



Simulations of Post-Merger Evolution: Mass Ejection

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Overview

1. Timescales and processes

2. Accretion disk evolution

3. Challenges

Neutron Star Mergers



RF & Metzger (2016)

Reference Numbers

$$\begin{split} M &\sim M_{\odot} = 2 \times 10^{30} \text{ kg} & r_g = \frac{GM}{c^2} \sim 1.5 \text{ km} \\ R &\sim 10 \text{ km} \\ t_{\rm dyn} &= 2\pi \left(\frac{R^3}{GM}\right)^{1/2} \sim 10^{-3} \text{ s} & \longrightarrow f_{\rm GW} \lesssim 1 \text{ kHz} \\ \tau_\nu &\sim \frac{R}{\ell_\nu} \sim \left(\frac{\varepsilon_\nu}{15 \text{ MeV}}\right)^2 \left(\frac{\rho}{10^{11} \text{ g cm}^{-3}}\right) & \text{Neutrinos} \\ \bar{\rho}_{\rm NS} &\sim \frac{M}{R^3} \sim 10^{15} \text{ g cm}^{-3} & \downarrow \\ \text{Nuclear physics} & kT_{\rm vir} \sim 100 \text{ MeV} \end{split}$$

Dynamical Phase: Merger

Rezzolla+ (2010)

Unequal mass NS-NS merger:



Phases:

- inspiral
- merger
- remnant + ejecta
- relativistic jet (?)

Large body of work:

MPA, Kyoto, Caltech-Cornell-CITA Princeton, Frankfurt, Trento, Stockholm, Illinois, Perimeter, etc.

Late-time dynamical ejecta



Rosswog+ (2014)

Disk evolution



also Popham+ (1999), Chen & Beloborodov (2003)

$$\begin{split} t_{\rm orb} &\simeq 3 R_{50}^{3/2} M_3^{-1/2} \mbox{ ms} \\ t_{\rm visc} &\simeq 1 \alpha_{0.03}^{-1} R_{50}^{3/2} M_3^{-1/2} \left(H/3R \right) \mbox{ s} \\ t_{\rm therm} &\simeq \frac{c_s^2}{v_K^2} t_{\rm visc} \lesssim t_{\rm visc} \end{split}$$

- Disk evolves on timescales long compared to the dynamical (orbital) time, due to viscous processes
- Weak interactions freeze-out as the disk spreads viscously: final Ye
- Gravitationally-unbound outflows driven by:
 - Neutrino heating (on thermal time) Ruffert & Janka (1999), Dessart+ (2009)
 - Viscous heating and nuclear recombination (on viscous time)

 ${E_lpha\over GM_{
m BH}/R}\simeq 1R_{600}M_3^{-1}$ Metzger+ (2009)

- MHD stresses

Kiuchi (2015), Siegel (2017)

Equations

$$\begin{array}{lll} \mbox{mass}\\ \mbox{conservation:} & \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho {\bf v}) = 0 & \rho: \mbox{ density} \\ \mbox{w: velocity} \\ \mbox{momentum}\\ \mbox{conservation:} & \frac{\partial {\bf v}}{\partial t} + ({\bf v} \cdot \nabla) {\bf v} + \frac{1}{\rho} \nabla p = -\nabla \Phi & + \frac{1}{\rho} \nabla \cdot \mathbb{T} \\ \mbox{gas}\\ \mbox{gas}\\ \mbox{grassure} & \mbox{gravity} & \mbox{angular mom}, \\ \mbox{transport} \\ \mbox{energy}\\ \mbox{conservation:} & \frac{De_{\rm int}}{Dt} - \frac{p}{\rho^2} \frac{D\rho}{Dt} = \frac{1}{\rho^2 \nu} \mathbb{T} : \mathbb{T} & + Q_{\nu, \rm abs} & -Q_{\nu, \rm em} \\ \mbox{viscous} & \mbox{neutrino} & \mbox{neutrino} \\ \mbox{heating} & \mbox{neutrino} & \mbox{neutrino} \\ \mbox{neutrino} & \mbox{neutrino} \\ \mbox{neutrino} & \mbox{neutrino} \\ \mbox{absorption} & \mbox{emission} \\ \mbox{EOS:} & p = p(\rho, e_{\rm int}, Y_e) & Y_e = \frac{n_e}{n} = \frac{n_e}{\rho/m_n} & Y_e: \mbox{electron fraction} \\ \end{array}$$

Equations

 $\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$ hydrodynamics: mass conservation: FLASH $\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla)\mathbf{v} + \frac{1}{\rho}\nabla p = -\nabla \Phi + \frac{1}{\rho}\nabla \cdot \mathbb{T}$ momentum pseudo-Newtonian conservation: gravity gas angular mom. gravity pressure transport $\frac{De_{\text{int}}}{Dt} - \frac{p}{\rho^2} \frac{D\rho}{Dt} = \frac{1}{\rho^2 \nu} \mathbb{T} : \mathbb{T} + Q_{\nu,\text{abs}} - Q_{\nu,\text{em}}$ energy a-viscosity conservation: neutrino neutrino viscous heating heating cooling neutrino lepton # $\frac{DY_e}{Dt} =$ leakage $\Gamma_{\nu, \mathrm{abs}}$ $+\Gamma_{\nu,\mathrm{em}}$ conservation: neutrino neutrino emission absorption lightbulb self-irradiation $Y_e = \frac{n_e}{n} = \frac{n_e}{\rho/m_n}$ EOS: $p = p(\rho, e_{\text{int}}, Y_e)$

Helmholtz EOS



Wind from remnant accretion disk

- Neutrino cooling shuts down as disk spreads on accretion timescale (~300ms)
- Viscous heating & nuclear recombination are unbalanced
- Fraction ~10-20% of initial disk mass ejected, ~1E-3 to 1E-2 solar masses
- Material is neutron-rich (Ye ~ 0.2-0.4)
- Wind speed (~0.05c) is slower than dynamical ejecta (~0.1-0.3c)

RF & Metzger (2013), MNRAS Just et al. (2015), MNRAS RF et al. (2015), MNRAS Setiawan et al. (2005)

Lee, Ramirez-Ruiz, & Lopez-Camara (2009)

Metzger (2009)

Effect of BH spin on disk wind



Hypermassive NS versus BH



Disk around HMNS



e.g. Metzger et al. (2018)

Interplay of disk wind and dynamical ejecta



Interplay of disk wind and dynamical ejecta



RF, Foucart, Kasen, Lippuner, et al. (2017)

Interplay of disk wind and dynamical ejecta



RF, Foucart, Kasen, Lippuner, et al. (2017)

SGRB jet: neutrino pair annihilation

Energy injection weak, and jet has trouble breaking out of dynamical ejecta for NS-NS mergers



Nucleosynthesis with Tracer Particles



M-R Wu, RF, Martinez-Pinedo & Metzger (2016)

- Nuclear network: ~7000 isotopes, include neutrino effects
- Non-spinning BH, parameter dependencies

Black Hole Accretion Disks



• Most sensitive to viscosity: expansion time vs weak interaction time

 Also sensitive to disk mass and degeneracy: neutrinos & equilibrium Ye M-R Wu, RF, Martinez-Pinedo & Metzger (2016)

- Not very sensitive to initial Ye
- See also Just et al. 2015

Black Hole Accretion Disks



M-R Wu, RF, Martinez-Pinedo & Metzger (2016)

HMNS disks



Lippuner, RF, Roberts, et al. (2017)

Pending Issues in Disk Modeling

1) Include magnetic fields (in progress)

Siegel & Metzger (2017), Nouri et al. (2017)

2) Improve neutrino transport: outflow composition

Just et al. (2015), Fujibayashi et al. (2017), Foucart et al. (2018)

3) Realistic initial conditions for magnetic field and matter, interplay with dynamical ejecta

HMNS evolution with viscosity



Angular velocity profile inside the HMNS changes from differential to uniform over ~10ms. Large release of energy, and driving of an outflow.

Previous work had not included viscosity for the internal HMNS evolution.

Shibata et al. (2017), Fujibayashi et al. (2017)

SGRB jet: MHD

Winding of toroidal field and buildup of magnetic pressure around rotation axis.



Ruiz+ (2016)

In BH-NS, it requires the NS to have an external poloidal field Paschalidis+ (2015)

 See also:
 Rezzolla+ (2011)
 Kiuchi+ (2015)
 Kawamura+ (2016)

Summary

1. Accretion disks formed in NS mergers can eject significant amounts of mass, reprocessed by neutrinos and hence with different composition than bulk of dynamical ejecta

2. Outstanding issues in post-NS-merger modeling: GRMHD, neutrino transport, and realistic initial conditions including dynamical ejecta

