



Axial nucleon form factor from a Bayesian analysis of the neutrino scattering data ISNET-6

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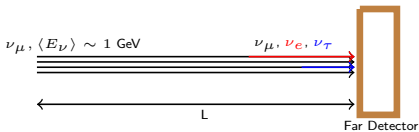


Motivation

- ▶ Neutrinos $\nu_e, \nu_\mu, \nu_\tau, \dots$, fundamental particles
- ▶ weakly interacting particles, neutral
- ▶ neutrino oscillations:
 - ▶ neutrino oscillation parameters, mass differences and mixing angles
 - ▶ CP -violation phase
 - ▶ mass hierarchy
- ▶ ...

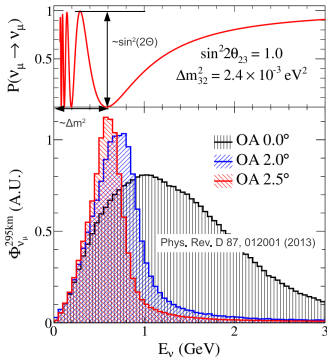
A need of better knowledge of ν -nuclei, ν -nucleon scattering cross sections (in particular for $E_\nu \sim 1$ GeV)

ν -A Interactions



$$P(\nu_\mu \rightarrow \nu_\tau) = \sin^2 \theta_{23} \left(\frac{\Delta m_{32}^2 L}{4 E_\nu} \right)$$

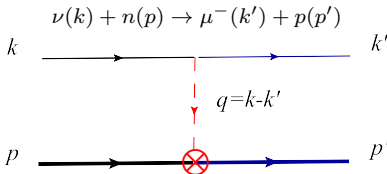
- Neutrino energy reconstructed from the analysis of QuasiElastic (QE) scattering events!



- accelerator source of neutrinos
→ $E_\nu \sim 1 \text{ GeV}$ but neutrino energy not monochromatic: E_ν known with some precision!

Axial Nucleon Form Factor

- Consider **Charged Current Quasielastic (CCQE)**



$$i\mathcal{M}_{Born} \approx i \frac{G_F \cos \theta_C}{\sqrt{2}} \cdot \bar{u}(k') \gamma^\mu (1 - \gamma_5) u(k) \times \bar{u}(p') \Gamma^\mu(q) u(p)$$

- q^μ - energy-momentum transfer
- Γ^μ has a V^μ -ector – A^μ -xial structure



- ▶ Vector part expressed in terms electromagnetic nucleon form factors

- ▶ *Axial*:

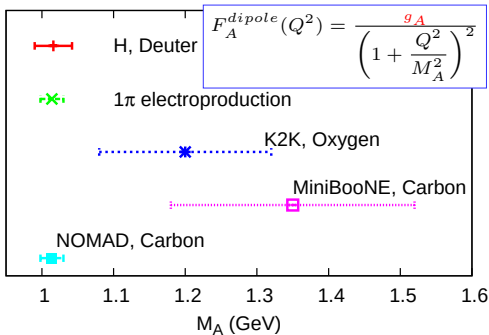
$$A^\mu = \left(\gamma_\mu \gamma_5 + \frac{q^\mu \gamma_5}{2M} \frac{4M^2}{m_\pi^2 + Q^2} \right) F_A(Q^2)$$

where

$$Q^2 = -q^\mu q_\mu \equiv -q^2 > 0$$

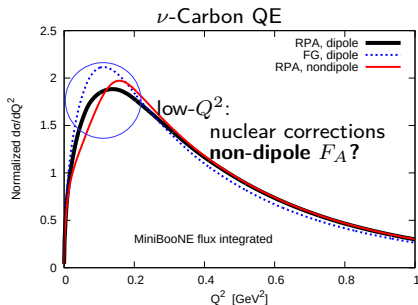
- ▶ Neutron β -decay asymmetry $F_A(0) = g_A = 1.2723 \pm 0.0023$. C. Patrignani et al., PDG, Chin. Phys. C40, 100001 (2016)
- ▶ F_A an input for the ν -nucleus scattering models
- ▶ ν -nucleon and ν -nucleus cross sections dominated by the axial contribution

Dipole Axial Nucleon Form Factor



The most reliable information about F_A :
bubble chamber data i.e. $\nu - d$ scattering
(ANL, BNL, and FNAL, CERN data)

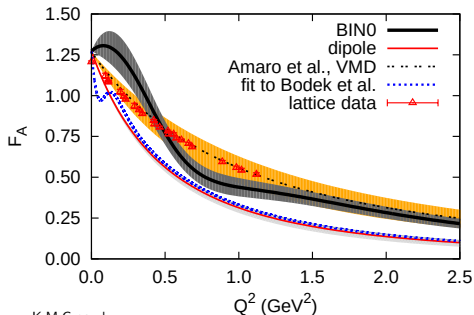
Recent measurements: ν -Carbon,
 ν -Oxygen, ... scattering: extraction of F_A
is nuclear model dependent!



recent studies: Bodek, et al.
EPJC53, 349 (duality
constrained); Adamuscin et al.
PRC78 035201 (monopole);
Amaro et al. PRD93, 053002
(VMD, two monopoles); Meyer et
al., PRD93, 113015 (z-expansion)

Non-dipole dependence of F_A ?

Electric and magnetic form factors of
the nucleon have non-dipole Q^2 -shape!



Goal

- ▶ analyze the bubble chamber ν -deuteron scattering data trying to find deviations of F_A from dipole shape
- ▶ obtain predictions of F_A with uncertainties

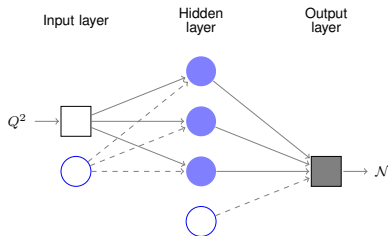
- ▶ F_A obtained from the data
- ▶ consider the class of flexible functions $\rightarrow$ feed-forward neural networks

$$F_A(Q^2) = F_A^{\text{dipole}}(Q^2) \times \mathcal{N}_M(Q^2; \{w_i\})$$

- ▶ choose the "best" fit using Bayesian algorithm
- * (used to extract G_E , G_M and 2γ correction from elastic ep scattering data)

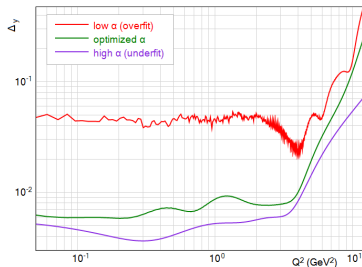
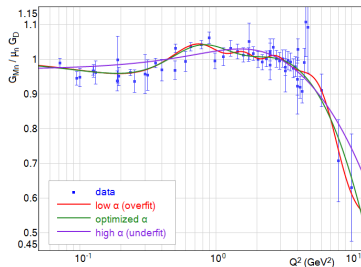
Main assumptions:

- (I) $F_A(Q^2)$ is assumed to be a continuous function of Q^2 in its validity domain
- (II) As $F_A(Q^2)$ is bounded, there must be a $C > 1$:
 $F_A(Q^2)/F_A^{\text{dipole}}(Q^2) < C$ in the whole Q^2 interval of $(0, 3)$ GeV²



Bias-variance trade off

- ▶ too complex fit reproduces the statistical noise and overfits the data
→ poor generalization
- ▶ too simple fit underfits the data
- * unrealistic uncertainties





Bayesian statistics embodies Occam's razor

Bayesian Framework for neural networks: MacKay, Neural Computation 4 (3), (1992) 415; Neural Computation 4 (3), (1992) 448

- ▶ consider function $\mathcal{N}$ (depending on W weights $\{w_i\}$)

$$P(\{w_i\}|\mathcal{D}, \mathcal{N}) = \frac{\overbrace{P(\mathcal{D}|\{w_i\}, \mathcal{N})}^{\text{likelihood}} \overbrace{P(\{w_i\}|\mathcal{N})}^{\text{prior}}}{\underbrace{P(\mathcal{D}|\mathcal{N})}_{\text{evidence}}}$$

- ▶ Evidence for model $\mathcal{N}$

$$P(\mathcal{D}|\mathcal{N}) = \int \prod_k dw_k P(\mathcal{D}|\{w_i\}, \mathcal{N}) P(\{w_i\}|\mathcal{N})$$

▶

$$P(\mathcal{N}|\mathcal{D}) = \frac{P(\mathcal{D}|\mathcal{N})P(\mathcal{N})}{P(\mathcal{D})}.$$

$P(\mathcal{D})$ - normalization

- ▶ Assume that $P(\mathcal{N}_1) = P(\mathcal{N}_2) = \dots$ then $P(\mathcal{D}|\mathcal{N})$ ranks models...

- ▶ Let consider the N independent measurements y_i with uncertainty σ_i
 $\mathcal{D} = \{(y_1(x_1), \sigma_1), \dots\}$
- ▶ likelihood (of the Gaussian form)

$$2 \ln \mathcal{P}(\mathcal{D} | \{w_i\}, \mathcal{N}) \sim -\chi_{ex}^2(\mathcal{D}, \{w_i\}), \quad \chi_{ex}^2(\mathcal{D}, \mathcal{N}) = \sum_{i=1}^N \frac{(y_i(x_i) - f(\mathcal{N}))^2}{\sigma_i^2}$$

- ▶ prior

$$\mathcal{P}(\{w_i\} | \alpha, \mathcal{N}) \sim \exp\left(-\frac{\alpha}{2} E_w\right), \quad E_w = \sum_{k=1}^W w_k^2$$

α regularizer (hyperparameter), W - number of weights (neural network parameters)

- ▶ Posterior

$$2 \ln P(\{w_i\} | \mathcal{D}, \mathcal{N}, \alpha) \sim -(\chi_{ex}^2(\mathcal{D}, \{w_i\}) + \alpha E_w) = -\mathcal{E}(\mathcal{D}, \{w_i\})$$

- ▶ $\mathcal{E}(\mathcal{D}, \{w_j\})$ has a minimum at $\{w_j\}_{MP}$
- ▶ α : overfit, underfit $\rightarrow$ optimal model



- Analytic calculations (in Gaussian approximation)

$$\mathcal{P}(\{w_i\} | \mathcal{D}, \alpha, \mathcal{N}) = \frac{\mathcal{P}(\mathcal{D} | \{w_i\}, \alpha, \mathcal{N}) \mathcal{P}(\{w_i\} | \alpha, \mathcal{N})}{\mathcal{P}(\mathcal{D} | \alpha, \mathcal{N})} \rightarrow$$

$$\mathcal{P}(\alpha | \mathcal{D}, \mathcal{N}) = \frac{\mathcal{P}(\mathcal{D} | \alpha, \mathcal{N}) \mathcal{P}(\alpha | \mathcal{N})}{\mathcal{P}(\mathcal{D} | \mathcal{N})} \rightarrow$$

$$\mathcal{P}(\mathcal{N} | \mathcal{D}) = \frac{\mathcal{P}(\mathcal{D} | \mathcal{N}) \mathcal{P}(\mathcal{N})}{\mathcal{P}(\mathcal{D})}$$

$$2\alpha_{MP} E_w = \sum_{i=1}^W \frac{\lambda_i}{\lambda_i + \alpha_{MP}} \equiv \gamma(\alpha_{MP})$$

λ_i - eigenvalues of $H_{kj} = \nabla_k \nabla_j \chi_{ex}$

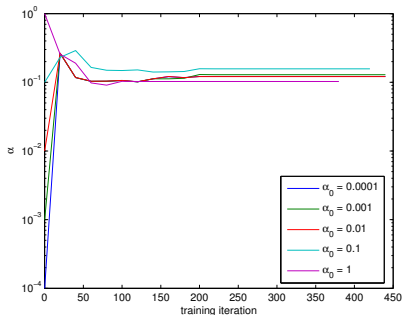
→ α_{MP} : α is iterated according to (12)

$$\alpha_{k+1} = \gamma(\alpha_k) / 2E_w$$

- Evidence approximation →

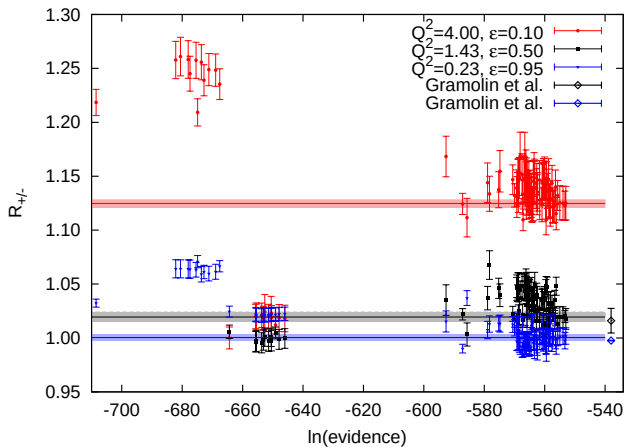
$$\begin{aligned} \ln \mathcal{P}(\mathcal{D} | \mathcal{N}) &\approx -\mathcal{E}(\mathcal{D}, \{w_j\}_{MP})/2 - \frac{1}{2} \ln |A| \\ &+ \frac{W}{2} \ln \alpha_{MP} - \frac{1}{2} \ln \frac{\gamma_{MP}}{2} + \text{combinatorial} \end{aligned}$$

where $A = \nabla_i \nabla_j S(\mathcal{D}, \{w_j\})$





elastic ep scattering data and 2γ exchange correction





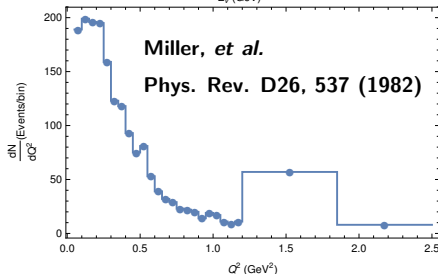
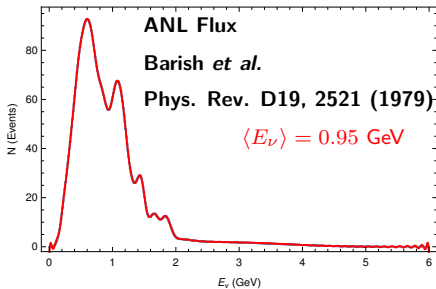
Bubble Chamber Data: $\nu_\mu + d \rightarrow \mu^- + p + p$

- ▶ **ANL** data: Argonne National Laboratory data (today's talk)

- * Distribution of events:
 $Q^2 \in (0.05, 2.5) \text{ GeV}^2$,

- **BNL**: Brookhaven National Laboratory, $\langle E_\nu \rangle \sim 1.56 \text{ GeV}$,
 $Q^2 \in (0.03, 2.8) \text{ GeV}^2$

- **FNAL** data:
 $\langle E_\nu \rangle \sim 22.6 \text{ GeV}$,
 $Q^2 \in (0., 3) \text{ GeV}^2$





- The least square-function:

$$\chi_{ex}^2 = \chi_{ANL}^2 + \chi_{g_A}^2$$

where

$$\chi_{ANL}^2 = \sum_{i=k}^{n_{ANL}} \frac{(N_i - pN_i^{th})^2}{N_i} + \left(\frac{1-p}{\Delta p} \right)^2$$

see D'Agostini, NIMA346, 306, Roe, arXiv:1506.09077, Baker and Cousins, NIM221 437

- Δp is a systematic uncertainty for $\#N$ of events
-

$$N_i^{th} = \int_0^\infty dE_\nu \frac{\frac{d\sigma}{dQ^2}(E_\nu, F_A)}{\sigma(E_\nu, F_A)} \frac{dN}{dE_\nu}, \quad \sigma(E_\nu, F_A) = \int_{min}^{max} \frac{d\sigma}{dQ^2}(E_\nu, F_A) dQ^2$$

see Meyer et al. (2017)

- $F_A(Q^2 = 0)$,

$$\chi_{g_A}^2 = \left(\frac{F_A(0) - g_A}{\Delta g_A} \right)^2$$

g_A and Δg_A from PDG

Model of $\nu_\mu + d \rightarrow \mu^- + p + p$

$$\frac{d\sigma}{dQ^2}(\nu_\mu n \rightarrow \mu^- p) = \frac{G_F^2 M}{8\pi E_\nu^2} \times [A(Q^2, E_\nu) F_A^2 + B(Q^2, E_\nu) F_A + C(Q^2, E_\nu)]$$

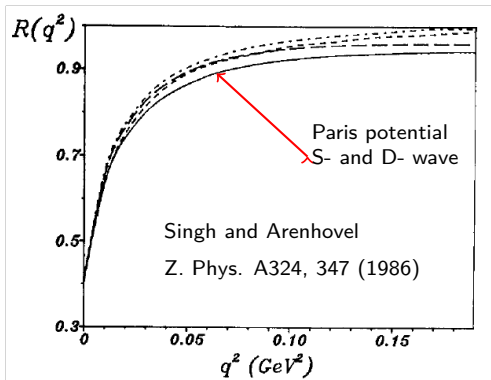
- Deuteron corrections:

$$\frac{d\sigma}{dQ^2} \rightarrow R(Q^2) \frac{d\sigma}{dQ^2}$$

where

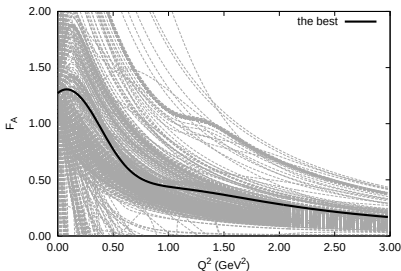
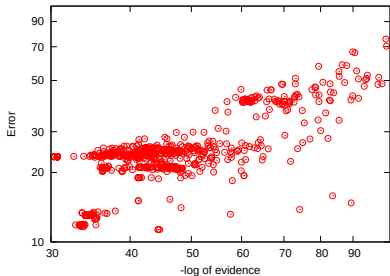
$$R(Q^2) = \frac{\frac{d\sigma}{dQ^2}(\nu_\mu d \rightarrow \mu^- pp)}{\frac{d\sigma}{dQ^2}(\nu_\mu n \rightarrow \mu^- p)} \quad (1)$$

- correction concerns first three bins of ANL data!



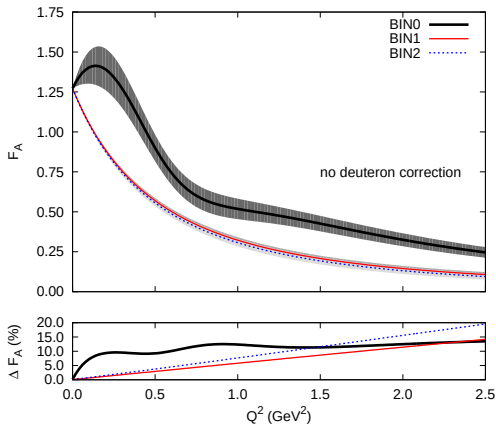
Analysis

- (i) consider network with $M = 1$;
- (ii) perform the network training;
- (iii) repeat steps (i)-(iii) for various $\{w_j\}_0$;
- (iv) choose the best fit according to the evidence;
- (v) repeat steps (i)-(iv) for $M = 2, \dots$;
- (vi) among the all the best fits choose the model with the highest evidence

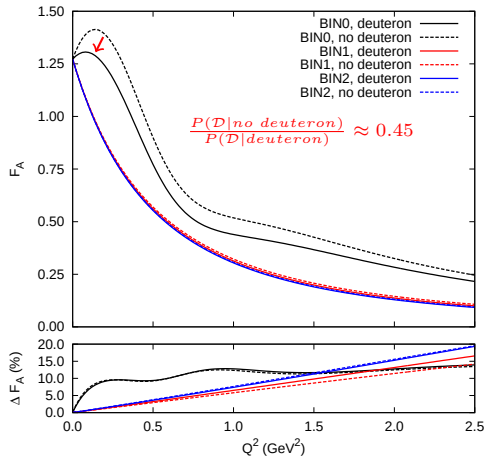


Analysis

- (i) BIN0: all ANL bins included
 - (ii) BIN k : where $k = 1$ or $k = 2$: ANL bins without the first k bins
- ▶ Unexpected F_A fit for all bins (BIN0)?!
 - ▶ The slope of F_A at $Q^2 = 0$ not consistent with other determinations...
 - ▶ fits of BIN1 and BIN2 data consistent with original ANL analysis w dipole par. (Miller et al.)!
 - ▶ similar result for BNL data, but not for FNAL data!(preliminary)



Impact of deuteron corrections

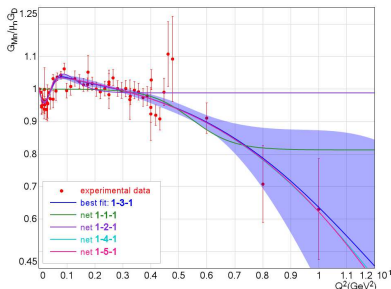


- ▶ problem with the first bin?
- ▶ deuteron correction important for the first bin

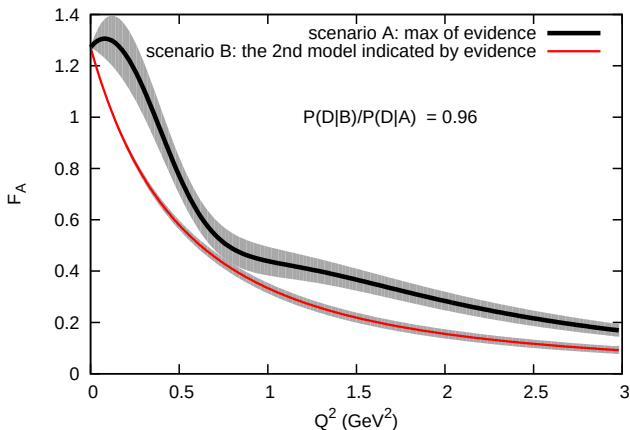
Problem with 1st bin?

Possible explanation:

- ▶ a low quality of the measurements at low- Q^2 due to low and not well understood efficiency
- ▶ an improper description of the nuclear corrections
- ▶ the actual value of the slope $dF_A(Q^2 = 0)/dQ^2$ might not be properly estimated because of the lack of very low- Q^2 data



two competitive scenarios: on the edge of dipole...



- ▶ small effects play a role...
- ▶ there is a tension between the first bin and constraint for F_A at $Q^2 = 0$
- ▶ assumption: $dF_A(Q^2 = 0)/dQ^2 < 0 \rightarrow$ the fit in agreement with dipole shape...



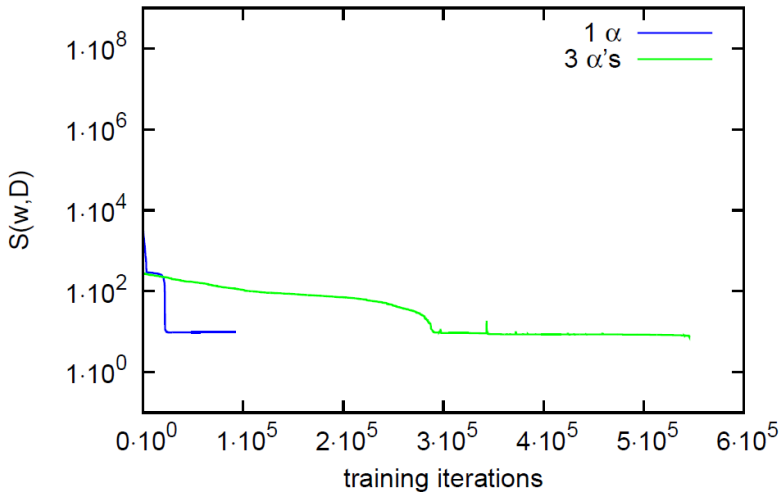
Conclusions

- ▶ The ANL data are not enough informative to study deviation of F_A from the dipole shape
- ▶ The fits of BIN1 and BIN2 in agreement with dipole fits
- ▶ Future: global analysis of ANL, BNL and FNAL data
- ▶ New measurements of ν -H, ν -D necessary

THIS IS THE END



GEN: 121



least-square vs. Poisson

	Fit	N_{fit}	$M_A [MeV]$	$r_A^2 [fm^2]$	bins
deuteron -2LL	52.7	1.11 ± 0.03	1032 ± 60	0.438 ± 0.051	49
deuteron χ^2	27.6	1.09 ± 0.03	1042 ± 60	0.430 ± 0.050	25

$$-2LL = 2 \sum_i^N \left[N_i^{th} - N_i + N_i \ln \left(\frac{N_i}{N_i^{th}} \right) \right] \quad (2)$$