

Efimov Physics with Fermions

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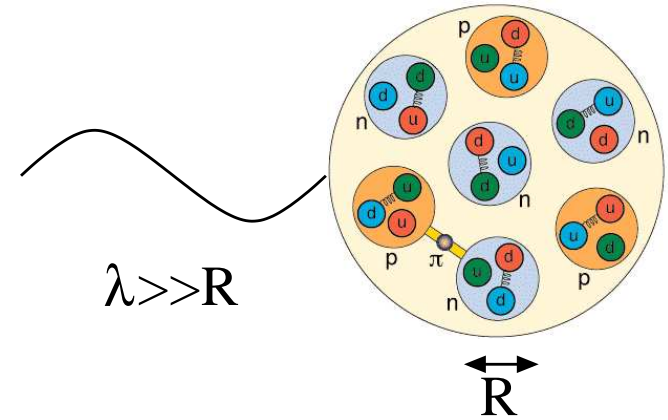
DFG

Collaborators: E. Braaten, D. Canham, D. Kang, L. Platter, ...

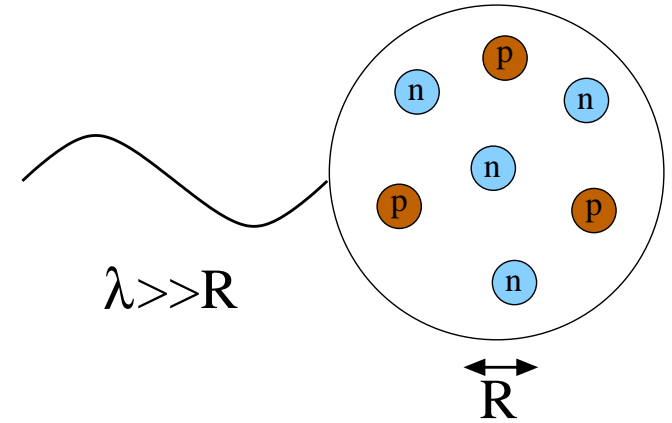
- Introduction
- Resonant Interactions and Efimov Physics
- Effective Field Theory for Large Scattering Length
- Applications
 - Cold atoms
 - Infrared limit cycle in QCD?
 - Halo nuclei
- Summary and Outlook

Review article: Braaten, HWH, Phys. Rep. **428** (2006) 259

- Separation of scales:
 $1/k = \lambda \gg R$
- Limited resolution at low energy:
→ expand in powers of kR

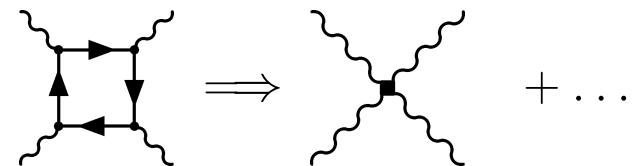


- Separation of scales:
 $1/k = \lambda \gg R$
- Limited resolution at low energy:
→ expand in powers of kR
- Short-distance physics not resolved
→ capture in low-energy constants using renormalization
→ include long-range physics explicitly
- Systematic, model independent → error estimates
- Classic example: light-light-scattering (Euler, Heisenberg, 1936)



Simpler theory for $\omega \ll m_e$:

$$\mathcal{L}_{QED}[\psi, \bar{\psi}, A_\mu] \rightarrow \mathcal{L}_{eff}[A_\mu]$$



- Painting at the limit of resolution of the human eye



G. Seurat, A Sunday on La Grande Jatte

- Resonant interaction \implies large scattering length a
- Natural expansion parameter: $\ell/|a|, k\ell, \dots$ ($\ell \sim r_e, l_{vdW}, \dots$)

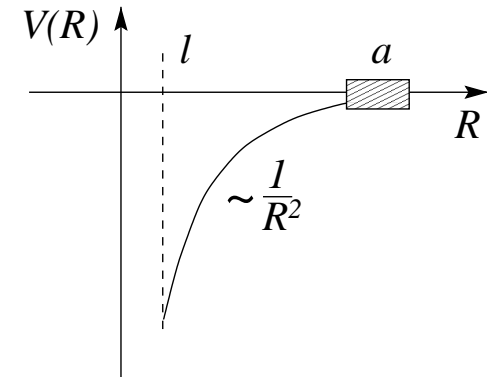
$$\text{Universal properties: } B_2 = \frac{\hbar^2}{ma^2} + \dots \quad (a < 0)$$

- Atomic physics:
 - ^4He : $a \approx 104 \text{ \AA} \gg r_e \approx 7 \text{ \AA} \sim l_{vdW} \longrightarrow B_d \approx 100 \text{ neV}$
 - Feshbach resonances \implies variable scattering length
- Nuclear physics: S -wave NN -scattering, halo nuclei, ...
 - $^1S_0, ^3S_1$: $|a| \gg r_e \sim 1/m_\pi \longrightarrow B_d \approx 2.2 \text{ MeV}$
 - ^6He : $2n$ separation energy $\approx 973 \text{ keV}$
- Particle physics: $X(3872)$ as a $D^0 D^{0*}$ molecule? ($J^{PC} = 1^{++}$)
$$m_X - (m_{D^0} + m_{D^{0*}}) = (-0.3 \pm 0.4) \text{ MeV}$$

(V. Efimov, Phys. Lett. **33B** (1970) 563)

- Three-body system with large scattering length a
- Hyperspherical coordinates: $R^2 = (r_{12}^2 + r_{13}^2 + r_{23}^2)/3$
- Schrödinger equation simplifies for $|a| \gg R \gg l$:

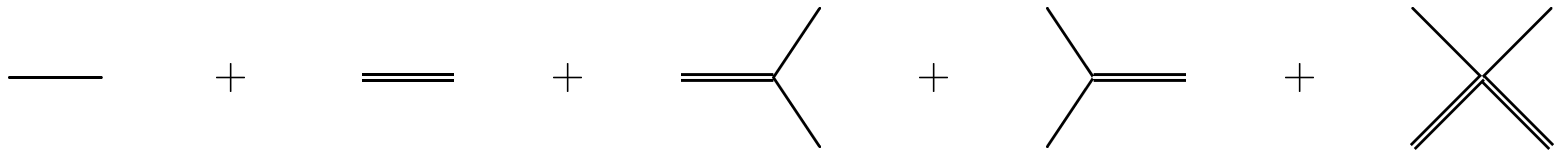
$$-\frac{\hbar^2}{2m} \left[\frac{\partial^2}{\partial R^2} + \frac{s_0^2 + 1/4}{R^2} \right] f(R) = \underbrace{-\frac{\hbar^2 \kappa^2}{m}}_E f(R)$$



- **Singular Potential:** renormalization required
- **Boundary condition at small R :** breaks scale invariance
 \implies **dependence of observables on 3-body parameter (and a)**
- **EFT formulation:** boundary condition \implies 3-body interaction

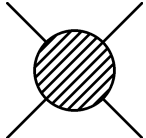
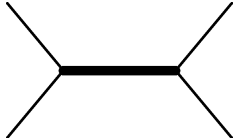
- Effective Lagrangian

$$\mathcal{L}_d = \psi^\dagger \left(i\partial_t + \frac{\vec{\nabla}^2}{2m} \right) \psi + \frac{g_2}{4} d^\dagger d - \frac{g_2}{4} (d^\dagger \psi^2 + (\psi^\dagger)^2 d) - \frac{g_3}{36} d^\dagger d \psi^\dagger \psi + \dots$$



- Interacting dimeron propagator \longrightarrow sum bubbles

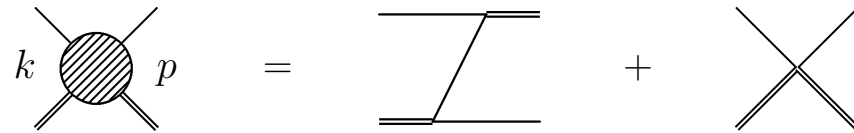


- Two-body amplitude $\mathcal{T}_2(k, k)$:  =  $\propto \frac{1}{1/a - ik} + \dots$

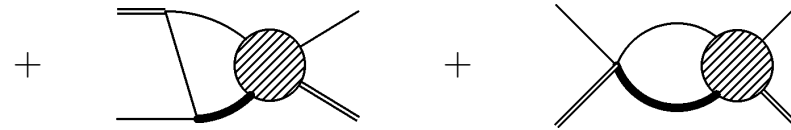
- Matching: $g_2 \longleftarrow a, B_d$

- RG fixed points of g_2 : $a = 0$ and $a = \infty$

- Higher order corrections: perturbation theory



- Three-body equation :

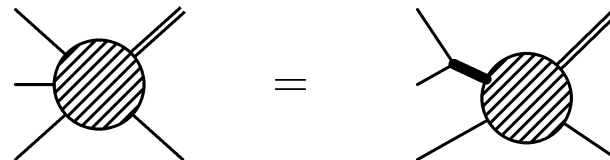


$$\mathcal{T}_3(k, p) = M(k, p) + \frac{4}{\pi} \int_0^\Lambda dq q^2 M(q, p) D_d(q) \mathcal{T}_3(k, q)$$

with $M(k, p) = \underbrace{F(k, p)}_{\text{1-atom exchange}} \underbrace{-\frac{g_3}{9g_2^2}}_{H(\Lambda)/\Lambda^2}$

($g_3 = 0, \Lambda \rightarrow \infty \rightarrow$ Skorniakov, Ter-Martirosian '57)

- Three-body recombination:



- Observables must be independent of regulator/cutoff Λ

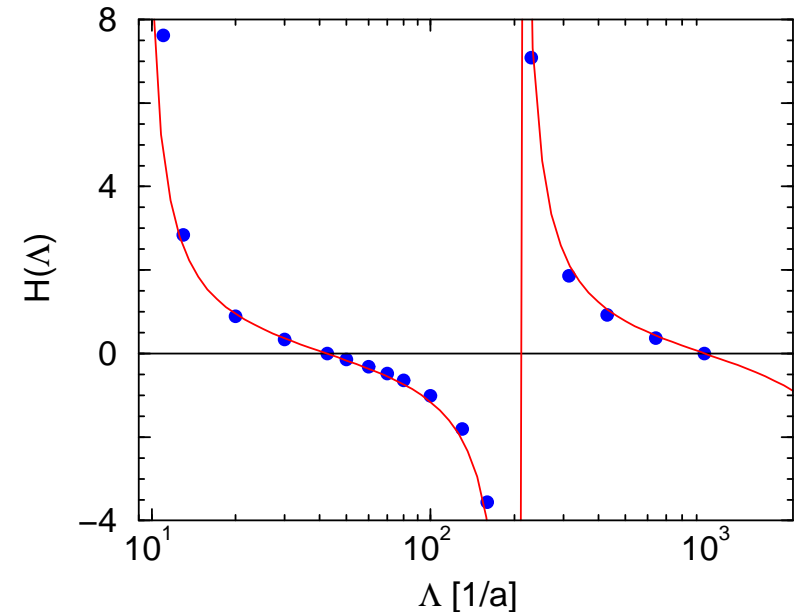
⇒ Running coupling $H(\Lambda)$

- $H(\Lambda)$ periodic: **limit cycle**

$$\Lambda \rightarrow \Lambda e^{n\pi/s_0} \approx \Lambda (22.7)^n$$

(cf. Wilson, 1971)

- Full scale invariance broken to discrete subgroup



$$H(\Lambda) = \frac{\cos(s_0 \ln(\Lambda/\Lambda_*) + \arctan(s_0))}{\cos(s_0 \ln(\Lambda/\Lambda_*) - \arctan(s_0))}, \quad s_0 \approx 1.00624$$

- **Limit cycle** \iff **Discrete scale invariance**
- **Matching:** $\Lambda_* \longleftarrow B_t, K_3, \dots \longrightarrow \kappa_*, a_*, a'_*$

- Similarity to Matryoschka doll
→ “Russian Doll Renormalization”

- Other examples

- $1/r^2$ potential in QM
- Field theory models
(Wilson, Glazek, LeClair et al.,...)
- Turbulence, earthquakes,
financial markets,...
(cf. Sornette, Phys. Rep. **297** (1998) 239)

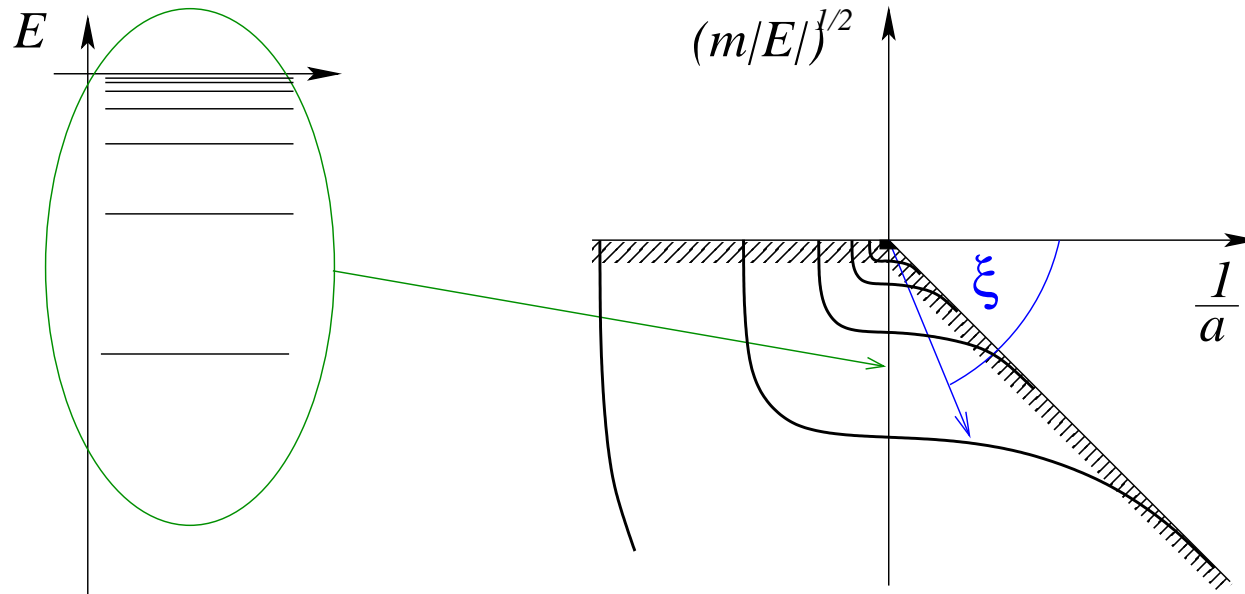
- Observable Consequences?

⇒ universal correlations, Efimov effect, log-periodic dependence on the scattering length,...



- Universal spectrum of three-body states

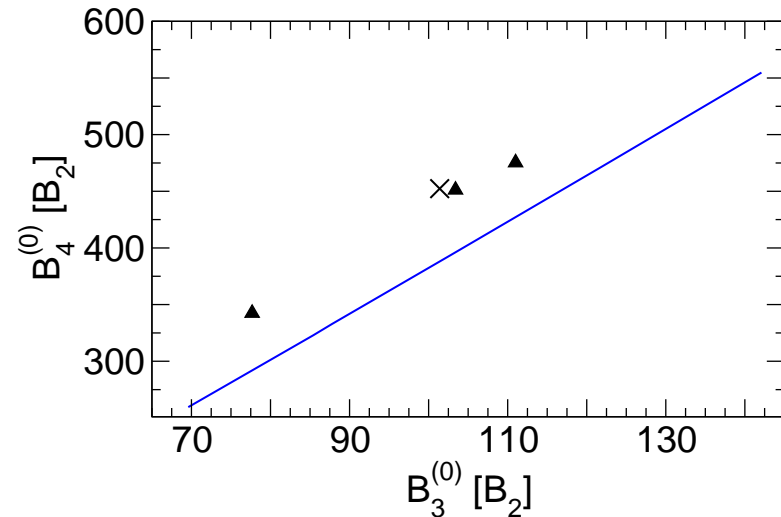
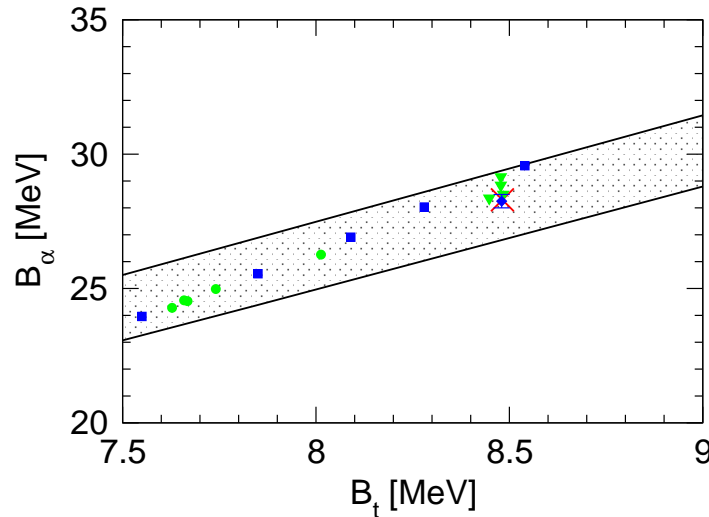
(V. Efimov, Phys. Lett. **33B** (1970) 563)



- Discrete scale invariance for fixed angle ξ
- **Geometrical spectrum** für $1/a \rightarrow 0$ (\Rightarrow “Efimov effect”)

$$B_3^{(n)} / B_3^{(n+1)} \xrightarrow{1/a \rightarrow 0} 515.035\dots$$

- 2 Parameters at LO \Rightarrow 3-body observables are correlated
 \Rightarrow Phillips line (Phillips, 1968)
- No four-body parameter at LO (Platter, HWH, Meißner, 2004)
 \Rightarrow 4-body observables are correlated \Rightarrow Tjon line



- Variation of 3-body parameter generates correlations
- Nuclear physics: Λ dependence of V_{low-k} (Bogner et al., 2004)
- Tjon line also at NLO (Kirscher et al., 2009)

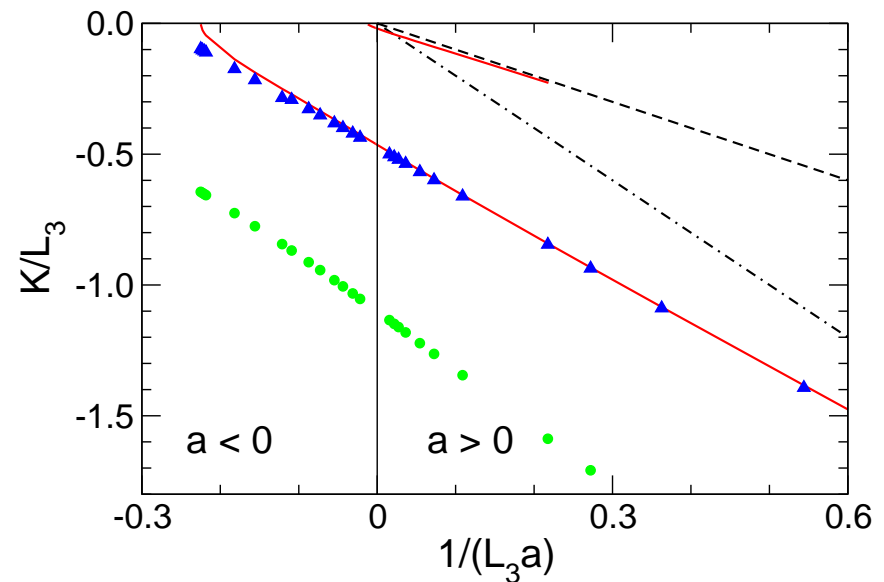
- Universal properties of 4-body system with large a
 - Bound state spectrum
 - Scattering Observables
- “Efimov-plot”: 4-body bound state spectrum as function of $1/a$

$$K = \text{sign}(E) \sqrt{m|E|}$$

$$B_4^{(0)} = 5B_3^{(0)} \quad (1/a \equiv 0)$$

$$B_4^{(1)} = 1.01B_3^{(0)}$$

(Platter, HWH, EPJA **32** (2007) 113)

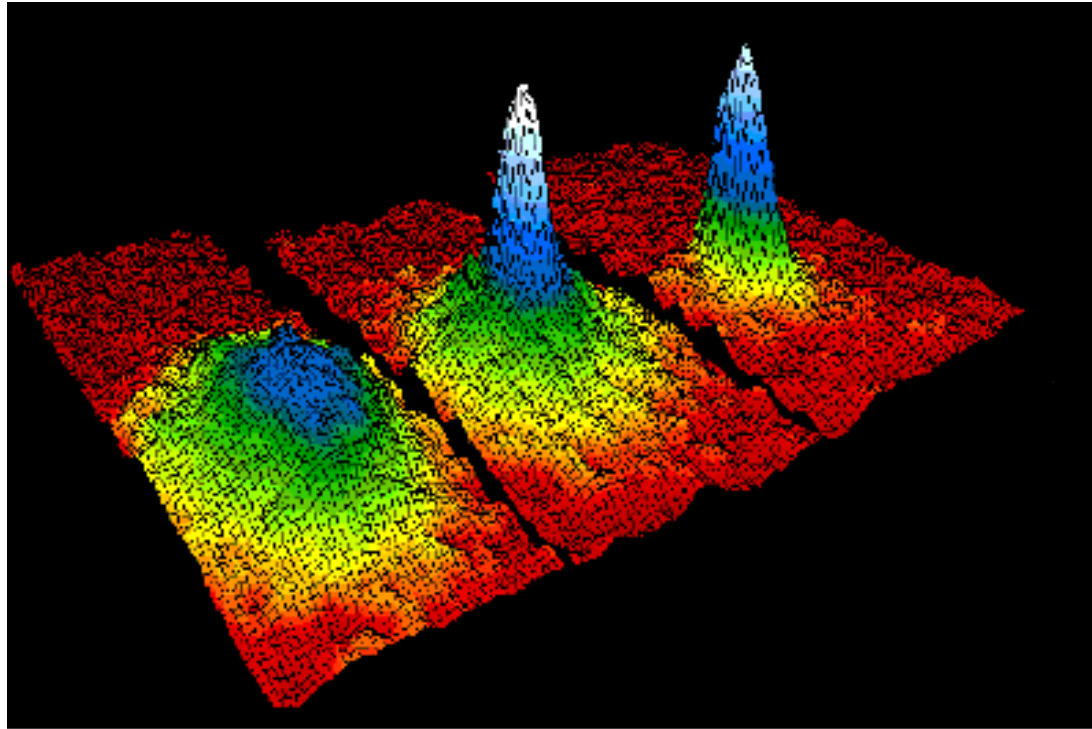


- Signature in Cs loss data

von Stecher, D’Incao, Greene, Nature Physics **5** (2009) 417

Ferlaino, Knoop, Berninger, Harm, D’Incao, Nägerl, Grimm, PRL **102** (2009) 140401

- Velocity distribution ($T = 400$ nK, 200 nK, 50 nK)

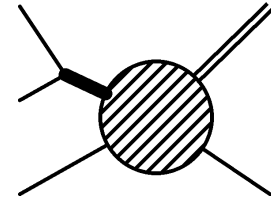


(Source: <http://jilawww.colorado.edu/bec/>)

- Few-body loss rates provide window on Efimov physics
- Variable scattering length via Feshbach resonances

- Recombination into weakly-bound dimer:

3 atoms \rightarrow dimer + atom \Rightarrow **loss of atoms**



- Recombination constant: $\dot{n}_A = -3 \alpha n_A^3$ (gas)

- Scattering length dependence for $a > 0$:

(Nielsen, Macek, 1999; Esry, Greene, Burke, 1999; Bedaque, Braaten, HWH, 2000)

$$\alpha \approx 67.1 \sin^2 [s_0 \ln(a\kappa_*) + 1.16] \frac{\hbar a^4}{m}, \quad s_0 \approx 1.00624..$$

- Alkali atoms form deeply-bound dimers
- Modification from deep dimers: Efimov states acquire width

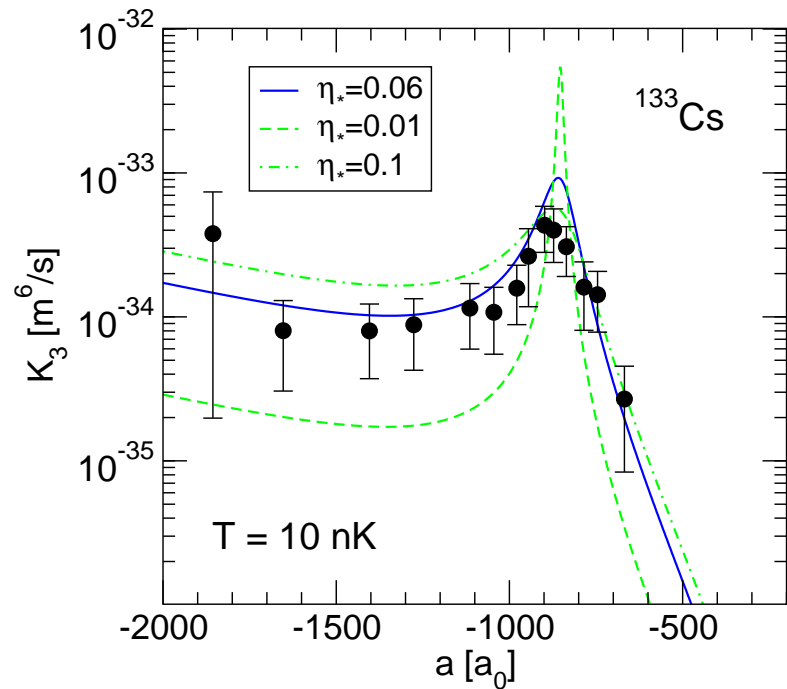
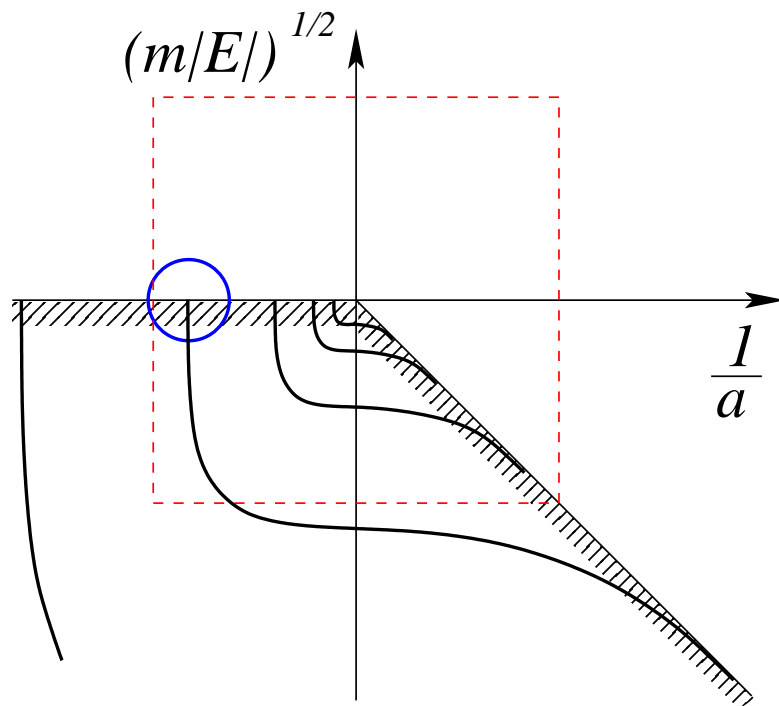
$$\implies \kappa_* \rightarrow \kappa_* \exp(i\eta_*/s_0) \quad (\text{Braaten, HWH, 2004})$$

- Recombination into deep dimers

- Experimental evidence for Efimov states in ^{133}Cs

(Kraemer et al. (Innsbruck), Nature **440** (2006) 315)

- Identification via 3-body recombination rate



$$a < 0: \quad a'_* = -875a_0$$

- Finite temperature effects small for $T \lesssim 200$ nK

(Braaten, HWH, Kang, Platter, 2008)

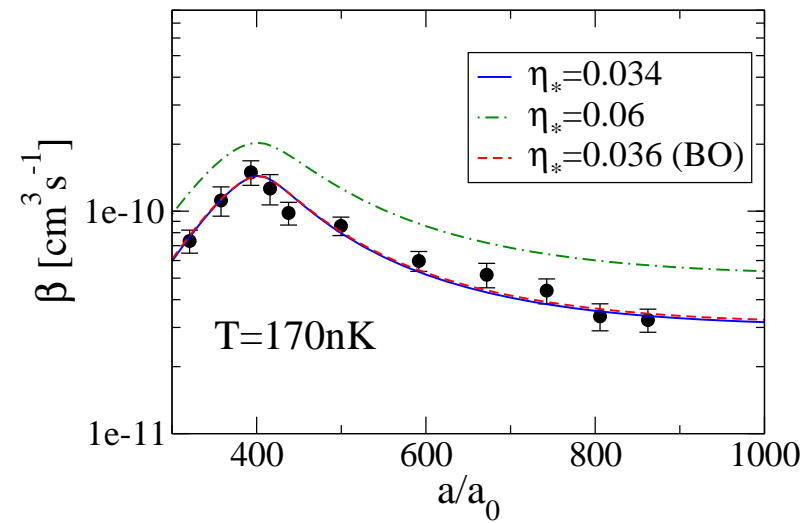
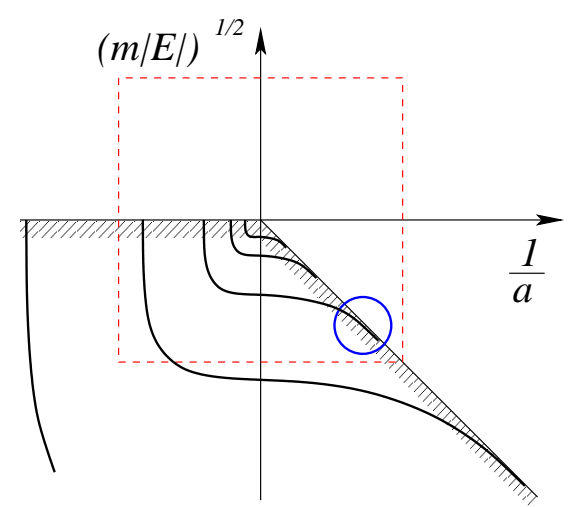
- Dimer Relaxation: atom + dimer \rightarrow atom + deep dimer (energetic)

- Relaxation constant:
$$\frac{dn_A}{dt} = \frac{dn_D}{dt} = -\beta n_A n_D$$

- Recent experiment: Knoop et al. (Innsbruck), Nature Physics **5** (2009) 227

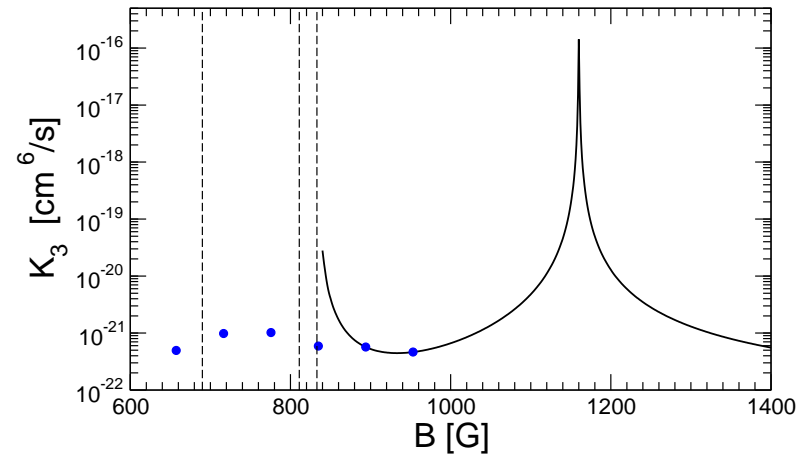
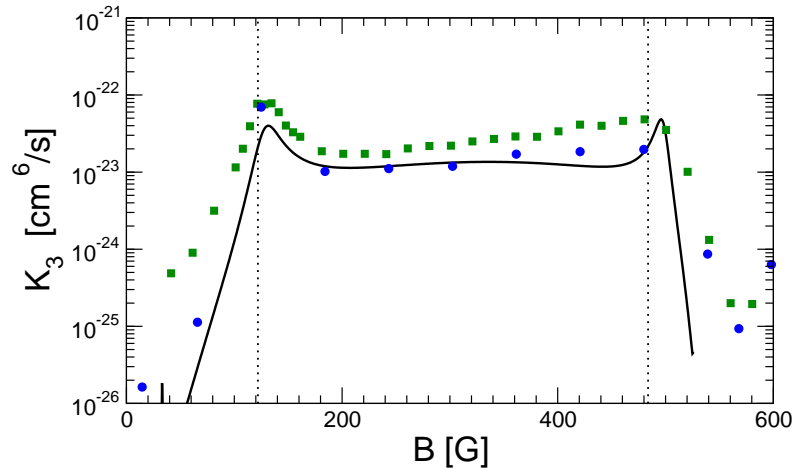
- Finite temperature $T \sim T_c$: Bose-Einstein average

- Include trap geometry:
$$\mu_i \longrightarrow \mu_i - m_i \bar{\omega}^2 r^2 / 2, \quad i = a, d$$



Helfrich, HWH, EPL **86** (2009) 53003

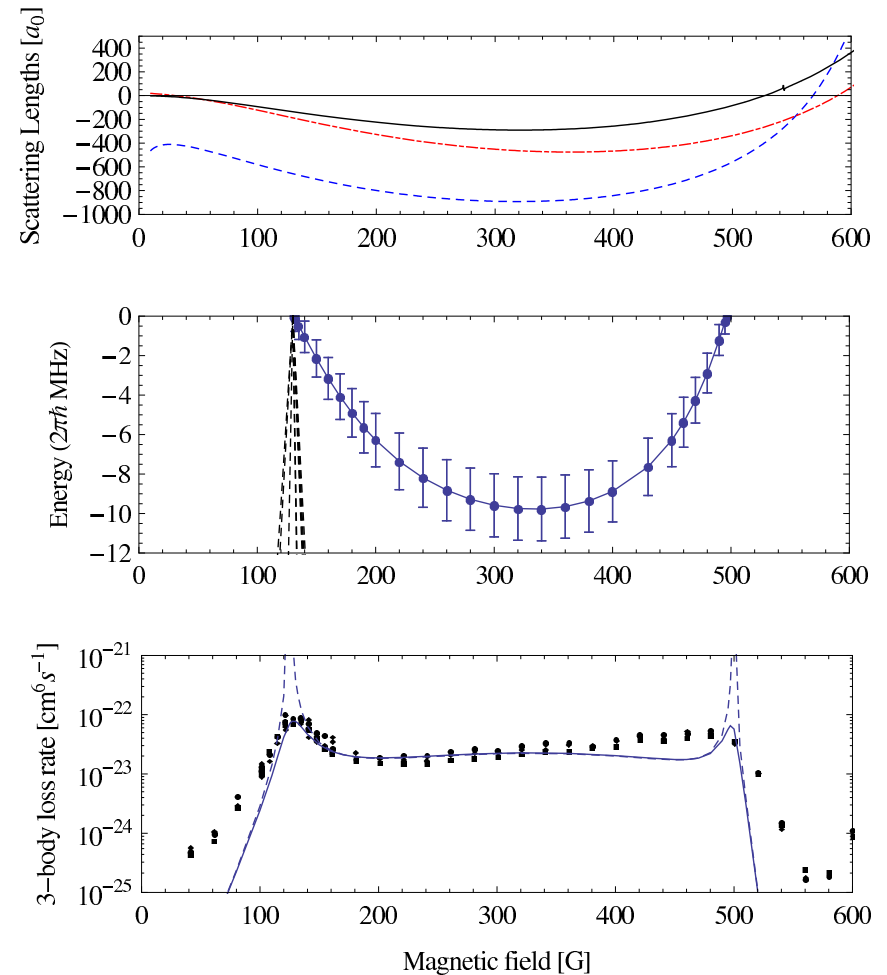
- Efimov effect for fermions $\Rightarrow \geq 3$ spin states
- Experimental evidence for Efimov states in ${}^6\text{Li}$
 - Ottenstein et al. (Heidelberg), Phys. Rev. Lett. **101** (2008) 203202
 - Huckans et al. (Penn State), Phys. Rev. Lett. **102** (2009) 165302



(Braaten, HWH, Kang, Platter, arXiv:0811.3578)

- Systematic normalization error: 70-90%
- Possible resonance around $B \approx 1160$ G
- Other approaches: Schmidt et al., arXiv:0812.1191; Naidon et al., arXiv:0811:4086

- **Solution of Faddeev equations in hyperspherical formalism**
(Naidon, Ueda, arXiv:0811:4086)
- **Functional renormalization group with explicit trion field**
(Schmidt, Floerchinger, Wetterich, arXiv:0812.1191)
- **Good agreement among different approaches**
(figure taken from Naidon, Ueda)



- Efimov resonances in a mixture of ^{41}K and ^{87}Rb atoms
(Barontini et al. (Florence), arXiv:0901:4584)
⇒ Connected K-Rb-Rb resonances for $a > 0$ and $a < 0$
- Efimov spectrum in ultracold potassium atoms near a Feshbach resonance (Zaccanti et al. (Florence), arXiv:0904.4453)
⇒ Observation of first two states of an Efimov spectrum
- Observation of universality in ultracold ^7Li three-body recombination (Gross et al. (Ramat-Gan), arXiv:0906.4731v1)

- Large scattering lengths in NN interaction imply:
QCD close to critical trajectory for infrared limit cycle
- Conjecture:

QCD can be tuned to critical trajectory by small changes in the up- and down-quark masses

(Braaten, HWH, 2003)

- Quark mass dependence of low-energy constants
 - Chiral EFT with explicit pions
 - **Extrapolation in quark masses possible**

Beane, Bedaque, Savage, van Kolck, 2002

Beane, Savage, 2003

Epelbaum, Meißner, Glöckle, 2003

- Gell-Mann-Oakes-Renner relation: $m_\pi^2 \propto (m_u + m_d)$
- Chiral NN potential: $V = V_{OPE} + V_{TPE} + V_{contact}$

$$V_{OPE} = \left(-\frac{g_A^2}{4F_\pi^2} + \dots \right) \boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2 \frac{(\boldsymbol{\sigma}_1 \cdot \mathbf{q})(\boldsymbol{\sigma}_2 \cdot \mathbf{q})}{\mathbf{q}^2 + m_\pi^2}$$

$$V_{contact} = \bar{C}_S + m_\pi^2 \bar{D}_S + (\bar{C}_T + m_\pi^2 \bar{D}_T) \boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2 + \dots$$

- $\bar{D}_{S,T}$ not known \longrightarrow dimensional analysis estimate:

(Epelbaum, Meißner, Glöckle, 2003)

$$-3.0 < F_\pi^2 \Lambda_\chi^2 \bar{D}_{S,T} < 3.0 \quad \longrightarrow \text{error band in extrapolation}$$

- Extrapolation of a_t, a_s, B_d as function of $m_\pi^2 \propto (m_u + m_d)$

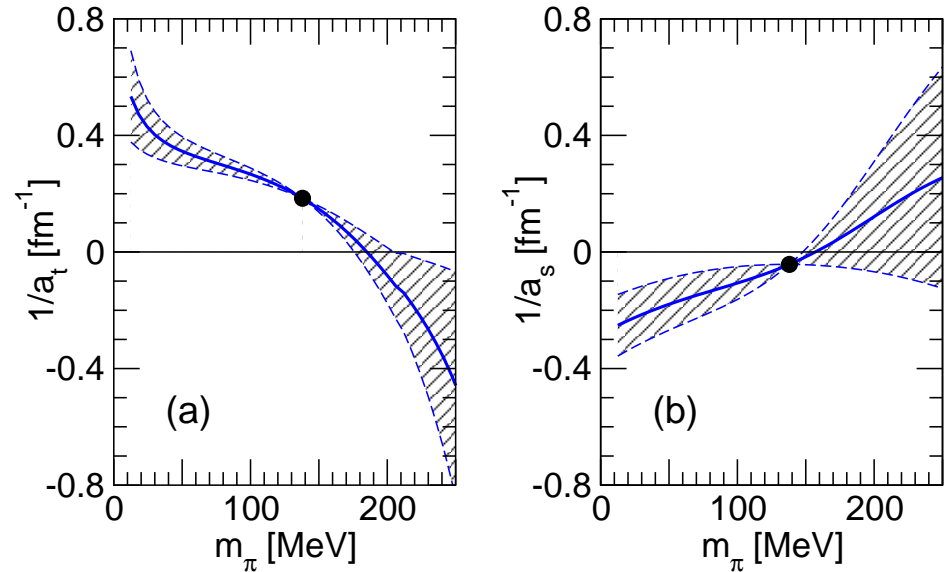
- Remember: $m_\pi^2 \propto (m_u + m_d)$

- Limit cycle \Rightarrow require

$$1/a_t = 1/a_s = 0$$

- Solutions exist for (NLO)

$$175 \text{ MeV} \lesssim m_\pi^{crit.} \lesssim 205 \text{ MeV}$$



(Epelbaum, Meißner, Glöckle, NPA **714** ('03) 535)

- Vary m_u and m_d independently, but:

only operators $\propto (m_u + m_d)$ at NLO

- Properties of limit cycle universal: study one of solutions in NLO

($m_\pi^{crit.} = 197.8577 \text{ MeV}$) (Epelbaum, HWH, Meißner, Nogga, 2006)

- Beane & Savage: similar results for $m_\pi > (m_\pi)_{phys}$

- Identification:

 - Efimov effect for triton

 - infinitely many excited states

- Calculate in chiral/universal EFT

- Properties:

 - $2N - N$ cluster structure close to threshold

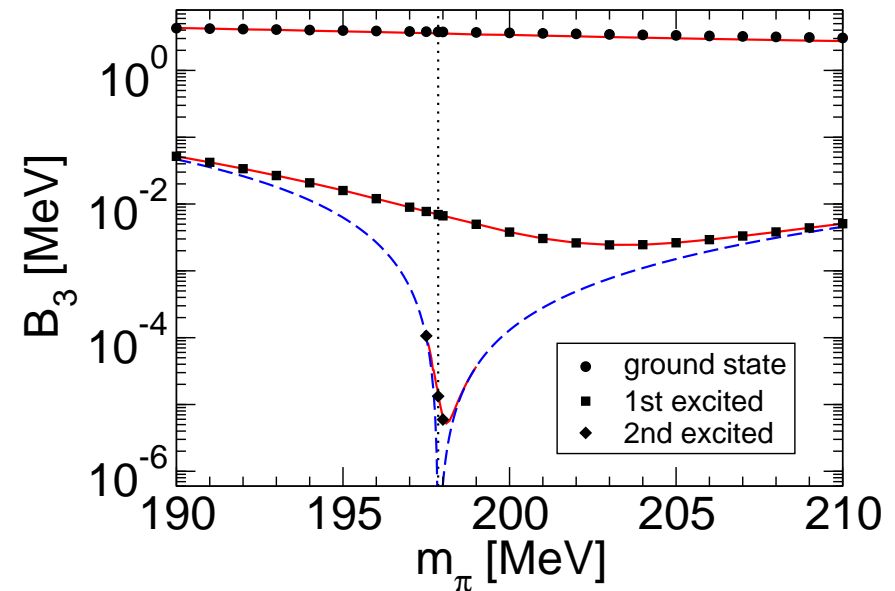
 - $2N - 3N$ overlap vanishes at critical point

- Can use chiral EFT for extrapolations

- Effect observable in Lattice QCD calculations?

 - (cf. Wilson, 2005)

- Variation of fundamental constants?



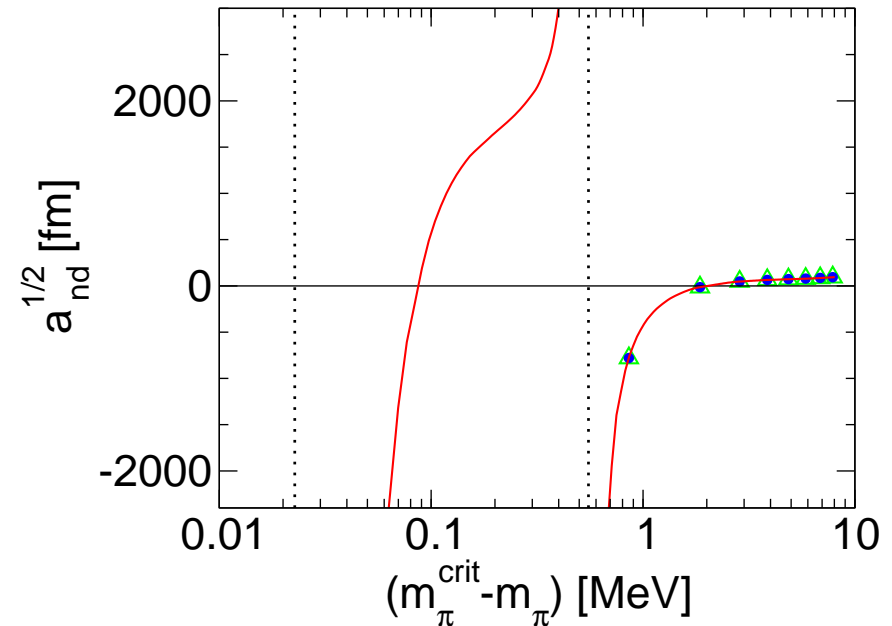
- “Matching” close to critical point:

chiral EFT \longleftrightarrow universal EFT

- Approaches complementary: e.g. scattering observables in universal EFT

- Divergences:

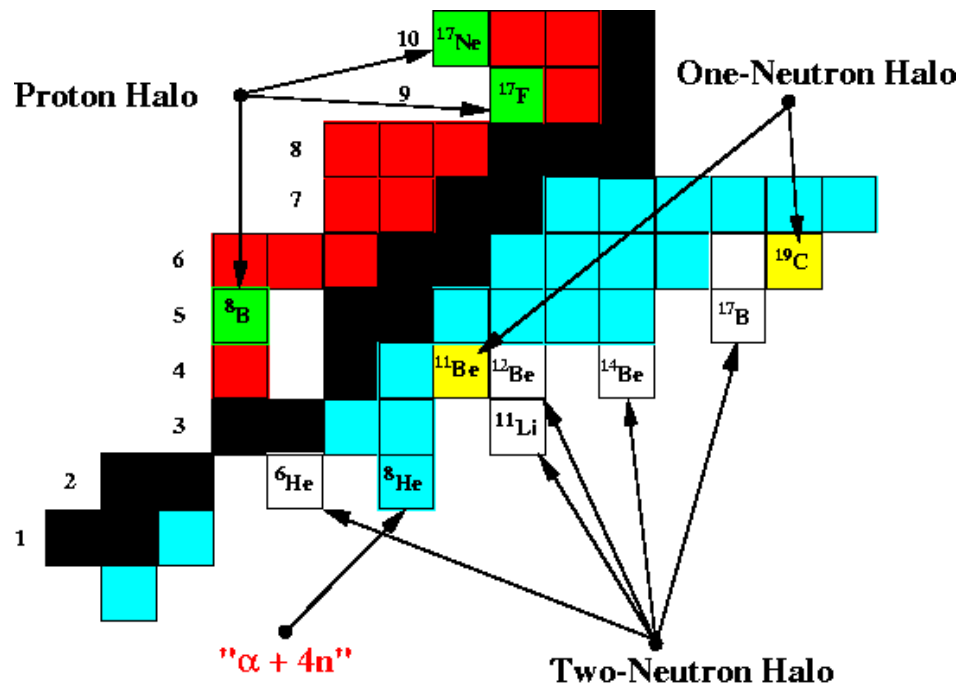
- at critical point
- when new Efimov states appear



- N2LO in universal theory (including effective range corrections)

(HWH, Phillips, Platter, 2007)

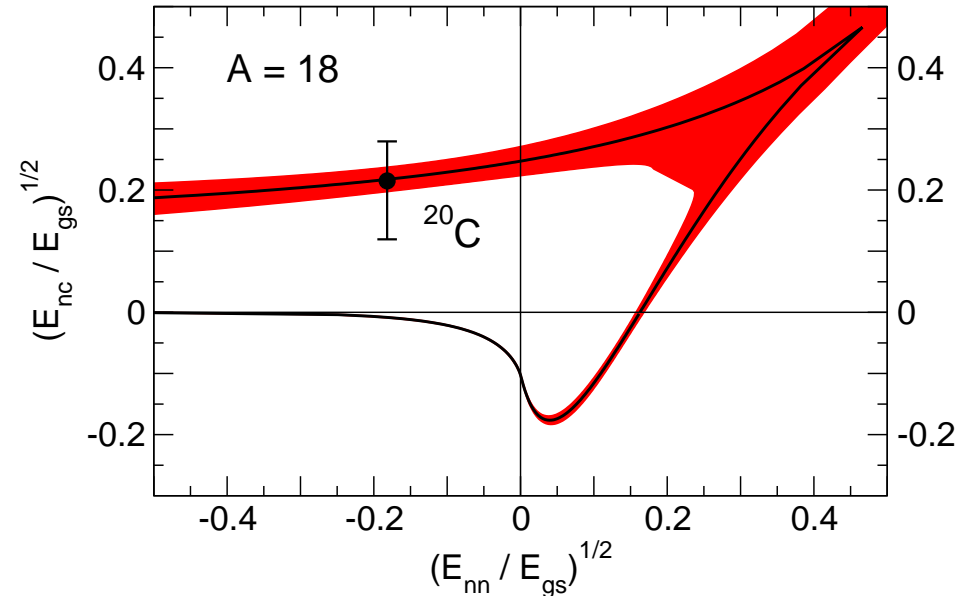
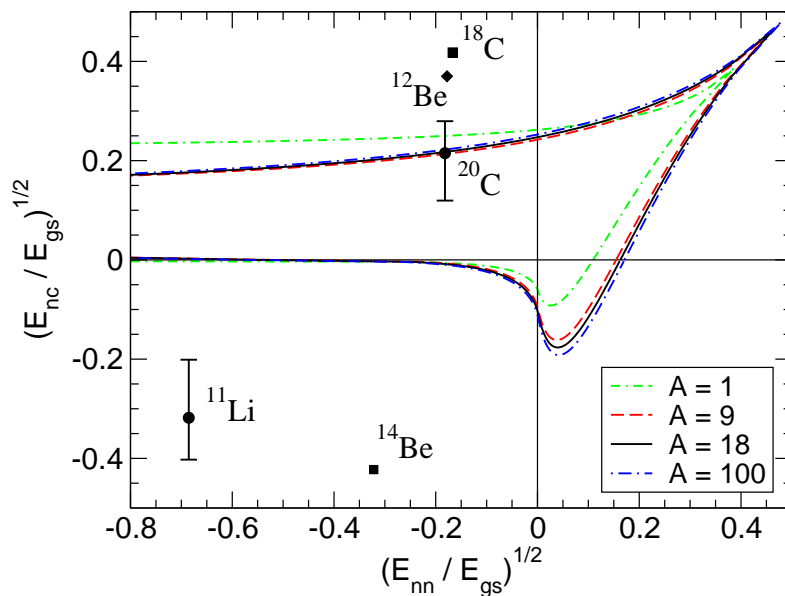
- Low separation energy of valence nucleons: $B_{valence} \ll B_{core}, E_{ex}$
 → close to “nucleon drip line” → **scale separation** → EFT



<http://www.nupecc.org>

- EFT for halo nuclei
 - $n\alpha$ -System (“ ^5He ”) (Bedaque, Bertulani, HWH, van Kolck, 2002)
 - $\alpha\alpha$ -System (“ ^8Be ”) (Higa, HWH, van Kolck, 2008)

- **Examples:** $^{14}\text{Be} \longleftrightarrow ^{12}\text{Be} + n + n$, $^{20}\text{C} \longleftrightarrow ^{18}\text{C} + n + n$
- **“Effective” 3-body system:** separation energy of valence nucleons small compared to binding energy of “core”
- **Efimov effect in halo nuclei?** \Rightarrow **excited states**



Canham, HWH, Eur. Phys. J. A **37** (2008) 367
 (cf. Amorim, Frederico, Tomio, 1997)

- Structure of halo nuclei \rightarrow matter form factors, radii

nucleus	B_{nnc} [keV]	B_{nc} [keV]	$\sqrt{\langle r_{nn}^2 \rangle}$ [fm]	$\sqrt{\langle r_{nc}^2 \rangle}$ [fm]
^{14}Be	1120	-200.0	4.1 ± 0.5	3.5 ± 0.5
^{20}C	3506	161	2.8 ± 0.3	2.4 ± 0.3
	3506	530	3.0 ± 0.7	2.5 ± 0.6
	3506	60	2.8 ± 0.2	2.3 ± 0.2
$^{20}\text{C}^*$	65 ± 6.8	60	42 ± 3	38 ± 3

Canham, HWH, Eur. Phys. J. A **37** (2008) 367

(cf. Yamashita, Tomio, Frederico, 2004)

- Input: TUNL Nuclear data evaluation project, ...

- Experiment: $^{14}\text{Be} \rightarrow \sqrt{\langle r_{nn}^2 \rangle} = (5.4 \pm 1.0)$ fm

(Marques et al., Phys. Rev. C **64** (2001) 061301)

- **Universality at large scattering length**
 - Limit cycle in three-body system \Leftrightarrow **Efimov physics**
 - Universal correlations (Phillips, Tjon line,...)
- **Applications in atomic, nuclear, and particle physics**
 - Cold atoms close to Feshbach resonance
 - Halo nuclei
 - Scattering properties of the $X(3872)$
- **Future directions:**
 - **Cold atoms:** $N \geq 4$, 2d-systems, ...
 - **Halo nuclei:** reactions, external currents, ...
 - **Hadronic molecules:** universal properties, Efimov effect?
 - **Three-nucleon system on the lattice:** limit cycle in “deformed” QCD?