



# Impact of vacuum quality in light sources: beam lifetime and ion instabilities

U. Iriso

ALBA - CELLS

# OUTLINE

## 1. Beam Lifetime

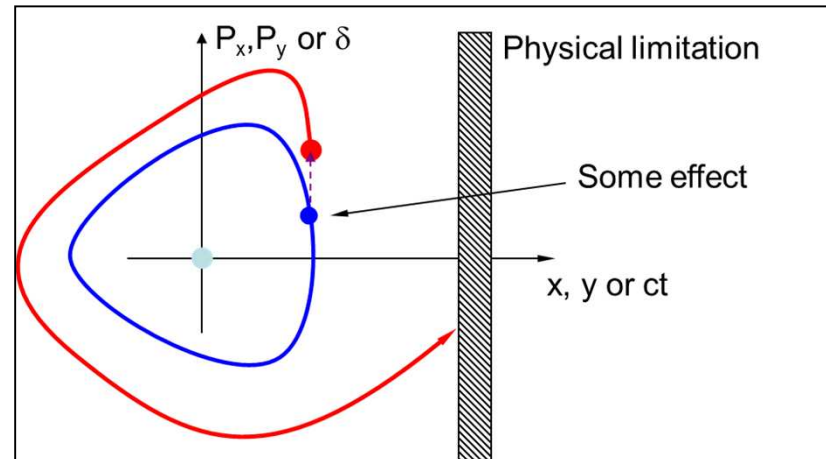
- a) Analytical description
- b) Observations in Light Sources

## 2. Ion Instabilities

- a) Analytical description
- b) Observations in Light Sources

## 3. Conclusions

# Lifetime: Particle Loss in LS



Main Lifetime contributors:

1. **Gas Lifetime:**

Interactions with residual gas nucleus or electrons (elastic or inelastic)

2. **Touschek effect:**

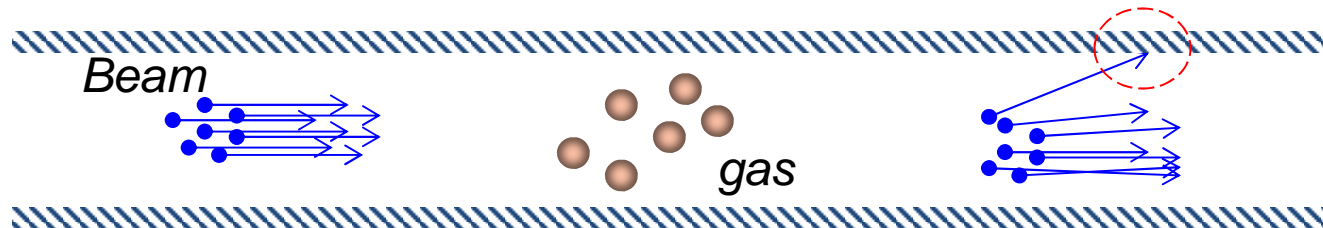
Interactions among beam particles with energy transfer.

3. **Quantum Lifetime:**

Emission of radiation quanta when particles are close to the aperture limitations. Usually negligible, since it is  $\sim 10^5$ h or more.

# Elastic Gas Scattering

Deflections caused by the **electrical force** of the residual gas atoms



Simplified Eqs\*:

$$\frac{1}{\tau_{\text{nuclei-elast}}} \approx \frac{2\pi r_e^2 c Z^2}{k_B T \gamma^2} \frac{\langle P_i \cdot \beta_{yi} \rangle}{A^2}.$$

$$\frac{1}{\tau_{\text{electron-elast}}} \approx \frac{2\pi r_e^2 c P Z}{k_B T \gamma} \frac{1}{\delta_{\text{acc}}}$$

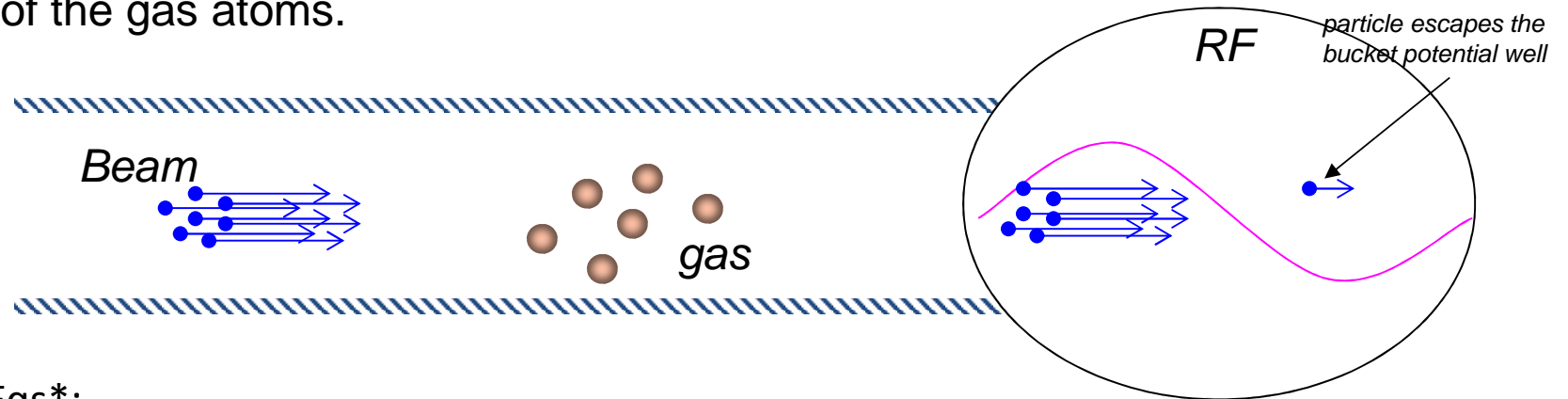
Note for LS, e-elast ~ [2-3]  
orders of magnitude larger

- Number of residual particles (pressure):
- Particle Charge:
- Beta function.
- Vacuum chamber:
- Beam energy:

$P \uparrow \rightarrow \tau \downarrow$   
 $Z \uparrow \rightarrow \tau \downarrow$   
 $\beta \uparrow \rightarrow \tau \downarrow$   
 $A \uparrow \rightarrow \tau \uparrow$   
 $\gamma \uparrow \rightarrow \tau \uparrow$

# Inelastic Gas Scattering

Energy loss caused by **radiation emission (mostly)** at the vicinity of the gas atoms.



Simplified Eqs\*:

$$\frac{1}{\tau_{\text{nuclei-inelast}}} \approx \frac{4r_e^2 c P Z^2}{137k_B T} \frac{4}{3} \ln \left( \frac{183}{Z^{1/3}} \right) \left[ \ln \left( \frac{1}{\delta_{\text{acc}}} \right) - \frac{5}{8} \right]$$

$$\frac{1}{\tau_{\text{electron-inelast}}} \approx \frac{4r_e^2 c P Z}{137k_B T} \frac{4}{3} \ln \left( \frac{2.5\gamma}{\delta_{\text{acc}}} - 1.4 \right) \left[ \ln \left( \frac{1}{\delta_{\text{acc}}} \right) - \frac{5}{8} \right]$$

Both shall be accounted for LS

1. Number of residual particles (pressure):
2. Charge of every particle:
3. RF acceptance.

$P \uparrow \rightarrow \tau \downarrow$   
 $Z \uparrow \rightarrow \tau \downarrow$   
 $\delta_{\text{acc}} \uparrow \rightarrow \tau \uparrow$

# Touscheck Lifetime (I)

- **Limiting effect in LS** (although not in the scope of this workshop)
- Due to energy exchange between particles within the bunch, which can bring the particles out of the energy acceptance
- It is not exponential, but asymptotic:

$$\frac{dN(t)}{N(t)dt} = -aN(t) \quad \longrightarrow \quad N(t) = \frac{N_0}{1+t/\tau}$$

- $\tau$  complex expression\*, but its main dependencies:

$$\tau_t \propto \frac{\sqrt{\epsilon_y} \sigma_z}{I_b} \delta_{acc}^3$$

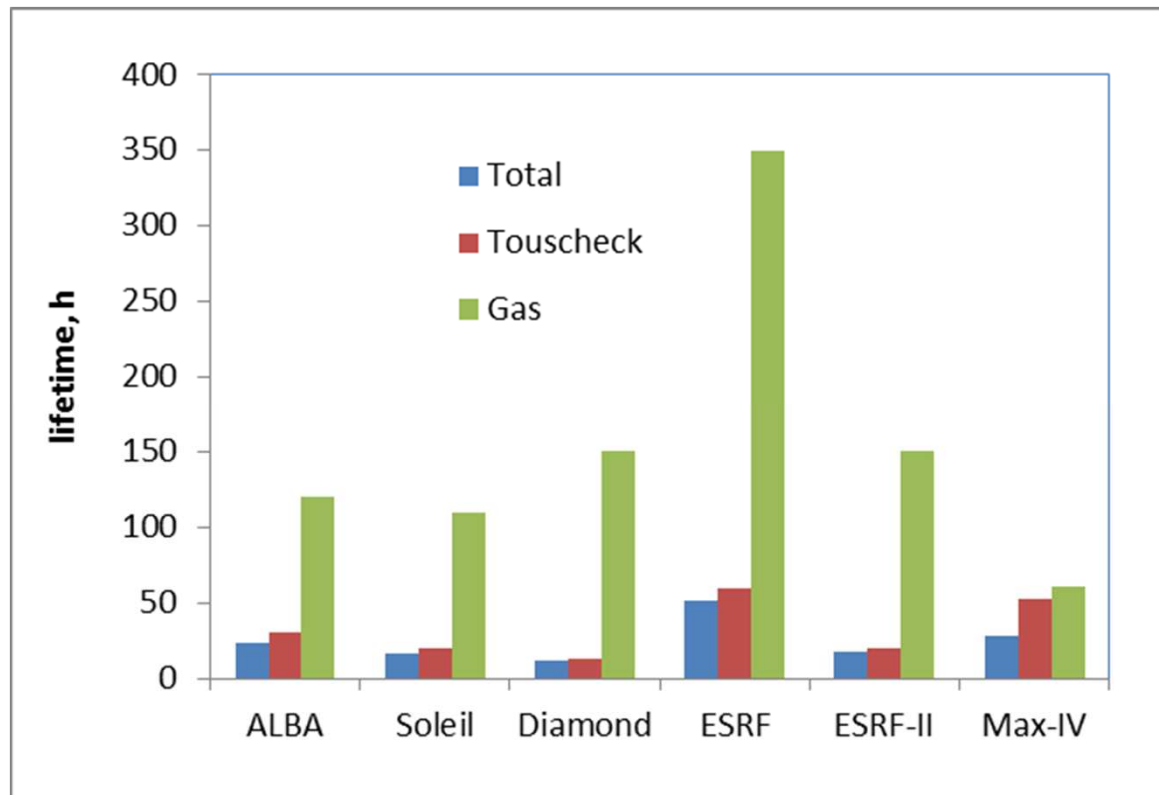
→ **Product  $\tau^* I_b = ct$**  , widely used in LS (instead of  $\tau$ )

- |  |   |
|--|---|
| 1. $V_{RF} \uparrow \rightarrow$ RF acceptance | $\delta_{acc} \uparrow \rightarrow \tau \uparrow$ |
| 2. Beam density (coupling, bunch length).      | $\rho \uparrow \rightarrow \tau \downarrow$       |
| 3. Beam bunch intensity:                       | $I_b \uparrow \rightarrow \tau \downarrow$        |

\*See Piwinsky Equation, Chao's book

# Lifetime in LS

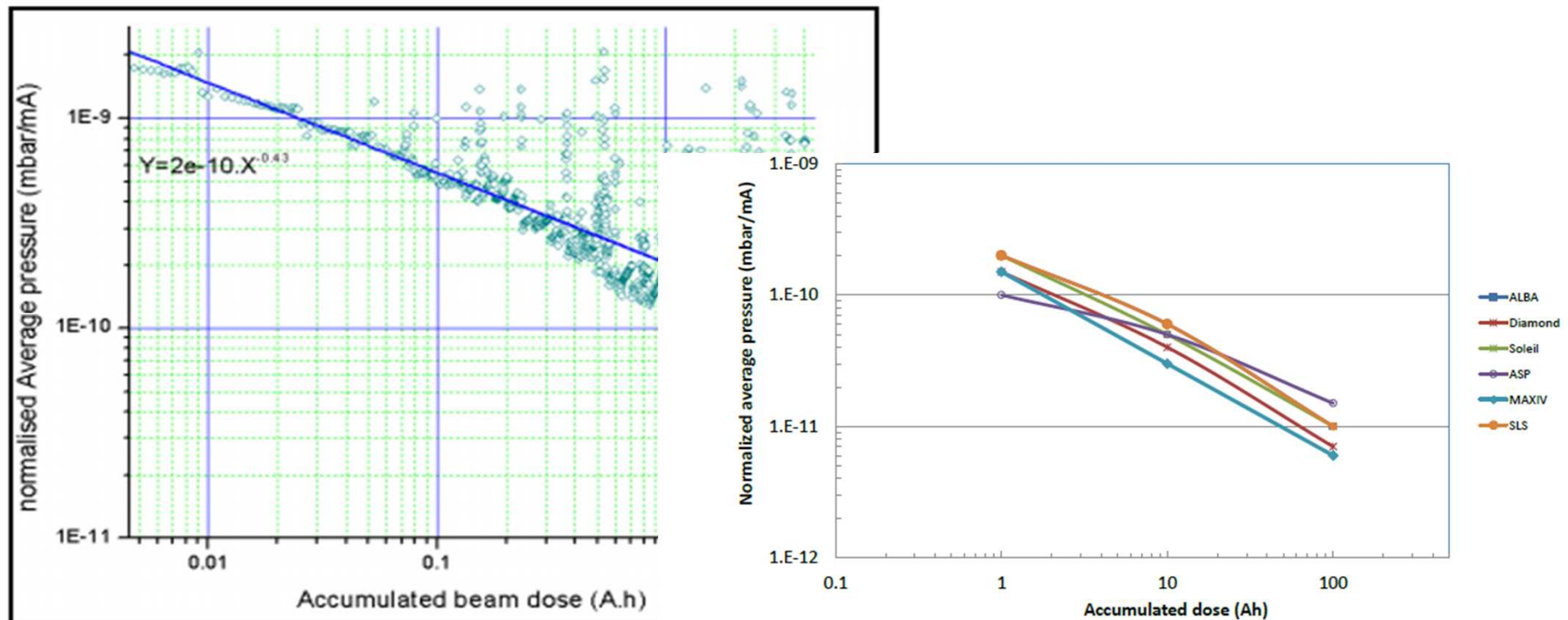
- Lifetime in LS is dominated by Touscheck effect, except in MAX-IV prediction@500mA
- Gas lifetime (including elastic & inelastic) only important at early commissioning phases or during installation of new IDs



Soleil: X.N. Gavalda, Phd Thesis  
 Diamond: I. Martin, priv. communications  
 ESRF: N. Carmignani, priv. communications  
 MAX-IV: S.C. Leeman, PRST-AB 12, 120701  
 (2009) – **500mA**

# Gas Lifetime in LS: commissioning

- Rule of thumb: Gas lifetime  $\sim 100\text{h}$ , when  $P\text{-avg} < 1\text{e-}9$ , which is usually achieved after  $\sim 100\text{ A}\cdot\text{h}$
- At ALBA, this was achieved after  $\sim 6$  months of operation



D. Einfeld, IPAC11

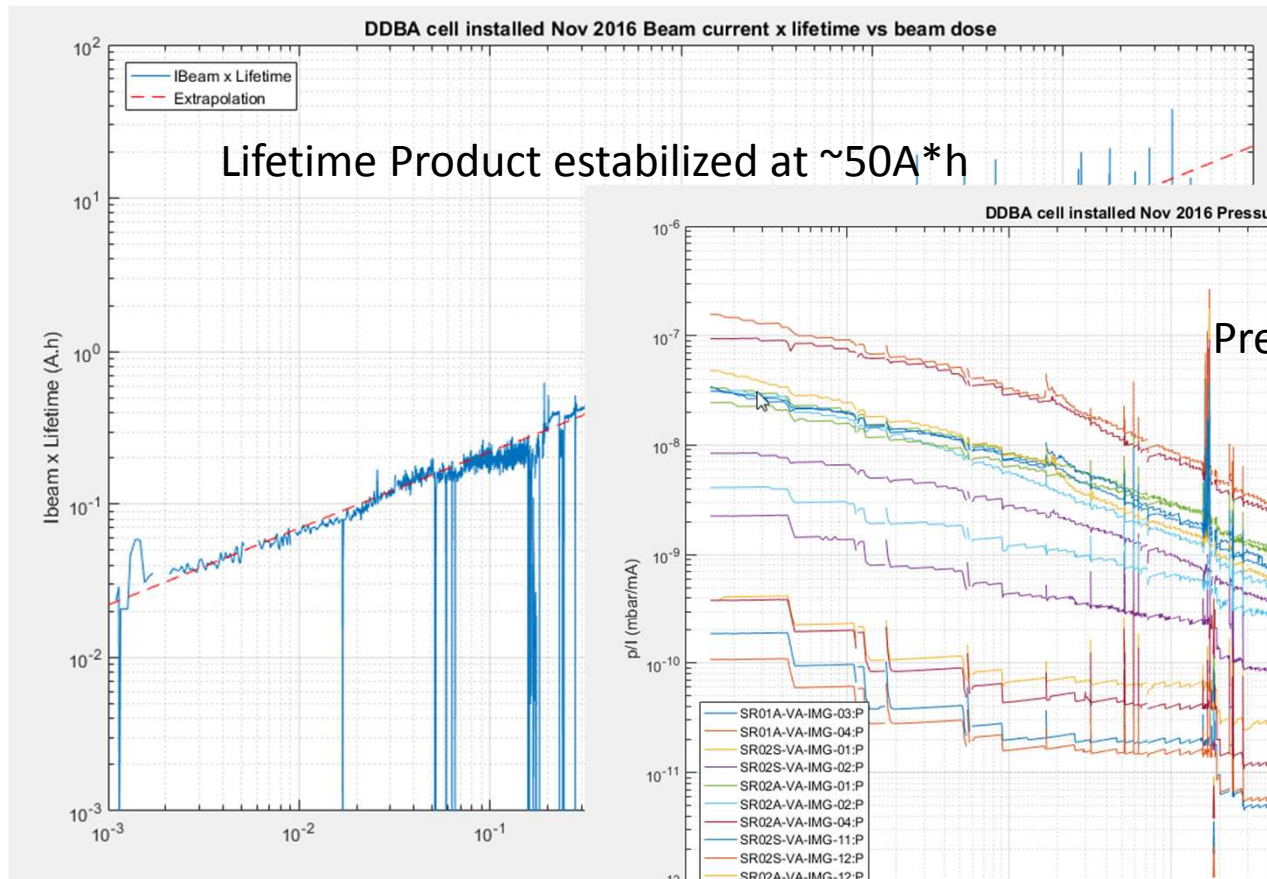
- ALBA: Raquel Monge, privet communication.
- Diamond: M P Cox et al, Commissioning of the diamond light source storage ring vacuum system, *Journal of Physics: Conference Series* 100 (2008) 092011
- Soleil: J.C. Besson, et al COMMISSIONING & OPERATION OF SOLEIL, WAO 2007. PSI - Scientific and Technical Report 2003 / Volume VI
- SLS: L. Schulz et al, STATUS REPORT OF THE SLS STORAGE RING VACUUM SYSTEM: EXPERIENCE AFTER TWO YEARS OF OPERATION
- ASP: E. Al-Dmour, VACUUM PERFORMANCE IN THE MOST RECENT THIRD GENERATION SYNCHROTRON LIGHT SOURCES, EPAC08.

E. Al-Dmour, XXIV SLS Workshop, 2016

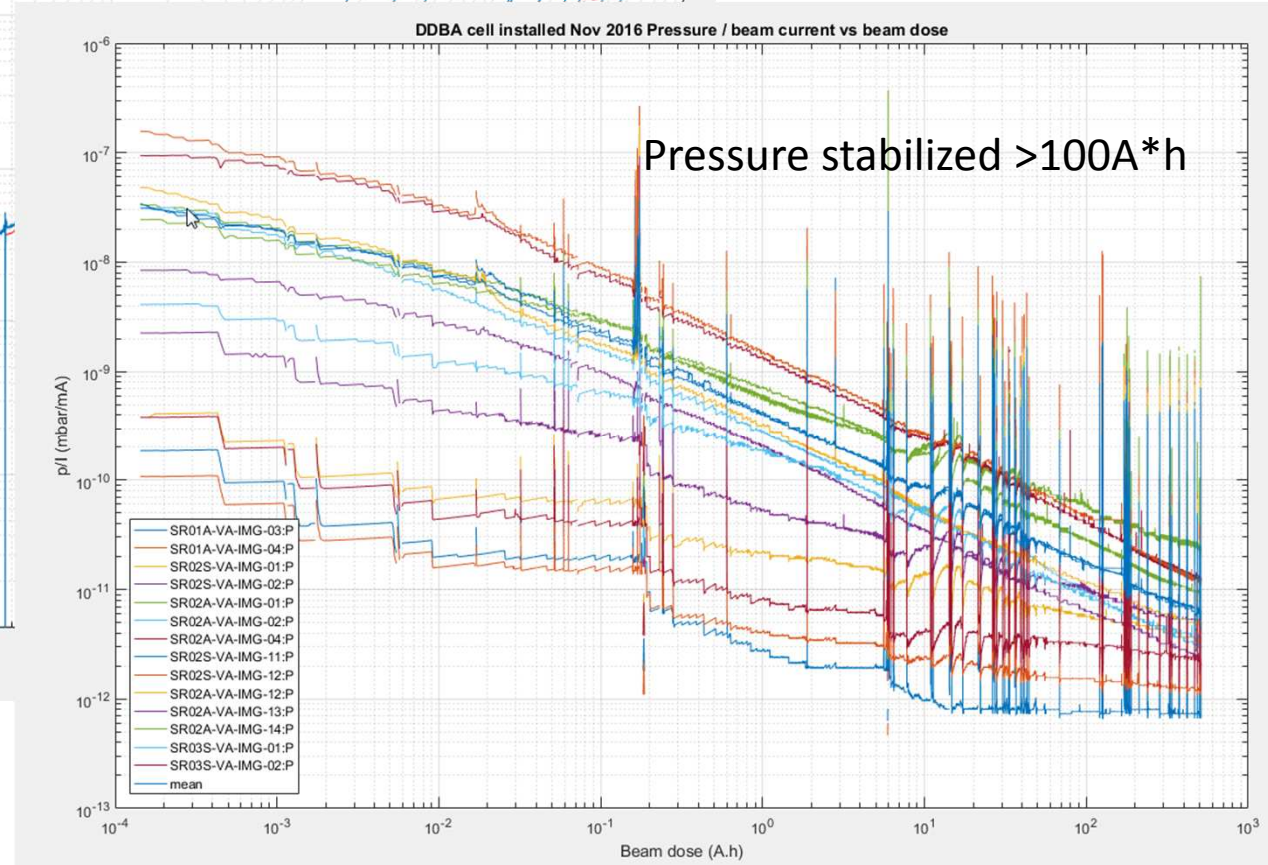


# Installation of new components

Example: Complete Cell exchange in Diamond

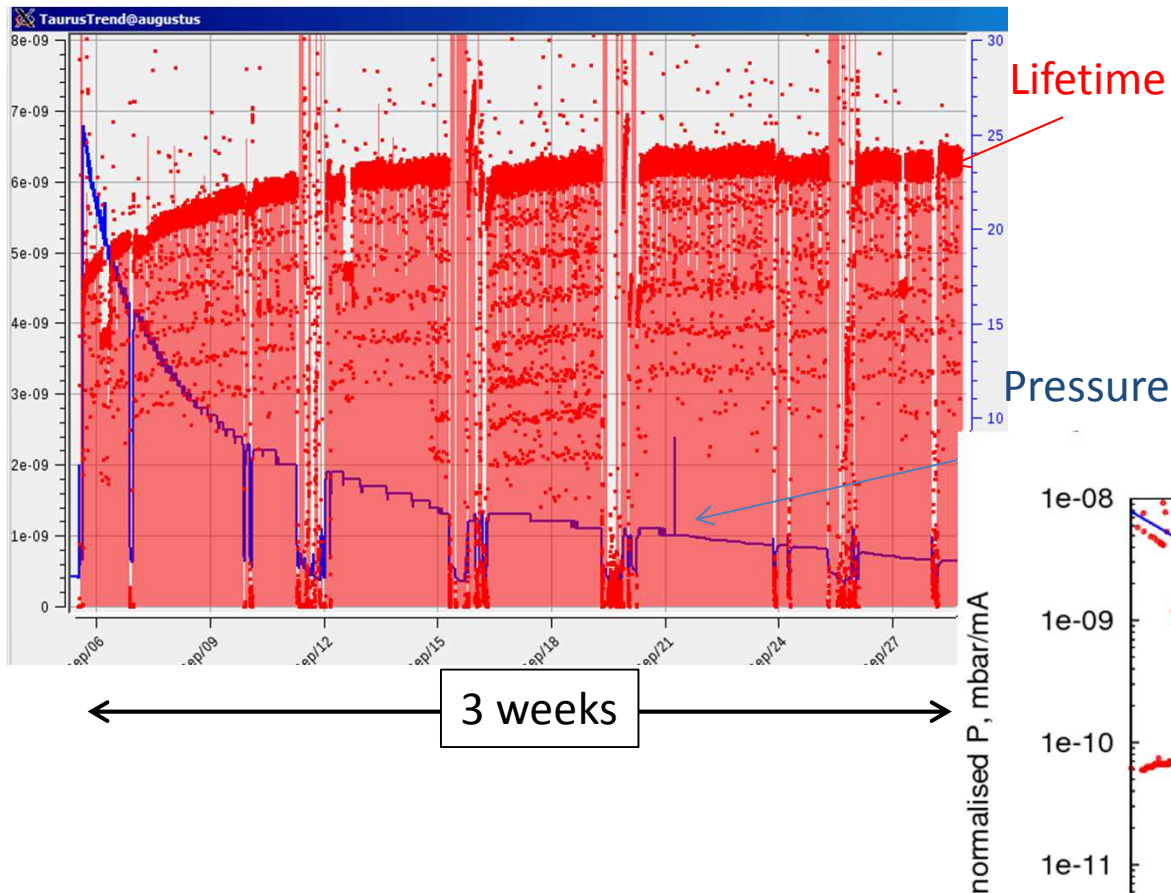


*Courtesy of M.Cox and I. Martin (Diamond)*

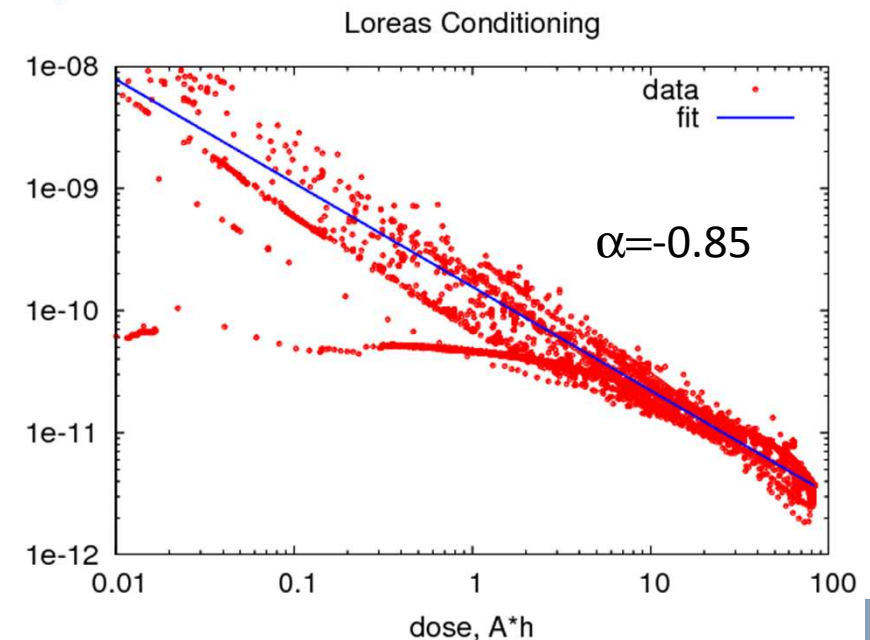


# Installation of new components

Example: New Insertion Device (Loreas) @ALBA



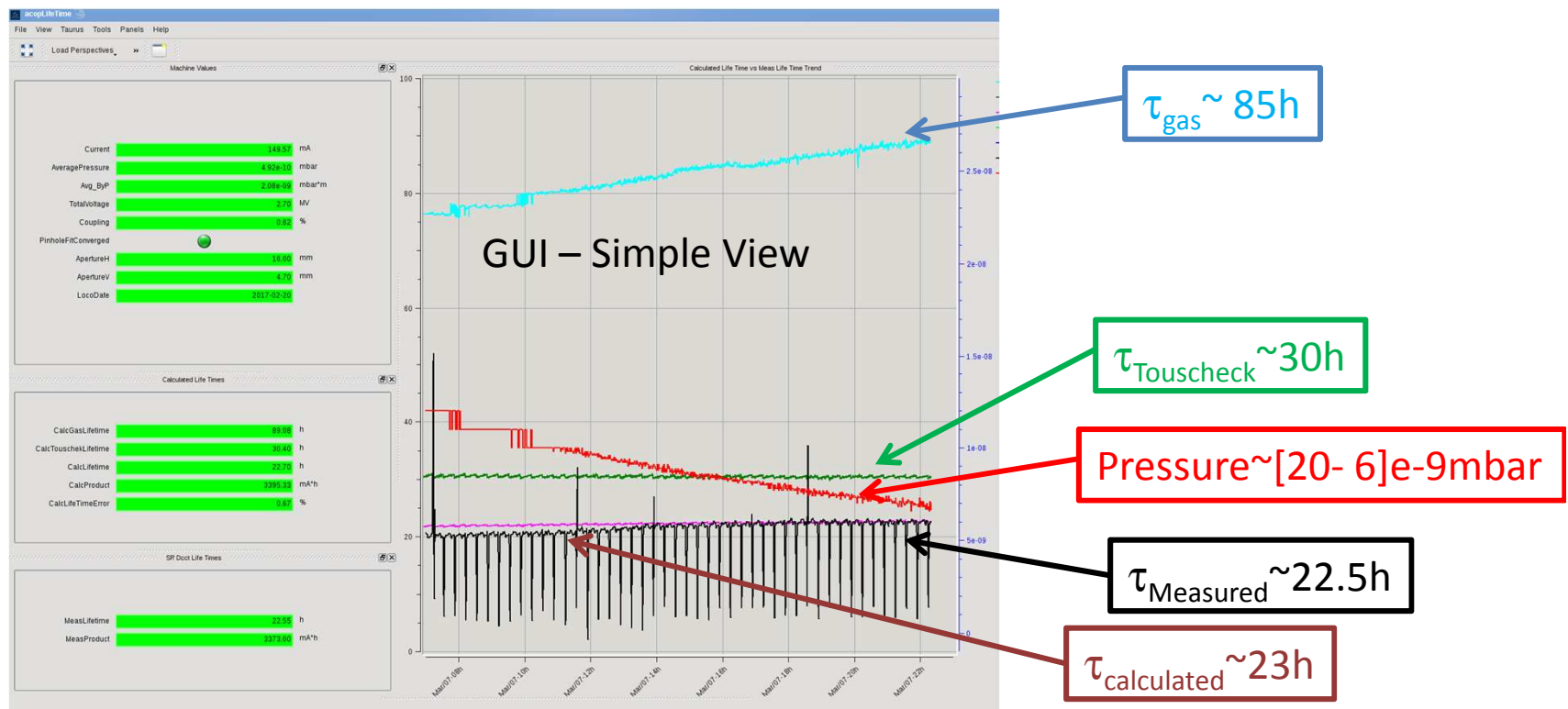
- ~3 weeks to decrease pressure  
~1order of magnitude
- ~2 weeks to stabilize lifetime  
(from 18h to 24h)
- Accumulated Dose: 80A\*h



# Lifetime Calculator GUI

- Used at ALBA Control Room to crosscheck machine performance
- $P@CCG$ ,  $\beta_{x,y}$  functions,  $\varepsilon_{xy}$ , cross sections, etc calculated on-line --> also  $\tau_{gas}$  &  $\tau_{Touscheck}$

Example during conditioning of new vac. chamber



Developed by M.Alvarez





# Lifetime Calculator GUI

- Used at ALBA Control Room to crosscheck machine performance
- $P@CCG$ ,  $\beta_{x,y}$  functions,  $\varepsilon_{xy}$ , cross sections, etc calculated on-line --> also  $\tau_{gas}$  &  $\tau_{Touscheck}$

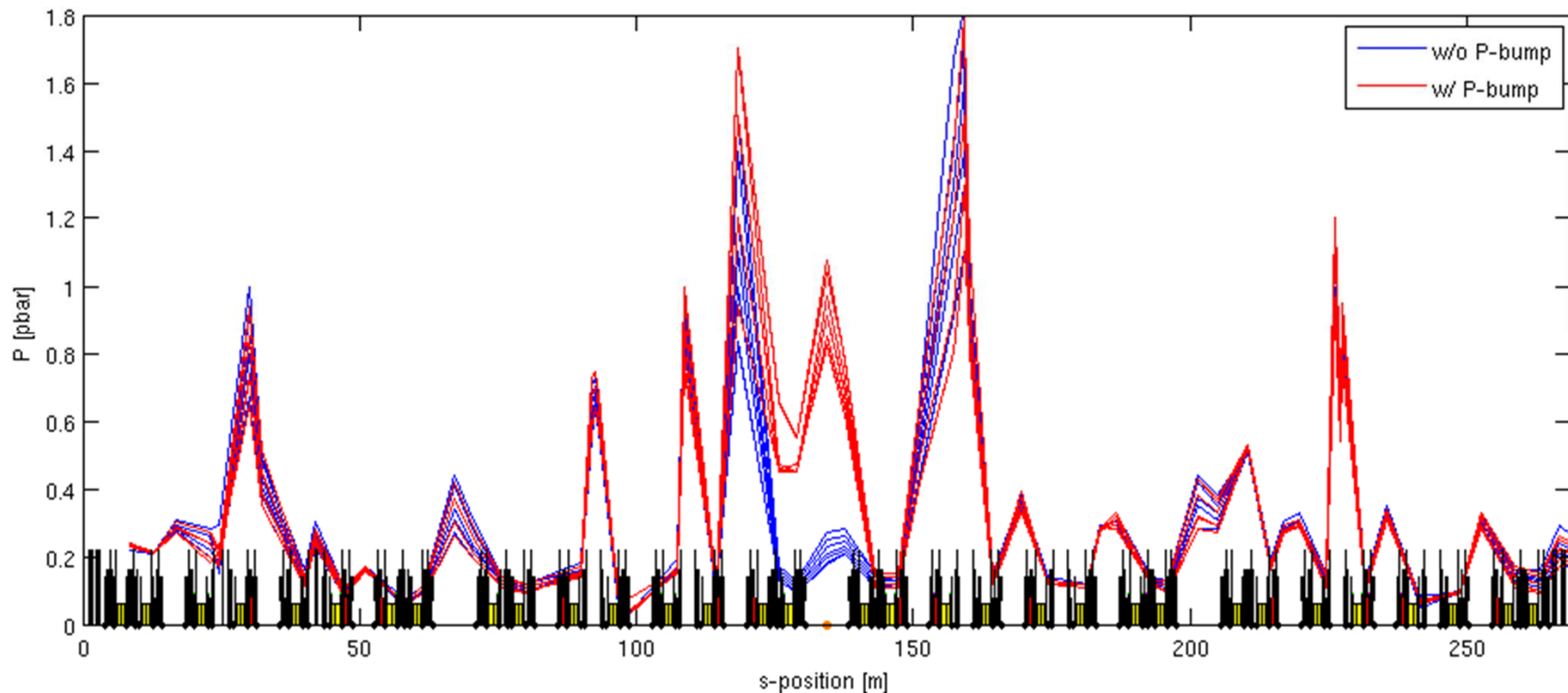
GUI – Expert View: allows you to control/measure all lifetime related params



Developed by M.Alvarez

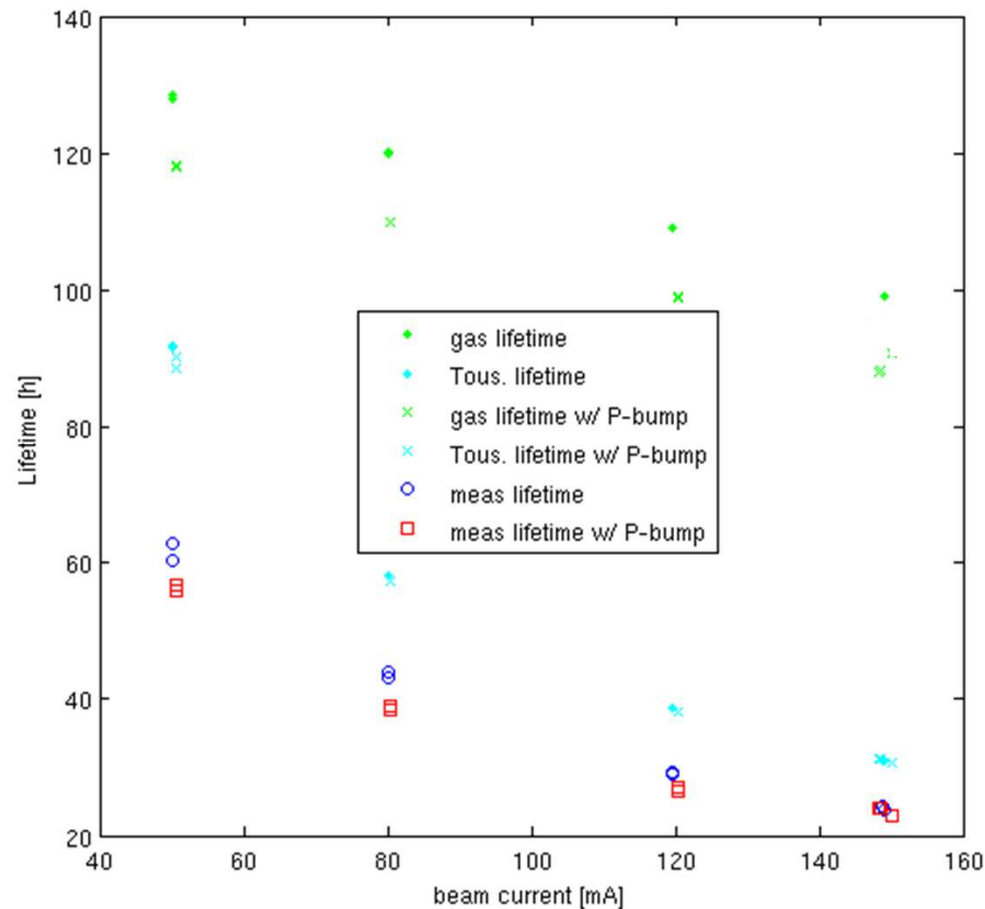
# Controlled pressure bump

- Experiment goal: change gas lifetime w.o. varying Touscheck
- Switch off Ion Pumps in one sector (out of 16).
- NEG pumping on  $\rightarrow$  pressure increase  $\sim 1$  order of magnitude in the Cell



# Controlled pressure bump

Lifetime comparison at different beam currents



At 50mA, the lifetime difference is only 10% .

At 150mA, the differences are barely noticeable, since Touscheck lifetime dominates



# OUTLINE

## 1. Beam Lifetime

- a) Analytical description
- b) Observations in Light Sources

## 2. Ion Instabilities

- a) Analytical Description
- b) Observations in Light Sources

## 3. Conclusions

# Ion Instabilities in LS

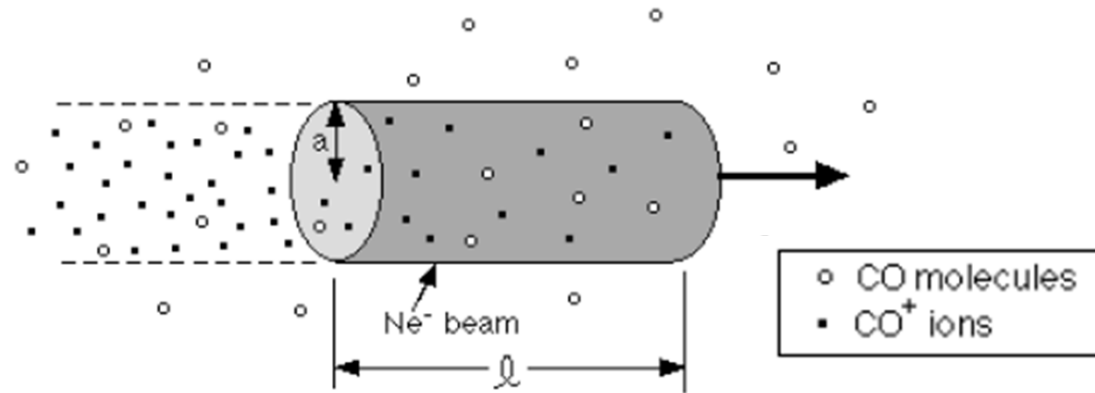
See G. Rumolo and R.  
Nagaoka talks

- Rest gas ionization produced by e-beam generates ions inside vac. chamber
- Ions are heavy particles and they usually feel only a sequence of attractive kicks from the bunches, which keep them confined near the beam core
- If the ion trapping condition is fulfilled, the ion remains oscillating around the beam
- The Ion Instabilities can be either:
  - “Ion Trapping Instability” (**ITI**): multi-turn ion accumulation that degrades beam quality (cured typically with abort gap)
  - “Fast Beam Ion Instability” (**FBII**), the ion production & accumulation occur within one turn, affecting mainly last bunches in the train.
- Main consequences: **emittance growth, tune shift, pressure rise...**



# Ion Instability: ion generation

- Rest gas ionization produced by e-beam generates ions inside vac. chamber



- Scattering ionization (depends on cross section  $\sigma_{\text{ion}}$ )

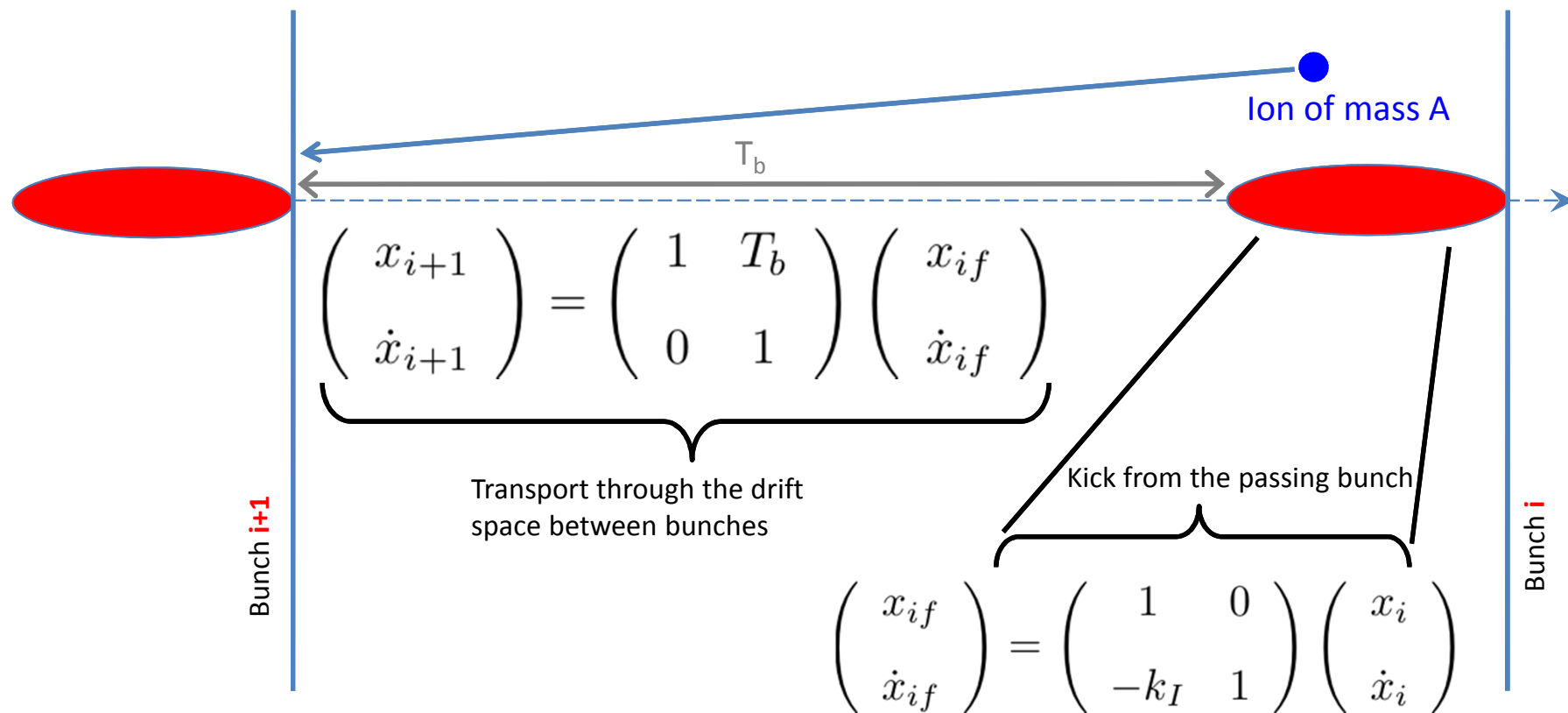
$$\lambda = \frac{N_b}{k_B T} \sum_{n=1}^N P_n \sigma_n \quad \sim \mathbf{20 \text{ ions/bunch/m}} \text{ (ALBA, } P=1\text{pbar, } \sigma_{\text{ion}} = 2\text{MB} - \text{CO})$$

- Field ionization could also happen above a certain threshold, but negligible in current LS – may be relevant for CLIC Linacs

# Ion Instability: Trapping Condition

$$\Delta \dot{x} = -\frac{2N_b r_p c}{A\sigma_x(\sigma_x + \sigma_y)}(x - \langle x \rangle) = -k_I(x - \langle x \rangle)$$

G. Rumolo - "Two stream Instabilities" – USPAS



# Ion Instability: Trapping Condition

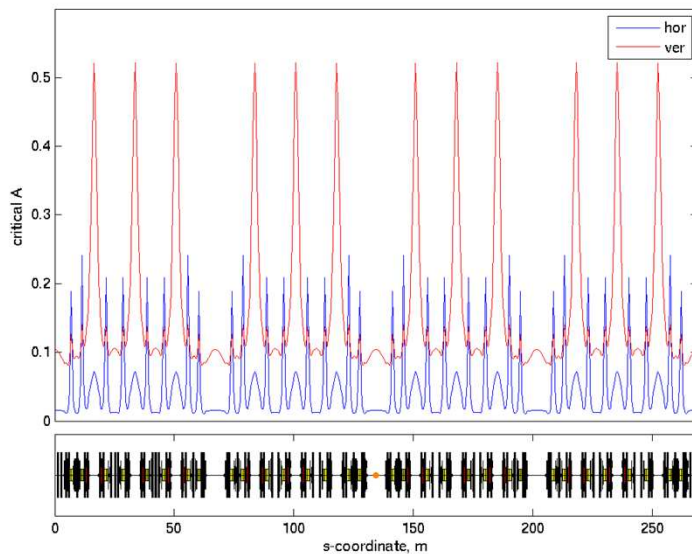
Ion Motion easily expressed in matrix notation:

$$\begin{pmatrix} x_{i+1} \\ \dot{x}_{i+1} \end{pmatrix} = \begin{pmatrix} 1 - k_I T_b & T_b \\ -k_I & 1 \end{pmatrix} \begin{pmatrix} x_i \\ \dot{x}_i \end{pmatrix} = \underline{\underline{\mathbf{A}}} \cdot \begin{pmatrix} x_i \\ \dot{x}_i \end{pmatrix}$$

Stability if  $\text{Tr}(\mathbf{A}) < 2$ :

$$A_{crit} > \frac{N_B r_p s_B}{2\sigma_y(\sigma_x + \sigma_y)}$$

- Assumed Gaussian beams
- Ions with atomic mass  $> A_{crit}$  are trapped!



ALBA @150mA:

- All ions can be trapped!!
- It does not affect ALBA due to use of abort gaps and low pressure ( $\sim 1e-10$  mbar)
- Dependence on lattice position through  $(\sigma_x, \sigma_y)$

# Ion Instability: Trapping Condition

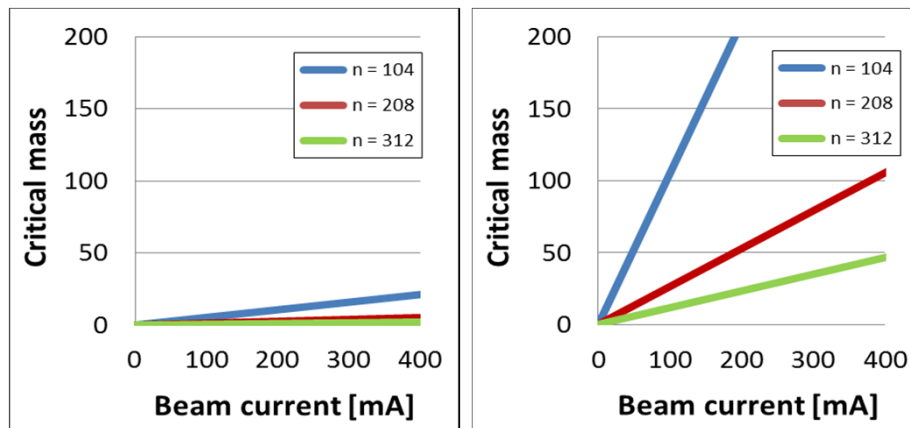
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$$A_{crit} > \frac{N_B r_p s_B}{2\sigma_y(\sigma_x + \sigma_y)}$$

- Assumed Gaussian beams
- Ions with atomic mass  $> A_{crit}$  are trapped!



Calculated with  $\varepsilon_x = 4$  nm (left) and  $0.2$  nm (right) with SOLEIL parameters (1% coupling)

R. Nagaoka talk

- For low-e rings,  $A_{crit}$  increases

# Ion Instability: Osc. Frequency

From 
$$\Delta \dot{x} = -\frac{2N_b r_p c}{A\sigma_x(\sigma_x + \sigma_y)} (x_i - \langle x \rangle) \quad \text{and} \quad \ddot{x} = \Delta x / T_B$$

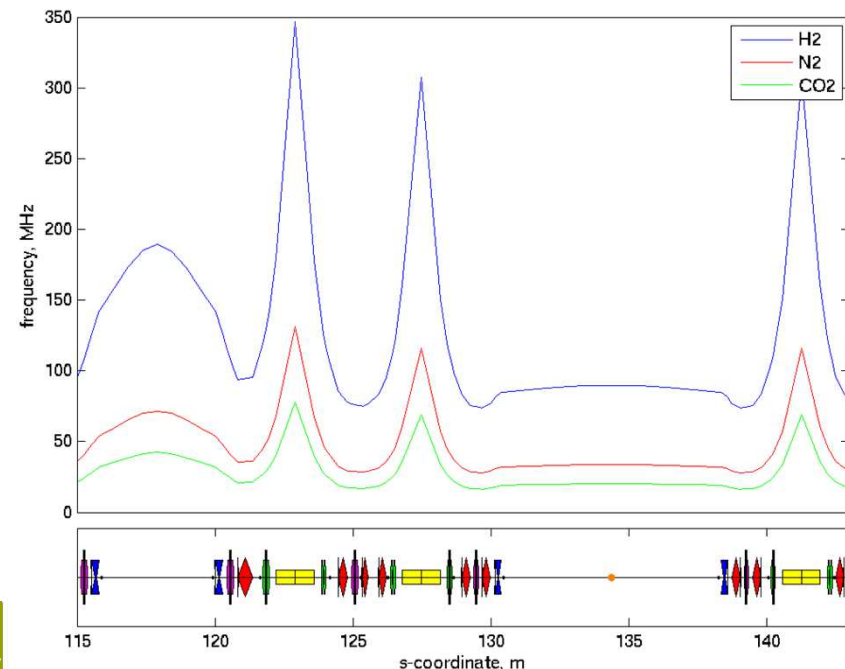
We get the ion **oscillation frequency**:

$$\omega^2_{ion} = \frac{2N_B c r_p}{A\sigma_y(\sigma_x + \sigma_y)T_B}$$

observable with spectrun analyser

ALBA Case @150mA in one sector:

- Inversely proportional to A
- Between  $\sim [10 - 300\text{MHz}]$
- Also depends on s-position



# Ion Instabilities: tune shift

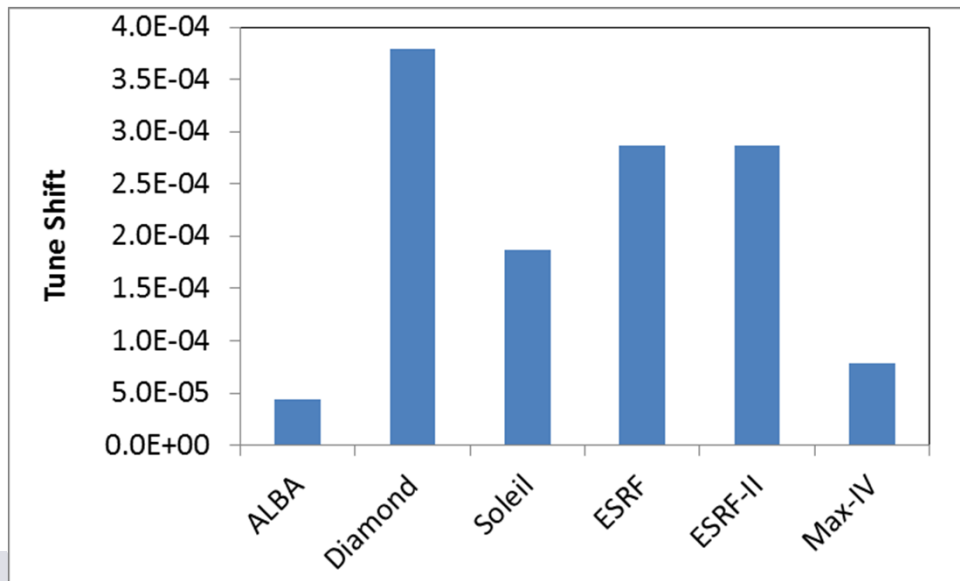
Induced tune shift due to focusing force:

$$\Delta\nu_\beta = + \frac{I_p r_e \eta}{2\pi e c (\beta_c)^3 \gamma B_f} \oint \frac{\beta_y ds}{\sigma_y (\sigma_x + \sigma_y)}$$

An easier simplification from above equation is:  $\Delta\nu_y = \frac{r_e \lambda_{ion} C \langle \beta_y \rangle}{4\pi \gamma \sigma_y (\sigma_x + \sigma_y)}$

Note:

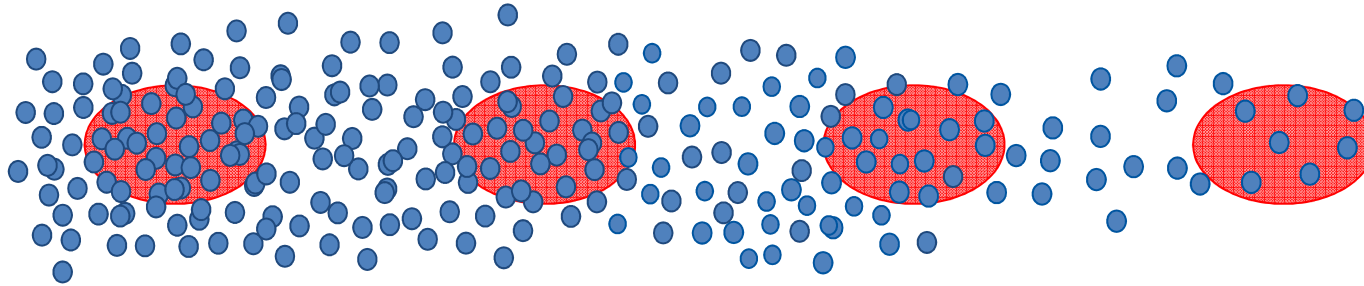
- positive tune shift
- For FBII,  $\lambda_{ion}$  depends on the bunch #, so the tune shift depends on the bunch 3



CO ions; P=1pbar;  $\sigma_{ion} = 2\text{MB}$   
Nominal operation conditions

Tune shift small, but measurable

# Fast Beam Ion Instability - FBII



- The ions accumulate along one bunch train
- Head and tail of the train are coupled through the ions
- Coupled motion between ions & e-beam: ions “keep memory” of the offset of the generating bunch and transfer this information to the following bunches.
- The driven oscillation is expected to be at a main frequency related to the ion oscillation frequency.

See analytical models at:

Raubenheimer *et al.* Phys. Rev. E 52, 5, 5487,  
Stupakov *et al.* Phys. Rev. E 52, 5, 5499

# Fast Beam Ion Instability - FBII

Analytical instability models\* allow to estimate rise times:

$$\frac{1}{\tau_e} \approx \frac{1}{3} \sqrt{\frac{2}{3}} \frac{c}{\Delta\omega_i/\omega_i} \beta_y k_y \quad \text{With } K_y \text{ the focusing strength} \quad k_y = \frac{\lambda_i r_e}{\gamma \sigma_y (\sigma_x + \sigma_y)}$$

- This gives rise times ~**2us** range at ALBA, although coherence effects along the ring can increase this rise times by 1-2 orders of magnitude
- Each case needs a careful analysis
- Computer simulation codes (FASTION ad PyHEADTAIL\*\*) can be used to carefully calculate these rise times.

\*Raubenheimer *et al.* Phys. Rev. E 52, 5, 5487,  
Stupakov *et al.* Phys. Rev. E 52, 5, 5499

\*\*L.Mether, "Numerical model of FBII", Proc. of HB2016  
A.Chatterjee et al, PRST-AB 18 064402 (2015)

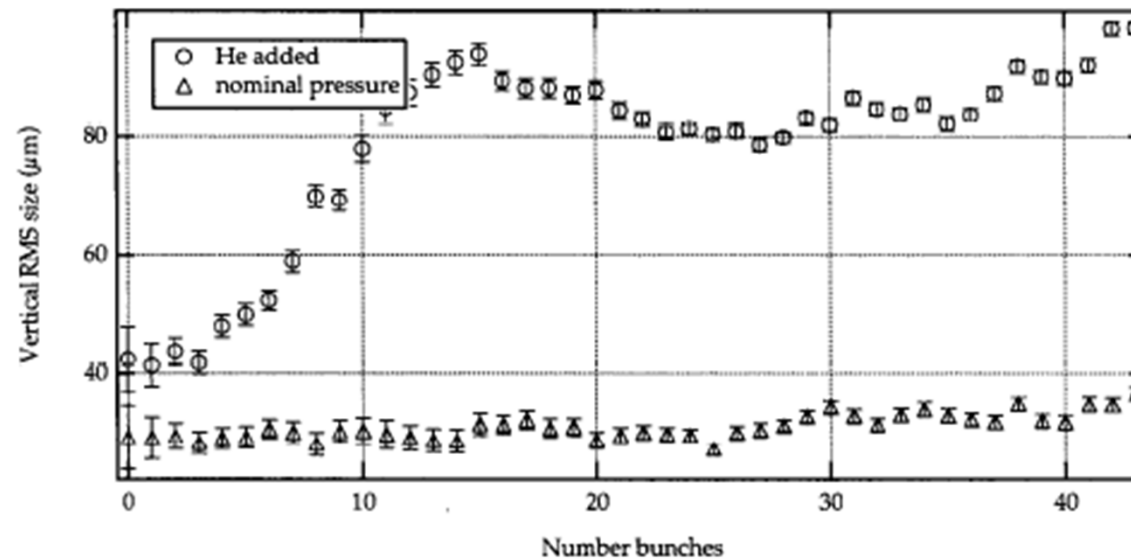


# Ion Instabilities Observations

- **Fast Beam Ion Instabilities** observed in LS are produced artificially induced by injecting gas into the vacuum chamber
- See examples at **ALS & CESR**

## ALS experiment\*:

- all Ion Pumps off
- Inject He: avg P from 0.25 to 80 nTorr



\*J.Byrd et al, PRL 79, 1 (1997)

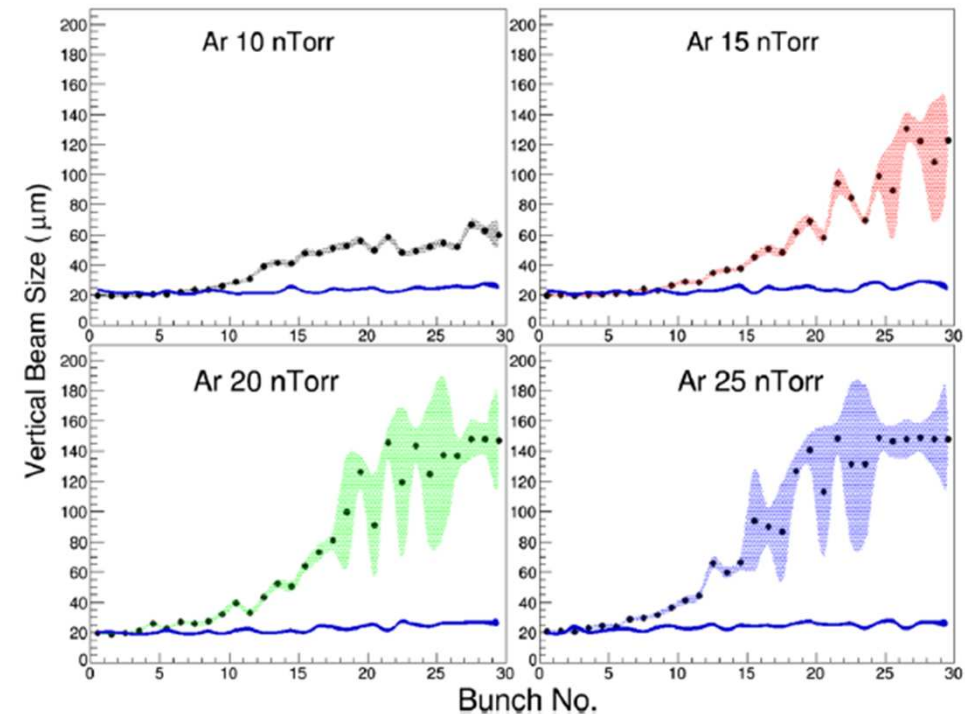
# Ion Instabilities Observations

- **Fast Beam Ion Instabilities** observed in LS are produced artificially induced by injecting gas into the vacuum chamber
- See examples at **ALS & CESR**

## CESR experiment\*:

- Localized pressure bump of  $\sim 10$  m
- Inject Ar and Kr: from 1 to 25 nTorr

- Beam Size observation
- BBB feedback can damp instability
- Observations match well with simulations using FASTION

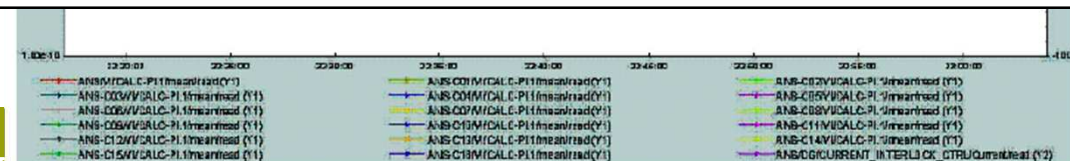


\*A.Chatterjee et al, PRST-AB 18 064402 (2015)

## R. Nagaoka, IPAC'10

- 

➔ Not Ion Trapping, but FBII



# Ion Instabilities Observations

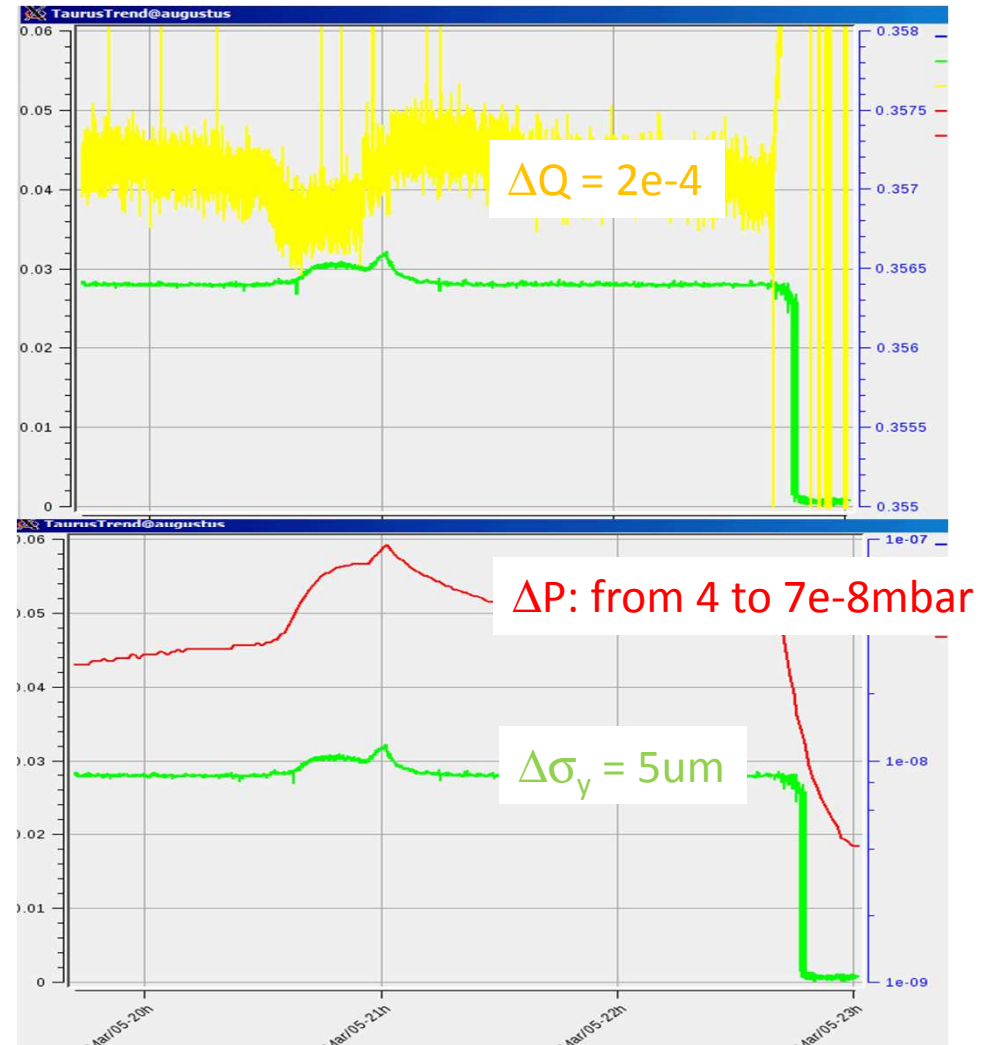
## Installation of CLIC SL Kicker @ALBA

(Data from 6/3/2017)

While conditioning the new vacuum chamber, injecting from **125mA** to **135mA** produced a:

- pressure rise:  $\sim 3e-8$ mbar
- beam blow up (by 20%)
- tune shift of  $2e-4$
- BBB feedback could not damp it

Compatible with ion instabilities in this case produced by larger atomic masses (see next)



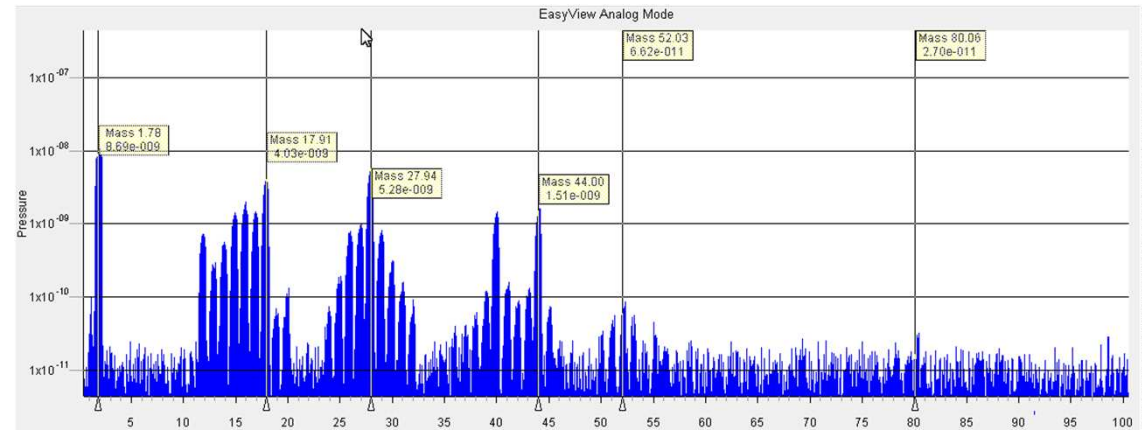
# Ion Instabilities Observations

## Installation of CLIC SL Kicker @ALBA

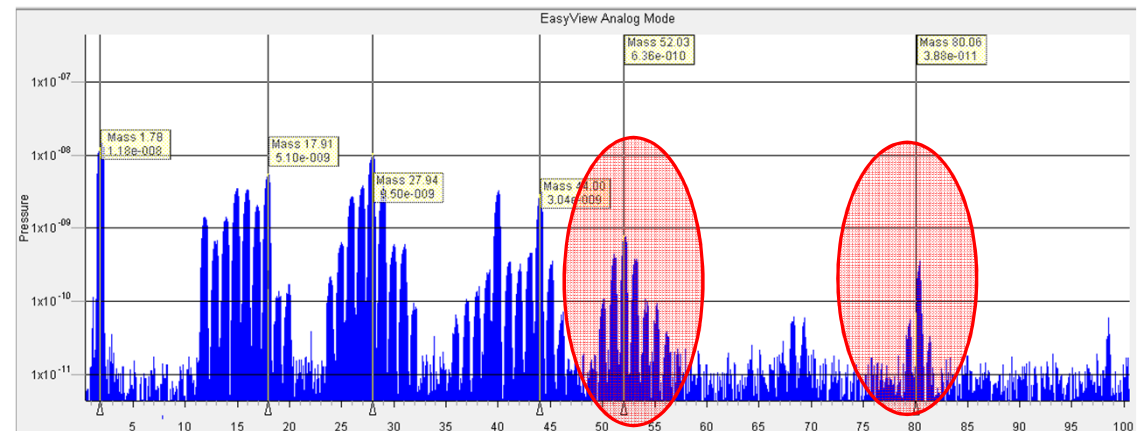
(Data from 6/3/2017)

- RGA data showed peaks at A=52 and A=80, likely due to desorption produced by Macor rings
- CLIC SL removed to guarantee operation

125 mA

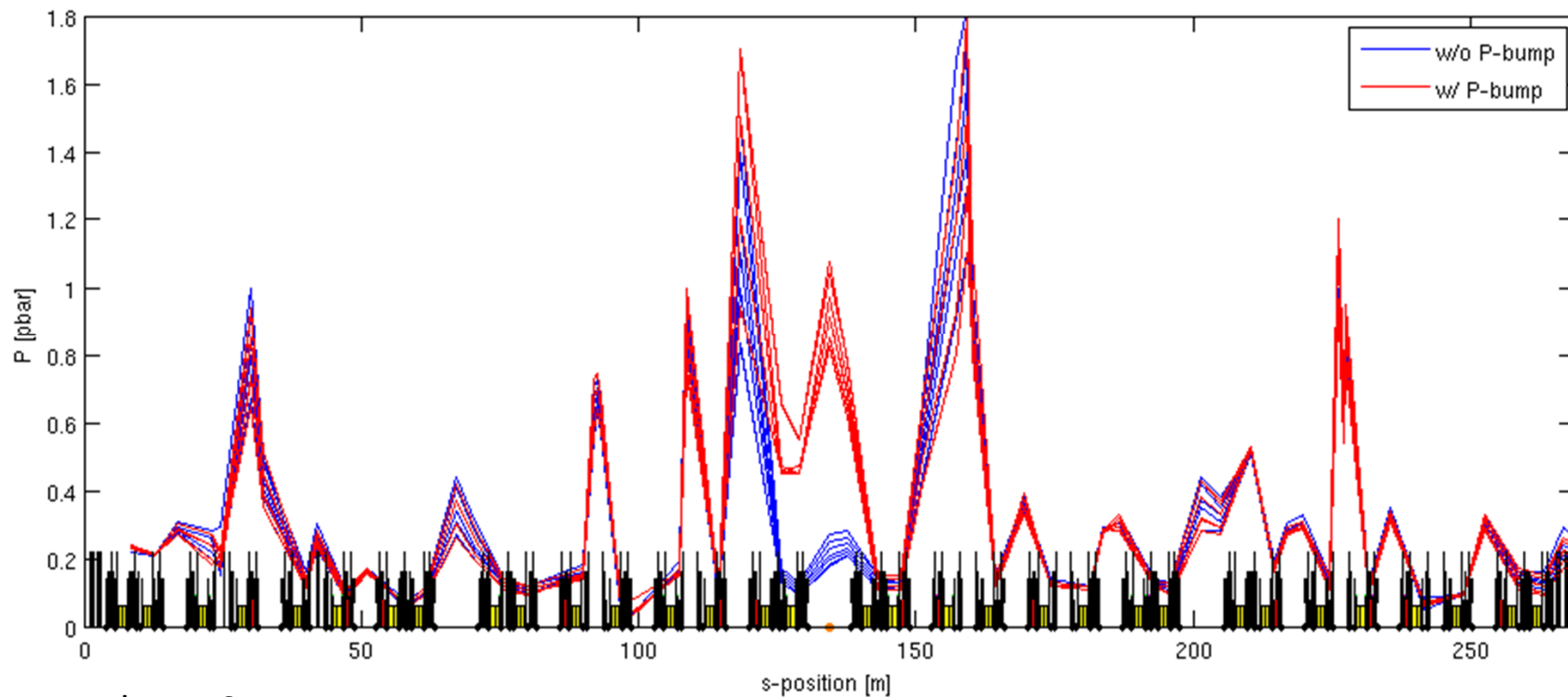


135 mA



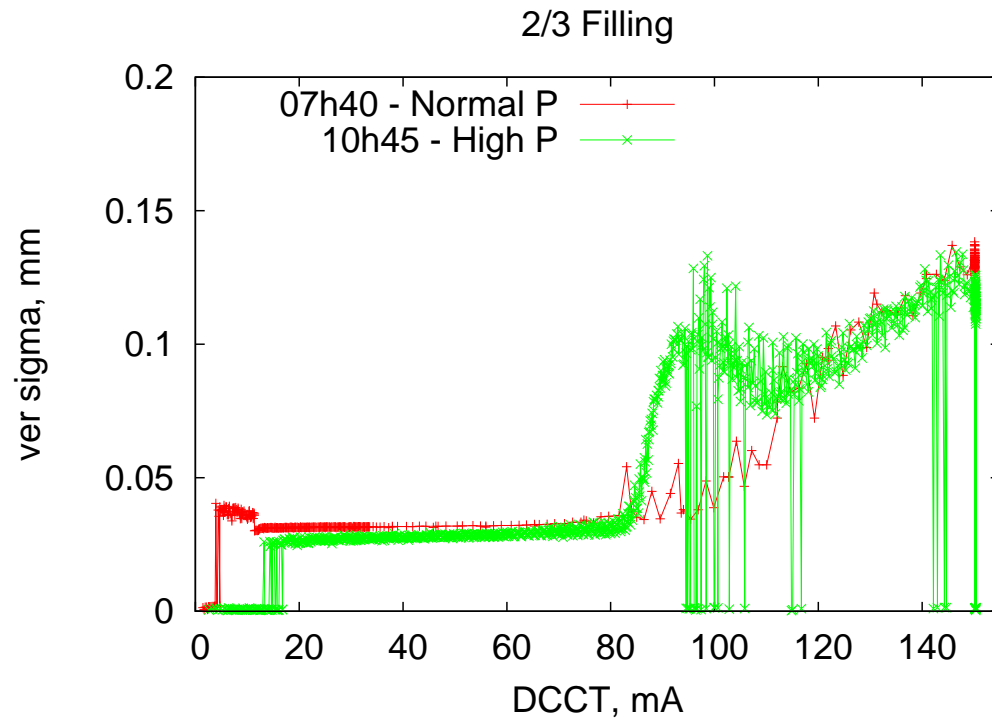
# Controlled pressure bump

- Switch off Ion Pumps in one sector (out of 16).
- NEG pumping on  $\rightarrow$  pressure increase  $\sim 1$  order of magnitude in the Cell
- Check instability thresholds w. & w.o. pressure bumps



# Controlled pressure bump

→ Beam size increase with increasing current with & w.o. pressure bump



- $\text{Xi} = (1,2)$ , threshold  $\sim 80\text{mA}$  for both cases
- But for normal pressure, the emittance increase is much violent. Co-existence of RW and Ions?
- At  $\sim 120\text{mA}$ , beam size follows the same trend as the case with normal P (larger Nb, ions “untrapped”?)
- Further studies going-on



# OUTLINE

## 1. Beam Lifetime

- a) Theoretical Equations
- b) Impact on current Light Sources

## 2. Ion Instabilities

- a) Analytical Equations
- b) Observations in Light Sources

## 3. Conclusions



# CONCLUSIONS

## As for the lifetime...

- Lifetime in LS is limited by Touscheck effect for machines with doses  $>100 \text{ A}^*\text{h}$
- Gas lifetime (elastic and inelastic)  $\sim 100\text{h}$ , while Touscheck lifetime  $\sim 20\text{h}$ .
- Gas lifetime only relevant in commissioning periods, or after installation of new components
- As a rule of thumb, need in  $\sim 100\text{A}^*\text{h}$  to recover  $P < 1\text{e-}9\text{mbar}$  ( $0.75\text{nTorr}$ )

# CONCLUSIONS

## As for the ion instabilities...

- Ion trapping condition is easily fulfilled in LS. However, observations of ion instabilities are rare due to a very good vacuum conditions ( $P_{avg} < 5 \times 10^{-10} \text{ mbar}$ ).
- Ion instabilities are only relevant in commissioning periods or after installation of new components in the vacuum chamber
- FBII have only been seen in artificially bad vacuum conditions (injecting Ar gas) and/or switching off Ion Pumps to bring  $P > 10 \text{ nTorr}$
- Ion instabilities may co-exist with other instabilities (RW), and active feedback may effectively damp them (CESR) or not (Soleil)
- Soleil is (so far) the only machine whose operation is limited by ion instabilities (FBII).
- Simulation codes exist (FASTION & PyHeadtail), but more benchmarking is appreciated



# Acknowledgements

Many colleagues from ALBA and other labs contributed to this talk:

M. Alvarez, T. Guenzel, Z. Marti, R. Monge (ALBA)

I.Martin, M.Cox (Diamond)

N.Carmignani (ESRF)

X. Nuel (Soleil)

J. Sundberg (Max-IV)

G. Rumolo (CERN)

<http://inspirehep.net/record/846682/files/PhysRevSTAB.12.pdf>

S.C. Leeman, Phys. Rev. ST Accel. Beams 12, 120701 (2009)

TABLE VII. Contributions to the total MAX IV 3 GeV storage-ring lifetime  $\tau$ . The results have been calculated for a “worst-case scenario”: four PMDWs and ten IVUs are installed in the storage ring while the total applied rf voltage is 1.5 MV which corresponds to an rf acceptance of only  $\delta_{\text{rf}} = 4.0\%$ .

	$\tau$ [h]
Elastic gas scattering	25.4
Inelastic gas scattering	53.1
Touschek scattering (with Landau cavities)	25.5
Total	10.3



## MACOR® Machinable Glass Ceramic

- is MACHINABLE with ordinary metal working tools
- allows FAST TURNAROUND, no post firing required
- holds TIGHT TOLERANCES, up to .0005"
- withstands HIGH TEMPERATURE, up to 1000°C (no load)
- is CLEAN, no outgasing and zero porosity

### Composition

MACOR Machinable Glass Ceramic is a white, odorless, porcelain-like (in appearance) material composed of approximately 55% fluorophlogopite mica and 45% borosilicate glass. It has no known toxic effects; however, the dust created in machining can be an irritant. This irritation can be avoided by good housekeeping and appropriate machining techniques. The material contains the following compounds:

	Approximate Weight %
Silicon - $\text{SiO}_2$	46%
Magnesium - $\text{MgO}$	17%
Aluminum - $\text{Al}_2\text{O}_3$	16%
Potassium - $\text{K}_2\text{O}$	10%
Boron - $\text{B}_2\text{O}_3$	7%
Fluorine - F	4%

# Fast Beam Ion Instability - FBII

- Analytical model available\*
  - Instability rise time equations, with & w.o. ion freq. spread along s

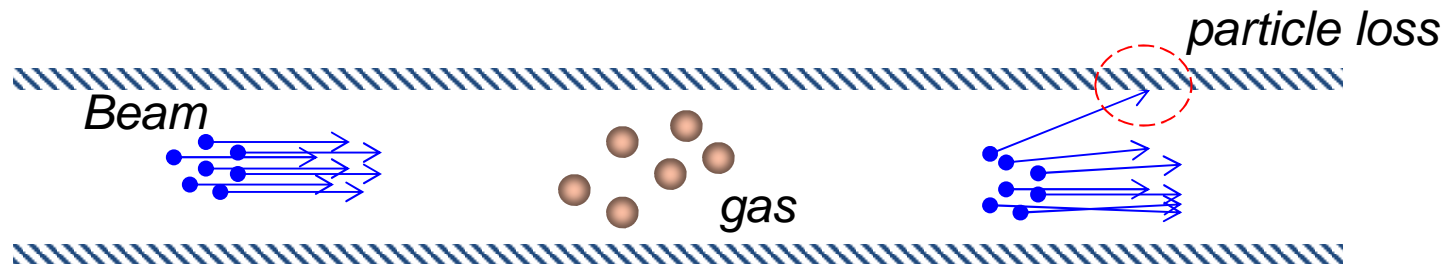
$$\frac{1}{\tau_F} = \frac{4d\sigma_{ion}\beta_y N_b^{3/2} n_b^2 r_p^{1/2} r_e s_b^{1/2} c}{3\sqrt{3}\gamma A^{1/2} \sigma_y^{3/2} (\sigma_x + \sigma_y)^{3/2}}$$

Lower  $\epsilon$  = faster instability

Raubenheimer *et al.* Phys. Rev. E 52, 5, 5487,  
Stupakov *et al.* Phys. Rev. E 52, 5, 5499

# Elastic Gas Scattering

Deflections caused by the **electromagnetic force** of the residual gas atoms. It can be due either by electrons or the nuclei (but the e- part is usually negligible for current LS)



Simplified Eqs:

$$\sigma_{el}(s) = 2\pi \frac{\beta_y(s)}{J_{y,\max}} \left( \frac{Zr_0}{\gamma} \right)^2$$

$$\tau_{el}^{-1} = \frac{c\beta}{kT} \langle P(s) \sigma_{el}(s) \rangle$$

Basic dependencies:

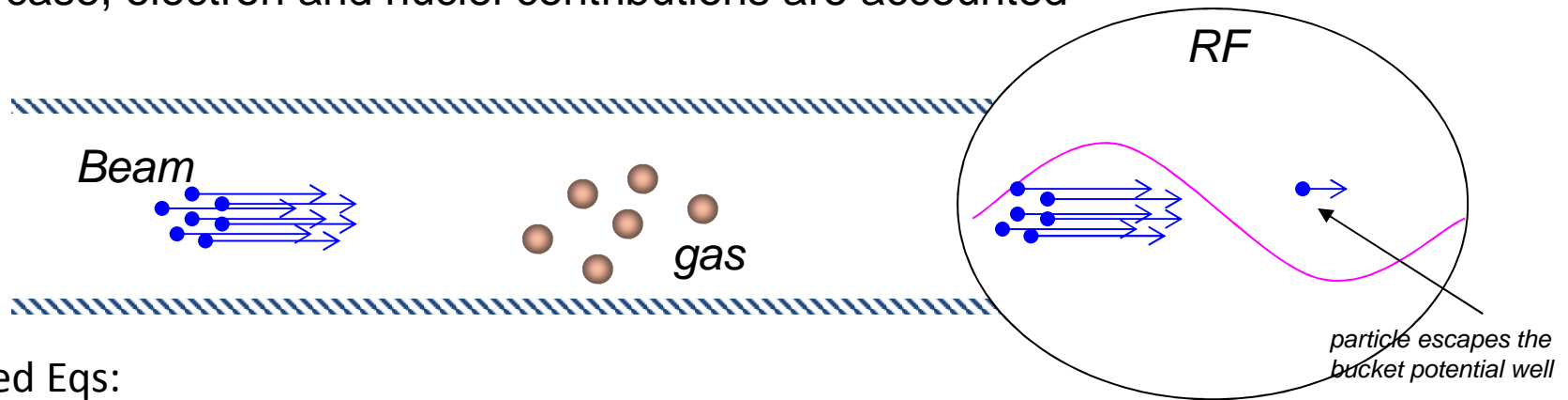
1. Number of residual particles (pressure):
2. Particle Charge:
3. Beam energy:
4. Beta function.
5. Beam size:

$$\begin{aligned} P \uparrow &\rightarrow \tau \downarrow \\ Z \uparrow &\rightarrow \tau \downarrow \\ \gamma \uparrow &\rightarrow \tau \uparrow \\ \beta \uparrow &\rightarrow \tau \downarrow \\ \sigma_y \uparrow &\rightarrow \tau \uparrow \end{aligned}$$

# Inelastic Gas Scattering

Energy loss caused by **radiation emission (mostly)** at the vicinity of the gas atoms.

In this case, electron and nuclei contributions are accounted



Simplified Eqs:

$$\sigma_{in}(s) = 4\alpha r_0^2 \left\{ \left[ \frac{4}{3} (\delta_{acc}(s) - \ln(\delta_{acc}(s))) + \frac{\delta_{acc}^2(s)}{2} \right] F(Z) + \frac{Z(Z+1)}{9} (\delta_{acc}(s) - \ln(\delta_{acc}(s))) \right\}$$

$$F(Z) = Z^2 \ln\left(\frac{183}{Z^{1/3}}\right) + Z \ln\left(\frac{1194}{Z^{1/3}}\right)$$

$$\tau_{in}^{-1} = \frac{c\beta}{kT} \langle P(s) \sigma_{in}(s) \rangle$$

Basic dependencies:

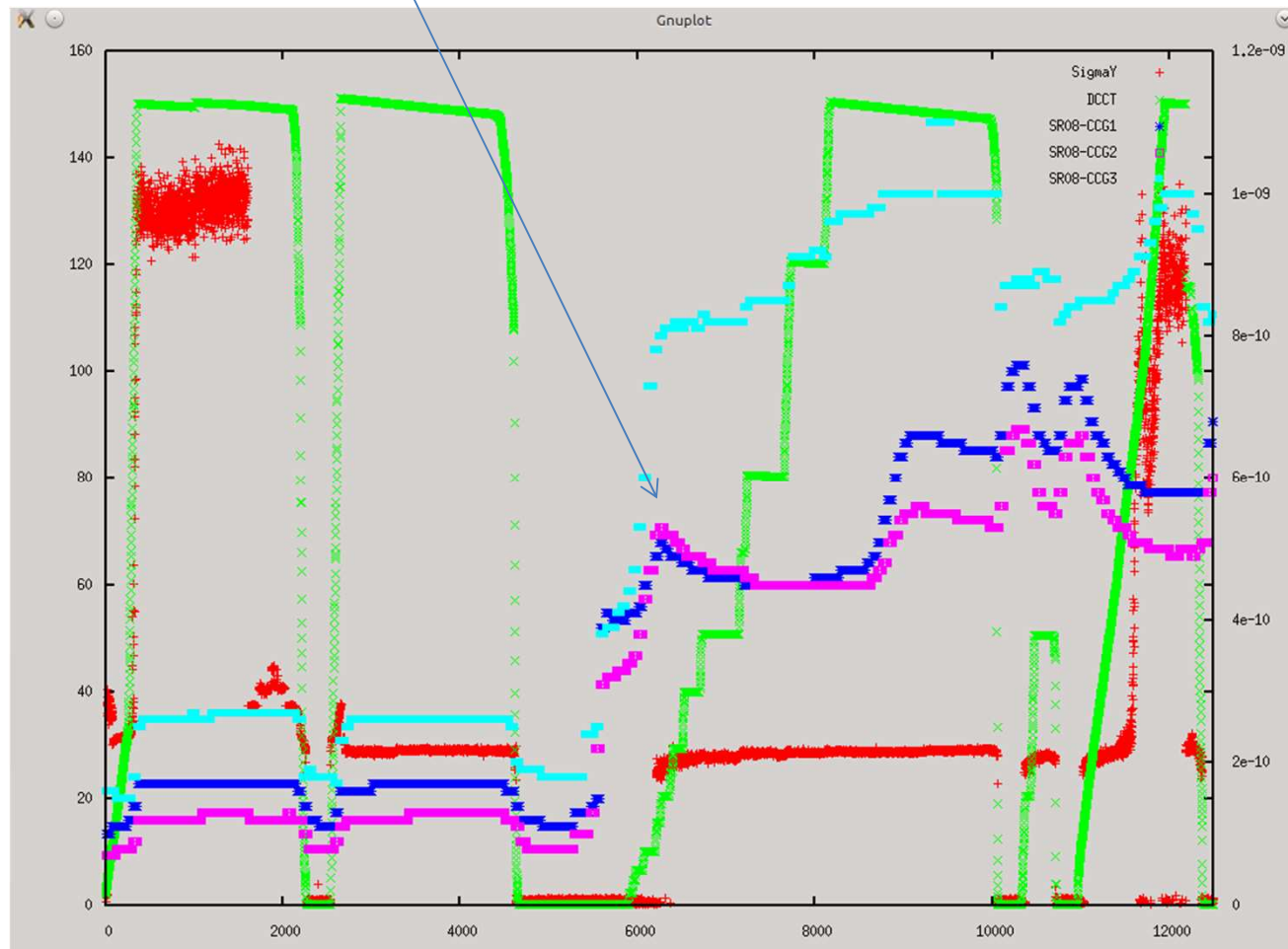
1. Number of rest gas atoms:
2. Charge of every particle:
3. RF acceptance.

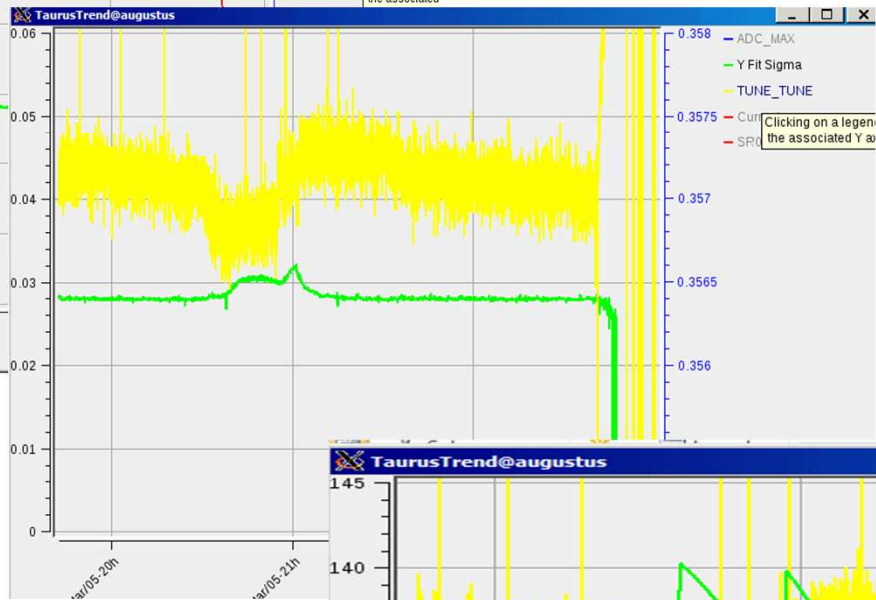
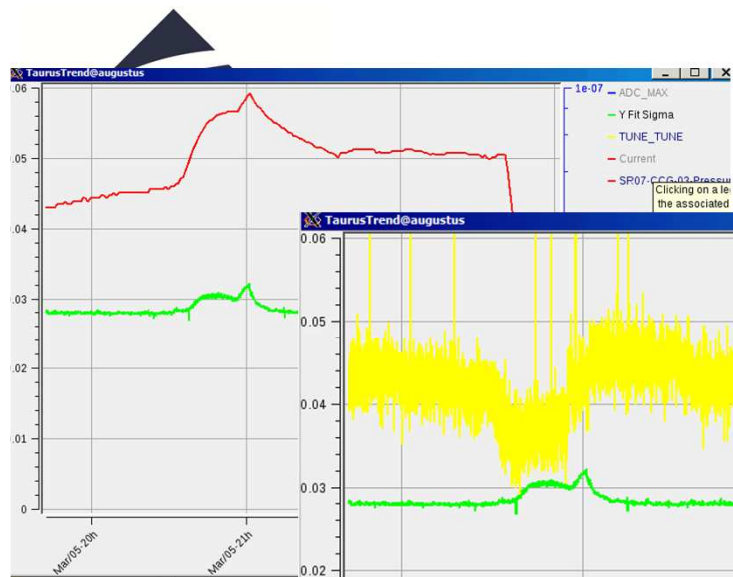
$$\begin{aligned} P \uparrow &\rightarrow \tau \downarrow \\ Z \uparrow &\rightarrow \tau \downarrow \\ \delta_{acc} \uparrow &\rightarrow \tau \uparrow \end{aligned}$$



Only SR08-CCG3 behaves “as expected” in the whole range.

For CCG1 and 2, after switching off certain pumps, there is a point in which both decrease even though the current increases: this (may) mean that after a certain P, you start the formation of a significant amount of gas ionization (or e-) which affects the pressure readings.





U. Iriso

B.Dynamics vs Vacuum

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