



Application of CLTDs* for the investigation of Z-yield distributions of fission fragments

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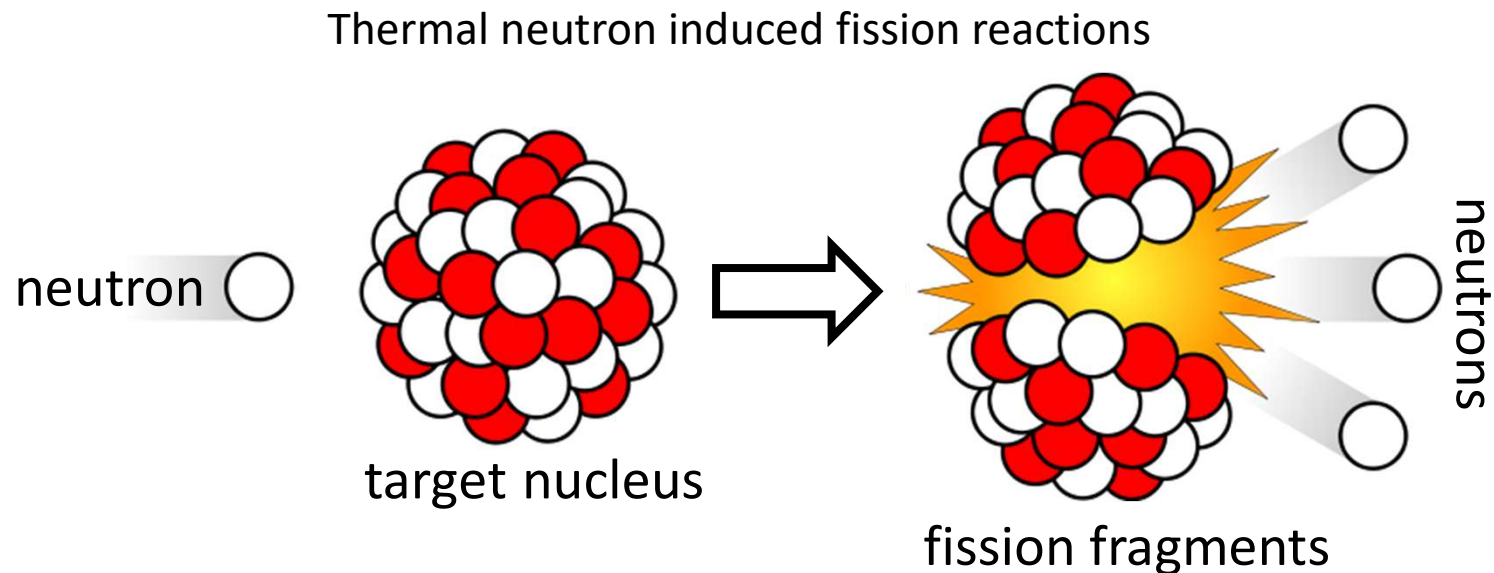
*CLTDs – Calorimetric Low Temperature Detectors

Motivation

motivation for studying fission fragment yields:

- Various nuclear energy applications
- better understanding of the nuclear fission process
- data relevant for reactor physics

Application of CLTDs for fission studies!



Outline

❖ Calorimetric Low Temperature Detectors:

- Concept
- Advantages
- Design

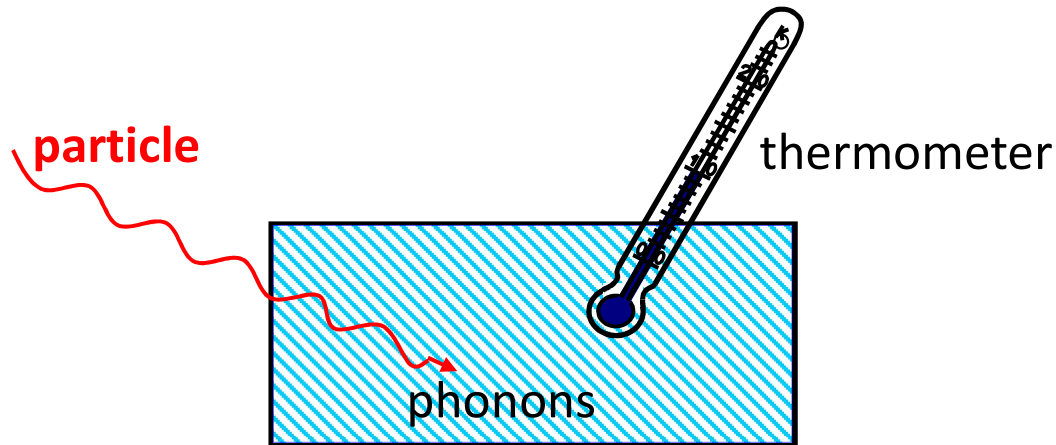
❖ Nuclear Charge Distribution Measurements:

- Experimental Set-up
- Preliminary results:
 - First time direct Z-yield determination for heavy fragment masses
 - Using passive absorber method
 - Proton odd-even effect in masses approaching symmetric fission
 - for ^{239}Pu & ^{241}Pu
 - Important for nuclear model description near scission
 - Precise ^{92}Rb & ^{96}Y yields
 - for ^{235}U , ^{239}Pu , ^{241}Pu (n_{th} , f)
 - Important for reactor antineutrino anomaly studies

CLTDs: Introduction

Calorimetric Detector → detection of heat

Idea: detection of particle energy independent of ionization processes



Interaction of ions with matter:

- Primary:

ionization → ionization detectors

- Secondary:

thermalization → calorimetric detectors

(conversion of energy to heat

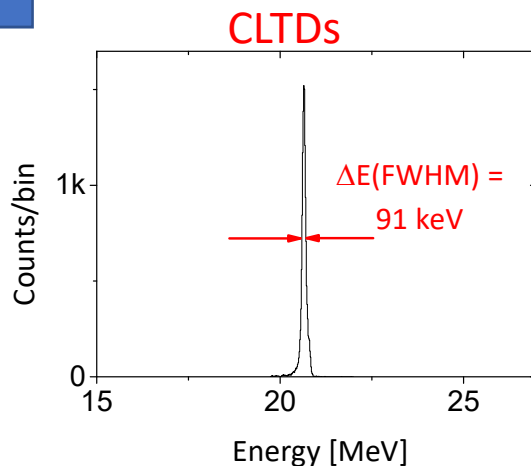
⇒ detection of thermal phonons)

Potential Advantages:

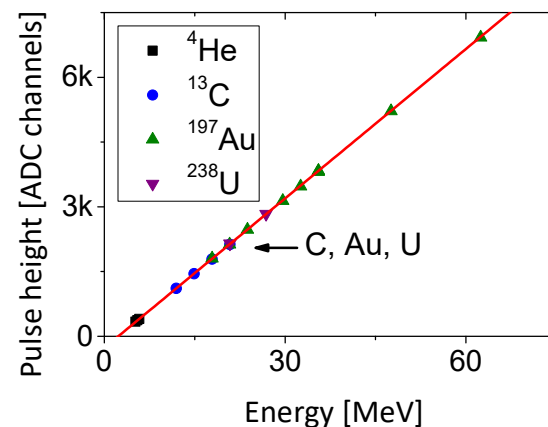
- energy resolution
- energy linearity
- detection threshold
- variety of detector materials

CLTDs: Advantages

Energy Resolution
(^{238}U @ 20.7 MeV)

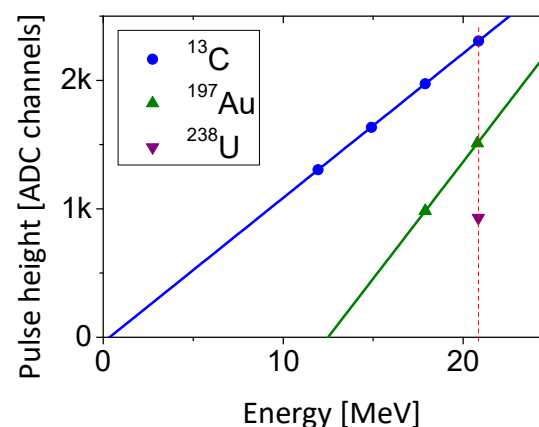
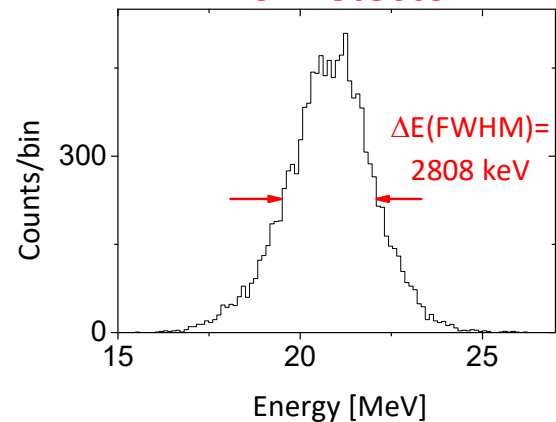


Pulse Height
Defect
(^{13}C , ^{197}Au , ^{238}U)



More complete energy detection
→ Better resolution and linearity

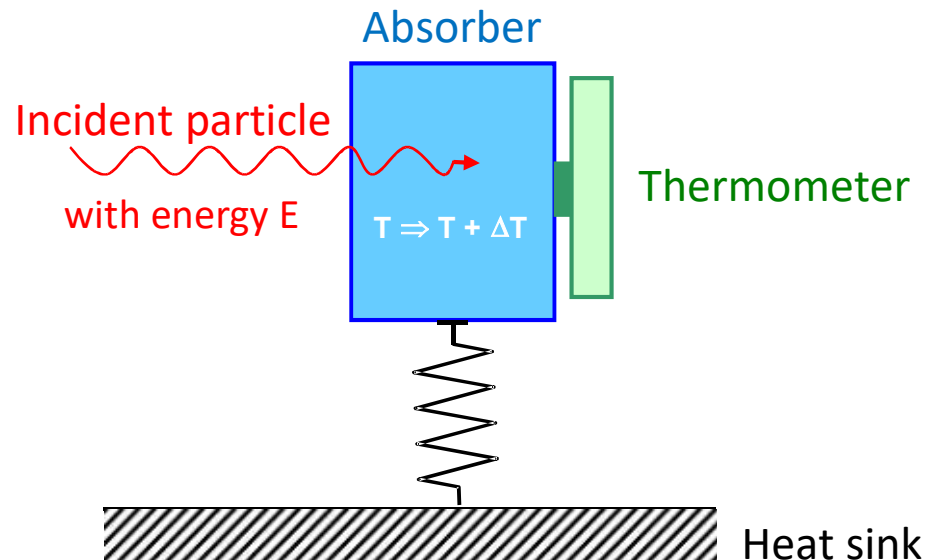
Si Detector



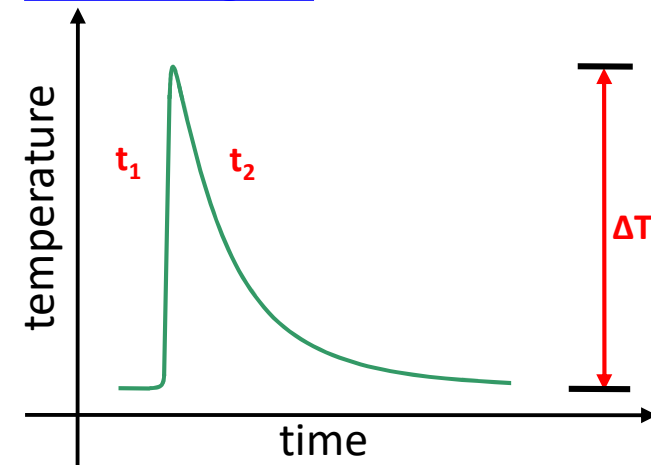
Losses at energy conversion:
Due to **dead layer**,
charge recombination

Plots from: Kraft-Bermuth et al., Rev. Sci. Instrum. 80, 103304 (2009)

CLTDs: Detection Principle



thermal signal:



Absorber:

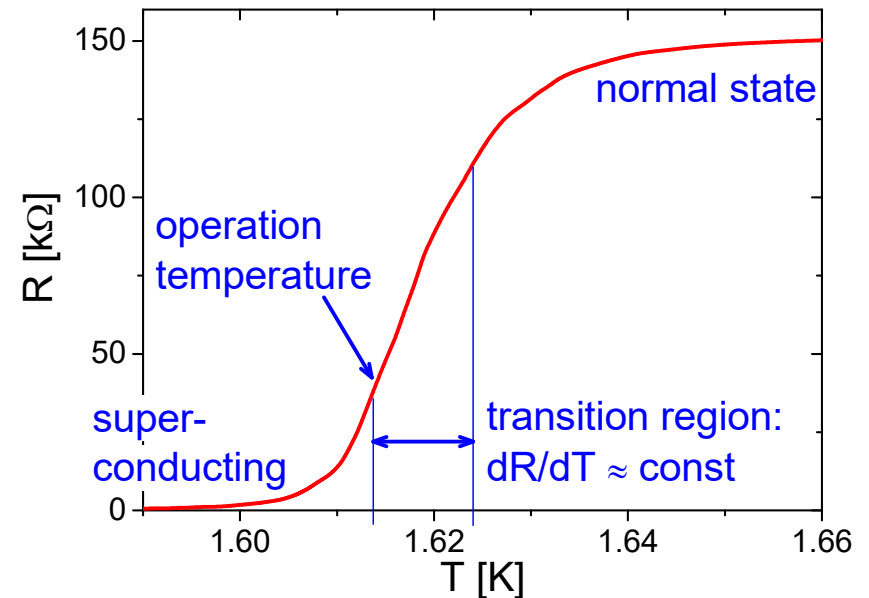
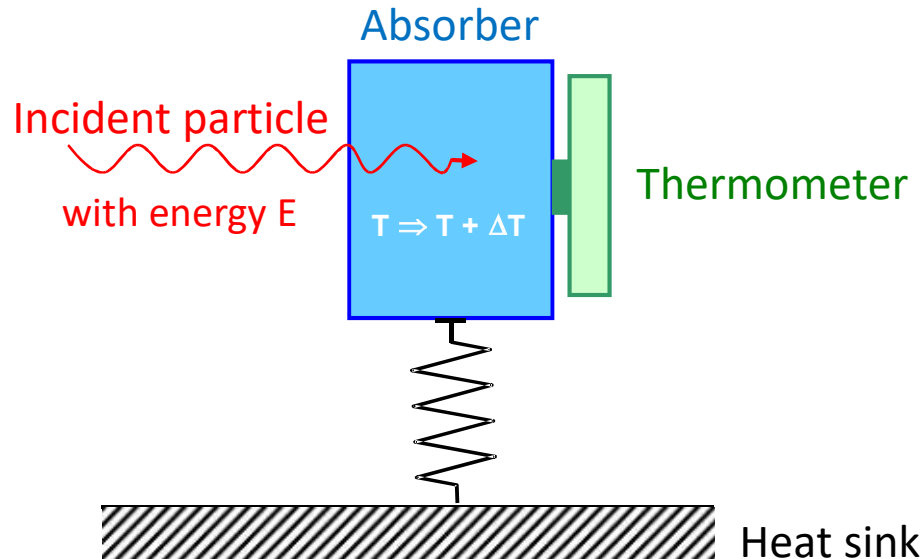
Particle energy is converted into heat,

$$\Delta T = \frac{E}{c \cdot m}$$

- High sensitivity for small specific heat c and small mass
- $c \sim (T^3/\Theta_D^3)$ from Debye law (for insulators, superconductors)
- **Operation at Low Temperatures**

Rise-Time	1 – 10 μ s
Decay-Time (C/k)	100 μ s – 10 ms

CLTDs: Detection Principle

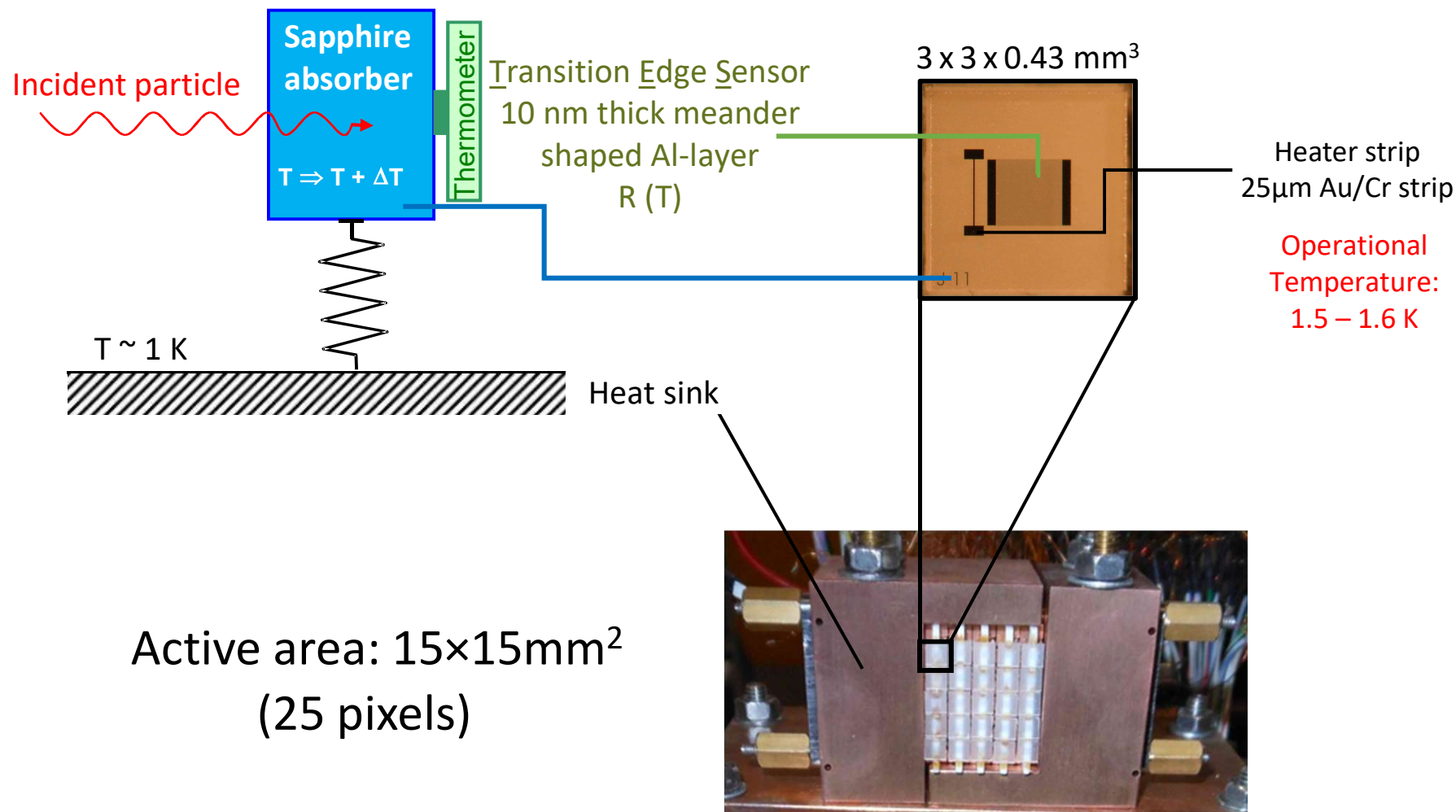


Thermometer:

Transition Edge sensors are used.

- Resistance (**superconductor**) biased with constant current.
- $\Delta T \rightarrow \Delta R(T) \rightarrow \Delta V$ voltage signal is detected.
- Detector temperature stabilized in **phase transition region** $\rightarrow$ high dR/dT .

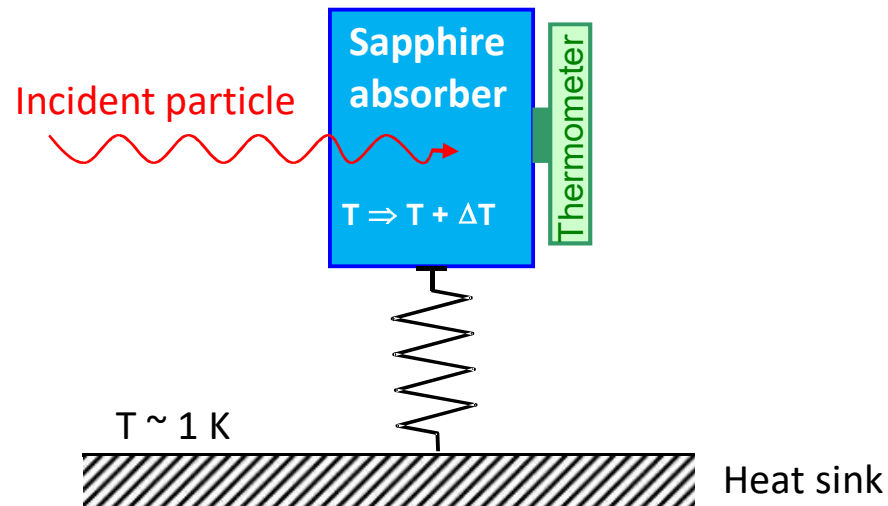
CLTDs: Current Detector Design



S. Kraft-Bermuth et al., Rev. Sci. Instrum. **80**, 2009

A. Echler et al., J. Low Temp. Phys. **167**, 2012,

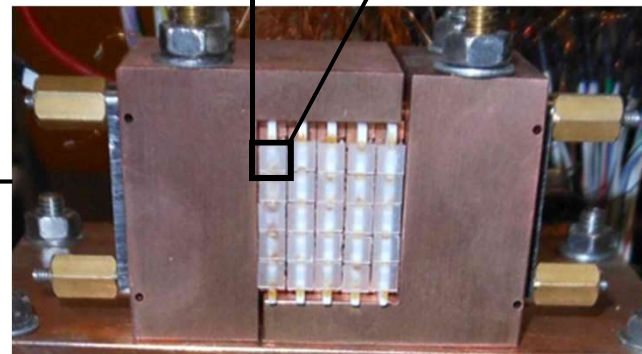
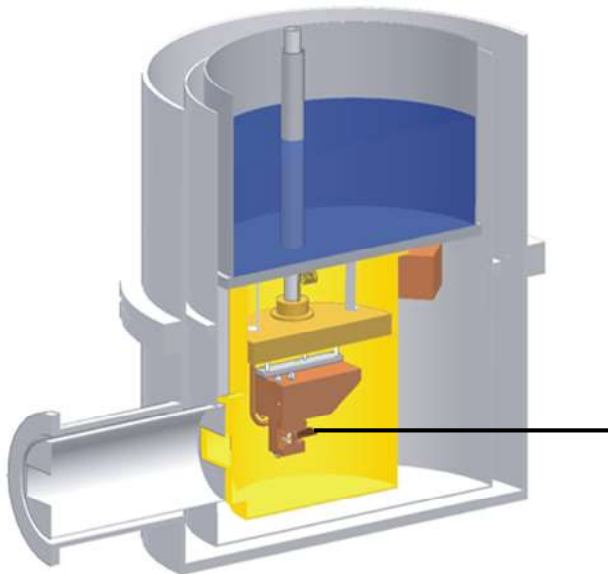
CLTDs: Current Detector Design



windowless ^4He bath cryostat



Operational
Temperature:
1.5 – 1.6 K



Z- distributions of fission fragments (n_{th}, f)

Methods for Z- separation:

➤ Radio chemistry

➤ γ -spectroscopy

⇒ restricted to particular nuclides

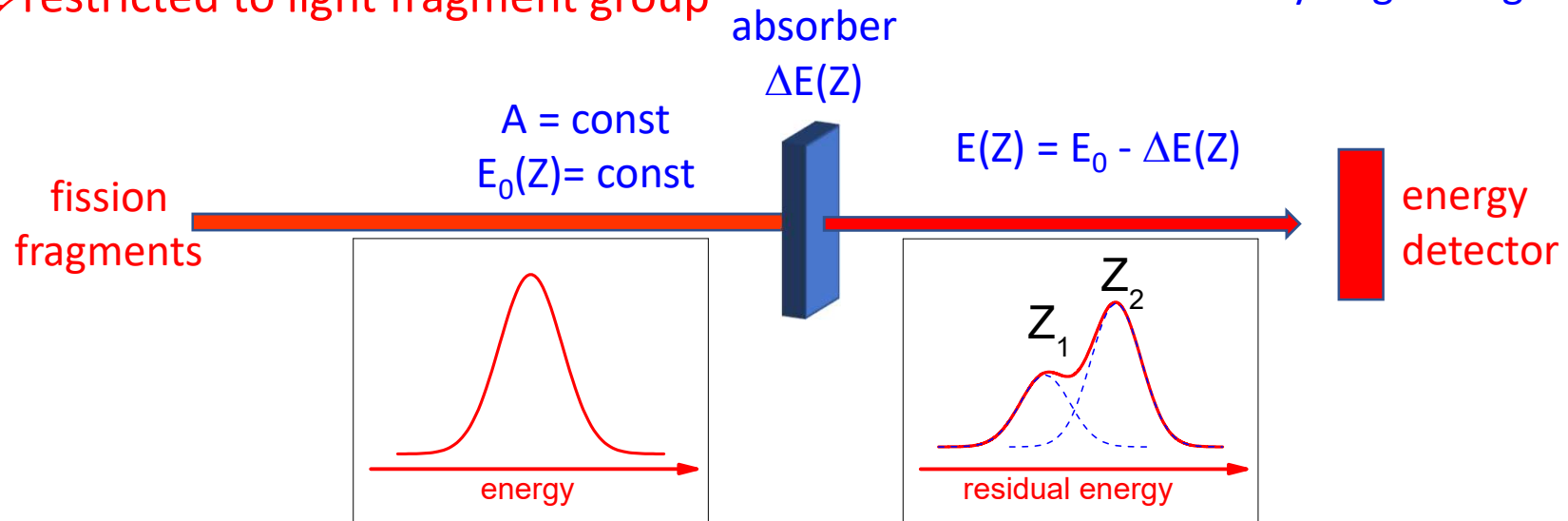
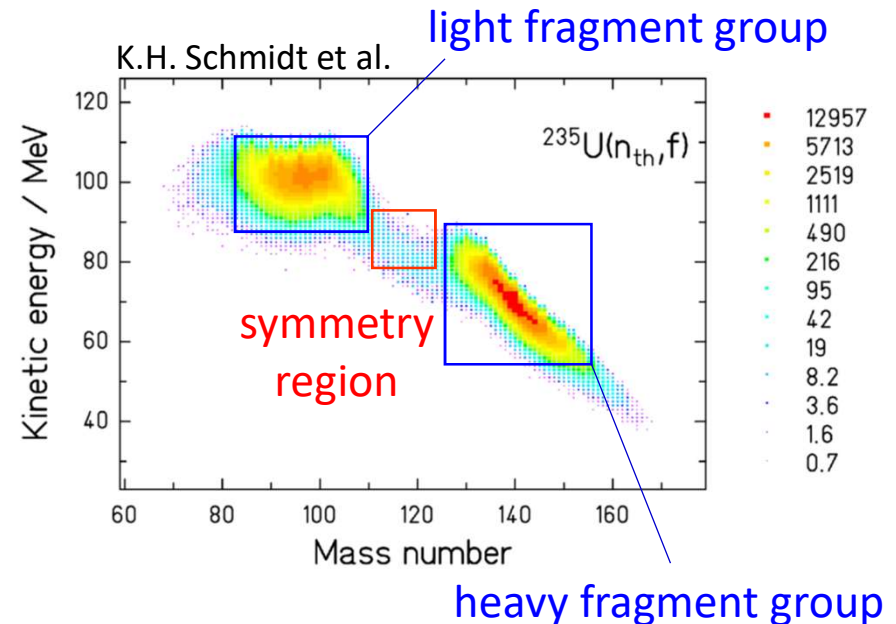
➤ Passive-absorber method

U. Quade et al., Nucl. Phys. A487 (1988) 1

best Z-resolution with:

- Parylene C absorbers
- ionization chamber

⇒ restricted to light fragment group



Z- distributions of fission fragments (n_{th}, f)

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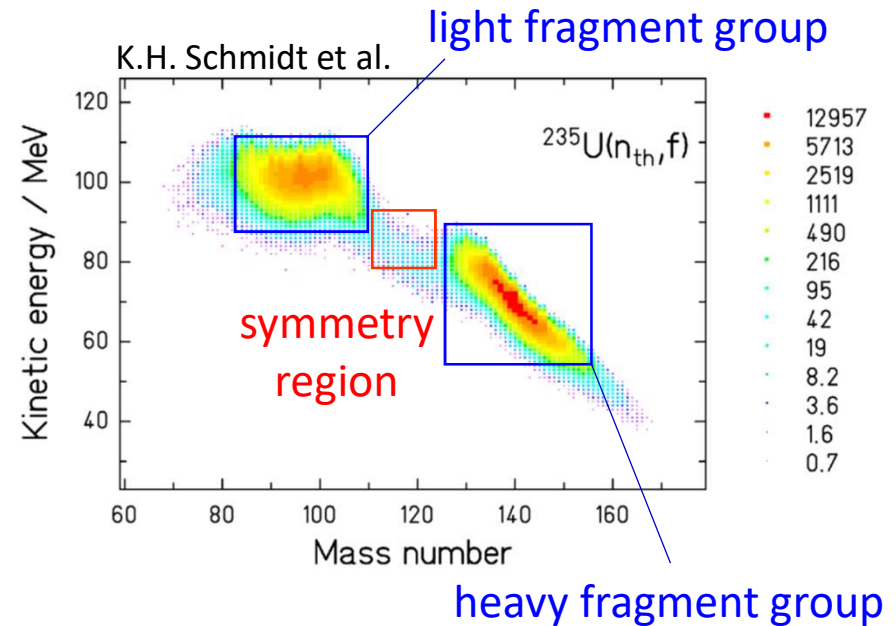
U. Quade et al., Nucl. Phys. A487 (1988) 1

best Z-resolution with:

- Parylene C absorbers

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⇒ restricted to light fragment group



Preliminary results from CLTDs using SiN absorber foils:

➤ First time possible to measure in the heavy mass region.

➤ New measurements in the light mass region approaching symmetric fission for $^{239}\text{P}(n_{th}, f)$ and $^{241}\text{Pu}(n_{th}, f)$.

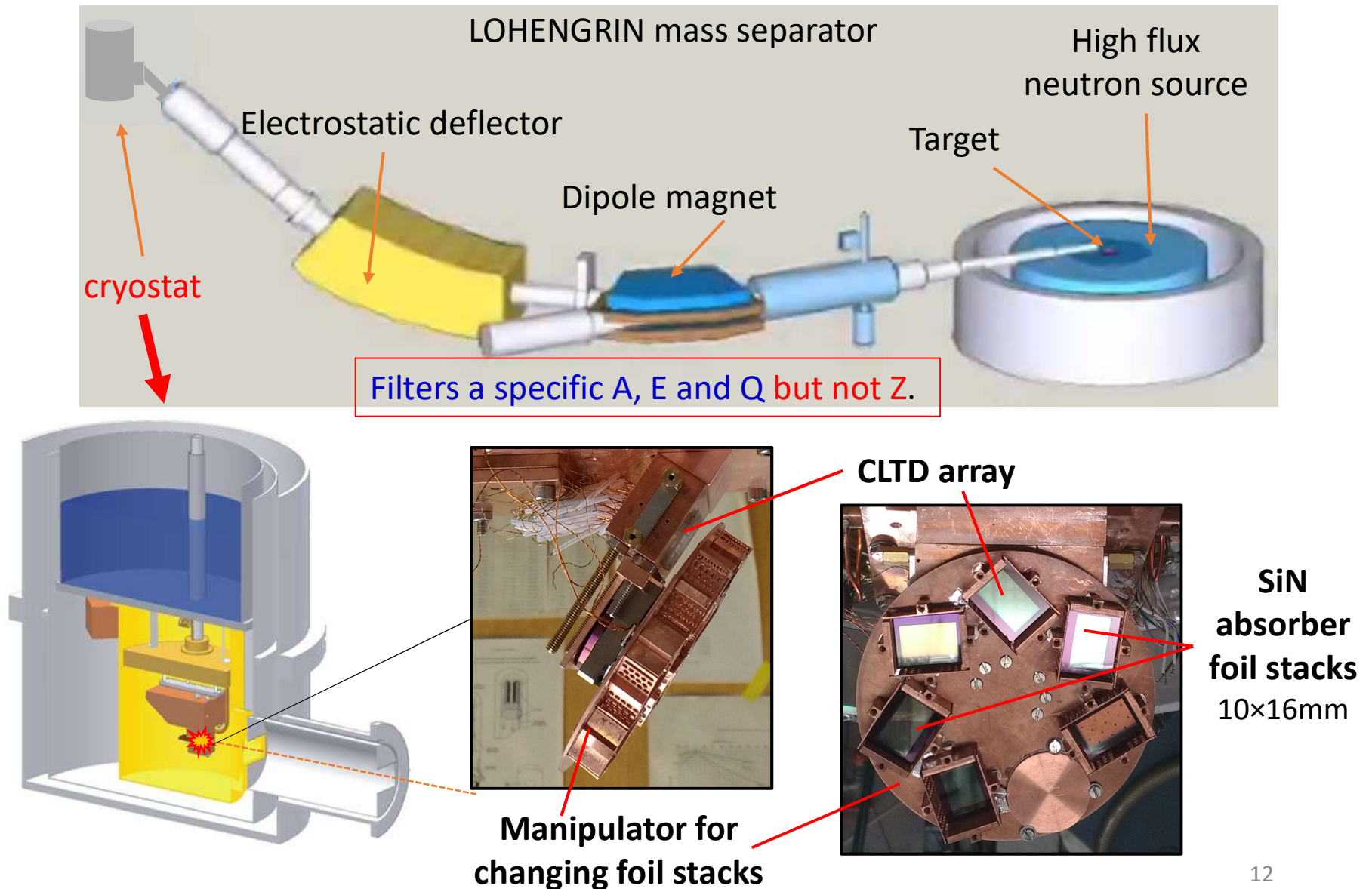
➤ Precise yield determination of ^{92}Rb , ^{96}Y .

SiN foils:

- Homogeneity
- High thermal stability
- Extreme hardness
- Chemical inertness

Investigation of Z- distributions of fission fragments (n_{th}, f)

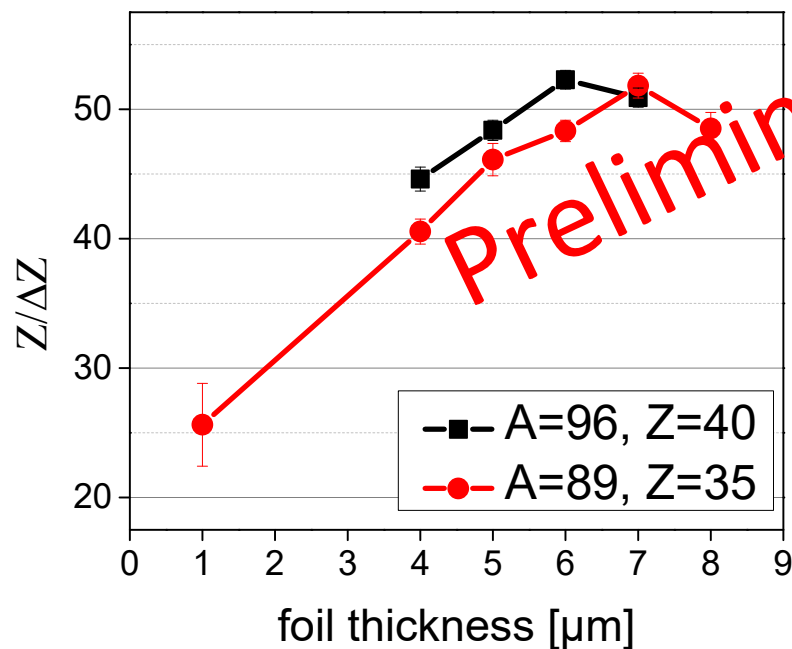
Experimental Set-up



Z-Resolution: Optimum thickness of SiN

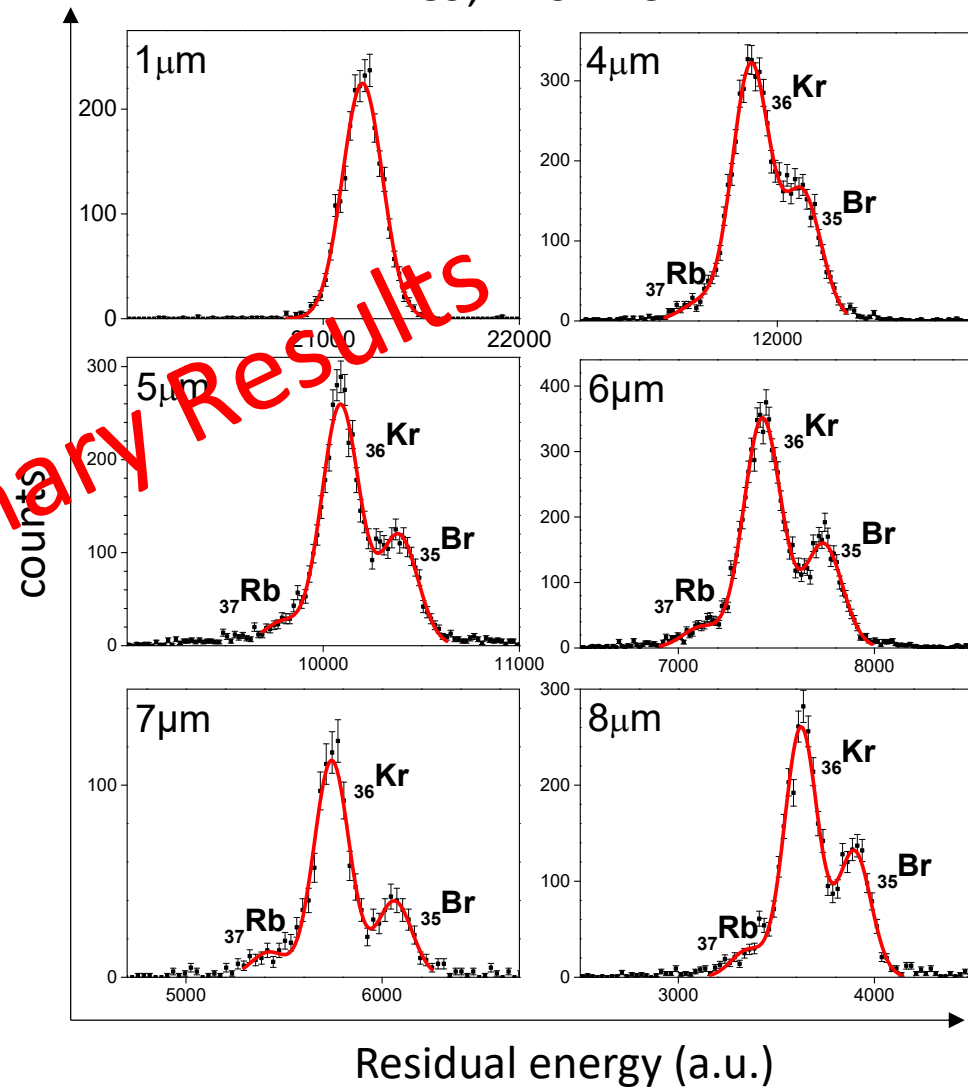
Systematic studies:
Z-Resolution vs SiN foil
thickness

$$Z/\Delta Z := Z \times \frac{\delta E(Z) - \delta E(Z - 1)}{FWHM}$$



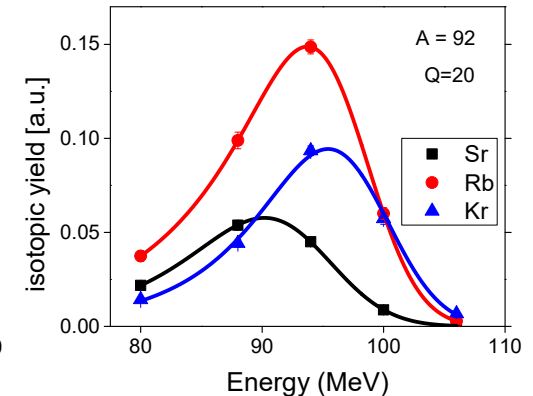
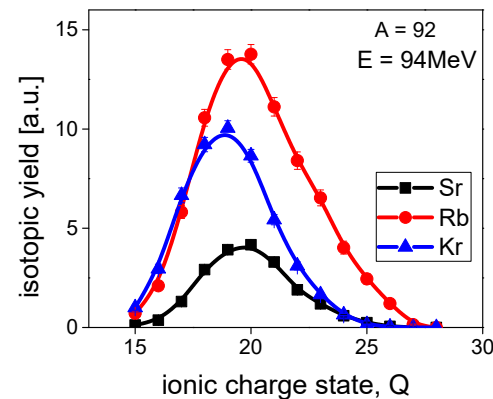
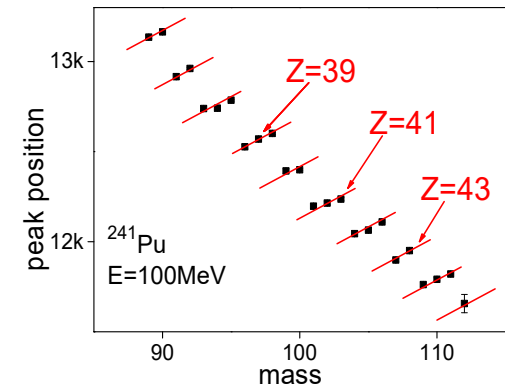
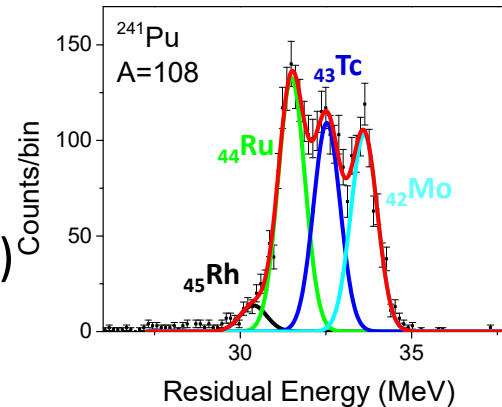
residual energy spectra

$A = 89, E = 94\text{MeV}$



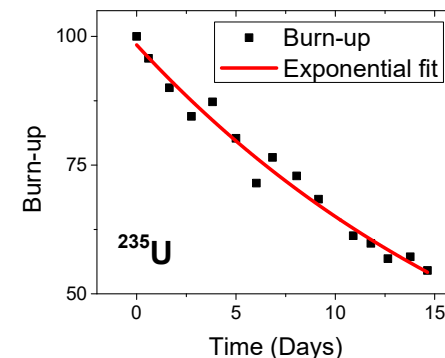
Procedure for determining Z-yields

- fit of spectrum
 - relative Z-Yield
 - line shape from Tandem experiment (for heavy masses)
- Z-Identification
- take into account
 - energy dependence
 - electronic charge state dependence (ns-isomers)
- absolute normalization
 - target burnup



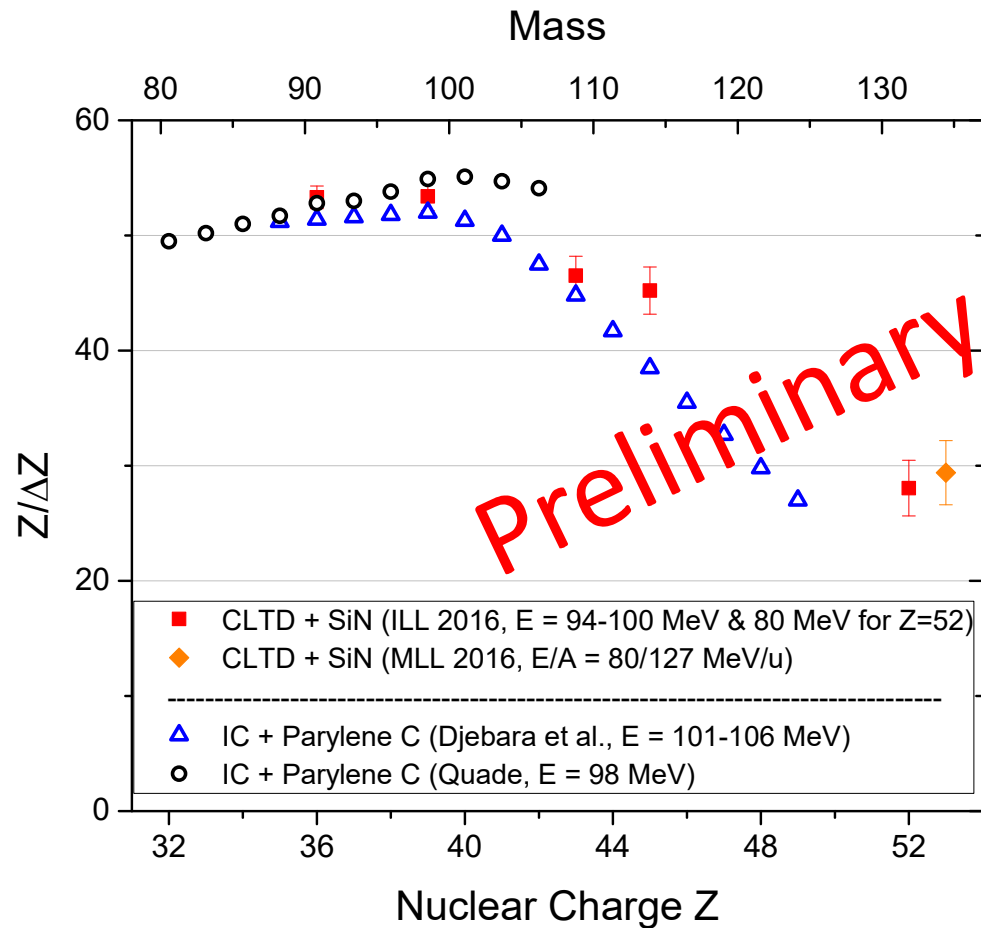
for Z-Yield of one mass:

⇒ upto 300 spectra to be analyzed



Resolution

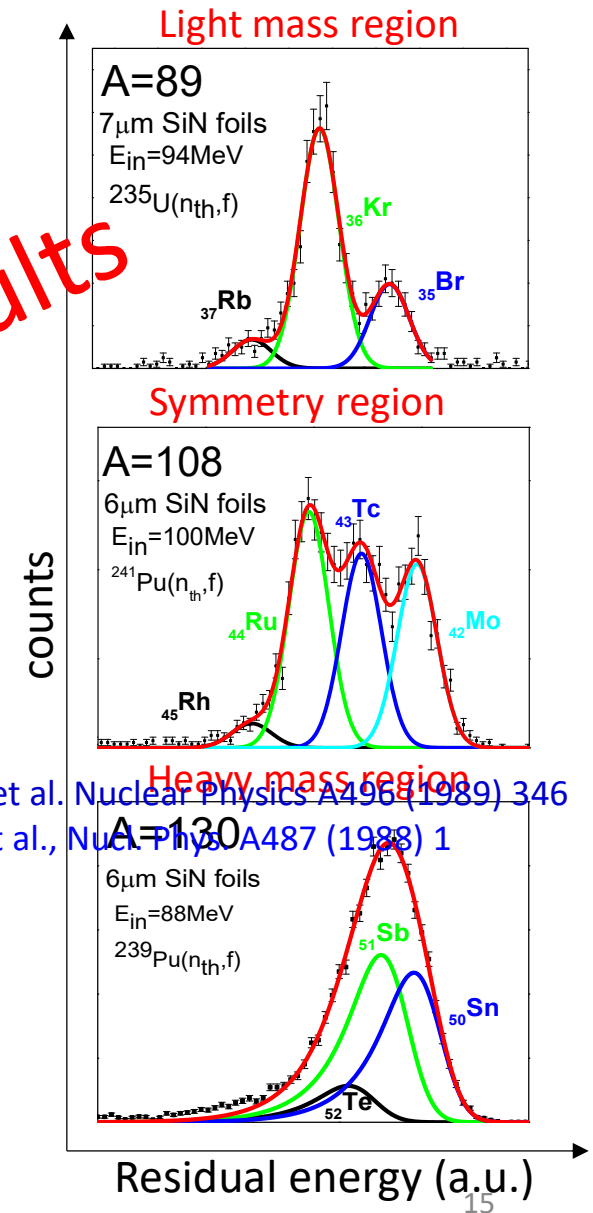
Quality of Z-Separation



CLTD+SiN: good resolution & linearity

→ Possibility to measure in symmetry & heavy mass region

Preliminary Results



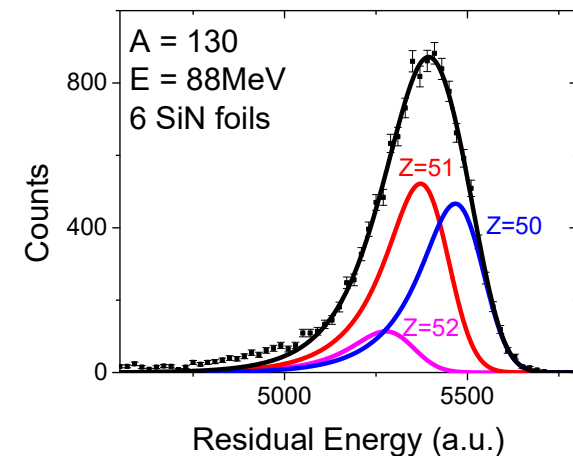
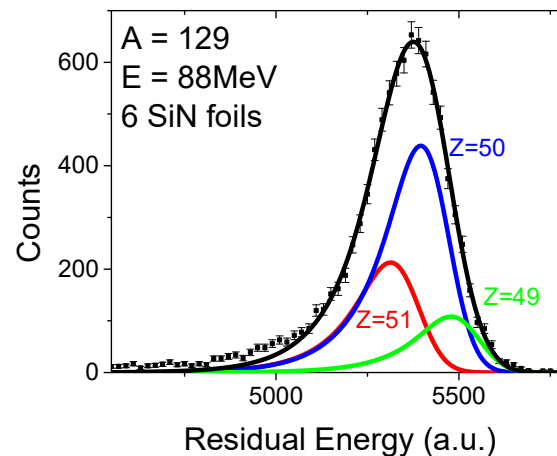
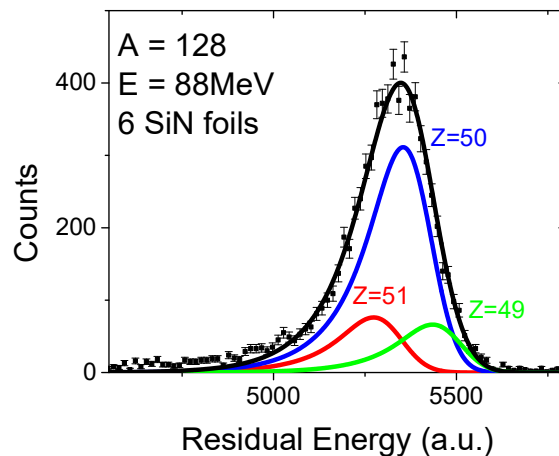
A. Djebara et al. Nuclear Physics A496 (1989) 346

U. Quade et al., Nucl. Phys. A487 (1988) 1

Heavy mass analysis

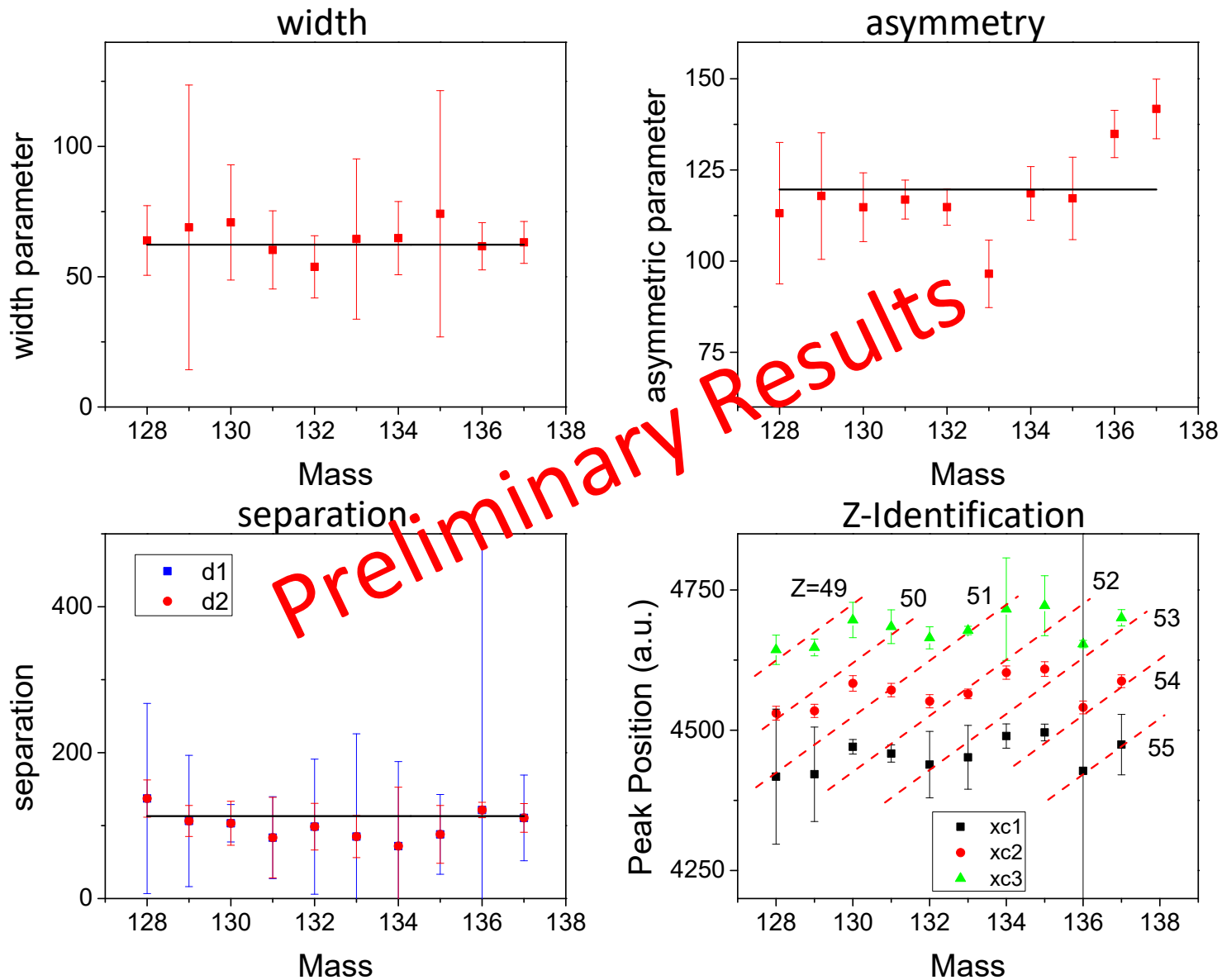
First direct Z-yield measurements in Heavy mass region:

- Standard method for unresolved peaks: Constrained fits with
 - good statistics (≈ 10000 counts)
 - good knowledge of response function
- With same set-up, measurements at similar mass and energy at MLL tandem tandem accelerator, allow to estimate the response function for the residual energy spectra.
- Stable fits for measured distributions.



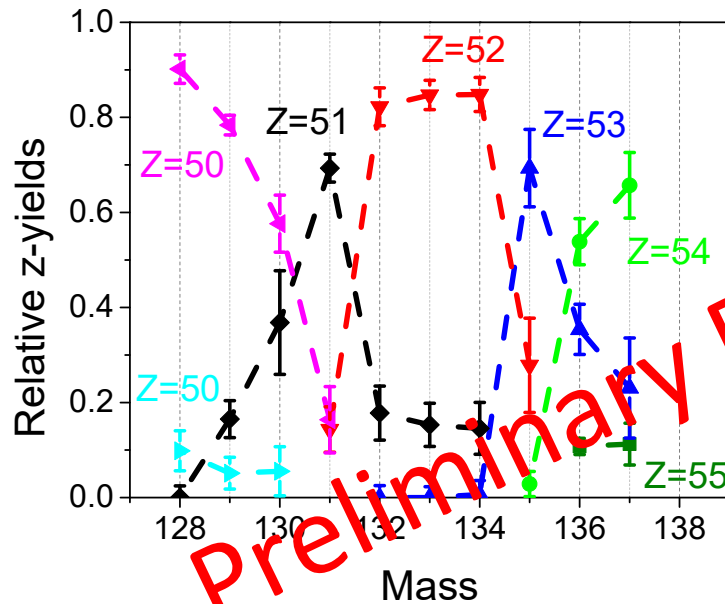
Preliminary Results

Heavy masses Z-yields Analysis

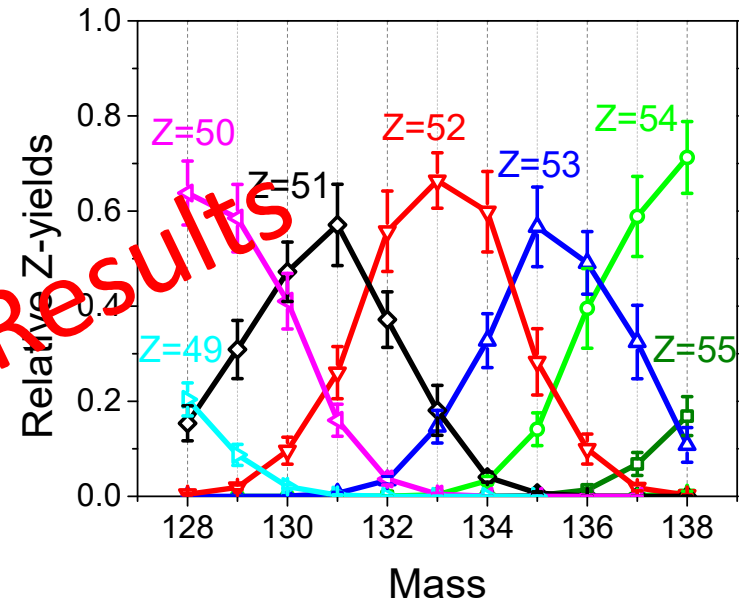


Heavy masses Z-yields: $^{239}\text{Pu}(n_{\text{th}}, f)$

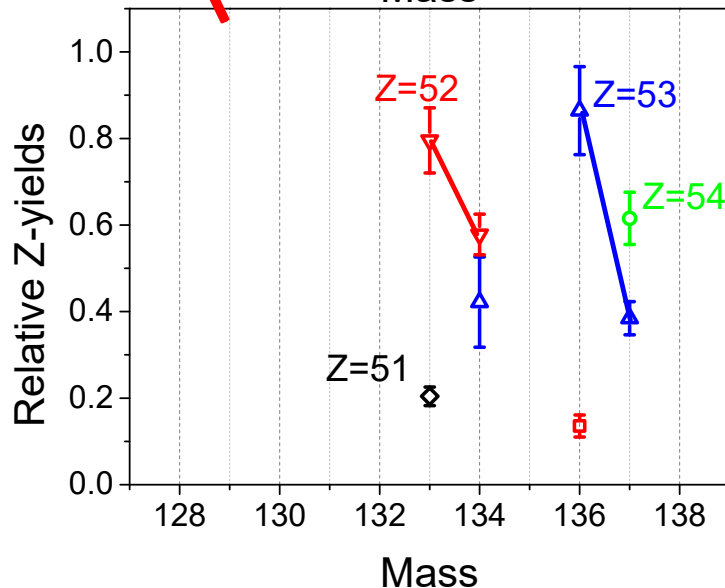
Our Expt.
Prelim. Results
 $E_{\text{LOH.}} = 80\text{MeV}$



Jeff Data base 3.1.1



Gamma
Spectroscopy*
Incl. all E & Q

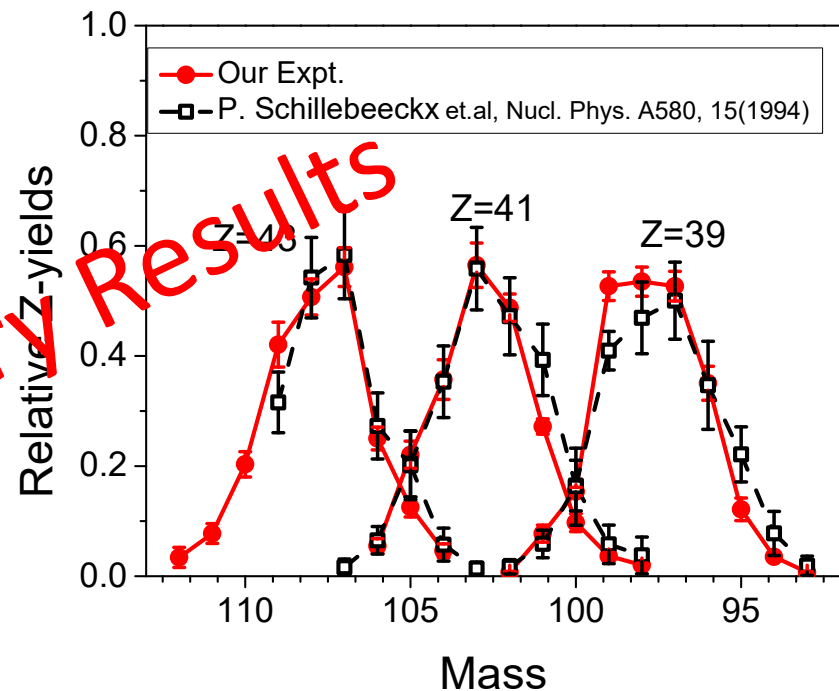
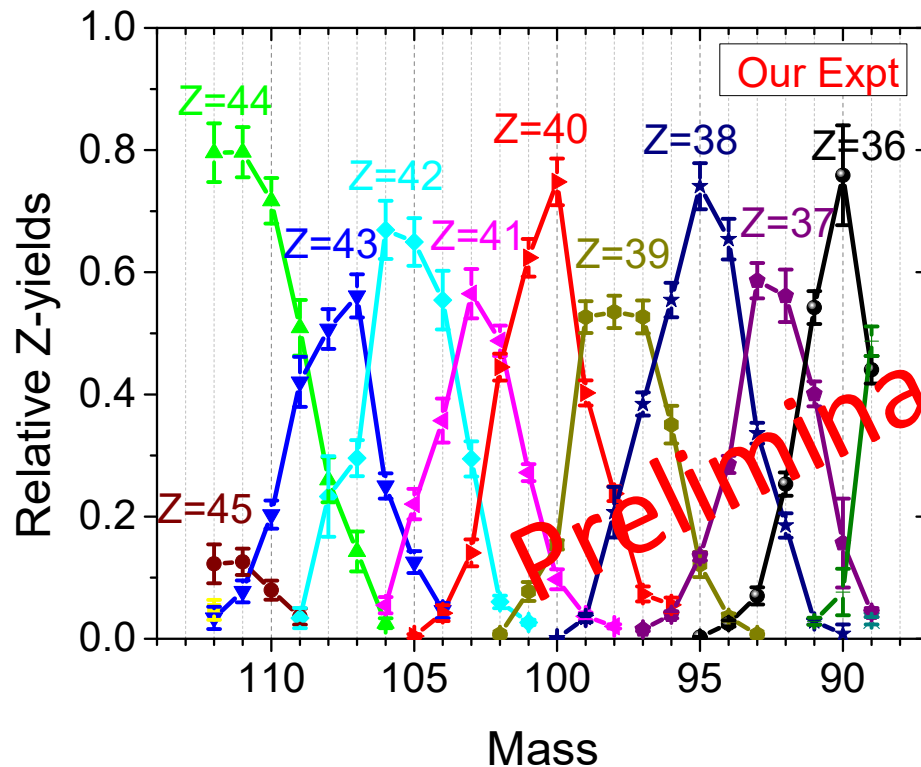


*A. Bail et al. Phys. Rev. C 84 (2011)
034605

$^{241}\text{Pu}(n_{\text{th}},f)$ Z-yield

First LOHENGRIN data on Z-yields for $A > 89$ were determined for $^{241}\text{Pu}(n_{\text{th}},f)$

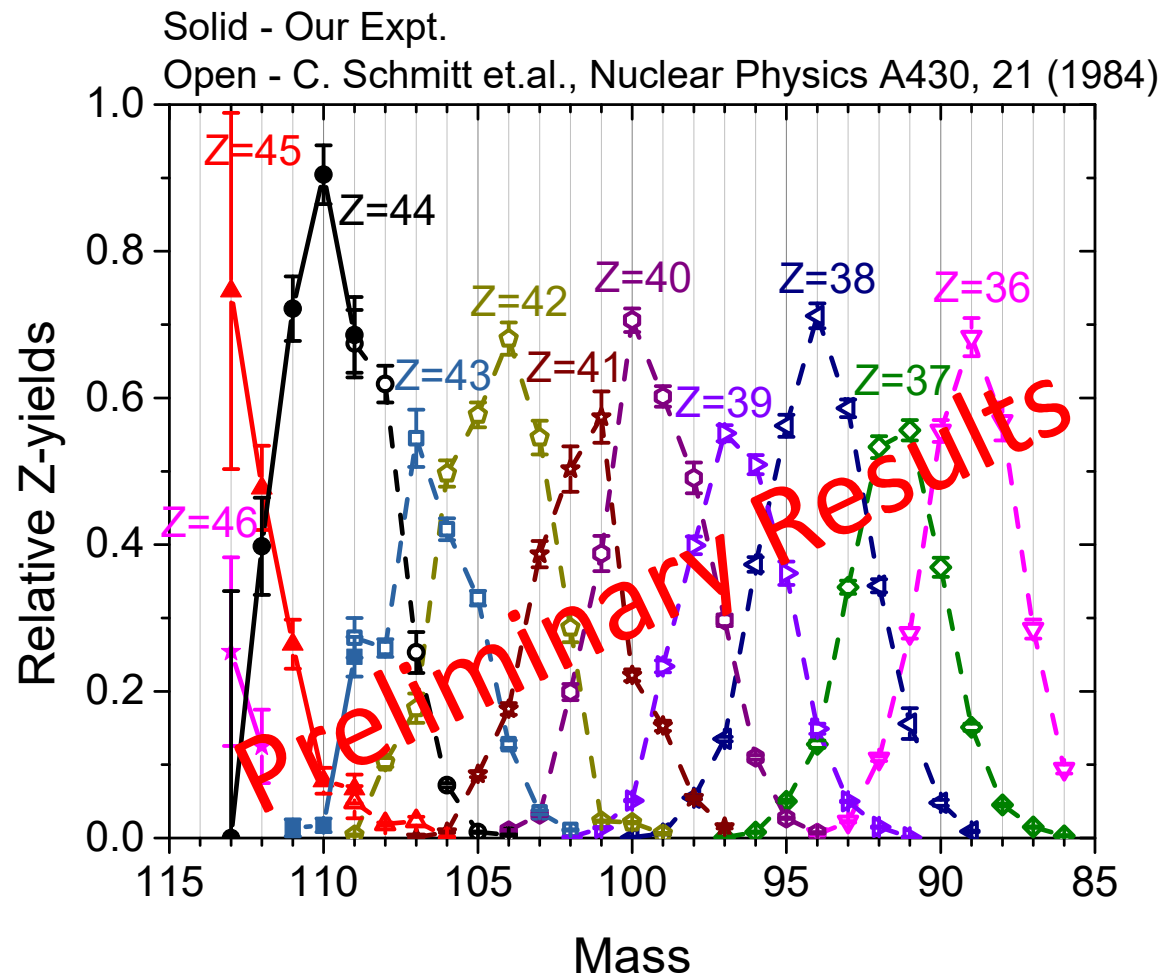
Good comparison with P.Schillebeeckx et.al. data – measurement with a different technique (Cosi-Fan-Tutte Spectrometer)



Prelim. Results

$^{239}\text{Pu}(n_{\text{th}},f)$ Z-yield

Also new data for $A = 109$ to 113 were determined for $^{239}\text{Pu}(n_{\text{th}},f)$



Prelim. Results

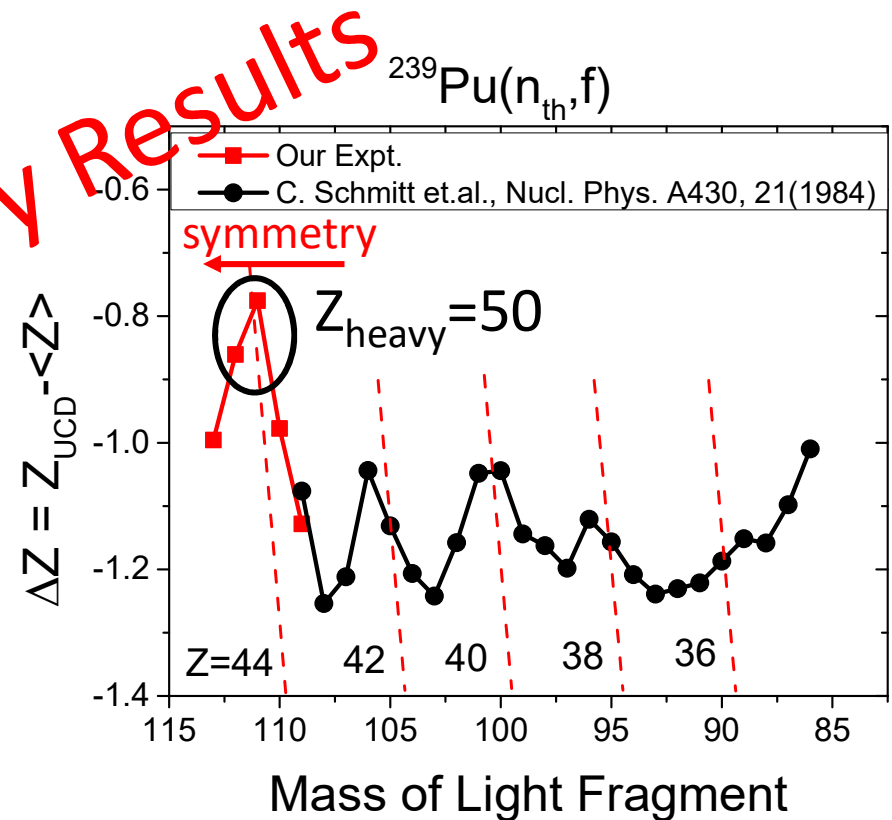
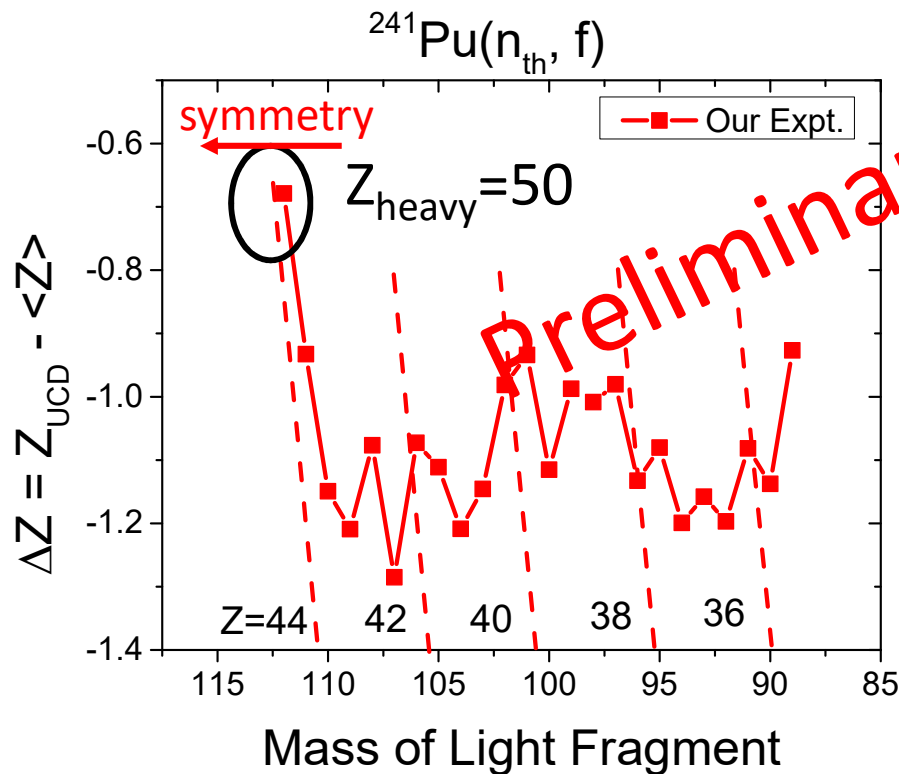
Odd-Even effect near symmetry

ΔZ = deviation from “democratic” distribution of neutron excess

$$\Delta Z = Z_{\text{UCD}} - \langle Z \rangle$$

$$Z_{\text{UCD}}(^{241}\text{Pu}) = \frac{94}{242}A$$

$$Z_{\text{UCD}}(^{239}\text{Pu}) = \frac{94}{24}A$$

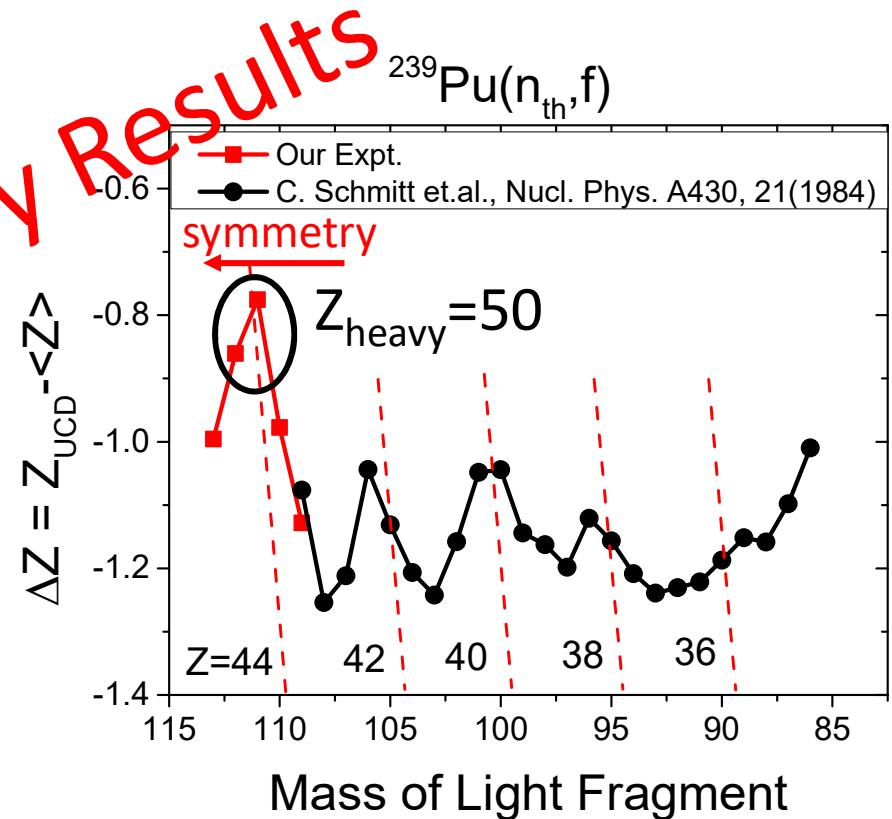
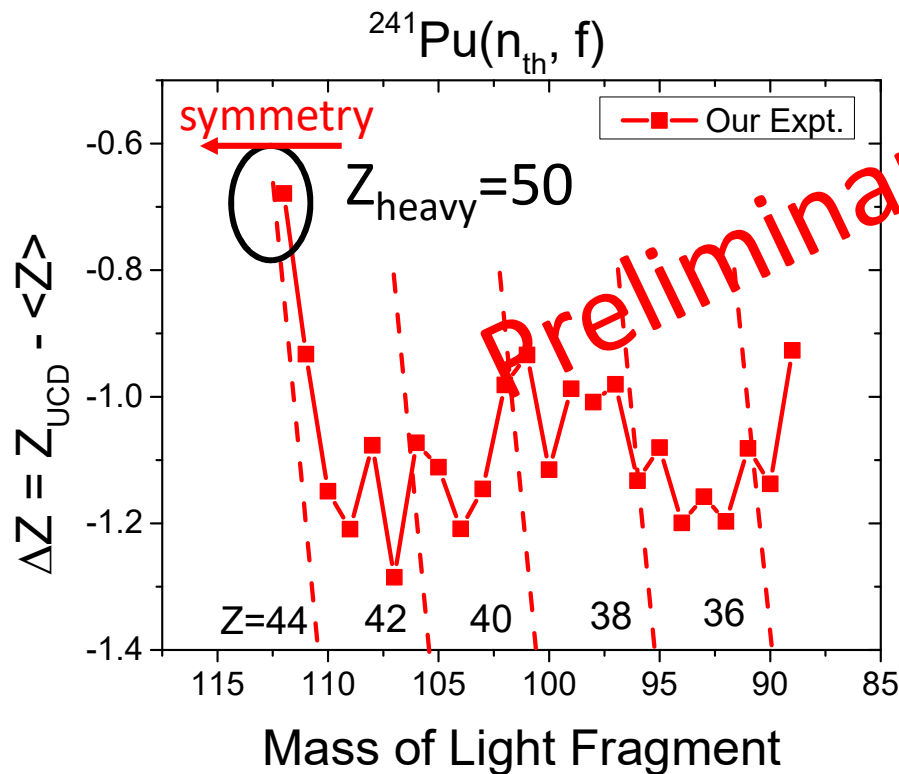


Prelim. Results

Odd-Even effect near symmetry

Rise in ΔZ towards corresponding heavy mass with closed shell, $Z=50$

For the first time in case of ^{239}Pu , we observe a drop in ΔZ beyond $Z_{\text{heavy}}=50$



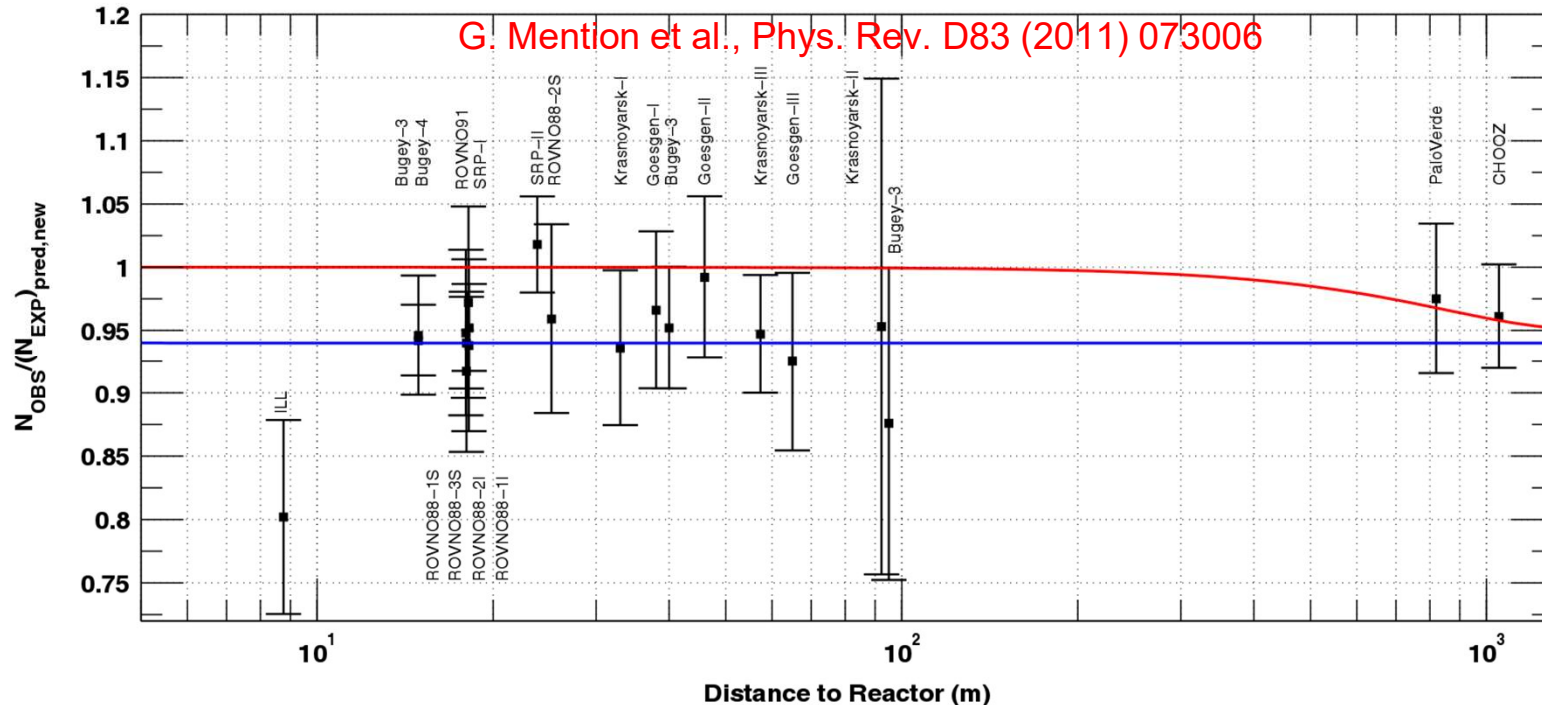
Prelim. Results

^{92}Rb & ^{96}Y yield

Motivation

Motivation for investigating Mass 92: The Reactor Antineutrino Anomaly

⇒ 5.7% deficit of measured antineutrinos
as expected from β -decay data of fission fragments



possible explanations: ⇒ 4th non standard „sterile“ neutrino
⇒ wrong Z-yields of fission fragments

^{92}Rb & ^{96}Y yield

A. A. Sonzogni et al., Phys. Rev. C91 (2015) 011301 (R)

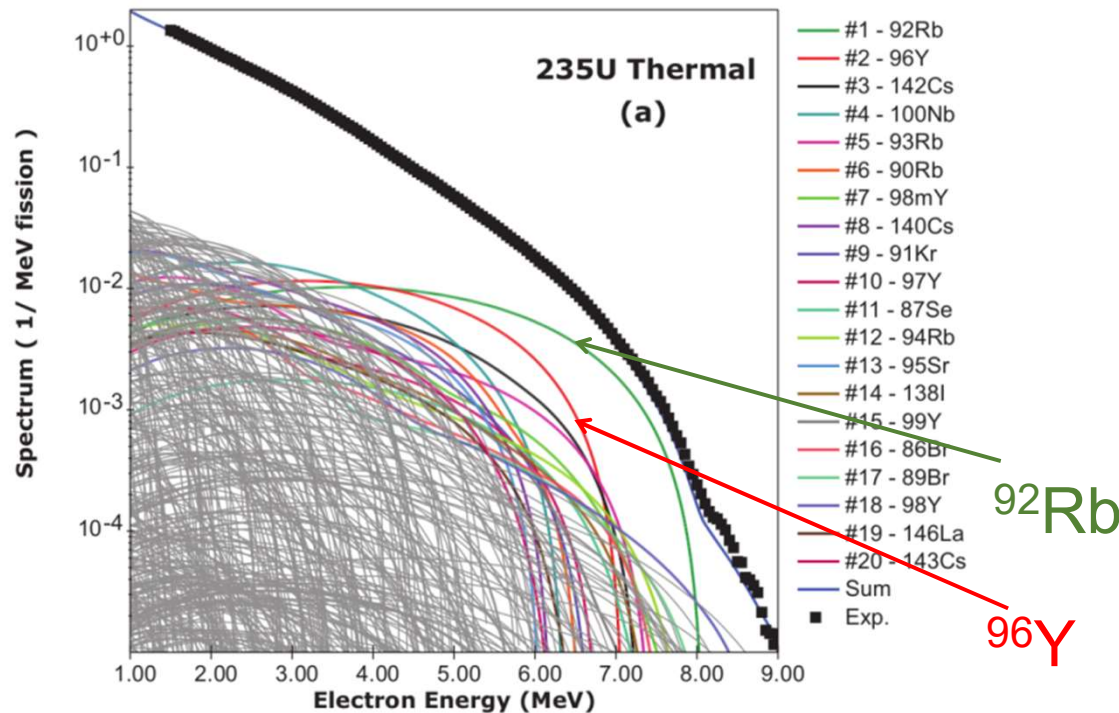


TABLE III. Same as Table II, except the strongest contributors from ^{235}U at 5.5 MeV are summarized.

Nuclide	Q_β (MeV)	GS BR (%)	$J_{gs}^\pi \rightarrow J_{gs}^\pi$	Contr. (%)
^{92}Rb	8.1	95.2(7)	$0^- \rightarrow 0^+$	21.6
^{96}Y	7.1	95.5(5)	$0^- \rightarrow 0^+$	14.5
^{142}Cs	7.3	56(5)	$0^- \rightarrow 0^+$	6.8
^{100}Nb	6.4	50(7)	$1^+ \rightarrow 0^+$	4.7
^{93}Rb	7.5	35(3)	$5/2^- \rightarrow 7/2^+$	4.6
^{90}Rb	6.6	33(4)	$0^- \rightarrow 0^+$	3.4
$^{98\text{m}}\text{Y}$	9.0	12(5) ^a	$(4,5) \rightarrow 4^+$	2.8
^{140}Cs	6.2	36(2)	$1^- \rightarrow 0^+$	2.4
^{91}Kr	6.8	18(3) ^b	$5/2^{(+)} \rightarrow (5/2^-)$	2.4

^aStrongest branch to a low-lying state is to a 4^+ , 1843-keV level.

^bStrongest branch is to a $(5/2^-)$, 109-keV level.

for ^{235}U target: contribution of ^{92}Rb to the β -spectrum at 5.5 MeV

⇒ JEFF 3.1.1 data basis: 21.6%

⇒ S.V.Tipnis et al. Phys. Rev. C 58, 905 (1998): 30.0%

A. A. Sonzogni et. al.:

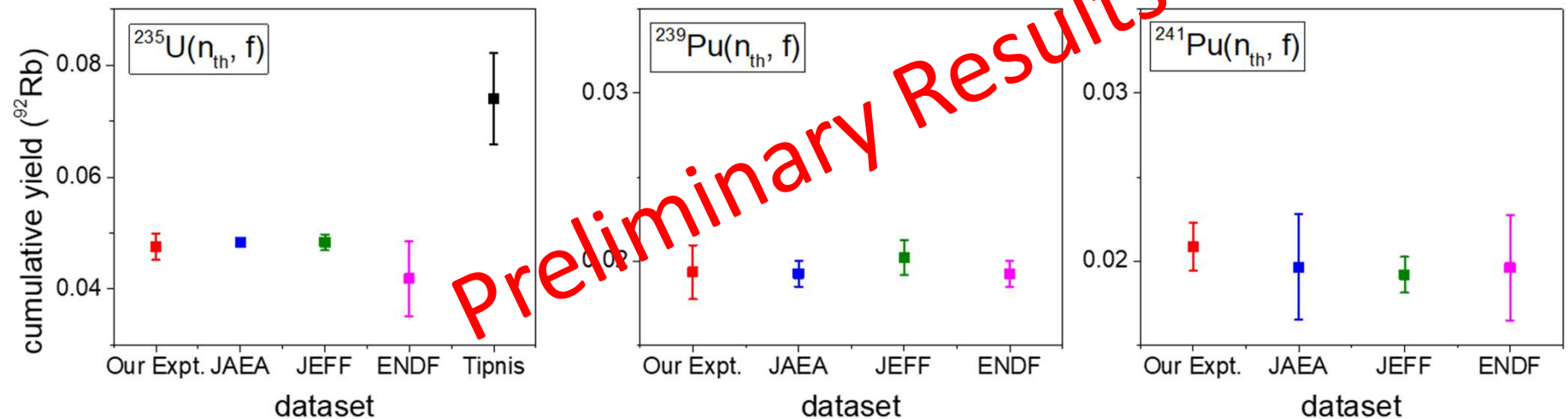
The ^{92}Rb cumulative fission yield following the thermal fission of ^{235}U definitely merits a new measurement.

^{92}Rb yield

Several measurements at different kinetic energies and ionic charges were made to precisely determine the cumulative yields

- Agreement with standard Data bases
- In case of ^{235}U , disagreement with data by S.V.Tipnis et al. Phys. Rev. C 58, 905 (1998)

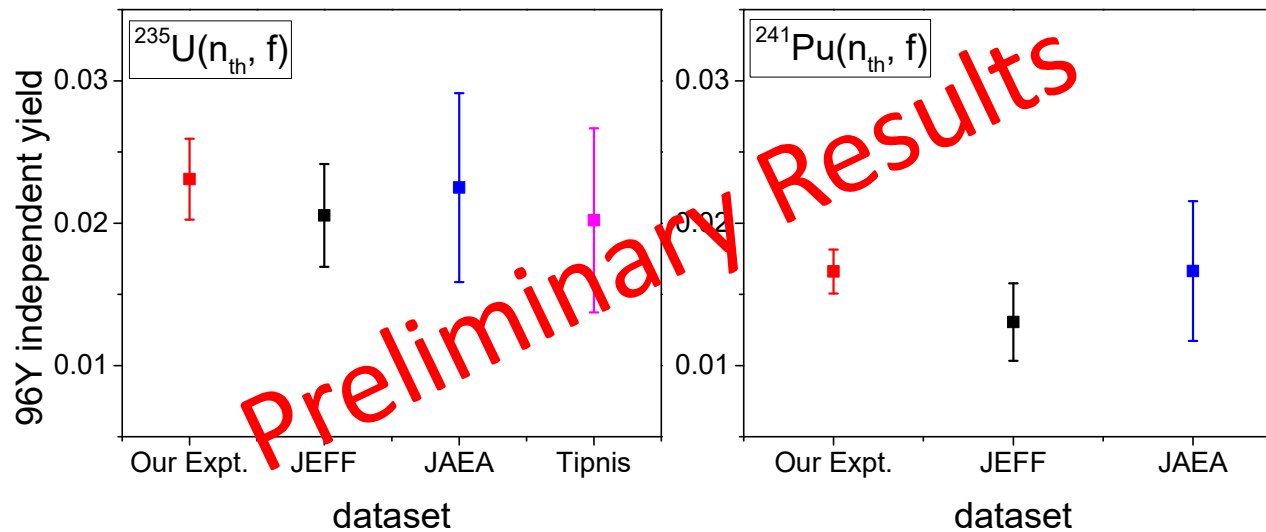
^{92}Rb



^{96}Y yield

- In agreement with the standard Data bases.

^{96}Y



S.V. Tipnis et al. Phys. Rev. C 58, 905 (1998)

Summary

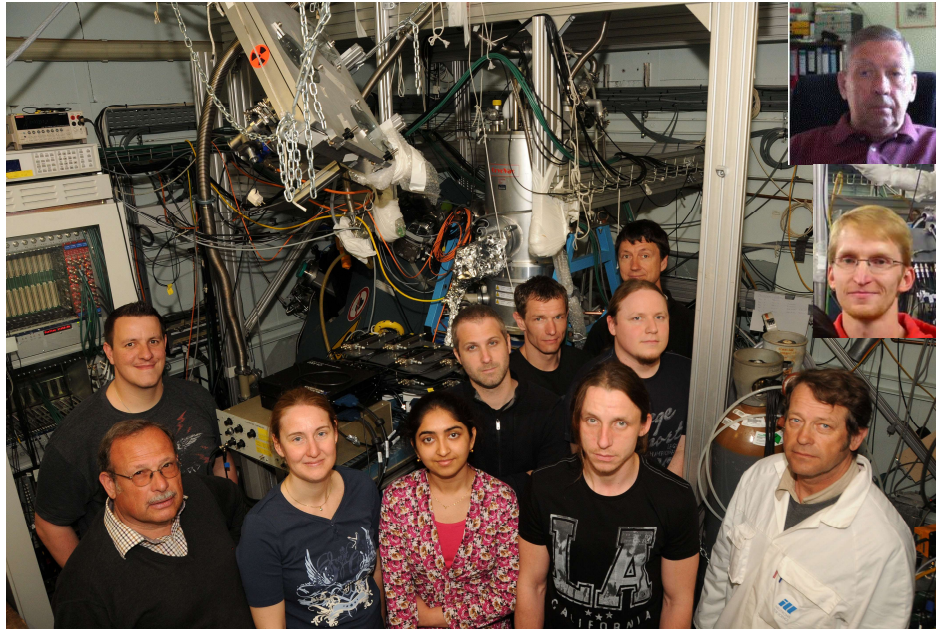
- Successful application of CLTDs for Z-yield measurements.
- CLTD + SiN: First time direct Z-yield measurements of heavy fragment masses.

A new method to measure also masses inaccessible with current technology e.g. gamma spectroscopy.

- Odd-even effect approaching symmetric fission for ^{239}Pu & ^{241}Pu
Important measurements for nuclear fission model description near scission

- Precise ^{92}Rb & ^{96}Y yields for ^{235}U , ^{239}Pu , ^{241}Pu (n_{th} , f)
Highest contributor to neutrino oscillations and reactor antineutrino anomaly studies

Thank you



S. Dubey^{1,2}, A. Echler^{1,2}, P. Egelhof^{1,2}, P. Grabitz^{1,2}, M. Mutterer¹, W. Lauterfeld², S. Stolte², S. Kraft-Bermuth³, P. Scholz³, S. Bishop⁴, J. Gomez⁴, A. Blanc⁵, U. Köster⁵, F. Gönnerwein⁶

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