Dibaryon candidates from Lattice QCD

~~Phase shift in strangeness interactions~~

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for HAL QCD Collaboration

Hadrons to Atomic nuclei

HAL

from Lattice QCD

HAL (Hadrons to Atomic nuclei from Lattice) QCD Collaboration

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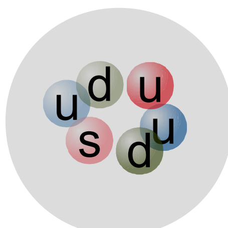
- ***What is dibaryon?***

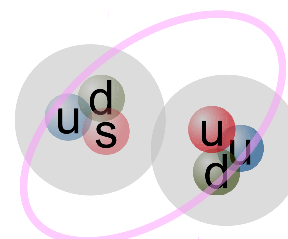
Bound and/or resonance two-baryon states

Tightly bound system



Compact exotic hadron

Molecular-like system



Deuteron-like behavior

Depend on the details of their interaction

Short range interactions

Repulsive core for general case of BB interactions?

Long range interactions

Strength of meson exchange contributions for each channel

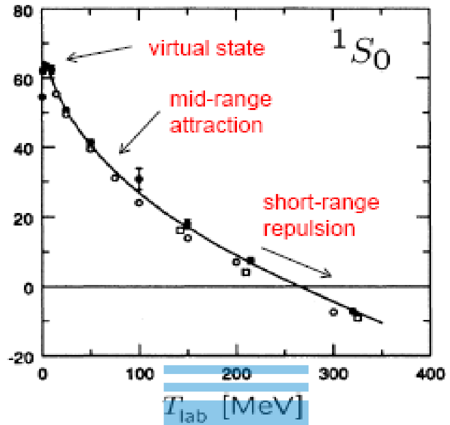
We need deep understandings of baryon-baryon interaction

How do we obtain the baryon force?

Phenomenological descriptions

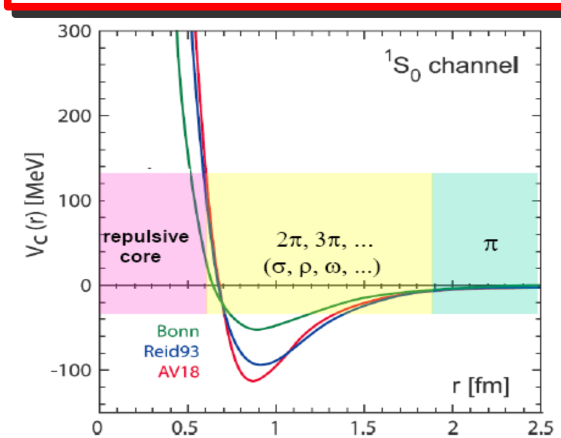
Traditional process to derive the BB interaction (potential)

Scattering observables



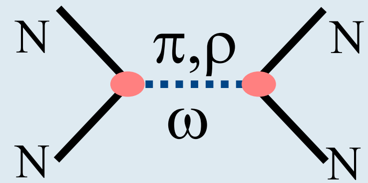
Model assumption

BB interaction (potential)



Meson exchange model

Described by hadron d.o.f.

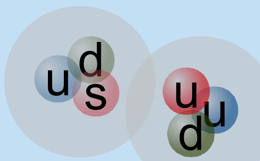


+ Phenomenological repul. core

H. Yukawa, PPMS17(1935)48
Th.A.Rijken et al, PTPS185(2010)14

Quark cluster model

Effective meson ex
+ quark anti-symmetrization

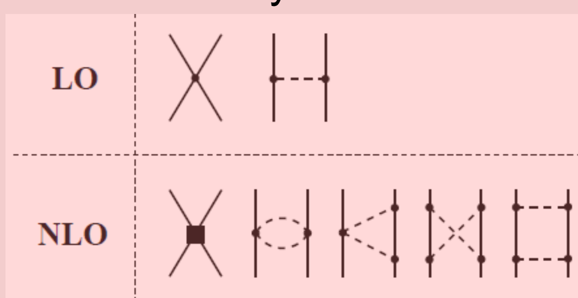


Quark Pauli effects
Color magnetic int.

M.Oka et al, PTPS137(2000)1
Y.Fujiwara et al, PPNP58(2007)439

Effective Field theory

Systematic calc. respecting with symmetry of QCD



Short range interaction is parametrized by contact term

E. Epelbaum et al, RMP81(2009)1773
R. Machleidt et al, PRept.503(2011)1
J. Haidenbauer et al, NPA954(2016)273

The models would be highly ambiguous if experimental data are scarce!

Clue to explore dibaryon candidates

We focus on the short range behavior of BB potential.

↔ Related to the tightly bound system.

In view of constituent quark cluster picture

- Short range interaction in between two baryons could be a result of Pauli principle and color-magnetic interaction for the quarks.

- Symmetry of constituent quarks
  - Assuming that all quarks are in s-orbit,
    Flavor SU(3) x Spin SU(2) x color SU(3)
    If totally anti-symmetric : Pauli allowed state
    If not : Pauli forbidden state

- Gluonic interaction between quarks at short range region
  - Gluon exchange contribution generates a color magnetic interaction

$$V_{OGE}^{CMI} \propto \frac{1}{m_{q1} m_{q2}} \langle \lambda_1 \cdot \lambda_2 \sigma_1 \cdot \sigma_2 \rangle f(r_{ij})$$

Dibaryon candidates

Possibility of dibaryon from short range interactions

- **For octet-octet system**

$$8 \otimes 8 = 1 \oplus 8_s \oplus 27 \oplus 8_a \oplus 10 \oplus \overline{10}$$

Spin 0 → **Pauli allowed**
Attractive color magnetic int. → **H-dibaryon** (Coupled $\Lambda\Lambda$–$N\Xi$–$\Sigma\Sigma$ system)

- **For decuplet-octet system**

$$10 \otimes 8 = 35 \oplus 8 \oplus 10 \oplus 27$$

Spin 2 → **Pauli allowed**
Attractive color magnetic int. → **N-Ω system** (Ground state of coupled $N\Omega$–$\Lambda\Xi^*$–$\Sigma\Xi^*$–$\Xi\Sigma^*$ system)

- **For decuplet-decuplet system**

$$10 \otimes 10 = 28 \oplus 27 \oplus 35 \oplus \overline{10}$$

Spin 3 → **Pauli allowed**
(Weakly) Attractive color magnetic int. → **$\Delta\Delta$ system**

Spin 0 → **Pauli allowed**
Repulsive color magnetic int. → **$\Omega\Omega$ system**

Dibaryon candidates

Some model calculations are performed for dibaryon candidates

- **H-dibaryon**
  - R.L.Jaffe PRL38(1977)
- **N-Ω system**
  - F.Wang et al. PRC51(1995)
  - Q.B.Li, P.N.Shen, EPJA8(2000)
- **ΔΔ and ΩΩ system**
  - F.J.Dyson,N-H.Xuong, PRL13(1964)
  - M.Oka, K.Yazaki, PLB 90(1980)
  - J. Haidenbauer, et al, nucl-th/1708.08071

- **Predicted B.E. and structures are highly depend on the model parameters.**
- **Some of them are still not found in experiments.**

We tackle this problem by Lattice QCD.

Baryon interaction from LQCD

BB interaction from QCD

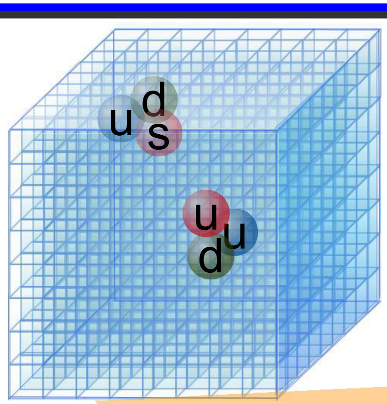
QCD is the fundamental theory of strong interactions

QCD Lagrangian

$$L_{QCD} = \bar{q}(i\gamma_\mu D^\mu - m)q + \frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu}$$

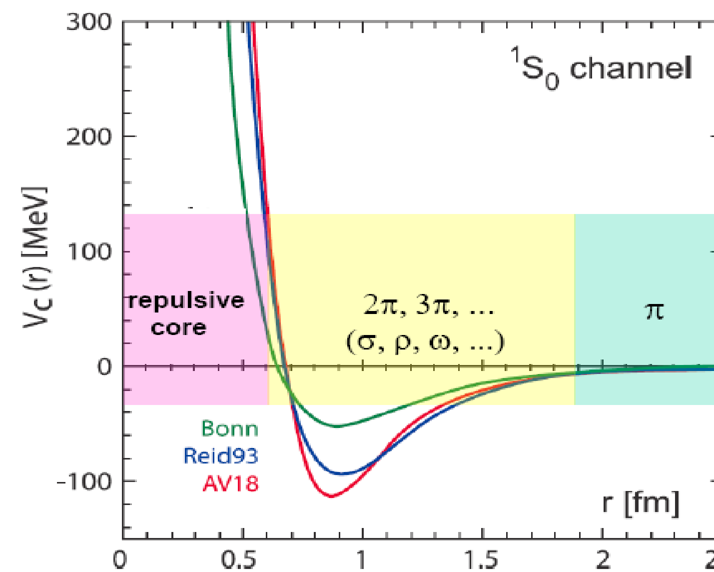
it is difficult to solve analytically the dynamics of quarks and gluons because of its non-perturbative nature at low-energies.

Lattice QCD simulation



- **Non-perturbative calculation.**
- **Hige computer resource is required.**
- **Independent of experimental situation.**

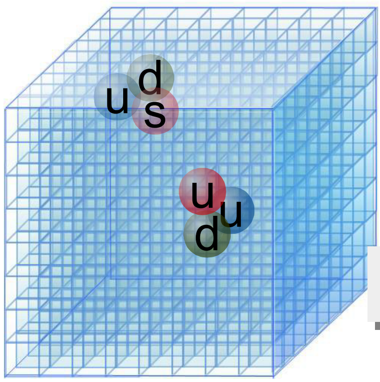
BB interaction (potential)



HAL QCD method

Hadron interaction from LQCD

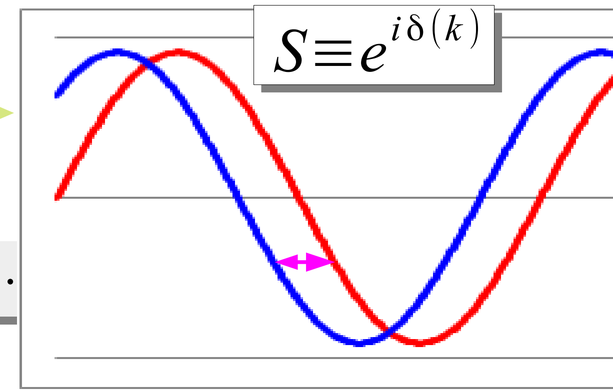
Lattice QCD simulation



HAL QCD method

Ishii, Aoki, Hatsuda, PRL**99** (2007) 022001

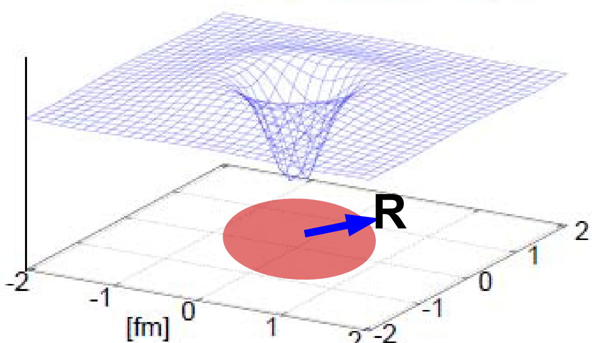
Scattering S-matrix



$$\langle 0 | B_1 B_2(t, \vec{r}) \bar{I}(t_0) | 0 \rangle = A_0 \Psi(\vec{r}, E_0) e^{-E_0(t-t_0)} + \cdots$$

$$\Psi(\vec{r}, E) e^{-Et} = \sum_{\vec{x}} \langle 0 | B_1(t, \vec{x}+\vec{r}) B_2(t, \vec{x}) | E \rangle$$

NBS wave function



$$(p^2 + \nabla^2) \Psi(\vec{r}, E) = 0, (r > R)$$

$$\Psi(\vec{r}, E) \simeq A \frac{\sin(pr + \delta(E))}{pr}$$

Phase shift is embedded in NBS w.f.

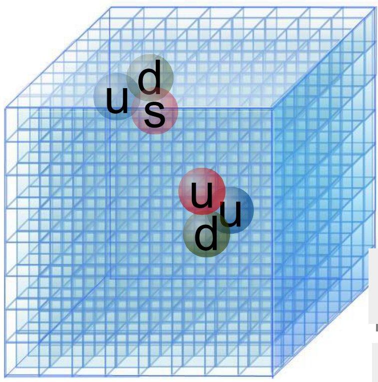
- **Inside of "interacting region"**

$$(p^2 + \nabla^2) \Psi^\alpha(E, \vec{x}) \equiv \int d^3 y \, U_\alpha^\alpha(\vec{x}, \vec{y}) \Psi^\alpha(E, \vec{y})$$

- **U(x,y) is faithful to the S-matrix.**
- **U(x,y) is not an observable.**
- **U(x,y) is energy independent but non-local.**

Hadron interaction (coupled-channel)

Lattice QCD simulation



HAL QCD method

S.Aoki et al [HAL] Proc. Jpn. Acad.,Ser.B,87 509

Scattering S-matrix

$$S(E)=\begin{pmatrix} \eta e^{2i\delta_1} & i\sqrt{1-\eta^2}e^{i(\delta_1+\delta_2)} \\ i\sqrt{1-\eta^2}e^{i(\delta_1+\delta_2)} & \eta e^{2i\delta_2} \end{pmatrix}$$

$$\langle 0|(B_1 B_2)^{\alpha}(t,\vec{r})\bar{I}(t_0)|0\rangle = A_0 \Psi^{\alpha}(\vec{r},E_0)e^{-E_0(t-t_0)}+\cdots$$

$$\langle 0|(B_1 B_2)^{\beta}(t,\vec{r})\bar{I}(t_0)|0\rangle = C_0 \Psi^{\beta}(\vec{r},E_0)e^{-E_0(t-t_0)}+\cdots$$

NBS wave function for each channel

$$\Psi^{\alpha}(\vec{r},E_i)e^{-E_i t} = \langle 0|(B_1 B_2)^{\alpha}(\vec{r})\,|E_i\rangle$$

$$\Psi^{\beta}(\vec{r},E_i)e^{-E_i t} = \langle 0|(B_1 B_2)^{\beta}(\vec{r})\,|E_i\rangle$$

Coupled-channel Schrödinger equation

$$\left(p_{\alpha}^2+\nabla^2\right)\Psi^{\alpha}(E,\vec{x}) \equiv \int d^3y\, U^{\alpha}_{\beta}(\vec{x},\vec{y})\Psi^{\beta}(E,\vec{y})$$

β: $\delta_{\alpha}, \delta_{\beta}, \eta$ ← $\Psi^{\alpha}_{sc}(\vec{r}), \Psi^{\beta}_{sc}(\vec{r})$

α: δ_{α} ← $\Psi^{\alpha}_{sc}(\vec{r})$

- **U(x,y) is faithful to the S-matrix beyond the threshold of channel β.**
- **U(x,y) is energy independent until the higher energy thershold opens.**
- **Derivative (velocity) expansion is used.**

S=-2 BB interaction

--- focus on the H-dibaryon ---

Keys to understand H-dibaryon state

A strongly bound state predicted by Jaffe in 1977 using MIT bag model.

H-dibaryon state is

- **SU(3) flavor singlet [uuddss], strangeness S=-2.**
- **spin and isospin equals to zero, and $J^P = 0^+$**

SU(3) classification

$$8 \times 8 = 27 + 8_S + 1 + \underline{10} + 10 + 8_A$$

▶ **Strongly attractive interaction is expected in flavor singlet channel.**

- **Strongly attractive Color Magnetic Interaction**
- **Flavor singlet channel is free from Pauli blocking effect**

	27	8	1	<u>10</u>	10	8
Pauli		forbidden	allowed		forbidden	
CMI	repulsive	repulsive	attractive	repulsive	repulsive	repulsive

M. Oka et al NPA464 (1987)

$J^p=0^+$, I=0

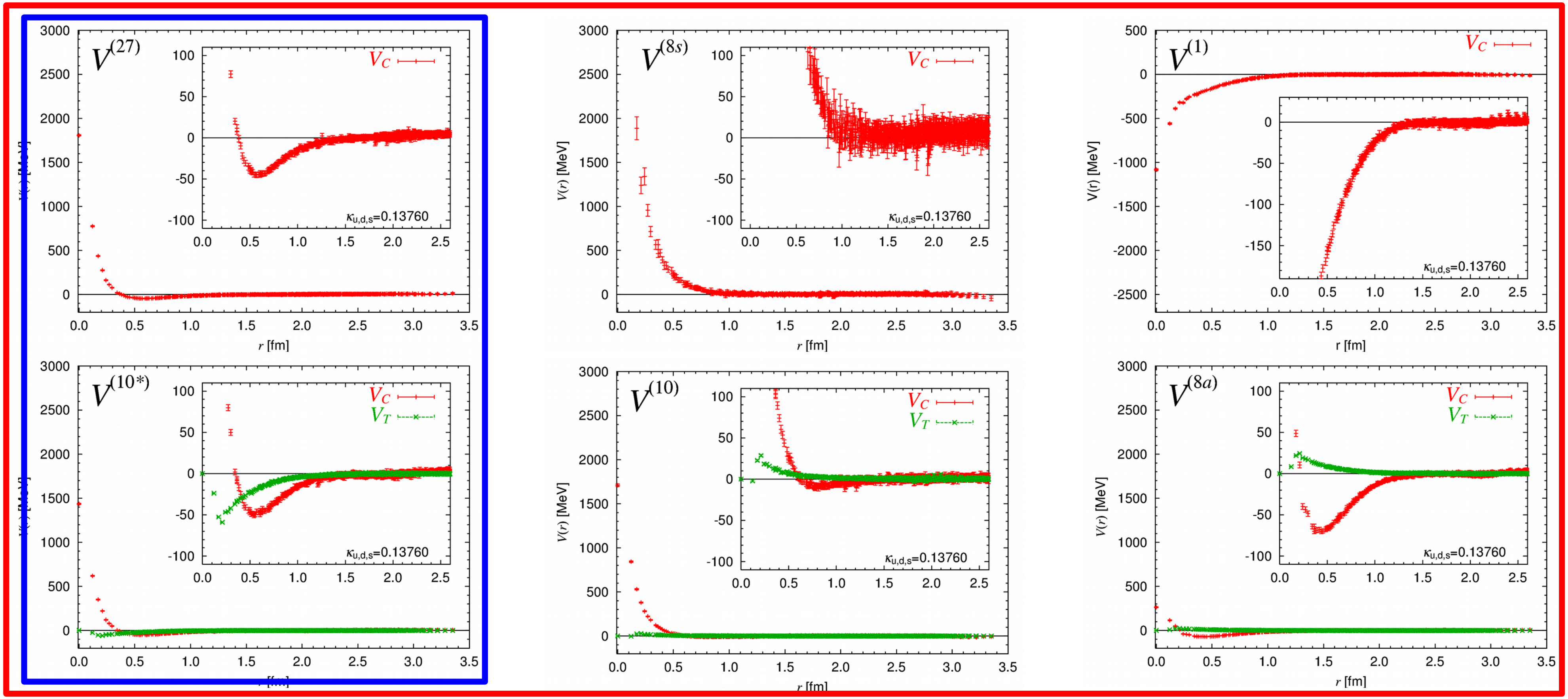
$$\begin{pmatrix} \Lambda\Lambda \\ N\Xi \\ \Sigma\Sigma \end{pmatrix} = \frac{1}{\sqrt{40}} \begin{pmatrix} -\sqrt{5} & -\sqrt{8} & \sqrt{27} \\ \sqrt{20} & \sqrt{8} & \sqrt{12} \\ \sqrt{15} & -\sqrt{24} & -1 \end{pmatrix} \begin{pmatrix} 1 \\ 8 \\ 27 \end{pmatrix}$$

B-B potentials in SU(3) limit

m_π= 469MeV

J=0

J=1



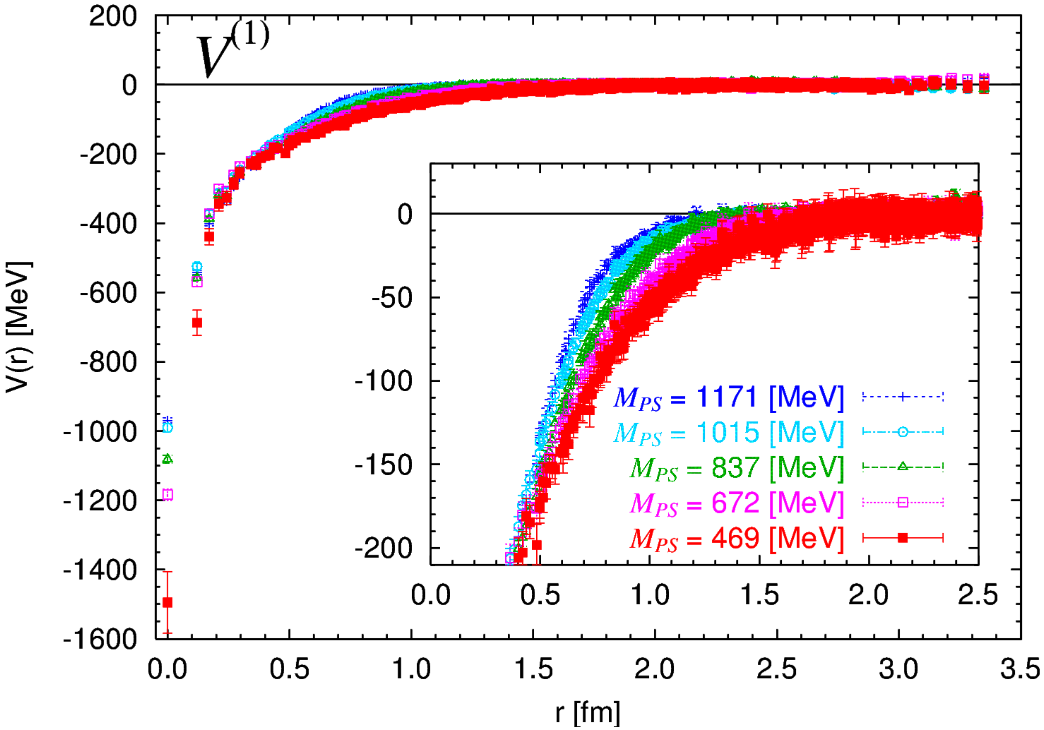
Two-flavors

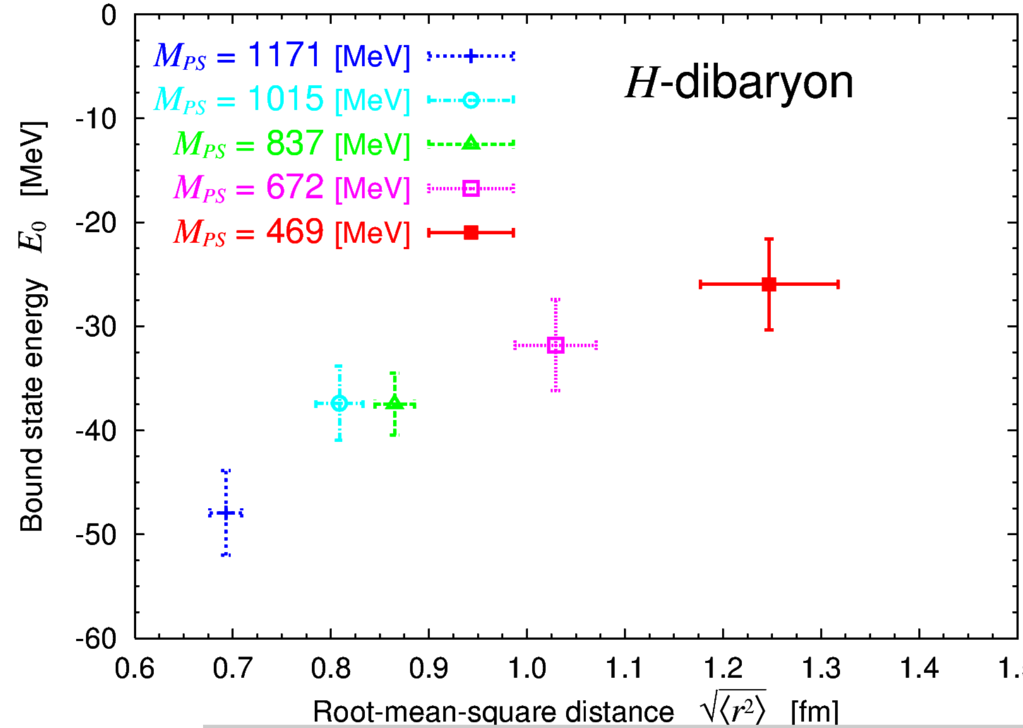
Three-flavors

- Quark Pauli principle can be seen at around short distances
  - No repulsive core in flavor singlet state
  - Strongest repulsion in flavor 8s state
- Possibility of bound H-dibaryon in flavor singlet channel.

H-dibaryon in SU(3) limit

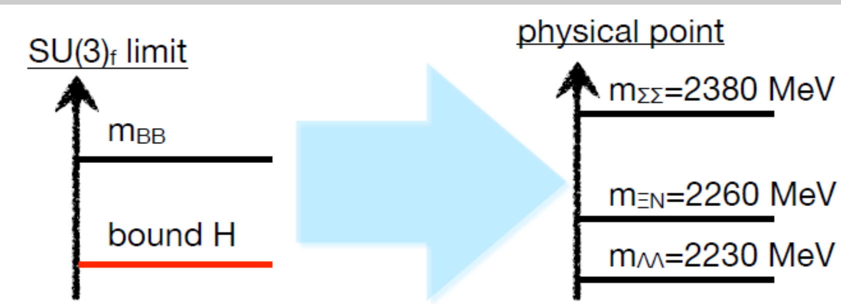
Strongly attractive interaction is expected in flavor singlet channel.





T.Inoue et al[HAL QCD Coll.] NPA881(2012) 28

m_B=1161MeV for M_{PS}=470MeV



- **Strongly attractive potential was found in the flavor singlet channel.**
- **Bound state was found in this mass range with SU(3) symmetry.**

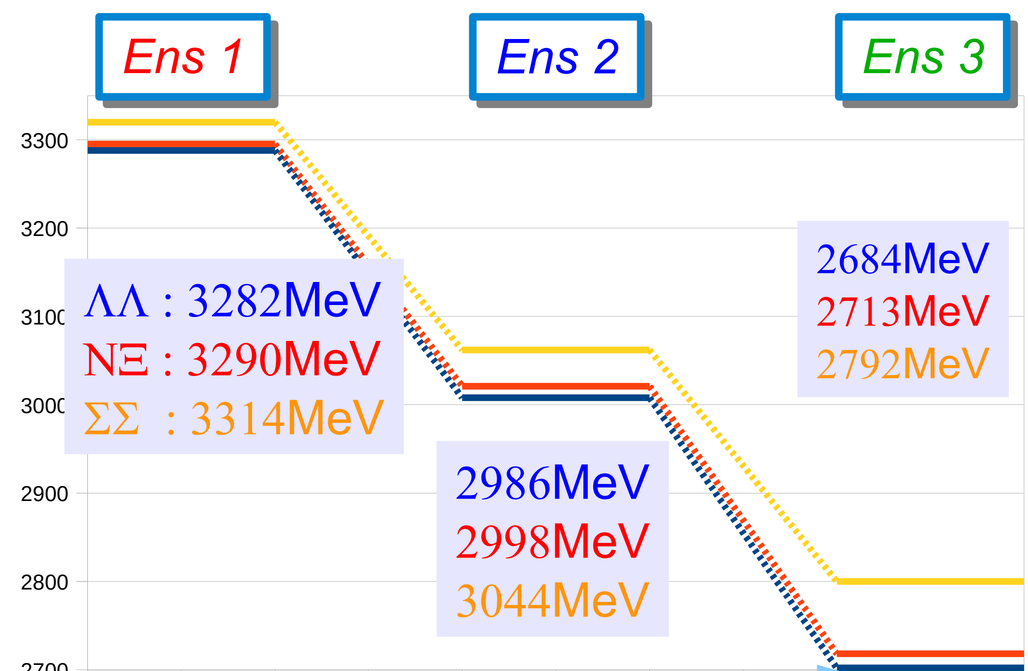
Go to the physical point simulation!

Numerical setup

- 2+1 flavor gauge configurations by PACS-CS collaboration.
  - RG improved gauge action & O(a) improved Wilson quark action
  - β = 1.90, a^{-1} = 2.176 [*GeV*], 32^3x64 lattice, *L* = 2.902 [*fm*].
  - κ_s = 0.13640 is fixed, κ_{ud} = 0.13700, 0.13727 and 0.13754 are chosen.
- Wall source is considered to produce S-wave B-B state.

In unit of MeV	*Ens 1*	*Ens 2*	*Ens 3*
π	701±1	570±2	411±2
K	789±1	713±2	635±2
m_π/m_K	0.89	0.80	0.65
N	1581±5	1398±7	1215±8
Λ	1641±5	1493±6	1342±6
Σ	1657±5	1522±7	1394±8
Ξ	1709±4	1600±5	1498±5

u,d quark masses lighter



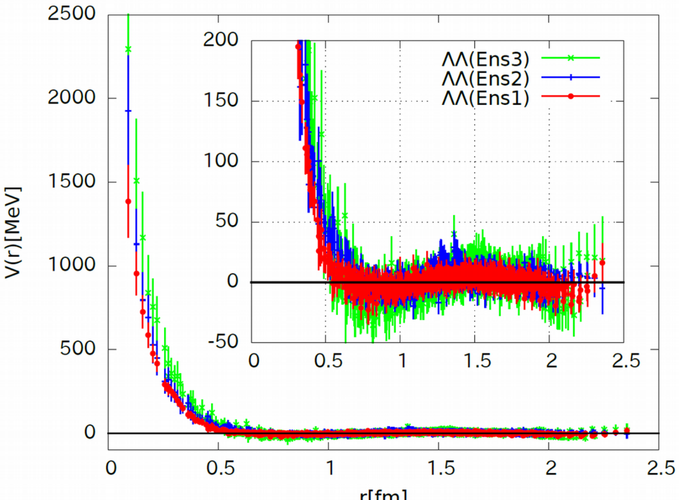
SU(3) breaking effects becomes larger

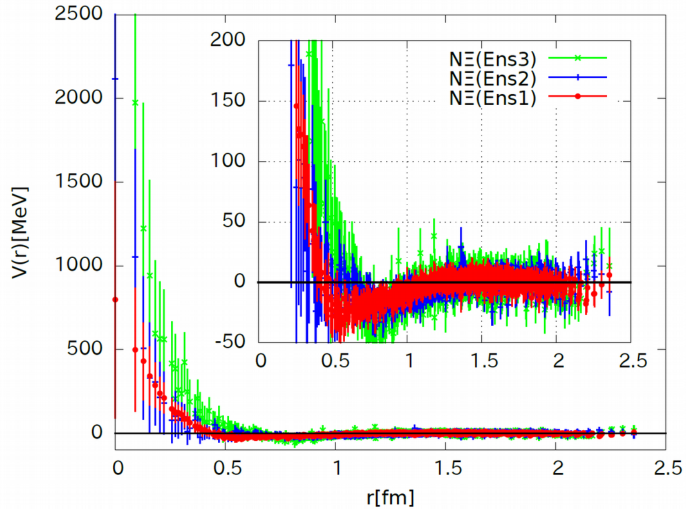
ΛΛ, NΞ, ΣΣ (I=0) $^1S_0$ channel

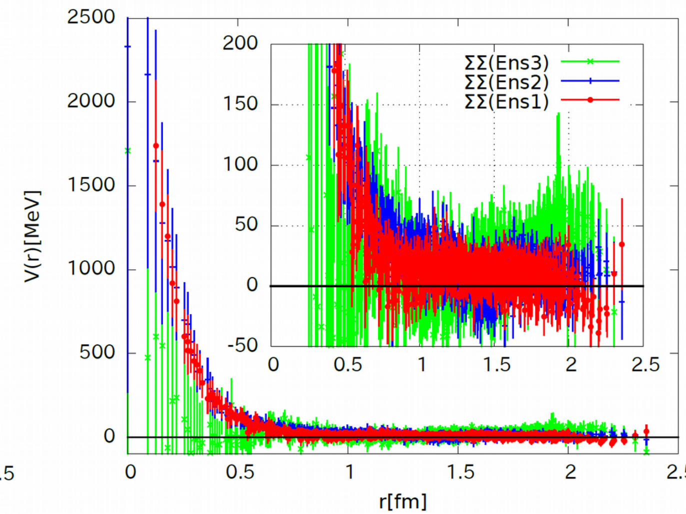
Esb1 : m_π= 701 MeV
Esb2 : m_π= 570 MeV
Esb3 : m_π= 411 MeV

- **N_f = 2+1 full QCD with L = 2.9fm**

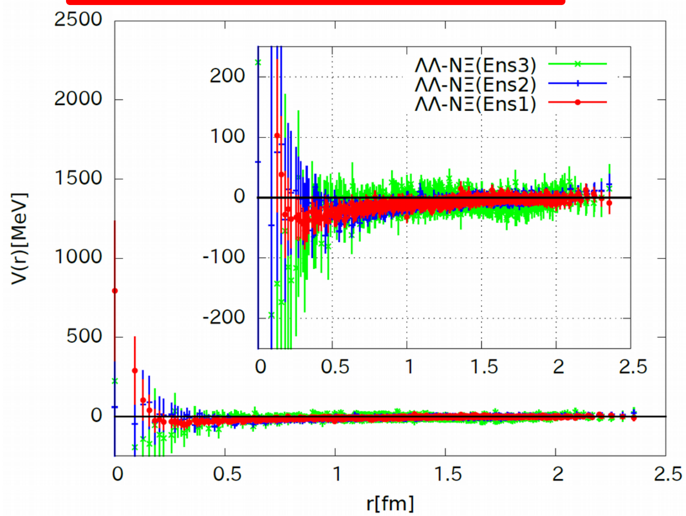
Diagonal elements

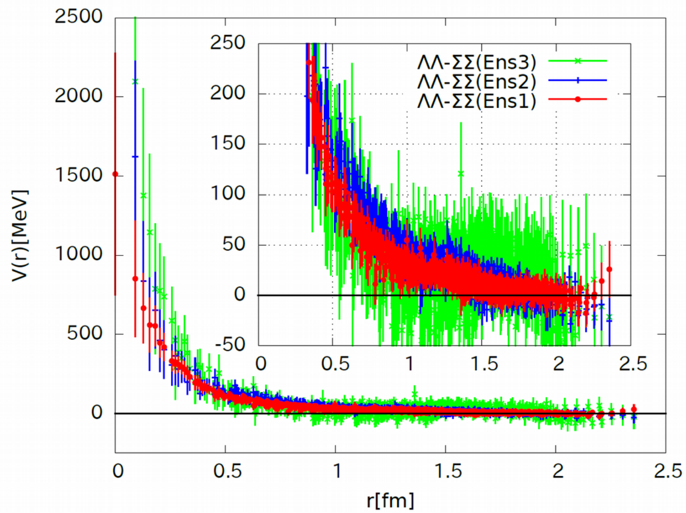


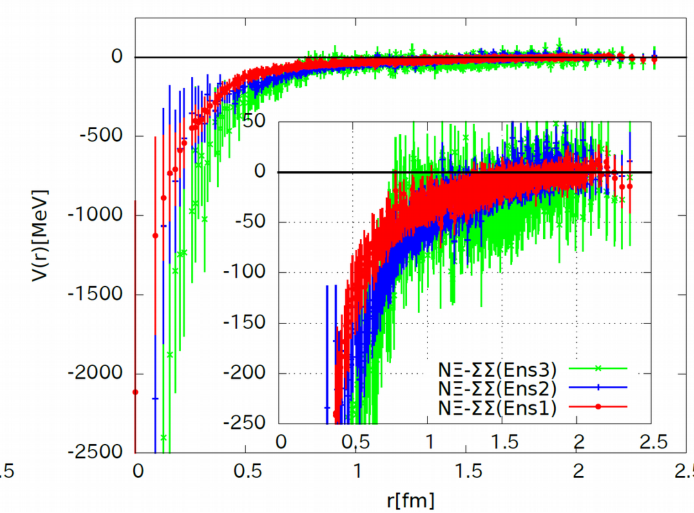




Off-diagonal elements



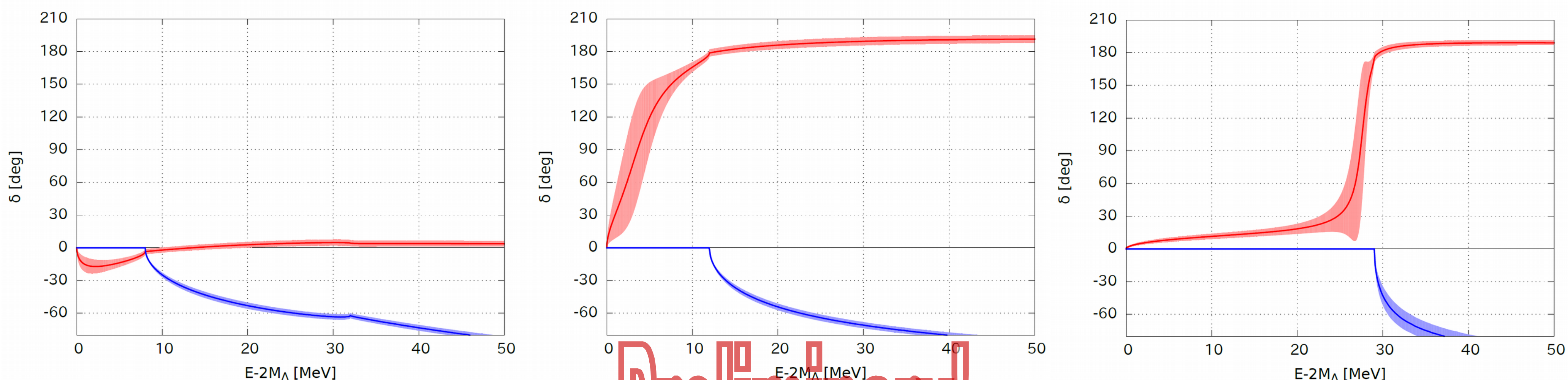




$\Lambda\Lambda$ and $N\Xi$ phase shifts

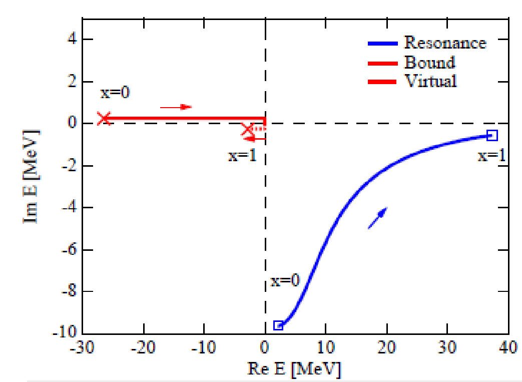
- $N_f = 2+1$ full QCD with L = 2.9fm

m_π = 700 MeV ***m_π = 570 MeV*** ***m_π = 410 MeV***



Preliminary!

- ***m_π = 700 MeV :*** bound state
- ***m_π = 570 MeV*** : resonance near $\Lambda\Lambda$ threshold
- ***m_π = 410 MeV*** : resonance near $N\Xi$ threshold..

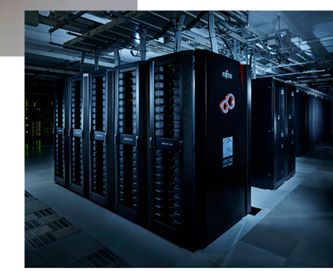


Y.Yamaguchi and T.Hyodo PRC 94 (2016) 065207

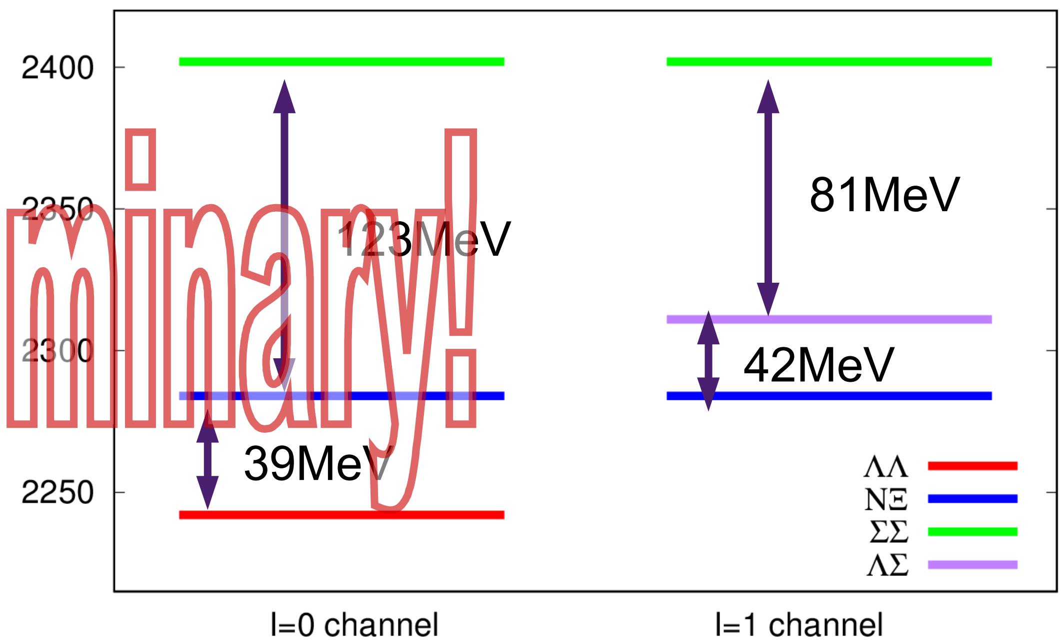
Go to the physical point simulation!

Numerical setup

- 2+1 flavor gauge configurations.
  - Iwasaki gauge action & O(a) improved Wilson quark action
  - a = 0.085 [fm], a^{-1} = 2.300 GeV.
  - 96^3x96 lattice, L = 8.21 [fm].
  - 414 confs x 84 sources x 4 rotations.
- Wall source is considered to produce S-wave B-B state.



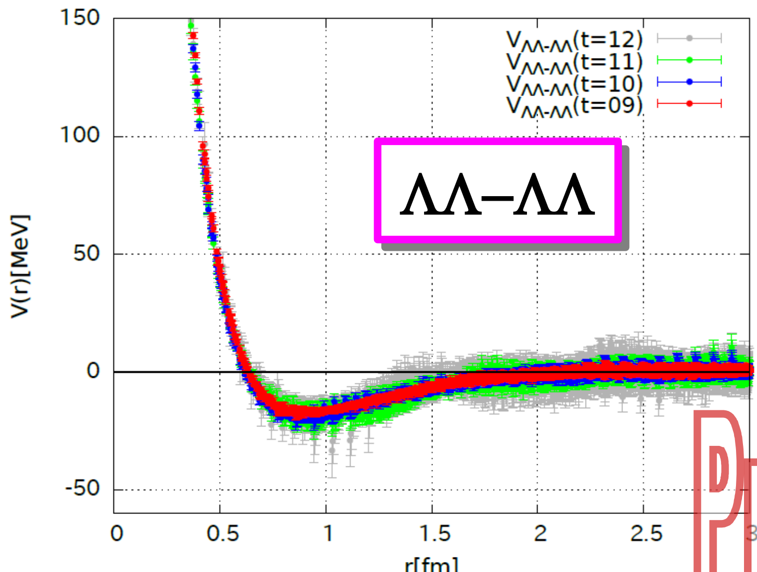
	Mass [MeV]
π	146
K	525
m_π/m_K	0.28
N	953±7
Λ	1123±3
Σ	1204±1
Ξ	1332±2

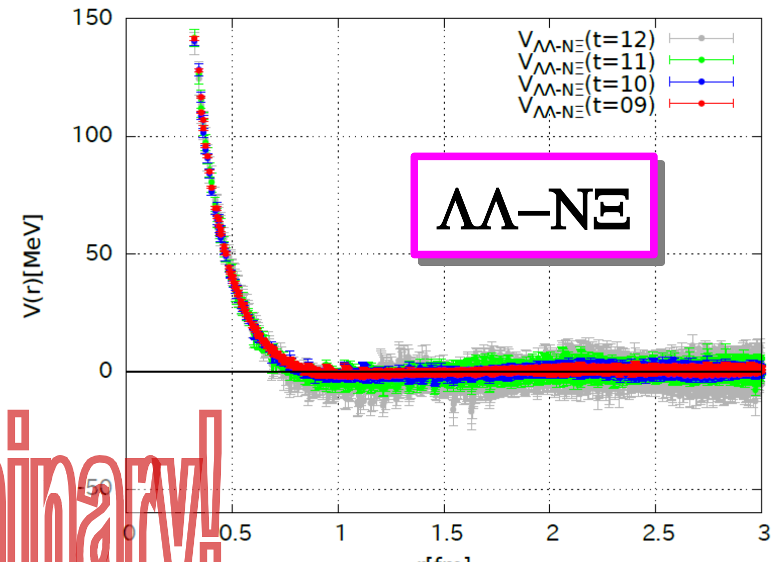


$\Lambda\Lambda$, $N\Xi$ (I=0) $^1S_0$ potential

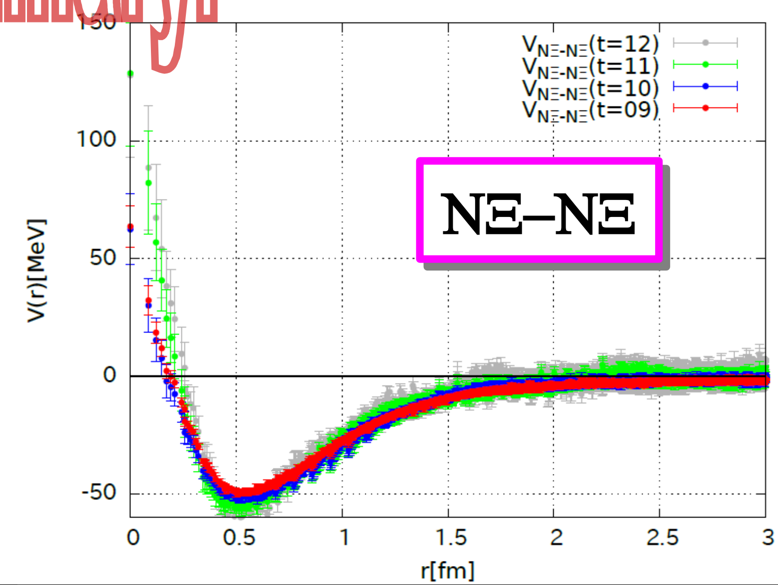
t=09
t=10
t=11
t=12

- N_f = 2+1 full QCD with L = 8.1fm, ***m_π = 146 MeV***





Preliminary!

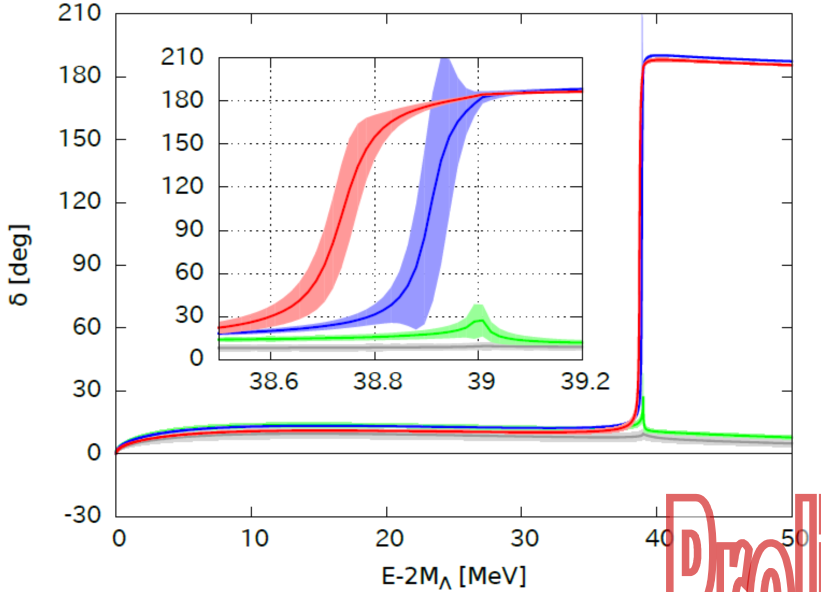


- **Potential calculated by only using $\Lambda\Lambda$ and $N\Xi$ channels.**
- **Long range part of potential is almost stable against the time slice.**
- **Short range part of $N\Xi$ potential changes as time t goes.**
- **$\Lambda\Lambda$–$N\Xi$ transition potential is quite small in r > 0.7fm region**

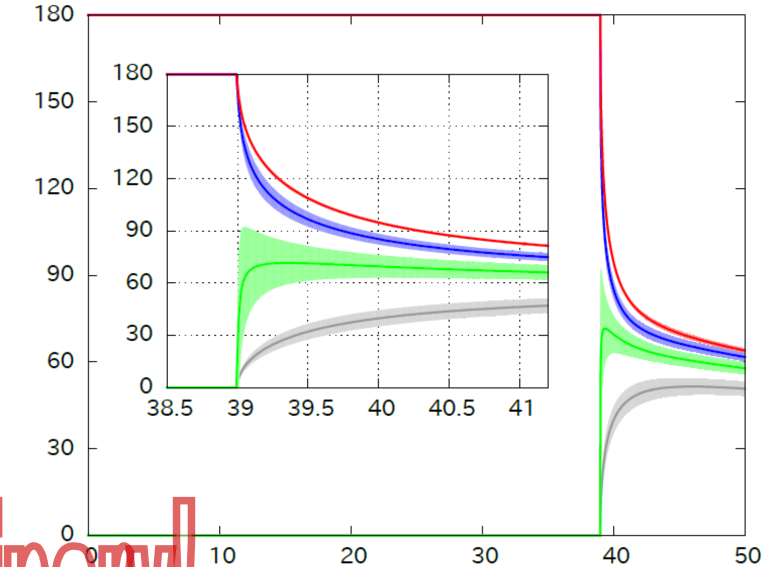
ΛΛ and NΞ phase shift and inelasticity

t=09
t=10
t=11
t=12

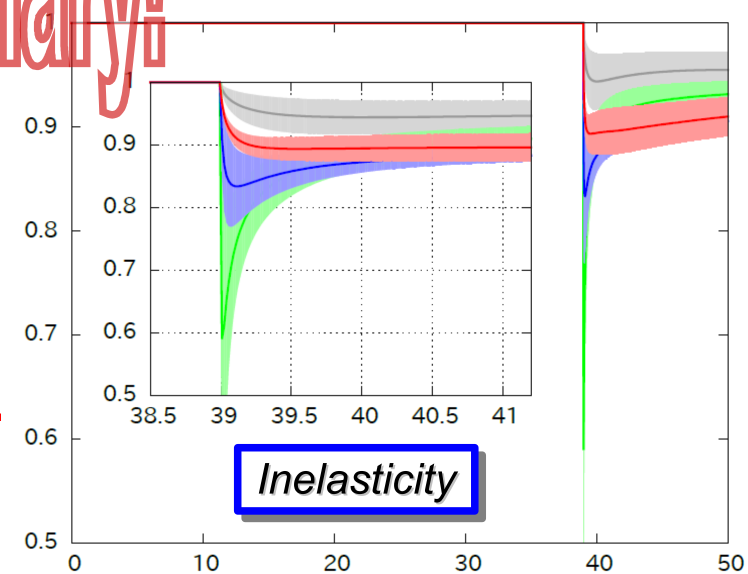
ΛΛ phase shift



NΞ phase shift



Preliminary!



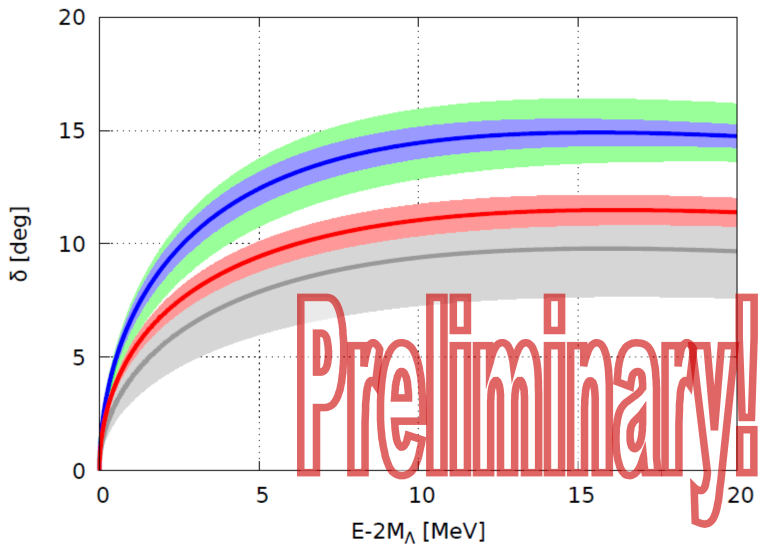
Inelasticity

- ΛΛ and NΞ phase shift is calculated
  by using 2ch effective potential.
- A sharp resonance is found
  below the NΞ threshold for t=9 and 10.

ΛΛ scattering length

t=09
t=10
t=11
t=12

Phase shift



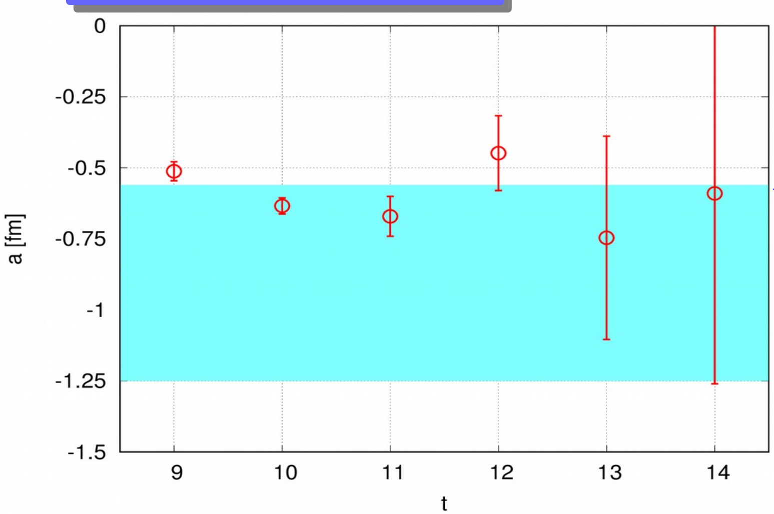
Λ		NLO			
		500	550	600	650
ΛΛ	a_{1S0}	−0.62	−0.61	−0.66	−0.70
	r_{1S0}	6.95	6.06	5.05	4.56

J. Haidenbauer et al, NPA954(2016)273

$a_{\Lambda\Lambda}=-0.821$fm **Y.Fujiwara et al, PPNP58(2007)439**

$a_{\Lambda\Lambda}=-0.97$fm **Th.A.Rijken et al, Few-Body Syst 54(2013)801**

Scattering length



$-1.25 < a_{\Lambda\Lambda} < -0.56$fm

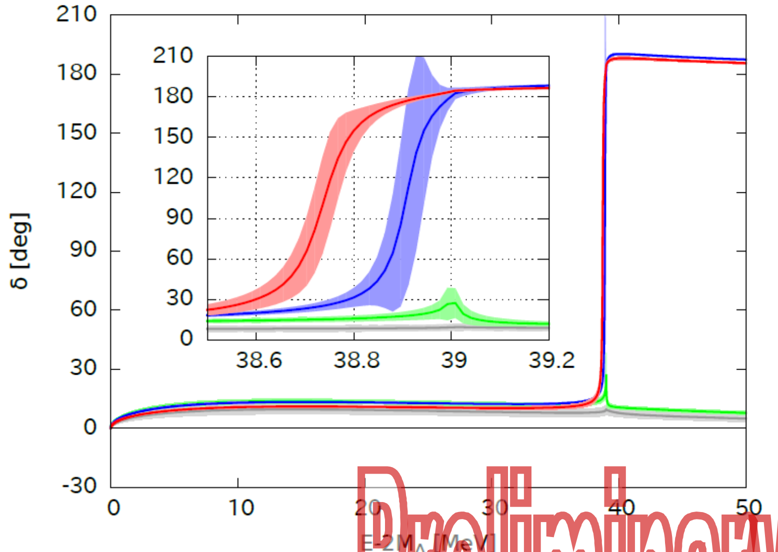
K.Morita et al, PRC91 (2015)024916

- **Scattering length in ΛΛ (I=0) is almost saturated.**
- **Attraction is a little bit weaker than the phenomenological values.**

ΛΛ and NΞ phase shift –comparison--

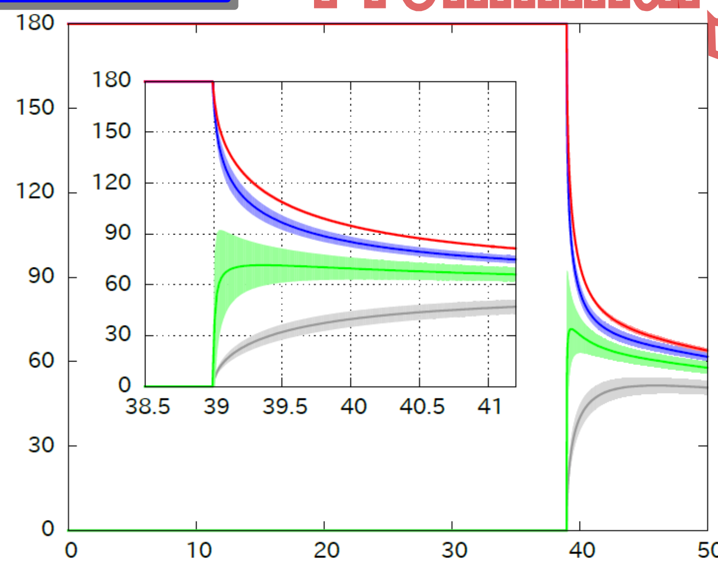
t=09
t=10
t=11
t=12

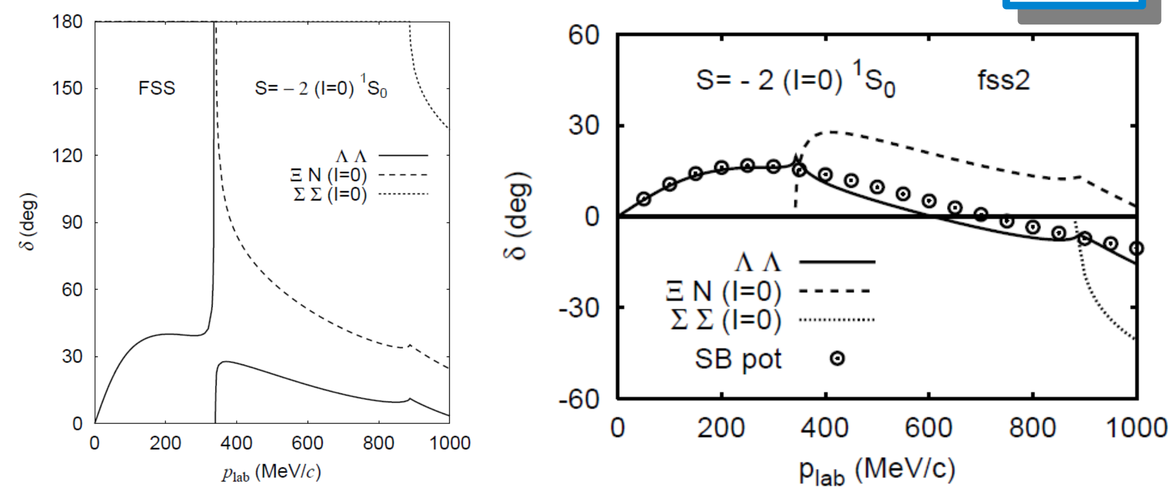
ΛΛ phase shift



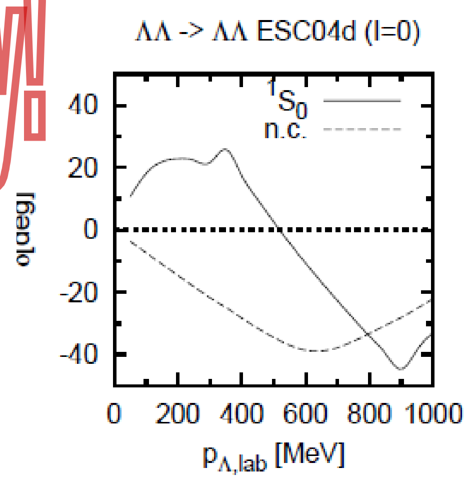
Preliminary!

NΞ phase shift

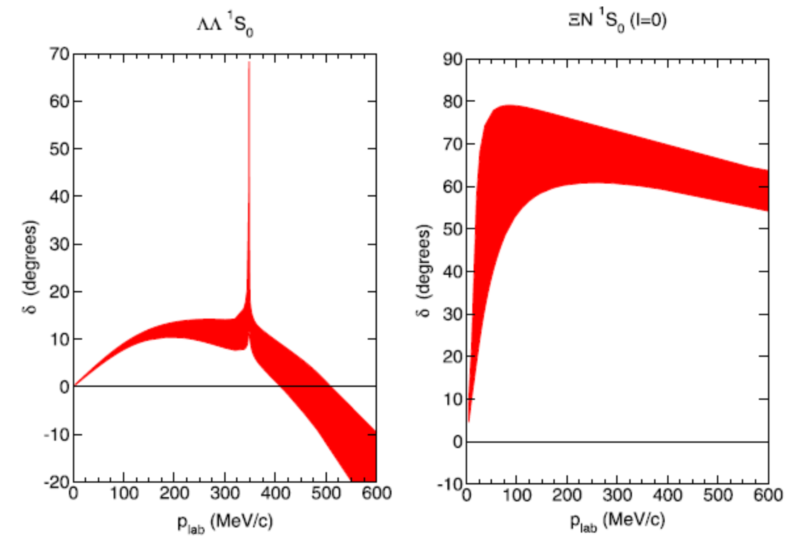




Y.Fujiwara et al, PPNP58(2007)439



Th.A. Rijken, nucl-th/060874



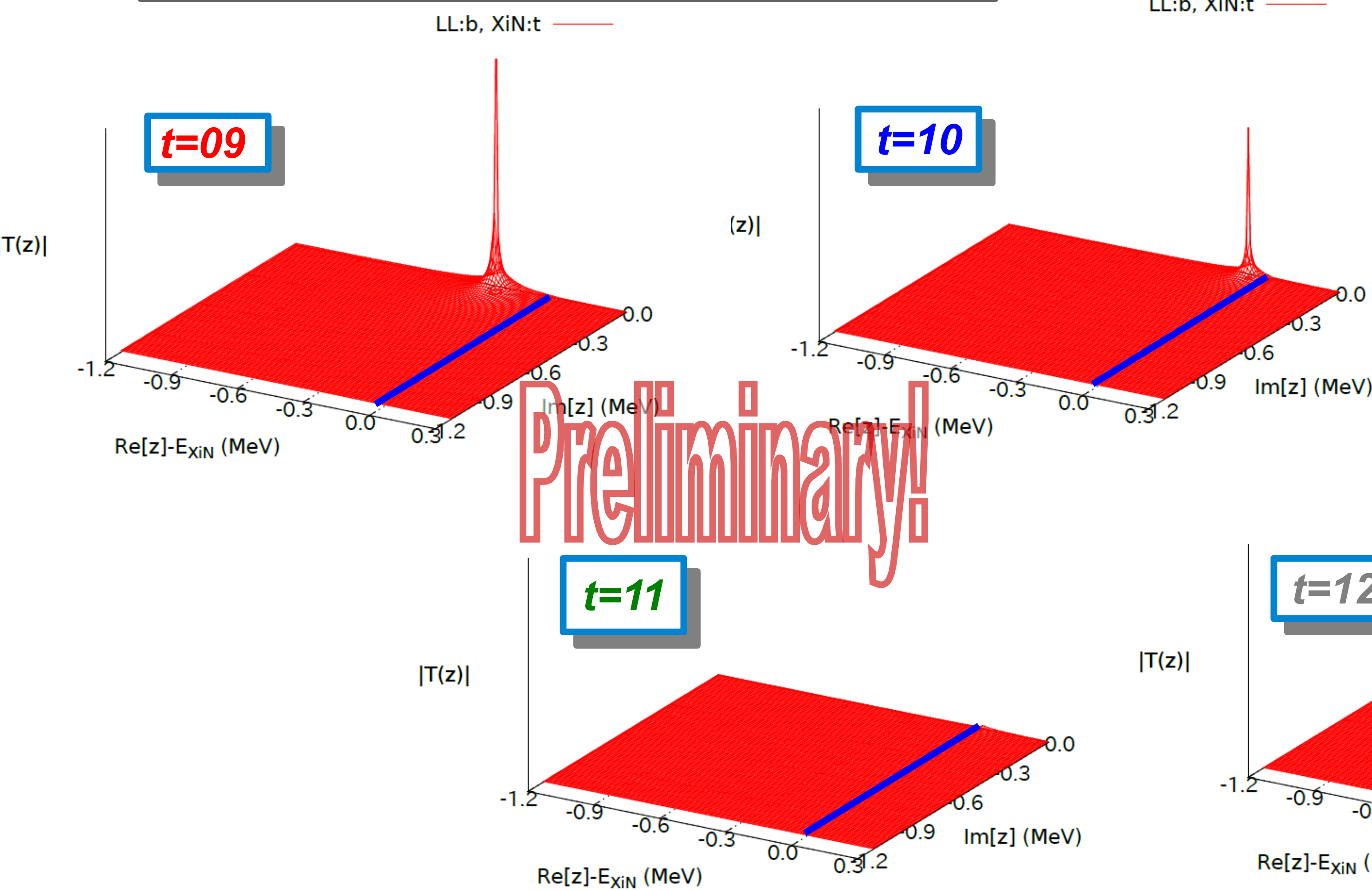
J. Haidenbauer et al, NPA954(2016)273

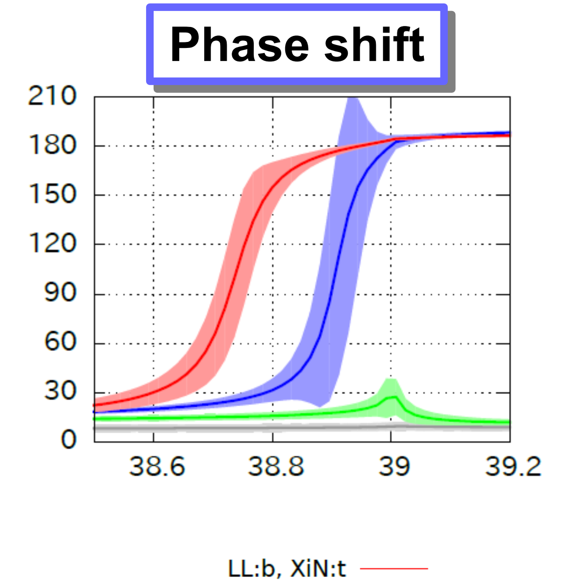
- Our results are compatible with the phenomenological ones.

Pole search

- N_f = 2+1 full QCD with L = 8.1fm, m_π = 146 MeV

T-matrix (ΛΛ : unphysical, NΞ : physical)





T-matrix pole is very close to the NX threshold

NΩ interaction

NΩ system

NΩ system from model calculations

- One of di-baryon candidate — T.Goldman et al PRL59(1987)627
- (Quasi-)Bound state is reported with J=2, I=1/2
  - Constituent quark model — M.Oka PRD38(1988)298
    - CMI does not contribute for this system because of no quark exchange between baryons.
    - Coupled channel effect is important.
  - Chiral quark model — Q.B.Li, P.N.Shen, EPJA8(2000)
    - Strong attraction yielded by scalar exchange

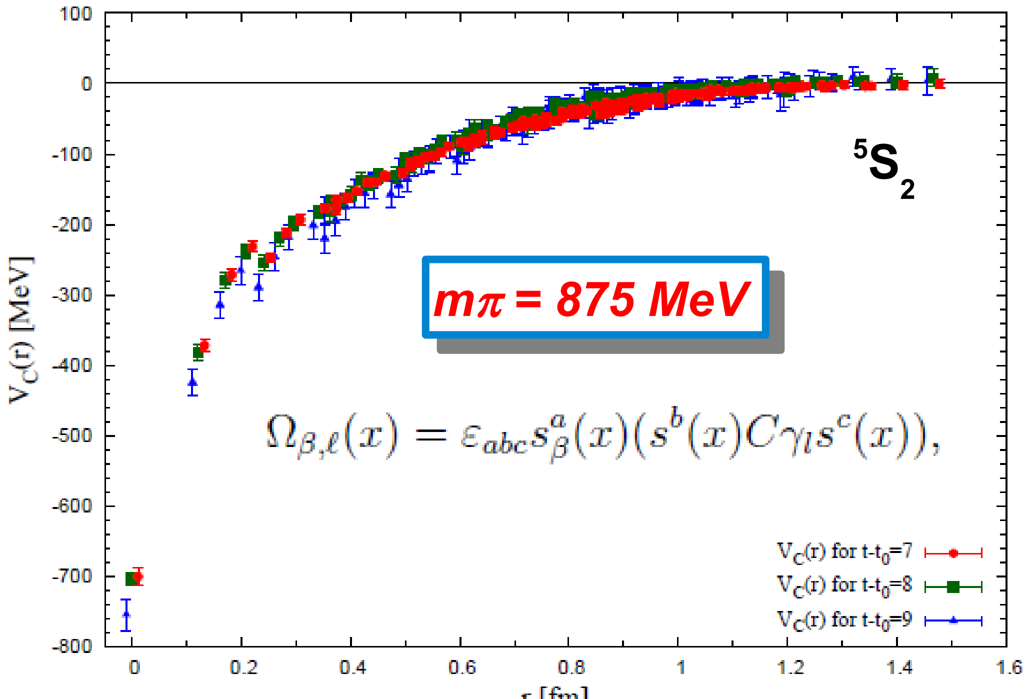
NΩ $J^p(I) = 2^+(1/2)$ is considered

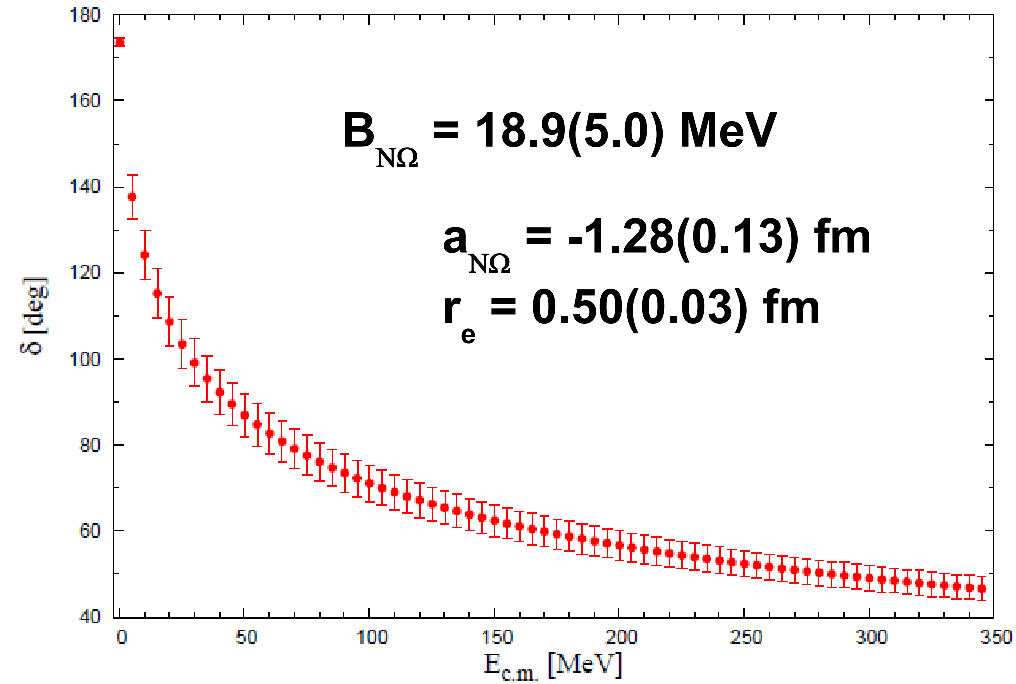
- Easy to tackle it by lattice QCD simulation
  - Lowest state in J=2 coupled channel
    - NΩ–ΛΞ*–ΣΞ*–ΞΣ*
  - Multi-strangeness reduces a statistical noise
  - Wick contraction is very simple

$N\Omega$ system $J^p(I) = 2^+(1/2)$

▶ $N_f = 2+1$ full QCD with L = 1.9fm

F.Etminan(HAL QCD), NPA928(2014)89





$N\Omega$ state cannot decay into $\Lambda\Xi$ (D-wave) state in this setup

- Strongly attractive S-wave effective potential in $J^p(I) = 2^+(1/2)$
- Measurement of strong $N\Omega$ attraction at RHIC and LHC is expected.

K.Morita et al PRC94(2016)031901

- Physical point result will be open soon by Iritani et al(HALQCD).

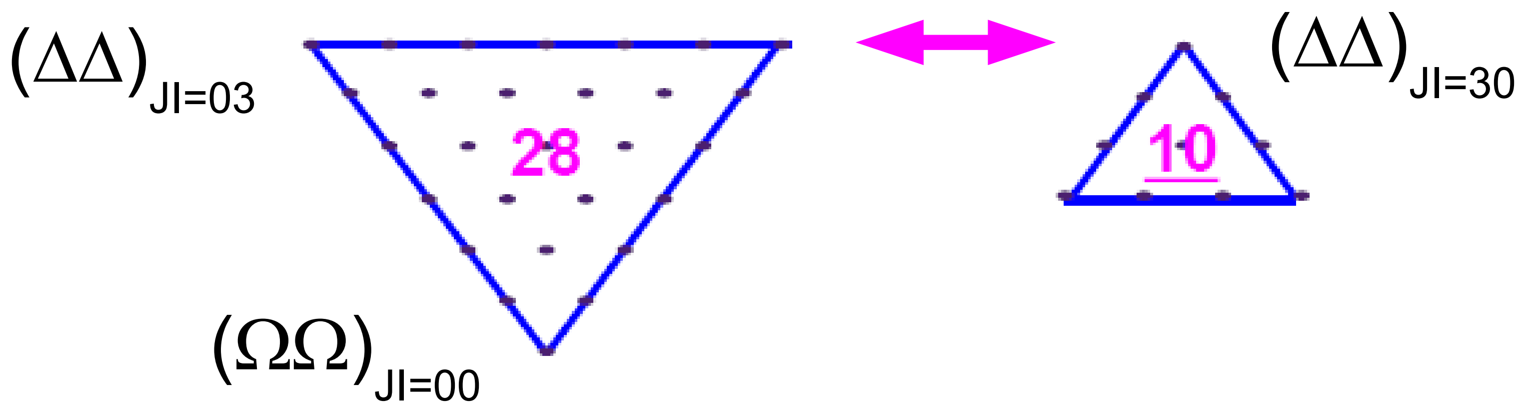
ΔΔ and ΩΩ interaction

Decuplet-Decuplet interaction

- **Flavor symmetry aspect**

Decuplet-Decuplet interaction can be classified as

$$10 \otimes 10 = 28 \oplus 27 \oplus 35 \oplus \overline{10}$$

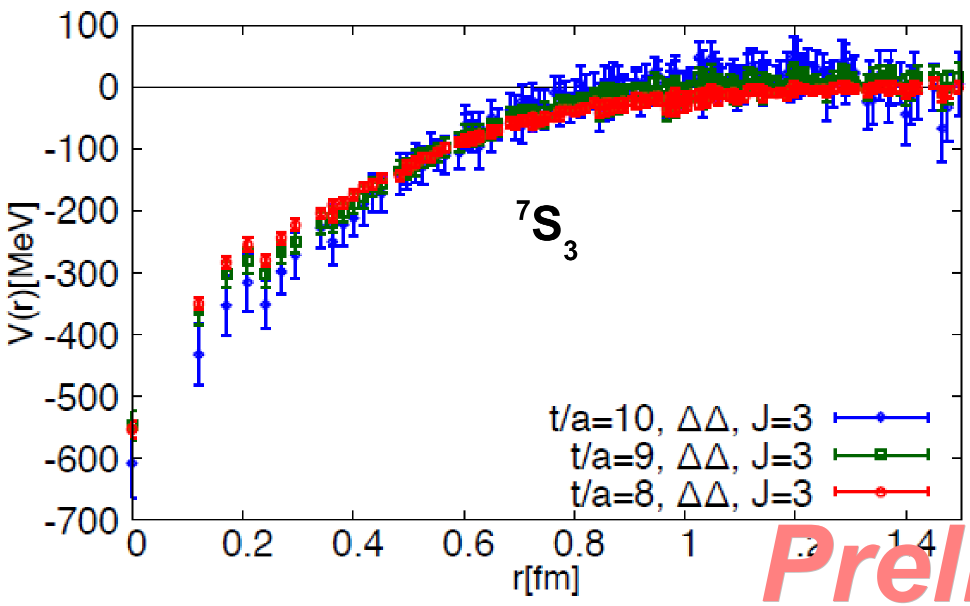


	28plet (0^+)	10*plet (3^+)
Pauli	allowed	allowed
CMI	repulsive	attractive

- **Δ–Δ(J=3) : Bound (resonance) state was found in experiment.**
- **Δ–Δ(J=0) [and Ω–Ω(J=0)] : Mirror of Δ–Δ(J=3) state**

Decuplet-Decuplet interaction in SU(3) limit

▶ N_f = 3 full QCD with L = 1.93fm, m_π = 1015 MeV m_Δ = 2225 MeV



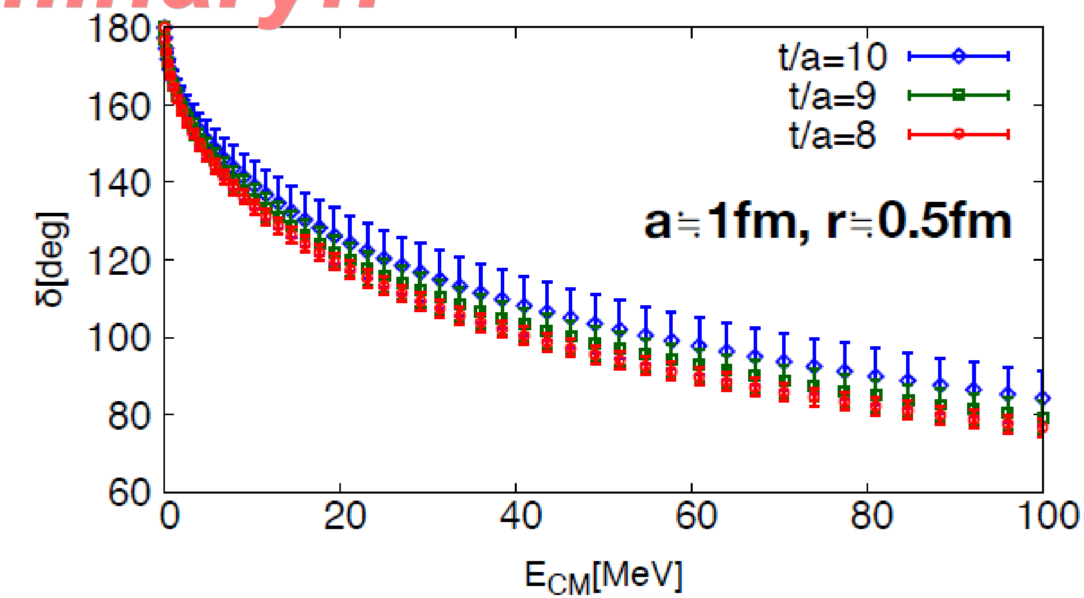
Δ–Δ(J=3) **10*plet**

- **Interaction in 10*plet [$J^P(I)=3^+(0)$] is strongly attractive.**
- **There is no repulsive core at short distances**

Preliminary!!

Bound ΔΔ state is found.

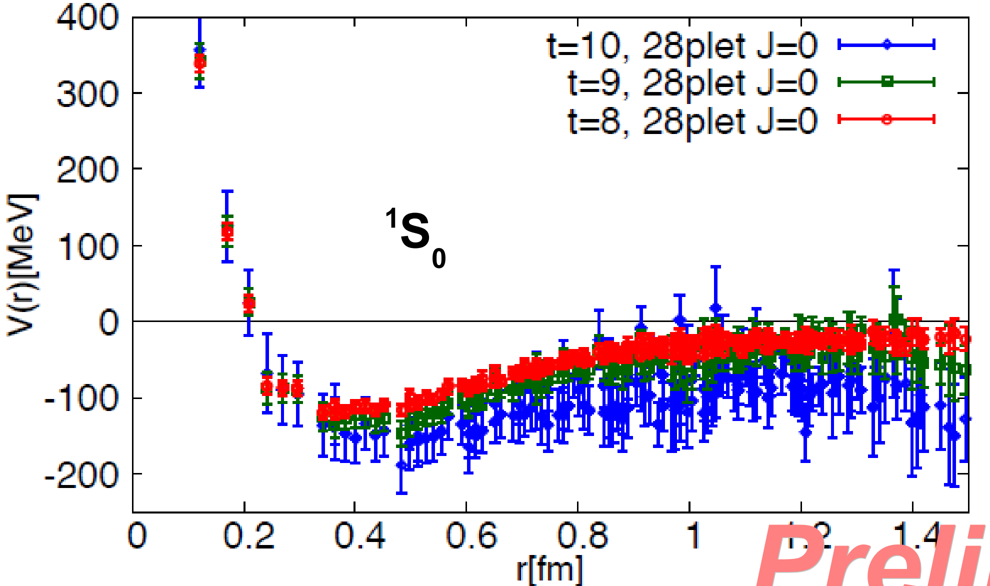
- **Decay to NN($^3D_3$) is neglected.**
- **Δ baryon can not decay into N+π in this lattice setup**



S. Gongyo, K. Sasaki (HAL QCD)

Decuplet-Decuplet interaction in SU(3) limit

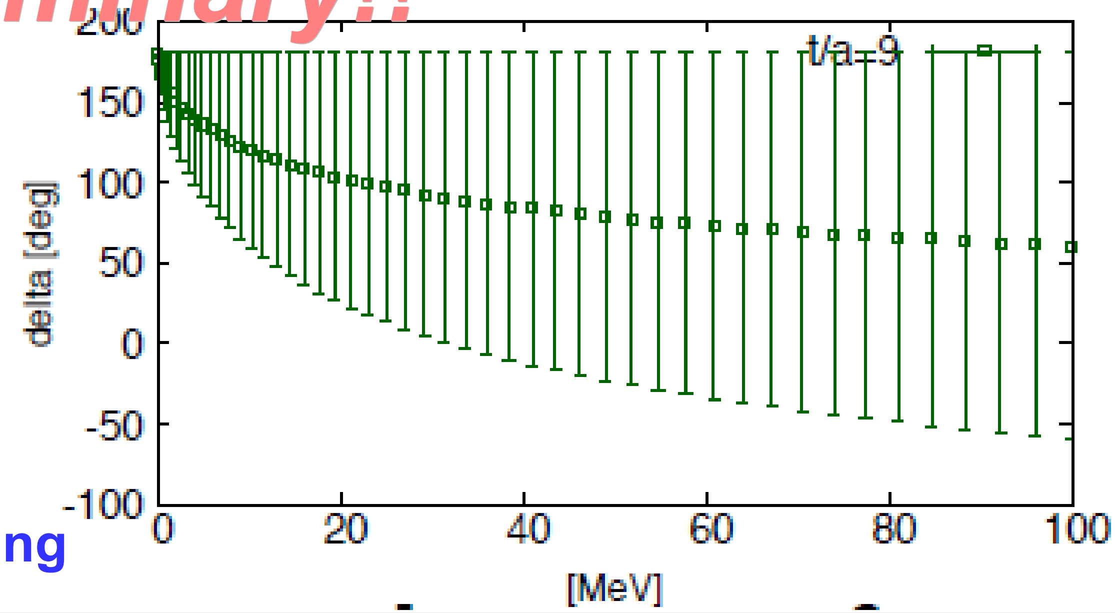
▶ N_f = 3 full QCD with L = 1.93fm, m_π = 1015 MeV m_Δ = 2225 MeV



Δ−Δ(J=0) **28plet**

- **Repulsive core is surrounded by attractive pocket in 28plet [ΔΔ $J^p(I)=0^+(3)$] and [ΩΩ $J^p(I)=0^+(0)$].**

Preliminary!!



Phase shift shows that the system is in the unitary limit.

- **Δ baryon can not decay into N+π in this lattice setup**

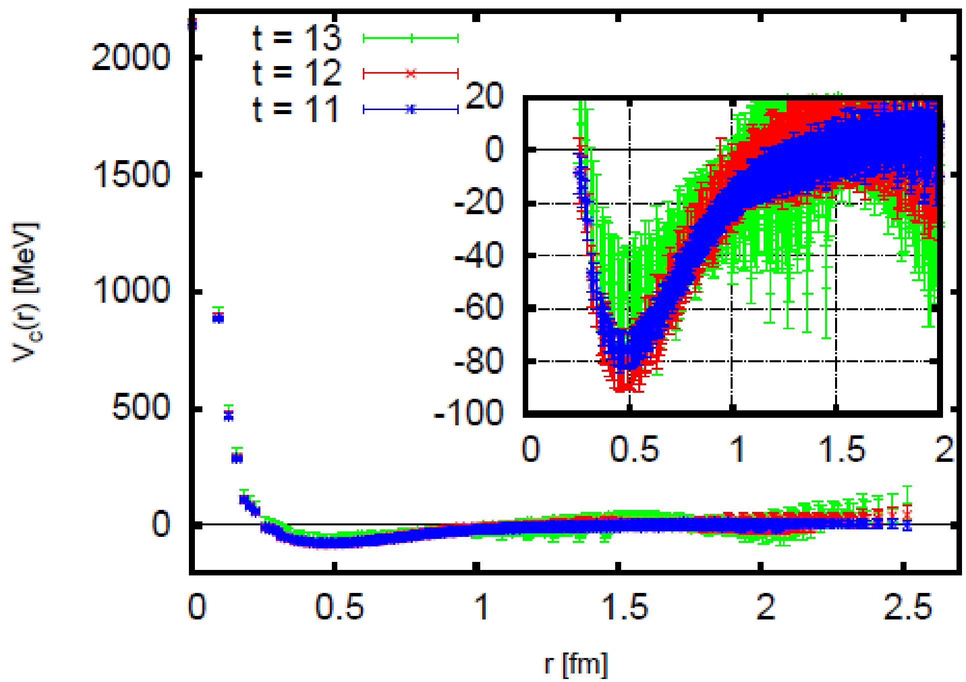
Go to the lighter quark mass region with consideration of SU(3) breaking

ΩΩ $J^p = 0^+$ state in unphysical region

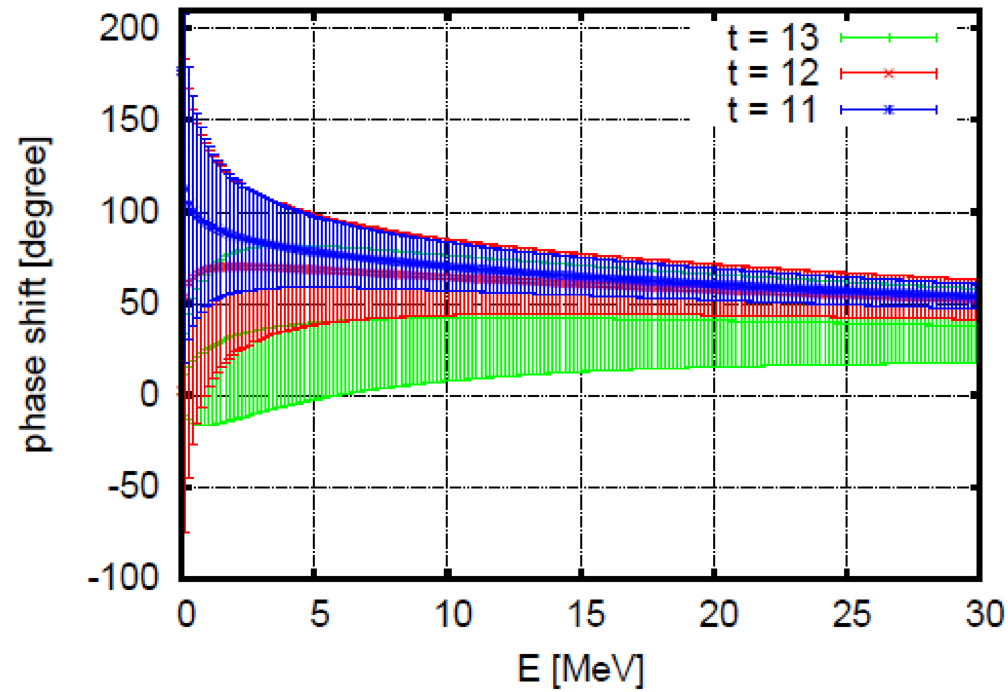
- N_f = 2+1 full QCD with L = 3fm, *m_π = 700 MeV* *m_Ω = 1966 MeV*

The ΩΩ state is stable against the strong interaction.

Potential



Phase shift

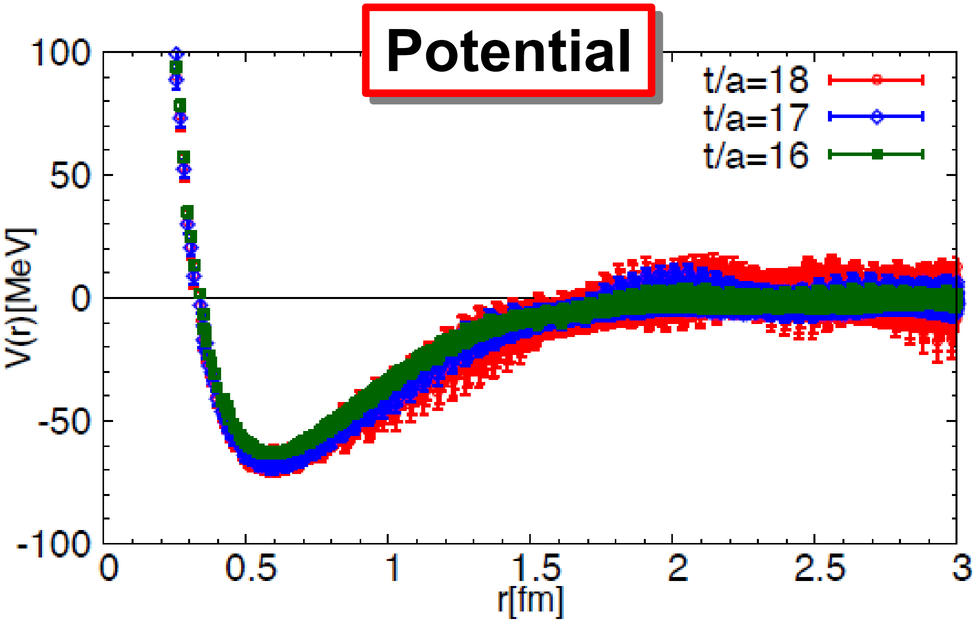


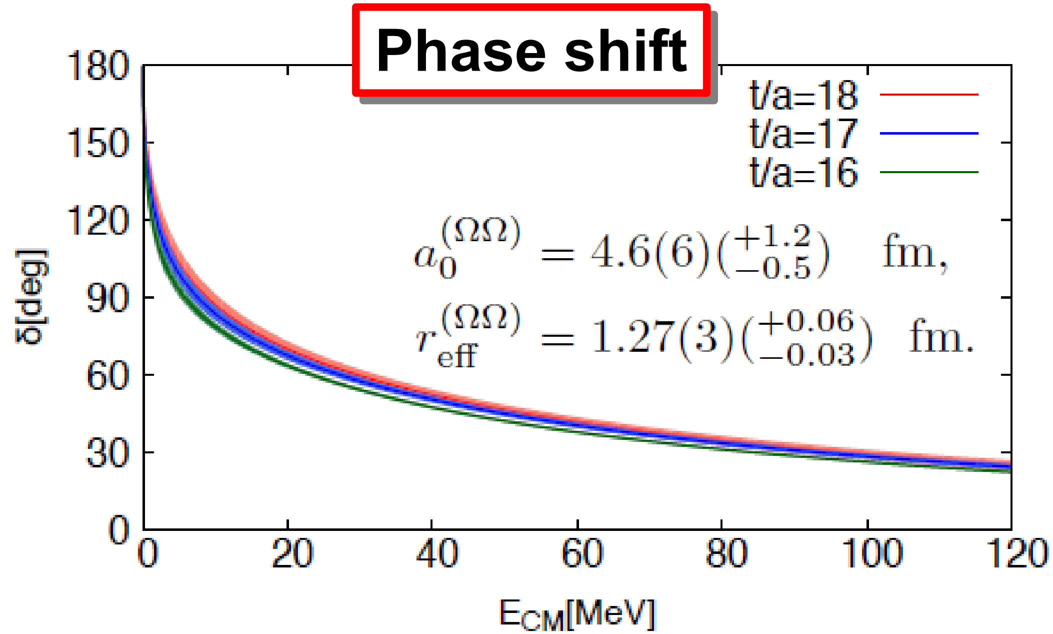
- Short range repulsion and attractive pocket are found.
- Potential is nearly independent on "t" within statistical error.
- The system may appear close to the unitary limit.

S. Gongyo, K. Sasaki (HAL QCD)arXiv:1709.00654

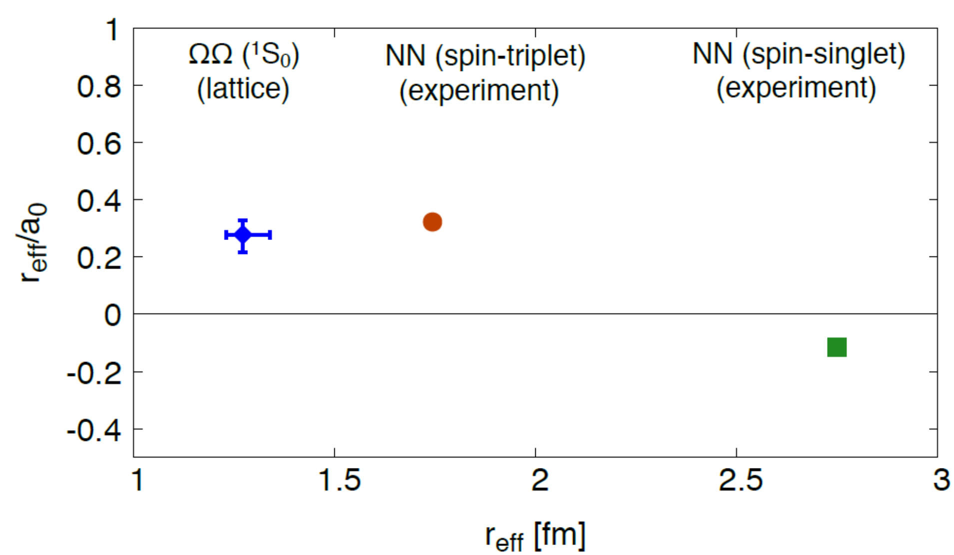
$\Omega\Omega$ $J^p(I) = 0^+(0)$ state near the physical point

- N_f = 2+1 full QCD with L = 8fm, m_π = 146 MeV m_Ω = 1712 MeV





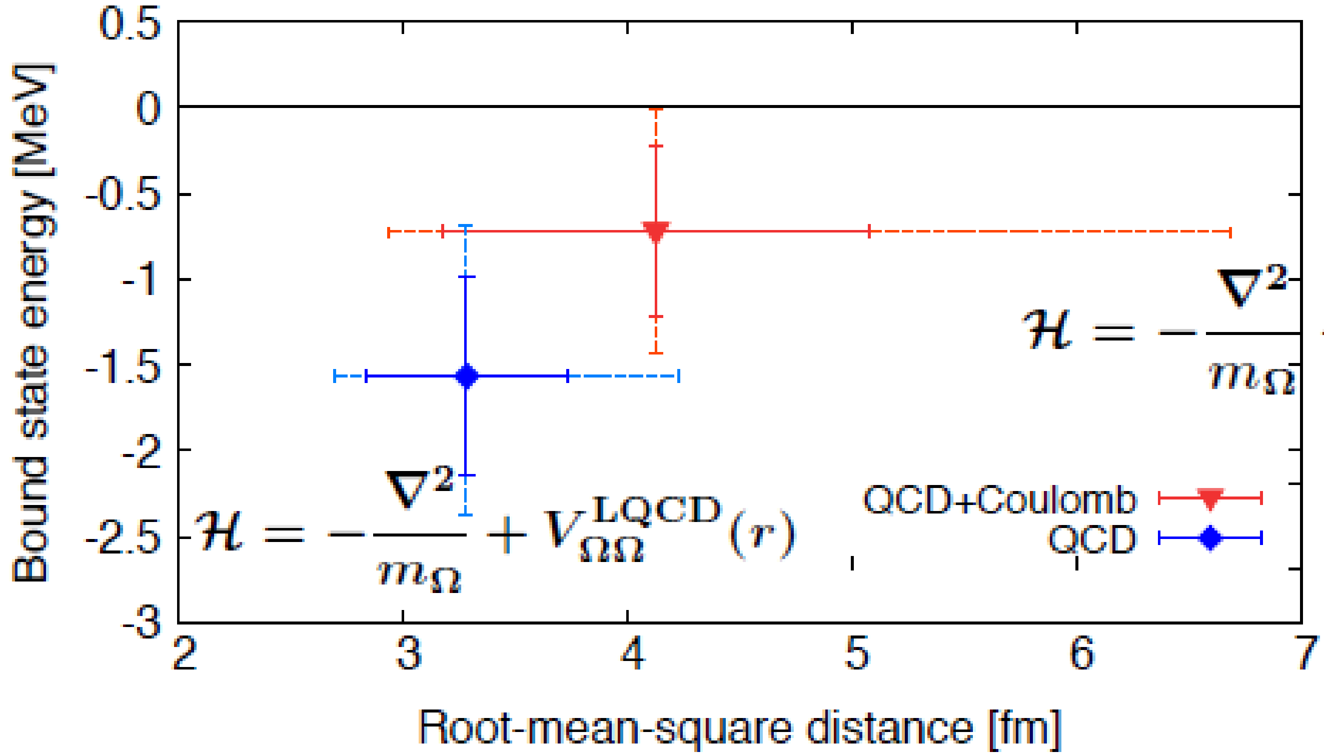
- Short range repulsion and attractive pocket is found.
- Calculated phase shift indicates a bound $\Omega\Omega$ state [Most strange dibaryon].
- Physical $\Omega\Omega$ state in $J^p(I) = 0^+(0)$ is very close to unitary region.



ΩΩ $J^p(I) = 0^+(0)$ state near the physical point

- **N_f = 2+1 full QCD with L = 8fm, *m_π = 146 MeV***

Binding energy and the Coulomb effect



$$\mathcal{H} = -\frac{\nabla^2}{m_\Omega} + V^{\rm LQCD}_{\Omega\Omega}(r) + \frac{\alpha}{r}$$

$$\mathcal{H} = -\frac{\nabla^2}{m_\Omega} + V^{\rm LQCD}_{\Omega\Omega}(r)$$

Most strange dibaryon appears (within 1σ)
even if Coulomb effect is taken into accound.

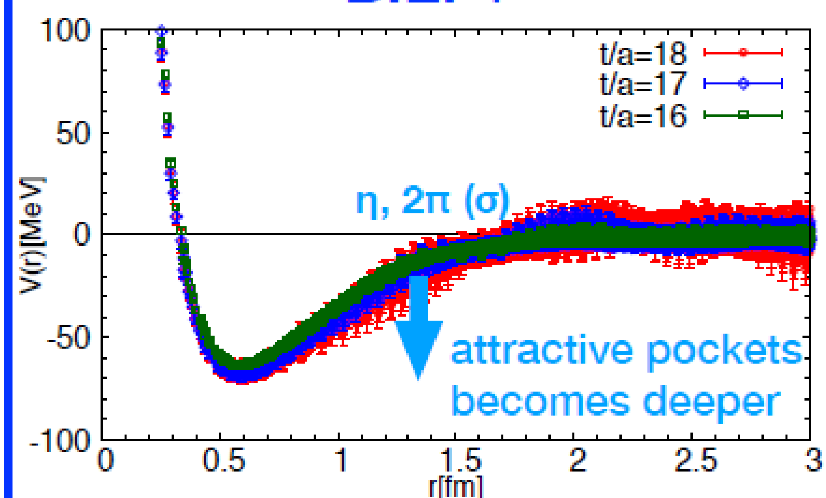
$$(B^{\rm (QCD)}_{\Omega\Omega}, B^{\rm (QCD+Coulomb)}_{\Omega\Omega}) = (1.6(6)\text{MeV}, 0.7(5)\text{MeV})$$

$\Omega\Omega$ $J^p(I) = 0^+(0)$ state at exact physical point

Conservative estimate at exact phys. pt.

m_π=146 MeV -> 135 MeV, m_Ω= 1712MeV -> 1672 MeV

B.E. ↑



V.S.

B.E. ↓

$$\mathcal{H} = -\frac{\nabla^2}{m_\Omega} + V_{\Omega\Omega}^{\mathrm{LQCD}}(r)$$

kinetic term becomes more dominant
-> B.E. is reduced

conservative estimate:
only change the mass of schroedinger eq.

$$(B_{\Omega\Omega}^{(\mathrm{QCD})}, B_{\Omega\Omega}^{(\mathrm{QCD+Coulomb})}) = (1.6(6)\mathrm{MeV}, 0.7(5)\mathrm{MeV})$$
$$\to (1.3(5)\mathrm{MeV}, 0.5(5)\mathrm{MeV})$$

These changes are well within errors

27

Summary

- We have investigated dibaryon candidate states from LQCD
  - H-dibaryon channel
    - We found a strong attraction in NΞ J=0 with I=0.
    - It is still difficult to conclude the fate of H-dibayon.
  - NΩ state with $J^p=2^+$
    - Interaction is strongly attractive and no short range repulsion.
    - It forms a bound state with about 20MeV B.E..
    - Physical point result will be open for ΩN channel.
  - ΔΔ and ΩΩ states
    - ΔΔ(I=0) have strongly attractive potential.
    - ΔΔ(I=3) and ΩΩ potential have repulsive core and attractive pocket.
    - Physical ΩΩ system in J=0 forms the most strange dibaryon (or unitary region...)

Backup

Lists of channels

I=0 states

Spin	BB channels			SU(3) representation		
$^1S_0$	$\Lambda\Lambda$	$N\Xi$	$\Sigma\Sigma$	1	8s	27
$^3S_1$	--	$N\Xi$	--	8a	--	--

I=1 states

Spin	BB channels			SU(3) representation		
$^1S_0$	$N\Xi$	--	$\Lambda\Sigma$	--	8s	27
$^3S_1$	$N\Xi$	$\Sigma\Sigma$	$\Lambda\Sigma$	8a	10	10*

I=2 states

Spin	BB channels	SU(3) representation		
$^1S_0$	$\Sigma\Sigma$	--	--	27
~~$^3S_1$~~				

Strong attraction (H-dibaryon)

Attraction

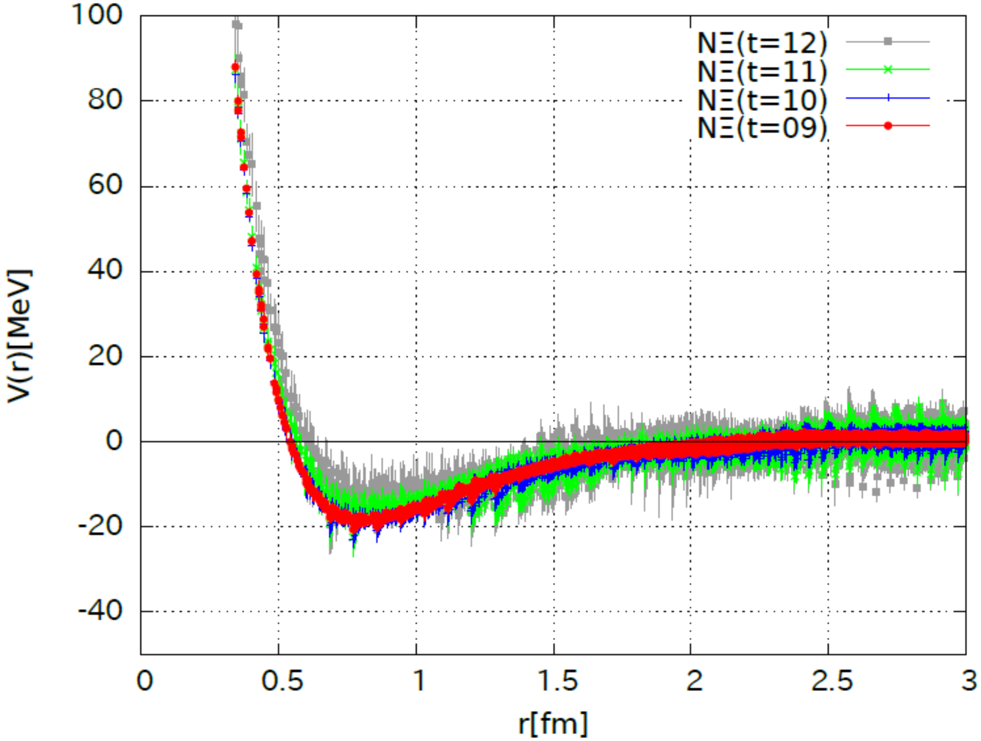
Strong repulsion

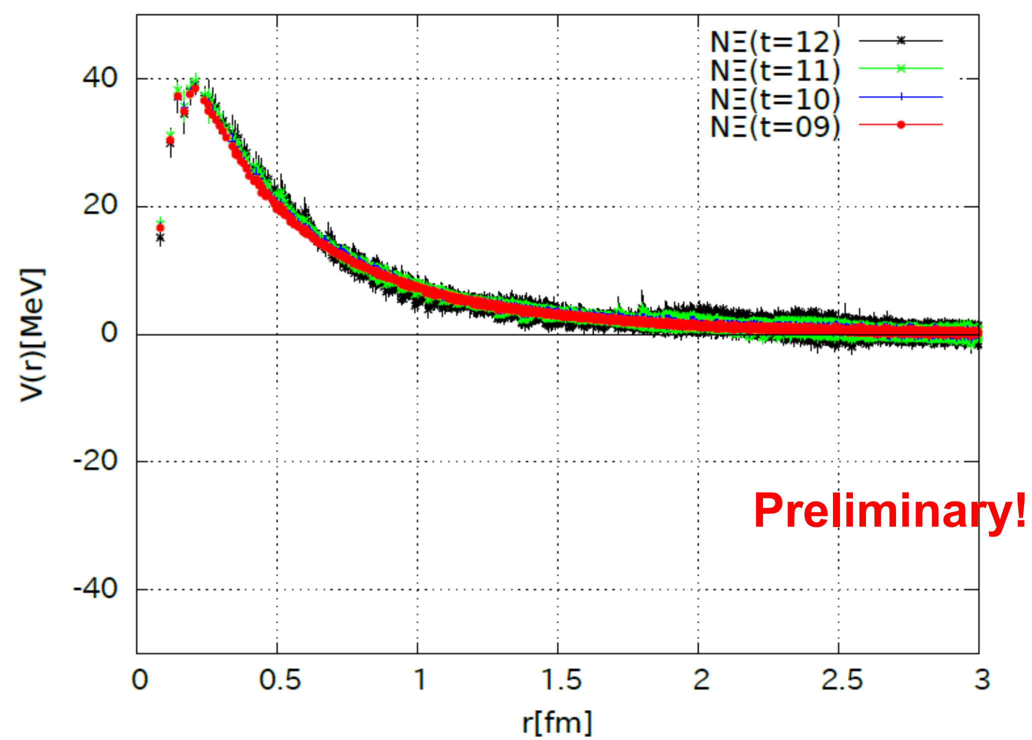
Repulsion

Similar to The NN potential

$N\Xi$ (I=0) $^3S_1$-$^3D_1$ channel

- N_f = 2+1 full QCD m_π=146MeV with L = 8.1fm





- NΞ (I=0,spin triplet) potential belongs to the 8plet in flavor SU(3) limit.
- The potential has an attractive pocket and repulsive core.
- Tensor potential is weaker than the phenomenological NN tensor potential.

$N\Xi$, $\Lambda\Sigma$, $\Sigma\Sigma$ (I=1) $^3S\text{-}D_1$ channel (central pot)

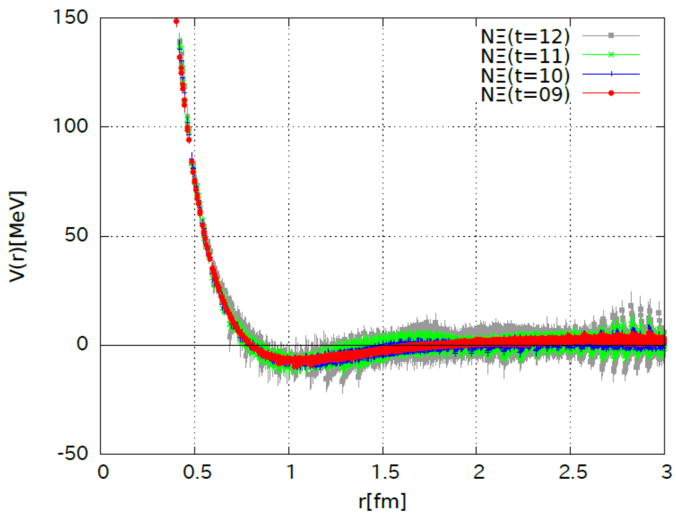
t=09
t=10
t=11
t=12

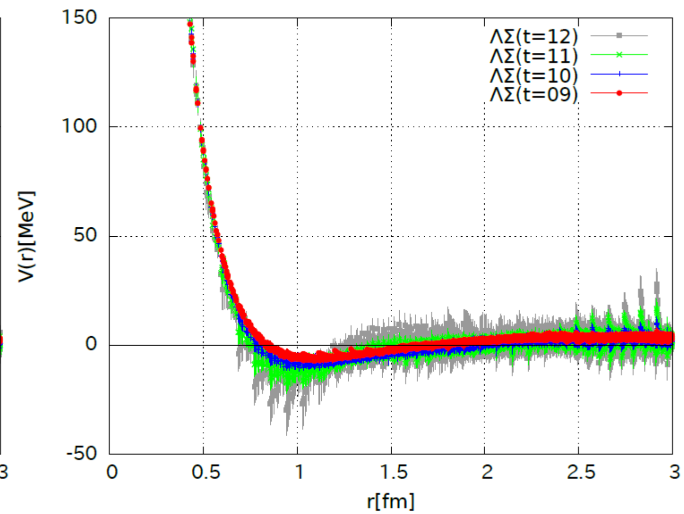
- **N_f = 2+1 full QCD m_π=146MeV with L = 8.12fm**

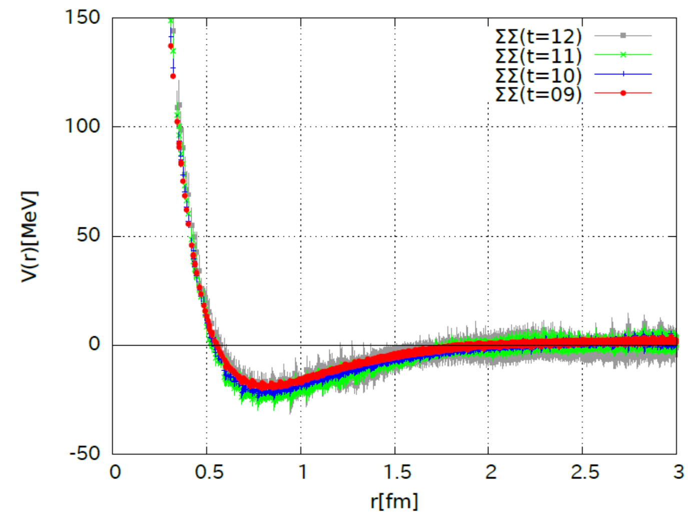
$$\begin{pmatrix} N\Xi \\ \Sigma\Lambda \\ \Sigma\Sigma \end{pmatrix} = \frac{1}{\sqrt{6}} \begin{pmatrix} -\sqrt{2} & -\sqrt{2} & \sqrt{2} \\ \sqrt{3} & -\sqrt{3} & 0 \\ 1 & 1 & \sqrt{4} \end{pmatrix} \begin{pmatrix} 8 \\ \overline{10} \\ 10 \end{pmatrix}$$

48src

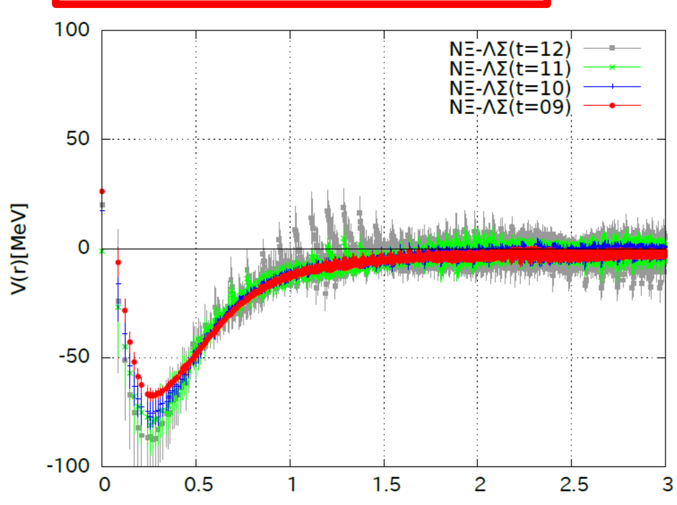
Diagonal elements

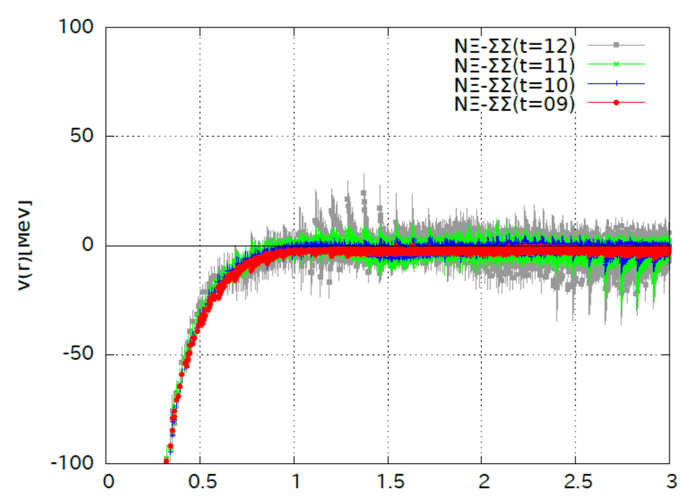


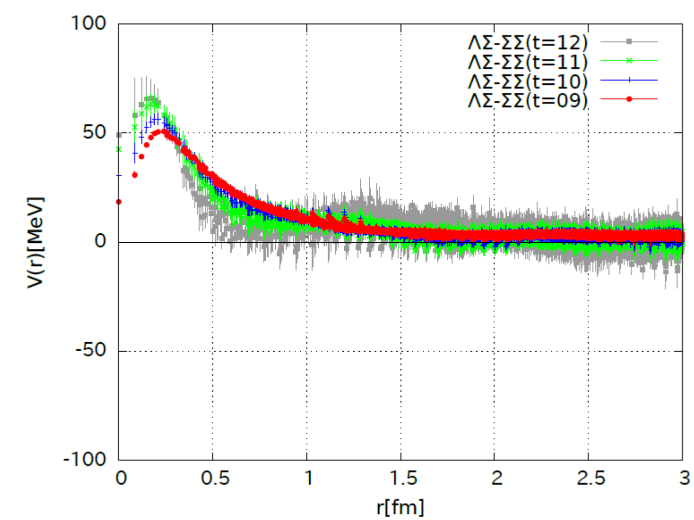




Off-diagonal elements







$N\Xi$, $\Lambda\Sigma$, $\Sigma\Sigma$ (I=1) $^3S\text{-}D_1$ channel (tensor pot)

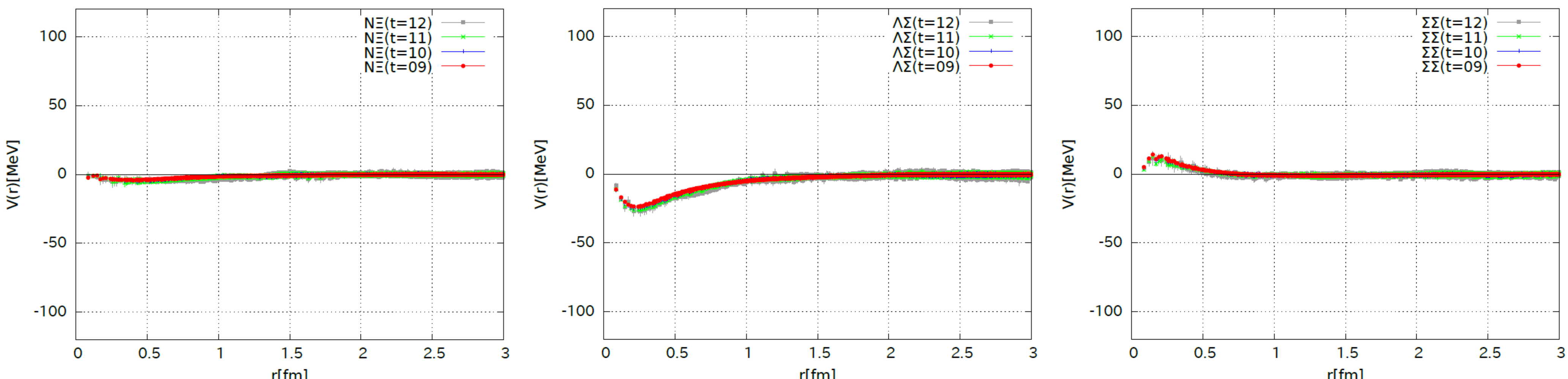
t=09
t=10
t=11
t=12

- N_f = 2+1 full QCD m_π=146MeV with L = 8.12fm

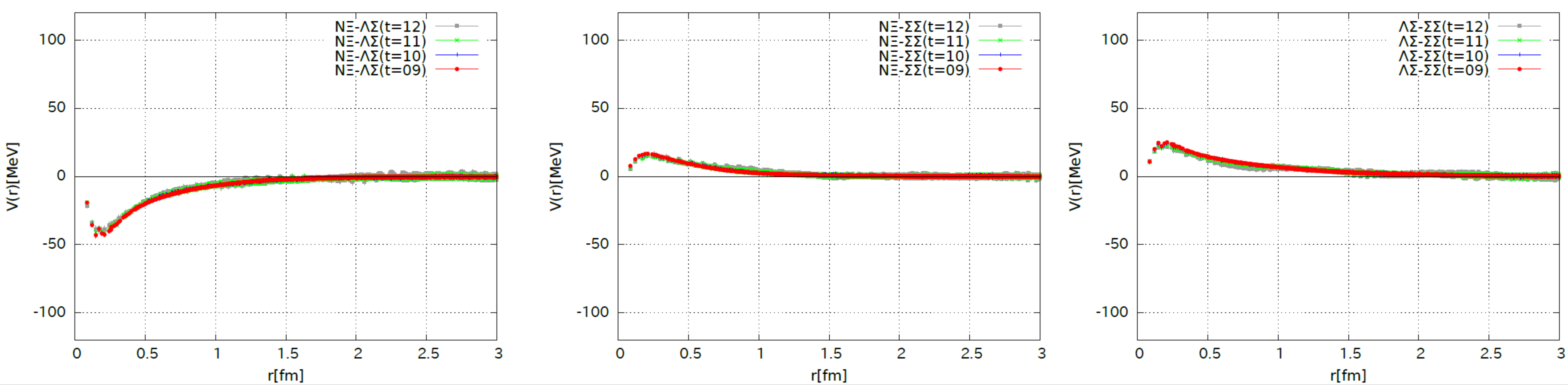
$$\begin{pmatrix} N\Xi \\ \Sigma\Lambda \\ \Sigma\Sigma \end{pmatrix} = \frac{1}{\sqrt{6}} \begin{pmatrix} -\sqrt{2} & -\sqrt{2} & \sqrt{2} \\ \sqrt{3} & -\sqrt{3} & 0 \\ 1 & 1 & \sqrt{4} \end{pmatrix} \begin{pmatrix} 8 \\ 10 \\ 10 \end{pmatrix}$$

48src

Diagonal elements



Off-diagonal elements



Baryon-baryon system with S=-2

Spin singlet states

Isospin	BB channels		
I=0	$\Lambda\Lambda$	$N\Xi$	$\Sigma\Sigma$
I=1	$N\Xi$	$\Lambda\Sigma$	——
I=2	$\Sigma\Sigma$	——	——

Spin triplet states

Isospin	BB channels		
I=0	$N\Xi$	——	——
I=1	$N\Xi$	$\Lambda\Sigma$	$\Sigma\Sigma$

Relations between BB channels and SU(3) irreducible representations

$8 \times 8 = 27 + 8_s + 1 + \underline{10} + 10 + 8_A$

$J^P=0^+$, I=0

$$\begin{pmatrix} \Lambda\Lambda \\ N\Xi \\ \Sigma\Sigma \end{pmatrix} = \frac{1}{\sqrt{40}} \begin{pmatrix} -\sqrt{5} & -\sqrt{8} & \sqrt{27} \\ \sqrt{20} & \sqrt{8} & \sqrt{12} \\ \sqrt{15} & -\sqrt{24} & -1 \end{pmatrix} \begin{pmatrix} 1 \\ 8 \\ 27 \end{pmatrix}$$

$J^P=1^+$, I=0

$$N\Xi \Leftrightarrow 8$$

$J^P=0^+$, I=1

$$\begin{pmatrix} N\Xi \\ \Sigma\Lambda \end{pmatrix} = \frac{1}{5} \begin{pmatrix} \sqrt{2} & -\sqrt{3} \\ \sqrt{3} & \sqrt{2} \end{pmatrix} \begin{pmatrix} 27 \\ 8 \end{pmatrix}$$

$J^P=0^+$, I=2

$$\Sigma\Sigma \Leftrightarrow 8$$

$J^P=1^+$, I=1

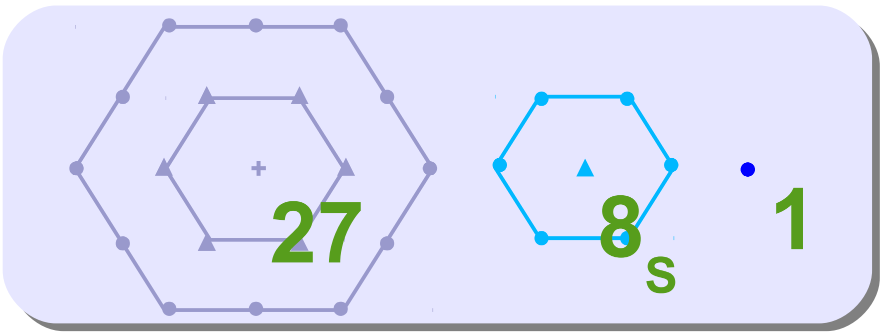
$$\begin{pmatrix} N\Xi \\ \Sigma\Lambda \\ \Sigma\Sigma \end{pmatrix} = \frac{1}{\sqrt{6}} \begin{pmatrix} -\sqrt{2} & -\sqrt{2} & \sqrt{2} \\ \sqrt{3} & -\sqrt{3} & 0 \\ 1 & 1 & \sqrt{4} \end{pmatrix} \begin{pmatrix} 8 \\ \overline{10} \\ 10 \end{pmatrix}$$

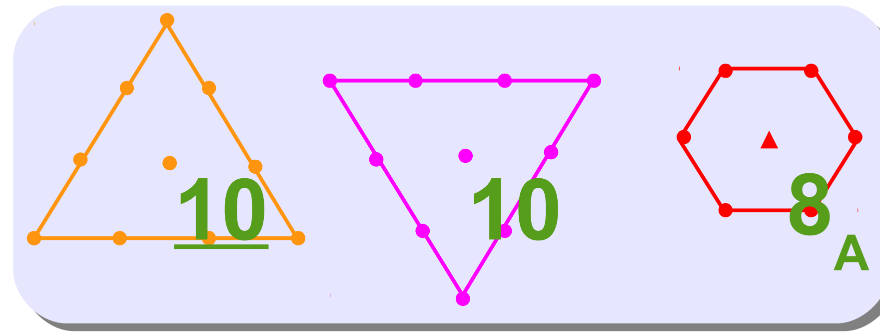
Features of flavor singlet interaction is integrated into the S=-2 $J^P=0^+$, I=0 system.

SU(3) feature of BB interaction

SU(3) classification

$8 \times 8 = 27 + 8_S + 1 + \underline{10} + 10 + 8_A$





In view of quark degrees of freedom

- Short range repulsion in BB interaction could be a result of **Pauli principle** and **color-magnetic interaction** for the quarks.

- For the s-wave BB system, **no repulsive core** is predicted in **flavor singlet state** which is known as **H-dibaryon** channel.

	Flavor symmetric states			Flavor anti-symmetric states		
	27	8	1	$\underline{10}$	10	8
Pauli	mixed	forbidden	allowed	mixed	forbidden	mixed
CMI	repulsive	repulsive	attractive	repulsive	repulsive	repulsive

Oka et al NPA464 (1987)

Works on H-dibaryon state

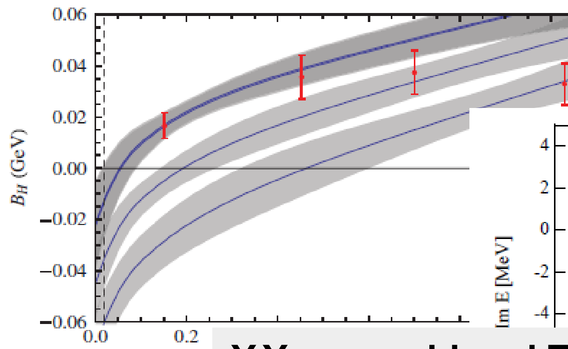
Theoretical status

Several sort of calculations and results
(bag models, NRQM, Quenched LQCD….)

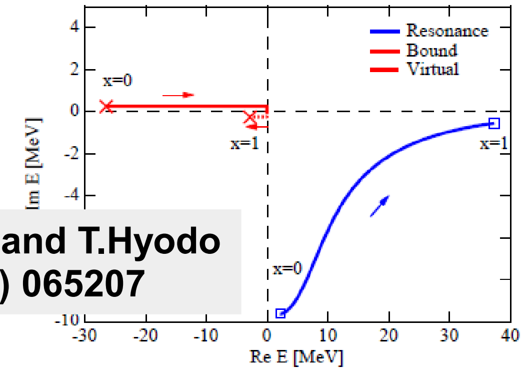
There were no conclusive result.

Chiral extrapolations of recent LQCD data

Unbound or resonance



P. E. Shanahan et al
PRL 107(2011) 092004



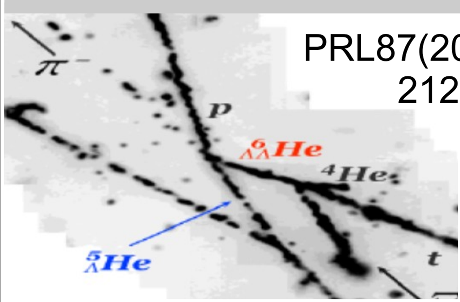
Y.Yamaguchi and T.Hyodo
PRC 94 (2016) 065207

Experimental status

"NAGARA Event"

K.Nakazawa et al
KEK-E176 & E373 Coll.

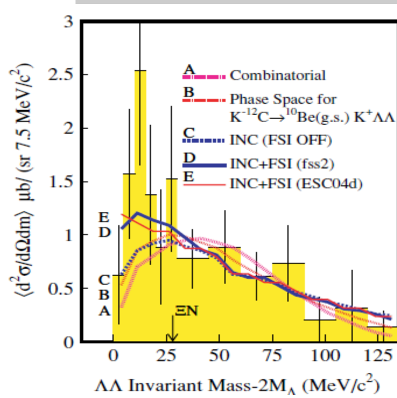
PRL87(2001) 212502



- Deeply bound dibaryon state is ruled out

"$^{12}C(K^-,K^+\Lambda\Lambda)$ reaction"

C.J.Yoon et al KEK-PS E522 Coll.



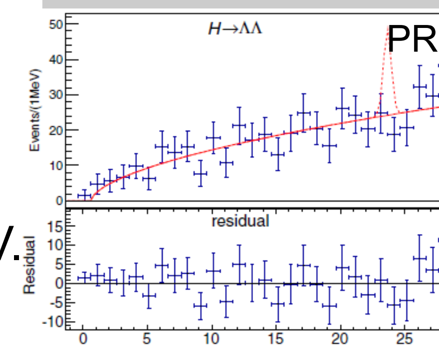
PRC75(2007) 022201(R)

- Significance of enhancements below 30 MeV.

Larger statistics
J-PARC E42

"Y(1S) and Y(2S) decays"

B.H. Kim et al Belle Coll.



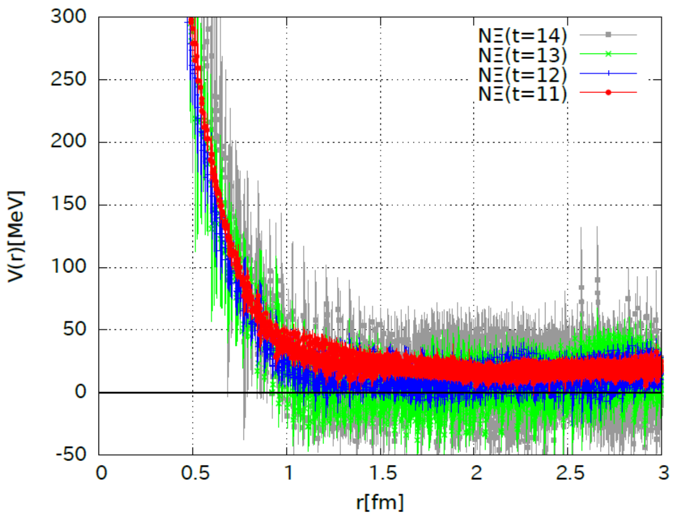
PRL110(2013) 222002

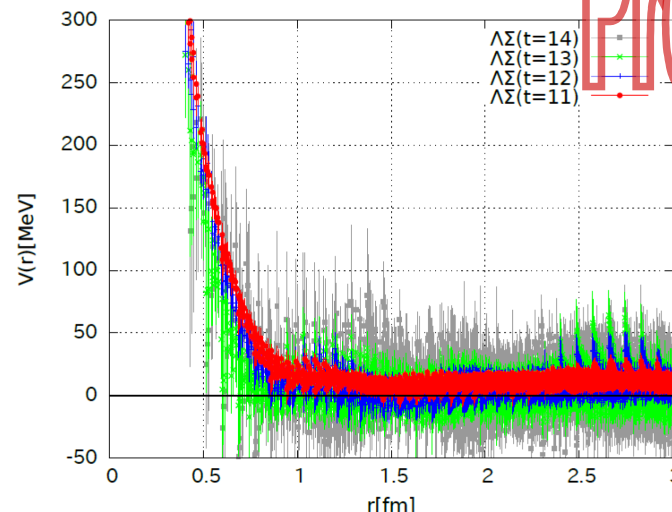
- There is no sign of near threshold enhancement.

$N\Xi$, $\Lambda\Sigma$ (I=1) $^1S_0$ channel

▶ N_f = 2+1 full QCD m_π=146MeV with L = 8.1fm

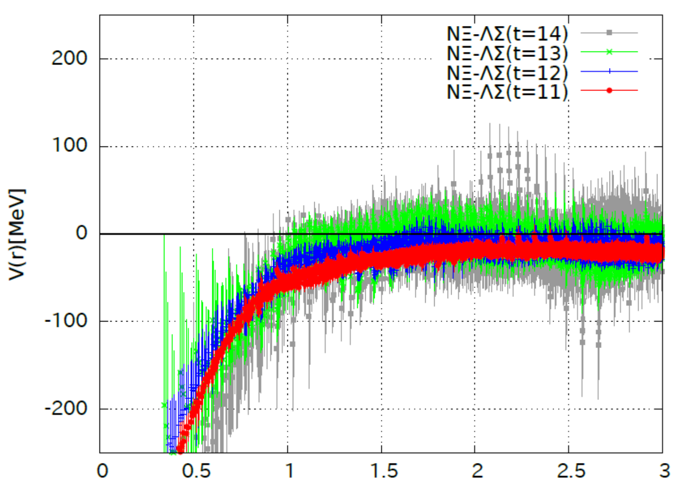
Diagonal elements





Preliminary!

Off-diagonal elements



$$\begin{pmatrix} N\Xi \\ \Sigma\Lambda \end{pmatrix} = \frac{1}{5}\begin{pmatrix} \sqrt{2} & -\sqrt{3} \\ \sqrt{3} & \sqrt{2} \end{pmatrix}\begin{pmatrix} 27 \\ 8 \end{pmatrix}$$

- Diagonal elements are repulsive in whole range.
- Diagonal $N\Xi$ potential is strongly repulsive.
- Potentials are not saturated in this time range.

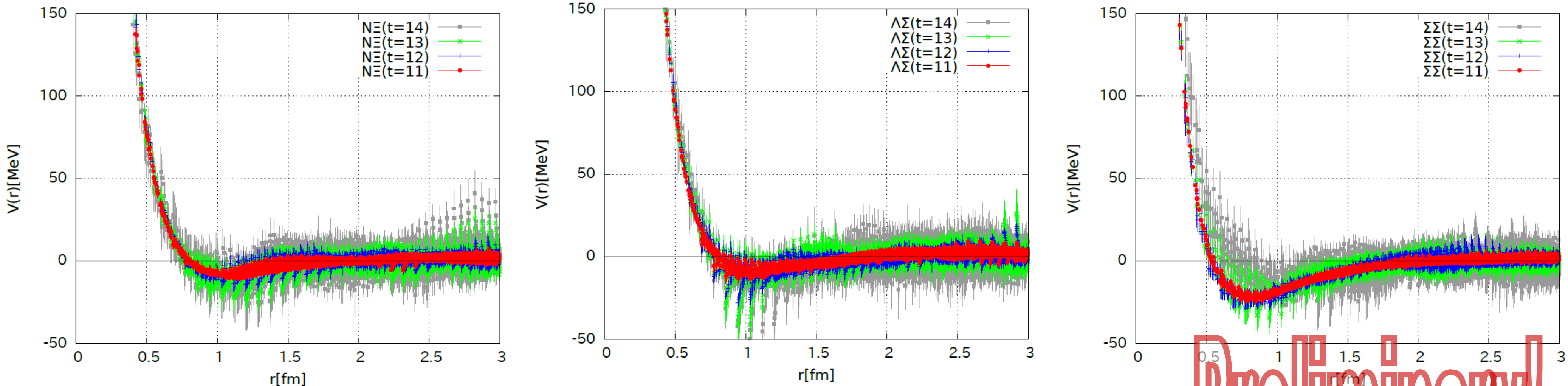
$N\Xi$, $\Lambda\Sigma$, $\Sigma\Sigma$ (I=1) $^3S_1$ channel (eff. pot.)

t=11
t=12
t=13
t=14

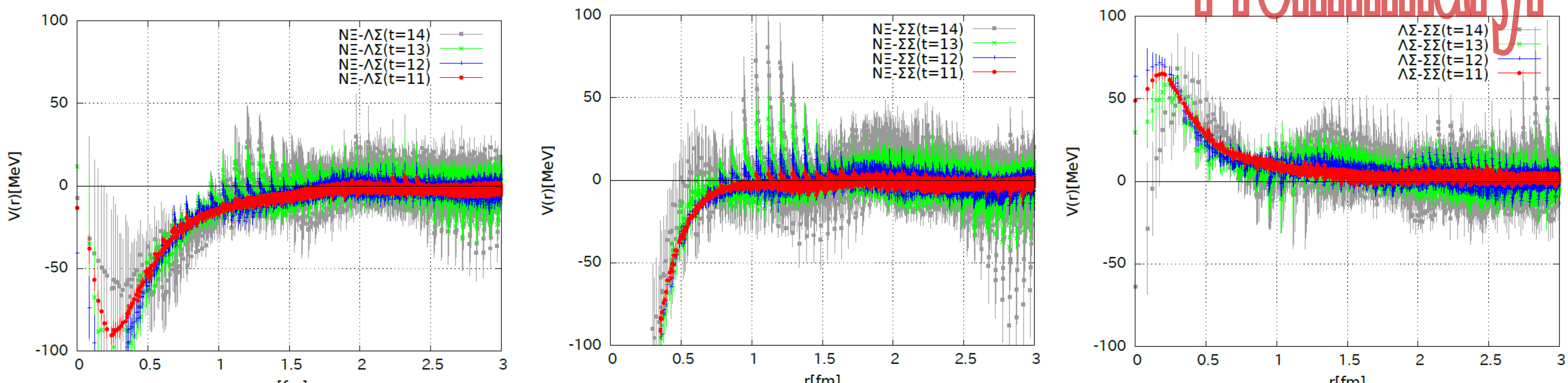
- N_f = 2+1 full QCD m_π=146MeV with L = 8.1fm

$$\begin{pmatrix} N\Xi \\ \Sigma\Lambda \\ \Sigma\Sigma \end{pmatrix} = \frac{1}{\sqrt{6}} \begin{pmatrix} -\sqrt{2} & -\sqrt{2} & \sqrt{2} \\ \sqrt{3} & -\sqrt{3} & 0 \\ 1 & 1 & \sqrt{4} \end{pmatrix} \begin{pmatrix} 8 \\ 10 \\ 10 \end{pmatrix}$$

Diagonal elements



Off-diagonal elements



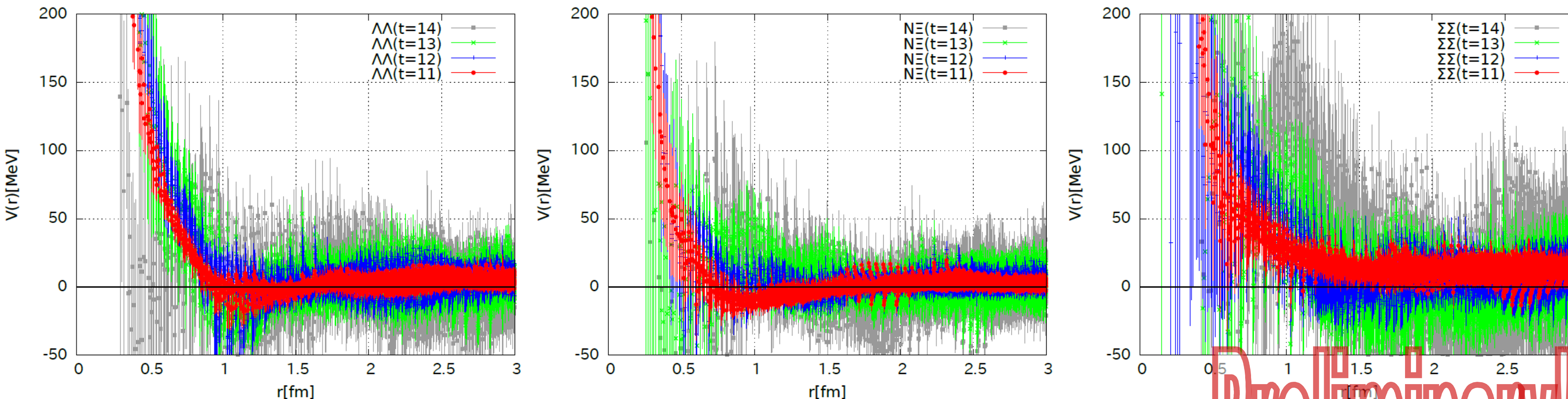
$\Lambda\Lambda$, $N\Xi$, $\Sigma\Sigma$ ($I=0$) $^1S_0$ channel

t=09
t=10
t=11
t=12

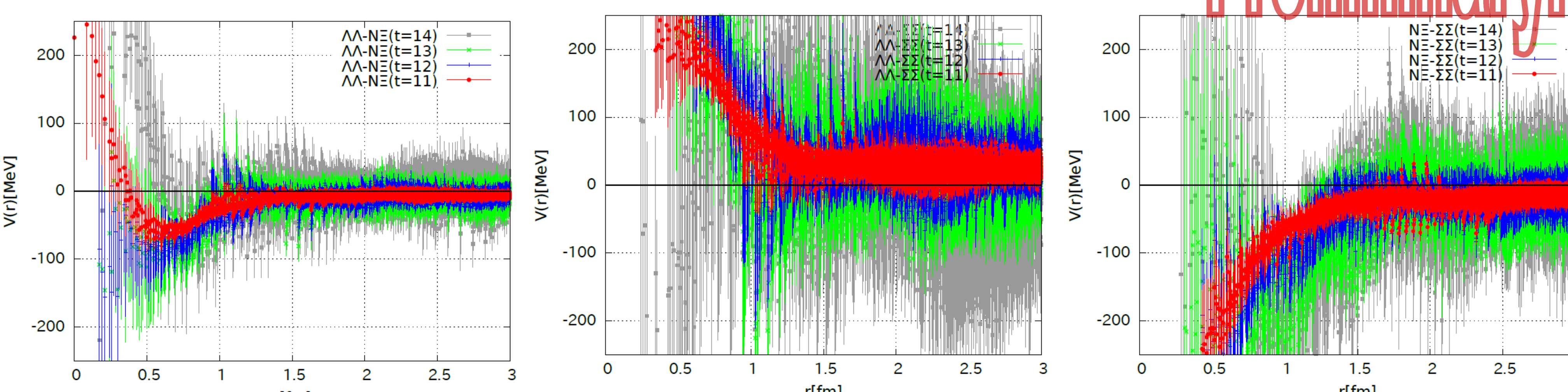
- N_f = 2+1 full QCD m_π=146MeV with L = 8.1fm

$$\begin{pmatrix}\Lambda\Lambda \\ N\Xi \\ \Sigma\Sigma\end{pmatrix} = \frac{1}{\sqrt{40}}\begin{pmatrix}-\sqrt{5} & -\sqrt{8} & \sqrt{27} \\ \sqrt{20} & \sqrt{8} & \sqrt{12} \\ \sqrt{15} & -\sqrt{24} & -1\end{pmatrix}\begin{pmatrix}1 \\ 8 \\ 27\end{pmatrix}$$

Diagonal elements



Off-diagonal elements



Comparison of potential matrices

Transformation of potentials
from the particle basis to the SU(3) irreducible representation (irrep) basis.

SU(3) Clebsh-Gordan coefficients

$$\begin{pmatrix} |1\rangle \\ |8\rangle \\ |27\rangle \end{pmatrix} = U \begin{pmatrix} |\Lambda\Lambda\rangle \\ |N\Xi\rangle \\ |\Sigma\Sigma\rangle \end{pmatrix}, \quad U \begin{pmatrix} V^{\Lambda\Lambda} & V^{\Lambda\Lambda}_{N\Xi} & V^{\Lambda\Lambda}_{\Sigma\Sigma} \\ V^{N\Xi}_{\Lambda\Lambda} & V^{N\Xi} & V^{N\Xi}_{\Sigma\Sigma} \\ V^{\Sigma\Sigma}_{\Lambda\Lambda} & V^{\Sigma\Sigma}_{N\Xi} & V^{\Sigma\Sigma} \end{pmatrix} U^t \rightarrow \begin{pmatrix} V_1 & & \\ & V_8 & \\ & & V_{27} \end{pmatrix}$$

In the SU(3) irreducible representation basis,
the potential matrix should be diagonal in the SU(3) symmetric configuration.

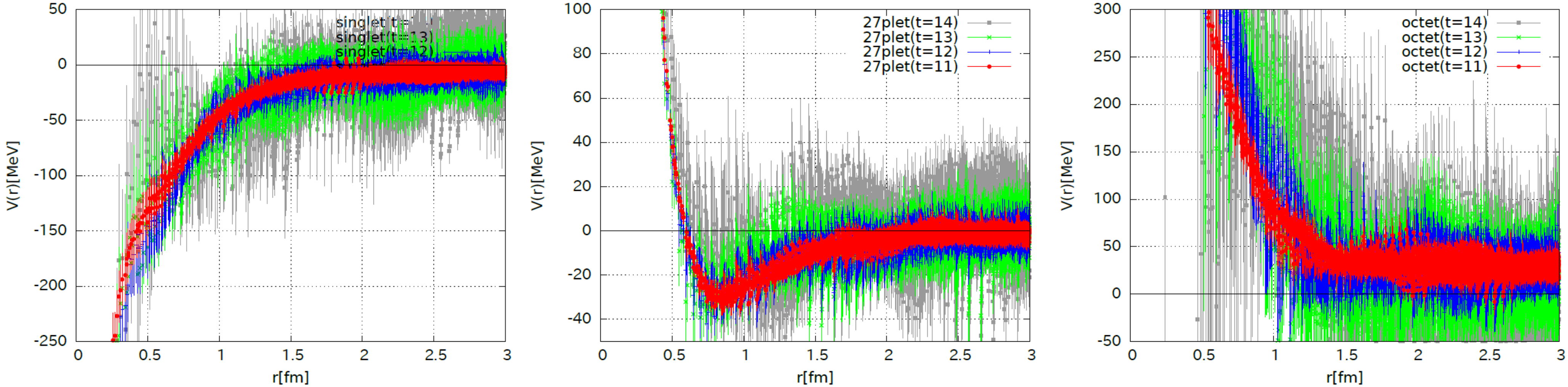
Off-diagonal part of the potential matrix in the SU(3) irrep basis
would be an effectual measure of the SU(3) breaking effect.

1, 8, 27plet (I=0) $^1S_0$ channel

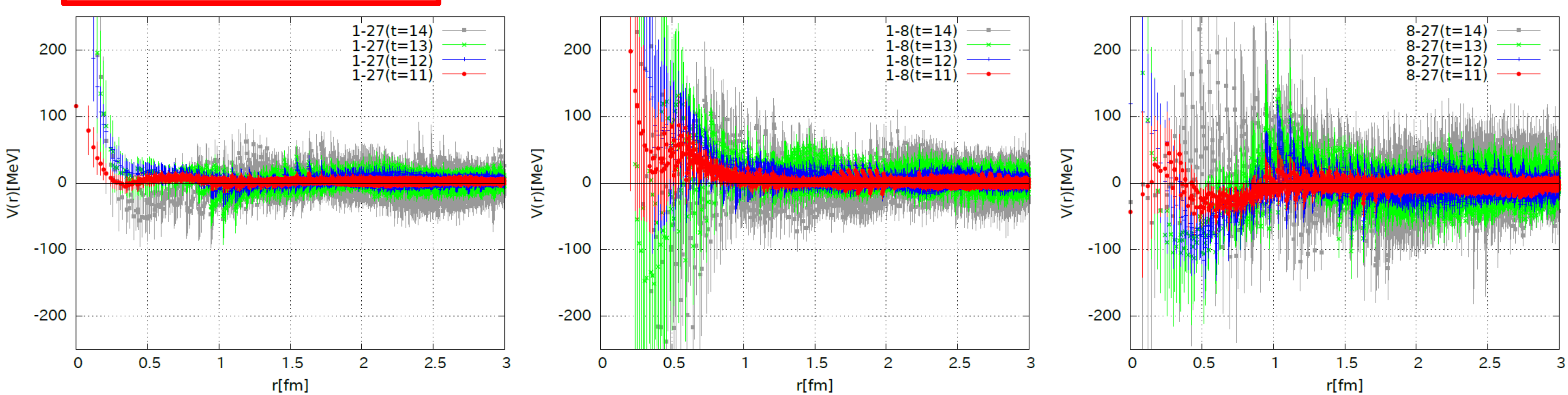
t=11
t=12
t=13
t=14

- N_f = 2+1 full QCD mπ=146MeV with L = 8.12fm

Diagonal elements



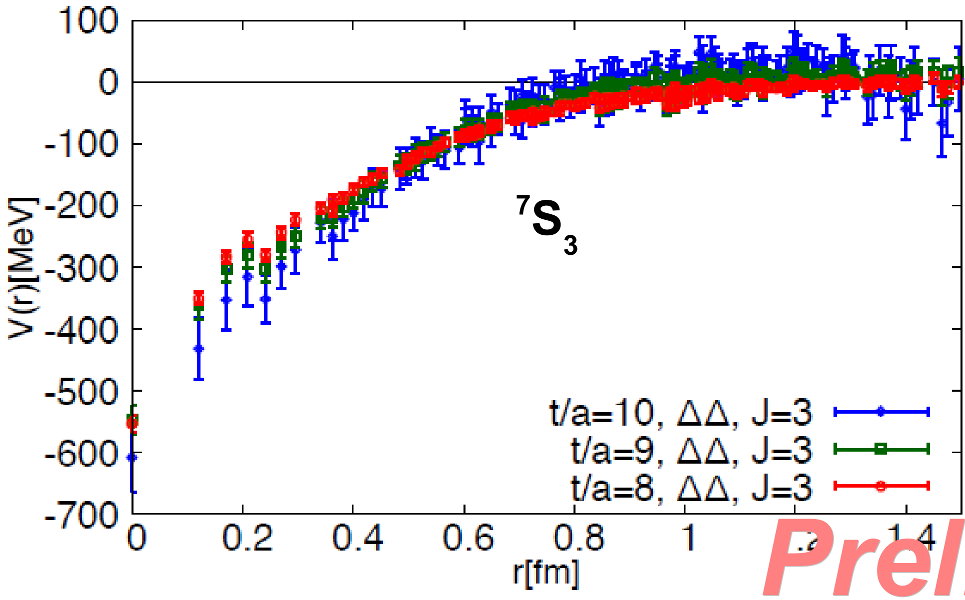
Off-diagonal elements



S. Gongyo, K. Sasaki (HAL QCD)

Decuplet-Decuplet interaction in SU(3) limit

▶ N_f = 3 full QCD with L = 1.93fm, m_π = 1015 MeV $m_\Delta \simeq 2225\text{MeV}$



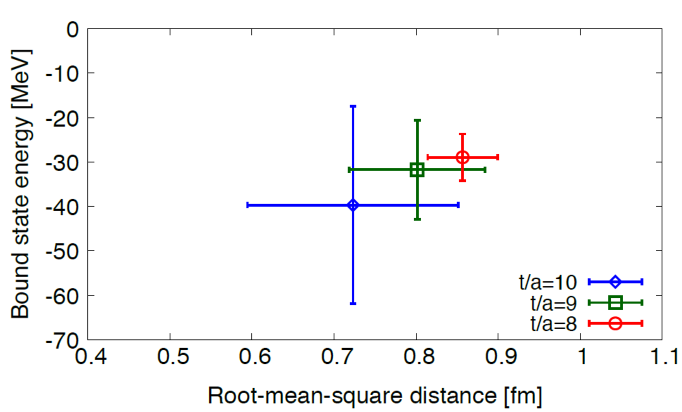
Δ–Δ(J=3) **10*plet**

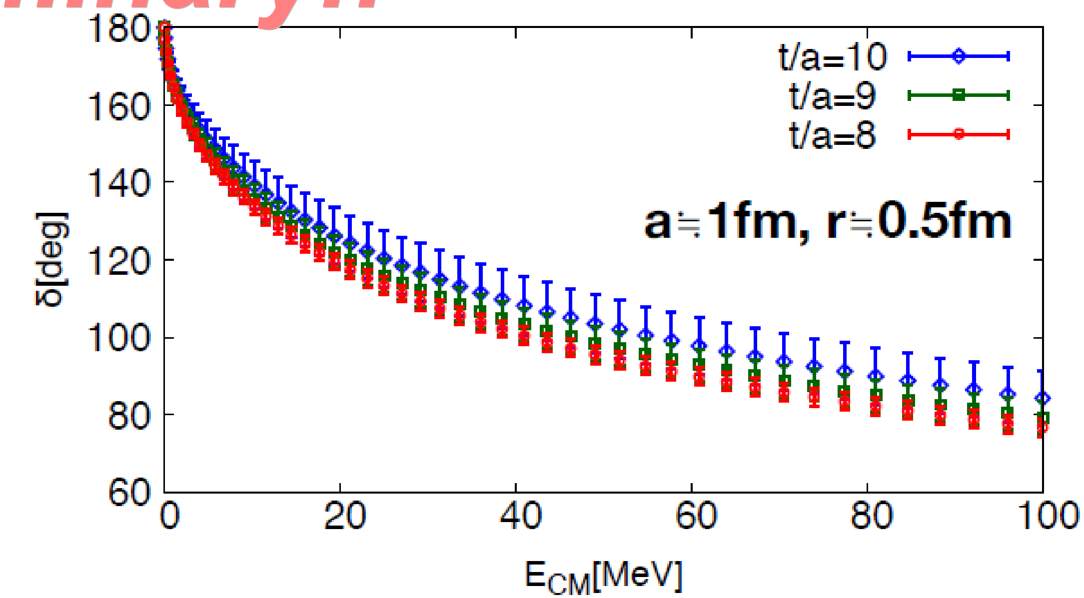
- **Interaction in 10*plet [$J^P(I)=3^+(0)$] is strongly attractive.**

- **Decay to NN($^3D_3$) is neglected.**
- **Δ baryon can not decay into N+π in this lattice setup**

Preliminary!!

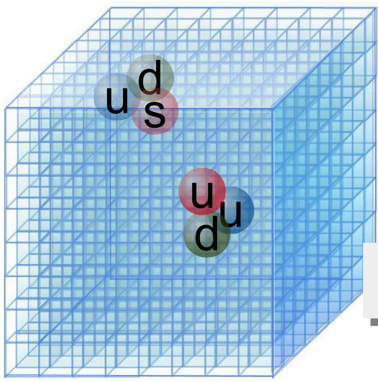
Bound ΔΔ state is found.



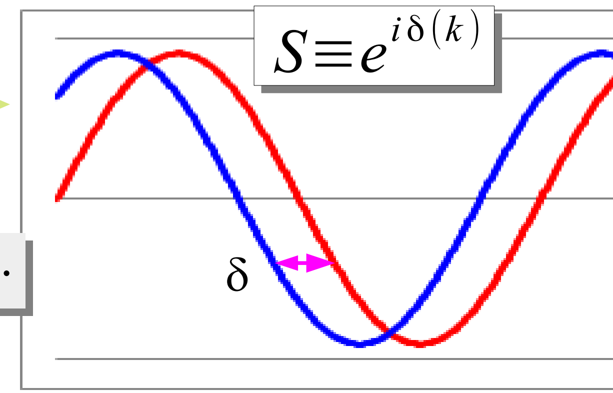


Hadron interaction from Lattice QCD

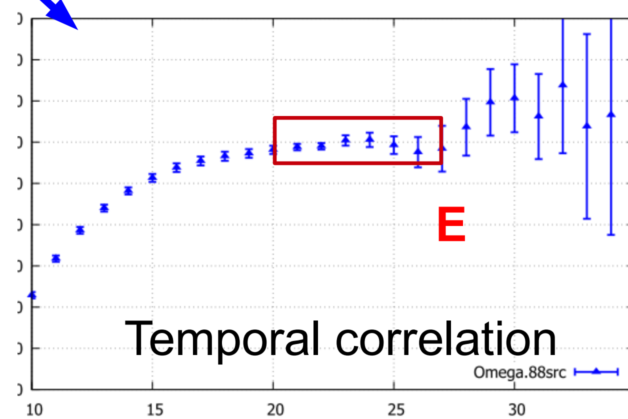
Lattice QCD simulation



Scattering S-matrix



$$\langle 0 | B_1 B_2(t, \vec{r}) \bar{I}(t_0) | 0 \rangle = A_0 \Psi(\vec{r}, E_0) e^{-E_0(t-t_0)} + \cdots$$



Lüscher's finite volume method

M. Lüscher, NPB354(1991)531

1. Measure the discrete energy spectrum, **E**
2. Put the **k** from E into the formula which connects **k** and **δ**

$$k_n \cot \delta(k_n) = \frac{4\pi}{L^3} \sum_{m \in \mathbb{Z}^3} \frac{1}{\vec{p}_m^2 - k_n^2}$$

Scattering phase shift

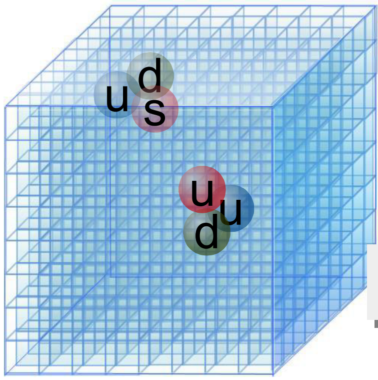
Eigen energy is extracted from plateau region

Be careful of mirage plateau

T. Iritani et al. (HAL), JHEP1610(2016)101

Hadron interaction from LQCD (coupled-channel)

Lattice QCD simulation



Scattering S-matrix

$$S(E)=\begin{pmatrix} \eta e^{2i\delta_1} & i\sqrt{1-\eta^2}\, e^{i(\delta_1+\delta_2)} \\ i\sqrt{1-\eta^2}\, e^{i(\delta_1+\delta_2)} & \eta e^{2i\delta_2} \end{pmatrix}$$

$$\langle 0|B_1 B_2(t,\vec{r})\bar{I}(t_0)|0\rangle = A_0 \Psi(\vec{r},E_0)\, e^{-E_0(t-t_0)} + \cdots$$

Lüscher's finite volume method

M. Lüscher, NPB354(1991)531

Two-channel S-matrix has 3-parameters

$$\delta_1(E),\quad \delta_2(E),\quad \eta(E)$$

Relations of parameters

- Assumption of Interaction
- Fixed form of K-matrix

These are related to the energy E by an eigenvalue equation (s-wave)

$$\cos\left(\Delta_1+\Delta_2-\delta_1^0-\delta_2^0\right)=\eta\cos\left(\Delta_1-\Delta_2-\delta_1^0+\delta_2^0\right)$$

$$\frac{1}{\tan\Delta_i}=\frac{4\pi}{k_i}\cdot\frac{1}{L^3}\sum_p \frac{1}{p^2-k_i^2}$$

Unlike the single channel case,
the number of equations is less than the number of parameters in S-matrix.

Extra-information (relation) is necessary to solve coupled channel scattering

Potential in HAL QCD method

We define potentials which satisfy Schrödinger equation

$$\left(p^2+\nabla^2\right)\Psi^\alpha(E,\vec{r}) \equiv \int d^3y\, \underline{U^\alpha_\alpha(\vec{x},\vec{y})}\, \Psi^\alpha(E,\vec{y})$$

Energy independent potential

$$\left(p^2+\nabla^2\right)\Psi^\alpha(E,\vec{r}) = K^\alpha(E,\vec{r})$$

$$\begin{aligned} K^\alpha(E,\vec{r}) &\equiv \int dE'\, K^\alpha(E',\vec{x}) \underline{\int d^3y\, \tilde{\Psi}^\alpha(E',\vec{y})\Psi^\alpha(E,\vec{y})} \\ &= \int d^3y \left[\underline{\int dE'\, K^\alpha(E',\vec{x})\, \tilde{\Psi}^\alpha(E',\vec{y})}\right] \Psi^\alpha(E,\vec{y}) \\ &= \int d^3y\, U^\alpha_\alpha(\vec{x},\vec{y})\, \Psi^\alpha(E,\vec{y}) \end{aligned}$$

We can define an energy independent potential but it is fully non-local.

This potential automatically reproduce the scattering phase shift

N. Ishii et al Phys. Lett. B**712**(2012)437

Time-dependent Schrödinger like equation

Start with the normalized four-point correlator.

$$R_I^{B_1 B_2}(t,\vec r) = F^{B_1 B_2}{}_I(t,\vec r)\, e^{(m_1+m_2)t}$$

$$= A_0\, \Psi(\vec r, E_0)\, e^{-(E_0-m_1-m_2)t} + A_1\, \Psi(\vec r, E_1)\, e^{-(E_1-m_1-m_2)t} + \cdots$$

$$\left(\frac{p_0^2}{2\mu}+\frac{\nabla^2}{2\mu}\right)\Psi(\vec r, E_0) = \int U(\vec r, \vec r\,')\,\Psi(\vec r\,', E_0)\, d^3r'$$

$$E_n - m_1 - m_2 \approx \frac{p_n^2}{2\mu}$$

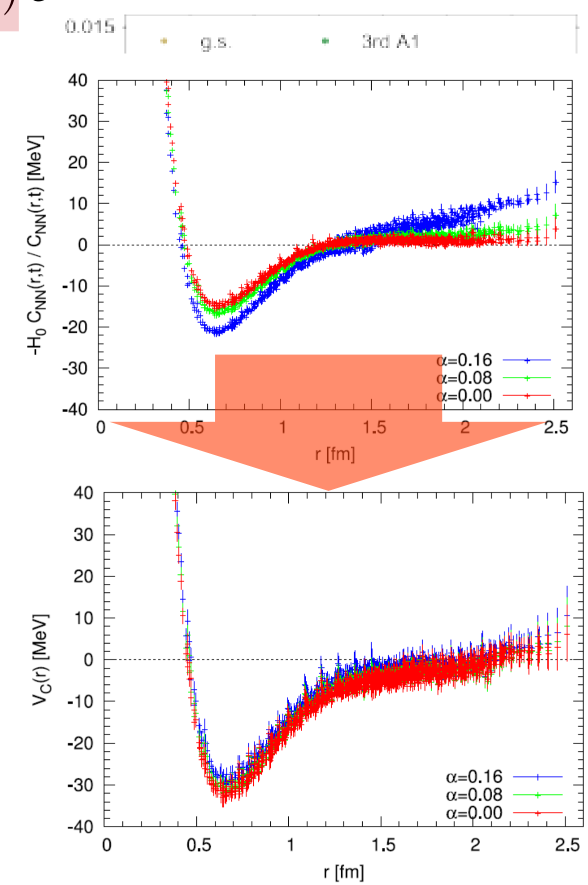
$$\left(\frac{p_1^2}{2\mu}+\frac{\nabla^2}{2\mu}\right)\Psi(\vec r, E_1) = \int U(\vec r, \vec r\,')\,\Psi(\vec r\,', E_1)\, d^3r'$$

A single state saturation is not required!!

$$\left(-\frac{\partial}{\partial t}+\frac{\nabla^2}{2\mu}\right) R_I^{B_1 B_2}(t,\vec r) = \int U(\vec r, \vec r\,')\, R_I^{B_1 B_2}(t,\vec r)\, d^3r'$$

Derivative (velocity) expansion of U

$$U(\vec r, \vec r\,') = \left[V_C(r) + S_{12} V_T(r)\right] + \left[\vec L\cdot\vec S_s V_{LS}(r) + \vec L\cdot\vec S_a V_{ALS}(r)\right] + O(\nabla^2)$$



BB interaction from NBS wave function

$$\left(-\frac{\partial}{\partial t}+\frac{\nabla^2}{2\mu}\right)R_I^{B_1B_2}(t,\vec{r})=\int U(\vec{r},\vec{r}\,')R_I^{B_1B_2}(t,\vec{r})d^3r'$$

Derivative (velocity) expansion of U is performed to deal with its nonlocality.

- For the case of oct-oct system,

$$U(\vec{r},\vec{r}\,') = \underbrace{\left[V_C(r)+S_{12}V_T(r)\right]}_{\text{Leading order part}}+\left[\vec{L}\cdot\vec{S}_s V_{LS}(r)+\vec{L}\cdot\vec{S}_a V_{ALS}(r)\right]+O(\nabla^2)$$

- For the case of dec-oct and dec-dec system,

$$U(\vec{r},\vec{r}\,') = \underbrace{\left[V_C(r)+S_{12}V_{T_1}(r)+S_{ii}V_{T_2}(r)+O(Spin\ op^3)\right]}_{\text{Leading order part}}+O(\nabla^2)$$

$$((\vec{r}\cdot\vec{S_1})^2-\frac{\vec{r}^2}{3}\vec{S_1}^2+(\vec{r}\cdot\vec{S_2})^2-\frac{\vec{r}^2}{3}\vec{S_2}^2)V_{T^2}(r)$$

$$\equiv\left[V_C^{eff}(r)\right]+O(\nabla^2)$$

We consider the effective central potential which contains not only the genuine central potential but also tensor parts.