Heavy Flavor in Nuclear Collisions

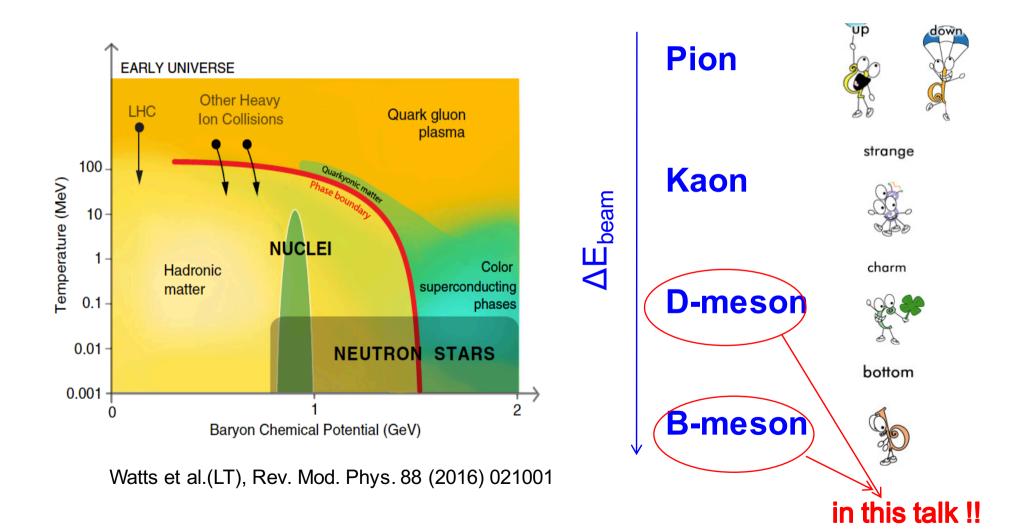


Laura Tolós



2nd EMMI Workshop: Anti-matter, hyper-matter and exotica production at the LHC 6-10 November 2017, Turin (Italy)

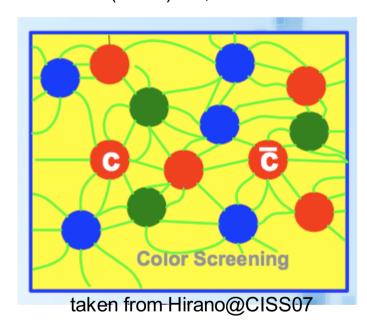
What happens to heavy hadrons at nonzero temperature and density?



Charm under Extreme Conditions

J/Ψ suppression

Gonin et al (NA50) '96, Matsui and Satz '86



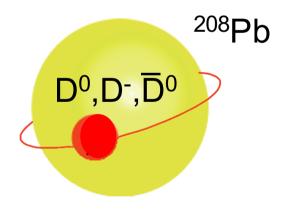
but also comover scattering

$$J/\Psi + \pi < -> D + \overline{D}$$

Capella, Armesto, Ferreiro, Vogt, Wang, Bratkovskaya, Cassing, Andronic..

D-mesic nuclei

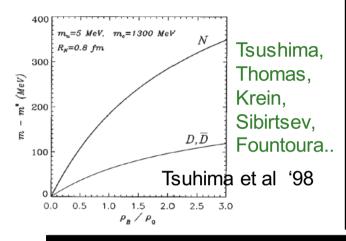
Tsushima et al '99, Garcia-Recio et al '10 Garcia-Recio et al '12 Yasui et al '12, Yamagata-Sekihara '16..



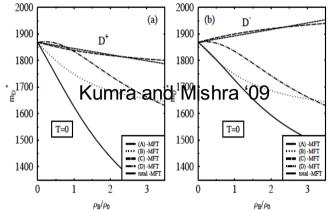
Theoretical models in matter

QMC model

exchange of ω , ρ , σ mesons among quarks in hadron bag



MF/RHF model



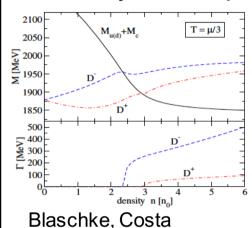
MF/RHF approach that includes charmed mesons

h Mazumdar,

Mishra. Kumar...

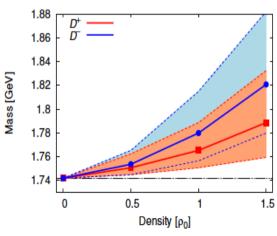
NJL model

SU(4) NJL model with Polyakov loop



and Kalinovsky '12

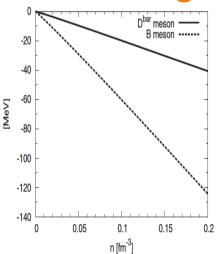
QCD sum-rule



Suzuki, Gubler and Oka '15

Gubler, Hayashigaki,
Hilger, Kaempfer, Leupold,
Nielsen, Navarra, Oka,
Suzuki, Thomas, Weise, ..
operator product
expansion for inmedium correlation
function and relate it
to the spectral
function

π exchange with HQ



heavy meson nucleon
interaction
mediated by
π exchange
with HQ

Yasui and Sudoh '13

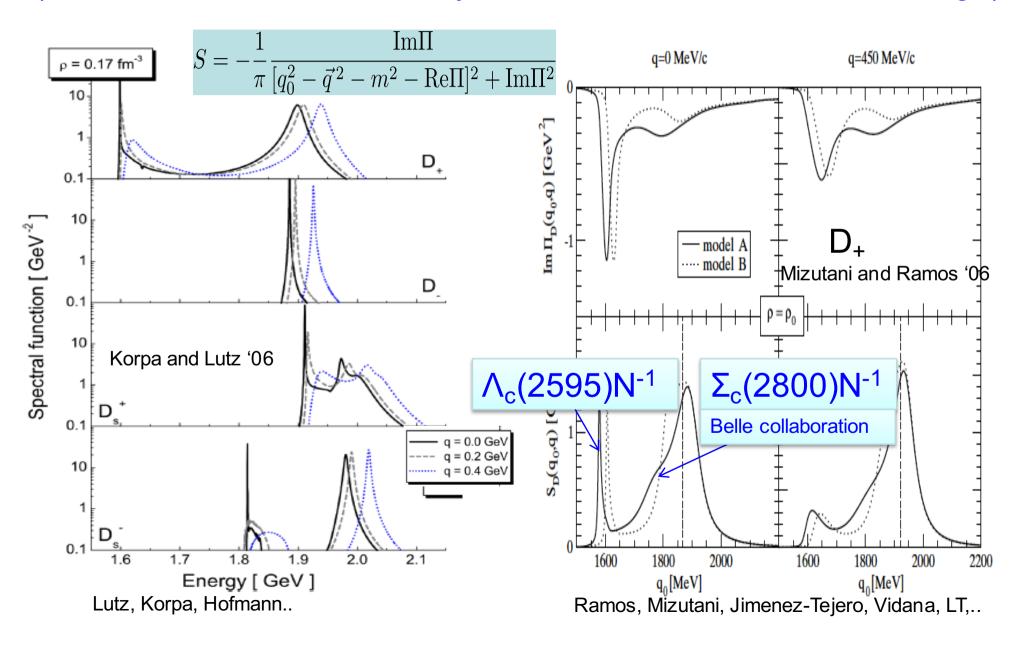
Unitarized theory in matter:

selfconsistent coupled-channel procedure

Unitarized theory in matter

selfconsistent coupled-channel procedure

(bare interaction saturated by t-channel vector-meson exchange)



Meson-baryon interaction with heavy quarks: Incorporate Heavy-Quark Spin Symmetry

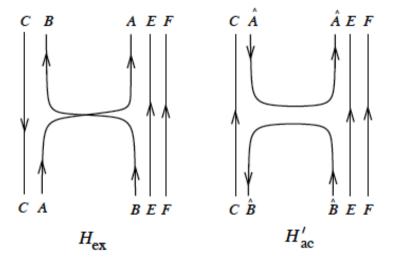
HQSS*: spin interactions vanish for infinitely massive quarks

*Isgur, Wise, Manohar, Neubert

To construct a model for four flavors for pseudoscalar and vector mesons as well as 1/2+ and 3/2+ baryons that incorporates HQSS in the charm sector: extended WT interaction that fulfills SU(6)xHQSS and it is consistent with chiral symmetry in the light sector

$$V = \frac{K(s)}{4f^2}H'_{WT}, \qquad H'_{WT} = H_{ex} + H'_{ac}.$$

K(s): depends on meson-baryon energy f: decay constant

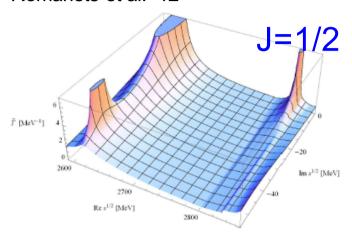


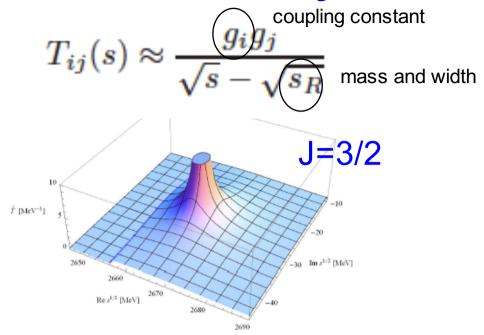
H_{ex}: exchange of quarks H'_{ac}: annhilation and creation of quarkantiquark pairs, corrected with HQSS constraints (only light quarks)

Spectroscopy of excited charmed baryons

Λ_{c} : C=1,S=0,I=0

Garcia-Recio et al.'09; Romanets et al. '12

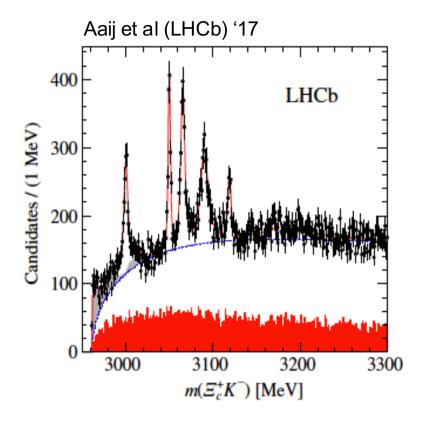




SU(8) irrep	SU(6) irrep	SU(3) irrep	M_R	Γ_R	Couplings to main channels	Status PDG	J
168	152,1	3*	2617.3	89.8	$g_{\Sigma_c \pi} = 2.3, g_{ND} = 1.6, g_{ND^*} = 1.4,$		1/2
168	15 _{2,1}	3*	2666.6	53.7		$\Lambda_c(2625) ****$	3/2
168	21 _{2,1}	3*2	2618.8	1.2	$g_{\Sigma_c^*\sigma} = 1.3$ $g_{\Sigma_c} = 0.2$ $g_{ND} = 3.5, g_{ND^*} = 5.6,$	$\Lambda_c(2595) ***$	1/2
120	21 _{2,1}	3_2^*	2828.4	0.8	$g_{\Lambda D_s} = 1.4, g_{\Lambda D_s} = 2.9, g_{\Lambda_c \eta} = 6.9$ $g_{ND} = 0.3, g_{\Lambda_c \eta} = 1.1, g_{\Xi_c K} = 1.6,$		1/2
					$g_{\Lambda D_s^*} = 1.1, g_{\Sigma_c \rho} = 1.1, g_{\Sigma_c^* \rho} = 1.0,$ $g_{\Xi_c^* K^*} = 0.8$		

- $\Lambda_c(2595)$ has large DN and D*N components
- Double-pole pattern for $\Lambda_c(2595)$, like for $\Lambda(1405)$
- Identification of $\Lambda_c(2625)$

Ω_{c} : C=1,S=-2,I=0

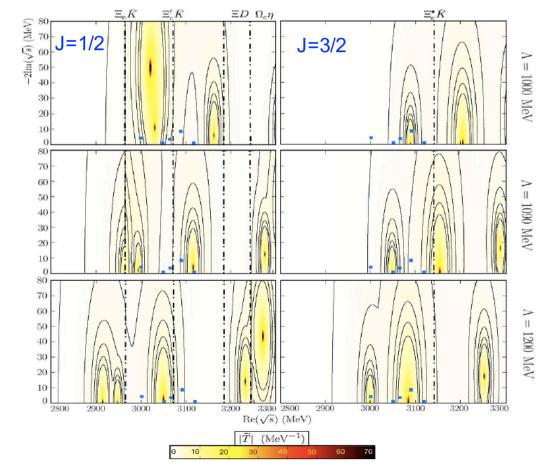


Resonance	Mass (MeV)	Γ (MeV)
$\Omega_c(3000)^0$	$3000.4 \pm 0.2 \pm 0.1^{+0.3}_{-0.5}$	$4.5 \pm 0.6 \pm 0.3$
$\Omega_c(3050)^0$	$3050.2 \pm 0.1 \pm 0.1^{+0.3}_{-0.5}$	$0.8\pm0.2\pm0.1$
	Viet.	<1.2 MeV, 95% C.L.
$\Omega_c(3066)^0$	$3065.6 \pm 0.1 \pm 0.3^{+0.3}_{-0.5}$	$3.5\pm0.4\pm0.2$
$\Omega_c(3090)^0$	$3090.2 \pm 0.3 \pm 0.5_{-0.5}^{+0.3}$	$8.7\pm1.0\pm0.8$
$\Omega_c(3119)^0$	$3119.1 \pm 0.3 \pm 0.9^{+0.3}_{-0.5}$	$1.1\pm0.8\pm0.4$
, ,	-4.7	<2.6 MeV, 95% C.L.

fine-tuning the subtraction constants (associated to a cutoff)

Λ=1090 MeV

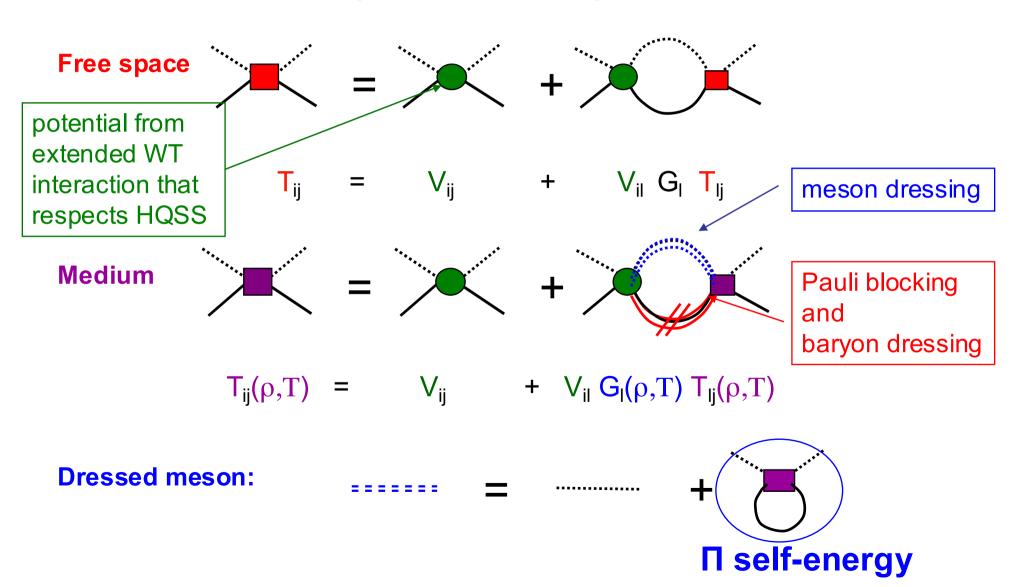
M_R (MeV)	Γ_R (MeV)	\boldsymbol{J}	M_R^{exp}	Γ_R^{exp}
2963.95	0.0	1/2		
2994.26	1.85	1/2	3000.4	4.5
3048.7	0.0	3/2	3050.2	0.8
3116.81	3.72	1/2	3119.1	1.1
3155.37	0.17	3/2	_	_



Nieves, Pavao and LT, in preparation

Charmed hadrons in matter

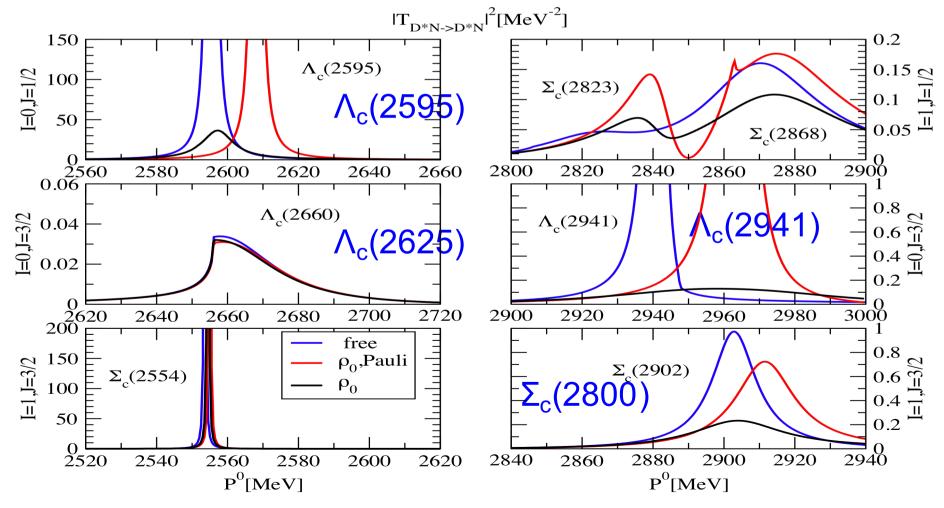
Unitarized theory in matter: selfconsistent coupled-channel procedure



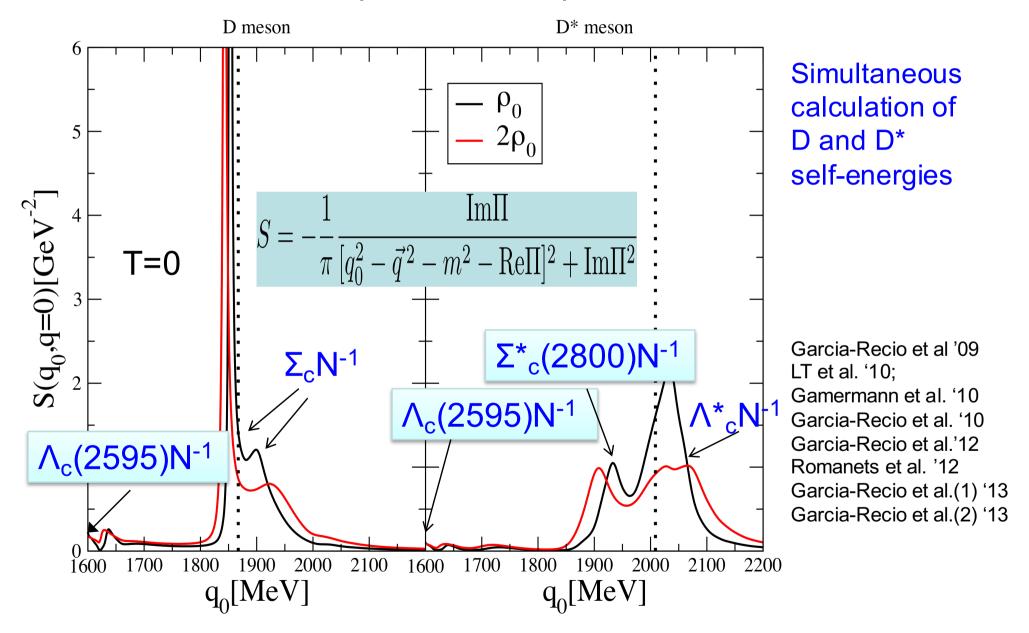
	Resonance	$I(J^P)$	${\bf Status}$	Mass [MeV]	$\Gamma \ [{ m MeV}]$
	$\Lambda_c(2595)$	$0(1/2^-)$	***	2592.25 ± 0.28	2.6 ± 0.6
	$\Lambda_c(2625)$	$0(3/2^-)$	***	2628.11 ± 0.19	< 0.97
C	$\Lambda_c(2765)$	$?(?^?)$	*	2766.6 ± 2.4	50
Ŏ	or $\Sigma_c(2765)$				
\Box	$\Lambda_c(2880)$	$0(5/2^+)$	***	2881.53 ± 0.35	5.8 ± 1.1
	$\Lambda_c(2940)$	$0(?^{?})$	***	2939.3 + 1.4 - 1.5	17 + 8 - 6
	$\Sigma_c(2800)^{++}$	$1(?^{?})$	***	2801 + 4 - 6	75 + 22 - 17
	$\Sigma_c(2800)^+$	$1(?^{?})$	***	2792 + 14 - 5	62 + 60 - 40
	$\Sigma_c(2800)^0$	$1(?^{?})$	***	2806 + 5 - 7	72 + 22 - 15

Dynamically-generated baryonic resonances in nuclear matter

LT, Garcia-Recio and Nieves '10 [α fitted to reproduce $\Lambda_c(2595)$ and analyze energies up to 3.5 GeV]



Unitarized theory in matter: selfconsistent coupled-channel procedure

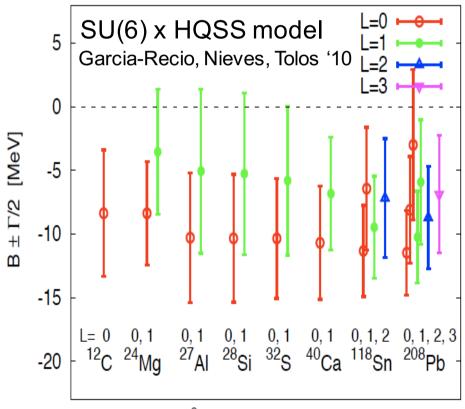


D-mesic nuclei

Solving Schroedinger or Klein-Gordon equation with

- potential from QMC model
- potential from SU(6) x HQSS model
- potential from π exchange model with HQS

 D^0 and D^+



D⁰-nucleus bound states

- SU(6) x HQSS model: weakly bound D⁰-nucleus states with important widths in contrast to previous QMC model

	²⁰⁸ Pb
State	$D^0(V^q_\omega)$
1 s	-96.2
1 <i>p</i>	-93.0
2 <i>s</i>	-88.5

Tsushima et al. '99 Krein et al '17 (review)

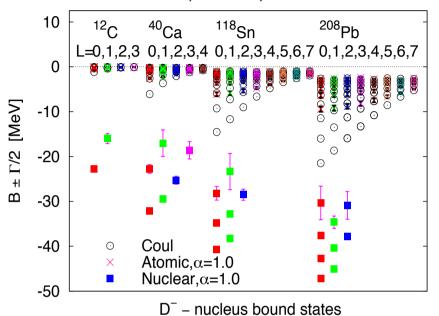
- SU(6) x HQSS model:

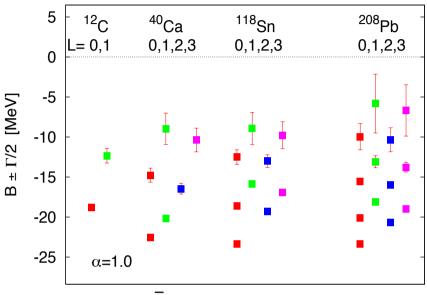
D⁺ does not bind

D^- and \overline{D}^0

SU(6) x HQSS model:

Garcia-Recio, Nieves, Salcedo and Tolos '12





 \overline{D}_0 – nucleus bound states

QMC model

Tsushima, Lu, Thomas, Saito and Landau '99 Krein, Thomas and Tsushima '17 (review)

²⁰⁸Pb

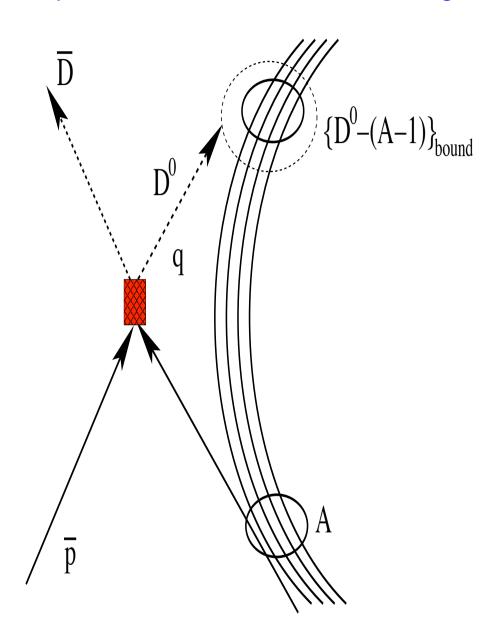
State	$D^-(\widetilde{V}^q_\omega)$	$D^-(V^q_\omega)$	$D^-(V^q_\omega$, no Coulomb)	$\bar{D}^0(\widetilde{V}^q_\omega)$	$\bar{D}^0(V^q_\omega)$
1 <i>s</i>	-10.6	-35.2	-11.2	Unbound	-25.4
1 <i>p</i>	-10.2	-32.1	-10.0	Unbound	-23.1
2 <i>s</i>	-7.7	-30.0	-6.6	Unbound	-19.7

π exchange model with HQS

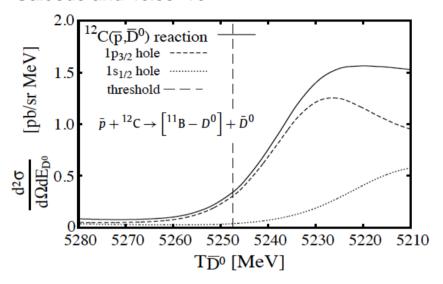
Yasui and Sudoh '13

⁴⁰ Ca	²⁰⁸ ₽b		
2.0 MeV	1.9 MeV		
5.9 MeV p-wave	9.2 MeV 7.3 MeV p-wave		
14.9 MeV	1 <u>8.7 MeV</u>		
28.2 MeV g.s.	27.0 MeV		
3-wave	32.8 MeV g.s. s-wave		

Experimental observation is, though, a difficult task (PANDA, J-PARC?)



Yamagata-Sekihara, Garcia-Recio, Nieves, Salcedo and Tolos '16



Large momentum transfer (about 1 GeV/c) makes any structure due to bound states not noticeable.

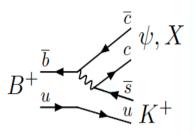
Need of reactions with lower momentum transfer, such as

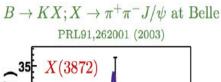
$$\bar{p} + p \to D^{*-} + D^{+},$$
 $D^{*-} + A_Z \to \pi^0 + [A_Z - D^{-}]_b$

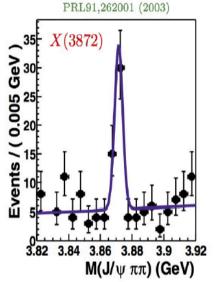
B decays: Belle, BaBar, LHCb, etc.

Charm resonances in matter: X(3872)

as D D* molecule







$\begin{array}{c c} X & \hat{g} & \hat{g} & X \\ \hline D^+ & X & \hat{g} & \hat{g} & X \\ D^{*-} & D^{*+} & D^{*+} \end{array}$	$ \begin{array}{c c} X & \begin{array}{c} & \begin{array}{c} & \begin{array}{c} & \\ & \\ & \end{array} \end{array} $ $ \begin{array}{c} & \begin{array}{c} & \\ & \\ & \end{array} \end{array} $ $ \begin{array}{c} & \begin{array}{c} & \\ & \\ & \end{array} \end{array} $ $ \begin{array}{c} & \\ & \\ & \end{array} $ $ \begin{array}{c} & \\ & \\ & \end{array} $ $ \begin{array}{c} & \\ & \\ & \\ & \end{array} $ $ \begin{array}{c} & \\ & \\ & \\ & \end{array} $ $ \begin{array}{c} & \\ & \\ & \\ & \\ & \end{array} $ $ \begin{array}{c} & \\ & \\ & \\ & \\ & \\ & \end{array} $ $ \begin{array}{c} & \\ & \\ & \\ & \\ & \\ & \\ & \\ & $
$\begin{array}{c} 1 \\ 0.8 \\ \hline \\ 0.6 \\ \hline \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ $	0.12 0.1 $P = 0.10\rho_0$ $\rho = 0.30\rho_0$ $\rho = 0.50\rho_0$ $\rho = 0.75\rho_0$ $\rho = 0.75\rho_0$ $\rho = 1.00\rho_0$ 0.02 0.04 0.02 $P_0 = 0.999$ 0.04 0.02 $P_0 = 0.999$ 0.04 0.08 $P_0 = 0.999$ 0.09 0.09 0.09 0.004 0.002 0.004 0.005 0.004 0.005 0.004 0.005 0.004 0.005 0.005 0.006 0.007 0.007 0.007 0.008 0.009 0.00

_		$P_0 = 0.999$		
_	ρ/ρ_0	E_{max} (MeV)	E_{qp} (MeV)	
	0.10	3873.3	3878.0	
	0.30	3870.8	_	
	0.50	3866.9	_	
	0.75	3860.3	_	
	1.00	3851.9	_	

- Spectral function: $S_X^{(\rho)}(E) = -\frac{1}{\pi} \text{Im} \Delta_{\rho}(E) =$ $\text{Im}\Pi_{\rho}(E)$ $\frac{1}{\pi} \frac{1}{(E^2 - m_0^2 - \text{Re}\Pi_{\rho}(E))^2 + (\text{Im}\Pi_{\rho}(E))^2}$
- Quasi-particle peak: $E_{\rm qp}^2 - m_0^2 - {\rm Re}\Pi(E_{\rm qp}) = 0.$

Albaladejo, Nieves and LT, in preparation

D meson propagation in dense hot matter

D-mesons: One of the cleanest probes of the early stages of the collision

Fokker-Planck equation (m _D >> m _{bath}, m _D >> T_{bath})

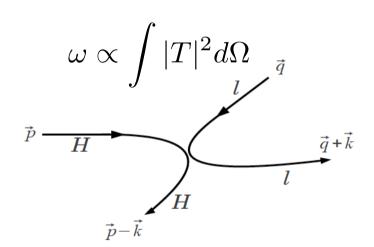
$$\frac{\partial f(t, \mathbf{p})}{\partial t} = \frac{\partial}{\partial p_i} \left[F_i(\mathbf{p}) f(t, \mathbf{p}) + \frac{\partial}{\partial p_j} \left[\Gamma_{ij}(\mathbf{p}) f(t, \mathbf{p}) \right] \right],$$

isotropic bath

$$F(p) = \int d\mathbf{k} w(\mathbf{p}, \mathbf{k}) \frac{k_i p^i}{p^2},$$

$$\Gamma_0(p) = \frac{1}{4} \int d\mathbf{k} w(\mathbf{p}, \mathbf{k}) \left[\mathbf{k}^2 - \frac{(k_i p^i)^2}{p^2} \right],$$

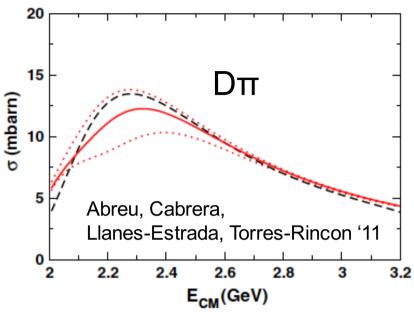
$$\Gamma_1(p) = \frac{1}{2} \int d\mathbf{k} w(\mathbf{p}, \mathbf{k}) \frac{(k_i p^i)^2}{p^2},$$



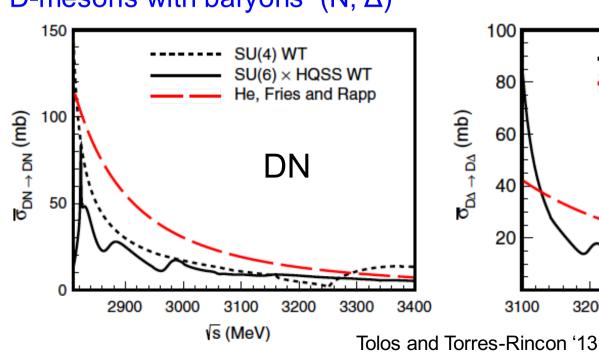
Previous works Laine '11; He, Fries, Rapp '11; Ghosh, Das, Sarkar, -eAlam '11
We need scattering amplitudes |T|² (first approximation: |T|² in free space)
Abreu, Cabrera, Llanes-Estrada, Torres-Rincon '11; LT and Torres-Rincon '13

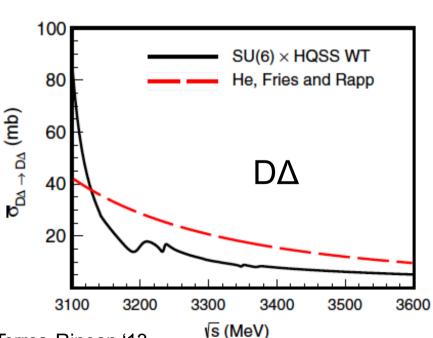
(Vacuum) cross sections for open charm with mesons and baryons

D-mesons with light mesons (π, K, K, η)



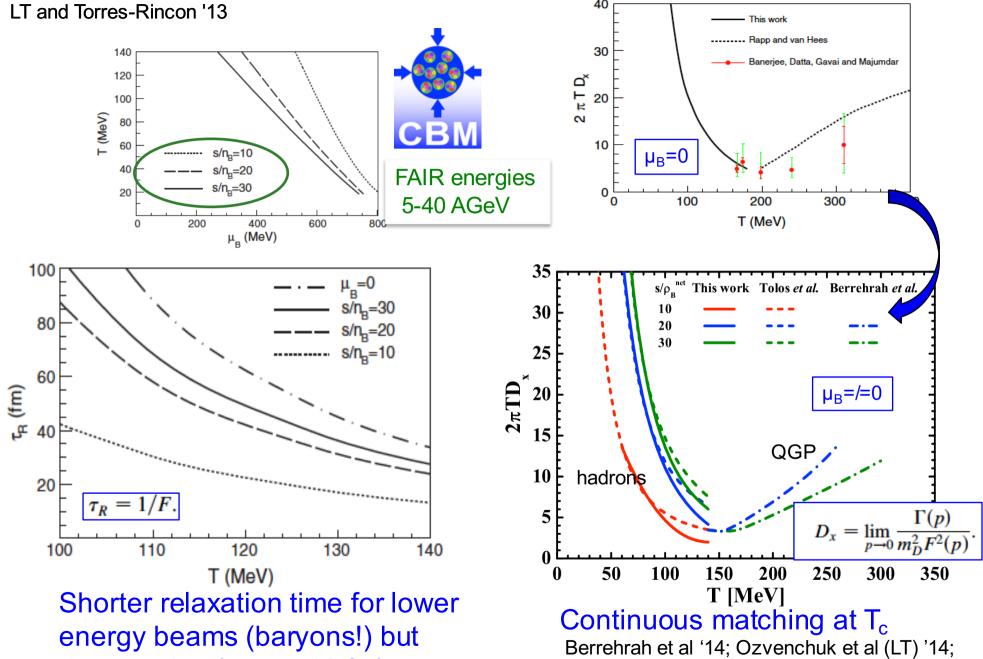
D-mesons with baryons (N, Δ)





Some results for CBM@FAIR energies

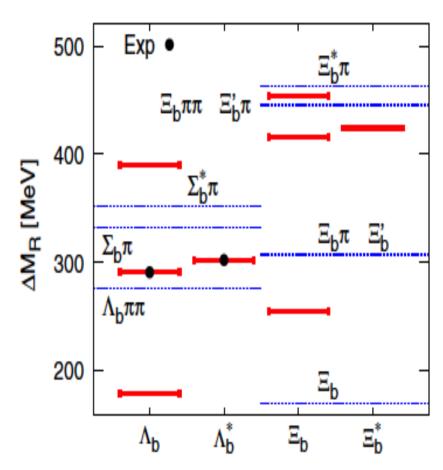
do not relax (τ_{fireball}~10 fm)



Song et al (LT) '15

Beauty under Extreme Conditions

Spectroscopy of excited beauty baryons



Garcia-Recio, Nieves, Romanets, Salcedo and LT '13 $\Lambda_b(5912)$ and $\Lambda^*_b(5920)$ found by LHCb* collaboration are described as mesonbaryon molecular states belonging to a HQSS doublet. New HQSS partners are predicted: $\Xi_b(6035)$ and $\Xi_b(6043)$

* Aaij et al (LHCb) '12

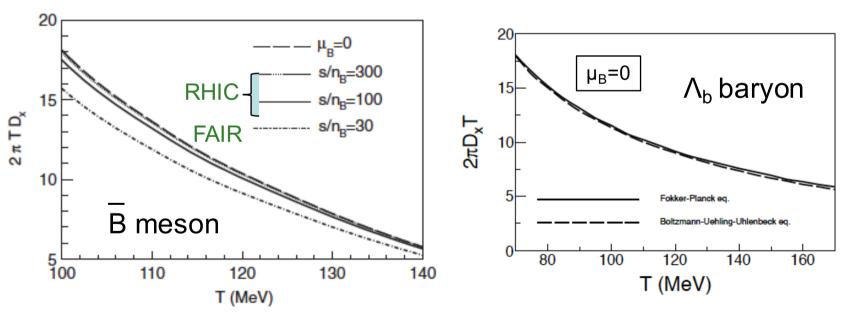
Beauty propagation in dense hot matter

Fokker-Planck equation

$$\frac{\partial f(t, \mathbf{p})}{\partial t} = \frac{\partial}{\partial p_i} \left[F_i(\mathbf{p}) f(t, \mathbf{p}) + \frac{\partial}{\partial p_j} [F_i(\mathbf{p}) f(t, \mathbf{p})] \right]_{\text{diffusion coefficient}}$$

Results from FAIR to RHIC energies (LHC?)

Torres-Rincon, LT and Romanets '14; LT, Torres-Rincon and Das '16; Song et al (LT) '16



- Results insensitive to trajectory for high s/n_B: prediction for behaviour of hadronic medium at RHIC energies
- Similar behaviour of diffusion coefficient for \overline{B} meson and Λ_b

Summary

> it is an exciting moment



- moving from the light to the heavy sector
- ➤ a lot of theoretical effort is needed (how to construct a reliable effective theory that implements the correct symmetries)
- ➤ but in close connection to experiments/lattice (how to provide feedback between theory and experiments/lattice: spectroscopy of excited states, spectra of meson-nucleus, transport coefficients,..)





