

Spectral Signature of the FFLO order parameter in one-dimensional optical lattices

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Haus Bergkranz, Riezlern, Austria

- Imbalanced Fermionic system: theory and experiments
- Mean-field approximation for 1D systems: lattice Bogoliubov-deGennes (BdG) formalism
- Radio-Frequency (RF) spectroscopy of a 1D imbalanced gas.

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Fermi systems with imbalanced spin population

Breakdown of superconductivity in a magnetic field:

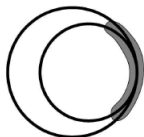
$$\mu_B H \approx \Delta \quad \text{Pauli limit}$$

Clogstone PRL, **9**, 266 (1962)

- FFLO (Fulde-Ferrel-Larkin-Ovchinnikov) state:
non-uniform order parameter



$$k_{F\uparrow} = k_{F\downarrow}, \quad Q = 0$$

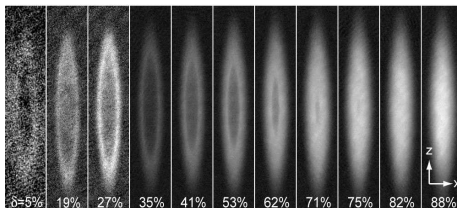


$$k_{F\uparrow} \neq k_{F\downarrow}, \quad Q = k_{F\uparrow} - k_{F\downarrow}$$

$$\Delta(\mathbf{r}) = \Delta_0 e^{i2\mathbf{q}\cdot\mathbf{r}} \quad (\text{single mode})$$

- FFLO observation in superconductors is still under debate:
 - H.A. Radovan *et.al*, Nature **425**, 51 (2003).
 - A. Bianchi *et.al*, PRL **91**, 187004 (2003).
 - K. Kakuyanagi *et.al*, PRL **44**, 047602 (2005).

$$\delta = \frac{N_{\uparrow} - N_{\downarrow}}{N_{\uparrow} + N_{\downarrow}}$$



M.W. Zwierlein *et.al*, Science **311**, 492 (2006)

M.W. Zwierlein *et.al*, Nature **442**, 54 (2006)

Y. Shin *et.al*, PRL **97**,030401 (2006).

Population imbalance \iff magnetic field in a superconductor.

How to detect FFLO in Fermi gases?

Till now, there is no clear manifestation of FFLO pairing mechanism in Fermi gases.

Proposal: using noise correlation measurement

Topic Review: Kun Yang, cond-mat/0603190

M. Rizzi *et.al*, Phys. Rev. B **77**, 245105 (2008)

Recent work: J.M. Edge and N.R. Cooper, arXiv:0906.1801
(collective modes)

Our proposal: using RF-spectroscopy to detect FFLO phase

Hubbard description for a trapped 1D Fermi gas

$$\mathcal{H} = -t \sum_{i,\sigma} \left(\hat{c}_{i,\sigma}^\dagger \hat{c}_{i+1,\sigma} + h.c \right) - U \sum_i \hat{n}_{i\uparrow} \hat{n}_{i\downarrow} \\ + \sum_{i,\sigma} (V_i^{\text{ext}} - \mu_\sigma) \hat{c}_{i,\sigma}^\dagger \hat{c}_{i,\sigma}$$

$$\hat{n}_{i\sigma} = \hat{c}_{i,\sigma}^\dagger \hat{c}_{i,\sigma}$$

$$V_i^{\text{ext}} = V_0 \left(i - \frac{L}{2} \right)^2 \quad \text{external trapping potential}$$

Validity of mean-field approx. in 1D



In 1D long-range order is absent (**Luttinger liquid paradigm**) so the mean-field approx is not valid as 3D. Specially the value of order parameter deviates from the exact solution.

F. Marsiglio, PRB **55**, 575 (1997)

Mean-field Hamiltonian for 1D interacting trapped Fermi gas

$$\begin{aligned}\mathcal{H}_{mf} = & -t \sum_{i,\sigma} \left(\hat{c}_{i,\sigma}^\dagger \hat{c}_{i+1,\sigma} + h.c \right) + \sum_i \left(\Delta_i \hat{c}_{i,\uparrow}^\dagger \hat{c}_{i,\downarrow}^\dagger + h.c \right) \\ & + \sum_{i,\sigma} \left\{ \left(V_i^{\text{ext}} - \mu_\sigma \right) - \underbrace{U \bar{n}_{i,\bar{\sigma}}}_{\text{Hartree term: neglect}} \right\} \hat{c}_{i,\sigma}^\dagger \hat{c}_{i,\sigma}\end{aligned}$$

$$\Delta_i \equiv -U \langle \hat{c}_{i,\downarrow} \hat{c}_{i,\uparrow} \rangle \quad \text{local pairing gap}$$

discrete Bogoliubov-deGennes (BdG) formalism

- Bogoliubov transformation

$$\hat{c}_{i,\sigma} = \sum_{\alpha} \left(u_{\alpha i \sigma} \hat{\gamma}_{\alpha \sigma} - \sigma v_{\alpha i \sigma}^* \hat{\gamma}_{\alpha \bar{\sigma}}^{\dagger} \right)$$

- diagonalized H

$$\mathcal{H}_{mf} = E_0 + \sum_{\alpha \sigma} E_{\alpha \sigma} \hat{\gamma}_{\alpha \sigma}^{\dagger} \hat{\gamma}_{\alpha \sigma}$$

- BdG equations

$$\sum_{j=1}^L \begin{pmatrix} \mathcal{H}_{ij}^{\sigma} & \Delta_{ij} \\ \Delta_{ij} & -\mathcal{H}_{ij}^{\bar{\sigma}} \end{pmatrix} \begin{pmatrix} u_{\alpha j \sigma} \\ v_{\alpha j \bar{\sigma}} \end{pmatrix} = E_{\alpha \sigma} \begin{pmatrix} u_{\alpha i \sigma} \\ v_{\alpha i \bar{\sigma}} \end{pmatrix}$$

$$\mathcal{H}_{ij}^{\sigma} = -t \delta_{i,j\pm 1} + (V_i^{\text{ext}} - \mu_{\sigma}) \delta_{ij}$$

$$\Delta_{ij} = \Delta_i \delta_{ij}$$

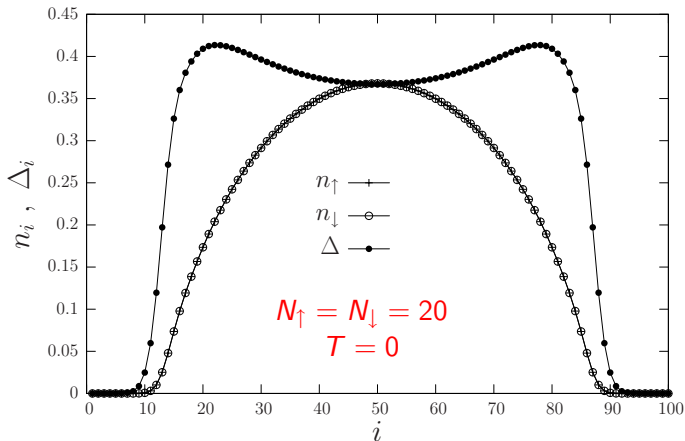
Self-consistent eqns. for spin-resolved density and gap

$$n_i^\sigma = \sum_{\alpha=1}^L \left[|u_{\alpha i \sigma}|^2 n_{\text{F}}(E_{\alpha \sigma}) + |v_{\alpha i \sigma}|^2 n_{\text{F}}(-E_{\alpha \bar{\sigma}}) \right]$$

$$\Delta_i = -U \sum_{\alpha=1}^L \left[u_{\alpha i \uparrow} v_{\alpha i \downarrow} n_{\text{F}}(E_{\alpha \uparrow}) - u_{\alpha i \downarrow} v_{\alpha i \uparrow} n_{\text{F}}(-E_{\alpha \downarrow}) \right]$$

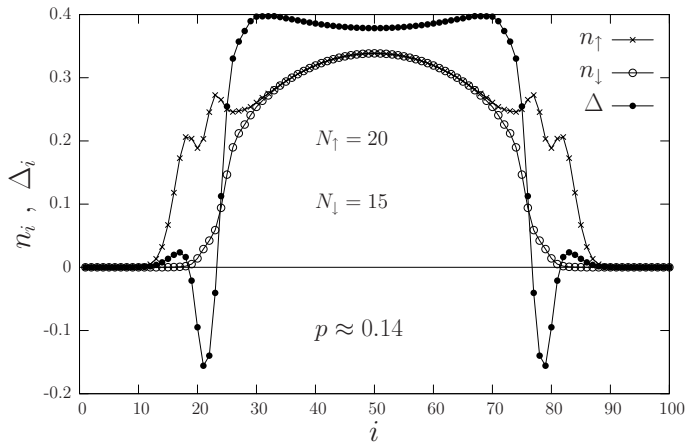
$$\sum_i n_i^\sigma = N_\sigma$$

Unpolarized gas



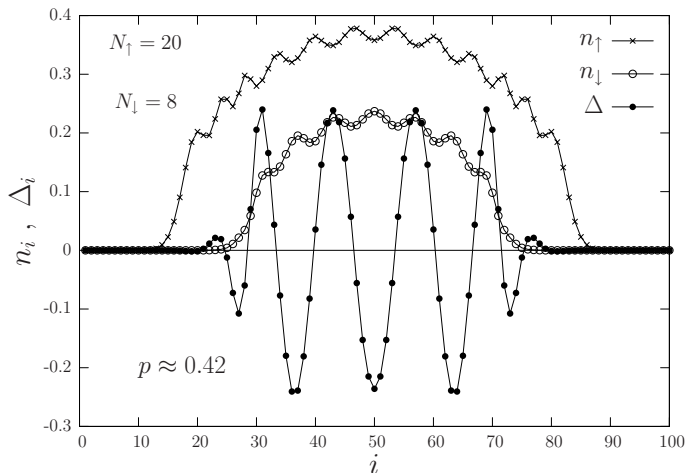
BCS-type of pairing: almost constant order-parameter.

Imbalanced gas



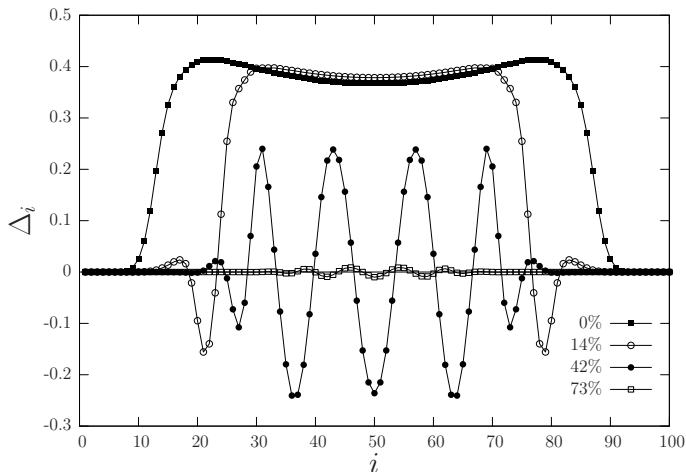
BCS pairing at the center and FFLO-type at the edges.

Imbalanced gas



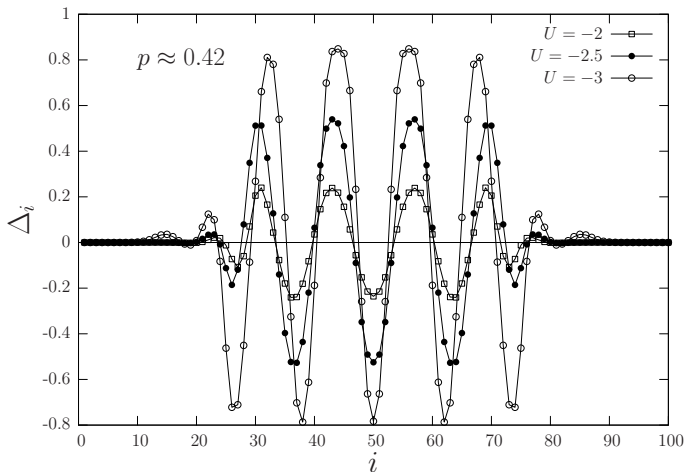
Minority component determines the extrema for order-parameter.

Order-parameter evolution



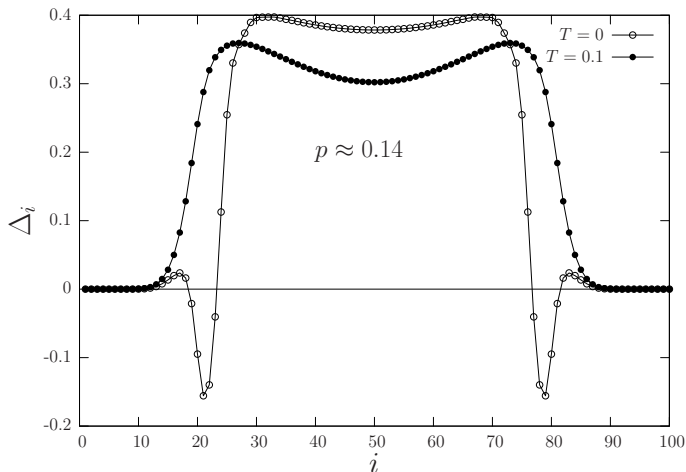
The oscillations become weaker and more confined upon increasing p .

Effect of interaction

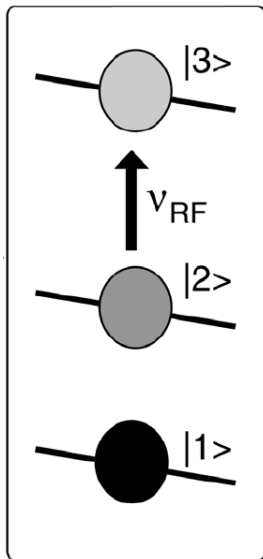


Stronger interaction enhance the oscillations.

Finite temperature effect



The oscillation disappear for high- T .



Original idea: using laser light to make a transition from an initially interactive state to a third empty state.

Goal: to detect BCS ground state and measure superconducting (-fluidity) order parameter.

P. Törma and P. Zoller, PRL **85**, 487 (2000)

Observation of pairing gap:

Experiment: Rudi Grimm group, *Science* **305**, 1128 (2004)

Theory: Päivi Törmä group, *Science* **305**, 1131 (2004)

Latest RF experiments

C.H. Schunck *et.al*, *Science* **316**, 867 (2007)

C.H. Schunck *et.al*, *Nature* **454**, 739 (2008)

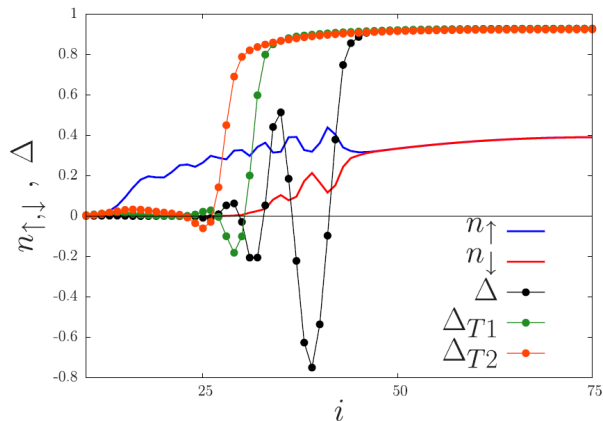
J.T. Stewart *et.al*, *Nature* **454**, 744 (2008)

A. Schirotzek *et.al*, *Phys.Rev.Lett* **102**, 230402 (2009)

$$\begin{aligned}
 J_{\uparrow/\downarrow}(\delta, K) = & - 2\pi \sum_{\alpha=1}^L \left[\left| \sum_{i=1}^L v_{\alpha i \downarrow/\uparrow} v_{Ki \uparrow/\downarrow}^{\text{non}} \right|^2 n_{\text{F}}(-E_{\alpha \downarrow/\uparrow}) \right. \\
 & \times \delta(E_{\alpha \downarrow/\uparrow} + \epsilon_K - \delta - \mu_{\uparrow/\downarrow}) \\
 & + \left| \sum_{i=1}^L u_{\alpha i \uparrow/\downarrow} v_{Ki \uparrow/\downarrow}^{\text{non}} \right|^2 n_{\text{F}}(E_{\alpha \uparrow/\downarrow}) \\
 & \left. \times \delta(E_{\alpha \uparrow/\downarrow} - \epsilon_K + \delta + \mu_{\uparrow/\downarrow}) \right]
 \end{aligned}$$

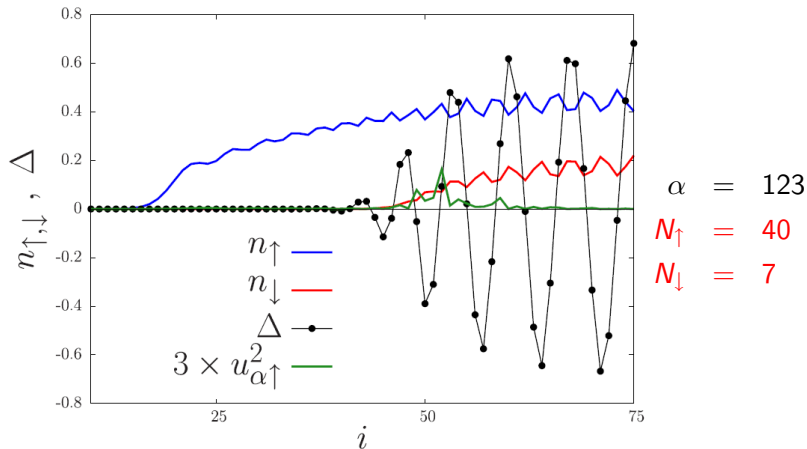
$$\delta = \omega_{\text{rf}} - (\omega_{\text{f}} - \omega_{\uparrow/\downarrow}) \quad \epsilon_K = \mu_{\uparrow/\downarrow} - E_{K \downarrow/\uparrow} \quad J_{\uparrow/\downarrow}(\delta) = \sum_{K=1}^L J_{\uparrow/\downarrow}(\delta, K)$$

density and gap profiles for more realistic system size

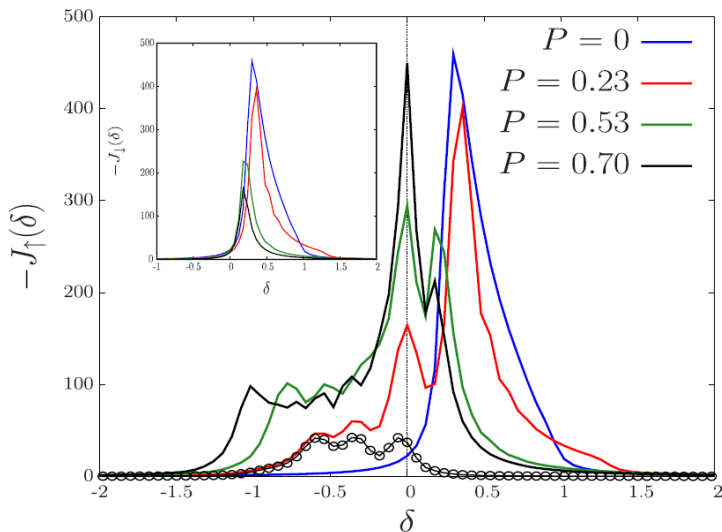


$$\begin{aligned} U &= -3 \\ V_0 &= 5 \times 10^{-4} \\ T_1 &= 0.075 \\ T_2 &= 0.1 \\ L &= 150 \\ N_{\uparrow} &= 40 \\ N_{\downarrow} &= 25 \end{aligned}$$

also the wave function for higher polarization

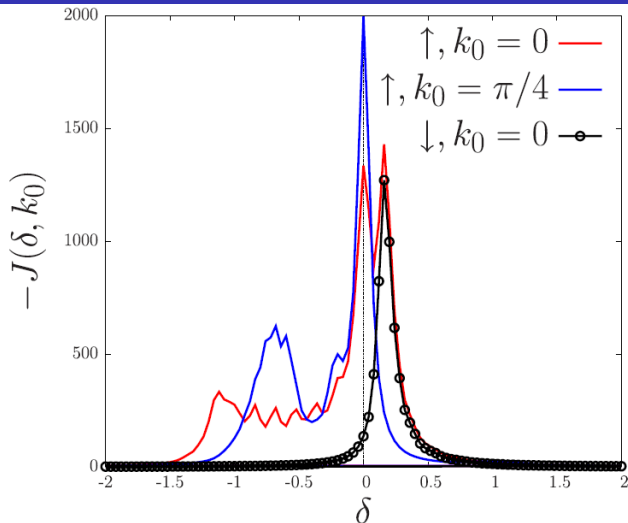


majority and minority spectra for different polarization



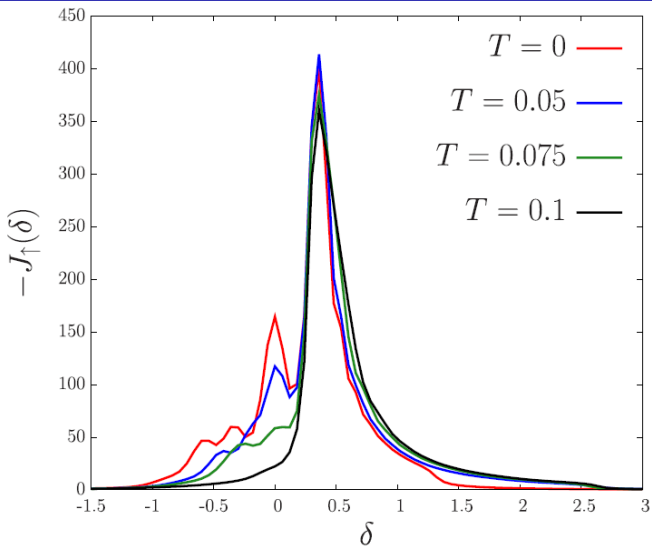
peaks at $\delta > 0$: paired atoms, peak at $\delta = 0$ unpaired majority.

Final state momentum-resolved spectra



$$J_{\uparrow/\downarrow}(\delta, k) = \sum_{K=1}^L \left| \sum_{i=1}^L e^{2(ki)} v_{Ki\uparrow/\downarrow}^{\text{non}} \right|^2 J_{\uparrow/\downarrow}(\delta, K)$$

T -dependence of majority component spectra



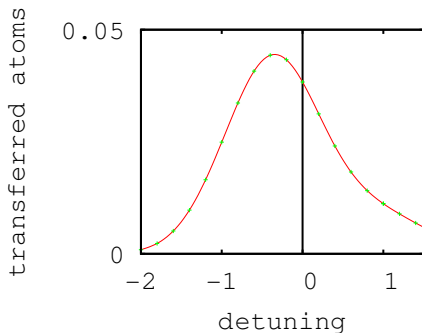
$T \gg \dots \implies$ vanish order parameter \implies disappear spectra at $\delta < 0$

validity of mean-field picture

comment about mean-field vs DMRG

R.A. Molina *et.al*, PRL, **102**, 168901 (2009)

M.R. Bakhtiari *et.al*, PRL, **102**, 168902 (2009)



calculated by Matrix Product States (MPS)

Summary and future plan

- FFLO signature of a 1D imbalanced gas at mean-field level
- finite temperature behavior of density and gap profile
- proposal on order parameter detection using RF spectra
- **Future plan:** FFLO with DMFT: higher D and arbitrary U
initial work: M. Snoek *et.al*, NJP **10**, 093008 (2008)
- **Q: Does FFLO observation make the life easier?**