



4d Phase Space Measurements

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Outline

- 4d rms emittance and eigen-emittances
- Motivation
- ROSE principle
- Error mitigation
- Results and error analysis
- Summary and outlook

- linear (4d), Hamiltonian beam line elements preserve :

- 4d rms emittance $E_{4d}^2 = \det \begin{bmatrix} \langle xx \rangle & \langle xx' \rangle & \langle xy \rangle & \langle xy' \rangle \\ \langle x'x \rangle & \langle x'x' \rangle & \langle x'y \rangle & \langle x'y' \rangle \\ \langle yx \rangle & \langle yx' \rangle & \langle yy \rangle & \langle yy' \rangle \\ \langle y'x \rangle & \langle y'x' \rangle & \langle y'y \rangle & \langle y'y' \rangle \end{bmatrix}$

- the two eigen-emittances

$$\varepsilon_1 = \frac{1}{2} \sqrt{-\text{tr}[(CJ)^2] + \sqrt{\text{tr}^2[(CJ)^2] - 16\det(C)}}$$

$$E_{4d} = \varepsilon_1 \cdot \varepsilon_2$$

$$\varepsilon_2 = \frac{1}{2} \sqrt{-\text{tr}[(CJ)^2] - \sqrt{\text{tr}^2[(CJ)^2] - 16\det(C)}}$$

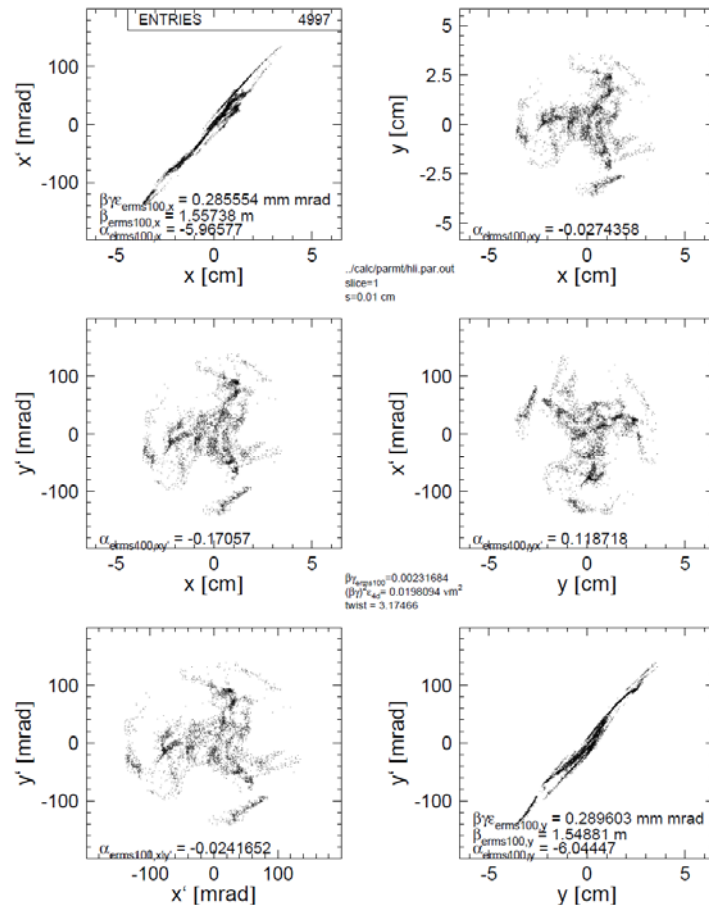
$$C = \begin{bmatrix} \langle xx \rangle & \langle xx' \rangle & \langle xy \rangle & \langle xy' \rangle \\ \langle x'x \rangle & \langle x'x' \rangle & \langle x'y \rangle & \langle x'y' \rangle \\ \langle yx \rangle & \langle yx' \rangle & \langle yy \rangle & \langle yy' \rangle \\ \langle y'x \rangle & \langle y'x' \rangle & \langle y'y \rangle & \langle y'y' \rangle \end{bmatrix} \quad J := \begin{bmatrix} 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & -1 & 0 \end{bmatrix}$$

- if, and only if there is no $x \leftrightarrow y$ correlation, i.e. $C = \begin{bmatrix} \langle xx \rangle & \langle xx' \rangle & 0 & 0 \\ \langle x'x \rangle & \langle x'x' \rangle & 0 & 0 \\ 0 & 0 & \langle yy \rangle & \langle yy' \rangle \\ 0 & 0 & \langle y'y \rangle & \langle y'y' \rangle \end{bmatrix}$
 - rms emittances = eigen-emittances
- if there is any coupling
 - rms emittances \neq eigen-emittances
 - coupling parameter $t = \frac{\epsilon_x \epsilon_y}{\epsilon_1 \epsilon_2} - 1 \geq 0$

- term „eigen-emittance“ is quite unknown, since generally beams are considered as uncoupled
- measured moments are just reliable, if they deliver reasonable eigen-emittances !

Motivation for 4d Diagnostics: ECR Source Performance Evaluation

4d distribution behind ECR source



- $\epsilon_{\text{rms},x} = 123 \text{ mm mrad}$
- $\epsilon_{\text{rms},y} = 125 \text{ mm mrad}$
- $E_1 = 17 \text{ mm mrad}$
- $E_2 = 231 \text{ mm mrad}$
- $\epsilon_{4d} = E_1 \cdot E_2 = 3927 \text{ (mm mrad)}^2$
- $\epsilon_{\text{rms},x} \cdot \epsilon_{\text{rms},y} = 15375 \text{ (mm mrad)}^2$
- $\epsilon_{\text{rms},x} \cdot \epsilon_{\text{rms},y} = 3.9 \epsilon_{4d}$
- to quantify the source performance, knowledge of $\epsilon_{\text{rms},x}$ and $\epsilon_{\text{rms},y}$ is not sufficient
- knowledge of E_1 and E_2 is required
- the correlation moments $\langle xy \rangle \dots$ etc are needed
- they give access to E_1 and E_2

Motivation: Envelopes along Solenoid Channel

initial beam:

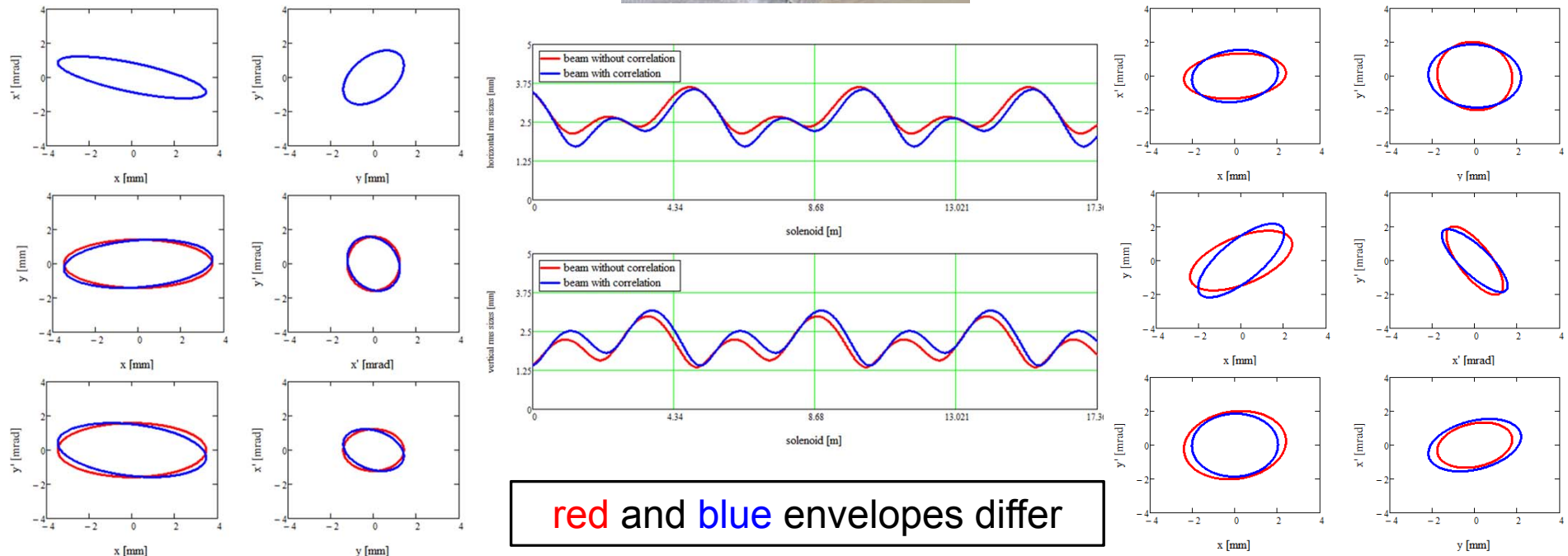
$$C_1 = \begin{bmatrix} 12.00 & -3.00 & 0.00 & 0.00 \\ -3.00 & 1.50 & 0.00 & 0.00 \\ 0.00 & 0.00 & 2.00 & 1.00 \\ 0.00 & 0.00 & 1.00 & 2.50 \end{bmatrix} \text{ uncorrelated}$$

$$C_2 = \begin{bmatrix} 12.00 & -3.00 & 1 & -1.5 \\ -3.00 & 1.50 & -0.5 & -0.35 \\ 1.00 & -0.50 & 2.00 & 1.00 \\ -1.50 & -0.35 & 1.00 & 2.50 \end{bmatrix} \text{ correlated}$$

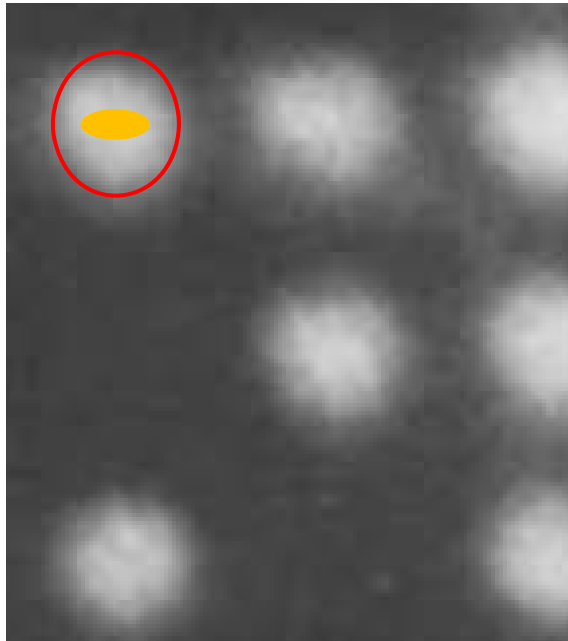


$$M_{sol} = \begin{bmatrix} C^2 & \frac{SC}{K} & SC & \frac{S^2}{K} \\ -KSC & C^2 & -KS^2 & CS \\ -SC & -\frac{S^2}{K} & C^2 & \frac{SC}{K} \\ KS^2 & -SC & -KSC & C^2 \end{bmatrix}$$

final beam:



Emittance Measurements: Pepper Pot (Praxis)



screen image size

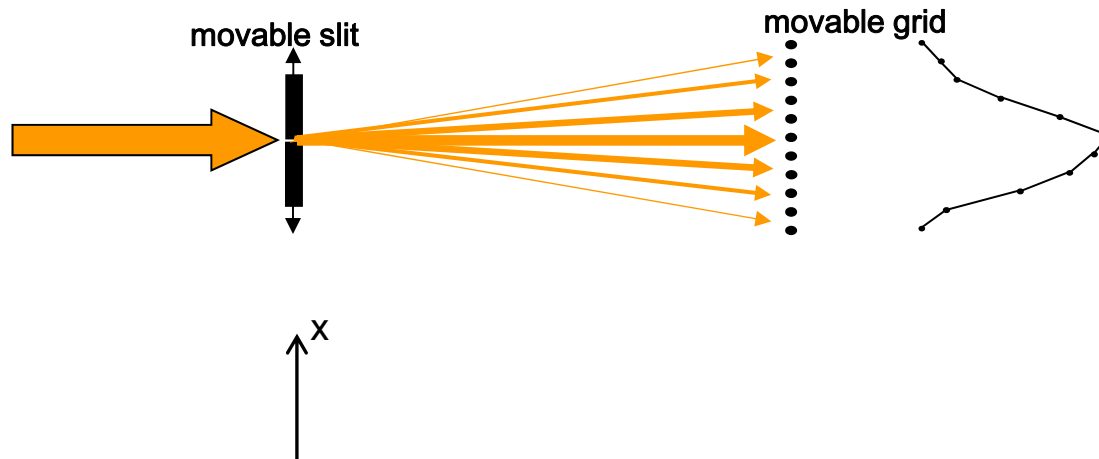
beamlet size

- to our knowledge, pepper pots do not work for ions with energies beyond ≈ 150 keV/u
- size of the beamlet image is dominated by screen resolution
- resolution limited by re-fluorescence
- excessive material research without success so far
- beamlet angles are always over-estimated
- GSI's energies are in MeV/u range
- pepper pots measure just few percent of the beam, as the mask dumps the main part

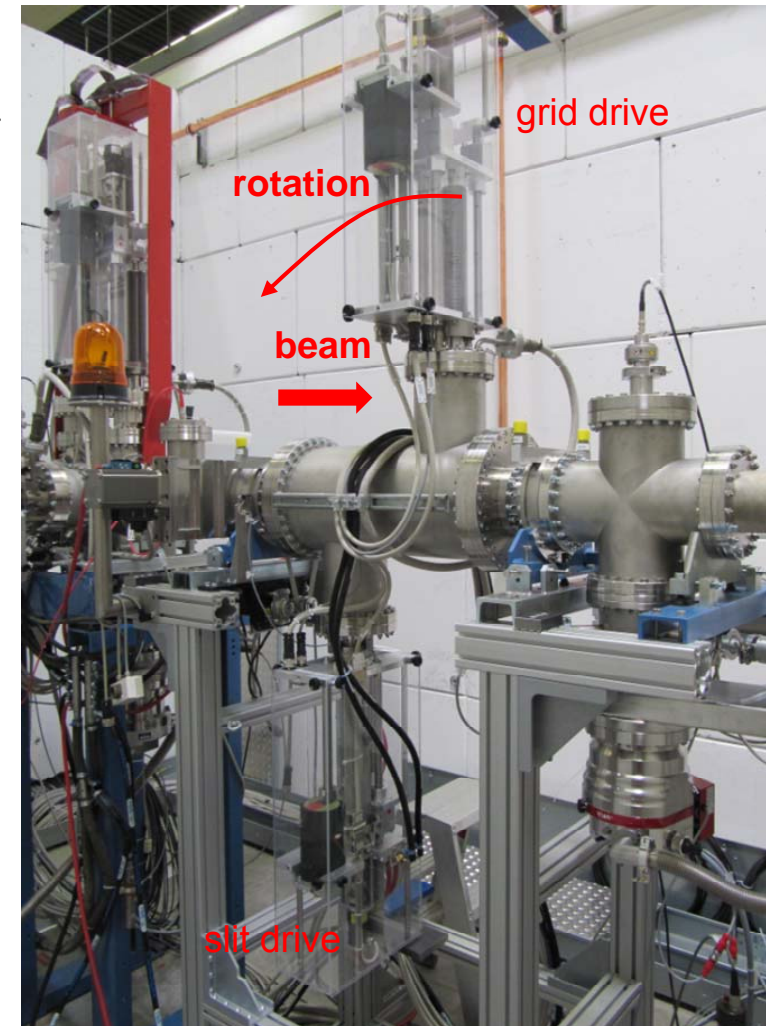
It seems that ion beam community misses a reliable tool for complete 4d transverse diagnostics

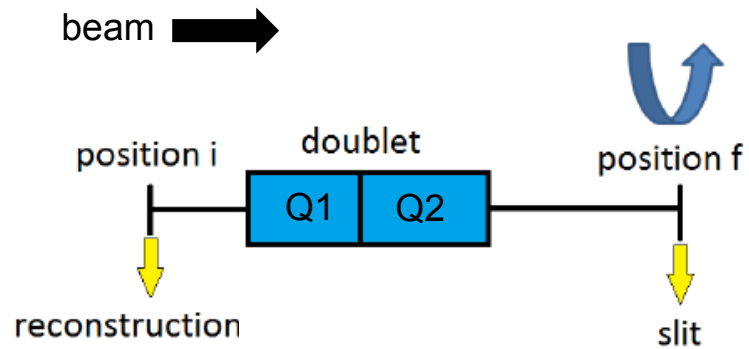
ROSE: Concept

ROtating System for Emittance Measurements

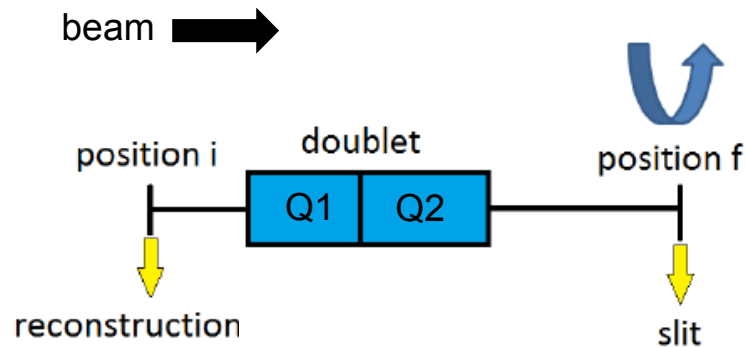


- slit & grid being mounted inside a rotatable chamber
- rotation gives access to $\langle xy \rangle$, $\langle xy' \rangle$, $\langle x'y \rangle$, $\langle x'y' \rangle$
- chamber is fixed during measurement
- chamber does not rotate during measurement





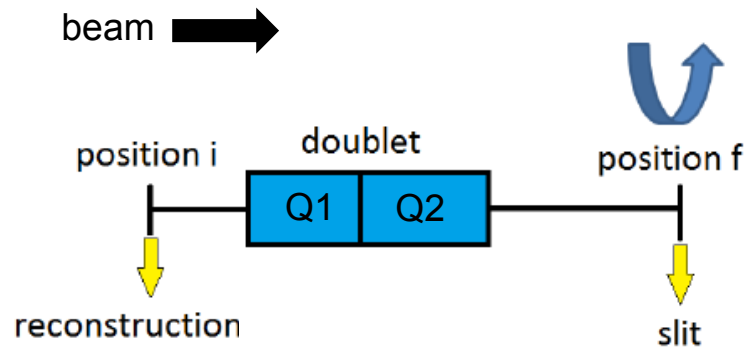
- beam moments to be determined at „reconstruction point“
- „reconstruction point“ and ROSE are separated by adjustable non-coupling element
- slit & grid of ROSE can be rotated simultaneously



four measurements are needed to determine the correlated moments:

1. at 0° with setting $Q1^a$, $Q2^a$
2. at θ with setting $Q1^a$, $Q2^a$
3. at 90° with setting $Q1^a$, $Q2^a$
4. at θ with setting $Q1^b$, $Q2^b$

doublet settings a & b deliver different envelopes

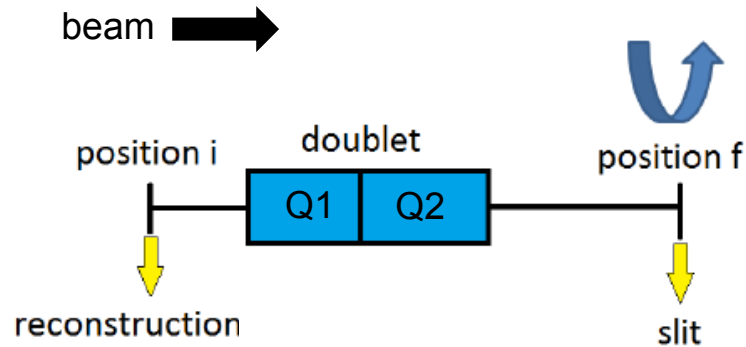


- beam transport to ROSE

$$\begin{bmatrix} x \\ x' \\ y \\ y' \end{bmatrix}_f^{a,b} = \begin{bmatrix} m_{11}^{a,b} & m_{12}^{a,b} & 0 & 0 \\ m_{21}^{a,b} & m_{22}^{a,b} & 0 & 0 \\ 0 & 0 & m_{33}^{a,b} & m_{34}^{a,b} \\ 0 & 0 & m_{43}^{a,b} & m_{44}^{a,b} \end{bmatrix} \begin{bmatrix} x \\ x' \\ y \\ y' \end{bmatrix}_i$$

- rotation of ROSE by θ is equivalent to rotation of beam by $-\theta$

$$\begin{bmatrix} x \\ x' \\ y \\ y' \end{bmatrix}_\theta^{a,b} = \begin{bmatrix} \cos\theta & 0 & \sin\theta & 0 \\ 0 & \cos\theta & 0 & \sin\theta \\ -\sin\theta & 0 & \cos\theta & 0 \\ 0 & -\sin\theta & 0 & \cos\theta \end{bmatrix} \begin{bmatrix} x \\ x' \\ y \\ y' \end{bmatrix}_f^{a,b}$$



- solving all equations finally gives an over-determined system of equations

$$\begin{bmatrix} \Gamma_{11} & \Gamma_{12} & \Gamma_{13} & \Gamma_{14} \\ \Gamma_{21} & \Gamma_{22} & \Gamma_{23} & \Gamma_{24} \\ \Gamma_{31} & \Gamma_{32} & \Gamma_{33} & \Gamma_{34} \\ \Gamma_{41} & \Gamma_{42} & \Gamma_{43} & \Gamma_{44} \\ \Gamma_{51} & \Gamma_{52} & \Gamma_{53} & \Gamma_{54} \\ \Gamma_{61} & \Gamma_{62} & \Gamma_{63} & \Gamma_{64} \end{bmatrix} \begin{bmatrix} \langle xy \rangle_i \\ \langle xy' \rangle_i \\ \langle x'y \rangle_i \\ \langle x'y' \rangle_i \end{bmatrix} = \begin{bmatrix} \Lambda_1 \\ \Lambda_2 \\ \Lambda_3 \\ \Lambda_4 \\ \Lambda_5 \\ \Lambda_6 \end{bmatrix}$$

- Γ -matrix includes quad strengths $Q1^{a,b}$ and $Q2^{a,b}$
- Λ -vector includes individual measurements of $\langle xx \rangle^{a,b,\theta}$, $\langle xx' \rangle^{a,b,\theta}$, and $\langle x'x' \rangle^{a,b,\theta}$
- $\langle \dots \rangle_i$ are the correlated beam moments to be determined

- final solution is

$$\begin{bmatrix} \langle xy \rangle_i \\ \langle xy' \rangle_i \\ \langle x'y \rangle_i \\ \langle x'y' \rangle_i \end{bmatrix} = (\Gamma^T \Gamma)^{-1} \Gamma^T \Lambda$$

- individually measured $\langle xx \rangle_{a,b,\theta}$, $\langle xx' \rangle_{a,b,\theta}$, and $\langle x'x' \rangle_{a,b,\theta}$ inhabit measurement errors
- errors will enter into the final result for $\langle \dots \rangle_i$
- final errors are minimized if **condition number K** of the matrix Γ is minimized

$$\begin{bmatrix} \langle xy \rangle_i \\ \langle xy' \rangle_i \\ \langle x'y \rangle_i \\ \langle x'y' \rangle_i \end{bmatrix} = (\Gamma^T \Gamma)^{-1} \Gamma^T \Lambda \quad \Gamma^\dagger = (\Gamma^T \Gamma)^{-1} \Gamma^T \quad \|\Gamma\|_2 := \sqrt{\sum_{i=1}^n \sum_{j=1}^k (\Gamma_{i,j})^2}, \quad \|\Gamma^\dagger\|_2 := \sqrt{\sum_{i=1}^k \sum_{j=1}^n (\Gamma_{i,j}^\dagger)^2}, \quad \kappa(\Gamma) := \|\Gamma\|_2 \|\Gamma^\dagger\|_2$$

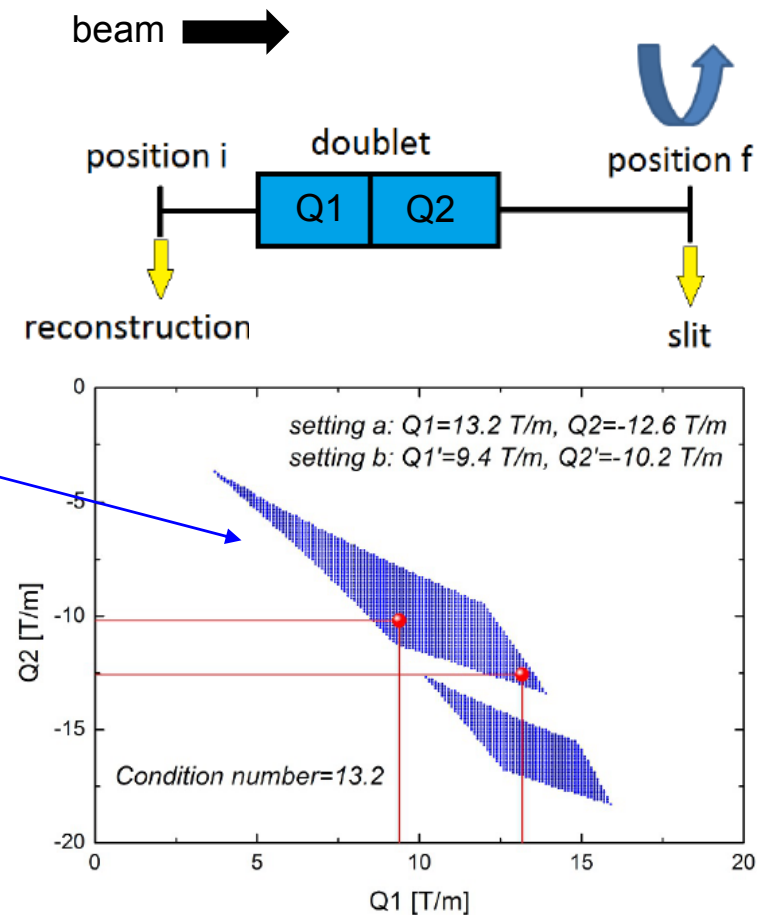
- K quantifies linear dependency of the matrix rows
- low linear dependency \rightarrow low $K \rightarrow$ errors contribute few to final result
- examples:

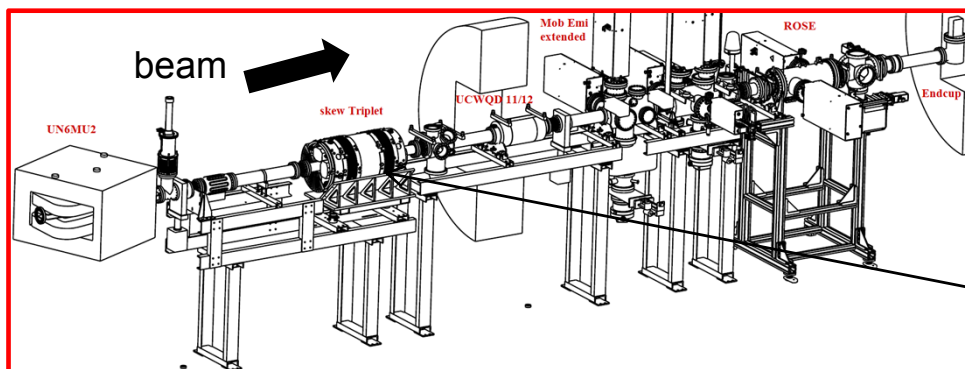
$$\begin{bmatrix} 5 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 2 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \rightarrow K = 5$$

$$\begin{bmatrix} 1 & 3 & 5 & 7 \\ 2.1 & 5.9 & 10.2 & 14.1 \\ 0.55 & 1.49 & 2.55 & 3.45 \\ 0.11 & 0.28 & 0.51 & 0.69 \end{bmatrix} \rightarrow K = 1.2 \times 10^{17}$$

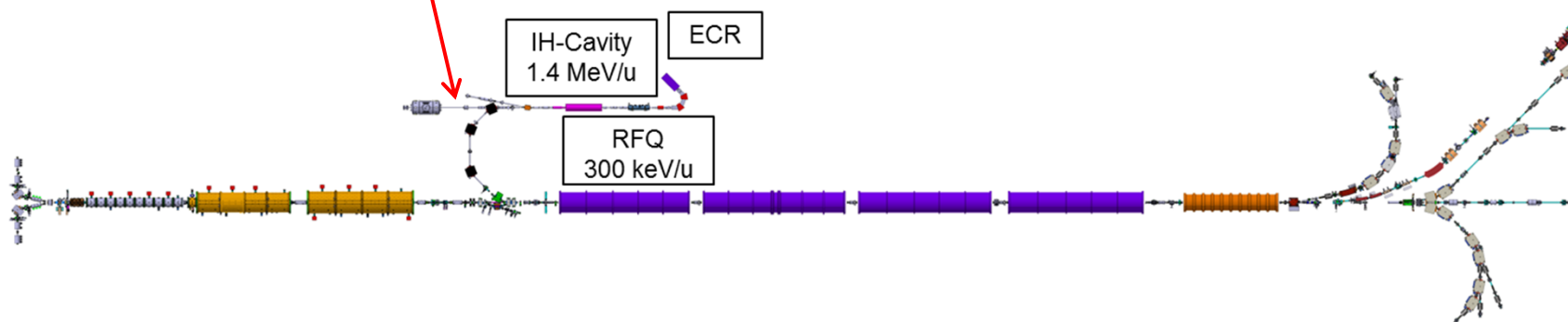
Measurement Error Mitigation

- do initial emittance measurements at 0° and 90° , i.e. in ver. and hor. plane
- backtransform to quad doublet entrance
- vary in brute force way two quad strengths $Q1$, $Q2$ and check for (analytically!) :
 - 100% transmission
 - reasonable beam size at slit
 - reasonable beam size at grid
- store settings $Q1$, $Q2$ that passed this test
- build all possible pairs of settings, check their K
- pick **pair with lowest K** for measurements:
 1. slit at 0° with setting $Q1^a$, $Q2^a$
 2. slit at θ with setting $Q1^a$, $Q2^a$
 3. slit at 90° with setting $Q1^a$, $Q2^a$
 4. slit at θ with setting $Q1^b$, $Q2^b$



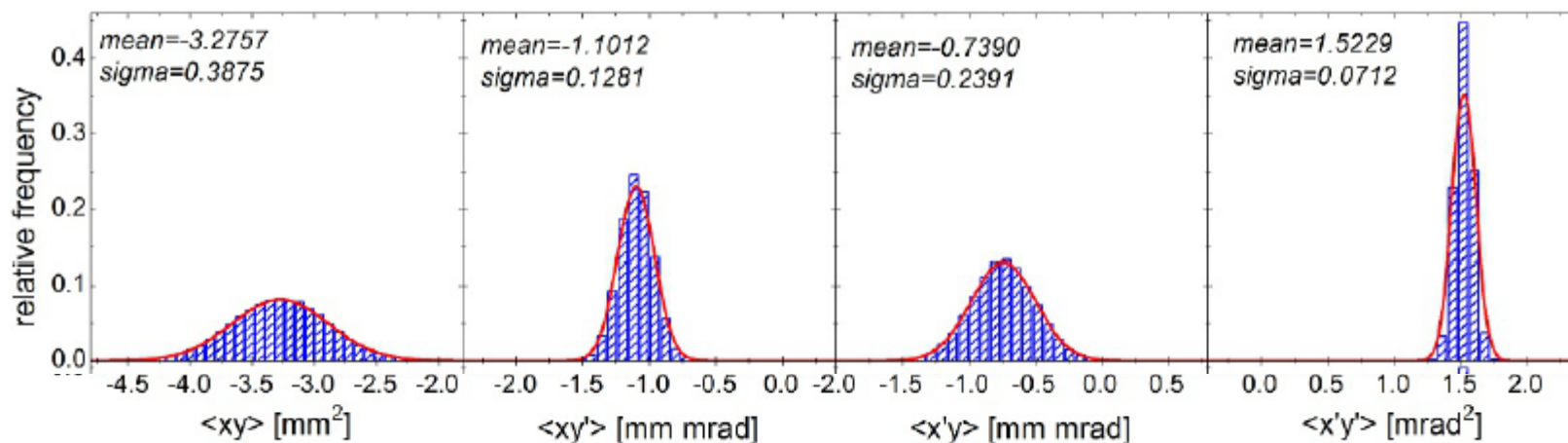


- measurements at exit of GSI's HLI
- 1.4 MeV/u
- low current of $^{83}\text{Kr}^{13+}$
- skew triplet to create $x \leftrightarrow y$ correlations
- full transmission



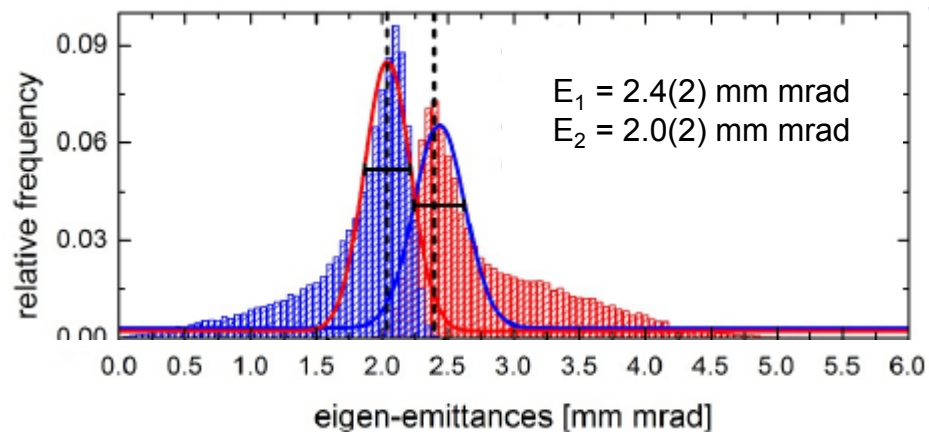
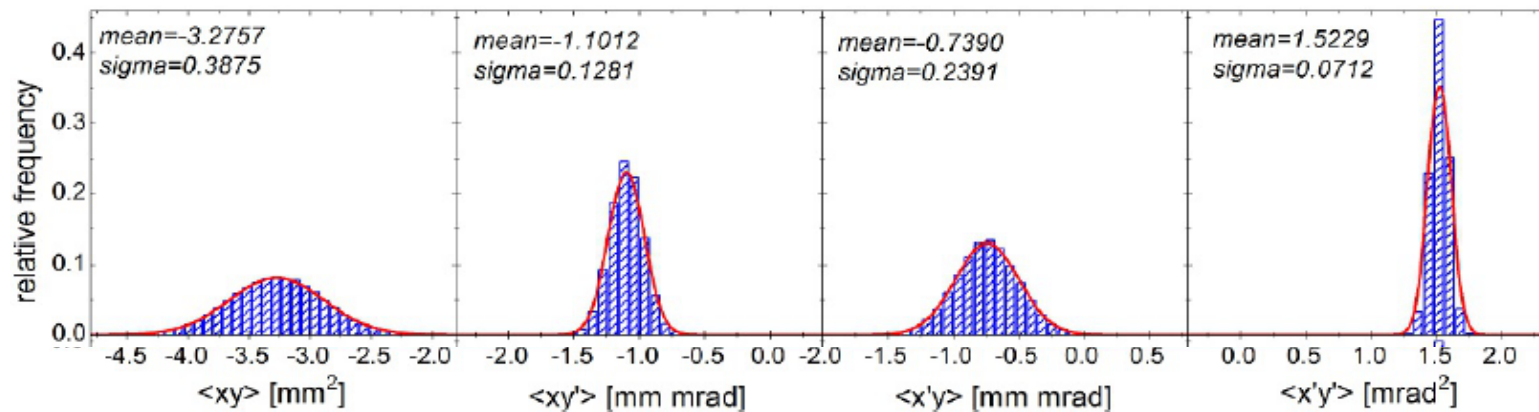
$$\begin{bmatrix} \langle xx \rangle & \langle xx' \rangle & \langle xy \rangle & \langle xy' \rangle \\ \langle x'x \rangle & \langle x'x' \rangle & \langle x'y \rangle & \langle x'y' \rangle \\ \langle yx \rangle & \langle yx' \rangle & \langle yy \rangle & \langle yy' \rangle \\ \langle y'x \rangle & \langle y'x' \rangle & \langle y'y \rangle & \langle y'y' \rangle \end{bmatrix} = \begin{bmatrix} 8.57 & -4.34 & -3.28 & -1.10 \\ -4.34 & 3.35 & -0.74 & 1.52 \\ -3.28 & -0.74 & 11.20 & -3.05 \\ -1.10 & 1.52 & -3.05 & 1.87 \end{bmatrix} \quad \text{in mm, mrad}$$

- each measured moment entering into the evaluation was varied randomly following a Gaussian distribution centered around its measured value ($3 \cdot \text{rms} = 10\%$)
- many sets of errors were used
- error bars for initial correlated moments were obtained



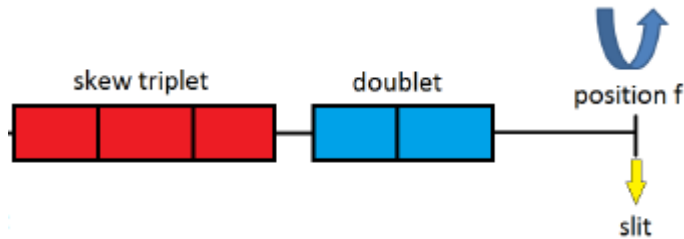
4d Emittance Measurements: Eigen-Emittances

$$\begin{bmatrix} \langle xx \rangle & \langle xx' \rangle & \langle xy \rangle & \langle xy' \rangle \\ \langle x'x \rangle & \langle x'x' \rangle & \langle x'y \rangle & \langle x'y' \rangle \\ \langle yx \rangle & \langle yx' \rangle & \langle yy \rangle & \langle yy' \rangle \\ \langle y'x \rangle & \langle y'x' \rangle & \langle y'y \rangle & \langle y'y' \rangle \end{bmatrix} = \begin{bmatrix} 8.57 & -4.34 & -3.28 & -1.10 \\ -4.34 & 3.35 & -0.74 & 1.52 \\ -3.28 & -0.74 & 11.20 & -3.05 \\ -1.10 & 1.52 & -3.05 & 1.87 \end{bmatrix} \quad \text{in mm, mrad}$$



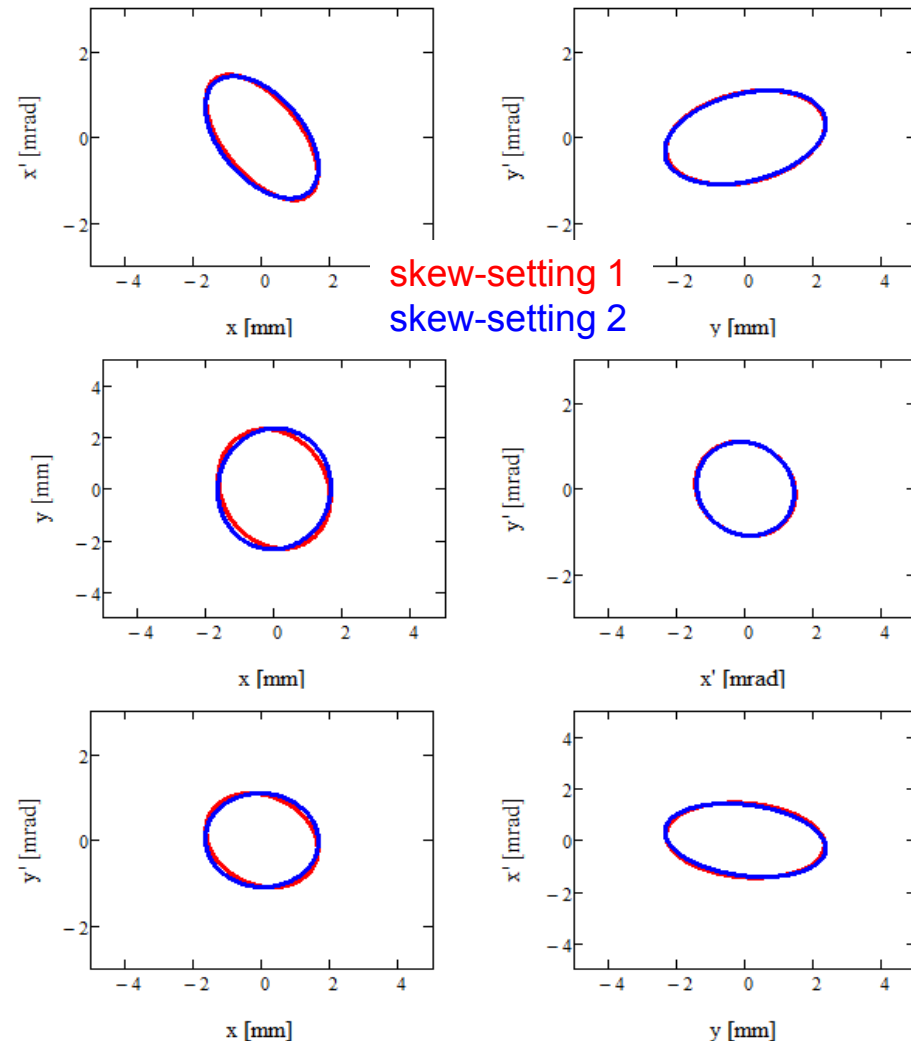
$$t = \frac{\epsilon_x \epsilon_y}{\epsilon_1 \epsilon_2} - 1 = 1.2(2)$$

4d Emittance Measurements: Check Reliability



- measurements were done for two different skew triplet settings
- both measurements were back-transformed to entrance of skew triplet
- back-transformations delivered very similar results
- the two measurements are consistent
- ROSE seems to be reliable

at skew triplet entrance:



- a new reliable device for full 4d transverse beam moments matrix measurements has been designed, built, and tested
- to our knowledge, it is the only reliable device working at ion energies beyond about 150 keV/u
- it was tested at the HLI of GSI at 1.4 MeV/u with $^{86}\text{Kr}^{13+}$
- device has a patent „*Deutsche Patentanmeldung Nr. 102015118017.0 eingereicht am 22.10.2015 beim Deutschen Patent- und Markenamt Titel der Patentanmeldung: Drehmodul für eine Beschleunigeranlage*“
- it will be installed:
 - behind the ECR test stand for source diagnostics
 - in the transfer channel to the synchrotron to detect/remove beam correlations
 - ...
- Literature:
 - Phys Rev. Accel. & Beams **19**, 072802 (2016)
 - Nucl. Instrum. & Meth. A 820 14, (2016)



rms-emittances defined through beam's second moments :

- a_i, b_i : two coordinates of particle i
- $\langle ab \rangle$: mean of product $a_i b_i$
- C is moment matrix (symmetric)

$$E_x^2 = \langle x x \rangle \langle x' x' \rangle - \langle x x' \rangle^2$$

$$C_x = \begin{bmatrix} \langle x x \rangle & \langle x x' \rangle \\ \langle x' x \rangle & \langle x' x' \rangle \end{bmatrix}, \quad E_x^2 = \det C_x$$

$$C_y = \begin{bmatrix} \langle y y \rangle & \langle y y' \rangle \\ \langle y' y \rangle & \langle y' y' \rangle \end{bmatrix}, \quad E_y^2 = \det C_y$$

linear transport from point_1 \rightarrow point_2 through matrices:

$$\begin{bmatrix} x \\ x' \end{bmatrix}_2 = M_x \begin{bmatrix} x \\ x' \end{bmatrix}_1 \quad M_x = \begin{bmatrix} m_{11} & m_{12} \\ m_{21} & m_{22} \end{bmatrix}, \quad \det M_x = 1$$

beam moments transport by matrix equation:

$$C_{x2} = M_x C_{x1} M_x^T$$

analogue in y

if x & y planes are not coupled

$$E_{4d}^2 = \det \begin{bmatrix} \langle xx \rangle & \langle xx' \rangle & \langle xy \rangle & \langle xy' \rangle \\ \langle x'x \rangle & \langle x'x' \rangle & \langle x'y \rangle & \langle x'y' \rangle \\ \langle yx \rangle & \langle yx' \rangle & \langle yy \rangle & \langle yy' \rangle \\ \langle y'x \rangle & \langle y'x' \rangle & \langle y'y \rangle & \langle y'y' \rangle \end{bmatrix}$$

transport of moments from 1 → 2 as usual :

$$\begin{bmatrix} x \\ x' \\ y \\ y' \end{bmatrix}_2 = \begin{bmatrix} m_{11} & m_{12} & m_{13} & m_{14} \\ m_{21} & m_{22} & m_{23} & m_{24} \\ m_{31} & m_{32} & m_{33} & m_{34} \\ m_{41} & m_{42} & m_{43} & m_{44} \end{bmatrix} \begin{bmatrix} x \\ x' \\ y \\ y' \end{bmatrix}_1, \quad \det M = 1$$

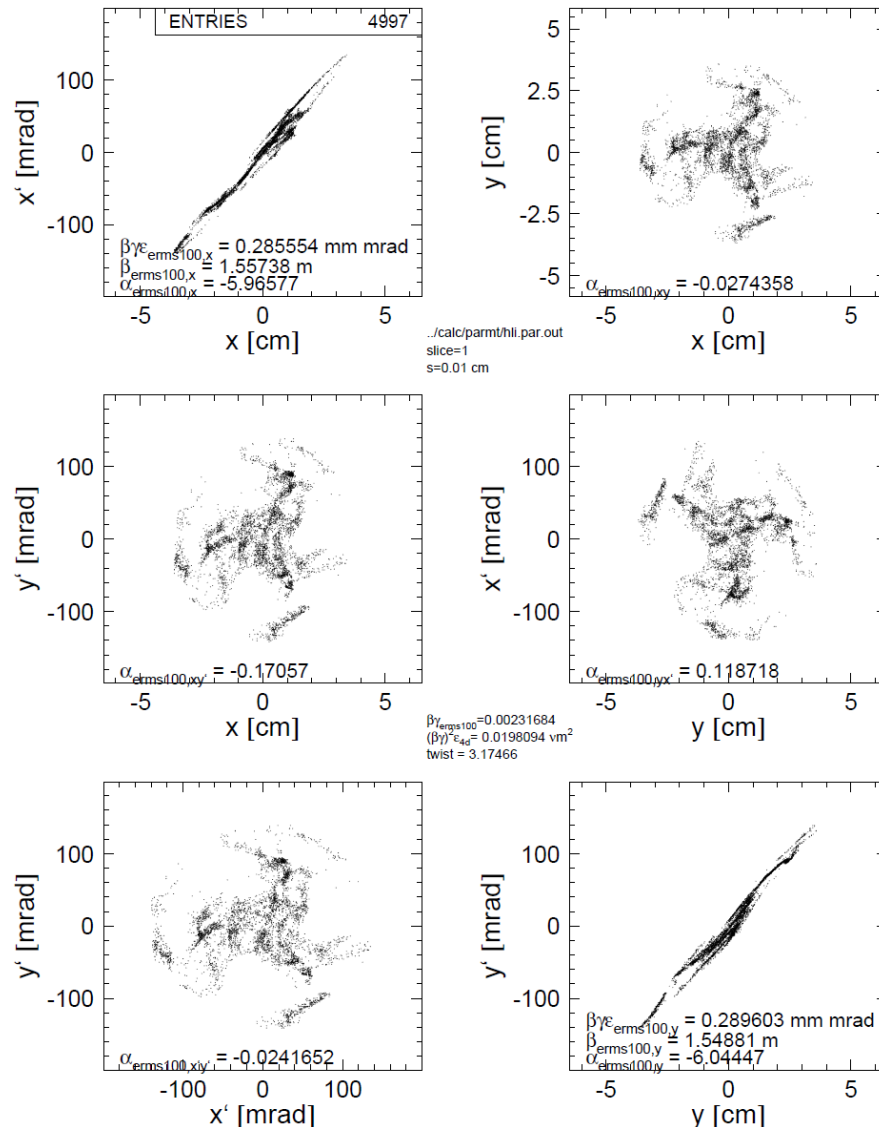
$$C_2 = M C_1 M^T$$

$$E_{4d}^2 = \det \begin{bmatrix} \langle xx \rangle & \langle xx' \rangle & 0 & 0 \\ \langle x'x \rangle & \langle x'x' \rangle & 0 & 0 \\ 0 & 0 & \langle yy \rangle & \langle yy' \rangle \\ 0 & 0 & \langle y'y \rangle & \langle y'y' \rangle \end{bmatrix} = (E_x \cdot E_y)^2$$

transport of moments from 1 → 2 as usual :

$$M = \begin{bmatrix} m_{11} & m_{12} & 0 & 0 \\ m_{21} & m_{22} & 0 & 0 \\ 0 & 0 & m_{33} & m_{34} \\ 0 & 0 & m_{43} & m_{44} \end{bmatrix}, \quad \det M = \det M_x \cdot \det M_y = 1 \cdot 1 = 1$$

$$C_2 = M C_1 M^T$$



4d-distribution behind ECR source

- $\epsilon_{rms,x} = 123 \text{ mm mrad}$
- $\epsilon_{rms,y} = 125 \text{ mm mrad}$
- $E_1 = 17 \text{ mm mrad}$
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- $\epsilon_{rms,4d} = E_1 \cdot E_2 = 3927 \text{ (mm mrad)}^2$
- $\epsilon_{rms,x} \cdot \epsilon_{rms,y} = 15375 \text{ (mm mrad)}^2$
- $\epsilon_{rms,x} \cdot \epsilon_{rms,y} = 3.9 \epsilon_{rms,4d}$

Motivation:

Envelopes along non-Coupling Focusing Elements



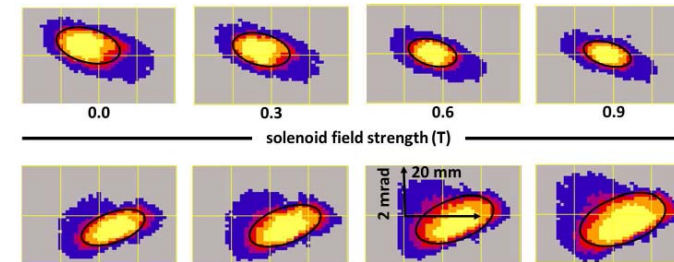
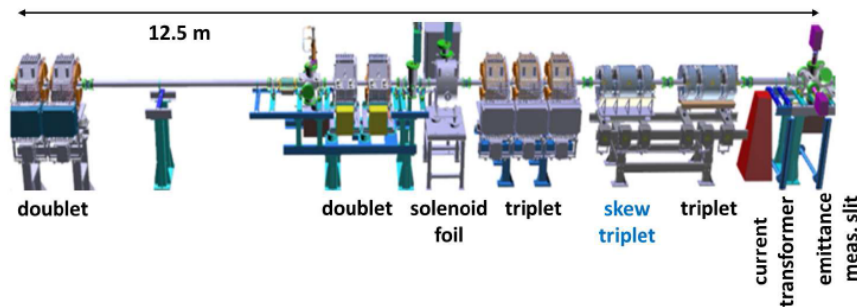
- if beam line comprises just elements that do not couple as $M = \begin{bmatrix} m_{11} & m_{12} & 0 & 0 \\ m_{21} & m_{22} & 0 & 0 \\ 0 & 0 & m_{33} & m_{34} \\ 0 & 0 & m_{43} & m_{44} \end{bmatrix}$,
the beam envelopes $\langle xx \rangle$ and $\langle yy \rangle$ do not depend on the initial correlated moments
- M has simply no access to $\langle xy \rangle$, $\langle xy' \rangle$... when applying $C_2 = M C_1 M^T$ to $C = \begin{bmatrix} \langle xx \rangle & \langle xx' \rangle & \langle xy \rangle & \langle xy' \rangle \\ \langle x'x \rangle & \langle x'x' \rangle & \langle x'y \rangle & \langle x'y' \rangle \\ \langle yx \rangle & \langle yx' \rangle & \langle yy \rangle & \langle yy' \rangle \\ \langle y'x \rangle & \langle y'x' \rangle & \langle y'y \rangle & \langle y'y' \rangle \end{bmatrix}$
- consequence for non-coupling lattices:
 - initial correlated moments can be ignored, even if they are non-zero
 - hor. / ver. envelopes can be calculated, if just the non-correlated moments are known, as $\langle xx \rangle$, $\langle xx' \rangle$,, $\langle y'y' \rangle$
- conventional emittance scanners are sufficient for beam dynamics layouts for non-coupling lattices

Motivation: Envelopes along Coupling Focusing Elements



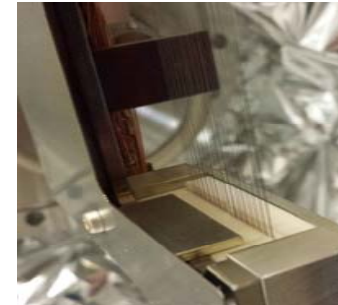
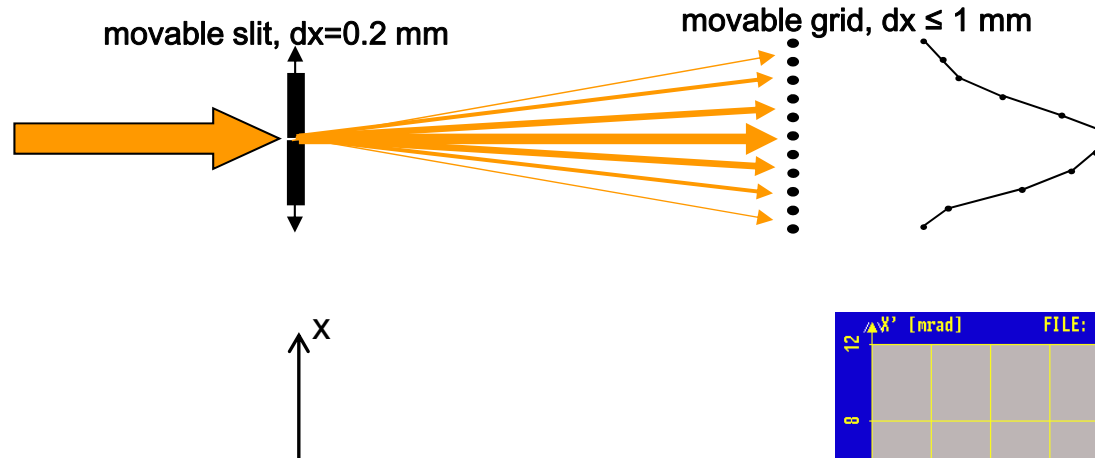
- if beam line comprises elements that do couple as $\begin{bmatrix} m_{11} & m_{12} & m_{13} & m_{14} \\ m_{21} & m_{22} & m_{23} & m_{24} \\ m_{31} & m_{32} & m_{33} & m_{34} \\ m_{41} & m_{42} & m_{43} & m_{44} \end{bmatrix}$,
the beam envelopes $\langle xx \rangle$ and $\langle yy \rangle$ do depend on the initial correlated moments
- M has access on $\langle xy \rangle$, $\langle xy' \rangle$... when applying $C_2 = M C_1 M^T$ to $C = \begin{bmatrix} \langle xx \rangle & \langle xx' \rangle & \langle xy \rangle & \langle xy' \rangle \\ \langle x'x \rangle & \langle x'x' \rangle & \langle x'y \rangle & \langle x'y' \rangle \\ \langle yx \rangle & \langle yx' \rangle & \langle yy \rangle & \langle yy' \rangle \\ \langle y'x \rangle & \langle y'x' \rangle & \langle y'y \rangle & \langle y'y' \rangle \end{bmatrix}$
- consequence for coupling beam lines:
 - initial correlated moments cannot be ignored, if they are non-zero
 - hor. / ver. envelopes can only be calculated, if also all correlated moments are known, as $\langle xy \rangle$, $\langle xy' \rangle$, $\langle x'y \rangle$, and $\langle x'y' \rangle$
- conventional emittance scanners are not sufficient for beam dynamics layouts for coupling lattices

Motivation: Diagnostics for Advanced EmTEX

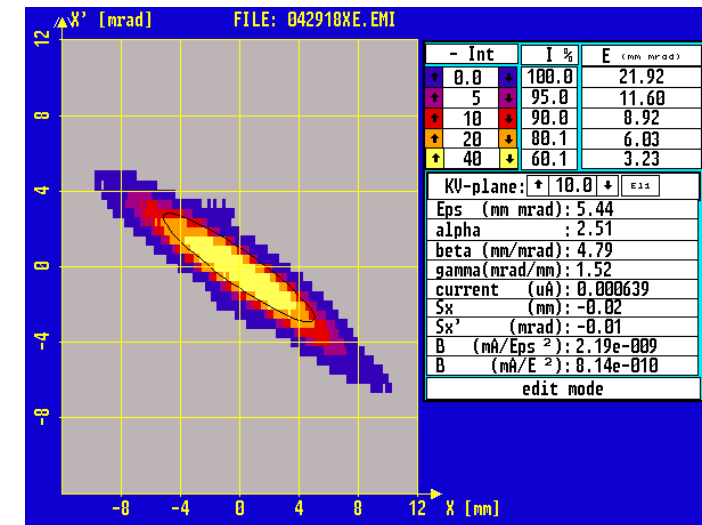


- EmTEX: Emittance Transfer Experiment with linear dynamics:
 - no space charge, low momentum spread, initially uncorrelated beam
 - was done just by knowing initial 2d-Twiss parameters before solenoid
 - coupling in solenoid, decoupling in skew triplet
 - beam in front of decoupling skew quads from pure linear 4d beam dynamics
- if advanced EmTEX shall be applied to high space charge beams:
 - linear 4d dynamics not applicable to deduce correlations in front of decoupling line
 - full 4d beam moments must be known in front of decoupling skew triplet
 - 4d diagnostics needed

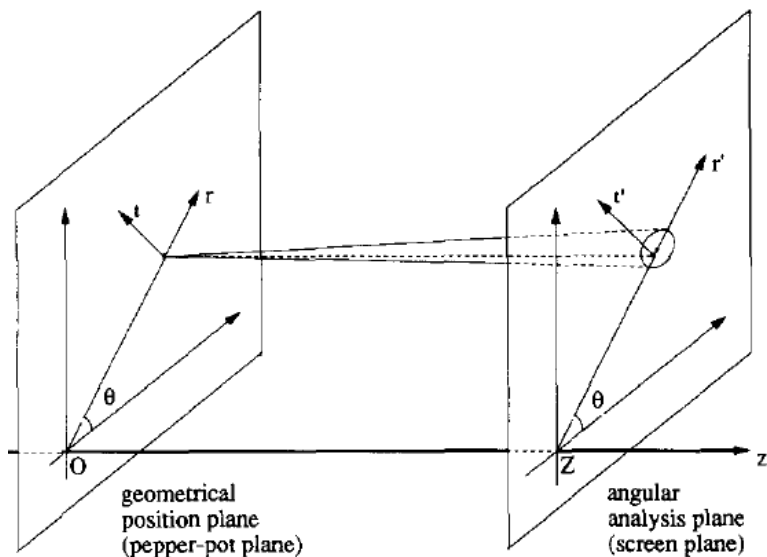
Emittance Measurements: Slit / Grid



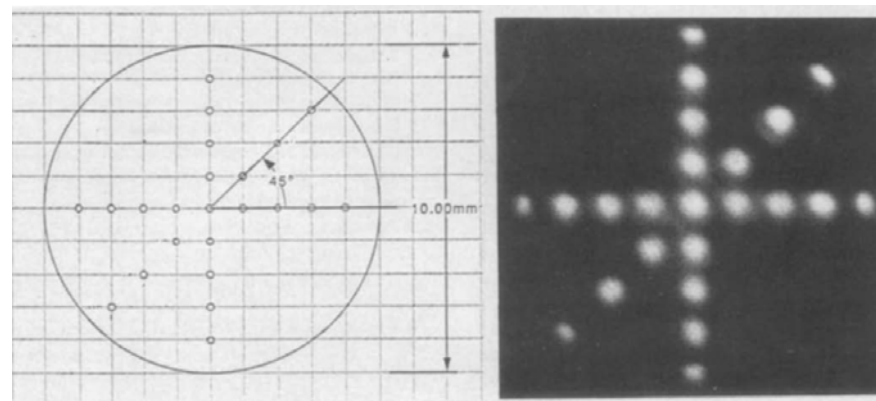
- slit position determines position x
- positions of wire and slit determine x'
- for each phase space pixel (x, x') wire current is recorded
- map of pixel intensities $I(x, x')$ is evaluated
- delivers $\langle xx \rangle$, $\langle xx' \rangle$, and $\langle x'x' \rangle$, i.e. the x -Twiss parameters



Emittance Measurements: Pepper-Pot (Concept)



J.G. Wang et al. NIM A 307, p. 190, (1991)



- hole position determines position x, y
- positions of hole and its image determine x', y'
- for each phase space pixel (x, x', y, y') screen light emission is recorded
- map of pixel intensities $I(x, x', y, y')$ is evaluated
- delivers 4d-Twiss parameters

- individually measured $\langle xx \rangle_{a,b,\theta}$, $\langle xx' \rangle_{a,b,\theta}$, and $\langle x'x' \rangle_{a,b,\theta}$ inhabit measurement errors
- errors will enter into the final result for $\langle \dots \rangle_i$
- final errors are minimized if „condition number K “ of the matrix Γ is minimized

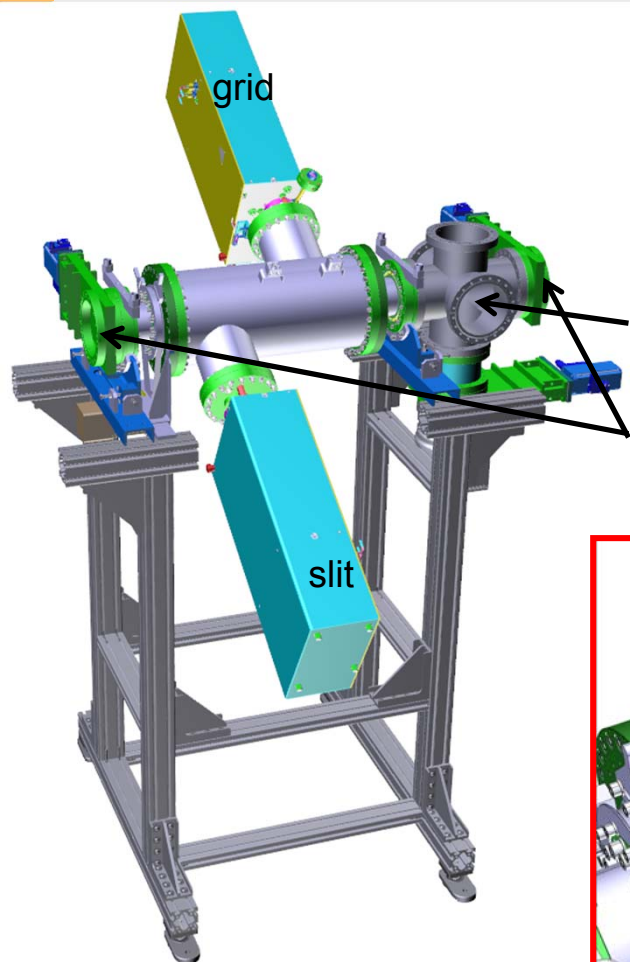
$$\begin{bmatrix} \langle xy \rangle_i \\ \langle xy' \rangle_i \\ \langle x'y \rangle_i \\ \langle x'y' \rangle_i \end{bmatrix} = (\Gamma^T \Gamma)^{-1} \Gamma^T \Lambda \quad \Gamma^\dagger = (\Gamma^T \Gamma)^{-1} \Gamma^T \quad \|\Gamma\|_2 := \sqrt{\sum_{i=1}^n \sum_{j=1}^k (\Gamma_{i,j})^2}, \quad \|\Gamma^\dagger\|_2 := \sqrt{\sum_{i=1}^k \sum_{j=1}^n (\Gamma_{i,j}^\dagger)^2}, \quad \kappa(\Gamma) := \|\Gamma\|_2 \|\Gamma^\dagger\|_2$$

- K quantifies linear dependency of the matrix rows
- low dependency \rightarrow low $K \rightarrow$ errors propagate less to final result
- examples:

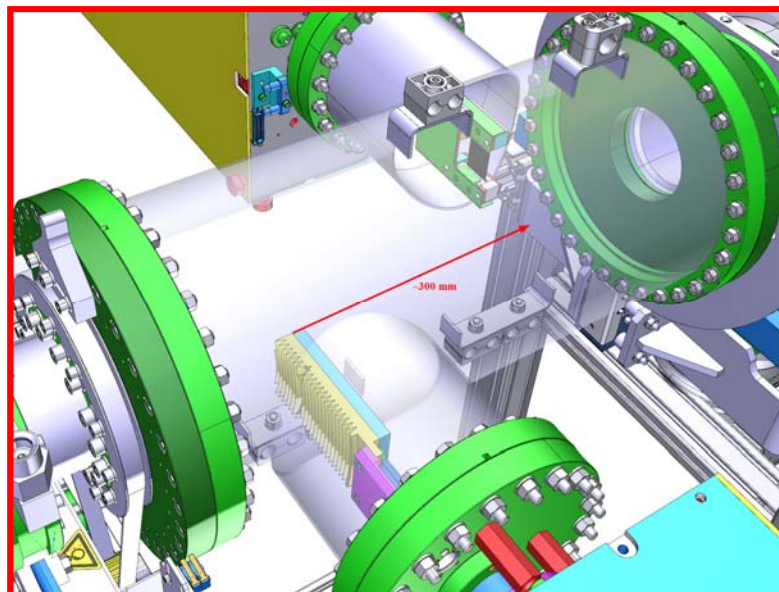
$$\begin{bmatrix} 5 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 2 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \rightarrow K = 5$$

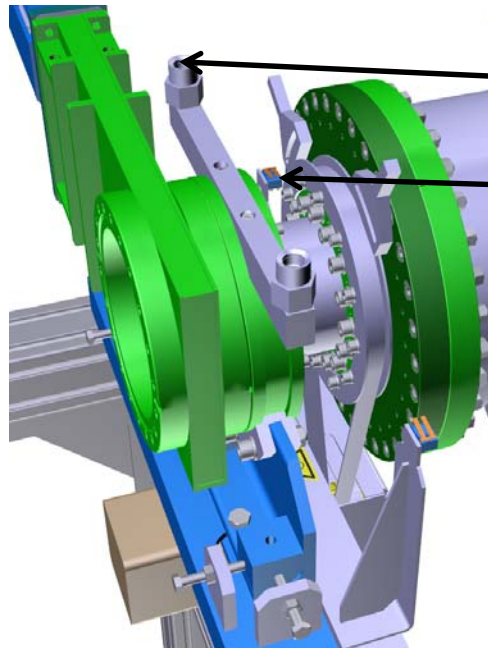
$$\begin{bmatrix} 1 & 3 & 5 & 7 \\ 2.1 & 5.9 & 10.2 & 14.1 \\ 0.55 & 1.49 & 2.55 & 3.45 \\ 0.11 & 0.28 & 0.51 & 0.69 \end{bmatrix} \rightarrow K = 1.2 \times 10^{17}$$

- accuracy of measurement is high for low K
- K is pure function of the transport matrix, i.e. of quadrupole settings a & b
- procedure to determine the two quad settings that minimize K :
 1. settings should provide 100% transmission up to the grid of ROSE
 2. settings should deliver lowest possible K

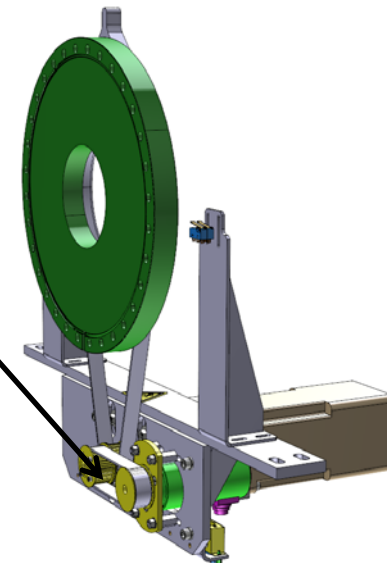
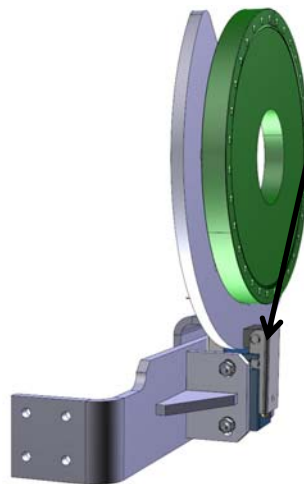


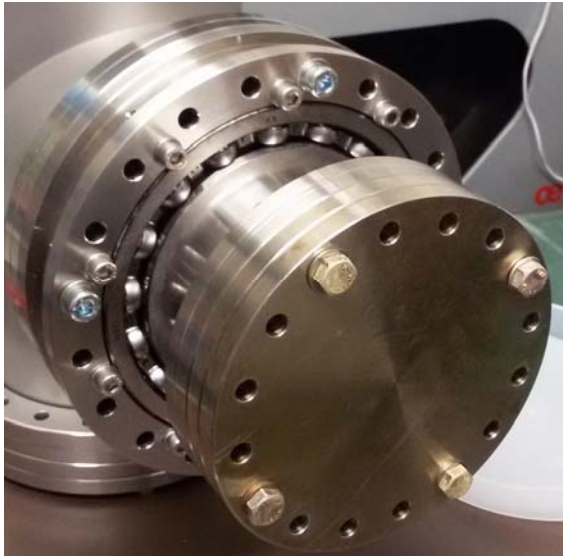
- slit & grid mounted on opposite sides to minimize torque
- cables wrapped around chamber
- fixed pumping chamber
- two gate valves (were not needed during rotation)



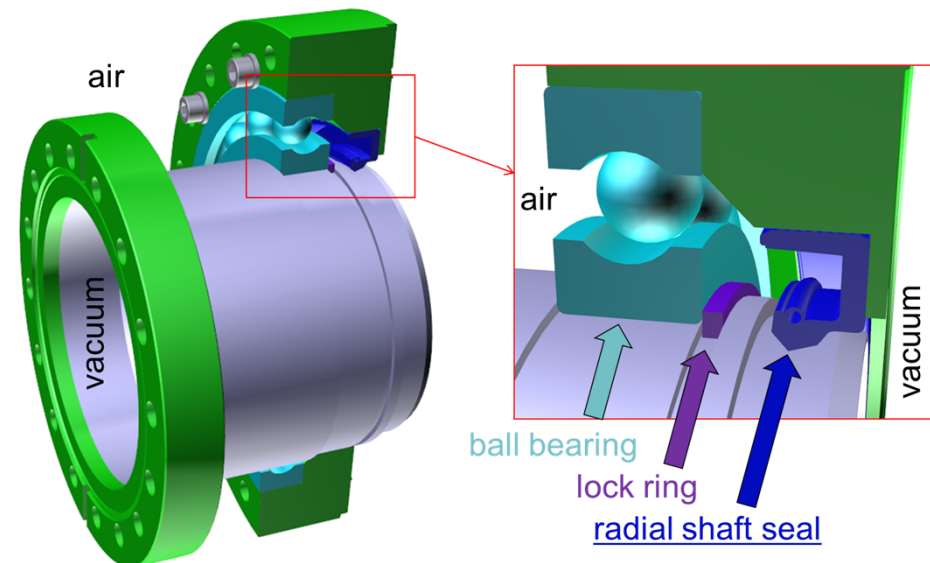
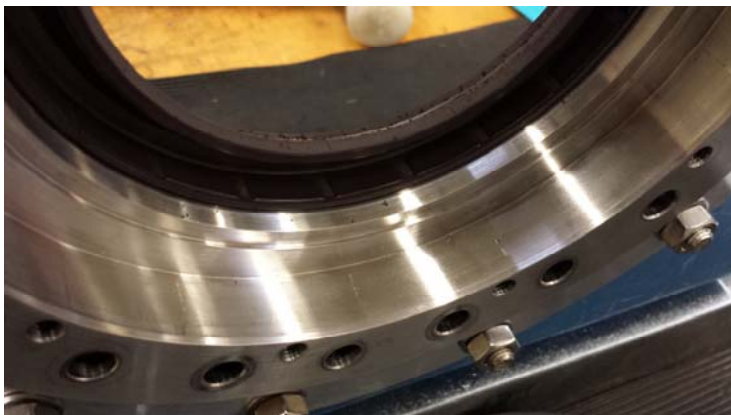


- fiducial points for alignment
- end switches separated by 180°
- disk brake (closed during measurements)
- motor driver with belts ($\delta\theta \leq 0.5^\circ$)
- rotation speed slowed actively down to $\approx 90^\circ/\text{min}$

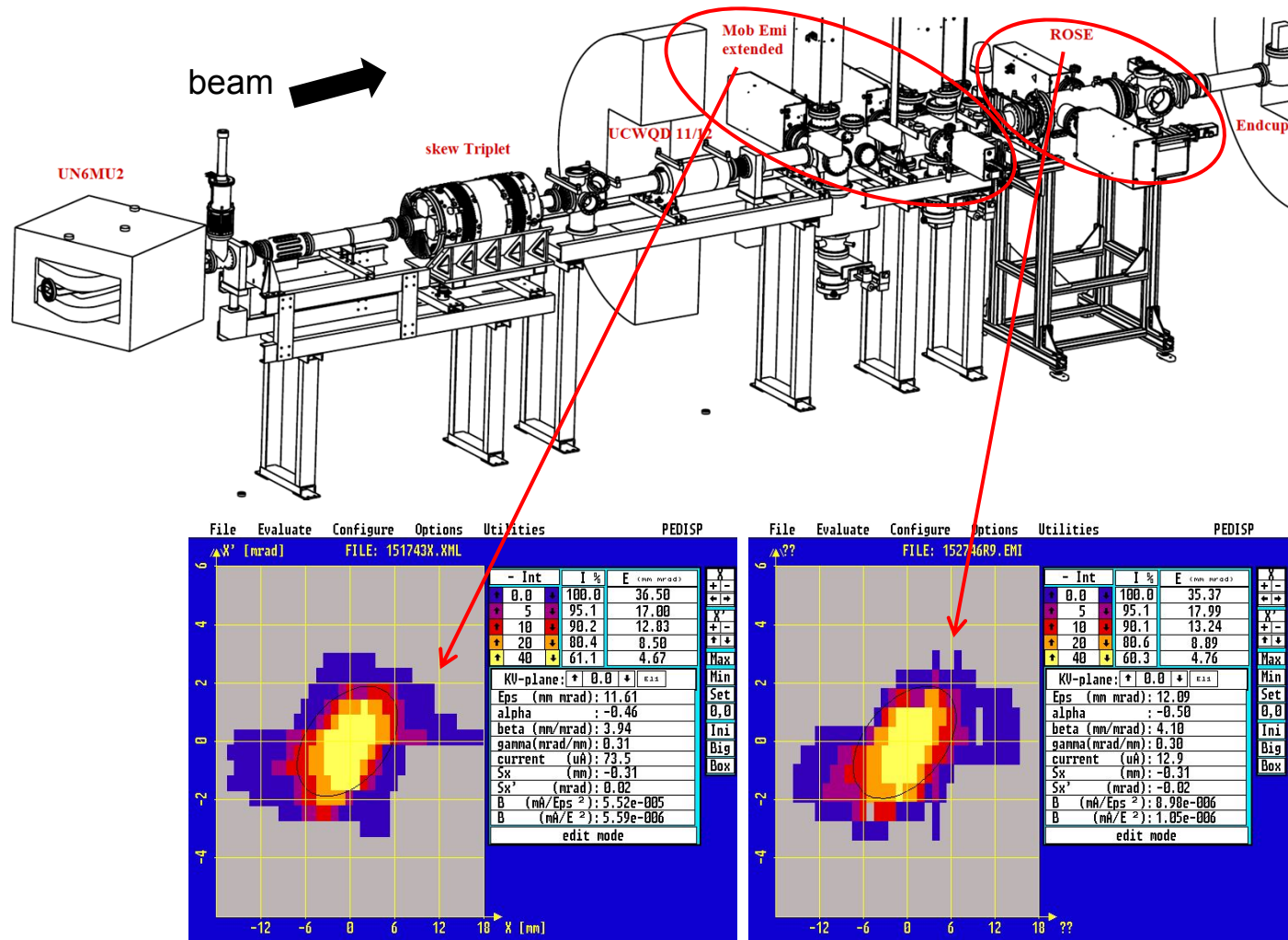




- static pressure $\approx 5 \times 10^{-8}$ mbar
- max. pressure during rotation $\approx 9 \times 10^{-8}$ mbar
- recovery time ≈ 1 min



ROSE ($\theta = 0^\circ$ and 90°) successfully benchmarked to well-tested „MobEmi“ – slit/grid device



4d Emittance Measurements

- initial measurements at 0° and 90° revealed at ROSE:

θ [deg]	ϵ_{rms} [mm mrad]	β [m]	α
0	3.1	4.6	-0.1
90	3.4	8.8	-2.4

- optimized quadrupole settings:
a: 13.2/-12.6 and b: 9.4/-10.2 T/m

- measurements:

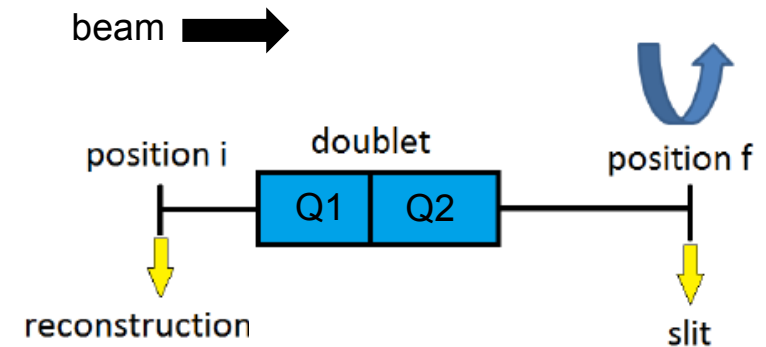
1. at 0° with setting Q1^a, Q2^a
2. at 30° with setting Q1^a, Q2^a
3. at 90° with setting Q1^a, Q2^a
4. at 30° with setting Q1^b, Q2^b

- evaluation:

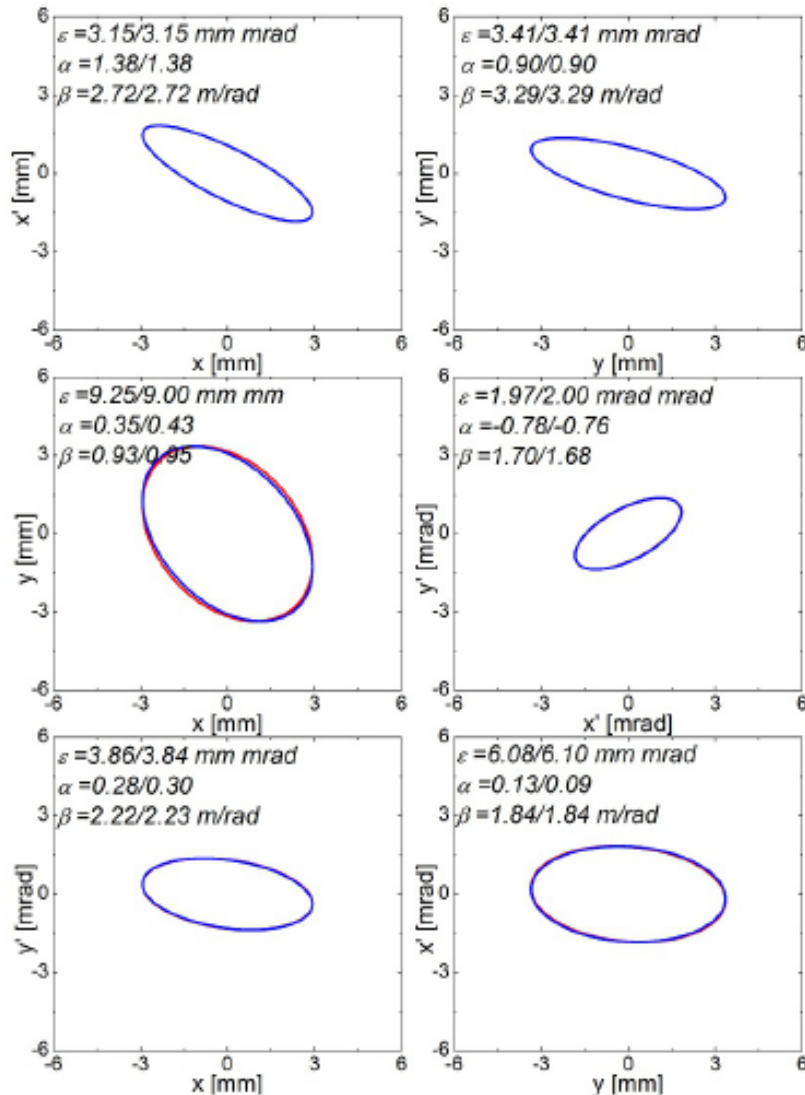
$$\begin{bmatrix} \langle xx \rangle & \langle xx' \rangle & \langle xy \rangle & \langle xy' \rangle \\ \langle x'x \rangle & \langle x'x' \rangle & \langle x'y \rangle & \langle x'y' \rangle \\ \langle yx \rangle & \langle yx' \rangle & \langle yy \rangle & \langle yy' \rangle \\ \langle y'x \rangle & \langle y'x' \rangle & \langle y'y \rangle & \langle y'y' \rangle \end{bmatrix} = \begin{bmatrix} 8.57 & -4.34 & -3.28 & -1.10 \\ -4.34 & 3.35 & -0.74 & 1.52 \\ -3.28 & -0.74 & 11.20 & -3.05 \\ -1.10 & 1.52 & -3.05 & 1.87 \end{bmatrix}$$

in mm, mrad

$$\begin{aligned} E_1 &= 2.4 \text{ mm mrad} \\ E_2 &= 2.0 \text{ mm mrad} \\ E_{4d} &= 4.8 \text{ mm}^2 \text{ mrad}^2 \end{aligned}$$



4d Emittance Measurements



check: do results change if six measurements are used instead of four?

- red ellipses from four measurements:
 1. at 0° with setting $Q1^a, Q2^a$
 2. at 30° with setting $Q1^a, Q2^a$
 3. at 90° with setting $Q1^a, Q2^a$
 4. at 30° with setting $Q1^b, Q2^b$
- blue ellipses from six measurements:
 1. at 0° with setting $Q1^a, Q2^a$
 2. at 30° with setting $Q1^a, Q2^a$
 3. at 90° with setting $Q1^a, Q2^a$
 4. at 0° with setting $Q1^b, Q2^b$
 5. at 30° with setting $Q1^b, Q2^b$
 6. at 90° with setting $Q1^b, Q2^b$

four measurements are sufficient

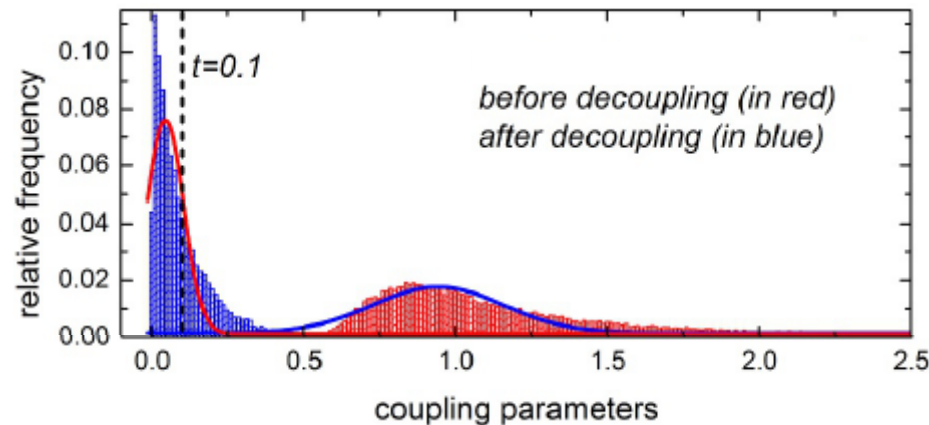
4d Emittance Measurements: Input for Decoupling



- decoupling reduces projected emittances, i.e. the usual hor. / ver. rms-emittances
- to decouple, the correlated moments must be known (EmTEx with space charge, ...)
- are the ROSE results sufficiently precise to determine the decoupling lattice?

4d Emittance Measurements: Input for Decoupling

- to answer that question, a virtual decoupling line was constructed (regular triplet and skew triplet for instance)
- the decoupling gradients are calculated from measured moments w/o errors
- moments with errors (from random error study) are transported through decoupling line
- the spectrum of final coupling parameters $t = \frac{\epsilon_x \epsilon_y}{\epsilon_1 \epsilon_2} - 1$ is evaluated:



**ROSE results are sufficiently precise
to perform decoupling**