

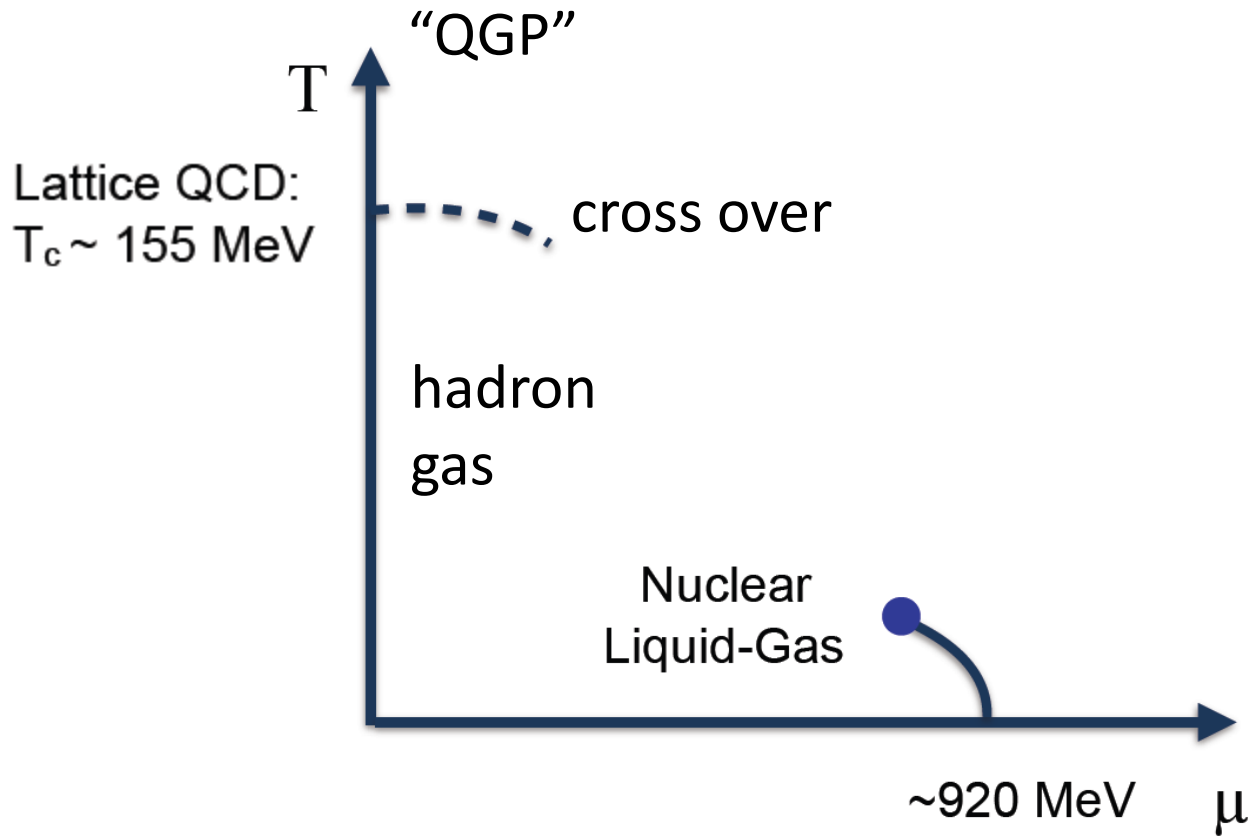
# Physics at high baryon density

Adam Bzdak

AGH University of Science and Technology, Kraków

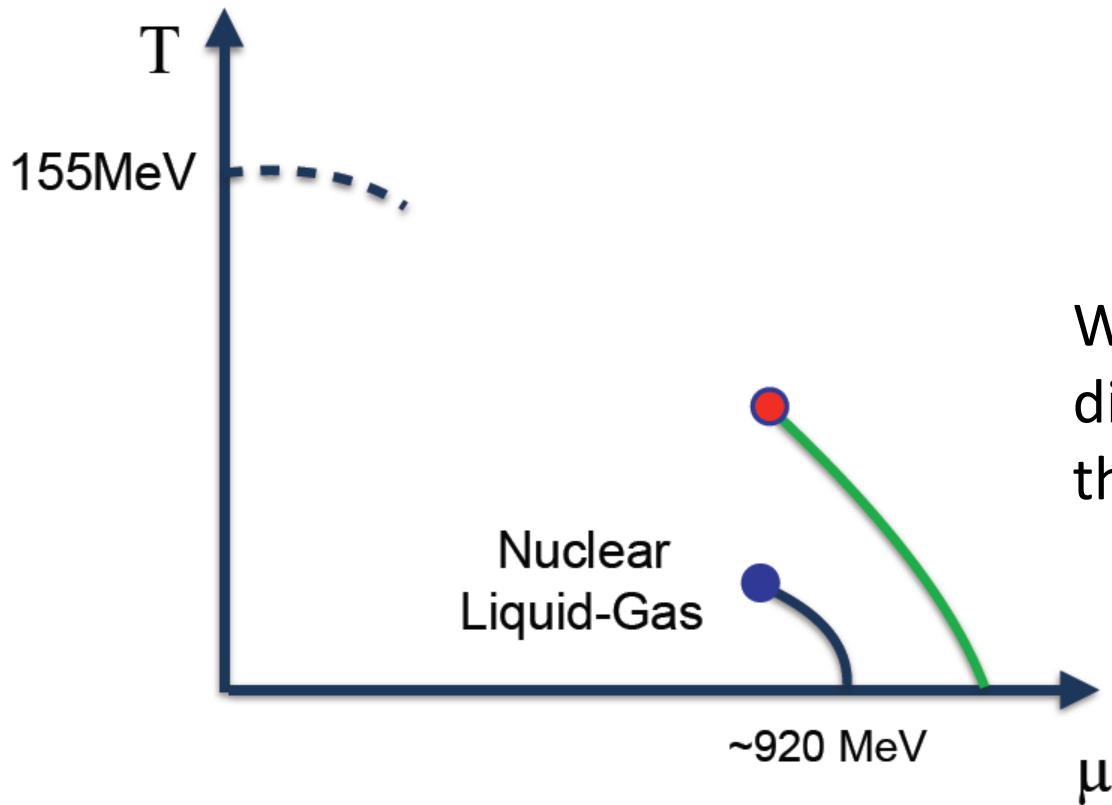


# What we really know about the QCD phase diagram



The rest is everybody's guess.

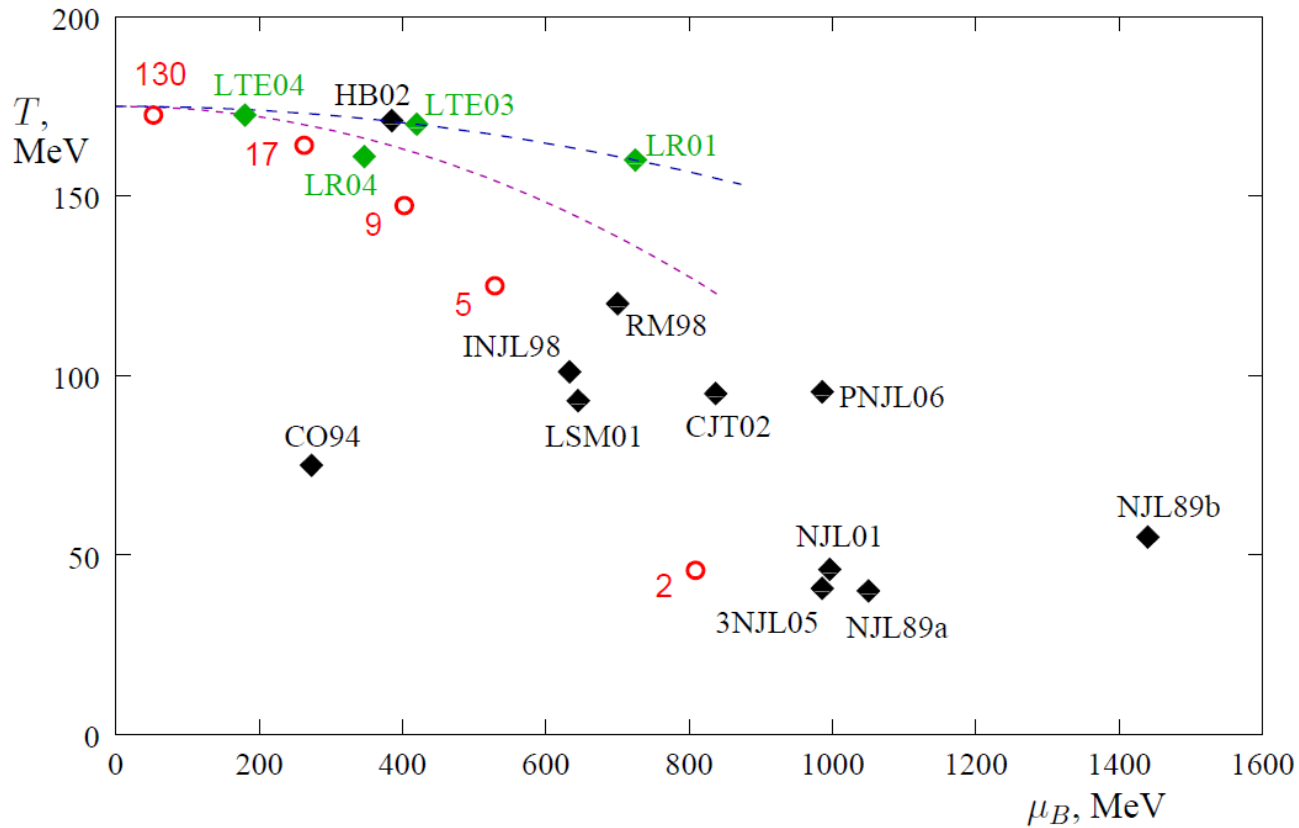
Usual expectation based on various effective models



What **exactly** distinguishes the two phases?

Are these models good?

Usually (hopefully) Nature is more interesting than our guesses



**Figure 4:** Comparison of predictions for the location of the QCD critical point on the phase diagram. Black points are model predictions: NJLa89, NJLb89 – [12], CO94 – [13, 14], INJL98 – [15], RM98 – [16], LSM01, NJL01 – [17], HB02 – [18], CJT02 – [19], 3NJL05 – [20], PNJL06 – [21]. Green points are lattice predictions: LR01, LR04 – [22], LTE03 – [23], LTE04 – [24]. The two dashed lines are parabolas with slopes corresponding to lattice predictions of the slope  $dT/d\mu_B^2$  of the transition line at  $\mu_B = 0$  [23, 25]. The red circles are locations of the freezeout points for heavy ion collisions at corresponding center of mass energies per nucleon (indicated by labels in GeV) – Section 5.

It is hard to expect any real theoretical progress (e.g., LQCD) in the near(est) future.

On the experimental side all we can do is to measure various fluctuation observables and hope to see some nontrivial energy or/and system-size dependence

see, e.g.,

Stephanov, Rajagopal, Shuryak, PRL (1998)

Stephanov, PRL (2009)

Skokov, Friman, Redlich, PRC (2011)

There are some intriguing results:

STAR, HADES

NA 49

Higher order  
net-proton cumulants

Intermittency in the transverse  
momentum phase space

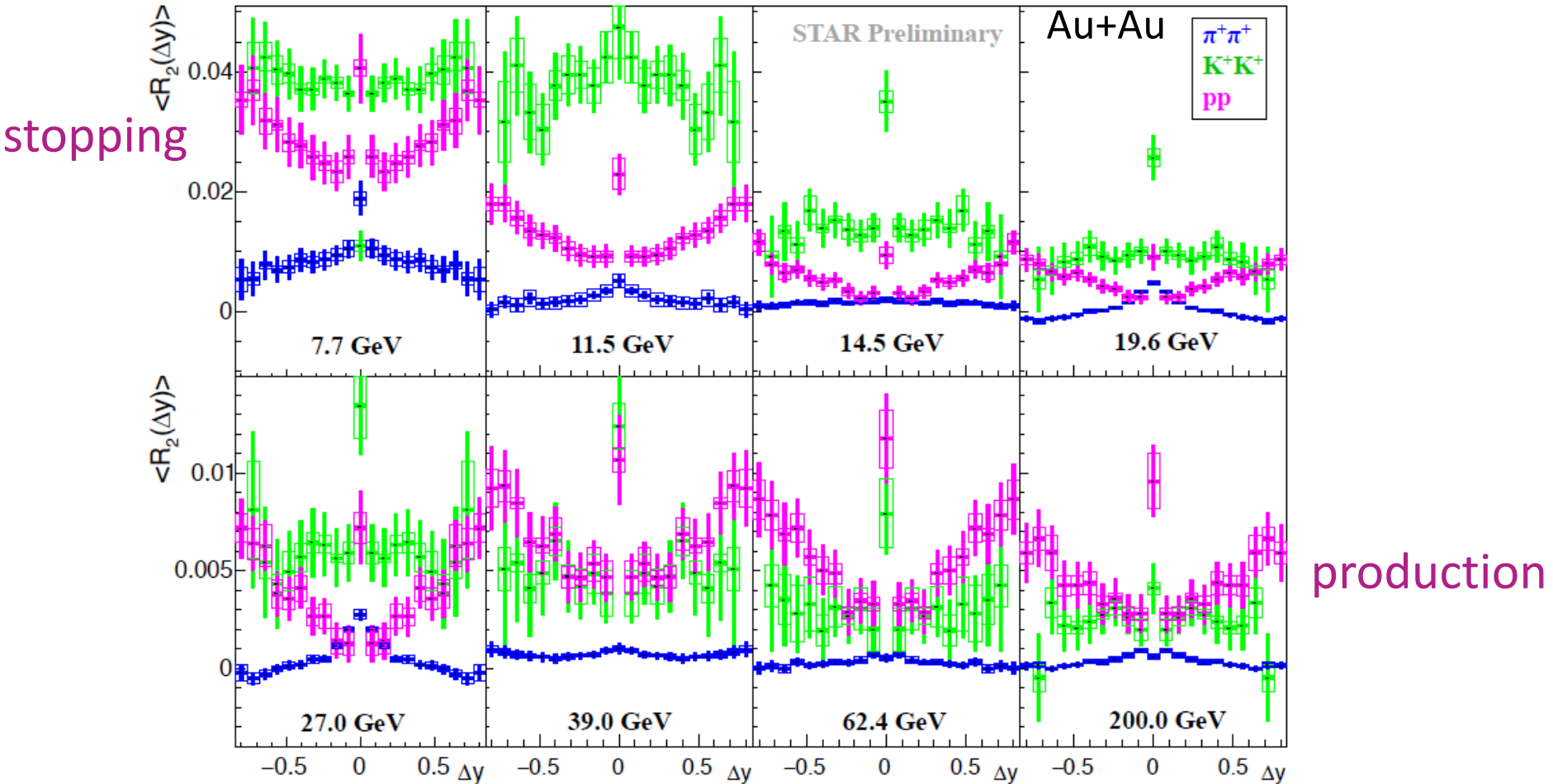
Proton  $v_1$  (STAR)

Strongly intensive variables

See, e.g., Eur. Phys. J. C75 (2015) 12, 587  
(1208.5292)

We do not understand basic baryon physics.

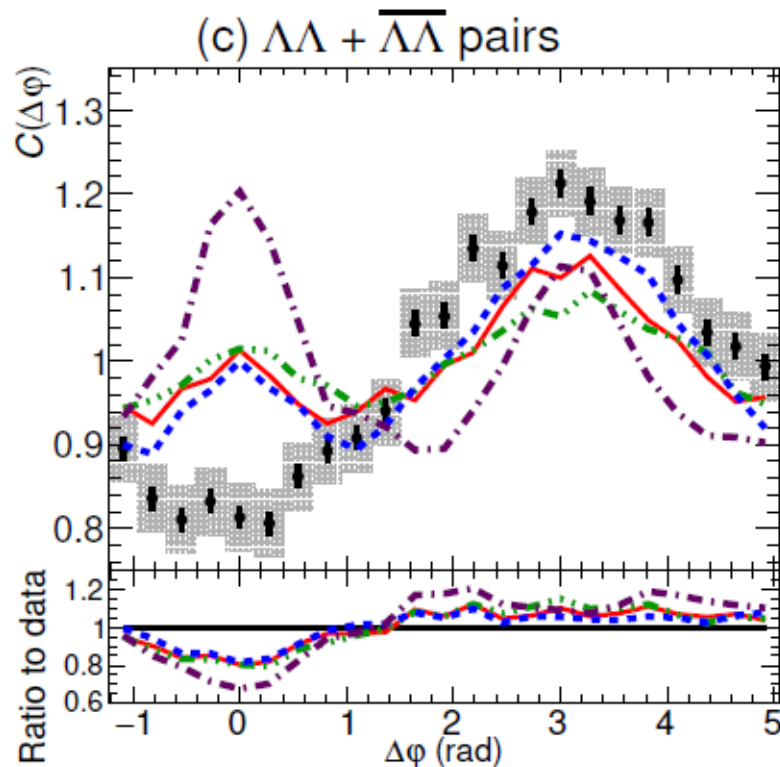
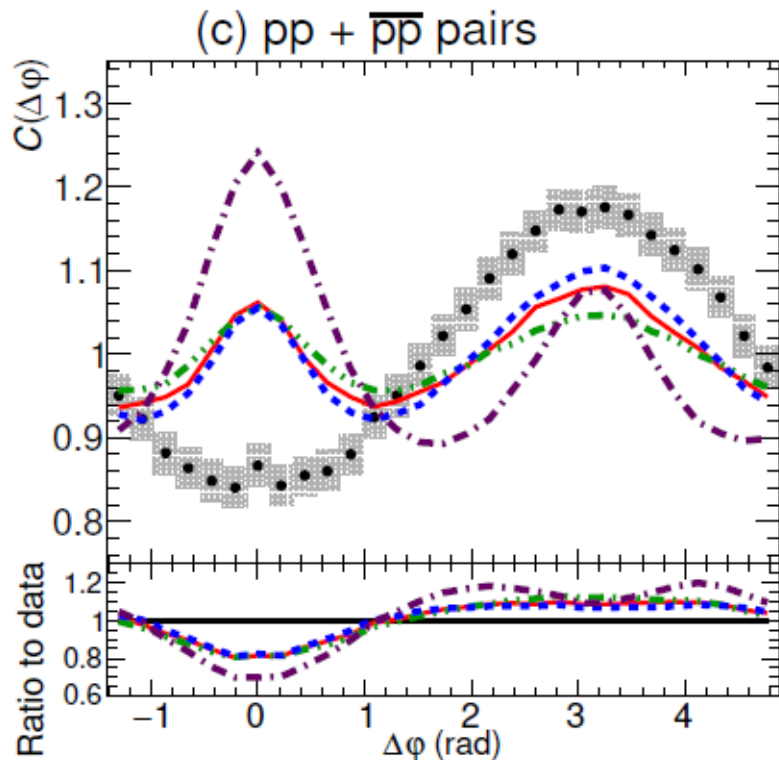
S. Jowzaee,  
QM2017



$$R_2(y_1, y_2) = -1 + \frac{\langle \rho_2(y_1, y_2) \rangle}{\langle \rho_1(y_1) \rangle \langle \rho_1(y_2) \rangle}$$

← Same event pair distributions

← Mixed event



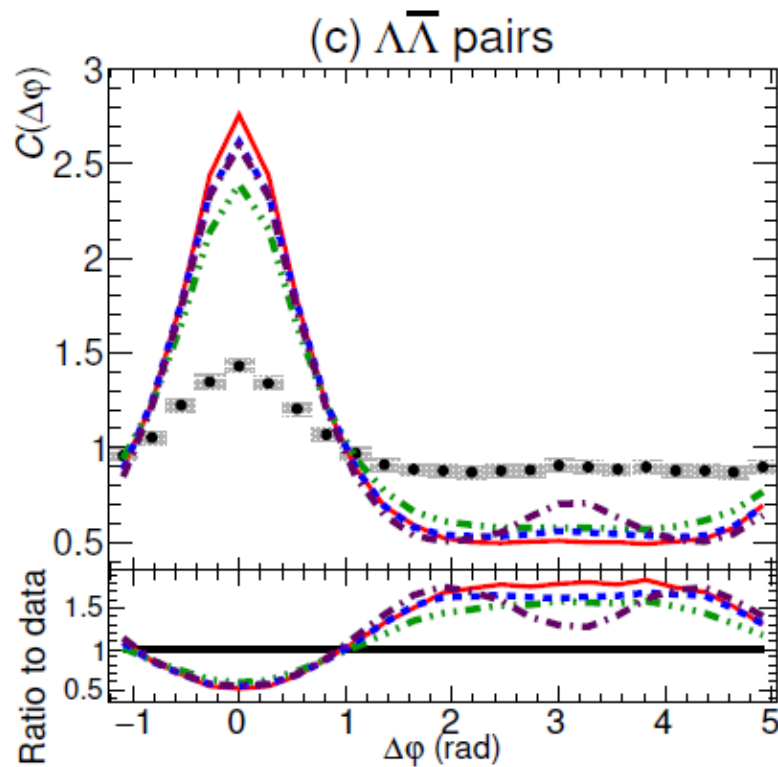
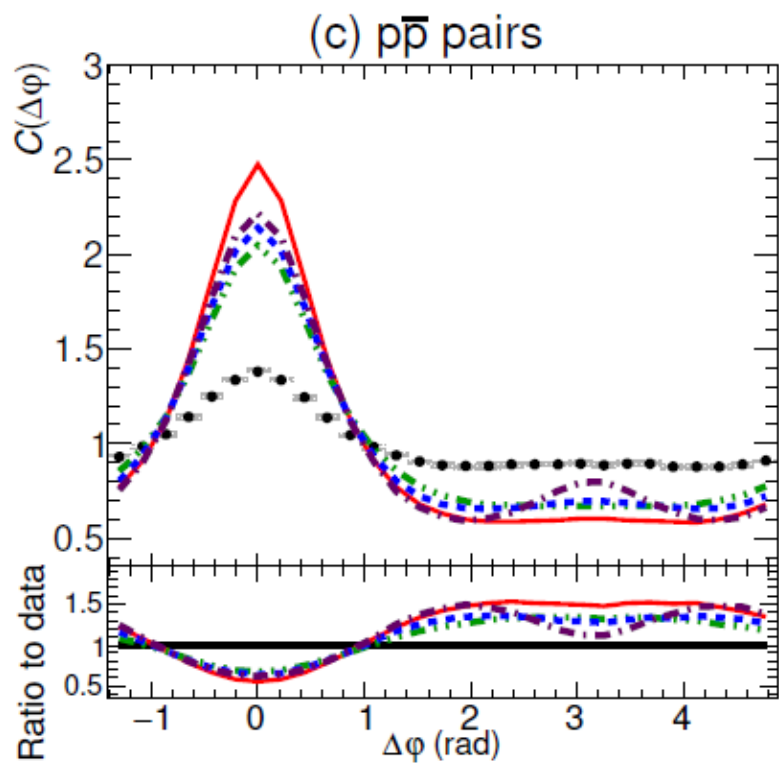
Baryons do not want to be close to each other in rapidity and azimuthal angle

And we want to create high baryon density

First seen by TPC/Two Gamma Collaboration in  $e+e-$  annihilation at 29-GeV, PRL 57, 3140 (1986).

- ↓ ALICE  $pp \sqrt{s} = 7 \text{ TeV}$ ,  $|\Delta\eta| < 1.3$
- PYTHIA6 Perugia-0
- PYTHIA6 Perugia-2011
- PYTHIA8 Monash
- PHOJET

## No surprise for baryon-antibaryon correlation





Proton stopping – is it obvious that we produce high baryon density?

At low energy protons are not produced. They are transferred from incoming nucleus.

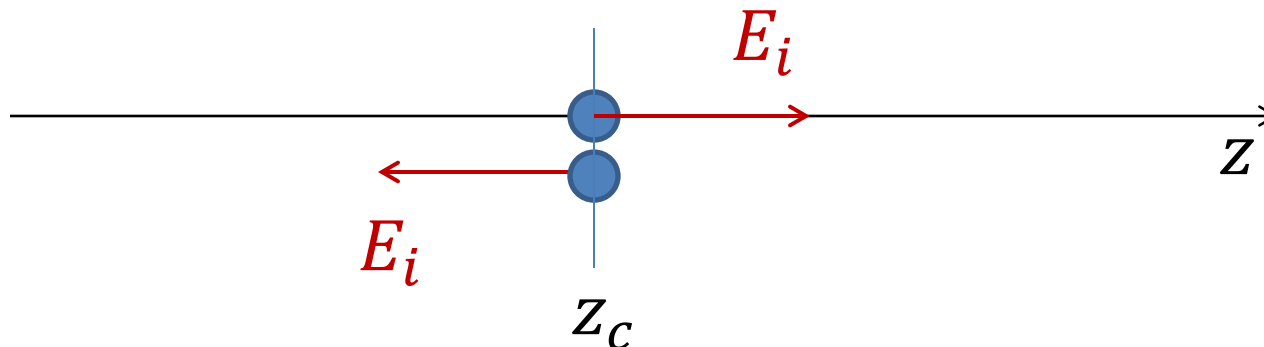
There is no infinite deceleration. It takes some time and length to slow down or stop a proton. Thus we expect (?) the stopped protons to be away from  $z = 0$ .

$$E_z = E_i - \sigma(z - z_c)$$

$E_i$  – initial energy;  $z_c$  – collision point;  $E_z$  – energy at a point  $z$

$\sigma$  – energy loss per unit length

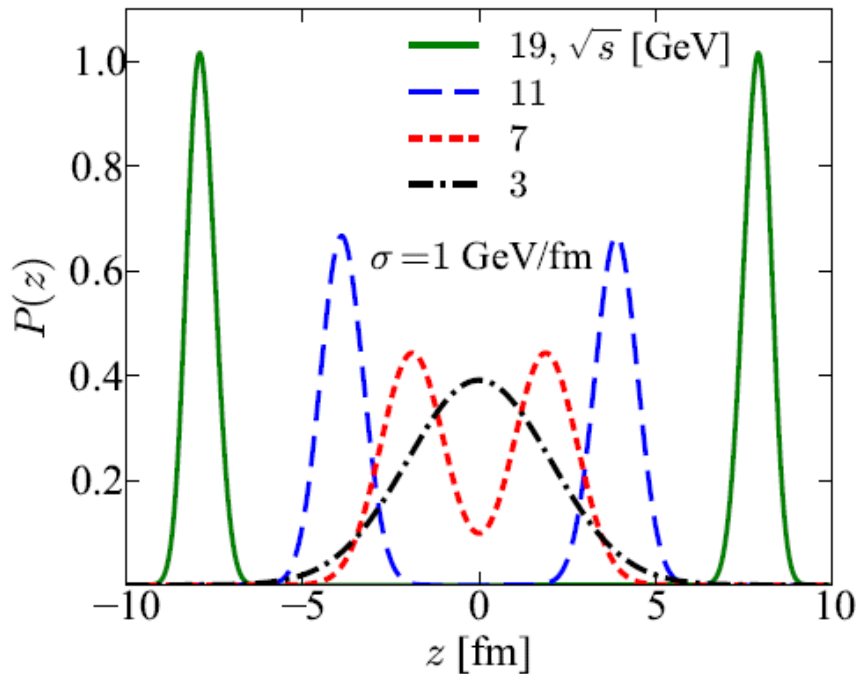
$$E_z \rightarrow M_t \cosh(y)$$



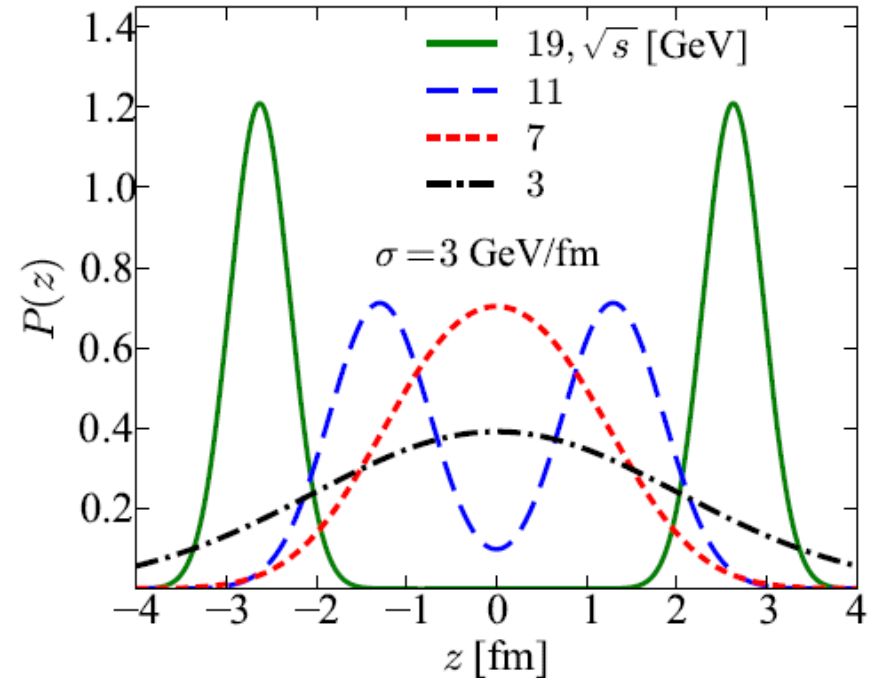
Now we need to average over the collision points  $z_c$  in A+A, and over the measured rapidity bin.

Here  $|y| < 1$ ,  $R_A = 6.5$  fm,  $P_t = 1$  GeV.

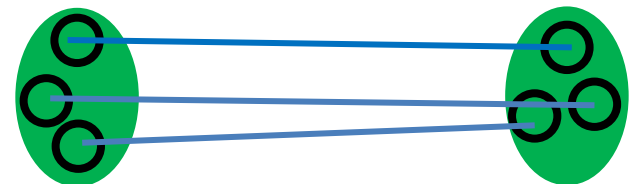
$$M_t^2 = M^2 + P_t^2$$



wounded nucleon model



wounded quark model



One would prefer to see a more uniform distribution in  $z$ .

We do not take into account resonances.

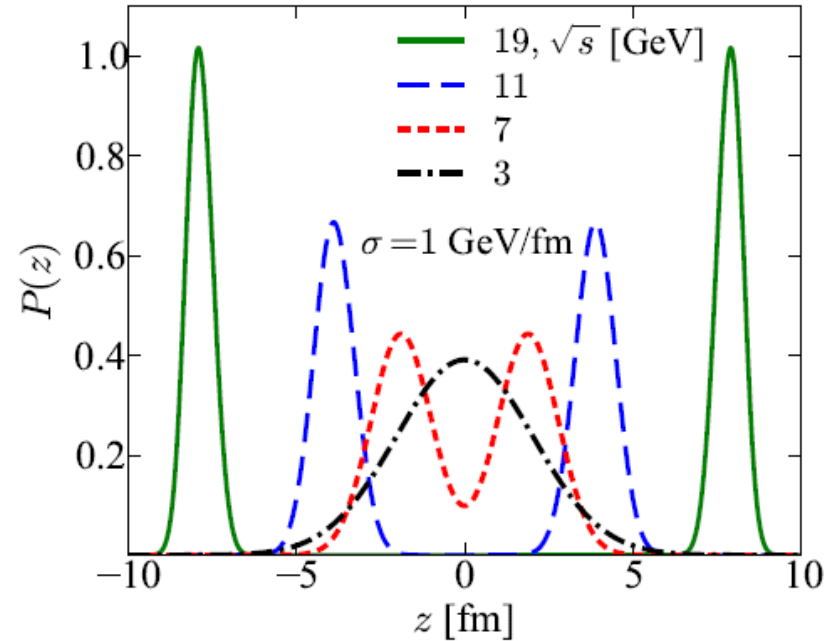
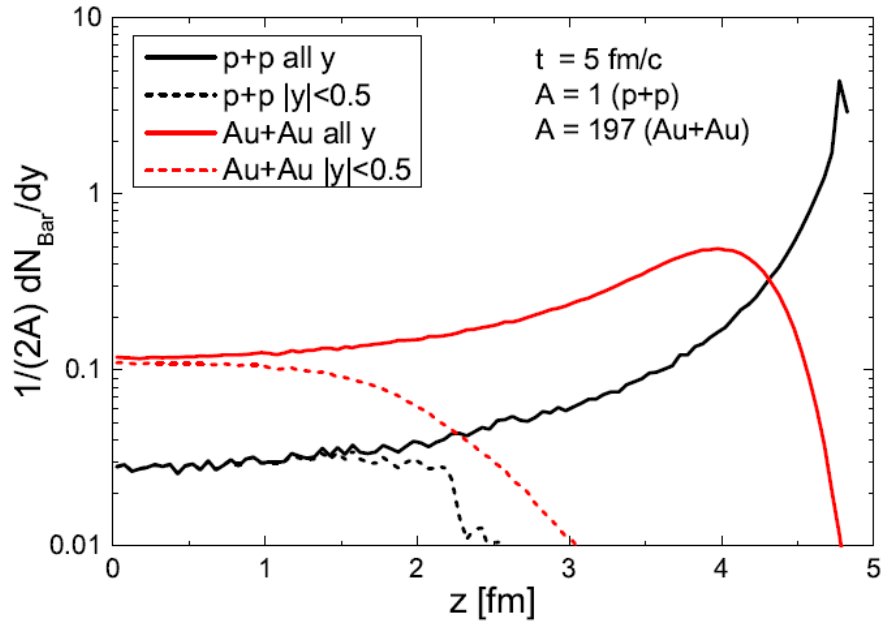
Real MC calculations are warranted but:

Y.Nara, H.Niemi, J.Steinheimer, H.Stoecker, 1611.08023

Also it has been pointed out that the usual description of baryon stopping in string models in a transport model assumes essentially an instant deceleration of the leading baryons because strings are immediately decay into hadrons with a formation time. In [59] it was argued that a constant deceleration would lead to a very different initial density profile in coordinate space, i.e. a smaller density at mid-rapidity, which may transform into a different time evolution of the particle flow.

# UrQMD for example

J. Steinheimer, INT-16-3

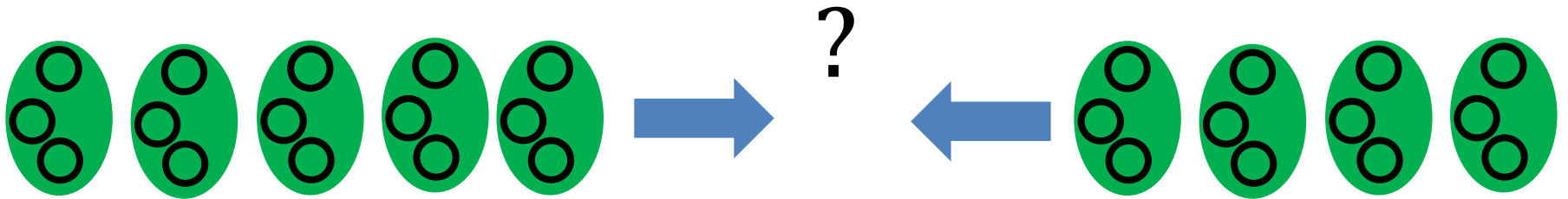


hopefully this is not right

Anyway, the question remains: what is the realistic  $z$  distribution of stopped protons? We need many baryons with similar  $y$  and  $z$

How to properly model baryon stopping?

Do baryons stop independently or maybe we have some multi-baryon correlations?

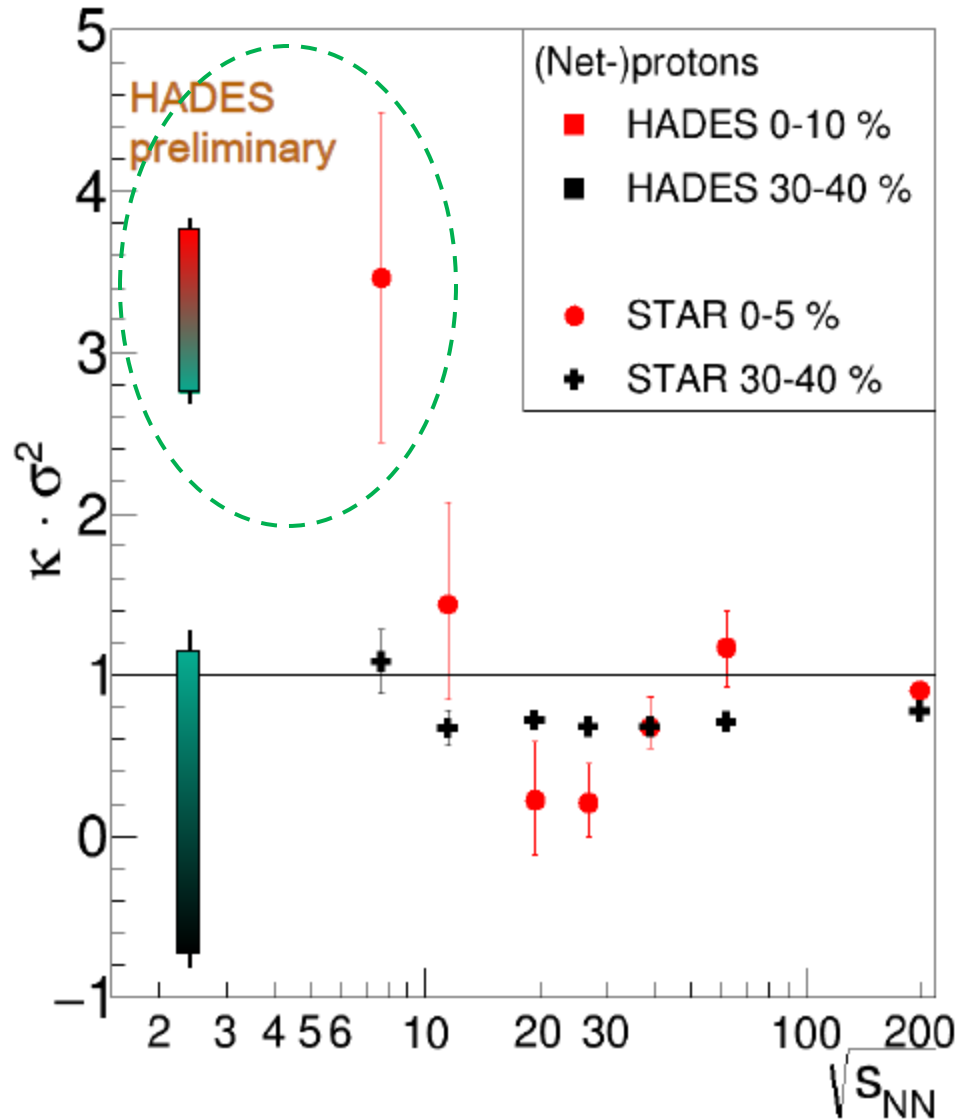


Important questions if we want to understand proton cumulants as measured by STAR and HADES.

STAR and possibly HADES see rather exciting multi-proton correlations at lower energies.

my notation

$$K_4/K_2$$



The signal at low (< 8 GeV) energy is surprisingly large

## Cumulants are not the best choice

$$K_2 = \langle (N - \langle N \rangle)^2 \rangle$$

$N$  – number of protons

$$K_3 = \langle (N - \langle N \rangle)^3 \rangle$$

we neglect anti-protons,  
good at low energies

$$K_4 = \langle (N - \langle N \rangle)^4 \rangle - 3\langle (N - \langle N \rangle)^2 \rangle^2$$

$$K_n = \langle N \rangle + \textit{physics}[2, \dots, n]$$

*physics* = two-, three-,  $n$ -particle  
correlation functions

for Poisson  $K_n = \langle N \rangle$ , (*physics* = 0)

We have

$$K_2 = \langle N \rangle + \mathbf{C}_2$$

$$K_3 = \langle N \rangle + 3\mathbf{C}_2 + \mathbf{C}_3$$

$$K_4 = \langle N \rangle + 7\mathbf{C}_2 + 6\mathbf{C}_3 + \mathbf{C}_4$$

cumulants mix  
correlation functions  
of different orders

For example:

$$\rho_2(y_1, y_2) = \rho(y_1)\rho(y_2) + \mathbf{C}_2(y_1, y_2)$$

$$\mathbf{C}_2 = \int \mathbf{C}_2(y_1, y_2) dy_1 dy_2$$

$$\mathbf{C}_2 = \langle N \rangle^2 \mathbf{c}_2$$



See, e.g.,

B. Ling, M. Stephanov, PRC 93 (2016) no.3, 034915

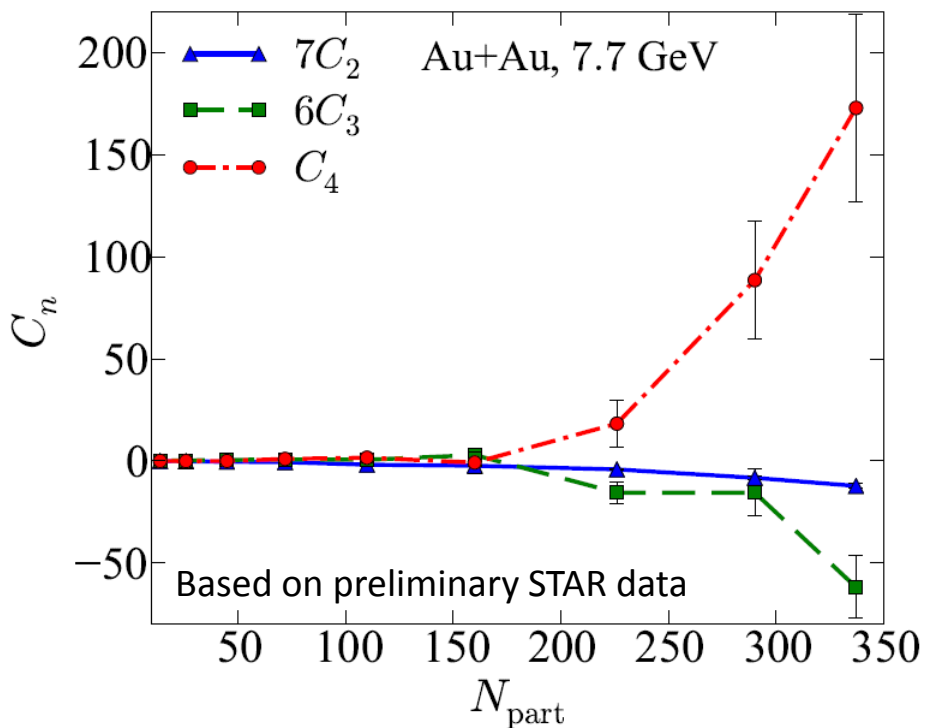
AB, V. Koch, N. Strodthoff, 1607.07375

we should  
measure this



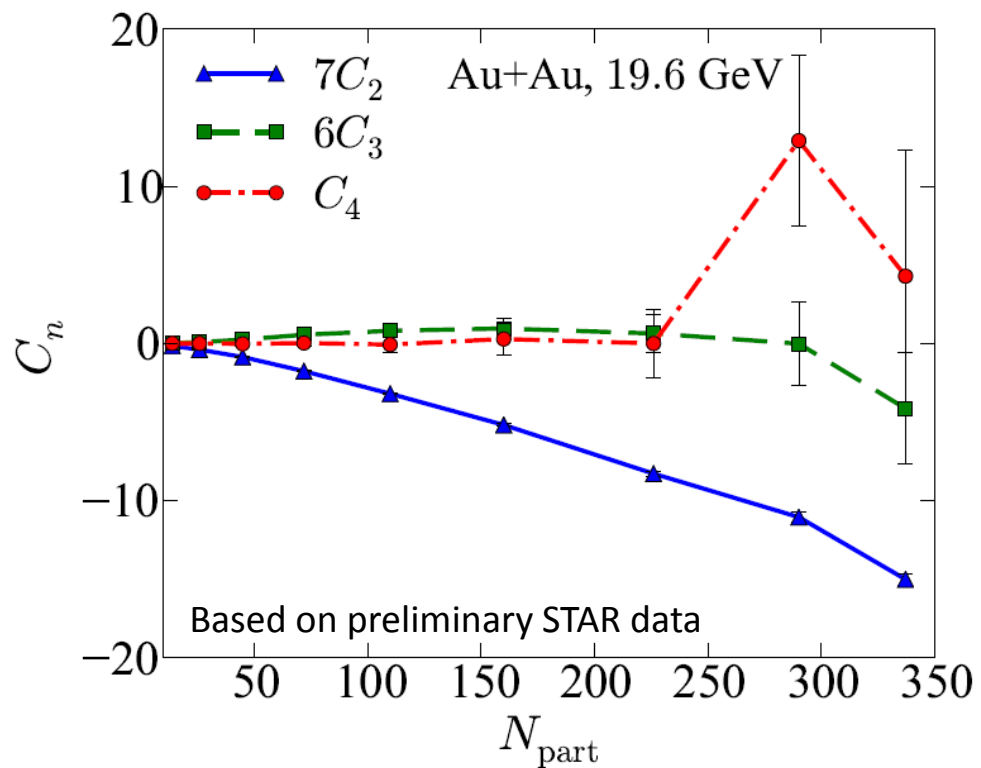
# Using preliminary STAR data we obtain $C_n$

central signal at 7.7 GeV is driven by large 4-particle correlations



$$C_4(7.7) \sim 170$$

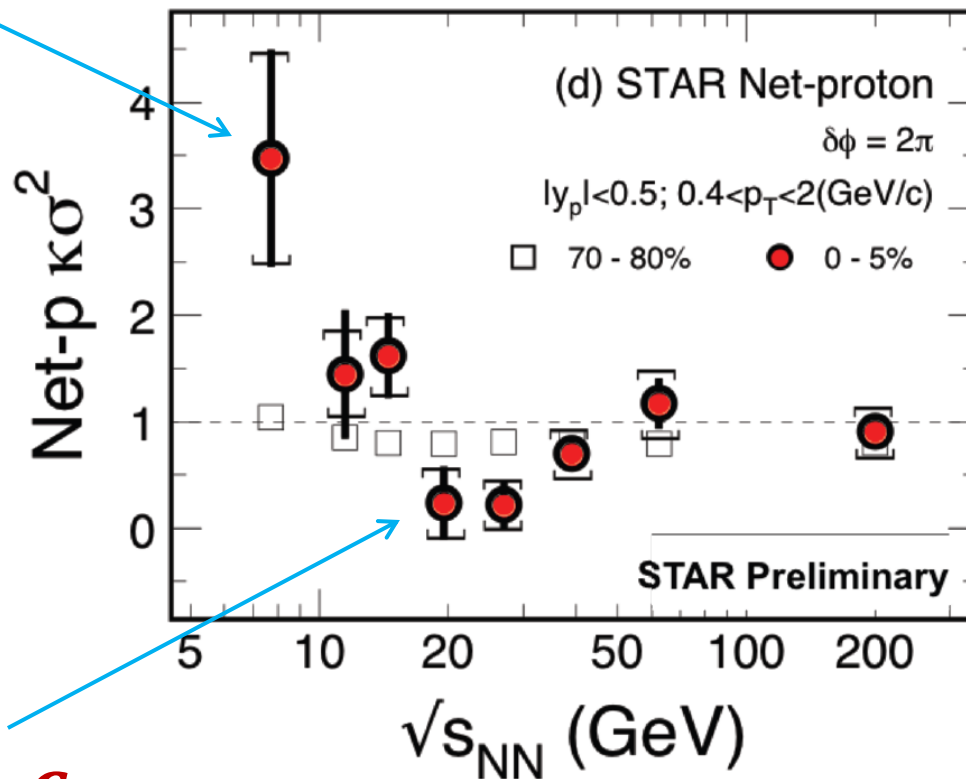
central signal at 19.6 GeV is driven by 2-particle correlations



$C_4$  and  $6C_3$  cancelation in most central coll.

here we see  $C_4$

X.Luo, N.Xu, 1701.02105



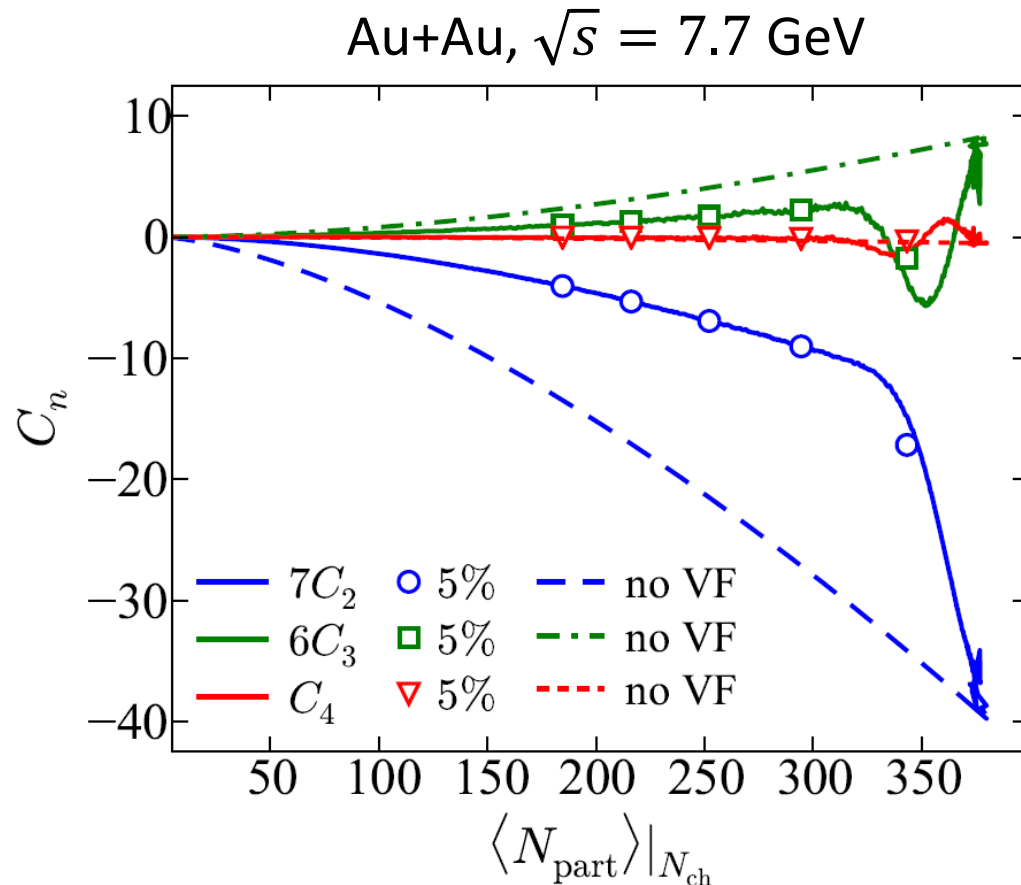
and here  $C_2$

e.g., baryon conservation

what about HADES ?

## Minimal model (MM) at low energies

- independent baryon stopping (baryon conservation by construction)
- $N_{\text{part}}$  fluctuations (volume fluctuation - VF)



STAR

$$C_4 \sim 170$$

$$6C_3 \sim -60$$

$$7C_2 \sim -15$$

we follow the STAR way (centrality etc.) as closely as possible

AB, V. Koch, V. Skokov, 1612.05128

See also: P. Braun-Munzinger, A. Rustamov, J. Stachel, NPA 960 (2017) 114

Let's put the STAR numbers in perspective.

Suppose that we have clusters (distributed according to Poisson) decaying always to 4 protons

$$C_k = \langle N_{cl} \rangle \cdot 4! / (4 - k)!$$



mean number  
of clusters

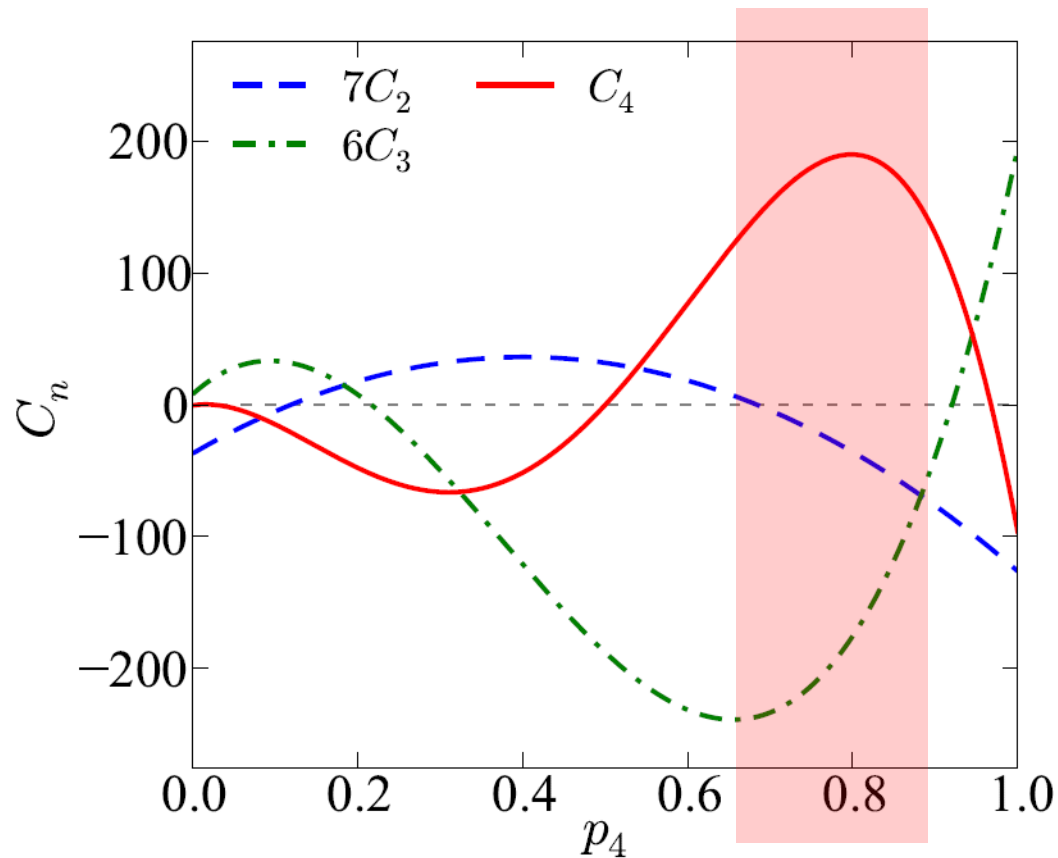
$$C_4 = \langle N_{cl} \rangle \cdot 24$$

To obtain  $C_4 \approx 170$  we need  $\langle N_{cl} \rangle \sim 7$ , it means 28 protons.  
STAR sees on average 40 protons in central collisions.

In this model  $C_2 > 0$  and  $C_3 > 0$  contrary to the STAR data

## Toy model:

- 16 protons stop in quartets with probability  $p_4$
- remaining protons stop independently with some small probability  $p_1 \sim 0.1$



qualitatively  
consistent  
with STAR

STAR

$$C_4 \sim 170$$

$$6C_3 \sim -60$$

$$7C_2 \sim -15$$

Take-home message:

$C_4$  (four-proton correlation function) observed by STAR (and likely by HADES) is larger by almost **three orders of magnitude** than the minimal model.

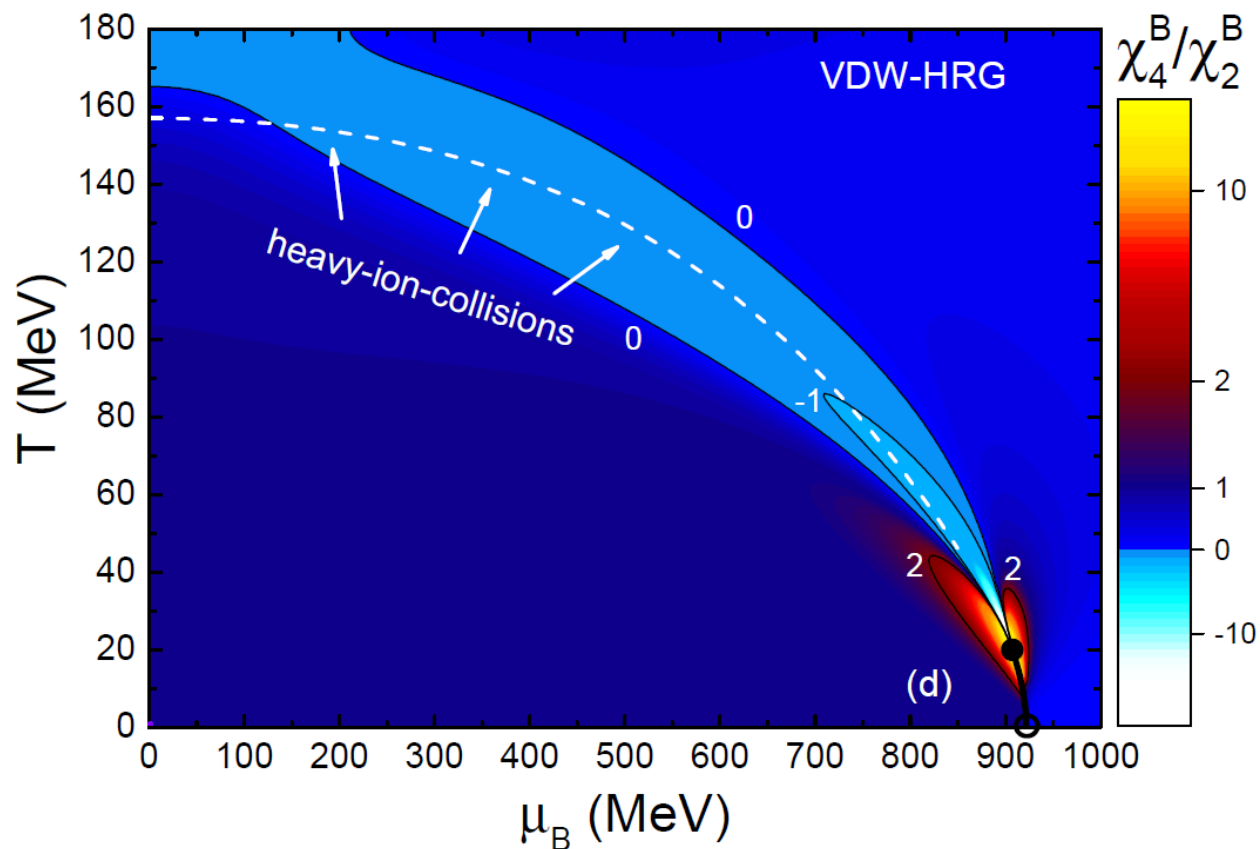
To explain  $C_4$  we need a strong source of multi-proton correlations.  
**Proton clusters?** Phase transition?

Similar story  
with 3-proton  
correlation  
function.



One needs to be careful with interpretations, for example the liquid-gas phase transition shines across a broad range of  $T$  and  $\mu$

HRG with attractive and repulsive Van der Waals interactions between (anti)baryons



## Conclusions

We do not know much about the QCD phase diagram.  
Real theoretical progress is not expected soon.

We need to make sure we produce high baryon density.  
The configuration space distribution of the final baryons (and other particles) requires careful study.

Multi-proton correlations from stopping, critical point, phase transition?

Very large 4-proton correlations at 7.7 GeV in central Au+Au collisions.

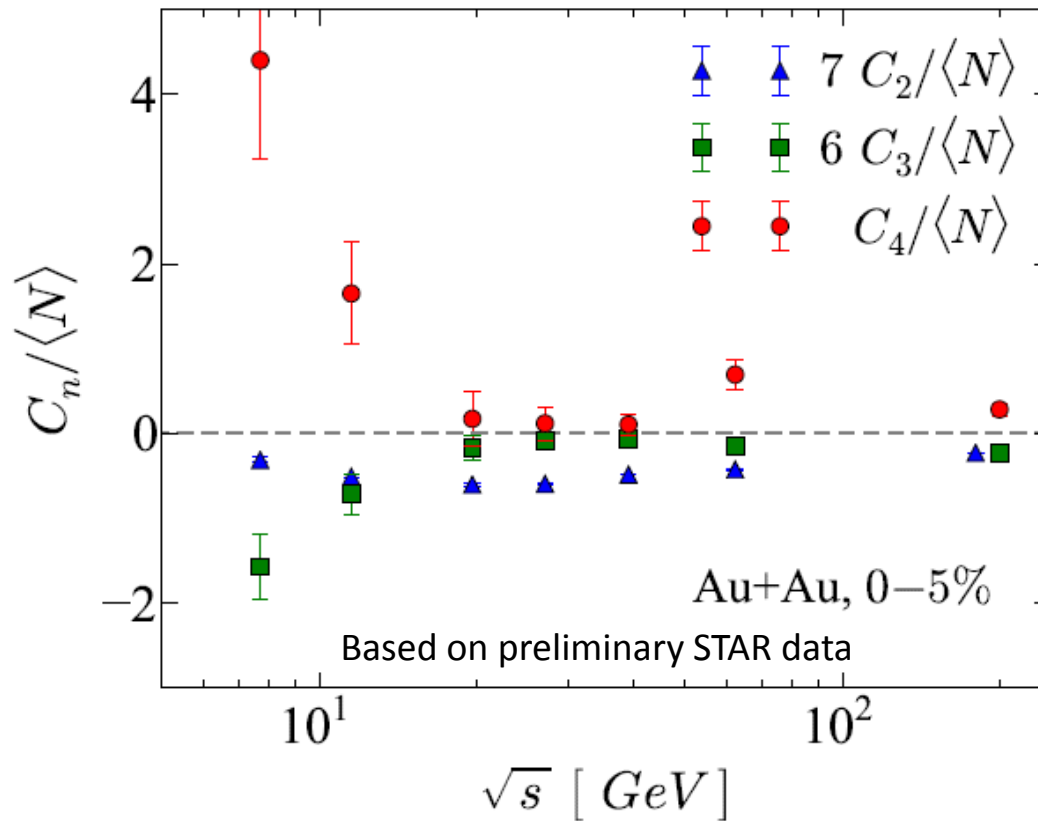
See backup: QCD exclusion plots, interesting centrality and rapidity dependence of the proton correlation functions



# Backup

## Exclusions plots

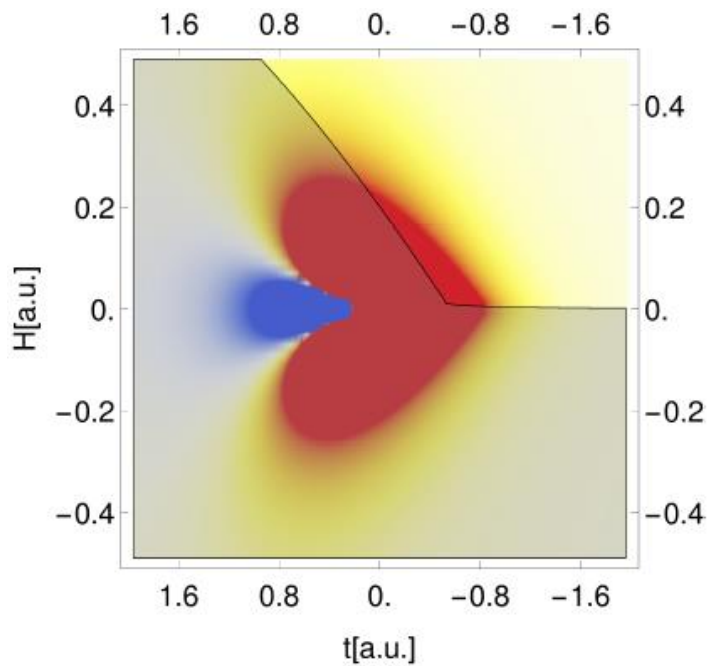
We propose to make the phase-diagram exclusion plots based on the **signs** of the correlation functions.



$C_4$  at 62 GeV !

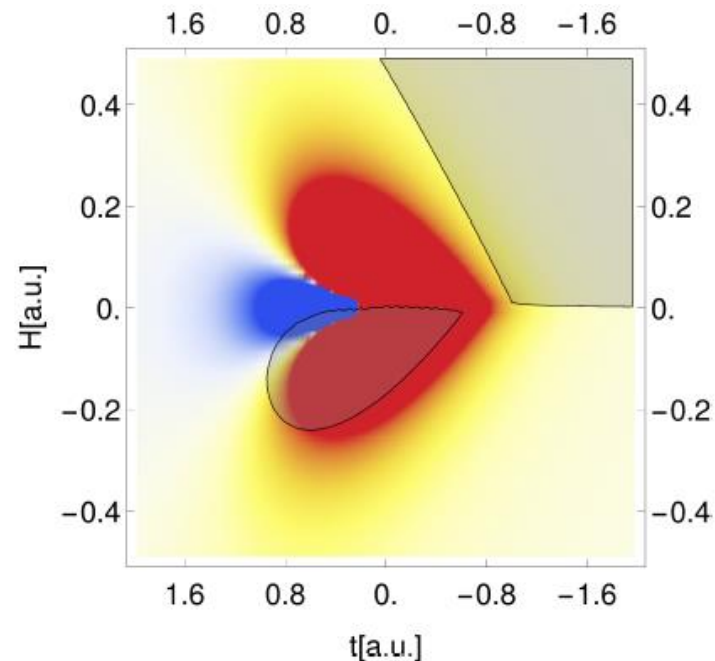
$$C_4 > 0, C_3 < 0 \quad C_2 < 0$$

# Exclusions plots based on the Ising model

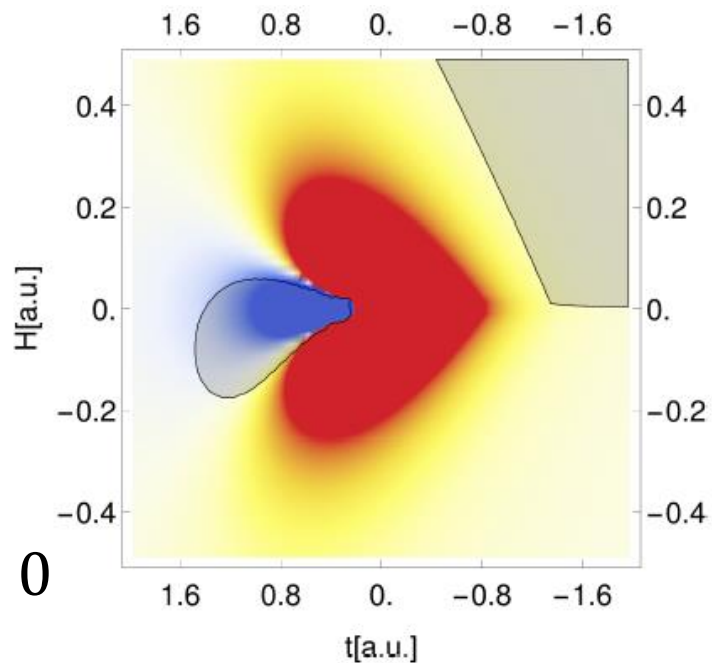


$$C_2 < 0$$

Clearly we need to use more realistic model with various effects included



$$C_3 < 0$$



$$C_4 > 0$$

## Observations

$$c_2 = \frac{\int \rho(y_1)\rho(y_2)c_2(y_1, y_2)dy_1dy_2}{\int \rho(y_1)\rho(y_2)dy_1dy_2}$$

$$K_2 = \langle N \rangle + \langle N \rangle^2 c_2$$

$$K_4 = \langle N \rangle + 7\langle N \rangle^2 c_2 + 6\langle N \rangle^3 c_3 + \langle N \rangle^4 c_4$$

Rapidity dependence:

**long-range** correlation

$$c_n(y_1, \dots, y_n) = c_n^0$$

$$c_n = c_n^0$$

$$K_2 = \langle N \rangle + c_2^0 \langle N \rangle^2, \quad \langle N \rangle \sim \Delta y$$

$$K_4 = \langle N \rangle + 7c_2^0 \langle N \rangle^2 + 6c_3^0 \langle N \rangle^3 + c_4^0 \langle N \rangle^4$$

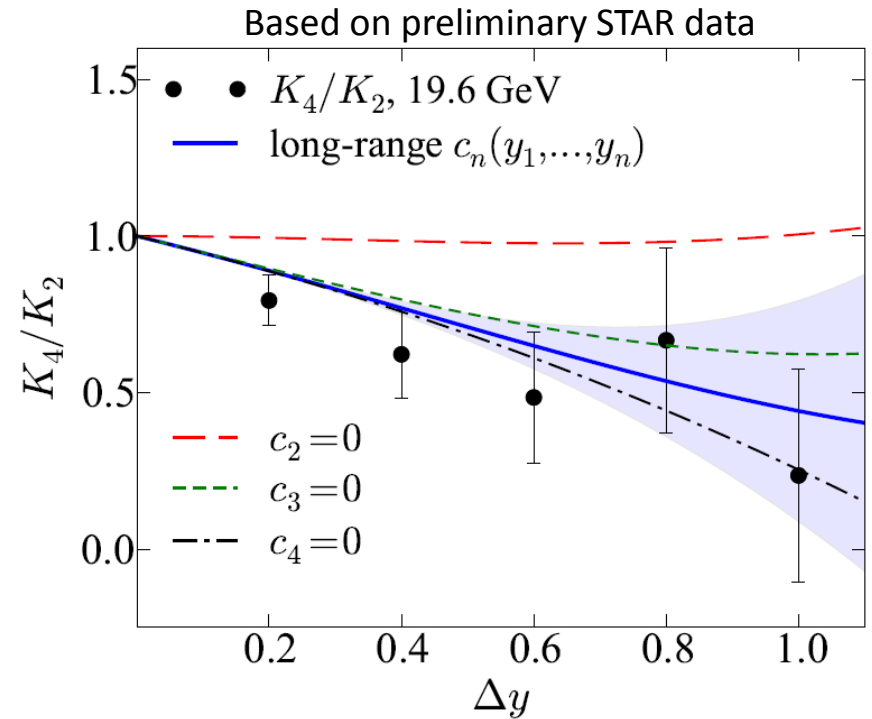
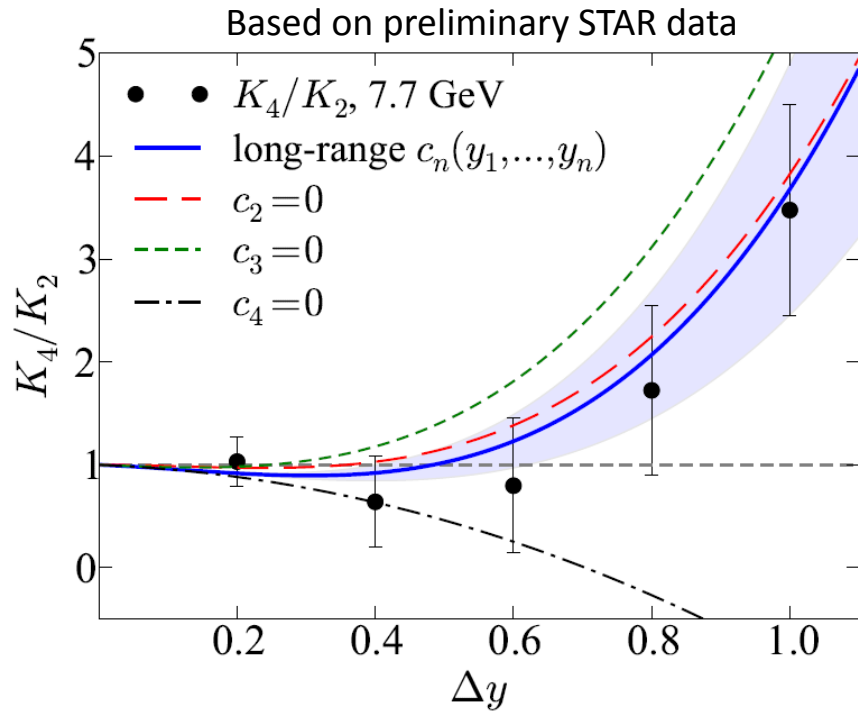
**short-range** correlation

$$c_2(y_1, y_2) = c_2^0 \delta(y_1 - y_2)$$

$$c_2 \sim 1/(\Delta y)$$

$$K_n \sim \Delta y$$

# Rapidity dependence consistent with long-range correlations



$|y| < 0.5$  is not particularly large

$$\rho_2(y_1, y_2) = \rho(y_1)\rho(y_2) + \mathbf{C}_2(y_1, y_2) \quad \text{correlation function}$$

$$\rho_2(y_1, y_2) = \rho(y_1)\rho(y_2)[1 + \mathbf{c}_2(y_1, y_2)] \quad \text{reduced correlation function}$$

$$\langle N(N - 1) \rangle = \langle N \rangle^2 + \langle N \rangle^2 \mathbf{c}_2$$

$$\mathbf{c}_2 = \frac{\int \rho(y_1)\rho(y_2)\mathbf{c}_2(y_1, y_2)dy_1dy_2}{\int \rho(y_1)\rho(y_2)dy_1dy_2}$$

and the second order cumulant

$$K_2 = \langle N \rangle + \underbrace{\langle N \rangle^2 \mathbf{c}_2}_{\mathbf{C}_2}$$

In the same way

$$\rho_3(y_1, y_2, y_3) = \rho(y_1)\rho(y_2)\rho(y_3)[1 + \mathbf{c}_2(y_1, y_2) + \cdots + \mathbf{c}_3(y_1, y_2, y_3)]$$

$$F_3 = \langle N(N-1)(N-2) \rangle = \langle N \rangle^3 + 3\langle N \rangle^2 \mathbf{c}_2 + \langle N \rangle^3 \mathbf{c}_3$$

$$\mathbf{c}_3 = \frac{\int \rho(y_1)\rho(y_2)\rho(y_3)\mathbf{c}_3(y_1, y_2, y_3)dy_1dy_2dy_3}{\int \rho(y_1)\rho(y_2)\rho(y_3)dy_1dy_2dy_3}$$

and the third order cumulant

$$K_3 = \langle N \rangle + 3\langle N \rangle^2 \mathbf{c}_2 + \langle N \rangle^3 \mathbf{c}_3$$

$3\mathbf{C}_2 \qquad \mathbf{C}_3$

Centrality dependence

$$C_2(y_1, y_2) = \rho(y_1)\rho(y_2)c_2(y_1, y_2)$$

$$K_4 = \langle N \rangle + 7C_2 + 6C_3 + C_4$$

$$C_k = \langle N \rangle^k c_k$$

$$K_4 = \langle N \rangle + 7\langle N \rangle^2 c_2 + 6\langle N \rangle^3 c_3 + \langle N \rangle^4 c_4$$

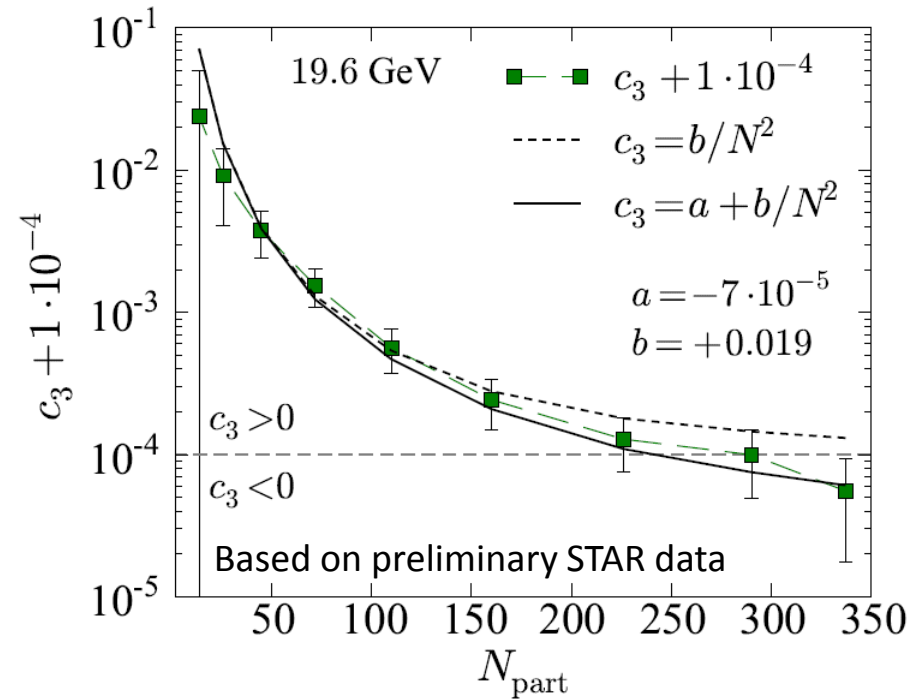
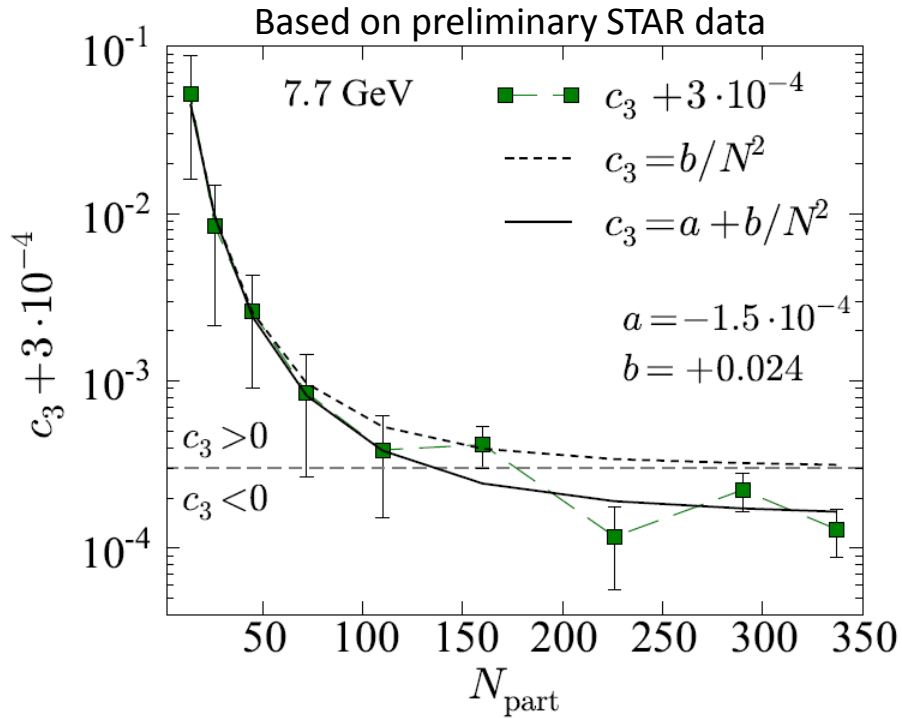
Suppose we have  $N_s$  **independent sources** of correlations  
(resonances, superposition of p+p etc.)

$$c_k \sim \frac{N_s}{N^k} \sim \frac{1}{N^{k-1}}$$



Using preliminary STAR data we obtain  $c_3$

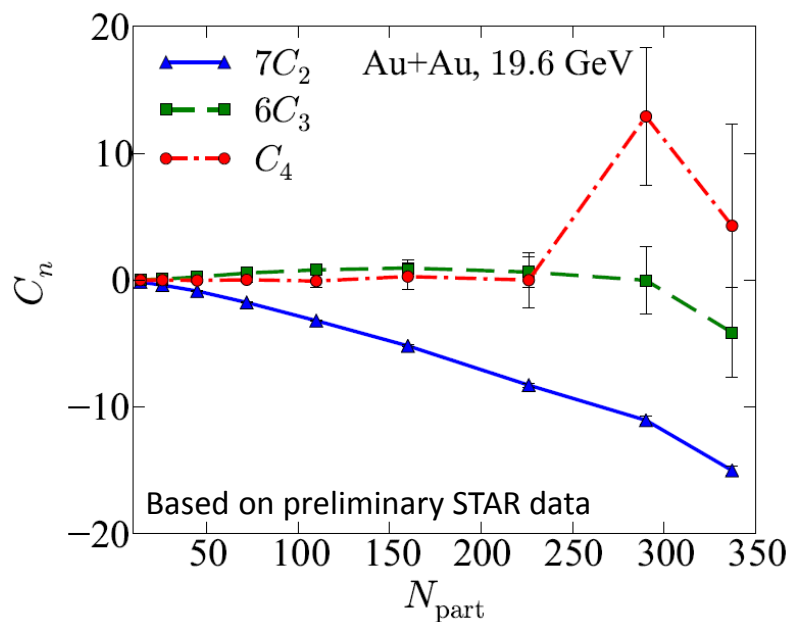
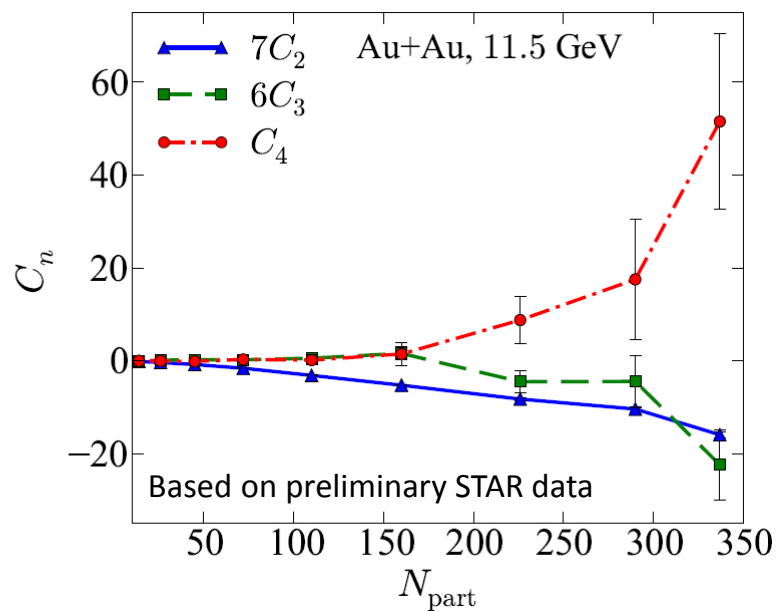
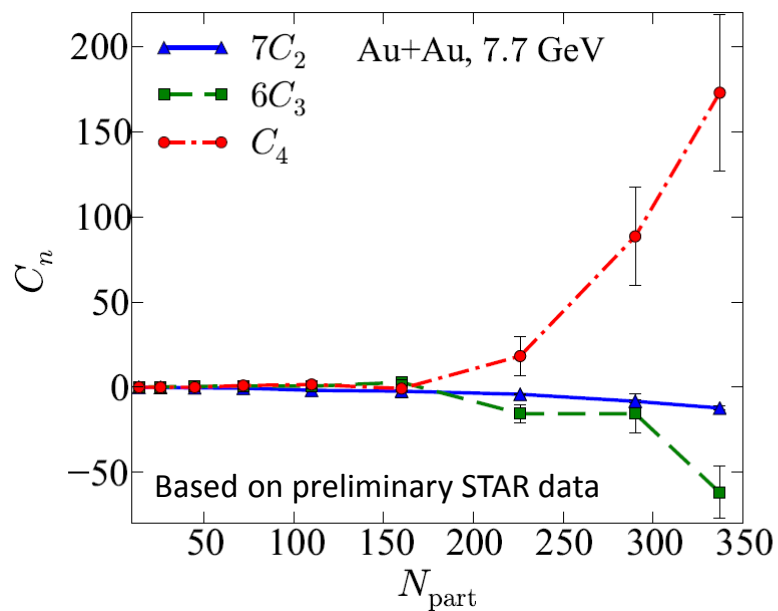
AB, V. Koch, N. Strodthoff,  
1607.07375



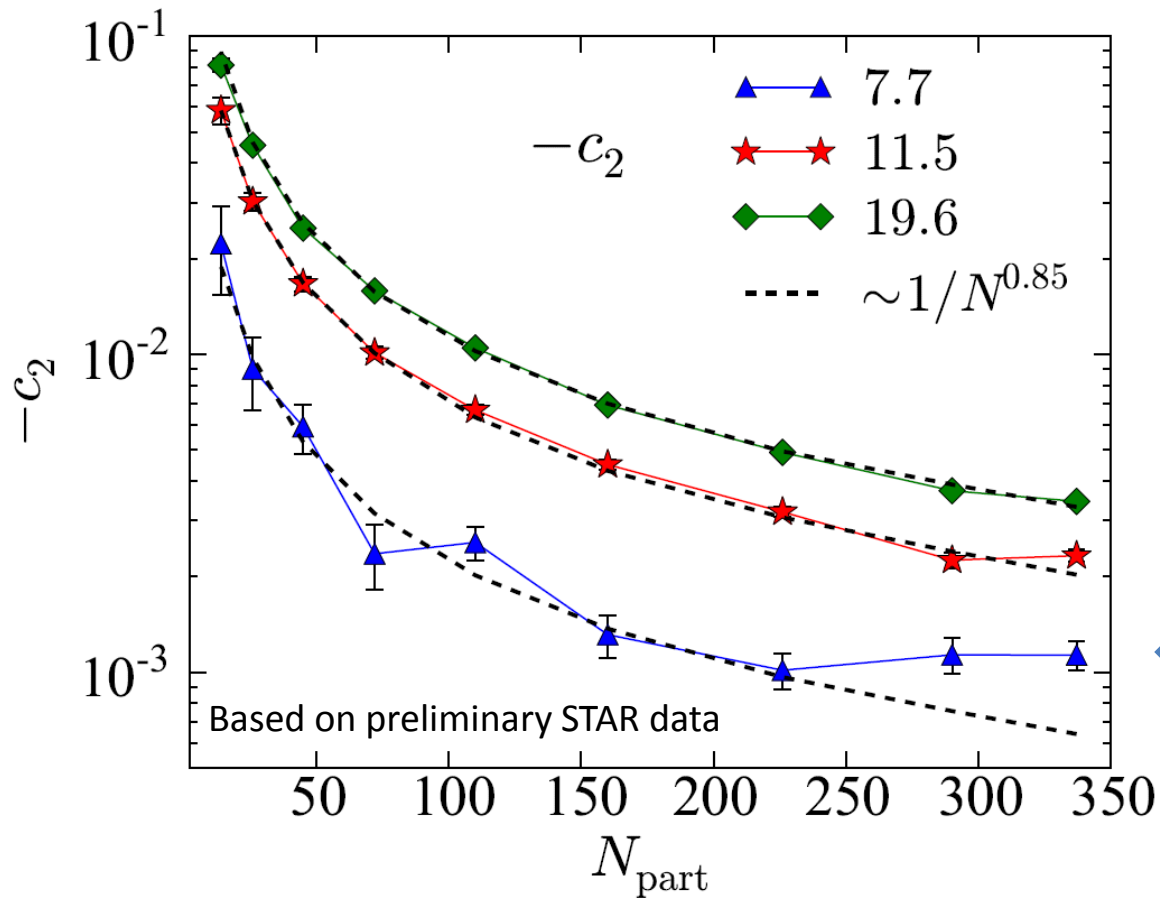
At 7.7 GeV we see  $1/N^2$  for small  $N_{\text{part}}$  then  $c_3$  changes sign and stays roughly constant...

Similar story for  $c_4$

# Comparison of 7.7, 11.5 and 19.6 GeV

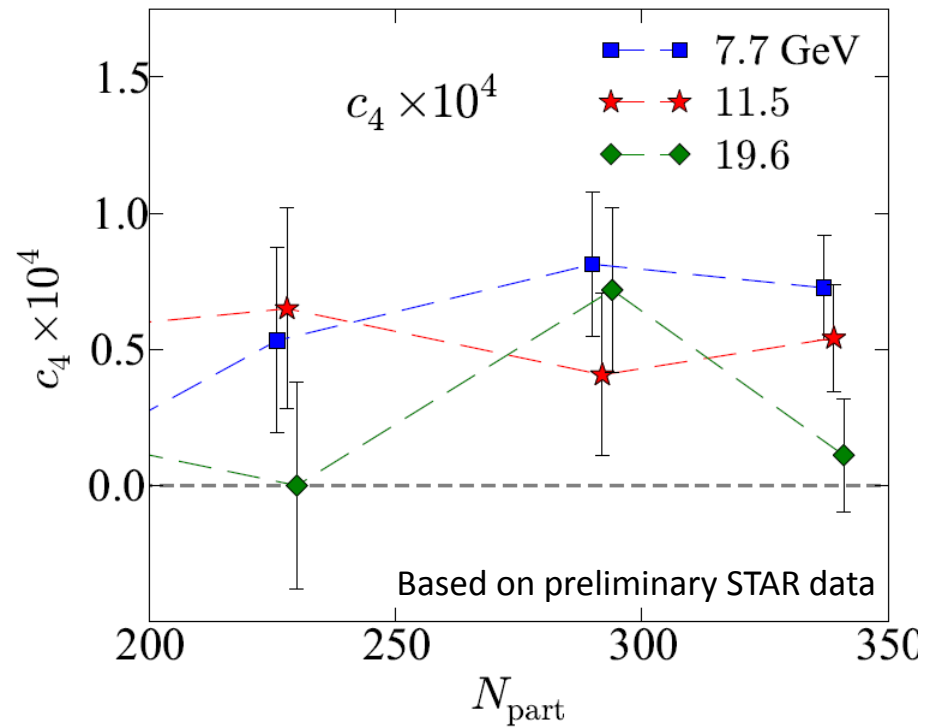
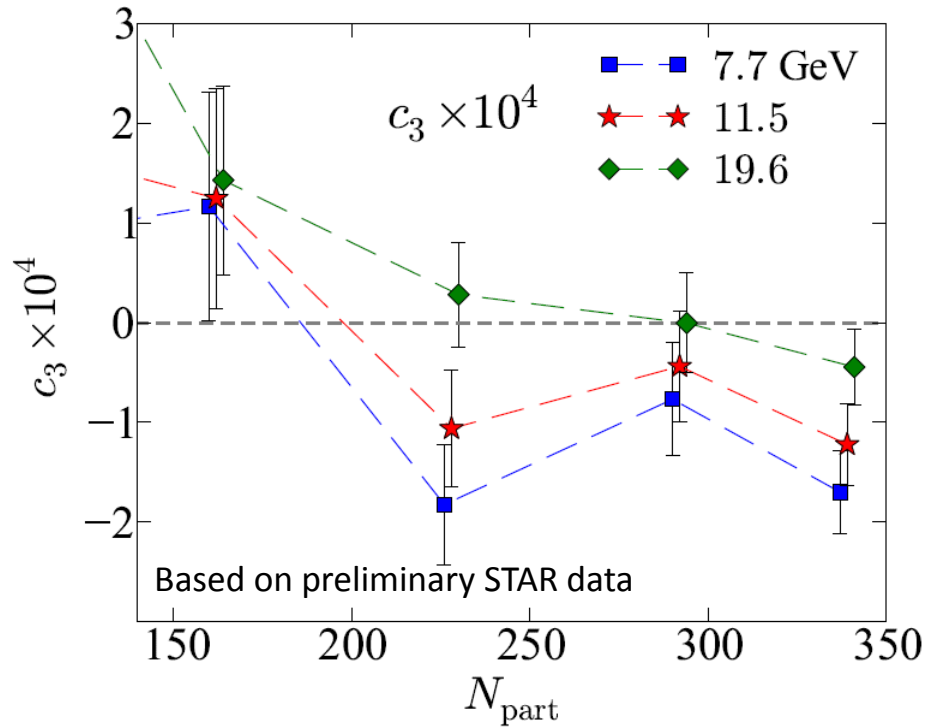


# results for $c_2$



central 7 GeV points are somehow special

# results for central $c_3$ and $c_4$



# results for $c_4$

