Direct observation of gravitational waves from the merger and inspiral of two black holes

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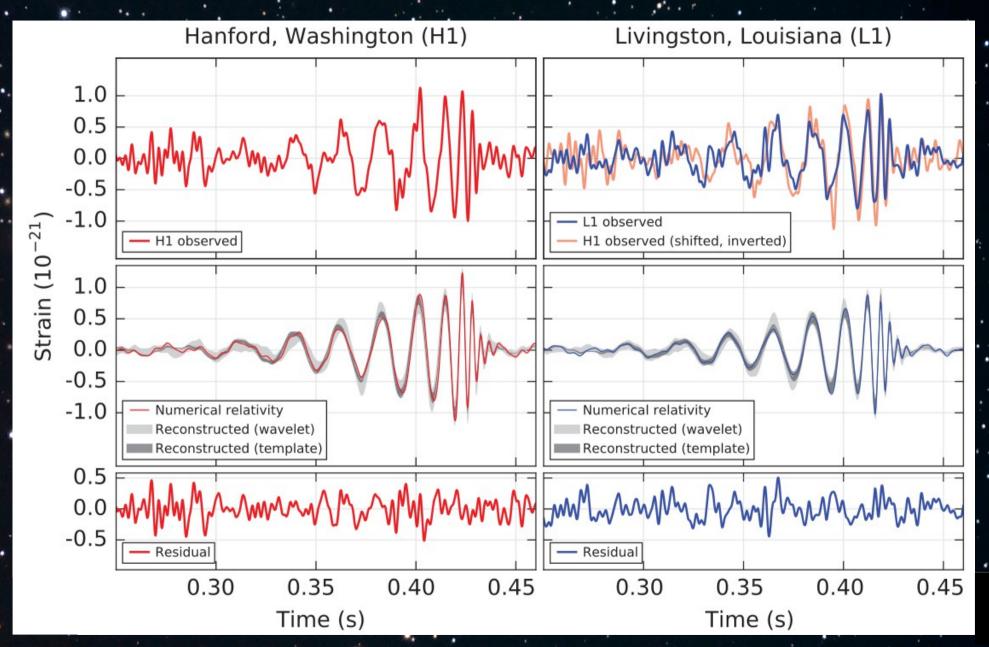
15th November 2016







What was seen 2015



Gravitational waves

What are gravitational waves?

Why are they detectable now?

What have we learnt?

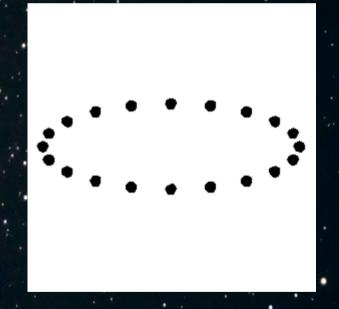
• Where are we going in the future?

What are gravitational waves?

$$G_{ab} = \frac{8 \pi G}{c^4} T_{ab}$$
 Einstein equation

Small linear perturbation $g_{ab} = \eta_{ab} + h_{ab}$

$$\nabla^2 \overline{h}_{ab} = 0$$
 Wave equation



$$Q_{ij} \equiv \int d^3x \, \rho \left(x_i x_j - \frac{1}{3} r^2 \delta_{ij} \right)$$

Einstein quadrupole formula

$$h_{ij} = \frac{2}{d_L} \frac{G}{c^4} \frac{d^2 Q_{ij}}{dt^2}$$

Einstein 1916

8 Sitzung der physikalisch-mathematischen Klasse vom 22. Juni 1916

Näherungsweise Integration der Feldgleichungen der Gravitation.

Von A. EINSTEIN.

Bei der Behandlung der meisten speziellen (nicht prinzipiellen) Probleme auf dem Gebiete der Gravitationstheorie kann man sich damit begnügen, die g_{sv} in erster Näherung zu berechnen. Dabei bedient man sich mit Vorteil der imaginären Zeitvariable $x_t = it$ aus denselben Gründen wie in der speziellen Relativitätstheorie. Unter »erster Näherung« ist dabei verstanden, daß die durch die Gleichung

$$q_{--} = -\delta_{--} + \gamma_{--} \tag{1}$$

definierten Größen γ_{sr} , welche linearen orthogonalen Transformationen gegenüber Tensorcharakter besitzen, gegen 1 als kleine Größen behandelt werden können, deren Quadrate und Produkte gegen die ersten Potenzen vernachlässigt werden dürfen. Dabei ist $\delta_{us}=1$ bzw. $\delta_{us}=0$, ie nachdem $\mu=v$ oder $\mu\neq v$.

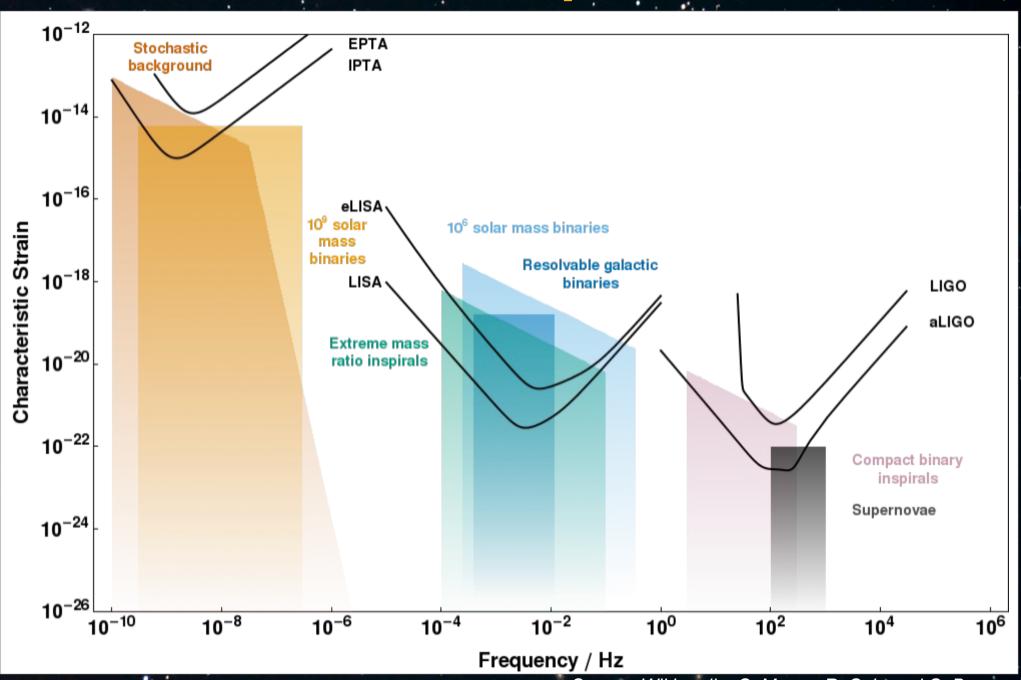
Wir werden zeigen, daß diese γ_{ω} in analoger Weise berechnet werden können wie die retardierten Potentiale der Elektrodynamik. Daraus folgt dann zunächst, daß sieh die Gravitationsfelder mit Liehtgeschwindigkeit ausbreiten. Wir werden im Anschluß an diese allgemeine Lösung die Gravitationswellen und deren Entstehungsweise untersuchen. Es hat sieh gezeigt, daß die von mir vorgeschlagene Wahl des Bezugssystems gemäß der Bedingung $g = |g_{\omega s}| = -1$ für die Berechnung der Felder in erster Näherung nicht vorteilhaft ist. Ich wurde hierauf aufmerksam durch eine briefliehe Mitteilung des Astronomen ne Strrex, der fand, daß man durch eine andere Wahl des Bezugssystems zu einem einfacheren Ausdruck des Gravitationsfeldes eines ruhenden Massenpunktes gelangen kann, als ieh ihn früher gegeben hatte'. Ich stütze mich daher im folgenden auf die allgemein invarianten Feldgleichungen.

"...so sieht man, daß A (die Ausstrahlung des Systems durch Gravitationswellen pro Zeiteneinheit) in allen nur denkbaren Fällen einen praktisch verschwindenden Wert haben muß."

A. Einstein, Sitzungsberichte der Königlich Preußischen Akademie der Wissenschaften (Berlin), 1916 688

Sitzungsber, XLVII, 1915, S. 833.

Full GW spectra

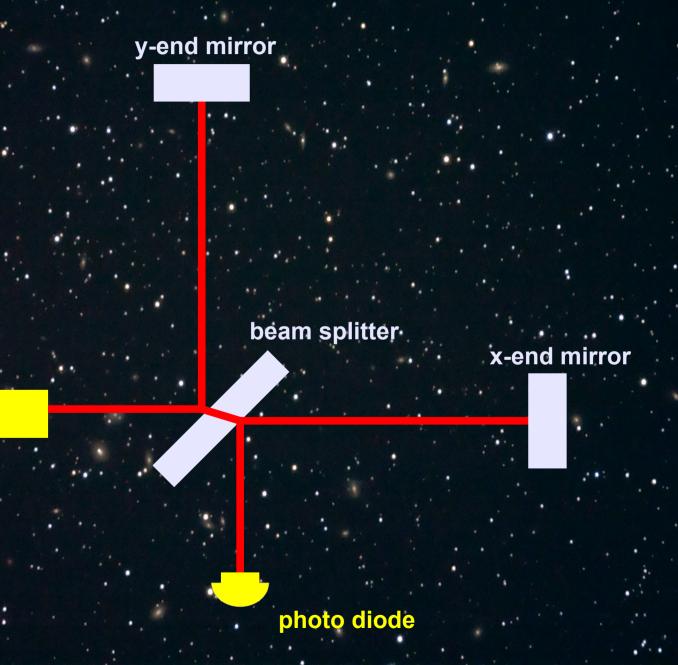


Source: Wikimedia; C. Moore, R. Cole and C. Berry



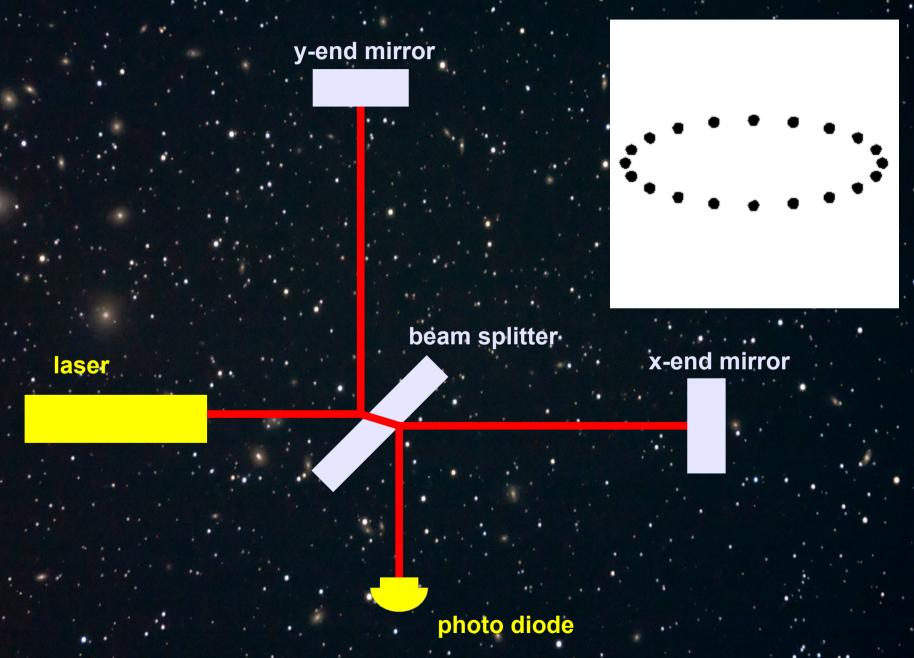
Source: LIGO Lab

Interferometers



laser

Interferometers



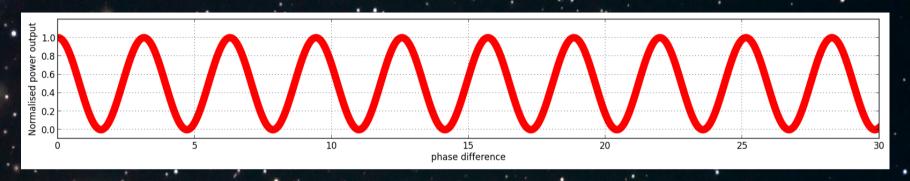
How LIGO really works

- Long arms: Earth's curvature over 4km is ~1m
- High vacuum: One trillionth atm, 10⁻⁹ torr in 10,000m³
- High power laser: 20W 1064 nm
 Nd:YAG (neodymium-doped yttrium aluminium garnet) (will be up to 200W)
- Higher power beams: Fabry-Perot cavities, 100kW, power and signal recycling

Reading between the lines

Interference pattern:

$$P_{out} = P_{max} \cos^2 \Delta \, \phi$$



$$\Delta \phi = \frac{\pi}{2} + B \frac{c \Delta T}{\lambda}$$

:Accumulated phase difference

Displacement sensitivity:

$$\Delta L = \frac{\lambda}{B} \sqrt{\frac{P_{out}}{P_{max}}}$$

What was seen II

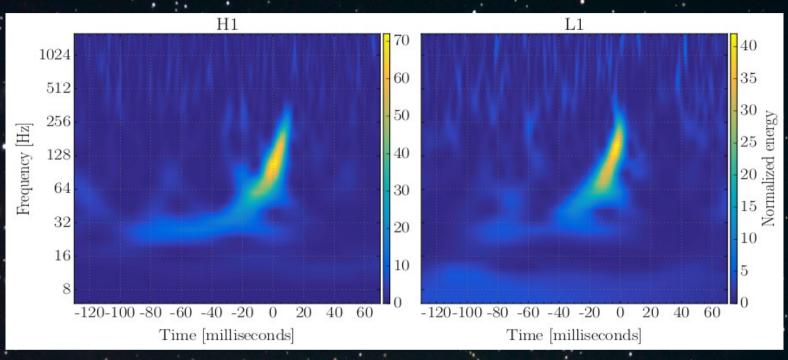


Fig. 10 Abbott et al. CQG33 (2016) 134001

Frequency ~30 Hz to ~250 Hz
Wavelength ~10,000 km to ~1,000 km
Visible duration ~ 0.1 secs
Increasing amplitude, increasing frequency = chirp
0.007 secs earlier in Livingston
The same signal in both detectors!

LVT151012

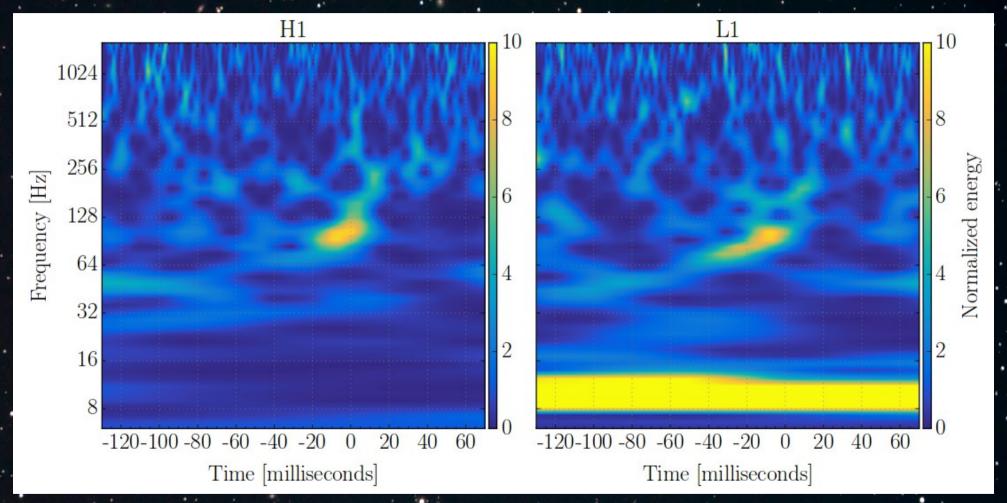
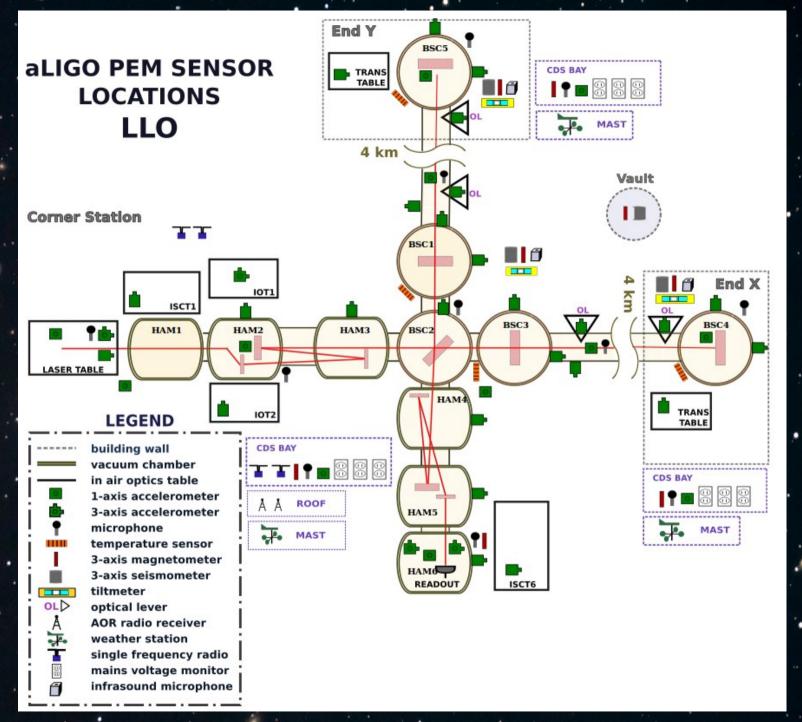


Fig. 13 Abbott et al. CQG33 (2016) 134001

False Alarm Rate, 1 per 2.3 years



http://pem.ligo.org/

Non-Gaussian transients

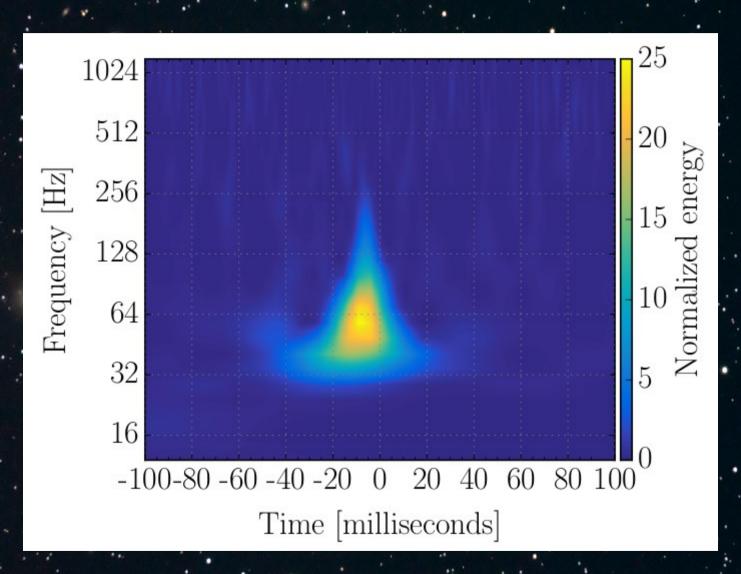


Fig. 3 Abbott et al. CQG33 (2016) 134001

Search significance

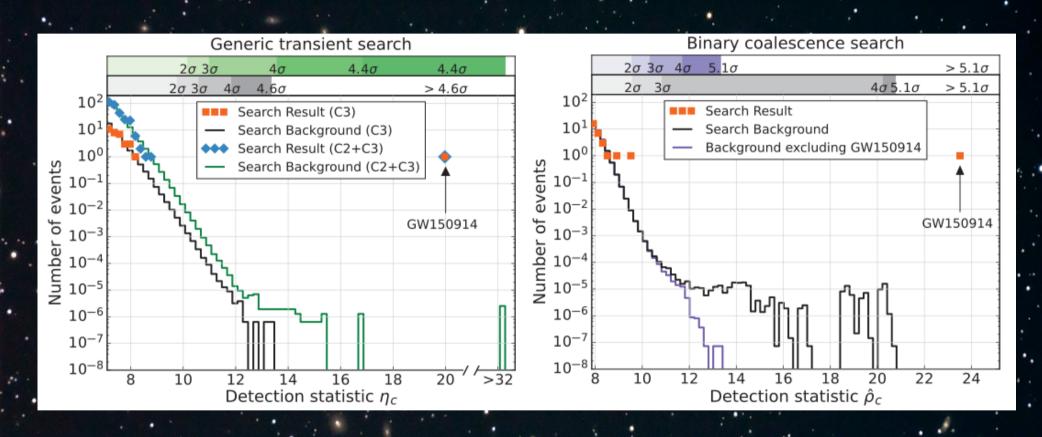


Fig 4 Abbott et al. PRL 116 (2016) 6, 061102

Binary inspiral

Quadrupole formula

$$\frac{dE_{GW}}{dt} = \frac{G}{c^5} \frac{d^3 Q_{ij}}{dt^3} \frac{d^3 Q_{ij}}{dt^3}$$

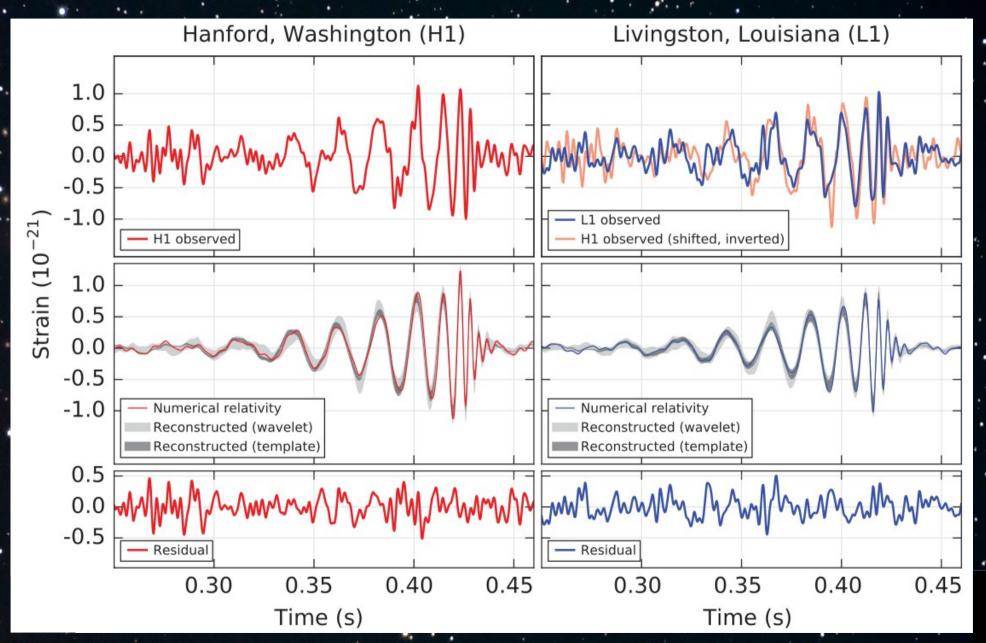
Newtonian order, chirp mass

$$f(t) = \frac{5^{3/8}}{8 \pi} \left(\frac{c^3}{GM_c} \right)^{5/8} (t_0 - t)^{-3/8}$$

Increase of frequency determined by "chirp mass"

$$M_c = \frac{(m_1 m_2)^{(3/5)}}{(m_1 + m_2)^{(1/5)}}$$

What was seen



Three binary black hole events

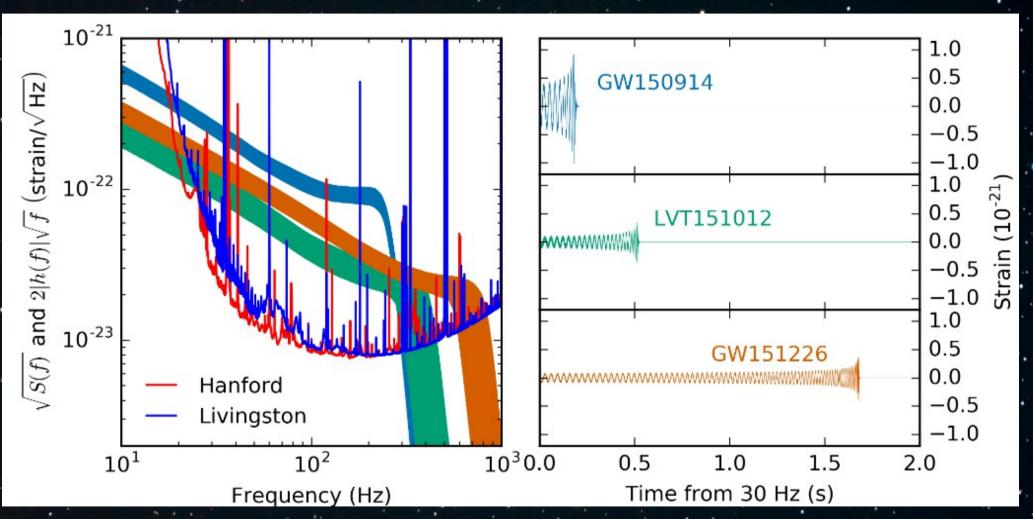


Fig 1 of Abbott et al. 1606.04856

Source parameters

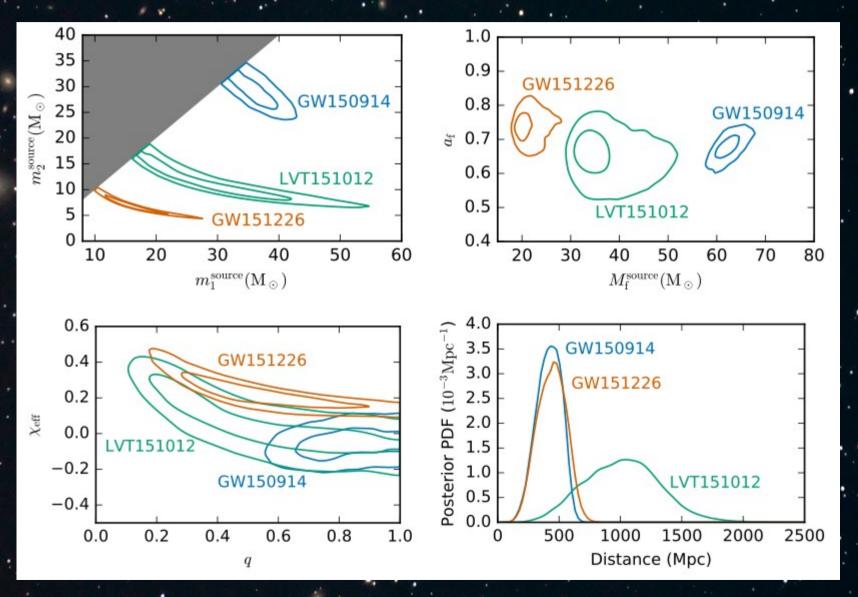
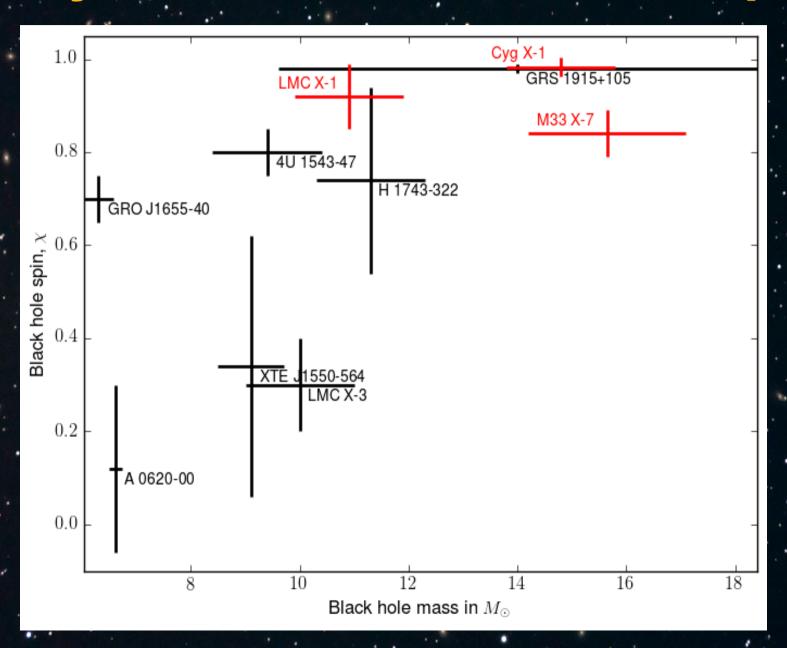
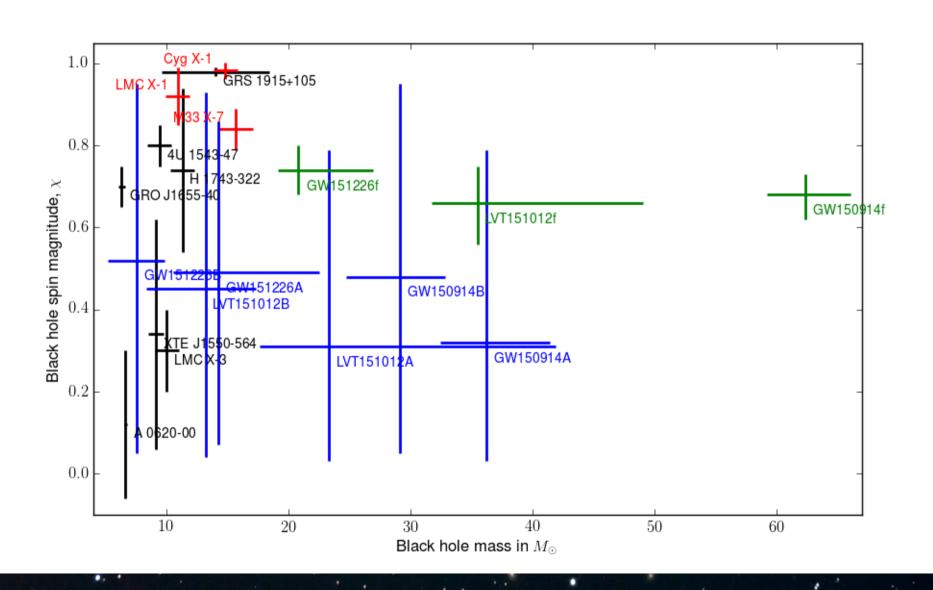


Fig 4 of Abbott et al. 1606.04856

X-ray binaries masses and spin



X-ray + GW masses and spins



Astrophysics

- Formation of heavy black holes direct collapse?
- Time to merge from 1AU by GW alone,
 ~ 100x age of universe common envelope?
- Formation of binary still open cluster or field?

• Peak energy flux 200^{+30}_{-20} solar masses per second

Signal consistency test

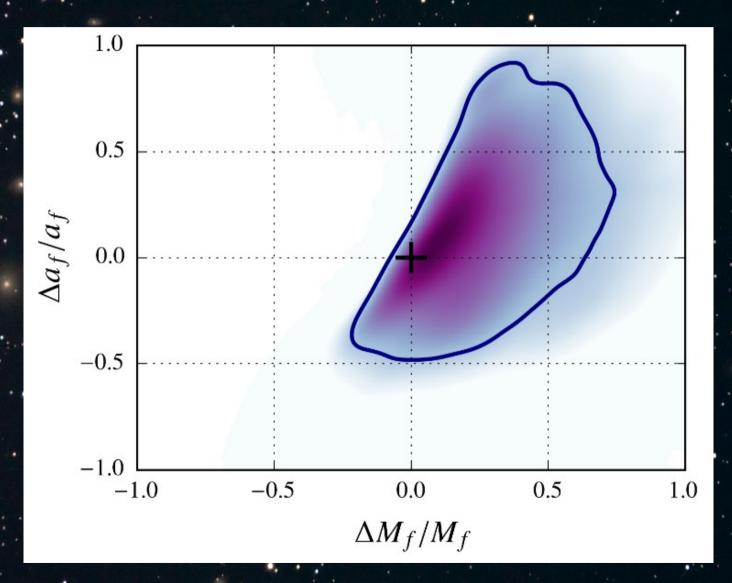
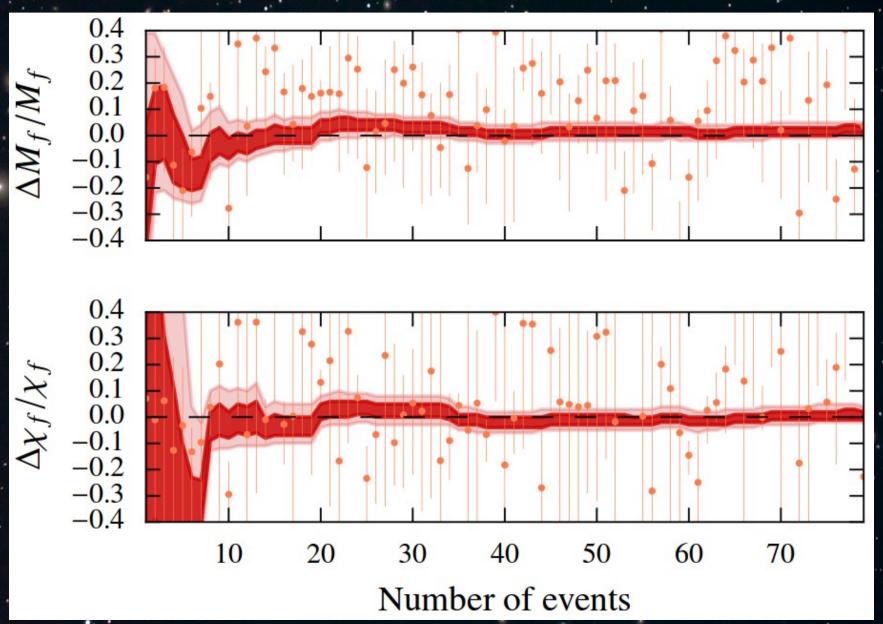


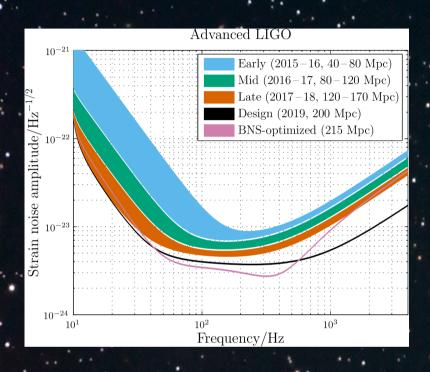
Fig. 3 (bot) of Abbott et al. PRL 16 (2016) 221101

IMR consistency going forward



Source: Archisman et al PRD94 (2016) 021101

Event rate estimates



Ten or more detections by the end of O2 is quite likely

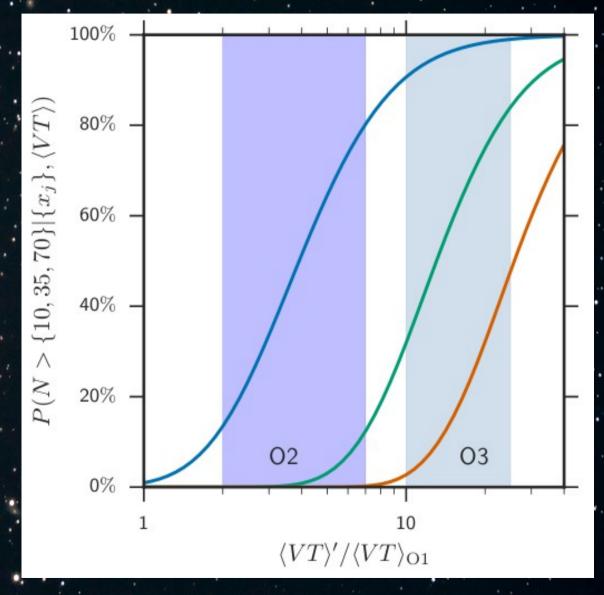
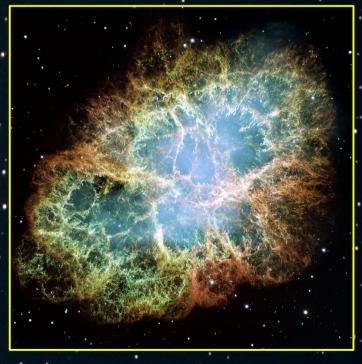


Fig 12 of Abbott et al. 1606.04856

Other potential aLIGO sources

- Neutron stars tidally disrupting
- Deformed rotating neutron stars
- Galactic supernovae
- Astrophysical background
- Cosmic strings
- First-order phase transitions
- Inflationary particle production
- Non-perturbative preheating
- Inflationary vacuum fluctuations



Source: NASA/HST

Squeezed light

- Heisenberg uncertainty in amplitude, phase
- Inject phase-locked squeezed vacuum state into output port
- Periodically poled potassium titanyl phosphate
- Hoped for ~30% gain in sensitivity

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PRL 117, 110801 (2016)

PHYSICAL REVIEW LETTERS

week ending SEPTEMBER 2010



Detection of 15 dB Squeezed States of Light and their Application for the Absolute Calibration of Photoelectric Quantum Efficiency

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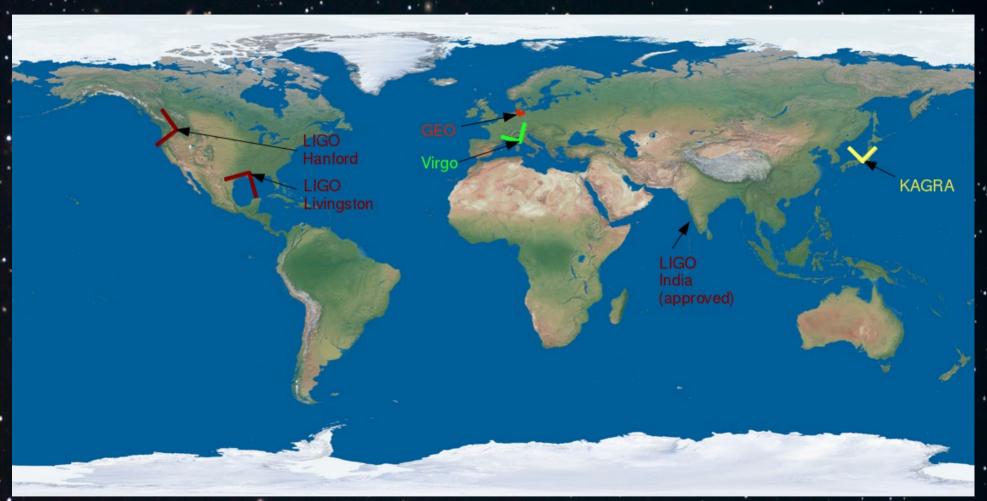
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Luruper Chaussee 149, 22761 Hamburg, Germany

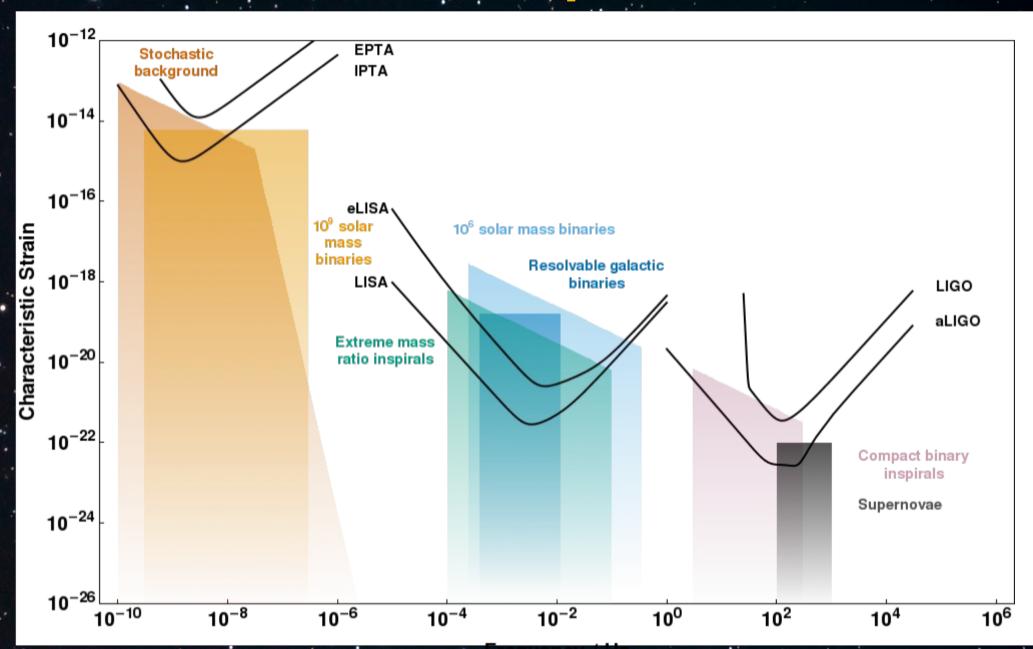
(Received 15 June 2016; published 6 September 2016)

Worldwide network



Source: Virgo/LAPP, T. Patterson

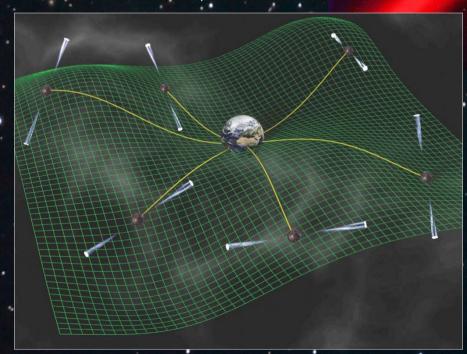
Full GW spectra

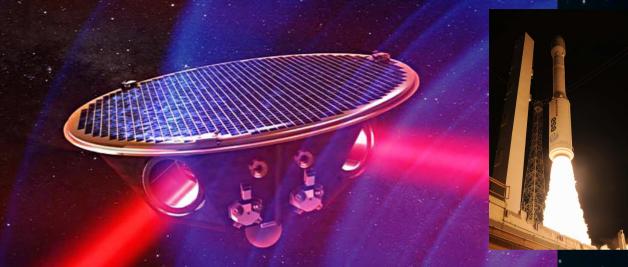


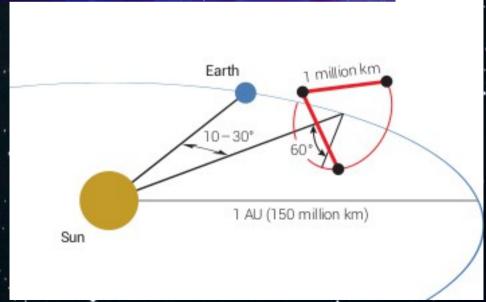
Pulsar Timing Arrays and LISA



Source: David Champion/MPIfR







Summary

- LIGO has detected gravitational waves
- Binary black hole systems exist
- Binary black holes merge
- The future is likely to bring more

Thank you

References and links

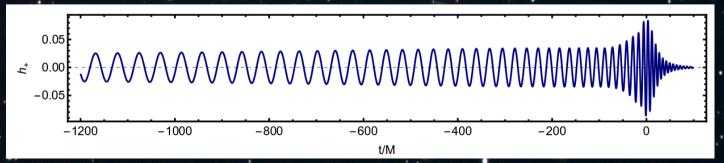
- Abbott et al. "The basic physics of the binary black hole merger GW150914" arXiv:1608.01940, Annalen der Physik
- Abbott et al. "Binary Black Hole Mergers in the first Advanced LIGO Observing Run" arXiv: 1606.04856, PRX
- Abbott et al. "Properties of the Binary Black Hole Merger GW150914" PRL 116 (2016) 241102
- LIGO Open Science Center: https://losc.ligo.org
- Logbooks: https://alog.ligo-la.caltech.edu/aLOG/
- Physical environment monitors: https://pem.ligo.org

Gravitational Waveforms

- Numerical relativity either finite differencing or spectral methods
- Effective One Body (EOBNR)
 maps two body problem to one body problem via effective
 Hamiltonian and calibrated to numerical simulations

IMRPhenom

combines post-Newtonian inspiral with phenomenological fit model of numerical simulations of late inspiral and merger, and quasi-analytical ringdown phase



Source: Khan et al. PRD 93 (2016) 044007

Average Sensitivity

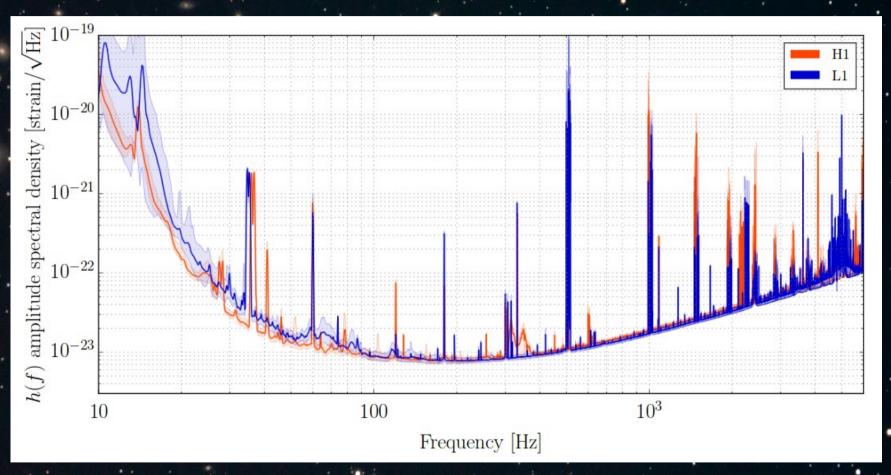


Fig. 1 Abbott et al. CQG33 (2016) 134001

Spectral lines

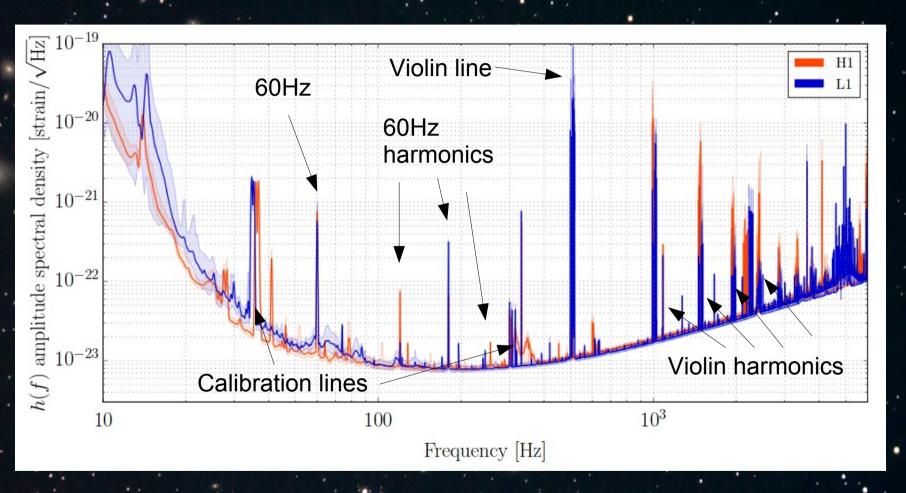


Fig. 1 Abbott et al. CQG33 (2016) 134001