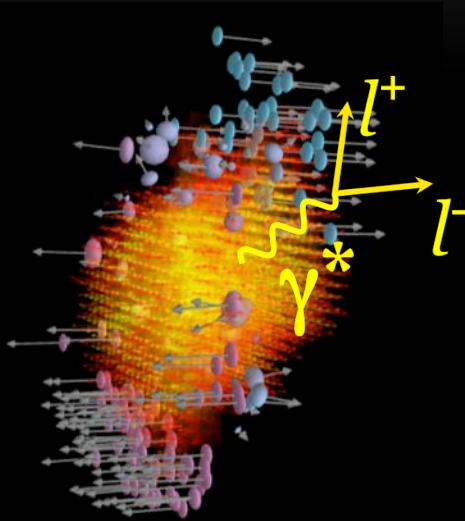


Constraining QCD Phase Boundary with Data from Heavy Ion Collisions

GSI, Darmstadt, Germany
February 12-14, 2018

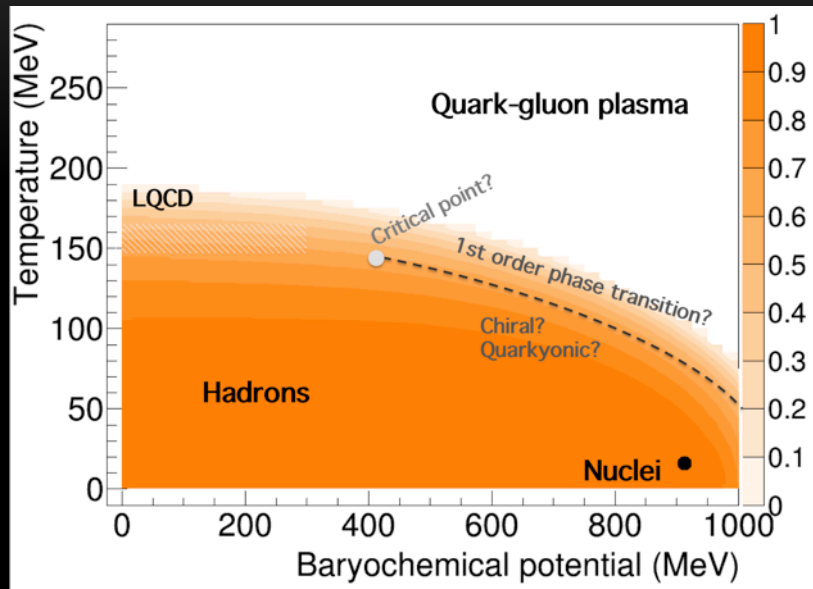


Exploring QCD phase diagram with dileptons

Tetyana Galatyuk
TU Darmstadt / GSI



Theoretical guidance



□ Vanishing μ_B , high T (Lattice QCD)

□ Crossover transition

□ $\epsilon_c \sim 1 \text{ GeV/fm}^3, T_c \sim 155 \text{ MeV}$

□ Large μ_B , moderate T (effective, Lattice QCD inspired models)

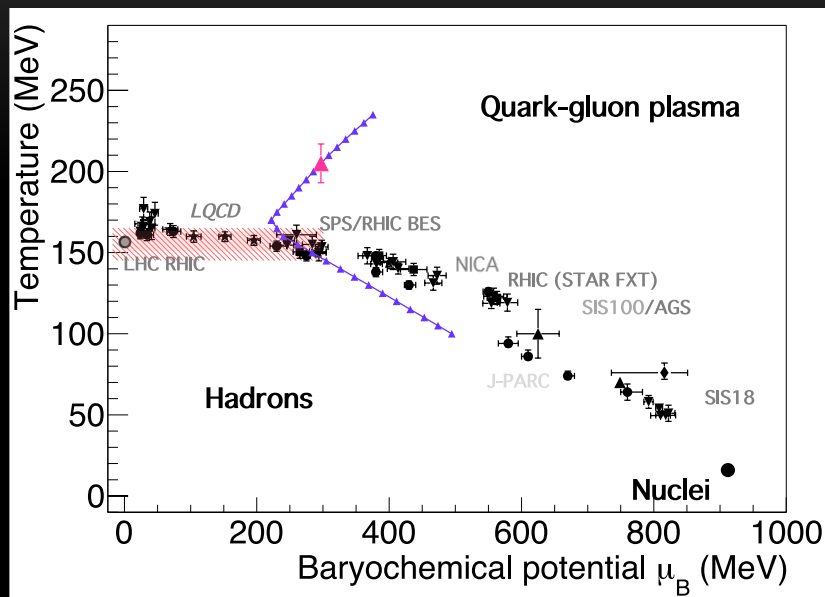
□ 1st order transition

□ QCD critical point

□ Melting of the condensate (order parameter $\langle 0 | \bar{q}q | 0 \rangle = \langle 0 | \bar{q}_L q_R + \bar{q}_R q_L | 0 \rangle \neq 0$)

$$\frac{\langle \bar{q}q \rangle_{T, \mu_B}}{\langle \bar{q}q \rangle_{T=0, \mu_B=0}} : \text{B.J. Schaefer and J. Wambach}$$

Searching for landmarks of the QCD matter phase diagram



□ Experimental approach:

- Probe with **highest precision** different regions of the QCD matter phase diagram

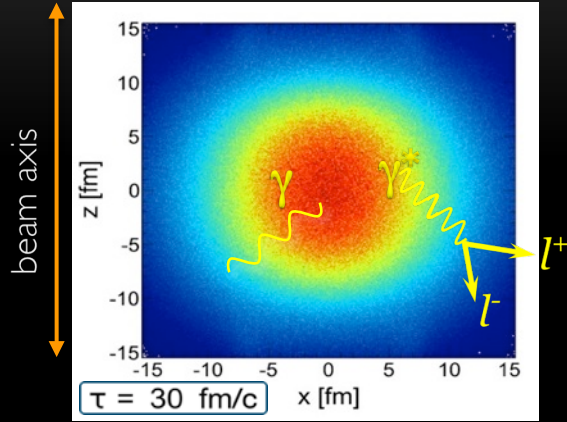
□ Observables:

- Flavor production (multi-strange, charm)
- **Emissivity of matter (dileptons)**
- Higher moments of e-b-e multiplicities (B, S, Q)

 O. Kaczmarek et al., PRD 83 (2011) 014504
S. Borsanyi et al. [Wuppertal-Budapest Coll.], JHEP 1009 (2010) 073
A. Andronic et al., arXiv:1710.09425 [nucl-th]
F. Becattini et al., Phys. Lett. B 764 (2017) 241
J. Cleymans et al., Phys. Rev. C 73 (2006) 034905
L. Adamczyk et al. [STAR Coll.], arXiv:1701.07065 [nucl-ex]
G. Agakishiev et al. [HADES Coll.], Eur. Phys. J. A 52 (2016) 178
X. Lopez et al. [FOPI Coll.], Phys. Rev. C 76 (2007) 052203
H. J. Specht et al. [NA60 Coll.], AIP Conf. Proc. 1322 (2010)

Electromagnetic radiation

Photons and lepton pairs probe the interior of fireballs – “PET” of the fireball

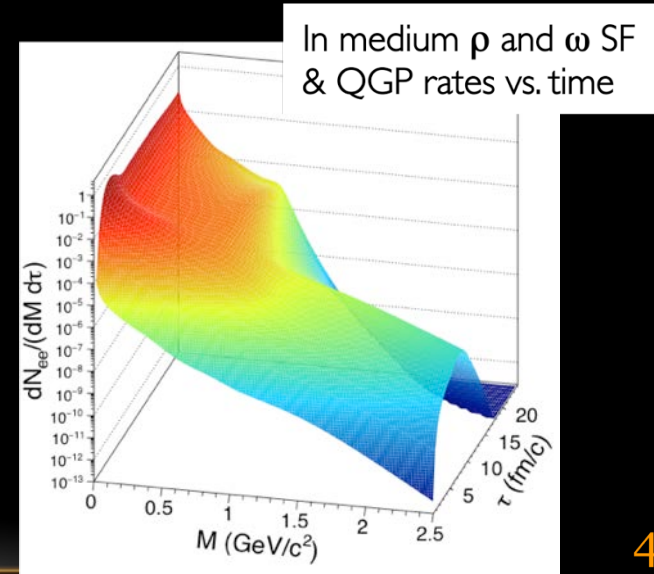


- EM radiation contains contributions from throughout the collision
- EM probes leave collision zone undisturbed
- Real γ characterized by transverse momentum (blue shifted)
- Lepton pairs carry extra information: invariant mass

The vector correlator is directly accessible in HIC

$$\frac{dN_{ll}}{d^4x d^4q} = \frac{-\alpha_{EM}^2}{\pi^3 M^2} f^B(q \cdot u; T) \text{Im} \Pi_{EM}^{\mu\nu}(M, q; \mu_B, T)$$

→ Unique direct access to in-medium spectral function



Electromagnetic correlator in vacuum

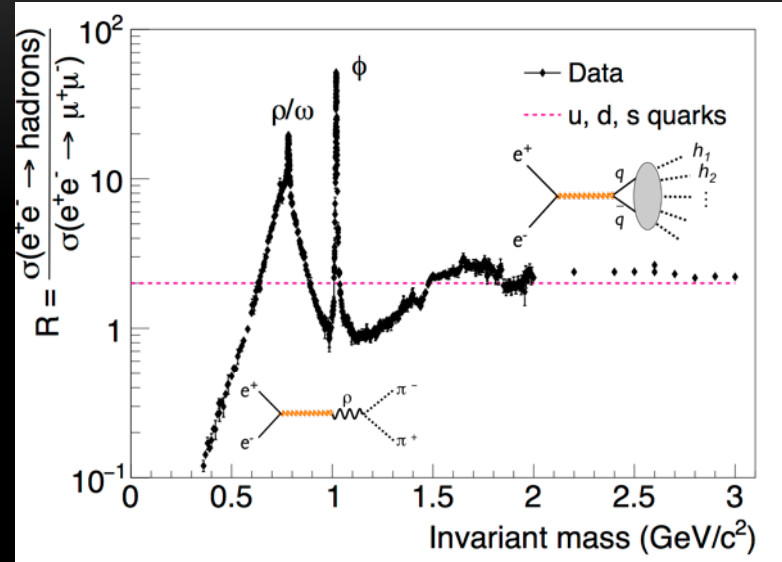
McLerran-Toimela formula

$$\frac{dN_{||}}{d^4x d^4q} = \frac{-\alpha_{EM}^2}{\pi^3 M^2} f^B(q \cdot u; T) \text{Im} \Pi_{EM}^{\mu\nu}(M, q; \mu_B, T)$$

$$R = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} \propto \frac{\text{Im} \Pi_{EM}^{vac}}{M^2}$$

- If $\frac{\text{Im} \Pi_{EM}^{\mu\nu}}{M^2} \sim \text{const.}$ thermal dilepton emission will follow a Bose distribution

→ Dileptons as thermometer



$$\text{Im} \Pi_{EM}^{vac}(M \leq 1.5 \text{ GeV}) = \sum_{\nu=\rho, \omega, \phi} \left(\frac{m_\nu^2}{g_\nu} \right)^2 \text{Im} D_\nu^{vac}(M)$$



Vector Meson Dominance: $J^P = 1^-$ for both γ^* and VM (with ρ playing a dominant role)

$$\text{Im} \Pi_{EM}^{vac}(M > 1.5 \text{ GeV}) = -\frac{M^2}{12\pi} \left(1 + \frac{\alpha_s(M)}{\pi} + \dots \right) N_c \sum_{q=u,d,s} (e_q)^2$$

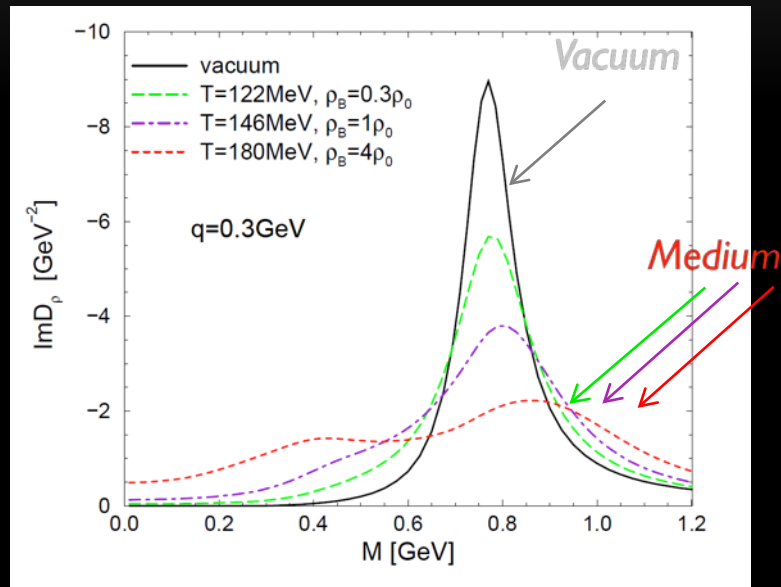


pQCD continuum

ρ meson in hot and dense medium

interacts with hadrons from heat bath $\rightarrow$ additional contributions to the ρ -meson self-energy

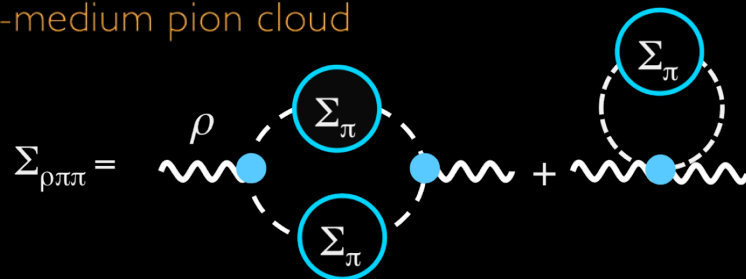
 R. Rapp and J. Wambach, Eur.Phys.J. A6 (1999)



Baryon effects important even at $\rho_{B,\text{tot}} = 0$:
sensitive to $\rho_{B,\text{tot}} = \rho_B + \rho_{\bar{B}}$ (ρ -N and ρ - $\bar{N}$
interactions identical)

$$D_\rho(M, q; \mu_B, T) = \frac{1}{\left[M^2 - m_\rho^2 - \underbrace{\Sigma_{\rho\pi\pi}}_{\text{In-medium pion cloud}} - \Sigma_{\rho B} - \Sigma_{\rho M} \right]}$$

In-medium pion cloud

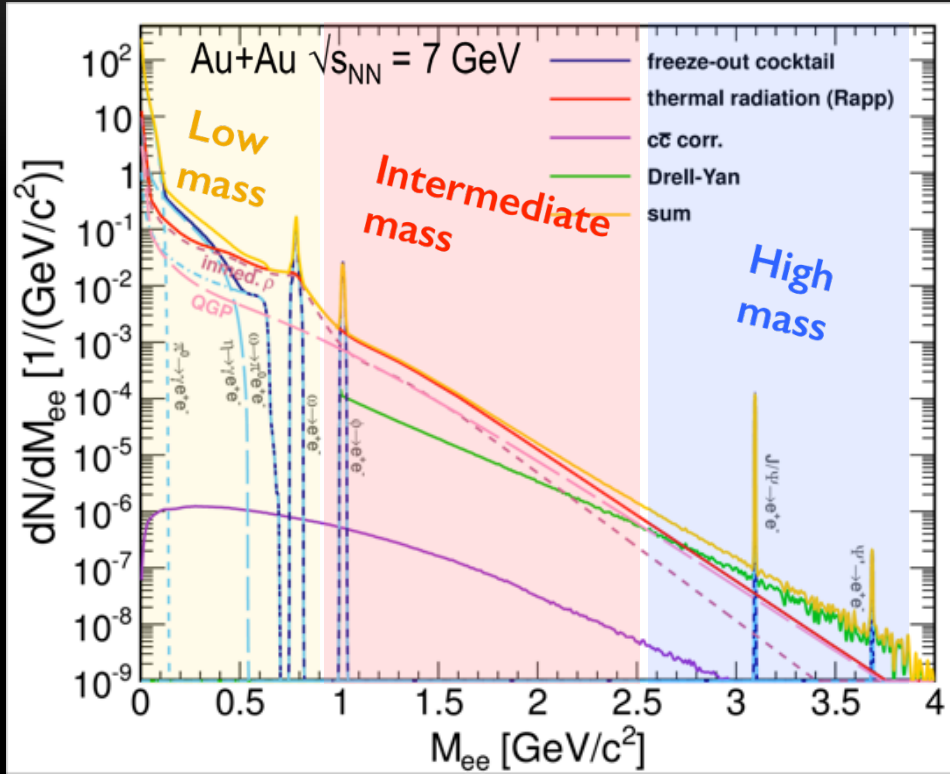


Direct ρ -hadron scattering

$$\Sigma_{\rho B, M} = \rho \circlearrowleft h = N, \pi, K$$

$R = \Delta, N(1520), a_1, K_1, \dots$

Characteristic regimes in invariant e^+e^- mass



- Drell-Yan: power-law $\sim M^n$
- Heavy-flavor: $c\bar{c} \rightarrow l^+l^-$
 - Becomes significant at RHIC and LHC
- Thermal radiation: $\sim \exp(-M/T)$
 - QGP – highest T , no flow
 - “ 4π annihilation”: $\pi a_1 \rightarrow l^+l^-$
 - In-medium ρ, ω – moderate T , flow
- Final state decays (hadron cocktail): $\pi^0, \eta \rightarrow \gamma e^+e^-$

$$M^2 = (P_{e^+} + P_{e^-})^2$$

$$p_t, v_2, \cos\theta$$

The experimental challenge ...

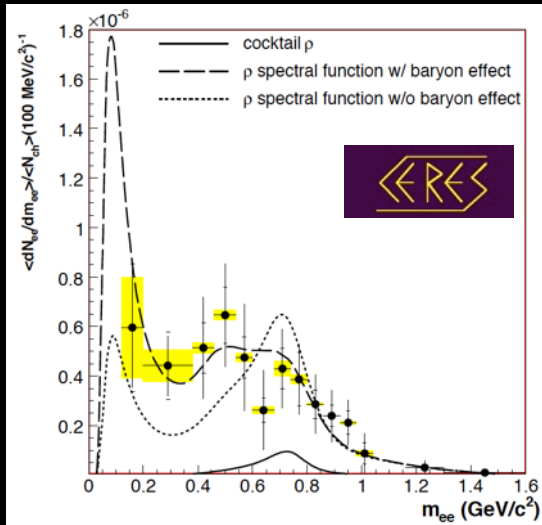
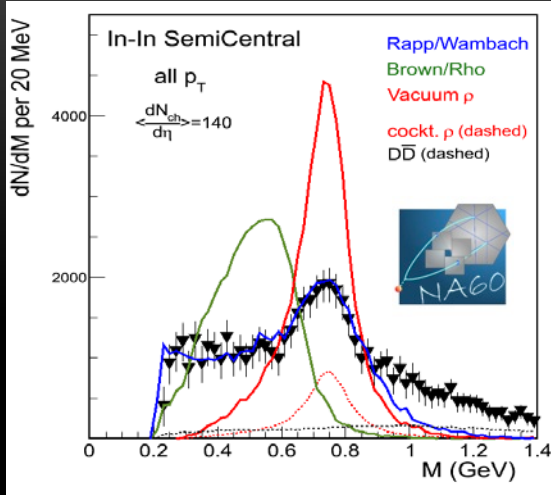
- ❑ Lepton pairs are rare probes ($\text{BR} < 10^{-4}$)
- ❑ at SIS energies sub-threshold vector meson production
 $\rightarrow M_r \times \Gamma_{ee}/\Gamma_{\text{tot}}$ decay per 10 mio events
- ❑ Large combinatorial background
 - ❑ in e^+e^- from Dalitz decays ($\pi^0 \rightarrow e^+e^-\gamma$) and conversion pairs (e^+e^-)
 - ❑ in $\mu^+\mu^-$: weak π, K decays
- ❑ Isolate the contribution to the spectrum from the dense stage
- ❑ Low-momentum coverage!



There is no such thing as a free lunch

DATA QUALITY

- ❑ The decisive parameters: Number of Interactions and Signal/Background
 - ❑ Range of B/S: 20 - 100 $\rightarrow B/S \gg 1$;
 - Effective sample size: $S_{\text{eff}} \sim I \times S/B$ reduction by factors of 20-100
 - Systematics: $\delta S_{\text{eff}}/S_{\text{eff}} = \delta B/B \times B/S$ $\delta B/B = 2 \dots 5 \times 10^{-2}$



Are narrow in-medium vector meson states with substantially shifted pole mass observed?

- Disfavors “dropping mass” scenario: $m_{had} \sim \langle \bar{q}q \rangle$
- Strongly supports in-medium broadening

$$\Sigma_{\rho B, M} = \text{Diagram with } \rho \text{ and } h = N, \pi, K \text{ lines and } R = \Delta, N(1520), a_1, K_1, \dots \text{ labels}$$



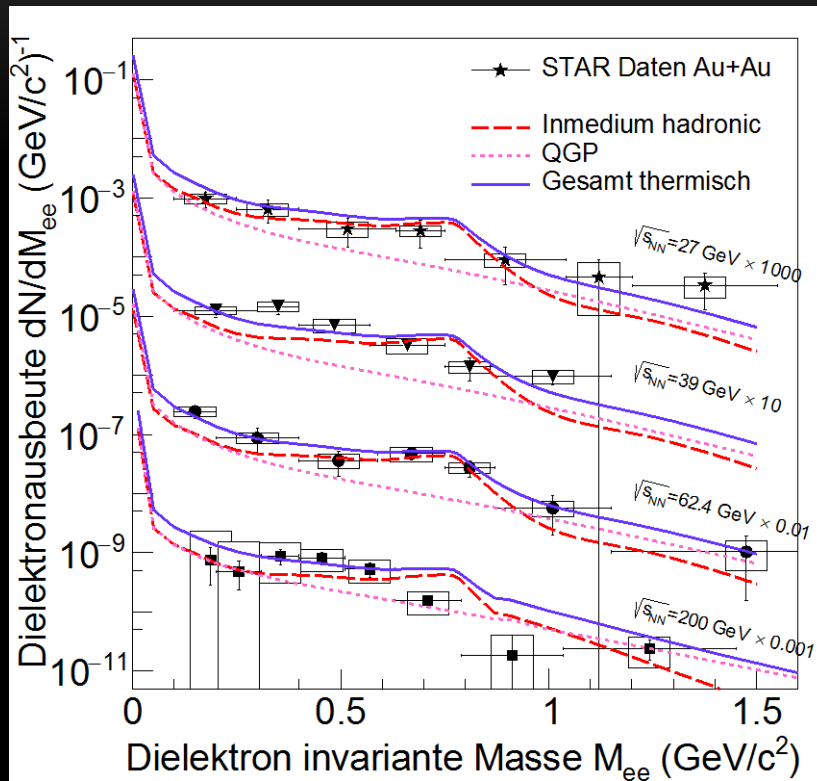
Data NA60: Phys. Rev. Lett. 96 (2006) 162302

Data CERES: Phys.Lett. B666 (2008)

Calculations: R.Rapp and H. van Hees, 2008

Dielectron mass spectra from STAR BES I

$Au+Au \sqrt{s_{NN}} = 19.6 - 200 \text{ GeV}$



Phys.Lett. B750 (2015) 64-71

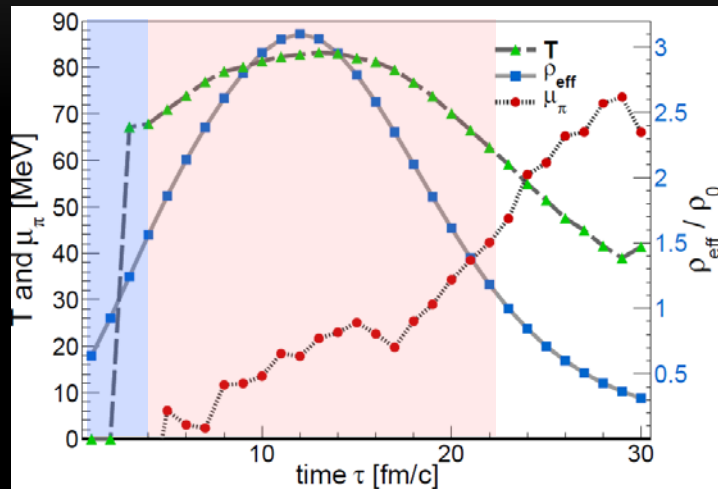
J. Butterworth et al., arXiv:1612.05484 [nucl-ex]

Model: Rapp/Wambach/Hees

- Isolation of the excess by subtracting the measured decay cocktail
- Acceptance corrected spectra
- In-medium **broadened ρ** spectral function consistently describes the low-mass electron-positron excess **for** all the energies **19.6-200 GeV**

Baryonic matter at few GeV beam energy $Au+Au \sqrt{s_{NN}} = 2.42 \text{ GeV}$

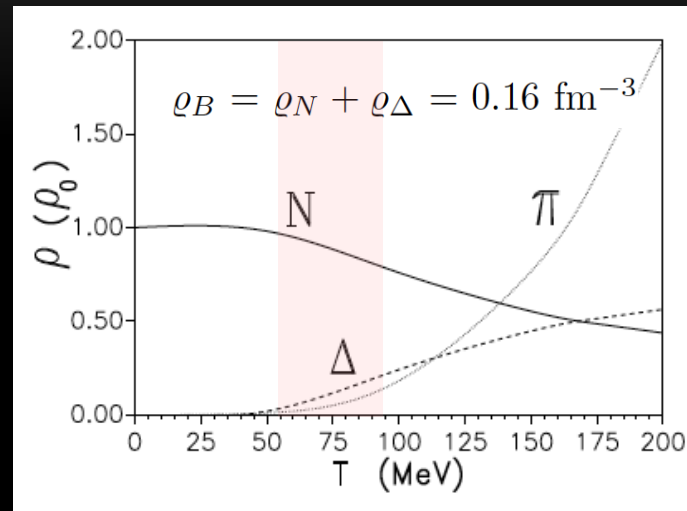
Central cell ($3 \times 3 \times 3 \text{ fm}^3$) thermodynamic properties from coarse graining UrQMD



TG, F. Seck et al., Eur. Phys. J. A 52 (2016) 131

- ❑ Long interpenetration times
- ❑ Comparatively long lifetime of the dense "fireball" ($\rho_{\text{max}} \approx 3 \rho_0$)

Composition of a hot $\pi\Delta N$ gas

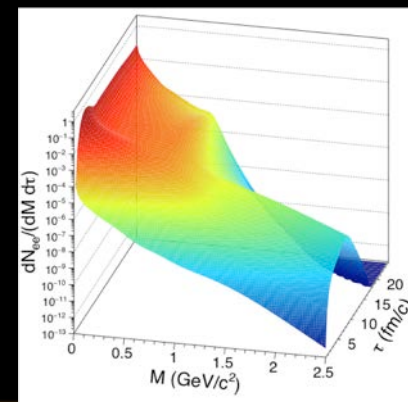
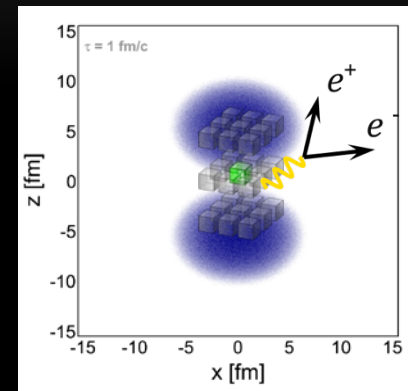
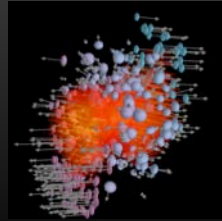
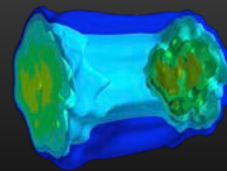


Rapp, Wambach, Adv. Nucl. Phys. 25 (2000)

- ❑ Moderate temperatures: $T < 90 \text{ MeV}$
- ❑ Baryon-dominated system throughout the evolution ($N_\pi / A_{\text{part}} \approx 10\%$)

Coarse-grained transport approach

- “Combine” the advantages of two descriptions: hydrodynamics & transport
- Simulate events with a transport model
 - ensemble average to obtain smooth space-time distributions
- Divide space-time evolution into 4-dimensional cells
 - $21 \times 21 \times 21$ space cells (1 fm^3), 30 time steps → $\sim 280 \text{ k cells}$
- Determine for each cell the bulk properties like T , μ_B , μ_π , collective velocity
 - parameterization of RW in-medium spectral function
- Sum up contributions of all cells



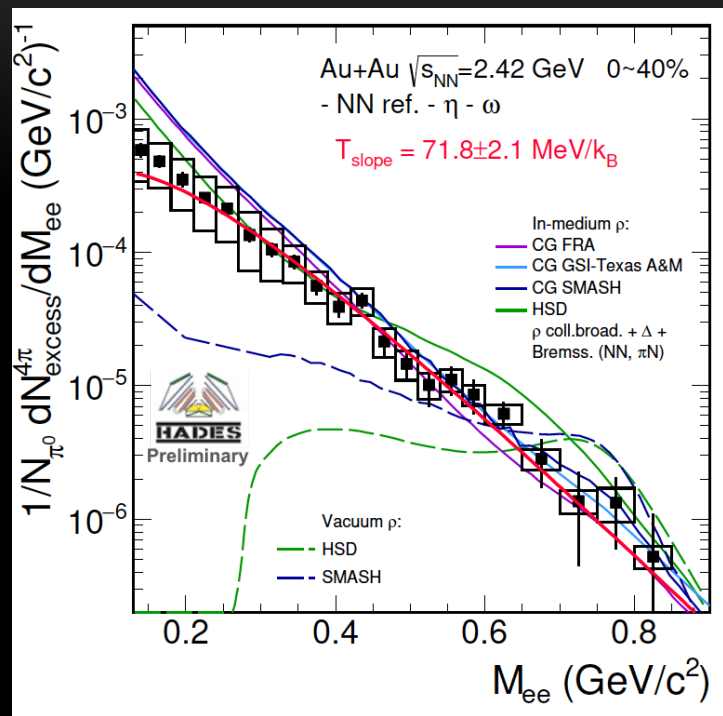
Huovinen et al., PRC 66 (2002) 014903

CG FRA Endres et al.: PRC 92 (2015) 014911

CG GSI-Texas A&M TG et al.: Eur.Phys.J. A52 (2016) no.5, 131

CG SMASH: J. Staudenmaier et al., arXiv:1711.10297v1

Dielectron mass spectra at $\sqrt{s_{NN}} = 2.42$ GeV



HADES, to appear soon

- Strong broadening of the in-medium ρ due to direct ρ -hadron scattering

$$\Sigma_{\rho B, M} = \begin{array}{c} \rho \\ \text{wavy line} \end{array} \begin{array}{c} \text{loop} \\ \text{with } R = \Delta, N(1520), a_1, K_1, \dots \\ \text{and } h = N, \pi, K \end{array}$$

- Thermal rates folded over coarse-grained UrQMD medium evolution works at low energies
- Supports baryon-driven medium effects at SPS and RHIC (LHC)!

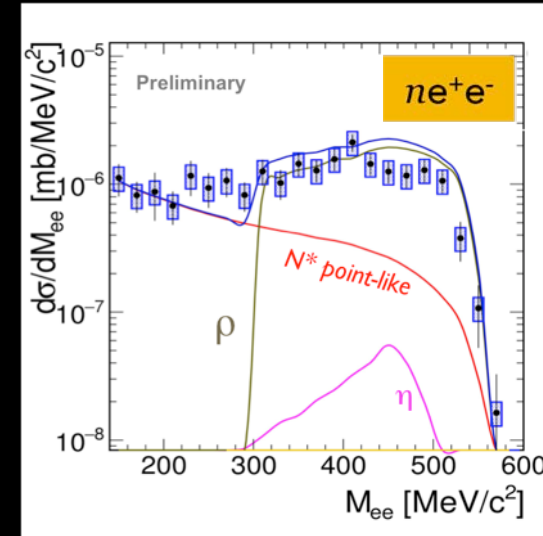
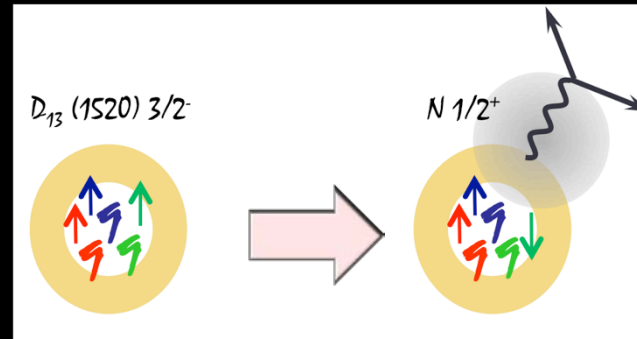
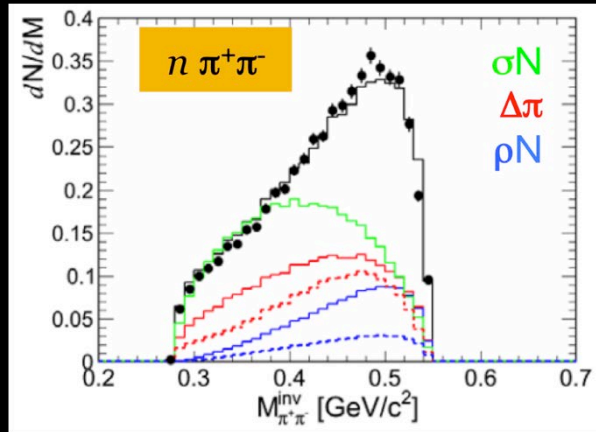
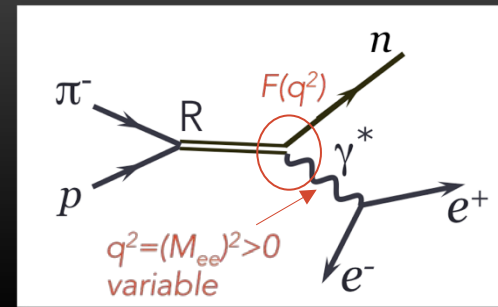
Robust understanding across QCD phase diagram



EM baryon-resonance decays

π^- (683,5 MeV/c) $p \rightarrow N^*$

- Worldwide unique combination of π beams (at low energies) with dilepton spectrometer
- Access to time-like etFF of baryons



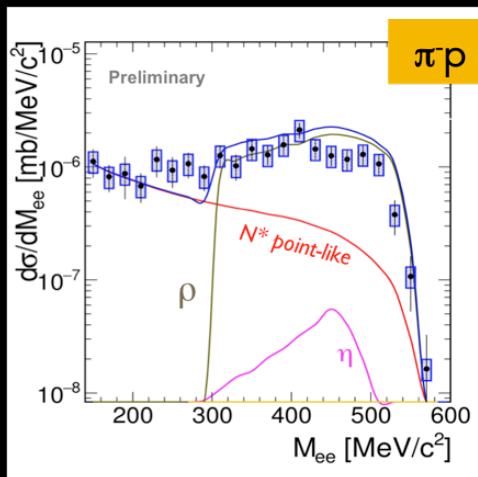
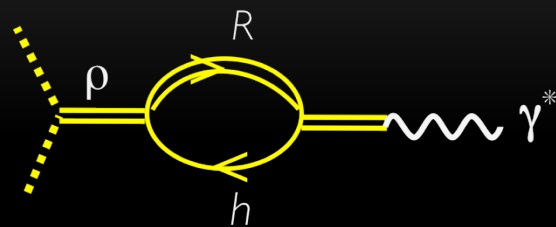
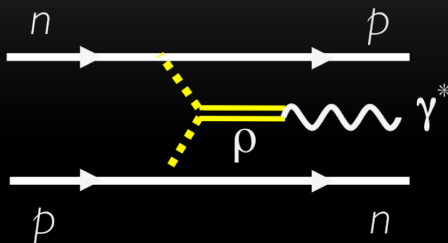
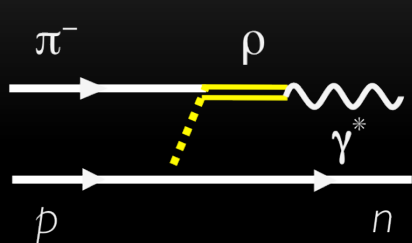
→ D_{13} (1520) is dominant in ρ production!

$$\frac{d\sigma}{dM_{ee}} = \frac{d\sigma}{dM_{\pi\pi}} c_p \left(\frac{m_\rho}{m_{ee}} \right)^3$$

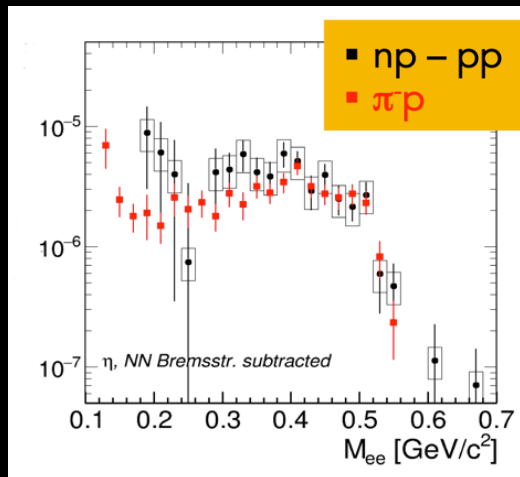
→ First measurement demonstrating the dominant role of intermediate ρ propagation in $N^* \rightarrow N e^+ e^-$ transition

The role of virtual pions in dilepton production

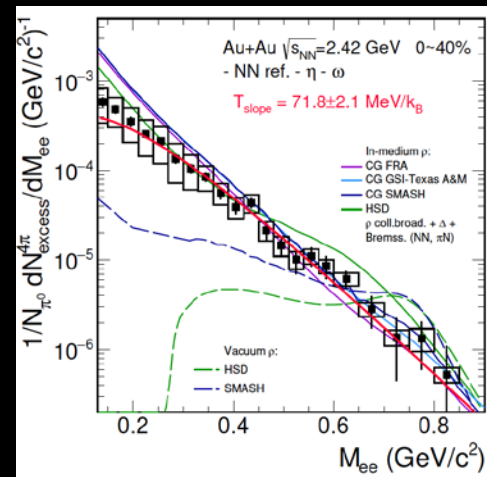
Three different collision systems, three surprises but likely the same underlying mechanism



HADES, in preparation



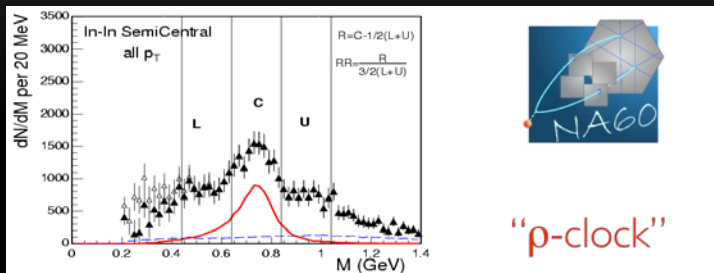
HADES Coll., EPJ A53 (2017) 149
HADES Coll., PRC 95 (2017) 065205



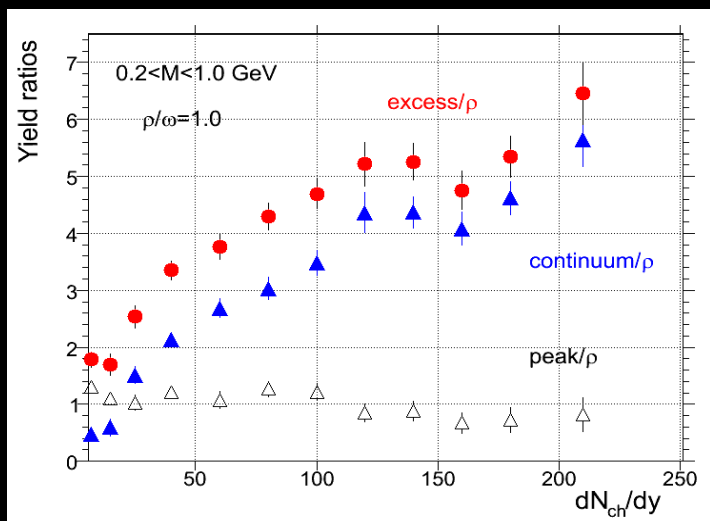
HADES, to appear soon

The dilepton clock

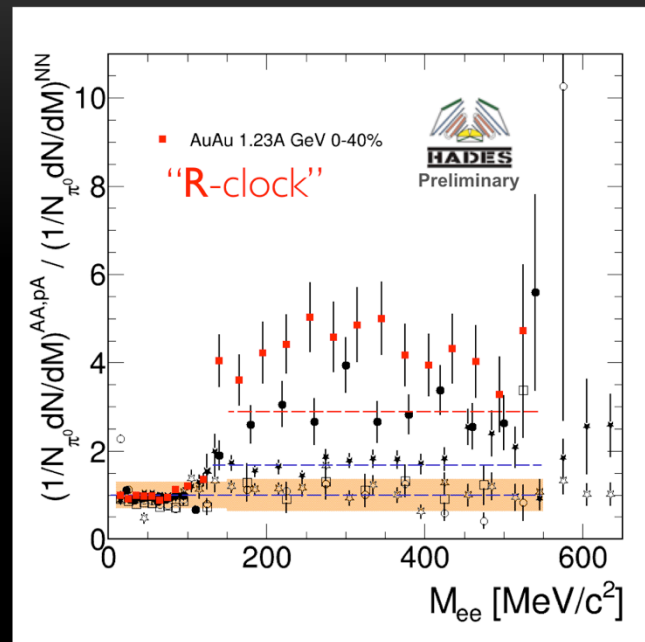
Centrality dependence of spectral shape



"p-clock"



System size dependence of excess



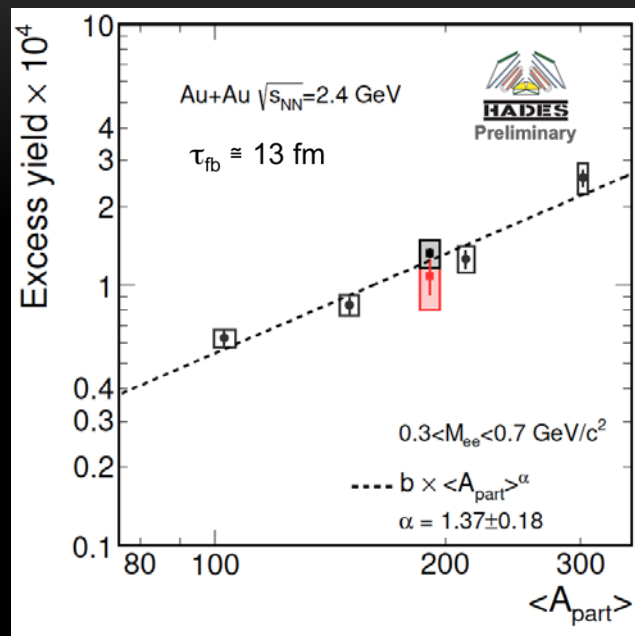
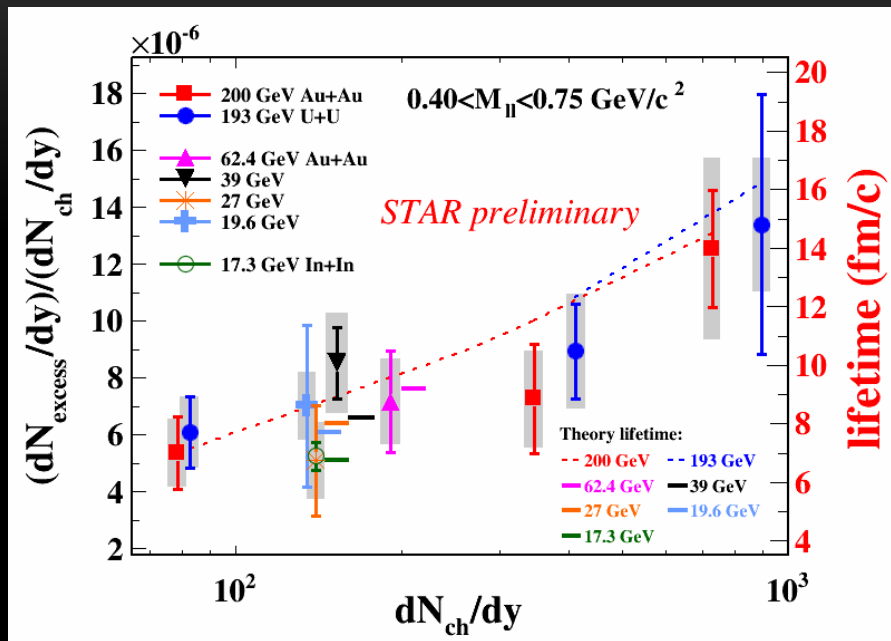
Rapid increase of relative yield reflect the number of ρ 's / R's regenerated in fireball



U.W. Heinz and K. S. Lee, PLB 259, 162 (1991)

H.W. Barz, B. L. Friman, J. Knoll and H. Schulz, PLB 254, 315 (1991)

The lifetime of the interacting fireball

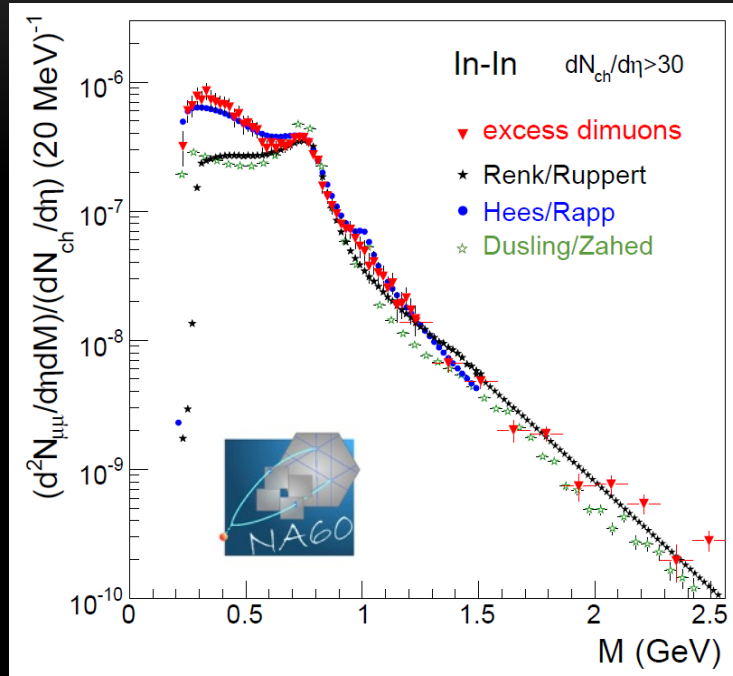


STAR data: arXiv:1612.05484 [nucl-ex]
Model: R. Rapp, H. van Hees, PLB 753 (2016) 586

- Normalized **excess yields** for the mass region $0.3 < M < 0.7$ GeV/c² is proportional to the **lifetime** of the interacting fireball
- Note: normalization to the number of pions should be done when going down with energy!

Dileptons as thermometer

Measurement of radiating source temperature



- Acceptance corrected excess yield
- $M_{\mu\mu} > 1 \text{ GeV}/c^2 \sim$ exponential fall-off - 'Planck-like'
- measurement of radiating source temperature

→ fit $\frac{dN}{dM} \sim M^{\frac{3}{2}} \times \exp\left(-\frac{M}{T}\right)$ to range

→ $M = 1.1 - 2.0 \text{ GeV}/c^2$: $\langle T \rangle_{\text{emitting source}} = 205 \pm 12 \text{ MeV}/k_B$

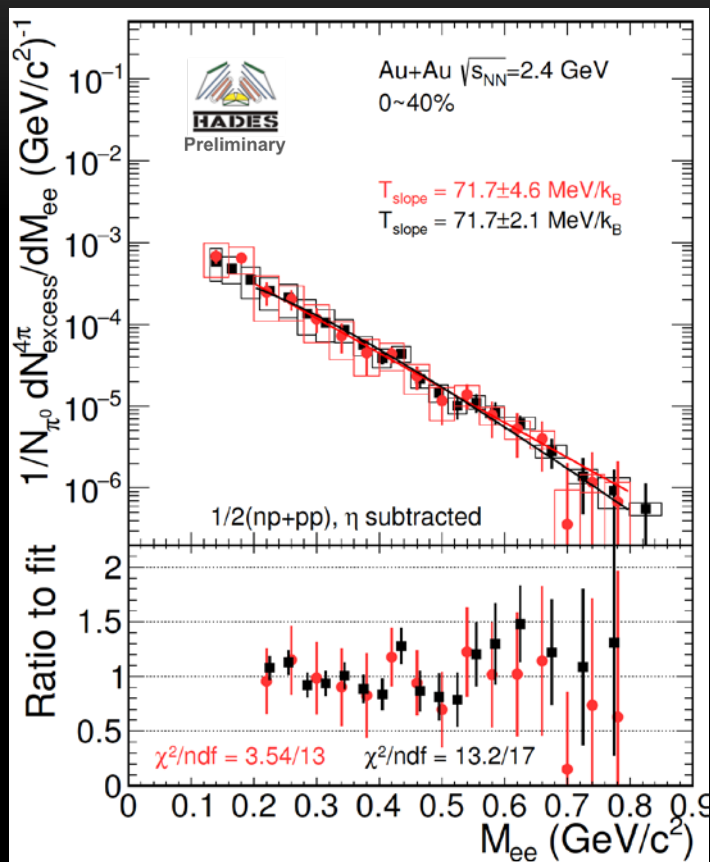
→ $M = 1.1 - 2.4 \text{ GeV}/c^2$: $\langle T \rangle_{\text{emitting source}} = 230 \pm 10 \text{ MeV}/k_B$



Eur. Phys. J. C 59 (2009) 607-623
CERN Courier 111 2009, 31-35
Chiral 2010, AIP Conf. Proc. (2010) 1322

" M_{ll} is the only Lorentz-invariant thermometer of the field"

Dileptons as thermometer at SIS18



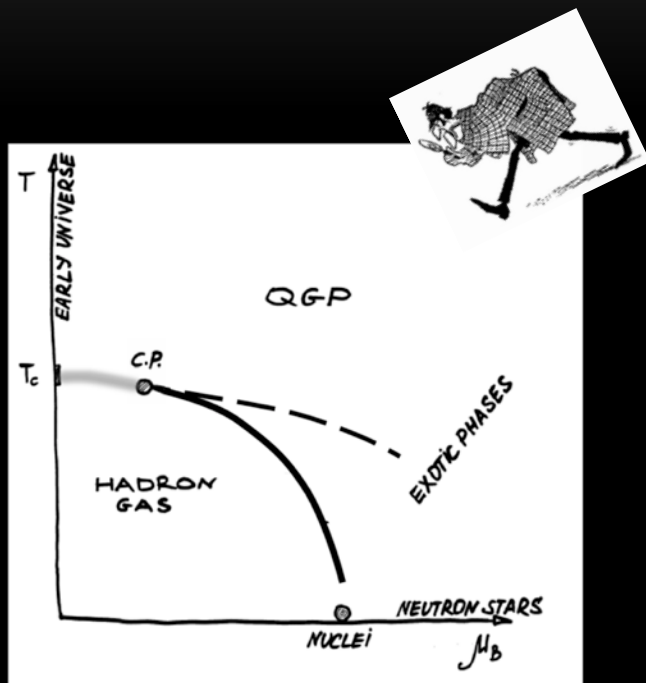
□ $T_{\text{initial}} \sim 3$ times lower at SIS18 compared to SPS

□ Shifts predominance of early emission from $M_{ee} > 1.2 \text{ GeV}/c^2$ at SPS to $M_{ee} > 0.4 \text{ GeV}/c^2$ at SIS

→ fit $\frac{dN}{dM} \sim M^{\frac{3}{2}} \times \exp\left(-\frac{M}{T}\right)$ to range $M=0.2-0.8 \text{ GeV}/c^2$

□ $\langle T \rangle_{\text{emitting source}} = 72 \pm 2 \text{ MeV}/k_B$

Dileptons and QCD phase diagram of matter

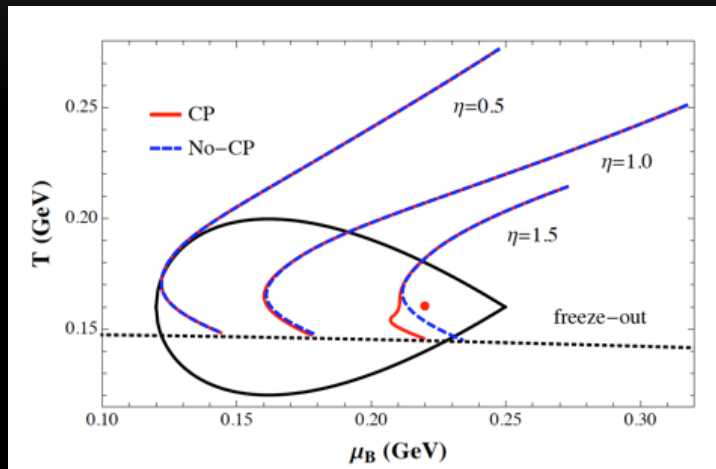


Excitation functions

- Fireball lifetime
- Emitting source temperature

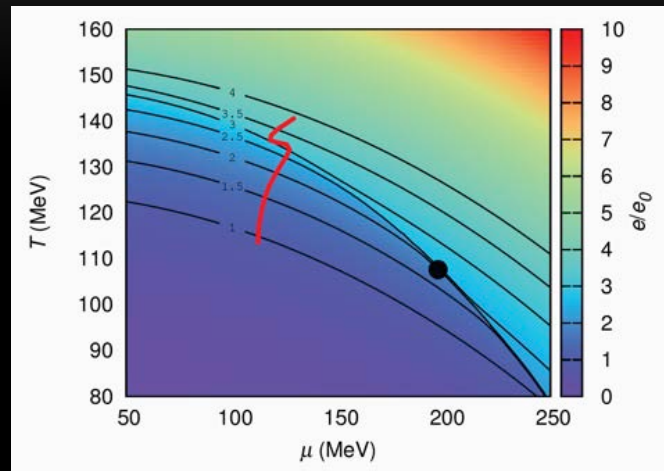
Mapping QCD phase diagram with dileptons

Hydrodynamic evolution trajectories



 A. Monnai, S. Mukherjee, Y. Yin: arXiv: 1606.00771

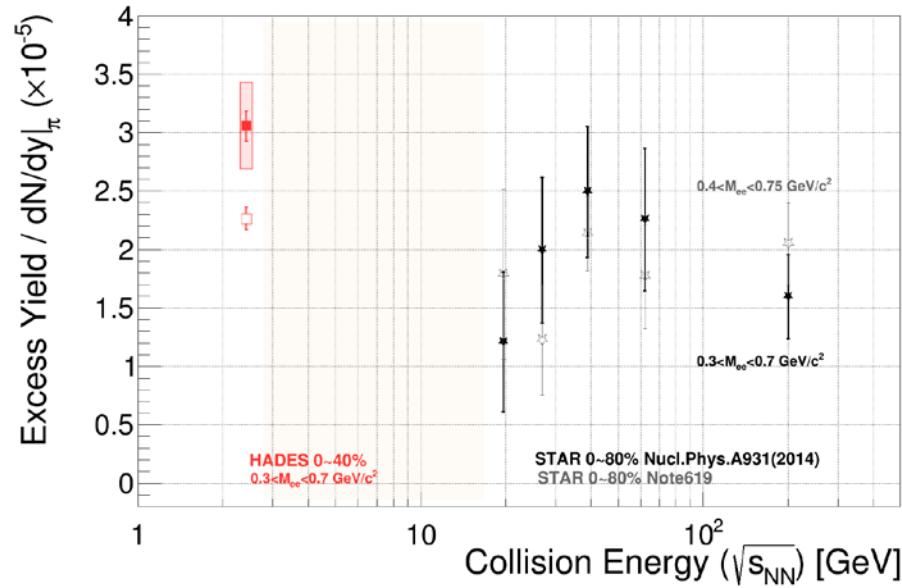
Event-averaged trajectory near the critical point (black dot)



 C. Herold, M. Nahrgang, Y. Yan and C. Kobdaj, PRC93 (2016)

What are the possible signatures in dilepton radiation?

Energy dependence of low-mass excess

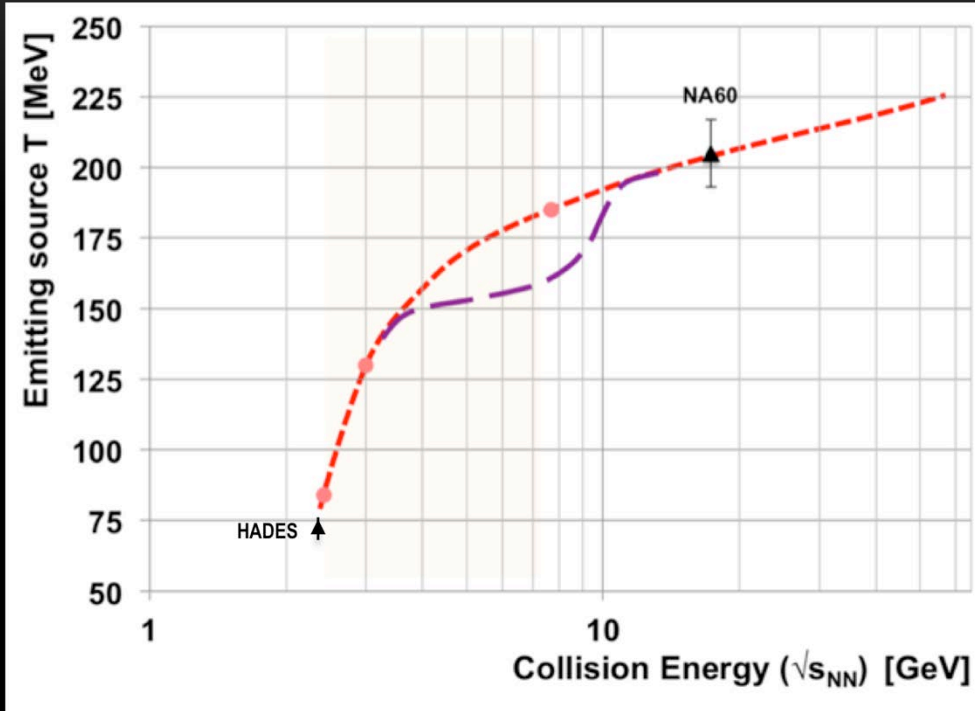


- Yield in low-mass window tracks fireball lifetime
- Search for **anomalous fireball lifetime** around phase transition (& critical point?)

- Quite moderate energy dependence
- Dilepton yield determined by interplay between temperature and $V \otimes \tau_{coll}$



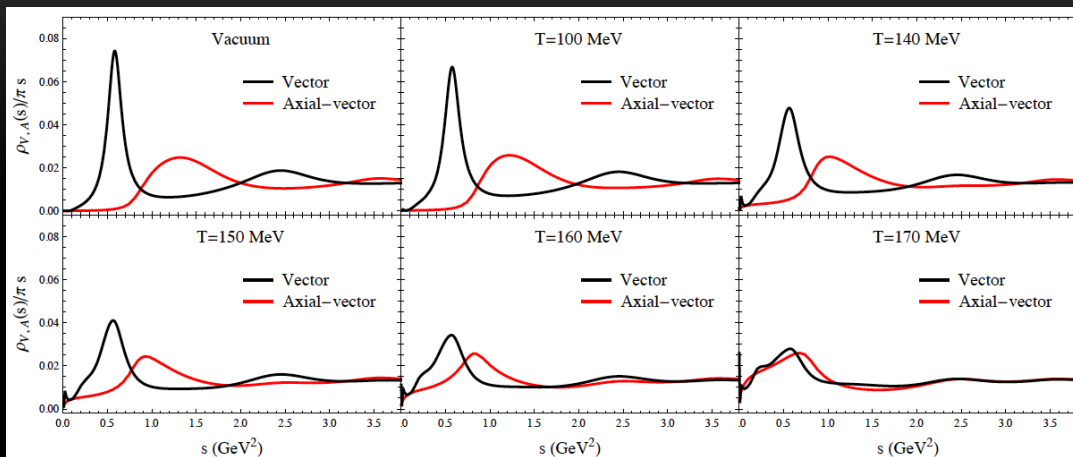
Energy dependence of intermediate mass slope



Dashed violet curve corresponds to a speculated shape with phase transition

- Measures the emitting source temperature (true, no blue shift)
- Measure T_s (note, $T_s < T_{\text{initial}}$) "caloric curve"
- Plateau around onset of deconfinement?
[see e.g. M. D'Agostino et al. NPA 749 (2005) 5533]
- Precision measurements are the key
→ CBM at FAIR

QCD and Weinberg sum rules in medium



Hohler & Rapp, Phys.Lett. B731 (2014)

Jung, Tripolt, von Smekal, Wambach, Phys.Rev. D95, 036020 (2017)

- Finite T vector and axial-vector spectral functions
- No baryon effects accounted for yet
- Chiral mass splitting “burns off”, resonances melt

CBM case: 4π processes: $\pi a_1 \rightarrow \gamma^* \rightarrow l^+ l^-$ (chiral mixing) is a dominant hadronic source in IMR; correlated charm contribution, Drell-Yan, QGP are suppressed

“You may say I’m a dreamer... but I’m not the only one”

		2016	2017	2018	2019	2020	2021	2022	2023	2024	2025
HADES	SIS18			2.5 GeV + pA, pA							
STAR BES	RHIC				8 – 20 GeV			200 GeV			
MPD	NICA				4 – 11 GeV						
CBM	SIS100										2 – 5 GeV
ALICE*	LHC	5.5 TeV					5.5 TeV				
	J-PARC**			pA 8 GeV							2 – 5 GeV
NA60+	SPS		LOI						6 – 17 GeV		
PHENIX	RHIC	200 GeV									
STAR	RHIC	200 GeV	pp(500)								
sPHENIX								200 GeV			
BM@N	Nuclotron		2 – 3 GeV								

* - ITS, 50kHz, lower field

** - Proposal to J-PARC → If approved, construction of HI injector and detectors in 10 years ?

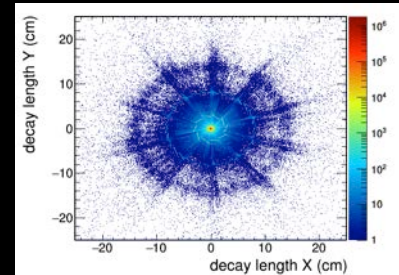
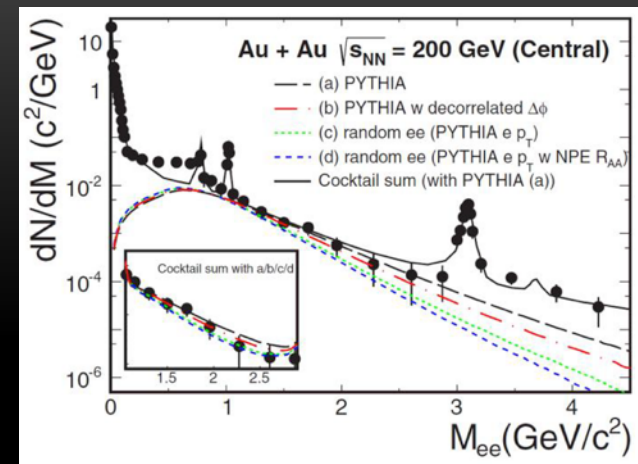
STAR dielectrons ongoing analysis

Au+Au $\sqrt{s_{NN}} = 200$ GeV

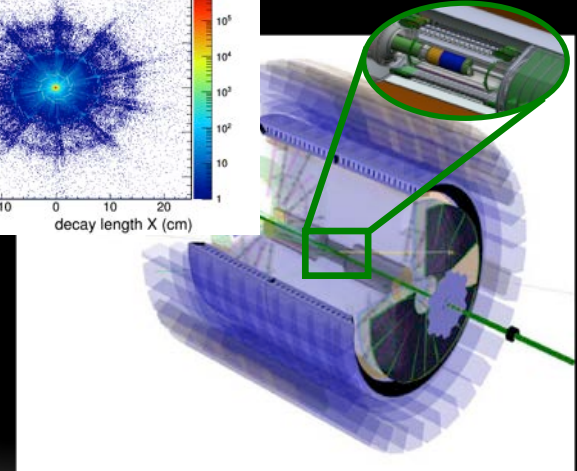
- $\mu_B \ll T$, i.e. vanishing net-baryon density
- Lattice QCD computations are most powerful
- Measure ρ spectral function and “calibrate” EM rates
- Extract fireball temperature (IMR)

Dielectron reconstruction with HFT

- Reduce uncertainties on charm contribution
- **Challenge:** γ conversion in HFT detector material
→ use excellent vertexing to reject it
(looks promising, statistics Runs 10,11,14,16 $\sim 4 \times 10^9$)



Heavy Flavor Tracker

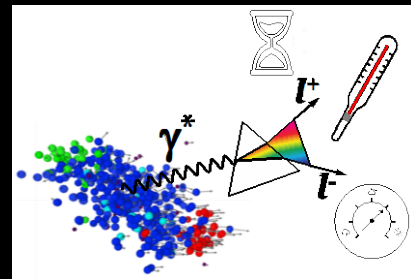


Résumé and prospects

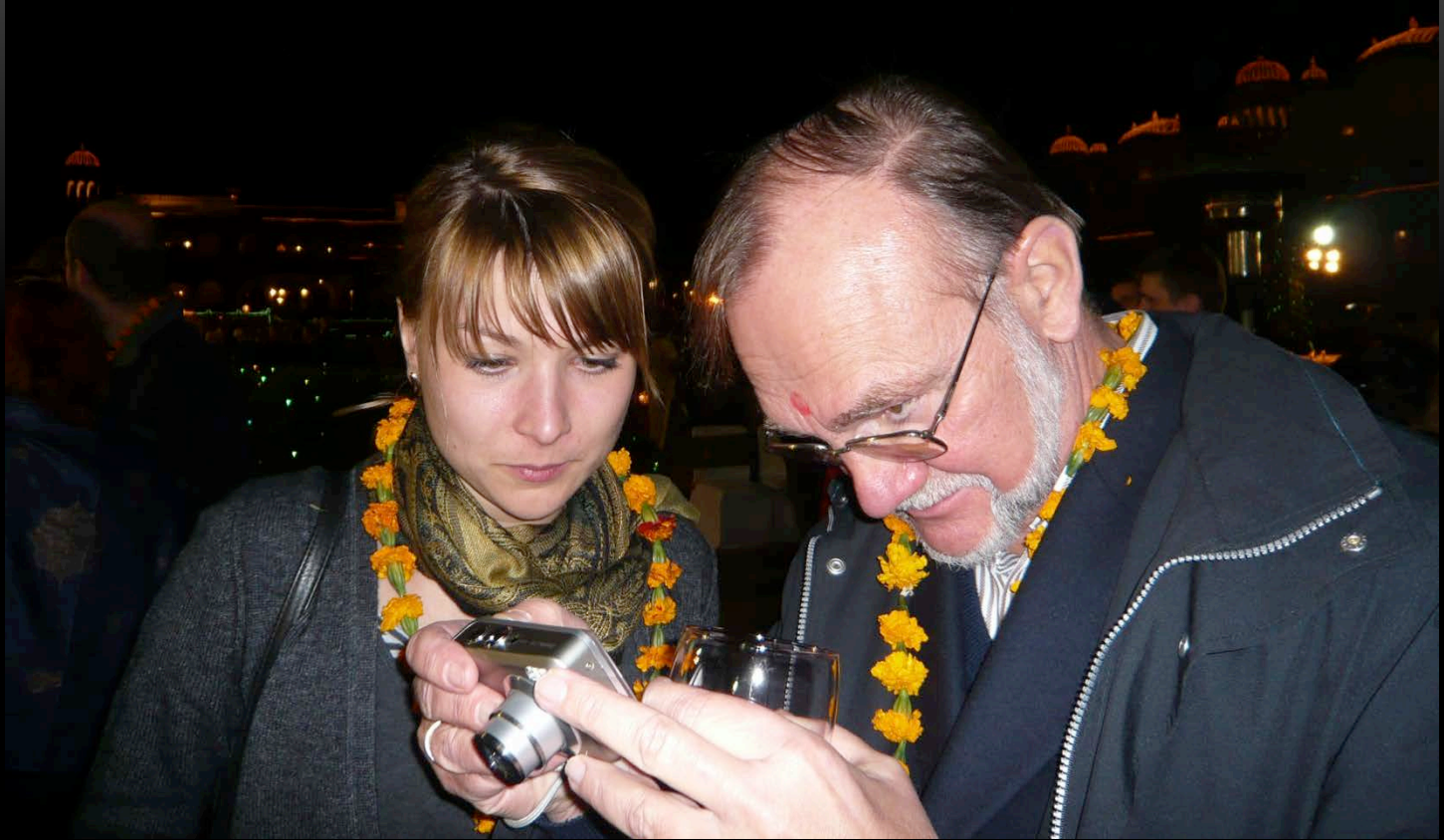
- **Unique** possibility of characterizing properties of hot and dense QCD matter with dileptons
- **Robust** understanding of low-mass dilepton excess radiation by ρ -baryon coupling (at top RHIC, RHIC BES, SPS and SIS18 energies)
- Enable unique measurements
 - Degrees of freedom of the medium
 - Restoration of chiral symmetry
 - Fireball lifetime
 - Emitting source temperature
- Future experiments allows for overlap and independent confirmation of results



There is no mission impossible



Thank you
for your attention!



There is a power that can take away life,
but never the memory of friends

Bonus slides

QCD and Weinberg sum rules

Weinberg Sum Rules...

 Weinberg '67, Das et al '67

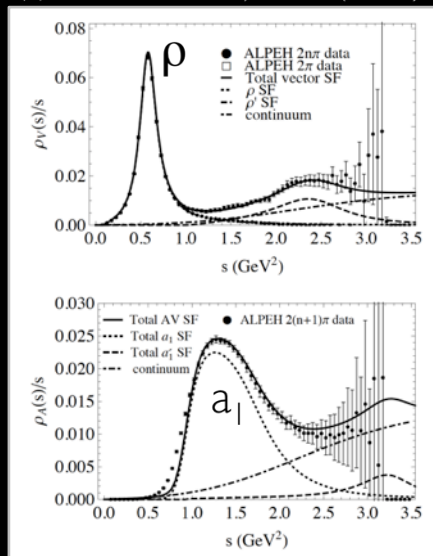
$$\int \frac{ds}{\pi} \frac{1}{s} (\rho_V - \rho_A) = f_\pi^2$$

$$\int \frac{ds}{\pi} (\rho_V - \rho_A) = -m_q \langle \bar{q}q \rangle$$

$$\int \frac{ds}{\pi} s (\rho_V - \rho_A) = c \alpha_s \langle (\bar{q}q)^2 \rangle$$

... accurately satisfied in vacuum

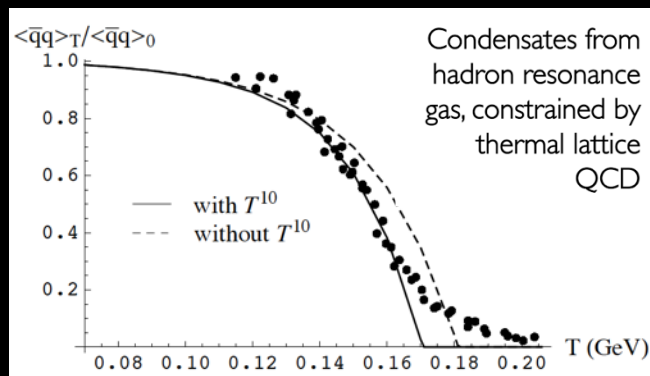
 Rapp et al, Annals Phys. 368 (2016)



ρ - a_1 mass splitting due to χ_s breaking ($\sim f_\pi \langle \bar{q}q \rangle$)

... remain valid in medium

 J. Kapusta, E. Suvayak '94



→ Test in-medium ρ - a_1 spectral function

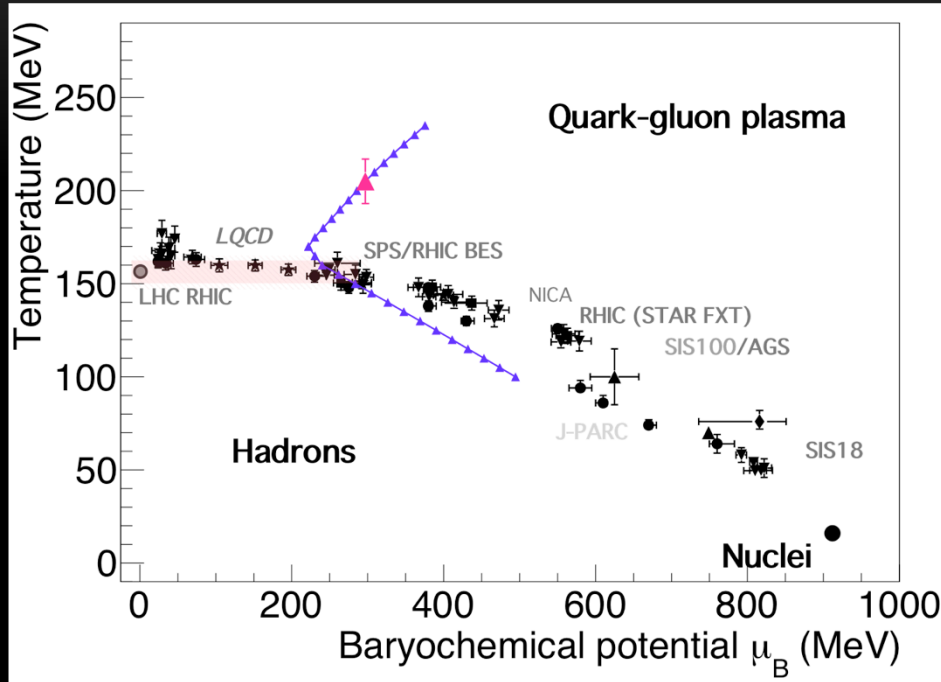
Components of EM probes

- Degrees of freedom of the medium:
 - Spectral function merges into QGP description
 - Direct evidence for transition hadrons → quarks & gluons
- Restoration of chiral symmetry
 - Condensates constrained by Lattice-QCD
 - Chiral mass splitting “burns off” → resonances melt
- Phenomenological tools → excitation functions
 - Fireball lifetime
 - Emitting source temperature
- Transport properties
 - Electric conductivity → probes soft limit of EM SF



$$\sigma_{EM}(T) = -e^2 \lim_{q_0 \rightarrow 0} \left[\frac{\partial}{\partial q_0} \text{Im} \Pi_{EM}(q_0, q = 0; T) \right]$$

Trajectories in the phase diagram



- Hadron production: well described by hadron gas in chemical and thermal equilibrium

$$n_i(\mu_B, T) = d_i \int \frac{d^3k}{(2\pi)^3} f(E_k; \mu_B, T)$$

- Dileptons - need to construct evolution:
 - Before up to earliest “formation” time
 $\tau_0 \leftrightarrow T_0 > T_{\text{chem}}$
 - After: down to thermal freezeout
 $\tau_f \leftrightarrow T_{f0} < T_{\text{chem}}$

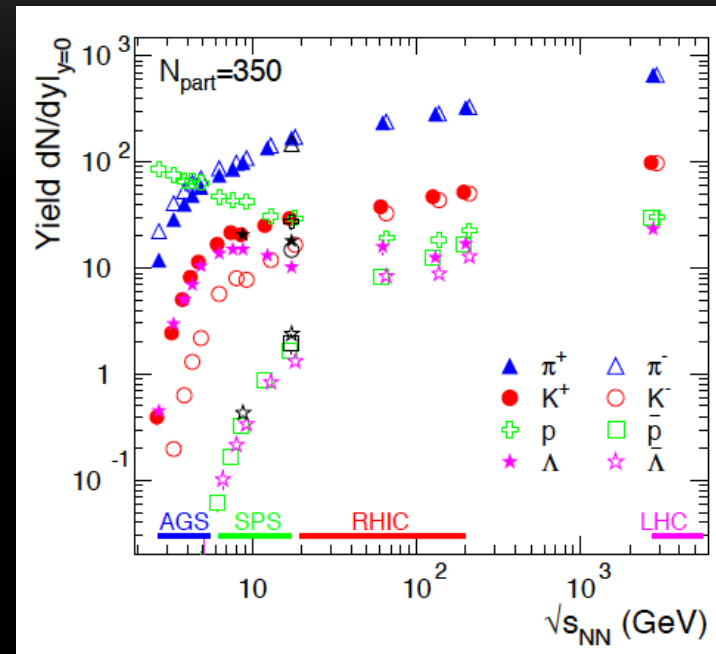
- Basic assumption: entropy (+baryon-number) conservation $\rightarrow$ fixes $T(\mu_B)$ in the phase diagram



From SPS to RHIC

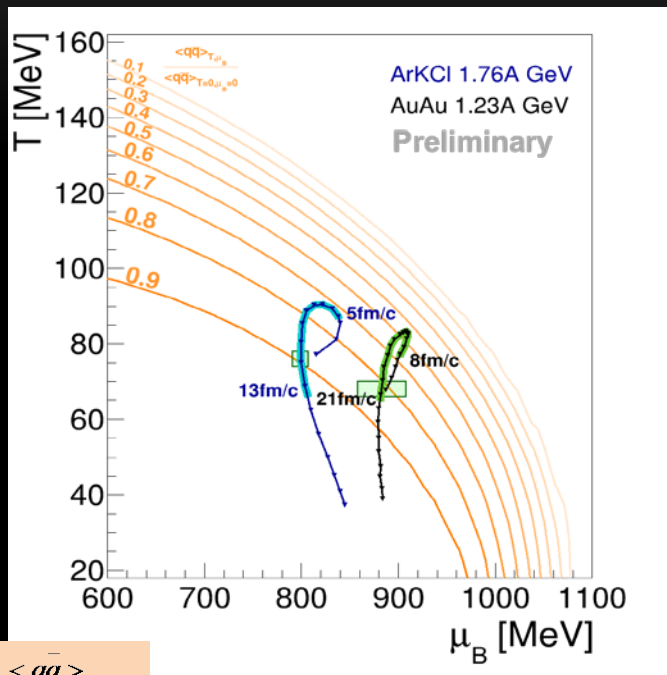
	SPS (Pb+Pb)	RHIC (Au+Au)
$dN(\bar{p})/dy$	6.2	20.1
produced baryons ($p, \bar{p}, n, \bar{n}$)	24.8	80.4
$p - \bar{p}$	33.5	8.6
participating nucleons $(p - \bar{p})A/Z$	85	21.4
total baryon number	110	102

- Although the NET baryon density is different at SPS and RHIC, baryon density is practically the same!
- Baryon effects important even at $\rho_{B,tot} = 0$: sensitive to $\rho_{B,tot} = \rho_B + \rho_{\bar{B}}$ (ρ -N and ρ - $\bar{N}$ interactions identical)
- Higher initial temperature at RHIC



A. Andronic, arXiv:1407.5003

HADES and QCD phase diagram of matter



- Chemical freeze-out from measured particle yields analyzed with SHM
- Trajectories extracted from inner cube of cells with coarse-grained UrQMD
- Time-window of dilepton emission → access to hot and dense stage of the heavy-ion collision

$T_{\text{max}} = 85 \text{ MeV}, \rho_{\text{max}} = 3 \rho_0$
 → Excitation of the vacuum
 (melting of condensate)
 matches spectral medium effects!

$$\frac{\langle \bar{q}q \rangle(T, \mu_B)}{\langle \bar{q}q \rangle} = 1 - \sum_h \frac{\rho_h^s \Sigma_h}{m_\pi^2 f_\pi^2}$$

$\frac{\langle \bar{q}q \rangle_{T, \mu_B}}{\langle \bar{q}q \rangle_{T=0, \mu_B=0}}$: B.J. Schaefer and J. Wambach



Ar+KCl data: Nucl. Phys. A931 (2014) c785