

Regulator dependence and isospin breaking

Sebastian König

**EMMI RRTF Workshop:
“The systematic treatment of the Coulomb interaction in few-body systems”**

GSI, Darmstadt

May 30, 2016



THE OHIO STATE UNIVERSITY

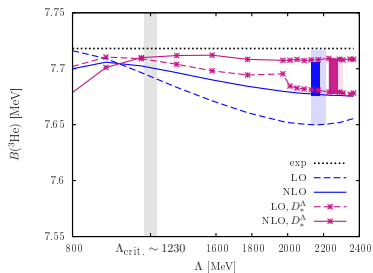
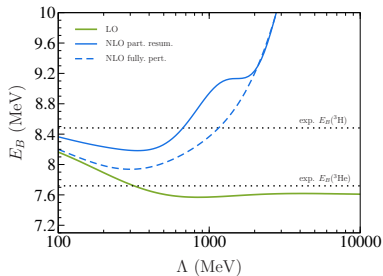


NUCLEI
Nuclear Computational Low-Energy Initiative

Motivation: counterterm controversy

Should there be a new counterterm at NLO!

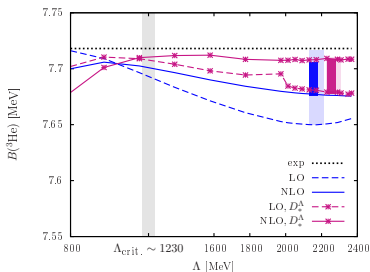
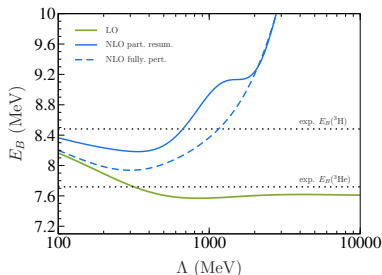
Yes! [Vanasse et al., PRC 89 \(2014\) 064003](#) No! [Kirscher+Gazit, PLB 755 \(2016\) 253](#)



Motivation: counterterm controversy

Should there be a new counterterm at NLO!

Yes! [Vanasse et al., PRC 89 \(2014\) 064003](#) No! [Kirscher+Gazit, PLB 755 \(2016\) 253](#)



Vanasse et al.

- momentum-space formalism
- sharp cutoff (non-local)
- can take Λ arbitrarily large

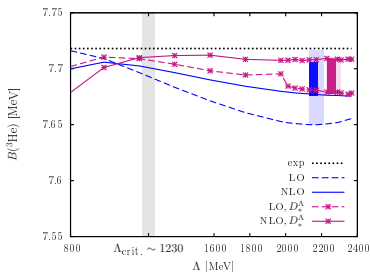
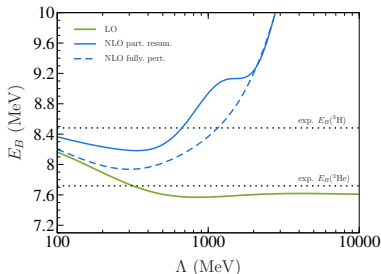
Kirscher and Gazit

- configuration-space (R)RGM
- Gaussian regulators (local)
- limited cutoff range

Motivation: counterterm controversy

Should there be a new counterterm at NLO!

Yes! [Vanasse et al., PRC 89 \(2014\) 064003](#) No! [Kirscher+Gazit, PLB 755 \(2016\) 253](#)



Vanasse et al.

- momentum-space formalism
- sharp cutoff (non-local)
- can take Λ arbitrarily large

Kirscher and Gazit

- configuration-space (R)RGM
- Gaussian regulators (local)
- limited cutoff range

power counting \leftrightarrow regulators?

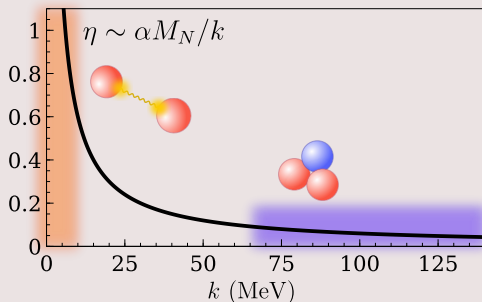
Coulomb corrections for nuclei



How much Coulomb should we really iterate?

Coulomb regimes

trinucleon binding momentum ~ 80 MeV ...



↪ should Coulomb not be a small perturbative correction?

Nonperturbative vs. perturbative and helium

$$\begin{aligned}
 \text{Diagram 1} &= \text{Diagram 1.1} + \text{Diagram 1.2} + \text{Diagram 1.3} + \text{Diagram 1.4} \times (\text{Diagram 1.1} + \text{Diagram 1.2} + \text{Diagram 1.3}) \\
 &+ \text{Diagram 1.5} \times (\text{Diagram 1.1} + \text{Diagram 1.3}) + \text{Diagram 1.6} \times (\text{Diagram 1.1} + \text{Diagram 1.3})
 \end{aligned}$$



$$\begin{aligned}
 \text{Diagram 2} &= \text{Diagram 2.1} + \text{Diagram 2.2} + \text{Diagram 2.3} \times (\text{Diagram 2.1} + \text{Diagram 2.2}) \\
 &+ \text{Diagram 2.4} \times (\text{Diagram 2.1} + \text{Diagram 2.3}) \\
 &+ \text{Diagram 2.5} \times (\text{Diagram 2.1} + \text{Diagram 2.2})
 \end{aligned}$$

$$\begin{aligned}
 \text{Diagram 3} &= \text{Diagram 3.1} + \text{Diagram 3.2} + \text{Diagram 3.3} \times (\text{Diagram 3.1} + \text{Diagram 3.2}) \\
 &+ \text{Diagram 3.4} \times (\text{Diagram 3.1} + \text{Diagram 3.2})
 \end{aligned}$$

- iterate $\mathcal{O}(\alpha)$ diagrams. . .
- get ${}^3\text{He}$ pole directly

Nonperturbative vs. perturbative and helium

$$\begin{aligned}
 \text{Diagram 1} &= \text{Diagram 1.1} + \text{Diagram 1.2} + \text{Diagram 1.3} + \text{Diagram 1.4} \times (\text{Diagram 1.1} + \text{Diagram 1.2} + \text{Diagram 1.3}) \\
 &+ \text{Diagram 1.5} \times (\text{Diagram 1.1} + \text{Diagram 1.3}) + \text{Diagram 1.6} \times (\text{Diagram 1.1} + \text{Diagram 1.3})
 \end{aligned}$$



$$\begin{aligned}
 \text{Diagram 2} &= \text{Diagram 2.1} + \text{Diagram 2.2} + \text{Diagram 2.3} + \text{Diagram 2.4} \\
 &+ \text{Diagram 2.5} \times (\text{Diagram 2.1} + \text{Diagram 2.2}) \\
 &+ \text{Diagram 2.6} \times (\text{Diagram 2.2} + \text{Diagram 2.4})
 \end{aligned}$$

$$\begin{aligned}
 \text{Diagram 3} &= \text{Diagram 3.1} + \text{Diagram 3.2} + \text{Diagram 3.3} \times (\text{Diagram 3.1} + \text{Diagram 3.2}) \\
 &+ \text{Diagram 3.4} \times (\text{Diagram 3.1} + \text{Diagram 3.2})
 \end{aligned}$$

- iterate $\mathcal{O}(\alpha)$ diagrams. . .
- get ${}^3\text{He}$ pole directly

nonperturbative!

Nonperturbative vs. perturbative and helium

$$\text{Diagram 1} = \text{Diagram 1a} + \text{Diagram 1b} + \text{Diagram 1c} + \text{Diagram 1d} \times (\text{Diagram 1a} + \text{Diagram 1b} + \text{Diagram 1c})$$

$$+ \text{Diagram 1e} \times (\text{Diagram 1a} + \text{Diagram 1c}) + \text{Diagram 1f} \times (\text{Diagram 1a} + \text{Diagram 1c})$$



$$\text{Diagram 2} = \text{Diagram 2a} + \text{Diagram 2b} + \text{Diagram 2c} + \text{Diagram 2d}$$

$$+ \text{Diagram 2e} \times (\text{Diagram 2a} + \text{Diagram 2c})$$

$$+ \text{Diagram 2f} \times (\text{Diagram 2b} + \text{Diagram 2d})$$

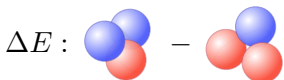
$$\text{Diagram 3} = \text{Diagram 3a} + \text{Diagram 3b} + \text{Diagram 3c} \times (\text{Diagram 3a} + \text{Diagram 3b})$$

$$+ \text{Diagram 3d} \times (\text{Diagram 3a} + \text{Diagram 3b})$$

nonperturbative!

- iterate $\mathcal{O}(\alpha)$ diagrams...
- get ${}^3\text{He}$ pole directly

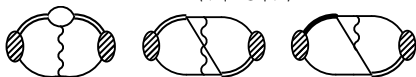
- use trinucleon wavefunctions
- fully perturbative in α !



$$\text{Diagram 5a} = \text{Diagram 5a1} + \text{Diagram 5a2} + \text{Diagram 5a3}$$

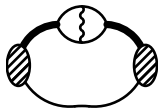
$$\text{Diagram 5b} = \text{Diagram 5b1} + \text{Diagram 5b2} + \text{Diagram 5b3}$$

$$\Delta E = \langle \psi | V_C | \psi \rangle$$



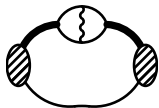
SK, Grießhammer, Hammer, J. Phys. G 42 045101 (2015)

Coulomb bubble divergence



- an additional diagram is logarithmically divergent. . .
- . . . but this divergence comes from the photon-bubble subdiagram!


Coulomb bubble divergence



- an additional diagram is logarithmically divergent. . .
- . . . but this divergence comes from the photon-bubble subdiagram!

Strategy

SK, Griethammer, Hammer, van Kolck, JPG **43** 055106 (2016), 1508.05085 [nucl-th]

- 1 isolate divergence: 
- 2 take the leading-order 1S_0 in the **unitarity limit!**
$$a_{1S_0} = -23.7 \approx \infty \rightsquigarrow 1/a_{1S_0} \approx 0$$
- 3 include divergent diagram together with finite a_{1S_0}

$$\text{---} \text{---} \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} \text{---} \text{---} = \text{finite}$$
The equation shows the sum of two diagrams. The first diagram is a photon bubble (a circle with a crack) connected to two external lines. The second diagram is a diamond-shaped regulator connected to two external lines. The result is labeled as 'finite'.

A new expansion

Take the leading-order singlet channel in the unitarity limit!

$$a_t = -23.4 \approx \infty !$$

$$\sigma_t = \sigma_t^{(0)} + \sigma_t^{(1)}$$

$$\boxed{\sigma_t^{(0)} - 2\Lambda/\pi = 0}$$

$$i\Delta_t^{(0)}(p_0, \mathbf{p}) \frac{-i}{\sigma_t + y_t^2 I_0(p_0, \mathbf{p})} \rightarrow \frac{-i}{\sqrt{\frac{\mathbf{p}^2}{4} - M_N p_0 - i\epsilon}}$$

A new expansion

Take the leading-order singlet channel in the unitarity limit!


$$a_t = -23.4 \approx \infty !$$


$$\sigma_t = \sigma_t^{(0)} + \sigma_t^{(1)}$$

$$\sigma_t^{(0)} - 2\Lambda/\pi = 0$$

NLO corrections

- scattering length: $\sigma_t^{(1)} \sim -1/a_t$
- effective range: $c_t^{(1)} = \frac{M_N r_0 t}{2}$


$$\sim -i\sigma_{t(pp)}^{(1)}$$


$$\sim -ic_t^{(1)}$$



$$i\Delta_t^{(1)}(p_0, \mathbf{p}) = i\Delta_t^{(0)}(p_0, \mathbf{p}) \times \left[-i\sigma_t^{(1)} - ic_t^{(1)} \left(p_0 - \frac{\mathbf{p}^2}{4M_N} \right) \right] \times i\Delta_t^{(0)}(p_0, \mathbf{p})$$

A new expansion

Take the leading-order singlet channel in the unitarity limit!


$$a_t = -23.4 \approx \infty !$$

$$\sigma_t = \sigma_t^{(0)} + \sigma_t^{(1)}$$


$$\sigma_t^{(0)} - 2\Lambda/\pi = 0$$

NLO corrections

- scattering length: $\sigma_t^{(1)} \sim -1/a_t$
- effective range: $c_t^{(1)} = \frac{M_N r_{0t}}{2}$



$\sim -i\sigma_{t(pp)}^{(1)}$



$\sim -ic_t^{(1)}$



$$i\Delta_t^{(1)}(p_0, \mathbf{p}) = i\Delta_t^{(0)}(p_0, \mathbf{p}) \times \left[-i\sigma_t^{(1)} - ic_t^{(1)} \left(p_0 - \frac{\mathbf{p}^2}{4M_N} \right) \right] \times i\Delta_t^{(0)}(p_0, \mathbf{p})$$

- new 1S_0 LO is **isospin-symmetric** and parameter-free
- **allows matching between perturbative and non-perturbative Coulomb regimes**

Divergence dissection

Vanasse, Egolf, Kerin, SK, Springer, PRC **89** 064003 (2014)

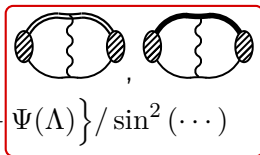
Look at structure of p - d three-body force:

$$H_{0,1}^{(\alpha)}(\Lambda) = h_I^{(\alpha)}(\Lambda) + h_{\kappa}^{(\alpha)}(\Lambda)$$

$$h_I^{(\alpha)}(\Lambda) = -\frac{3\pi(1+s_0^2)}{16} \left\{ \frac{1}{12}(r_{p-p} - r_t)\Lambda [1 - \dots] \right\}$$

+ various terms $\sim \log \Lambda$, all $\propto (r_{p-p} - r_t)$ or $\propto (\gamma_{p-p} - \gamma_t)$ $\left. \right\} / \sin^2(\dots)$

$$h_{\kappa}^{(\alpha)} = -\frac{\sqrt{3}\kappa\pi(1+s_0^2)}{48} \left\{ \text{various terms} \sim \log \Lambda \right.$$



that do not vanish in the isospin limit + $\Psi(\Lambda)$ $\left. \right\} / \sin^2(\dots)$

- pieces associated with $r_{p-p} \neq r_t$ and $\alpha\rho_d, \alpha r_t$
- these have been relegated to a higher order in the new scheme!

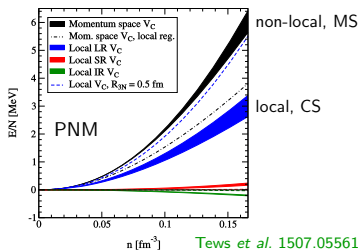
SK, Griebhammer, Hammer, van Kolck, JPG **43** 055106 (2016), 1508.05085 [nucl-th]

- but an otherwise remaining log-divergence is absent!

power counting \leftrightarrow regulators?

power counting \leftrightarrow regulators?

similar issues in chiral EFT!



pionless EFT can study this question in a clean scenario without other complications!

- chiral EFT with Weinberg counting does not absorb all cutoff dependence in counterterms
- pionless EFT formulated with RG invariance
- different schemes can still give different runnings for LECs

Pionless two- and three-body regularization

Momentum-space overview

- typically used in pionless EFT with dibaryons:
dim. reg. (PDS) in two-body sector, cutoff in three-body equation

$$\Delta(p) \sim \frac{1}{\underbrace{-1/a}_{=\sigma-\mu} + \sqrt{\frac{\mathbf{p}^2}{4} - M_N p_0 - i\epsilon}} \quad , \quad T(E; k, p) = K(E; k, p) + \int_0^\Lambda dq^2 K(E; k, q) D(E; q) T(E; q, p)$$

- as long as everything converges, this should not matter

Bedaque *et al.*, NPA **676** 357 (2000)

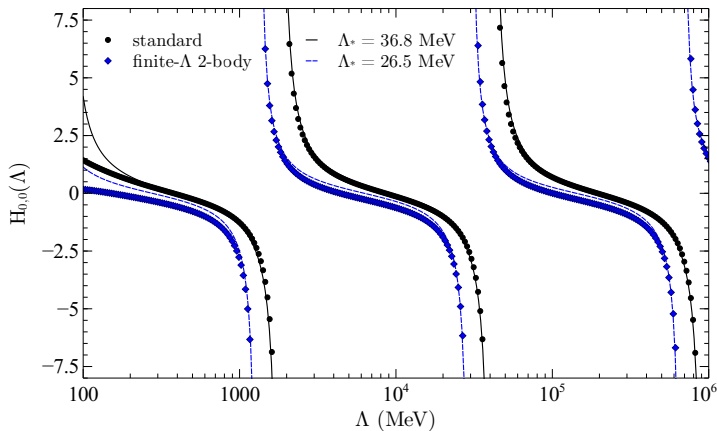
- **to check this:** use cutoff regulator for two-body system,
keep finite- Λ corrections

Vanasse, Coulomb RRTF Part I (Jan. 2016)

- Faddeev calculations use Gaussian regulator functions for two-
and three-body sector

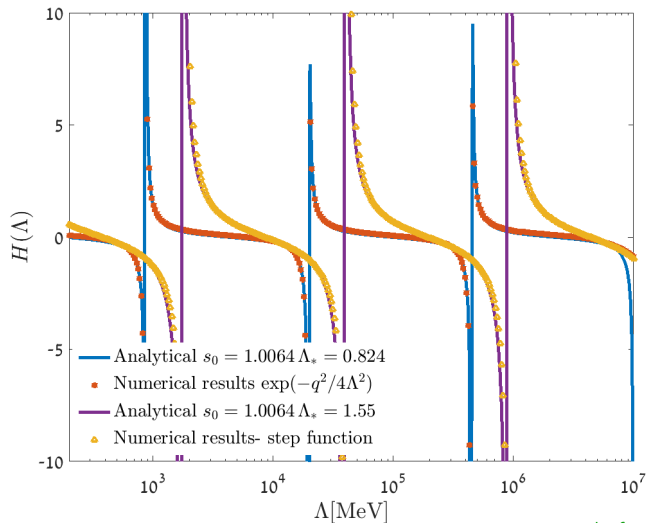
Platter *et al.*, PRA **70** 052101 (2004), PLB **607** 254 (2005), ...

Three-body force with two-body cutoff regulator



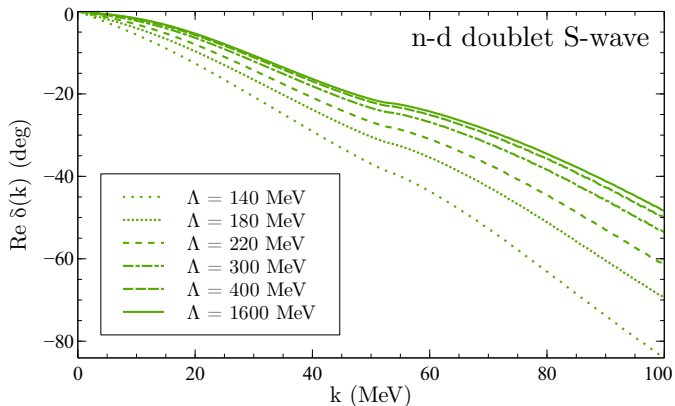
$$\Delta(p_0, \mathbf{p}) \sim \frac{1}{-\gamma + \sqrt{\frac{\mathbf{p}^2}{4} - M_N p_0 - i\epsilon} - \frac{2}{\pi} \sqrt{\frac{\mathbf{p}^2}{4} - M_N p_0 - i\epsilon} \arctan \left(\sqrt{\frac{\mathbf{p}^2}{4} - M_N p_0 - i\epsilon/\Lambda} \right)}$$

Three-body force with Gaussian regulators

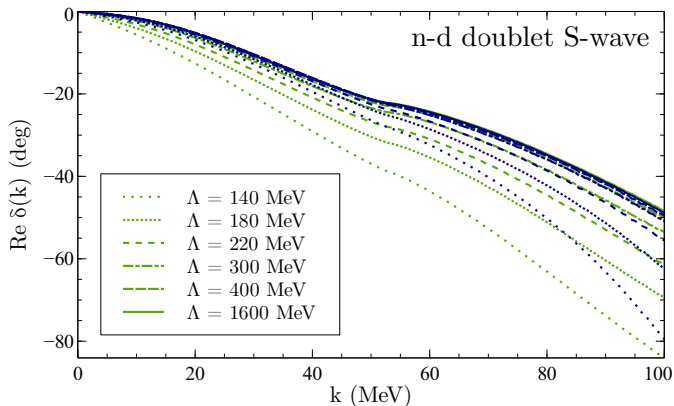


plot from H. DeLeon, Jan. 2016

nd doublet phase shift with two-body cutoff regulator



nd doublet phase shift with two-body cutoff regulator



$$\Delta(p_0, \mathbf{p}) \sim \frac{1}{-\gamma + \sqrt{\frac{\mathbf{p}^2}{4} - M_N p_0 - i\epsilon} - \frac{2}{\pi} \sqrt{\frac{\mathbf{p}^2}{4} - M_N p_0 - i\epsilon} \arctan \left(\sqrt{\frac{\mathbf{p}^2}{4} - M_N p_0 - i\epsilon} / \Lambda \right)}$$

Divergence dissection

Vanasse, Egolf, Kerin, SK, Springer, PRC **89** 064003 (2014)

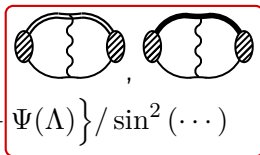
Look at structure of p - d three-body force:

$$H_{0,1}^{(\alpha)}(\Lambda) = h_I^{(\alpha)}(\Lambda) + h_{\kappa}^{(\alpha)}(\Lambda)$$

$$h_I^{(\alpha)}(\Lambda) = -\frac{3\pi(1+s_0^2)}{16} \left\{ \frac{1}{12}(r_{p-p} - r_t)\Lambda [1 - \dots] \right\}$$

+ various terms $\sim \log \Lambda$, all $\propto (r_{p-p} - r_t)$ or $\propto (\gamma_{p-p} - \gamma_t)$ $\left. \right\} / \sin^2(\dots)$

$$h_{\kappa}^{(\alpha)} = -\frac{\sqrt{3}\kappa\pi(1+s_0^2)}{48} \left\{ \text{various terms} \sim \log \Lambda \right.$$



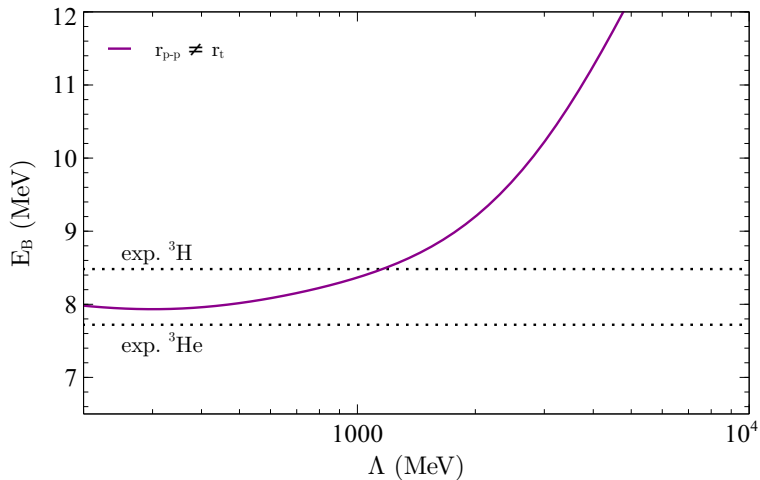
that do not vanish in the isospin limit + $\Psi(\Lambda)$ $\left. \right\} / \sin^2(\dots)$

- pieces associated with $r_{p-p} \neq r_t$ and $\alpha\rho_d, \alpha r_t$
- these have been related to a higher order in the new scheme!

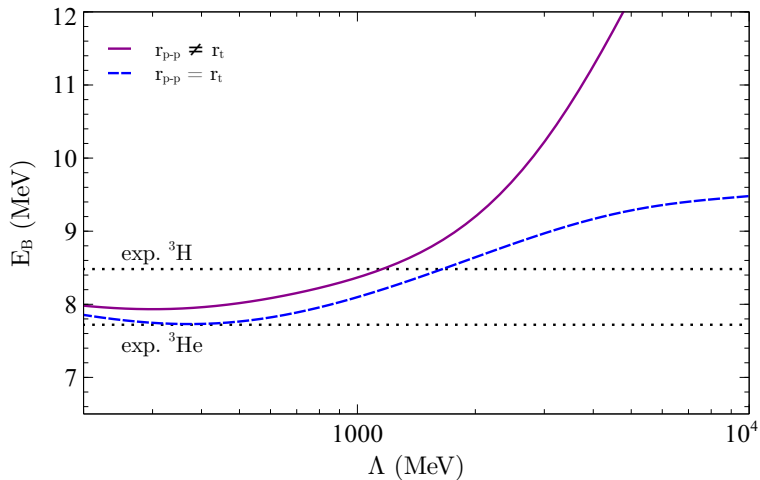
SK, Griebhammer, Hammer, van Kolck, JPG **43** 055106 (2016), 1508.05085 [nucl-th]

- **but an otherwise remaining log-divergence is absent!**

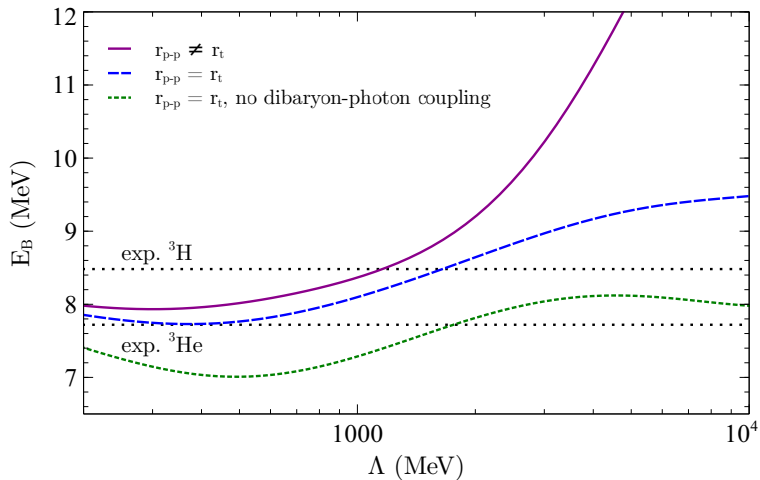
NLO divergence contributions



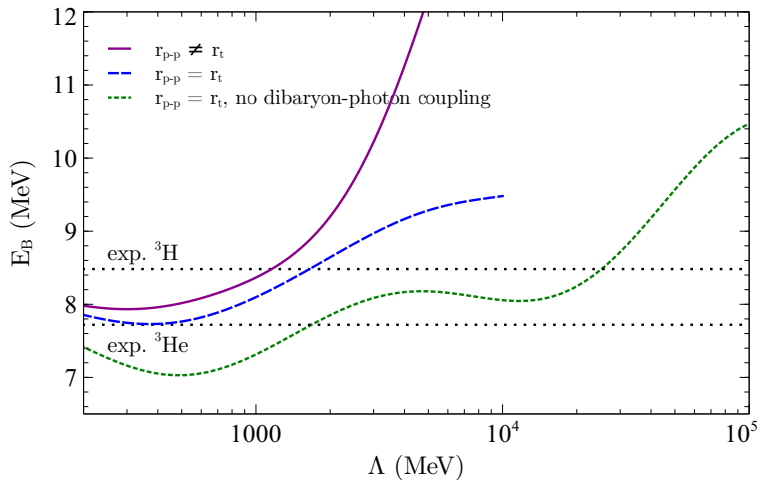
NLO divergence contributions



NLO divergence contributions



NLO divergence contributions



NLO divergence from isospin breaking (no Coulomb)

- fit three-nucleon force with isospin-symmetric 1S_0 dibaryon
- use this in calculation with separate n - n channel

