

Secondary sources generated by short intense laser pulses: high energy particles and radiation

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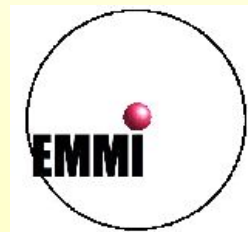


ExtreMe Matter Institute EMMI

Workshop on Plasma Physics
with Intense Laser and Heavy Ion Beams

May 14 - 15, 2009, JIHT, Moscow, Russia

Organized by JIHT RAS and EMMI



Outline

- **Introduction**
 - *What can we do with PW lasers and where we are?*
- **Laser wakefield electron acceleration by PHELIX**
 - *in guiding structures*
 - *In gas jets*
- **Ion sources by optical field ionization**
 - *residual ion energy*
 - *ponderomotive acceleration*
- **X-ray radiation using solid and nano-structured targets**
 - *mechanisms of hot electrons production*
 - *optimization of K_α x-ray yield*
- **Powerful sources of terahertz radiation**
 - *terahertz radiation in magnetized plasma*
 - *low-frequency transition radiation*
- **Mega-gauss magnetic fields in laser plasmas**

➤ *What can we do with PW lasers...?*

from the **Summary**

of the Franco-Russian ISTC Workshop on the physics with **PETAL**,

March 9 - 11, 2009, Le Barp, France

(**PETAL** is a 3.5 kJ, 0.5 – 10 ps, **PetaWatt-class laser** which is under construction in the Aquitaine)

Fundamental physics:

- ✚ **nuclear physics with lasers** (*production of isomeric nuclei*),
- ✚ **physics of high intensity laser plasma interaction** (*strong magnetic and electric fields, nonlocal energy transport*),
- ✚ **high energy density physics** (*development of new, more detailed equations of state, studies of transport and optical properties of non ideal plasma, modeling of atomic transitions and opacities*),
- ✚ **laser driven melting and ablation** (*especially with sub-ps laser pulses*).

It was mutually acknowledged that the collaboration in the domain of development of

secondary sources

of energetic charged particles, neutrons and radiation
is absolutely necessary for all envisioned applications of PETAL

➤ *What can we do with PW lasers...?*

secondary sources ***of high energy particles and radiation***

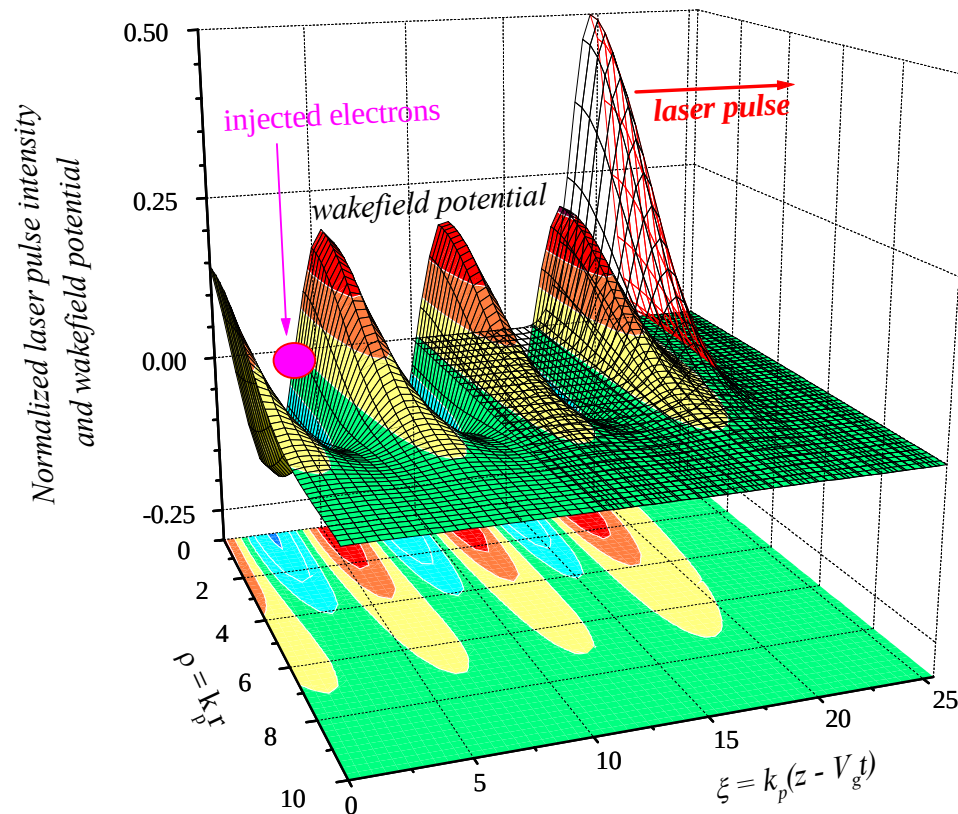
The particular subjects could be:

- + production of quasi-monoenergetic and stable source of protons
(e.g., with the energy ~ 200 MeV for medical applications),**
- + laser wake field electron acceleration,**
- + new sources of X-ray radiation using solid and nano-structured targets,**
- + powerful sources of THz radiation,**
- + laser-driven gamma-ray sources**

➤ laser wake field electron acceleration



Laser-Plasma Acceleration of electrons with record gradients ~ 10 GV/m



The result of optimization of laser and plasma parameters
for the monoenergetic acceleration of a short electron bunch in the wakefield
generated by the 50 TW, 100 fs laser pulse in the plasma channel



Expected potential of laser – plasma acceleration of electrons

Electric field of plasma wave (with phase velocity $\sim c$, $\lambda_p = 2\pi c / \omega_p$):

$$E_p [\text{V/m}] \approx 10^2 \alpha (n_e [\text{cm}^3])^{1/2} \propto \gamma_g^{-1} = \omega_p / \omega_0$$

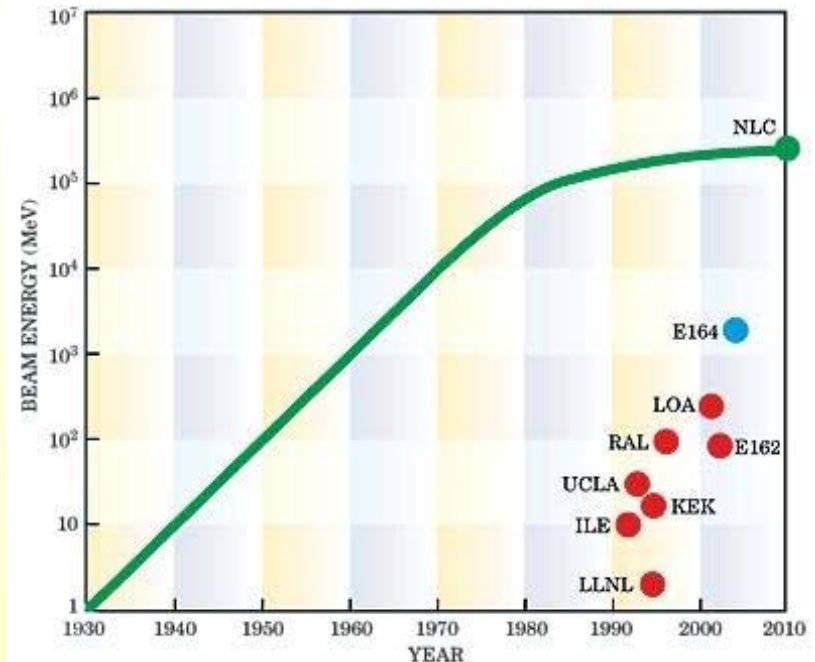
$\alpha = \delta n / n_0$ – plasma wave amplitude; at $\alpha = 0.3 \div 1.0$, $n_e = 10^{17} \div 10^{18} \text{ cm}^{-3}$:

$$E_p = 10 \div 100 \text{ GV/m}$$

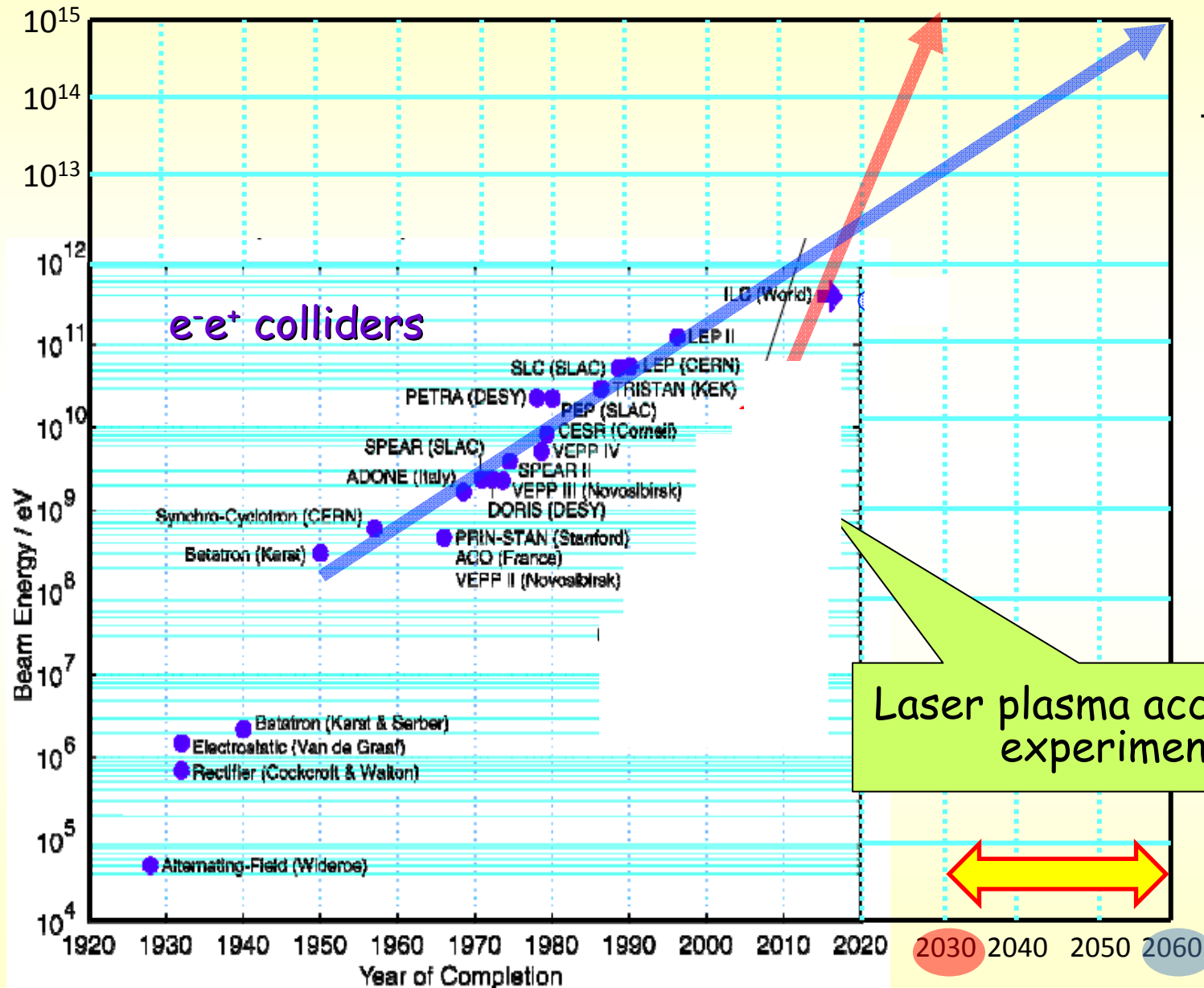
maximum of accelerating gradient
in traditional accelerators (RF linac):

$$E_{\text{RF}} \sim 10 - 100 \text{ MV/m}$$

Exponential growth of “the Livingston curve” began tapering off around 1980



When can we reach 1 PeV?: Suzuki Challenge(2)



(Suzuki,2009)

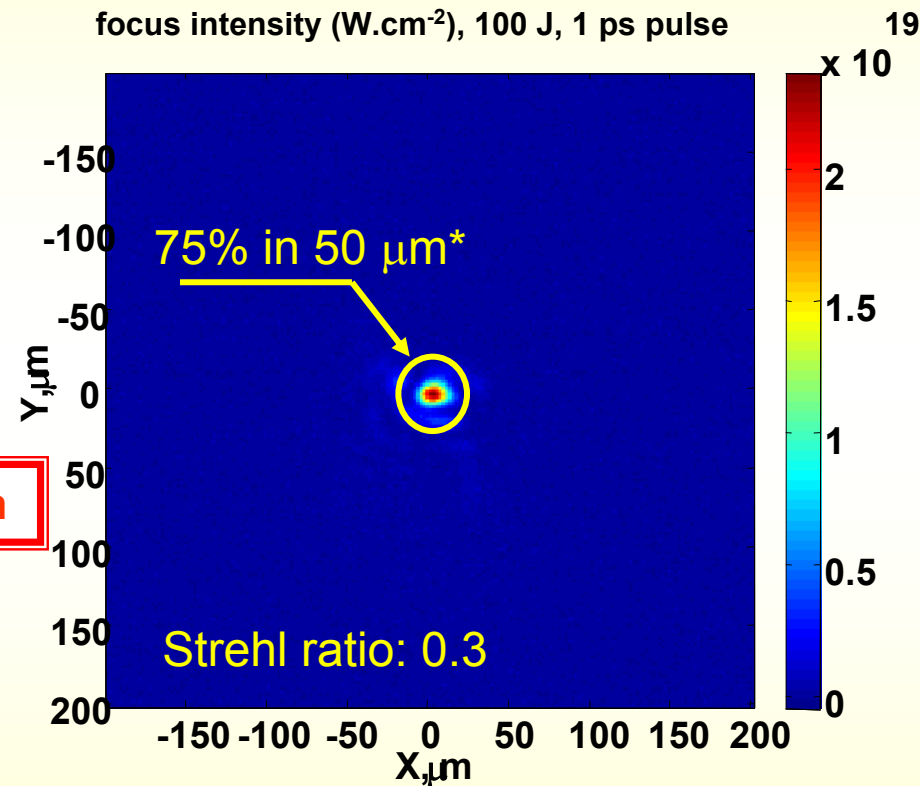
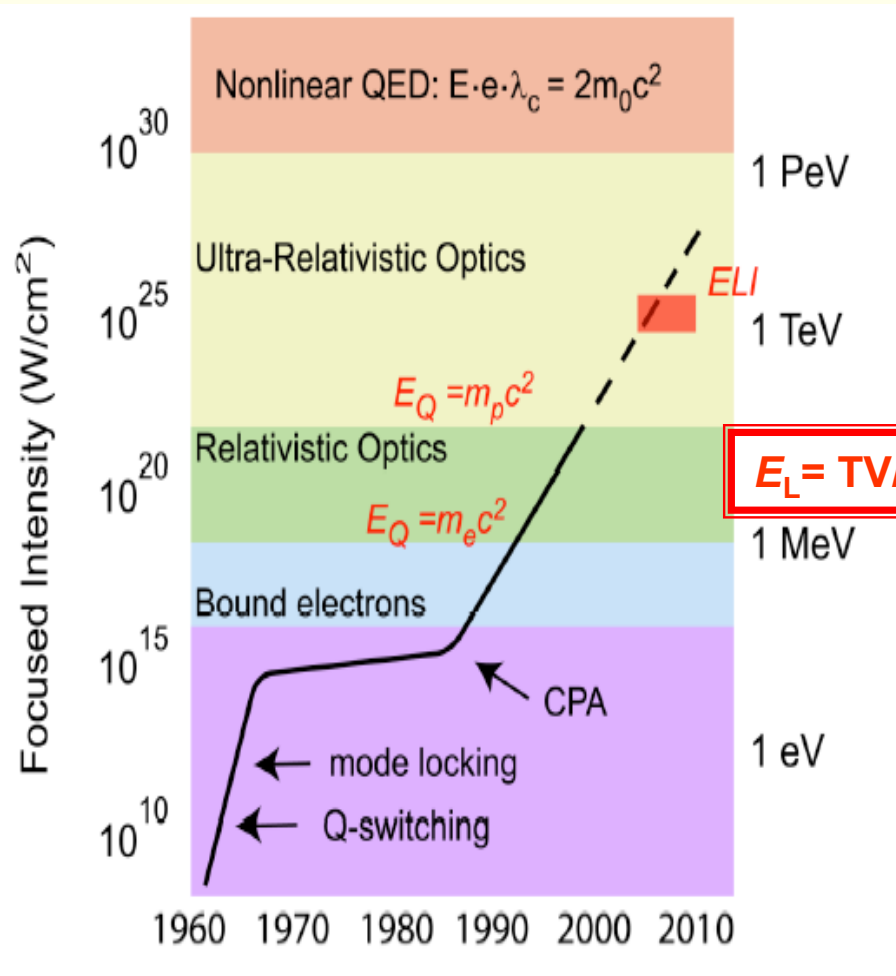
Toshi Tajima



ELI Grand Challenges

Paris, April
27th-28th 2009

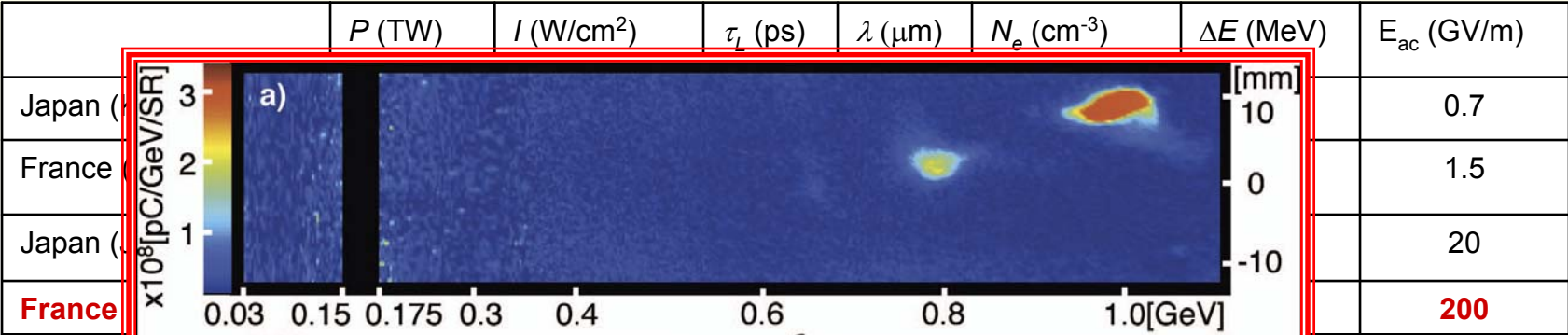
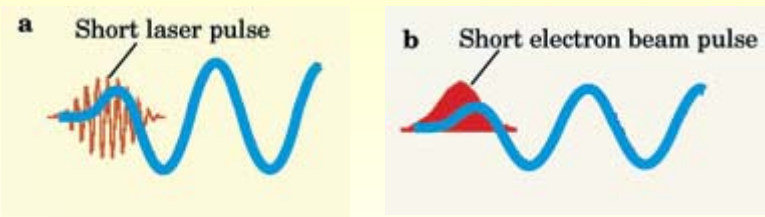
where we are with PHELIX intensity



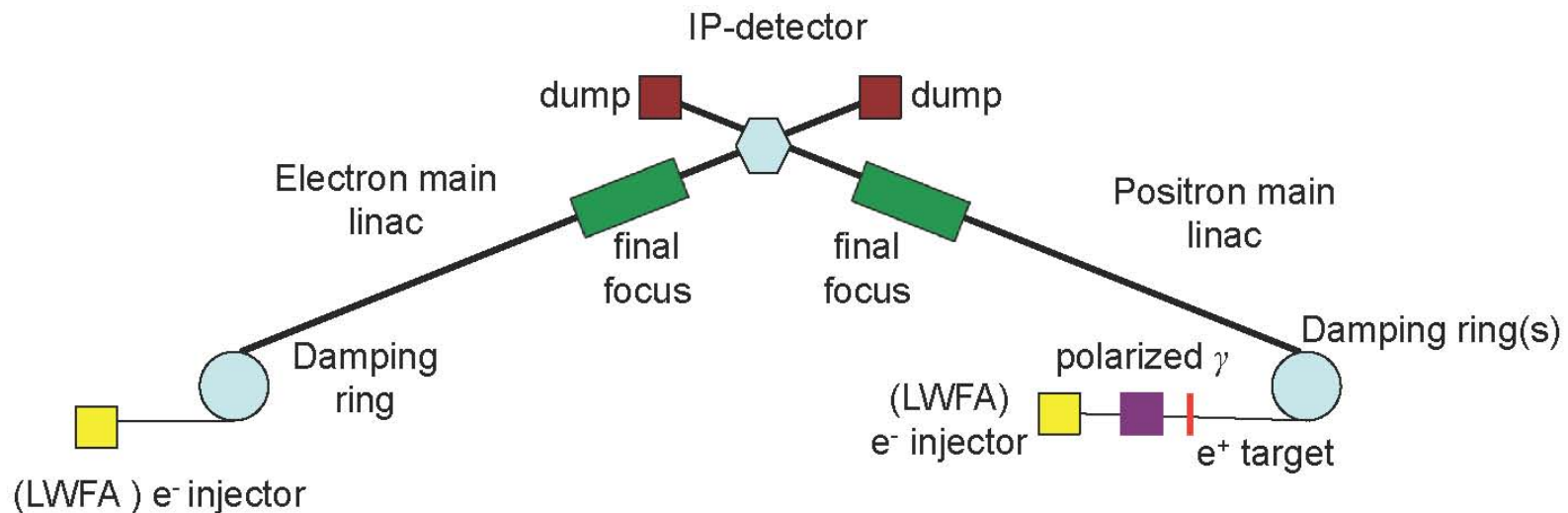
Parameters and results of some experiments

for standard LWFA scheme

$$c\tau_L \approx \lambda_p/2$$



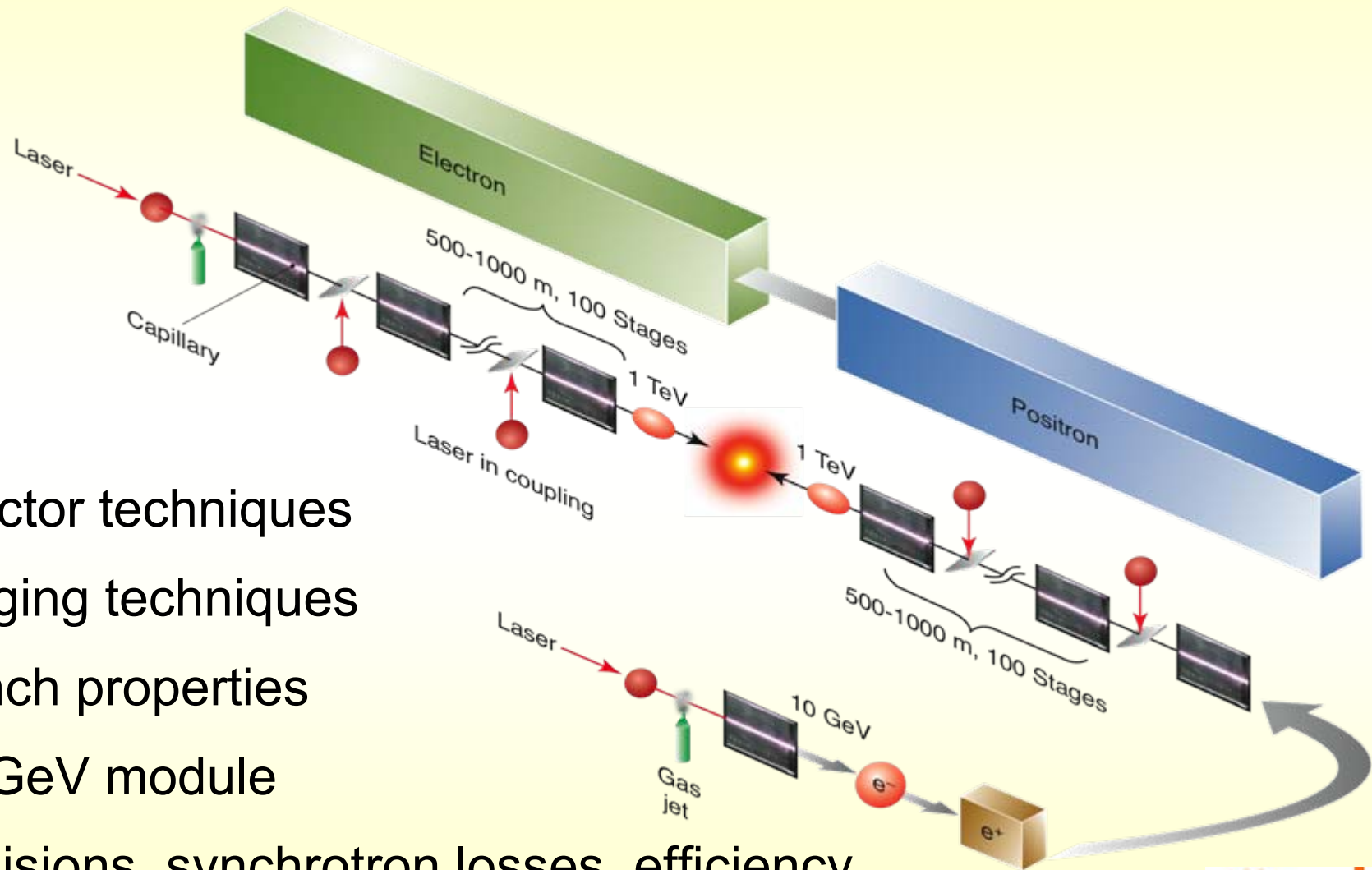
Electron-Positron Linear Collider



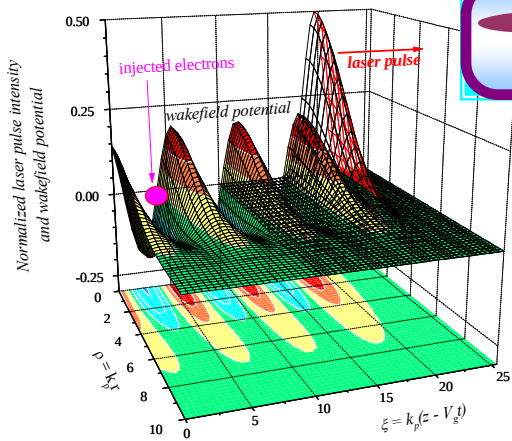
Conventional technology:

- Current generation of future linear collider designs based on existing technology (e.g., ILC): $E_{\text{cm}} \sim 0.5 \text{ TeV}$; gradient $\sim 0.03 \text{ GV/m}$; $\sim 30 \text{ km}$ ($\sim$ multi- $\$B$).
- Higher energy collider with existing technology: $5 \text{ TeV} \rightarrow >100 \text{ km}$, $> \text{tens of } \$B$

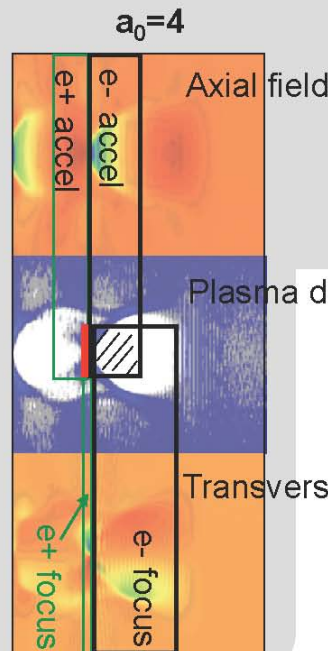
Laser plasma accelerator based concept for a Laser Plasma Linear Collider



- Injector techniques
- Staging techniques
- Bunch properties
- 10 GeV module
- Collisions, synchrotron losses, efficiency

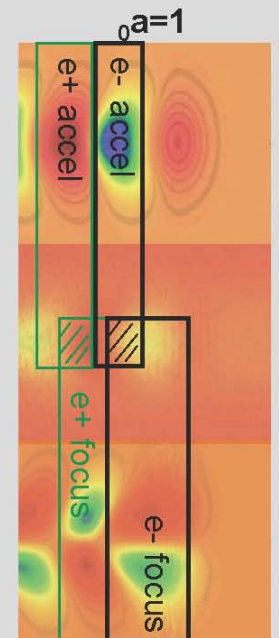


Linear & blowout regimes: e+/e- acceleration



- Blowout regime
 - high field
 - very asymmetric
 - focuses e-
 - defocuses e+

- Quasilinear
 - linear: symmetric e+/e-
 - high a_0 desired for gradient
 - too high enters bubble: ~~e+~~
 - $a_0 \sim 1-2$ good compromise





Physical Restrictions on the Monoenergetic Acceleration of Electrons/Positrons in Laser-Plasma Accelerators

- ❖ *Dephasing of electron/positron bunches;*
- ❖ *Diffraction of the laser pulse;*
- ❖ *Depletion of the laser pulse;*
- ❖ *Electron/positron bunches loading effect*
- ❖ *Focusing/defocusing of electron/positron bunches*
- ❖ *Nonlinear laser pulse distortions*

Physical Restrictions on the Energy of Accelerated Electrons in Laser-Plasma Accelerators

Dephasing of electron/positron bunches

Maximum length of acceleration:

$$l_a \leq L_{deph} = \frac{\lambda_p}{2(c - v_{ph})} c = \gamma_g^2 \lambda_p, \quad \gamma_g = \frac{\omega_0}{\omega_p} = \sqrt{\frac{n_c}{n_e}}$$

Energy gain: $\Delta W_e = e E_a l_a \cong \frac{e \Delta \varphi}{\lambda_p / 2} l_a = 2 m c^2 \gamma_g^2 \alpha$

Laser Pulse Diffraction:

$$L_{diff} = \pi Z_R = \frac{2 \pi^2 r_L^2}{\lambda}, \quad a_0^2 \propto P_L / r_L^2 \quad \alpha \propto a_0^2$$

for LWFA $\Delta W_{diff} [\text{MeV}] = 960 \frac{\kappa \lambda [\mu\text{m}]}{\tau [\text{fs}]} P [\text{TW}]$

Main Processes Accompanying the Propagation of Short Intense Laser Pulses in Gases

Typical parameters

Laser pulse:

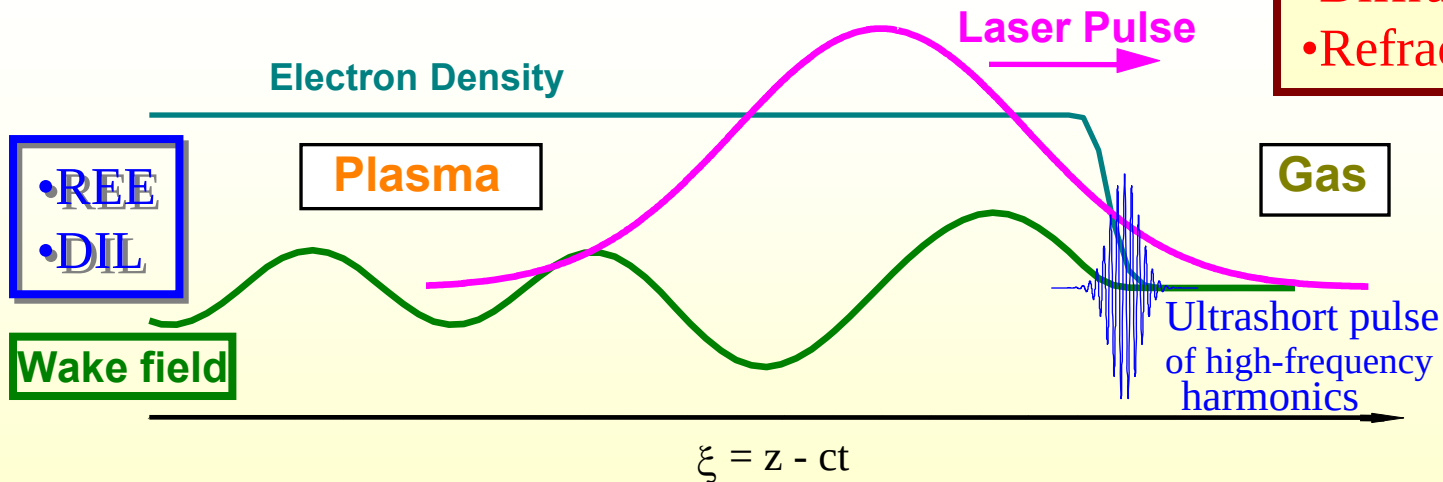
$I_{\max} > 10^{15} \text{ W/cm}^2$, FWHM $< 1 \text{ ps}$, $\lambda \approx 1 \text{ mkm}$

Gas:

comparatively light atoms $Z < 10$, $n_{\text{at}} < 10^{19} \text{ cm}^{-3}$

Laser Pulse Distortions:

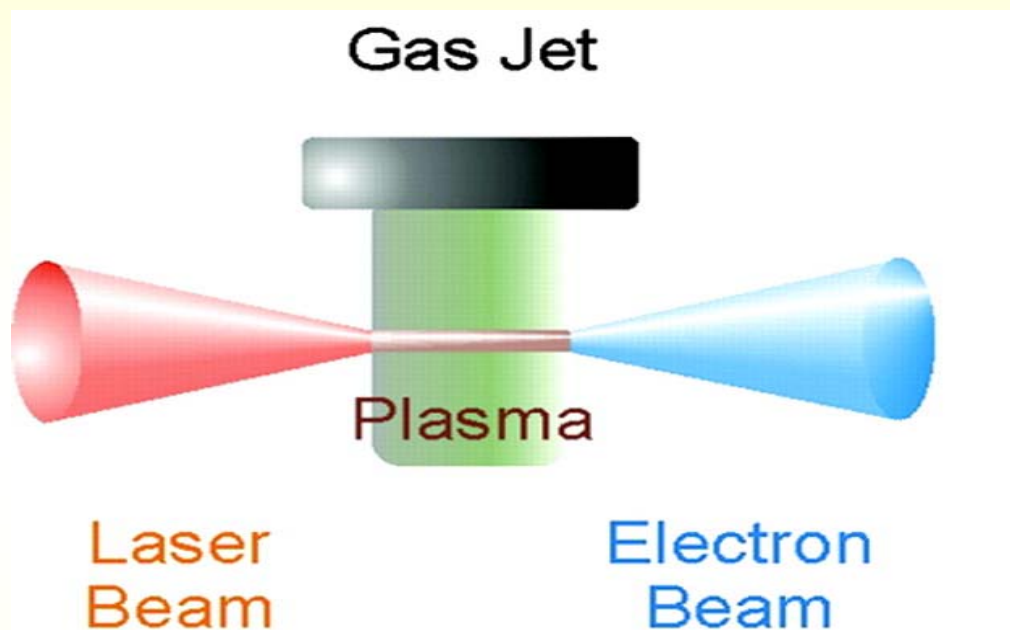
- LASRS
- RS-F & S-M
- Filamentation
- Diffraction
- Refraction



Electron Acceleration in the



Focus



Laser Parameters:

$$\lambda_0 = 1.053 \quad \tau_{\text{FWHM}} = 300 \text{ fsec}$$
$$w_0 = 100 \text{ } \mu\text{m}$$
$$I_L = 2 \times 10^{18} \text{ W/cm}^2, \quad E_L = 190 \text{ J}$$

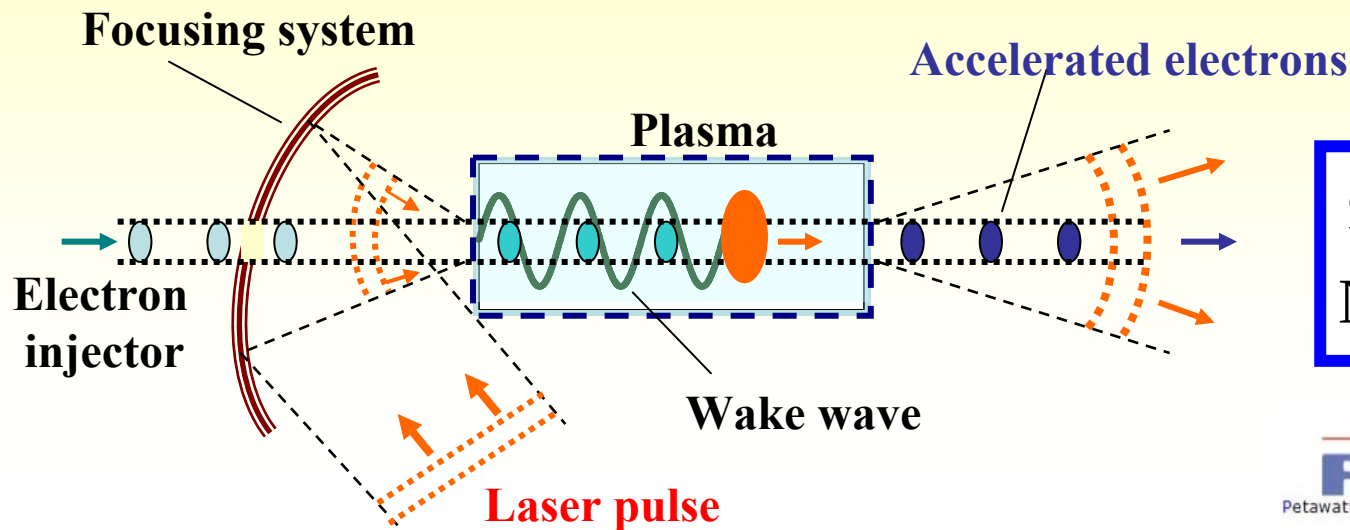
Plasma Density:

$$n_e = 2.5 \cdot 10^{16} \text{ cm}^{-3}$$
$$\gamma_g = 200$$

$$N_{e\text{-bunch}} \approx 10^9$$
$$Q_b \approx 200 \text{ pC}$$

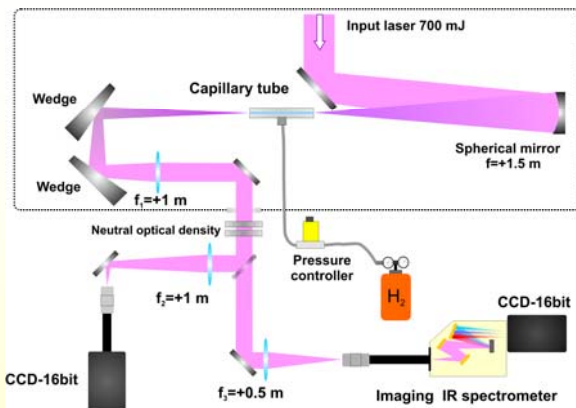
Injected Electrons with $E_{\text{inj}} = \gamma_g mc^2$ are accelerated to $E \sim 0.8 \text{ GeV}$,
Length of acceleration $l_a \sim L_{\text{diff}} \sim 10 \text{ cm}$

Laser Wakefield Electron Acceleration



$$Q_b \leq 100 \text{ pC}$$

$$N_{\text{e-bunch}} \approx 10^9$$



B. Cros, et al. Schematic view of the experiment at the Lund Laser Center

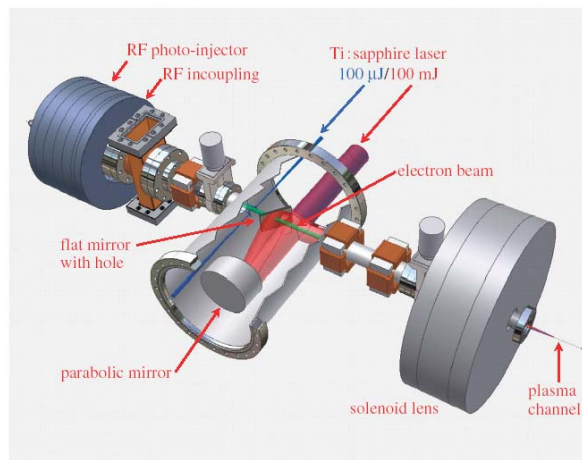


Figure 1. Layout of experiment planned at TU-Eindhoven on LWA with external injection.

$$\lambda_0 = 1.053 \mu\text{m}$$

$$\tau_{\text{FWHM}} = 300 \text{ fs}$$

$$w_0 = 200 \mu\text{m}$$

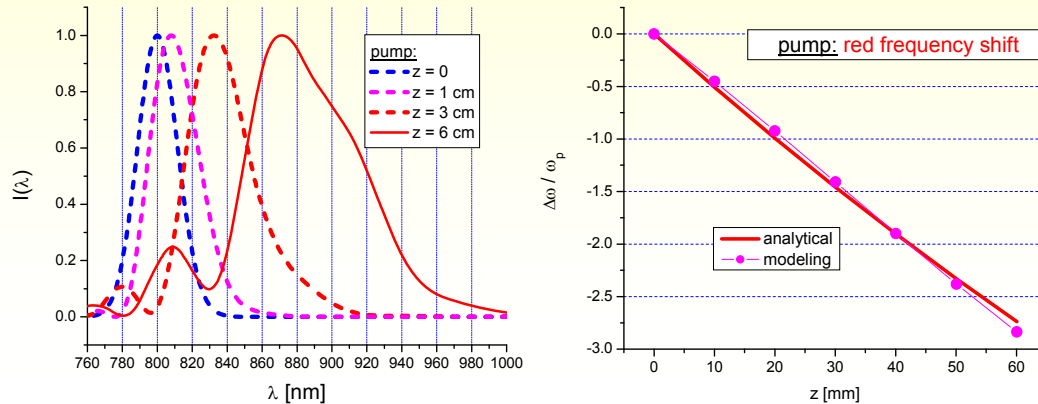
$$E_L = 150 \text{ J}$$

$$n_e = 1.7 \cdot 10^{16} \text{ cm}^{-3}$$

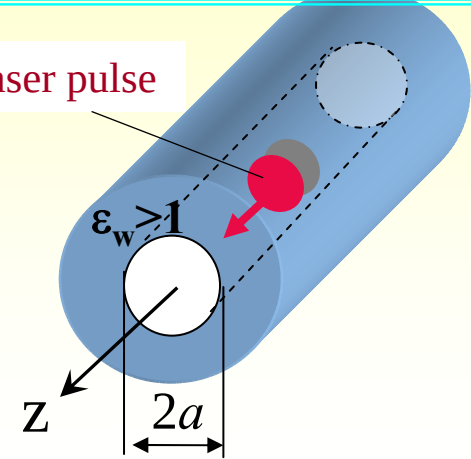
$$E \sim 2.7 \text{ GV/m}$$

$$L_{\text{trap}} \sim 10 \text{ cm}, L_{\text{acc}} \sim 1 \text{ m}$$

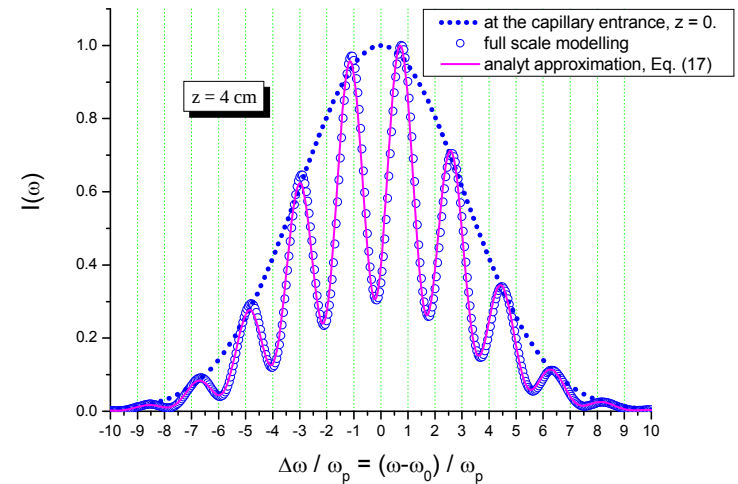
Frequency Shift of the Pump Pulse



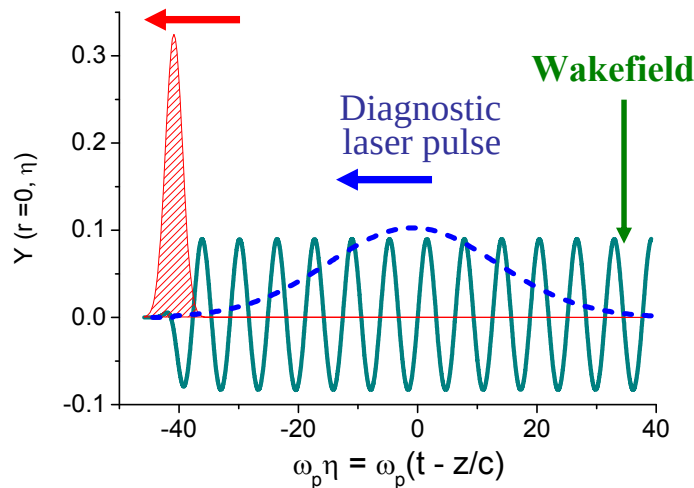
Laser pulse



Spectrum Modulation of the Probe Pulse

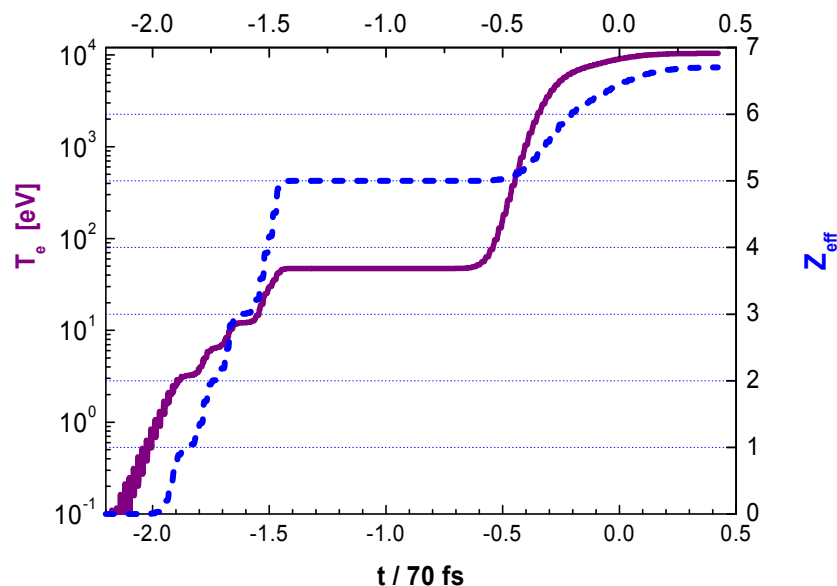


Short intense laser pulse



$$\bar{A}(z) = C_r \frac{k_0}{k_p^2 z} \left[\sin \left(\frac{1}{4} \frac{\Omega^2 \beta}{1 + \beta^2} \right) \right]^{-1} [\bar{I}(z) - \bar{I}(z = 0)]$$

Multi-group Residual Electron Energy Distribution of many-electron atoms



laser pulse parameters:
 $I_L = 10^{19} \text{ W/cm}^2$, $\tau_L = 70 \text{ fs}$,
 $\lambda = 0.78 \mu\text{m}$
Nitrogen gas

Energy of electron groups and relaxation times ($N_2 = 1.5 \times 10^{19} \text{ W/cm}^2$)

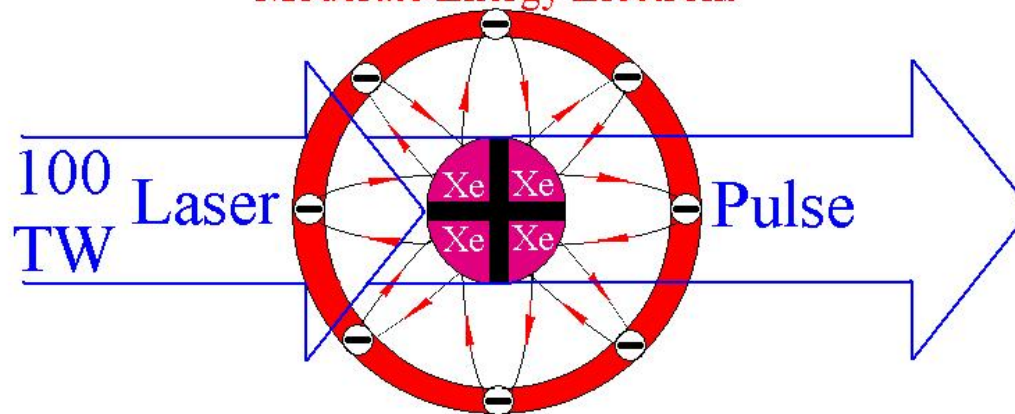
k	1	2	3	4	5	6	7
$Q_{e,k}$ [eV]	3.3	9.7	23.4	79.5	120.	39×10^3	39×10^3
$t_{ee,k}$	60 fs	150 fs	380 fs	1.8 ps	2.6 ps	13 ns	13 ns

A large amount of ions with low temperature (less than 100 eV), low momentum spread and narrowed charge state distribution can be produced by petawatt-class short-pulse laser

Low Temperature Xe^{+26} Ions Laser Source

Residual ion energy is the least unavoidable ion energy after ionization of gases by a short intense laser pulse

Moderate Energy Electrons



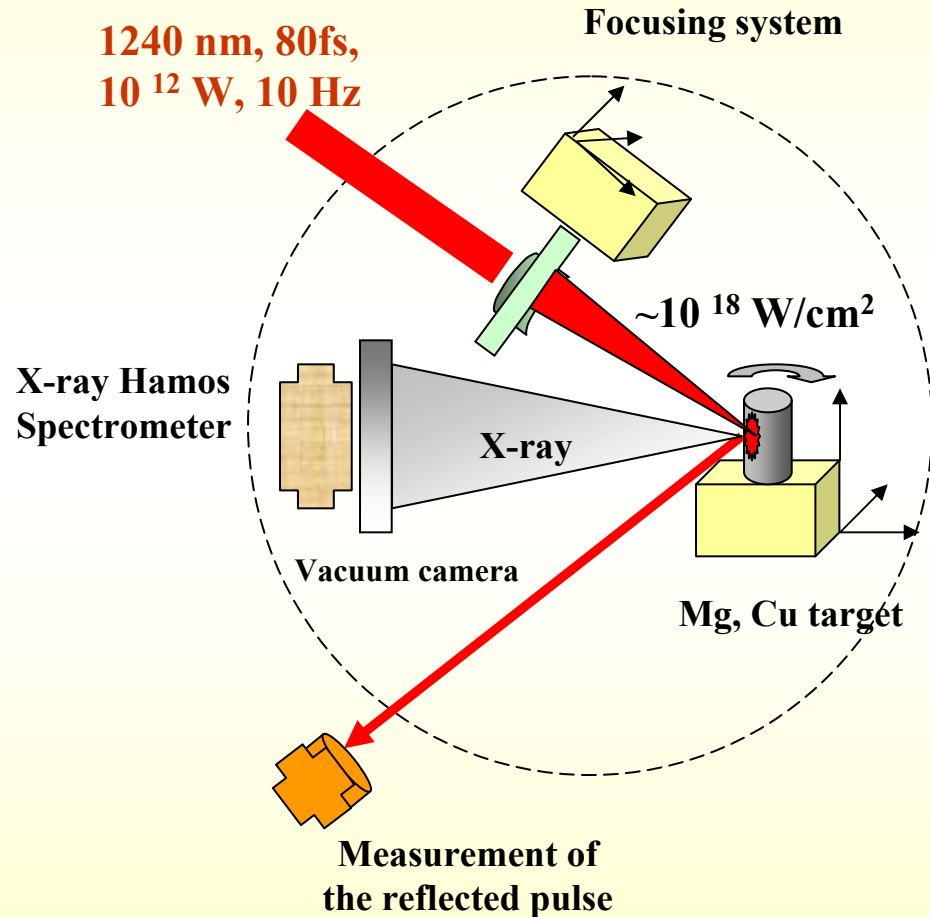
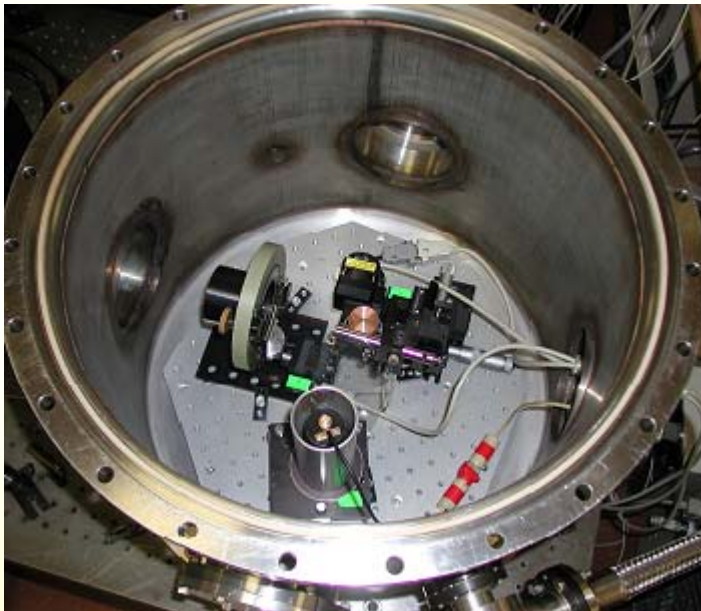
Moderate and Low Energy Ionization Electrons are Caught by Ambipolar Electric Field and Produce Undesirable:

- 1) Ion Heating, 2) Collisional Ionization, 3) Recombination and ponderomotive acceleration of ions by laser pulse

➤ new sources of X-ray radiation using solid and nano-structured targets

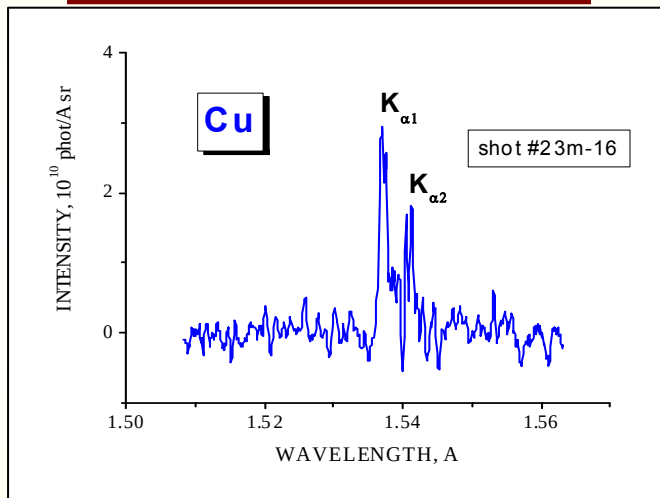


X-ray radiation of laser plasma

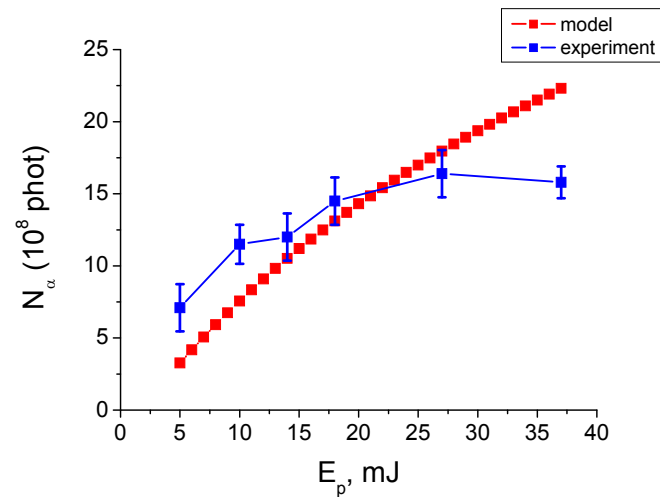


K_α radiation from Cu laser-produced plasmas
 Cr:Forsterite laser system, ω_0 , p -polarization,
 $E_L = 30$ mJ, $\tau_L = 80$ fs, $10^{16} \div 10^{17}$ W/cm²

Single laser short experiment



Optimization of K_α x-ray yield



The model calculation of K_α yield against pulse energy in comparison with experimental data received at IHED laser facility:

$\lambda = 1.24$ μ m, $\tau_p = 80$ fs,
 $d_f = 10$ μ m, $\theta = 45^\circ$, Fe target

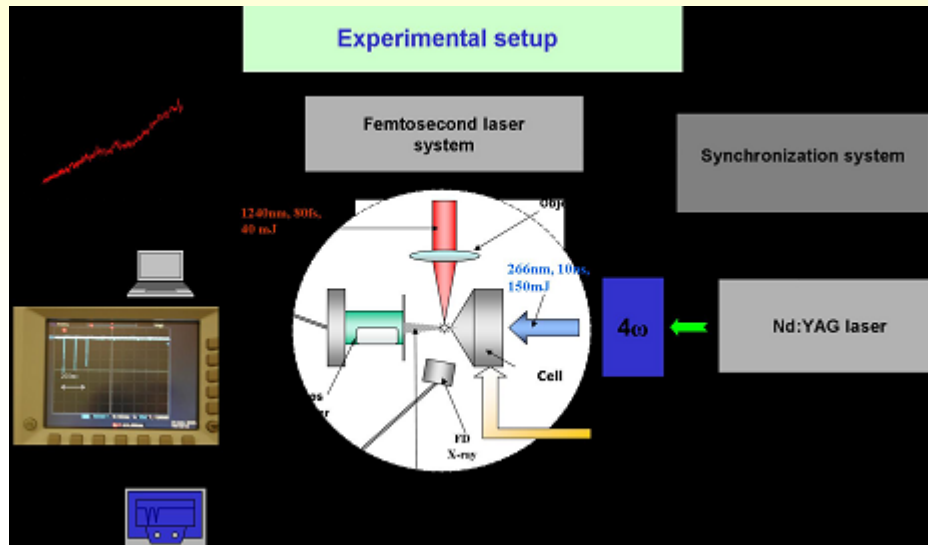
$$E_0 = E_L [1 + (1 - f)^{1/2}] \sin \theta$$

$$T_{h,keV} = 7.6 I_{L,16} \lambda_{\mu m}^2 \alpha^2 \sin^2 \theta$$

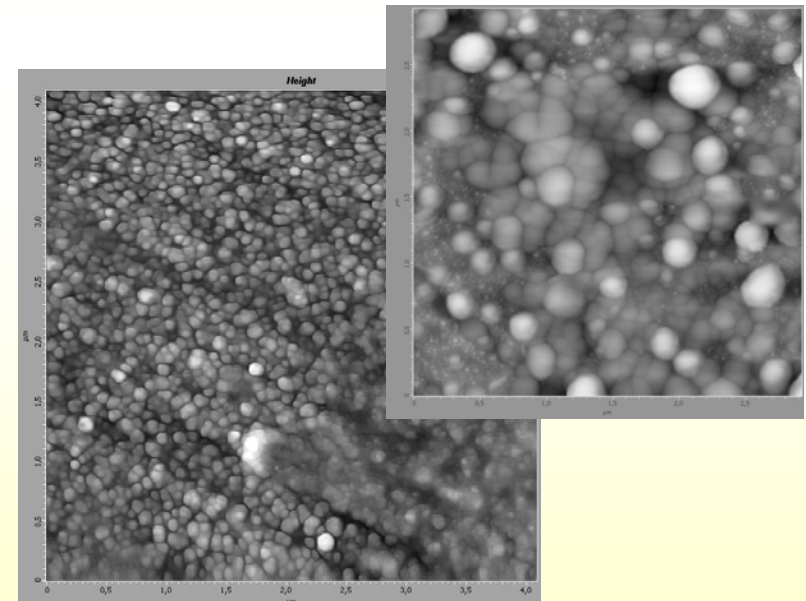
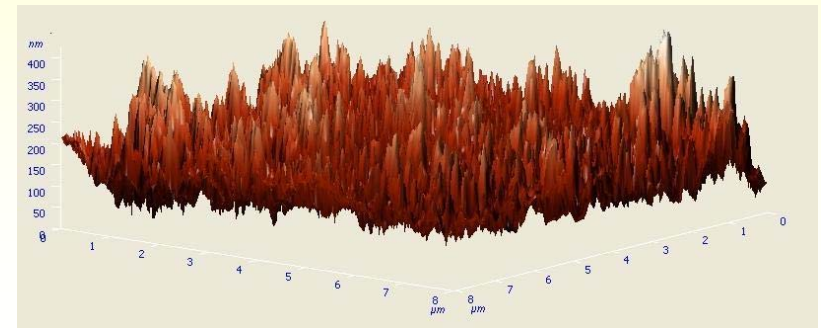
E_0 – driving, E_L – laser field, f – absorption, $\eta = 1.57$, $I_{L,16} = I_L / 10^{16}$ W/cm² – laser intensity

There is a range of laser pulse parameters in which K_α yield may be described by vacuum heating mechanism

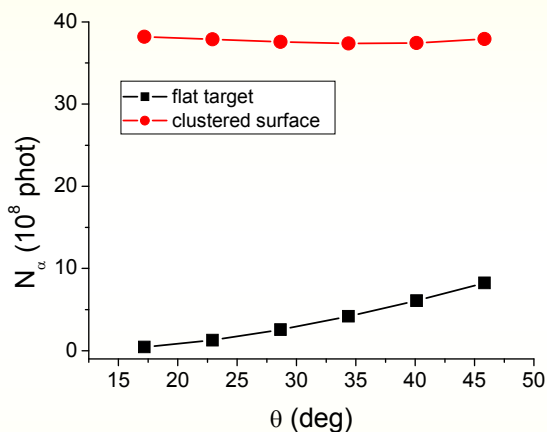
nano-structured targets



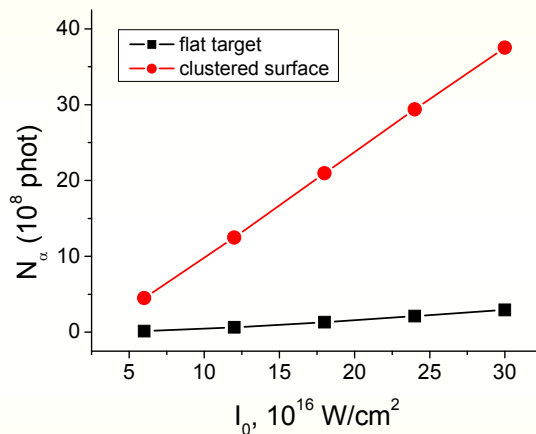
Fe clusters on Cu target



Fe clusters on Cu target



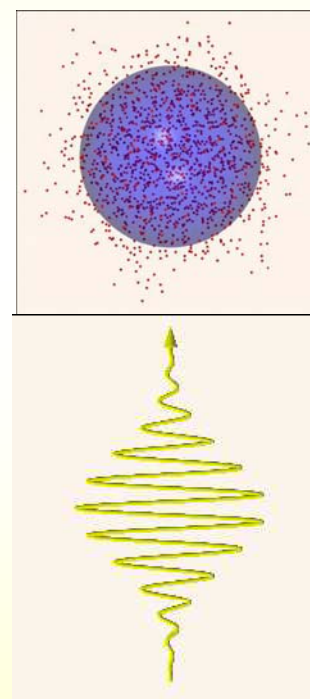
Dependence of K_α yield on angle of incidence for peak intensity $3 \times 10^{17} \text{ W/cm}^2$



Dependence of K_α yield on peak intensity for angle of incidence 30°

$$T_{h,keV} = 68.41 \times I_{L,16} \lambda_{\mu m}^2 \cos^2 \theta_1$$

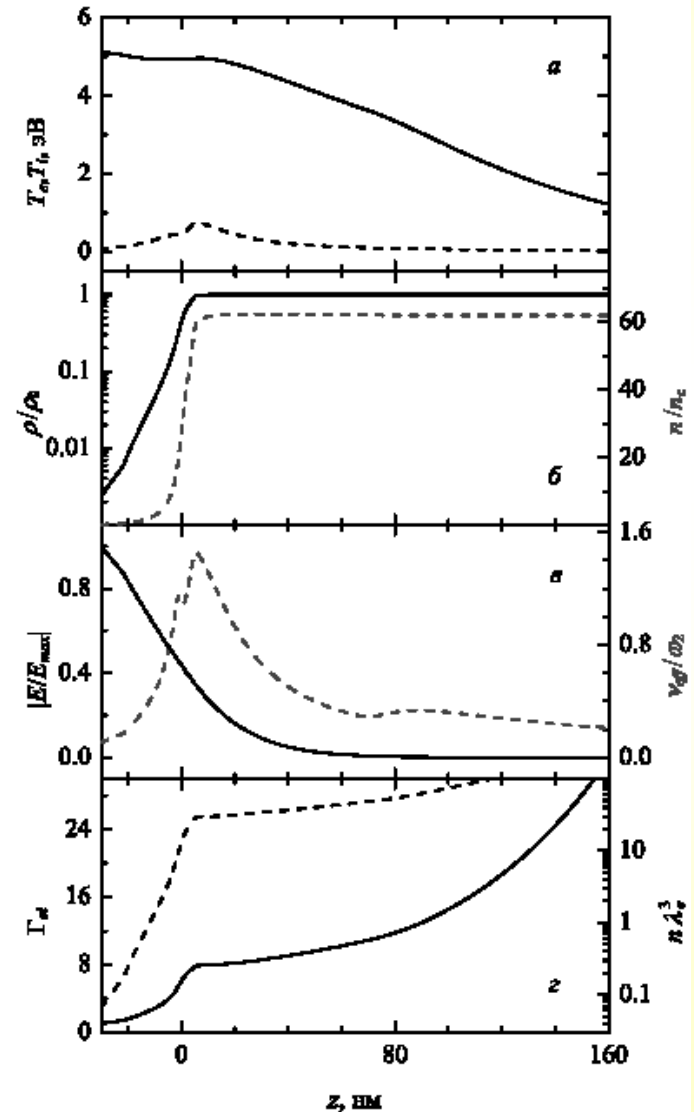
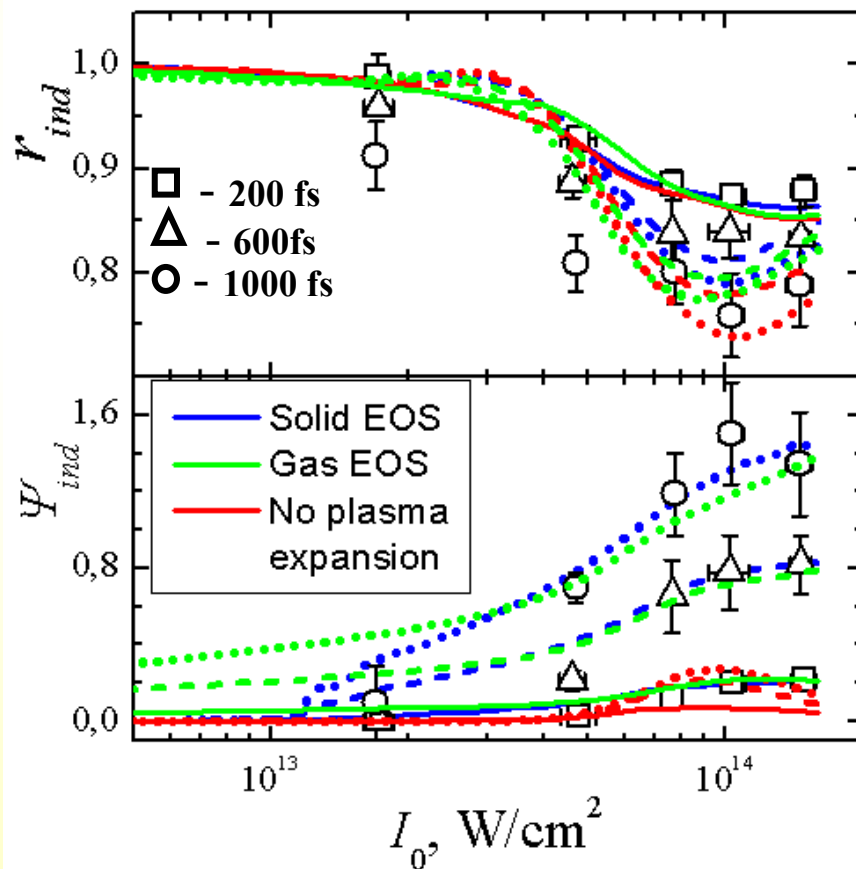
$$E_0 = 3E_L \cos \theta_1$$



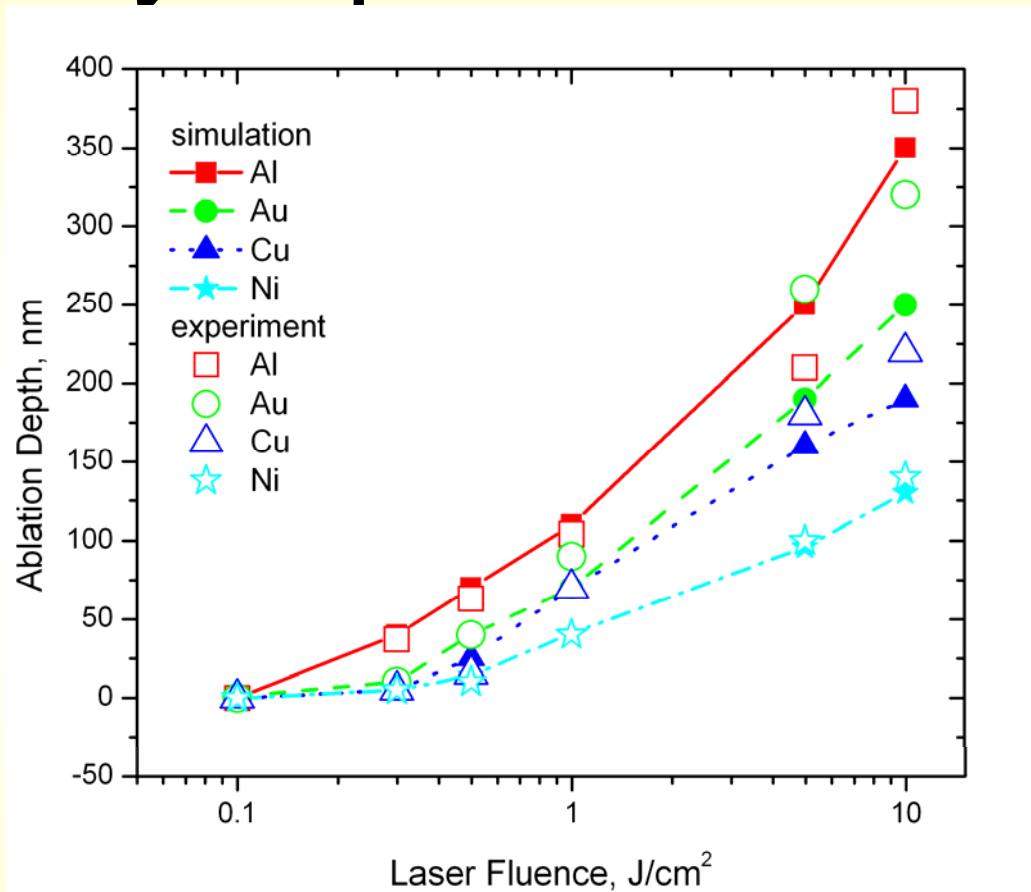
Enhancement of driving electric field at a cluster surface and favorable conditions for K_α photons to escape from the wafer lead to considerable enhancement of K_α yield

Femtosecond optical diagnostics and hydrodynamic simulation of Al plasma

Laser pulse: $\tau_L = 100$ fs; $I < 2 \cdot 10^{14}$ W/cm²



Ablation of metal targets by fs pulses

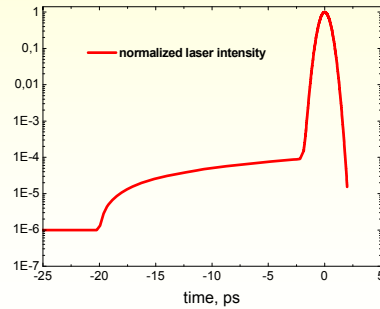
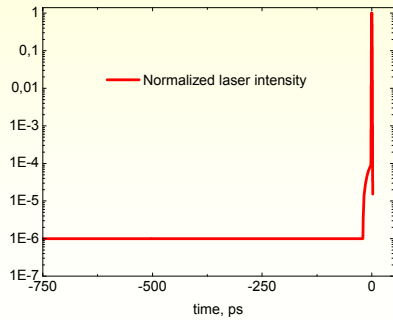


M. E. Povarnitsyn *et al.* // PRB **75**, 235414 (2007).

M. E. Povarnitsyn *et al.* // Appl. Surf. Sci. **253**, 6343 (2007)

M. E. Povarnitsyn *et al.* // Appl. Surf. Sci. (in press, 2008)

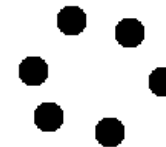
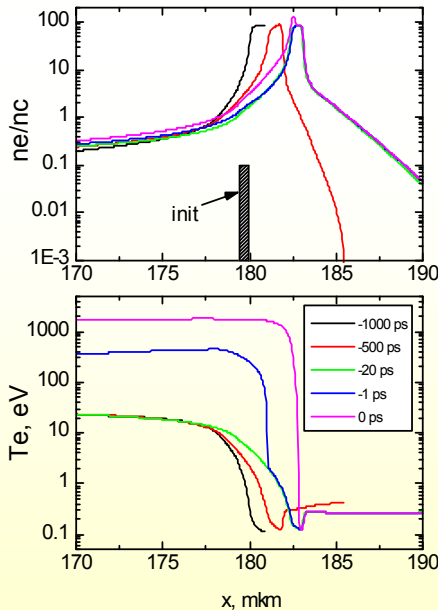
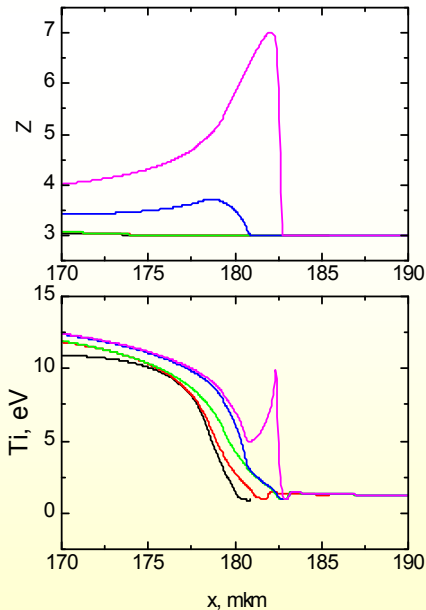
Pilot modeling of the nanostructure dynamics



Interaction with Cu wires (n_e/n_{cr})

000 fs

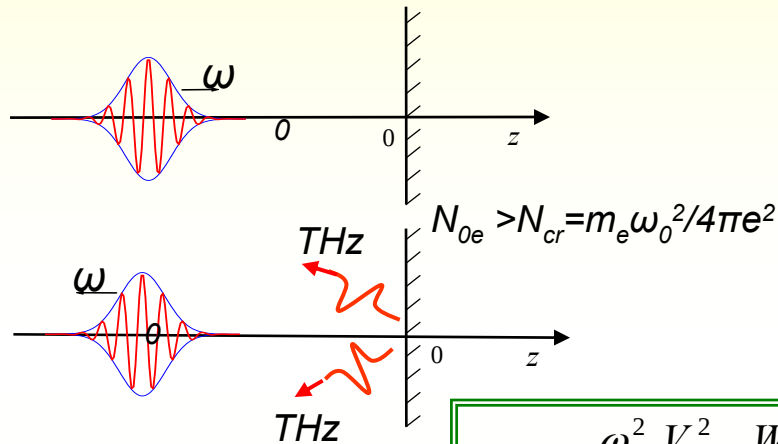
10^{16} W/cm^2



$$\lambda = 1.053 \text{ μm}, I_0 = 10^{17} \text{ W/cm}^2, \tau = 500 \text{ fs}$$

produced plasma with the density of order critical one will expand
over distances of the order of 5 μm before the main pulse.

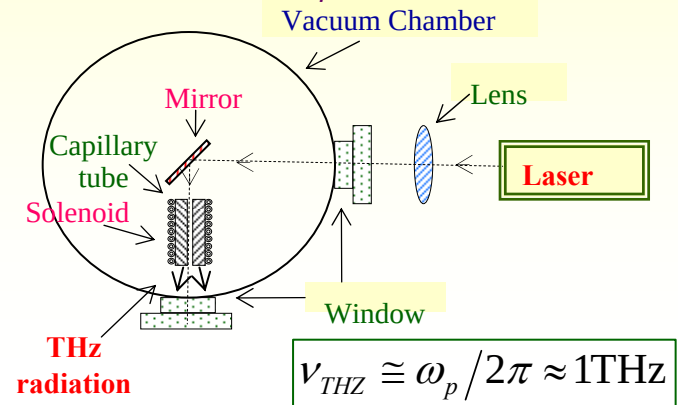
SiO₂ aerogel, $\rho=0.02$ g/cm³



$$\nu_{THz} \cong 1/\pi\tau \approx 1\text{THz}$$

$$W_{THz} \cong \frac{\omega_0^2}{\omega_p^2} \frac{V_E^2}{c^2} \frac{W_L}{\omega_p^2 \tau^2}$$

Magnetic field ($\omega_p \gg \Omega_e = eB_0/m_e c$)



$$\frac{dW_{THz}}{dz} = \frac{0.5}{\omega_0^2 \tau^2} \left(1 + \frac{2}{k_p^2 R_L^2} \right) \frac{V_E^2}{4c^2} \frac{W_L}{L}$$

Laser: $\lambda_0=1.053\mu\text{m}$; $\tau_{FWHM}=0.5\text{ps}$; $I_L=2 \cdot 10^{18} \text{ W/cm}^2$; $R_L=50 \mu\text{m}$ ($W_L=100\text{J}$)

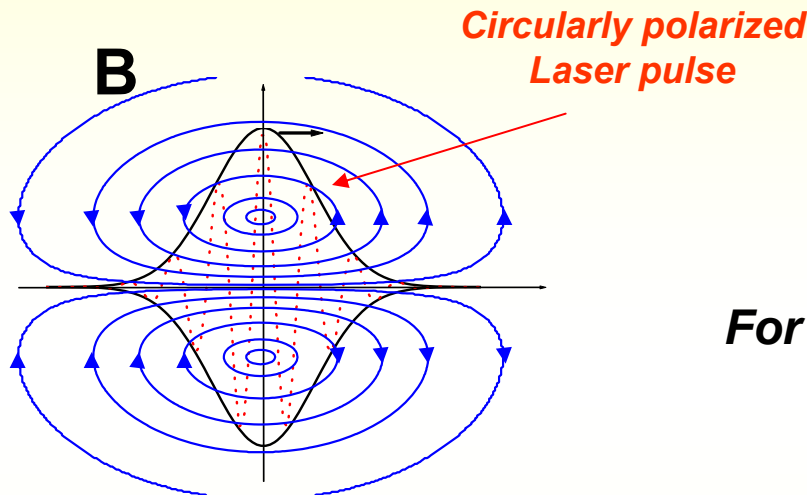
Plasma: $N_{0e}=6 \cdot 10^{21} \text{ cm}^{-3}$; $\omega_p \tau=1200$

THz: $P_{THz} \approx 30 \text{ MW}$; $W_{THz} \approx 10 \mu\text{J}$;
 $W_{THz}/W_L \approx 10^{-7}$

Plasma: $N_{0e}=1.4 \cdot 10^{16} \text{ cm}^{-3}$; $\omega_p = 6.7 \cdot 10^{12} \text{ s}^{-1}$

THz: $W_{THz}/W_L \approx 3 \cdot 10^{-4} z[\text{cm}]$; $z=10\text{cm}$:

$P_{THz} \approx 10\text{GW}$; $W_{THz} \approx 0.3\text{J}$; $W_{THz}/W_L \approx 3 \cdot 10^{-3}$



$$\frac{eB_z}{m_e c \omega_0} = \frac{\omega_p^2}{2\omega_0^2} \frac{V_E^2 / c^2}{(1 + V_E^2 / c^2)}$$

**For relativistic laser intensities, $V_E / c > 1$
at critical density, $\omega_p \approx \omega_0$**

$$\Omega_e \approx \omega_0$$

Laser: $\lambda_0 = 1.053 \mu\text{m}$; $\tau_{FWHM} = 0.5 \text{ ps}$ ($\tau = 0.3 \text{ ps}$); $I_L = 3 \cdot 10^{19} \text{ W/cm}^2$; $R_L = 15 \mu\text{m}$
($W_L = 100 \text{ J}$; $P_L = 300 \text{ TW}$)

gas: $N_{0e} = 3 \cdot 10^{19} \text{ cm}^{-3} \rightarrow \mathbf{B_z} \approx 3 \text{ M}\Gamma\text{c}$; aerogel: $N_{0e} \sim 10^{21} \text{ cm}^{-3} \rightarrow \mathbf{B_z} \approx 100 \text{ MGauss}$

$$\Omega = 1.8 \cdot 10^7 B_0 [\text{G}] s^{-1} \quad \omega_p = 6.7 \cdot 10^{12} s^{-1} \quad \Omega = 0.1 \omega_p \Rightarrow B_0 = 3.7 \cdot 10^4 \text{ G}$$

People who say it cannot be done should not interrupt those who are doing it.

George Bernard Shaw

1929



LHC, 2009



Size x 10^5

Energy x 10^9

Helmholtz Alliance

Extremes of Density and Temperature: Cosmic Matter in the Laboratory

ExtreMe Matter Institute EMMI

Workshop on Plasma Physics
with Intense Laser and Heavy Ion Beams

May 14 - 15, 2009, JIHT, Moscow, Russia



Organized by:

Joint Institute for High Temperature, Russia and
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Thank

You for attention!

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