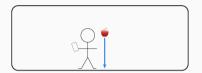
# Quantifying Errors from Chiral Effective Field Theory

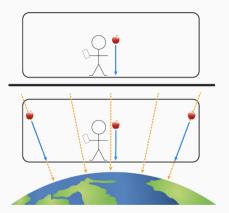
Iordan Melendez<sup>1</sup> Dick Eurnstahl<sup>1</sup> Daniel Phillips<sup>2</sup> Matt Pratola<sup>1</sup> Sarah Wesolowski<sup>3</sup> October 11. 2018 Posterior THE OHIO STATE UNIVERSITY True value <sup>1</sup>The Ohio State University an <sup>2</sup>Ohio University **BUQEYE** Collaboration Nuclear Computational Low-Energy Initiative <sup>3</sup>Salisbury University U.S. DEPARTMENT OF NSF ENERGY

Based on: • S. Wesolowski, R. J. Furnstahl, J. A. Melendez, D. Phillips, arXiv:1808.08211; & J. A. Melendez, S. Wesolowski, and R. J. Furnstahl Phys. Rev. C **96**, 024003, Editors' Suggestion

# **Physical Motivation**



F = -mg

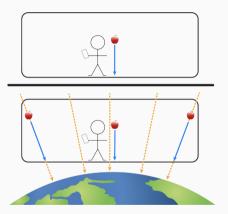


F = -mg

Grav. force (large distances):

 $F = -\frac{GMm}{r^2}$ 

The laws look quite different!



F = -mg

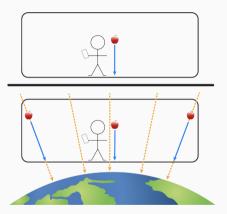
Grav. force (large distances):



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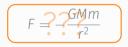
Connected via series expansion about radius of Earth R:

$$F \approx -mg + 2mg\left(\frac{r-R}{R}\right) - 3mg\left(\frac{r-R}{R}\right)^2 + \mathcal{O}\left[\left(\frac{r-R}{R}\right)^3\right]$$



F = -mg

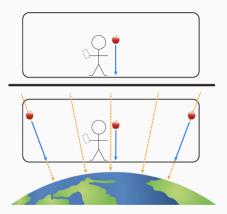
#### Grav. force (large distances):



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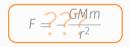
Can fit unknown parameters to data  $\Rightarrow$  inverse problem!

$$F \approx a_0 + a_1 \left(\frac{r-R}{R}\right) + a_2 \left(\frac{r-R}{R}\right)^2 + \mathcal{O}\left[\left(\frac{r-R}{R}\right)^3\right]$$



F = -mg

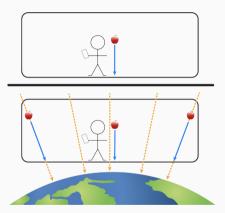
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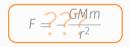
Use prior info from physics:

$$F \approx mg\left\{a_0' + a_1'\left(\frac{r-R}{R}\right) + a_2'\left(\frac{r-R}{R}\right)^2 + \mathcal{O}\left[\left(\frac{r-R}{R}\right)^3\right]\right\}$$



F = -mg

#### Grav. force (large distances):

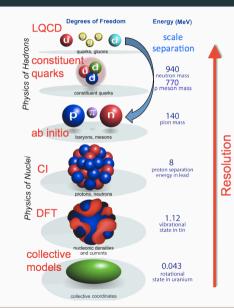


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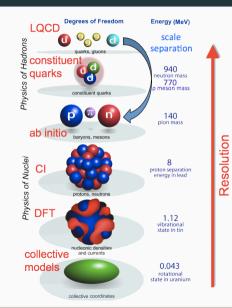
Propagate full uncertainty

$$F \approx mg \left\{ a'_0 + a'_1 \left( \frac{r-R}{R} \right) + a'_2 \left( \frac{r-R}{R} \right)^2 + \mathcal{O} \left[ \left( \frac{r-R}{R} \right)^3 \right] \right\}$$

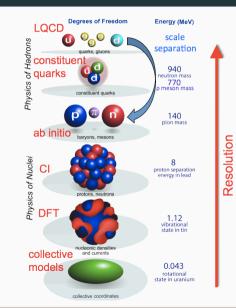
• There is interesting physics at all scales



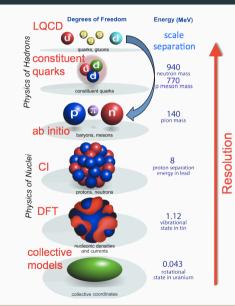
- There is interesting physics at all scales
- $\cdot\,$  Nuclear physics spans lengths from  $10^{-15}\text{--}10^9\,\text{m}$

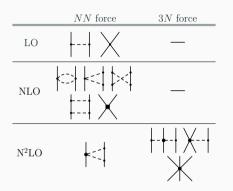


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- Fine details at one level of analysis do not affect the physics at a coarser level of analysis

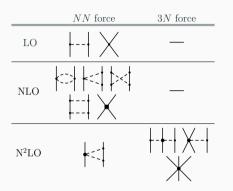


- There is interesting physics at all scales
- $\cdot\,$  Nuclear physics spans lengths from  $10^{-15}\text{--}10^9\,\text{m}$
- Fine details at one level of analysis do not affect the physics at a coarser level of analysis
- Start simple  $\rightarrow$  add corrections to reach desired precision.





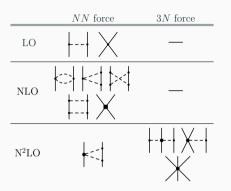
- $\cdot$  An expansion in the nuclear force
- $\cdot$  Ordered by increasing factors of small parameter Q
- $\cdot\,$  Truncation  $\rightarrow$  main source of uncertainty
- + Force convergence  $\neq$  prediction convergence
- $\cdot\,$  The debate on the "best" expansion is ongoing



#### We want to

- $\cdot$  An expansion in the nuclear force
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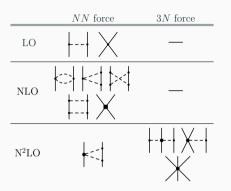
• Fit unknown parameters  $\vec{a}$ , or low-energy constants, with discrepancy  $\delta y_{th}$ 



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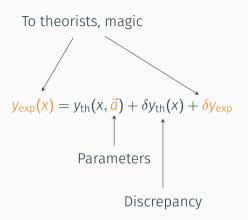


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- Fit unknown parameters  $\vec{a}$ , or low-energy constants, with discrepancy  $\delta y_{th}$
- $\cdot$  Quantify uncertainty in predictions (aka observables)  $y_{\rm th}$
- Test existing EFTs, uncover physics

$$y_{\exp}(x) = y_{th}(x, \vec{a}) + \delta y_{th}(x) + \delta y_{\exp}(x)$$



$$\chi^{2} \text{ fit}$$

$$y_{\exp}(x) = y_{\text{th}}(x, \vec{a}) + \delta y_{\exp}(x)$$

rigorous fit  

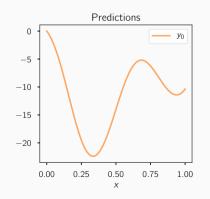
$$y_{exp}(x) = y_{th}(x, \vec{a}) + \delta y_{th}(x) + \delta y_{exp}$$

Full Prediction  

$$y_{exp}(x) = \overline{y_{th}(x, \vec{a}) + \delta y_{th}(x)} + \delta y_{exp}$$

 $y_{\exp}(x) = y_{th}(x, \vec{a}) + \underbrace{\delta y_{th}(x)}_{Can we build this?}$ Can we use it?

• Theoretical predictions could look like the following

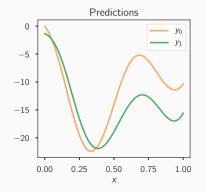


### $\{y_0\}$

 $y_0 \rightarrow LO$ 

• Theoretical predictions could look like the following

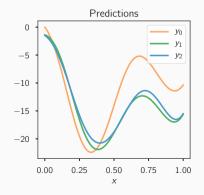
 $\{\mathbf{y_0}, y_1\}$ 





• Theoretical predictions could look like the following

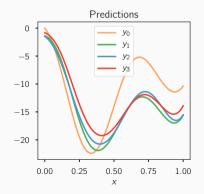
 $\{ \textbf{y_0}, y_1, \textbf{y_2} \}$ 



$$y_0 \rightarrow LO$$
  
 $y_1 \rightarrow NLO$   
 $y_2 \rightarrow N^2LC$ 

• Theoretical predictions could look like the following

# $\{y_0, y_1, y_2, y_3\}$



$$y_0 \to LO$$
  

$$y_1 \to NLO$$
  

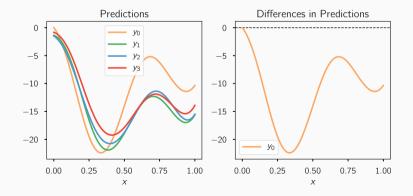
$$y_2 \to N^2 LO$$
  

$$\vdots$$
  

$$y_k \to N^k LO$$

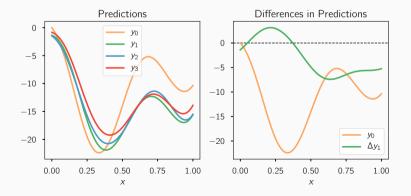
- Theoretical predictions could look like the following
- One can change variables for convenience/insight.

$$y_0 = y_0$$



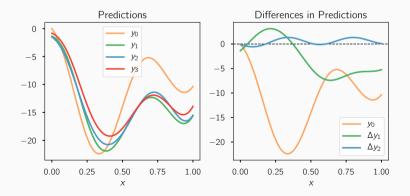
- Theoretical predictions could look like the following
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$$y_1 = \mathbf{y}_0 + \Delta y_1$$



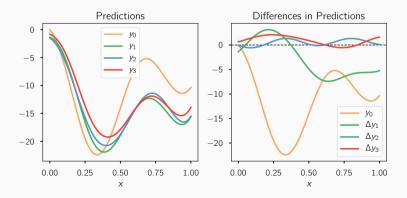
- Theoretical predictions could look like the following
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$$y_2 = y_0 + \Delta y_1 + \Delta y_2$$



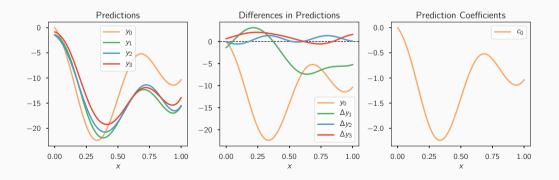
- Theoretical predictions could look like the following
- One can change variables for convenience/insight.
- $\Delta y_n = y_{ref}c_nQ^n$

$$y_3 = y_0 + \Delta y_1 + \Delta y_2 + \Delta y_3$$



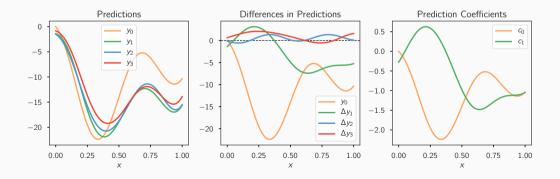
- $\cdot$  Theoretical predictions could look like the following
- One can change variables for convenience/insight.
- $\Delta y_n = y_{ref}c_nQ^n$

$$y_0 = y_{\rm ref} \left[ {\rm C}_0 Q^0 \right]$$



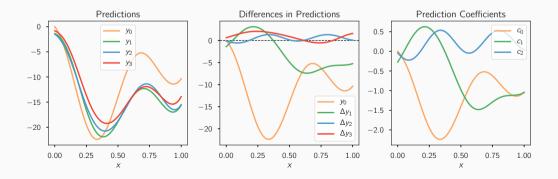
- Theoretical predictions could look like the following
- One can change variables for convenience/insight.
- $\Delta y_n = y_{\text{ref}} c_n Q^n$

$$y_1 = y_{ref} \left[ \frac{c_0 Q^0 + c_1 Q^1}{c_0 Q^0} + \frac{c_1 Q^1}{c_0 Q^0} \right]$$



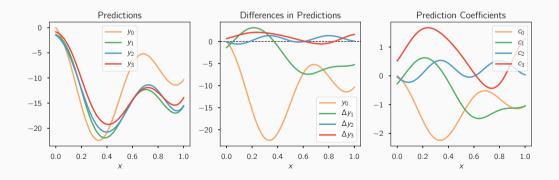
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$$y_2 = y_{ref} \left[ \frac{C_0 Q^0 + c_1 Q^1 + c_2 Q^2}{1 + c_2 Q^2} \right]$$



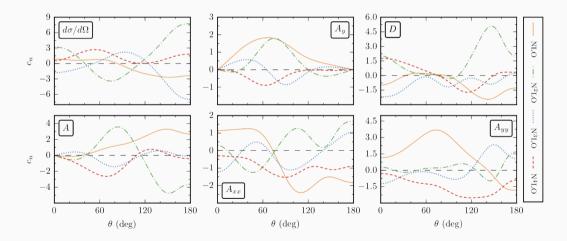
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$$y_3 = y_{\text{ref}} \left[ c_0 Q^0 + c_1 Q^1 + c_2 Q^2 + c_3 Q^3 \right]$$



#### **Real Life**

#### Coefficients from NN scattering look like our toy model!



# Statistical Model

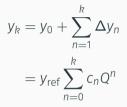
• Decompose prediction

$$y_k = y_0 + \sum_{n=1}^k \Delta y_n$$

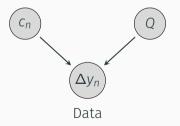


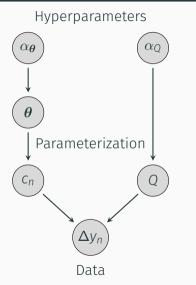
Data

Decompose prediction

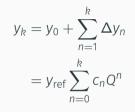


#### Parameterization



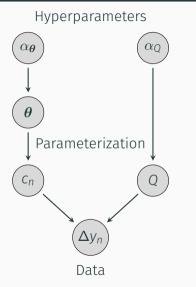


Decompose prediction

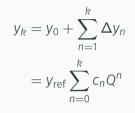


• Put priors on  $c_n$  (and Q)

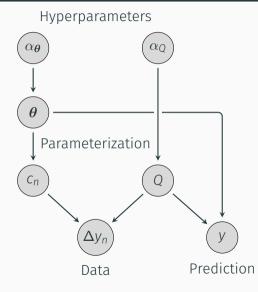
 $\operatorname{pr}(c_n \mid \boldsymbol{\theta}) \stackrel{\text{iid}}{\sim} \mathcal{GP}(\mu, \sigma^2 R_\ell)$ 



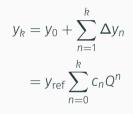
Decompose prediction



- Put priors on  $c_n$  (and Q)
  - $\operatorname{pr}(c_n \mid \boldsymbol{\theta}) \stackrel{\text{iid}}{\sim} \mathcal{GP}(\mu, \sigma^2 R_\ell)$
- $\cdot$  Learn heta and Q

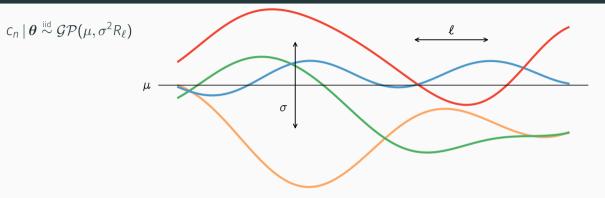


 $\cdot$  Decompose prediction

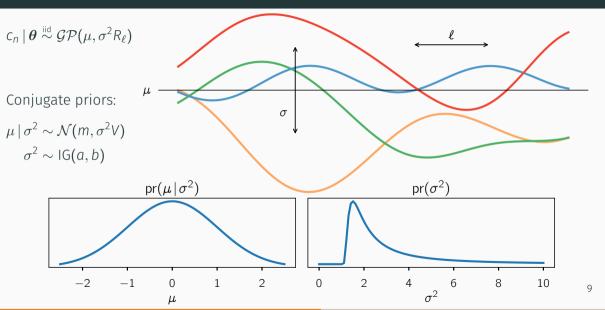


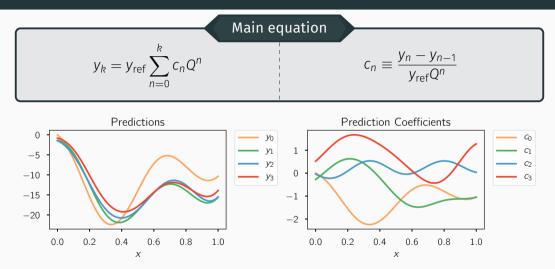
- Put priors on  $c_n$  (and Q)
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- $\cdot$  Learn heta and Q
- Predict pr(y | D)

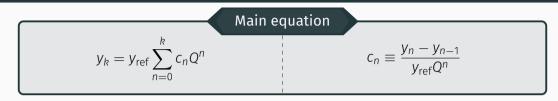
## Gaussian Process Priors on Observable Coefficients

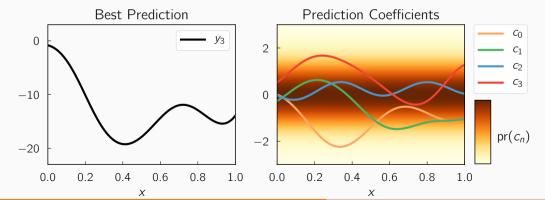


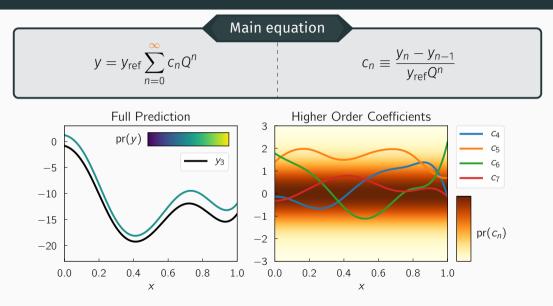
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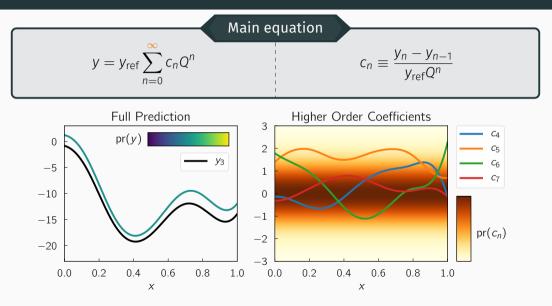


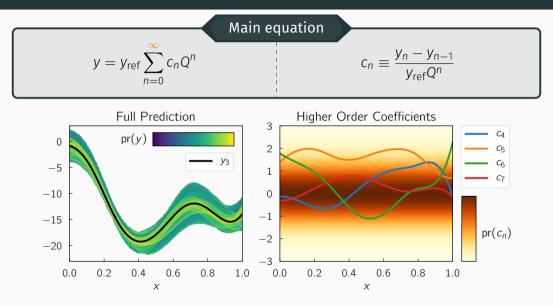












Remember the goal:

 $y_{\exp}(x) = y_{\mathrm{th}}(x, \vec{a}) + \delta y_{\mathrm{th}}(x) + \delta y_{\mathrm{exp}}$ 

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Our convergence assumptions

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$$\delta y_{\text{th}}(x) = y_{\text{ref}} \sum_{n=k+1}^{\infty} c_n Q^n$$

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Gaussian sum rules

$$a\mathcal{N}(\mu_1, \Sigma_1) + b\mathcal{N}(\mu_2, \Sigma_2) = \mathcal{N}(a\mu_1 + b\mu_2, a^2\Sigma_1 + b^2\Sigma_2)$$

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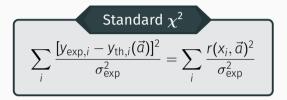
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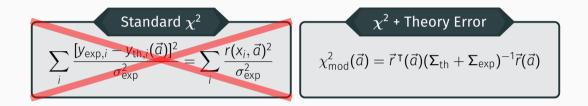
Gaussian sum rules

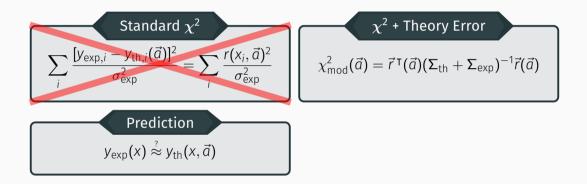
$$a\mathcal{N}(\mu_1, \Sigma_1) + b\mathcal{N}(\mu_2, \Sigma_2) = \mathcal{N}(a\mu_1 + b\mu_2, a^2\Sigma_1 + b^2\Sigma_2)$$

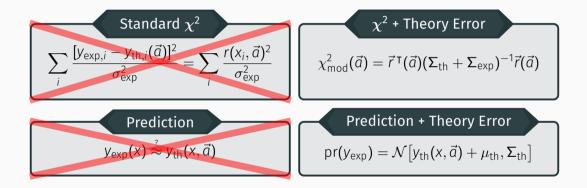
Discrepancy Distribution

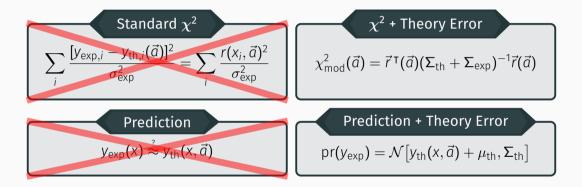
$$\operatorname{pr}(\delta y_{\mathrm{th}} | \boldsymbol{\theta}) = \mathcal{GP}(\mu_{\mathrm{th}}, \boldsymbol{\Sigma}_{\mathrm{th}}) = \mathcal{GP}\left(\frac{\mu Q^{k+1}}{1-Q}, y_{\mathrm{ref}}^2 \frac{\sigma^2 Q^{2(k+1)}}{1-Q^2} R_\ell\right)$$



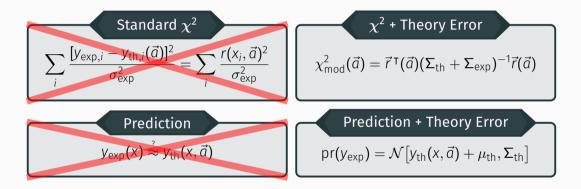








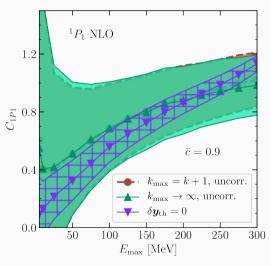
• Gaussian process correlations propagate via  $\Sigma_{th}$  matrix (computed once!)



- Gaussian process correlations propagate via  $\Sigma_{th}$  matrix (computed once!)
- $\cdot \,$  Different correlation assumptions  $\rightarrow$  different results!

#### What You Get for Free: Max Energy Insensitivity

- $\cdot$  y axis: posterior median  $\pm$  1 $\sigma$
- x axis: max energy of data in fit

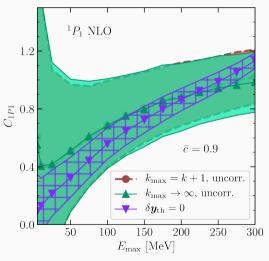


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- x axis: max energy of data in fit
- $\cdot$  *Q*, and hence  $\delta y_{\mathrm{th}}$ , grows with energy

$$\delta y_{\rm th} = y_{\rm ref} \sum_{n=k+1}^{k_{\rm max}} c_n Q^n$$

- This weights high energy data less!
- Stabilizes LEC fit as a function of E

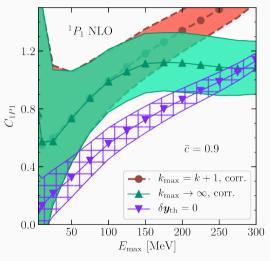


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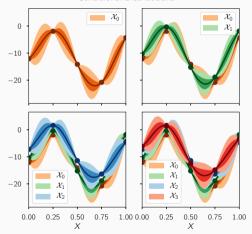
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- This weights high energy data less!
- Stabilizes LEC fit as a function of E
- Correlation assumptions can lead to different results

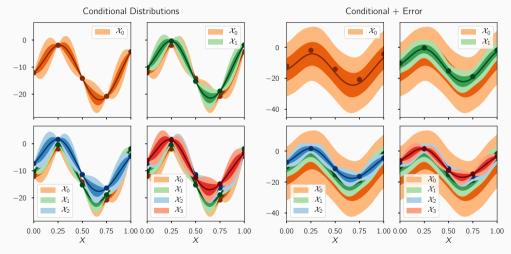


### Quantifying Truncation Uncertainty



Conditional Distributions

#### Quantifying Truncation Uncertainty



This model permits mostly analytic calculation of evidence

$$\operatorname{pr}(\mathcal{D} \mid \ell, Q) = \frac{\Gamma(a^*)}{\Gamma(a)} \frac{b^a}{(b^*)^{a^*}} \sqrt{\frac{|V^*|}{|V|} \frac{|2\pi R_\ell|^{-(k+1)/2}}{|Q|^{k(k+1)/2}}}$$

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Important for model comparison and for posteriors:

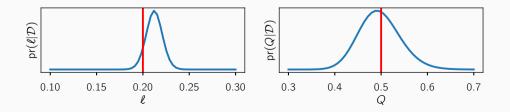
$$pr(\ell \mid \mathcal{D}, Q) \propto pr(\mathcal{D} \mid \ell, Q) pr(\ell)$$
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This model permits mostly analytic calculation of evidence

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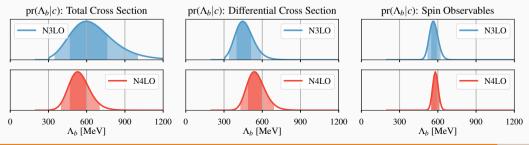
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Here, 
$$Q \propto \frac{1}{\Lambda_b}$$



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— Albert Einstein

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Does our model refer to reality? How can we check?

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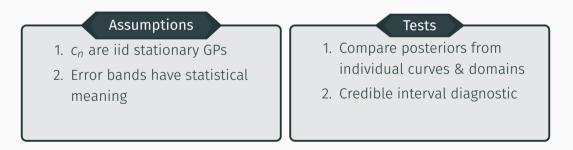
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Assumptions	Tests
1. <i>c<sub>n</sub></i> are iid stationary GPs	1. Compare posteriors from individual curves & domains

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# Model Checking

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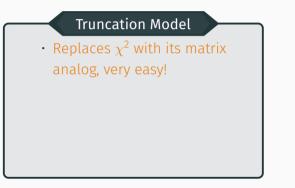
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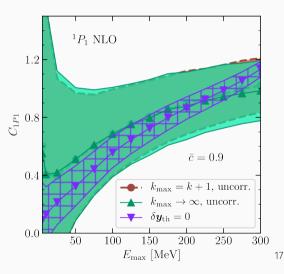


$$\chi^2_{\rm mod}(\vec{a}) = \vec{r}^{\mathsf{T}}(\vec{a})(\Sigma_{\rm th} + \Sigma_{\rm exp})^{-1}\vec{r}(\vec{a})$$

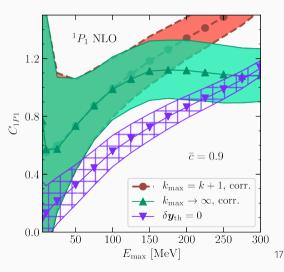
- Replaces  $\chi^2$  with its matrix analog, very easy!
- Full error can be propagated

$$pr(y_{exp}) = \mathcal{N}[y_{th}(x, \vec{a}) + \mu_{th}, \Sigma_{th}]$$

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#### Todo List

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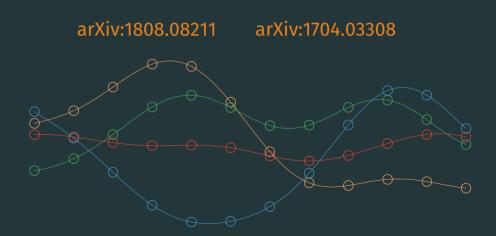
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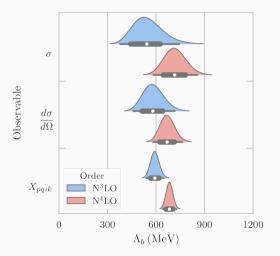
Suggestions welcome!

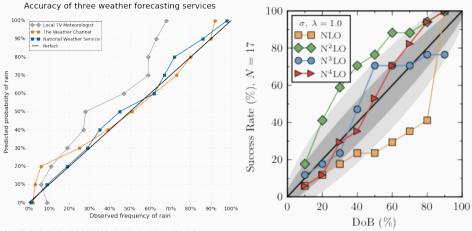
# Thank you!



## **Uncorrelated Posteriors**

#### Assumes that the variance of the $c_n$ is independent at each point





Source: "The Signal and the Noise" by Nate Silver | Author: Randy Olson (randalolson.com / @randal\_olson)