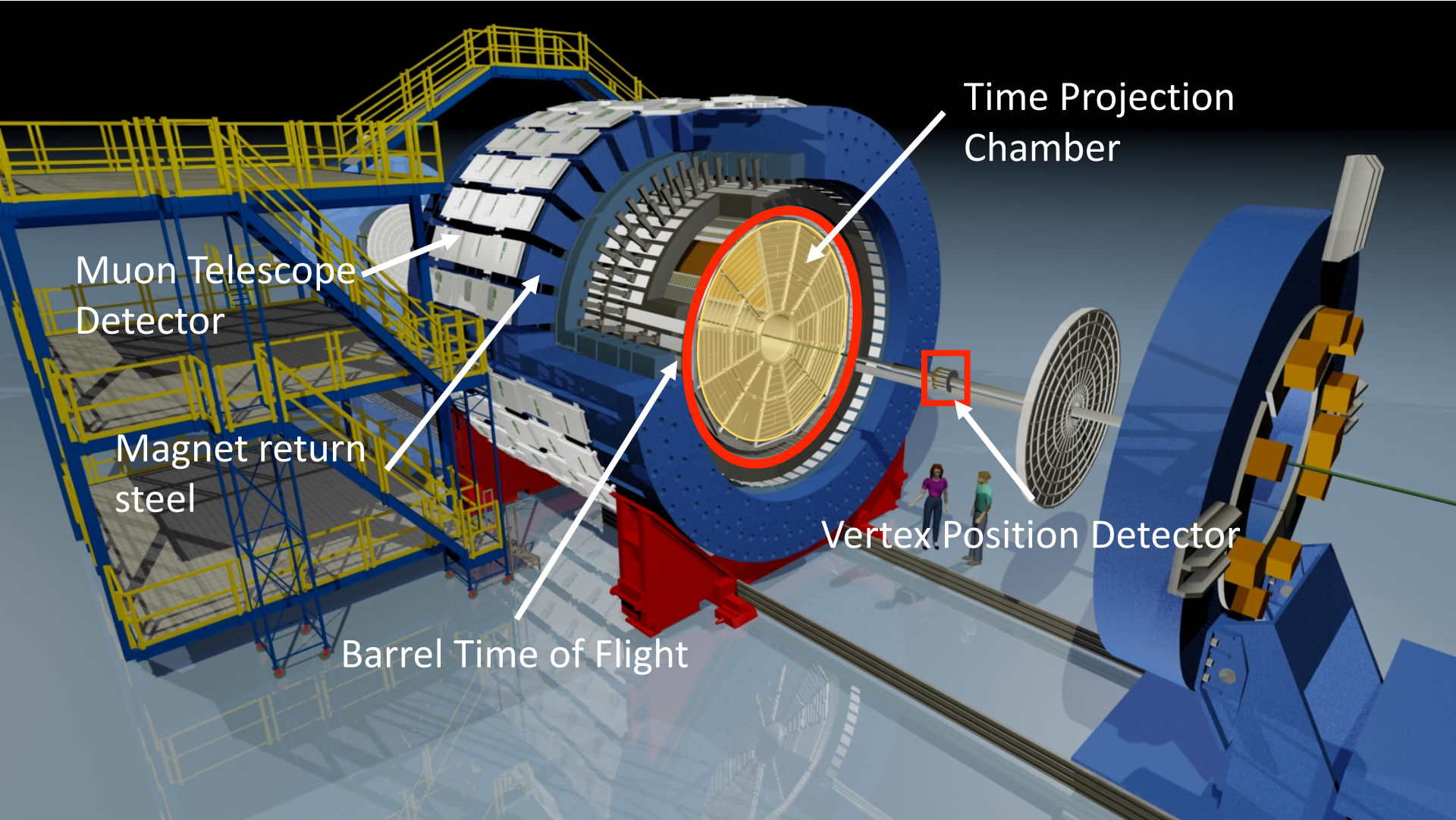
- Full azimuthal acceptance
- 18 to 85 degrees in Polar
- e^+e^- par acceptance : 35%
- Mass resolution : 2% (ρ , ω)

Physics Program:

- Excitation function for low-mass dielectron pairs
- Excitation function for (multi-)strange baryon/mesons

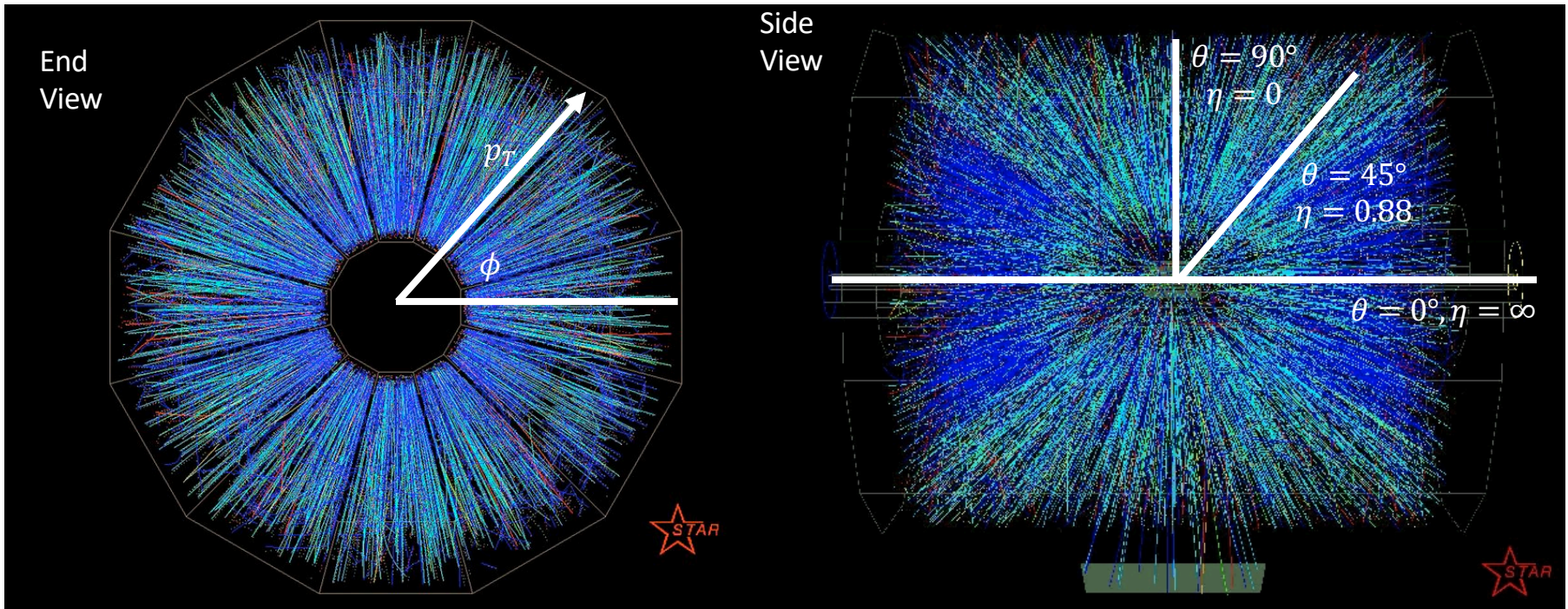






Heavy Ion Collision : Common Terms

18



Transverse Momentum:

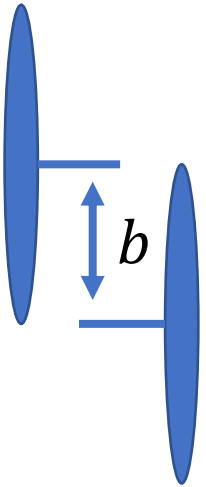
$$p_T = \sqrt{p_x^2 + p_y^2}$$

Rapidity: $y = \frac{1}{2} \ln \left(\frac{E + P_z}{E - P_z} \right)$
Additive under Lorentz boost

Pseudo-Rapidity (y for $m=0$) :

$$\eta = \frac{1}{2} \ln \left(\frac{P + P_z}{P - P_z} \right) = -\ln \left(\tan \frac{\theta}{2} \right)$$

Centrality in HICs

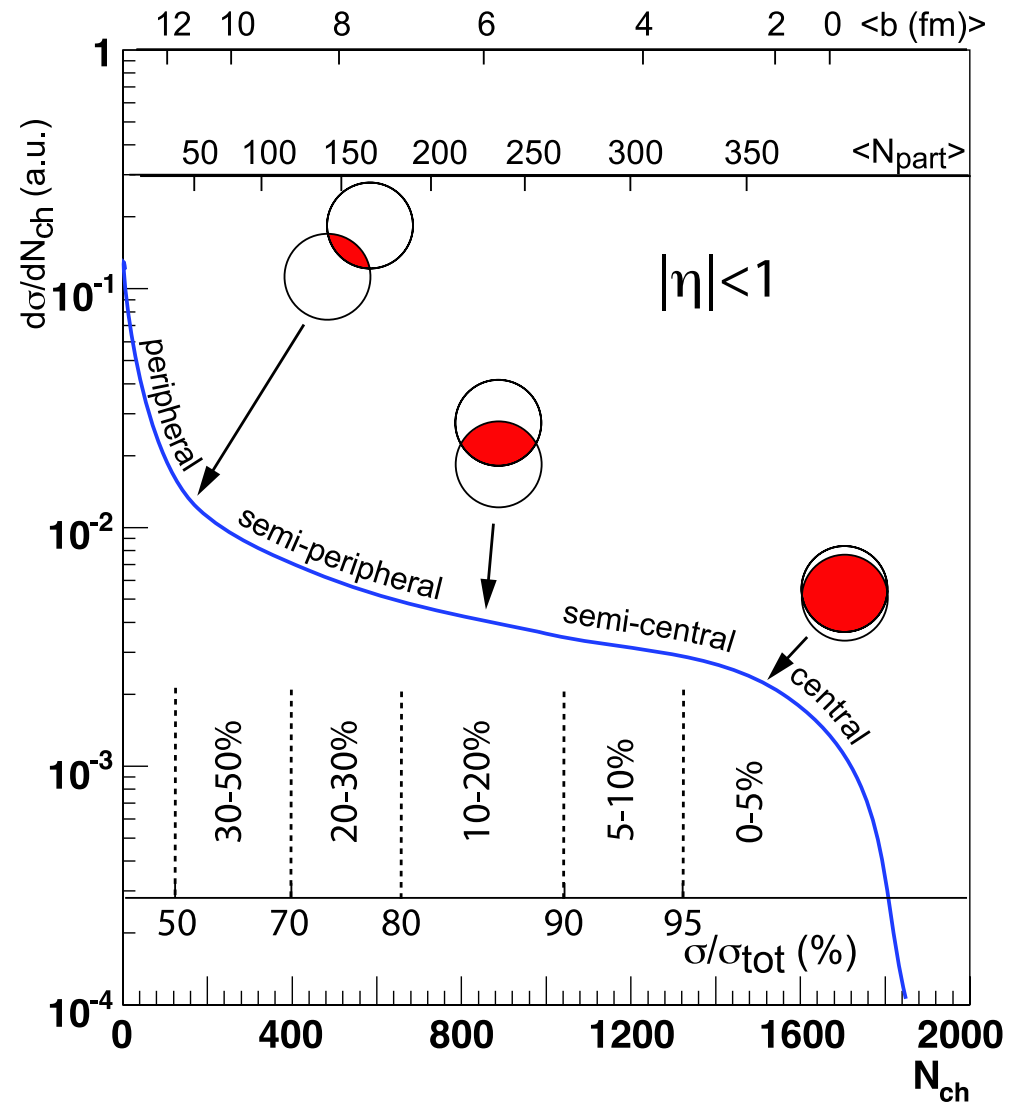


We cannot measure impact parameter (b) directly!

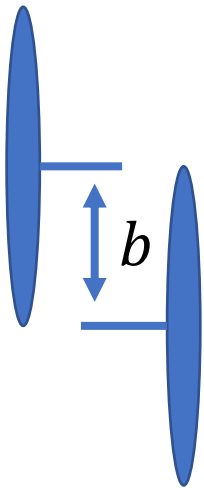
Instead we measure charged particle multiplicity (N_{ch}) and correlate to b

Same technique is used for:

- N_{part} - number of participating nucleons
- N_{coll} - number of binary collisions between nucleons



Centrality in HICs

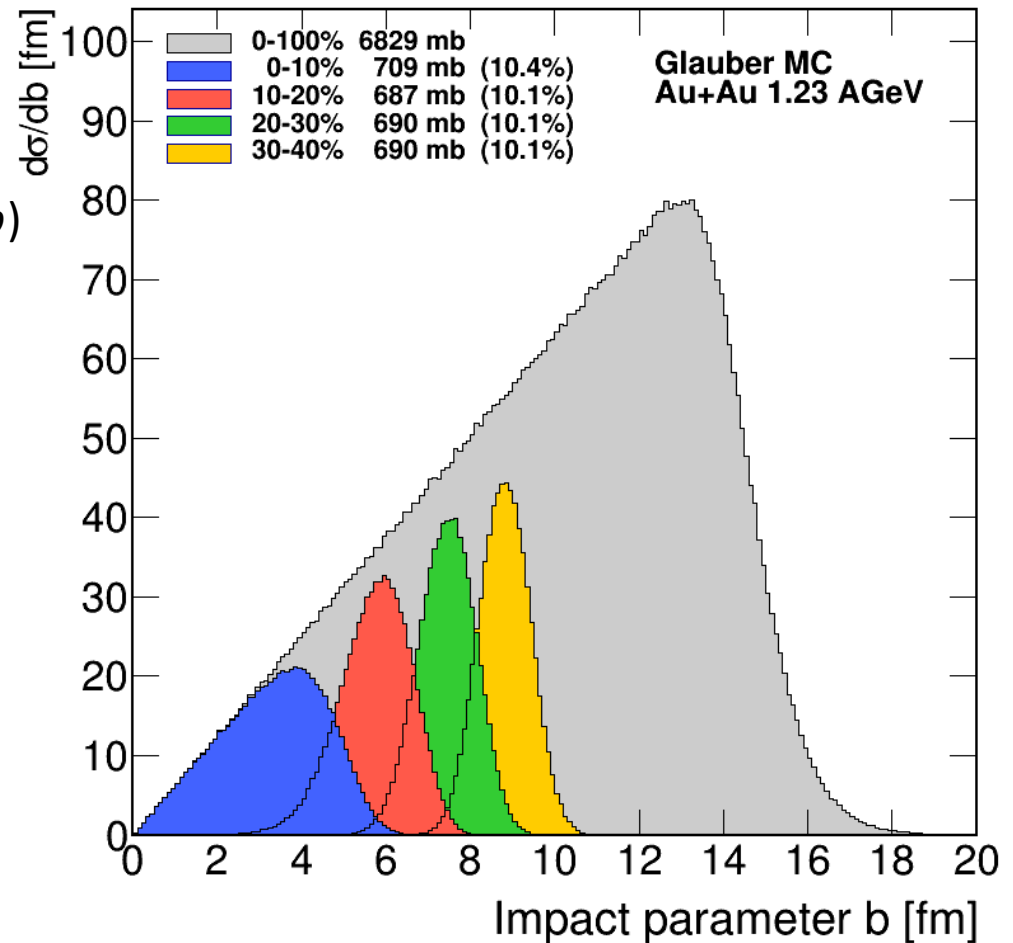


We cannot measure impact parameter (b) directly!

Instead we measure charged particle multiplicity (N_{ch}) and correlate to b

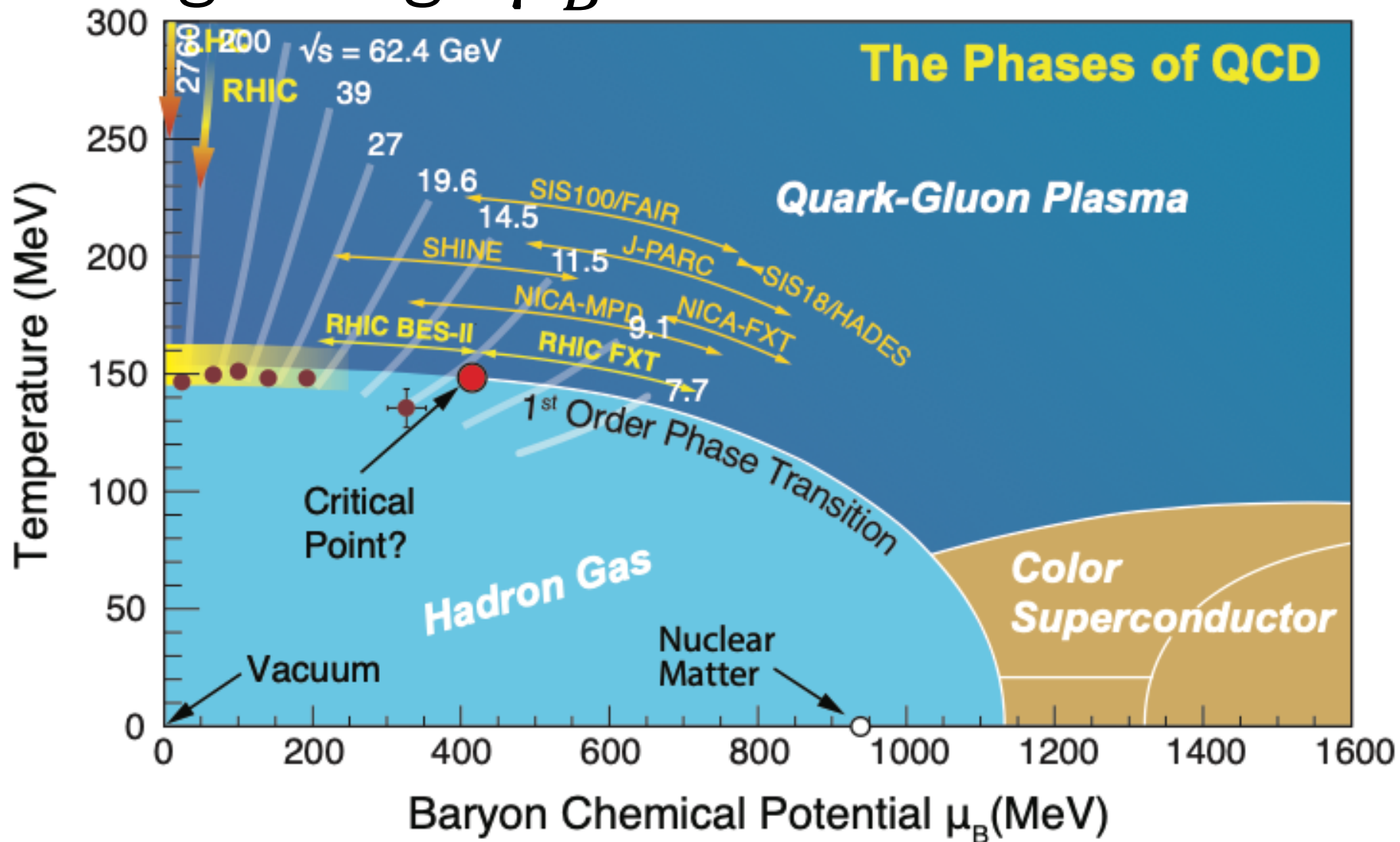
Same technique is used for:

- N_{part} - number of participating nucleons
- N_{coll} - number of binary collisions between nucleons



Going to high μ_B

21



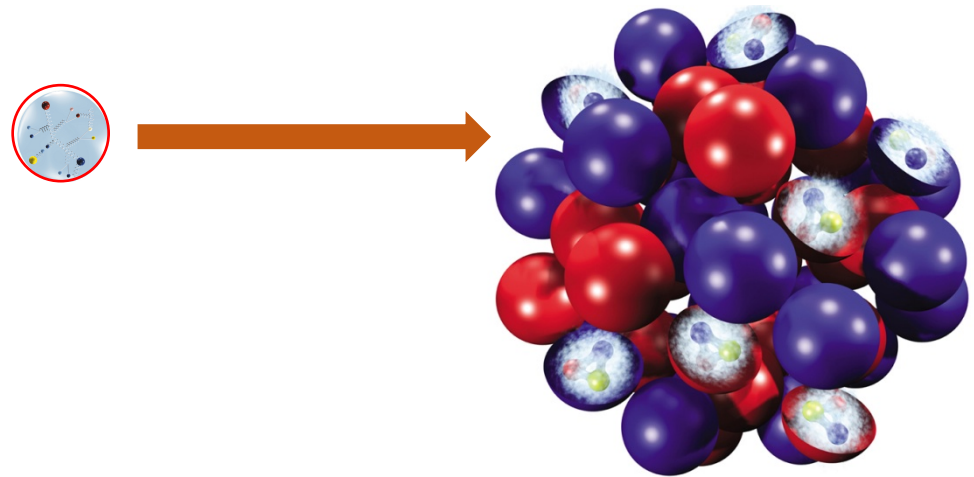
Effect of Collision Energy on μ_B

22

Beam energy effects deBroglie wavelength of partons.

Wavelength determines the scale observed by colliding parton :

- Whole nucleus
- Individual nucleons
- Individual partons



Low collision energy

High collision energy

Nucleons are opaque
Valence quarks are stopped
Excess quarks = high μ_B
Baryon dominated

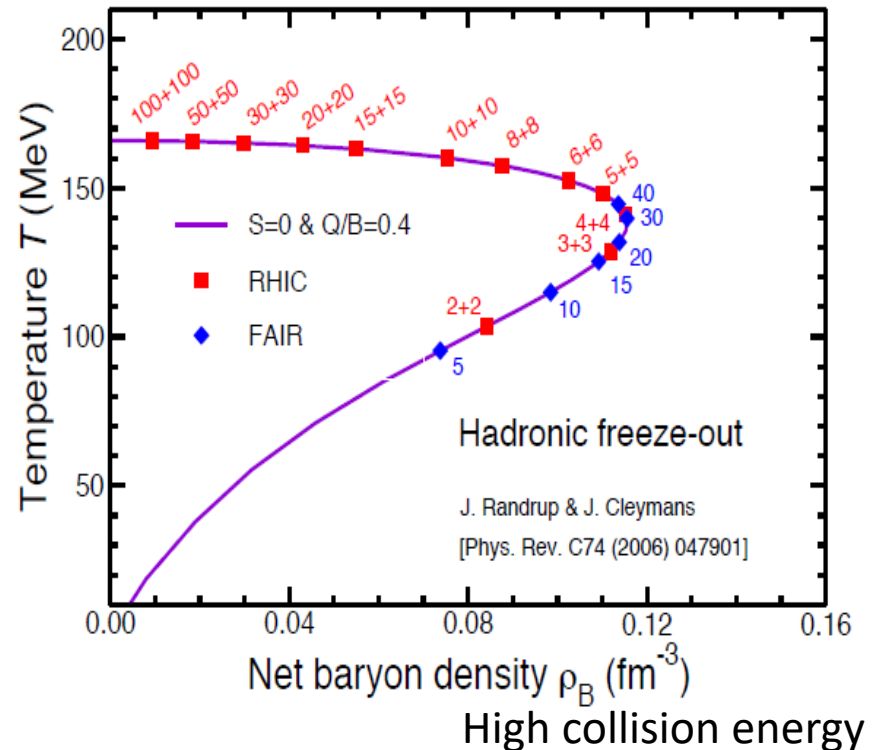
Nucleons are transparent
Valence quarks pass through
Equal quark / anti-quarks = low μ_B
Meson dominated

Effect of Collision Energy on μ_B

Beam energy effects deBroglie wavelength of partons.

Wavelength determines the scale observed by colliding parton :

- Whole nucleus
- Individual nucleons
- Individual partons



Low collision energy

High collision energy

Nucleons are opaque
Valence quarks are stopped
Excess quarks = high μ_B
Baryon dominated

Nucleons are transparent
Valence quarks pass through
Equal quark / anti-quarks = low μ_B
Meson dominated

Where in the phase diagram?

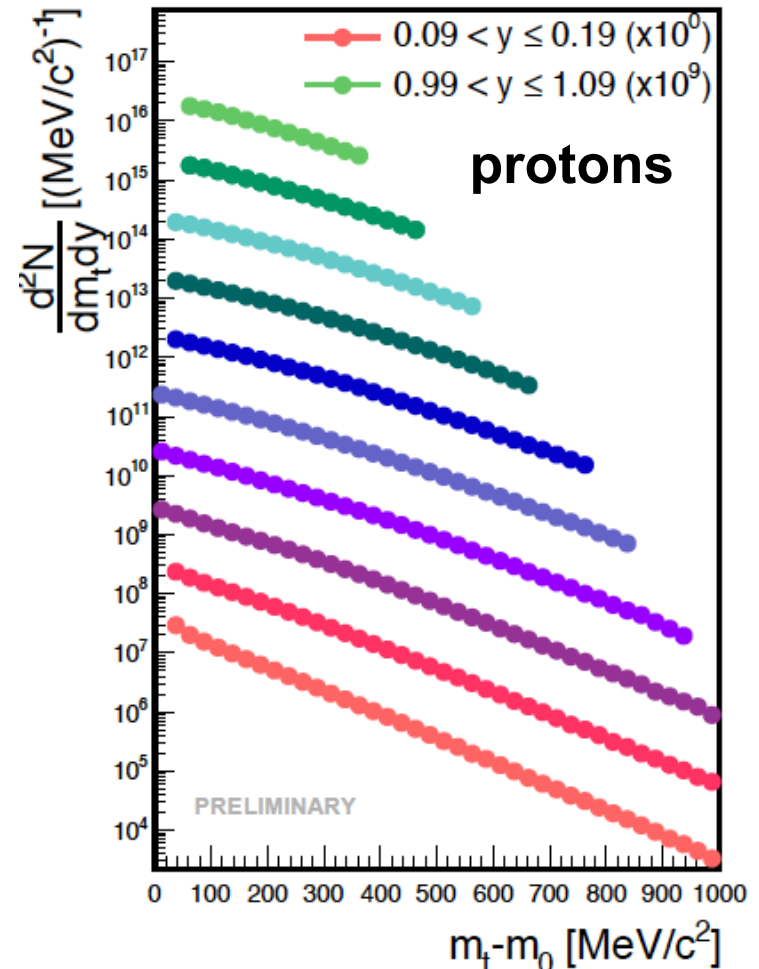
Experimentally how do we measure where we are in the phase diagram for a given:

- Collision energy
- Centrality
- Rapidity range

Using a model which assumes statistical equilibrium (Statistical Hadronization Models):
Relate the measured particle yields (π , K, p, Λ , Ξ , ϕ , Ω) and their ratios, to the location in the phase diagram

$$\ln Z^{GC}(T, V, \{\mu_i\}) = \sum_{\text{species } i} \frac{g_i V}{(2\pi)^3} \int d^3p \ln \left(1 \pm e^{-\beta(E_i - \mu_i)} \right)^{\pm 1}$$

$$\mu_i = B_i \mu_B + S_i \mu_S + Q_i \mu_Q$$



Step 1 : Measure Particle Spectra

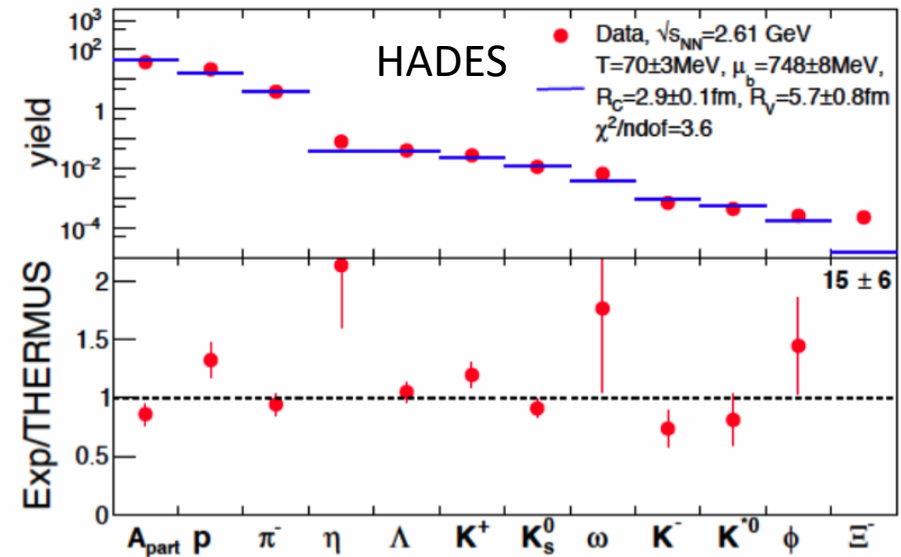
Where in the phase diagram?

25

Experimentally how do we measure where we are in the phase diagram for a given:

- Collision energy
- Centrality
- Rapidity range

Using a model which assumes statistical equilibrium (Statistical Hadronization Models):
Relate the measured particle yields (π , K , p , Λ , Ξ , ϕ , Ω) and their ratios, to the location in the phase diagram



$$\ln Z^{GC}(T, V, \{\mu_i\}) = \sum_{\text{species } i} \frac{g_i V}{(2\pi)^3} \int d^3p \ln \left(1 \pm e^{-\beta(E_i - \mu_i)} \right)^{\pm 1}$$

$$\mu_i = B_i \mu_B + S_i \mu_S + Q_i \mu_Q$$

Step 2 : Fit with SHM

Where in the phase diagram?

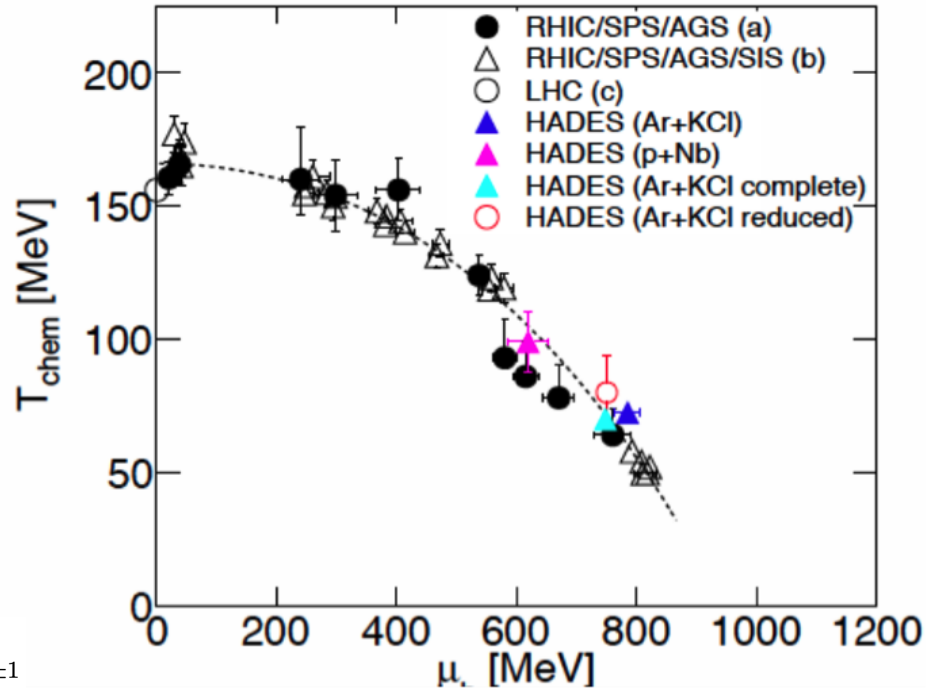
Experimentally how do we measure where we are in the phase diagram for a given:

- Collision energy
- Centrality
- Rapidity range

Using a model which assumes statistical equilibrium (Statistical Hadronization Models):
Relate the measured particle yields (π , K, p, Λ , Ξ , ϕ , Ω) and their ratios, to the location in the phase diagram

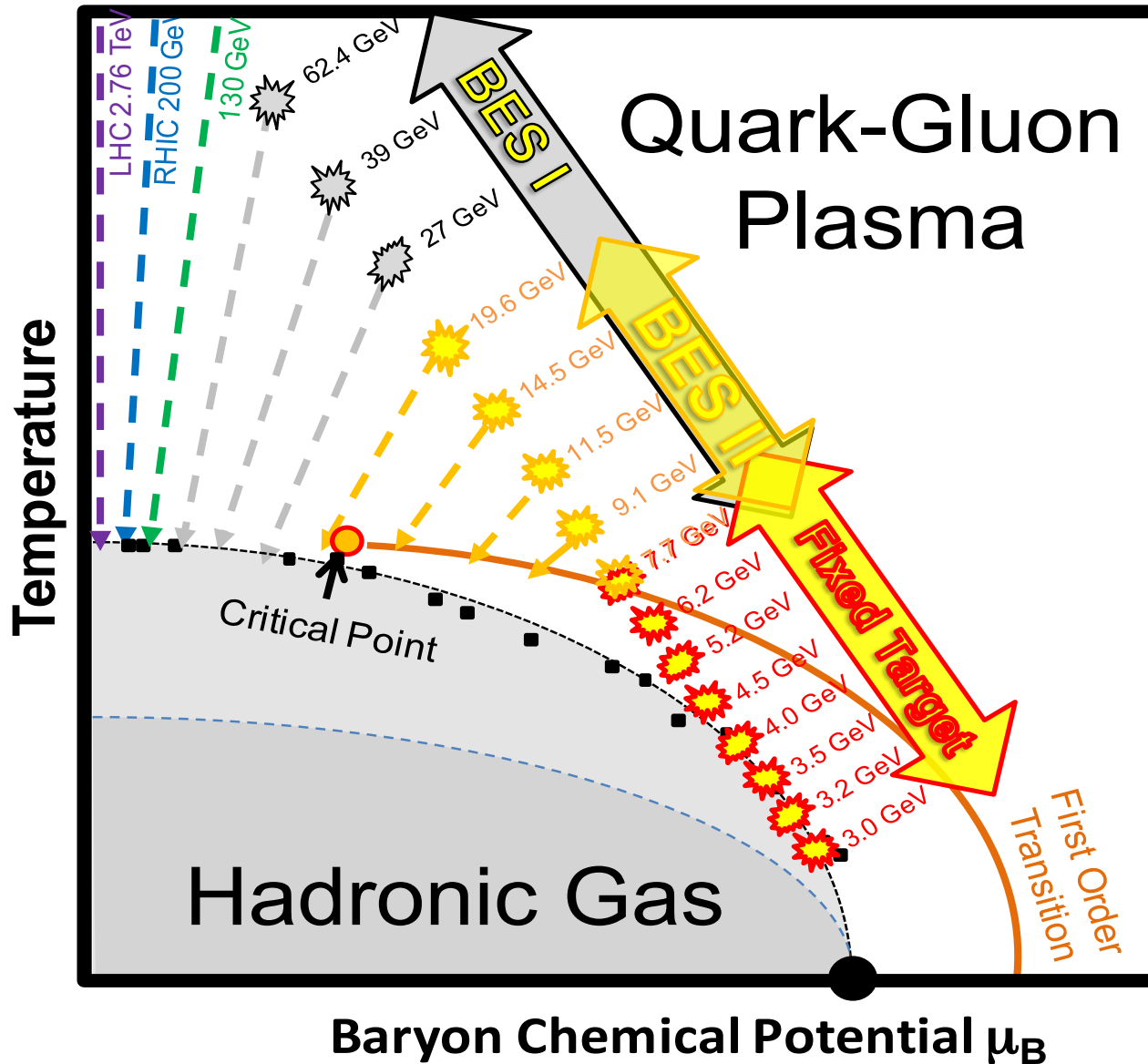
$$\ln Z^{GC}(T, V, \{\mu_i\}) = \sum_{\text{species } i} \frac{g_i V}{(2\pi)^3} \int d^3p \ln \left(1 \pm e^{-\beta(E_i - \mu_i)} \right)^{\pm 1}$$

$$\mu_i = B_i \mu_B + S_i \mu_S + Q_i \mu_Q$$



Step 3 : Get T_{chem} (chemical freeze-out temp) and μ_B

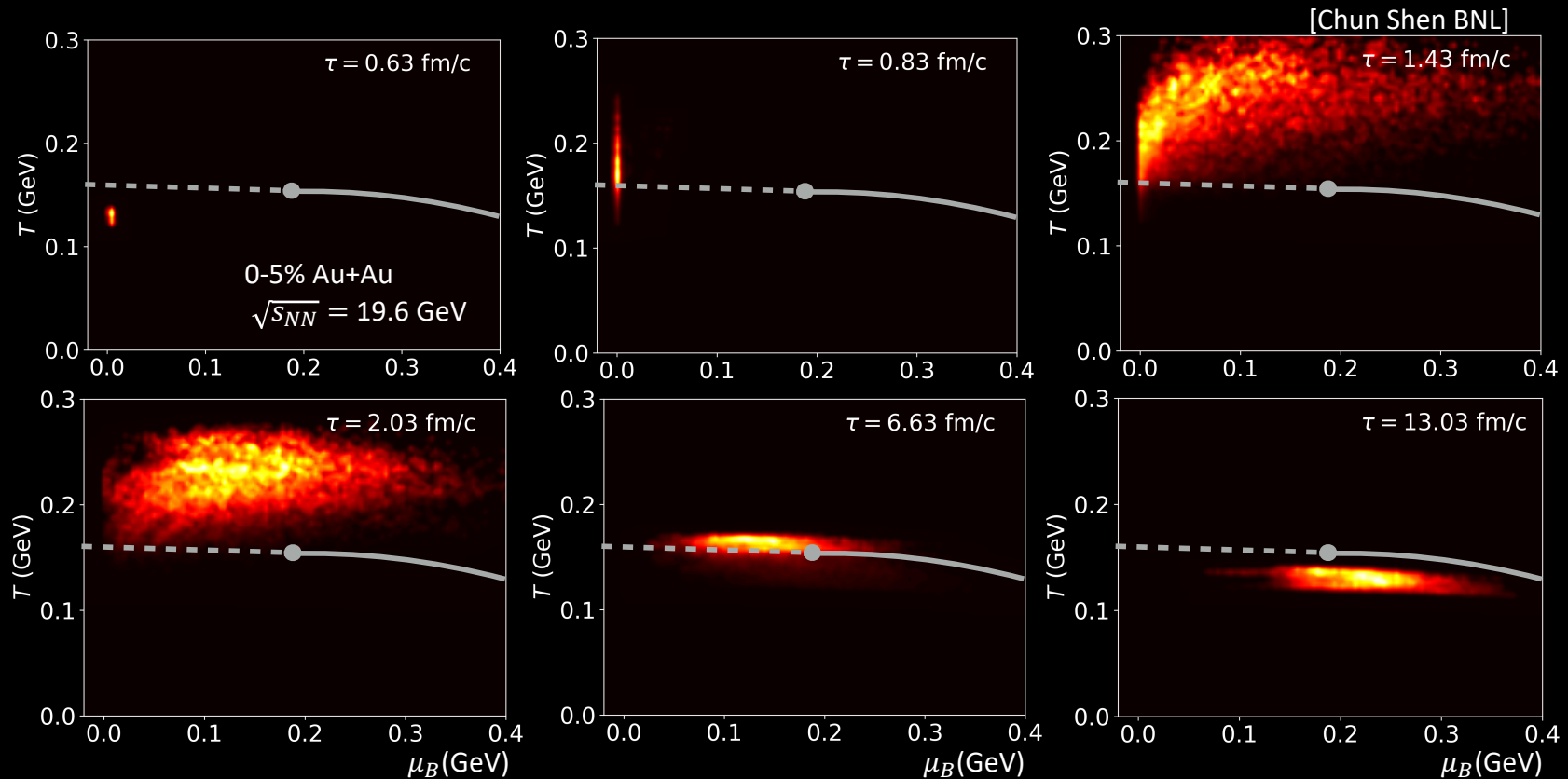
RHIC Beam Energy Scan



	Energy (GeV)	Chemical Potential μ_B	Pred. Temp. (MeV)
LHC	2760.0	2	166.0
RHIC	200.0	24	165.9
RHIC	130.0	36	165.8
RHIC	62.4	73	165.3
RHIC	39.0	112	164.2
RHIC	27.0	156	162.6
RHIC	19.6	206	160.0
SPS	17.3	229	158.6
RHIC	14.5	262	156.2
SPS	12.4	299	153.1
RHIC	11.5	316	151.6
SPS	8.8	383	144.4
RHIC	7.7	422	139.6
SPS	7.7	422	139.6
SPS	6.4	476	131.7
AGS	4.7	573	114.6
AGS	4.3	602	108.8
AGS	3.8	638	100.6
AGS	3.3	686	88.9
AGS	2.7	752	70.4
SIS	2.3	799	55.8

Trajectory through Phase Diagram of a HIC

28

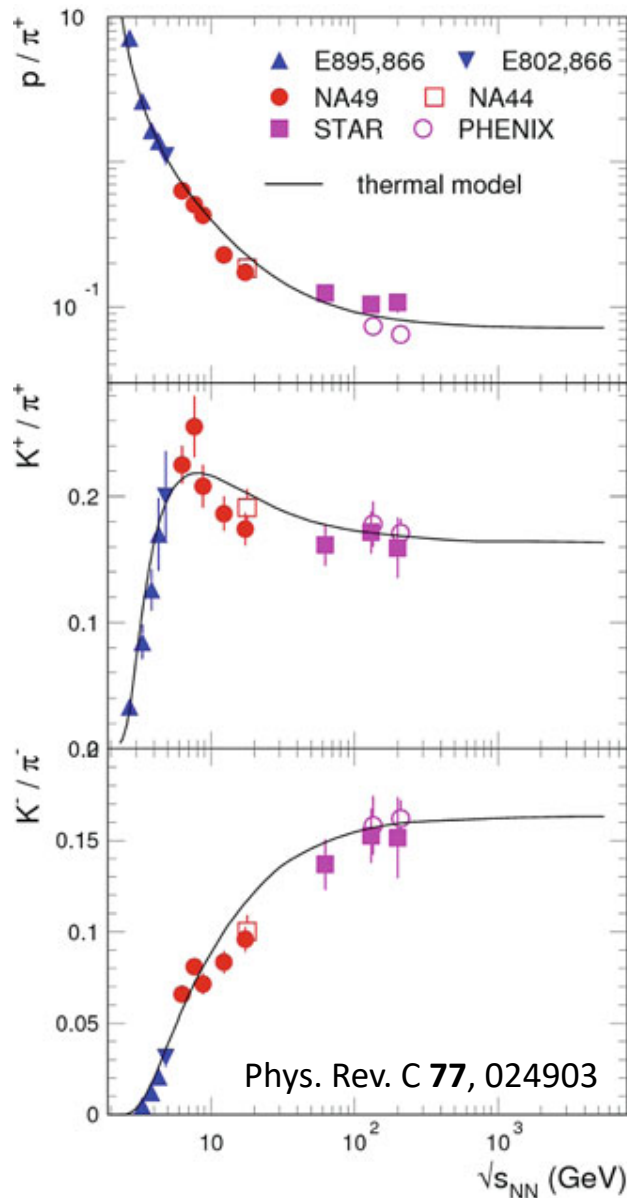


Keep in mind:

Path through the phase diagram is not as trivial and localized as we often depict

Particle Ratio

29



Steep decrease in p/π^+ : decreasing μ_B and transition from baryon dominated to meson dominated

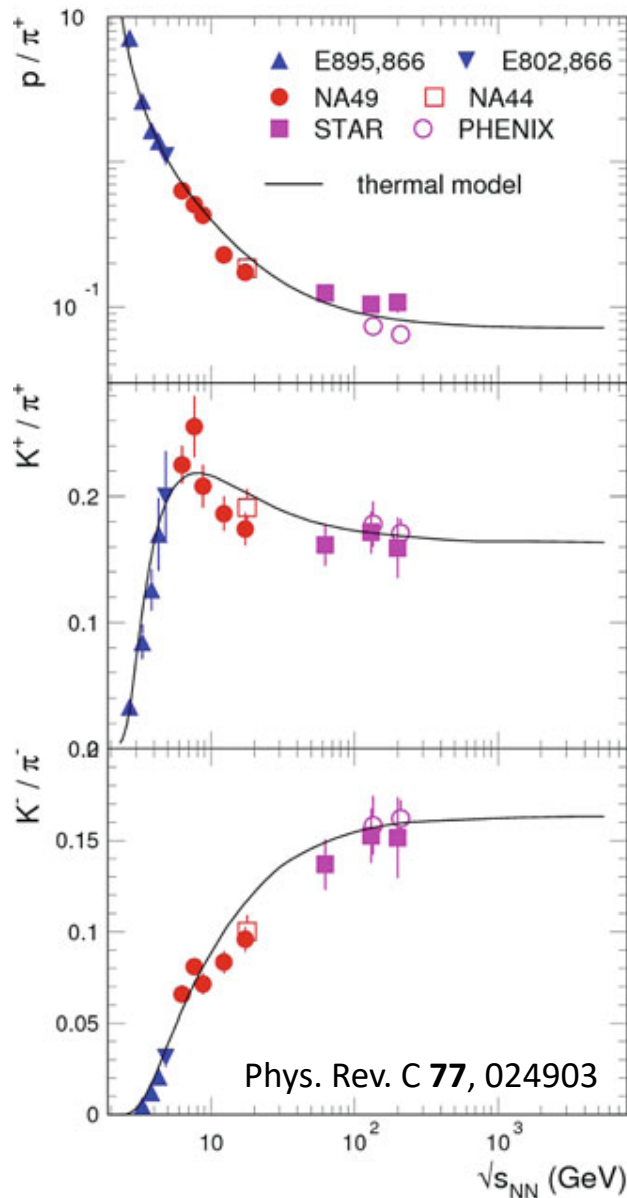
Peak in K^+/π^+ : competing effects. 1) decreasing light quark fraction, 2) increasing $\bar{q}q$ production ($u \bar{s}$) in kaon

➤ Interpreted as indication of the onset of deconfinement

Peak in K^-/π^- : Both are produced particles. Steep increase corresponds to the rapid rise in temperature

Particle Ratio

30

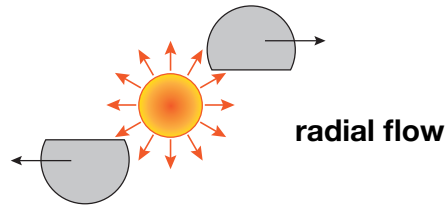


Something interesting is happening
for $\sqrt{s_{NN}} < \sim 20$ GeV

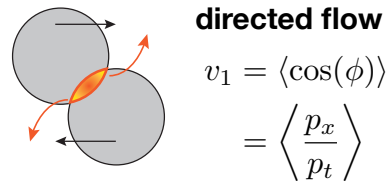
Now, Lets look for the 1st order
Phase Transitions and Critical point

Measuring the Flow

31



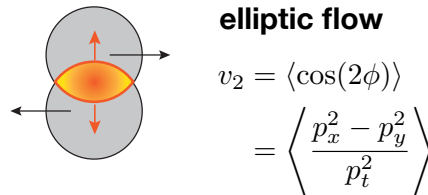
Look at the collective behavior of the medium



$$v_1 = \langle \cos(\phi) \rangle$$

$$= \left\langle \frac{p_x}{p_t} \right\rangle$$

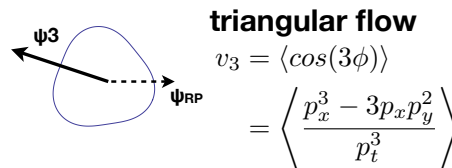
$$\frac{dN}{d\varphi} \propto 1 + \sum_{n=1}^{\infty} 2v_n \cos [n(\varphi - \Psi_n)]$$



$$v_2 = \langle \cos(2\phi) \rangle$$

$$= \left\langle \frac{p_x^2 - p_y^2}{p_t^2} \right\rangle$$

Where ψ is the azimuthal angle and Ψ_n is the n^{th} order event plane



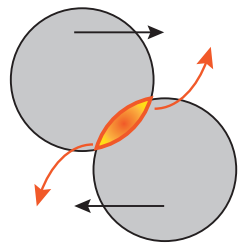
$$v_3 = \langle \cos(3\phi) \rangle$$

$$= \left\langle \frac{p_x^3 - 3p_x p_y^2}{p_t^3} \right\rangle$$

Behruz Kardan

Proton v1

32



directed flow

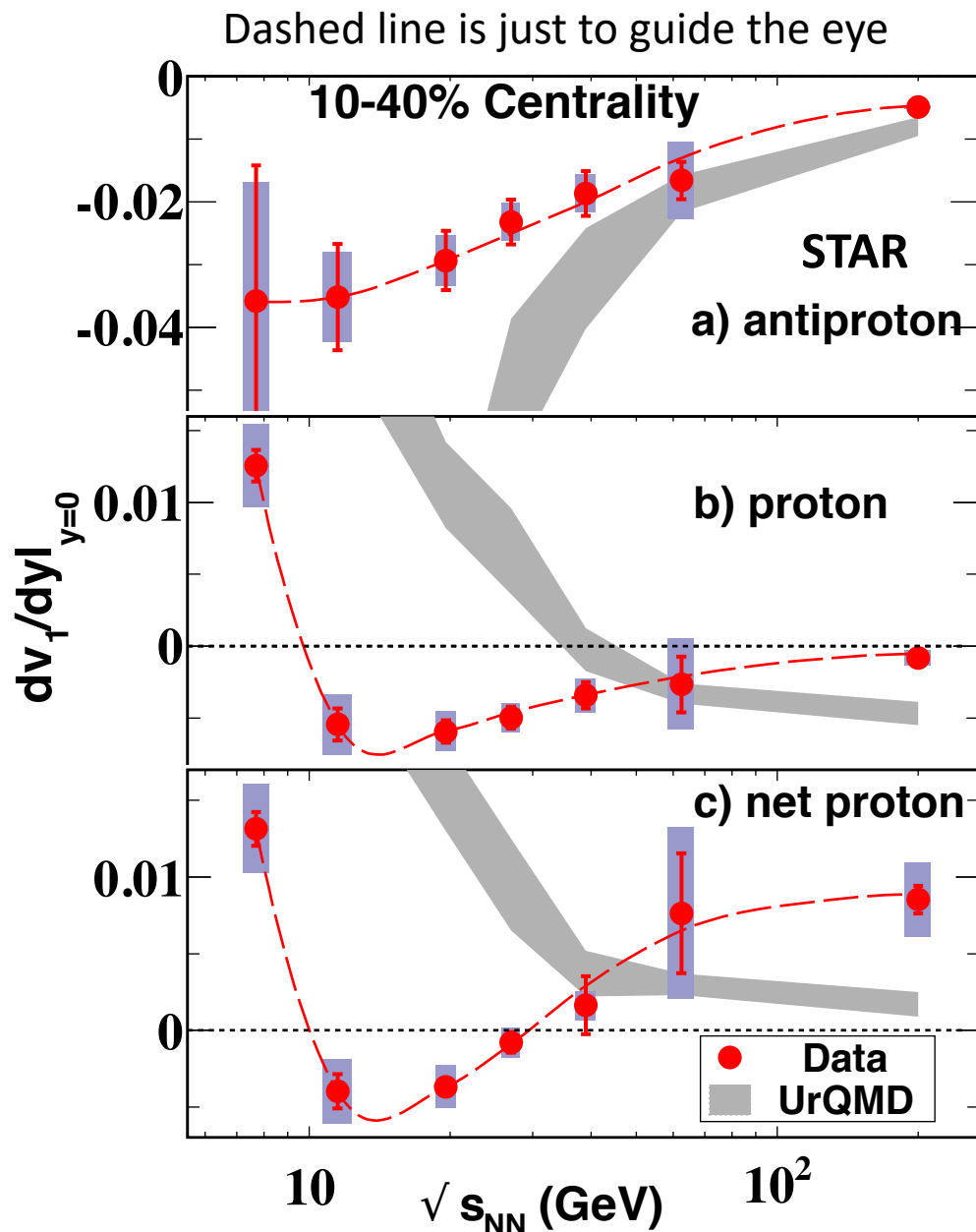
$$v_1 = \langle \cos(\phi) \rangle$$

$$= \left\langle \frac{p_x}{p_t} \right\rangle$$

v_1 is Sensitive to the compressibility

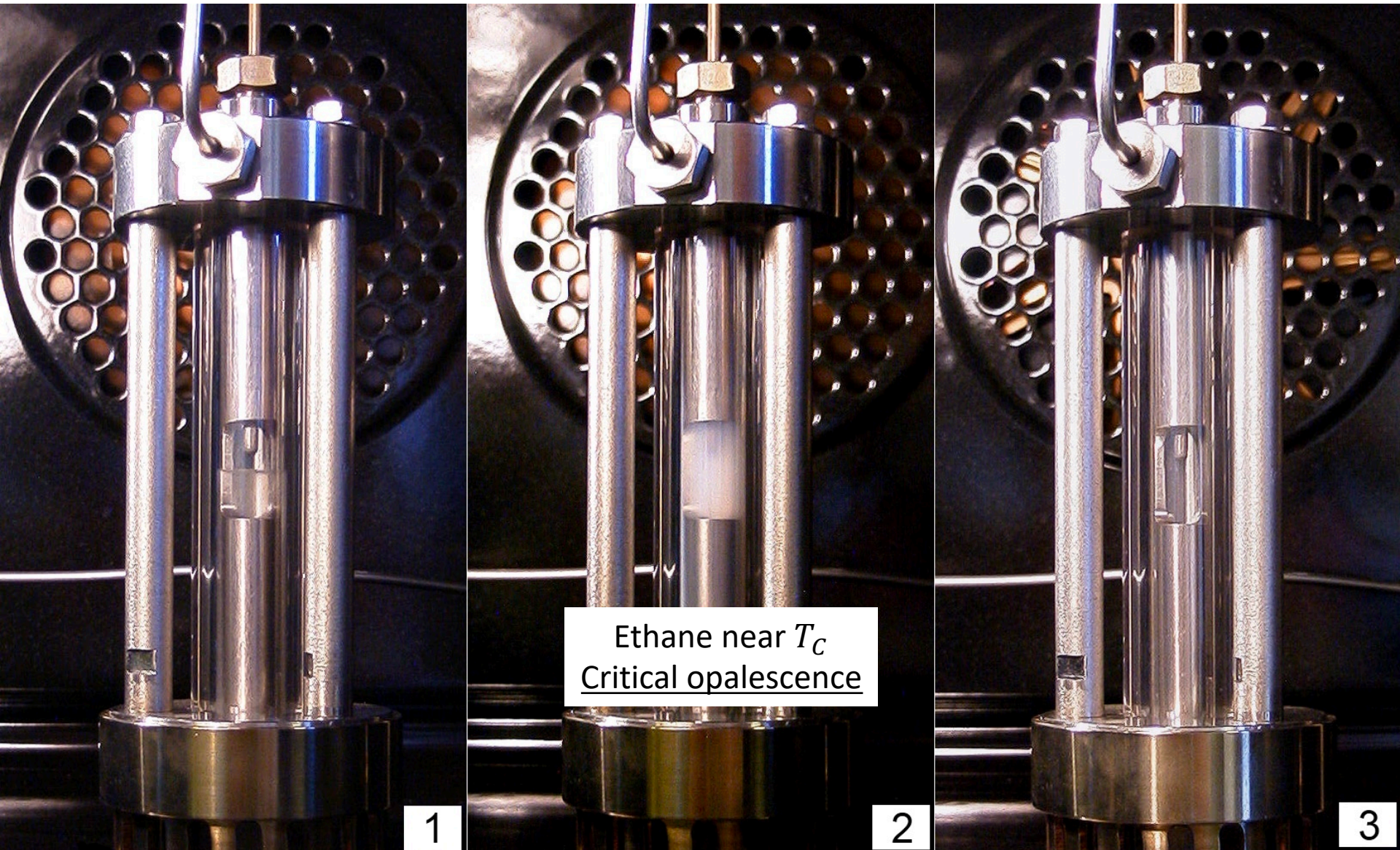
First order phase transition produces a **spinoidal region** – region of phase coexistence = softest point in EoS

Minimum in dv_1/dy is consistent with predictions for a softest point in EoS due to 1st order phase transition



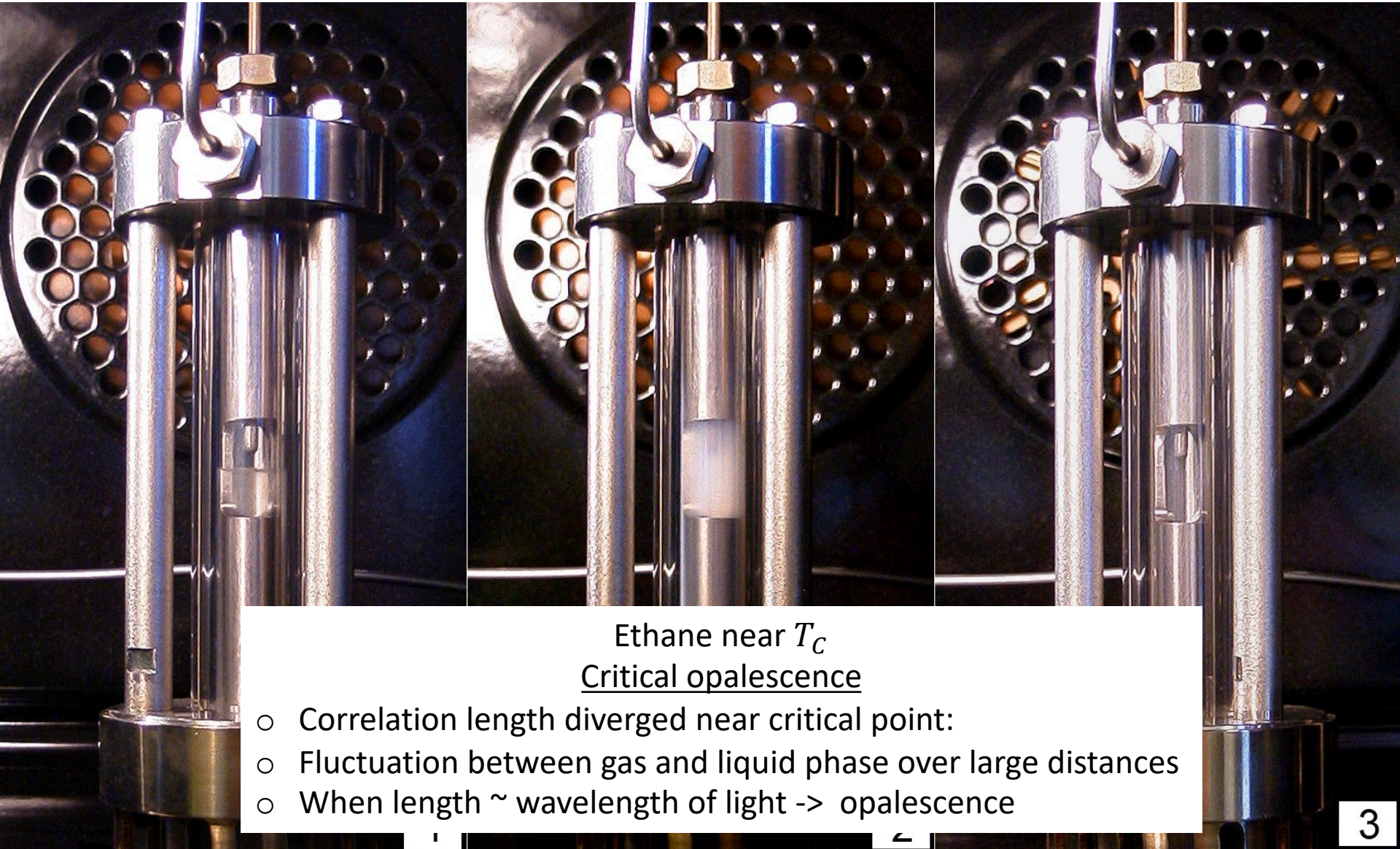
Fluctuations and Criticality

33



Fluctuations and Criticality

34



Ethane near T_c Critical opalescence

- Correlation length diverged near critical point:
- Fluctuation between gas and liquid phase over large distances
- When length \sim wavelength of light \rightarrow opalescence

3

Fluctuations in QCD

Q: Fluctuations of what? A: Conserved quantities

Charge(C), Baryon Number(B), Strangeness(S)

- Look for event-by-event fluctuations in a fixed volume
- Corresponds to divergence in susceptibilities

$$\chi_{lmn}^{BSQ} = \frac{\partial^{l+m+n}(P/T^4)}{\partial(\mu_B/T)^l \partial(\mu_S/T)^m \partial(\mu_S/T)^n} \quad \mu_i = \langle(\delta N)^i\rangle$$

$$\delta N = N - \langle N \rangle$$

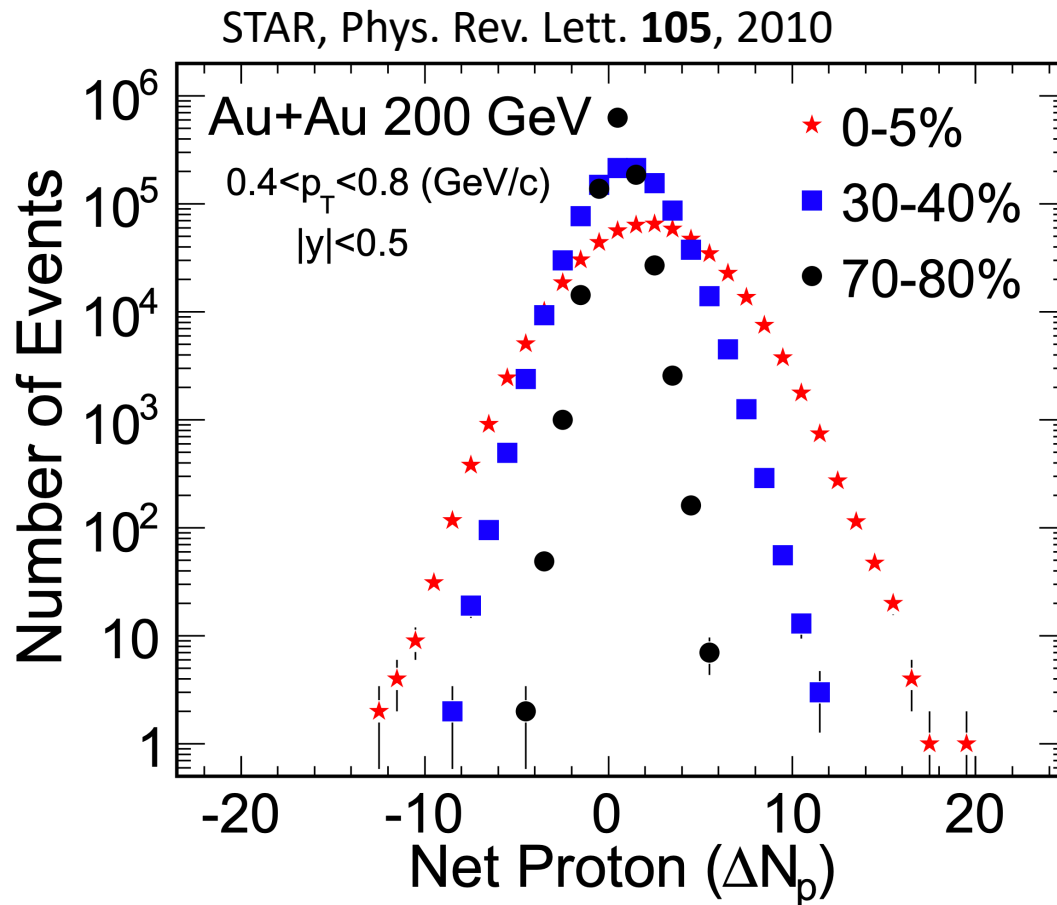
$$\begin{array}{llllll} M & = & K_1 & = & \mu & = & \langle N \rangle & = & VT^3 \cdot \chi_1 \\ \sigma^2 & = & K_2 & = & \mu_2 & = & \langle (\delta N)^2 \rangle & = & VT^3 \cdot \chi_2 \\ S & = & K_3/\sigma^3 & = & \mu_3/\sigma^3 & = & \langle (\delta N)^3 \rangle / \sigma^3 & = & VT^3 \cdot \chi_3 / (VT^3 \cdot \chi_2)^{3/2} \\ \kappa & = & K_4/\sigma^4 & = & (\mu_4 - 3\mu_2^2)/\mu_2^2 & = & \langle (\delta N)^4 \rangle / \sigma^4 - 3 & = & (VT^3 \cdot \chi_4) / (VT^3 \cdot \chi_2)^2 \end{array}$$

Susceptibilities are directly accessible in lattice QCD!

BUT – cannot predict behavior near Critical Point

Fluctuations in QCD

36



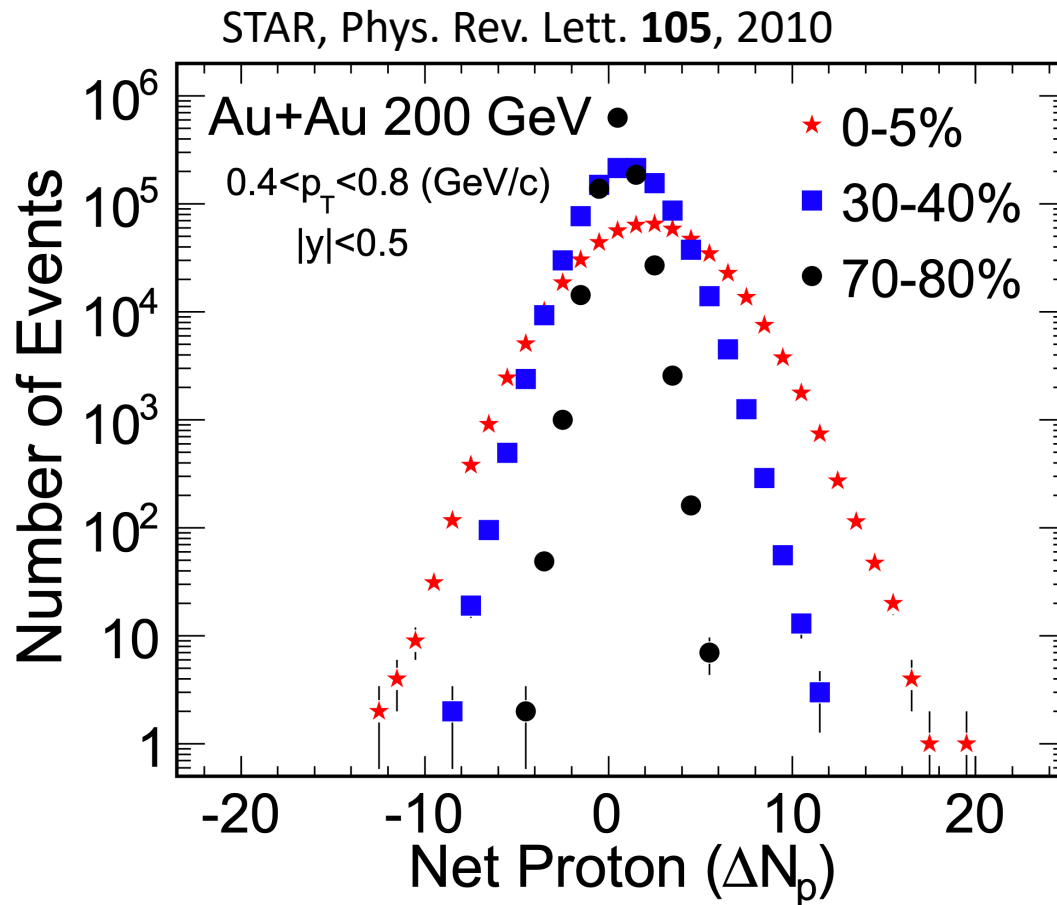
Experimentally, we cannot measure baryon number **directly**

Proxy : net-protons = $p - \bar{p}$

Also, we don't know the volume...

Fluctuations in QCD

37



Experimentally, we cannot measure baryon number **directly**

Proxy : net-protons = $p - \bar{p}$

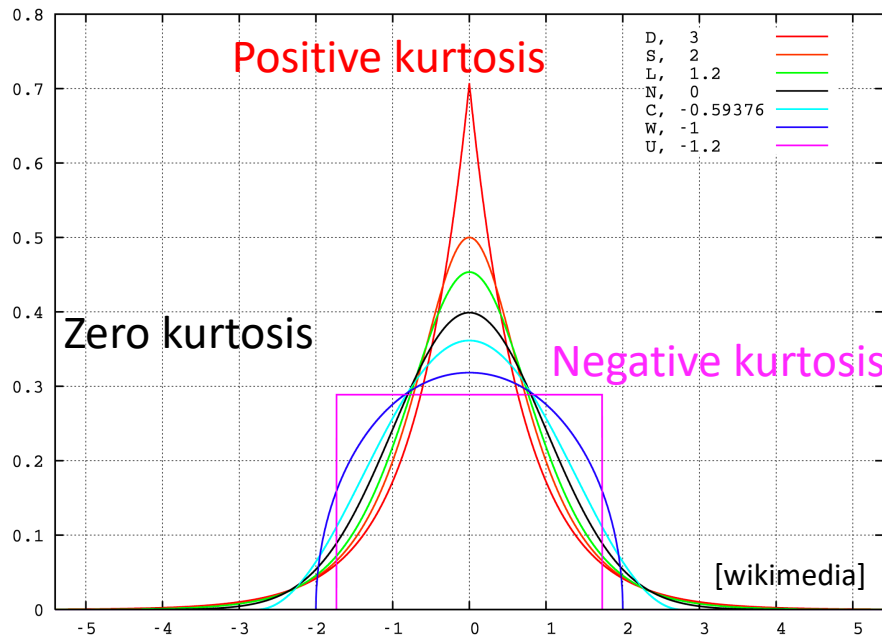
Also, we don't know the volume...

Make ratios or products that cancel volume term:

$$\kappa \sigma^2 = \frac{\chi_B^{(4)}}{\chi_B^{(2)}/T^2}$$

Independent of volume!

Fluctuations in QCD



Experimentally, we cannot measure baryon number **directly**

Proxy : net-protons = $p - \bar{p}$

Also, we don't know the volume...

Make ratios or products that cancel volume term:

Very sensitive to the “tails” of distribution:
Very statistics hungry

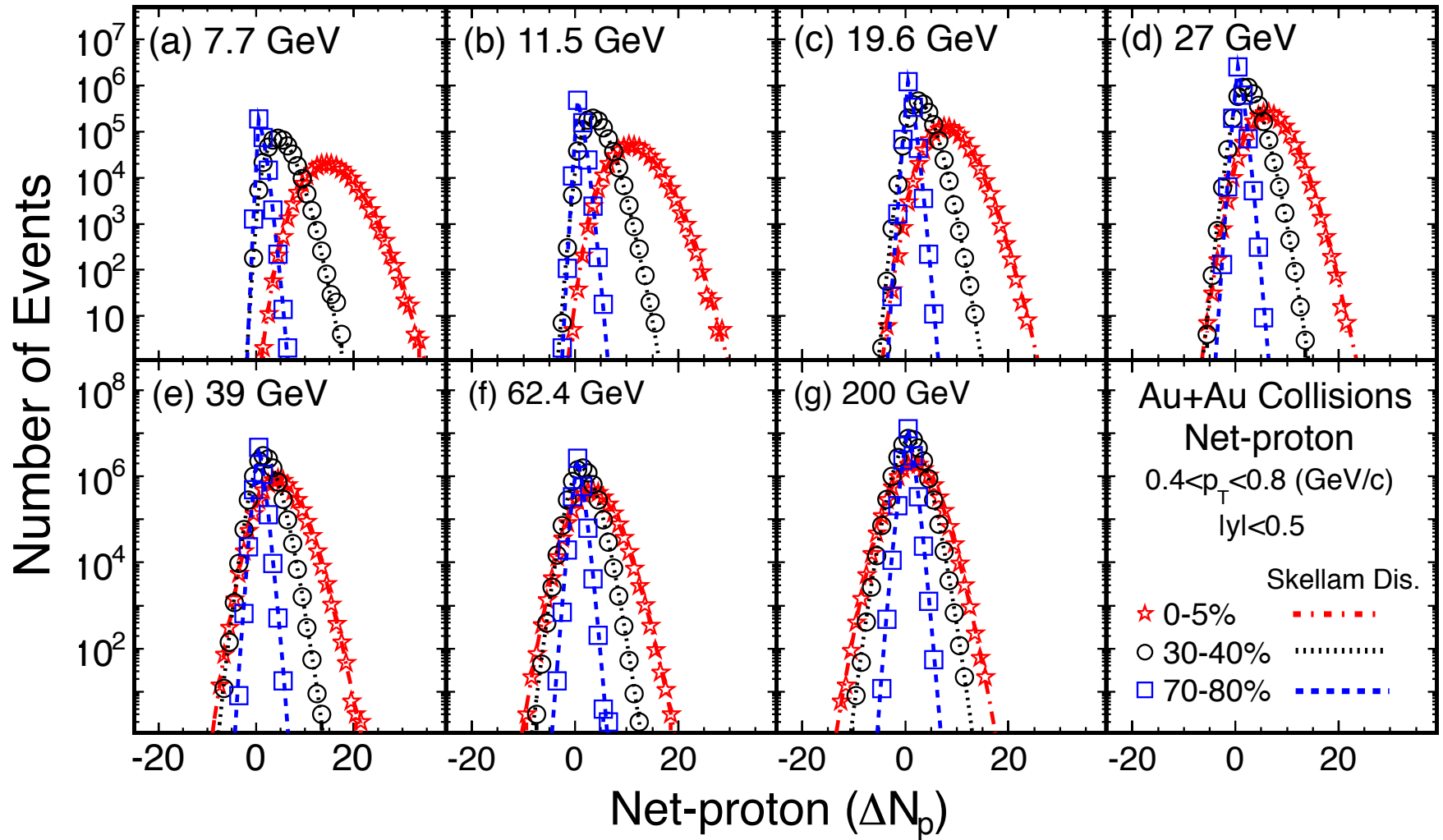
$$\kappa \sigma^2 = \frac{\chi_B^{(4)}}{\chi_B^{(2)} / T^2}$$

Independent of volume!

Measuring Fluctuations

39

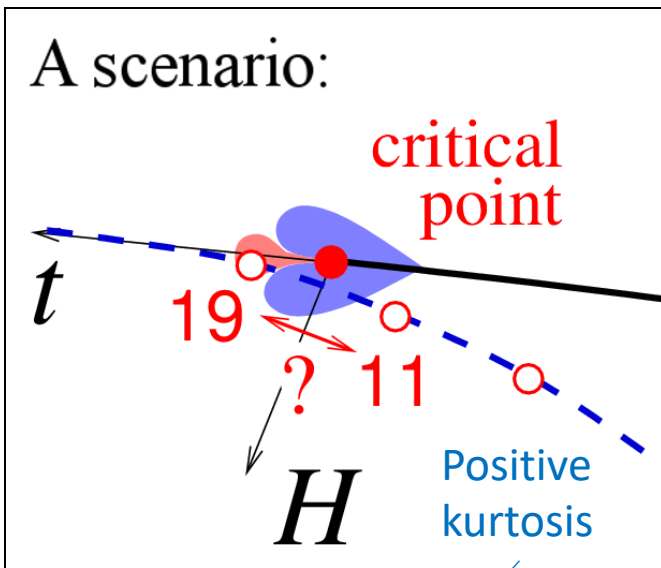
[STAR@RHIC, PRL 112 (2014)]



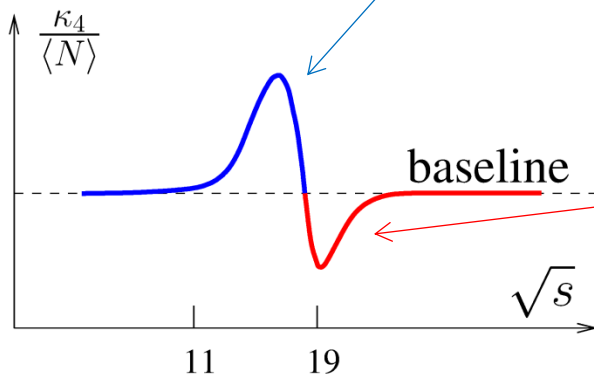
Search for the Critical Point

40

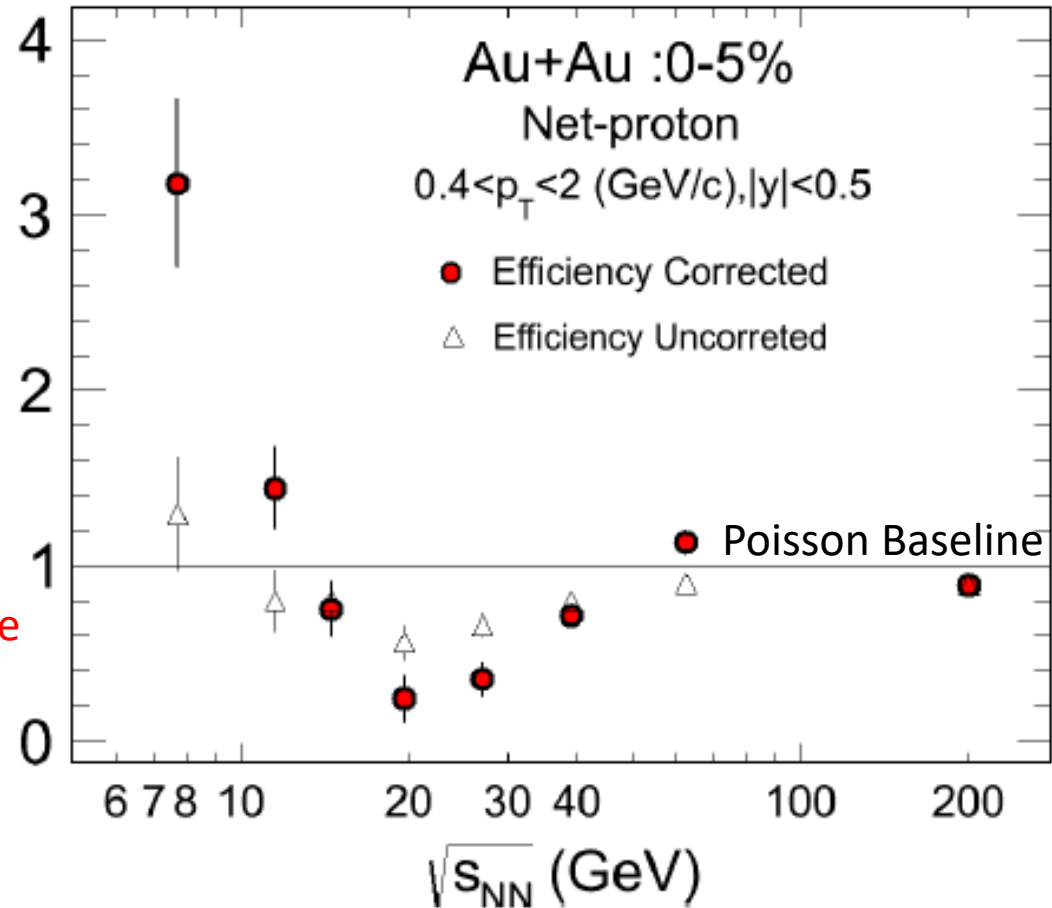
A scenario:



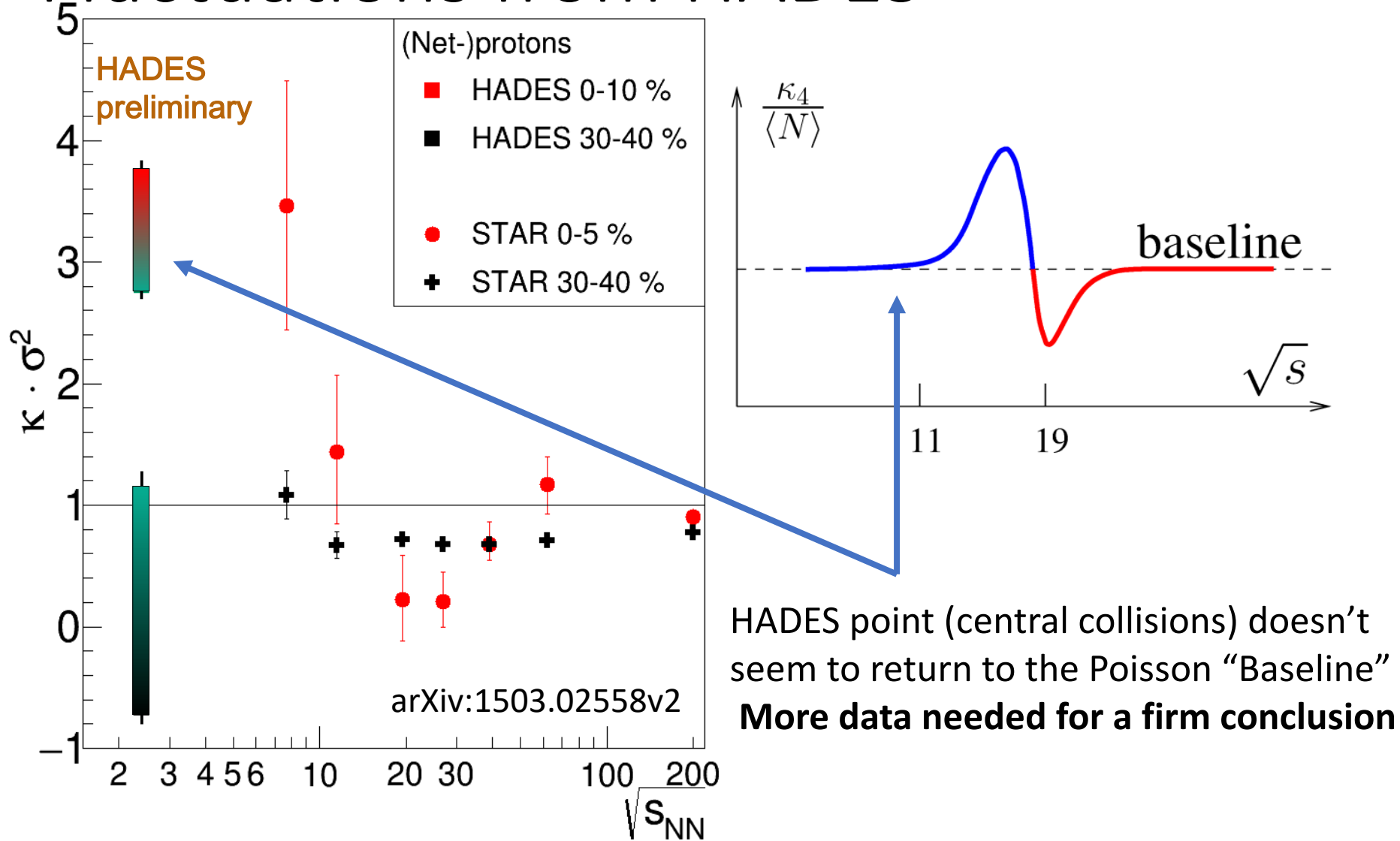
M. Stephanov



<http://arxiv.org/abs/1503.02558v2>



Fluctuations from HADES



red/black = unfolding (preferred method) + vol. flucs. corr.

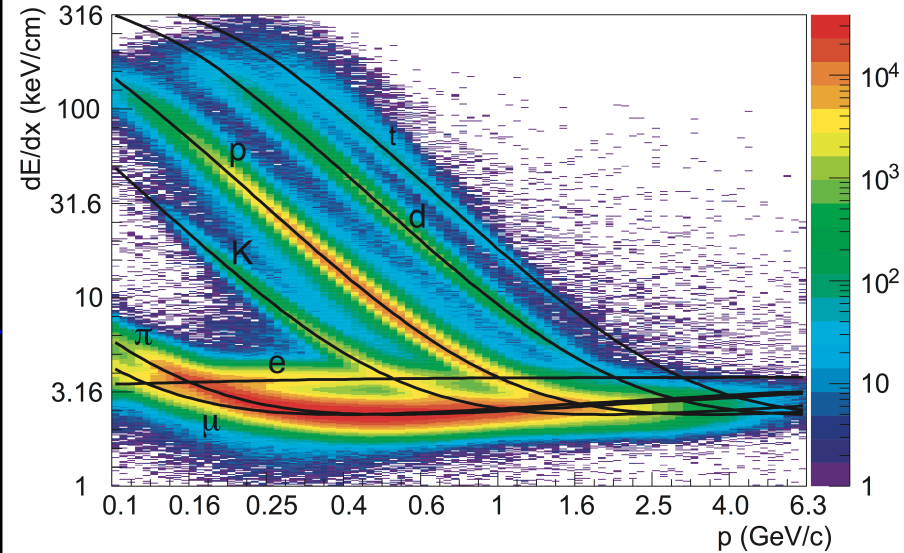
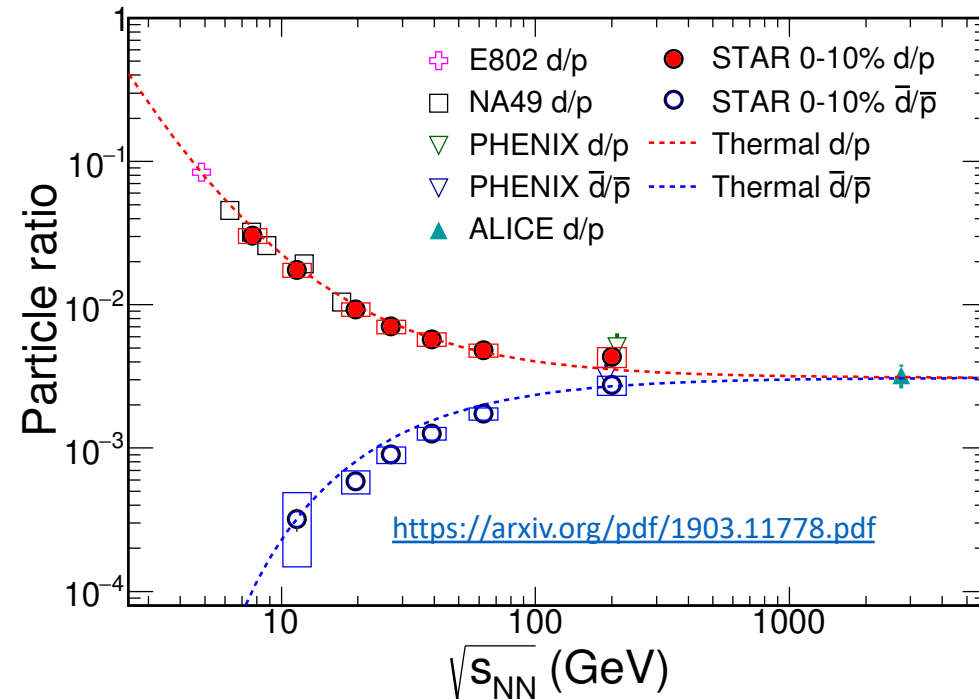
green = evt-by-evt eff correction of factorial moments + vol. flucs. corr.

5/24/19

J. D. Brandenburg : FAIRNESS 2019

Experimental Complications

42



Substantial fraction of protons are bound in heavy fragments : d, t, He, etc.

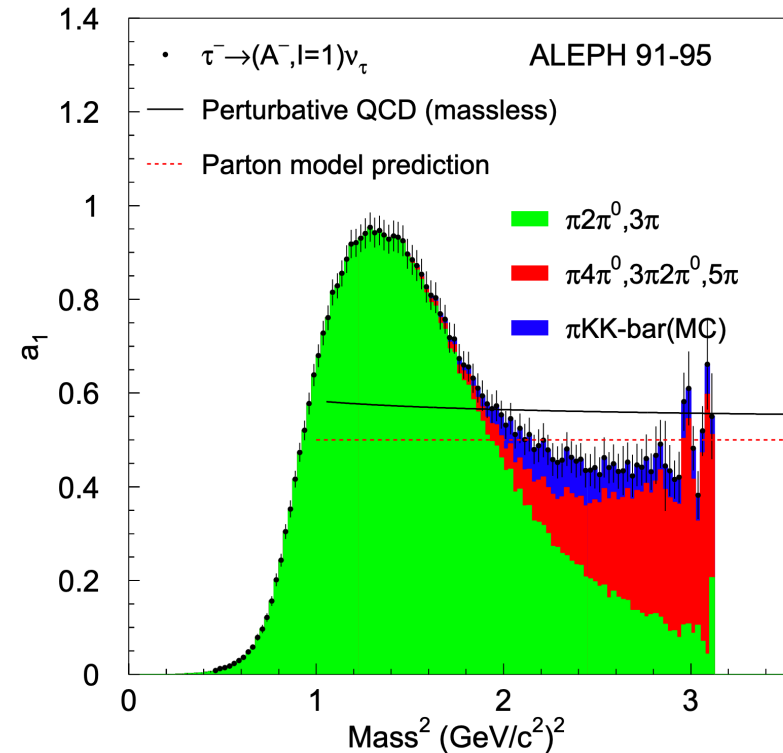
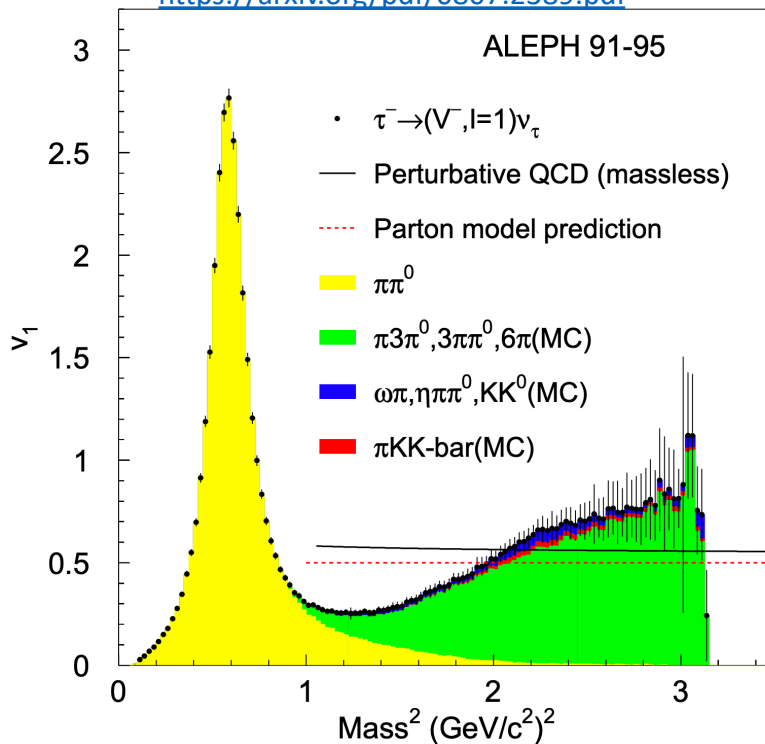
○ How does this effect the baryon number fluctuations measured through protons?

Particle identification techniques do not give 100% purity/efficiency

Search for Chiral Symmetry Restoration

43

<https://arxiv.org/pdf/0807.2389.pdf>



Mass difference of light quark states : $\frac{u\bar{u}-d\bar{d}}{\sqrt{2}}$

Mass (MeV/c²)

π^0

135

ρ^0

770

a_1

~1260

Chiral symmetry restoration : mass of chiral partners will become the same

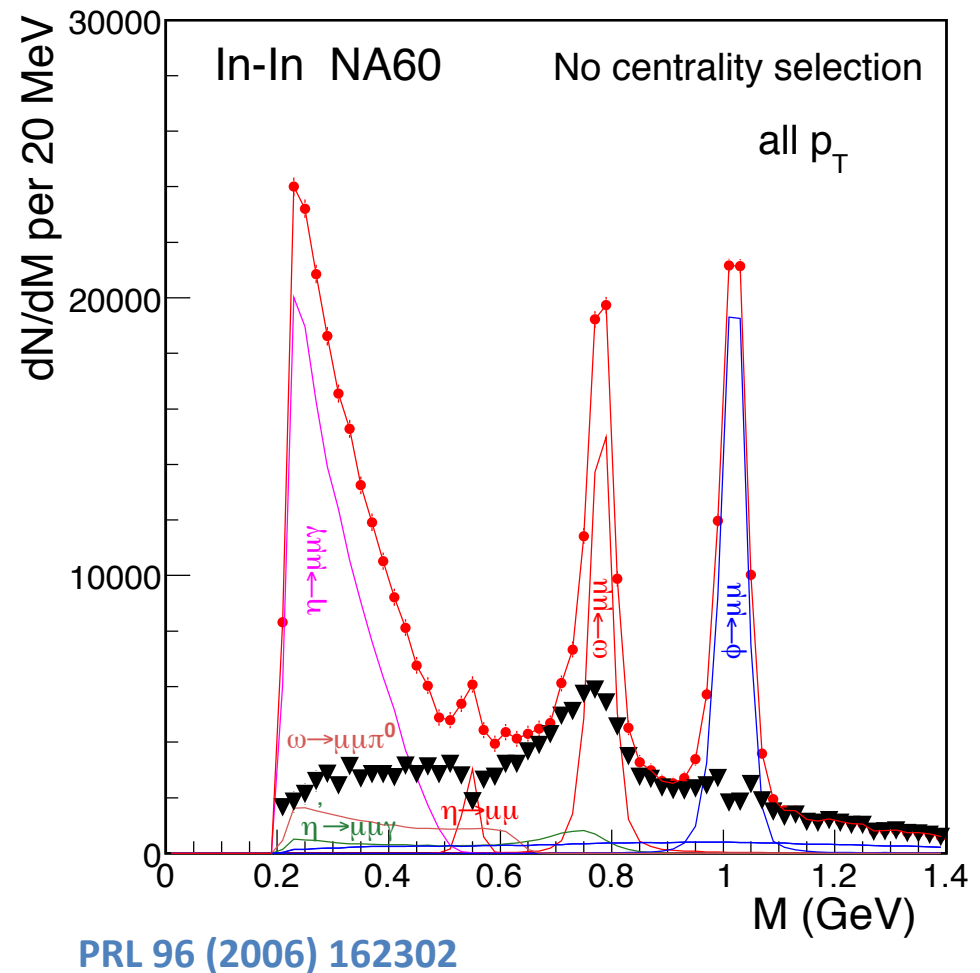
Search for Chiral Symmetry Restoration 44

NA60 Measured $\mu^+\mu^-$
in In+In collisions at
 $\sqrt{s_{NN}} = 17.3$ GeV

Extremely high precision

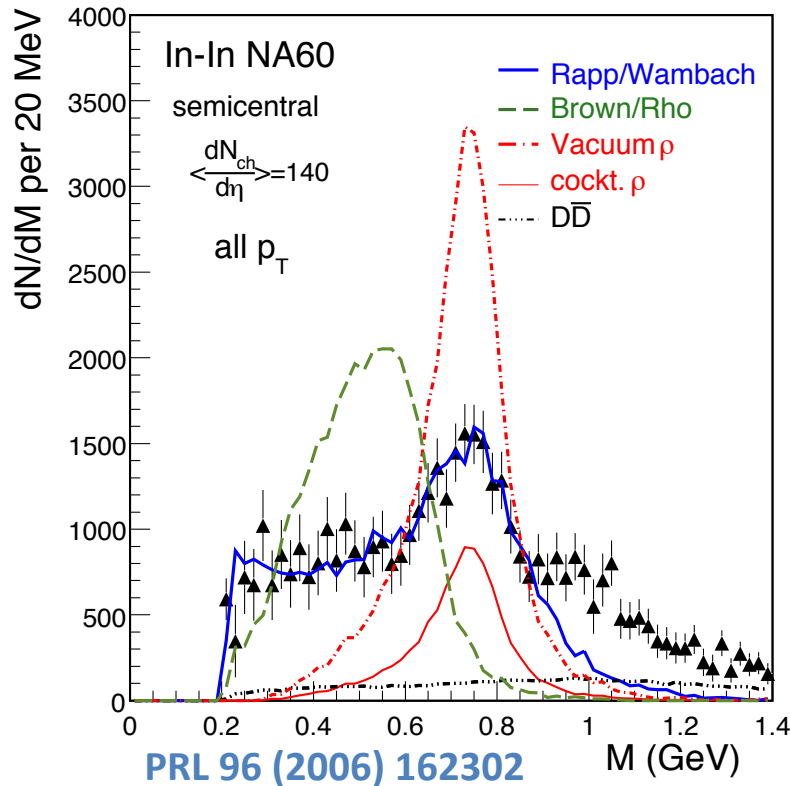
→ Isolate ρ^0 meson

→ Thermal dileptons



NA60 Thermal Dileptons

45

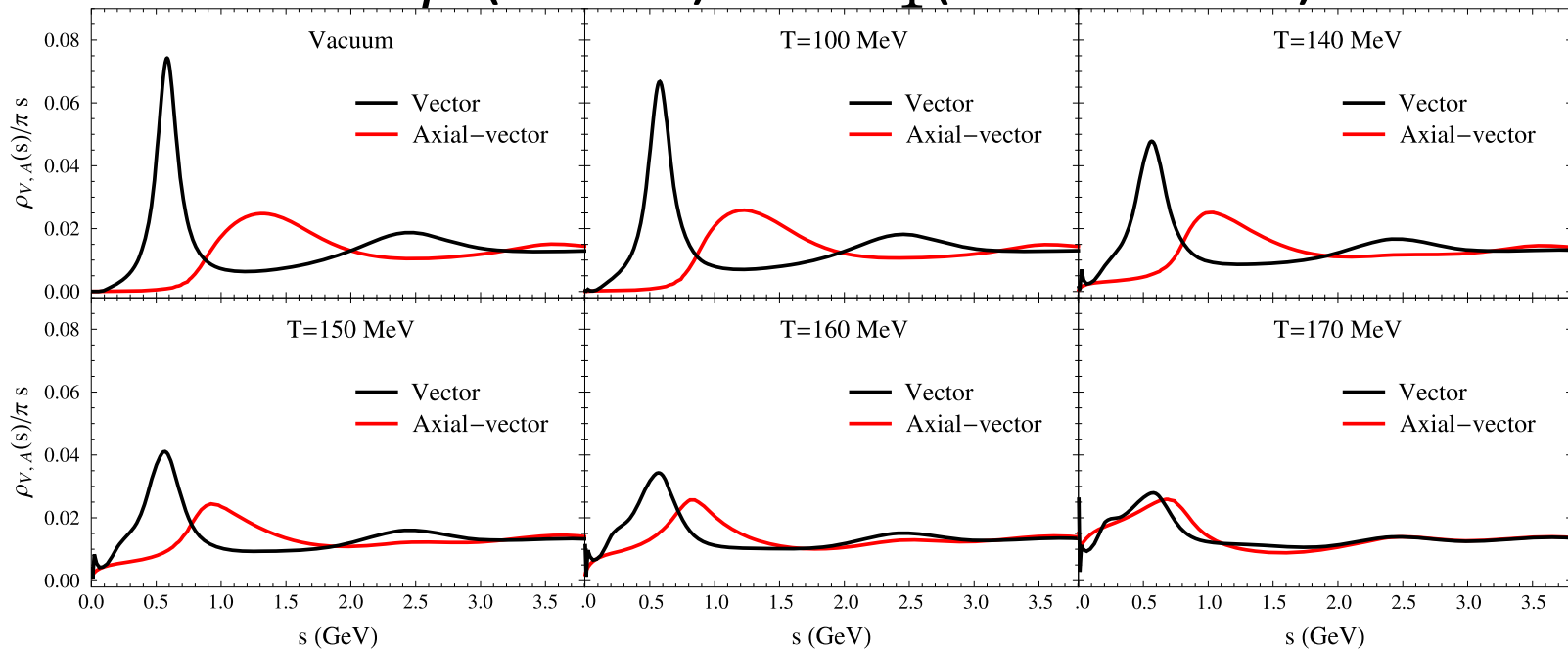


- ρ -meson broadening/melting through medium interactions
- Possible link to **chiral symmetry restoration**

How does low mass excess evolve with collision energy?

Chiral Symmetry Restoration

- Phenomenological description of mass splitting between ρ (vector) and a_1 (axial-vector)

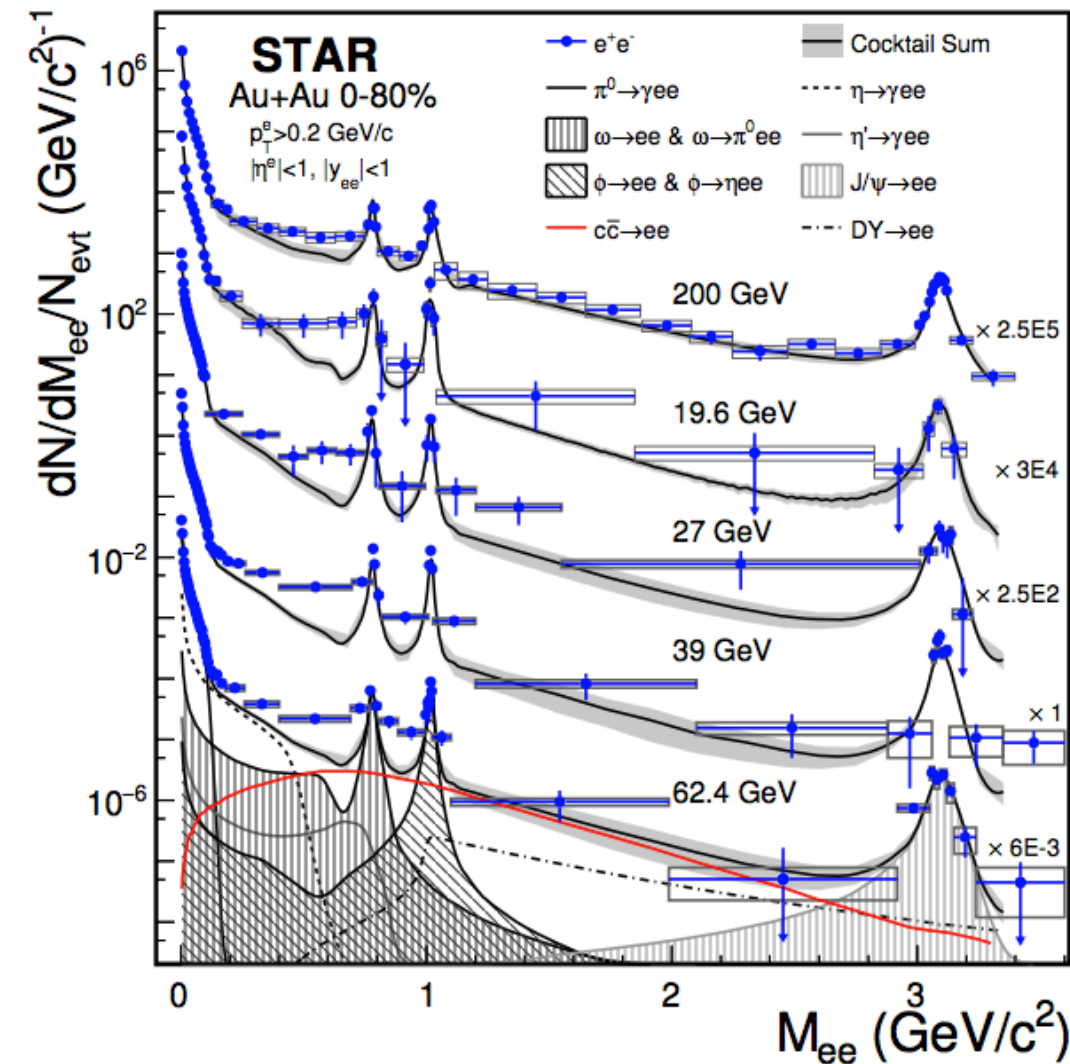


- Use experimentally verified ρ spectral function(SF) \rightarrow search for a_1 SF satisfying QCD & Weinberg sum rules

$$(\text{QCD}) \quad \frac{1}{M^2} \int ds \frac{\rho_{V/A}(s)}{s} e^{-s/M^2} = \sum_n C_n \langle O_n \rangle \quad (\text{Weinberg}) \quad \int ds (\rho_V - \rho_A) s^n = f_n$$

Systematic study of e^+e^-

47

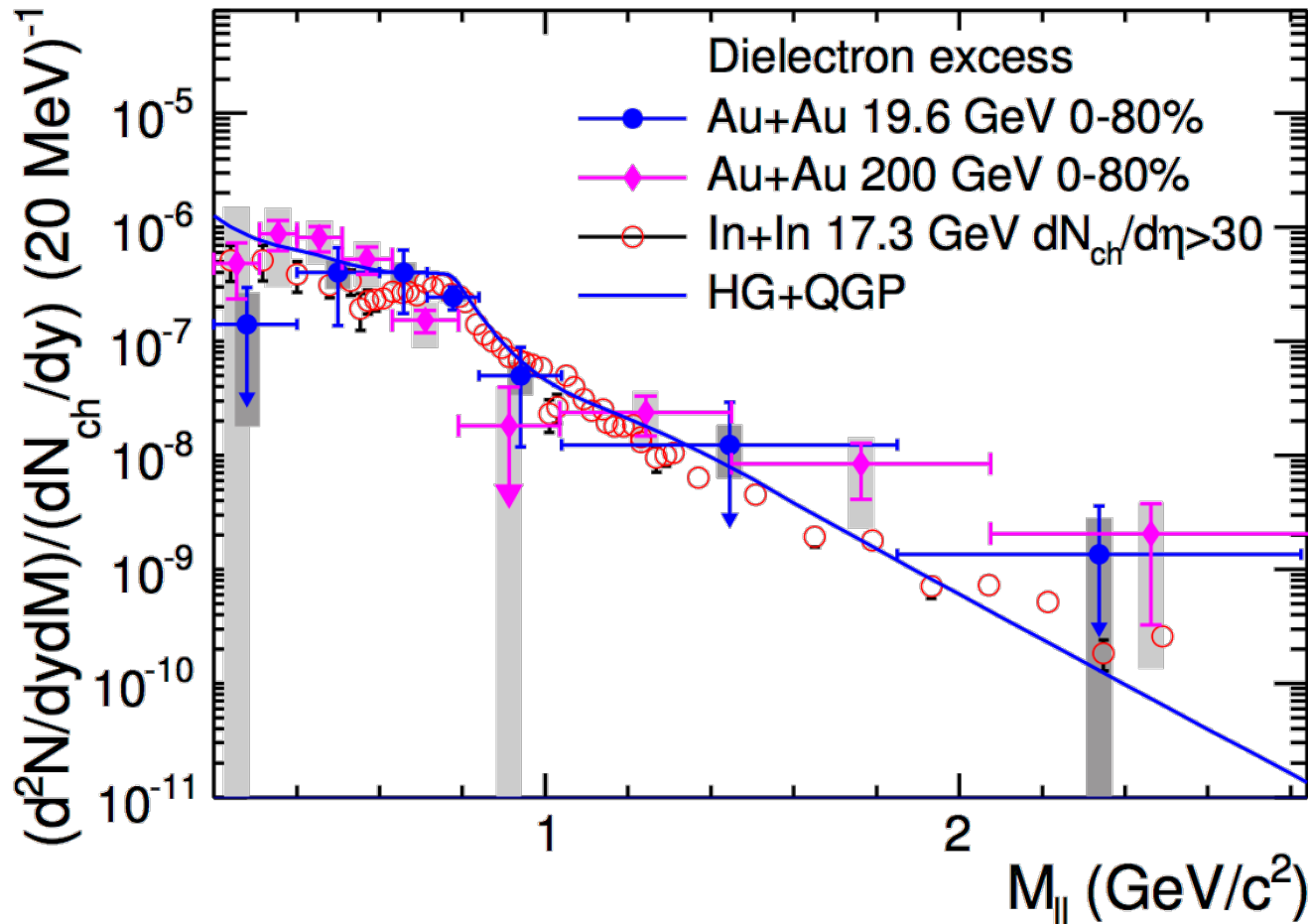


STAR measurements of e^+e^- in BES 1

- Away from ρ – spectra well described by “hadronic cocktail”
- Sustained excess observed in ρ meson region

The Melted ρ

48



Model includes :

ρ^0 “melting” from chiral symmetry restoration

QGP radiating like a thermal black-body

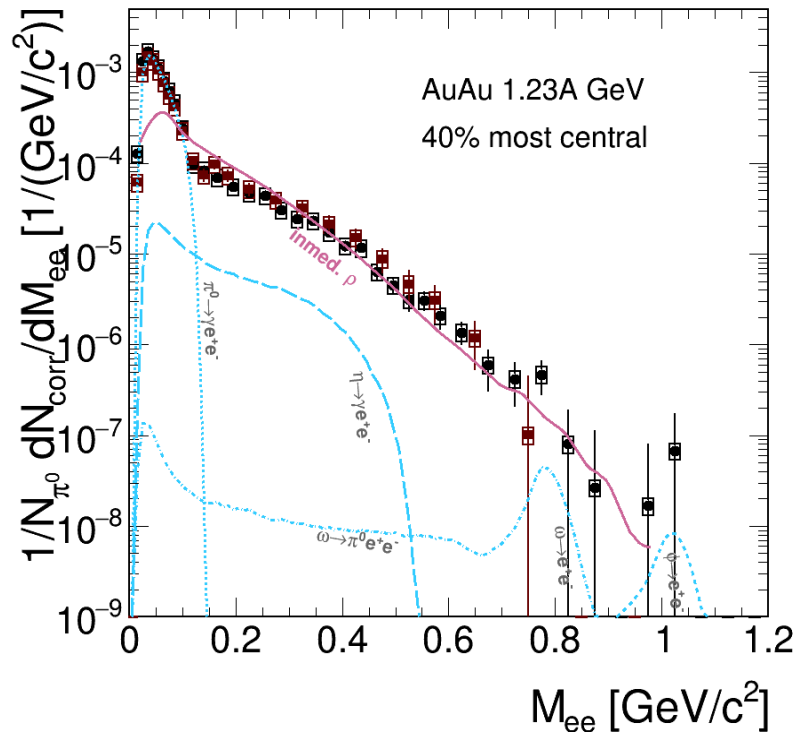
NA60 & STAR excess yields are well described by melted ρ +QGP model

Strong experimental evidence that ρ -meson “melts” →

Hohler & Rapp: “Is ρ -meson melting compatible with chiral restoration?”

HADES High Precision Measurements

49



Significant excess observed over expected sources

But $T \approx 80 \text{ MeV} < T_C$!

At high μ_B the Chiral transition may proceed at lower T than deconfinement transition

$$T_C^{\chi} < T_C^{QGP}$$

Excess is well described by the same in-medium broadened ρ as for higher energies

HADES in CBM: Measure the $\rho - a_1$ mixing

Summary

Heavy Ion Collisions gives us the tool we need to systematically exploring the QCD Phase Diagram

1st Order Phase Transition

Net-Proton v_1 shows double sign change → softest point in EoS (phase coexistence region)

Critical End Point

Fluctuations in conserved quantities → access susceptibilities
 $\kappa\sigma^2$ for net-p shows possible deviation from Poisson Baseline

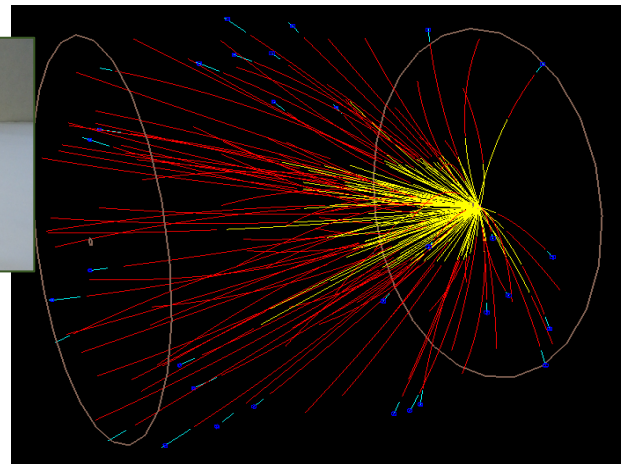
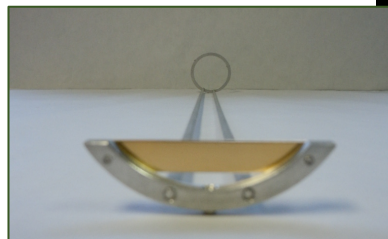
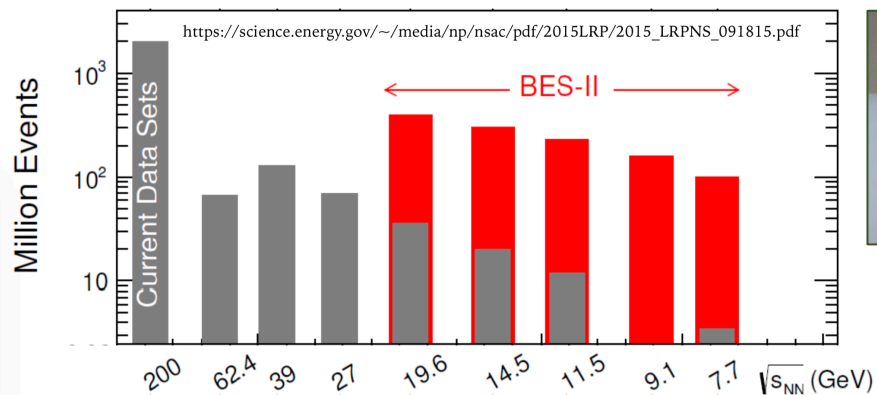
Chiral Symmetry Restoration

Melting of the ρ^0 meson may be linked to Chiral Symmetry restoration
Model consistently describes excess from HADES to RHIC energies

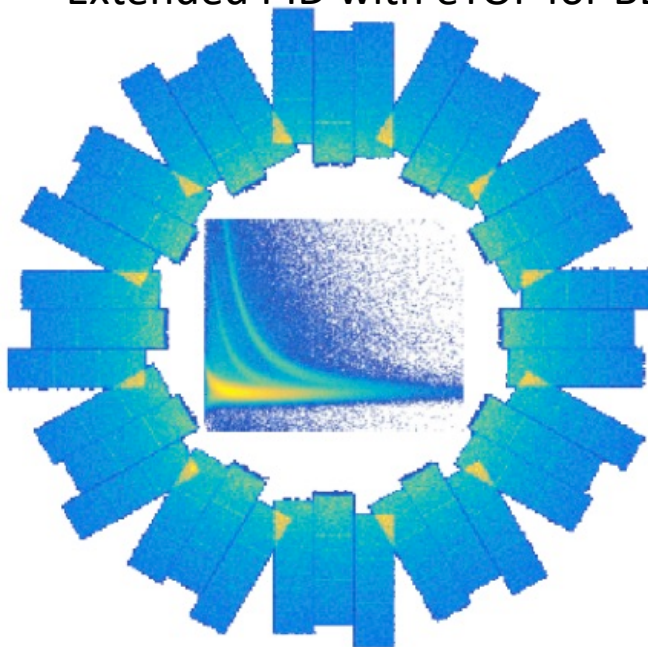
Conclusive results need more data → Outlook

RHIC Beam Energy Scan II

51



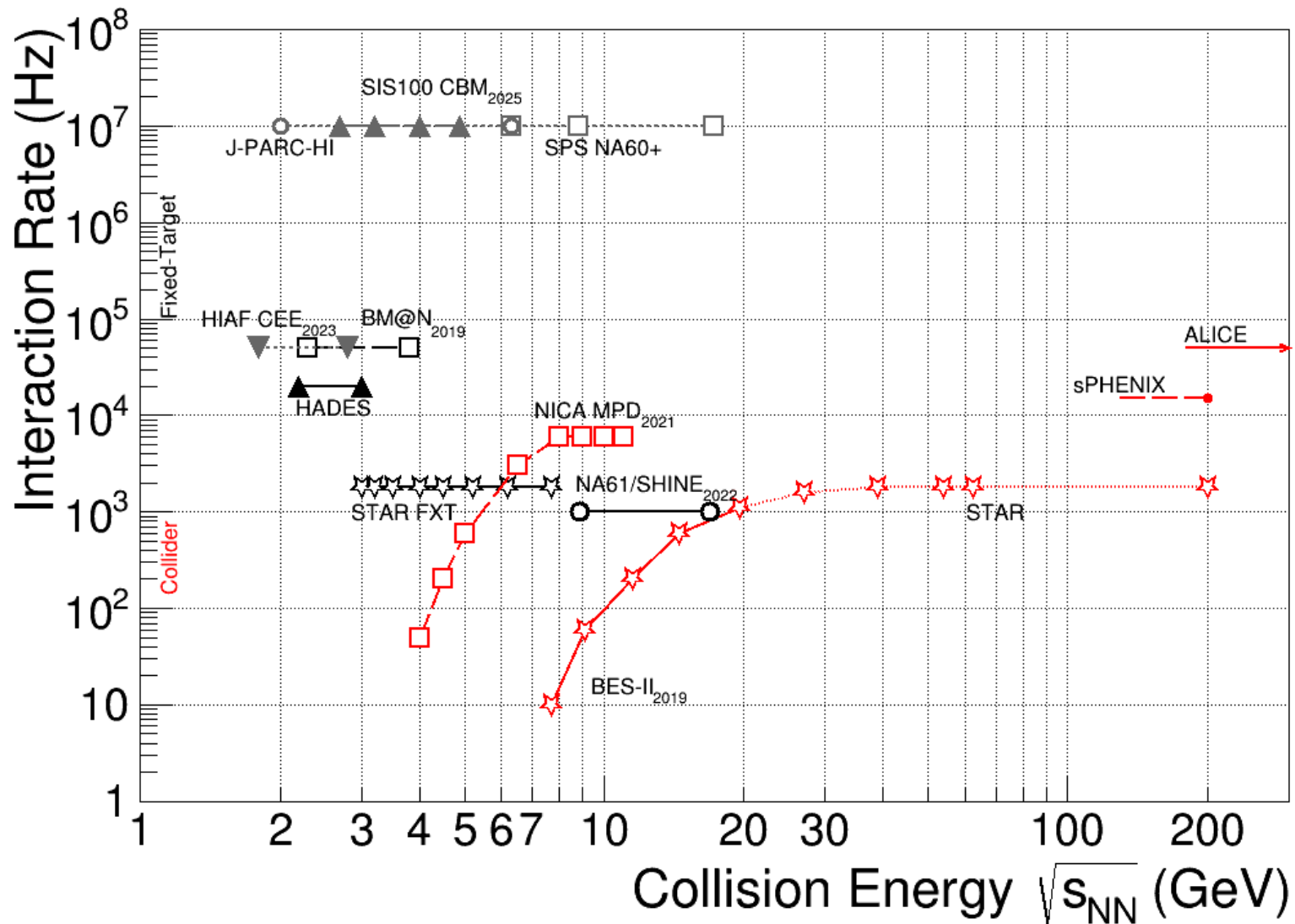
Extended PID with eTOF for BESII



Successful collaboration with CBM!

The Next Era of high- μ_B Experiments

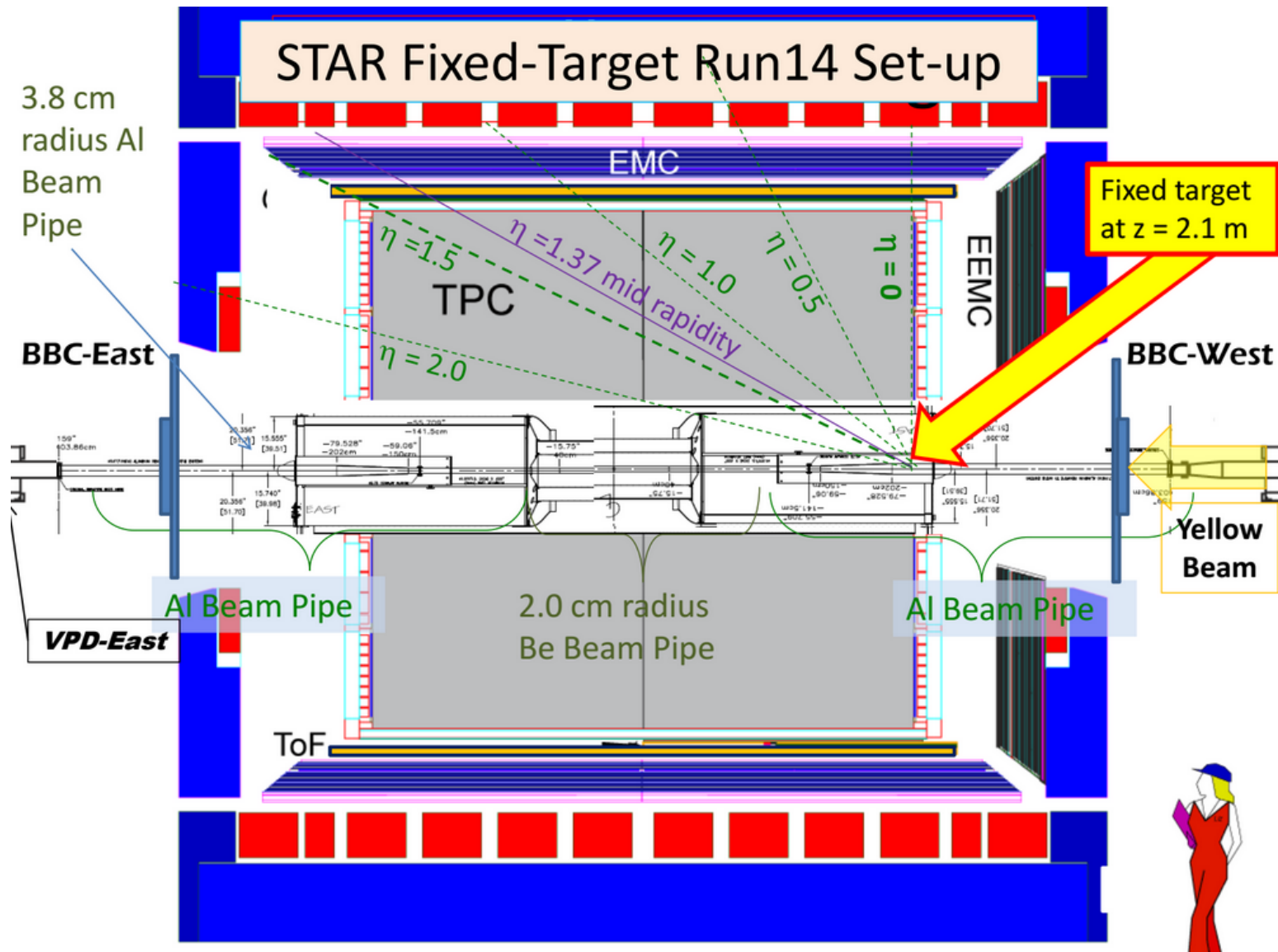
52



Thank you!

STAR Fixed Target

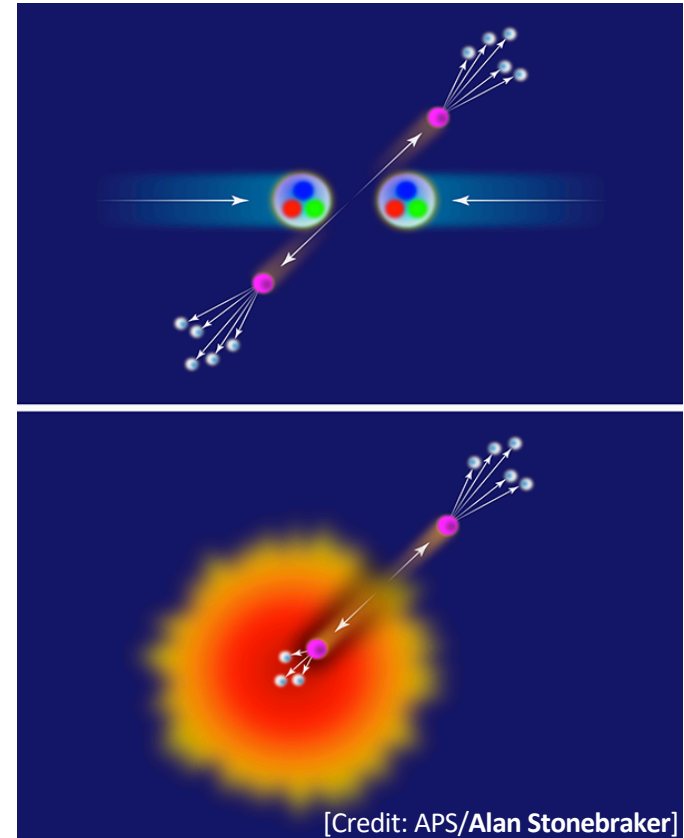
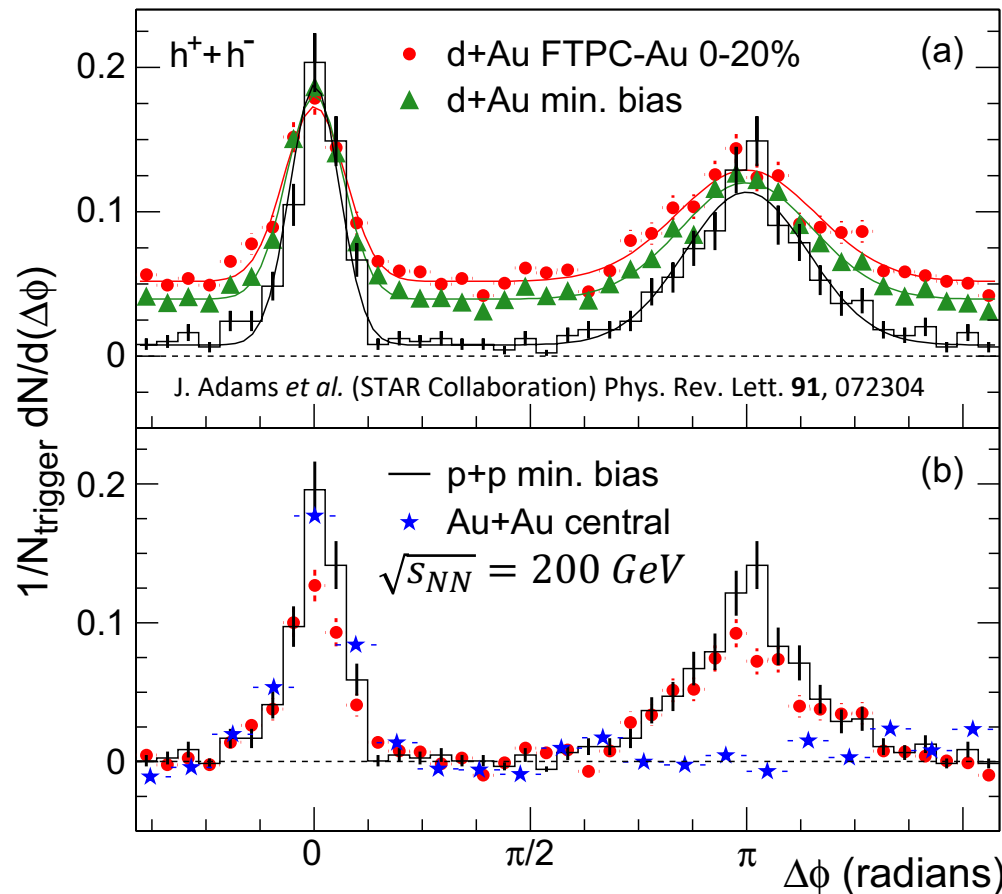
55



Setting the Scene

Jet Quenching in the Quark Gluon Plasma

57

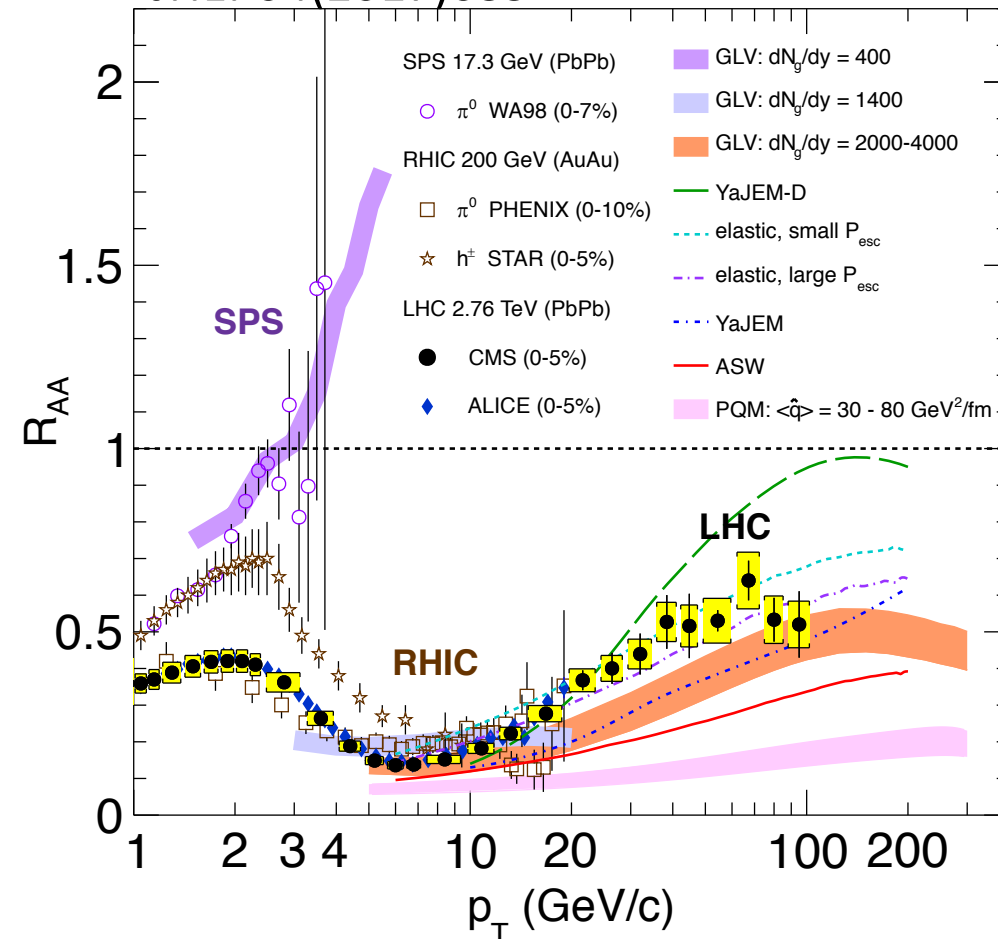


- High- p_T partons (quarks and gluons) – forebears of jets – lose energy via strong interactions with the partonic medium
- Clear signature of the Quark Gluon Plasma observed by all 4 RHIC experiments in 2003 (cover of PRL!)

Jet Quenching in the Quark Gluon Plasma

58

JHEP04(2017)039



Quantify jet quenching:

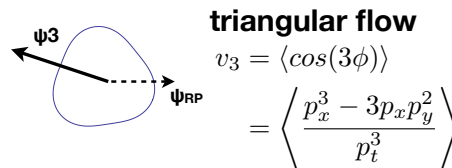
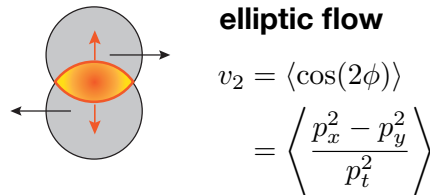
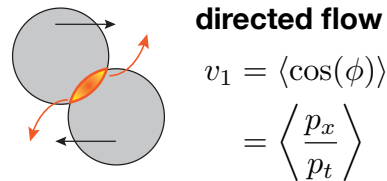
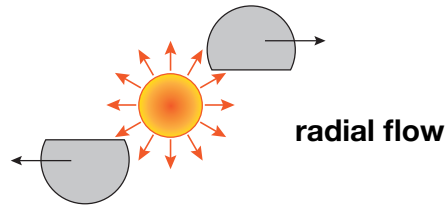
$$R_{AA} = \frac{\text{Yield in AA Collisions}}{\text{Yield in pp Collisions} \times \text{\# "binary" collisions in AA}}$$

Significant Jet quenching ($R_{AA} \ll 1$) has identified as a clear signature of QGP formation

Clear evidence for QGP formation at RHIC and LHC

Finding the QGP

59



Behruz Kardan

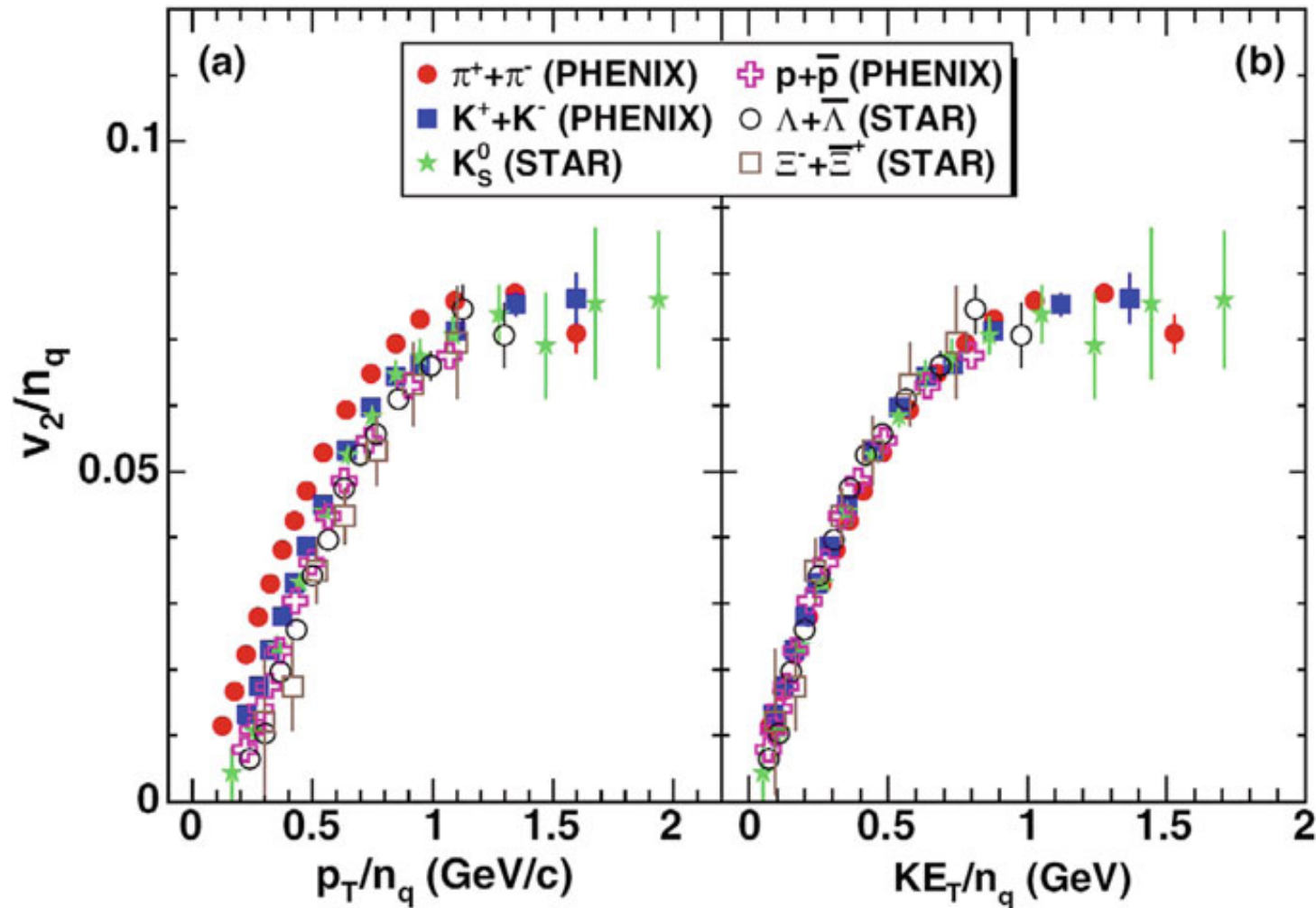
Look at the collective behavior of the medium

$$\frac{dN}{d\varphi} \propto 1 + \sum_{n=1}^{\infty} 2v_n \cos [n(\varphi - \Psi_n)]$$

Where ψ is the azimuthal angle and Ψ_n is the n^{th} order event plane

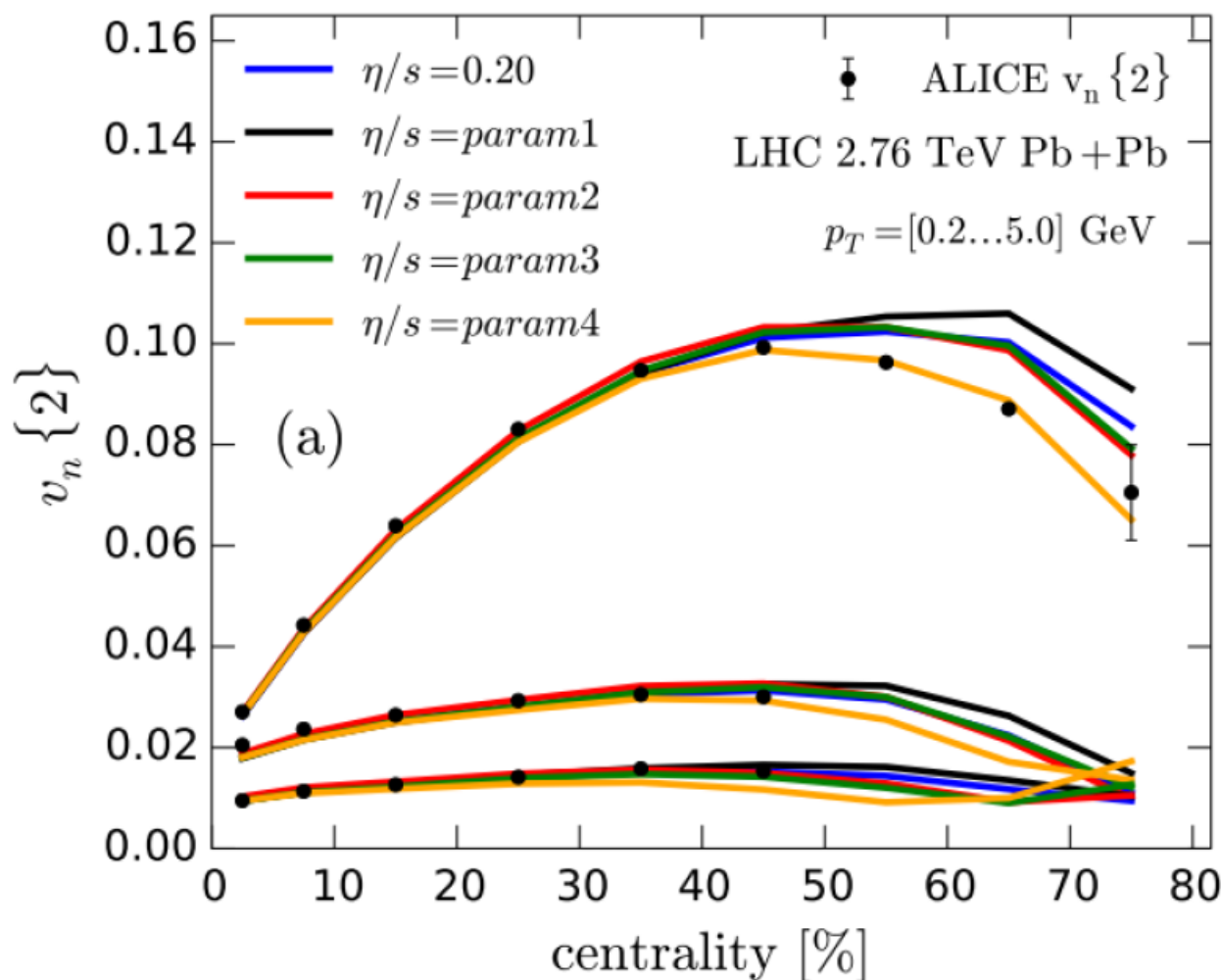
Finding the QGP

○ How do we know the medium is partonic?



QGP : Ideal Gas or Perfect Liquid?

61



Fit and Bayesian Analysis

give most likely value of

$$\eta/s \approx 0.07^{+0.05}_{-0.04} + c(T - T_c)$$

Note, c is mostly

unconstrained [0, 1.6] /GeV

Average η/s strikingly close
the theoretical limit from
holography:

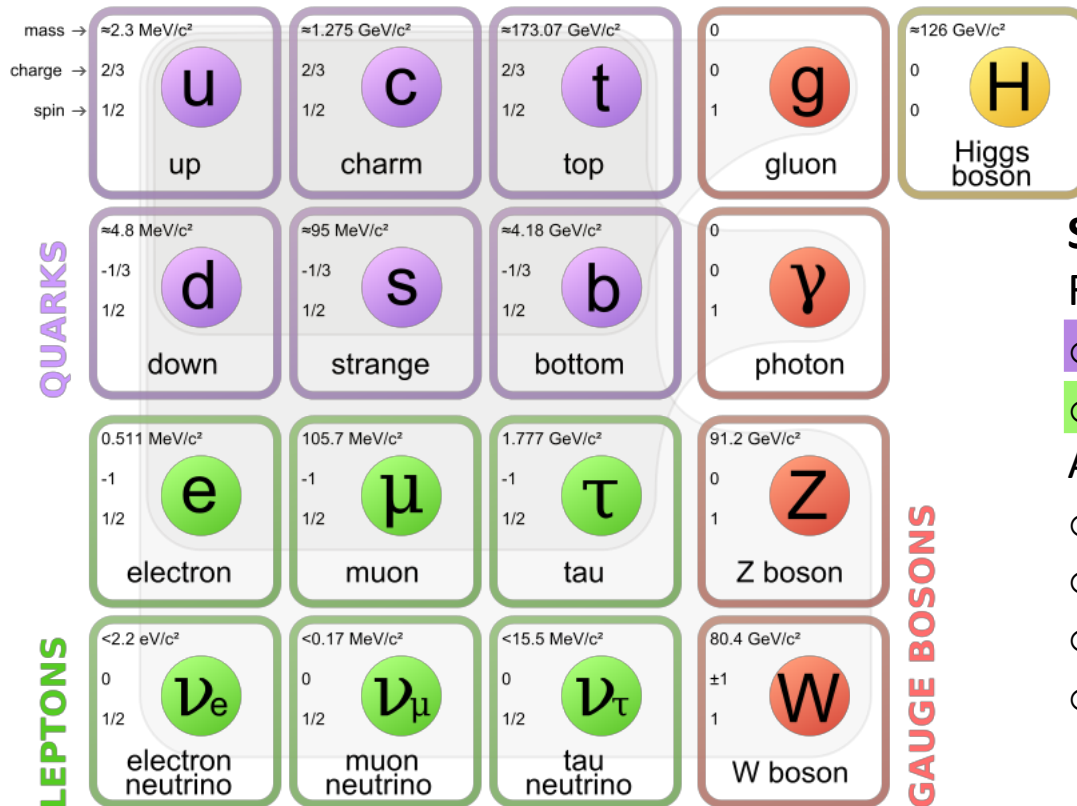
$$(1/4\pi \approx 0.08)$$

If QGP were a weakly
interacting gas of quarks
and gluons : expect
isotropic distribution

Hydrodynamic flow converts spatial anisotropies into momentum anisotropy

Why Study QCD Matter?

62



Standard Model of Physics

Fundamental building blocks of matter:

- Quarks
- Leptons

And their fundamental interactions:

- Electromagnetic
- Weak interaction
- **Strong interaction**
- Gravity (not yet!)

But normal matter isn't just bare quarks

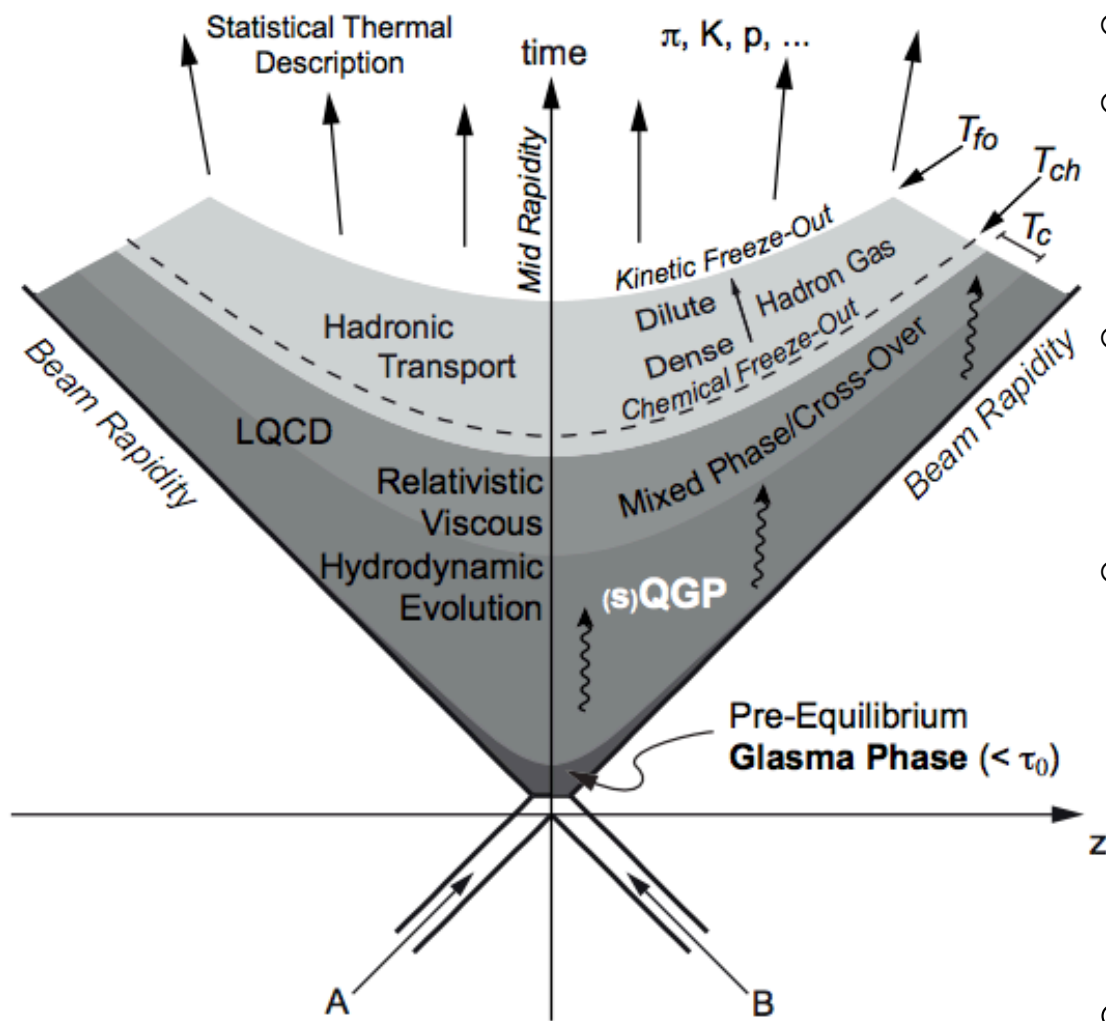
We study QCD to understand the emergence and organization of “normal” matter from quarks and gluons

Life of a Heavy Ion Collision

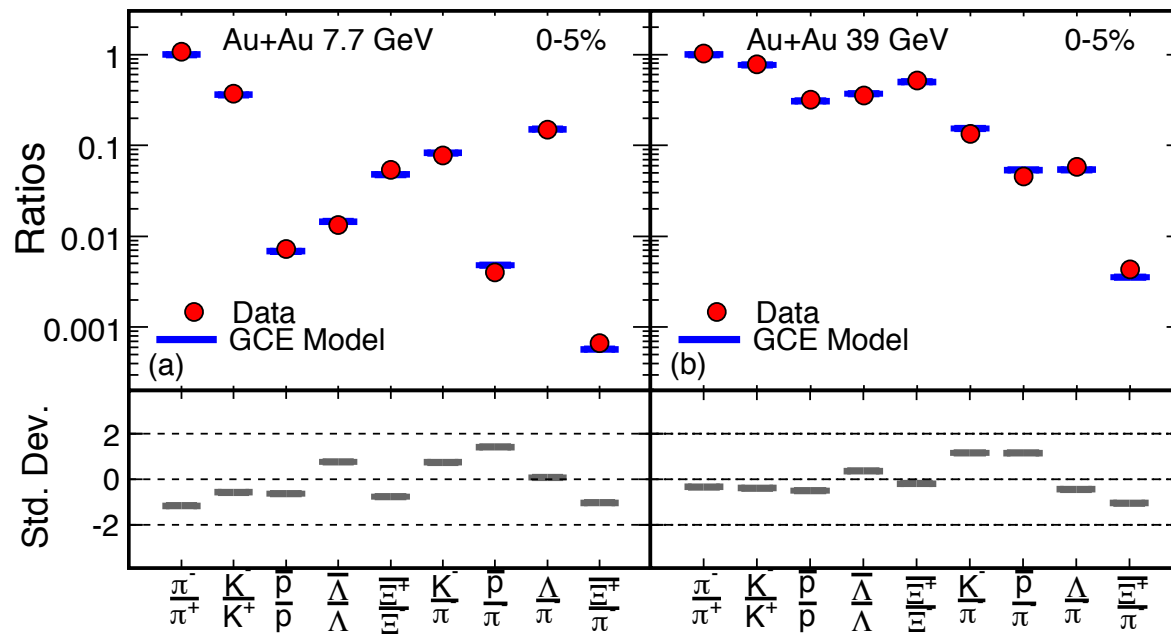
U.W. Heinz, (2004).

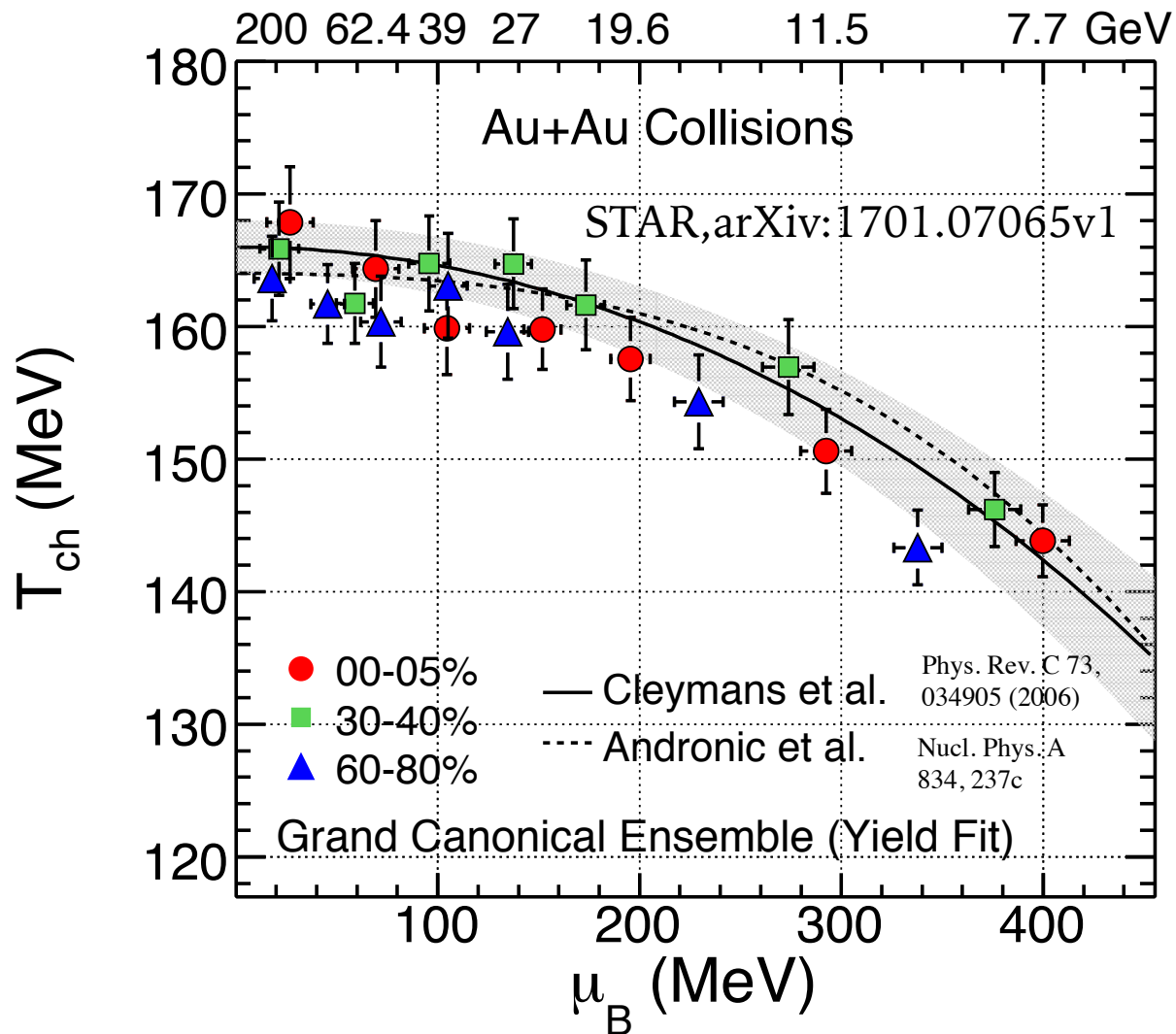
63

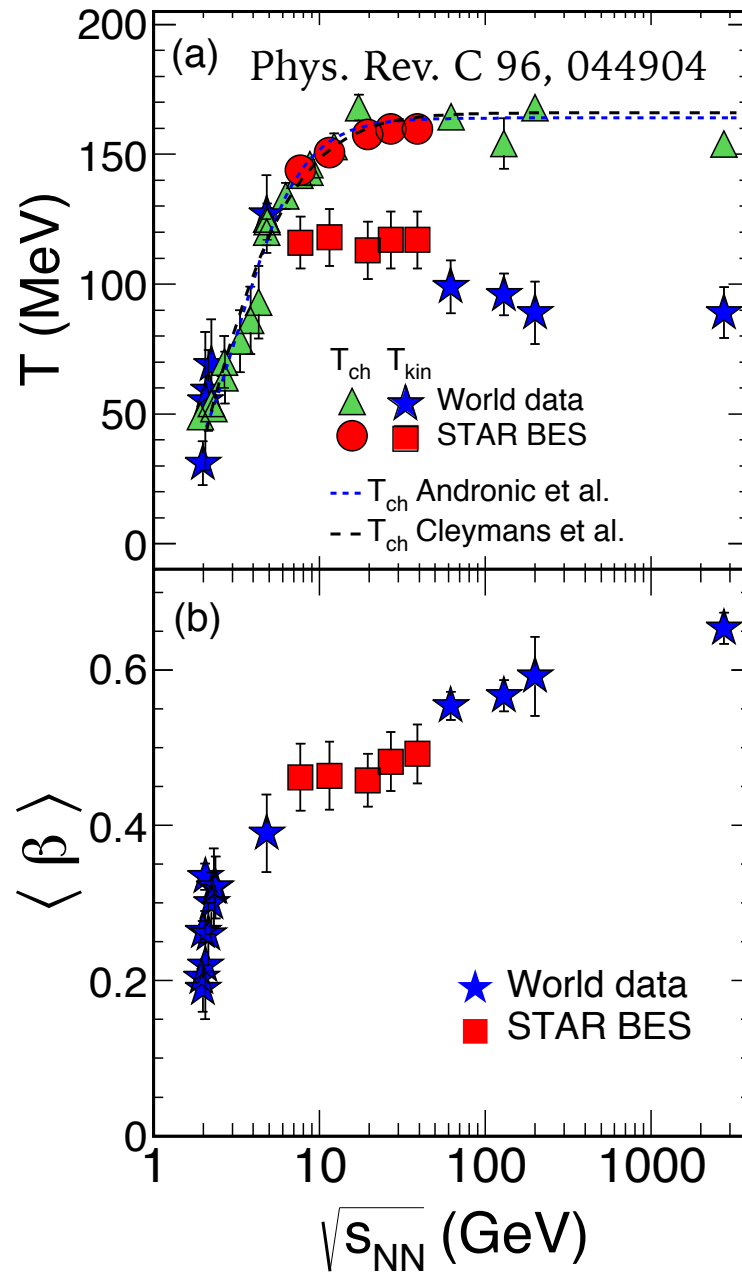
<http://arxiv.org/abs/hep-ph/0407360>.



- $\tau = 0$: initial collision
- $\tau < \sim 1$ fm/c : pre-equilibrium – not much is known about the dynamics. Assume fast approach to thermal equilibrium
- $1 < \tau < \approx 15$ fm/c : QGP is formed if $\varepsilon > \varepsilon_c, T > T_c$
 - Viscous hydrodynamic expansion and cooling
- $\tau > 10 - 15$ fm/c : Freeze-out
 - T drops below T_c
 - Chemical freeze-out : inelastic scatterings cease - particle species are fixed
 - Kinetic freeze-out : elastic scatterings cease – particle momenta are fixed
- Free stream to detectors



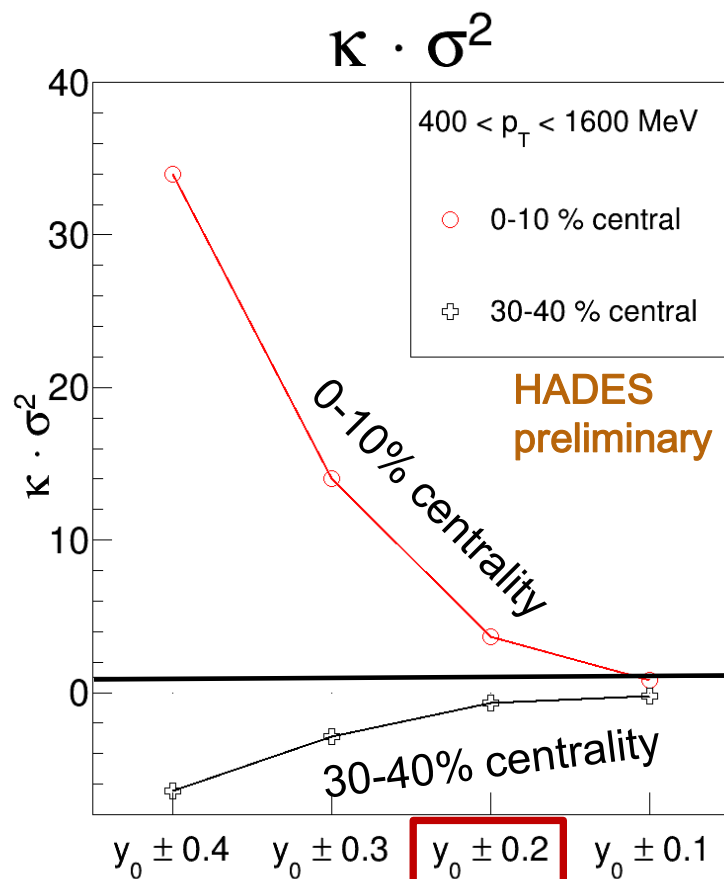




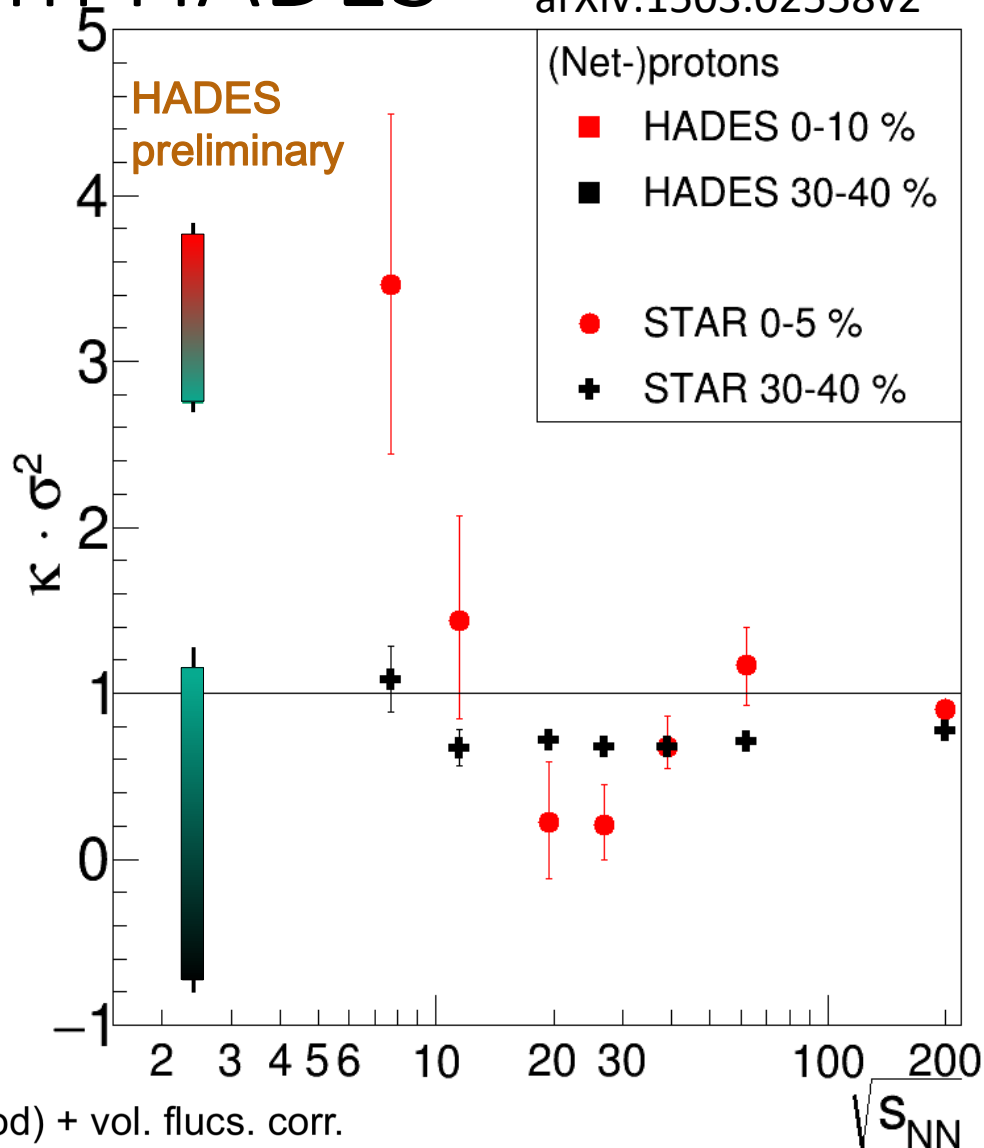
Fluctuations from HADES

67

arXiv:1503.02558v2



Demonstrate Poisson
limit for small volume

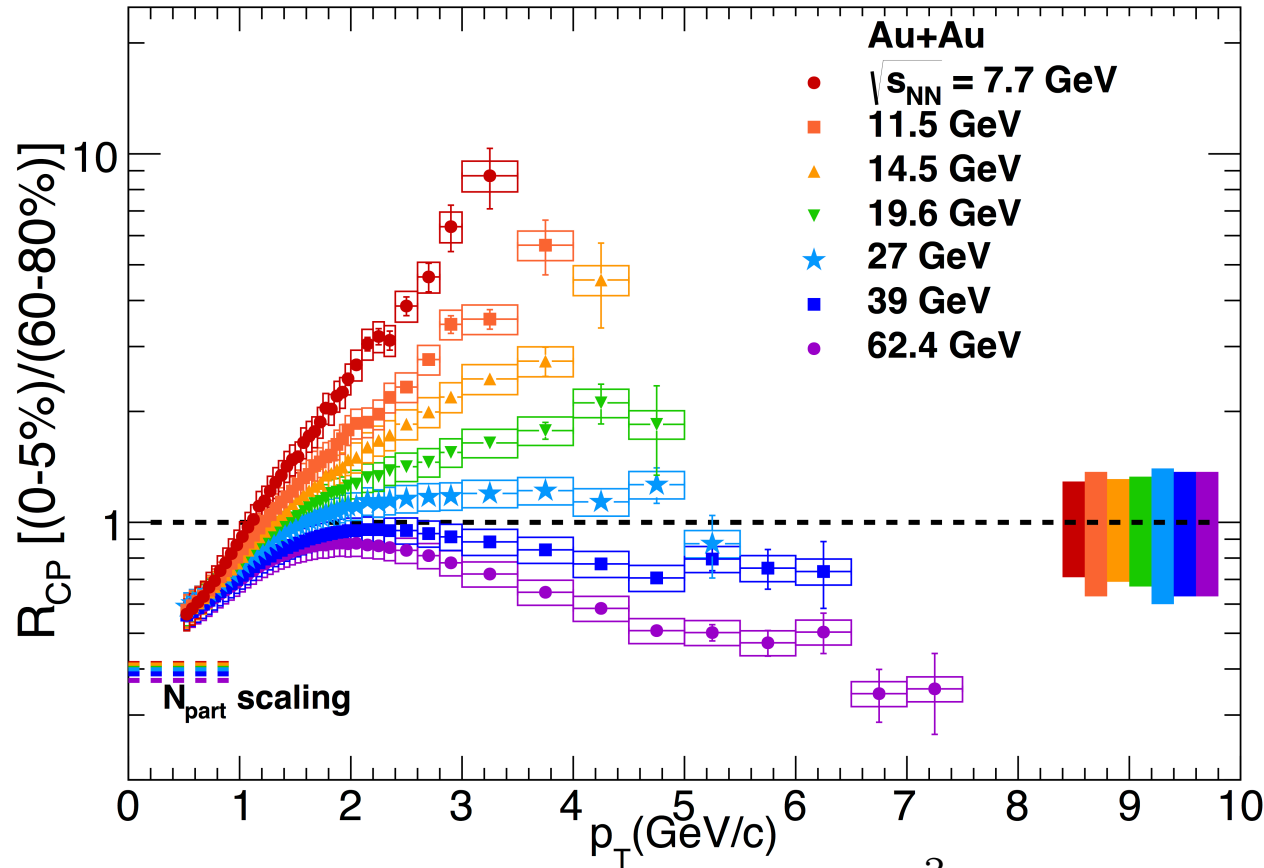


red/black = unfolding (preferred method) + vol. flucs. corr.

green = evt-by-evt eff correction of factorial moments + vol. flucs. corr.

Nuclear Modification Factor in BES 1

68



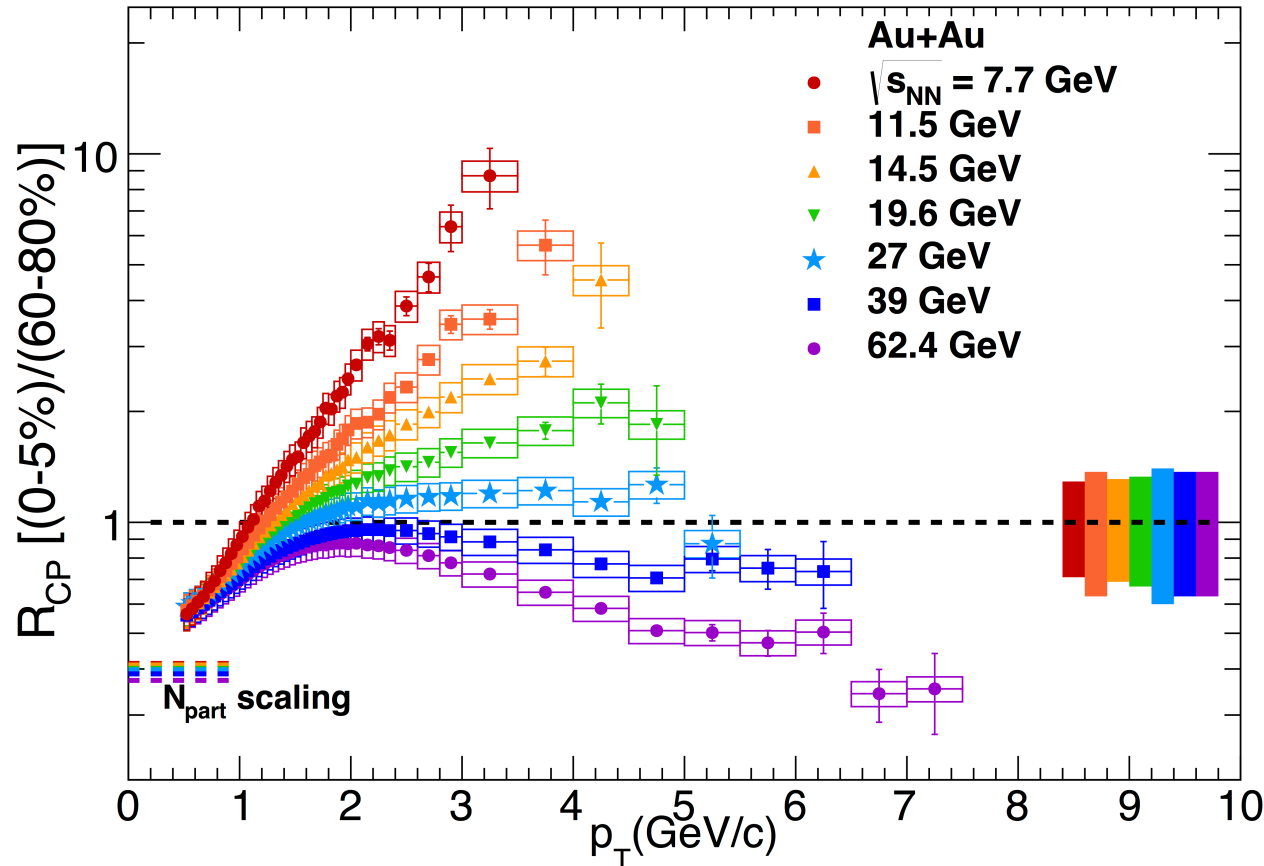
Smooth transition from suppression ($R_{CP} < 1$) at highest energies to enhancement ($R_{CP} > 1$) at lowest energies

Turn off of the QGP?

$$R_{CP} = \frac{\langle N_{coll} \rangle_{\text{Peripheral}}}{\langle N_{coll} \rangle_{\text{Central}}} \frac{\left(\frac{d^2 N}{dp_T d\eta} \right)_{\text{Central}}}{\left(\frac{d^2 N}{dp_T d\eta} \right)_{\text{Peripheral}}}$$

Nuclear Modification Factor in BES 1

69

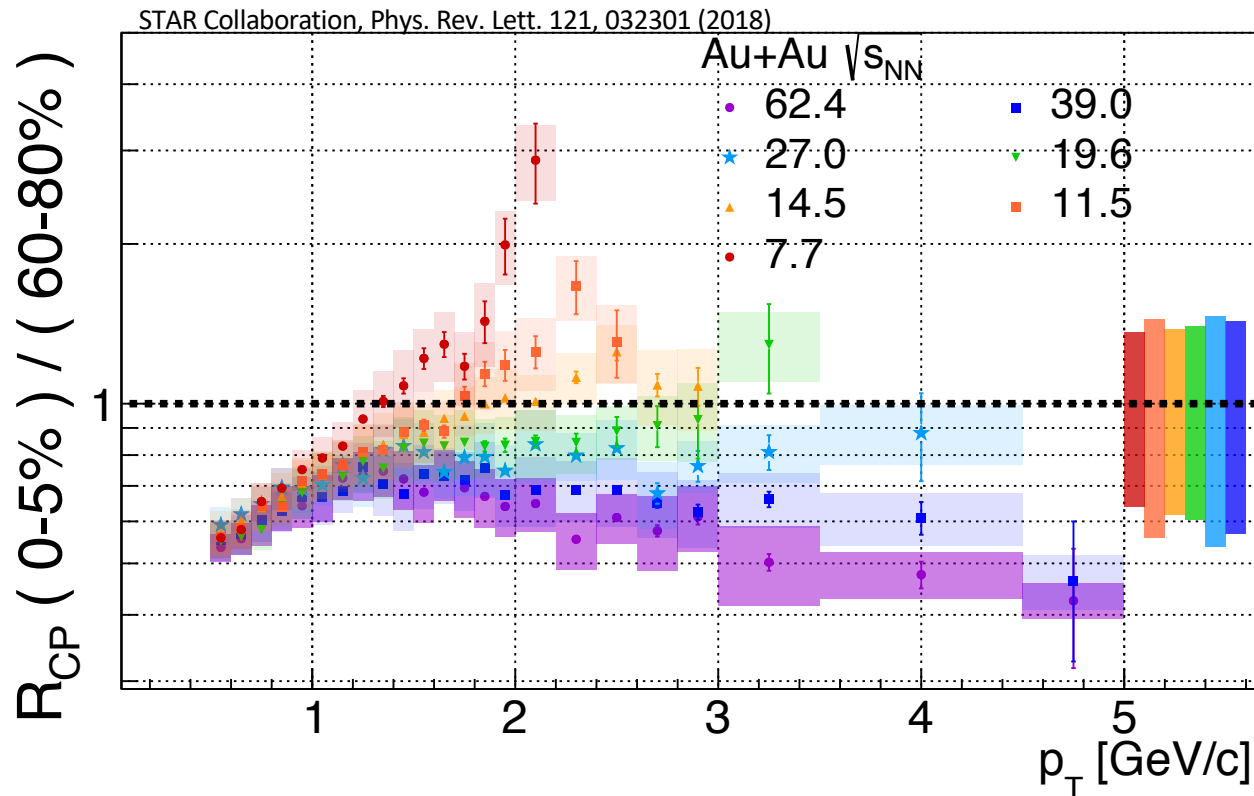


What can we say from this measurement?

- $R_{CP} \ll 1$ (at high p_T) : QGP is definitely formed
- If $R_{CP} \geq 1$: QGP may be formed or may not - competing effects

R_{CP} for Identified Hadrons

70



- Pions are more sensitive probes of suppression than heavy particles
- No sign of suppression above ~ 19.6 GeV