# Chemical freeze-out conditions from net-kaon fluctuations at RHIC

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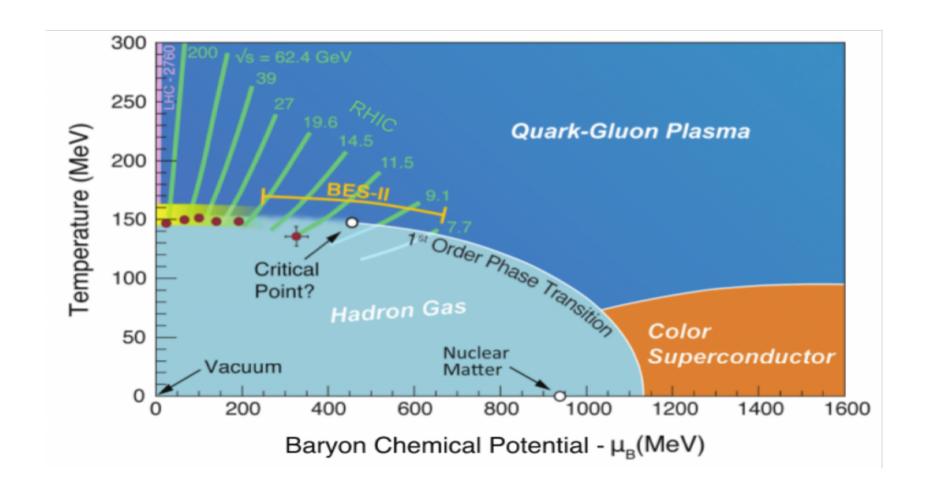


## QCD Phase Diagram



The different phases of QCD matter can be understood by studying the characteristics of the phase diagram

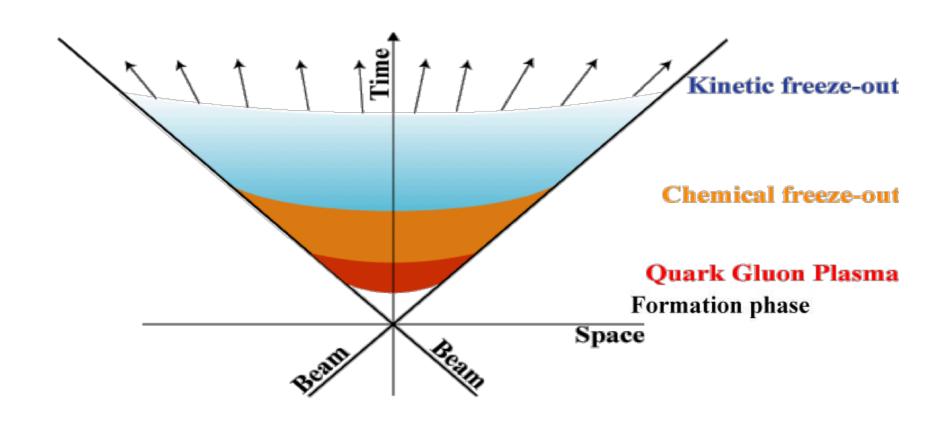
- ▶ QGP is formed at large T,  $\mu_B$ ; ordinary hadronic matter at small T,  $\mu_B$
- Crossover transition at T ~ 155MeV; possible first order phase transition
  - Search for the critical point with the Beam Energy Scan (BES)



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## Evolution of a heavy-ion collision



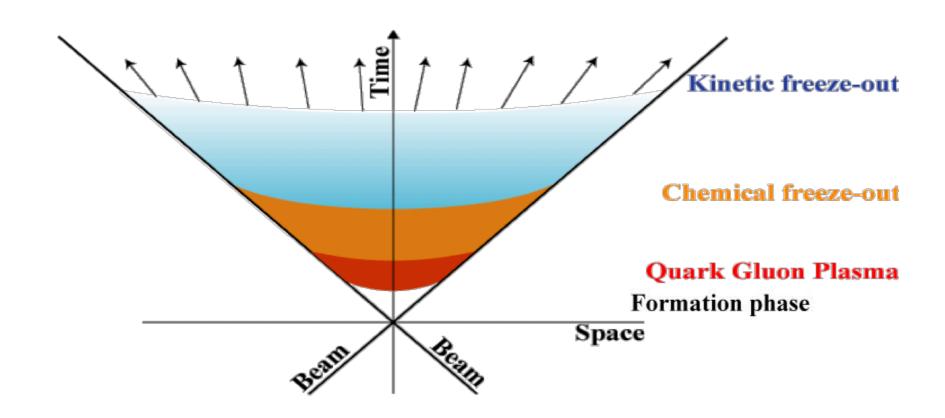


- Chemical freeze-out: inelastic collisions cease; the chemical composition is fixed (yields and fluctuations)
- Kinetic freeze-out: elastic collisions cease; spectra and correlations are fixed

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## Evolution of a heavy-ion collision





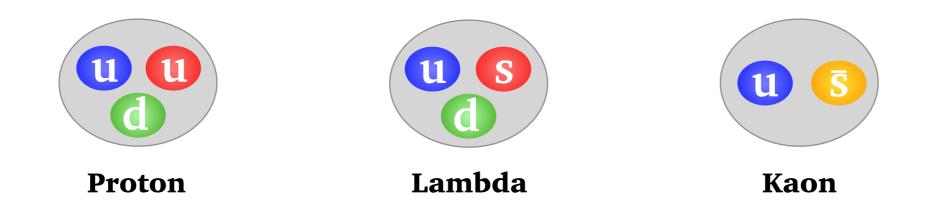
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## Chemical freeze-out in heavy-ion collisions



Is there a **flavor hierarchy** in the chemical freeze-out temperature?



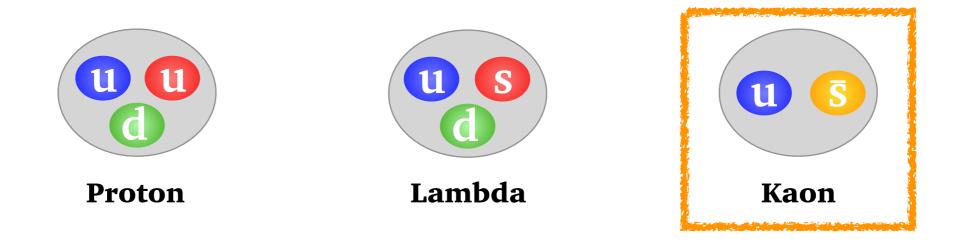
Utilize Hadron Resonance Gas (HRG) model to directly compare with experiment and extract chemical freeze-out parameters

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## Chemical freeze-out in heavy-ion collisions



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Utilize Hadron Resonance Gas (HRG) model to directly compare with experiment and extract chemical freeze-out parameters

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#### Hadron Resonance Gas Model



In the low temperature regime, the system is well-described by a gas of hadrons:

- Grand Canonical Ensemble
- Interacting gas of ground-state hadrons
  - Treat as non-interacting system of resonant states
- Adaptable to match experimental conditions
- List of particles from the Particle Data Group (PDG)

$$\frac{P}{T^4} = \frac{1}{VT^3} \sum_{i} \ln Z_i(T, V, \vec{\mu})$$

where: 
$$\ln Z_i^{M/B} = \mp \frac{V d_i}{(2\pi)^3} \int d^3p \, \ln(1 \mp \exp\left[-\left(\epsilon_i - \mu_a X_a^i\right)/T\right])$$

energy 
$$\epsilon_i = \sqrt{p^2 + m_i^2}$$

conserved charges  $\vec{X}_i = (B_i, S_i, Q_i)$ 

degeneracy  $d_i$ , mass  $m_i$ , volume V

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## Hadron Resonance Gas Model



hadron	$m_{\it i}$ (GeV)	$d_i$	$B_i$	$s_i$	$I_i$	hadron	$m_{\it i}$ (GeV)	$d_i$	$B_i$	$s_i$	$I_i$
$\pi$	0.140	3	0	0	1	N (1535)	1.530	4	1	0	1/2
K	0.496	2	0	1	1/2	π <sub>1</sub> (1600)	1.596	9	0	0	1
$\overline{K}$	0.496	2	0	-1	1/2	Δ (1600)	1.600	16	1	0	3/2
η	0.543	1	0	0	0	Λ (1600)	1.600	2	1	-1	0
ρ	0.776	9	0	0	1	△ (1620)	1.630	8	1	0	3/2
$\omega$	0.782	3	0	0	0	$\eta_2$ (1645)	1.617	5	0	0	0
$K^*$	0.892	6	0	1	1/2	N (1650)	1.655	4	1	0	1/2
$\overline{K}^*$	0.892	6	0	-1	1/2	ω (1650)	1.670	3	0	0	0
N	0.939	4	1	0	1/2	Σ (1660)	1.660	6	1	-1	1
$\eta'$	0.958	1	0	0	0	Λ (1670)	1.670	2	1	-1	0
$f_0$	0.980	1	0	0	0	Σ (1670)	1.670	2	1	-1	1
$a_0$	0.980	3	0	0	1	ω <sub>3</sub> (1670)	1.667	7	0	0	0
$\phi$	1.020	3	0	0	0	π <sub>2</sub> (1670)	1.672	15	0	0	1
Λ	1.116	2	1	-1	0	$\Omega^-$	1.672	4	1	-3	0
$h_1$	1.170	3	0	0	1	N (1675)	1.675	12	1	0	1/2
$\Sigma$	1.189	6	1	-1	1	φ (1680)	1.680	3	0	0	0
$a_1$	1.230	9	0	0	1	K* (1680)	1.717	6	0	1	1/2
$b_1$	1.230	9	0	0	1	K * (1680)	1.717	6	0	-1	1/2
Δ	1.232	16	1	0	3/2	N (1680)	1.685	12	1	0	1/2
$f_2$	1.270	5	0	0	0	ρ <sub>3</sub> (1690)	1.688	21	0	0	1
$K_1$	1.273	6	0	1	1/2	Λ (1690)	1.690	4	1	-1	0
$\overline{K}_1$	1.273	6	0	-1	1/2	三(1690)	1.690	8	1	-2	1/2
$f_1$	1.285	3	0	0	1	ρ (1700)	1.720	9	0	0	1
$\eta$ (1295)	1.295	1	0	0	0	N (1700)	1.700	8	1	0	1/2
$\pi$ (1300)	1.300	3	0	0	1	$\Delta(1700)$	1.700	16	1	0	3/2

List made by C. Ratti

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## Fluctuations of Conserved Charges



$$\chi_{lmn}^{BSQ} = \frac{\partial^{\,l+m+n} p/T^4}{\partial (\mu_B/T)^l \partial (\mu_S/T)^m \partial (\mu_Q/T)^n}.$$

mean :  $M = \chi_1$ 

variance :  $\sigma^2 = \chi_2$ 

skewness :  $S = \chi_3/\chi_2^{3/2}$ 

kurtosis :  $\kappa = \chi_4/\chi_2^2$ 

The volume-independent ratios are:

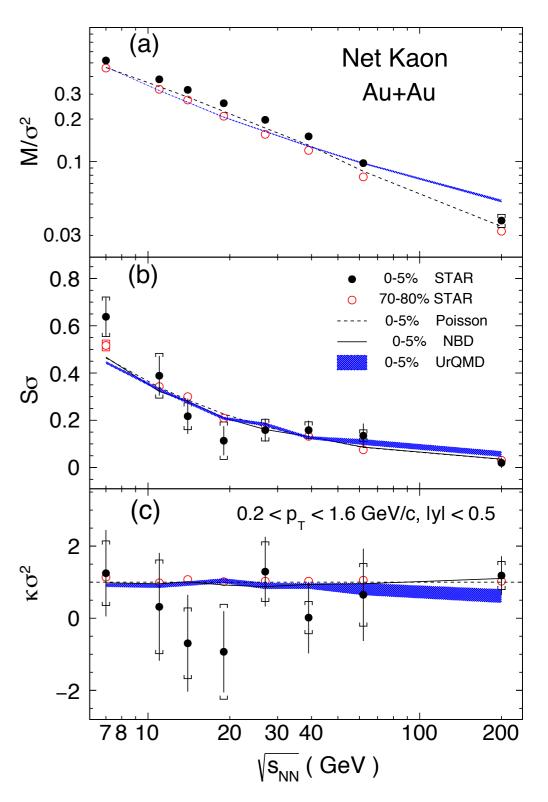
$$S\sigma = \chi_3/\chi_2$$

$$\kappa \sigma^2 = \chi_4/\chi_2$$

$$M/\sigma^2 = \chi_1/\chi_2$$

$$S\sigma^3/M = \chi_3/\chi_1$$

Directly compare HRG Model to experiment to identify chemical freeze-out conditions!



STAR Collaboration (Adamczyk, L. et al.) Phys.Lett. B 785 (2018) 551-560 arXiv:1709.00773 [nucl-ex]

## Fluctuations of Conserved Charges



Kaon susceptibilities in the HRG Model:

$$\chi_{n}^{\text{net-}K} = \sum_{i \in \text{HRG}} \frac{(Pr_{i \to \text{net-K}})^{n}}{T^{3-(n-1)}} \frac{S_{i}^{1-n}d_{i}}{4\pi^{2}} \frac{\partial^{n-1}}{\partial \mu_{S}^{n-1}} \times \left\{ \int_{-0.5}^{0.5} dy \int_{0.2}^{1.6} dp_{T} \frac{p_{T}\sqrt{p_{T}^{2} + m_{i}^{2}} \operatorname{Cosh}[y]}{(-1)^{B_{i}+1} + \exp\left((\operatorname{Cosh}[y]\sqrt{p_{T}^{2} + m_{i}^{2}} - (B_{i}\mu_{B} + S_{i}\mu_{S} + Q_{i}\mu_{Q})\right)/T)} \right\}$$

- ☑ Limits of integration correspond to the acceptance cuts from the experiment
- oxdots Strangeness neutrality imposed on  $\mu_B$ ,  $\mu_Q$ ,  $\mu_S$  matches the experimental conditions:

$$\langle n_s \rangle = 0$$
  $\langle n_O \rangle = 0.4 \langle n_B \rangle$ 

☑ Include feed-down from resonances by utilizing the branching ratios:

$$Pr_{i \to K_{net}} = Br_{i \to K_{net}} n_i(K_{net})$$

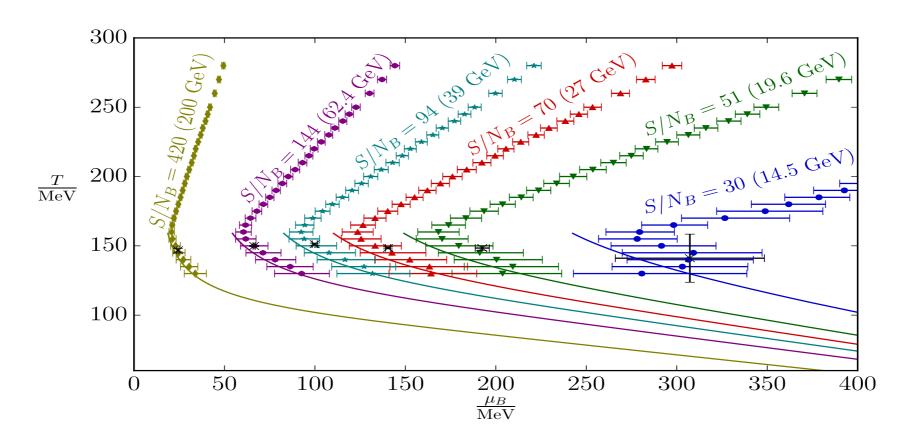
R. Bellwied, JS, et al., Phys. Rev. C99 (2019) 034912

## Lattice QCD Isentropes



In order to extract the  $\{T_f, \mu_{B,f}\}$ , we utilize the isentropic trajectories from Lattice QCD

- ► S/N<sub>B</sub> is conserved
- ► Shows the path of the system across the phase diagram



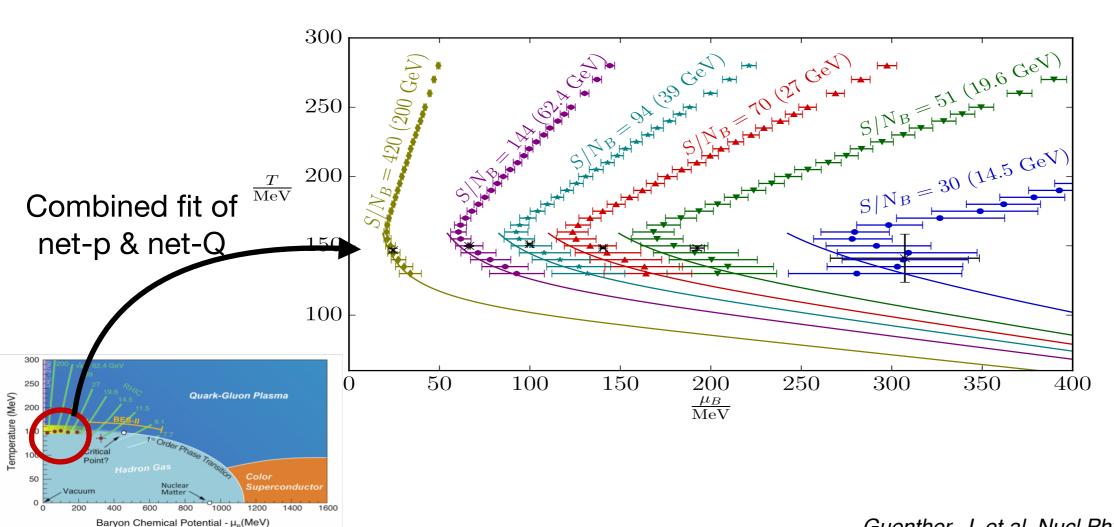
Guenther, J. et al. Nucl. Phys. A967 (2017) 720-723

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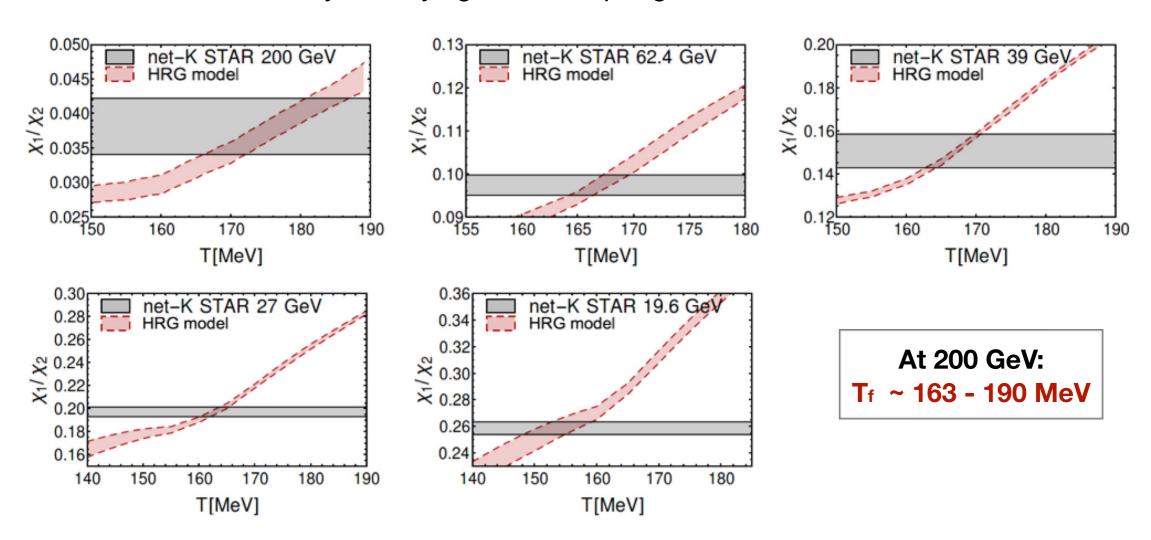
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#### Kaon fluctuations



Calculate  $\chi_1/\chi_2$  along the isentropes corresponding to the five highest energies of the Beam Energy Scan at RHIC

Extract Tf by identifying the overlap regions



R. Bellwied, JS, et al., Phys. Rev. C99 (2019) 034912

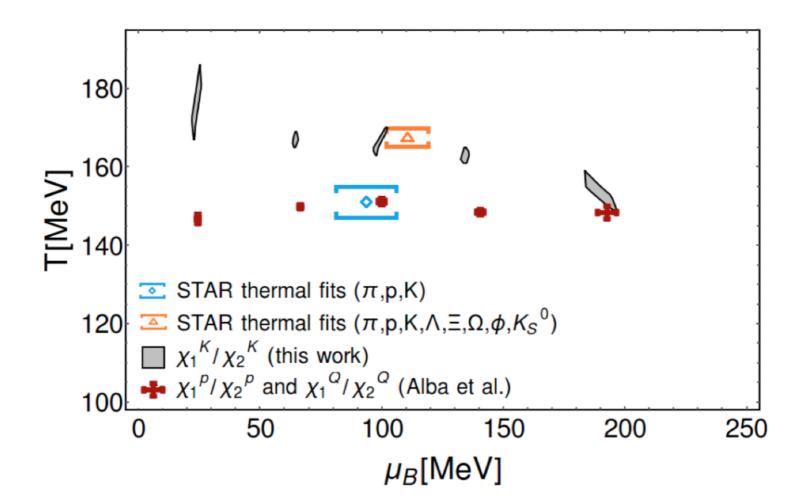
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## Freeze-out parameters for net-kaons



#### Compare the freeze-out parameters for net-K to:

- light freeze-out (combined fit of net-proton and net-electric charge)
- thermal fits from the experiment

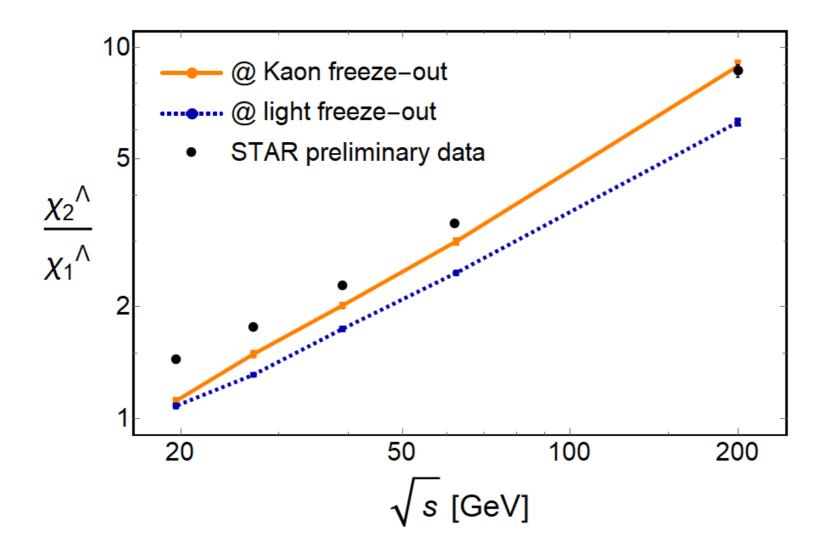


R. Bellwied, JS, et al., Phys. Rev. C99 (2019) 034912

## Lambda predictions



Calculate fluctuations for net- $\Lambda$  using the kaon and light hadron freezeout parameters and compare with preliminary experimental data



R. Bellwied, JS, et al., Phys. Rev. C99 (2019) 034912

#### Freeze-out with different PDG lists

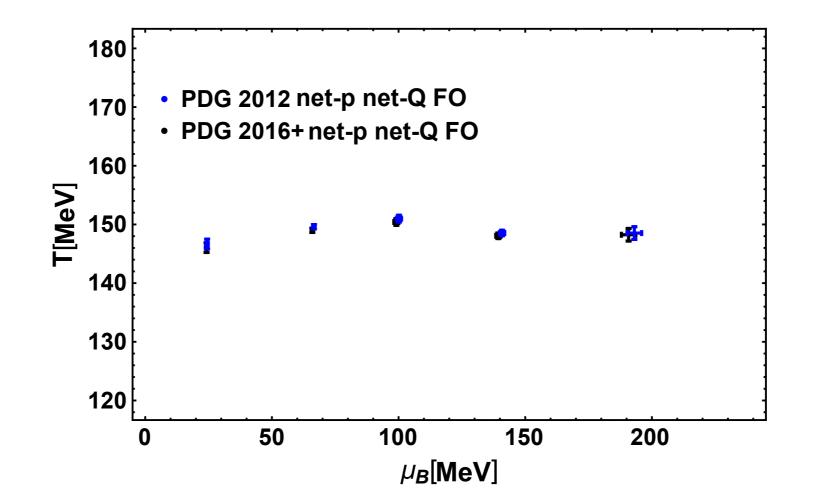


Recall the pressure in the HRG Model:  $\frac{P}{T^4} = \frac{1}{VT^3} \sum_i \ln Z_i(T,V,\vec{\mu})$ 

Different particle lists will yield different results for the freeze-out parameters

PDG2012: 524 particles

PDG2016+: 739 particles



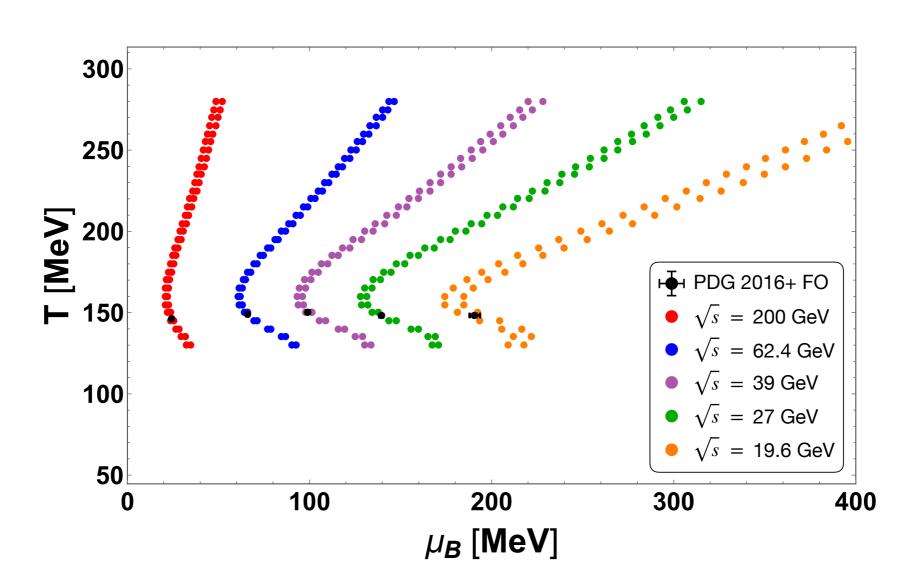
## Lattice QCD Isentropes with 2016+



Calculate new isentropes from freeze-out parameters determined with PDG 2016+

$$\frac{s}{T^3} = \left. \frac{1}{T^3} \frac{\partial p}{\partial T} \right|_{\mu_i}$$

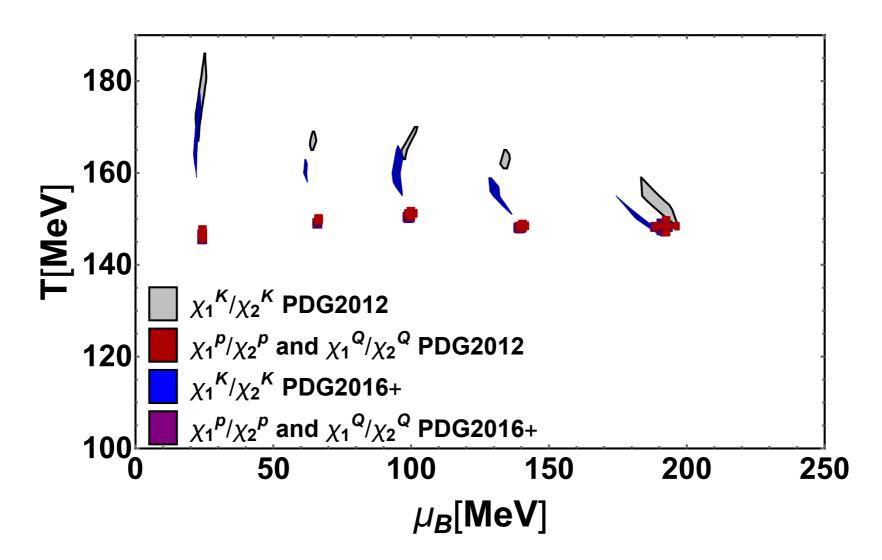
$$\frac{n_i}{T^3} = \frac{1}{T^3} \left. \left( \frac{\partial p}{\partial \mu_i} \right) \right|_{T,\mu_i}$$



### Kaon freeze-out for PDG2012 and PDG2016+



Compare the freeze-out parameters for the kaons and light particles for the different lists in order to determine the effect of the number of resonant states:



Even with the inclusion of more states in the HRG Model, there is evidence of a flavor hierarchy

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#### Conclusions



• The net-kaon fluctuation data from the STAR collaboration cannot be reproduced in the HRG model by using the freeze-out parameters obtained from the combined fit of  $\chi^p_{\ 1}/\chi^p_{\ 2}$  and  $\chi^Q_{\ 1}/\chi^Q_{\ 2}$ .

 At the highest collision energy, the kaons freeze-out above T=163 MeV, about 10-15 MeV higher than the light hadrons.

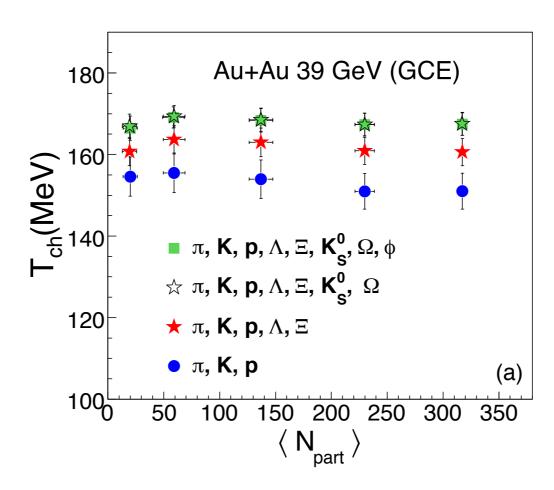
 Confirmation of a freeze-out flavor-hierarchy by the inclusion of more resonances in the HRG Model.

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## Back-up slides

## Is there a flavor hierarchy for chemical freeze-out?





- Chemical freeze-out: inelastic collisions cease; the chemical composition is fixed (yields and fluctuations)
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