### Central Depression of Nucleonic Densities Trend analysis in the nuclear DFT approach

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## Physical motivation

- Superheavy nuclei with vanishing/reduced central density (bubble/semi-bubble) expected already in the 1940s.
- Effect attributed to the electrostatic repulsion of protons.
- Central depression of light nuclei also expected in <sup>34</sup>Si, <sup>46</sup>Ar and <sup>206</sup>Hg. Here depression is due to low occupation of s-orbitals.
- Proton density distributions probed through electron scattering. PREX and CREX can probe neutron densities of <sup>208</sup>Pb and <sup>48</sup>Ca.

Questions:

- What is the mechanism behind central depression in both light and heavy nuclei?
- Can we gain information about nuclear matter properties from density distributions?



Decharge et al., Nucl. Phys. A 716 (2003) 55 2 / 15

#### Nucleonic Densities

- Proton central depression for <sup>34</sup>Si, <sup>302</sup>Og and <sup>472</sup>164. Neutron central depression only for <sup>472</sup>164.
- <sup>34</sup>Si is very sensitive to EDF. (EDFs share about the same bulk properties).
- Central depression in <sup>34</sup>Si is large for SLy6 parametrization predicting large shell gap between 0d<sub>5/2</sub> and 1s<sub>1/2</sub>.
- In general, all models agree in heavy and superheavy nuclei.



Radial Skyrme DFT densities of <sup>34</sup>Si, <sup>48</sup>Ca, <sup>208</sup>Pb, <sup>302</sup>Og, and <sup>472</sup>164. Shaded areas mark spread of results obtained with SV-min, SLy6, and UNEDF1.

#### Quantification of central depression

Modified Helm model:

• Parametrize density as:

$$\rho_H(r; R_d, \sigma) = \rho_h(r; R_d) * \rho_G(r; \sigma),$$

$$\rho_h(r; R_d) = (1 + wr^2)\rho_0 \Theta(R_d - r)$$

 $\rho_G(r; \sigma)$ : Gaussian with width  $\sigma$ .  $R_d$ : Diffraction radius.

Quantification with

$$\bar{w} = w \cdot R_d^2$$

which can be obtained from shift of first and second zero in form factor.

• Robust with respect to shell fluctuations.

Differences of densities:

Quantification with

$$\bar{\rho}_{t,c} = (\rho_{t,av} - \rho_{t,c})/\rho_{t,av}$$

$$\begin{split} t &= (n,p) \\ \rho_{t,\mathrm{av}} &= N_t / (4/3\pi R_d^3): \text{ Average} \\ \text{density up to diffraction radius.} \\ \rho_{t,\mathrm{c}}: \text{ Central density.} \end{split}$$

- Straightforward way to quantify central depression.
- Sensitive to oscillations due to shell effects.

## Coefficient of Determination (CoD)

Question: How well determined is an observable by one model parameter?

- $\bullet$  Assessment through statistical analysis around the  $\chi^2$  minimum of EDF parametrization.
- Coefficient of determination:

$$R_{x,y}^2 = \left(\frac{\operatorname{cov}(x,y)}{\sigma_x \sigma_y}\right)^2$$

where cov(x, y) is the covariance of x and y and  $\sigma_x$  is the standard deviation of x.

- Gives values between 0 and 1.
- Can be evaluated for correlations between two model parameters, two observables or one observable and one model parameter.
- CoD does not give any information about the rate of change (slope).

#### Central Depression

- <sup>34</sup>Si proton central depression increases dramatically without pairing (large Z=14 subshell closure).
- \$\bar{w}\_p\$ larger than \$\bar{w}\_n\$, flat up to Z=90.
   Increases in the N=184 isotonic chain.
- $\rho_{p,c}$  increasing for N=82 and 126 isotonic chains, but flat for N=184 chain.
- Dip at <sup>208</sup>Pb due to full occupation of 2s orbit.
- Correlation of central density with Coulomb energy only significant for N=184 chain.



Proton central depression  $\bar{w}_{\rho}$  (a), central density  $\rho_{t,c}$  (b), and CoD between  $E_{\rm Coul}$  and  $\rho_{\rho,c}$  (c) predicted by SV-min.  $\times$  marks <sup>34</sup>Si values obtained without pairing.

### Correlation Analysis

- <sup>48</sup>Ca typical for nuclei lighter than <sup>208</sup>Pb, <sup>302</sup>Og for heavier nuclei.
- CoD between measures for central depression is low for <sup>48</sup>Ca because of shell effects, but high for <sup>302</sup>Og.
- CoD between observable and model parameters insignificant for <sup>48</sup>Ca, but large in <sup>302</sup>Og especially with E/A, J, L, a<sub>surf</sub>.

 $\rightarrow$  Central depression in nuclei below  $^{208}\text{Pb}$  is governed by shell effects. For superheavy nuclei CoDs reveal strong correlation with LDM parameters and Coulomb.



Matrices of CoD for SV-min parameters and selected observables in  $^{48}$ Ca (upper triangle) and  $^{302}$ Og (lower triangle).

#### Multiple Correlations

- LDM: (E/A,  $\rho_{eq}$ , K, J, L,  $a_{surf}$ ,  $a_{surf,s}$ )
- bulk: (E/A, ρ<sub>eq</sub>, K, J, L)
- sym: (J, L)
- Is+pair:  $(C_t^{\rho \nabla J}, V_{\text{pair},t}, \rho_{\text{pair}}), t=(n,p)$
- $E_{\rm Coul}$  almost fully determined by LDM.
- Surface effects important for N=82 and for large Z due to competition between Coulomb and surface tension.
- $\bar{\rho}_{\rm p,c}$  largely determined by LDM, especially for superheavy nuclei.
- Symmetry energy very important for heavy and superheavy nuclei, where Coulomb determines the central depression.



Multiple correlation coefficients with  $E_{\text{Coul}}$  (a) and  $\bar{\rho}_{p,c}$  (b) in heavy nuclei.

#### Localization for superheavy nuclei



Jerabek, B.S., Schwerdtfeger, Nazarewicz PRL 120, 053001 (2018)

 Localization: High value: One shell dominates. Low value: shells overlapping.
 0.5: Fermi gas limit.

- Neutron and proton localization function for <sup>132</sup>Sn, <sup>302</sup>Og, and <sup>472</sup>164.
- Going heavier, shell structure fades and extrema become fainter in NLF. Transition to Fermi gas.
- Effect is stronger for Neutrons, because of more particles occupying the same space (higher level density).

#### Charge radius isotopic shifts



- Measurement of nobelium charge radius isotopic shifts at GSI.
- DFT can describe isotopic shifts of charge radii for heavy nuclei.
- Simultaneously, DFT predicts appreciable semi-bubbles.
- Simple droplet models fail to predict isotopic shifts.

## Conclusions

#### Light and medium mass nuclei:

- Central depression in <sup>34</sup>Si highly sensitive to the input used (interactions, EDFs).
- No strong correlation of density distribution of <sup>48</sup>Ca with specific parameter.

 $\rightarrow$  Density distributions of light nuclei below  $^{208}\text{Pb}$  especially  $^{34}\text{Si}$  are governed by shell effects and carry little information about nuclear matter parameters.

#### Heavy nuclei:

- Central depression in heavy nuclei above lead is mainly driven by electrostatic repulsion.
- Since the Coulomb term has a large isovector component, the symmetry energy is strongly correlated to central depression.
- Heavy nuclei are closer to the leptodermous limit (determined by LDM parameters) and density distribution can be related better to nuclear matter parameters.
- Predicted densities are consistent with newest charge radii measurements.



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#### Density Functional Theory with Skyrme Energy Density Functionals

• Nucleonic part of the energy density functional (EDF):

$$\mathcal{E}_{\mathrm{Sk}} = \sum_{t=0,1} \left( C_t^{\rho}(\rho_0) \rho_t^2 + C_t^{\rho \Delta \rho} \rho_t \Delta \rho_t + C_t^{\tau} \rho_t \tau_t + C_t^{\rho \nabla J} \rho_t \nabla J_t \right) \quad (t=0: \text{ isoscalar; } t=1: \text{ iso}$$

$$\mathcal{E}_{\mathrm{pair}} = \sum_{q=p,n} rac{V_{\mathrm{pair},q}}{2} \left[ 1 - rac{
ho_0}{
ho_{\mathrm{pair}}} 
ight] ilde{
ho}^2$$

- Some constants can be expressed through nuclear matter properties:
  - $\rho_{eq}$  Equilibrium density.
  - E/A Energy-per-nucleon at equilibrium.
    - K Incompressibility
  - m\*/m Effective mass characterizing the dynamical isoscalar response.
    - J Symmetry energy.
    - L Slope of symmetry energy.

- κ Thomas-Reiche-Kuhn sum-rule enhancement characterizing the dynamical isovector response.
- asurf Surface energy coefficient
- a<sub>surf,s</sub> Surface-symmetry energy coefficient.
- $C_t^{\rho \nabla J}$  Spin-orbit parameters.
- $V_{\text{pair},q}$ ,  $\rho_{\text{pair}}$  Pairing parameters.

Multiple correlation coefficient (MMC)

Question: How well determined is an observable by a group of model parameters?

- CoDs are not an additive quantity due to correlations between model parameters.
- Multiple correlation coefficient:

$$R_{a,x}^2 = \boldsymbol{c}^T (R_{a,a})^{-1} \boldsymbol{c}$$

where  $R_{a,a}$  contains correlations between the model parameters of group a,  $c = (R_{a_1,x}, R_{a_2,x}, ...)$  contains correlations between the observable and single group members.

- Value of 0.3 means that 30% of the variance of x is predictable from group a.
- If **a** contains all model parameters,  $R_{a,x}^2 = 1$  always.

## Correlation Analysis

- Isoscalar and isovector densities correlate weakly with nuclear matter properties and other model parameters for nuclei below <sup>208</sup>Pb.
- Central densities in nuclei above lead carry information on nuclear matter properties, especially the symmetry energy.



Correlations of central densities and model parameters.