

# Predicting and Discovering True Muonium

Henry Lamm

w. Yao Ji and Naveen Raman

EXA2017 - Vienna, Austria

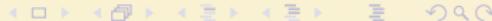


Universität Regensburg

September 12, 2017

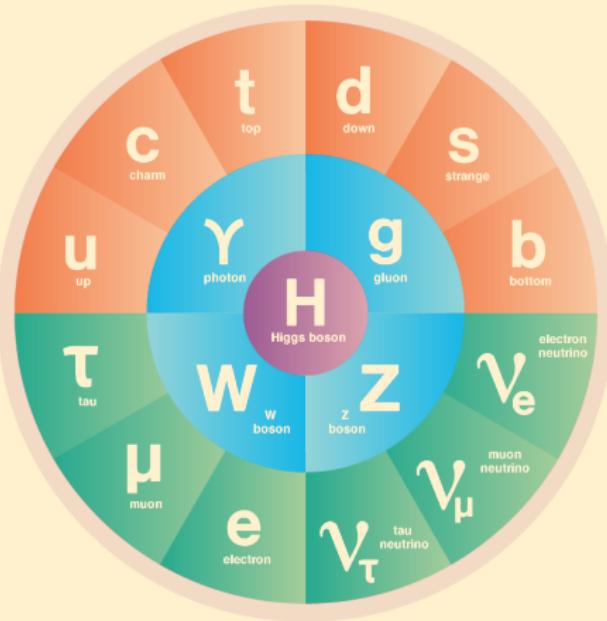


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MARYLAND

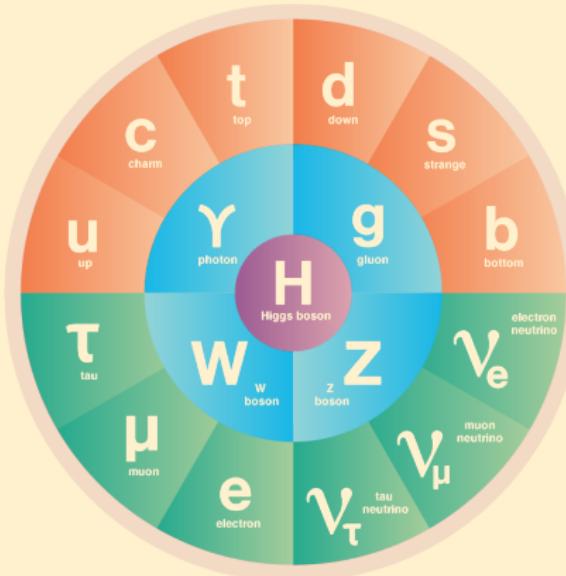


# Outline

- 1 The Muon Problem & ( $\mu^+\mu^-$ )
- 2 Predicting
- 3 Discovering
- 4 Conclusions

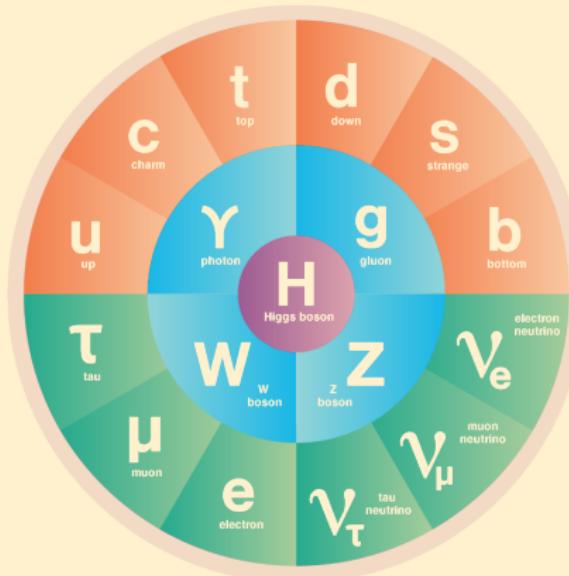


# Who ordered that?



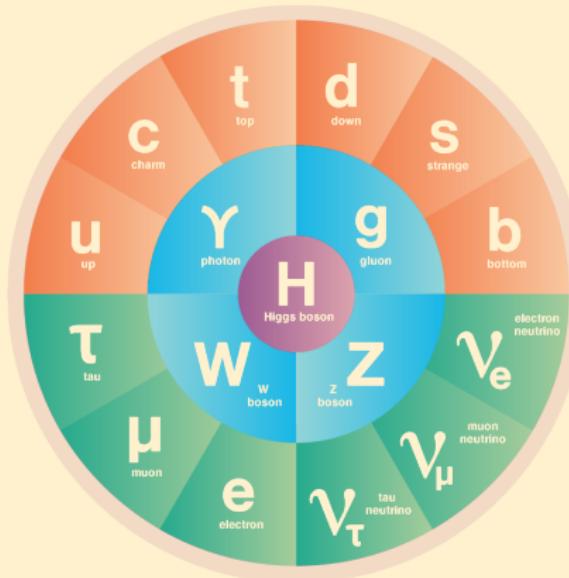
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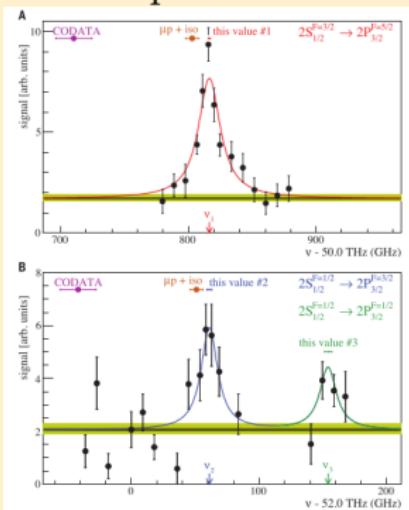
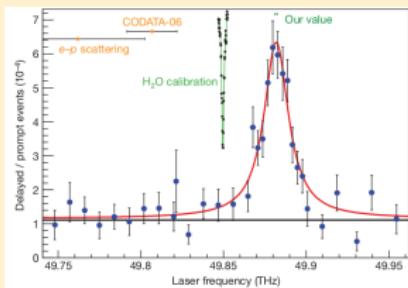
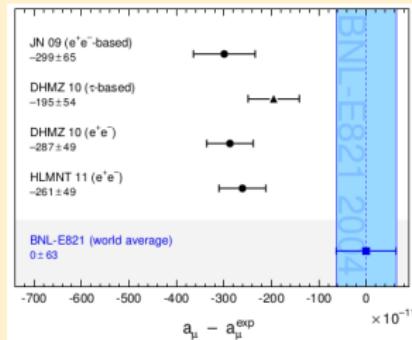


$\nu_\ell$  require modification beyond Standard Model, so at some level lepton universality is violated

# What's the deal with muon physics?<sup>123</sup>



**mu-on problem:** (n) the curious observation that discrepancies exist in muon sector when compared to other leptons

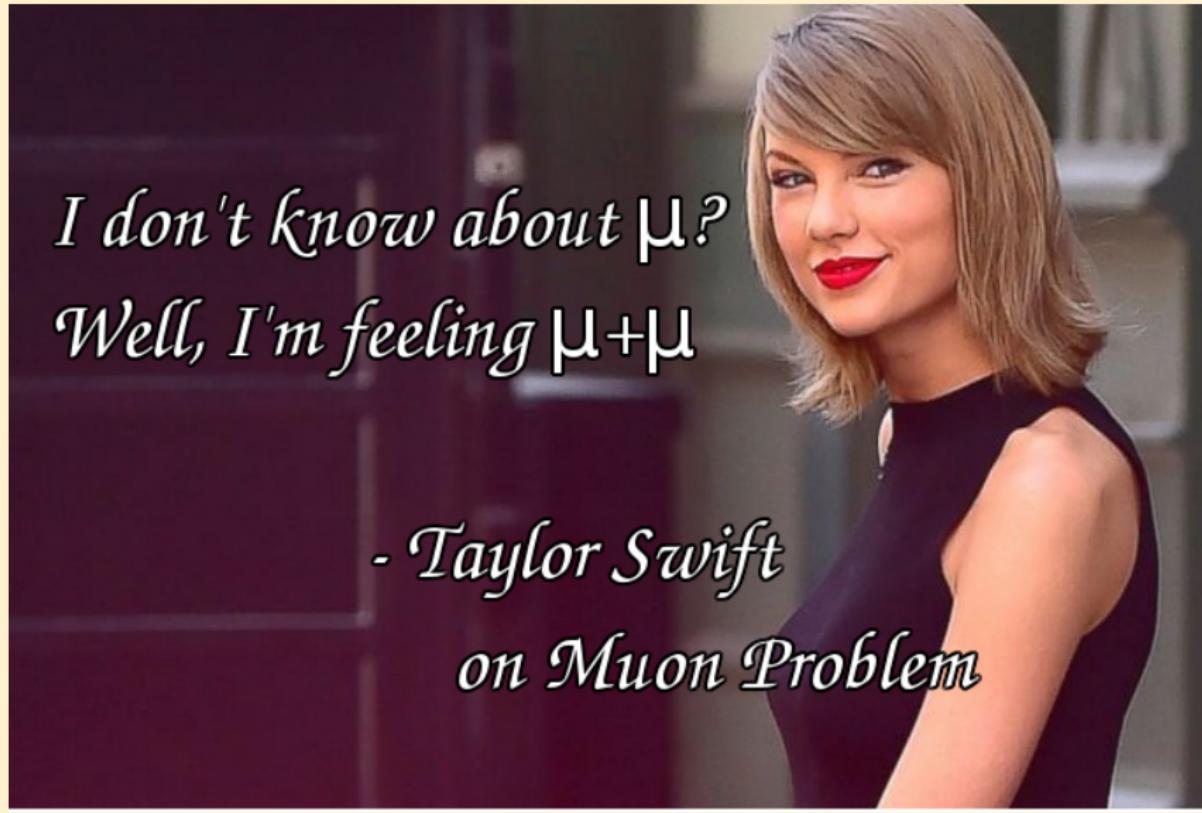


<sup>1</sup>G. Bennett et al. "Final Report of the Muon E821 Anomalous Magnetic Moment Measurement at BNL". In: *Phys. Rev. D73* (2006), p. 072003. arXiv: hep-ex/0602035 [hep-ex].

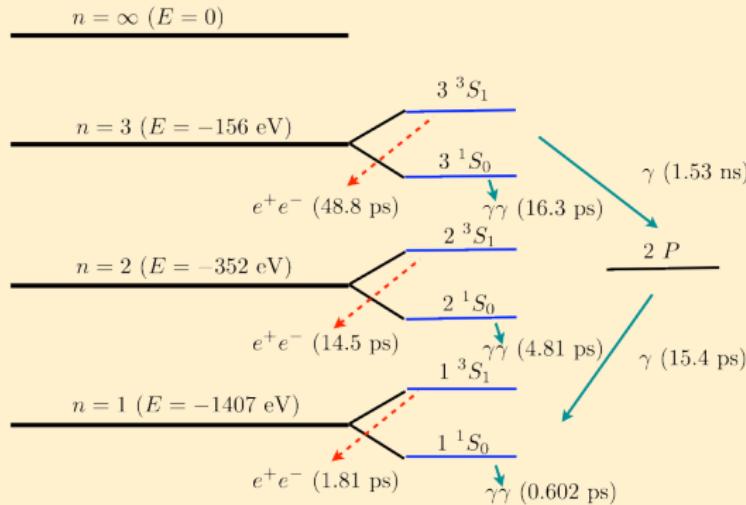
<sup>2</sup>A. Antognini et al. "Proton Structure from the Measurement of  $2S - 2P$  Transition Frequencies of Muonic Hydrogen". In: *Science 339* (2013), pp. 417–420.

<sup>3</sup>R. Pohl et al. "Laser spectroscopy of muonic deuterium". In: *Science 353.6300* (2016), pp. 669–673.

# Building a new streetlight: True Muonium



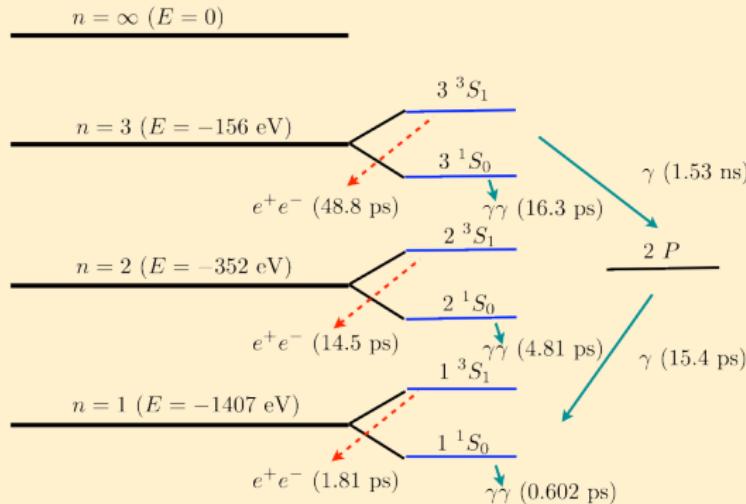
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- $(\mu^+\mu^-)$  proposed in 1961, still undetected

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# Wait? I can bind muons?

The true muonium ( $\mu^+ \mu^-$ ) and true tauonium ( $\tau^+ \tau^-$ ) [and the much more difficult to produce “mu-tauonium” ( $\mu^\pm \tau^\mp$ )] bound states are not only the heaviest, but also the most compact pure QED systems [the ( $\mu^+ \mu^-$ ) Bohr radius is 512 fm]. The relatively rapid weak decay of the  $\tau$  unfortunately makes the observation and study of true tauonium more difficult, as quantified below. In the case of true muonium the proposed production mechanisms include  $\pi^- p \rightarrow (\mu^+ \mu^-) n$  [6],  $\gamma Z \rightarrow (\mu^+ \mu^-) Z$  [6],  $eZ \rightarrow e(\mu^+ \mu^-) Z$  [12],  $Z_1 Z_2 \rightarrow Z_1 Z_2 (\mu^+ \mu^-)$  [13] (where  $Z$  indicates a heavy nucleus), direct  $\mu^+ \mu^-$  collisions [7],  $\eta \rightarrow (\mu^+ \mu^-) \gamma$  [14], and  $e^+ e^- \rightarrow (\mu^+ \mu^-)$  [15]. In addition, the properties of true muonium have been studied in a number of papers [9,16,17].

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- ...but many channels suggested<sup>4</sup>

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...but clearly wrong are Muononium and Muonic Muonium

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# ...because if you did<sup>8</sup>910111213141516171819

<sup>8</sup>A. Badertscher et al. “An Improved Limit on Invisible Decays of Positronium”. In: *Phys. Rev.* D75 (2007), p. 032004. arXiv: [hep-ex/0609059 \[hep-ex\]](#).

<sup>9</sup>J. Jaeckel and S. Roy. “Spectroscopy as a test of Coulomb’s law: A Probe of the hidden sector”. In: *Phys. Rev.* D82 (2010), p. 125020. arXiv: [1008.3536 \[hep-ph\]](#).

<sup>10</sup>B. Batell, D. McKeen, and M. Pospelov. “New Parity-Violating Muonic Forces and the Proton Charge Radius”. In: *Phys. Rev. Lett.* 107 (2011), p. 011803. arXiv: [1103.0721 \[hep-ph\]](#).

<sup>11</sup>B. Batell, M. Pospelov, and A. Ritz. “Multi-lepton Signatures of a Hidden Sector in Rare B Decays”. In: *Phys. Rev.* D83 (2011), p. 054005. arXiv: [0911.4938 \[hep-ph\]](#).

<sup>12</sup>D. Tucker-Smith and I. Yavin. “Muonic hydrogen and MeV forces”. In: *Phys. Rev.* D83 (2011), p. 101702. arXiv: [1011.4922 \[hep-ph\]](#).

<sup>13</sup>J. Kopp, L. Michaels, and J. Smirnov. “Loopy Constraints on Leptophilic Dark Matter and Internal Bremsstrahlung”. In: *JCAP* 1404 (2014), p. 022. arXiv: [1401.6457 \[hep-ph\]](#).

<sup>14</sup>A. H. Gomes, A. Kostelecký, and A. J. Vargas. “Laboratory tests of Lorentz and CPT symmetry with muons”. In: *Phys. Rev.* D90.7 (2014), p. 076009. arXiv: [1407.7748 \[hep-ph\]](#).

<sup>15</sup>S. G. Karshenboim. “HFS interval of the  $2s$  state of hydrogen-like atoms and a constraint on a pseudovector boson with mass below  $1 \text{ keV}/c^2$ ”. In: *Phys. Rev.* A83 (2011), p. 062119. arXiv: [1005.4875 \[hep-ph\]](#).

<sup>16</sup>S. Karshenboim. “Precision physics of simple atoms and constraints on a light boson with ultraweak coupling”. In: *Phys. Rev. Lett.* 104 (2010), p. 220406. arXiv: [1005.4859 \[hep-ph\]](#).

<sup>17</sup>S. Karshenboim and V. Flambaum. “Constraint on axion-like particles from atomic physics”. In: *Phys. Rev.* A84 (2011), p. 064502. arXiv: [1110.6259 \[physics.atom-ph\]](#).

<sup>18</sup>S. G. Karshenboim, D. McKeen, and M. Pospelov. “Constraints on muon-specific dark forces”. In: *Phys. Rev.* D90.7 (2014), p. 073004. arXiv: [1401.6154 \[hep-ph\]](#).

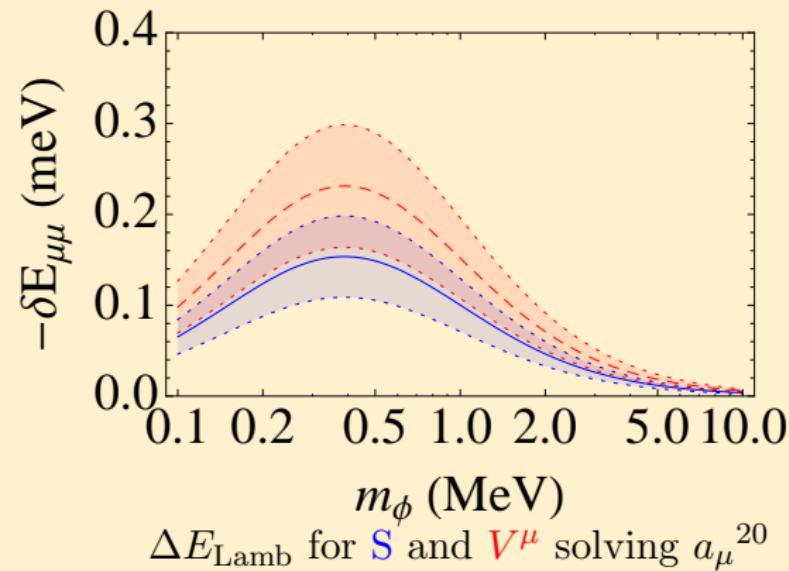
<sup>19</sup>H. Lamm. “Applying Bayesian Inference to Galileon Solutions of the Muon Problem”. In: *Phys. Rev.* D94.11 (2016), p. 115007. arXiv: [1609.07520 \[hep-ph\]](#).

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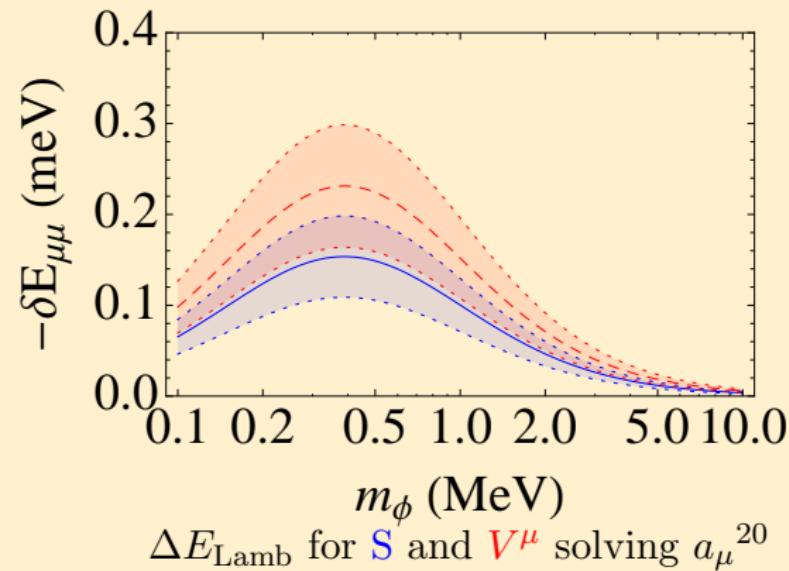
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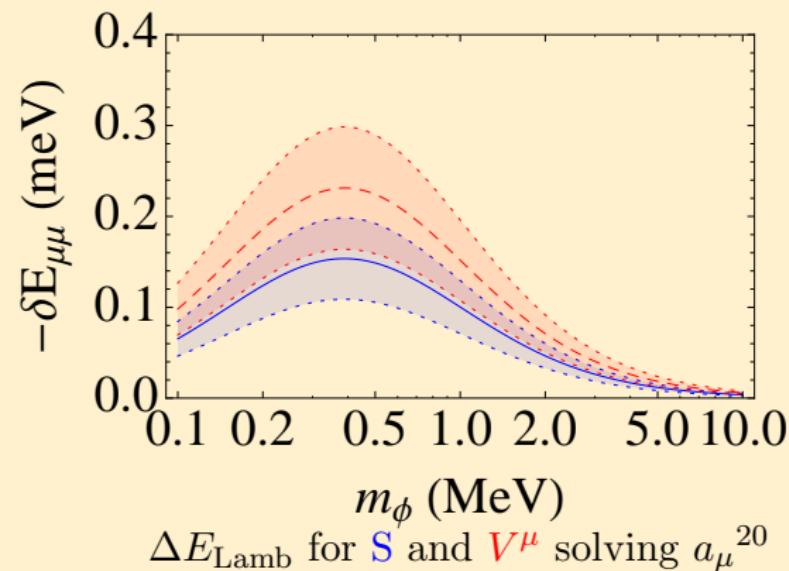
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Annihilation channel through Fierz identity enhances suppressed interactions

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$$\delta a_\mu \approx 2 \times 10^{-9} \left( \frac{2 \text{TeV}}{M_x/g} \right)^2 \quad \delta E \approx \frac{g^2}{M_x^2} |\psi(0)|^2 = \mathbf{100 \text{ MHz}} \left( \frac{1 \text{TeV}}{M_x/g} \right)^2$$

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<u>Hydrogen</u>	<u><math>\mu\text{H}</math></u>	<u>Positronium</u>	<u><math>\mu^+\mu^-</math></u>
$\equiv 2^2P_{3/2}$	$\equiv 2^2P_{3/2}$	$2^3S_1$	$2^3P_2$
$\equiv 2^2P_{1/2}$	$\equiv 2^2P_{1/2}$	$2^3P_2$ $2^1P_1$ $2^3P_1$	$2^3P_1$ $2^3P_0$
$\equiv 2^2S_{1/2}$	$\equiv 2^2S_{1/2}$	$2^3P_0$ $2^1S_0$	$2^3S_1$ $2^1S_0$

# Goal: $\mathcal{O}(100 \text{ MHz})$ true muonium predictions

Obs.	$E[\text{MHz}]$ 2000	$E[\text{MHz}]$ 2017
Lamb	$1.37(5) \times 10^7$	$1.3813(14) \times 10^7$
$1s - 2s$	$2.55(5) \times 10^{11}$	—
hfs	$42329604(800)_{\text{had}}(1200)_{\text{miss}}$	$42329435(16)_{\text{had}}(90)_{\text{ind}}(700)_{\mu}$

- Unlike Positronium  $e, had, \tau$  loops contribute appreciably  
 $\frac{\alpha m_\mu}{m_e} \approx 1.5$ .

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<sup>21</sup>HL and NR, in prep.

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- Technology exist for these computations diagram by diagram

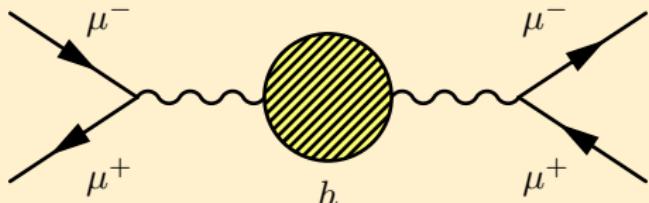
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# Improving the hfs prediction

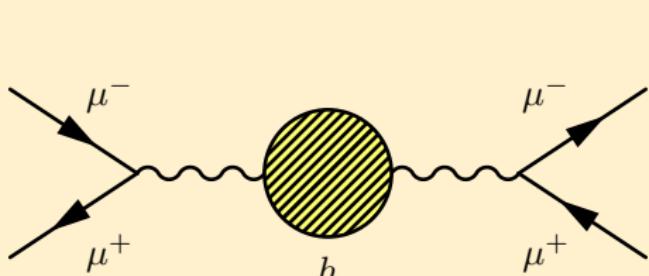
# Improving the hfs prediction



$$\Delta E_{1,\text{hfp}} = \frac{m_\mu \alpha^5}{n^3 \pi} \left[ m_\mu^2 \int_{4m_\pi^2}^{\infty} ds \frac{R(s)}{3s(4m_\mu^2 - s)} \right]$$

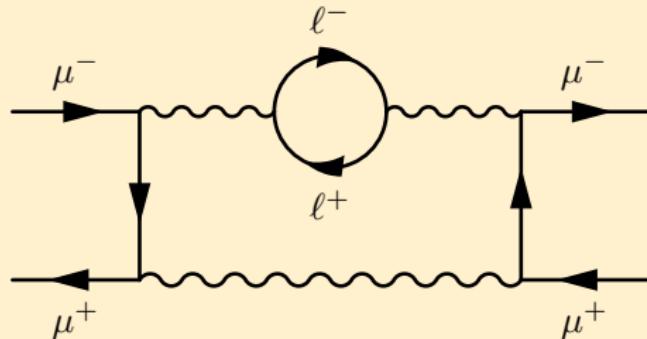
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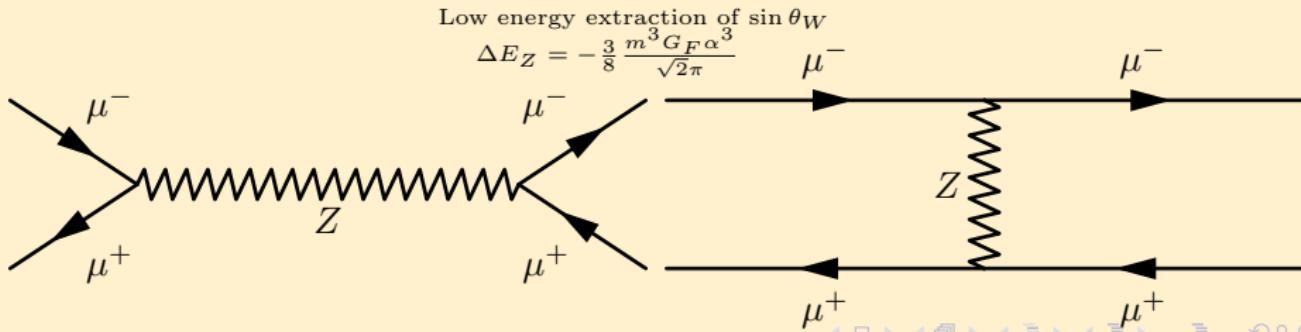
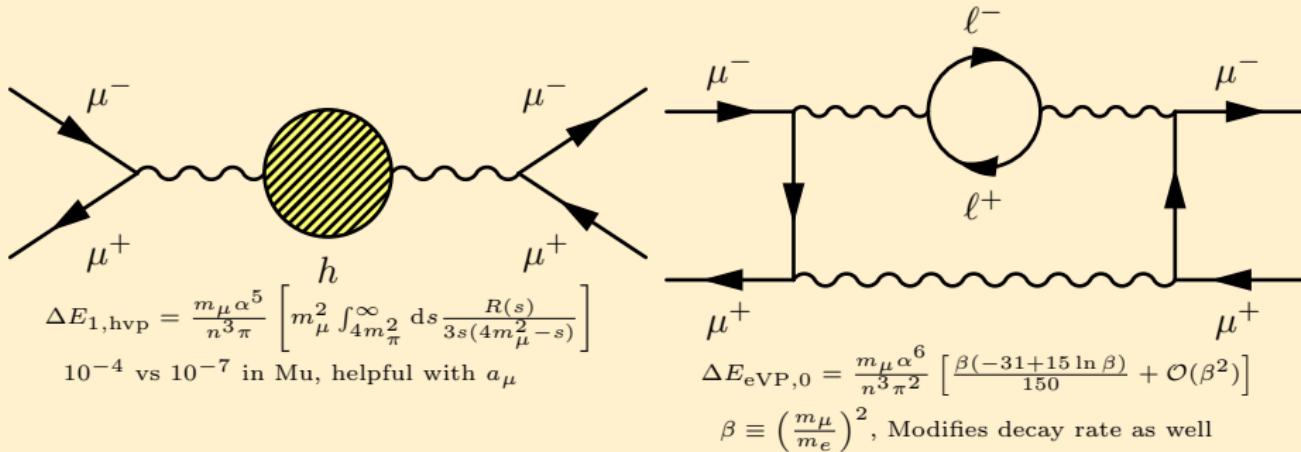
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$$\Delta E_{\text{eVP},0} = \frac{m_\mu \alpha^6}{n^3 \pi^2} \left[ \frac{\beta(-31 + 15 \ln \beta)}{150} + \mathcal{O}(\beta^2) \right]$$

$\beta \equiv \left( \frac{m_\mu}{m_e} \right)^2$ , Modifies decay rate as well

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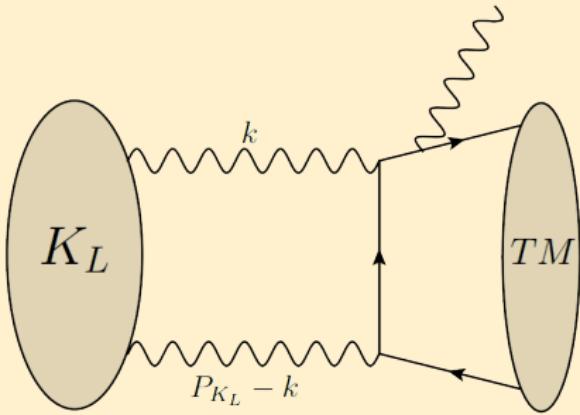
# The current state of hfs in true muonium

$\mathcal{O}(m\alpha^n)$	$C_n$	$\delta E_{\text{hfs}}^{1s} [\text{MHz}]$
$m\alpha^4$	$\frac{7}{12}$	42260692
$\frac{m\alpha^5}{\pi}$	$-\frac{1}{2} \ln(2) - \frac{8}{9}$	-207904
$\frac{m\alpha^5}{\pi} _\mu$	0.81650(9)	1611.04258
$\frac{m\alpha^6}{\pi^2}$	$-\frac{52}{32} \zeta(3) + \left(\frac{221}{24} \ln(2) - \frac{5197}{576}\right) \zeta(2) + \frac{1}{2} \ln(2) + \frac{1367}{648}$	-1515
$\frac{m\alpha^6}{\pi^2} _\mu$	350.572533*	1607.07059, 1701.04362
$m\alpha^6 \ln\left(\frac{1}{\alpha}\right)$	$\frac{5}{24}$	3954
$\frac{m\alpha^7}{\pi^3}$	160*	145*(90)
$\frac{m\alpha^7}{\pi} \ln\left(\frac{1}{\alpha}\right)$	$-\frac{17}{3} \ln(2) + \frac{217}{90}$	-67
$\frac{m\alpha^7}{\pi} \ln^2\left(\frac{1}{\alpha}\right)$	$-\frac{7}{8}$	-190
$\frac{m^3 G_F \alpha^3}{\sqrt{2}\pi}$	$-\frac{3}{8}$	1502.03841
$\frac{m\alpha^7}{\pi^3} _\mu$	-5.324248*	1507.07841
Total		42329435(16)(90)(700)

$|_\mu$  indicates true muonium corrections missing from positronium, which depend upon  $m_\mu/m_e$ . Errors consists of (1) hadronic model dependence, (2) missing mass-independent QED contributions at  $\mathcal{O}(m\alpha^7)$ , (3) missing mass-dependent QED contributions at  $\mathcal{O}(m\alpha^6)$ . \* indicate partial terms.

# Outline

- 1 The Muon Problem &  $(\mu^+ \mu^-)$
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- 3 Discovering
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# Time for the bad news



“... we choose to  
do these things  
not because  
they are easy  
but because  
they are hard.”

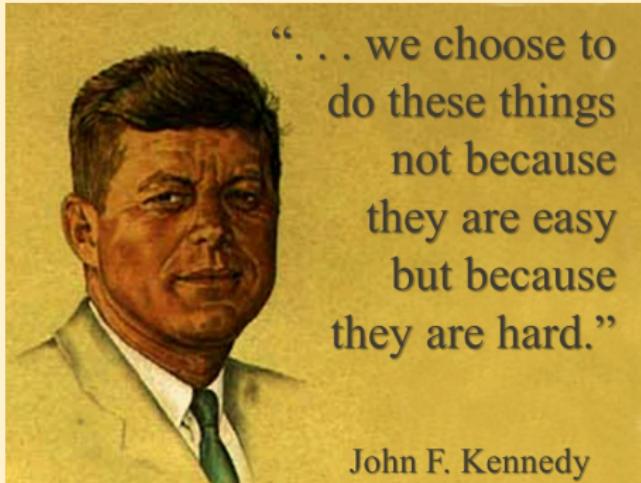
John F. Kennedy

# Time for the bad news



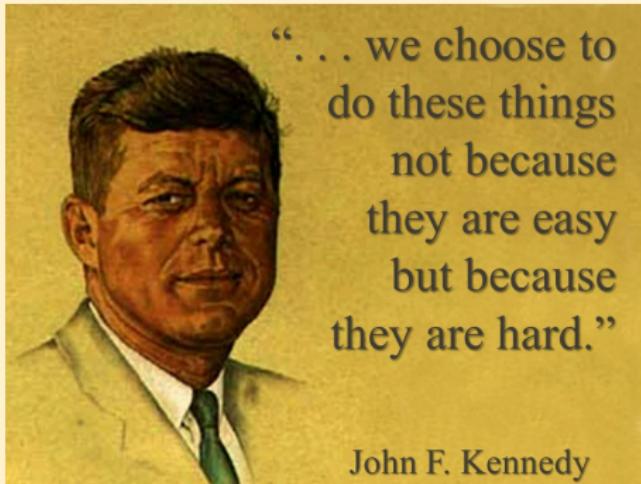
- Analogous to pion production, but  $\propto \left(\frac{\alpha_{em}}{\alpha_S}\right)^3$

# Time for the bad news



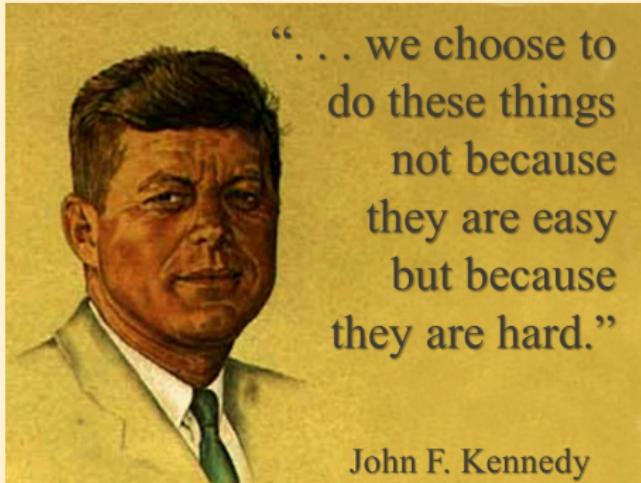
- Analogous to pion production, but  $\propto \left(\frac{\alpha_{em}}{\alpha_S}\right)^3$
- Lifetime is  $\tau_{(\mu^+ \mu^-)} = n^3 \text{ ps} \ll \tau_{\pi^0} = 26 \text{ ns}$

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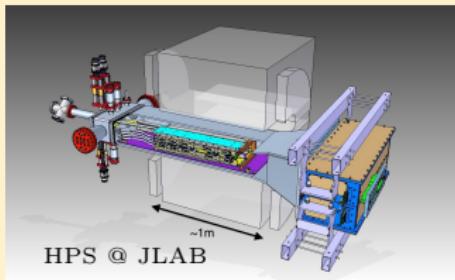
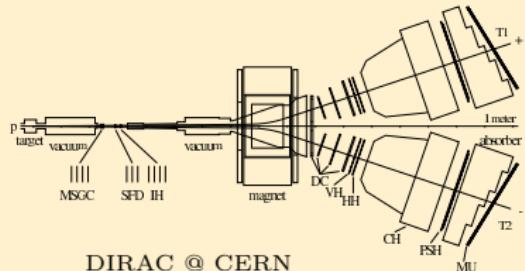
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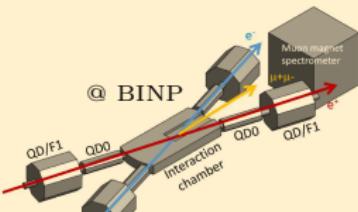
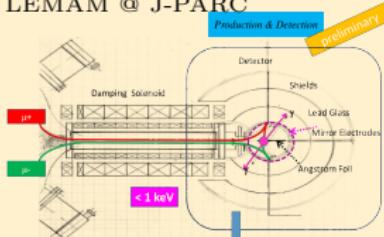


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- ...but true muonium is likely to be **relativistic**
- Truly an **intensity** frontier proposition

# Near future experimental efforts to detect possible

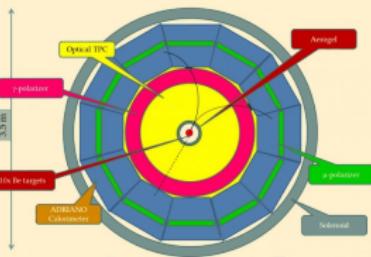


LEMAM @ J-PARC

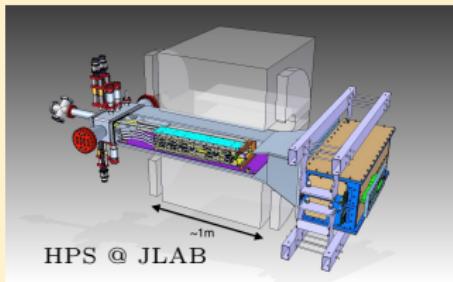
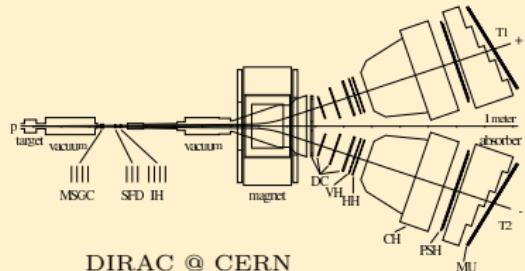


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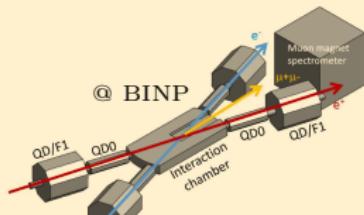
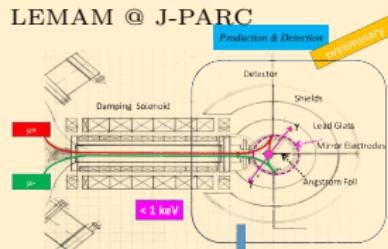
REDTOP @ Fermilab



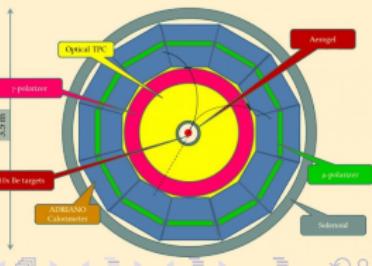
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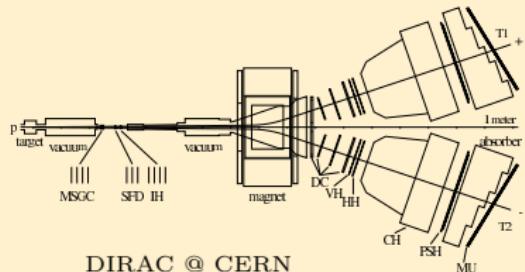
- Exist: DIRAC, HPS are fixed targets



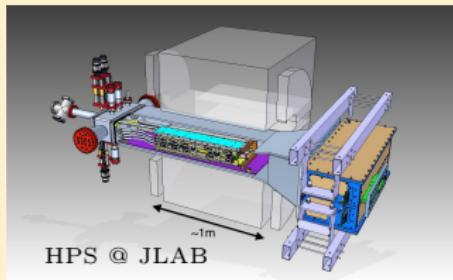
REDTOP @ Fermilab



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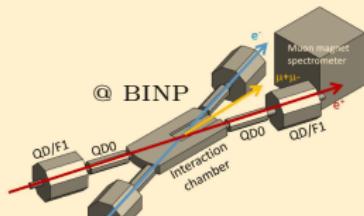
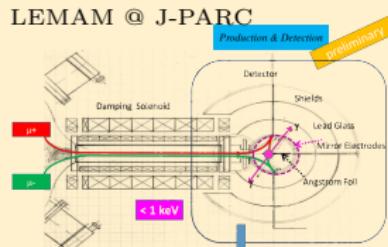


DIRAC @ CERN



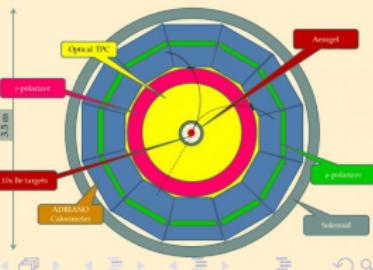
HPS @ JLAB

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- Proposed: LEMAM is  $\mu^+\mu^-$  collider

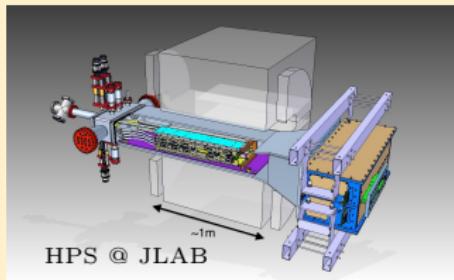
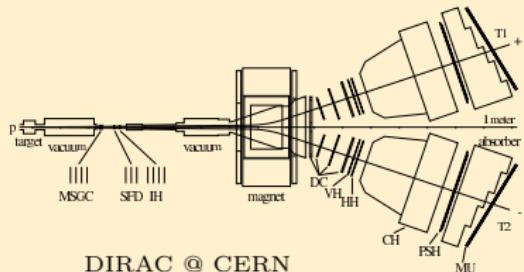


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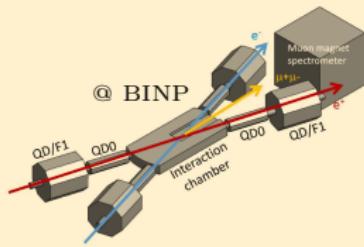
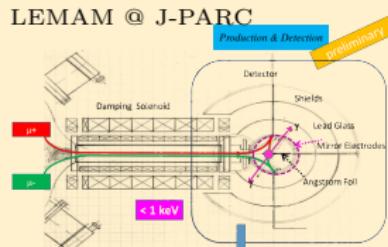
REDTOP @ Fermilab



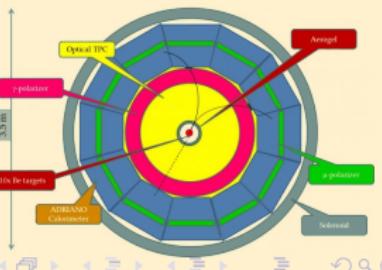
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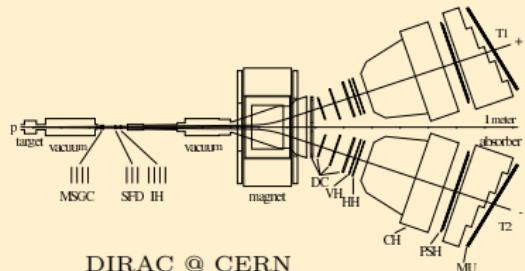
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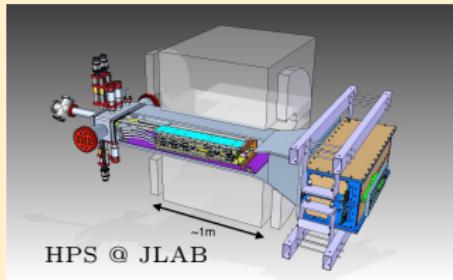
REDTOP @ Fermilab



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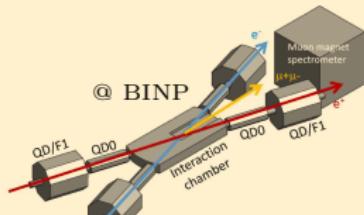
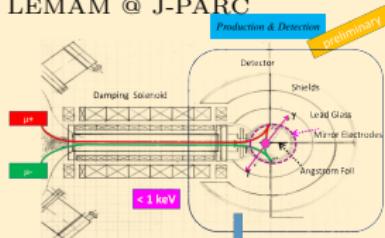


DIRAC @ CERN



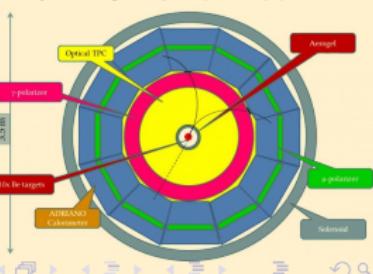
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- Proposed: **REDTOP is a  $\eta/\eta'$  factory**

LEMAM @ J-PARC



See Anton's Talk: Wed 16:20

REDTOP @ Fermilab



# Rare meson decay may open the path

Proposals for rare decay searches, like REDTOP, would be capable of detecting true muonium+ $\gamma$ . Large  $\mathcal{O}(\alpha)$  corrections and  $F_{\gamma^*}(q^2)$  dependence needed

$$\frac{\mathcal{B}(\eta \rightarrow \gamma(\mu^+ \mu^-))}{\mathcal{B}(\eta \rightarrow \gamma\gamma)} = 1.476(5)_{\text{stat}}(4)_{\text{sys}} \times 10^{-9}$$

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$$\frac{\mathcal{B}(K_L \rightarrow \gamma(\mu^+ \mu^-))}{\mathcal{B}(K_L \rightarrow \gamma\gamma)} \approx 1.26(2)_{\text{model}} \times 10^{-9}$$

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<sup>24</sup>YJ and HL, in prep.

<sup>25</sup>Y. Ji and H. Lamm. “Discovering True Muonium in  $K_L \rightarrow (\mu^+ \mu^-)\gamma$ ”. In: (2017). arXiv: 1706.04986 [hep-ph].

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REDTOP:  $2 \times 10^{12} \eta \rightarrow \mathbf{80 \text{ events!}^{24}}$  ( $\epsilon = 10\%$ )

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Near-future experiments could also probe  $\eta'$ ,  $K_L$ <sup>25</sup>,  $K_S$ , ...

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Discovery would constrain on  $F_{\gamma^*}(q^2)$

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# Outline

- 1 The Muon Problem &  $(\mu^+ \mu^-)$
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# The future looks bright

- Your favorite **new physics** can be severely constrained with true muonium
- A theoretical program is underway to reduce the **uncertainty** in observables to  $\mathcal{O}(100 \text{ MHZ})$
- Contributions possible to  $a_\mu$  and  $\sin \theta_W$
- Most promising experimental outlook in 56 years!