

The Hypertriton as an Efimov State

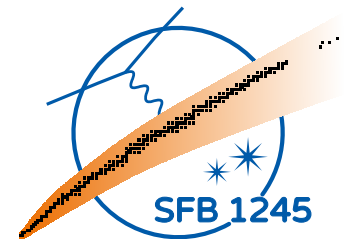
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“Anti-matter, hyper-matter and exotica production at the LHC” Turin, Italy, Nov. 6-10, 2017

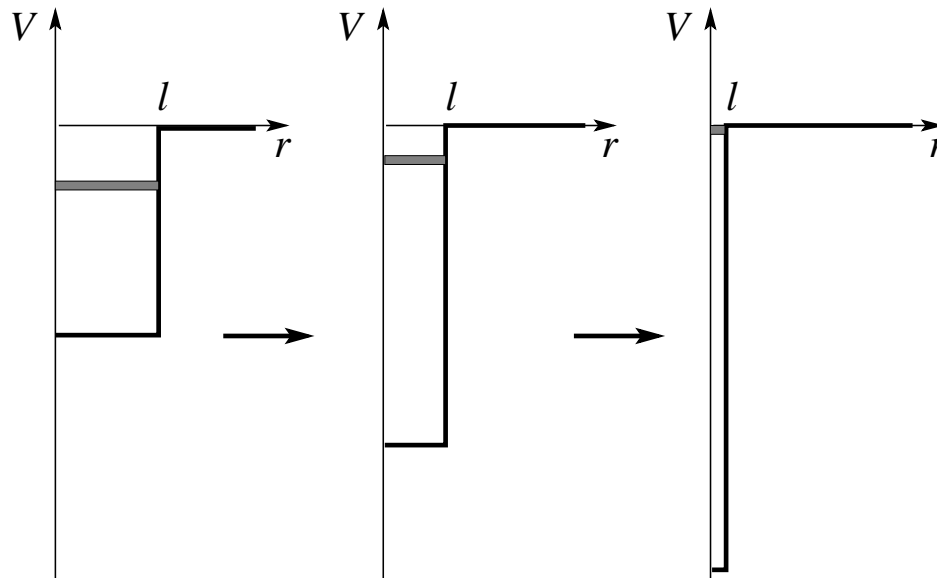
- Threshold bound states and the unitary limit
- Limit cycles and Efimov physics
- Hypertriton
- Factorization in break-up and recombination
- Summary and outlook

Collaborators: F. Hildenbrand, ...

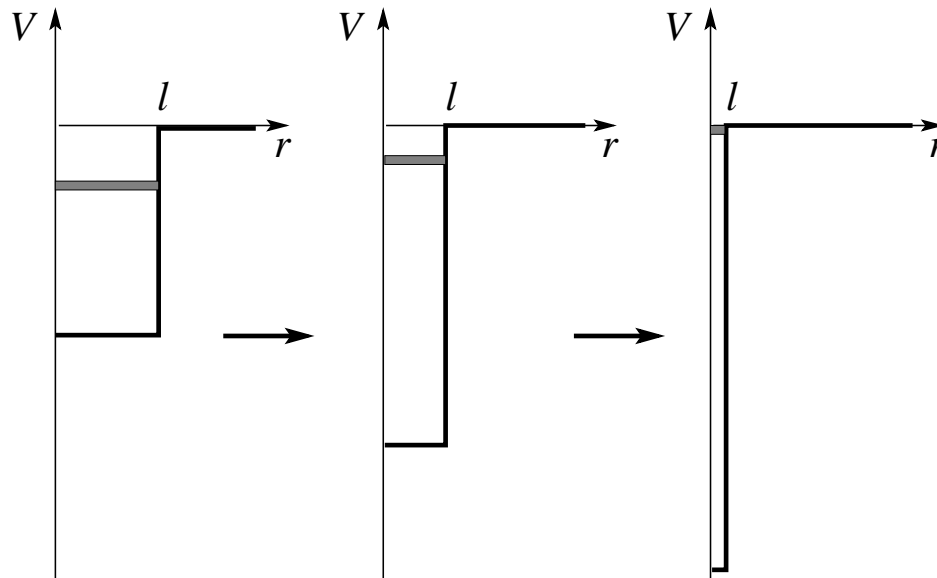
Review articles: HWH, Ji, Phillips, J. Phys. G **44** (2017) 103002

Braaten, HWH, Phys. Rep. **428** (2006) 259

- Consider system with short-ranged, resonant interactions
- Unitary limit: $a \rightarrow \infty$, $\ell \sim r_e \rightarrow 0$ (cf. Bertsch problem, 2000)



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- Use as starting point for description of few-body properties
 - Large scattering length: $|a| \gg \ell \sim r_e, l_{vdW}, \dots$
 - Natural expansion parameter: $\ell/|a|, k\ell, \dots$
 - **Universal dimer** with energy $B_2 = -1/(ma^2)$ ($a > 0$)
size $\langle r^2 \rangle^{1/2} \sim a \implies$ **halo state**

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- But Efimov effect in certain channels

EFT for the Unitary Limit

- Effective Lagrangian

(Kaplan, 1997; Bedaque, HWH, van Kolck, 1999)

$$\mathcal{L}_{eff} = \text{---} + \text{==} + \text{==} \text{---} + \text{---} \text{==} + \text{---} \text{---} + \dots$$

- 2-body amplitude:

$$\text{---} = \text{==} + \text{---} \text{---} + \text{---} \text{---} \text{---} + \dots$$

- 2-body coupling g_2 near fixed point ($1/a = 0$)

\Rightarrow **scale and conformal invariance** \iff **unitary limit**

(Mehen, Stewart, Wise, 2000; Nishida, Son, 2007; ...)

- 3-body amplitude:

$$\text{---} \text{---} \text{---} = \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---}$$

$g_3(\Lambda) \Rightarrow$ **limit cycle**

\Rightarrow **discrete scale inv.**

$$+ \text{---} \text{---} \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} \text{---} \text{---}$$

Λ

Three-Body Force: Limit Cycle

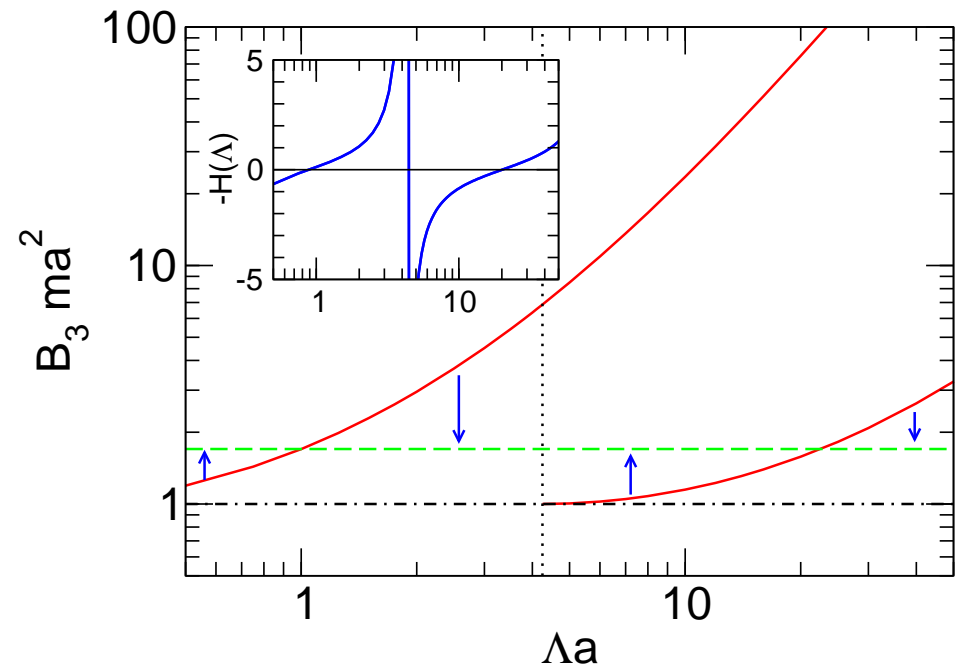
- RG invariance \implies running coupling $H(\Lambda) = g_3 \Lambda^2 / (9g_2^2)$

- $H(\Lambda)$ periodic: **limit cycle**

$$\Lambda \rightarrow \Lambda e^{n\pi/s_0} \approx \Lambda (22.7)^n$$

(cf. Wilson, 1971)

- **Anomaly:** scale invariance broken to discrete subgroup



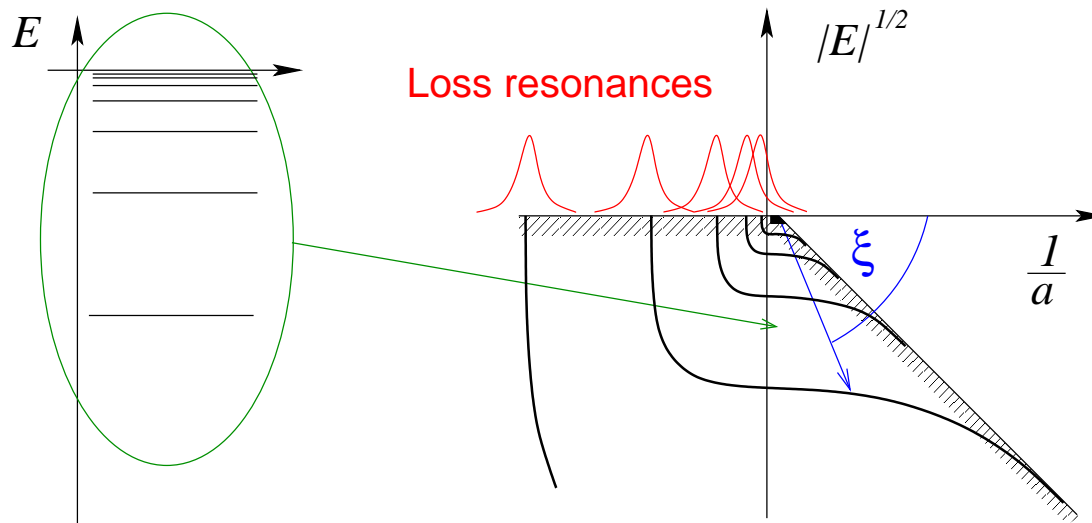
$$H(\Lambda) \approx \frac{\cos(s_0 \ln(\Lambda/\Lambda_*) + \arctan(s_0))}{\cos(s_0 \ln(\Lambda/\Lambda_*) - \arctan(s_0))}, \quad s_0 \approx 1.00624$$

(Bedaque, HWH, van Kolck, 1999)

- **Limit cycle** \iff **Discrete scale invariance** \iff **Efimov physics**

Limit Cycle: Efimov Physics

- Universal spectrum of three-body states (Efimov, 1970)

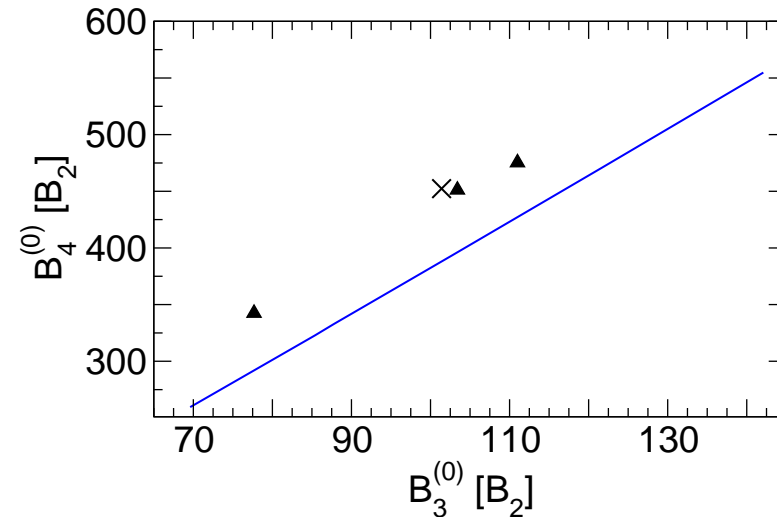
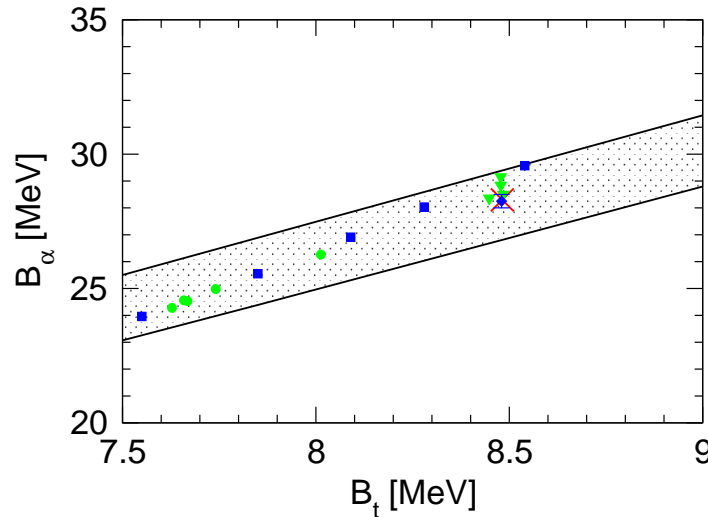


- Discrete scale invariance for fixed angle ξ
- Geometrical spectrum for $1/a \rightarrow 0$

$$B_3^{(n)} / B_3^{(n+1)} \xrightarrow{1/a \rightarrow 0} e^{2\pi/s_0} = 515.035\dots$$

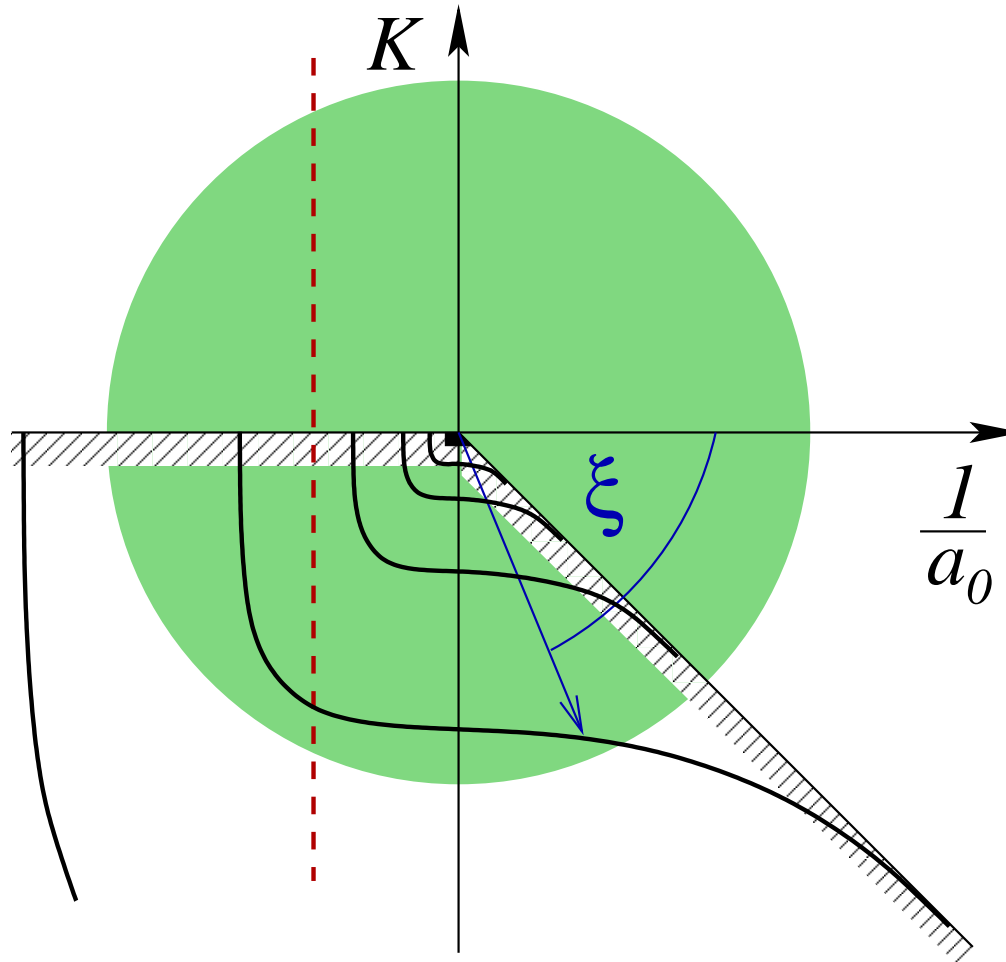
- Ultracold atoms \implies variable scattering length \implies loss resonances
- Nuclei \implies universal correlations and scaling relations

- 2 Parameters at LO \Rightarrow 3-body observables are correlated
 \Rightarrow Phillips line (Efimov, Tkachenko, 1985; Bedaque, HWH, van Kolck, 2000)
- No four-body parameter at LO (Platter, HWH, Meißner, 2004)
 \Rightarrow 4-body observables are correlated \Rightarrow Tjon line

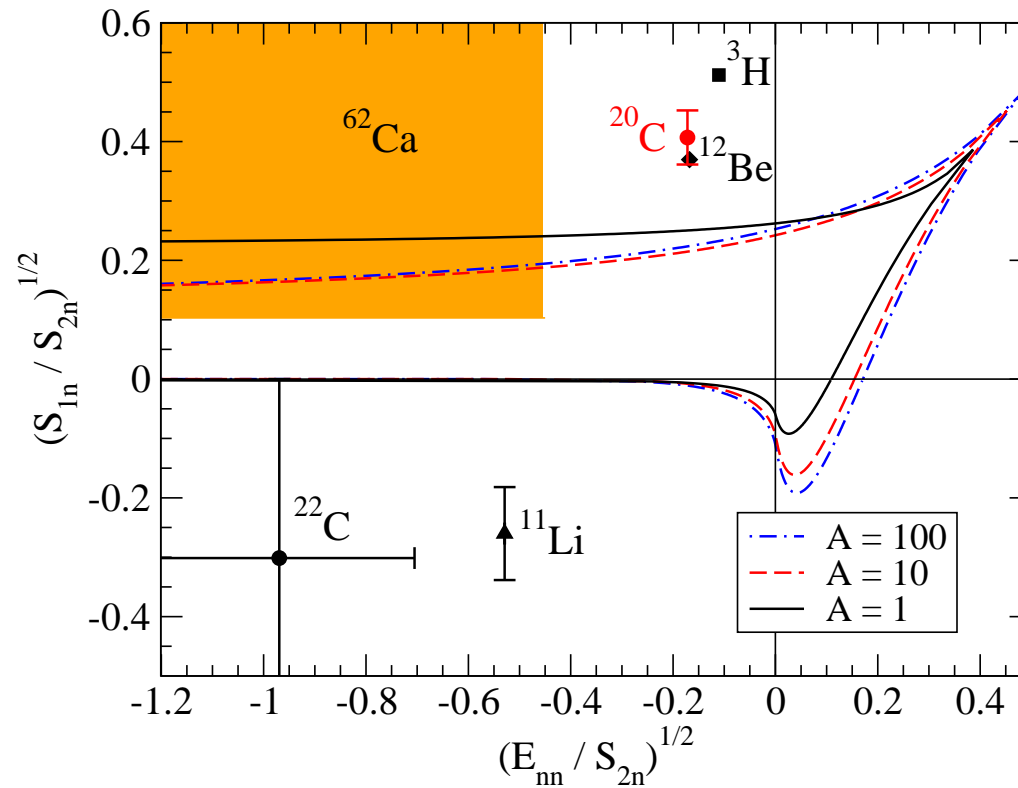


- Variation of 3-body parameter generates correlations
- RG-evolved interactions: Λ dependence traces correlations
(cf. Nogga, Bogner, Schwenk, 2004)

- Window of universality



- Efimov effect in halo nuclei? (Fedorov, Jensen, Riisager, 1994)
 \implies excited states obeying scaling relations
- Correlation plot: $E_{nn} \leftrightarrow S_{1n}$ (Amorin, Frederico, Tomio, 1997)



adapted from Canham, HWH, Eur. Phys. J. A **37** (2008) 367

- Hypertriton

- $np\Lambda$ bound state with $J^P = \frac{1}{2}^+$, $I = 0$
- Λd separation energy: $B^\Lambda = 0.13 \pm 0.05$ MeV
- total binding energy: $B_3^\Lambda = 2.35$ MeV

- EFT for large scattering lengths

⇒ shallow hypertriton follows naturally

- Leading order EFT ⇒ S-wave interactions

- ${}^3S_1(NN) + \Lambda \longrightarrow a_d \sim 1/\gamma_d$
- ${}^3S_1(\Lambda N) + N \longrightarrow a_3 \sim 1/\gamma_3$
- ${}^1S_0(\Lambda N) + N \longrightarrow a_1 \sim 1/\gamma_1$

- Scattering lengths large compared to interaction range

($NN \rightarrow \pi$ -exchange, $\Lambda N \rightarrow 2\pi$ -exchange)

- ΛN system unbound
- (Old) effective range analyses inconclusive (few data at relatively high energies)

$$0 > a_1 > -15 \text{ fm}$$

$$0 < r_1 < 15 \text{ fm}$$

$$-0.6 \text{ fm} > a_3 > -3.2 \text{ fm}$$

$$2.5 \text{ fm} < r_3 < 15 \text{ fm}$$

- Extractions using hyperon-nucleon potentials

$$a_1 \approx -2.9 \text{ fm}, \quad a_3 \approx -1.6 \text{ fm}, \quad \gg R \sim 1/(2m_\pi)$$

(chiral EFT: Haidenbauer et al., Nucl. Phys. A **915** (2013) 24)

- Characteristic three-body momentum

$$\gamma_3^\Lambda \sim 2\sqrt{|MB_3^\Lambda - \gamma_d^2|/3} \approx 14 \text{ MeV} \ll \sqrt{m_\Lambda(m_\Sigma - m_\Lambda)} \approx 300 \text{ MeV}$$

\Rightarrow $\Lambda\Sigma$ conversion is short range \implies three-body force

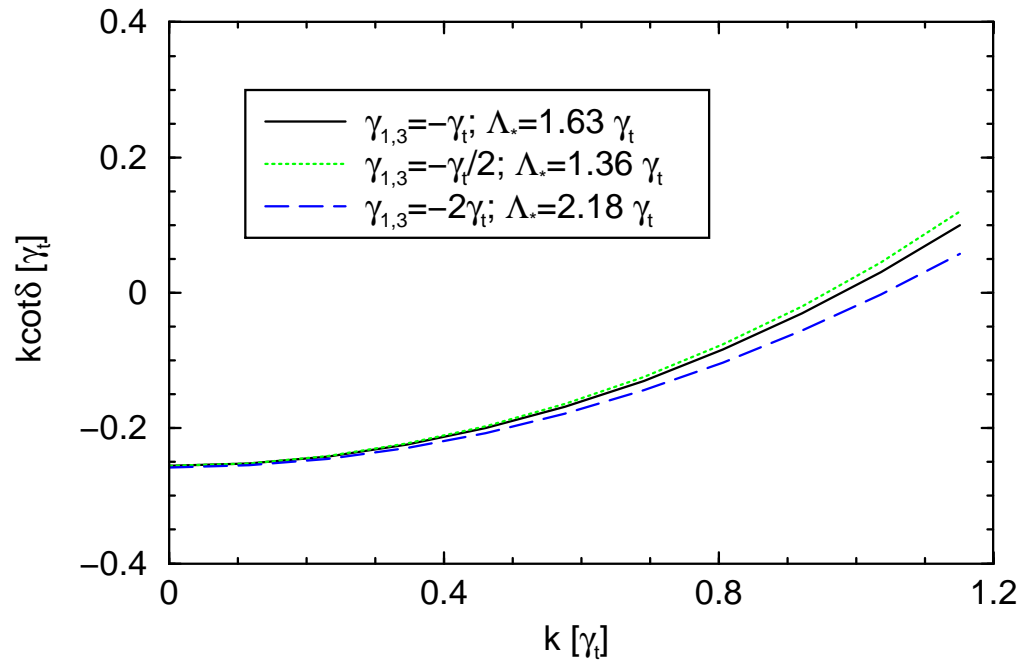
Integral Equations for Hypertriton



$$\begin{aligned}
 \overline{\overline{T_A}} &= \overline{\overline{T_B}} \text{ (with 3)} + \overline{\overline{T_C}} \text{ (with 1)} \\
 \overline{\overline{T_B}} \text{ (with 3)} &= \text{ (diagram)} + \overline{\overline{T_A}} \text{ (with 3)} + \overline{\overline{T_B}} \text{ (with 3)} + \overline{\overline{T_C}} \text{ (with 3)} \\
 \overline{\overline{T_C}} \text{ (with 1)} &= \text{ (diagram)} + \overline{\overline{T_A}} \text{ (with 1)} + \overline{\overline{T_B}} \text{ (with 1)} + \overline{\overline{T_C}} \text{ (with 1)}
 \end{aligned}$$

HWH, Nucl. Phys. **A705** (2002) 173; Hildenbrand, HWH, in preparation

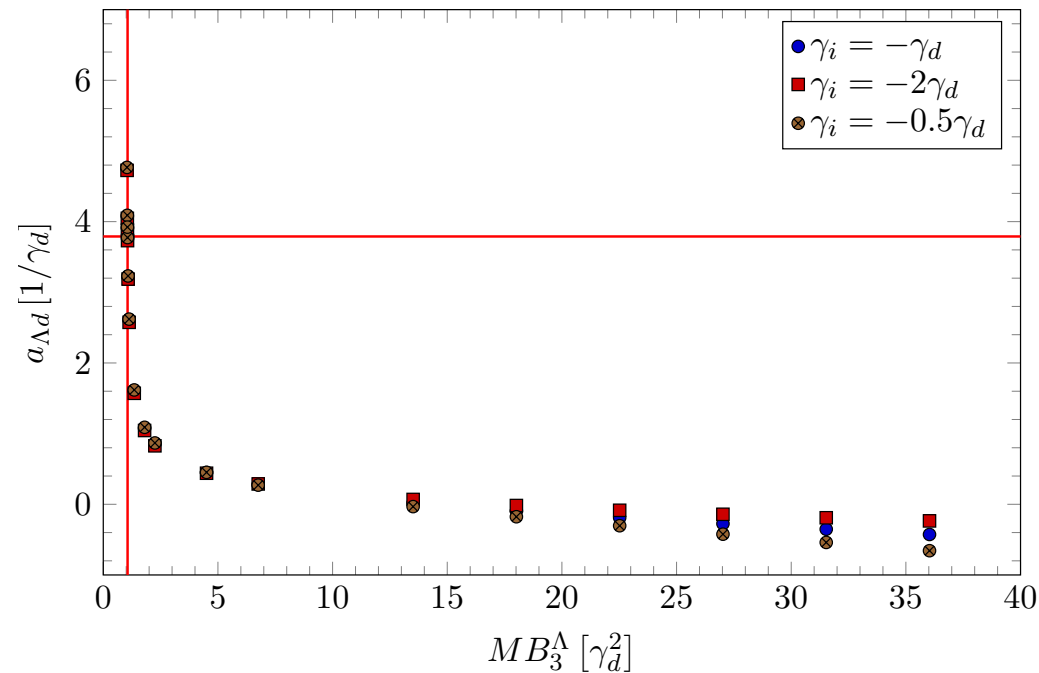
- **Strong cutoff dependence** \implies renormalize with Λ_{np}
three-body force (cf. triton, bosons)
- **Limit cycle** with $s_0 = 1.0076$ (unequal masses)
- **Scaling factor:** $\exp \pi / s_0 \approx 22.60$
- **Corrects error in original publication**
- **No room for excited states....**



HWH, Nucl. Phys. **A705** (2002) 173

- Exact value of γ_i not determined by B_3^Λ
- Phase shifts independent of γ_i \iff shallowness of hypertriton
- Low-energy parameters:

$$a_{\Lambda d} = (16.8^{+4.4}_{-2.4}) \text{ fm} \quad \text{and} \quad r_{\Lambda d} = (2.3 \pm 0.1) \text{ fm}$$

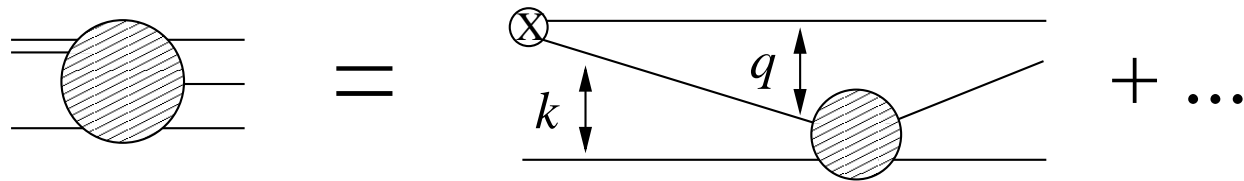


Hildenbrand, HWH, in preparation

- Correlation between hypertriton triton binding energy and $S = 1/2$ Λd scattering length (cf. Phillips '68)
- Sensitivity to specific values of γ_i only for deeper binding
- Hypertriton wave function can also be extracted
- Production in heavy ion collisions?

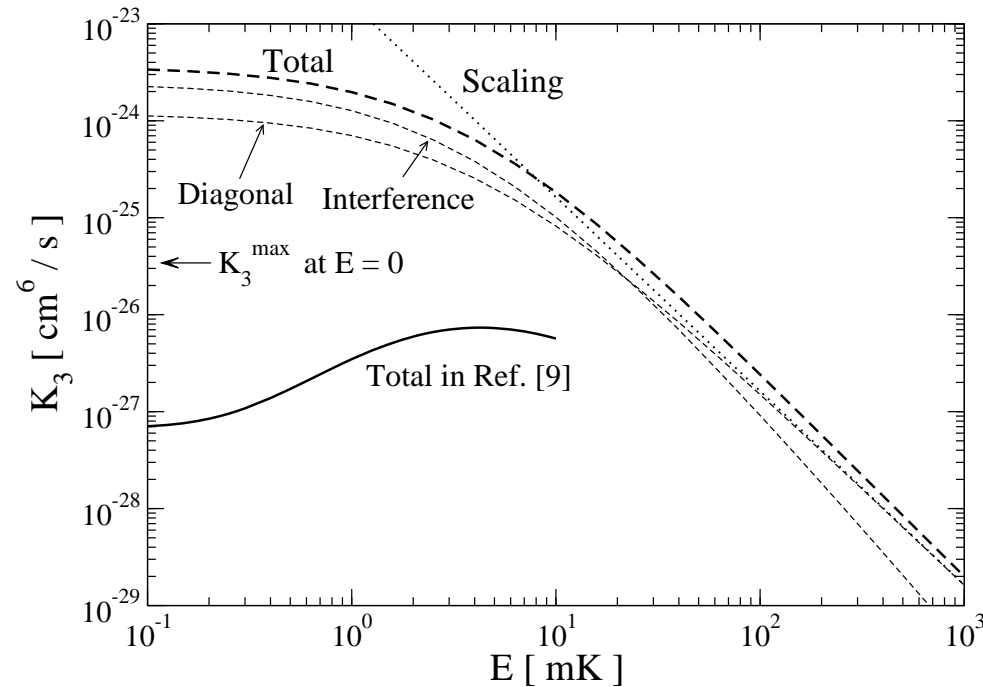
- Breakup and recombination reactions of shallow atomic bound states at higher energy E
(Braaten, Zhang, Phys. Rev. A **73**, 042707 (2006))
- Factorization of amplitude if typical momenta between atom and bound state satisfy $k \gg 1/a \sim q$, $E \sim k^2/m$

$$\mathcal{T} \approx \frac{4(\pi/a)^{1/2}}{q^2 + 1/a^2} \times \frac{8\pi/m}{(\frac{1}{2}r_e k^2 + \dots) - ik}$$



- Final state interaction suppressed
- Works well for breakup reactions of ^4He atoms
- Compare with exact calculation for recombination

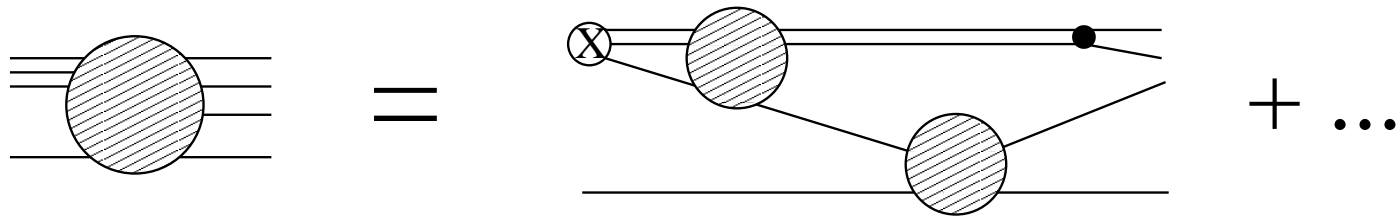
- Three-body recombination rate coefficient K_3



(Braaten, Zhang, Phys. Rev. A **73**, 042707 (2006))

- Interference terms important at low energies ($B_2 = 1.6$ mK)
- Higher orders in atom-atom scattering amplitude

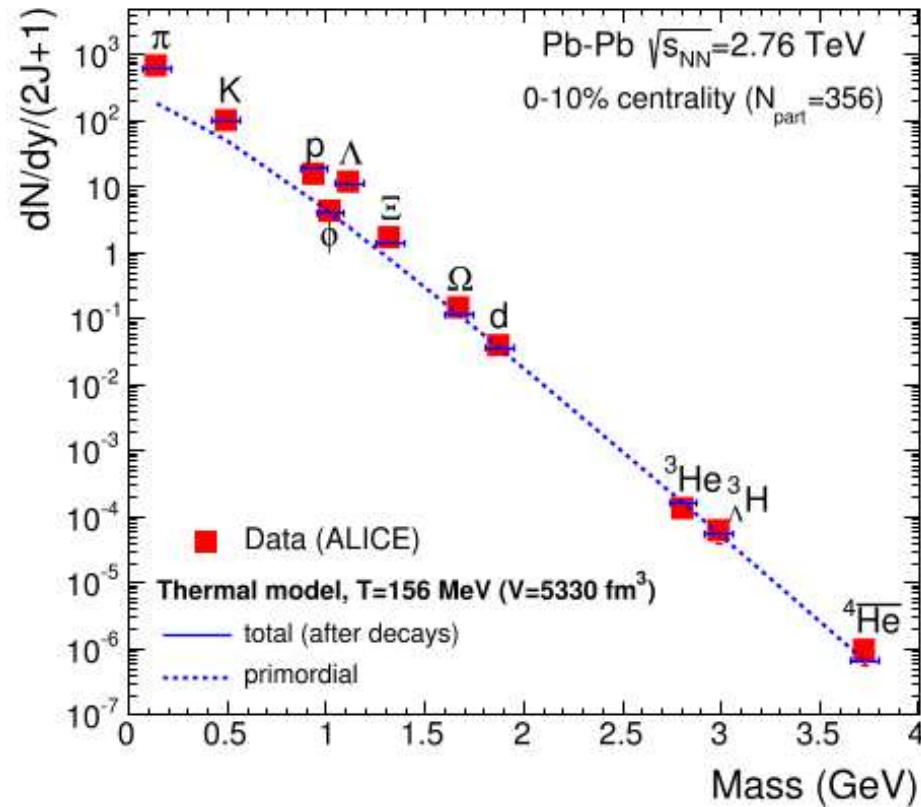
- Extension to three-body bound states



- Requires momentum space wave function of trimers
 - ⇒ three-body problem requires numerical solution
 - ⇒ dependence on a and Λ_*
- Extension to higher-body states possible
- Only bound state calculations with $N - 1$ bodies required
- Closure relation for sum over all (initial/final) states

$$\int d^3r \sum_{X_{N-1}} |\langle X_{N-1} | \psi(r) | M \rangle|^2 = \int d^3r \langle M | \psi^\dagger \psi(r) | M \rangle = N$$

- Application to heavy ion collisions?



Andronic et al. (2016)

- Short-distance factor cancels in ratios?

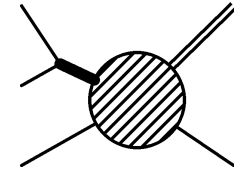
- Effective field theory for unitary limit
- Universal aspects of (Discrete) Scale Invariance \Leftrightarrow Efimov physics
 - Effective field theory for threshold states
 -
- Applications in atomic, nuclear, and particle physics
 - Cold atoms close to Feshbach resonance
 - Few-body nuclei: triton, **hypertriton**, halo nuclei, ...
 - Hadronic molecules: $X(3872)$, ...
- Factorization for breakup and recombination reactions
 - Application to production of weakly-bound objects in heavy ion collisions?

Additional Slides



- Three-body recombination:

3 atoms \rightarrow dimer + atom \Rightarrow **loss of atoms**



- Recombination constant: $\dot{n}_A = -K_3 n_A^3$

- K_3 has log-periodic dependence on scattering length

(Nielsen, Macek, 1999; Esry, Greene, Burke, 1999; Bedaque, Braaten, HWH, 2000)

- Deep dimers: Efimov trimers acquire width \Rightarrow **resonances**

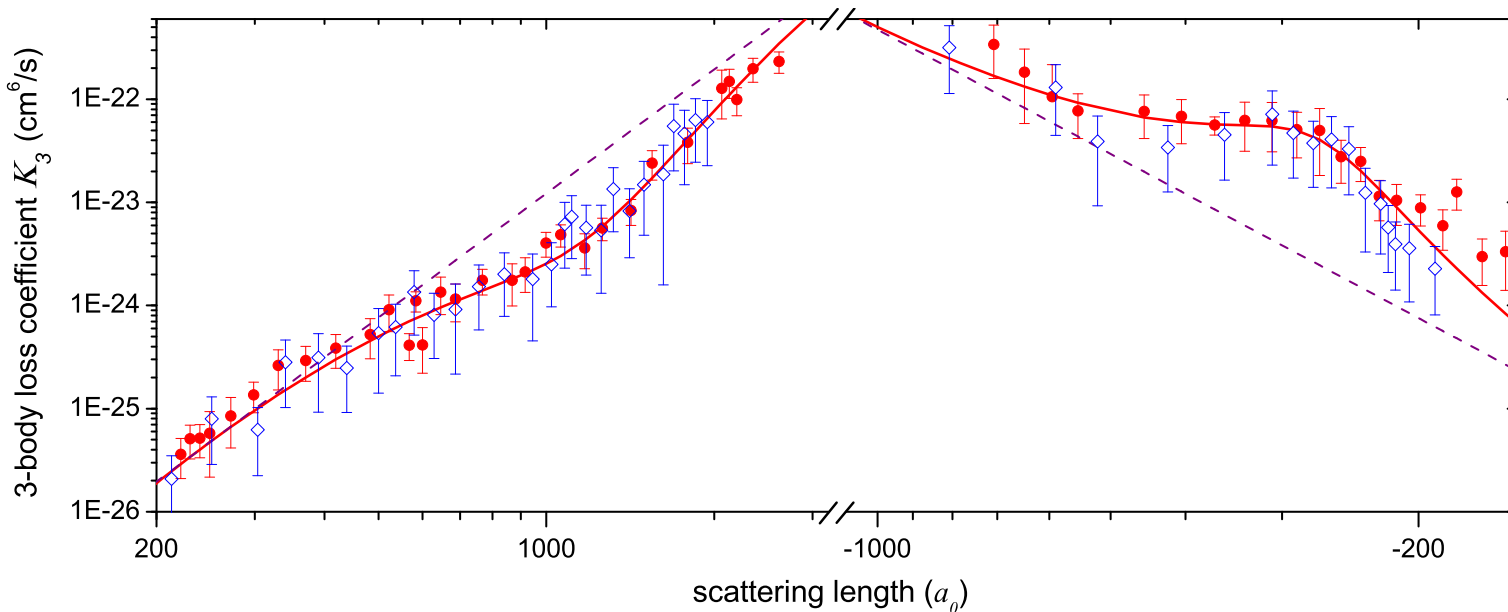
- Loss term in short distance b.c.: $\Lambda_* \rightarrow \Lambda_* \exp^{i\eta_*/s_0}$

- Universal line shape of recombination resonance ($a < 0$, $T = 0$)

$$K_3^{deep} = \frac{64\pi^2(4\pi - 3\sqrt{3}) \coth(\pi s_0) \sinh(2\eta_*)}{\sin^2 [s_0 \ln(a/a_-)] + \sinh^2 \eta_*} \frac{\hbar a^4}{m}, \quad s_0 \approx 1.00624..$$

and other observables ...

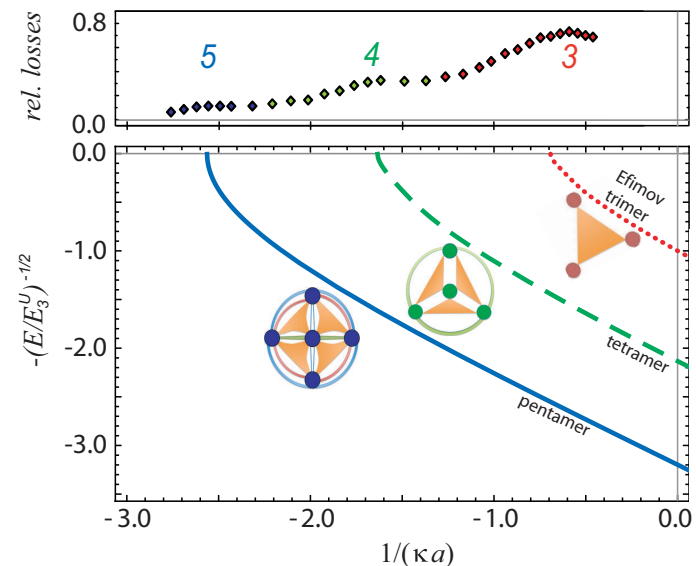
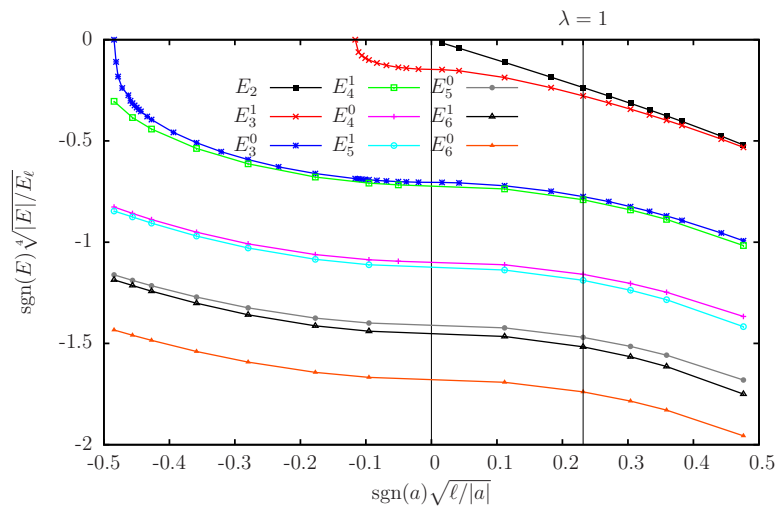
- First experimental evidence in ^{133}Cs (Krämer et al. (Innsbruck), 2006)
now also ^6Li , ^7Li , ^{39}K , $^{41}\text{K}/^{87}\text{Rb}$, $^6\text{Li}/^{133}\text{Cs}$
- Example: Efimov spectrum in $^6\text{Li}/^{133}\text{Cs}$ mixture
(Gross et al. (Bar-Ilan Univ.), Phys. Rev. Lett. **105** (2010) 103203)



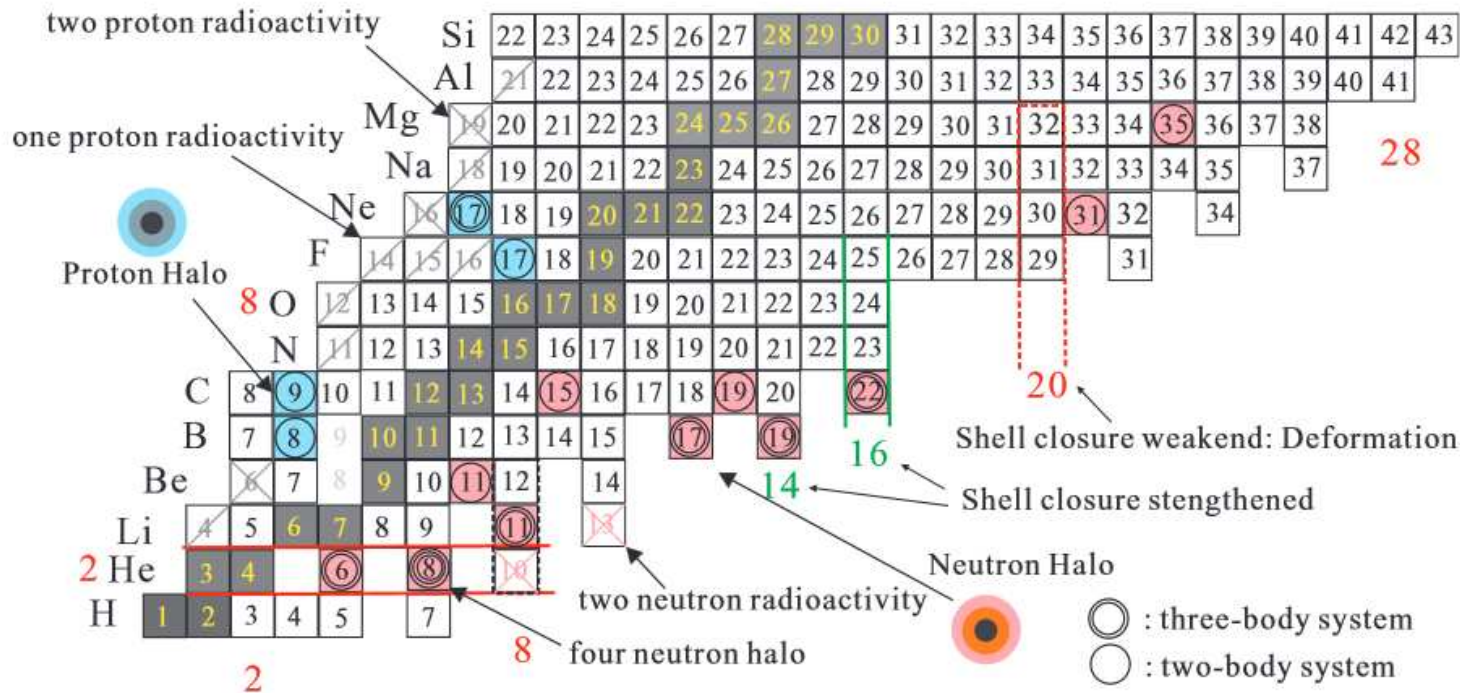
- Van der Waals tail determines $a_-/l_{vdW} \approx -10$ ($\pm 15\%$)
(Wang et al., 2012; Naidon et al. 2012, 2014; ...)
... but not η_* ...

Universal Tetramers and Beyond

- **Universal tetramers:** $B_4^{(0)} = 4.610(1) B_3$, $B_4^{(1)} = 1.00227(1) B_3$
(Platter, HWH, 2004, 2007; von Stecher et al., 2009; Deltuva 2010-2013)
- **Two tetramers attached to each trimer**
- **Universal states up to $N = 16$ calculated**
(von Stecher, 2010, 2011; Gattobigio, Kievsky, Viviani, 2011-2014)
- **Observation up to $N = 5$ in Cs losses** (Grimm et al. (Innsbruck), 2009, 2013)



- Low separation energy of valence nucleons: $B_{valence} \ll B_{core}, E_{ex}$
 → close to “nucleon drip line” → **scale separation** → EFT

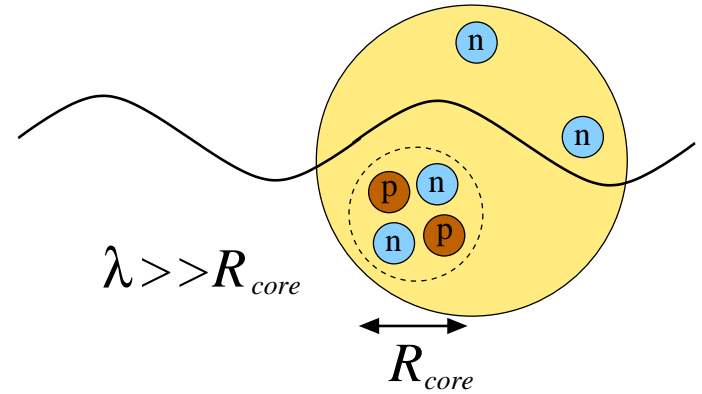


C.-B. Moon, Wikimedia Commons

- EFT for halo nuclei

(Bertulani, HWH, van Kolck, 2002; Bedaque, HWH, van Kolck, 2003; ...)

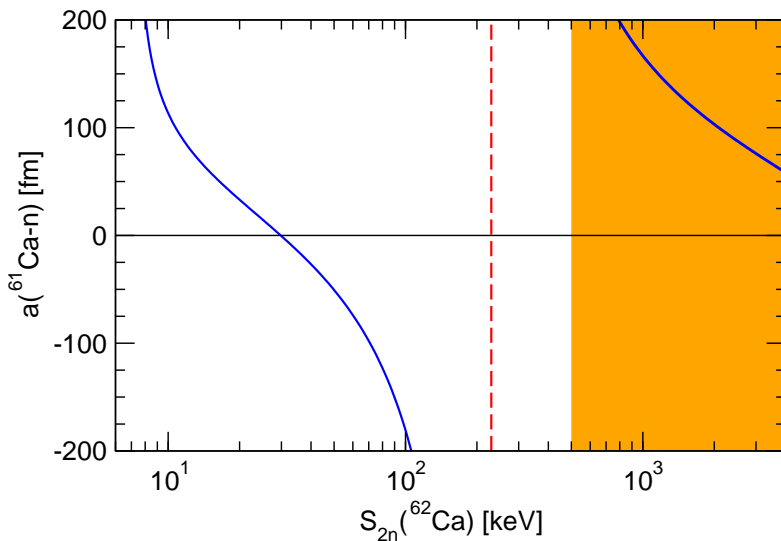
- Scales: $E \sim p^2/(2\mu) \sim 1/(2\mu R^2)$
- Separation of scales:
 $1/k = \lambda \gg R_{core}$
- Limited resolution at low energy:
→ expand in powers of kR_{core}
- Short-distance physics not resolved
→ capture in low-energy constants using renormalization
→ include long-range physics explicitly
- Systematic, model independent \implies universal properties
- Very low energies: only short-range physics \implies pionless EFT
- Exploit cluster substructures \implies Halo EFT



(G. Hagen, P. Hagen, HWH, Platter, Phys. Rev. Lett. **111** (2013) 132501)

- From many to few: emergence of halo degrees of freedom
- Coupled cluster calculations of ^{60}Ca and ^{61}Ca using chiral N2LO two-body force and schematic three-body force:

⇒ ^{61}Ca is a weakly bound S-wave state (or virtual state)



- Prospects for excited Efimov states in ^{62}Ca :

$$S_{\text{deep}} = 1/(\mu_{cn} r_{cn}^2) \approx 500 \text{ keV}$$

scaling factor $\lambda_0 \approx 16$

⇒ possible if $S_{2n} \gtrsim 230 \text{ keV}$