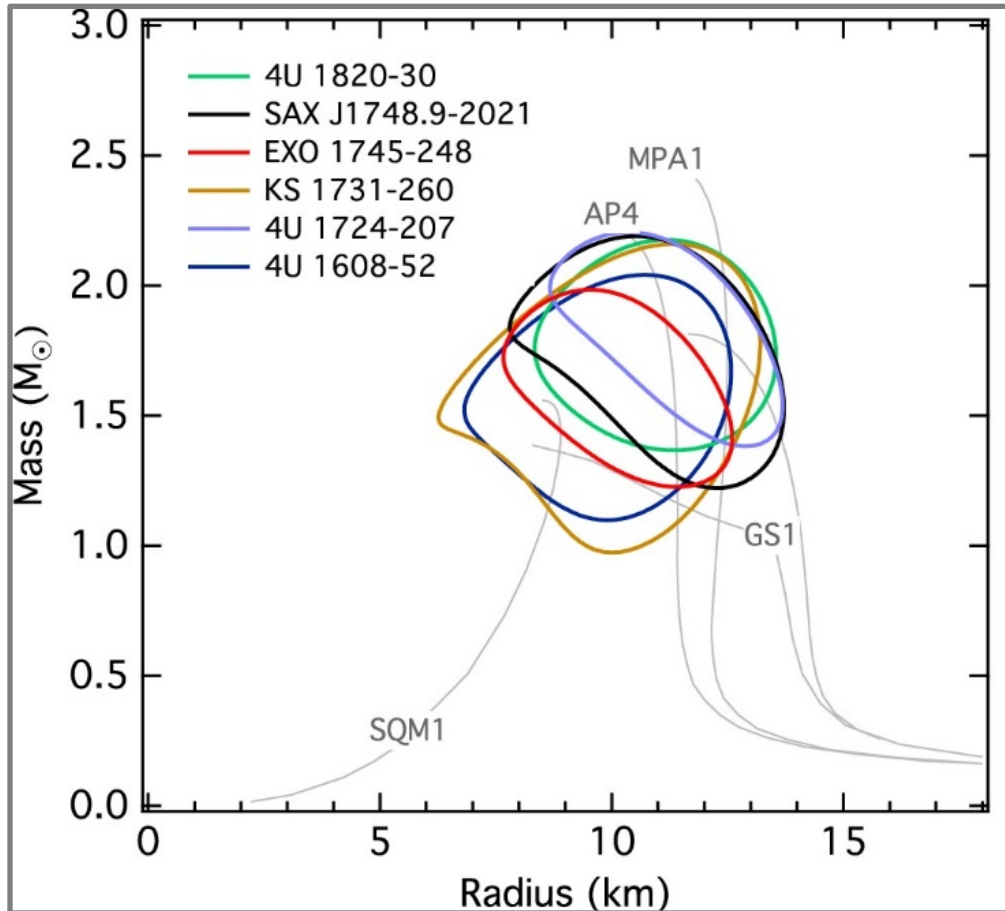
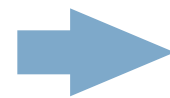


- Neutron Stars and what we learn from Femtoscopy
  - The Hyperon Puzzle
  - How can Femtoscopy help?
- Correlation functions
  - Theoretical description of the correlation function
  - Experimental correlation functions
- Results of the Analysis in the p-p collision system at 7 TeV
  - pp and  $\Lambda$ p Correlations
  - $\Lambda\Lambda$  Correlations
  - Future prospects with CATS
- Summary and Outlook





- Observation from X-Ray measurements of binary Systems:
  - Neutron star masses  $m > m_{\odot}$  and radii between 5 – 15 km
- From Shapiro Delay Measurements of binary systems: (P.Demorest et. Al, Nature 467 (2010), pp. 1081-1083)
  - Neutron star  $m = (1.97 \pm 0.04)m_{\odot}$
- Large nucleon chemical potential ( $\rho > \rho_0$ )



Hyperon production becomes energetically possible? e.g.  $\Lambda, \Xi, \Sigma$

Ozel et al. 2015 (arXiv:1505.05155v2)



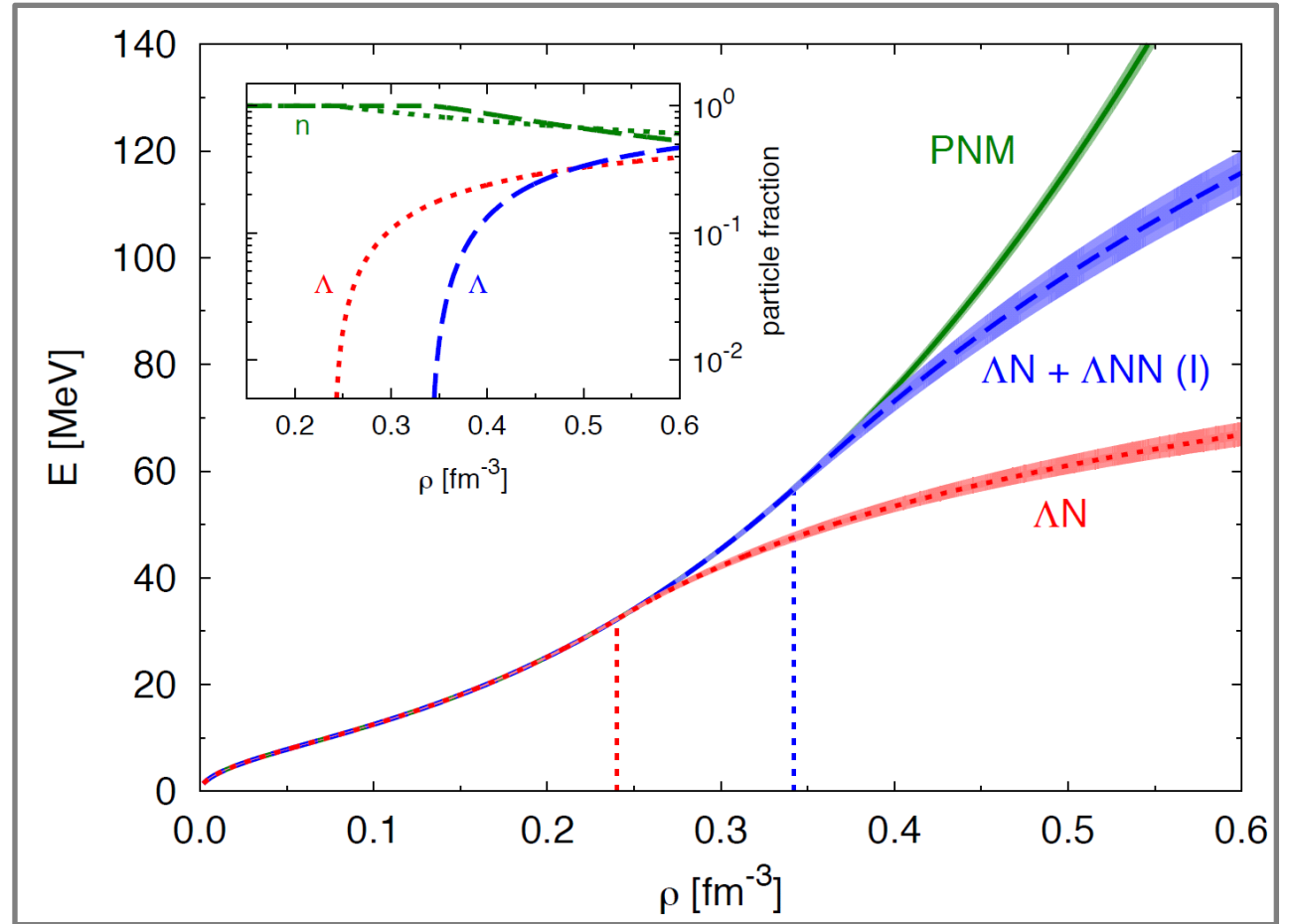


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# The Equation of State (EoS)



- Interaction potential between hyperons and neutron matter governs the onset of production and behavior of the EoS
- Quantum Monte Carlo study of pure nuclear matter and hyper neutron matter
  - Attractive  $\Lambda N$  interaction softens EoS
  - Repulsive  $\Lambda NN$  interaction stiffens EoS



D. Lonardoni, A. Lovato, S. Gandolfi, F. Pederiva Phys. Rev. Lett. 114, 092301 (2015)

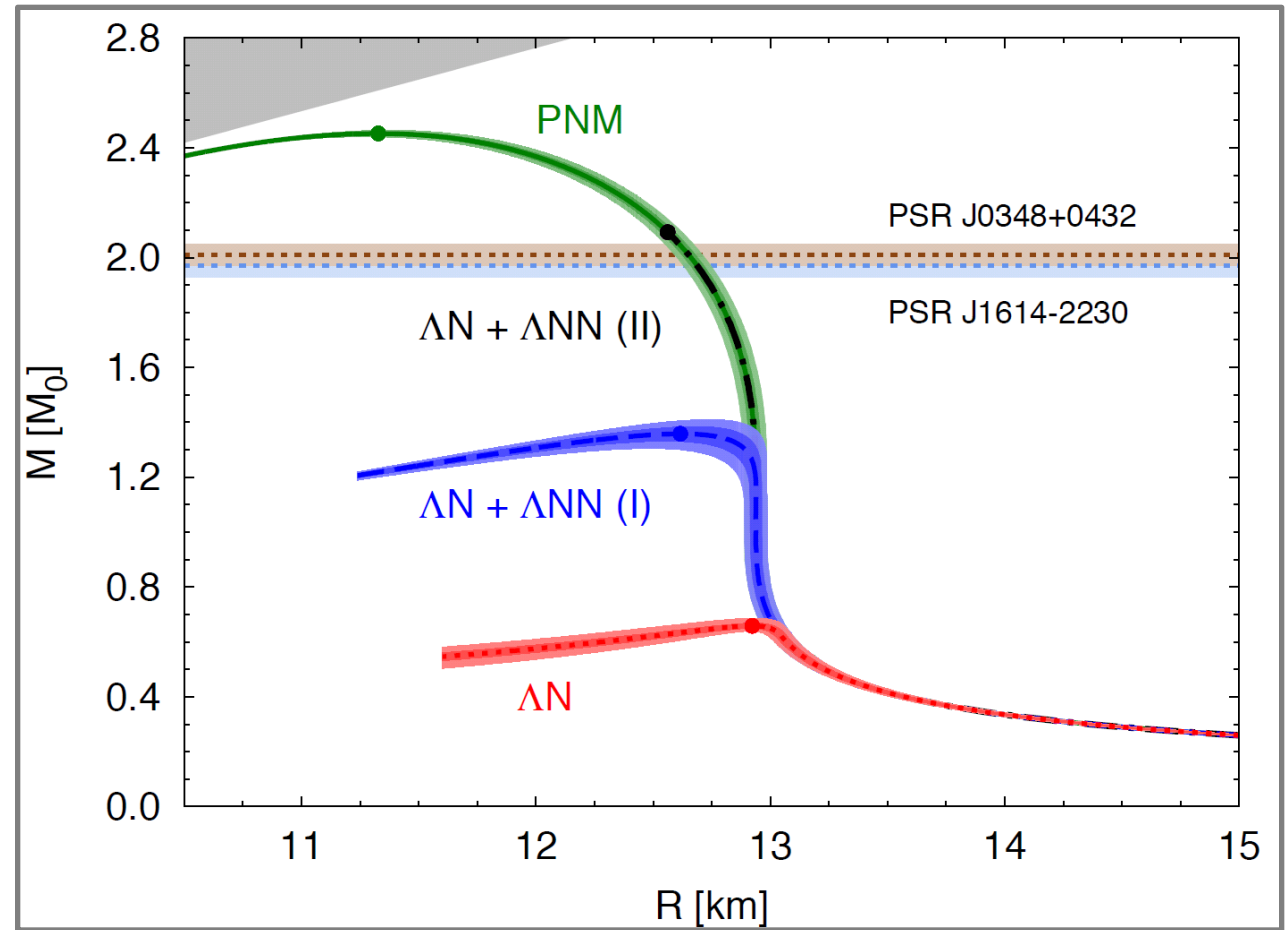


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# Mass – Radius Relation of Neutron Stars

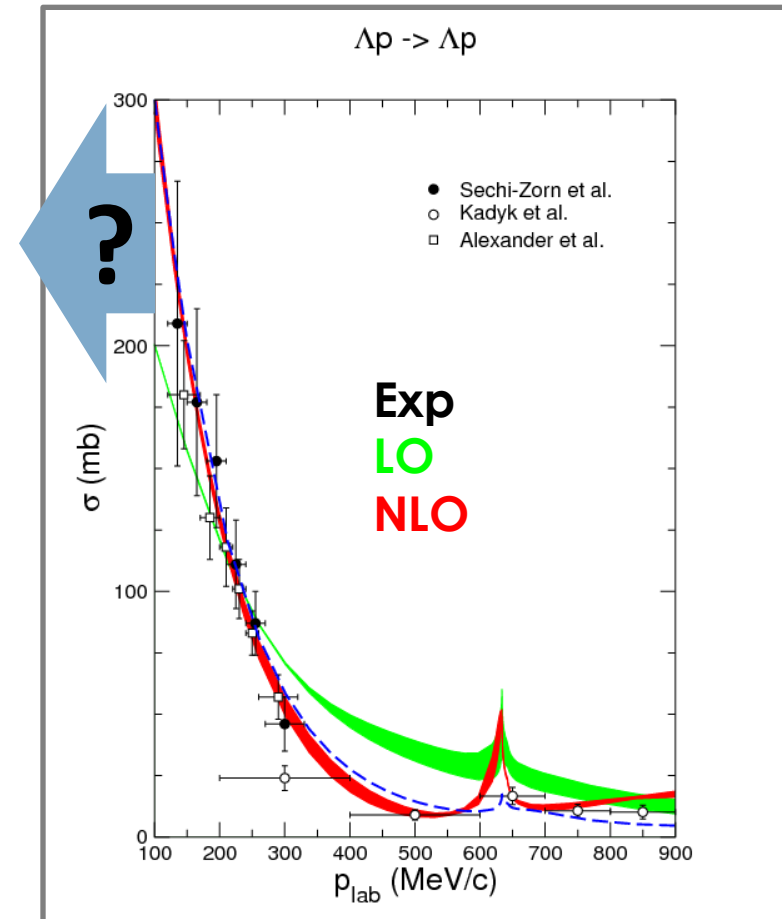


- EoS allows to solve the Tolman-Oppenheimer-Volkoff Equations
  - Mass Radius Relation
- Repulsive 3 body interaction shifts onset of hyperon production to larger densities and allows for higher masses
  - Not well constrained
- To understand the role of hyperons in neutron stars more constraints on the hyperon-neutron force are needed



D. Lonardoni, A. Lovato, S. Gandolfi, F. Pederiva Phys. Rev. Lett. 114, 092301 (2015)

- Data from scattering experiments from 1968 and 1971 in bubble chambers
  - $K^- + p \rightarrow \Sigma^0 + \pi^0, \Sigma^0 \rightarrow \Lambda + \gamma$
  - Production threshold for  $\Lambda$ 's :  $p \gtrsim 100$  MeV
- Different type of measurement needed to obtain constraints at low momentum
- Can we use Femtoscopic measurements?

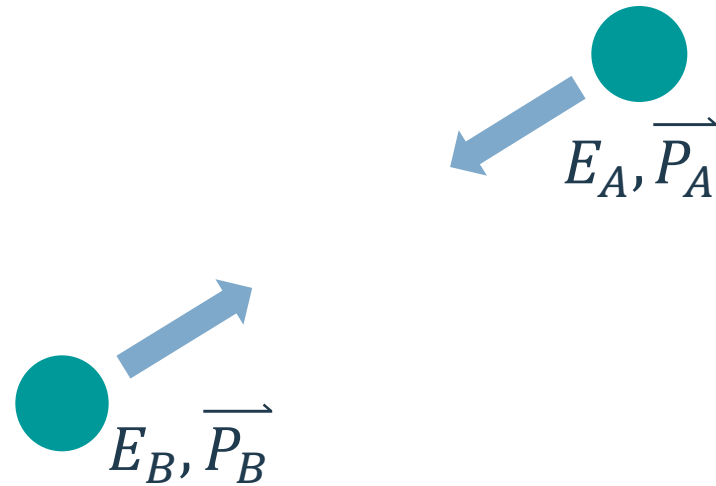


LO: H. Polinder, J.H., U. Meißner, NPA 779 (2006) 244  
 NLO: J.Haidenbauer., N.Kaiser, et al., NPA 915 (2013) 24



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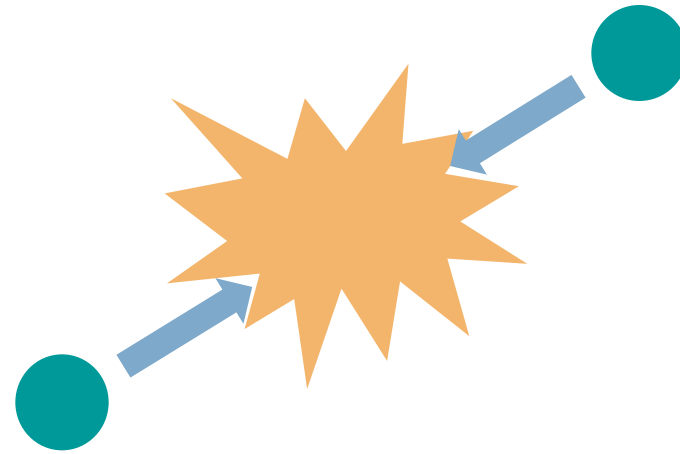
# Nuclear Collisions





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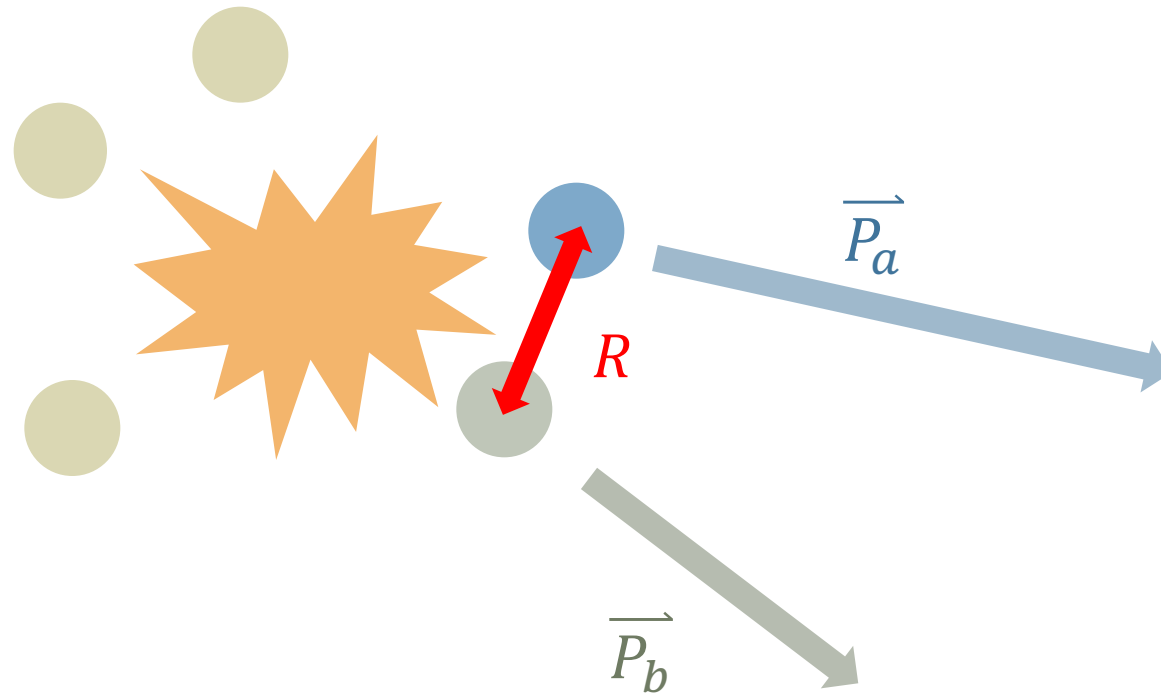
# Nuclear Collisions





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# Particle Production

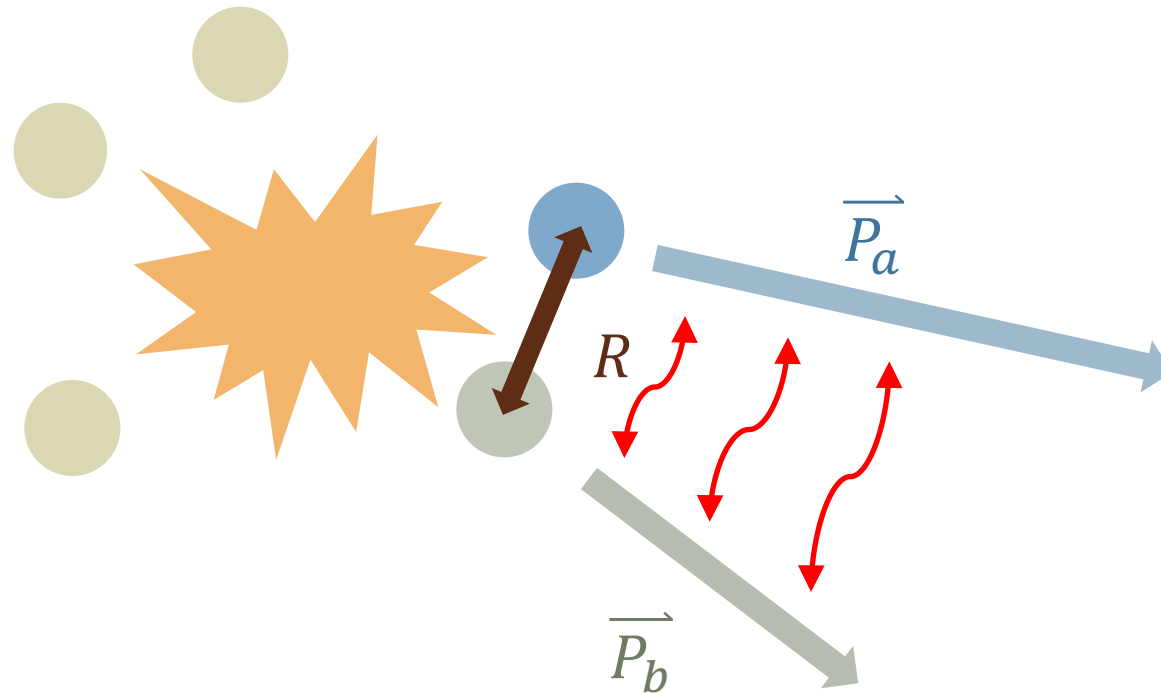






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# Particle Propagation





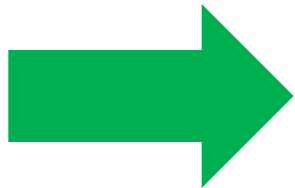
Theoretical formulation of the Correlation Function:

$$C(\vec{p}_a, \vec{p}_b) = \frac{P(\vec{p}_a, \vec{p}_b)}{P(\vec{p}_a)P(\vec{p}_b)} = C(k) = \int S(\vec{r}, k) |\psi(\vec{r}, k)|^2 d\vec{r}$$

Source



Wave Function



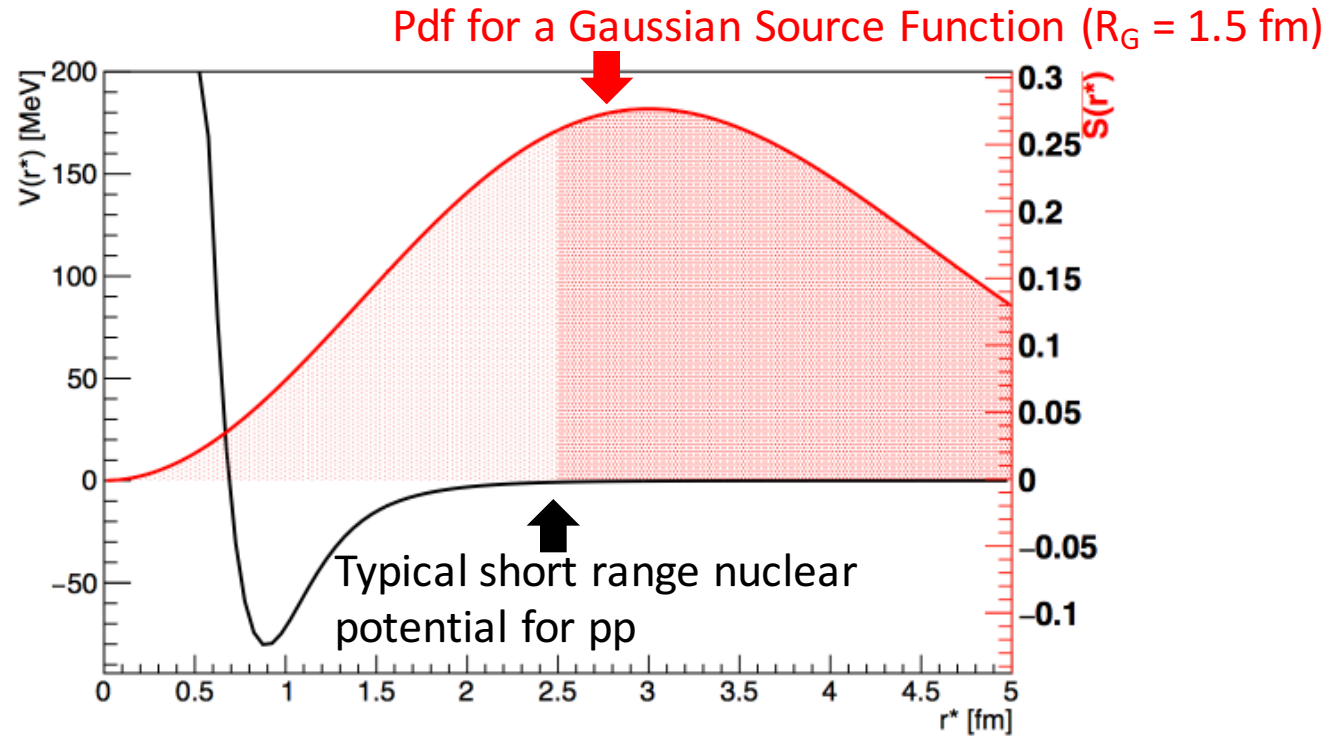
If we know about the **source**, we can learn about the **interaction**

Fix the parametrization of the source by fitting the pp and  $\Lambda p$  Correlation Function simultaneously to test different models of the  $p\Lambda$  interaction



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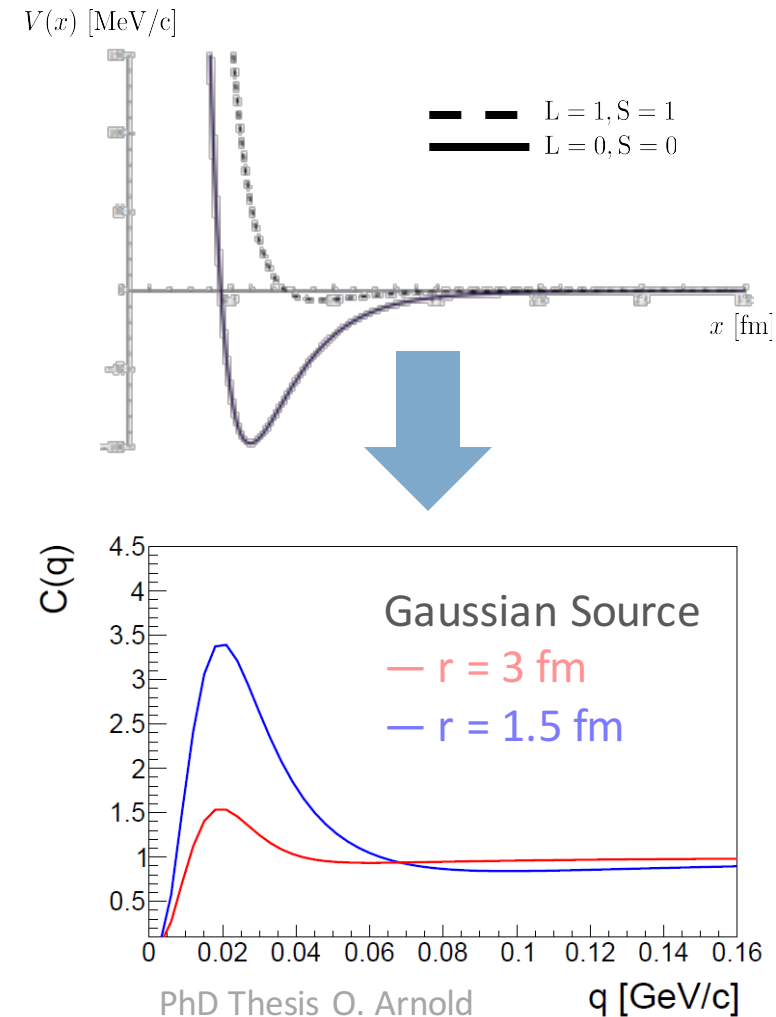
# Small Source $\rightarrow$ Repulsive Core



- Governed by:
  - Coulomb Interaction
  - Strong Interaction
  - Quantum Statistics
- Koonin Fit Function
  - Assumes a **Gaussian source** of size  $R_G$
$$C(k) = \int dr^3 \phi_{rel}^2(r, k) \exp\left(-\frac{r^2}{4R_G^2}\right)$$

S. E. Koonin, Phys. Lett. B 70 (1977) 43  
S. Pratt et al., Nucl. Phys. A 566 (1994) 103c

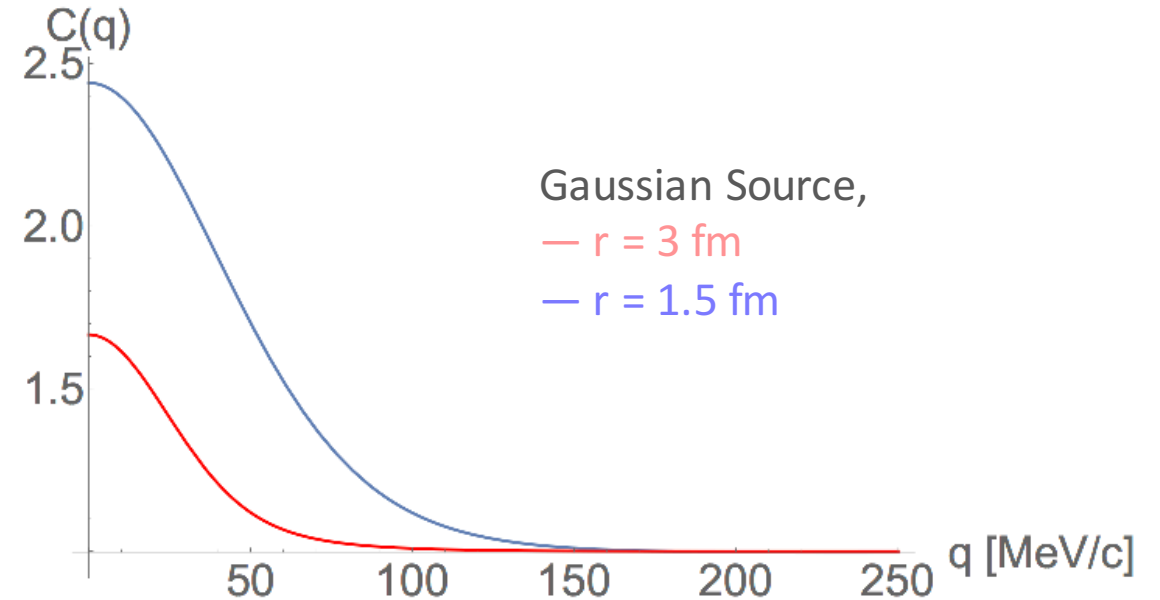
  - $\phi_{rel}$  from solving the Schrödinger Equation with the **known potentials** for the Coulomb and Strong interaction





# The $\Lambda p$ Correlation Function

- Governed by:
  - Strong Interaction
  - No Coulomb Interaction
- Lednický model
  - Assumes a **Gaussian source** of size  $R_G$
  - Based on the effective Range expansion
  - The interaction is modeled using the **scattering length** ( $f_0$ ) and the **effective range** ( $d_0$ )

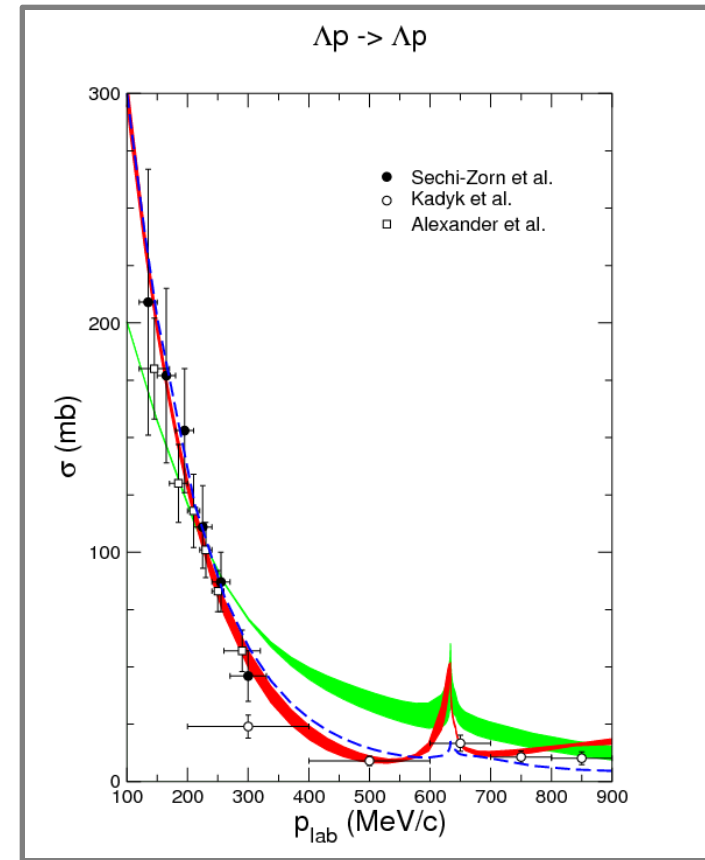


R. Lednický and V. L. Lyuboshits, *Sov. J. Nucl. Phys.* **35**, 770 (1982), [*Yad. Fiz.*35,1316(1981)].

$$C(k) = 1 + \sum_S \rho_S \left[ \frac{1}{2} \left| \frac{f^S(k)}{R_G^{\Lambda p}} \right|^2 \left( 1 - \frac{d_0^S}{2\sqrt{\pi}R_G^{\Lambda p}} \right) + 2 \frac{\mathcal{R}f^S(k)}{\sqrt{\pi}R_G^{\Lambda p}} F_1(QR_G^{\Lambda p}) - \frac{\mathcal{I}f^S(k)}{R_G^{\Lambda p}} F_2(QR_G^{\Lambda p}) \right]$$

- Governed by:
  - Strong Interaction
  - No Coulomb interaction
- Lednický model
  - Assumes a **Gaussian source** of size  $R_G$
  - Based on the effective Range expansion
  - The interaction is modeled using the **scattering length** ( $f_0$ ) and the **effective range** ( $d_0$ )

R. Lednický and V. L. Lyuboshits, *Sov. J. Nucl. Phys.* **35**, 770 (1982), [*Yad. Fiz.*35,1316(1981)].

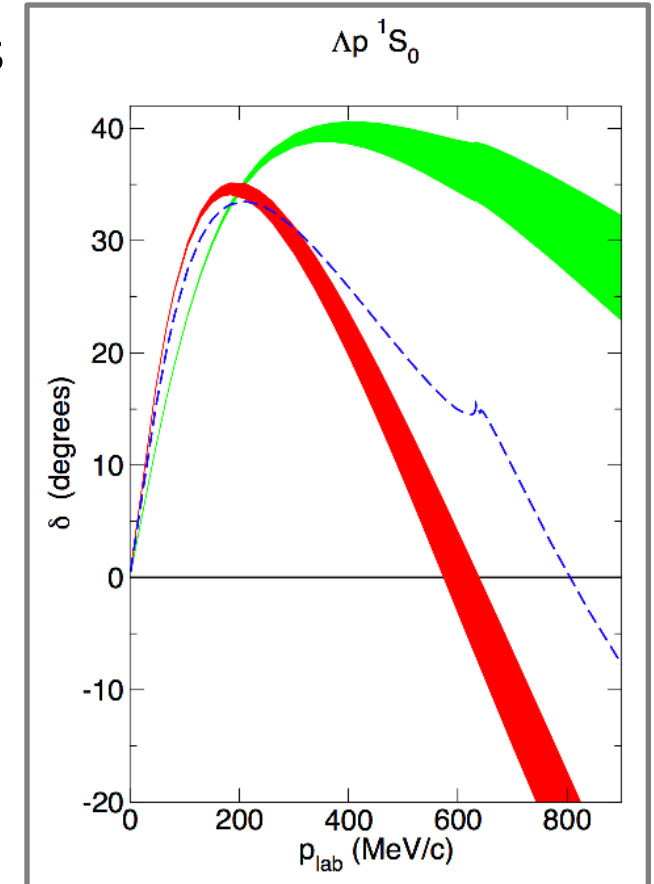


An expression containing  $f_0$

$$C(k) = 1 + \sum_S \rho_S \left[ \frac{1}{2} \left| \frac{f^S(k)}{R_G^{\Lambda p}} \right|^2 \left( 1 - \frac{d_0^S}{2\sqrt{\pi}R_G^{\Lambda p}} \right) + 2 \frac{\mathcal{R}f^S(k)}{\sqrt{\pi}R_G^{\Lambda p}} F_1(QR_G^{\Lambda p}) - \frac{\mathcal{I}f^S(k)}{R_G^{\Lambda p}} F_2(QR_G^{\Lambda p}) \right]$$

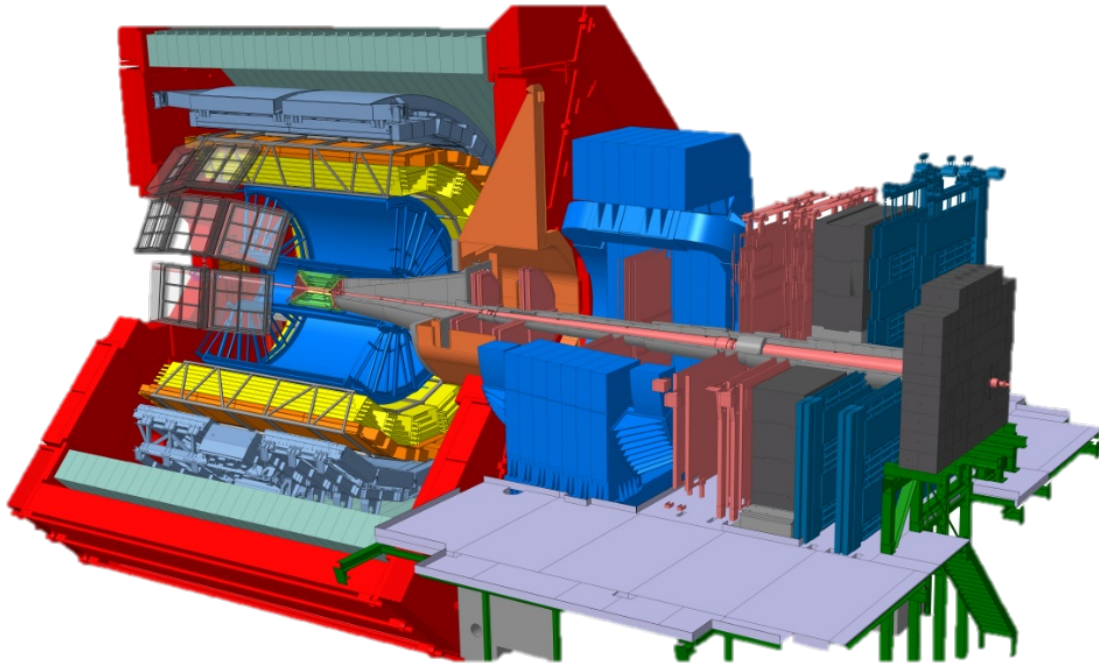
- Calculations based on an chiral effect field theory calculations
- Discrimination of the two models not possible due to the lack of data
- NLO predicts a repulsive core of the interaction potential

	NLO						LO
$\Lambda$ [MeV]	450	500	550	600	650	700	600
$a_s^{\Lambda p}$	-2.90	-2.91	-2.91	-2.91	-2.90	-2.90	-1.91
$r_s^{\Lambda p}$	2.64	2.86	2.84	2.78	2.65	2.56	1.40
$a_t^{\Lambda p}$	-1.70	-1.61	-1.52	-1.54	-1.51	-1.48	-1.23
$r_t^{\Lambda p}$	3.44	3.05	2.83	2.72	2.64	2.62	2.13



Nucl. Phys. A915 (2013) 24-58,  
arXiv:1304.5339

pp 7 TeV



- Tracking and Particle Identification (PID) of charged particles
  - $|\eta| < 0.9$
  - $2\pi$  coverage in the azimuth
- Inner Tracking System (ITS)
- Time Projection Chamber (TPC)
- Time of Flight (TOF) System





$$C(\vec{p}_a, \vec{p}_b) = \frac{P(\vec{p}_a, \vec{p}_b)}{P(\vec{p}_a)P(\vec{p}_b)} = \mathcal{N} \frac{N_{SE}(k)}{N_{ME}(k)}$$



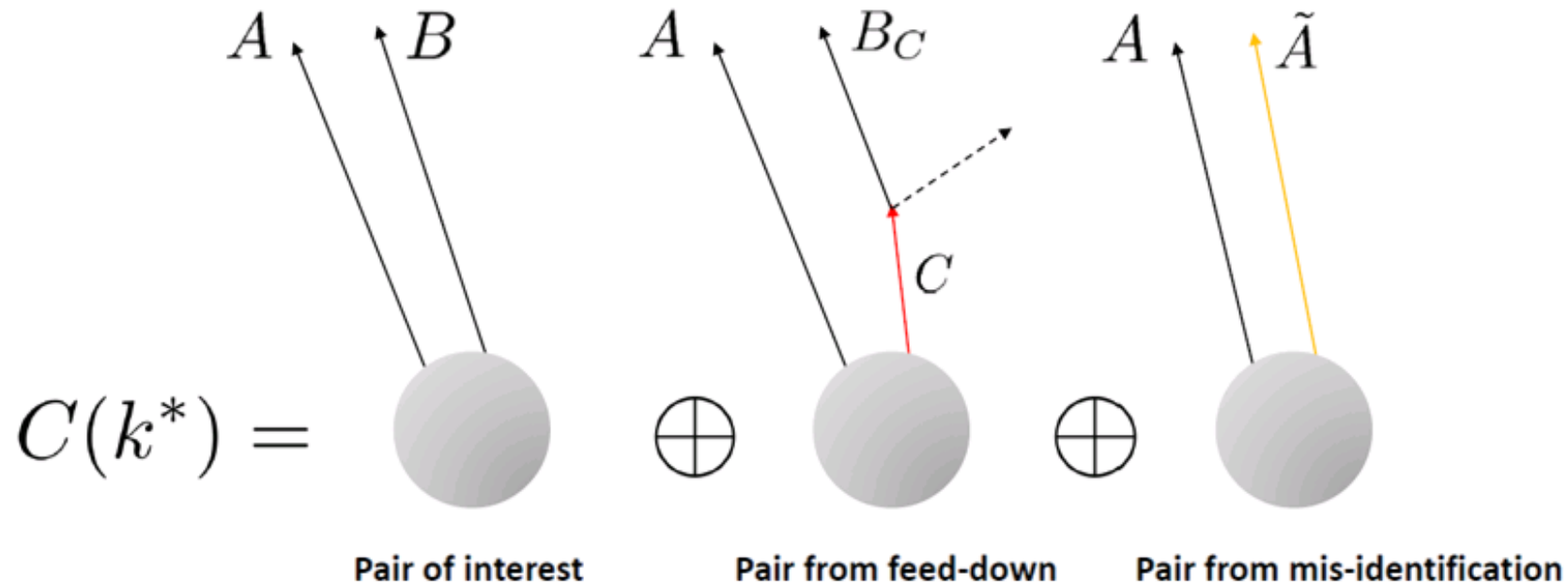


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# The Experimental Correlation Function



$$C(\vec{p}_a, \vec{p}_b) = \frac{P(\vec{p}_a, \vec{p}_b)}{P(\vec{p}_a)P(\vec{p}_b)} = \mathcal{N} \frac{N_{SE}(k)}{N_{ME}(k)}$$



$$C(\vec{p}_a, \vec{p}_b) = \frac{P(\vec{p}_a, \vec{p}_b)}{P(\vec{p}_a)P(\vec{p}_b)} = \mathcal{N} \frac{N_{SE}(k)}{N_{ME}(k)} = \lambda_1 C_1(k) + \lambda_2 C_2(k) + \dots$$

Correlation function of interest

Contributions from impurities, secondaries etc.

$\lambda$  Parameters can be related to measured single-particle quantities, e.g. Purity  $\mathcal{P}$  or feed-down fractions  $f$

(PhD Thesis of O. Arnold - [https://www.das.ktas.ph.tum.de/DasDocs/Public/PhD\\_Theses/arnold-oliver\\_thesis.pdf](https://www.das.ktas.ph.tum.de/DasDocs/Public/PhD_Theses/arnold-oliver_thesis.pdf))





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# Decomposition of the Correlation Functions



$$\{pp\} = pp + p_{\Lambda}p + p_{\Lambda}p_{\Lambda} + p_{\Sigma^+}p + p_{\Sigma^+}p_{\Sigma^+} + p_{\Lambda}p_{\Sigma^+} + \tilde{p}p + \tilde{p}\tilde{p}.$$

$$\begin{aligned} \{p\Lambda\} = & p\Lambda + p\Lambda_{\Xi^-} + p\Lambda_{\Xi^0} + p\Lambda_{\Sigma^0} + p_{\Lambda}\Lambda + p_{\Lambda}\Lambda_{\Xi^-} \\ & + p_{\Lambda}\Lambda_{\Xi^0} + p_{\Lambda}\Lambda_{\Sigma^0} + p_{\Sigma^+}\Lambda + p_{\Sigma^+}\Lambda_{\Xi^-} + p_{\Sigma^+}\Lambda_{\Xi^0} + p_{\Sigma^+}\Lambda_{\Sigma^0} \\ & + \tilde{p}\Lambda + p\tilde{\Lambda} + \tilde{p}\tilde{\Lambda} \end{aligned}$$

Secondary contributions from measurements

Pair	Percentage %
$pp$	75
$p_{\Lambda}p$	16
$p_{\Lambda}p_{\Lambda}$	1
$p_{\Sigma^+}p$	6
$p_{\Sigma^+}p_{\Sigma^+}$	0
$p_{\Sigma^+}p_{\Lambda}$	0
$\tilde{p}p$	2
$\tilde{p}\tilde{p}$	0

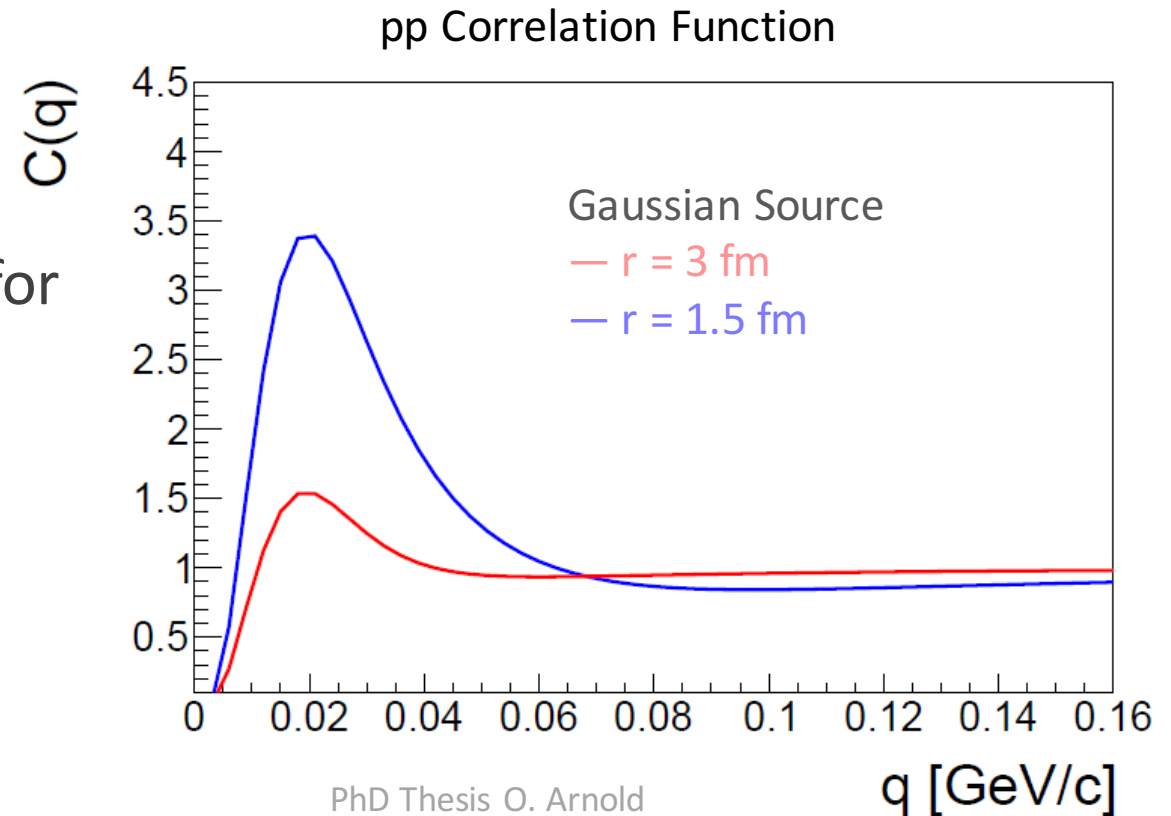
Pair	Percentage %
$p\Lambda$	49
$p\Lambda_{\Xi^-}$	10
$p\Lambda_{\Xi^0}$	10
$p\Lambda_{\Sigma^0}$	16
$p_{\Lambda}\Lambda$	5
$p_{\Lambda}\Lambda_{\Xi^-}$	1
$p_{\Lambda}\Lambda_{\Xi^0}$	1
$p_{\Lambda}\Lambda_{\Sigma^0}$	2
$p_{\Sigma^+}\Lambda$	0
$p_{\Sigma^+}\Lambda_{\Xi^-}$	0
$p_{\Sigma^+}\Lambda_{\Xi^0}$	1
$p_{\Sigma^+}\Lambda_{\Sigma^0}$	2
$\tilde{p}\Lambda$	1
$p\tilde{\Lambda}$	2
$\tilde{p}\tilde{\Lambda}$	0



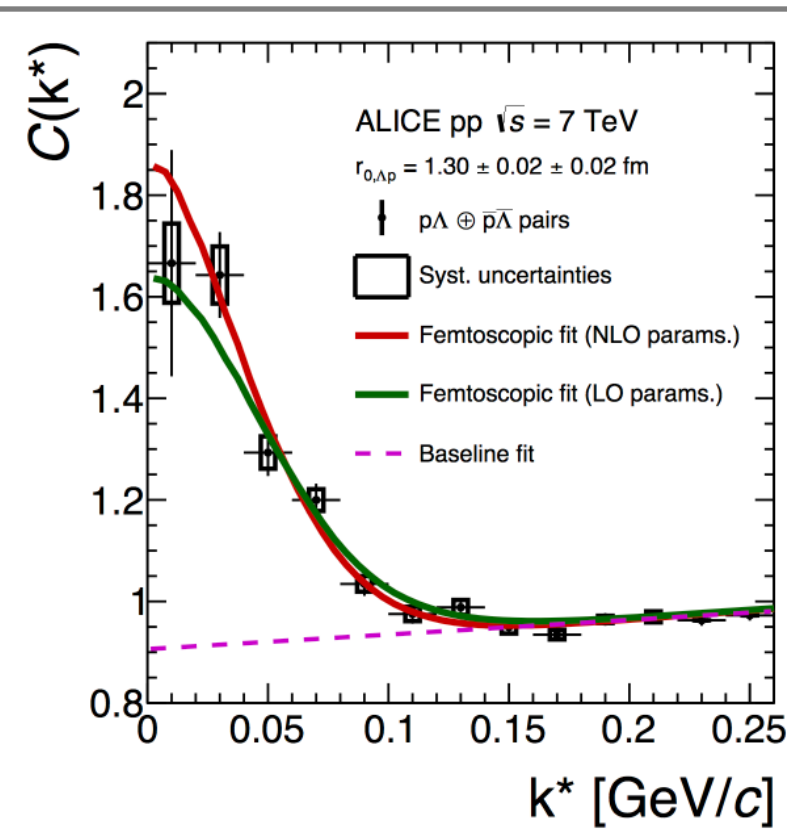
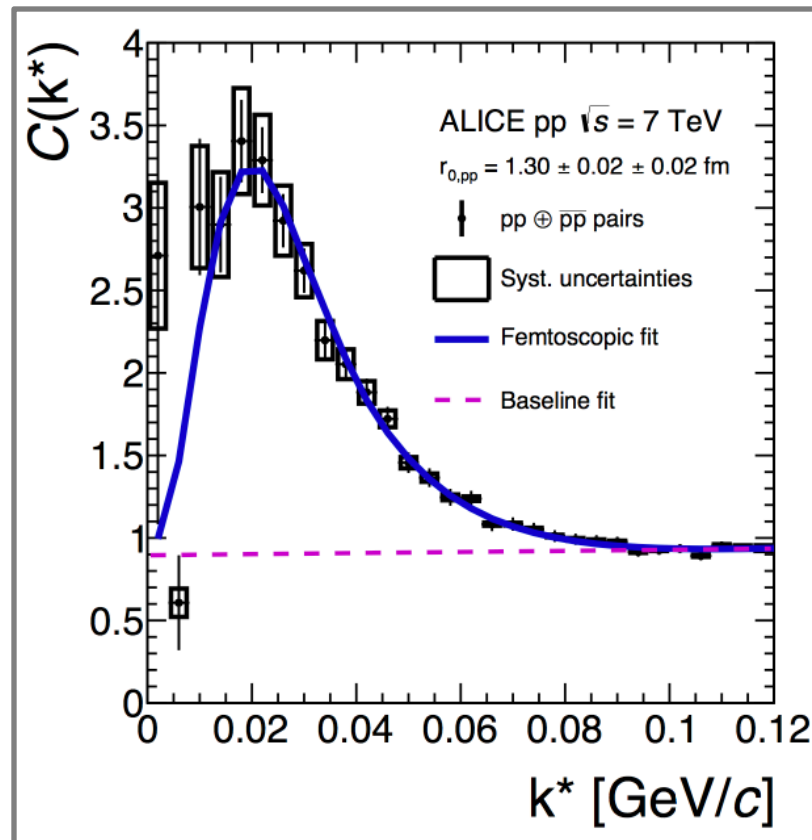
- Mini Jet Background not present for baryon baryon pairs

(see J.Adam et al., Europ. Phys. Jour. C, Aug. 17, 77:569)

- Evolution of the system is better understood compared to Pb-Pb
  - Same freeze out times and source size for all particles
- Small source sizes
  - Stronger signal
  - Sensitivity to the shape of the potential

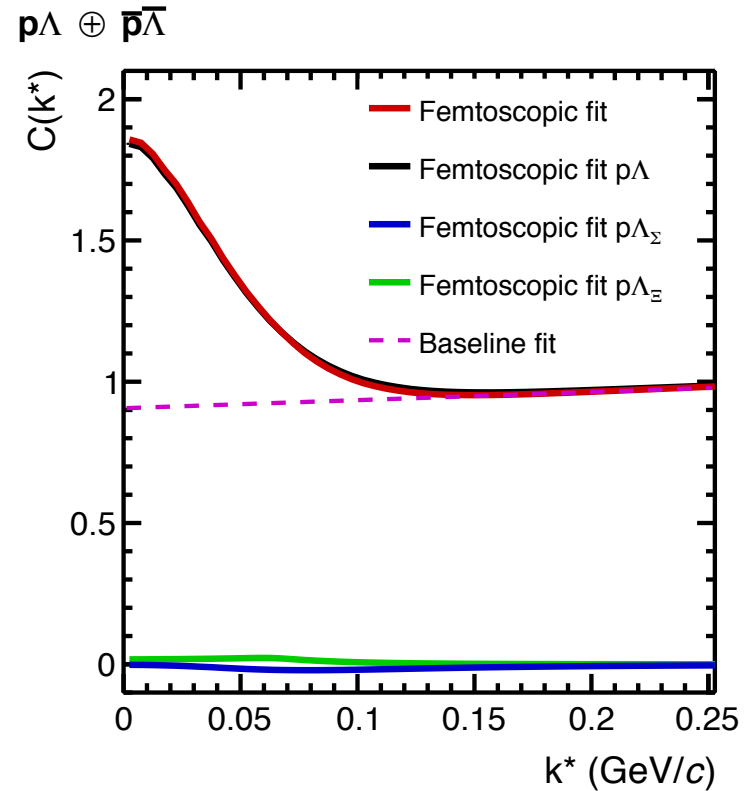
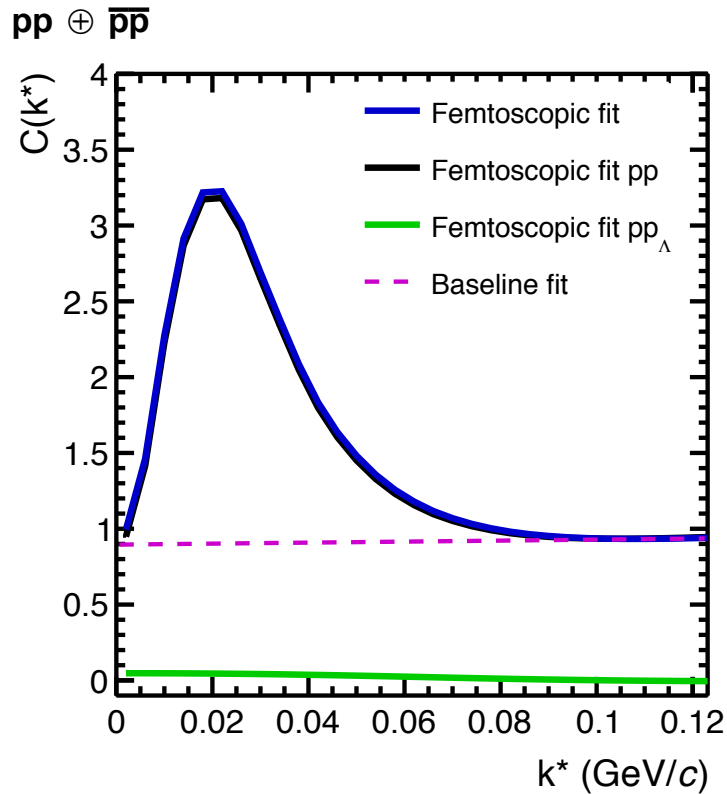


- **NLO** and **LO** Parameters show different features at low  $k^*$
- More statistics are required to distinguish between the two calculations
  - LHC Run 2
  - Other collision systems e.g. p-Pb





# Fit Decomposition



? Any Sensitivity to the repulsive core of the potential?  
 Included in pp not in Λp currently.

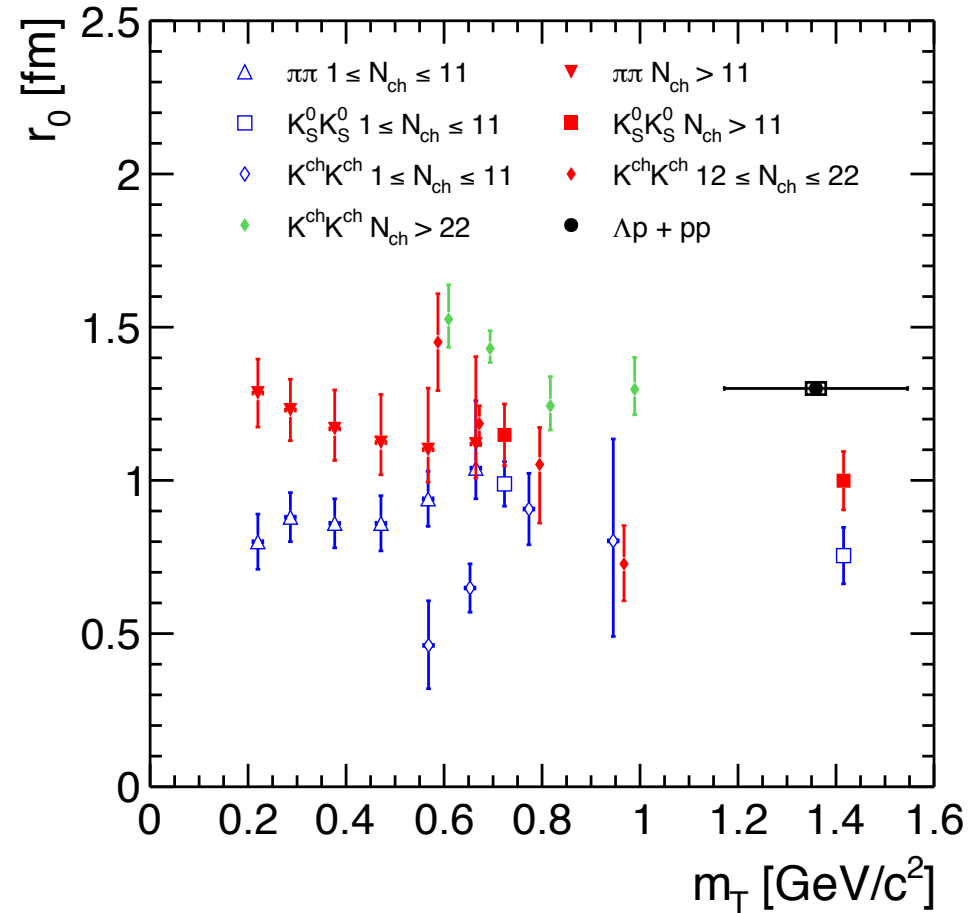


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# Radius



ALICE Phys. Lett. B717 (2012) nn. 151–161. Neutral Kaons femtoscopy  
ALICE Phys. Rev. D87.5 (20 Charged Kaons femtoscopy

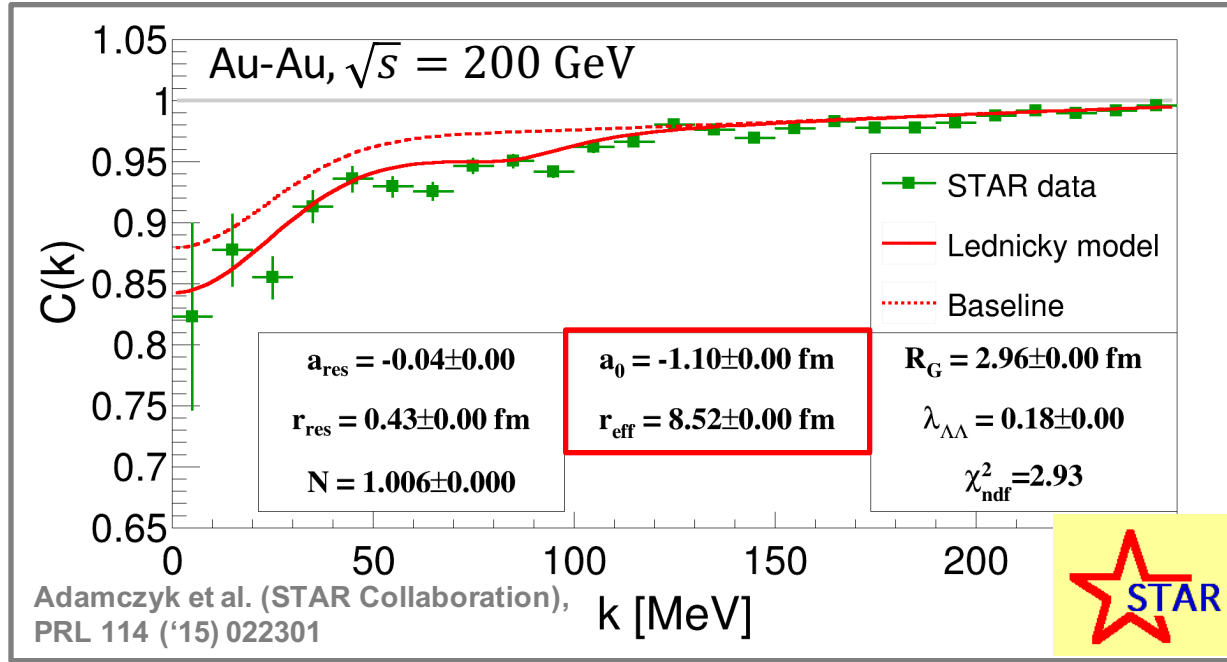






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# The $\Lambda\Lambda$ Correlation Function

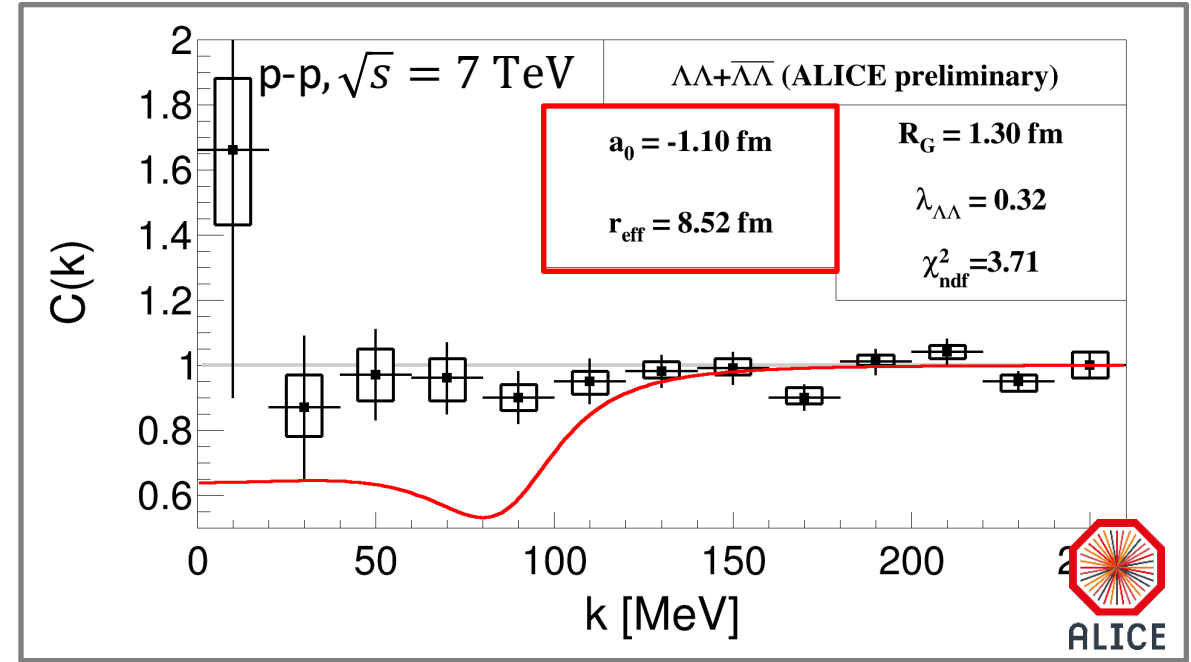
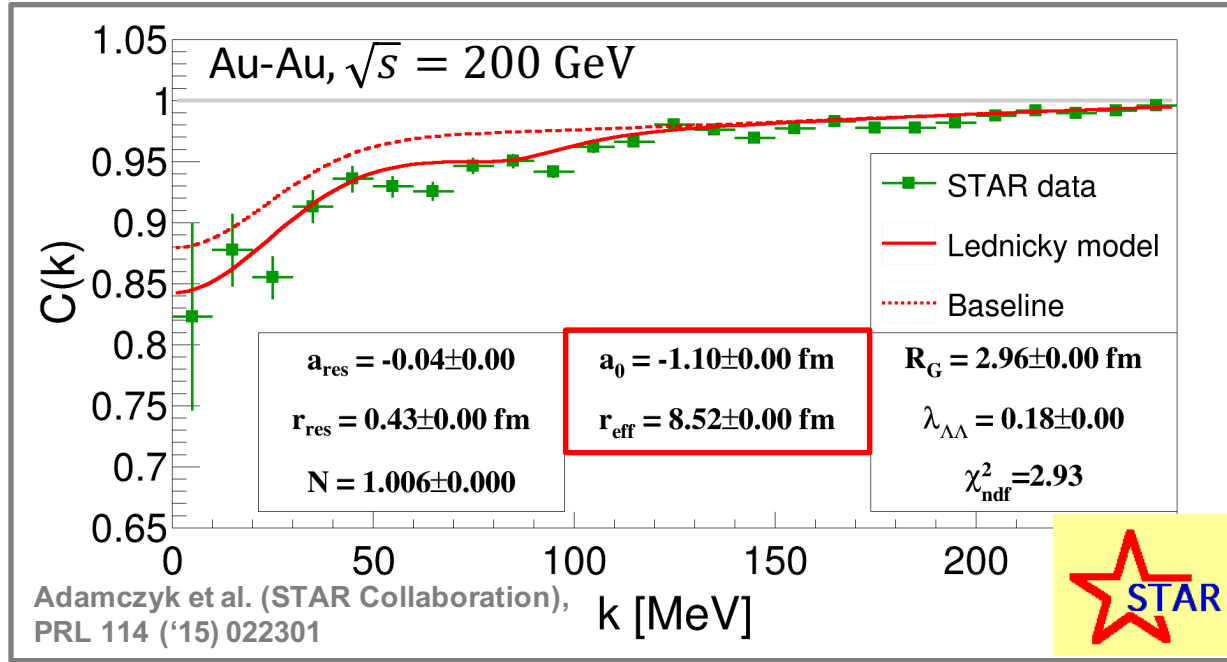


- Analysis of the  $\Lambda\Lambda$  correlation function in Au-Au Collisions (STAR):
  - Fit of the Lednický Model with an additional residual term as a free parameter
  - Results in a slightly repulsive interaction ( $a_0 < 0$  fm)
- Modeled correlation function does not describe ALICE Data with an adjusted source



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# The $\Lambda\Lambda$ Correlation Function

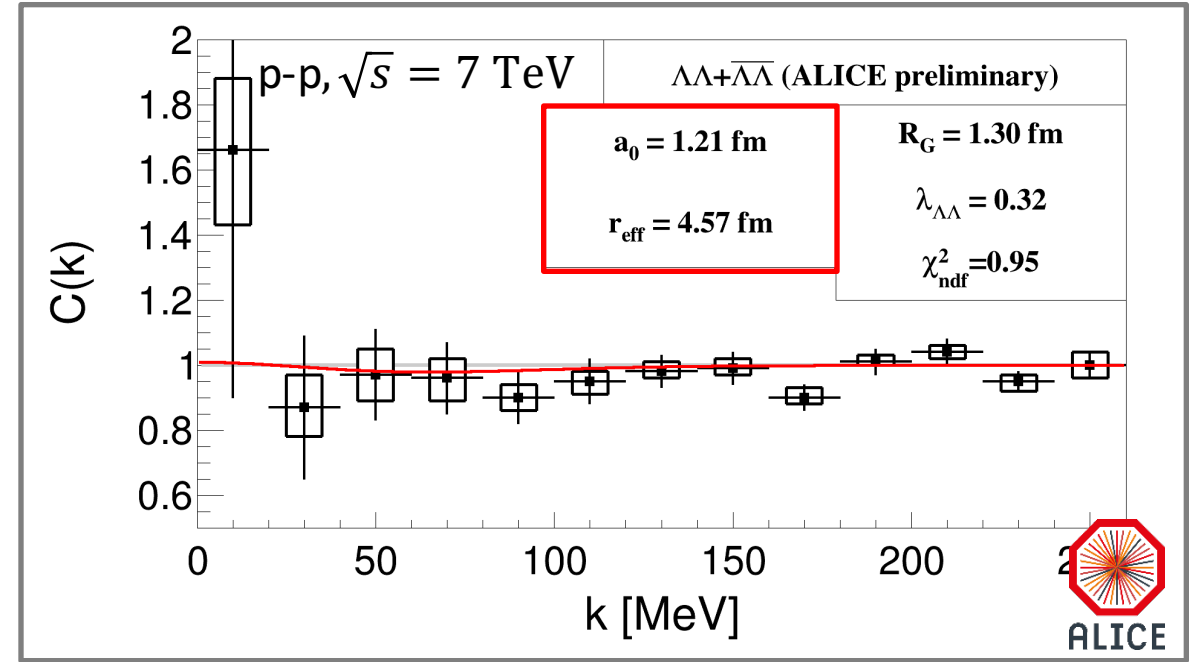
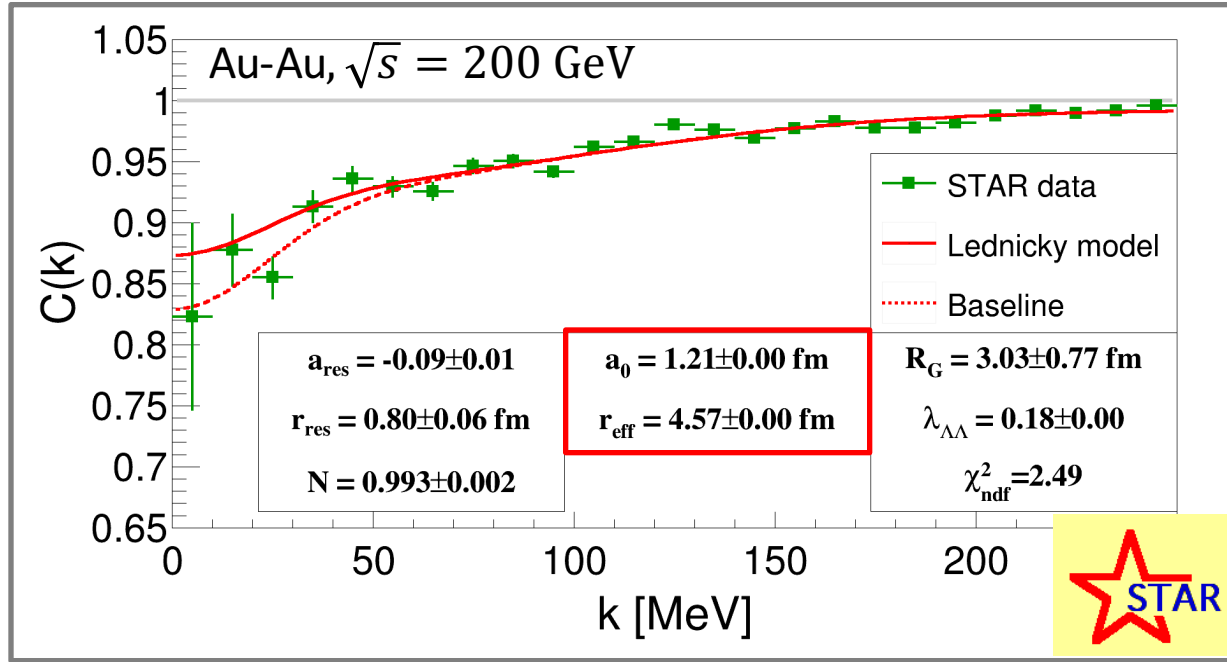


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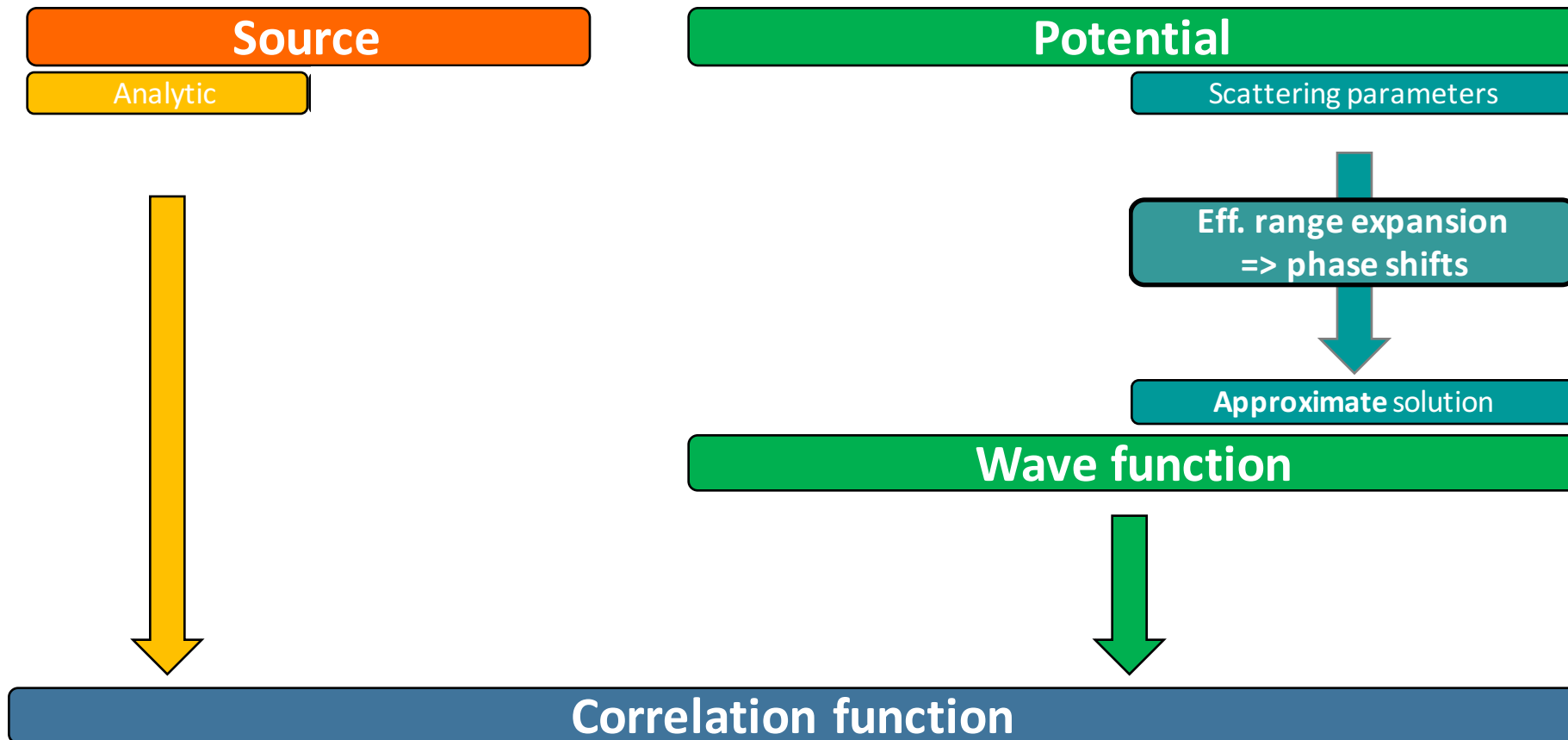


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# The $\Lambda\Lambda$ Correlation Function



- Reanalysis of the data (K. Morita et al., T.Furumoto, AO PRC91('15) 024916) finds an attractive interaction ( $a_0 > 1.25$  fm)
- An attractive interaction allows to describe both STAR and ALICE data
- Larger statistics of Run 2 could yield a conclusive answer





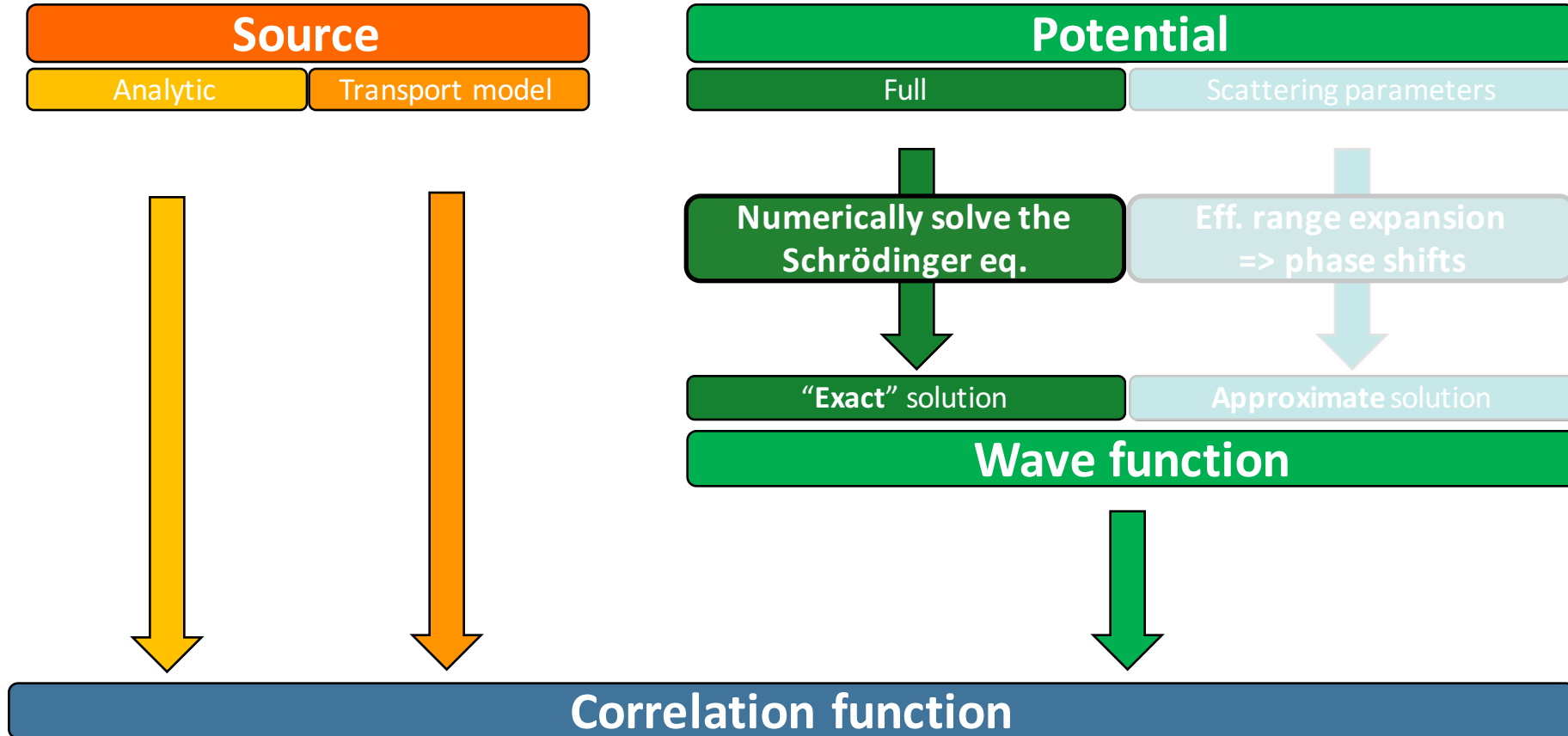
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# CATS – Correlation Analysis Tool Using the Schrödinger Equation



<http://www.denseandstrange.ph.tum.de/index.php?id=78>

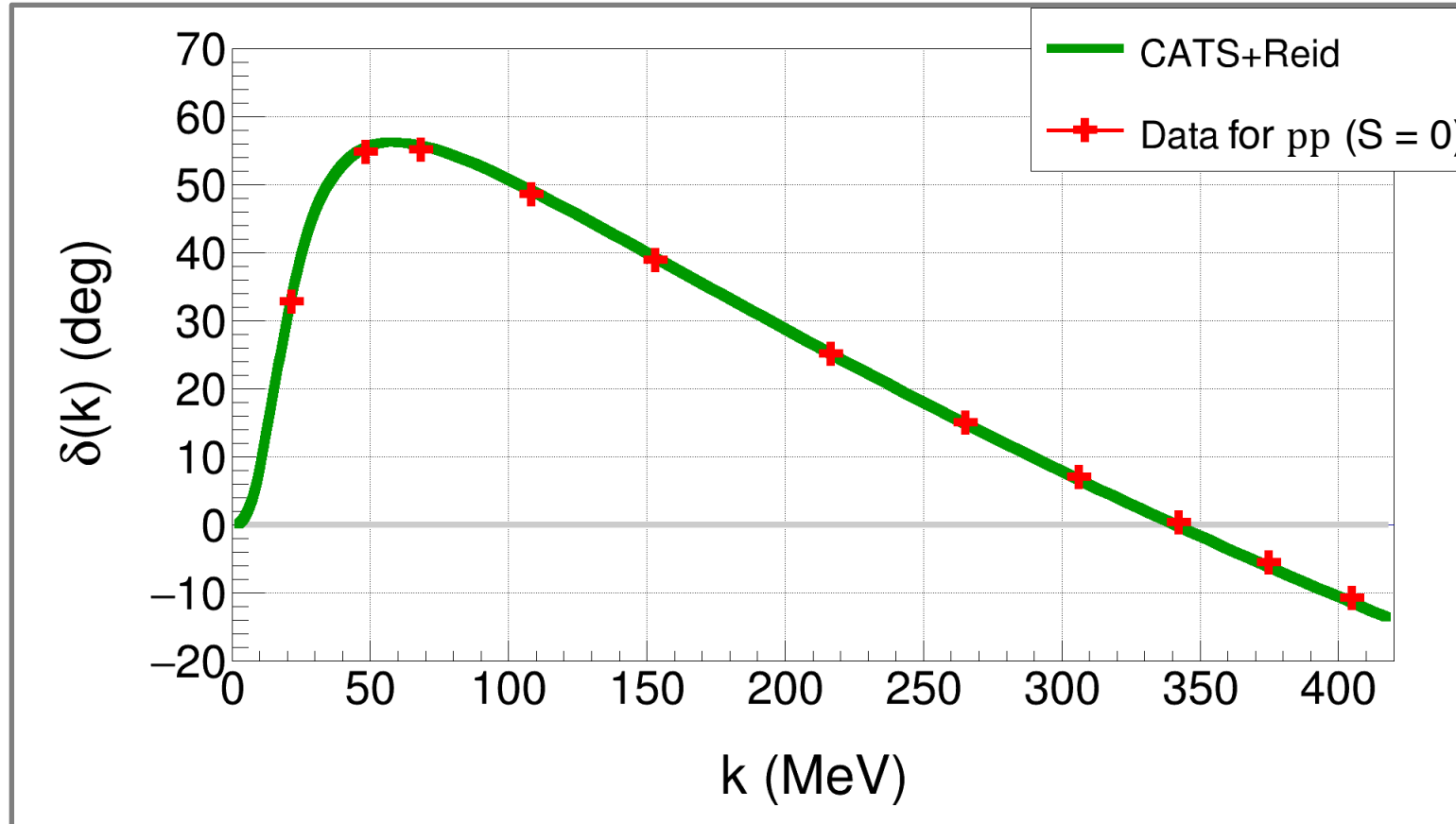
$$C(k) = \int S(\vec{r}, k) |\psi(\vec{r}, k)|^2 d\vec{r} \xrightarrow{k \rightarrow \infty} 1$$



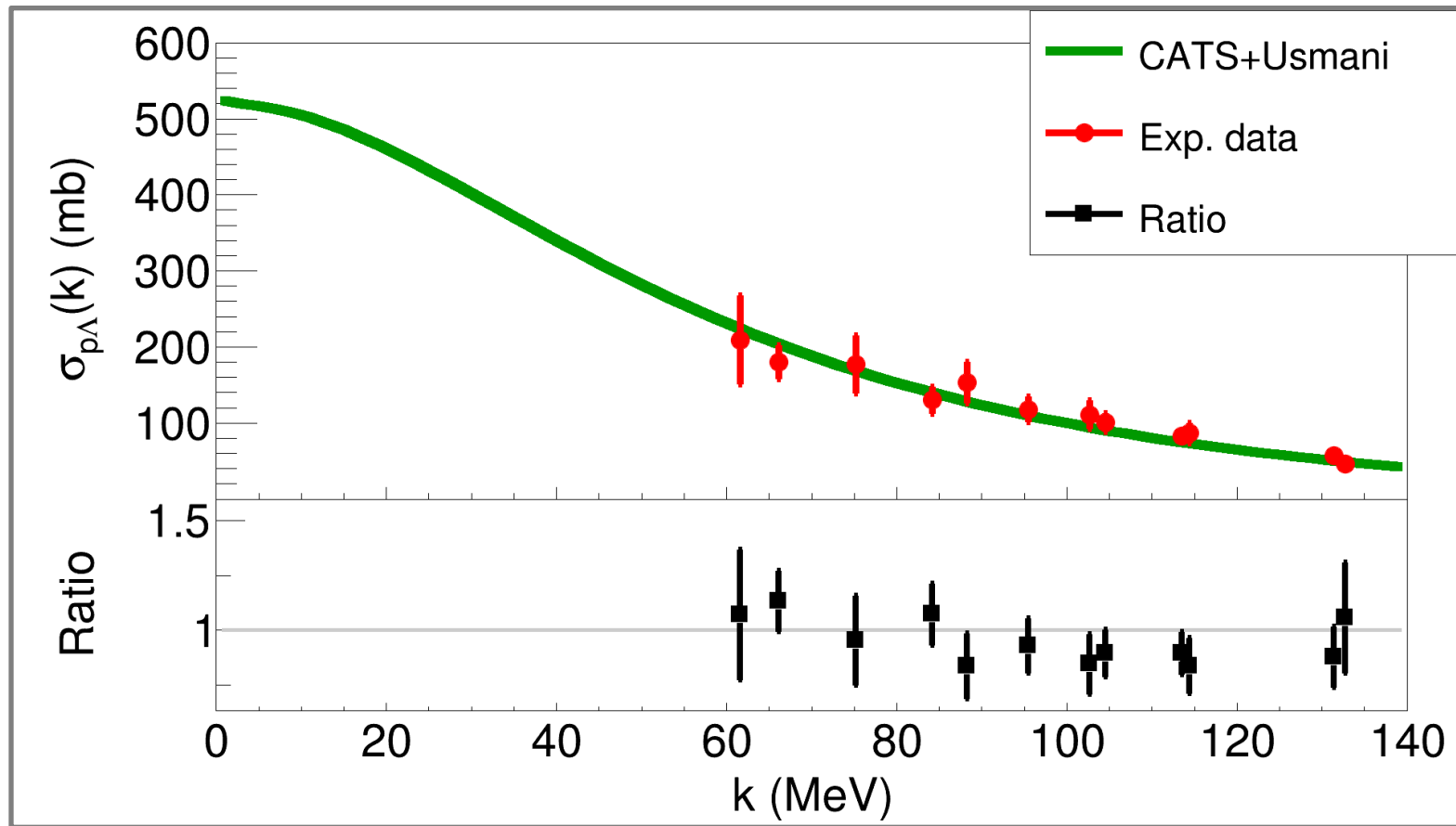


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# Phase Shift calculations for pp ( $S = 0$ )



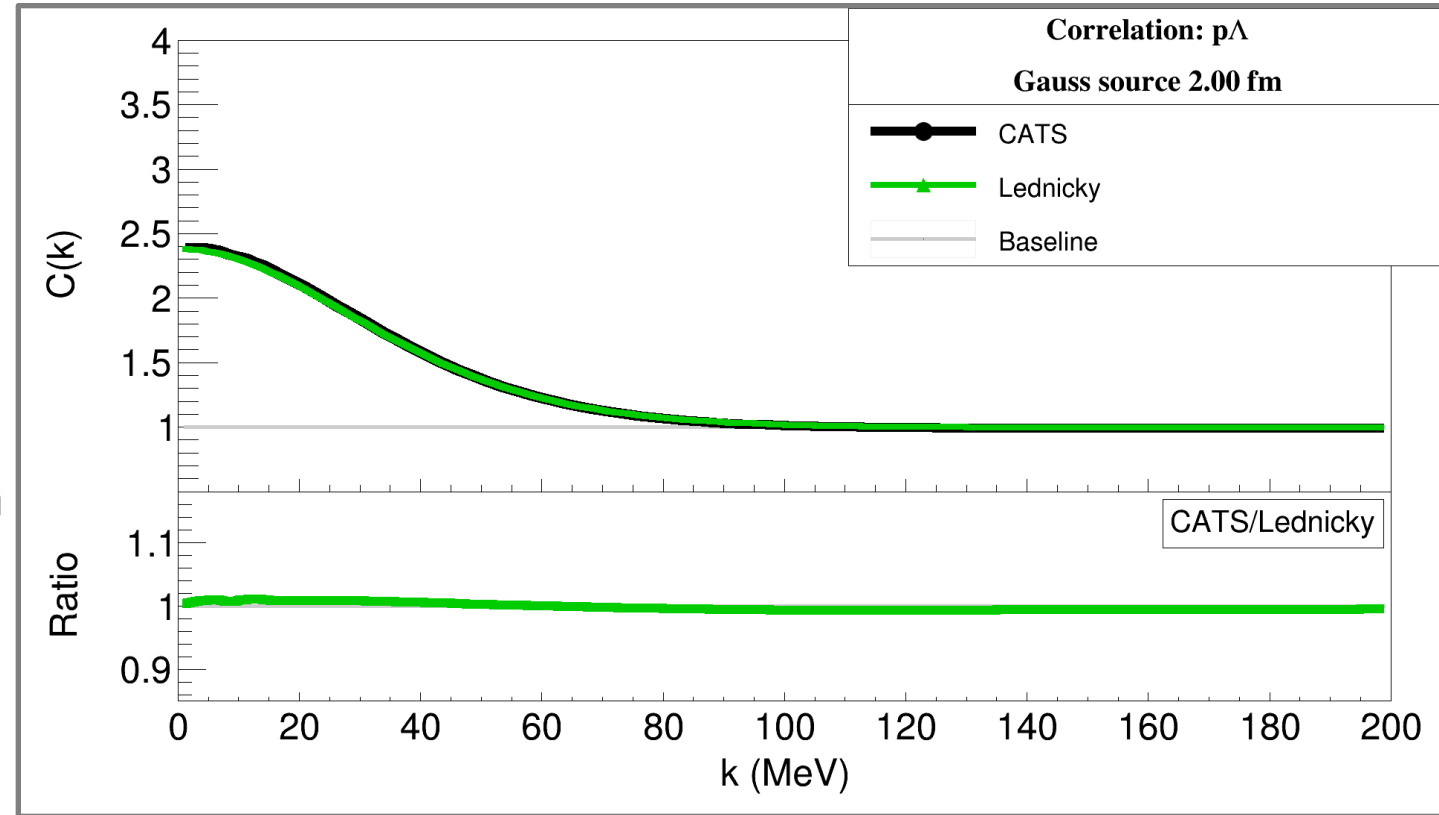
- Using CATS phase shifts  $\delta_l$  can be extracted for any given potential
  - Comparison to global data



- $\Lambda p$  cross sections for any given potential can be extracted
  - Comparison to global data

- **CATS:**
  - Usmani Potential for the interaction
  - Gaussian Source
- **Lednicky:**
  - Scattering parameters and effective range obtained from the Usmani Potential

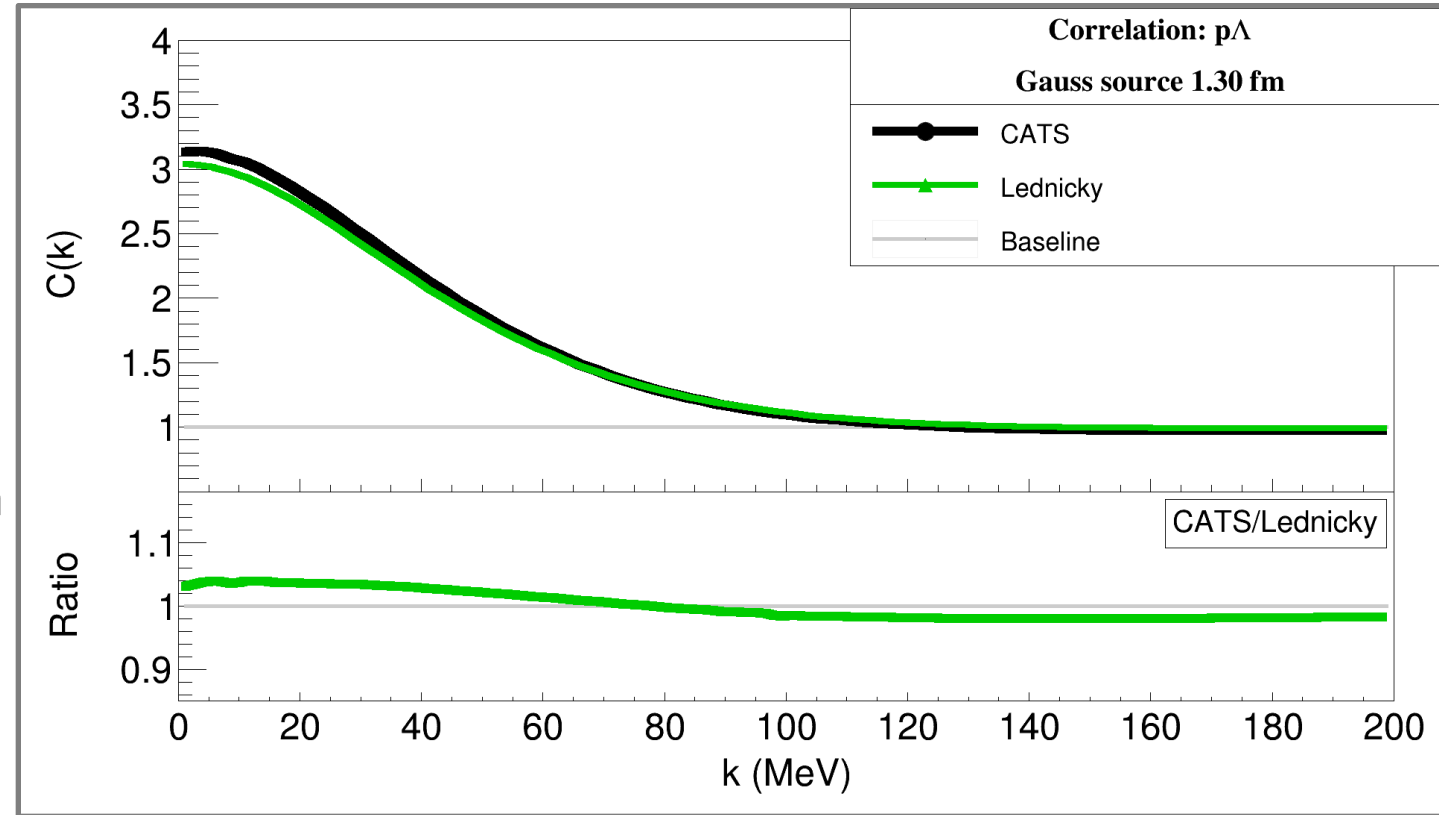
Wang and Pratt, Phys. Rev. Lett. 83, 3138, 1999



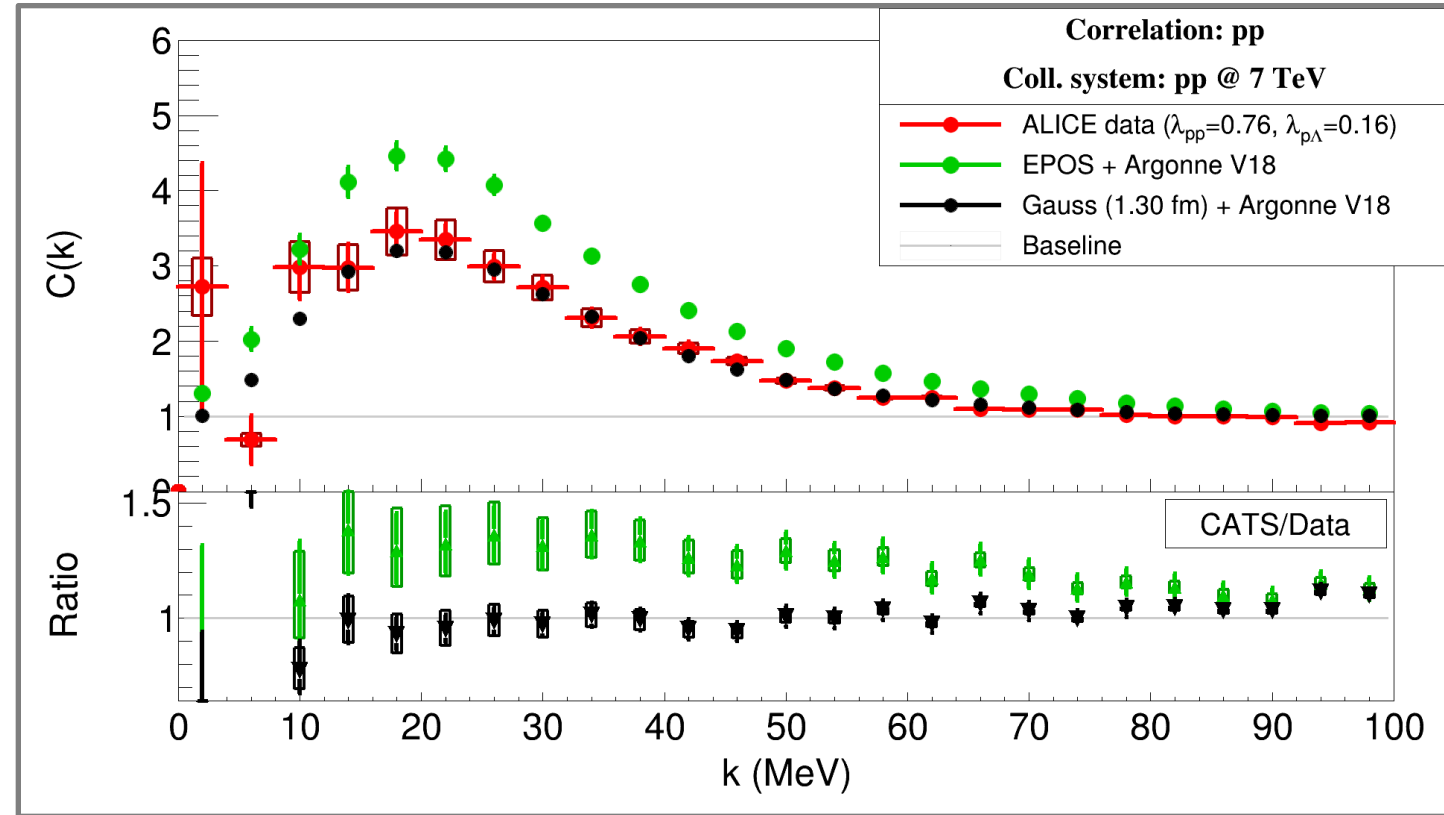


- **CATS:**
  - Usmani Potential for the interaction
  - Gaussian Source
- **Lednicky:**
  - Scattering parameters and effective range obtained from the Usmani Potential

Wang and Pratt, Phys. Rev. Lett. 83, 3138, 1999
- Agreement within 4% for a source size  $R_G = 1.30$  fm
  - Corrections for small source sizes in the Lednicky Model play a role

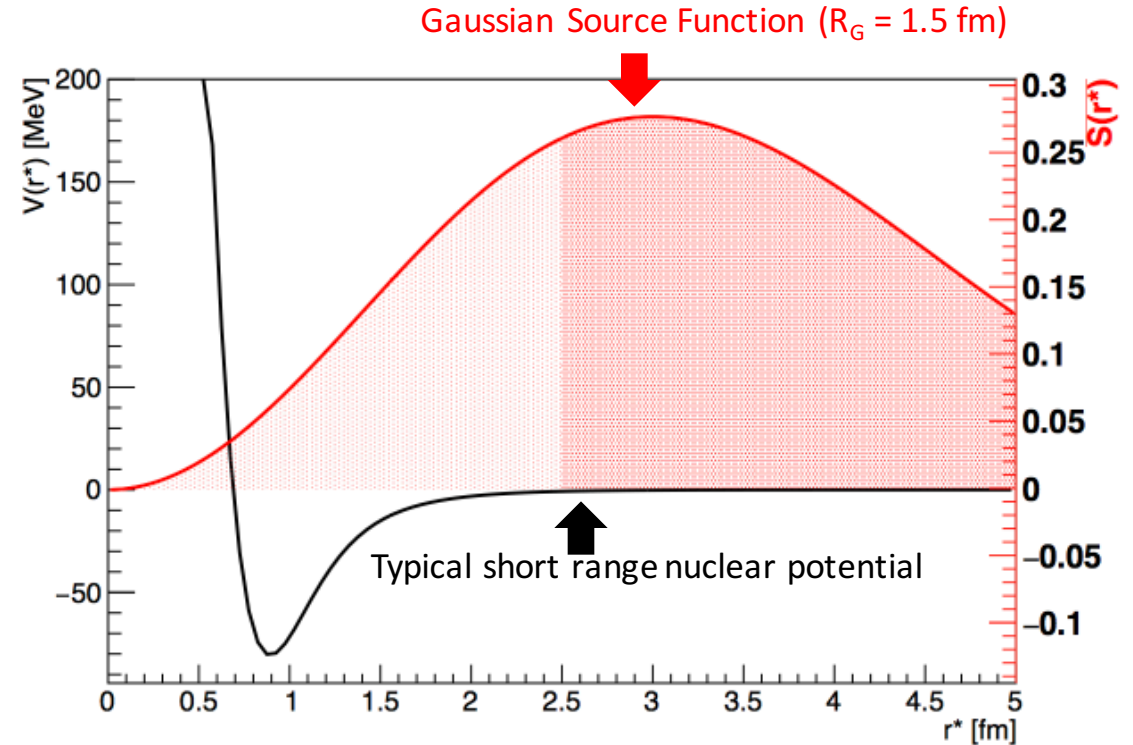


- CATS allows to describe **experimental data** of our ALICE analysis
- Input 1:
  - Argonne V18 Potential
  - Gaussian Source function with  $R_G = 1.30$  fm
- Input 2:
  - Argonne V18 Potential
  - Source Distribution from EPOS
- EPOS not suitable to describe the source



Argonne V18 potential and data: Phys. Rev. C 51, 38

- Reduce degrees of freedom by fitting different collision systems (pp 7 and 13 TeV and pPb at 5 TeV) at different energies at the same time
- **Scattering experiments:** intermediate region outside the core of the potential
- **Femtoscopy:** pairs are also produced in the region inside the core
  - Access to the shape of the potential?



- Femtoscopy in small systems is feasible
- New method to calculate different contributions to the total correlation function based on single particle properties
- Modelling of the correlation function with CATS
- Analysis of Run 2 Data in p-p at 13 TeV and p-Pb Collisions at 5 TeV ongoing
  - Additionally obtain the  $\Sigma$  and  $\Xi$  Correlation Function
- Universal and Robust Femto Analysis Tool
  - Fit the correlation function of various systems simultaneously in combination with CATS

**UrFAT**

Universal &  
robust  
Femtoscopy  
Analysis  
Tool





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# Thank you for your attention!

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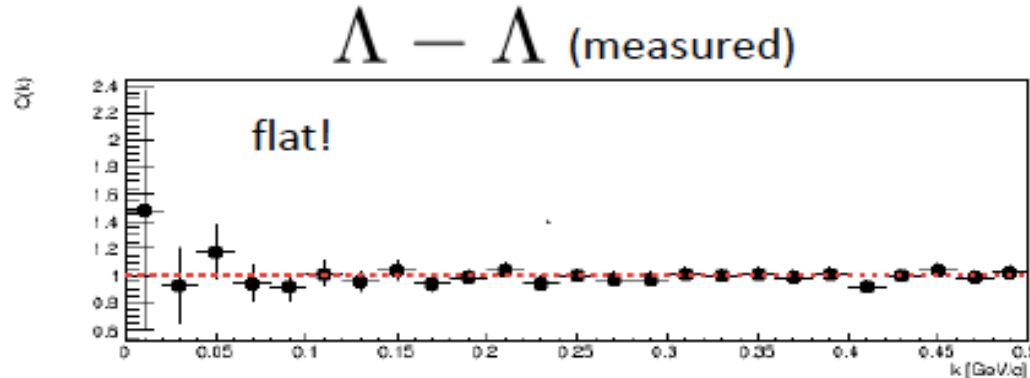
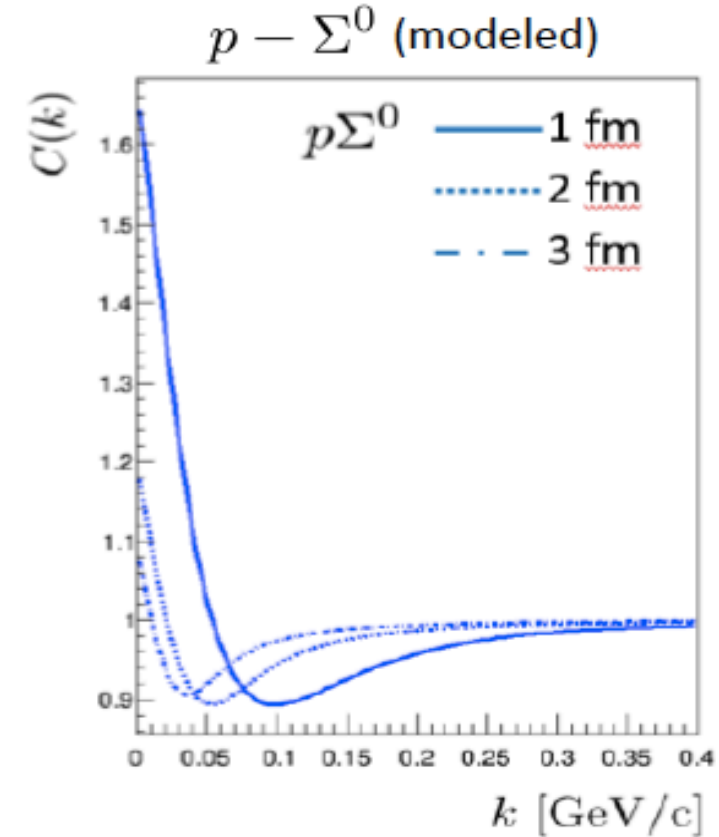
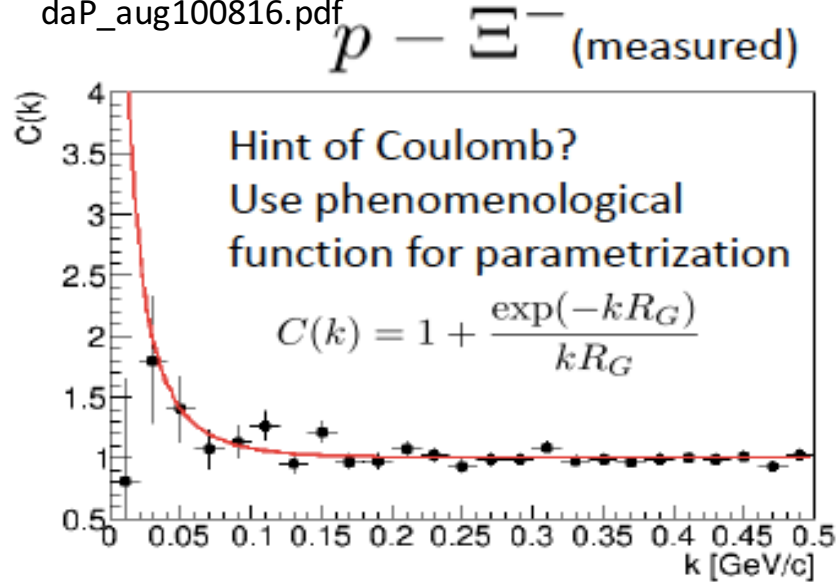


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Study of the theoretical prediction for the  $p\Sigma^0$  correlation in

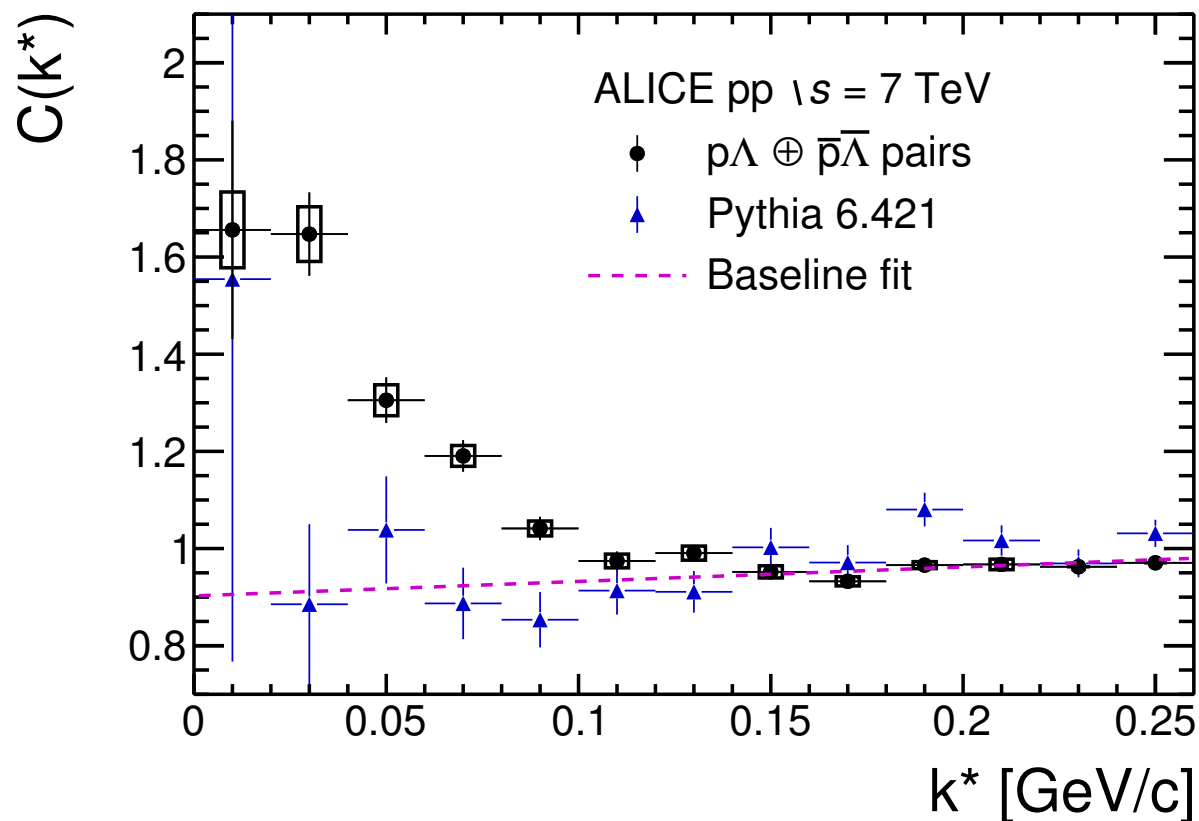
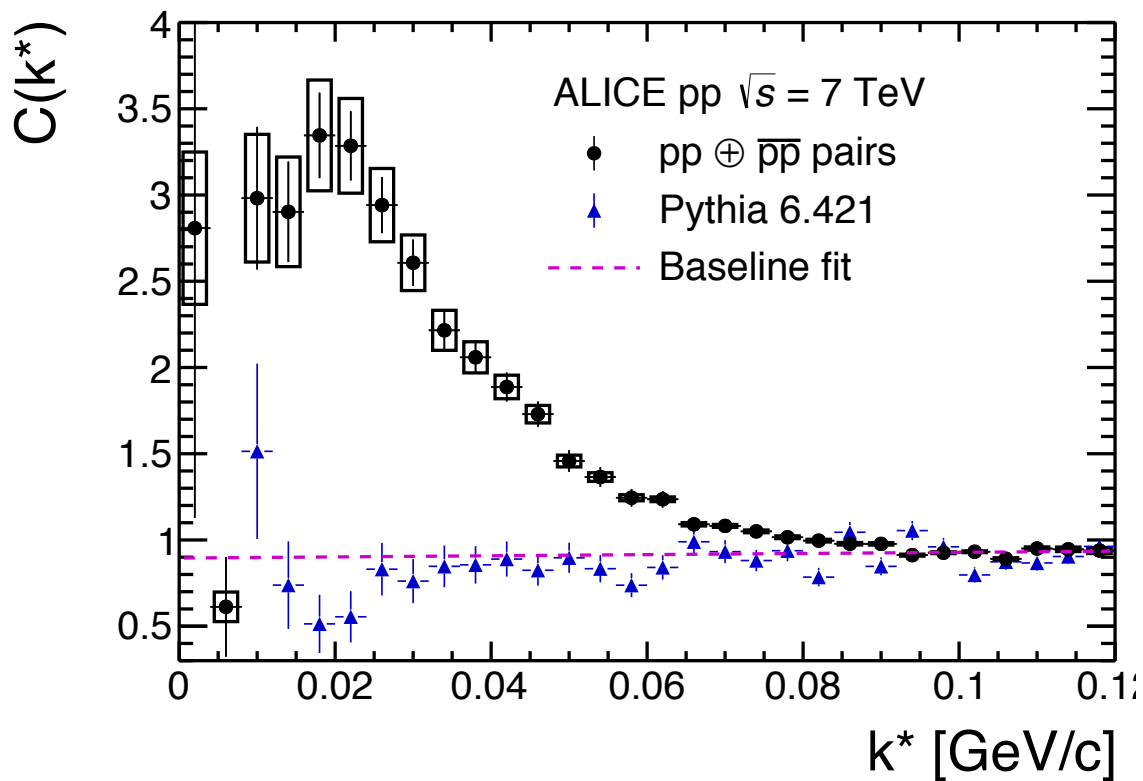
[https://indico.cern.ch/event/562339/contributions/2271695/attachments/1322381/1983690/oarnold\\_Update\\_LambdaP\\_aug100816.pdf](https://indico.cern.ch/event/562339/contributions/2271695/attachments/1322381/1983690/oarnold_Update_LambdaP_aug100816.pdf)



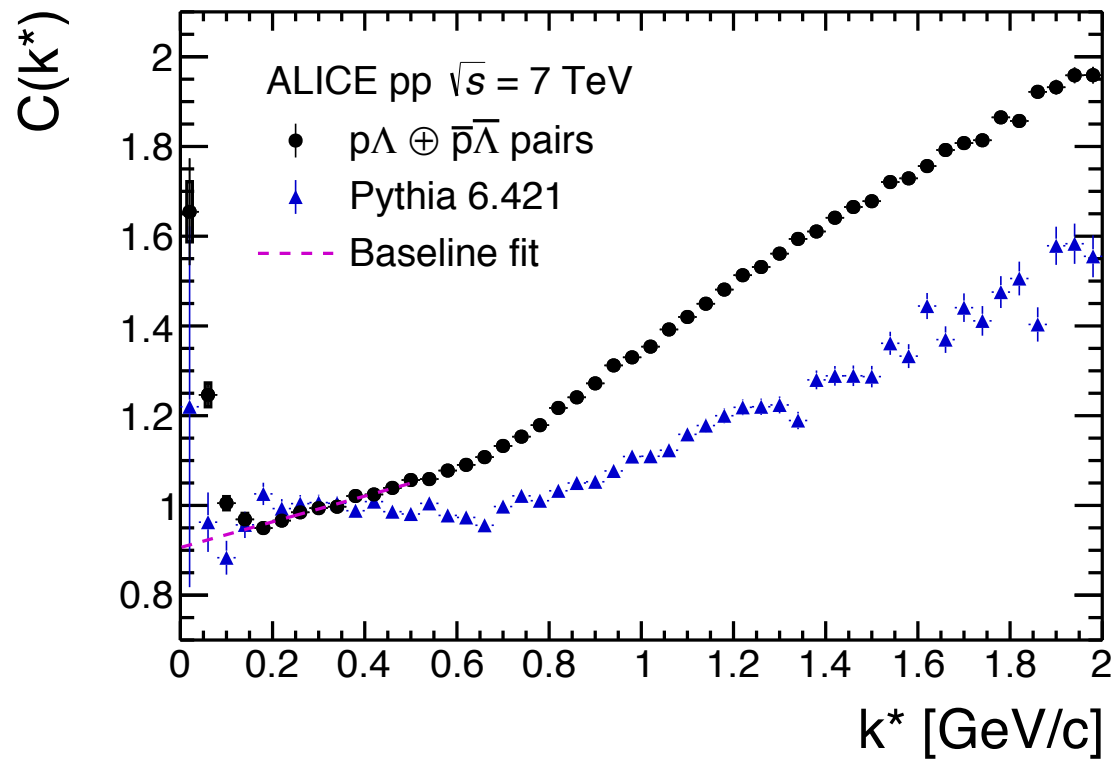
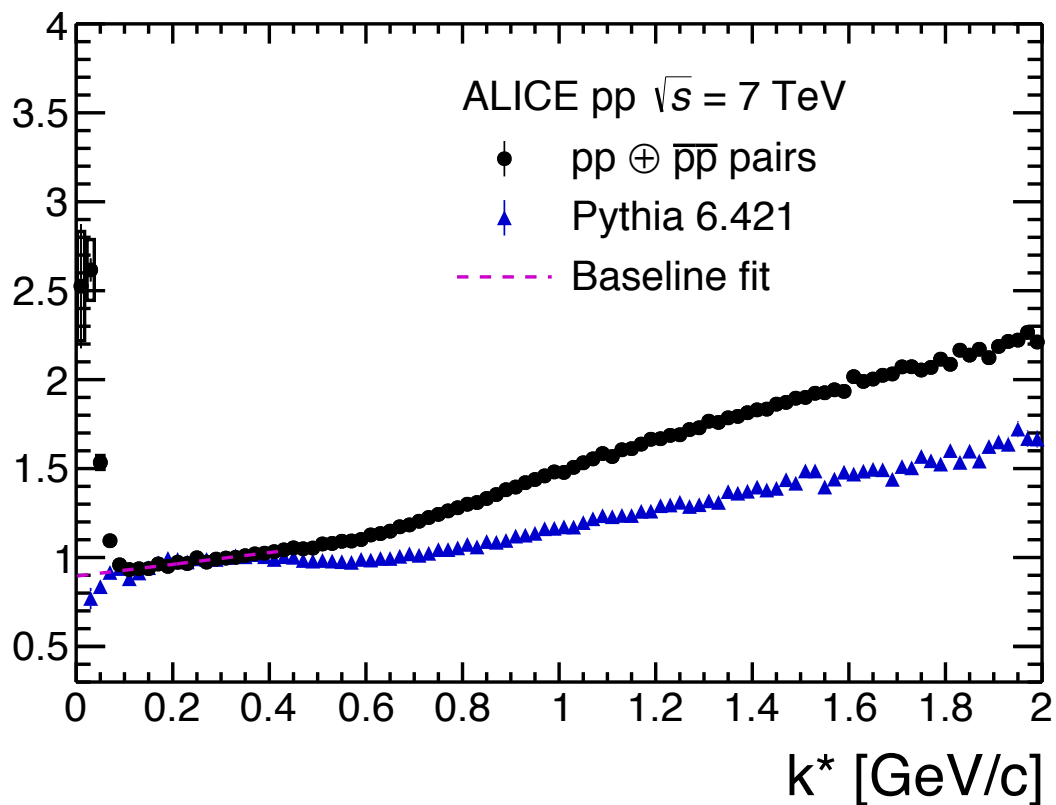
arXiv:0704.3290

<http://inspirehep.net/record/I283372>

4

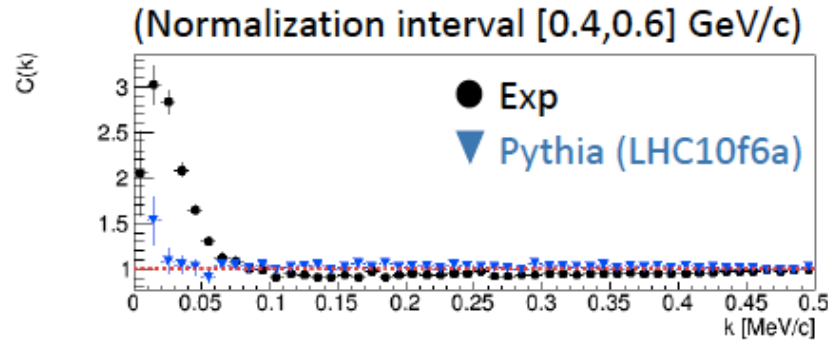




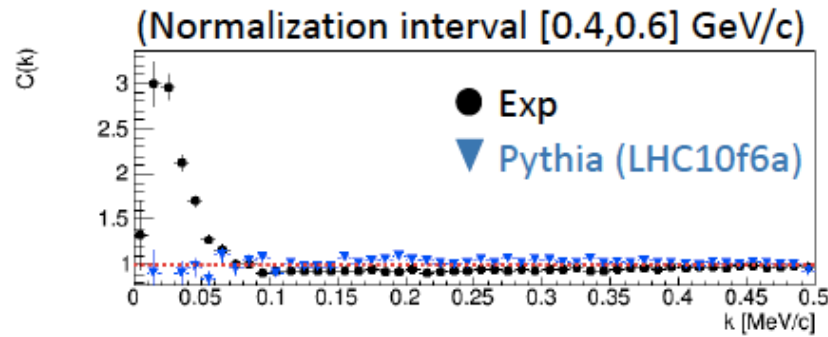


# Minijets Background

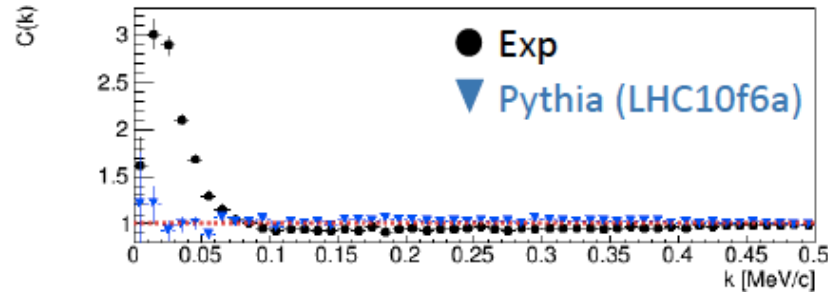
Correlations:



$pp$

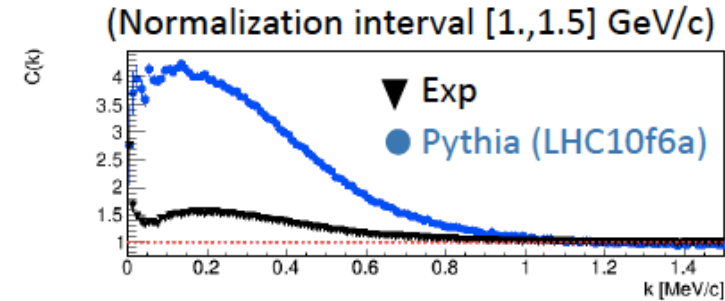


$p\bar{p}$

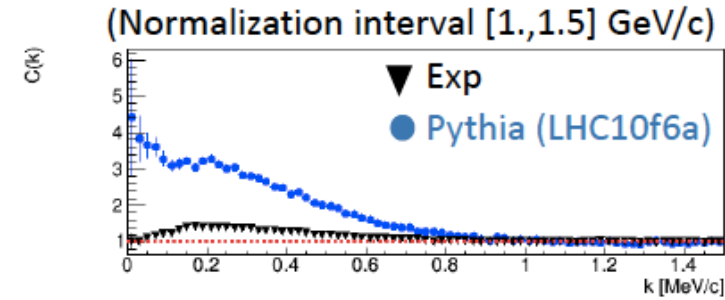


$pp + p\bar{p}$

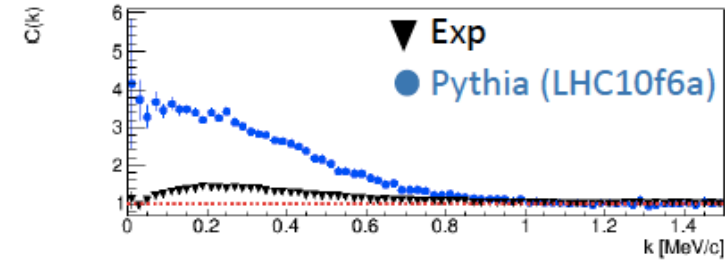
$n$  Correlations:



$p\bar{p}\bar{p}$



$p\bar{\Lambda}$



$p\bar{\Lambda}$

