production of loosely bound states at the QCD phase boundary – 'snowballs in hell'

- introduction and perspective
- the hadron resonance gas
- (u,d,s) hadron production, Lattice QCD and the QCD phase structure
- loosely bound states
- comments on coalescence models

EMMI workshop on
Anti-Matter, Hypermatter and Exotica

Torino, Italy
Nov. 6 -10, 2017
pbm

ISOQUANT
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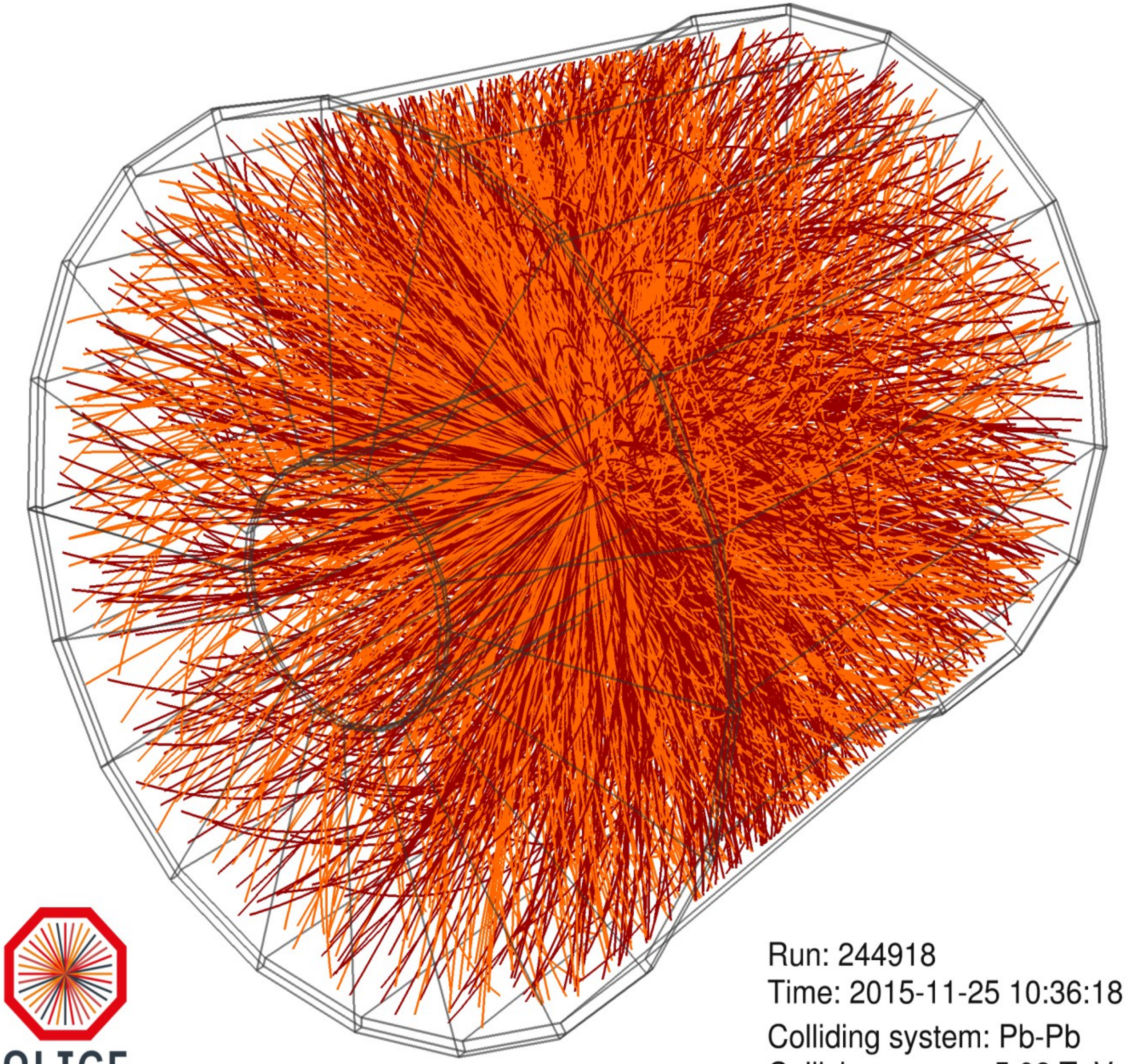
GSI

phenomenology results obtained in collaboration with
Anton Andronic,
Krzysztof Redlich, and Johanna Stachel
for a recent review see
Andronic, pbm, Redlich, Stachel, arXiv :1710.09425

first PbPb collisions at LHC at √s = 5.02 A TeV



ALICE

Run1: 3 data taking campaigns
pp, pPb, Pb—Pb
> 145 publications

Run2

2015: 13 TeV pp
Pb—Pb run
in November 2015

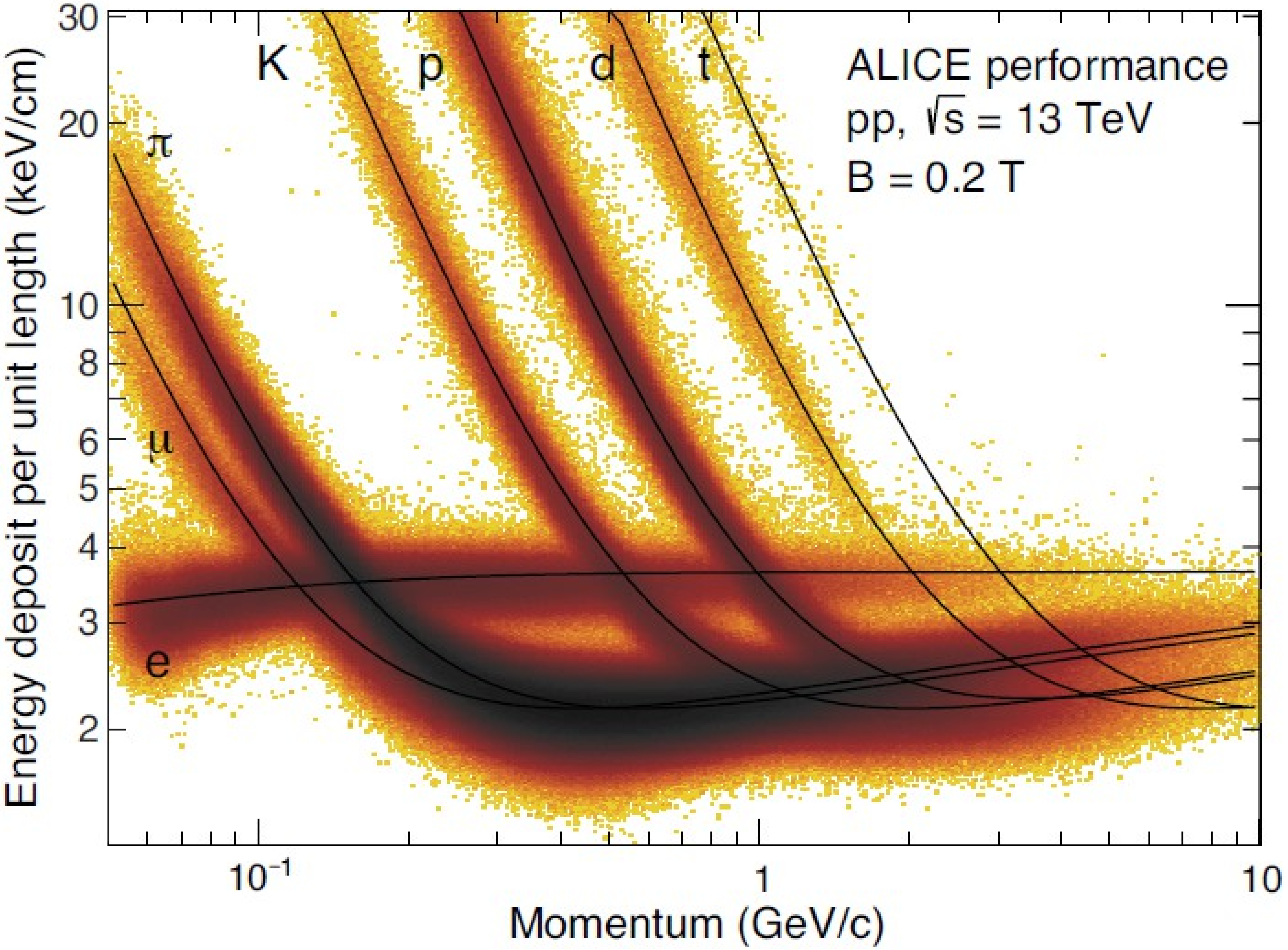
2016: 13 TeV pp
+ pPb 5 TeV and 8 TeV

2017: pp running at 13 and 5 TeV

2018: pp + Pb—Pb running

particle identification with the ALICE TPC

from 50 MeV to 50 GeV



hadron production and the QCD phase boundary

part 1: the hadron resonance gas

duality between hadrons and quarks/gluons (I)

Z: full QCD partition function

all thermodynamic quantities derive from QCD partition functions

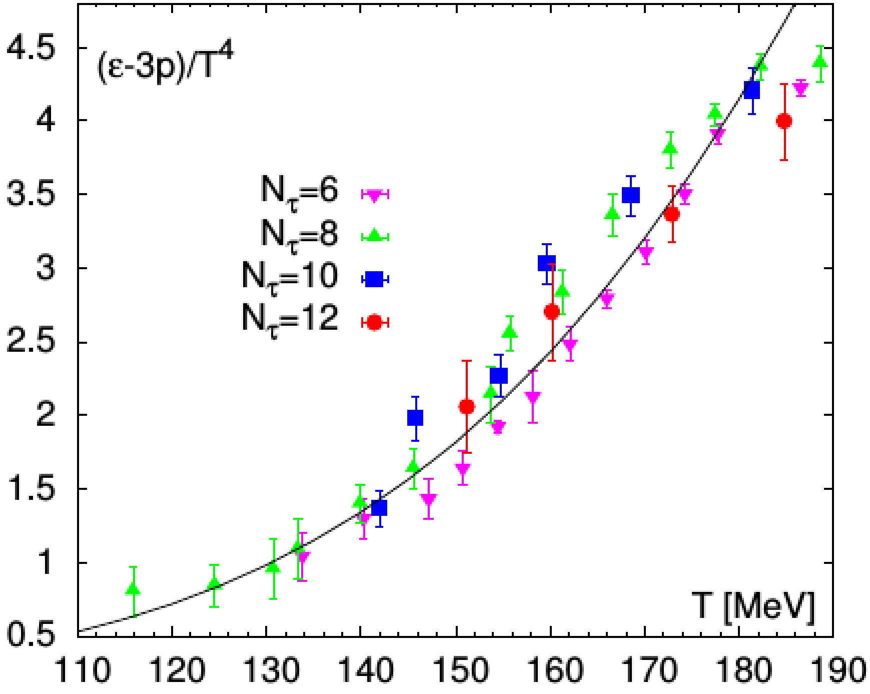
for the pressure we get:

$$\frac{p}{T^4} = \frac{1}{T^3}\frac{\partial \ln Z(V,T,\mu)}{\partial V}$$

comparison of trace anomaly from LQCD
Phys.Rev. D90 (2014) 094503
HOTQCD coll.

with hadron resonance gas prediction
(solid line)

LQCD: full dynamical quarks with realistic pion mass

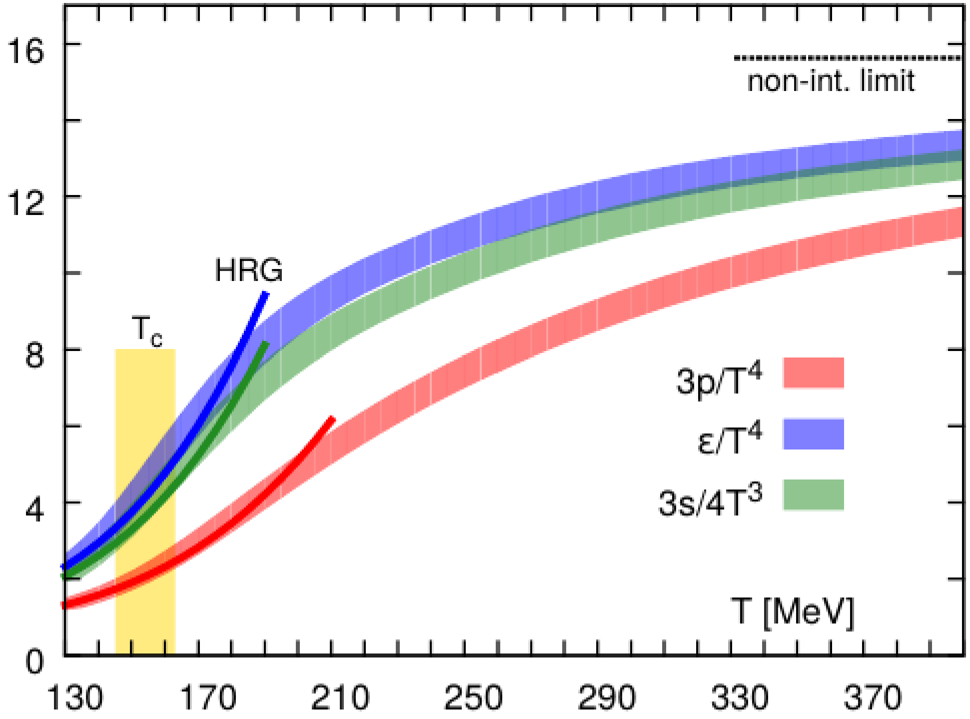


duality between hadrons and quarks/gluons (II)

comparison of equation of state from LQCD
Phys.Rev. D90 (2014) 094503
HOTQCD coll.

with hadron resonance gas predictions (colored lines)

essentially the same results also from Wuppertal-Budapest coll.
Phys.Lett. B730 (2014) 99-104



pseudo-critical temperature

$$T_c = (154 \pm 9)\ \mathrm{MeV}$$

$$\varepsilon_{crit} = (340 \pm 45)\ \mathrm{MeV/fm}^3$$

$$\varepsilon_{nucl} = 450\ \mathrm{MeV/fm}^3$$

duality between hadrons and quarks/gluons (III)

in the dilute limit T < 165 MeV:

$$\ln Z(T,V,\mu) \approx \sum_{i\in\ mesons} \ln \mathcal{Z}^{M}_{M_i}(T,V,\mu_Q,\mu_S) + \sum_{i\in\ baryons} \ln \mathcal{Z}^{B}_{M_i}(T,V,\mu_b,\mu_Q,\mu_S)$$

where the partition function of the hadron resonance model is expressed in mesonic and baryonic components. The chemical potential μ reflects then the baryonic, charge, and strangeness components $\mu = (\mu_b, \mu_Q, \mu_S)$.

hadron production and the QCD phase boundary

part 2: analysis with the statistical hadronization model

statistical hadronization model of particle production and QCD

partition function Z(T,V) contains sum over the full hadronic mass spectrum and is fully calculable in QCD

for each particle i, the statistical operator is:

$$\ln Z_i = \frac{V g_i}{2\pi^2} \int_0^\infty \pm p^2 \mathrm{d}p \ln[1 \pm \exp(-(E_i - \mu_i)/T)]$$

particle densities are then calculated according to:

$$n_i = N_i/V = -\frac{T}{V}\frac{\partial \ln Z_i}{\partial \mu} = \frac{g_i}{2\pi^2} \int_0^\infty \frac{p^2 \mathrm{d}p}{\exp[(E_i - \mu_i)/T] \pm 1}$$

from analysis of all available nuclear collision data we now know the energy dependence of the parameters T, mu_b, and V over an energy range from threshold to LHC energy and can confidently extrapolate to even higher energies

in practice, we use the full experimental hadronic mass spectrum from the PDG compilation (vacuum masses) to compute the 'primordial yield'

comparison with measured hadron yields needs evaluation of all strong decays

implementation

$$n_i = N_i/V = -\frac{T}{V}\frac{\partial \ln Z_i}{\partial \mu} = \frac{g_i}{2\pi^2}\int_0^\infty \frac{p^2 \mathrm{d}p}{\exp[(E_i-\mu_i)/T]\pm 1}$$

Latest PDG hadron mass spectrum ...quasi-complete up to m=2 GeV;
our code: 555 species (including fragments, charm and bottom hadrons)

for resonances, the width is considered in calculations

Minimize: $\chi^2 = \sum_i \frac{(N_i^{exp}-N_i^{therm})^2}{\sigma_i^2}$

N_i hadron yield, σ_i experimental uncertainty (stat.+syst.)

$\Rightarrow (T,\, \mu_B,\, V)$

canonical treatment whenever needed (small abundances)

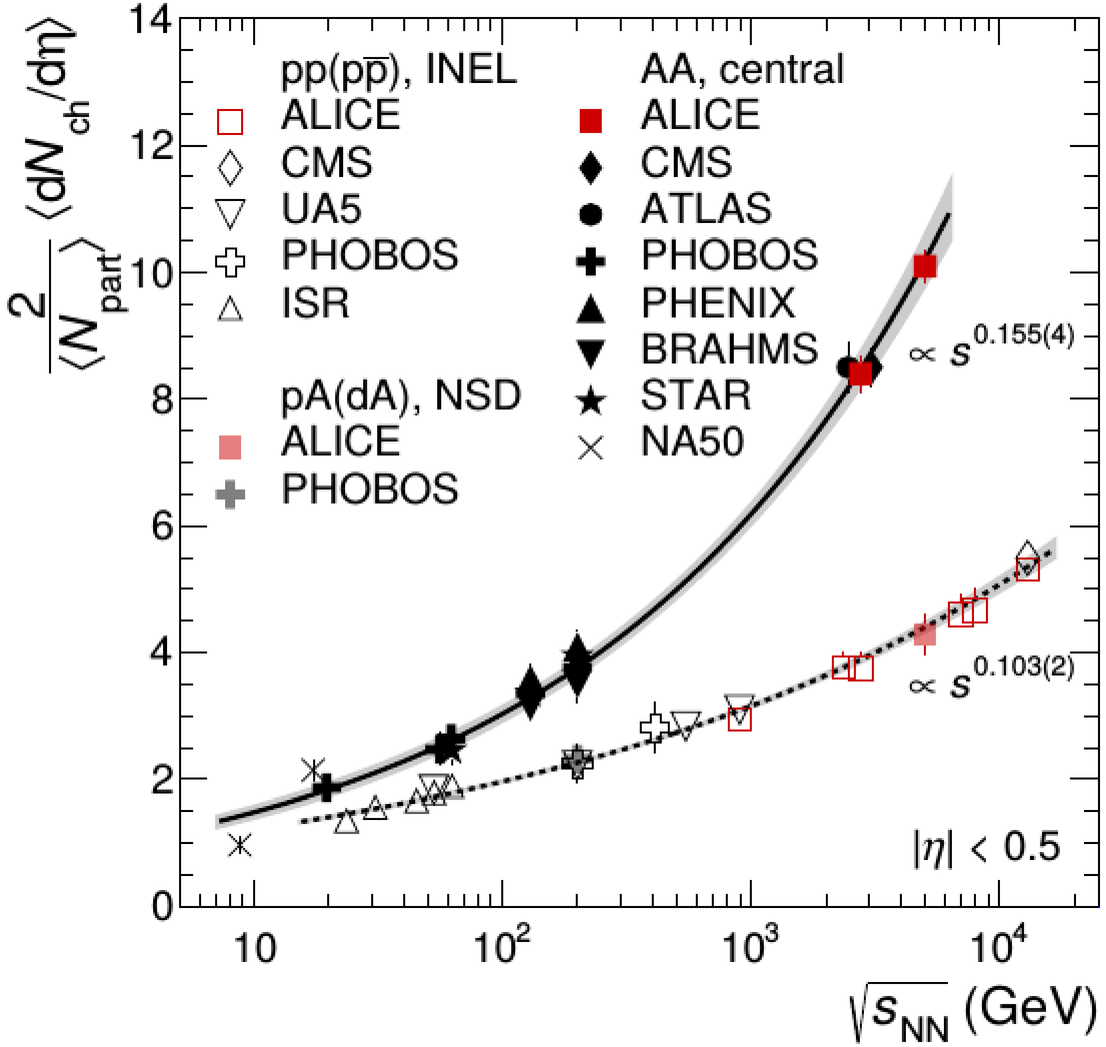
energy dependence of hadron production in central Pb-Pb (Au-Au) collisions

total number of hadrons produced

2.76 TeV $N_{had} = 25800$

5.02 TeV $N_{had} = 32300$

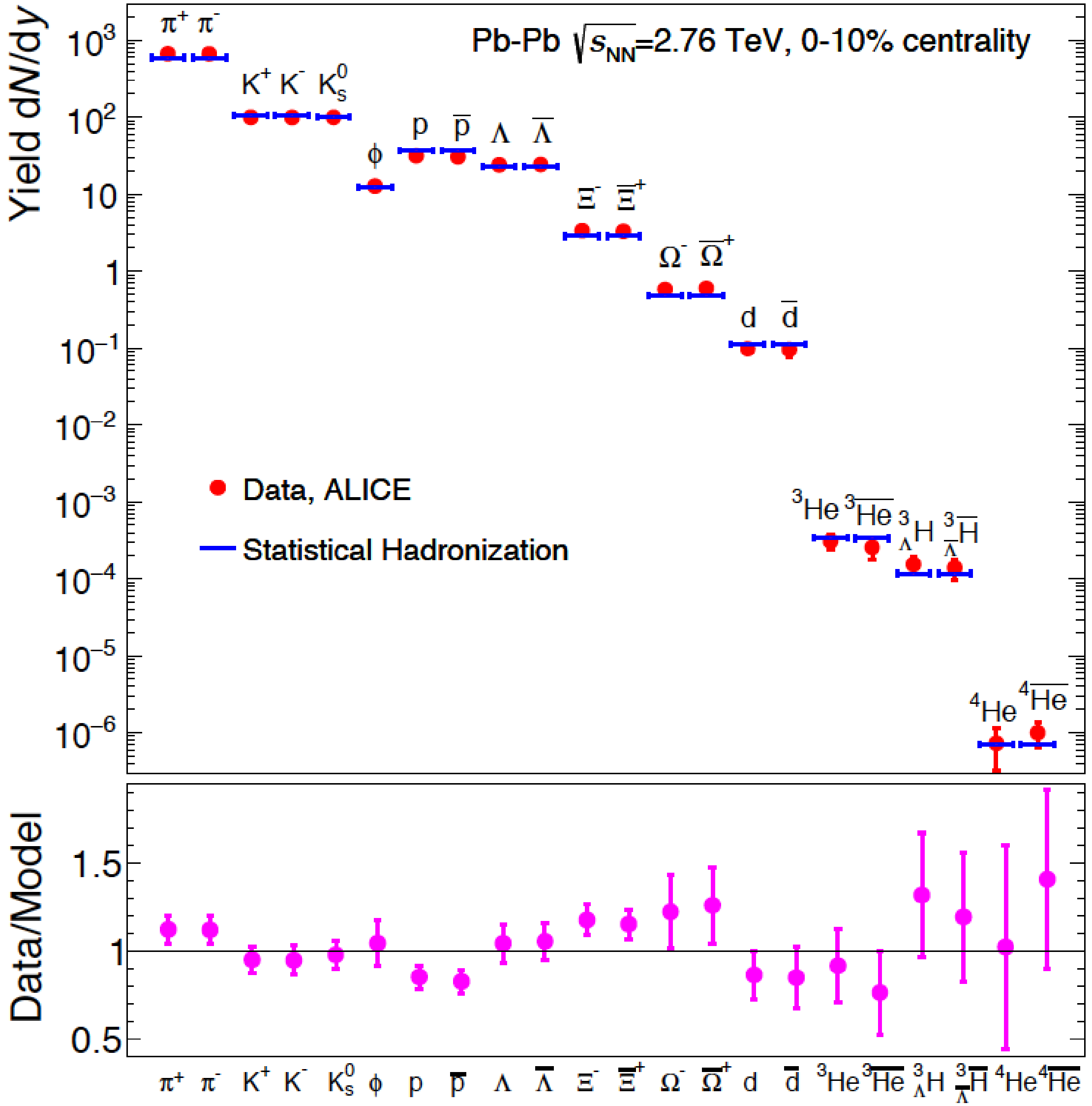
fireball with 'macroscopic' number of produced particles



data from LHC run1 and run2

ALICE coll., Phys.Rev.Lett. 116 (2016) no.22, 222302

July 2017 update: excellent description of ALICE@LHC data



fit includes loosely bound systems such a deuteron and hypertriton

hypertriton is bound-state of (Λ,p,n), Λ separation energy about 130 keV

size about 10 fm, the **ultimate halo nucleus,**

produced at T=156.5 MeV. close to an Efimov state

proton discrepancy 2.8 sigma

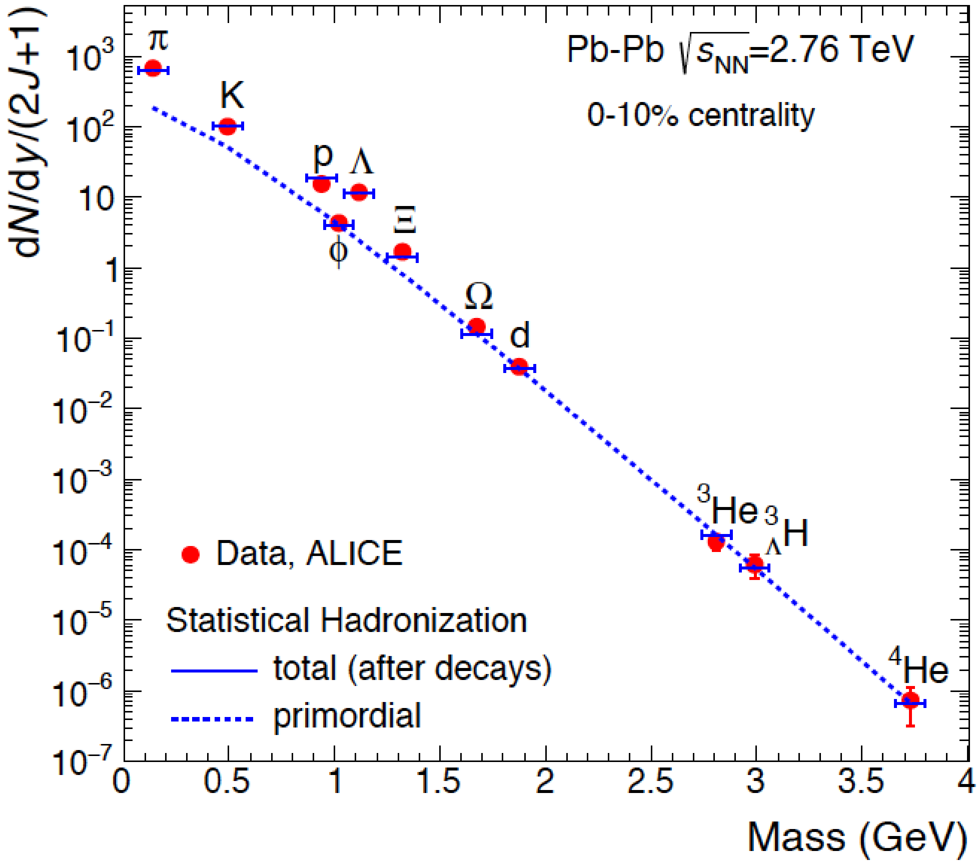
Xi discrepancy?

Andronic, pbm, Redlich, Stachel,
arXiv :1710.09425

J. Stachel, A. Andronic, P. Braun-Munzinger and K. Redlich, Confronting LHC data with the statistical hadronization model, J.Phys.Conf.Ser.**509** (2014) 012019, arXiv:1311.4662 [nucl-th].

excellent agreement over 9 orders of magnitude

agreement over 9 orders of magnitude with QCD statistical operator prediction



exponential decrease with mass and common temperature T = 159 MeV
of yields for light nuclei predicted from the thermal phenomenology discussed above

production near the phase boundary

yield of light nuclei predicted in: pbm, J. Stachel, J.Phys. G28 (2002) 1971-1976,
J.Phys. G21 (1995) L17-L20

a note on the chemical freeze-out temperature

T_{chem} = 156.5 ± 1.5 MeV from fit to all particles

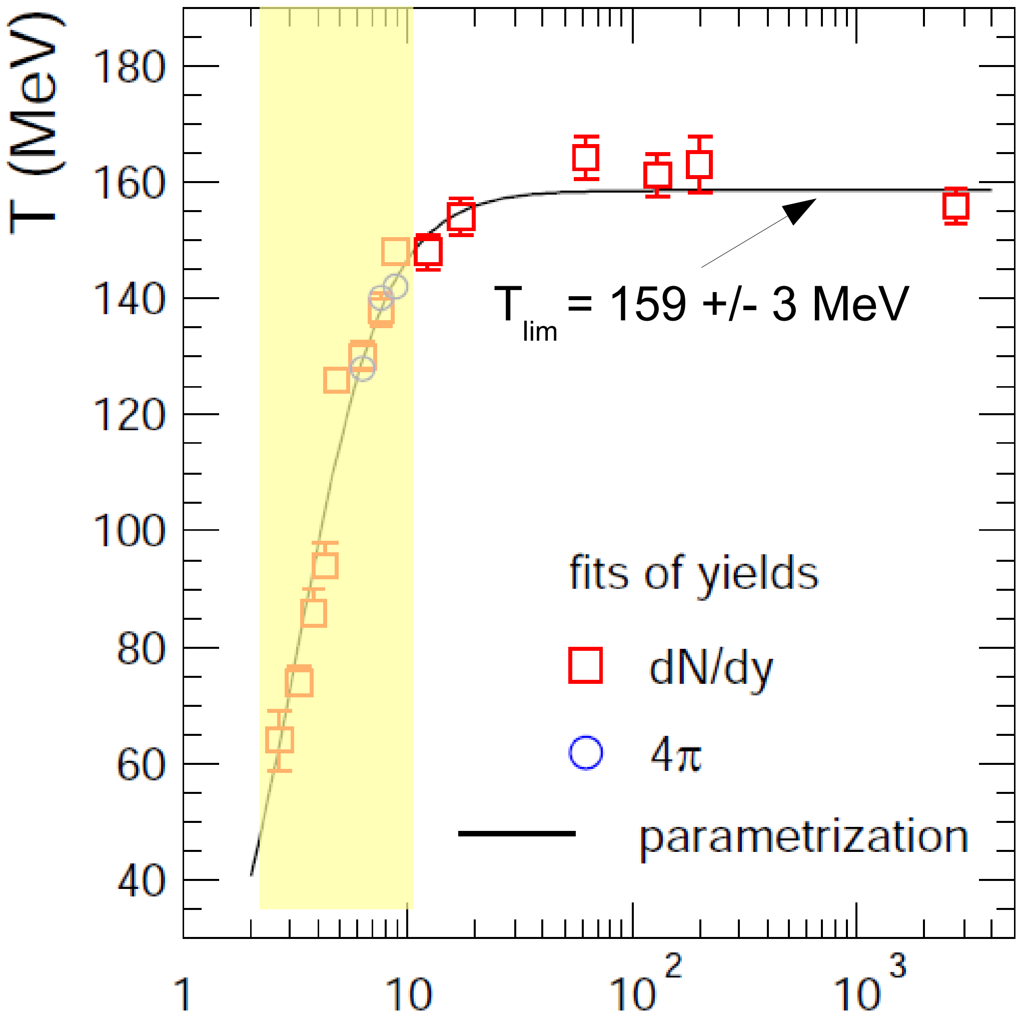
there is an additional uncertainty because of the poorly known hadronic mass spectrum for masses > 2 GeV

for d, 3He, hypertriton and alpha, there is very little feeding from heavier states and none from high mass states in the hadronic mass spectrum, for these particles the temperature T_{nuc} can be determined 'on the back of an envelope' :

T_{nuc} = 159 ± 5 MeV, independent of hadronic mass spectrum

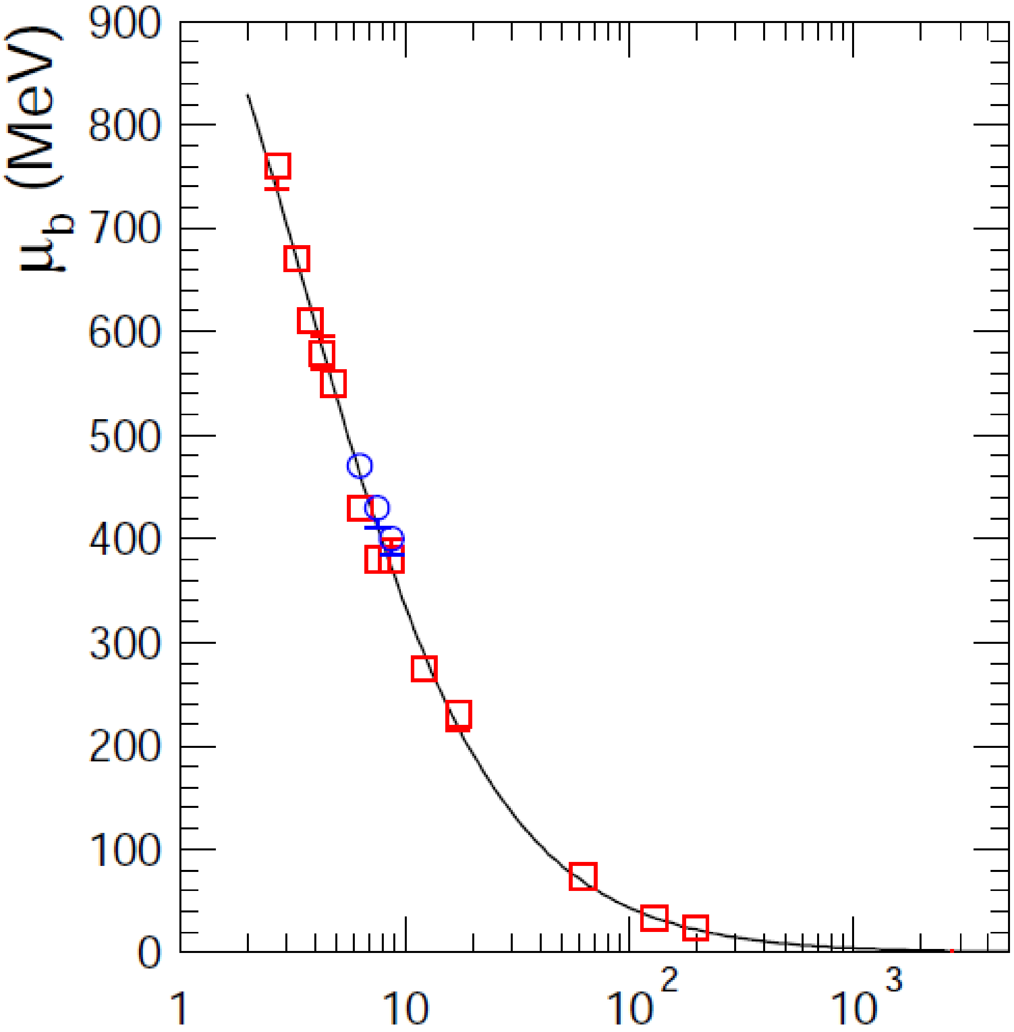
energy dependence of temperature and baryo-chemical potential

energy range from LHC down to threshold (FAIR)



T_{lim} = 159 +/- 3 MeV is maximum hadronic temperature

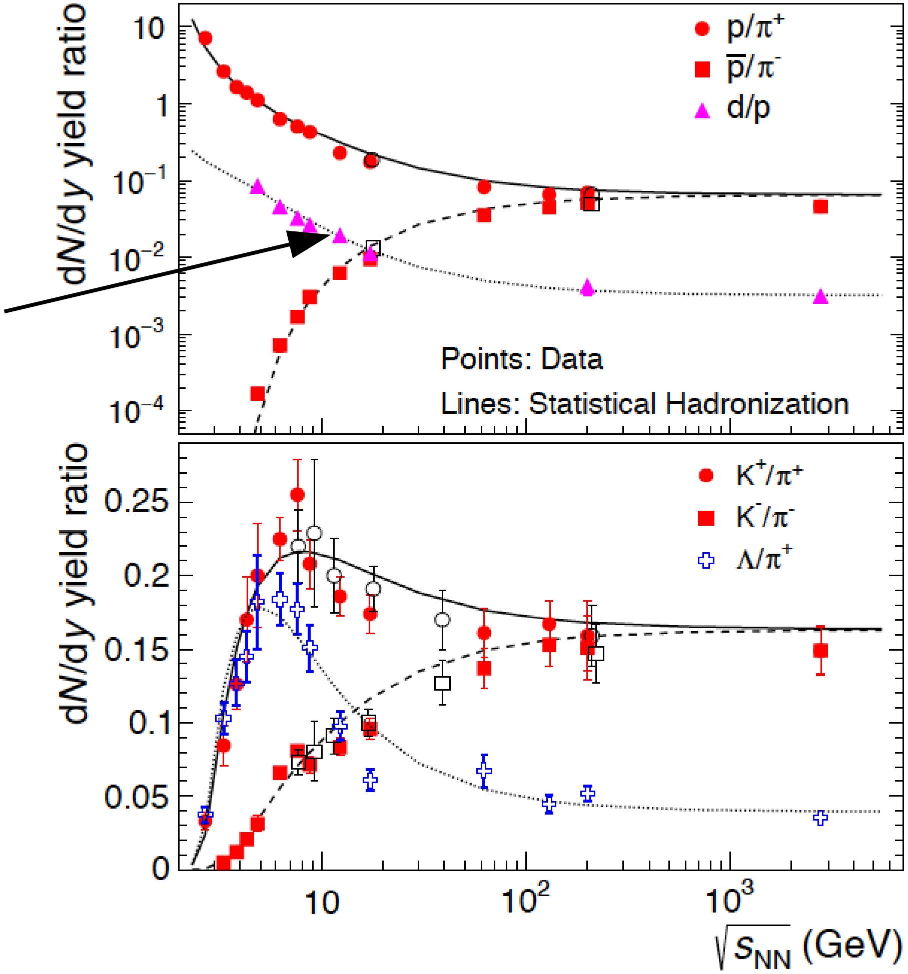
is phase boundary ever reached for $\sqrt{s_{NN}}$ < 10 GeV?



T_c = 154 +/- 9 MeV from lattice

energy dependence of hadron production described quantitatively

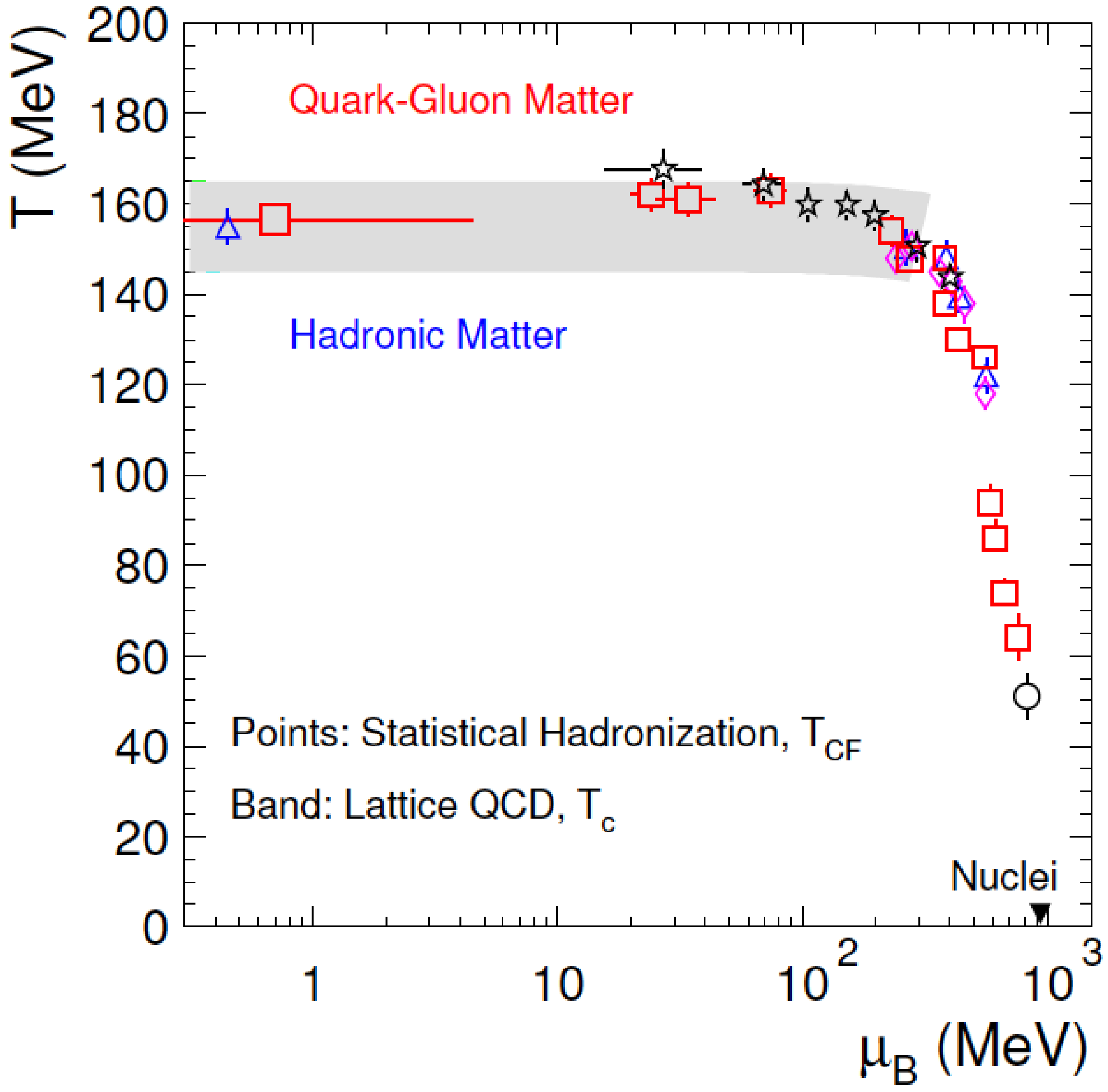
energy dependence of d/p ratio quantitatively described, no new parameters



together with known energy dependence of charged hadron production in Pb-Pb collisions we can predict yield of all hadrons at all energies with < 10% accuracy

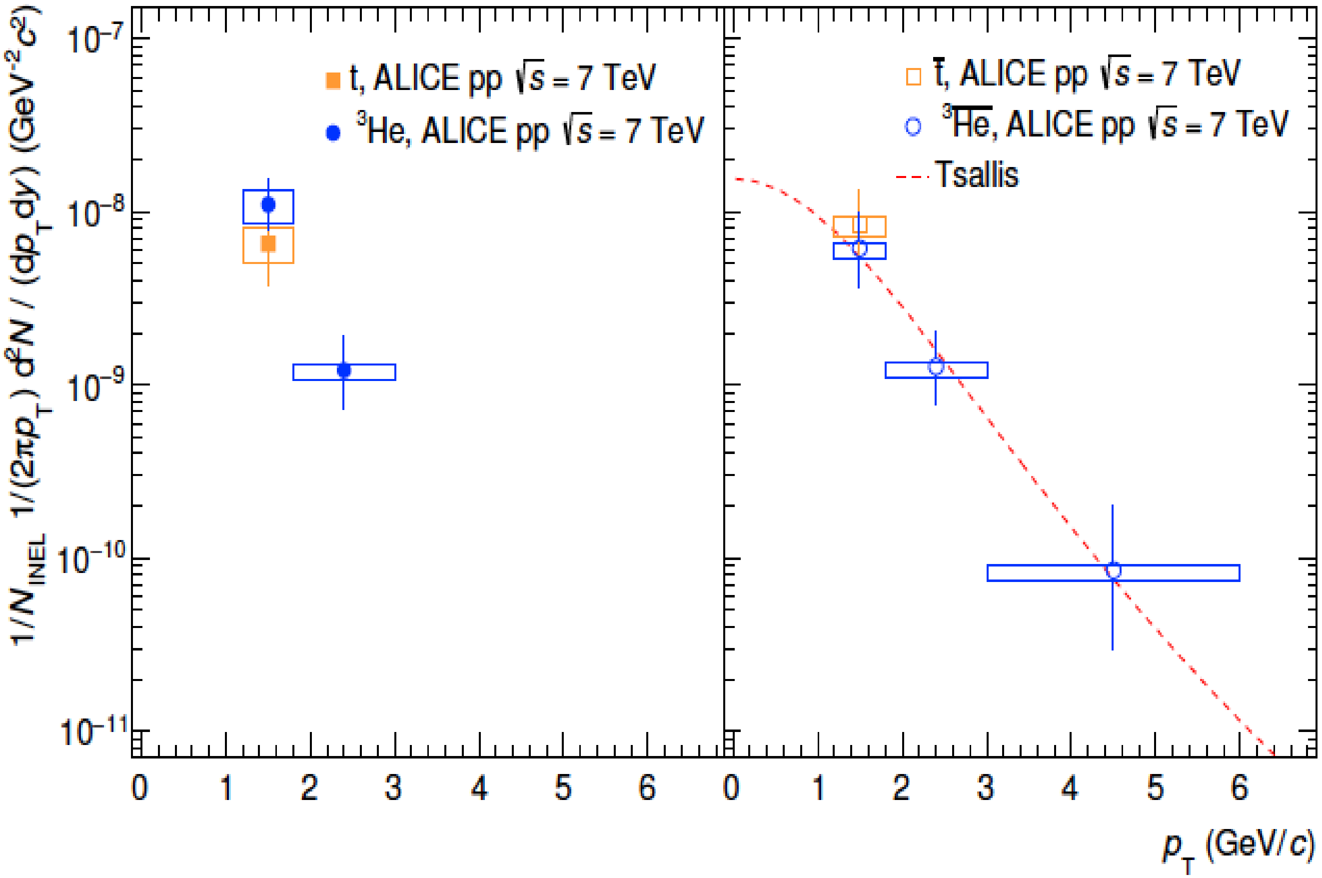
no new physics needed to describe K+/pi+ ratio including the 'horn'

the QGP phase diagram, LQCD, and hadron production data



quantitative agreement of chemical freeze-out parameters with LQCD predictions for baryo-chemical potential < 300 MeV

the LHC is a 'gluon collider' – isospin plays no role in particle production



3He = t, p=n, and anti-particles

Systematic uncertainties in statistical hadronization model

in general, not easy to estimate

from analysis of uncertainties in mass spectrum, and in branching ratios, and considering the Boltzmann suppression, we get:

$\Delta T \leq 5$ MeV at $\mu_b = 0$ and $T = 156$ MeV

now loosely bound objects

The Hypertriton

mass = 2990 MeV, binding energy = 2.3 MeV

Lambda sep. energy = 0.13 MeV

molecular structure: (p+n) + Lambda

2-body threshold: (p+p+n) + pi- = ^{3}He + pi-

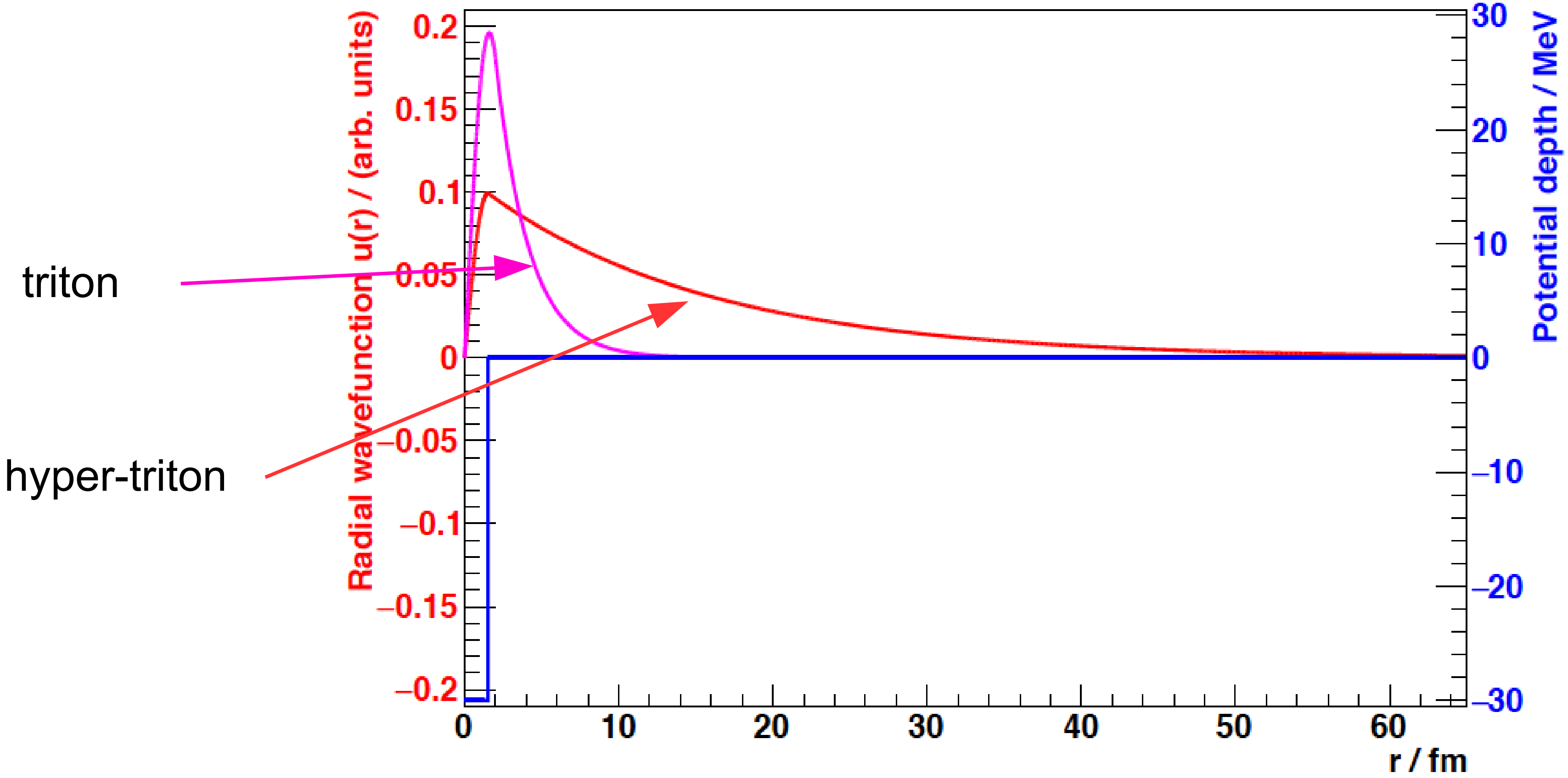
rms radius = $(4\ \mathrm{B.E.}\ M_{\mathrm{red}})^{-1/2}$ = 10.3 fm =
rms separation between d and Lambda

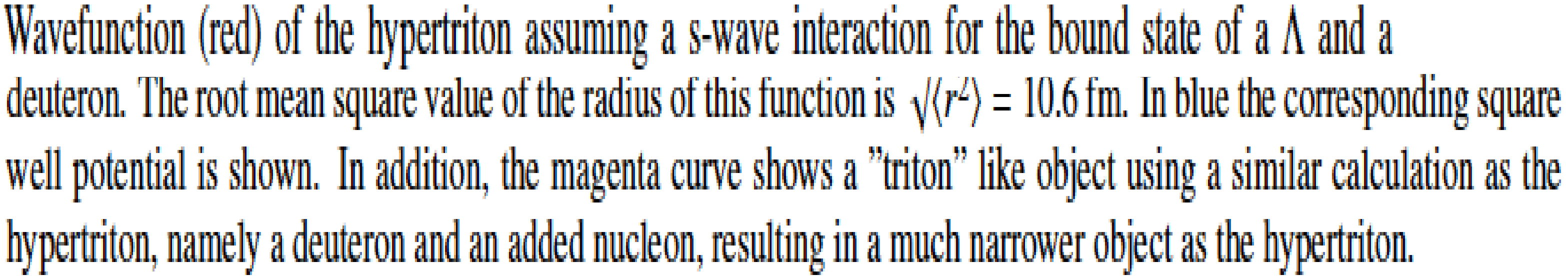
in that sense: hypertriton = (p n Lambda) =
(d Lambda) is the ultimate halo state

yet production yield is fixed at 156 MeV temperature
(about 1000 x separation energy.)

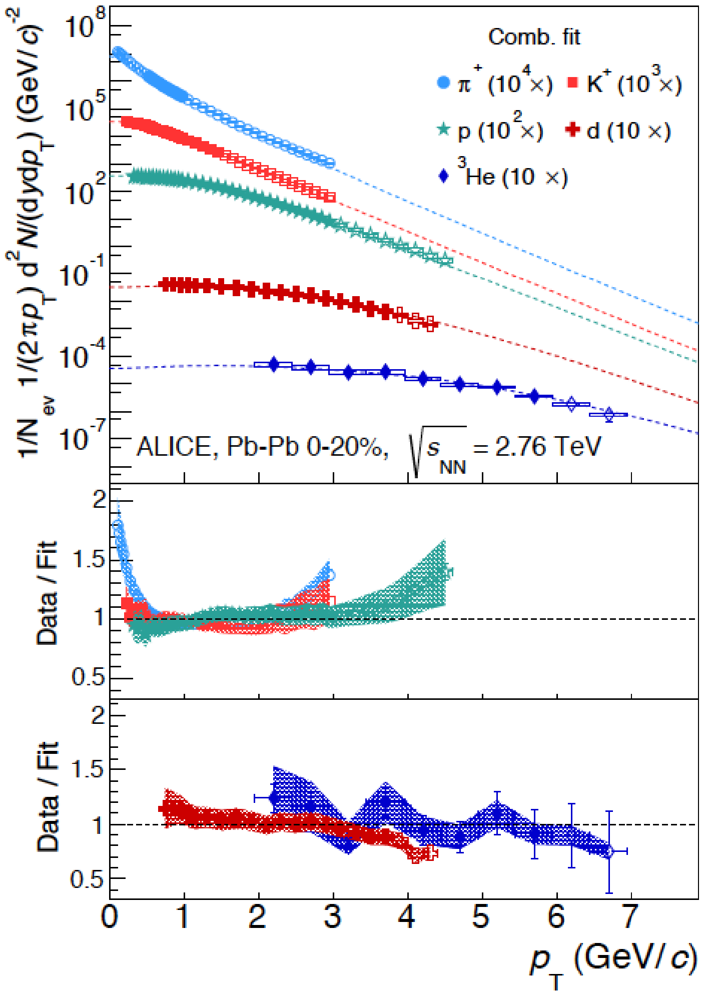
wave function of the hyper-triton – schematic picture

figure by Benjamin Doenigus, August 2017

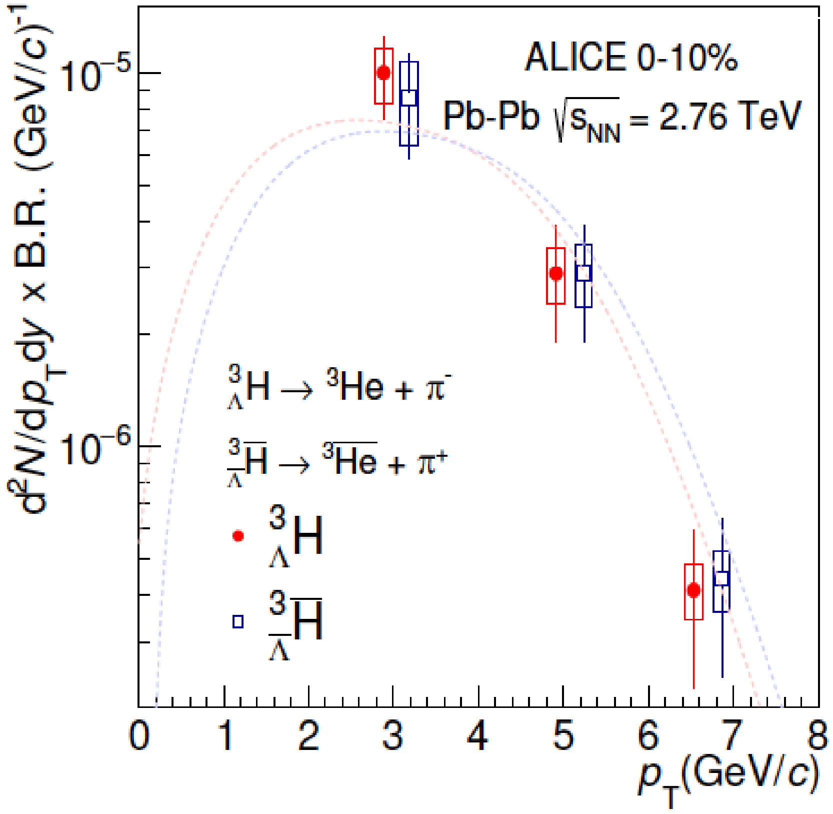



Wavefunction (red) of the hypertriton assuming a s-wave interaction for the bound state of a Λ and a deuteron. The root mean square value of the radius of this function is $\sqrt{\langle r^2 \rangle} = 10.6$ fm. In blue the corresponding square well potential is shown. In addition, the magenta curve shows a "triton" like object using a similar calculation as the hypertriton, namely a deuteron and an added nucleon, resulting in a much narrower object as the hypertriton.

light nuclei flow with same fluid velocity as pions, kaons, and protons

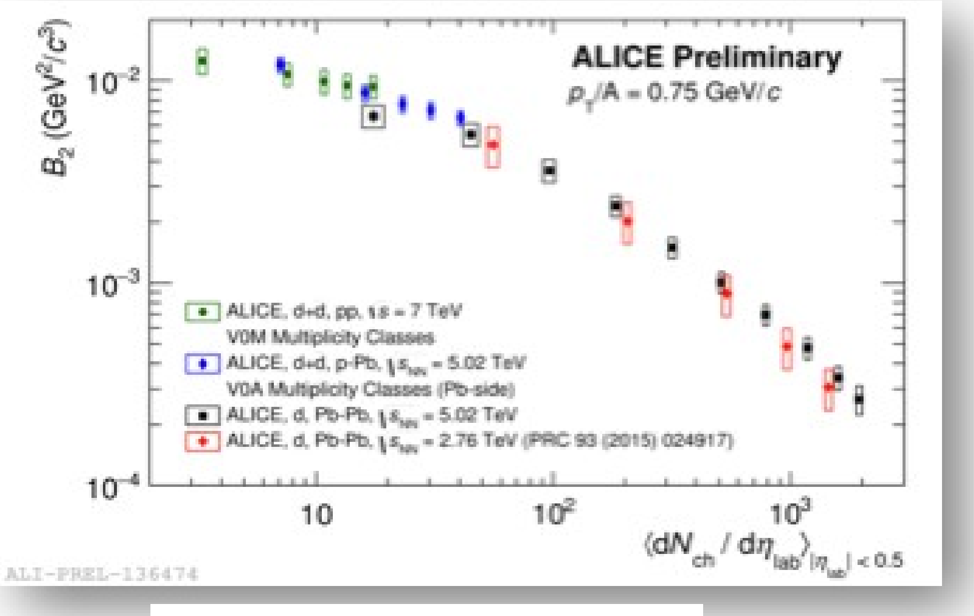


even hyper-triton flows with same common fluid velocity



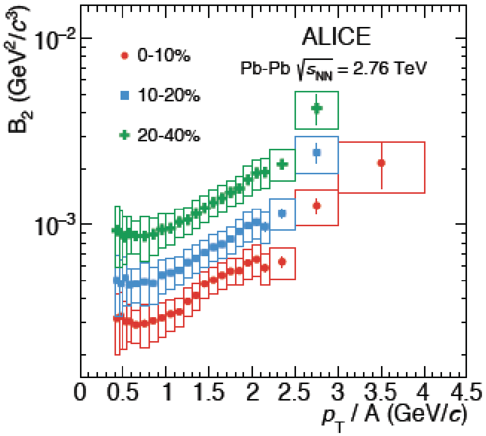
is coalescence approach an alternative?

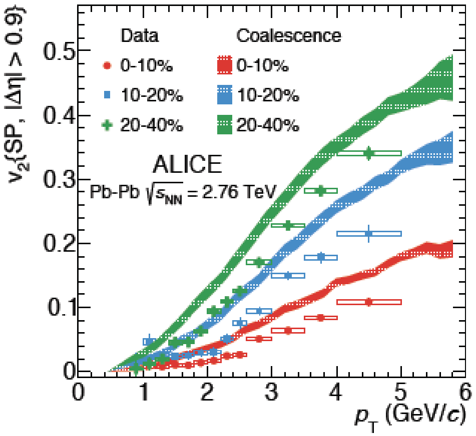
$$E_i \frac{d^3N_i}{dp_i^3} = B_A \left(E_p \frac{d^3N_p}{dp_p^3}\right)^A \qquad B_A = \left(\frac{4\pi}{3}p_0^3\right)^{A-1}\frac{M}{m^A}$$

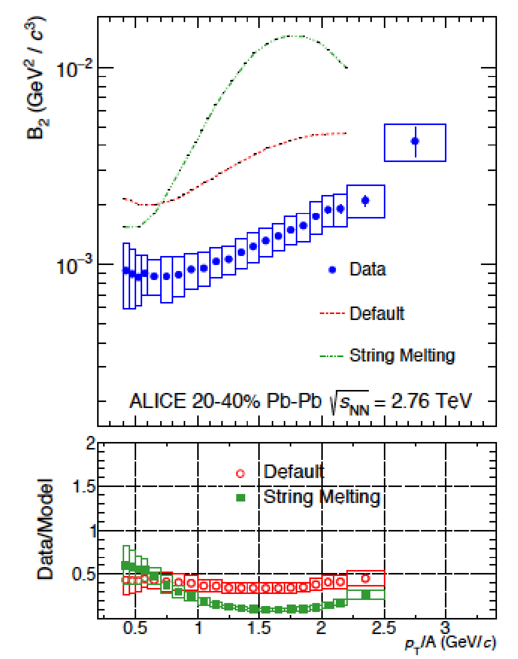


centrality and p_T dependence of coalescence parameter not understood and not well reproduced by models such as AMPT

ALICE: arXiv:1707.07304







coalescence approach, general considerations for loosely bound states

- production yields of loosely bound states is entirely determined by mass, quantum numbers and fireball temperature.

- hyper-triton and 3He have very different wave functions but essentially equal production yields.

- energy conservation needs to be taken into account when forming objects with baryon number A from A baryons.

- delicate balance between formation and destruction; maximum momentum transfer onto hyper-triton before it breaks up: $\Delta Q_{max} < 20$ MeV/c, typical pion momentum p_pi = 250 MeV/c, typical hadronic momentum tranfer > 100 MeV/c

- hyper-triton interaction cross section with pions or nucleons at thermal freeze-out is of order $\sigma > 70\ \mathrm{fm}^2$. For the majority of hyper-tritons to survive, the mfp λ has to exceed 15 fm → density of fireball at formation of hyper-triton $n < 1/(\lambda\,\sigma) = 0.001/\mathrm{fm}^3$. Completely inconsistent with formation at kinetic freeze-out, where $n \approx 0.05$

a possible way out

Quark Model Spectroscopy

Why does the quark model work so well?

Why do M and B body plans dominate?

Why don't multibaryons make one big bag?

Frank Wilczek, QM2014 introductory talk

hypothesis:
all nuclei and hyper-nuclei are formed as compact multi-quark states at the phase boundary. Then slow time evolution into hadronic respresentation.

Andronic, pbm, Redlich, Stachel, arXiv :1710.09425

How can this be tested?

precision measurement of spectra and flow pattern for light nuclei and hyper-nuclei

a major new opportunity for ALICE Run3
and for CBM/NICA/JPARC/NA61

summary

- statistical hadronization model is effective tool to understand the phenomenology of hadron production in relativistic nuclear collisions from SIS to LHC energy
- deeply rooted in duality 'hadrons – quarks' near QCD phase boundary
- present precision is at the 10% level, mostly limited by incomplete knowledge of hadron mass spectrum and related branching ratios for decays
- measurements from ALICE at the 5% accuracy level shows deviations for protons and cascades at the 2 – 3 sigma level → need to be followed up
- yields of light nuclei and hyper-nuclei successfully predicted → maybe produced as quark bags?
- coalescence approach not well suited for loosely bound states
- statistical hadronization approach also applies to the heavy quark sector – not covered here

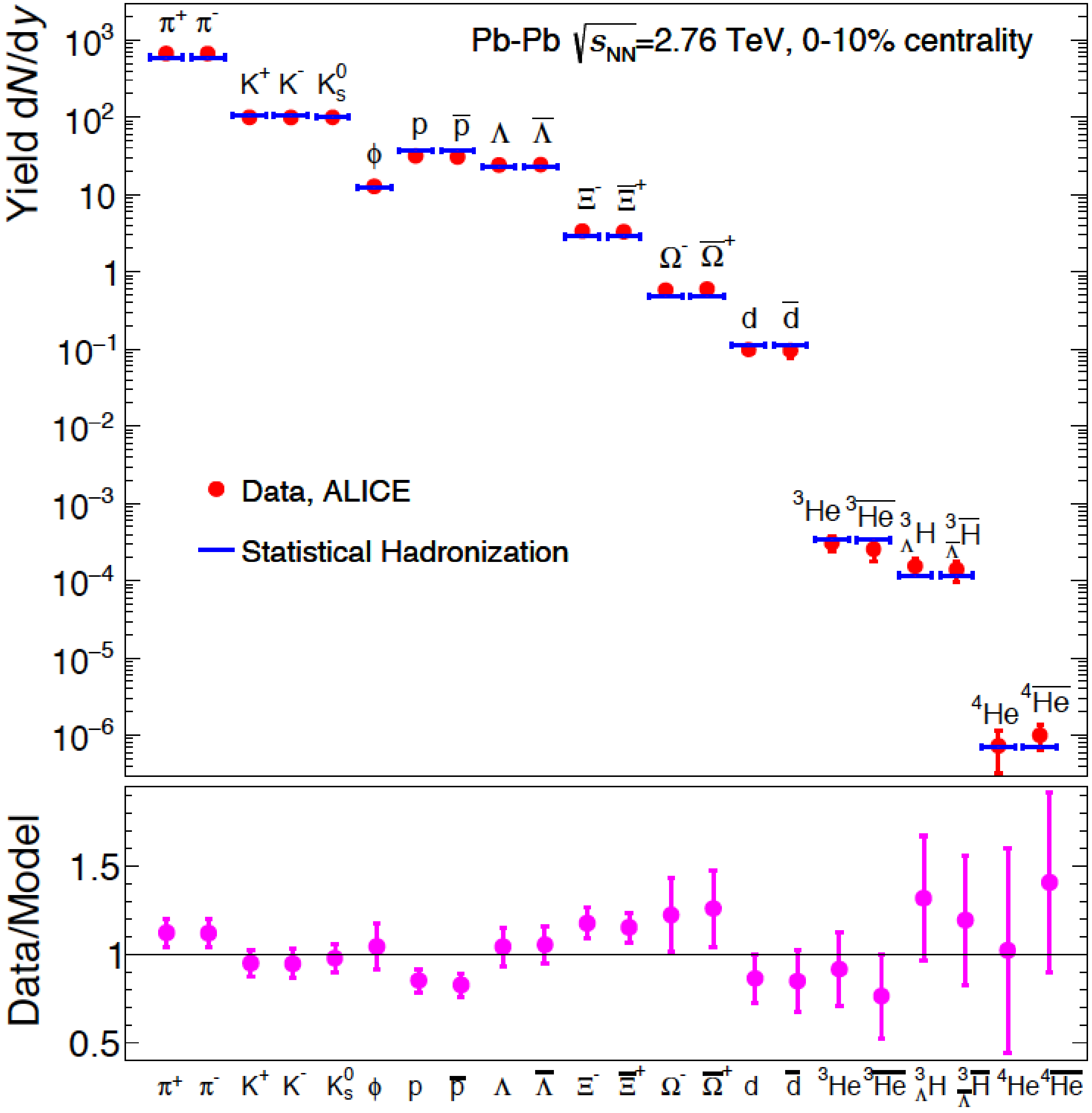
key results:
experimental location of QCD phase boundary for $\mu_b < 300$ MeV:
$T_c = 156 \pm 5$ MeV
new insight into hadronization

open issues and questions

- why vacuum masses near phase boundary?
- transition from canonical to grand canonical regime
- are higher moments more sensitive to thermal parameters?
- incomplete hadron mass spectrum?
- uncertainty from statistical hadronization model

thermal fit with statistical hadronization model uses vacuum masses for all hadrons!

fit includes loosely bound systems such a
deuteron and hypertriton
hypertriton is bound-state of (Λ,p,n), Λ separation energy about 130 keV
size about 10 fm, the **ultimate halo nucleus,**
produced at T=156 MeV. close to an Efimov state

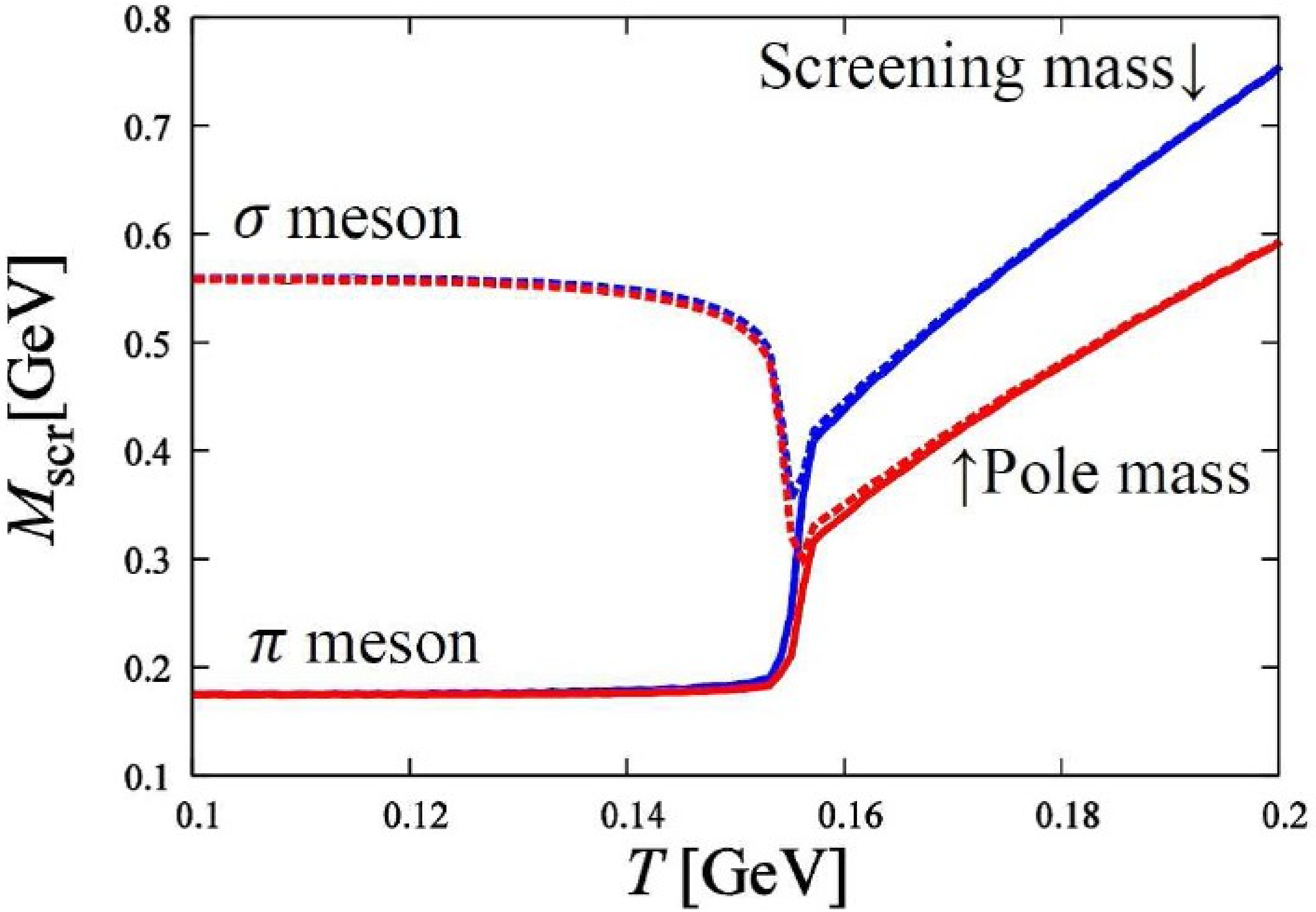


proton discrepancy 2.8 sigma

Xi discrepancy?

J. Stachel, A. Andronic, P. Braun-Munzinger and K. Redlich, Confronting LHC data with the statistical hadronization model, J.Phys.Conf.Ser.**509** (2014) 012019, arXiv:1311.4662 [nucl-th].

temperature dependence of meson masses in a NJL model



Mesonic correlation functions at finite temperature and density in the Nambu-Jona-Lasinio model with a Polyakov loop

H. Hansen, W.M. Alberico (INFN, Turin & Turin U.), A. Beraudo (Saclay, SPhT), A. Molinari, M. Nardi (INFN, Turin & Turin U.), C. Ratti (ECT, Trento & INFN, Trento). Sep 2006. 26 pp.

Phys.Rev. D75 (2007) 065004

If the pion mass would be 300 MeV near T_c this would have drastic consequences, especially if nucleon mass is unchanged, see below

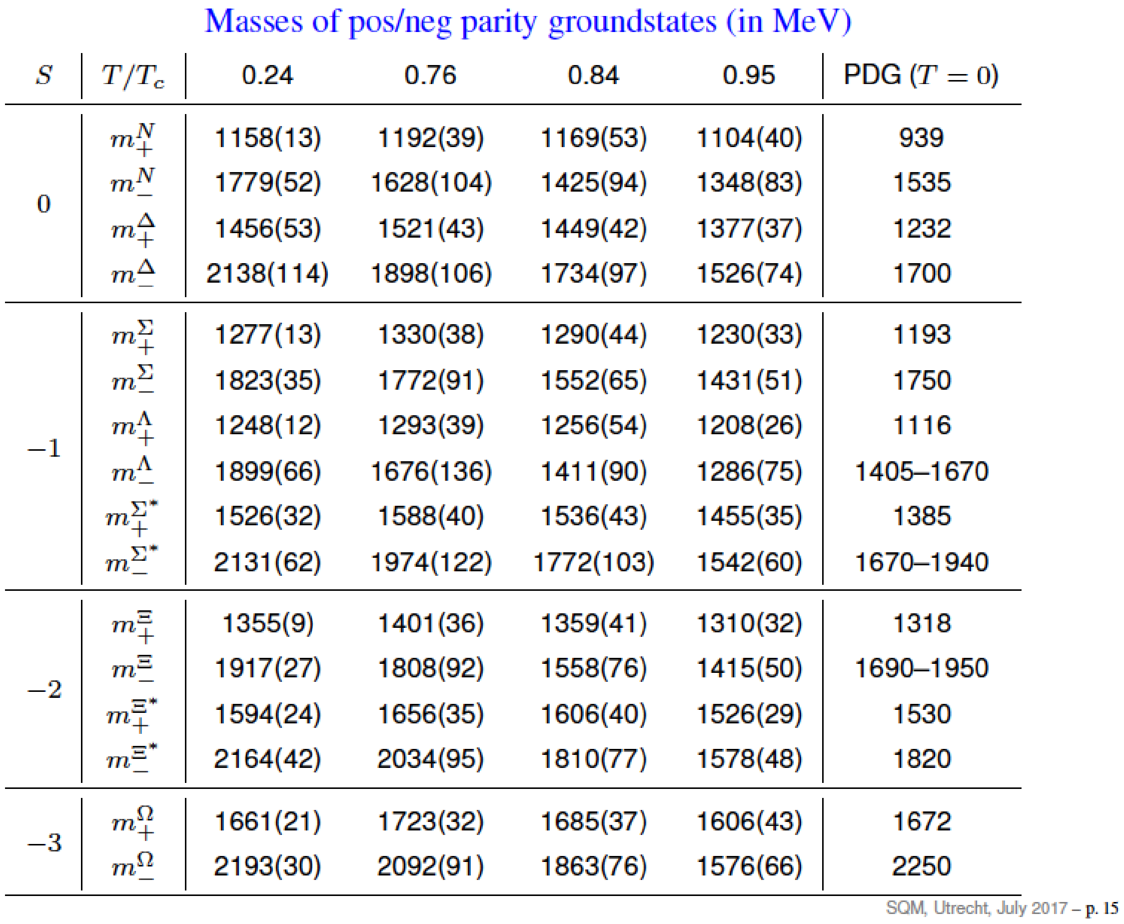
also: changing masses near $T_c = T_{chem}$ would invalidate the chemical freeze-out picture as it implies a dense hadronic phase below T_c
strong interactions are needed to bring masses back on the mass shell and adjust particle numbers

From G. Aarts, SQM2017

`PRD 92 (2015) 014503, arXiv:1502.03603 [hep-lat]`

`JHEP 06 (2017) 034, arXiv:1703.09246 [hep-lat]`

`in preparation`



Masses of pos/neg parity groundstates (in MeV)

S	T/T_c	0.24	0.76	0.84	0.95	PDG ($T=0$)
0	m_+^N	1158(13)	1192(39)	1169(53)	1104(40)	939
	m_-^N	1779(52)	1628(104)	1425(94)	1348(83)	1535
	m_+^Δ	1456(53)	1521(43)	1449(42)	1377(37)	1232
	m_-^Δ	2138(114)	1898(106)	1734(97)	1526(74)	1700
-1	m_+^Σ	1277(13)	1330(38)	1290(44)	1230(33)	1193
	m_-^Σ	1823(35)	1772(91)	1552(65)	1431(51)	1750
	m_+^Λ	1248(12)	1293(39)	1256(54)	1208(26)	1116
	m_-^Λ	1899(66)	1676(136)	1411(90)	1286(75)	1405–1670
	$m_+^{\Sigma^*}$	1526(32)	1588(40)	1536(43)	1455(35)	1385
	$m_-^{\Sigma^*}$	2131(62)	1974(122)	1772(103)	1542(60)	1670–1940
-2	m_+^Ξ	1355(9)	1401(36)	1359(41)	1310(32)	1318
	m_-^Ξ	1917(27)	1808(92)	1558(76)	1415(50)	1690–1950
	$m_+^{\Xi^*}$	1594(24)	1656(35)	1606(40)	1526(29)	1530
	$m_-^{\Xi^*}$	2164(42)	2034(95)	1810(77)	1578(48)	1820
-3	m_+^Ω	1661(21)	1723(32)	1685(37)	1606(43)	1672
	m_-^Ω	2193(30)	2092(91)	1863(76)	1576(66)	2250

SQM, Utrecht, July 2017 – p. 15

change of baryon masses near T_c

From G. Aarts, SQM2017

`PRD 92 (2015) 014503, arXiv:1502.03603 [hep-lat]`

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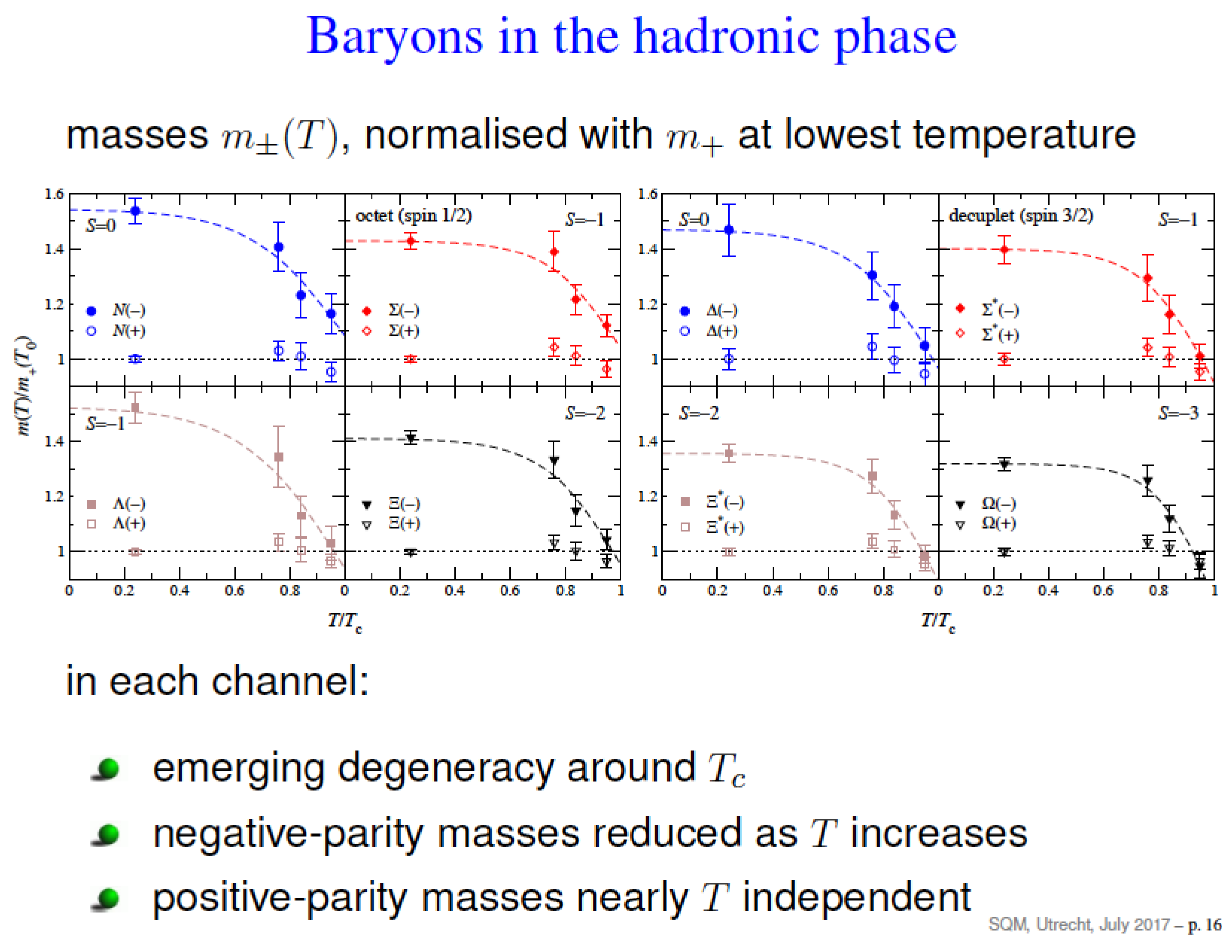

Baryons in the hadronic phase

masses $m_\pm(T)$, normalised with m_+ at lowest temperature

in each channel:

- emerging degeneracy around T_c
- negative-parity masses reduced as T increases
- positive-parity masses nearly T independent

SQM, Utrecht, July 2017 – p. 16

but negative parity baryons all lie higher up in the mass distribution
→ small effects on statistical hadronization results … to be tested