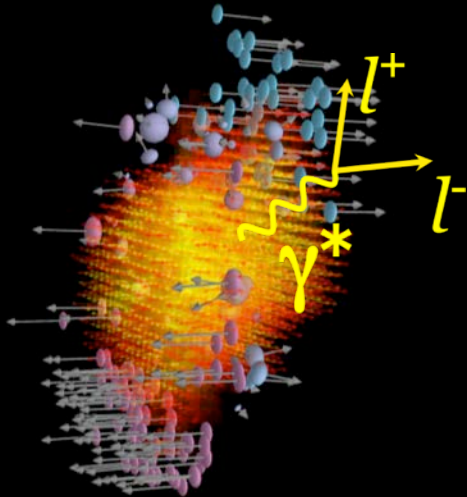




WE-Heraeus Physics School

QCD – Old Challenges and
New Opportunities

Bad Honnef, Sept 24–30, 2017



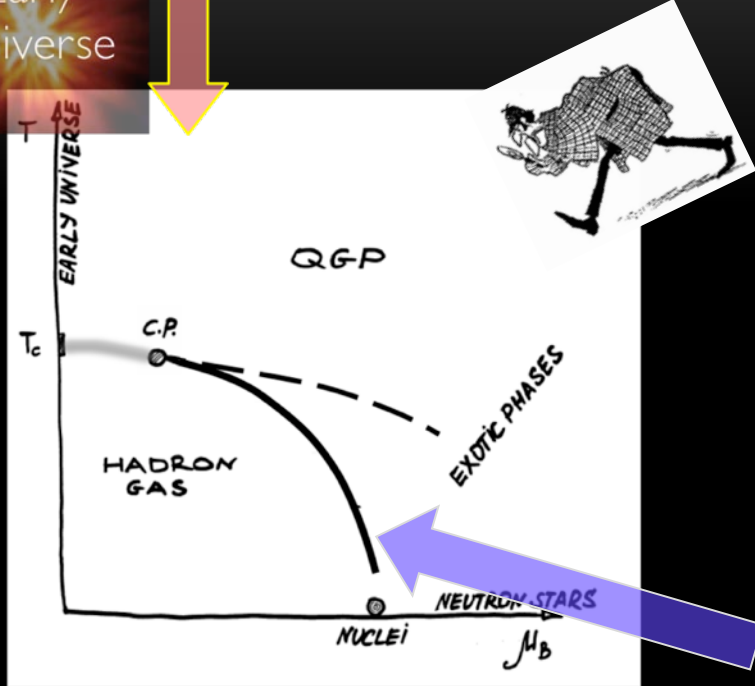
Experimental search for signals of chiral symmetry restoration in heavy-ion collisions

Tetyana Galatyuk

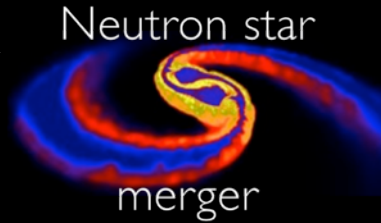
Technische Universität Darmstadt /
GSI Helmholtzzentrum für Schwerionenforschung

Exploring the phase diagram of QCD matter

Early Universe



- What are the properties of matter under extreme temperatures and densities?
- Where are the phase boundaries located?
- Is there a critical point?
- Where are the limits of hadronic existence?



merger

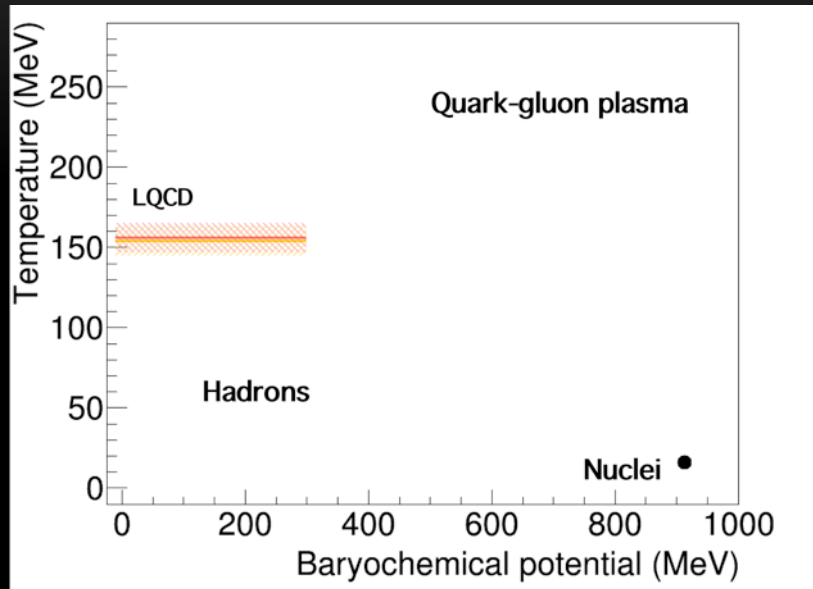
Rosswog



supernova

NASA, SN1987A

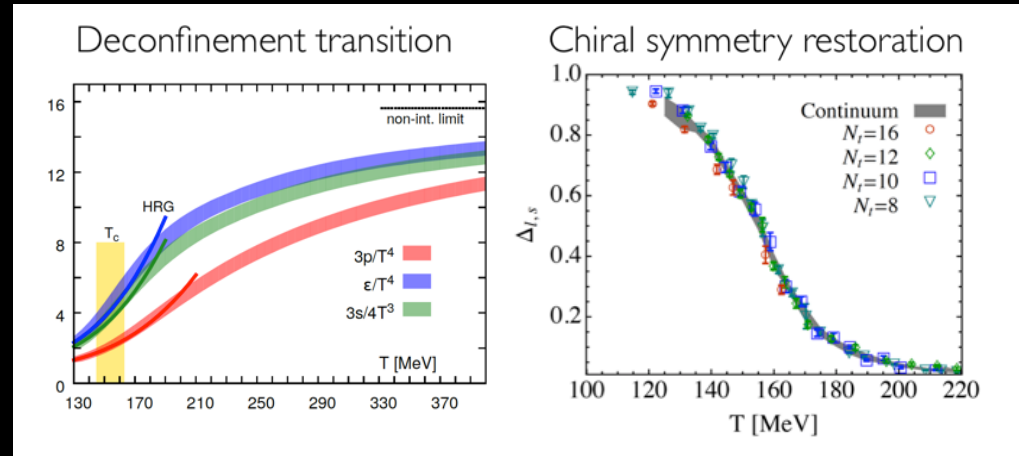
Theoretical guidance



□ Vanishing μ_B , high T (Lattice QCD)

□ Crossover transition

□ $\varepsilon_c \sim 1 \text{ GeV}/\text{fm}^3, T_c \sim 155 \text{ MeV}$

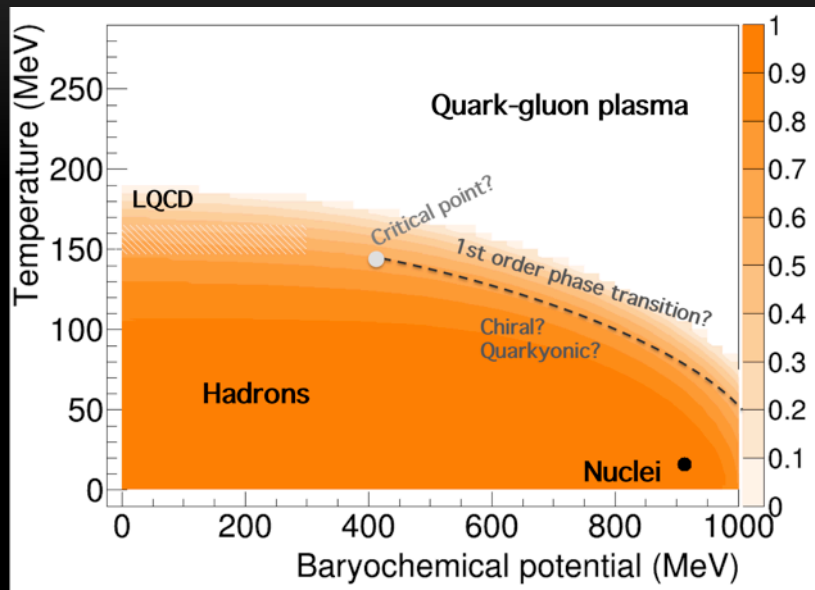


O. Kaczmarek et al., PRD 83 (2011) 014504

S. Borsanyi et al. [Wuppertal-Budapest Coll.], JHEP 1009 (2010) 073

A. Bazavov et al. [Hot QCD Coll.], PRD90 (2014) 094503

Theoretical guidance



□ Vanishing μ_B , high T (Lattice QCD)

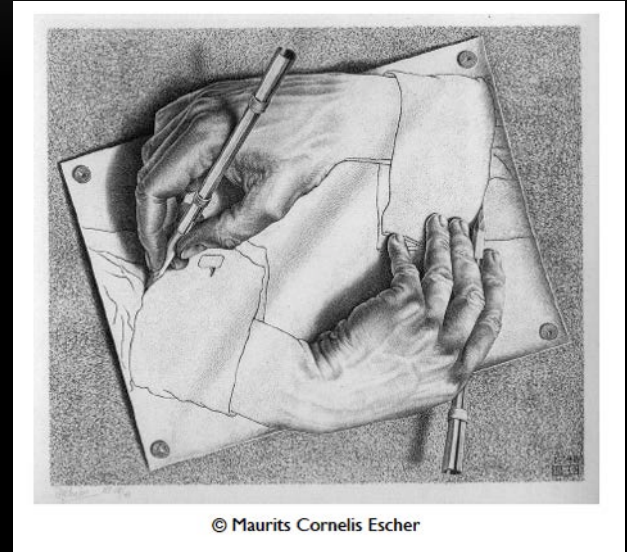
- Crossover transition
- $\varepsilon_c \sim 1 \text{ GeV}/\text{fm}^3, T_c \sim 155 \text{ MeV}$

□ Large μ_B , moderate T (effective, Lattice QCD inspired models)

- 1st order transition
- QCD critical point
- Melting of the condensate (order parameter $\langle 0 | \bar{q}q | 0 \rangle = \langle 0 | \bar{q}_L q_R + \bar{q}_R q_L | 0 \rangle \neq 0$)

$$\frac{\langle \bar{q}q \rangle_{T, \mu_B}}{\langle \bar{q}q \rangle_{T=0, \mu_B=0}} : \text{B.J. Schaefer and J. Wambach}$$

Chiral symmetry in QCD



Chiral Symmetry in QCD: Vacuum

$$\mathcal{L}_{QCD} = \bar{q} (i\not{\partial} + g\not{A} - \hat{m}_q) q - \frac{1}{4} G_{a\mu\nu}^2$$

Current quark masses: $m_u \approx m_d \approx 5-10$ MeV

Chiral $SU(2)_V \times SU(2)_A$ transformation up to $\mathcal{O}(m_q)$, L_{QCD} invariant under

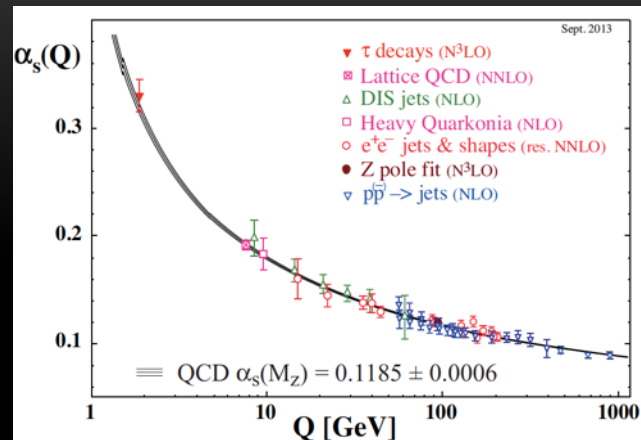
$$q \mapsto R_V(\vec{\alpha}_V) q = \exp(-i\vec{\alpha}_V \cdot \vec{\tau} / 2) q$$

$$q \mapsto R_A(\vec{\alpha}_A) q = \exp(-i\gamma_5 \vec{\alpha}_A \cdot \vec{\tau} / 2) q$$

Rewrite L_{QCD} using $q_{L,R} = (1 \pm \gamma_5) / 2 q$:

$$\mathcal{L}_{QCD} = (\bar{u}_L, \bar{d}_L) i\not{D} \begin{pmatrix} u_L \\ d_L \end{pmatrix} + (\bar{u}_R, \bar{d}_R) i\not{D} \begin{pmatrix} u_R \\ d_R \end{pmatrix} + \mathcal{O}(m_q) - \frac{1}{4} G_{a\mu\nu}^2$$

$$q_{L,R} \mapsto \exp(-i\vec{\alpha}_{L,R} \cdot \vec{\tau} / 2) q_{L,R}$$



Invariance under
isospin and “handedness”

Chiral symmetry breaking

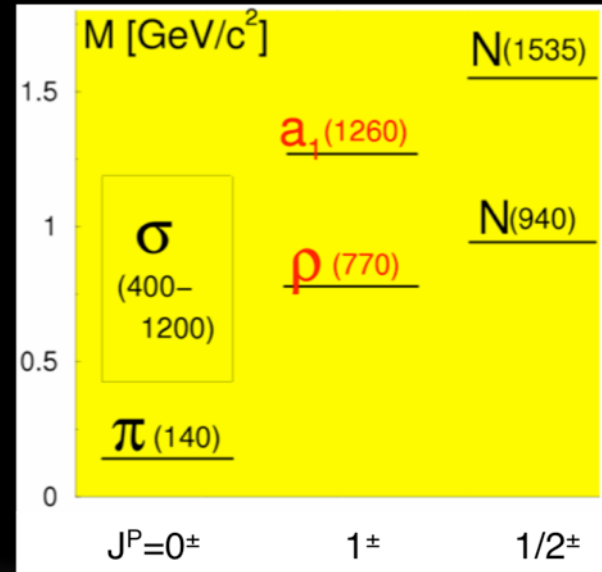
Chiral symmetry is **explicitly broken** by the finite masses of the current (u, d, s) quarks

On top of this, chiral symmetry is **spontaneously broken**: strong $q\bar{q}$ attraction \rightarrow Chiral Condensate fills QCD vacuum $\langle \bar{q}q \rangle = \langle \bar{q}_L q_R + \bar{q}_R q_L \rangle \approx -[(229 \pm 9) \text{MeV}]^3$ an order parameter of chiral symmetry

- \rightarrow Mass generation
- \rightarrow $\langle \bar{q}q \rangle$ is not an observable!


But: **hadronic excitations reflect spontaneous breaking of chiral symmetry:**

- “massless” Goldstone bosons $\pi^{0,\pm}$
(explicit breaking: $f_\pi^2 m_\pi^2 = m_q \langle \bar{q}q \rangle$)
- “chiral partners” split: $\Delta M \approx 0.5 \text{GeV}$



QCD and Weinberg sum rules

Weinberg Sum Rules...


 Weinberg '67, Das et al '67

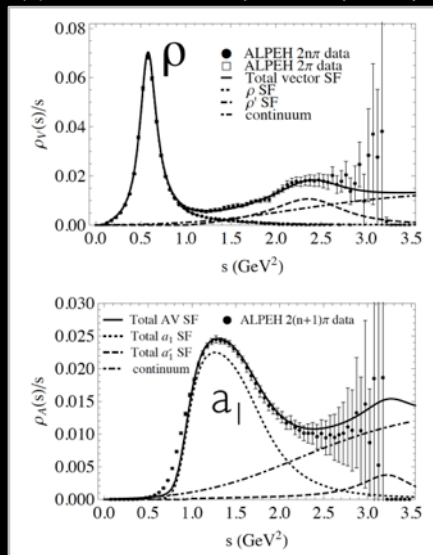
$$\int \frac{ds}{\pi} \frac{1}{s} (\rho_V - \rho_A) = f_\pi^2$$

$$\int \frac{ds}{\pi} (\rho_V - \rho_A) = -m_q \langle \bar{q}q \rangle$$


$$\int \frac{ds}{\pi} s (\rho_V - \rho_A) = c\alpha_s \langle (\bar{q}q)^2 \rangle$$

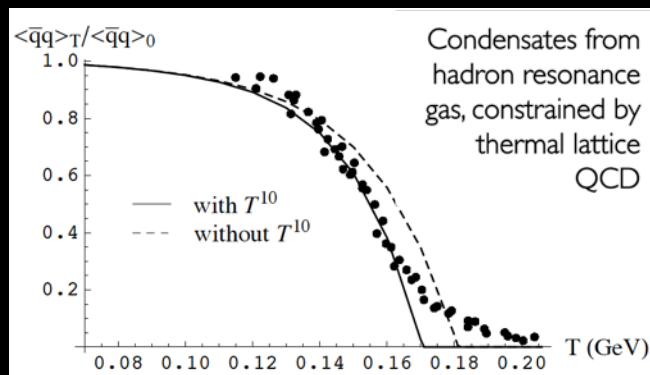
... accurately satisfied in vacuum

 Rapp et al, *Annals Phys.* 368 (2016)



... remain valid in medium

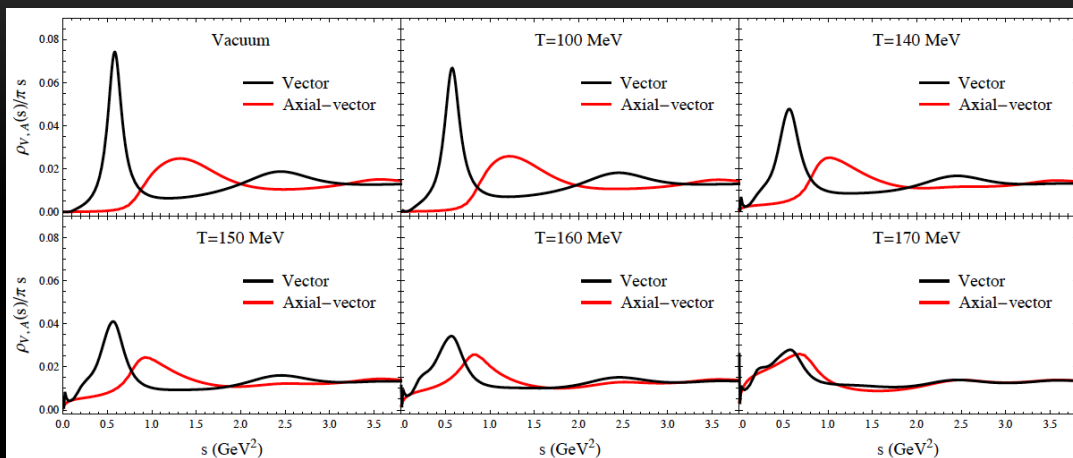
 J. Kapusta, E. Surya '94



ρ - a_1 mass splitting due to χ_s breaking ($\sim f_{\pi T} \langle \bar{q}q \rangle$)

→ Test in-medium ρ - a_1 spectral function

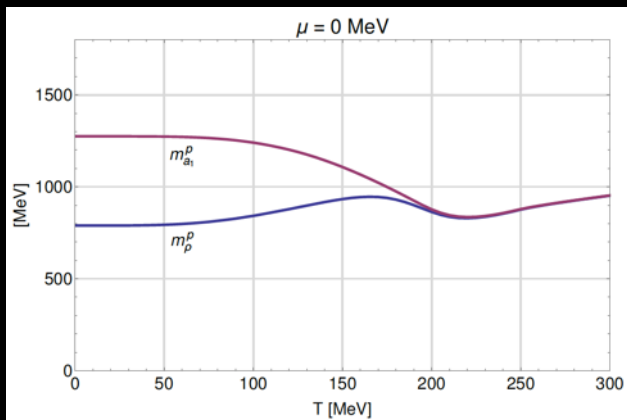
QCD and Weinberg sum rules in medium



- Finite T vector and axial-vector spectral functions
- No baryon effects accounted for yet
- Chiral mass splitting “burns off”, resonances melt



Hohler & Rapp, *Phys.Lett. B731* (2014)



- Degeneracy of hadronic chiral partners at finite T from Functional Renormalization Group (FRG)



Jung, Tripolt, von Smekal, Wambach, *Phys.Rev. D95, 036020* (2017)

Experimental approach: high energy heavy-ion collisions



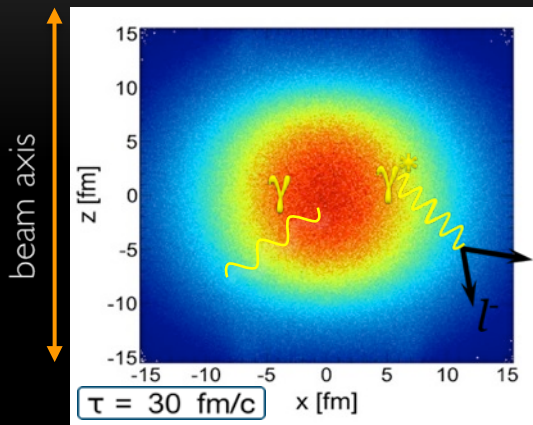
Systematic experimental measurements (E_{beam}, A)

→ Extract numbers that might be related to the QCD phase diagram

→ Objective: use dileptons to probe the nature of Strongly Interacting Matter

Electromagnetic radiation

Photons and lepton pairs probe the interior of fireballs – “PET” of the fireball

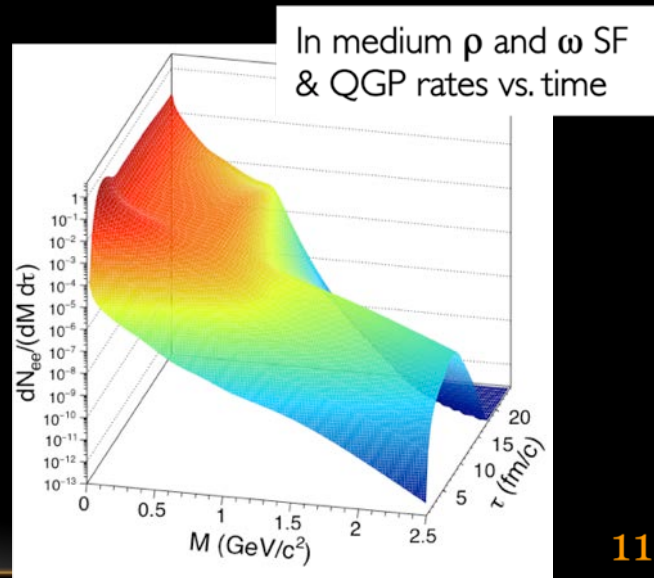


- EM radiation contains contributions from throughout the collision
- EM probes **leave collision zone undisturbed**
- Real γ characterized by transverse momentum
- Lepton pairs carry extra information: invariant mass

The vector correlator is directly accessible in HIC

$$\frac{dN_{ll}}{d^4x d^4q} = \frac{-\alpha_{EM}^2}{\pi^3 M^2} f^B(q \cdot u; T) \text{Im} \Pi_{EM}^{\mu\nu}(M, q; \mu_B, T)$$

→ Unique direct access to in-medium spectral function



Electromagnetic correlator in vacuum

McLerran-Toimela formula

$$\frac{dN_{||}}{d^4x d^4q} = \frac{-\alpha_{EM}^2}{\pi^3 M^2} f^B(q \cdot u; T) \text{Im} \Pi_{EM}^{\mu\nu}(M, q; \mu_B, T)$$

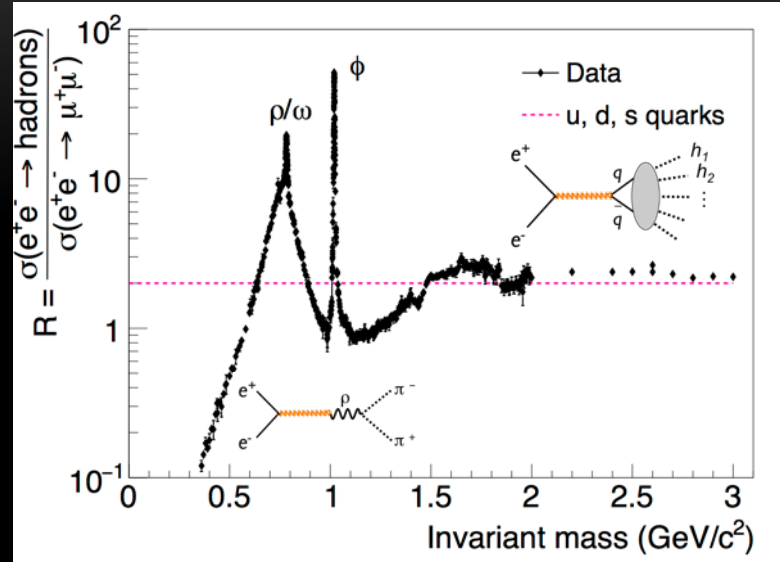
$$R = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} \propto \frac{\text{Im} \Pi_{EM}^{vac}}{M^2}$$

- If $\frac{\text{Im} \Pi_{EM}^{\mu\nu}}{M^2} \sim \text{const.}$ thermal dilepton emission will follow a Bose distribution

→ Intermediate mass range as thermometer

$$\text{Im} \Pi_{EM}^{vac}(M \leq 1.5 \text{ GeV}) = \sum_{v=\rho, \omega, \phi} \left(\frac{m_v^2}{g_v} \right)^2 \text{Im} D_v^{vac}(M)$$

$$\text{Im} \Pi_{EM}^{vac}(M > 1.5 \text{ GeV}) = -\frac{M^2}{12\pi} \left(1 + \frac{\alpha_s(M)}{\pi} + \dots \right) N_c \sum_{q=u,d,s} (e_q)^2$$



→ Vector Meson Dominance: $J^P = 1^-$ for both γ^* and VM (with ρ playing a dominant role)

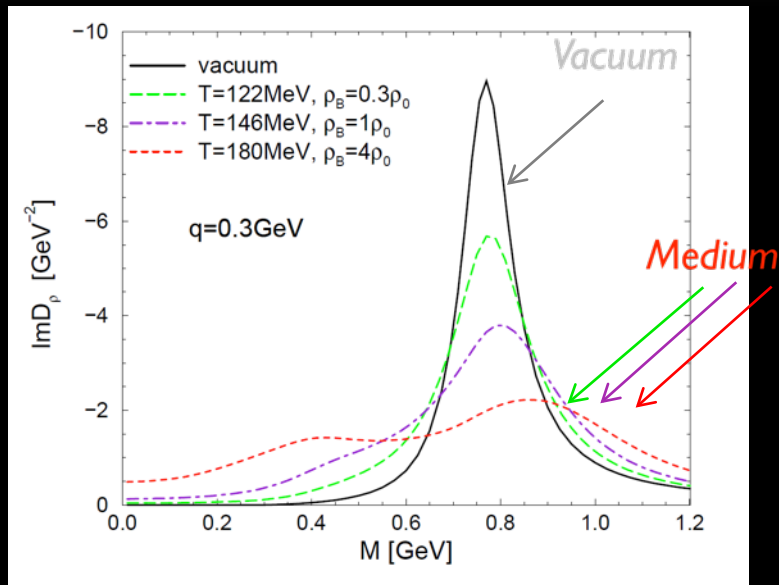
→ pQCD continuum

ρ meson in hot and dense medium

interacts with hadrons from heat bath \rightarrow

additional contributions to the ρ -meson self-energy in the medium

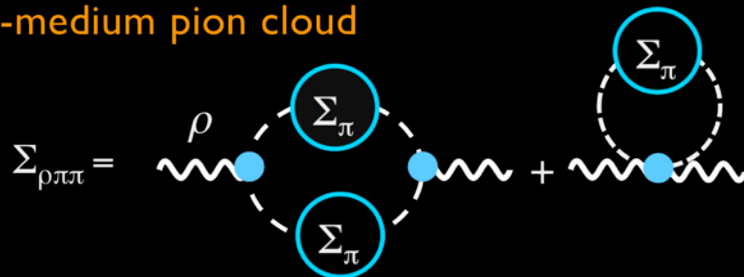
 R. Rapp and J. Wambach, Eur.Phys.J. A6 (1999)



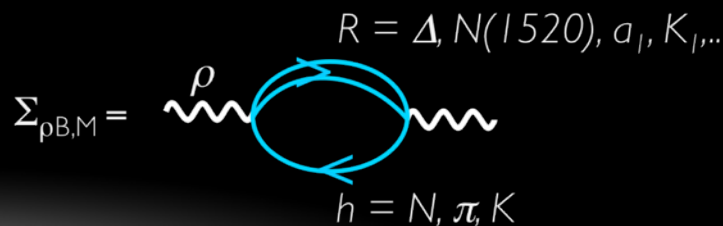
The ρ spectral function **strongly broadens** in the medium because the ρ couples to baryons!

$$D_\rho(M, q; \mu_B, T) = \frac{1}{\left[M^2 - m_\rho^2 - \Sigma_{\rho\pi\pi} - \Sigma_{\rho B} - \Sigma_{\rho M} \right]}$$

In-medium pion cloud



Direct ρ -hadron scattering



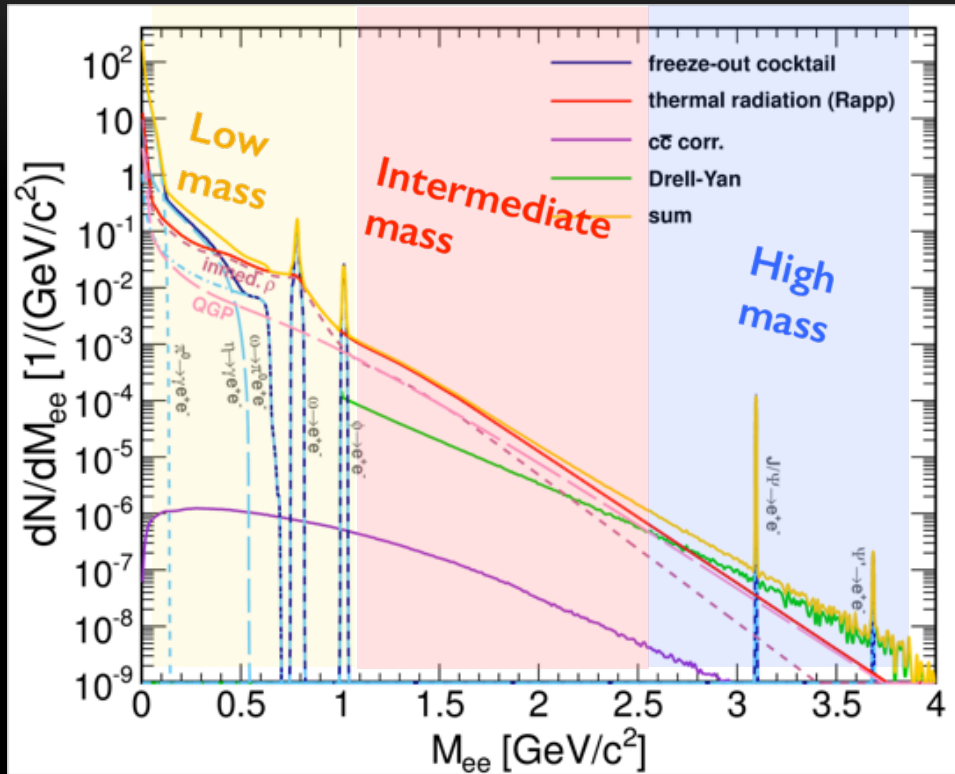
Components of EM probes

- Degrees of freedom of the medium:
 - Spectral function merges into QGP description
 - Direct evidence for transition hadrons \rightarrow quarks & gluons
- Restoration of chiral symmetry
 - Condensates constrained by Lattice-QCD
 - Chiral mass splitting “burns off” \rightarrow resonances melt
- Phenomenological tools \rightarrow excitation functions
 - Fireball lifetime
 - Emitting source temperature
- Transport properties
 - Electric conductivity \rightarrow probes soft limit of EM SF



$$\sigma_{EM}(T) = -e^2 \lim_{q_0 \rightarrow 0} \left[\frac{\partial}{\partial q_0} \text{Im} \Pi_{EM}(q_0, q = 0; T) \right]$$

Characteristic regimes in invariant e^+e^- mass



- Drell-Yan: power-low $\sim M^n$
- Heavy-flavor: $c\bar{c} \rightarrow l^+l^-$
- Thermal radiation: $\sim \exp(-M/T)$
 - QGP – highest T, no flow
 - “4 π annihilation”: $\pi a_1 \rightarrow l^+l^-$
 - In-medium ρ , ω – moderate T, flow
- Final state decays (hadron cocktail): $\pi^0, \eta \rightarrow \gamma e^+e^-$

$$M^2 = (P_{e^+} + P_{e^-})^2$$

The experimental challenge ...

- ❑ Lepton pairs are rare probes ($BR < 10^{-4}$)
- ❑ at SIS energies sub-threshold vector meson production
→ $M_r \times \Gamma_{ee}/\Gamma_{tot}$ decay per 10 mio events
- ❑ Large combinatorial background
 - ❑ in e^+e^- from Dalitz decays ($\pi^0 \rightarrow e^+e^-\gamma$) and conversion pairs (e^+e^-)
 - ❑ in $\mu^+\mu^-$: weak π, K decays
- ❑ Isolate the contribution to the spectrum from the dense stage
- ❑ Low-momentum coverage!

DATA QUALITY

- ❑ The decisive parameters: Number of Interactions and Signal/Background
 - ❑ Range of B/S: 20 - 100 → $B/S \gg 1$;
 - Effective sample size: $S_{eff} \sim 1 \times S/B$
 - Systematics: $\delta S_{eff}/S_{eff} = \delta B/B \times B/S$



There is no such thing as a free lunch



reduction by factors of 20-100

$$\delta B/B = 2 \dots 5 \times 10^{-2}$$

Italian Artist Sven Sachsaber looked for a needle in haystack (November 2014)

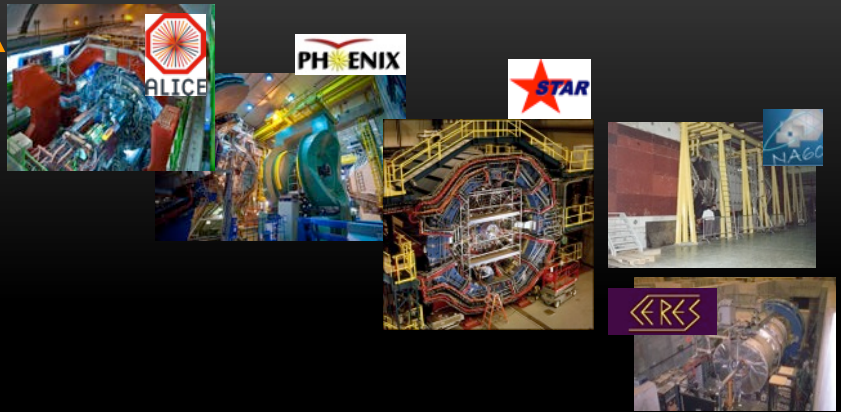


... but after 2 days, he found it !

... he has gone through periods of doubt
and serious discouragement

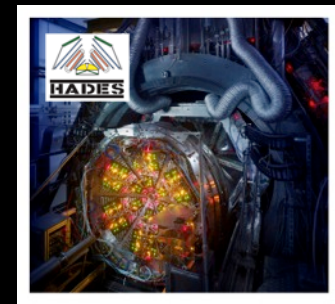


T



Highly interesting results from
LHC, RHIC BES, SPS, SIS18

→ Lepton pairs as true messengers of the dense phase

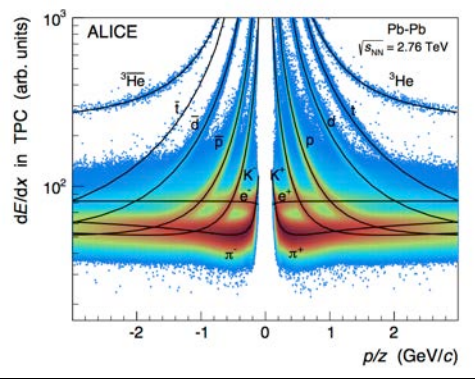


μ_B

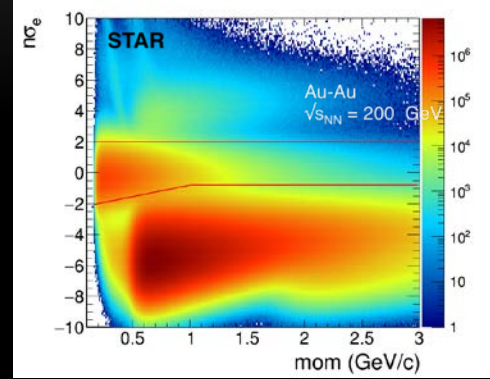
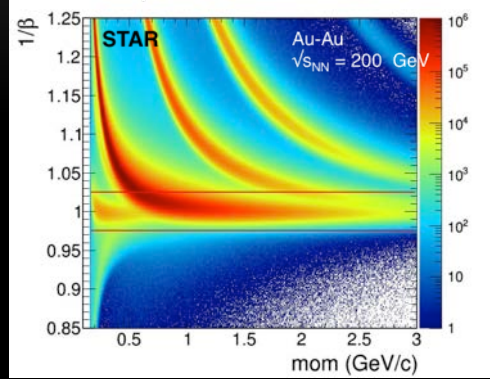
Particle identification

Electron identification by means of: momentum, dE/dx , velocity, RICH information

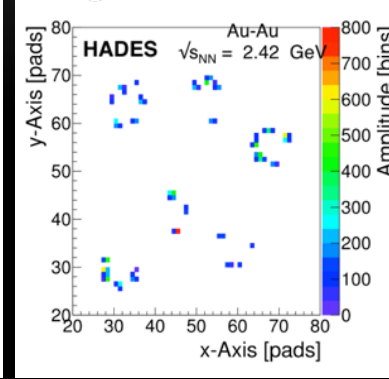
dE/dx in TPC



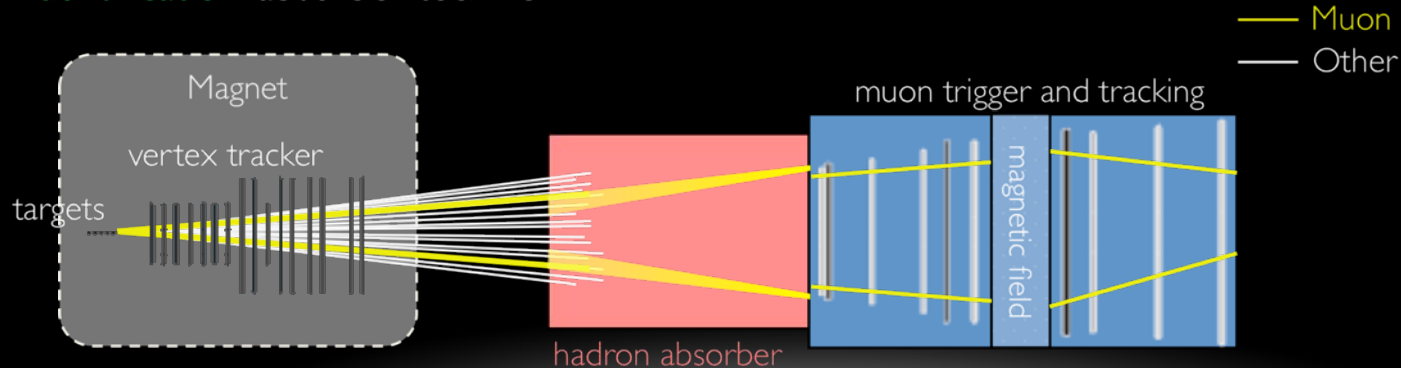
Velocity in ToF



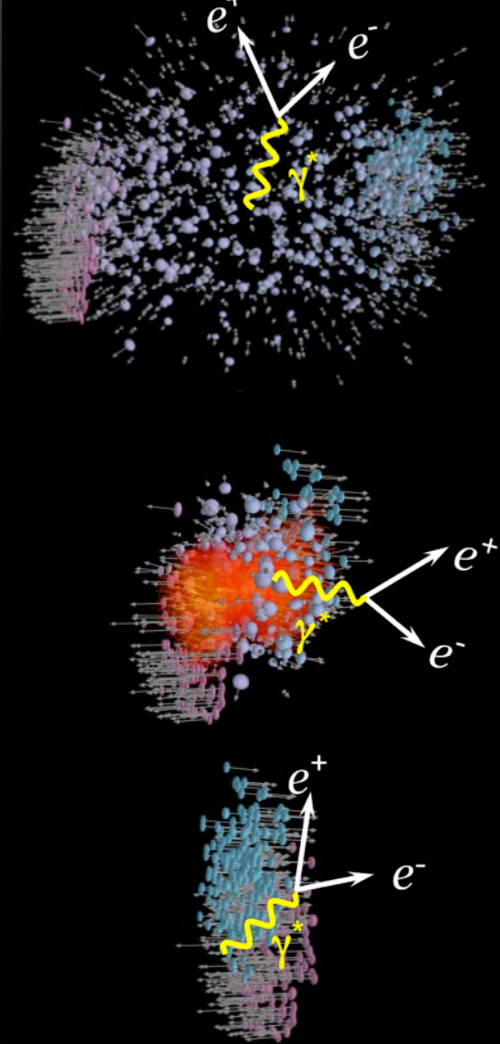
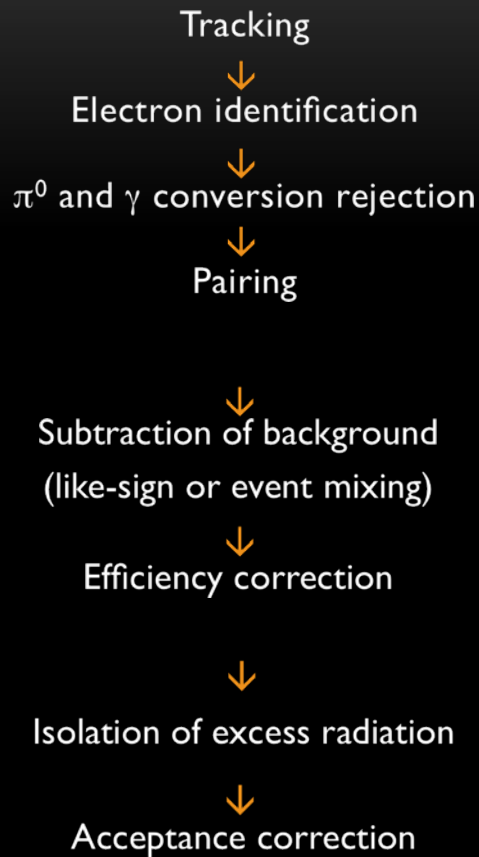
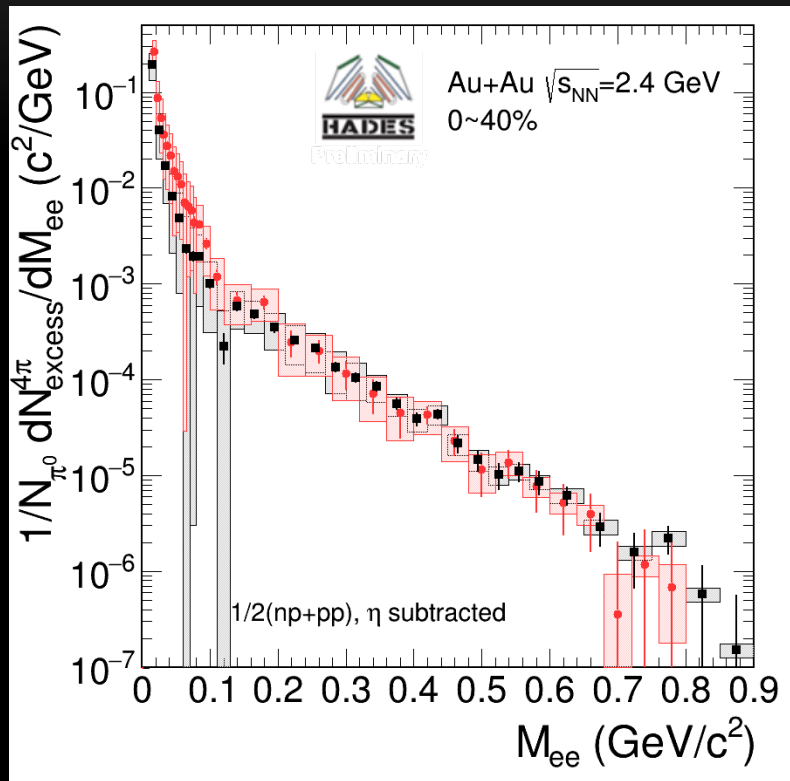
Rings in RICH



Muon identification: absorber technic



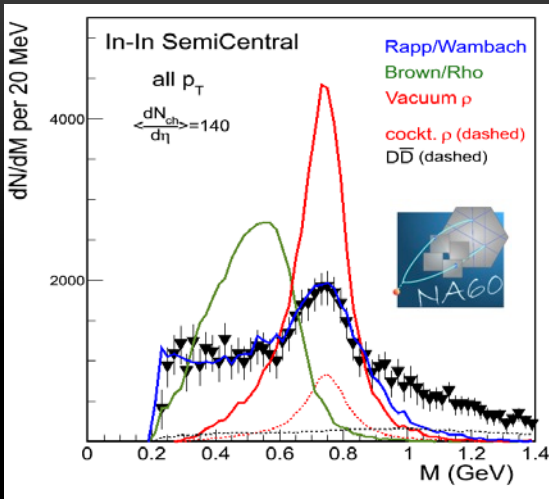
Just few steps ;)



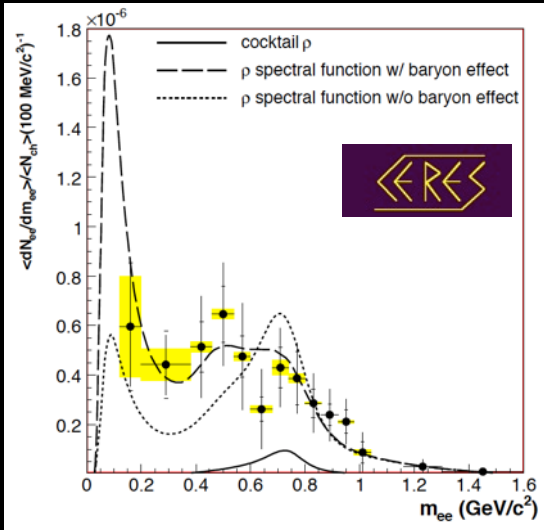
Dileptons as spectrometer



Are narrow in-medium vector meson states with substantially shifted pole mass observed?



- ❑ Disfavors “dropping mass” scenario: $m_{had} \sim \langle \bar{q}q \rangle$
- ❑ Strongly supports in-medium broadening



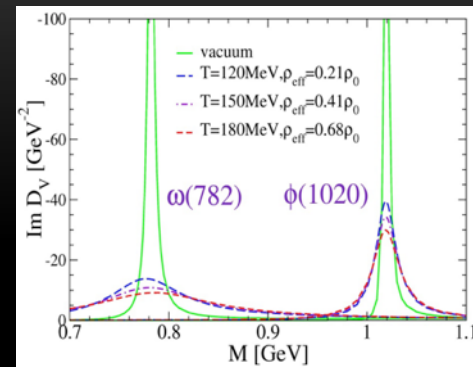
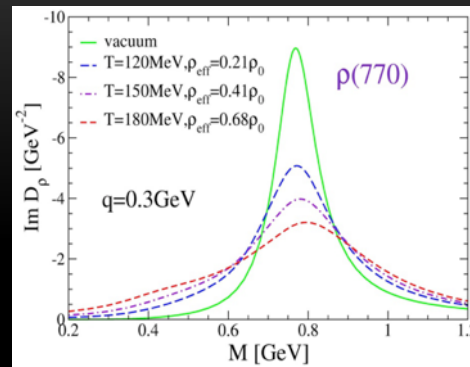
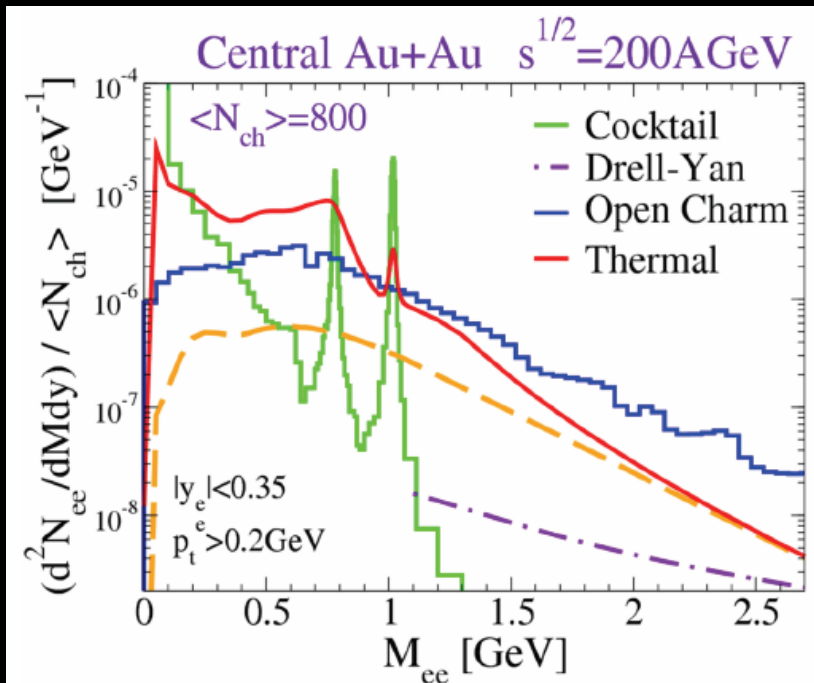
$$\Sigma_{\rho B, M} = \rho \text{ (loop) } h = N, \pi, K$$

$R = \Delta, N(1520), a_1, K_1, \dots$



Data NA60: Phys. Rev. Lett. 96 (2006) 162302
 Data CERES: Phys.Lett. B666 (2008)
 Calculations: R.Rapp and H. van Hees, 2008

Predictions for RHIC

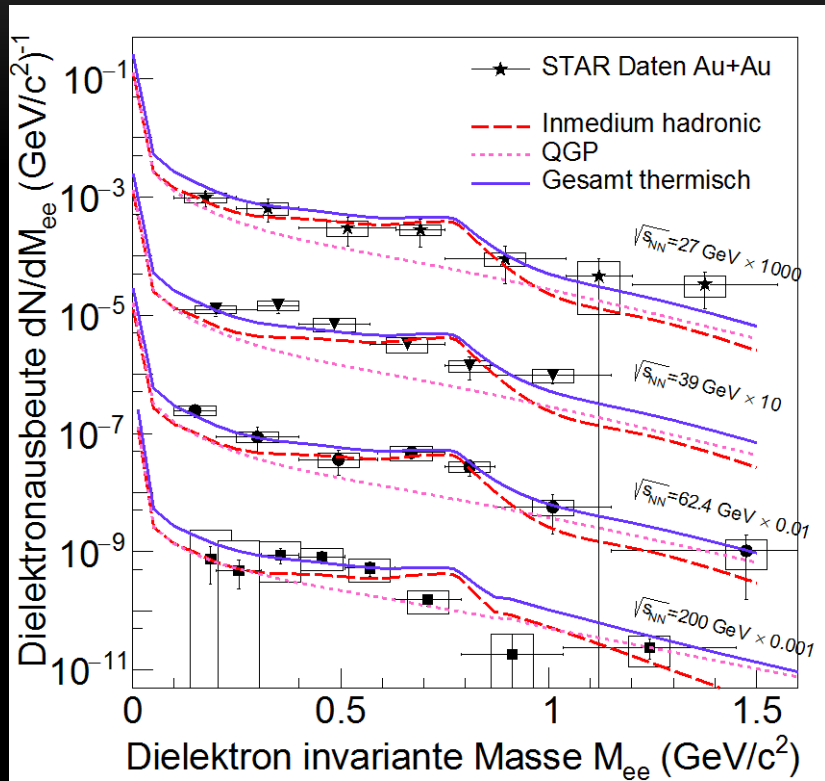


In-medium modifications of vector mesons persists (R. Rapp)

Open charm contribution becomes significant

Dielectron mass spectra from STAR BES I

Au+Au $\sqrt{s_{NN}} = 19.6 - 200 \text{ GeV}$



Phys.Lett. B750 (2015) 64-71

J. Butterworth et al., arXiv:1612.05484 [nucl-ex]

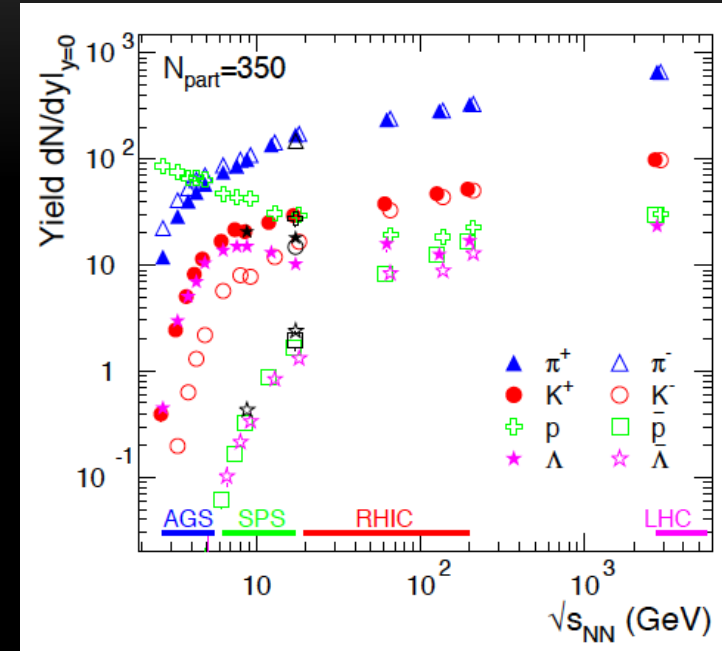
Model: Rapp/Wambach/Hees

- Isolation of the excess by subtracting the measured decay cocktail
- Acceptance corrected spectra
- In-medium **broadened ρ** spectral function consistently describes the low-mass electron-positron excess **for all the energies 19.6-200 GeV**

From SPS to RHIC

	SPS (Pb+Pb)	RHIC (Au+Au)
$dN(\bar{p})/dy$	6.2	20.1
produced baryons (p, \bar{p} , n, \bar{n})	24.8	80.4
$p - \bar{p}$	33.5	8.6
participating nucleons $(p - \bar{p})A/Z$	85	21.4
total baryon number	110	102

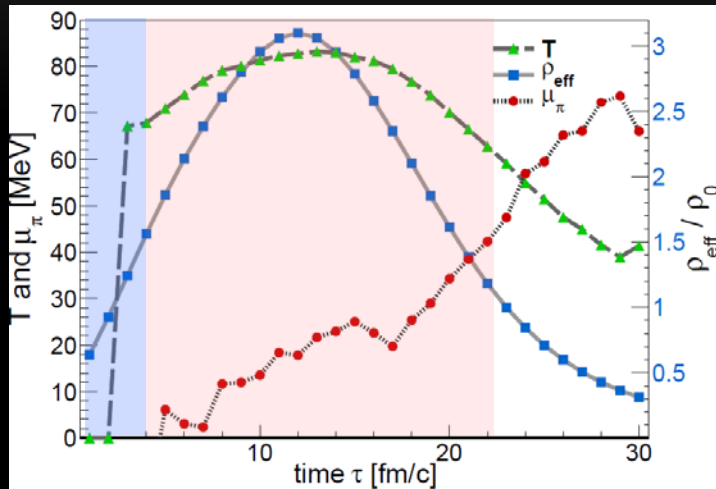
- Although the NET baryon density is different at SPS and RHIC, baryon density is practically the same!
- Baryon effects important even at $\rho_{B,tot} = 0$: sensitive to $\rho_{B,tot} = \rho_B + \rho_{\bar{B}}$ (ρ -N and ρ - \bar{N} interactions identical)
- Higher initial temperature at RHIC




A. Andronic, arXiv:1407.5003

Baryonic matter at few GeV beam energy $Au+Au \sqrt{s_{NN}} = 2.42 \text{ GeV}$

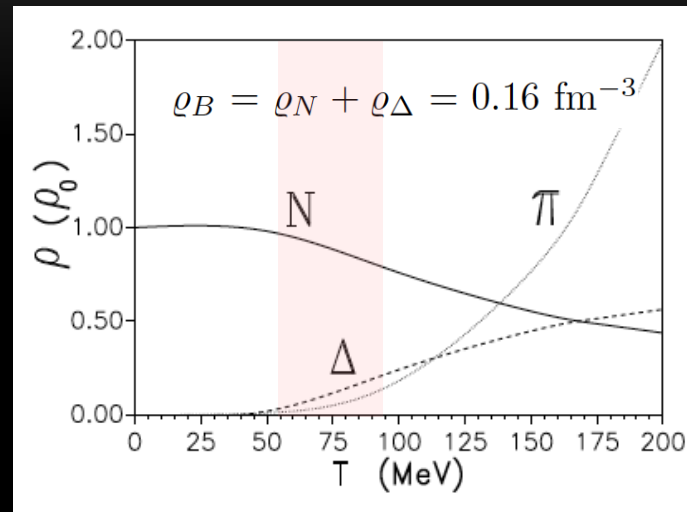
Central cell ($3 \times 3 \times 3 \text{ fm}^3$) thermodynamic properties from coarse graining UrQMD




 TG, F. Seck et al., *Eur. Phys. J. A* 52 (2016) 131

- ❑ Long interpenetration times
- ❑ Comparatively long lifetime of the dense "fireball" ($\rho_{\text{max}} \approx 3 \rho_0$)

Composition of a hot $\pi\Delta N$ gas

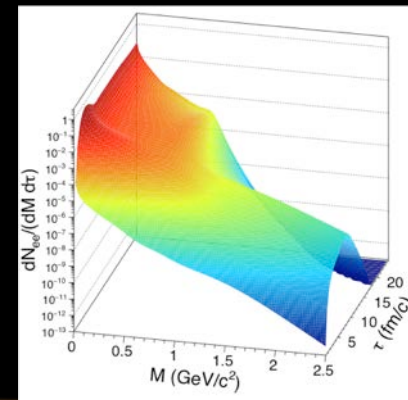
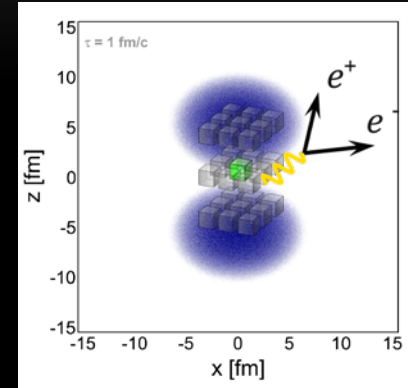
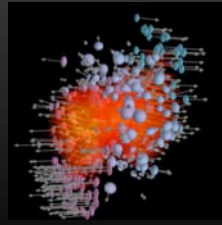
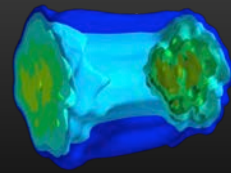


 Rapp, Wambach, *Adv. Nucl. Phys.* 25 (2000)

- ❑ Moderate temperatures: $T < 90 \text{ MeV}$
- ❑ Baryon-dominated system throughout the evolution ($N_{\pi}/A_{\text{part}} \approx 10\%$)

Coarse-grained transport approach

- “Combine” the advantages of two descriptions: hydrodynamics & transport
- Simulate events with a transport model
 - ensemble average to obtain smooth space-time distributions
- Divide space-time evolution into 4-dimensional cells
 - $21 \times 21 \times 21$ space cells (1 fm^3), 30 time steps → ~ 280 k cells
- Determine for each cell the bulk properties like T , μ_B , μ_π , collective velocity
 - parameterization of RW in-medium spectral function
- Sum up contributions of all cells

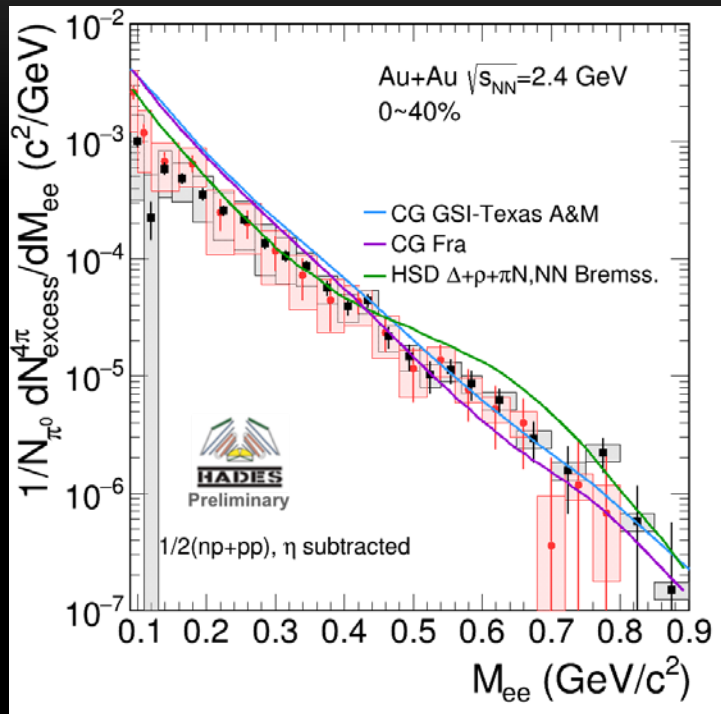


Huovinen et al., PRC 66 (2002) 014903

CG FRA Endres et al.: PRC 92 (2015) 014911

CG GSI-Texas A&M TG et al.: Eur.Phys.J. A52 (2016) no.5, 131

Dielectron mass spectra at $\sqrt{s_{NN}} = 2.42$ GeV



HADES., collaboration review

- Strong broadening of the in-medium ρ due to direct ρ -hadron scattering

$$R = \Delta, N(1520), a_1, K_{1,1}, \dots$$

$$\Sigma_{\rho B, M} = \text{wavy line } \rho \text{ loop } \text{wavy line } h = N, \pi, K$$

- Thermal rates folded over coarse-grained UrQMD medium evolution works at low energies
- Supports baryon-driven medium effects at UrHIC (SPS and RHIC)!

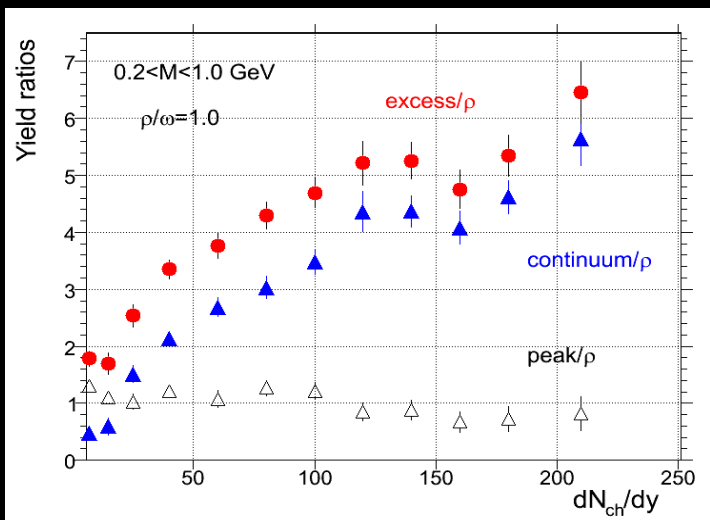
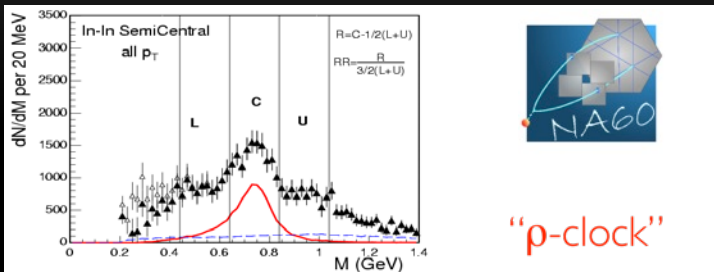
Robust understanding across QCD phase diagram

Dileptons as chronometer

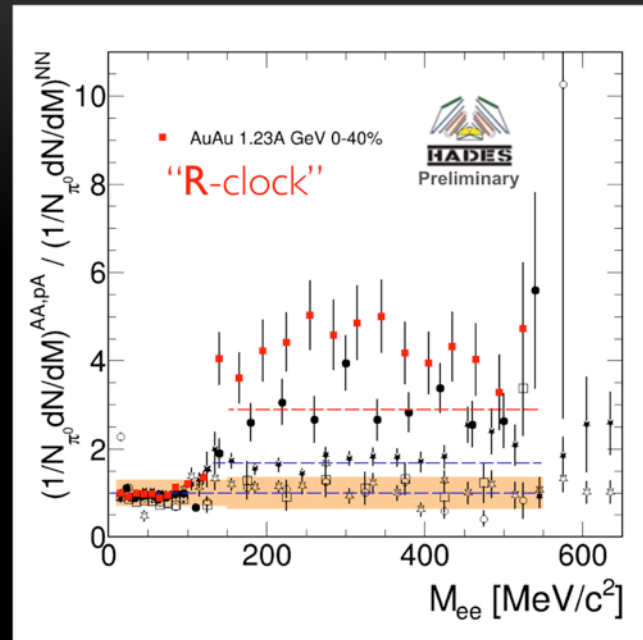


The dilepton clock

Centrality dependence of spectral shape



System size dependence of excess

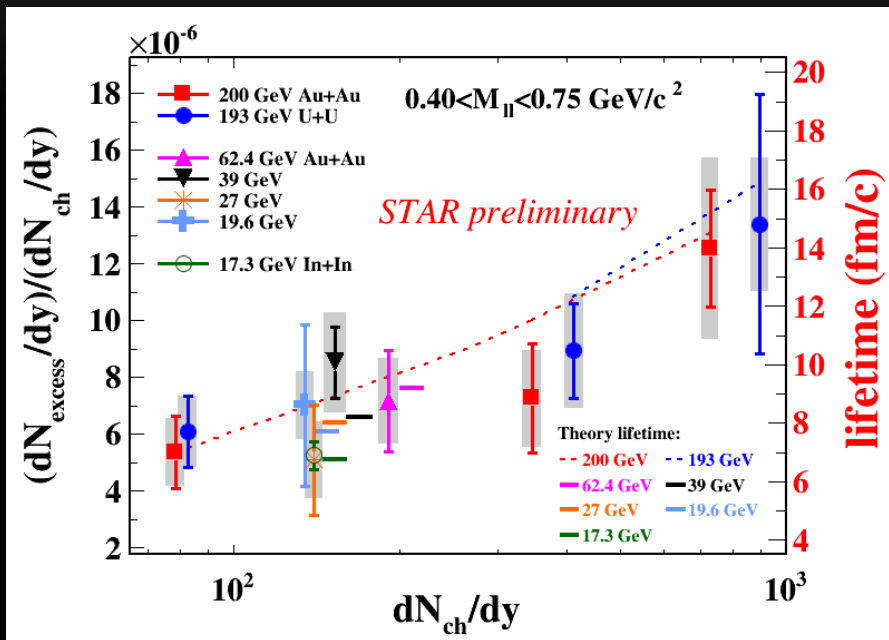


Rapid increase of relative yield reflect the number of ρ 's / R's regenerated in fireball



U.W. Heinz and K. S. Lee, PLB 259, 162 (1991)
H.W. Barz, B. L. Friman, J. Knoll and H. Schulz, PLB 254, 315 (1991)

The lifetime of the interacting fireball



- Normalized **excess yields** for the mass region $0.3 < M < 0.7 \text{ GeV}/c^2$ is proportional to the **lifetime** of the interacting fireball
- Note: normalization to the number of pions should be done when going down with energy!



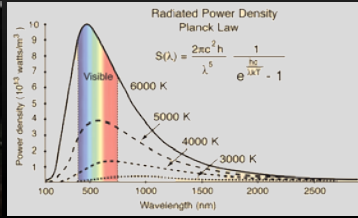
STAR data: arXiv:1612.05484 [nucl-ex]

Model: R. Rapp, H. van Hees, PLB 753 (2016) 586



Spectral shape of the light \rightarrow
temperature of the emitting object

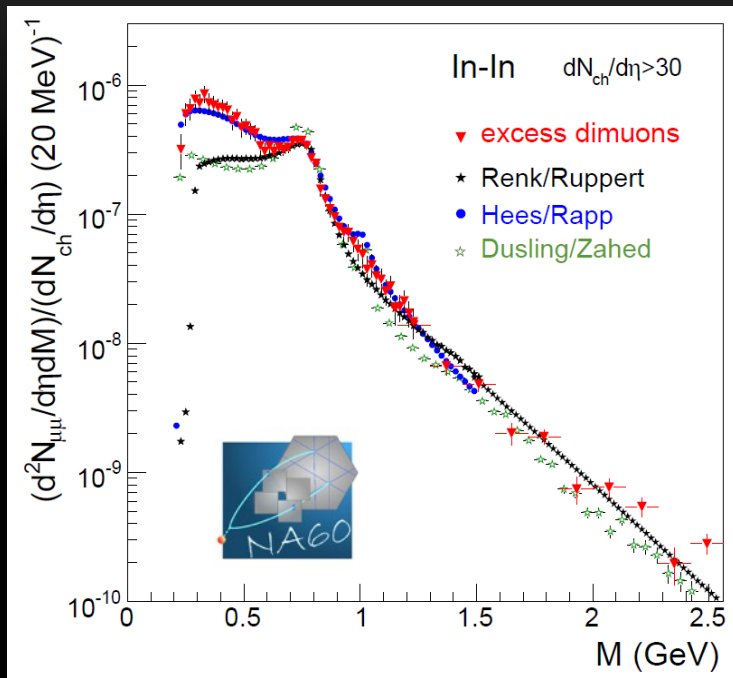
$$E = \hbar\nu = \frac{\hbar c}{\lambda}$$



Dileptons as thermometer



Measurement of radiating source temperature at SPS



- Acceptance corrected excess yield
- $M_{\mu\mu} > 1 \text{ GeV}c^2 \sim$ exponential fall-off - 'Planck-like'
- measurement of radiating source temperature

→ fit $\frac{dN}{dM} \sim M^3 \times \exp\left(-\frac{M}{T}\right)$ to range

→ $M=1.1-2.0 \text{ GeV}c^2$: $\langle T \rangle_{\text{emitting source}} = 205 \pm 12 \text{ MeV}k_B$

→ $M=1.1-2.4 \text{ GeV}c^2$: $\langle T \rangle_{\text{emitting source}} = 230 \pm 10 \text{ MeV}k_B$



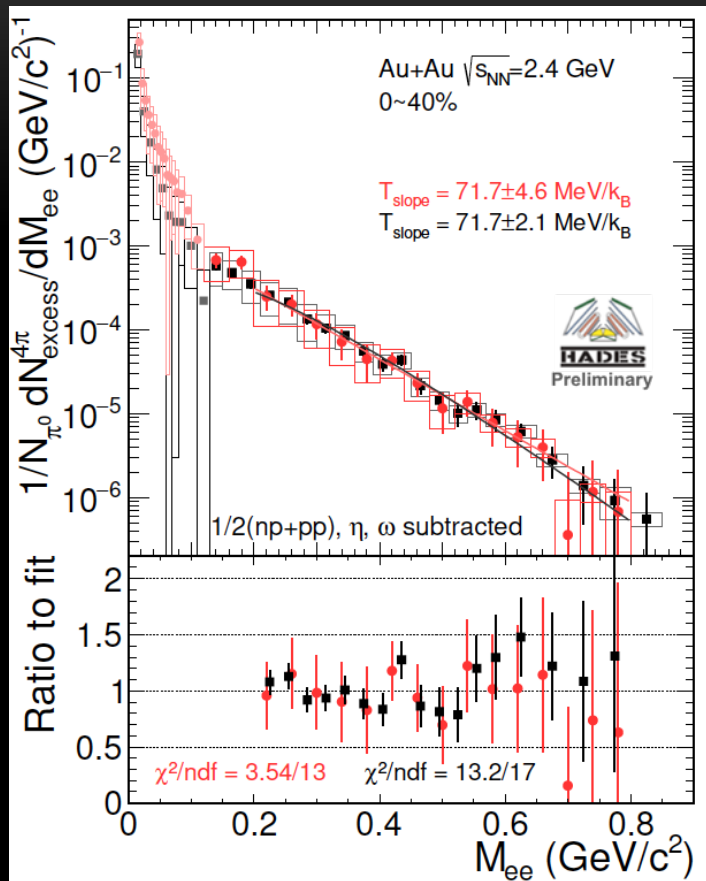
Eur. Phys. J. C 59 (2009) 607-623

CERN Courier 111 2009, 31-35

Chiral 2010, AIP Conf. Proc. (2010) 1322

M is the only Lorentz-invariant thermometer of the field

Virtual photon emission – fireball thermometer at SIS18



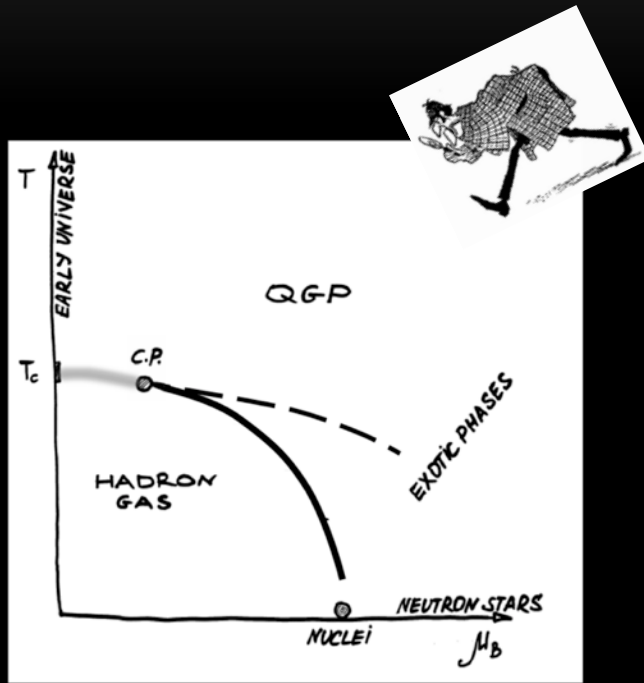
□ $T_{\text{initial}} \sim 3$ times lower at SIS18 compared to SPS

□ Shifts predominance of early emission from $M_{ee} > 1.2 \text{ GeV}/c^2$ at SPS to $M_{ee} > 0.4 \text{ GeV}/c^2$ at SIS

→ fit $\frac{dN}{dM} \sim M^3 \times \exp\left(-\frac{M}{T}\right)$ to range $M=0.2-0.8 \text{ GeV}/c^2$

□ $\langle T \rangle_{\text{emitting source}} = 72 \pm 2 \text{ MeV}/k_B$

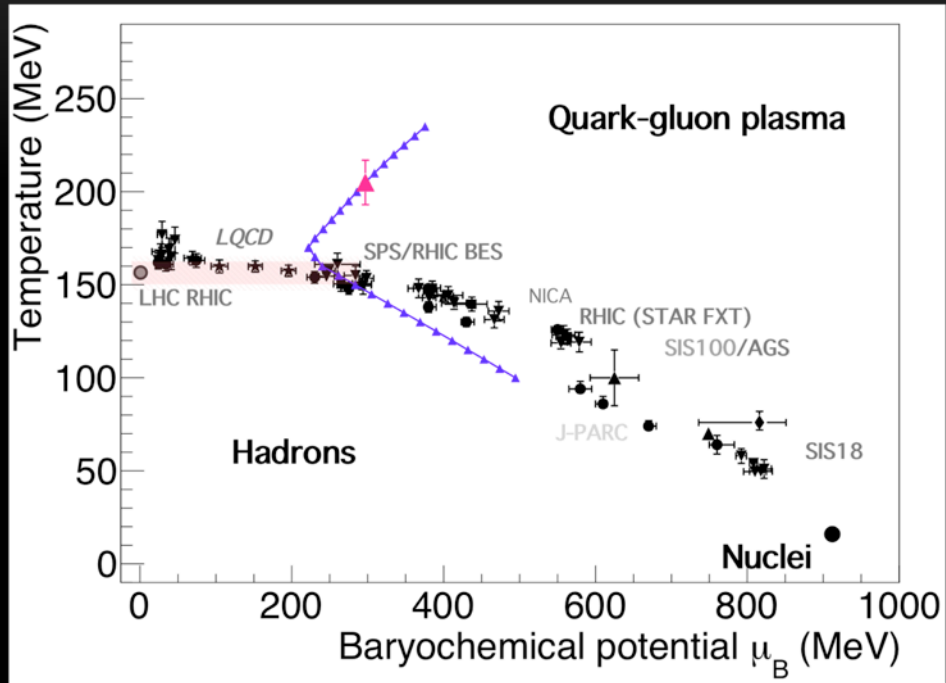
Dileptons and QCD phase diagram of matter



Excitation functions

- Fireball lifetime
- Emitting source temperature

Trajectories in the phase diagram



- Hadron production: well described by hadron gas in chemical and thermal equilibrium

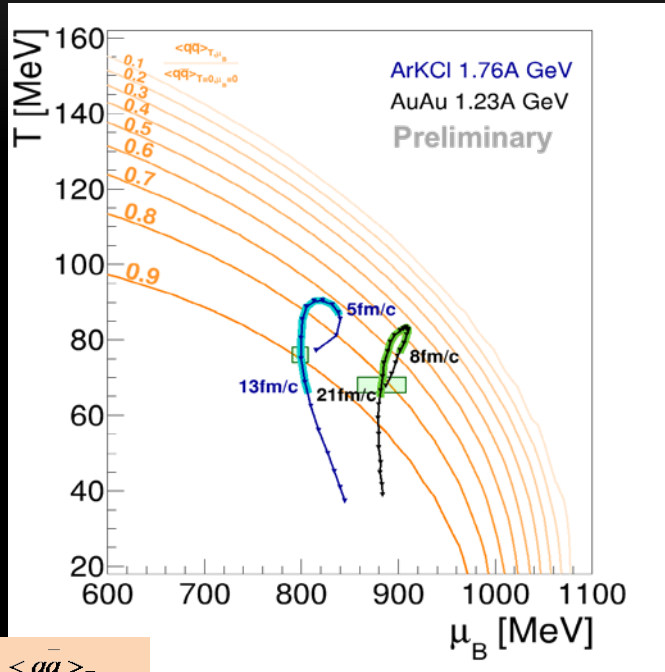
$$n_i(\mu_B, T) = d_i \int \frac{d^3k}{(2\pi)^3} f(E_k; \mu_B, T)$$

- Dileptons - need to construct evolution:
 - Before up to earliest “formation” time $\tau_0 \leftrightarrow T_0 > T_{\text{chem}}$
 - After: down to thermal freezeout $\tau_f \leftrightarrow T_{f_0} < T_{\text{chem}}$

- Basic assumption: entropy (+baryon-number) conservation \rightarrow fixes $T(\mu_B)$ in the phase diagram



HADES and QCD phase diagram of matter




- Chemical freeze-out from measured particle yields analyzed with SHM
- Trajectories extracted from inner cube of cells with coarse-grained UrQMD
- Time-window of dilepton emission \rightarrow access to hot and dense stage of the heavy-ion collision

$T_{\max} = 85 \text{ MeV}, \rho_{\max} = 3 \rho_0$
 \rightarrow Excitation of the vacuum
 (melting of condensate)
 matches spectral medium effects!

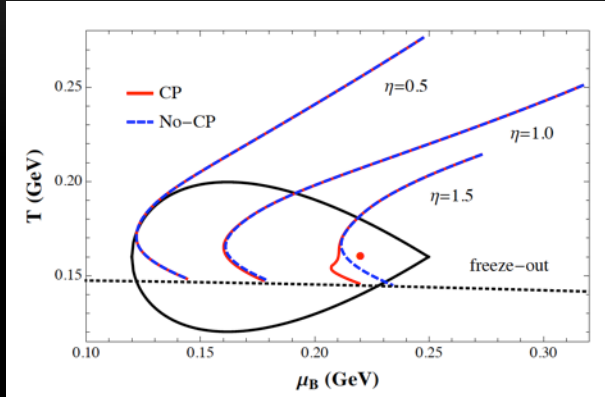
$$\frac{\langle \langle \bar{q}q \rangle \rangle (T, \mu_B)}{\langle \bar{q}q \rangle} = 1 - \sum_h \frac{\rho_h^s \Sigma_h}{m_\pi^2 f_\pi^2}$$

$\langle \bar{q}q \rangle_{T, \mu_B}$
 $\langle \bar{q}q \rangle_{T=0, \mu_B=0}$: B.J. Schaefer and J. Wambach

 Ar+KCl data: Nucl. Phys. A931 (2014) c785

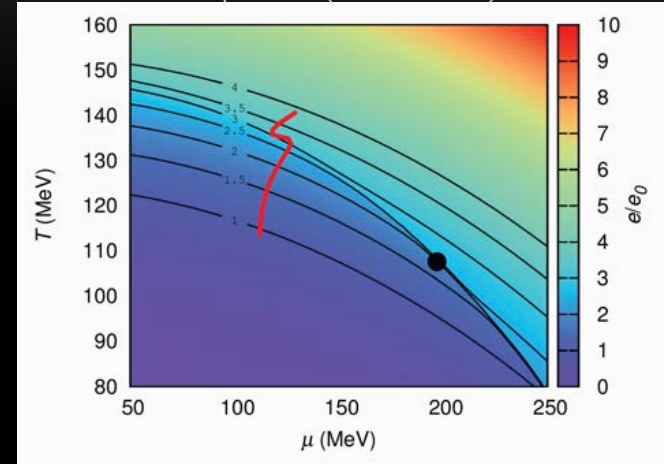
Mapping QCD phase diagram with dileptons

Hydrodynamic evolution trajectories



A. Monnai, S. Mukherjee, Y. Yin: arXiv: 1606.00771

Event-averaged trajectory near the critical point (black dot)



C. Herold, M. Nahrgang, Y. Yan and C. Kobdaj, PRC93 (2016)

Diverging bulk viscosity at QCD critical point

$$\zeta \sim \tau_{\Pi} \sim \tau_{\sigma} \sim \xi^3$$

Bulk viscosity

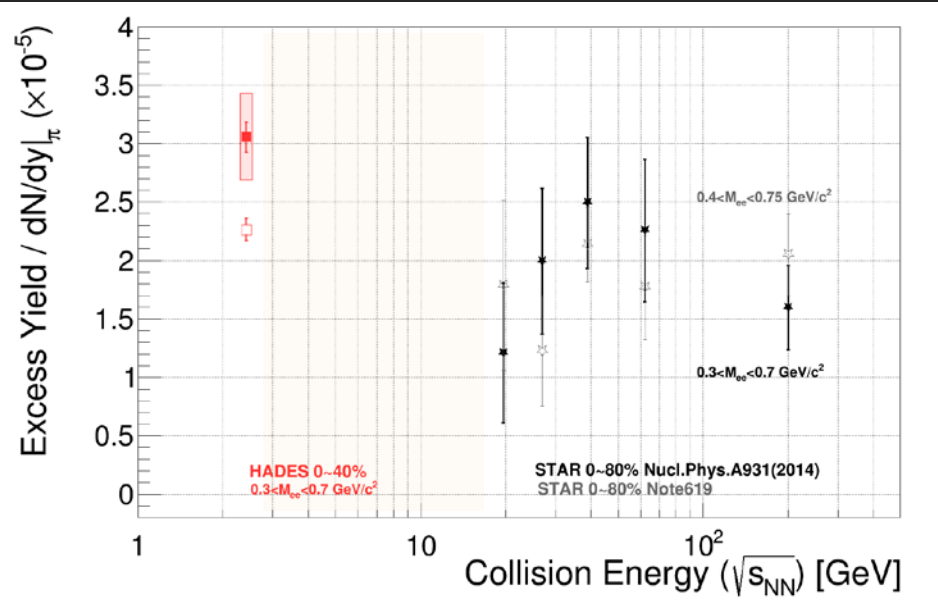
Bulk relaxation time

Relaxation time for the critical mode

Correlation length³

What are the possible signatures in dilepton radiation?

Energy dependence of low-mass excess



□ Yield in low-mass window tracks fireball lifetime

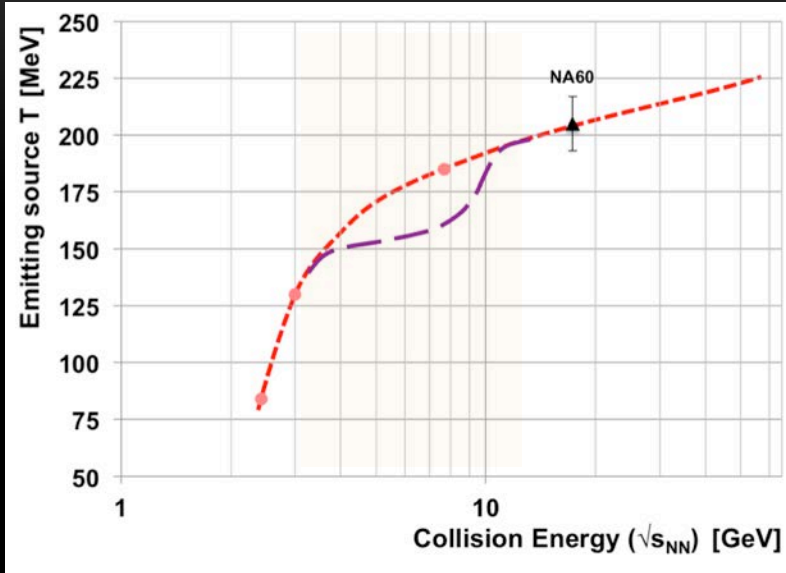
→ Search for **anomalous fireball lifetime** around phase transition & critical point

- 2019 - STAR at RHIC BES II
- 2024 - CBM at FAIR

- Quite moderate energy dependence
- Dilepton yield determined by interplay between temperature and $V \otimes \tau_{coll}$



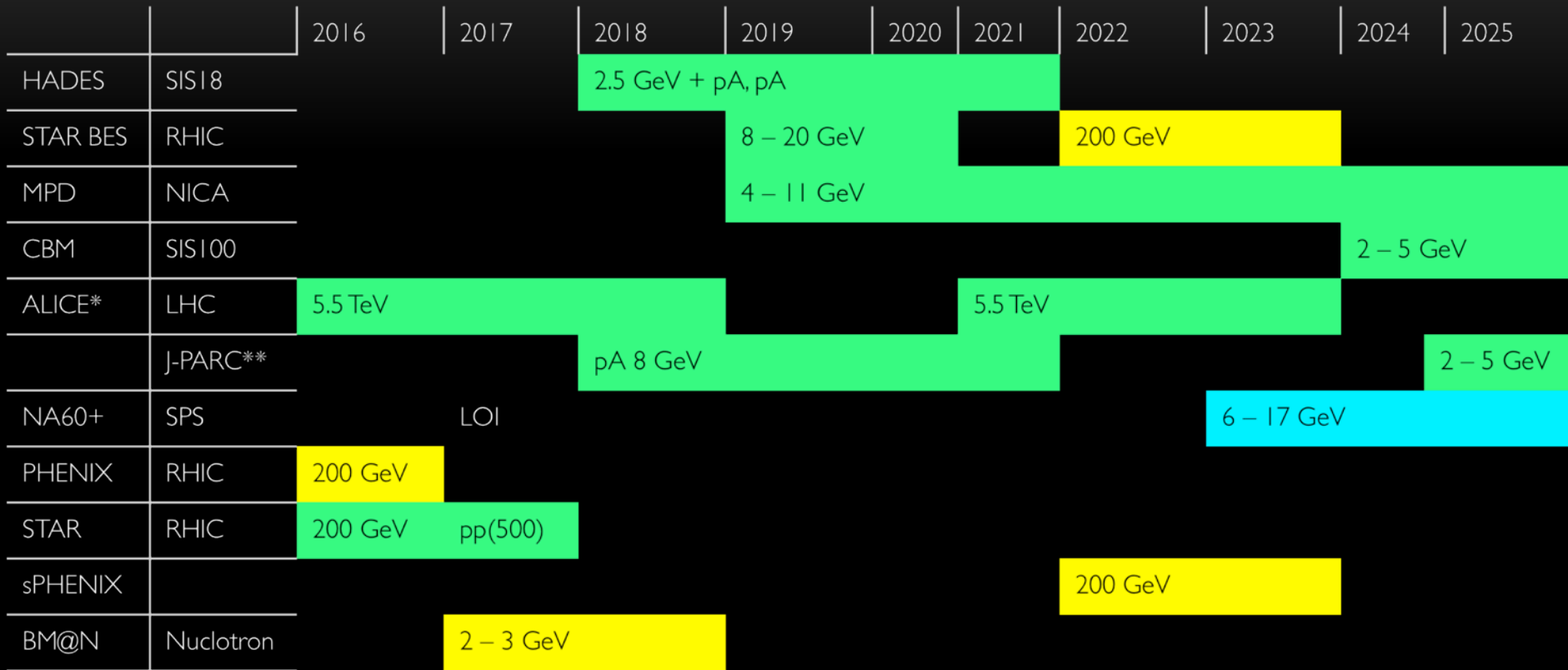
Energy dependence of intermediate mass slope



Dashed violet curve corresponds to a speculated shape with phase transition

- Measures the emitting source temperature (true, no blue shift)
- Measure T_s (note, $T_s < T_{\text{initial}}$) "caloric curve"
- Plateau around onset of deconfinement?
[see e.g. M. D'Agostino et al. NPA 749 (2005) 5533]
- Precision measurements are the key
- 2024 - CBM at FAIR

“You may say I’m a dreamer... but I’m not the only one”



* - ITS, 50kHz, lower field

** - Proposal to J-PARC → If approved, construction of HI injector and detectors in 10 years?

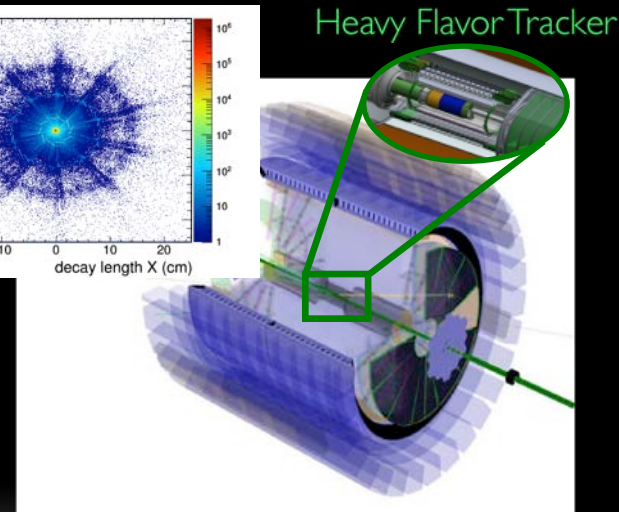
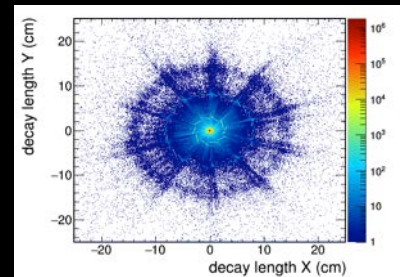
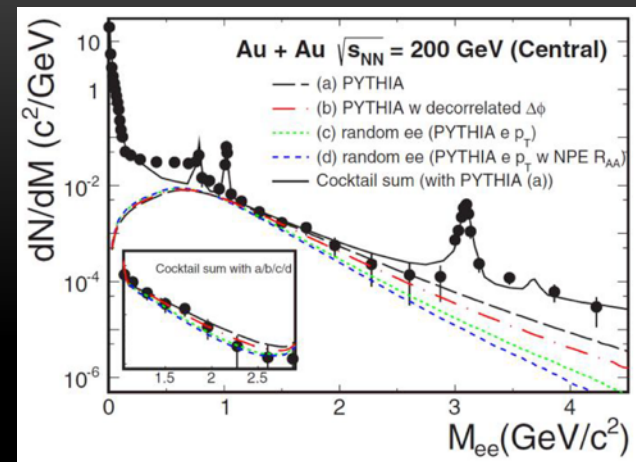
STAR dielectrons ongoing analysis

Au+Au $\sqrt{s_{NN}} = 200$ GeV

- $\mu_B \ll T$, i.e. vanishing net-baryon density
- Lattice QCD computations are most powerful
- Measure ρ spectral function and “calibrate” EM rates
- Extract fireball temperature (IMR)

Dielectron reconstruction with HFT

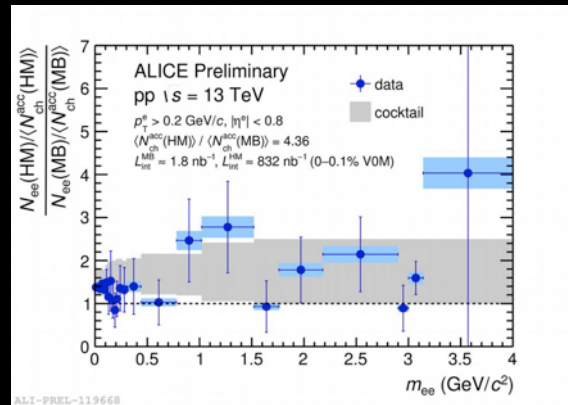
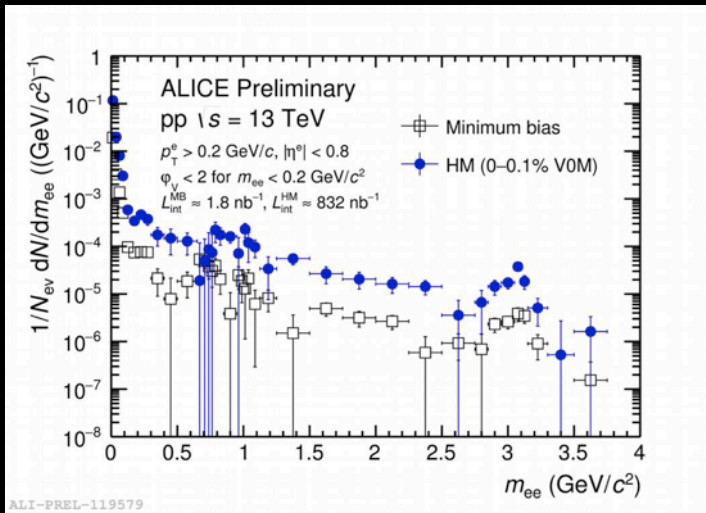
- Reduce uncertainties on charm contribution
- **Challenge:** γ conversion in HFT detector material
→ use excellent vertexing to reject it
(looks promising, statistics Runs 10, 11, 14, 16 $\sim 4 \times 10^9$)



pp at $\sqrt{s} = 13$ TeV: High-Multiplicity studies

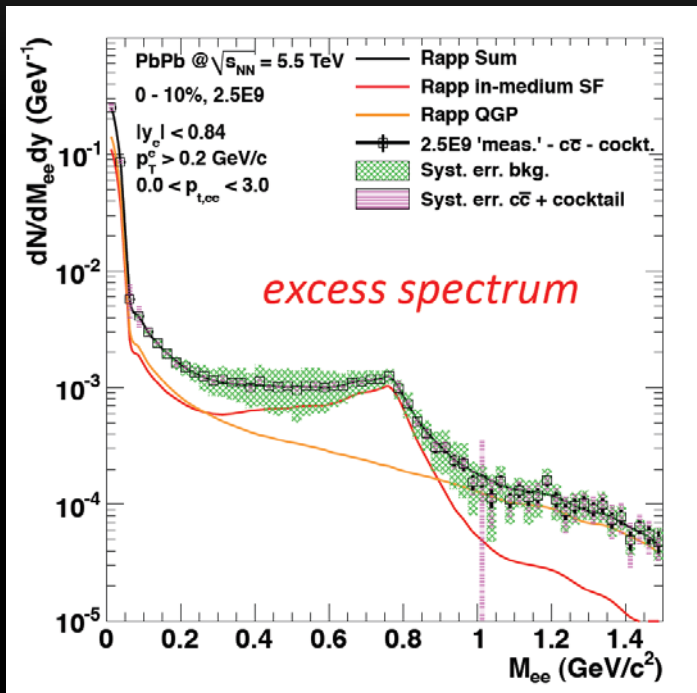
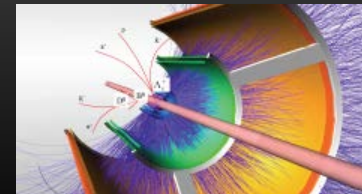
- New (or heavy-ion like) phenomena in High-Multiplicity (HM) pp events:
 - Production/destruction of ρ mesons
 - Multiplicity scaling of light hadrons, open heavy flavours and direct photons
- Thermal radiation? in HM events

→ So far: compare dielectron yield (uncorrected) in MB and HM pp collisions



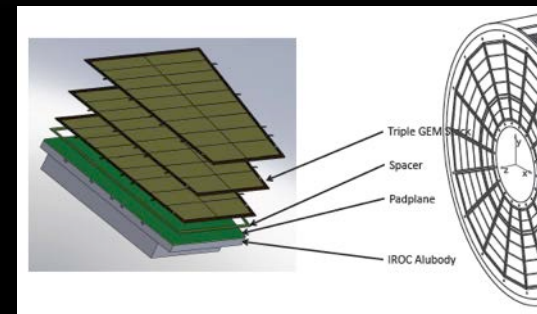
x 5 more pp data recorded in 2016

Future: Pb+Pb from ALICE

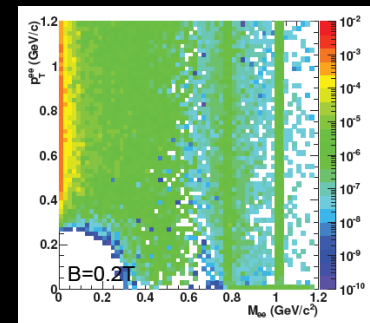


New ITS
 → charm

New TPC readout chambers with GEMs
 → 50kHz, statistics

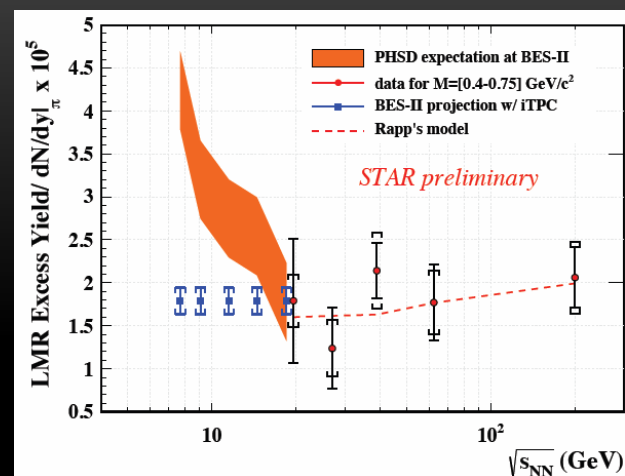


Low magnetic field run
 → acceptance



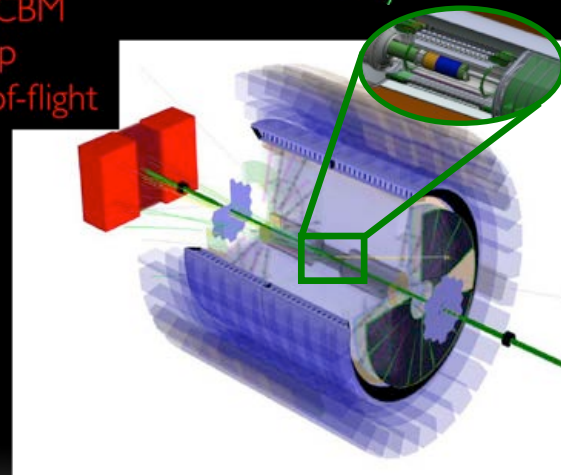
STAR dielectrons Beam Energy Scan II 2019 - 2020

- ❑ Improve statistics for existing low energy samples
- ❑ Quantify lifetime and baryon density dependence of the ρ spectral function
- ❑ Disentangle various model calculations
- ❑ Inner TPC upgrade
- ❑ Installation of eToF (CBM Phase-0)
- ❑ Collision energies 7.7, 9.1, 11.5, 14.5, 19.6 GeV

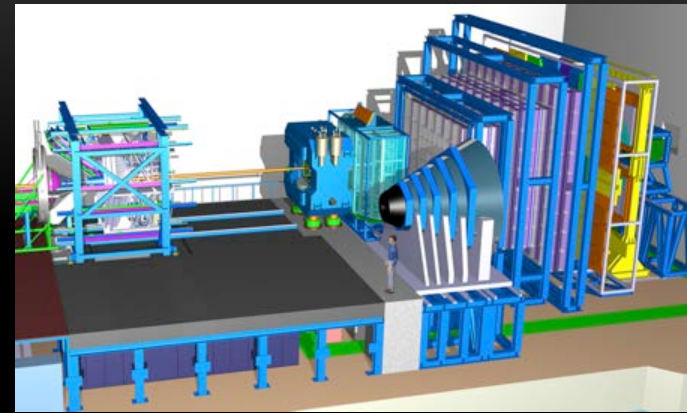
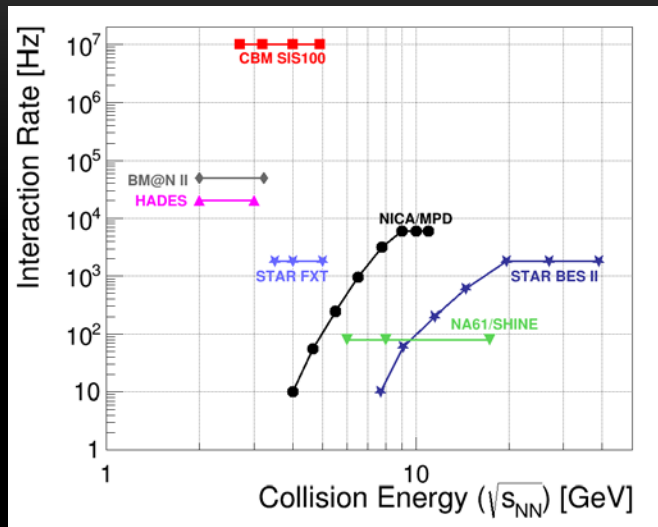


STAR/CBM
end-cap
Time-of-flight

Heavy Flavor Tracker

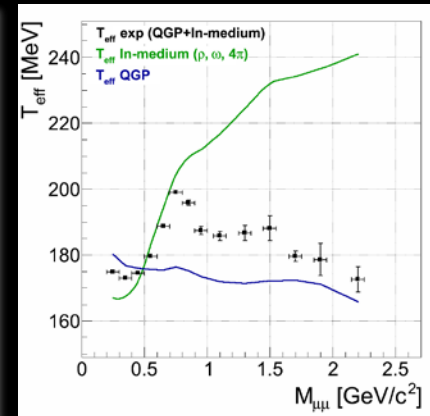
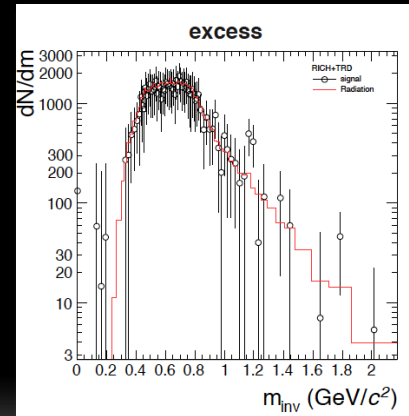


Compressed Baryonic Matter experiment at FAIR



CBM Collab., EPJA, 53 3 (2017) 60

- CBM will play a **unique role** in the exploration of the QCD phase diagram in the region of high net-baryon densities
- 4π processes: $\pi a_1 \rightarrow \gamma^* \rightarrow l^+l^-$ (chiral mixing) is a dominant hadronic source in IMR: correlated charm contribution, Drell-Yan, QGP are suppressed

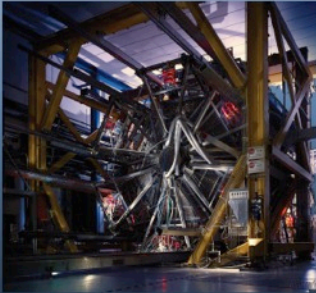


”If you are out to describe the truth,
leave elegance to the tailor” (A. Einstein)



Proposal for experiments at
SIS18 during FAIR Phase-0

The HADES Collaboration



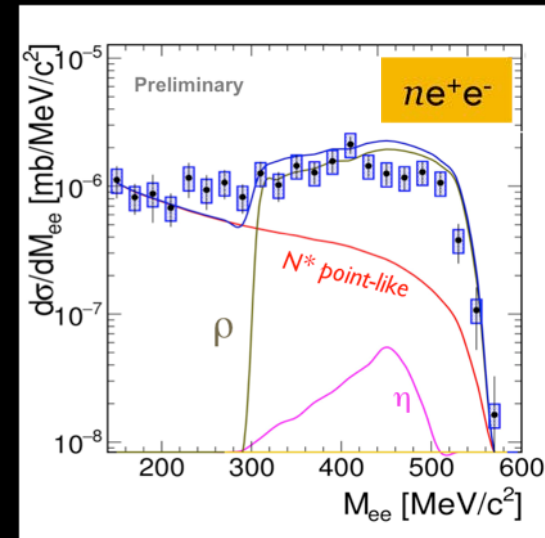
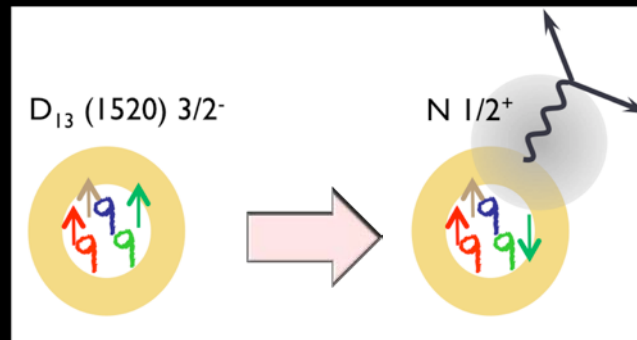
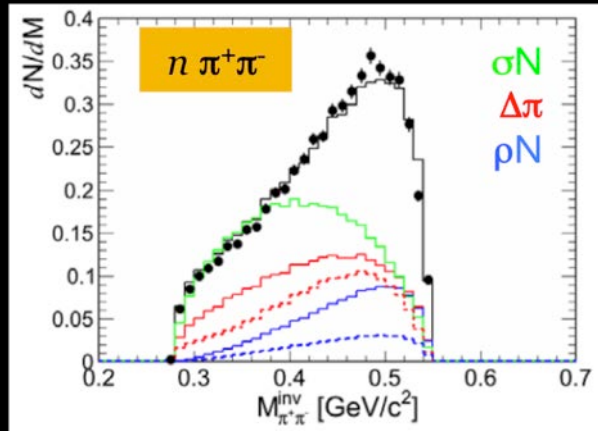
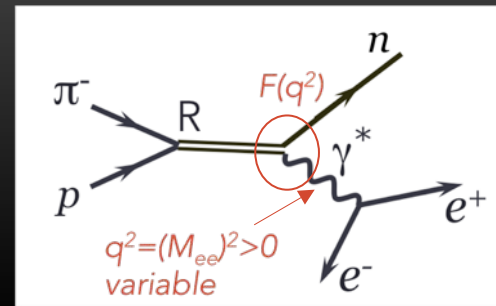
Properties of hadron resonances
and baryon rich matter

- HADES at SIS18 and SIS100
 - High statistics $\pi+p$, $\pi+A$, $p+A$, $A+A$
- Results in elementary collisions provide an **important baseline** for current and future explorations in HIC

EM baryon-resonance decays

π^- (683,5 MeV/c) $p \rightarrow N^*$

- Worldwide unique combination of π beams (at low energies) with dilepton spectrometer
- Access to time-like etFF of baryons



$\rightarrow D_{13}$ (1520) is dominant in ρ production!

$$\frac{d\sigma}{dM_{ee}} = \frac{d\sigma}{dM_{\pi\pi}} c_p \left(\frac{m_p}{m_{ee}} \right)^3$$

\rightarrow First measurement demonstrating the dominant role of intermediate ρ propagation in $N^* \rightarrow N e^+ e^-$ transition

Résumé and prospects

- ❑ **Unique** possibility of characterizing properties of hot and dense QCD matter with dileptons
- ❑ **Robust** understanding of low-mass dilepton excess radiation by ρ -baryon coupling (at top RHIC, RHIC BES, SPS and SIS18 energies)
- ❑ Enable unique measurements
 - ❑ **Degrees of freedom** of the medium
 - ❑ Restoration of **chiral symmetry**
 - ❑ **Fireball lifetime**
 - ❑ Emitting source **temperature**
- ❑ Future experiments allows for overlap and independent confirmation of results



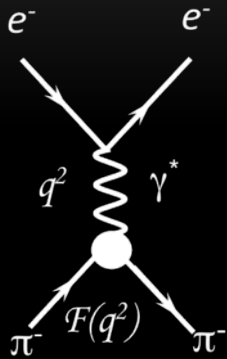
There is no mission impossible

**Thank you
for your attention!**

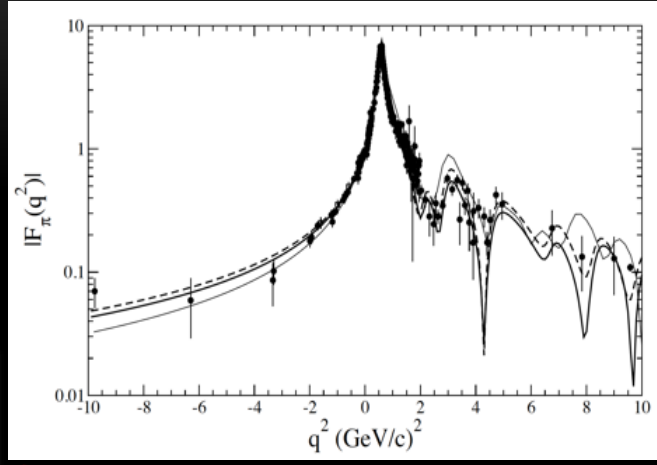
Bonus slides

How do photons couple to hadrons?

space-like photons

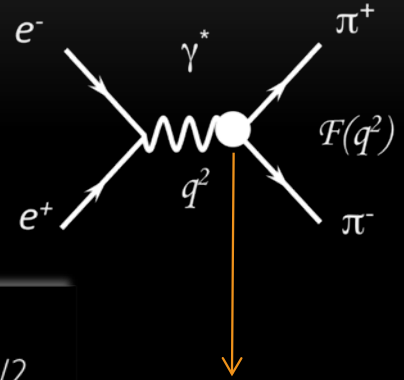


$q^2 < 0$
 $\Delta p \Delta x \geq \hbar/2$



De Melo et al., Phys. Rev. D73 (2006) 070413

time-like photons



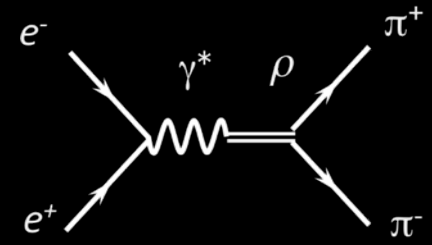
$q^2 > 0$
 $\Delta E \Delta t \geq \hbar/2$

$$F(q^2) = \frac{d\sigma/dq^2}{\left(d\sigma/dq^2\right)_{point\ like}}$$

$q^2 = (\Delta E)^2 - (\Delta p)^2$

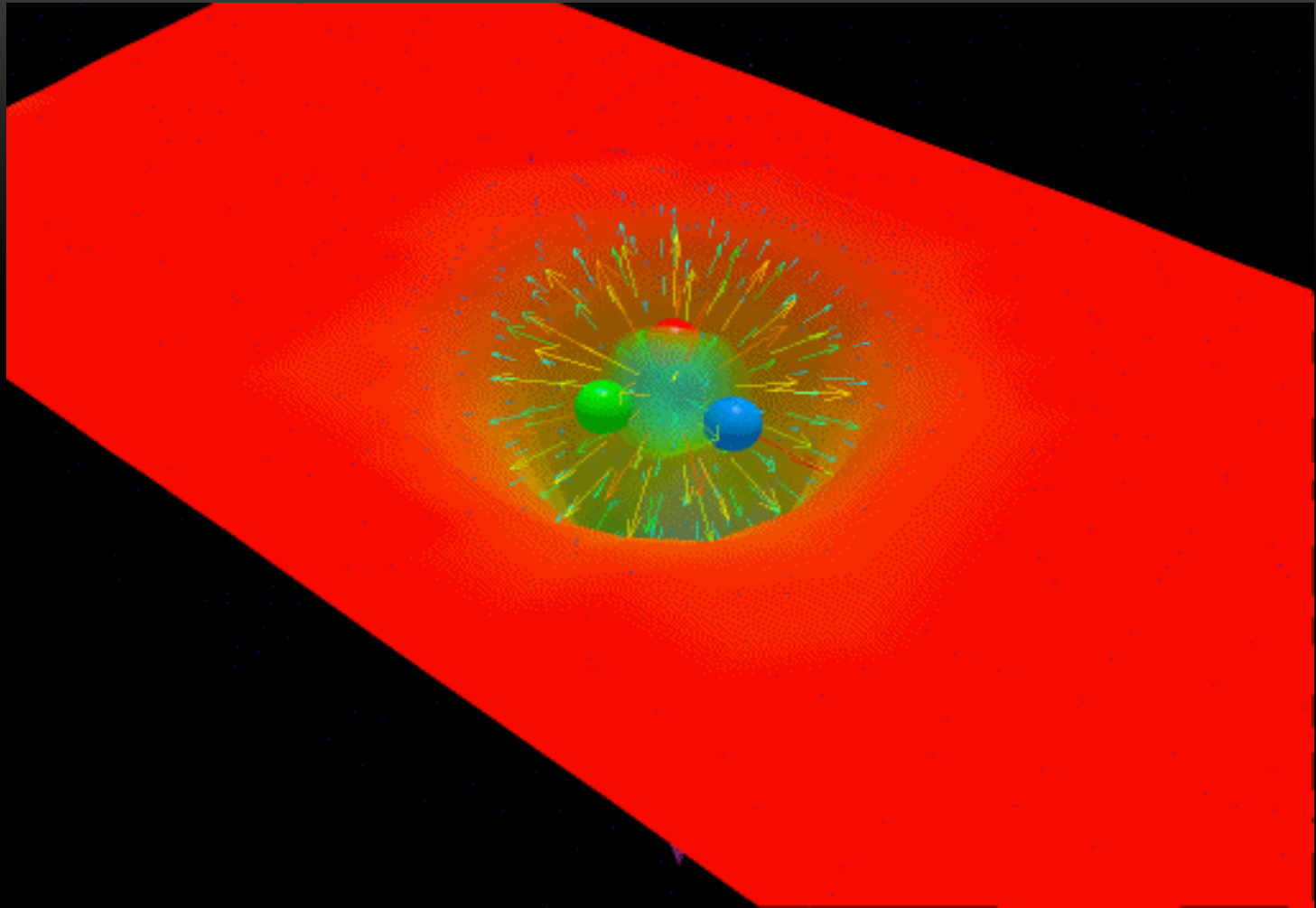
→ Form factor
 → Squared 4-momentum

$J^P = 1^-$ for both γ^* and Vector Meson



- Strong coupling of γ^* to Vector Meson
 → Vector Meson Dominance model
- Observable: vector mesons (ρ , ω , ϕ)

Nucleon
on Lattice QCD



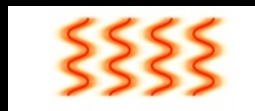
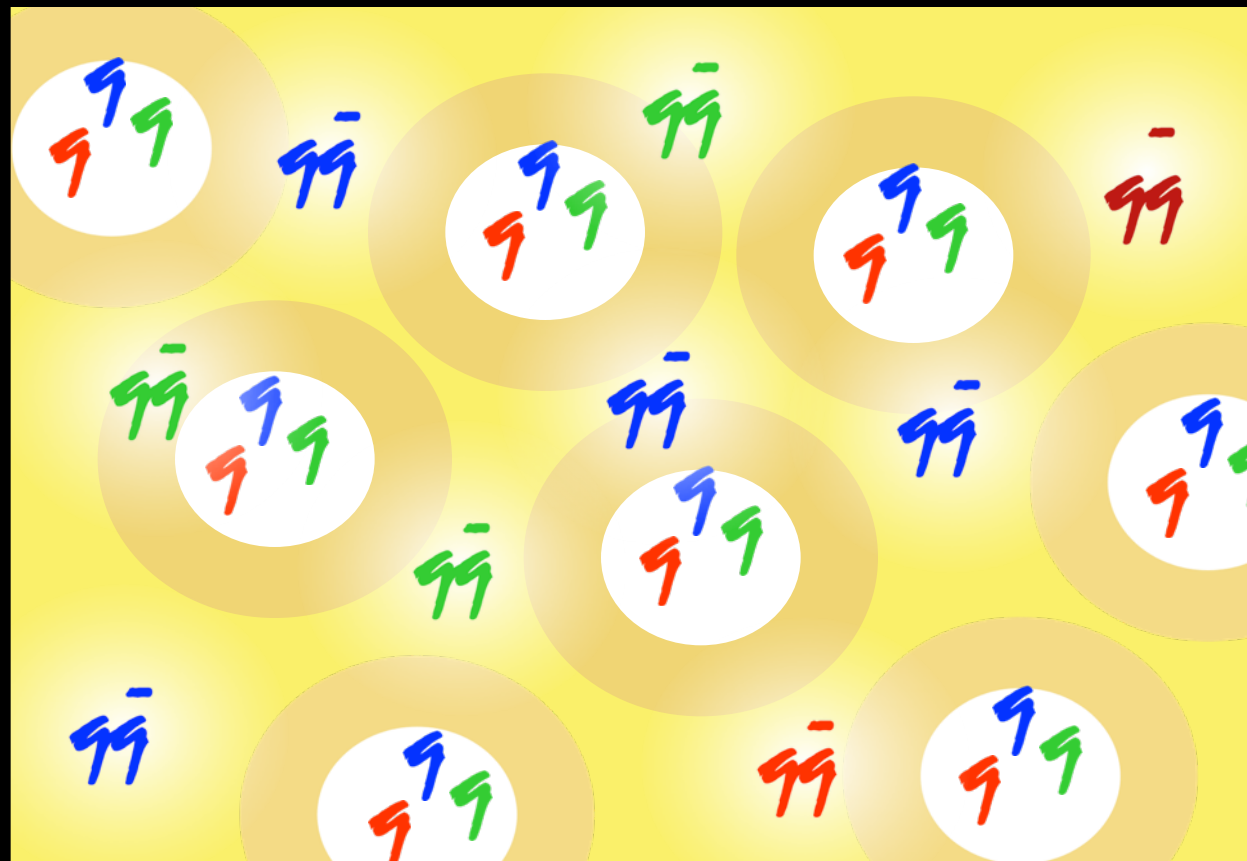
Swiss Cheese Picture of Nuclear Matter



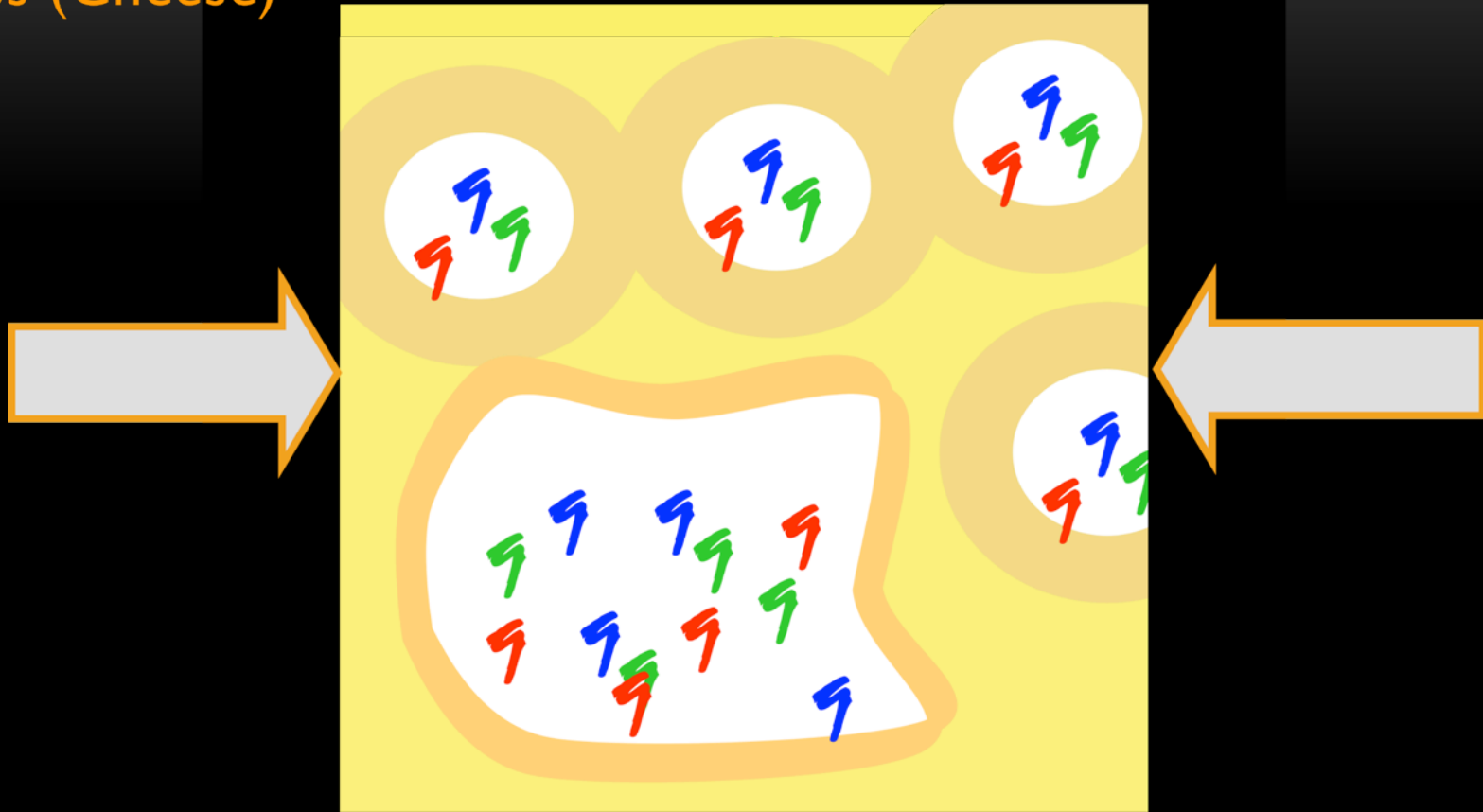
QCD condensates

Meson Cloud

Swiss (Cheese) Fondue



More Swiss (Cheese) Raclette



Dileptons as barometer

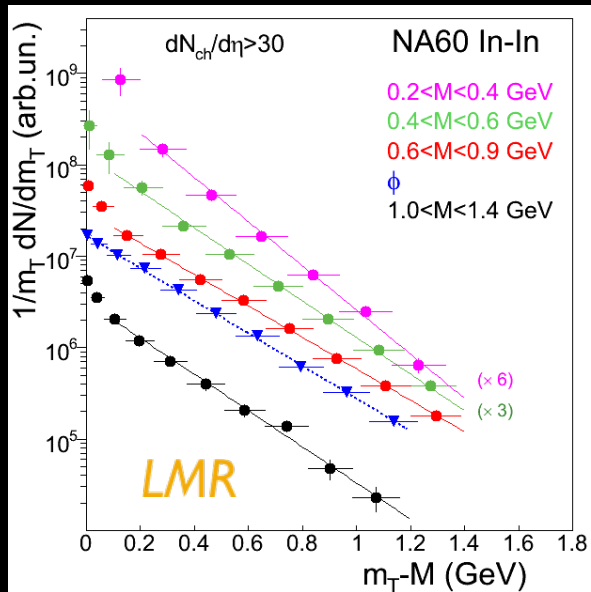


Transverse mass distributions of excess

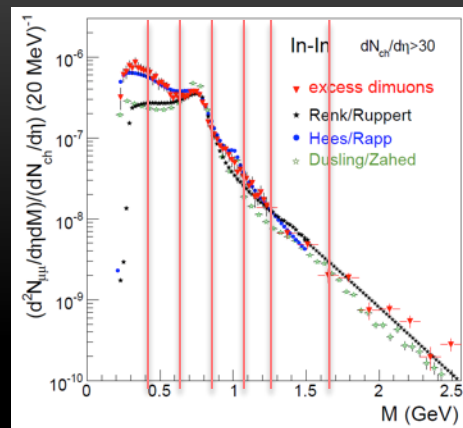
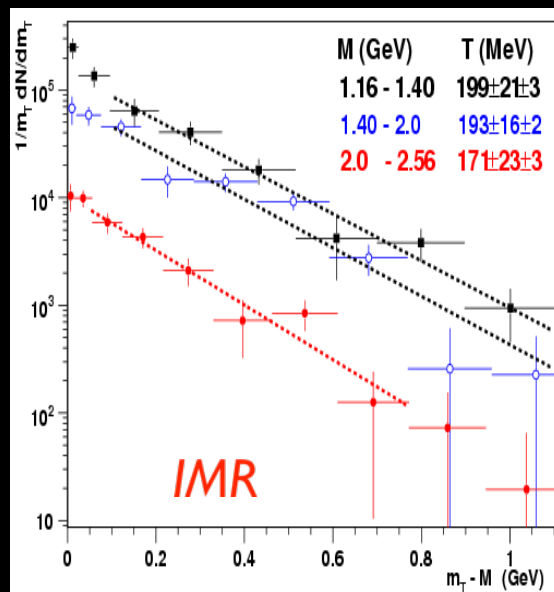
- For each bin of $\mu^+\mu^-$ project transverse mass spectrum: $m_T = (p_T^2 + M^2)^{1/2}$



Phys. Rev. Lett. 100 (2008) 022302

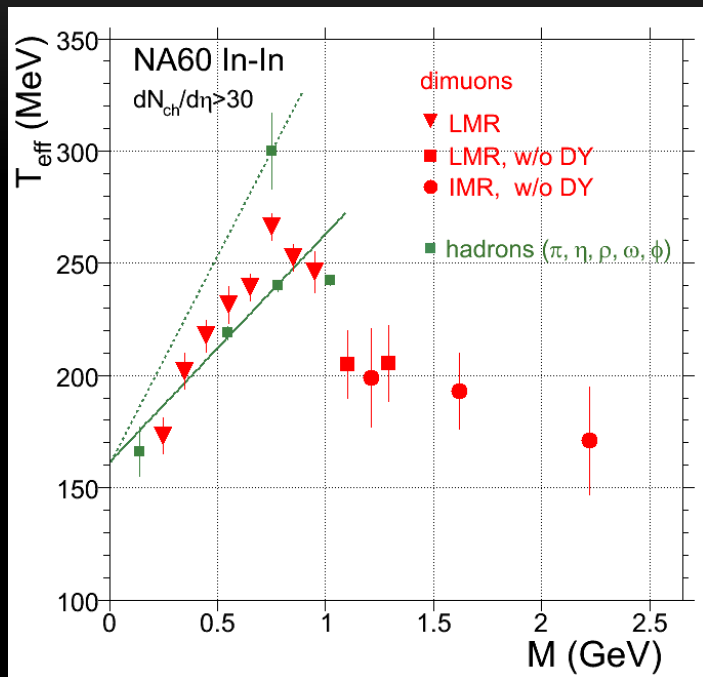



Eur. Phys. J. C 59 (2009) 607



- m_T spectra exponential for $m_T M > 0.1$ GeV; < 0.1 GeV ??
- Fit with $1/m_T dN/dm_T \sim \exp(-m_T/T_{eff})$
- Extract T_{eff} as a function of mass

The rise and fall of T_{eff} of thermal dimuons



 Phys. Rev. Lett. 100 (2008) 022302

- $M < 1 \text{ GeV}/c^2$
 - Strong, almost linear rise of T_{eff} with dimuon mass
 - Follows trend set by hadrons
- $M > 1 \text{ GeV}/c^2$
 - Drop of T_{eff} by $\sim 50 \text{ MeV}$
 - followed by an almost flat plateau

What can we learn from m_{T} spectra?

→ Radial Flow

→ Origin of dileptons

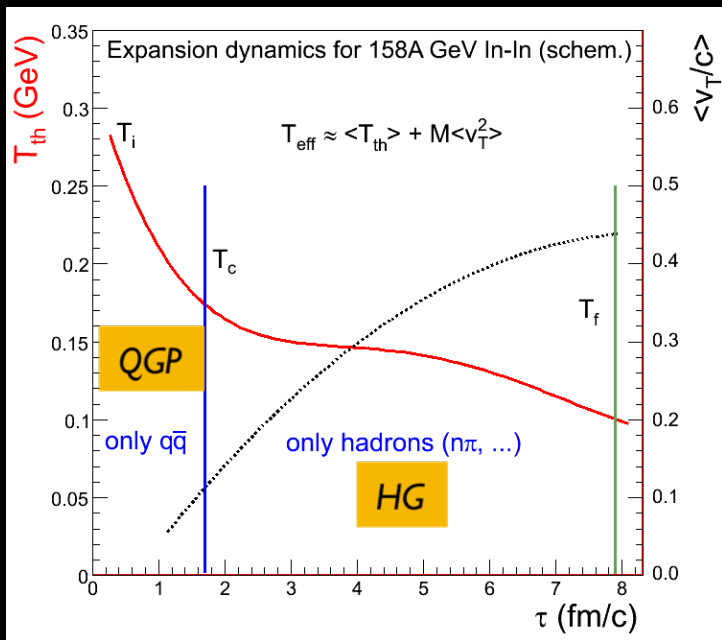
Interpretation of the dilepton m_T (p_T) spectra

T - dependence of **thermal distribution** of “mother” hadrons/partons

m - dependent **collective radial flow** (v_T) of “mother” hadrons/partons

(p_T - dependence of **spectral function**; dispersion relation)

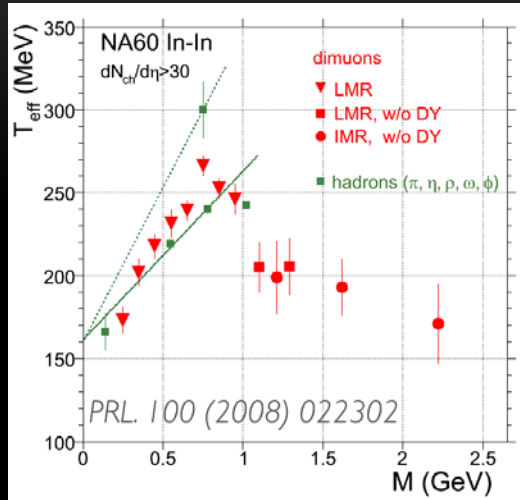
$$p_T = p_T^{th} + m v_T \quad \rightarrow \quad T_{eff} \sim T_f + \frac{1}{2} m \langle v_T^2 \rangle$$



- Hadron p_T spectra: determined at T_{final} (restricted information)
- Dilepton p_T spectra: superposition from all fireball stages
 - early emission \rightarrow high T , low v_T
 - late emission \rightarrow low T , high v_T
- Final spectra from space-time folding over T - v_T history from $T_{initial} \rightarrow T_{final}$
note: small flow in the QGP phase

\rightarrow Handle on emission region, i.e. nature of emitting source

Quantifying the average temperatures $\langle T_{th} \rangle$



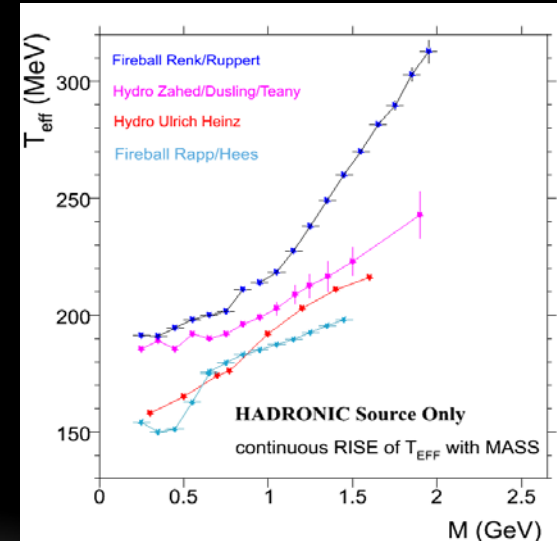
□ $M < 1 \text{ GeV}/c^2$

- extrapolate T_{eff} to $M=0$ (zero flow)
- $\langle T_{th} \rangle = 130 - 140 < T_c = 170 \text{ (MeV)}$
- *hadronic phase*

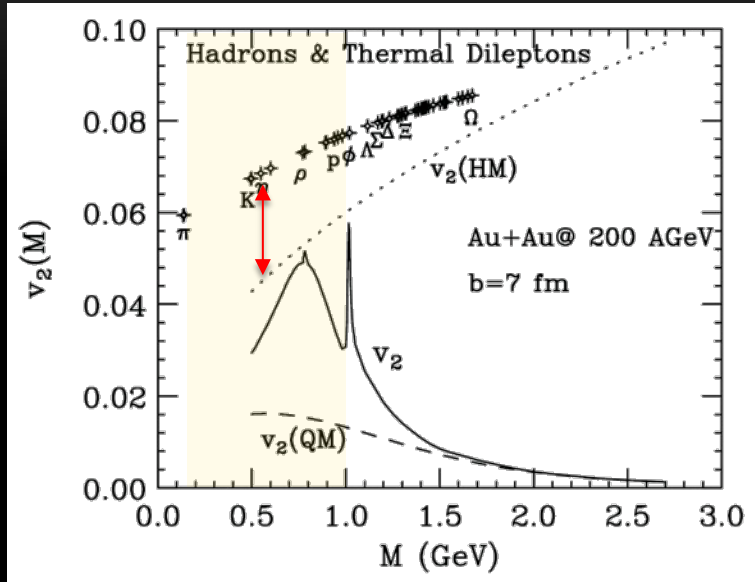
□ $M > 1 \text{ GeV}/c^2$

- T_{eff} independent of mass, negligible flow
- $\langle T_{th} \rangle \sim 200 \text{ MeV} > T_c = 170 \text{ (MeV)}$
- *partonic phase* ($T_{\text{initial}} \sim 250 \text{ MeV}$)

- Dominance of partons for $M > 1 \text{ GeV}/c^2$: support from theory:
 - Hadronic sources alone ($2\pi + 4\pi + a|\pi$) cannot produce a discontinuity
 - Continuous rise of T_{eff} with mass



Azimuthal anisotropy of virtual photons



R. Chatterjee et al., PRC 75 (2007), 054909

- Very clean tool to diagnose the collective expansion dynamics, i.e. origin of the electromagnetic emission source
- v_2 vs T_{eff} : T_{eff} is superposition of T and v_T
- $M < 1$ GeV/ c^2 : v_2 is large
→ late emissions → hadronic matter
- $M > 1$ GeV/ c^2 : v_2 is small
→ early emission → partonic matter

So far:

- STAR v_2 of inclusive e^+e^- (not of excess e^+e^-)
- HADES v_2 of excess radiation (coming soon)