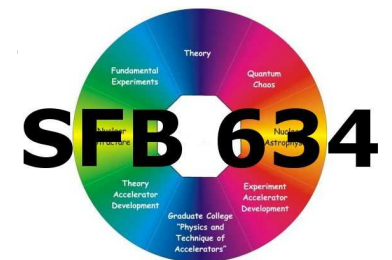


Mean-Field and Pairing Calculations in the UCOM Framework

Heiko Hergert

Institut für Kernphysik, TU Darmstadt



Overview

- Reminder: UCOM and SRG Basics
- Hartree-Fock and Perturbation Theory with V_{UCOM}
- Pairing in the UCOM Framework
 - Hartree-Fock-Bogoliubov & Projection
 - Quasiparticle RPA
- Conclusions

Current Approaches

Phenomenology

binding energies, radii
nuclear matter

QCD

Current Approaches

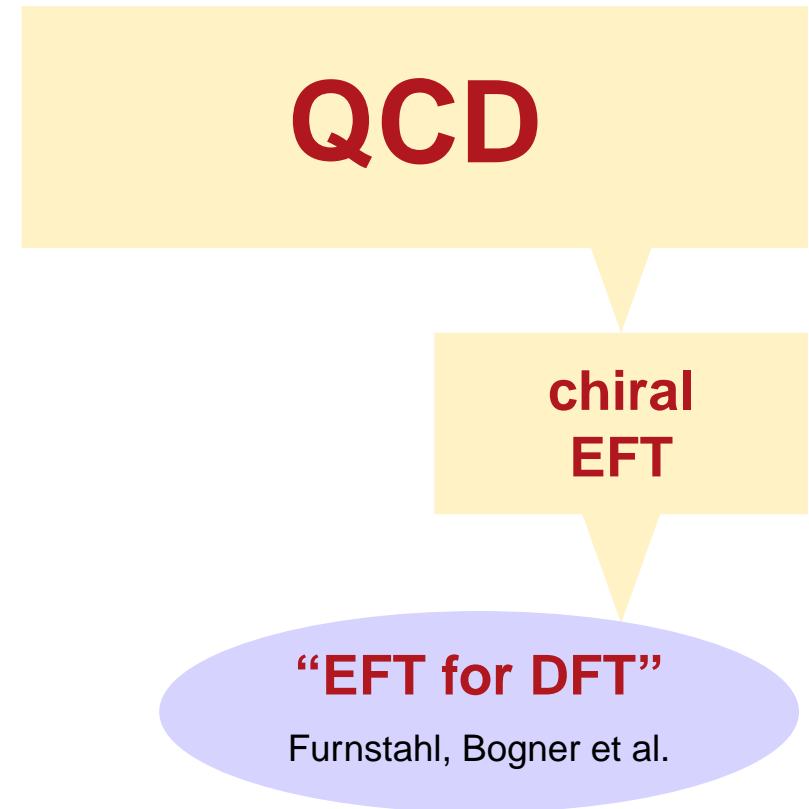
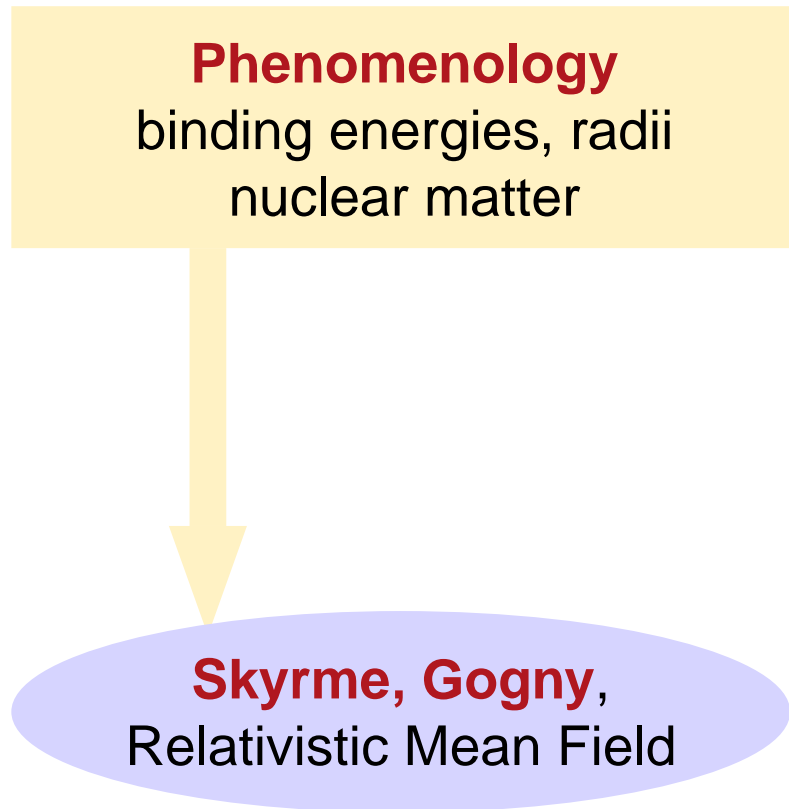
Phenomenology
binding energies, radii
nuclear matter



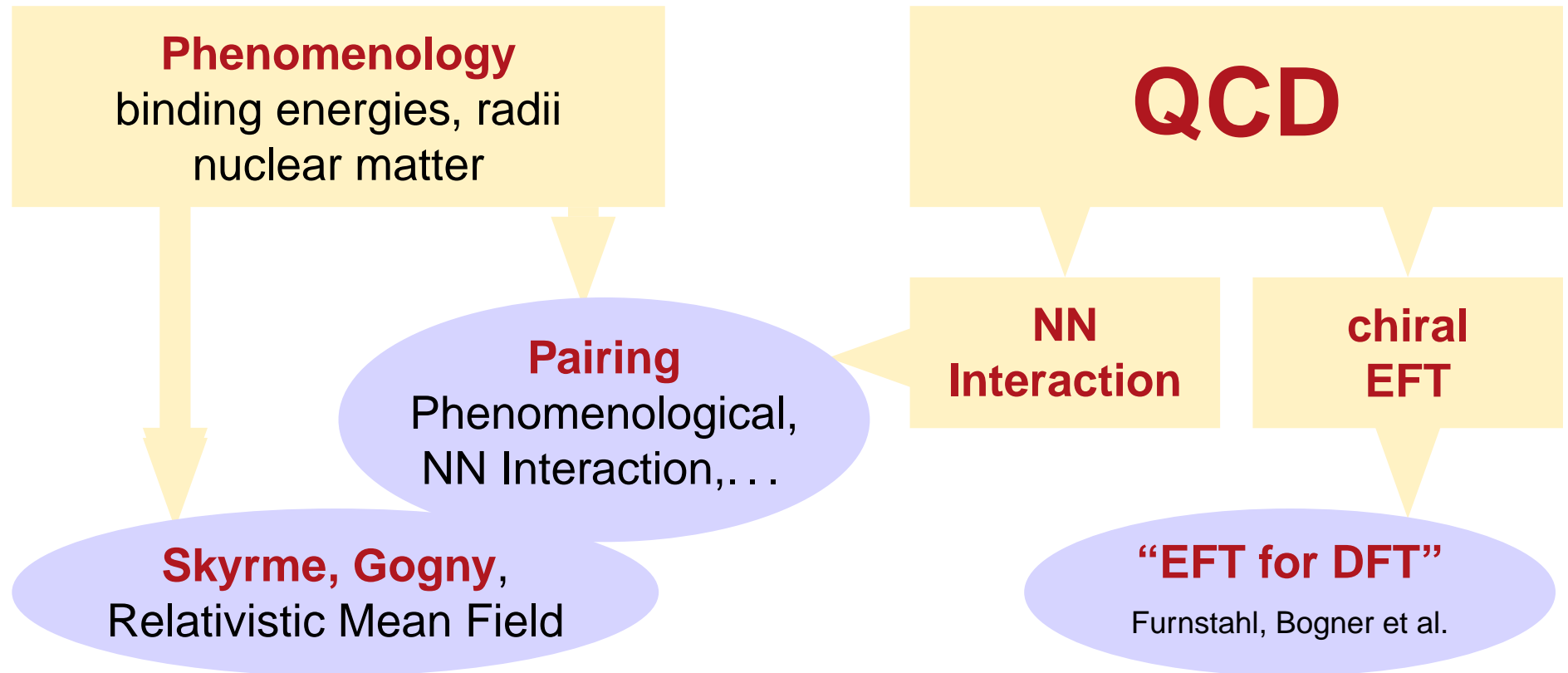
Skyrme, Gogny,
Relativistic Mean Field

QCD

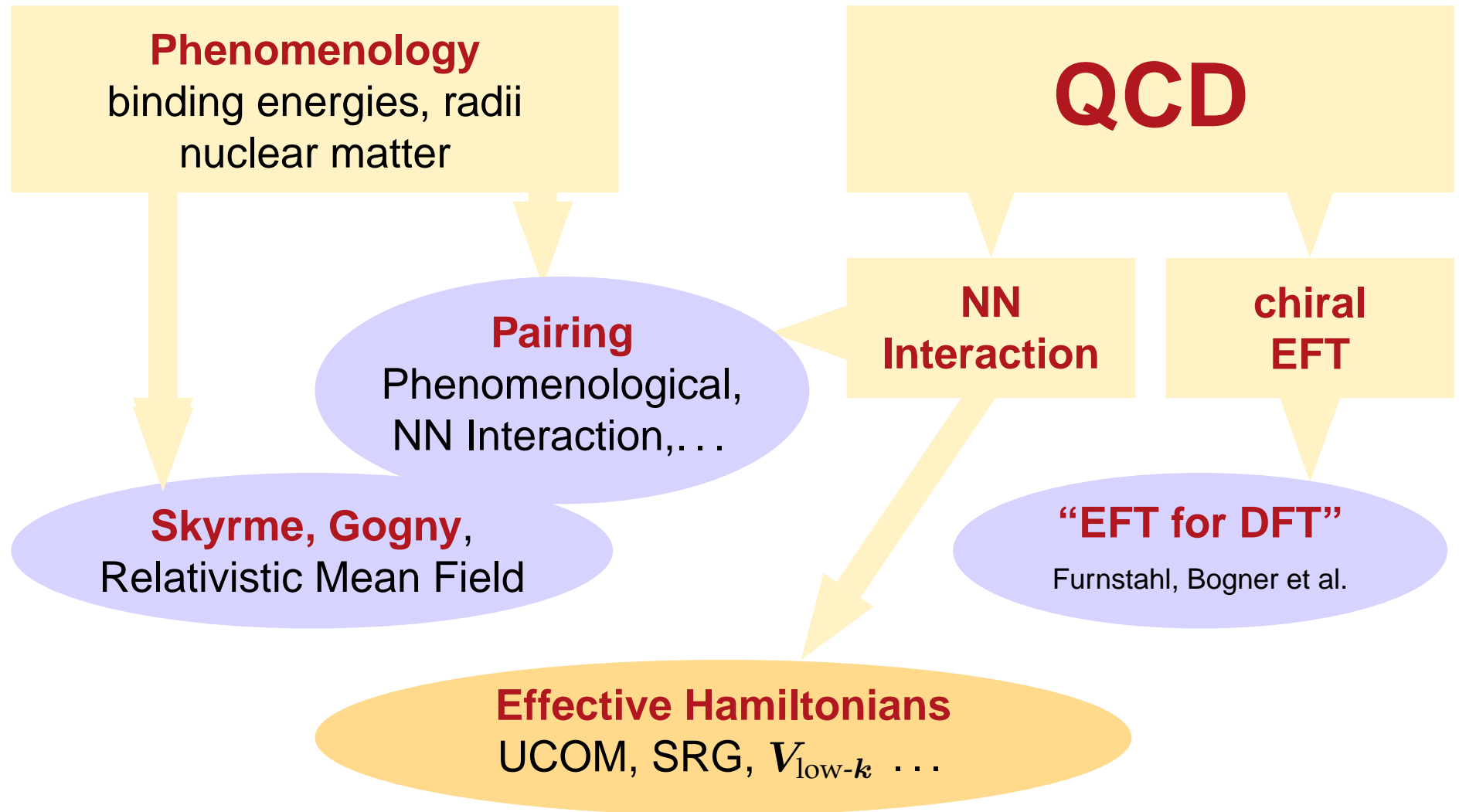
Current Approaches



Current Approaches



Current Approaches



Issues of DFT

- ✗ description of **exotic nuclei** and **spectroscopic observables** is lacking
- ✗ inconsistent treatment of **particle-hole** and **particle-particle** channel treated if separate forces are used
- ✗ local density approximation, gradient expansion essentially **uncontrolled**
- ✗ technical & conceptual difficulties due to **non-analytic terms**

☞ Dobaczewski et al., Phys. Rev. **C76**, 054315 (2007)

Duguet, Lacroix, Bender et al., **arXiv**:0809.2041, 0809.2045, 0809.2049

- ✗ fits of phenomenological functionals obscure underlying physics

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- ✗ fits of phenomenological functionals obscure underlying physics

consistent Hamiltonian approach
avoids most problems

UCOM and SRG Basics

Modern Effective Interactions

phase-shift equivalent interaction from unitary transformation of the Hamiltonian

Unitary Correlation Operator Method

■ transformed Hamiltonian

$$\tilde{H} = C_r^\dagger C_\Omega^\dagger H C_\Omega C_r$$

■ central correlations: radial shift

$$C_r = \exp\left(-i \sum_{i<j} g_{r,ij}[s(r_{ij})]\right)$$

■ tensor correlations: angular shift

$$C_\Omega = \exp\left(-i \sum_{i<j} g_{\Omega,ij}[\vartheta(r_{ij})]\right)$$

Similarity Renormalization Group

■ transformed Hamiltonian

$$\tilde{H}(\alpha) = C^\dagger(\alpha) H C(\alpha)$$

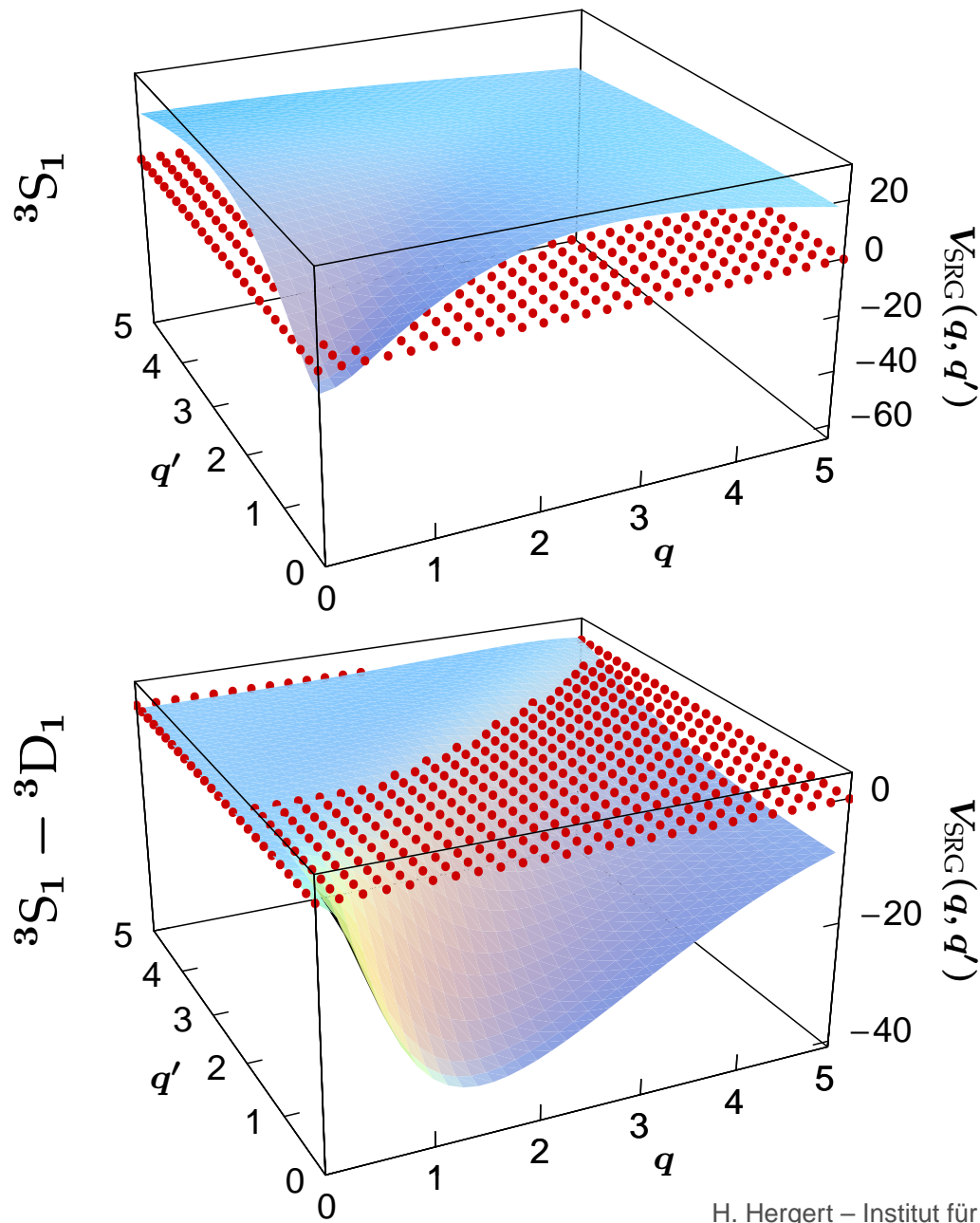
■ evolution via RG flow equation

$$\frac{d}{d\alpha} \tilde{H}(\alpha) = [\eta(\alpha), \tilde{H}(\alpha)]$$

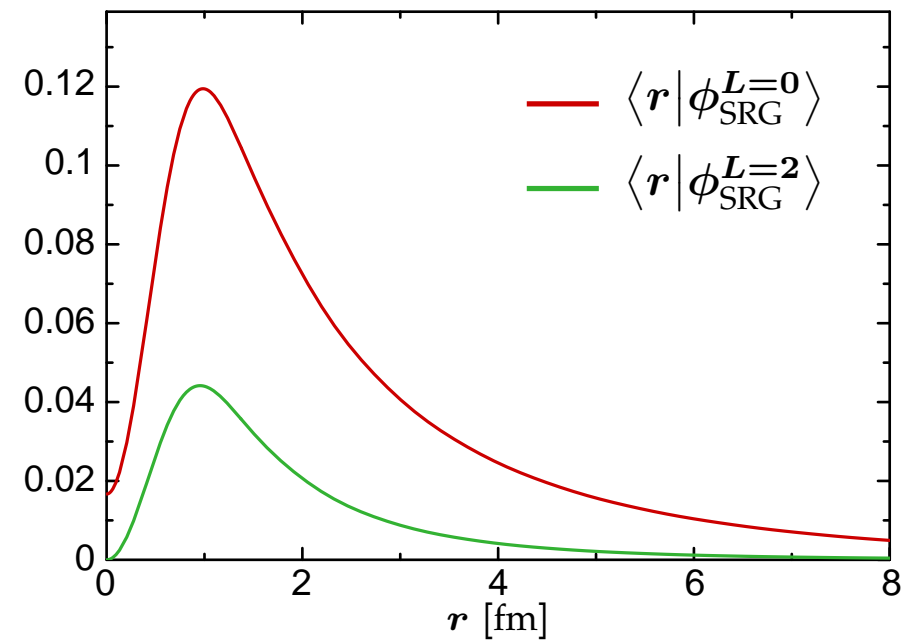
■ dynamical generator

$$\eta(\alpha) = \frac{1}{2\mu} [\vec{q}^2, \tilde{H}(\alpha)]$$

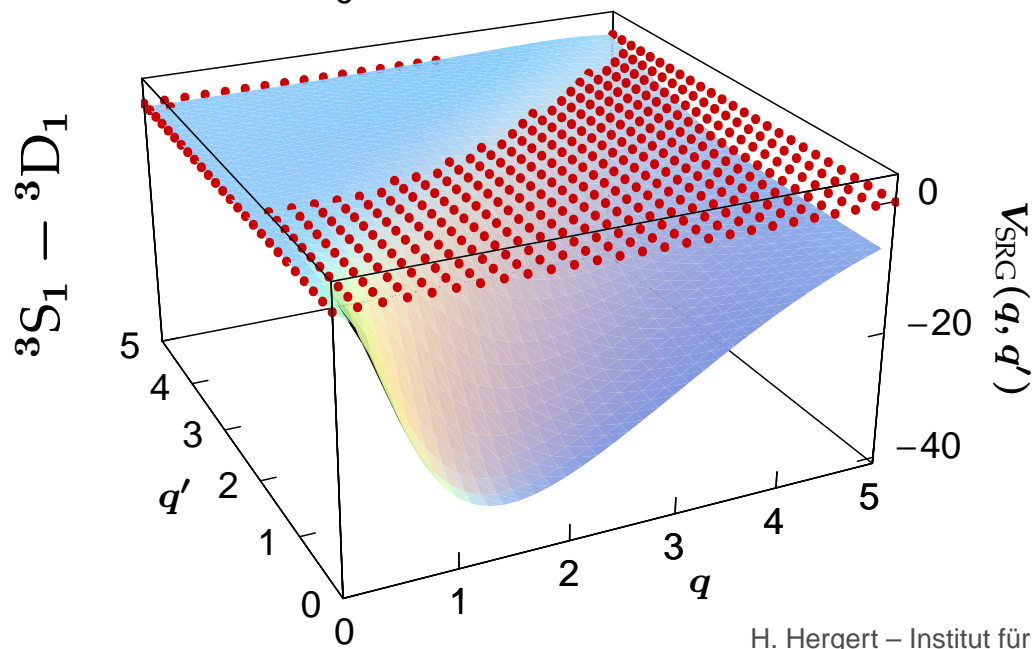
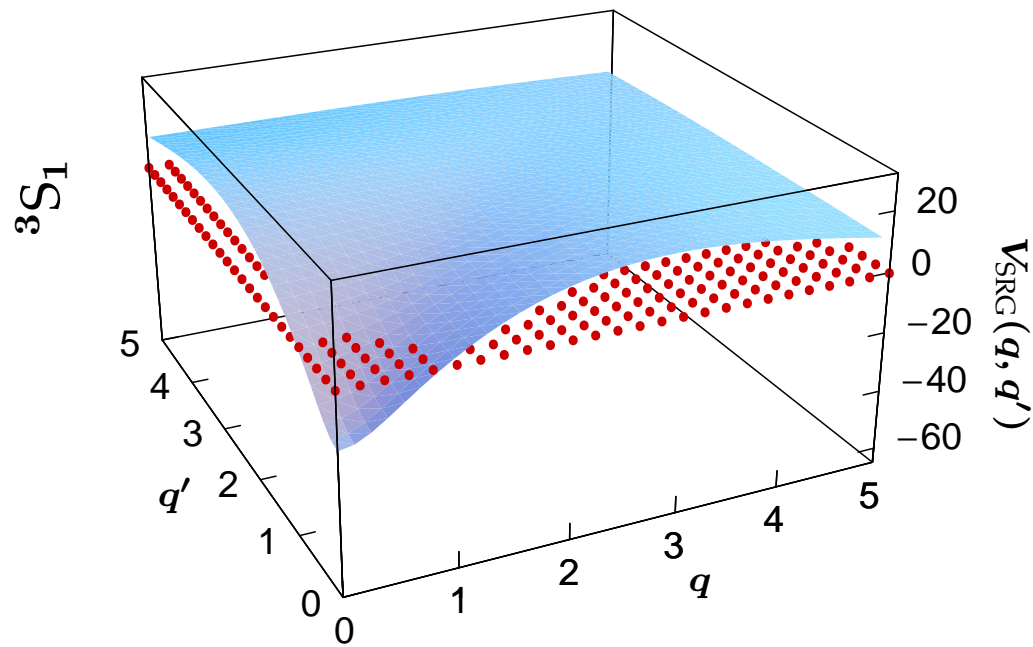
SRG Evolution: The Deuteron



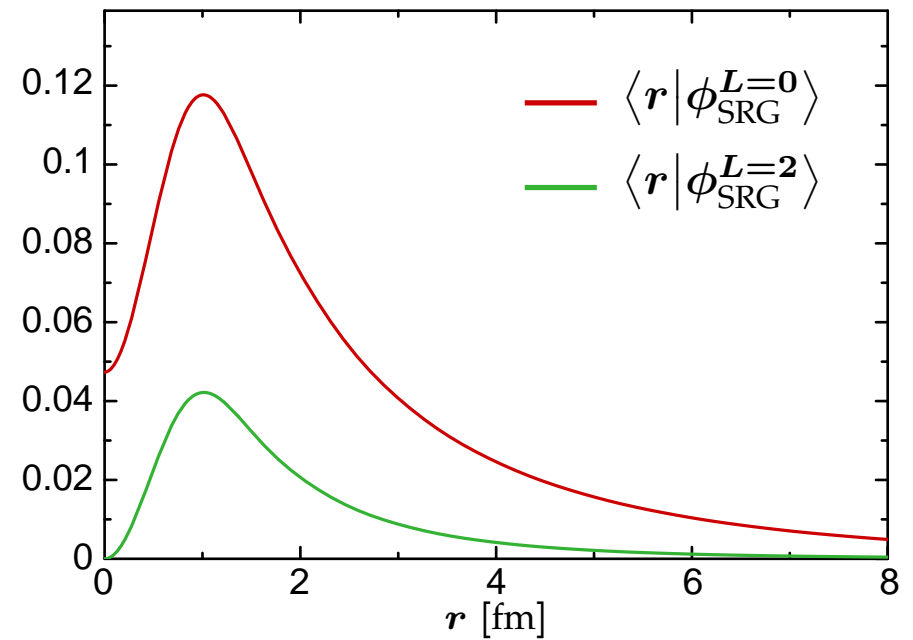
Argonne V18



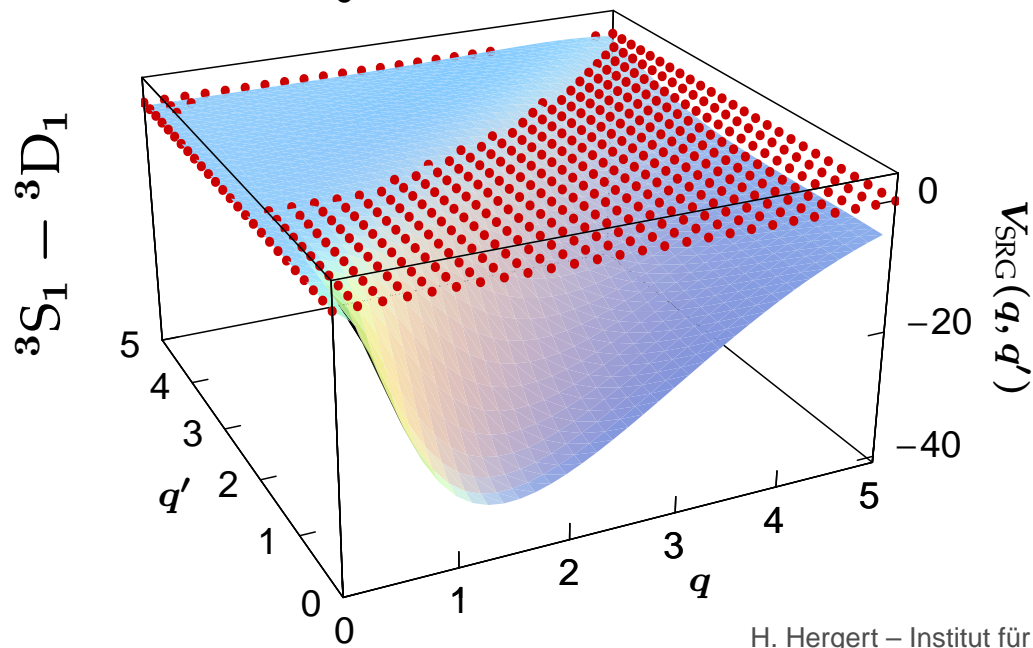
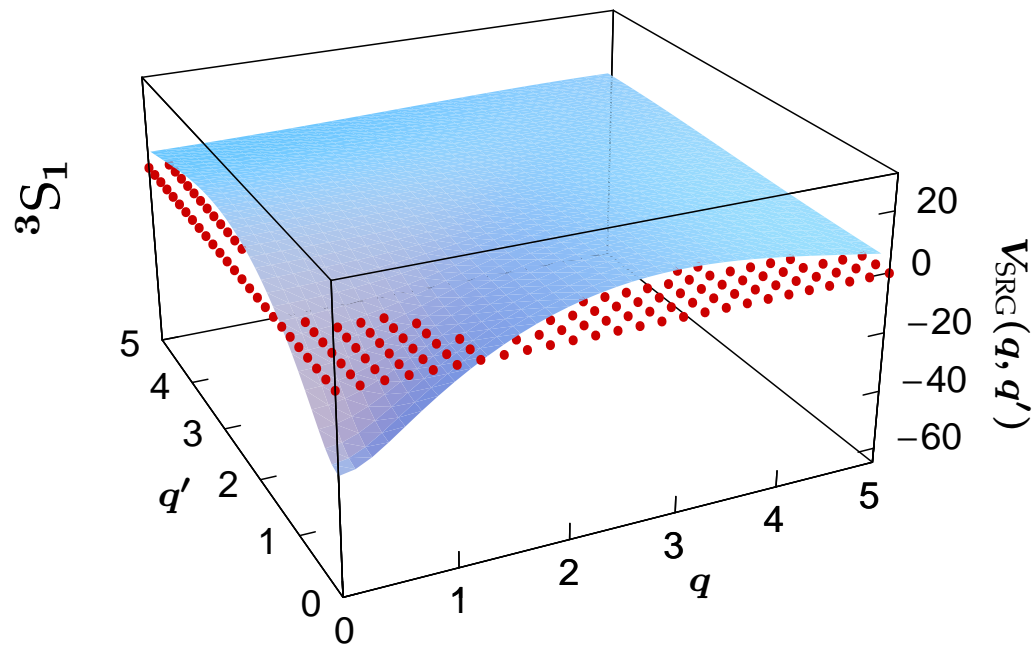
SRG Evolution: The Deuteron



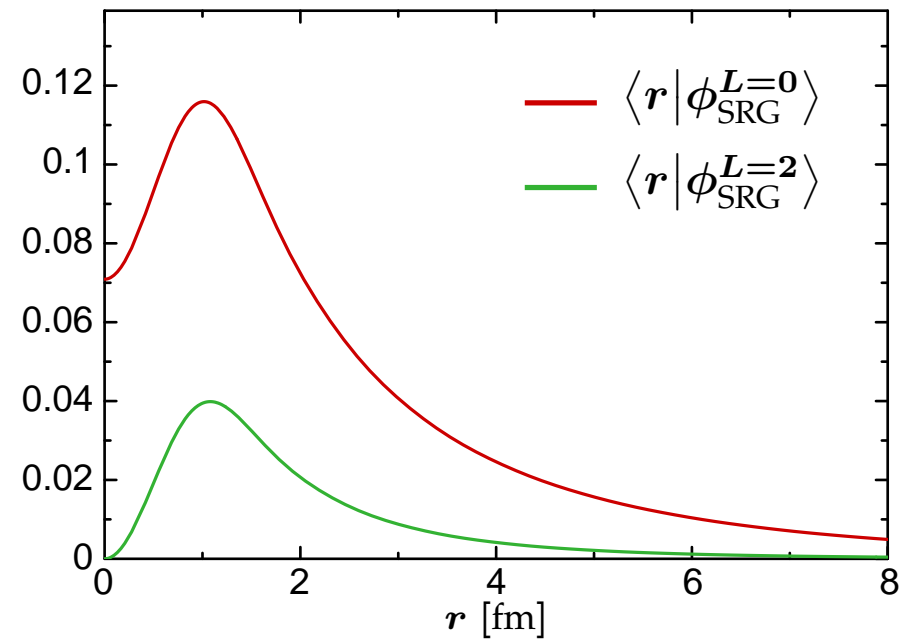
$$\alpha = 0.0004 \text{ fm}^4$$



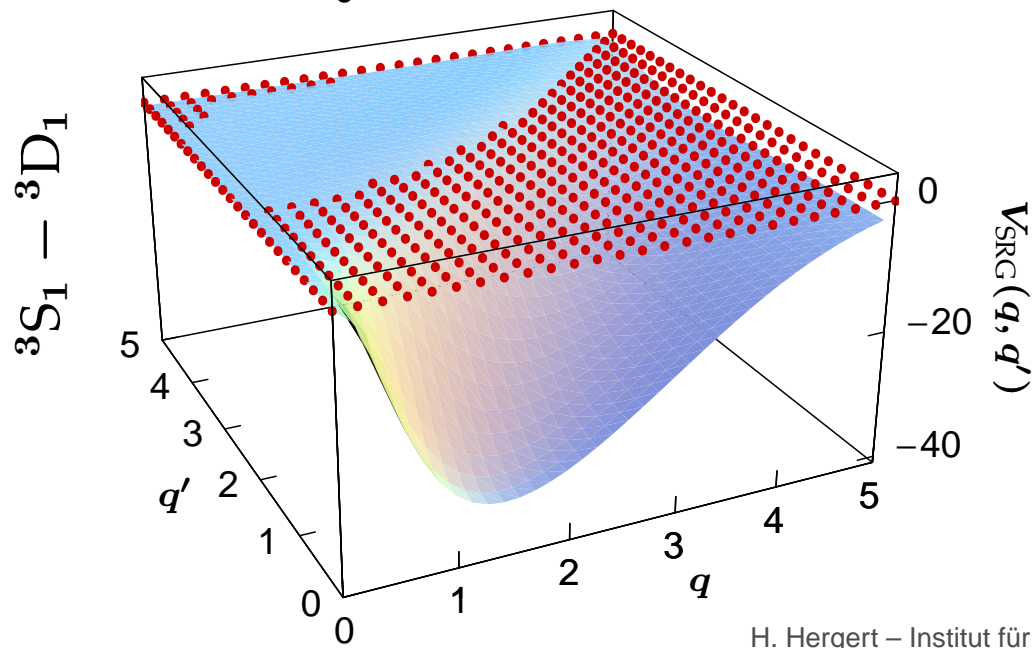
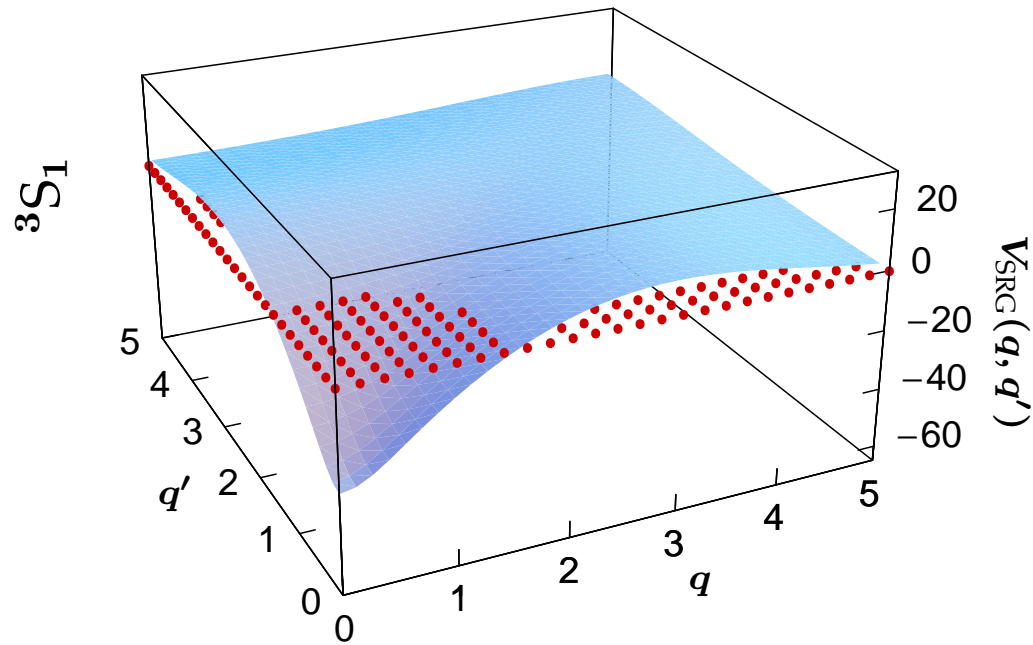
SRG Evolution: The Deuteron



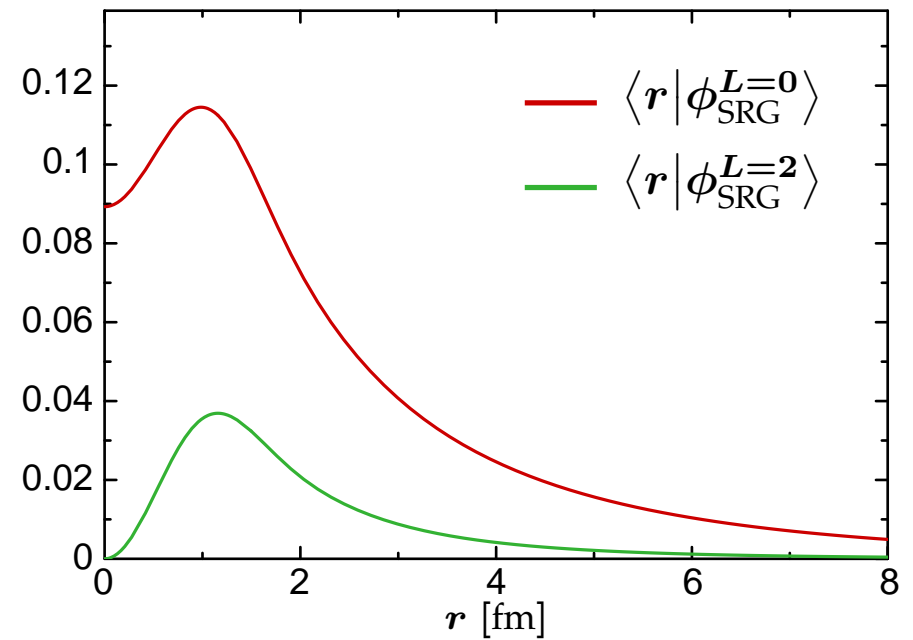
$$\alpha = 0.0010 \text{ fm}^4$$



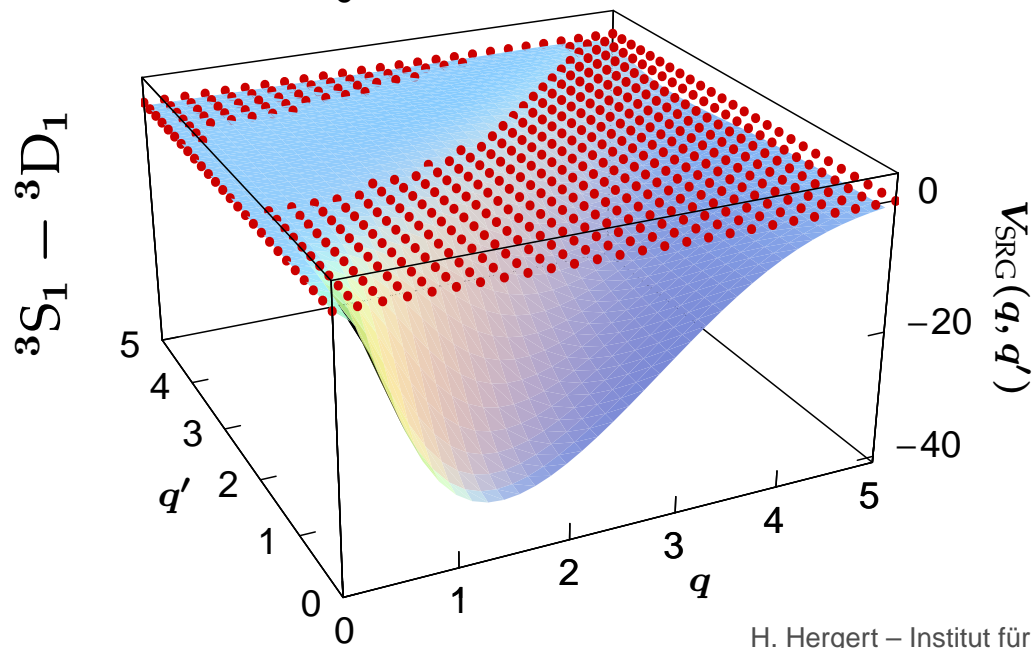
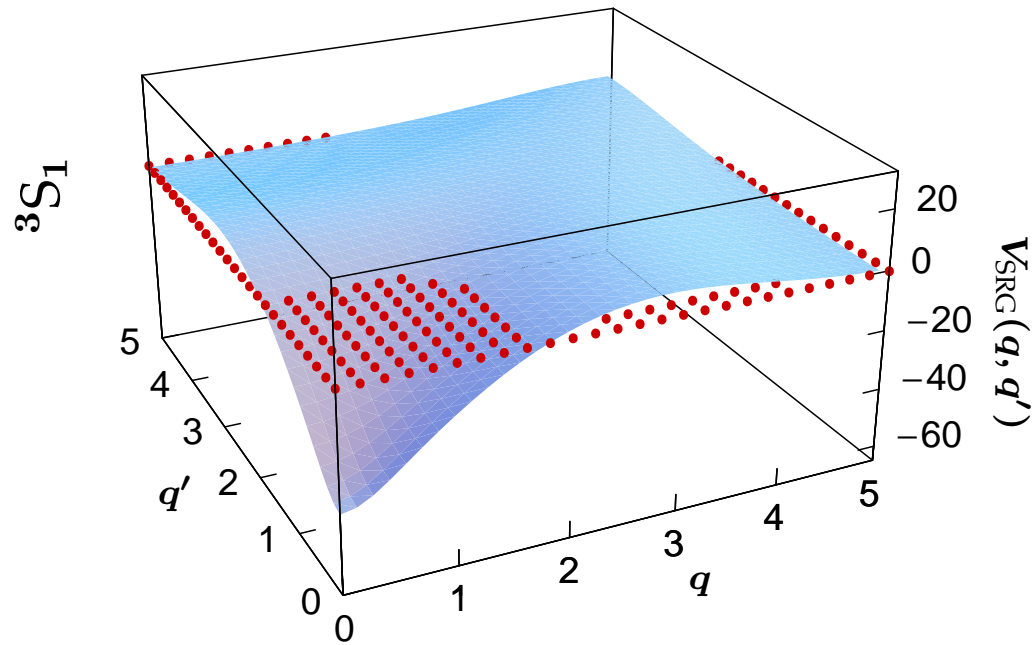
SRG Evolution: The Deuteron



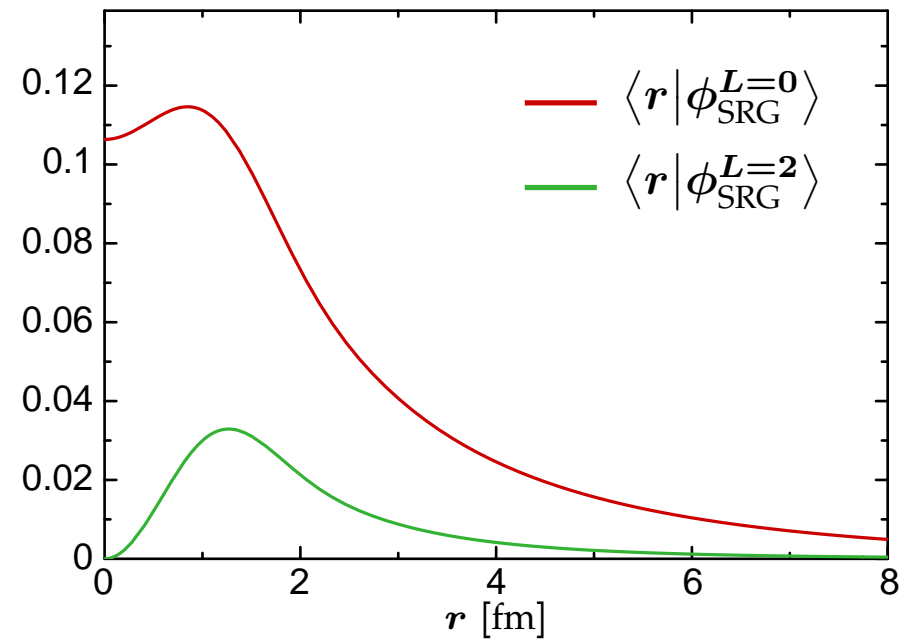
$$\alpha = 0.0020 \text{ fm}^4$$



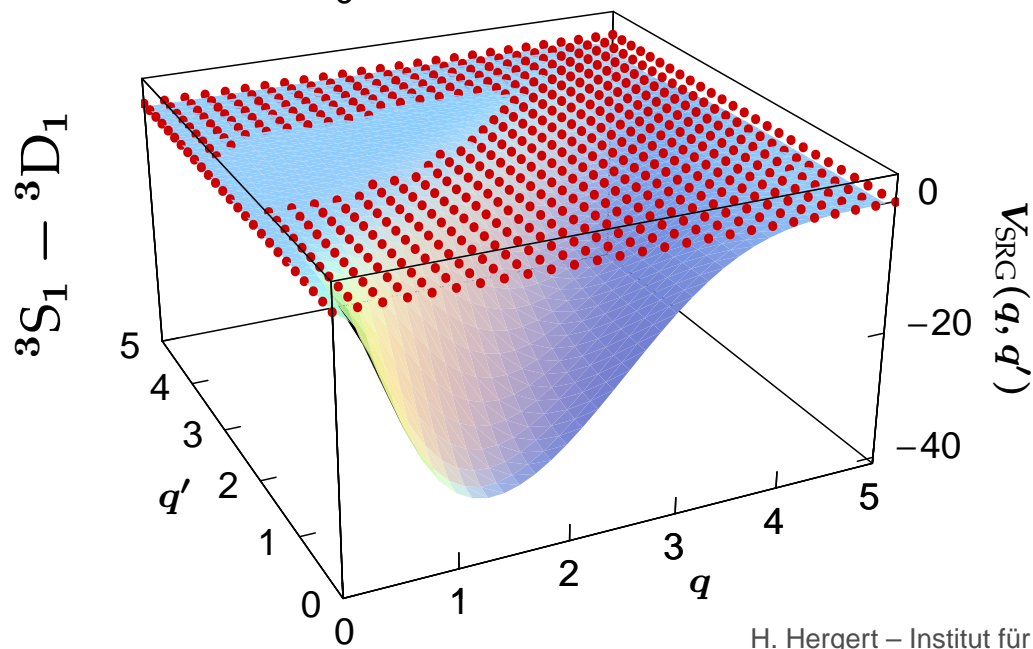
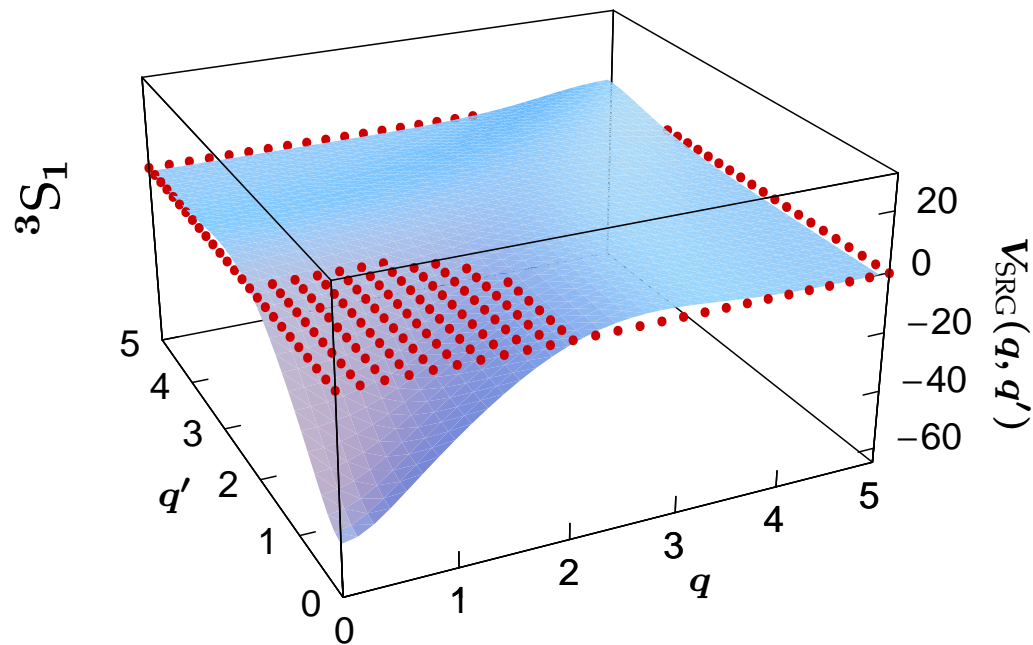
SRG Evolution: The Deuteron



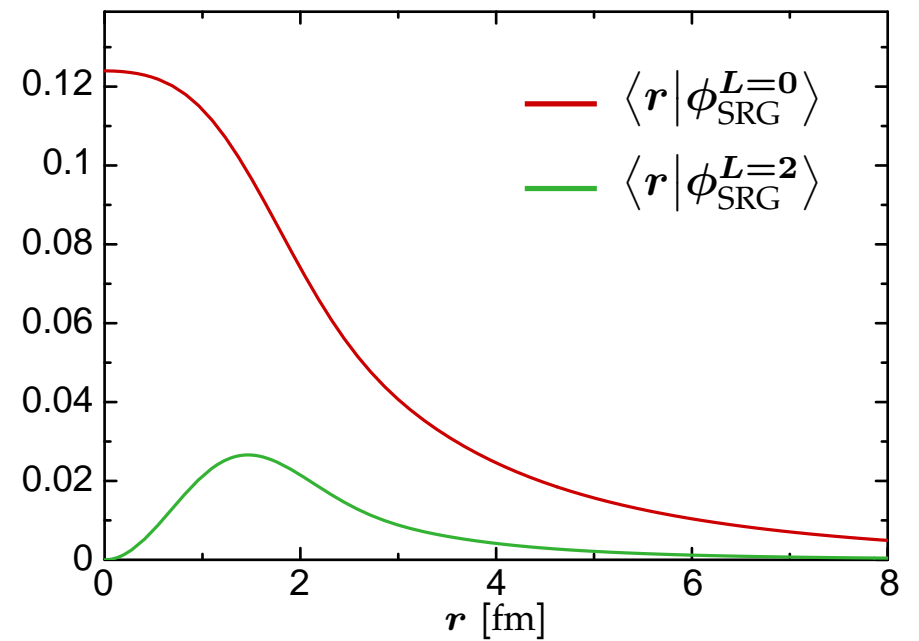
$$\alpha = 0.0040 \text{ fm}^4$$



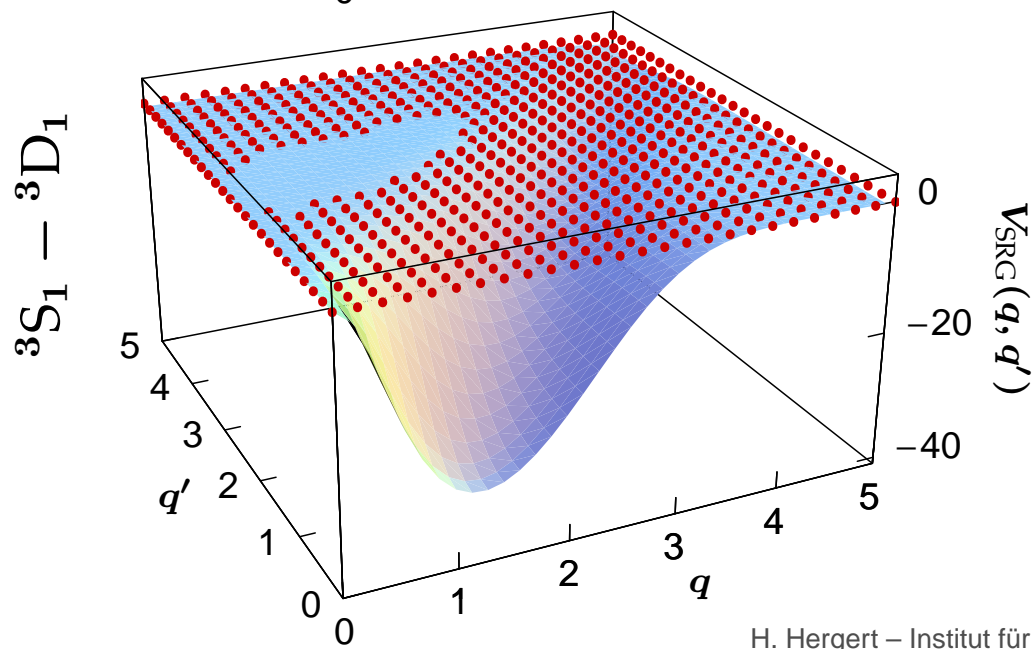
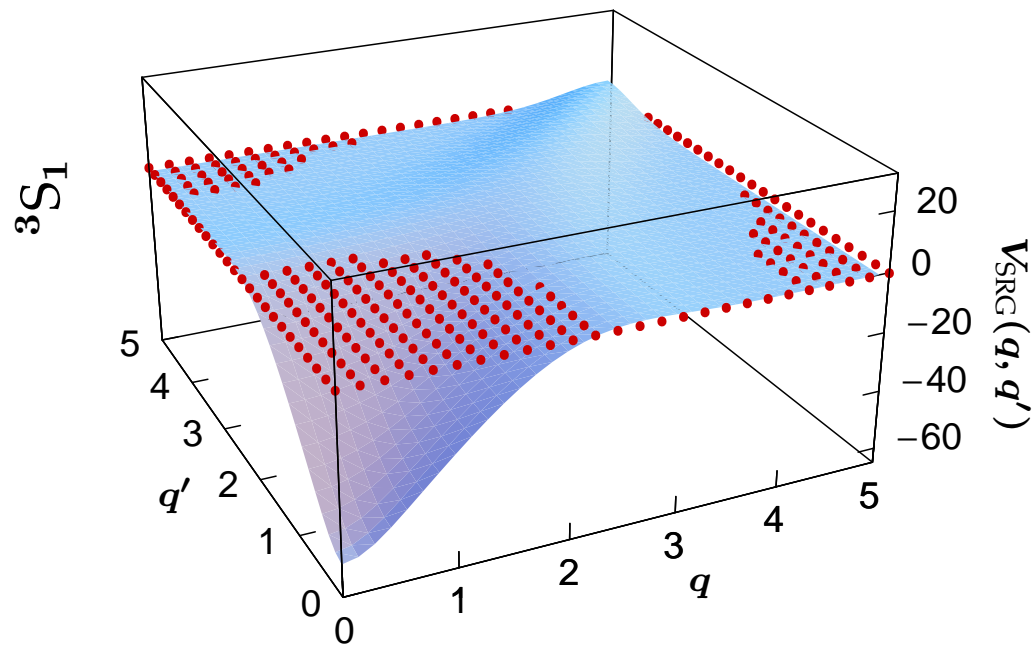
SRG Evolution: The Deuteron



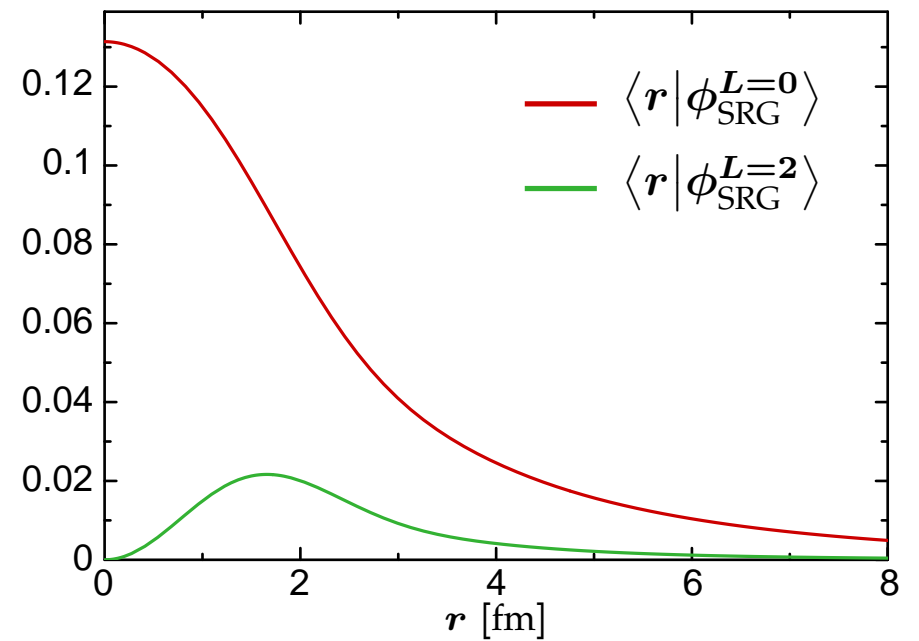
$$\alpha = 0.0100 \text{ fm}^4$$



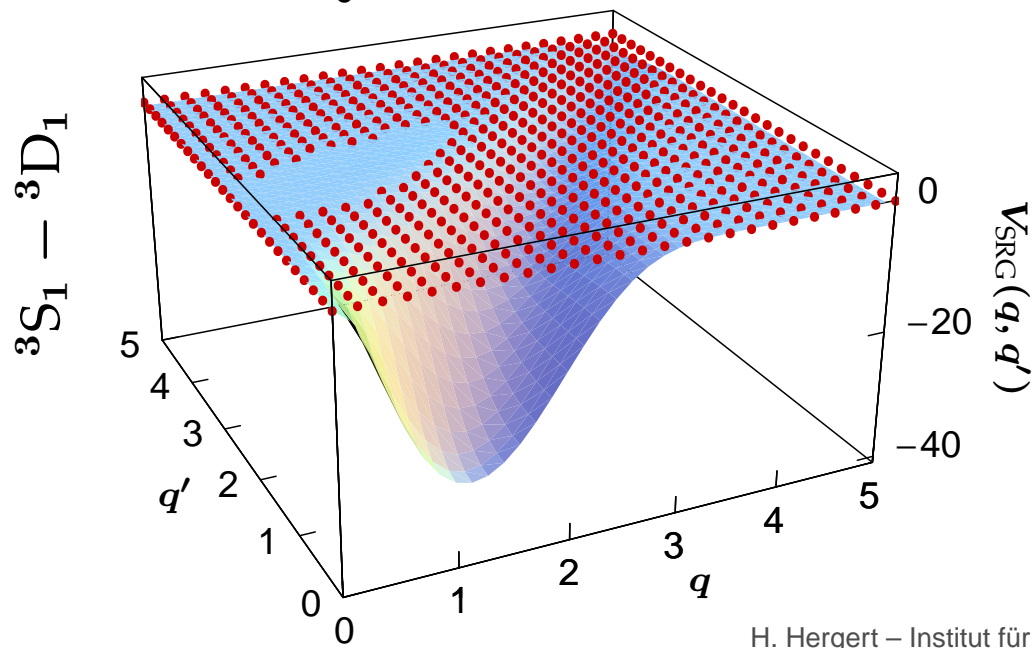
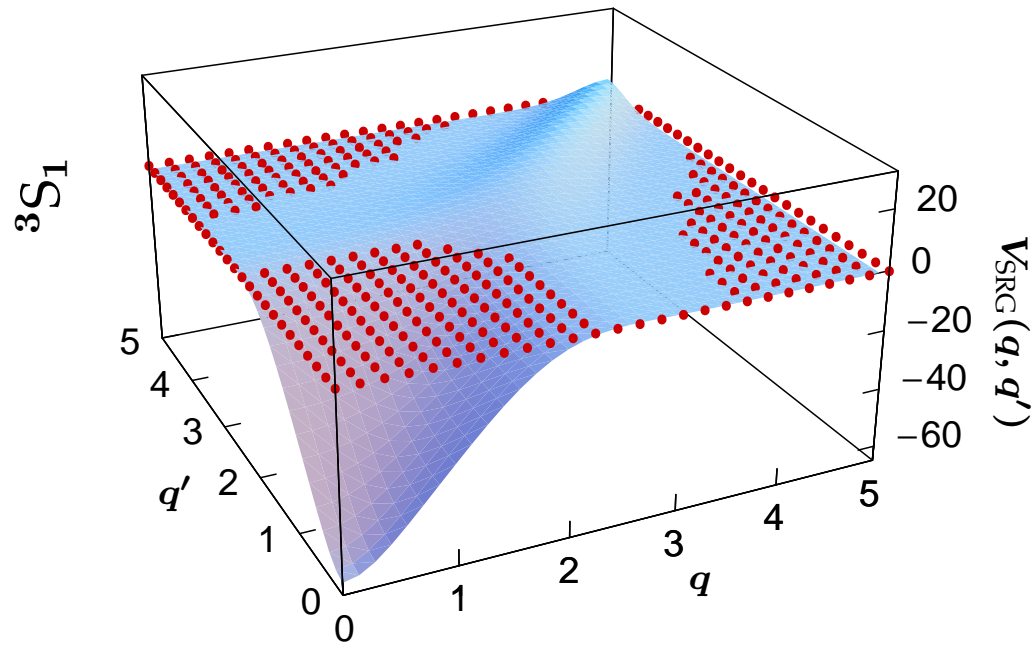
SRG Evolution: The Deuteron



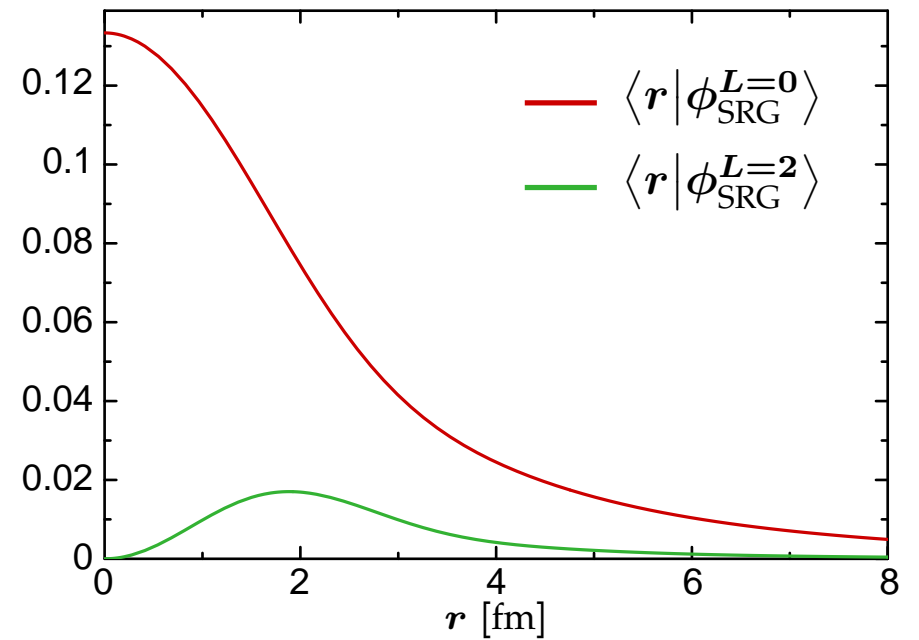
$$\alpha = 0.0200 \text{ fm}^4$$



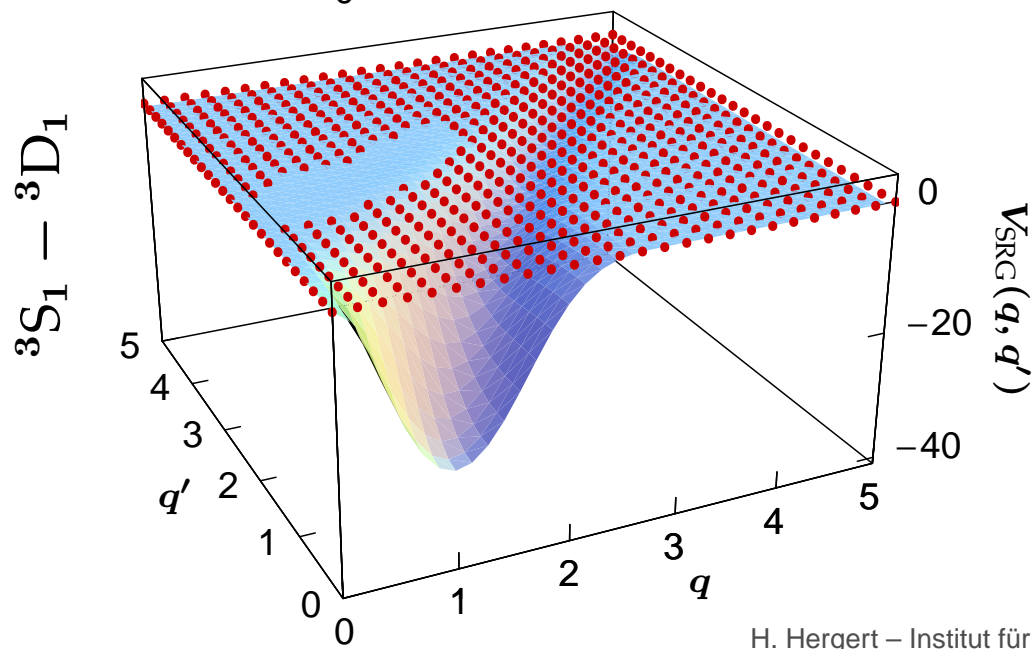
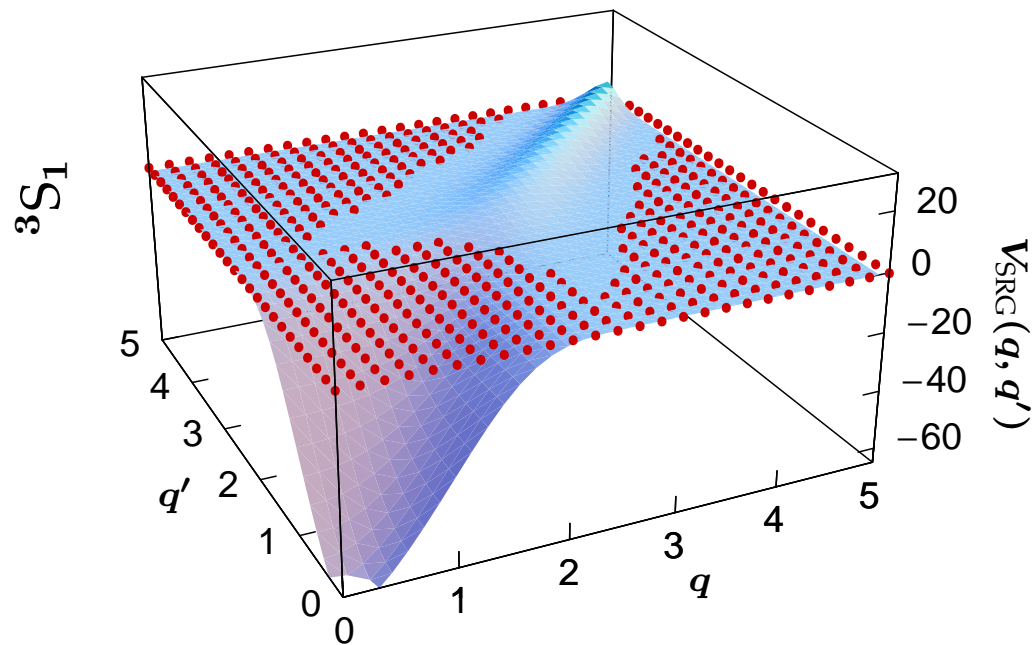
SRG Evolution: The Deuteron



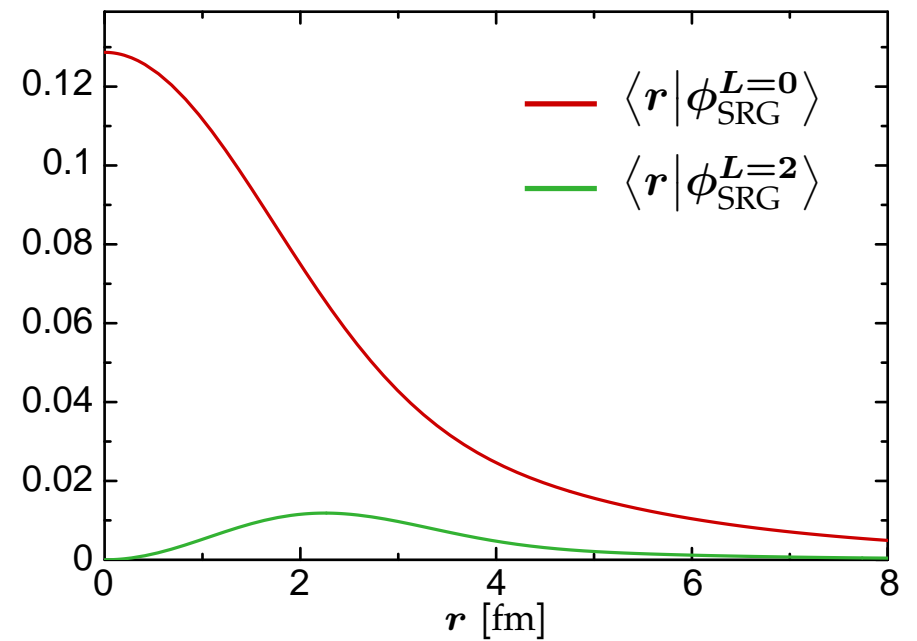
$$\alpha = 0.0400 \text{ fm}^4$$



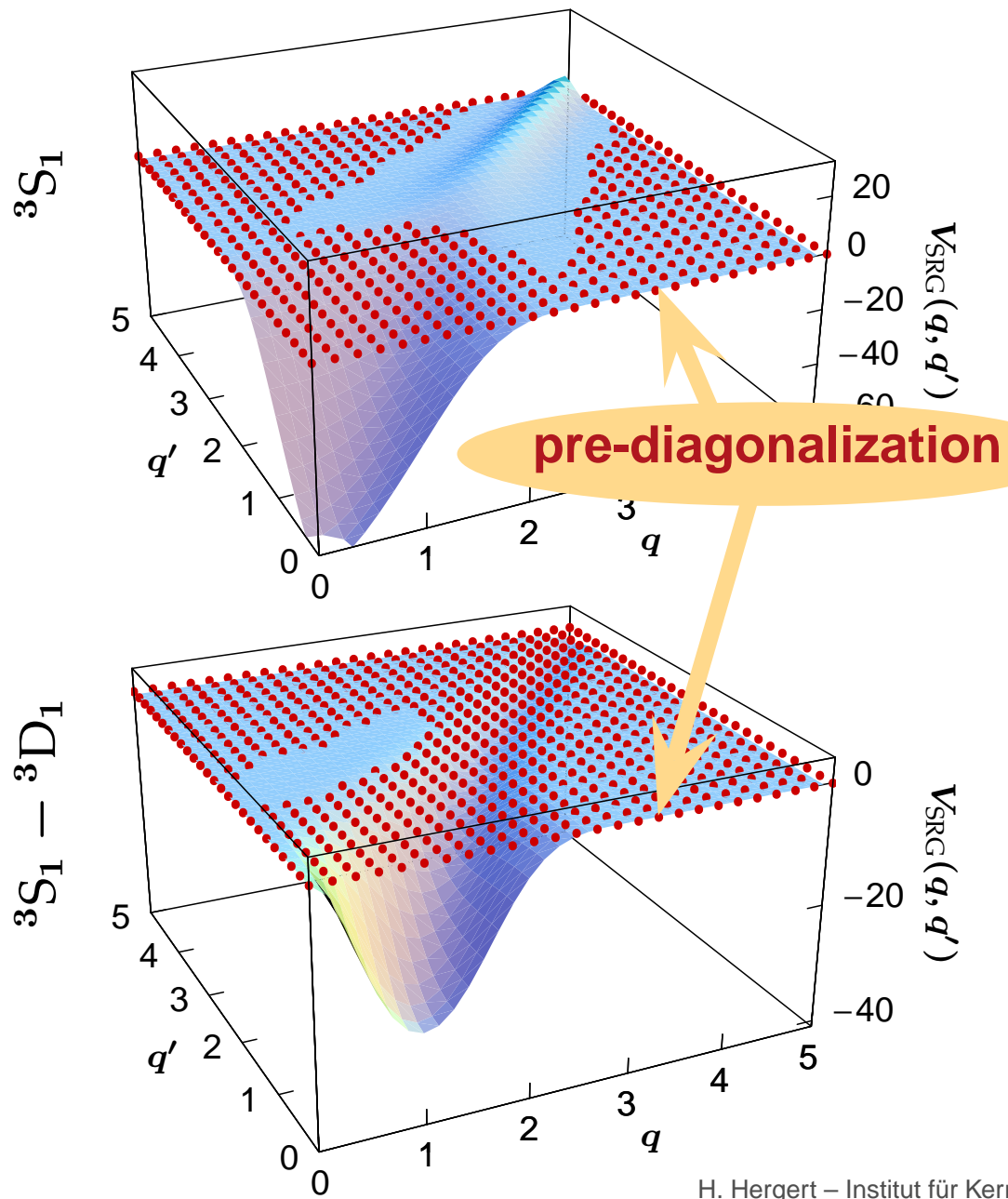
SRG Evolution: The Deuteron



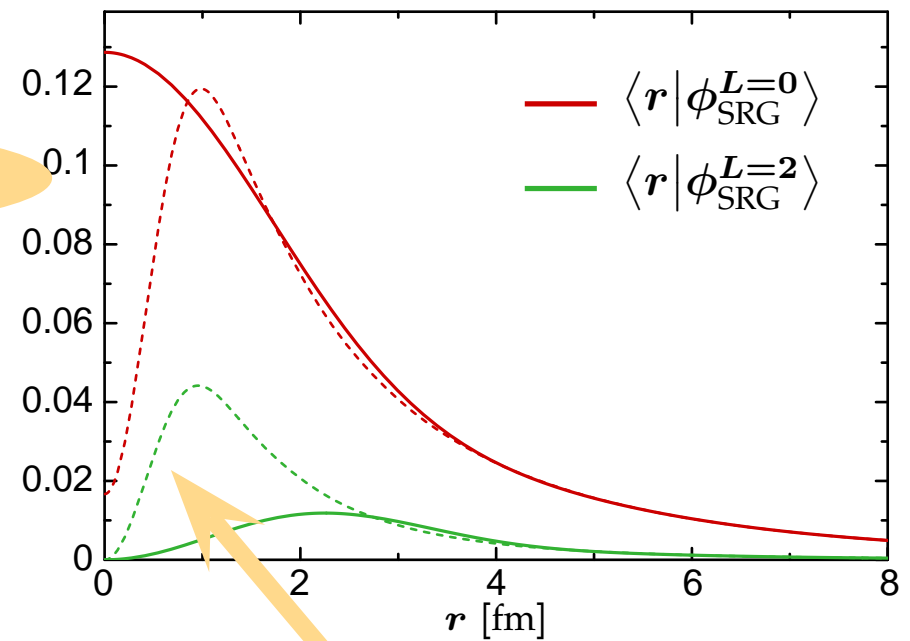
$$\alpha = 0.1000 \text{ fm}^4$$



SRG Evolution: The Deuteron

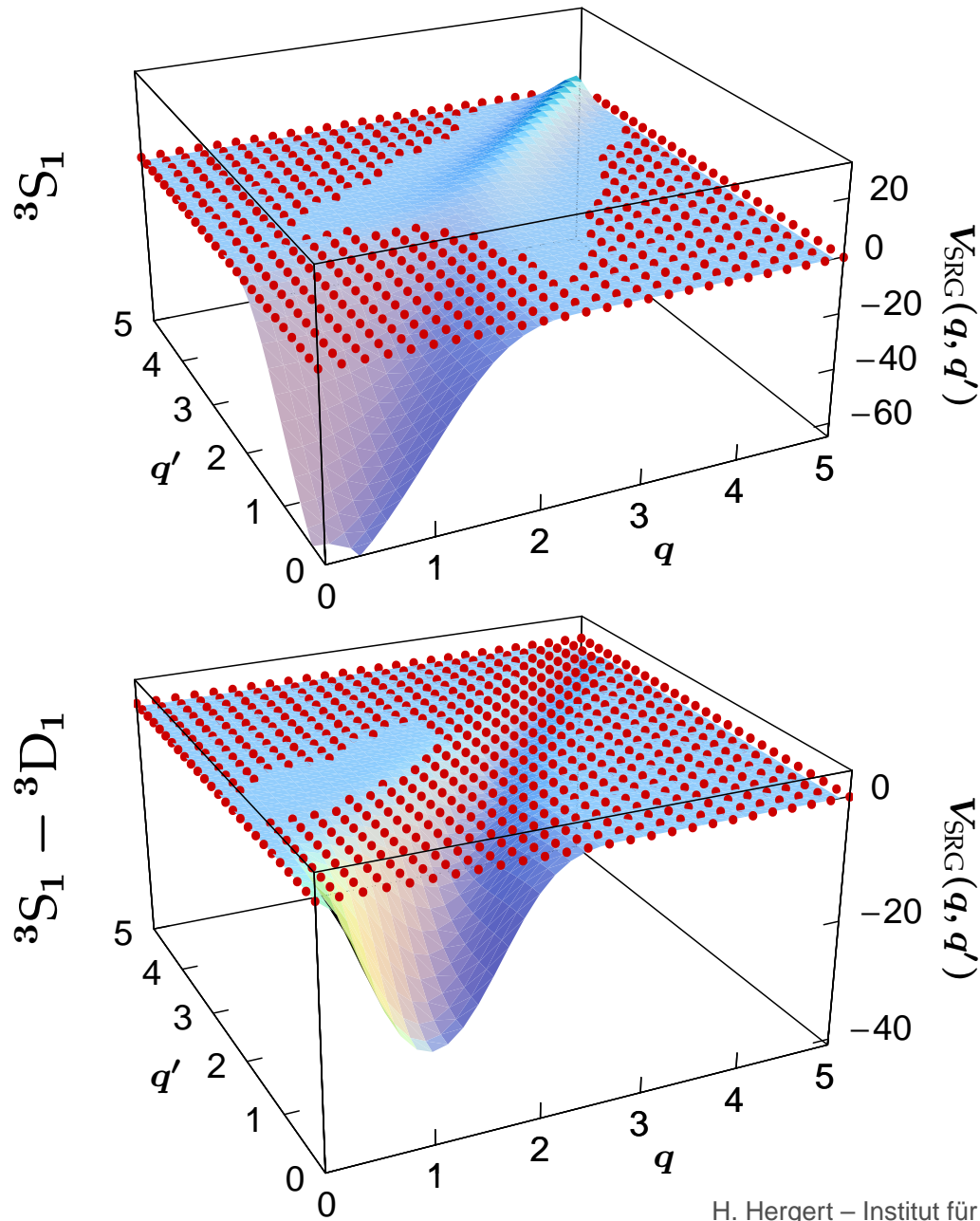


$$\alpha = 0.1000 \text{ fm}^4$$

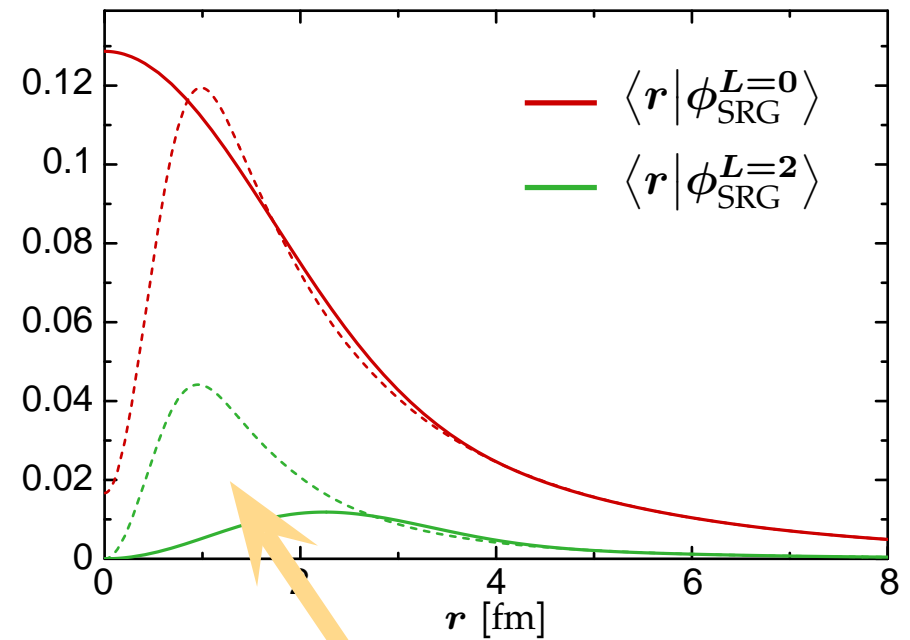


elimination of
short-range
correlations

SRG-Evolution: The Deuteron

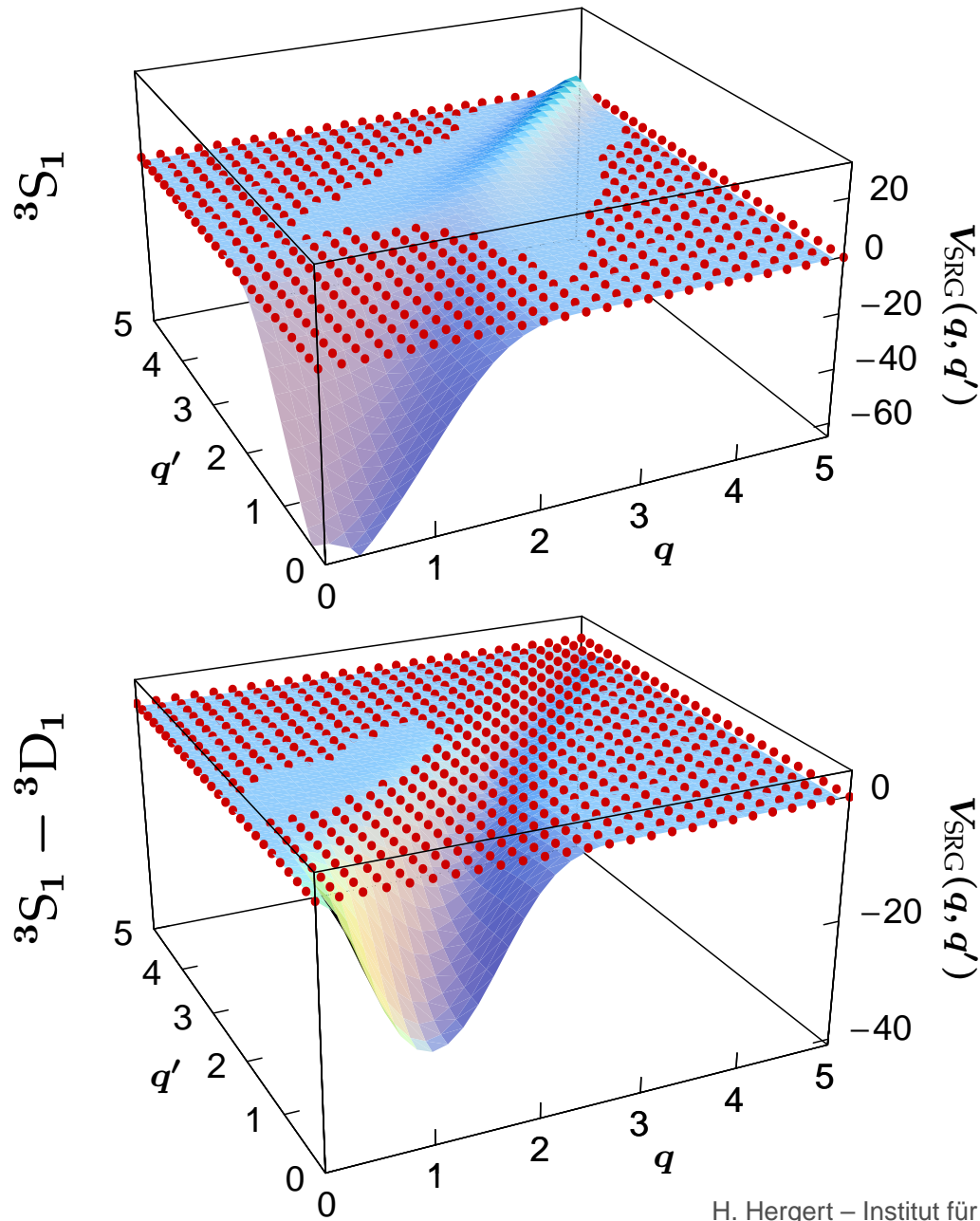


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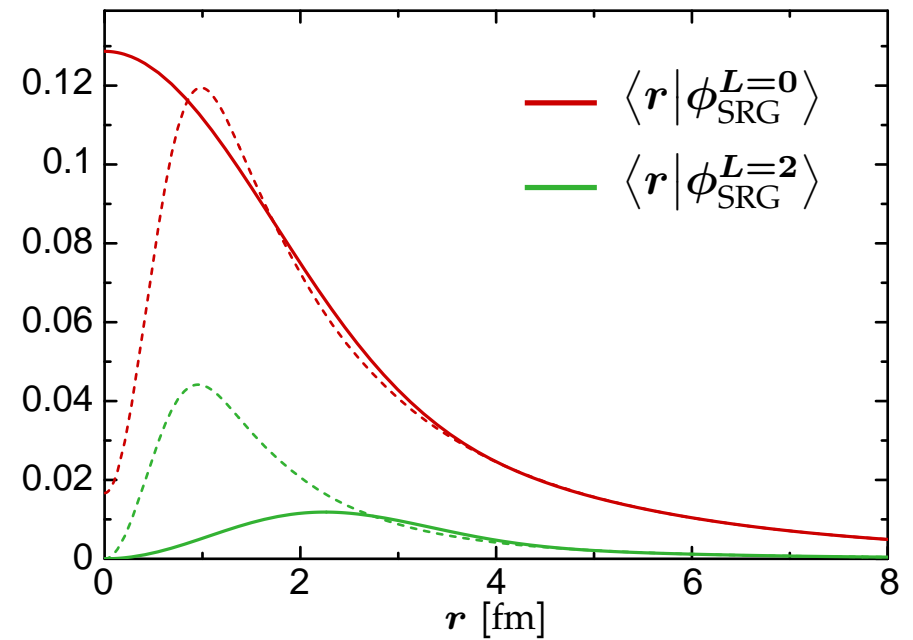


**extract UCOM
correlation functions**

SRG-Evolution: The Deuteron

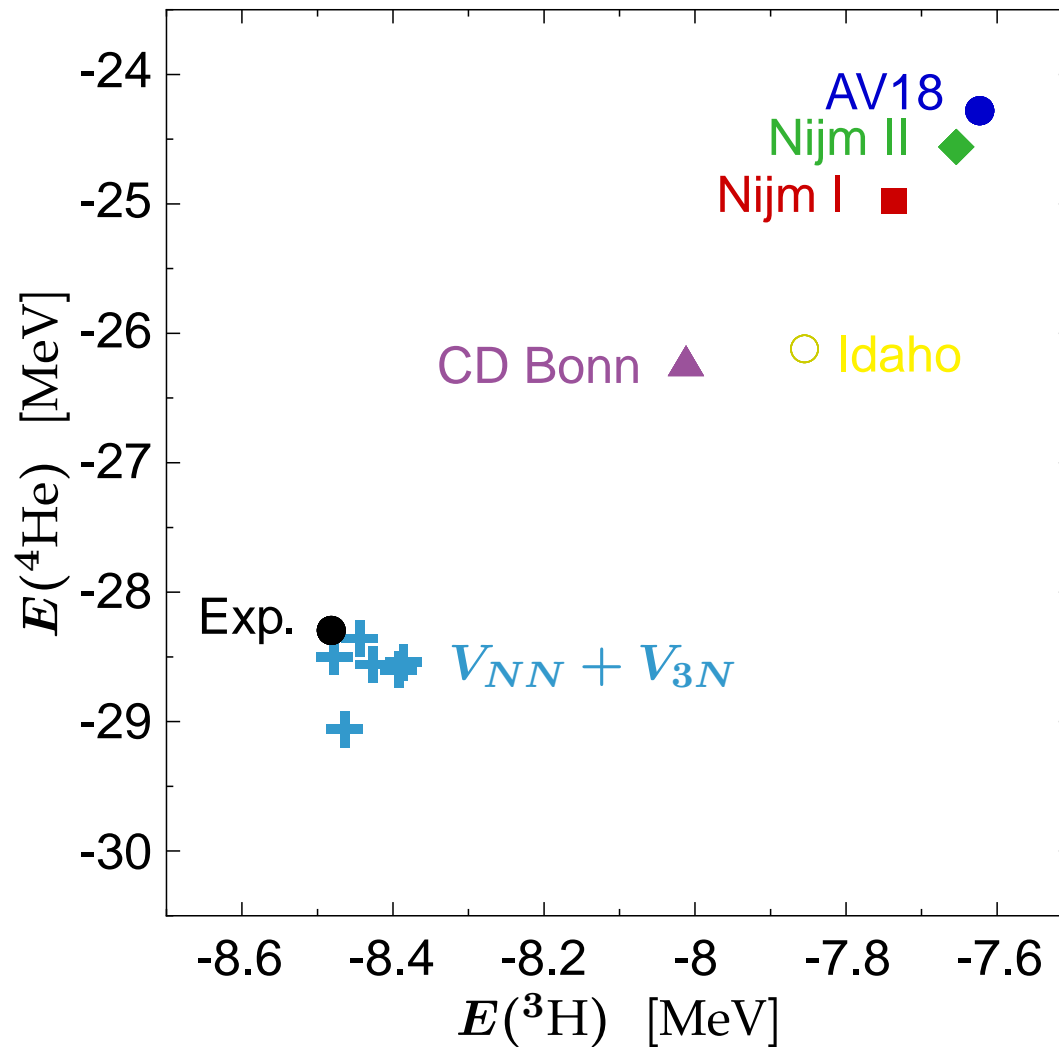


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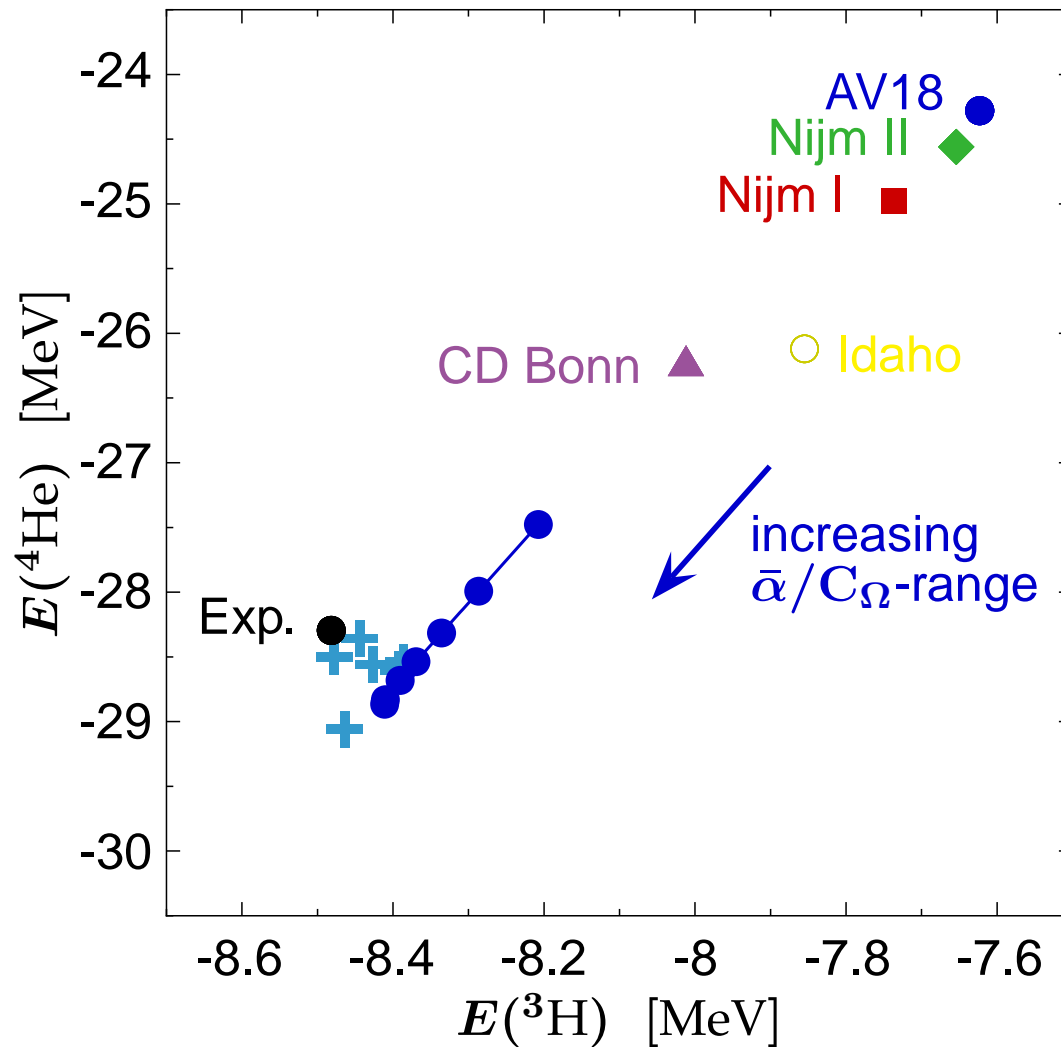
$$V_{\text{UCOM}} \neq V_{\text{SRG}}!$$

Tjon Line



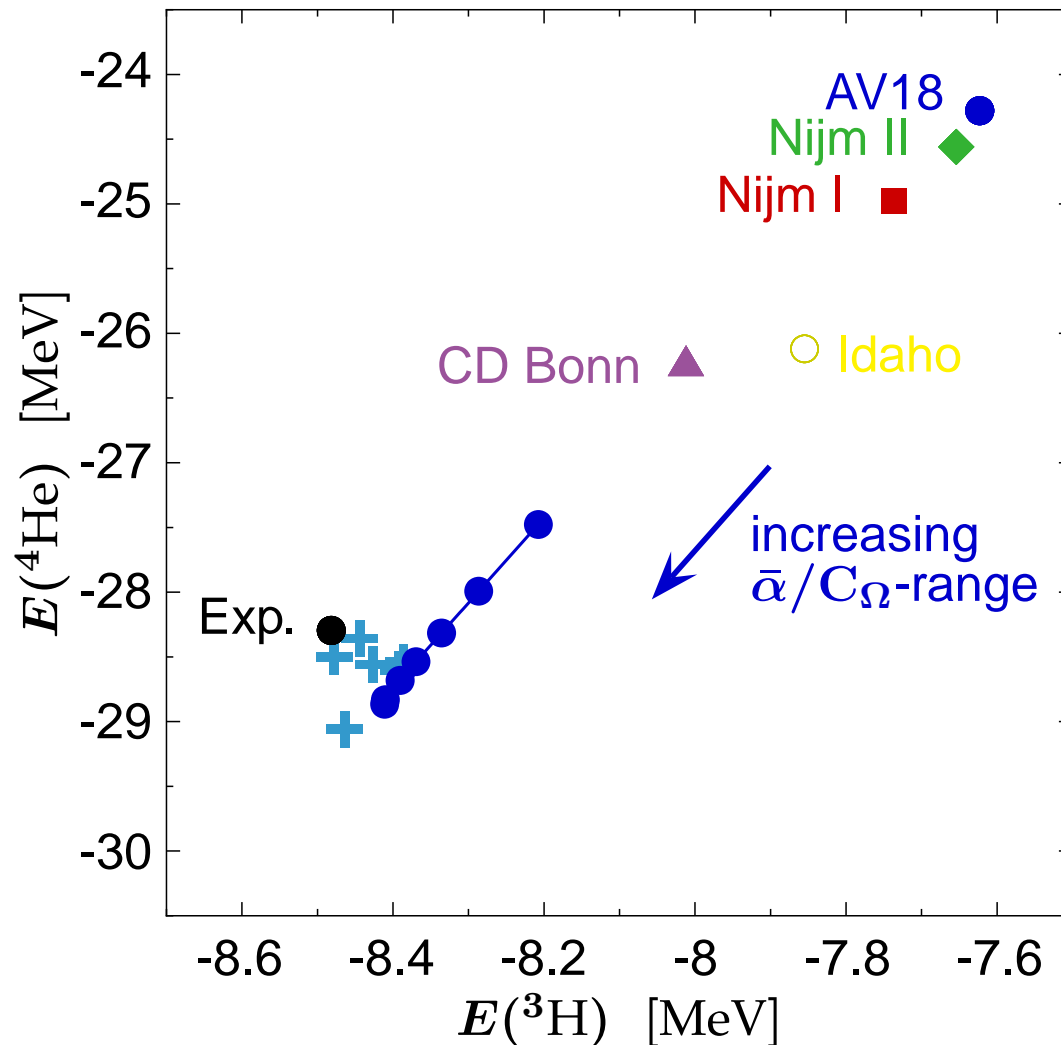
- **Tjon line:** $E({}^4\text{He})$ vs. $E({}^3\text{H})$ for phase-shift equivalent NN-interactions

Tjon Line



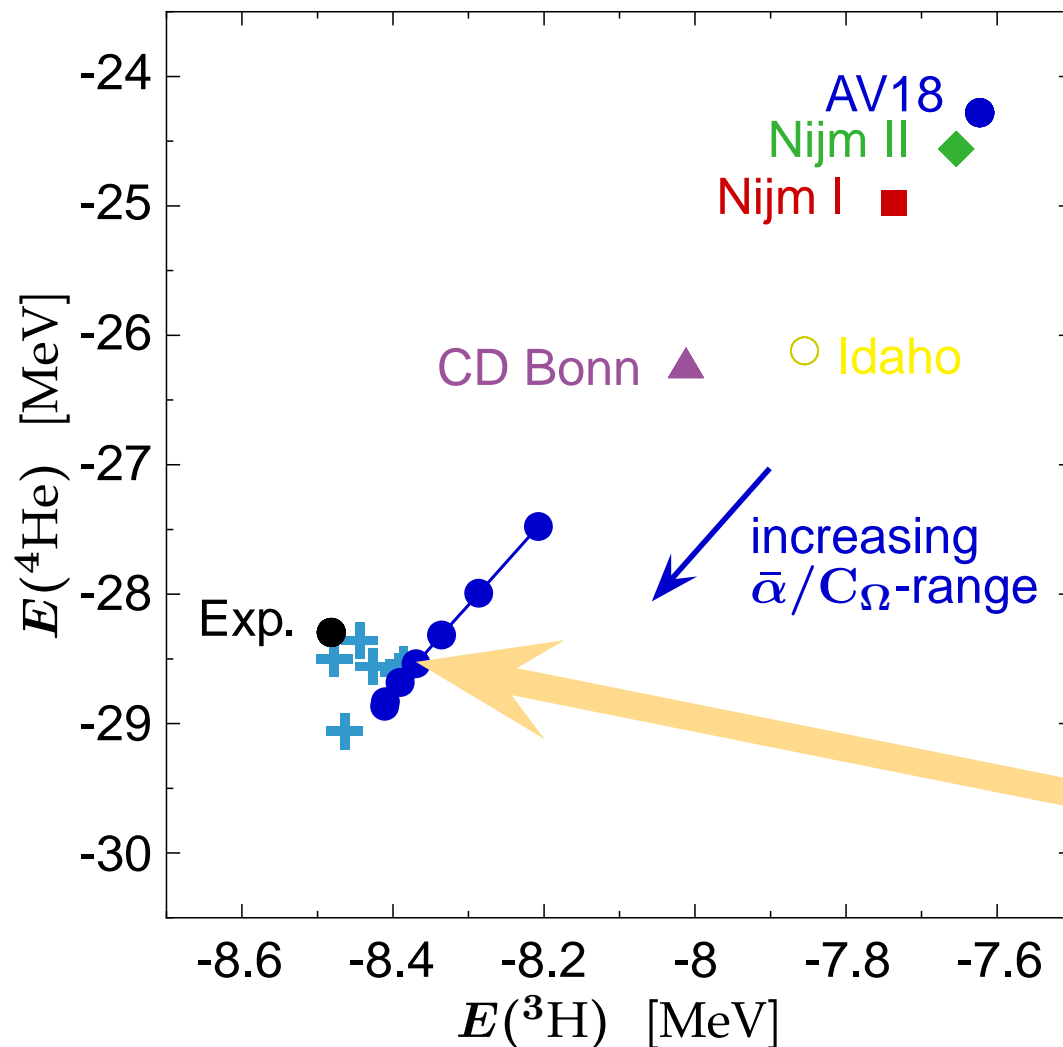
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Tjon Line



- **Tjon line:** $E({}^4\text{He})$ vs. $E({}^3\text{H})$ for phase-shift equivalent NN-interactions
- use $\bar{\alpha}$ / range of C_Ω
 - test dependence of V_{UCOM}
 - tune contributions of **net 3N force**

Tjon Line



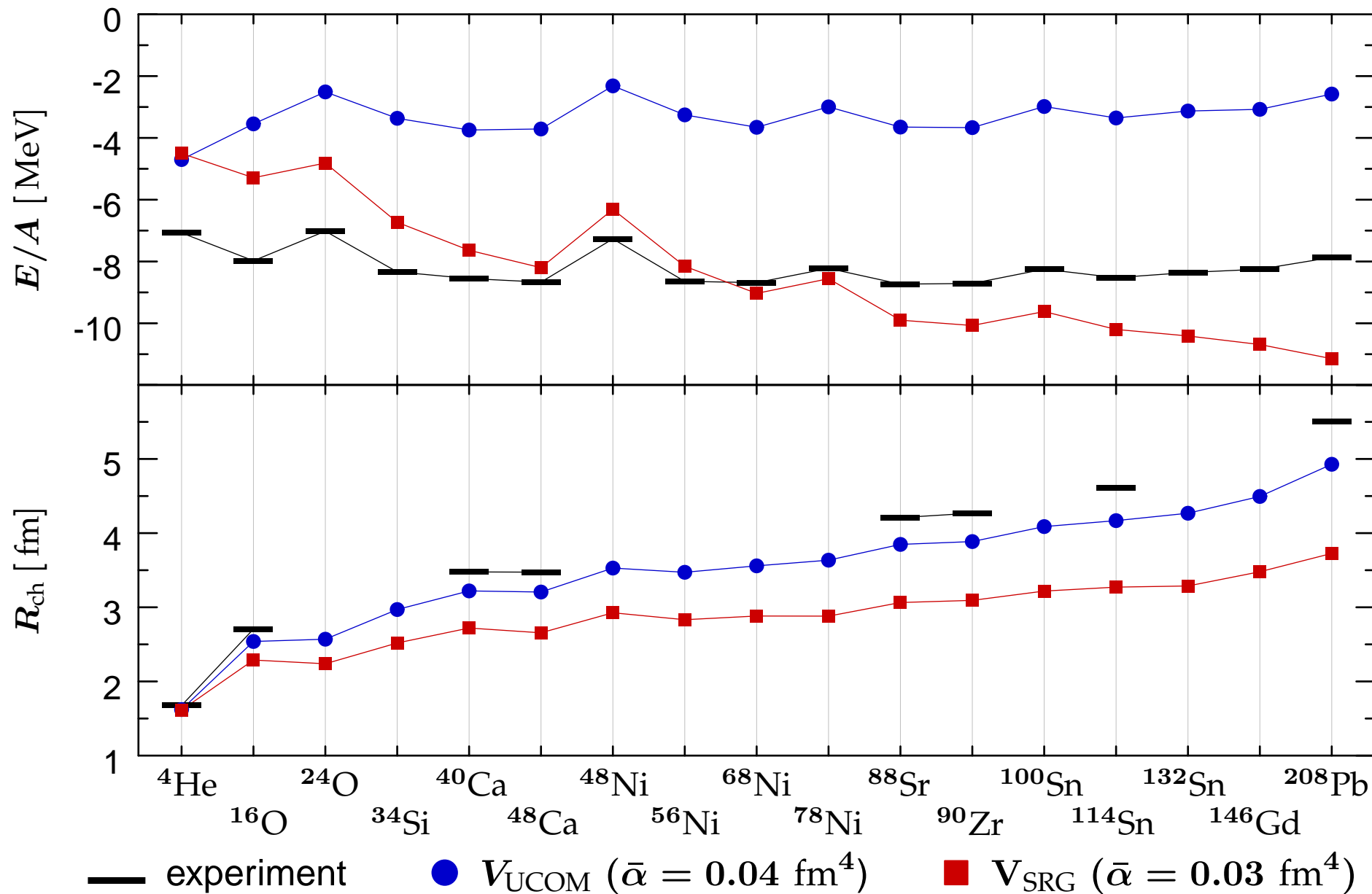
- **Tjon line:** $E(^4\text{He})$ vs. $E(^3\text{H})$ for phase-shift equivalent NN-interactions

- use $\bar{\alpha}$ / range of C_Ω to
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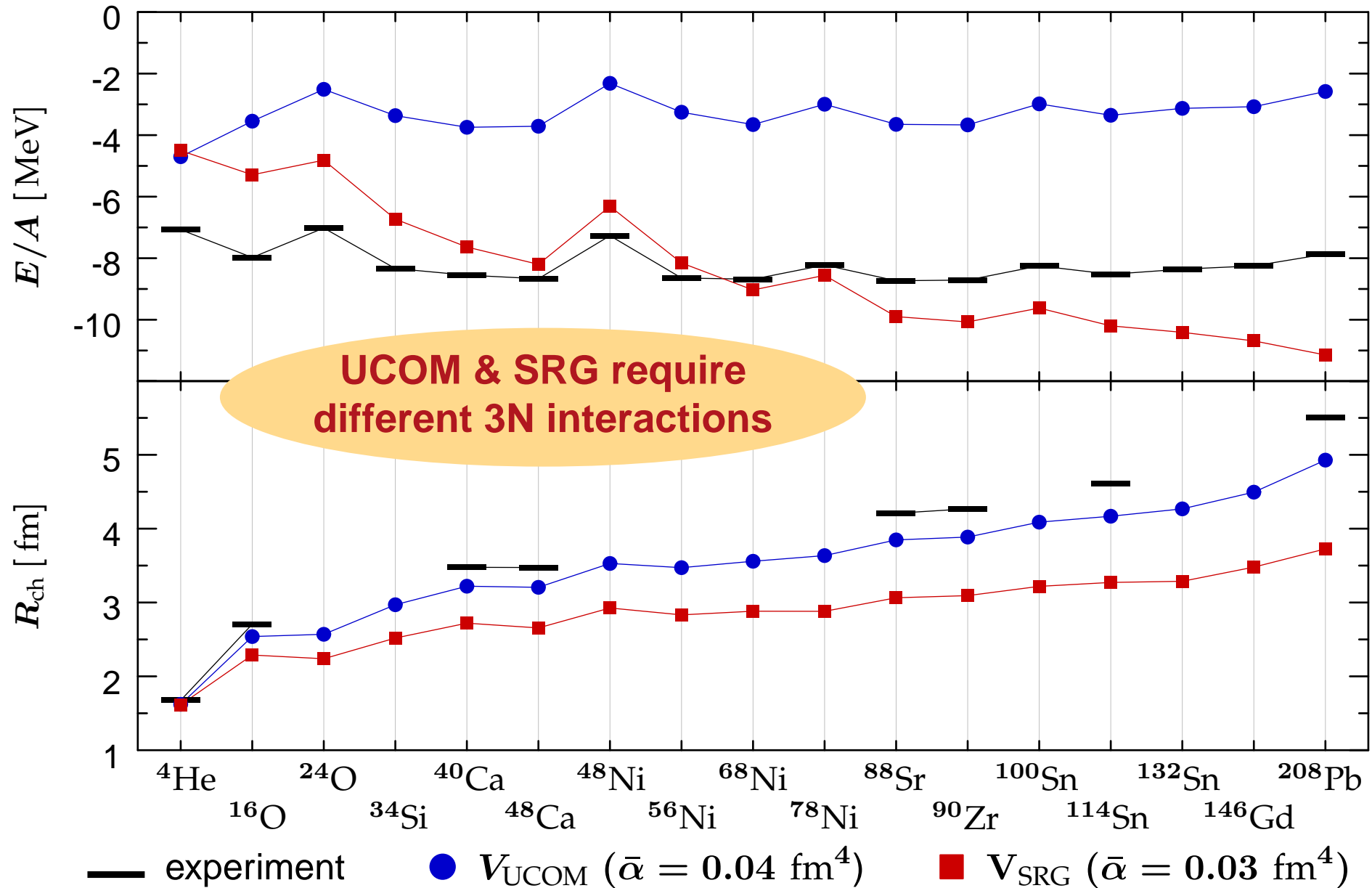
minimal net 3N interaction
use V_{UCOM} with
 $\bar{\alpha} = 0.04 \text{ fm}^4$

Hartree-Fock and Many-Body Perturbation Theory

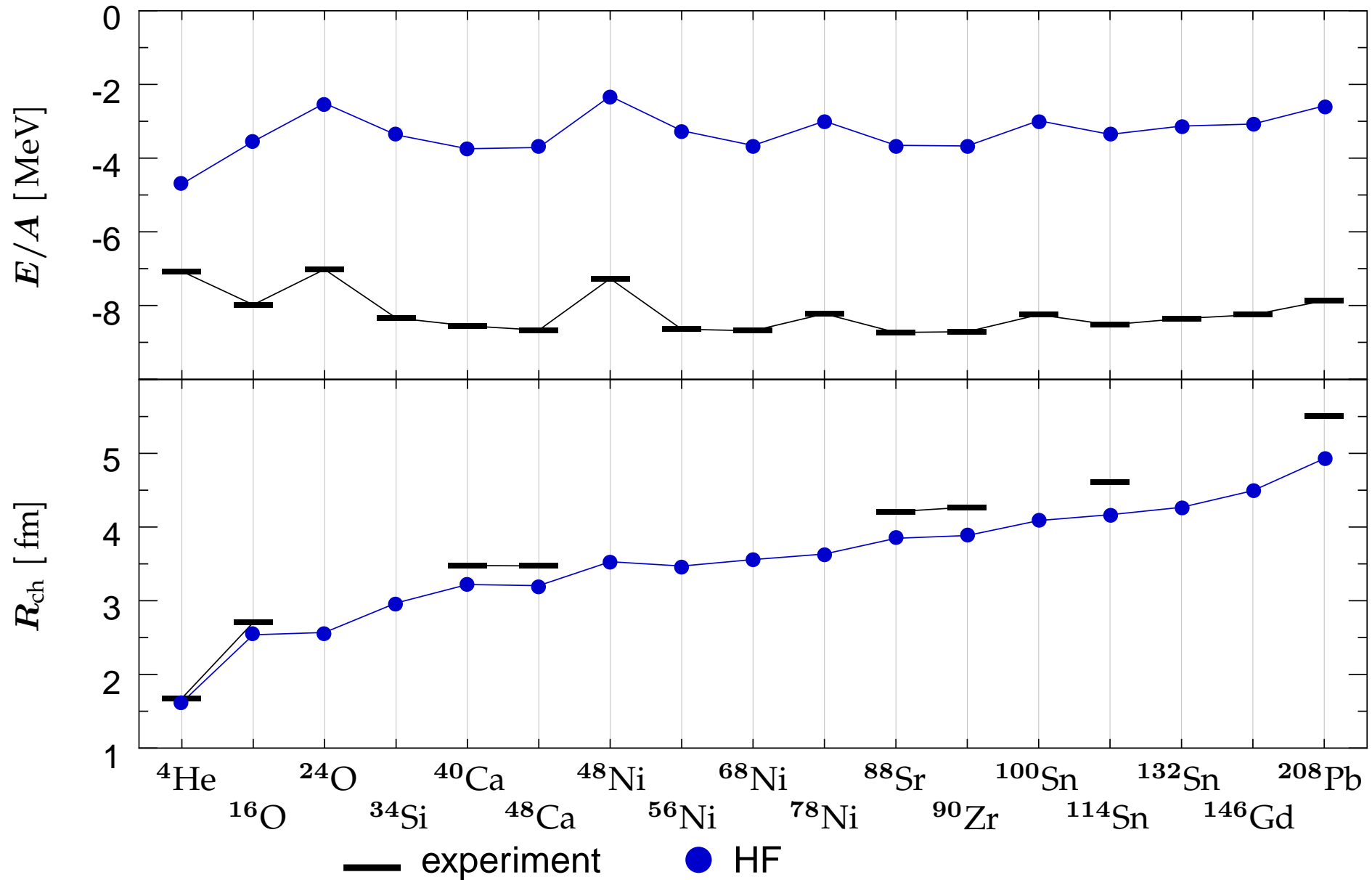
Hartree-Fock: UCOM vs. SRG



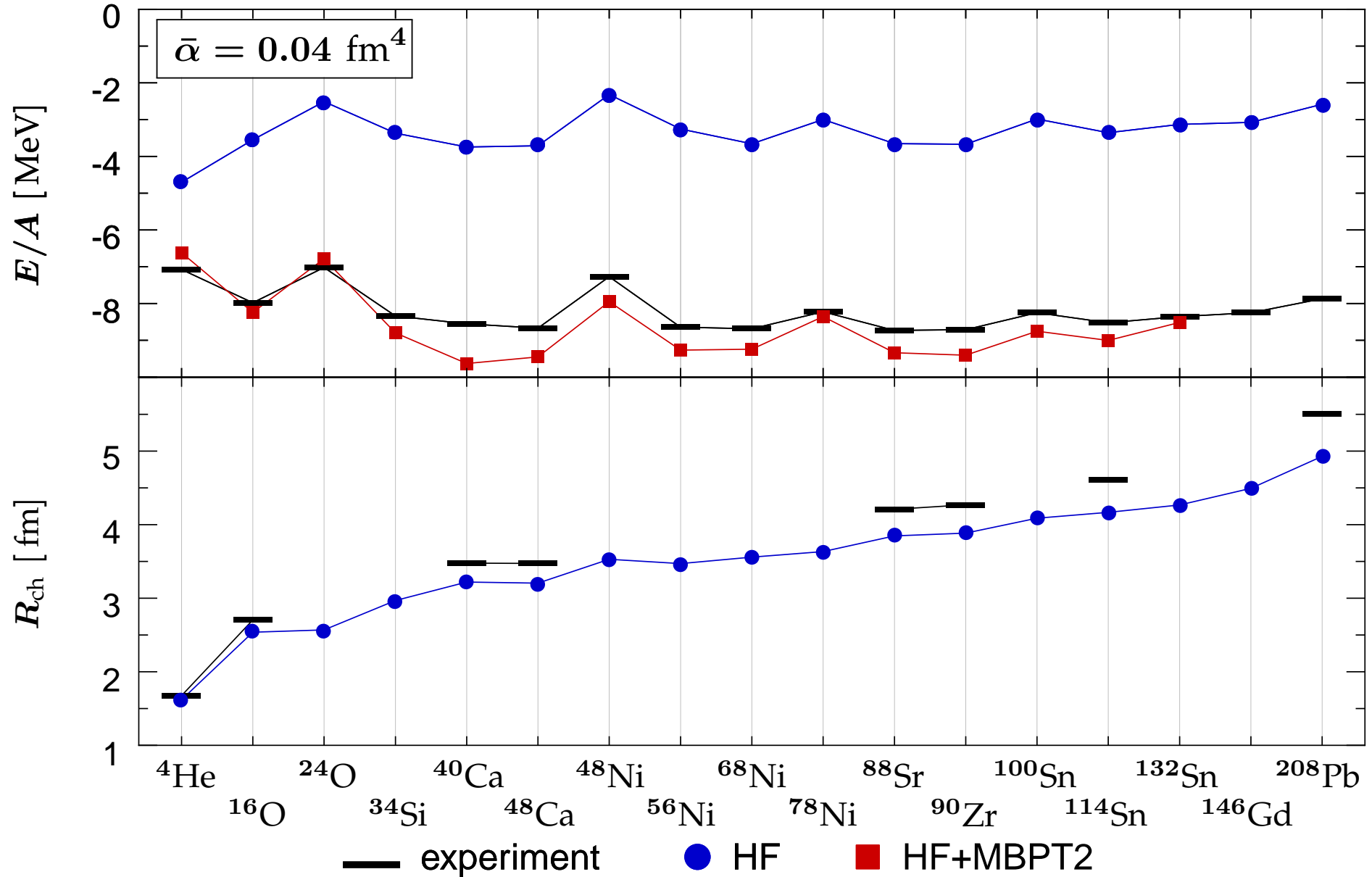
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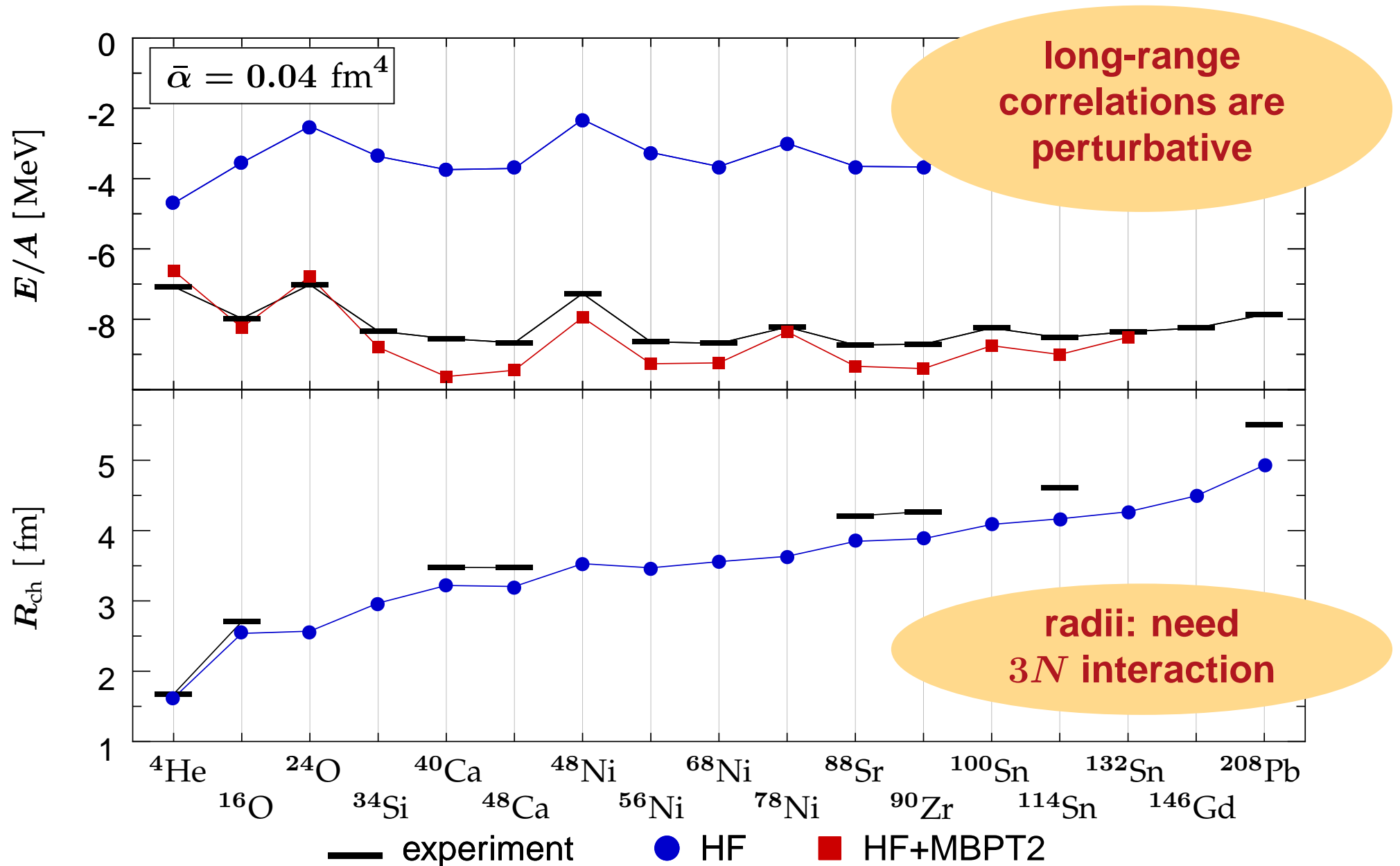
Hartree-Fock & Perturbation Theory



Hartree-Fock & Perturbation Theory



Hartree-Fock & Perturbation Theory



Pairing in the UCOM Framework

Hartree-Fock-Bogoliubov

HFB Theory Overview

Bogoliubov Transformation

$$\beta_k^\dagger = \sum_q U_{qk} c_q^\dagger + V_{qk} c_q$$

$$\beta_k = \sum_q U_{qk}^* c_q + V_{qk}^* c_q^\dagger$$

where

$$\{\beta_k, \beta_{k'}\} \stackrel{!}{=} \{\beta_k^\dagger, \beta_{k'}^\dagger\} \stackrel{!}{=} 0$$

$$\{\beta_k, \beta_{k'}^\dagger\} \stackrel{!}{=} \delta_{kk'}$$

HFB Densities & Fields

$$\rho_{kk'} \equiv \langle \Psi | c_{k'}^\dagger c_k | \Psi \rangle = (V^* V^T)_{kk'}$$

$$\kappa_{kk'} \equiv \langle \Psi | c_{k'} c_k | \Psi \rangle = (V^* U^T)_{kk'}$$

$$\Gamma_{kk'} = \sum_{qq'} \left(\frac{2}{A} \bar{t}_{\text{rel}} + \bar{v} \right)_{kq', k'q} \rho_{qq'}$$

$$\Delta_{kk'} = \sum_{qq'} \left(\frac{2}{A} \bar{t}_{\text{rel}} + \bar{v} \right)_{kk', qq'} \kappa_{qq'}$$

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$$\Delta_{kk'} = \sum_{qq'} \left(\frac{2}{A} \bar{t}_{\text{rel}} + \bar{v} \right)_{kk', qq'} \kappa_{qq'}$$

Energy

$$E[\rho, \kappa, \kappa^*] = \frac{\langle \Psi | \mathbf{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle} \equiv \frac{1}{2} (\text{tr } \Gamma \rho - \text{tr } \Delta \kappa^*)$$

HFB Equations

$$(\mathcal{H} - \lambda \mathcal{N}) \begin{pmatrix} U \\ V \end{pmatrix} \equiv \begin{pmatrix} \Gamma - \lambda & \Delta \\ -\Delta^* & -\Gamma^* + \lambda \end{pmatrix} \begin{pmatrix} U \\ V \end{pmatrix} = E \begin{pmatrix} U \\ V \end{pmatrix}$$

Gap Definitions

Gaps from Experiment

- **binding energy differences**; e.g.

$$\Delta^{(3)}(N) = (-1)^N \frac{1}{2} (E(N+1) - 2E(N) + E(N-1))$$

Common Definitions of Theoretical Gap

- lowest canonical state: Δ_μ for state with minimal canonical E_μ
- “correlated” average (\sim averaged **pairing energy**)

$$\langle \Delta \rangle = \frac{\sum_\mu \Delta_\mu u_\mu v_\mu}{\sum_\mu u_\mu v_\mu}$$

Gap Definitions

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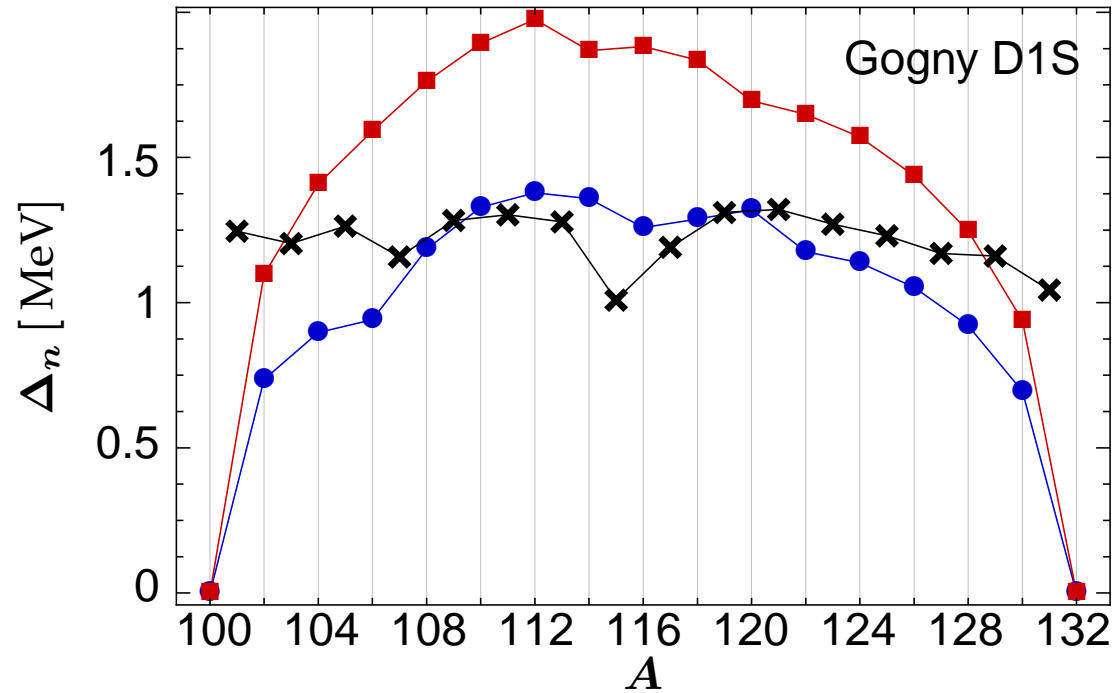
Common Definitions of Theoretical Gap

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✗ Such theoretical gaps are not observables!

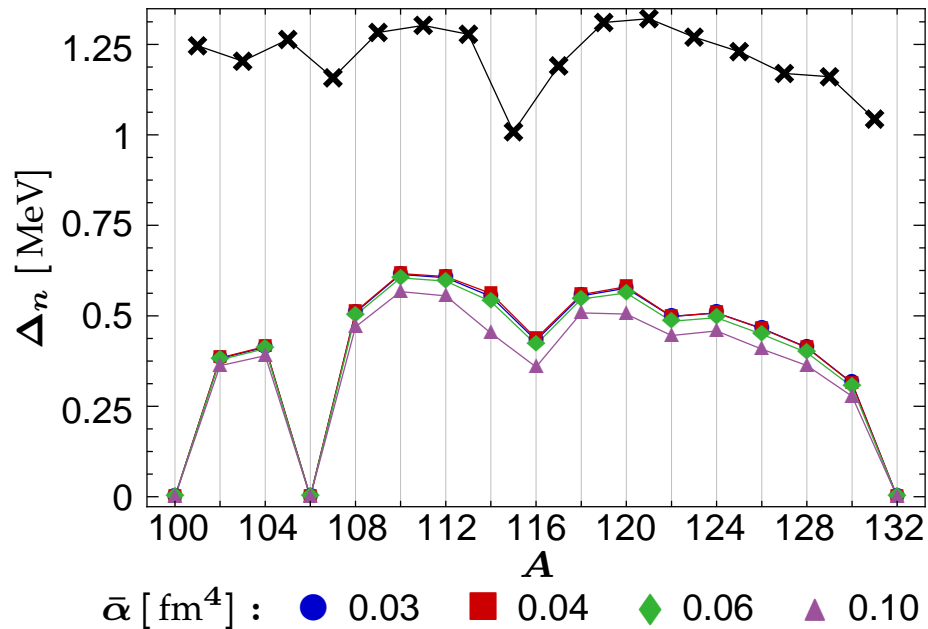
Center-of-Mass Correction



— exp. — full T_{int} — one-body approximation

$$T_{\text{int}} = \frac{2}{A} \sum_{i < j} \frac{\vec{q}_{ij}^2}{2\mu} = \underbrace{\left(1 - \frac{1}{A}\right) \sum_i \frac{\vec{p}_i^2}{2m}}_{\text{one-body}} - \underbrace{\frac{1}{Am} \sum_{i < j} \vec{p}_i \cdot \vec{p}_j}_{\text{two-body}}$$

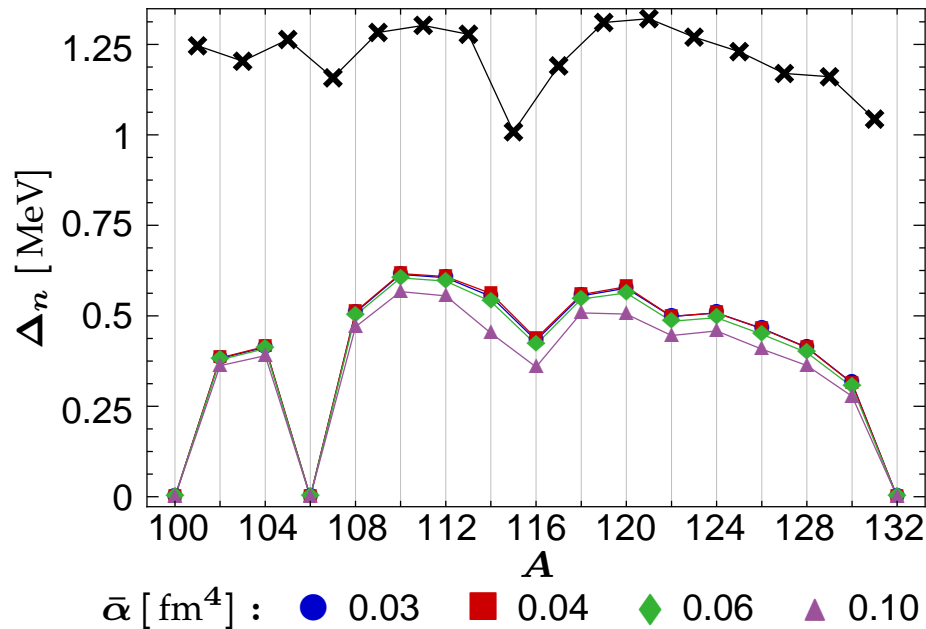
V_{UCOM} as a Pairing Force



variation of $\bar{\alpha}$

- small effect due to **stability** of 1S_0 matrix elements
- residual reduction of gaps through **partial wave mixing** by Talmi transformation

V_{UCOM} as a Pairing Force



variation of $\bar{\alpha}$

- small effect due to **stability** of 1S_0 matrix elements
- residual reduction of gaps through **partial wave mixing** by Talmi transformation

- 1S_0 gap in nuclear matter essentially determined by **phase shift**: expect similar gaps for phase-shift equivalent interactions in finite nuclei

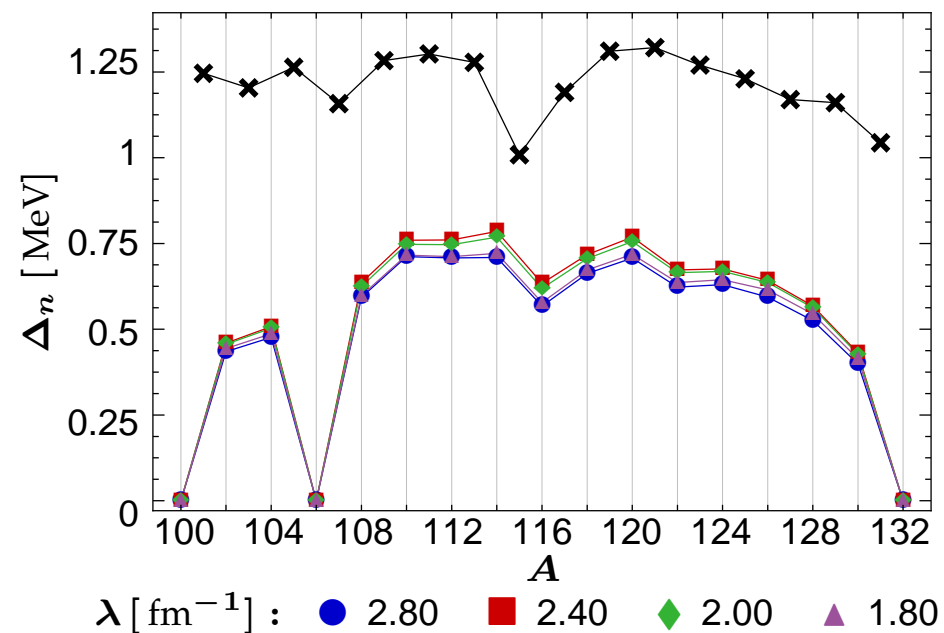
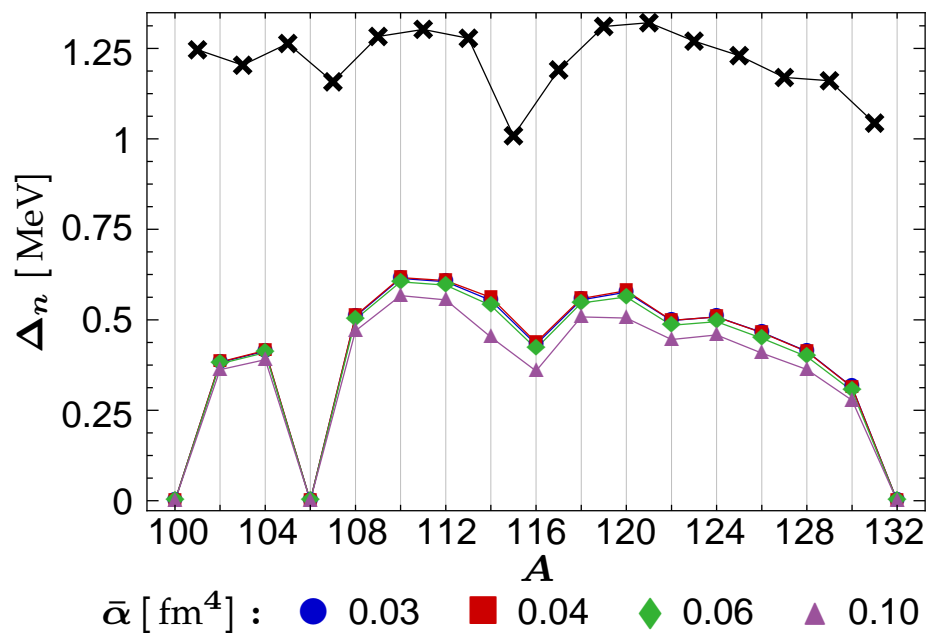
✓ Gogny D1S/SLy4 + AV14 pairing

(e.g. F. Barranco et al., EPJ A21 (2004), 57)

✗ $\sim 50\%$ smaller than SLy4 + $V_{\text{low-}k}$ study by Lesinski & Duguet

(arXiv: 0809.2895)

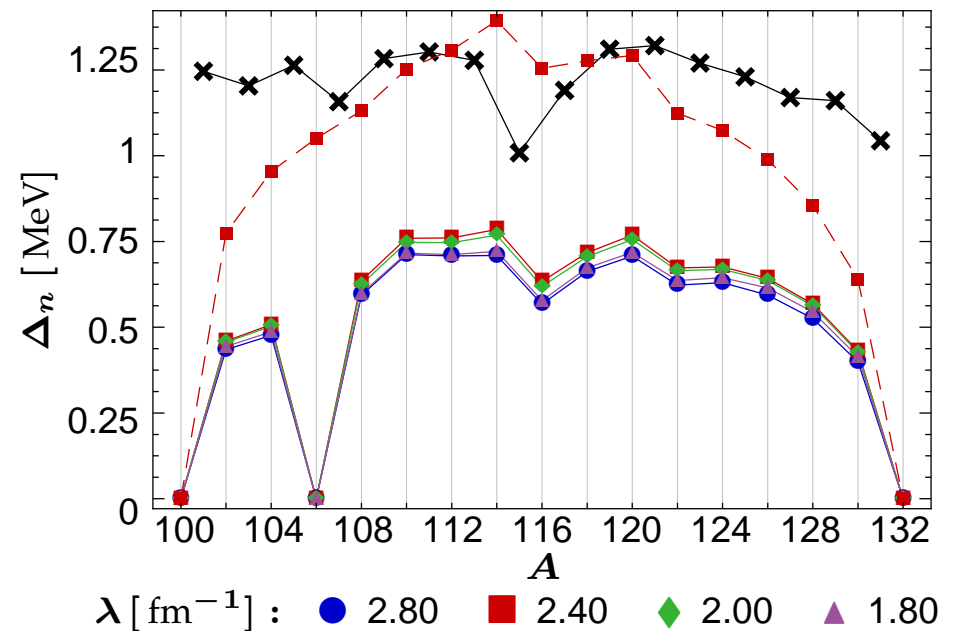
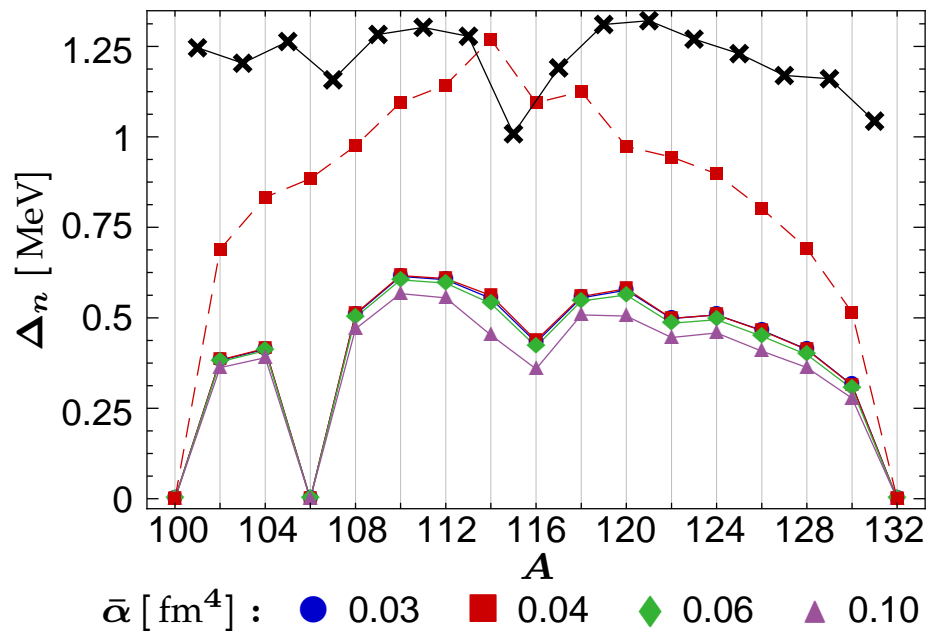
V_{UCOM} vs. V_{SRG}



■ compare with V_{SRG} (similar properties as $V_{\text{low-}k}$)

👉 center-of-mass treatment ?

V_{UCOM} vs. V_{SRG}

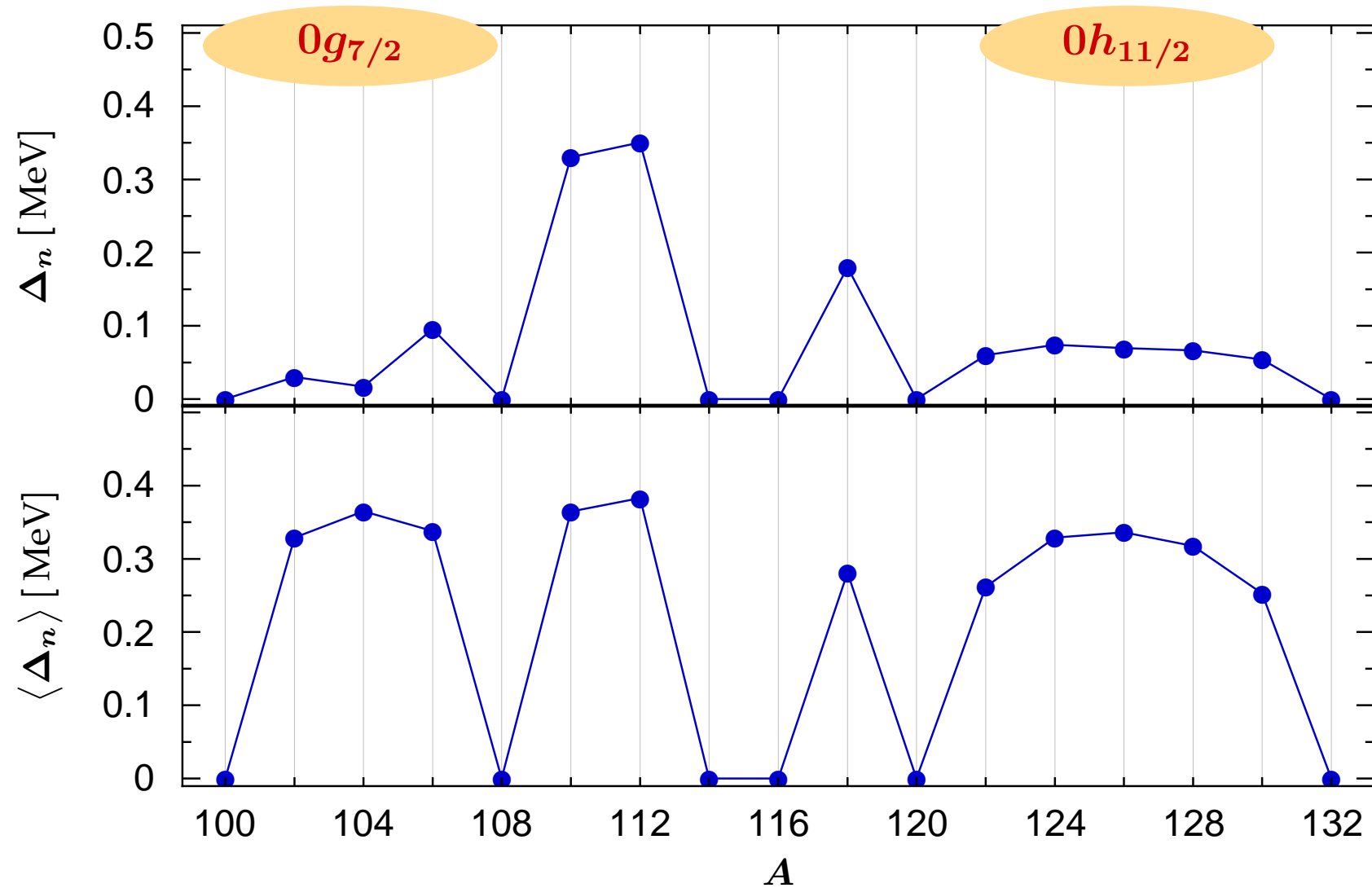


- compare with V_{SRG} (similar properties as $V_{\text{low-}k}$)
- ✓ different center-of-mass treatment explains discrepancy

Pairing in the UCOM Framework

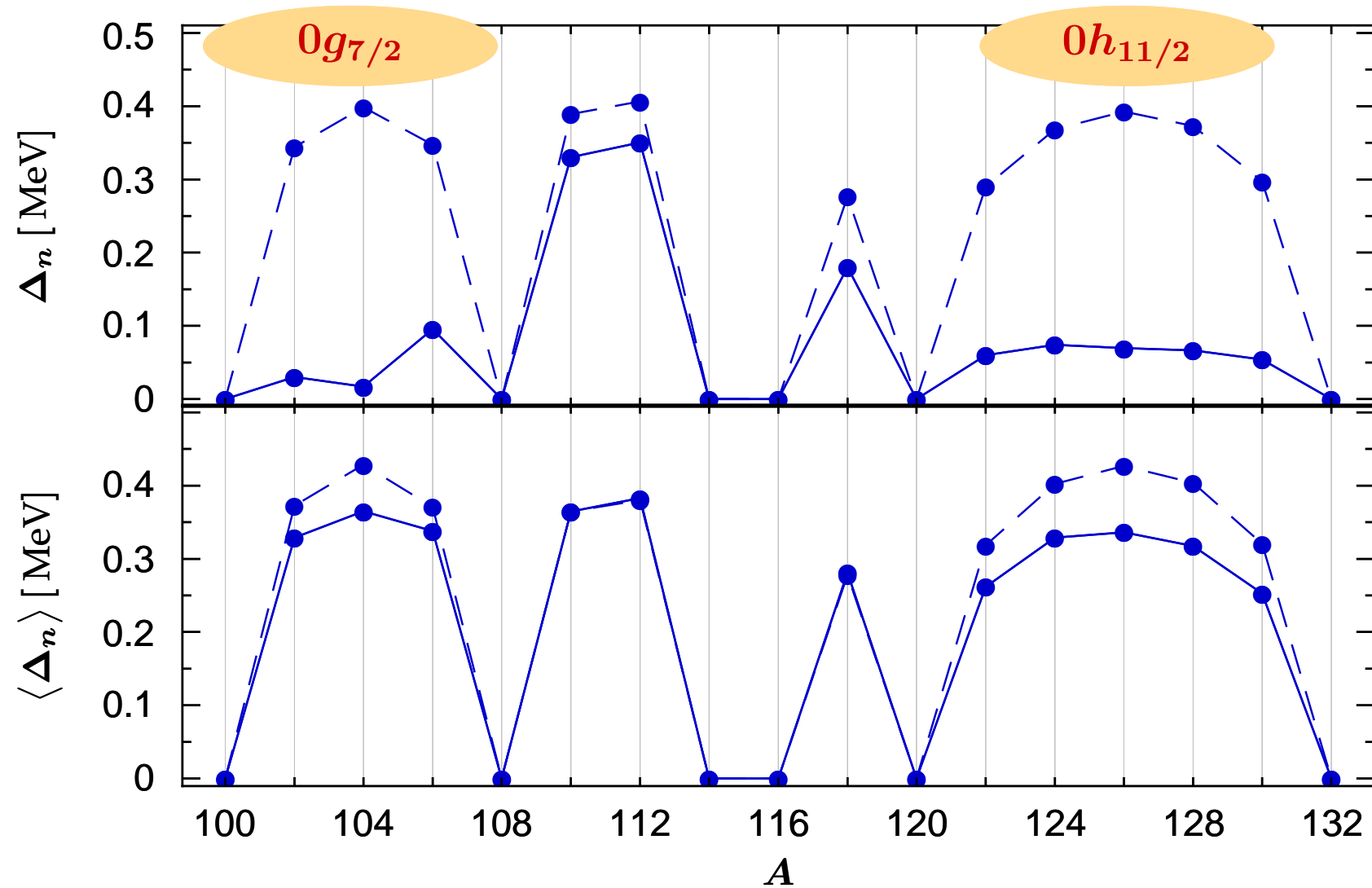
Fully Self-Consistent Hartree-Fock-Bogoliubov

Non-Central Interactions



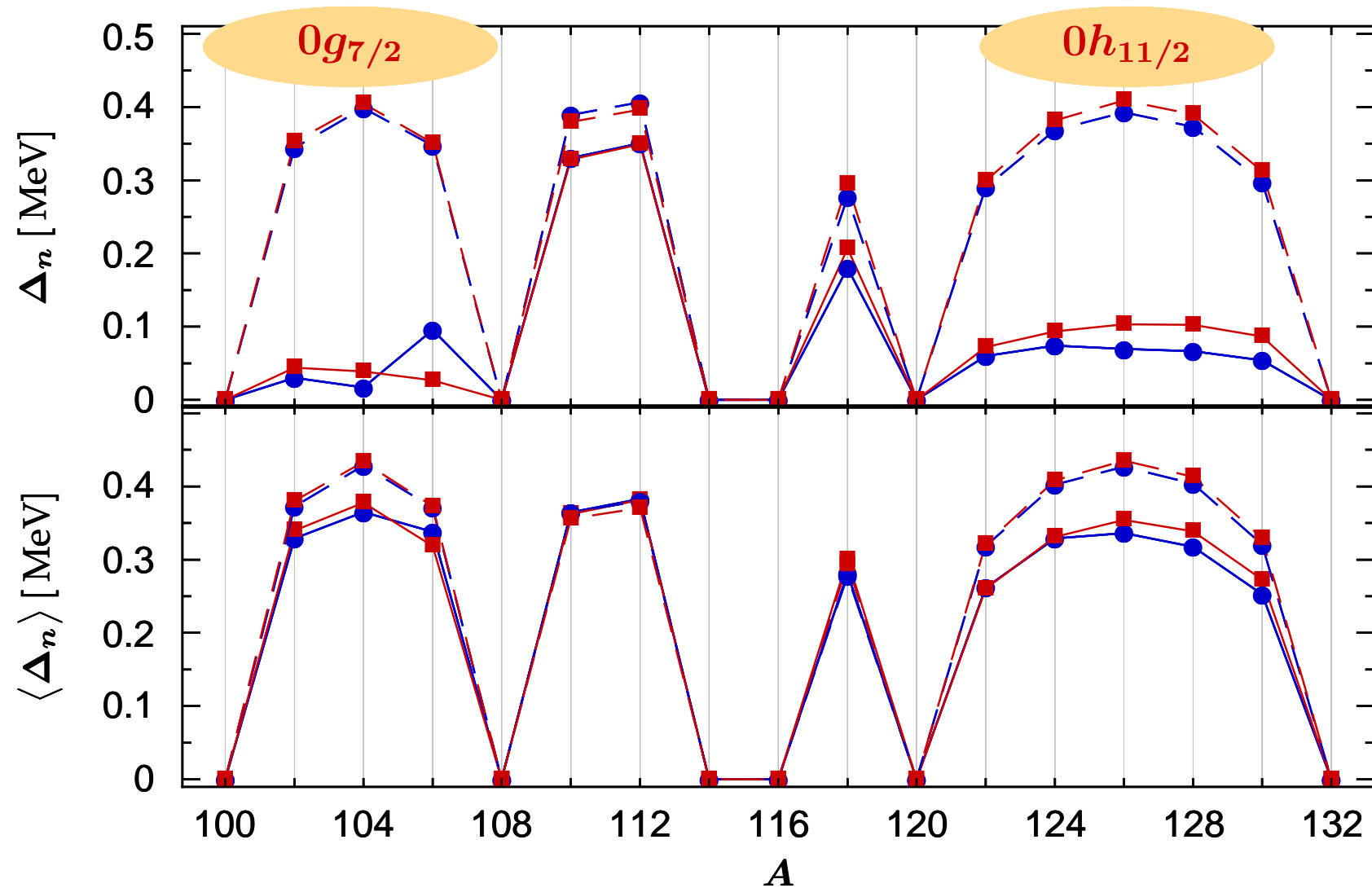
$\bar{\alpha}$ [fm⁴] : ●— (0.04, full)

Non-Central Interactions



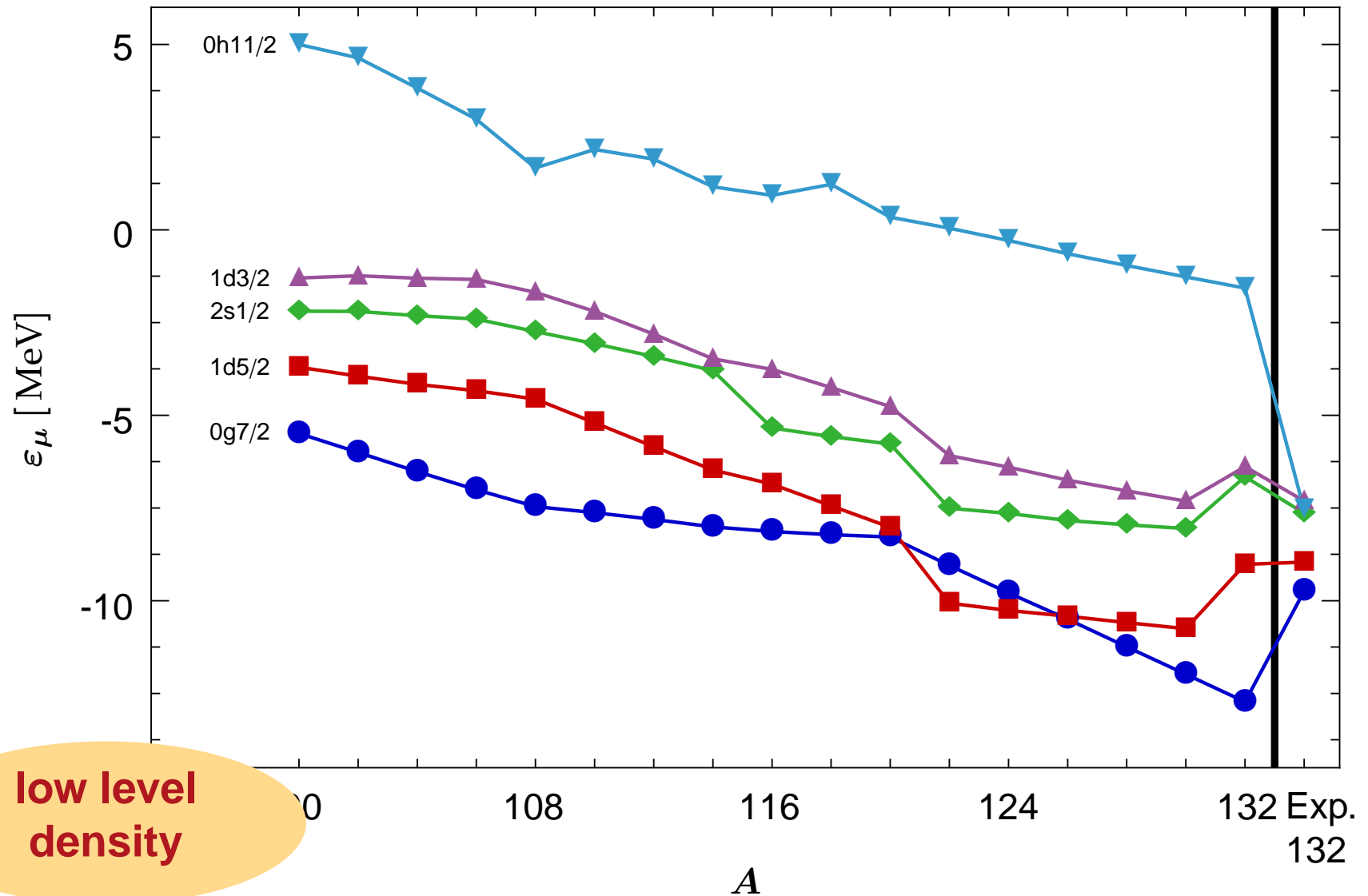
$\bar{\alpha}$ [fm⁴] : ●— (0.04, full) ●— (0.04, 1S_0)

Non-Central Interactions

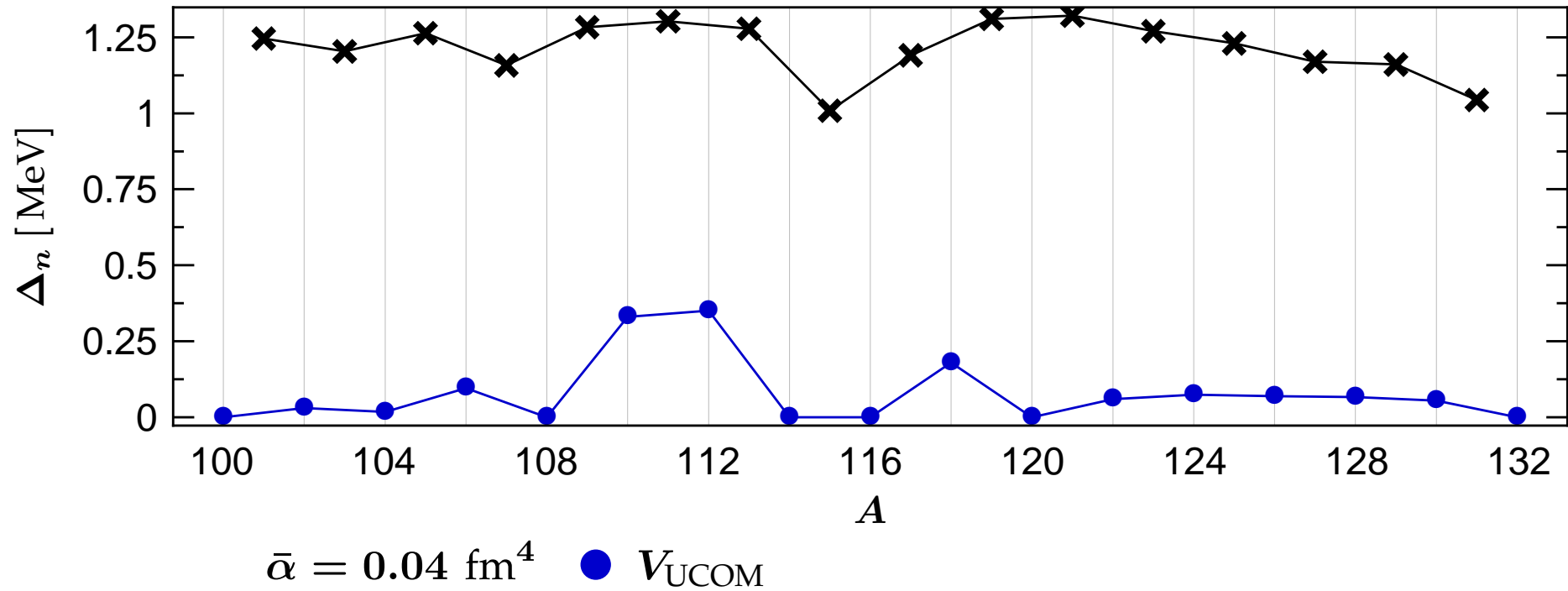


$\bar{\alpha}$ [fm⁴] : ●— (0.04, full) ●- (0.04, 1S_0) ■— (0.10, full) ■- (0.10, 1S_0)

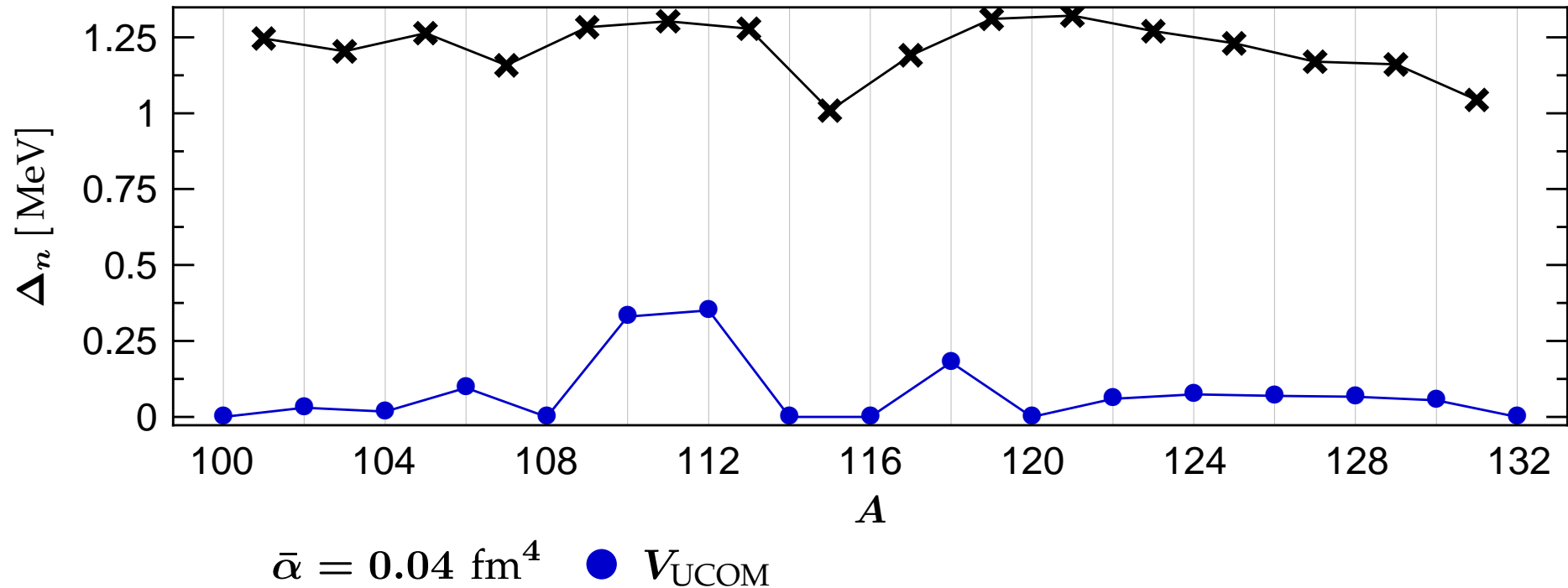
Canonical Single-Particle Spectra



Gaps



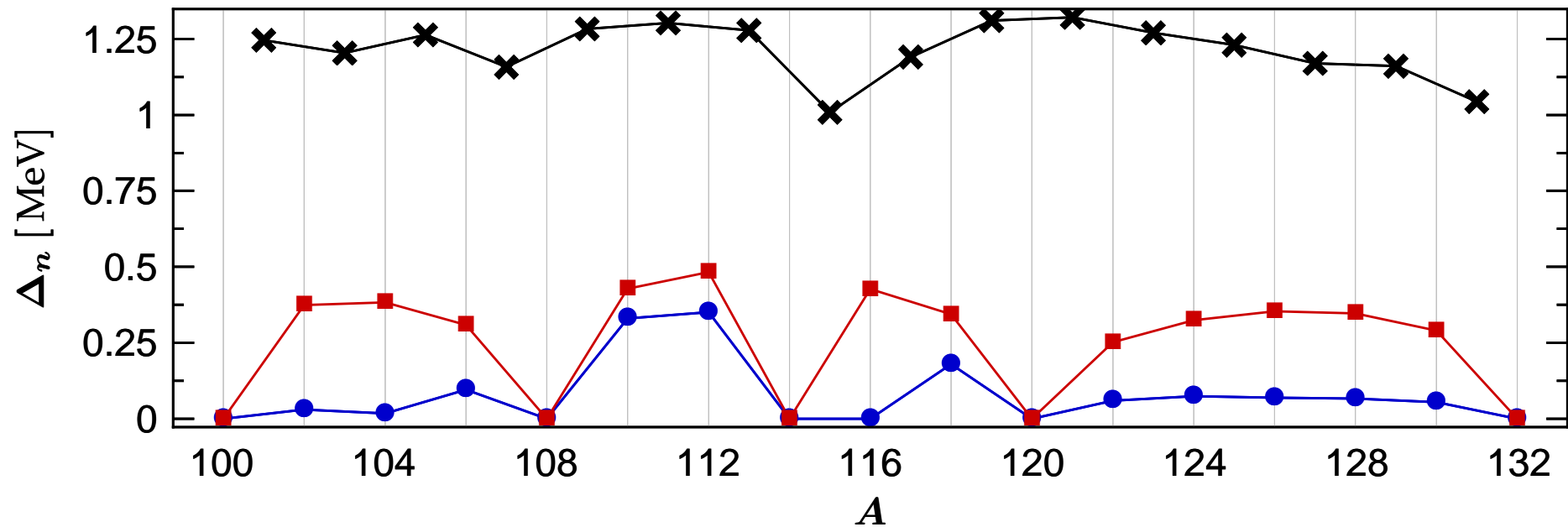
Gaps



- pairing is suppressed due to low level density at Fermi surface
- suppression of theoretical gap in valence shells with high j
- **linear density dependence:**

$$V_{\rho} = \frac{C_{3N}}{6} (1 + P_{\sigma}) \rho \left(\frac{1}{2} (\vec{r}_1 + \vec{r}_2) \right) \delta^3 (\vec{r}_1 - \vec{r}_2)$$

Gaps



$\bar{\alpha} = 0.04 \text{ fm}^4$ ● V_{UCOM} ■ $+ V_{\rho} (C_{3N} = 1.2 \text{ GeV fm}^6)$

- pairing is suppressed due to low level density at Fermi surface
- suppression of theoretical gap in valence shells with high j
- **linear density dependence:**

$$V_{\rho} = \frac{C_{3N}}{6} (1 + P_{\sigma}) \rho \left(\frac{1}{2} (\vec{r}_1 + \vec{r}_2) \right) \delta^3 (\vec{r}_1 - \vec{r}_2)$$

Particle Number Projection

Projected Energy

$$E(N_0) = \frac{\langle \Psi | \mathbf{H} \mathbf{P}_{N_0} | \Psi \rangle}{\langle \Psi | \mathbf{P}_{N_0} | \Psi \rangle} = \frac{1}{2\pi \langle \mathbf{P}_{N_0} \rangle} \int_0^{2\pi} d\phi \langle \Psi | \mathbf{H} e^{i\phi(\mathbf{N} - N_0)} | \Psi \rangle$$

Particle Number Projection

Variation of Projected Energy

$$\delta E(N_0) = \frac{1}{2\pi \langle \mathbf{P}_{N_0} \rangle} \int_0^{2\pi} d\phi \langle e^{i\phi(N-N_0)} \rangle \left\{ \delta \langle \mathbf{H} \rangle_\phi - \left(E(N_0) - \langle \mathbf{H} \rangle_\phi \right) \delta \log \langle e^{i\phi \mathbf{N}} \rangle \right\}$$

$$\langle \mathbf{H} \rangle_\phi \equiv \langle \mathbf{H} e^{i\phi \mathbf{N}} \rangle / \langle e^{i\phi \mathbf{N}} \rangle$$

Particle Number Projection

Variation of Projected Energy

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- ✓ Structure of **HFB equations is preserved!**
- ✓ manageable computational effort for variation after projection (VAP)
- ✓ implement with care: **subtle cancellations between divergences of direct, exchange, and pairing terms**

Particle Number Projection

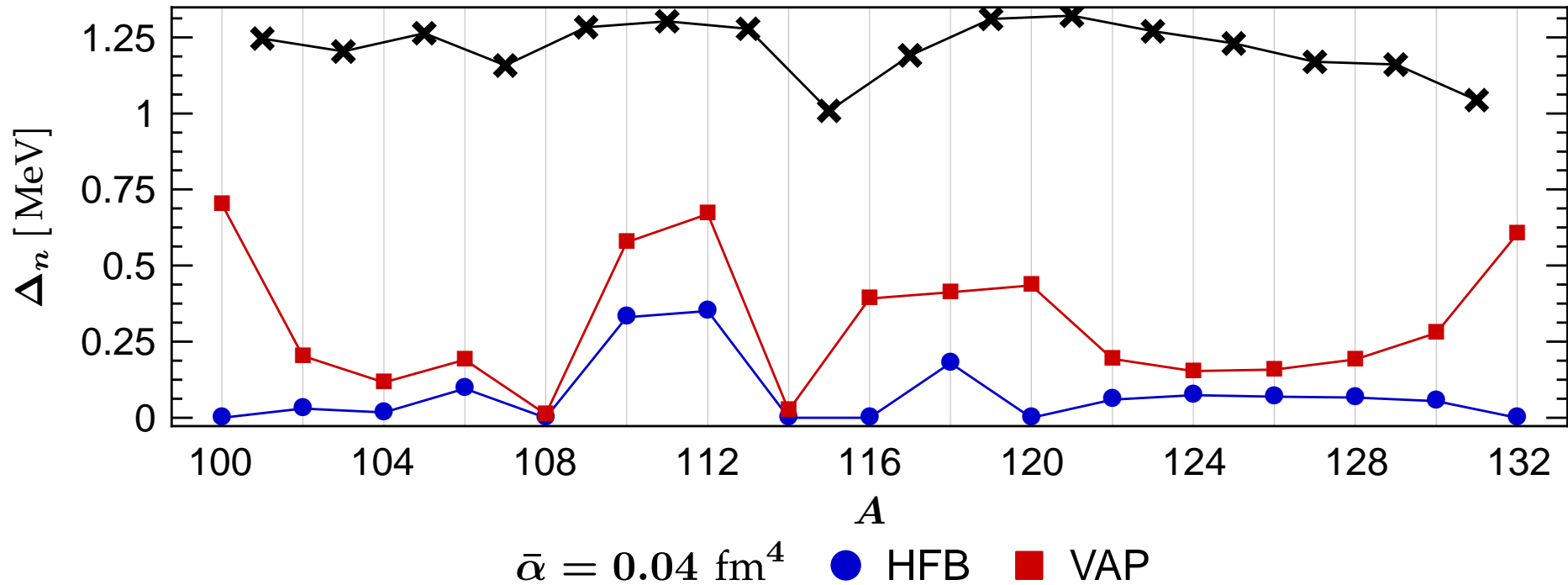
Variation of Projected Energy

$$\delta E(N_0) = \frac{1}{2\pi \langle P_{N_0} \rangle} \int_0^{2\pi} d\phi \langle e^{i\phi(N-N_0)} \rangle \left\{ \delta \langle \mathbf{H} \rangle_\phi - \left(E(N_0) - \langle \mathbf{H} \rangle_\phi \right) \delta \log \langle e^{i\phi N} \rangle \right\}$$
$$\langle \mathbf{H} \rangle_\phi \equiv \langle \mathbf{H} e^{i\phi N} \rangle / \langle e^{i\phi N} \rangle$$

- ✓ Structure of **HFB equations is preserved!**
- ✓ manageable computational effort for variation after projection (VAP)
- ✓ implement with care: **subtle cancellations between divergences of direct, exchange, and pairing terms**
- ✗ **density-dependent interaction:**
complex transition density has **poles** (serious problem for projection methods, GCM, ...)

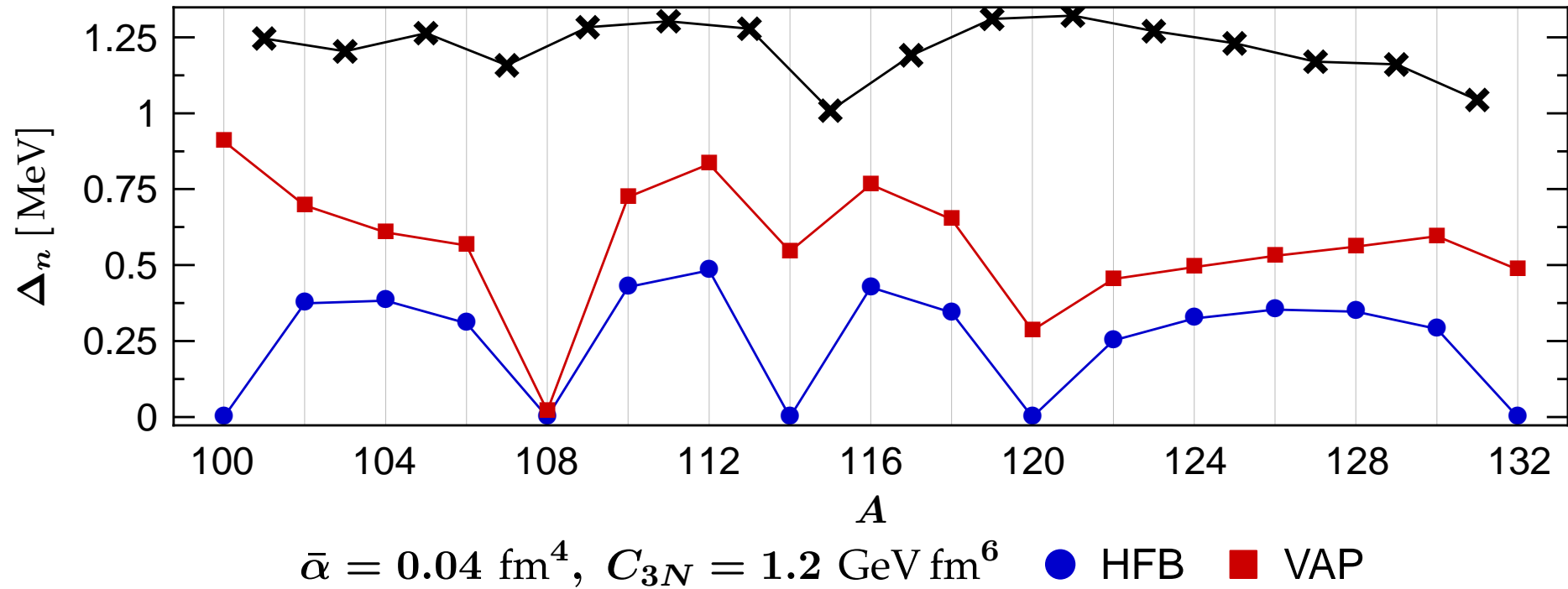
👉 Duguet, Lacroix, Bender et al., [arXiv:0809.2041](https://arxiv.org/abs/0809.2041), [0809.2045](https://arxiv.org/abs/0809.2045), [0809.2049](https://arxiv.org/abs/0809.2049)

PNP: Gaps

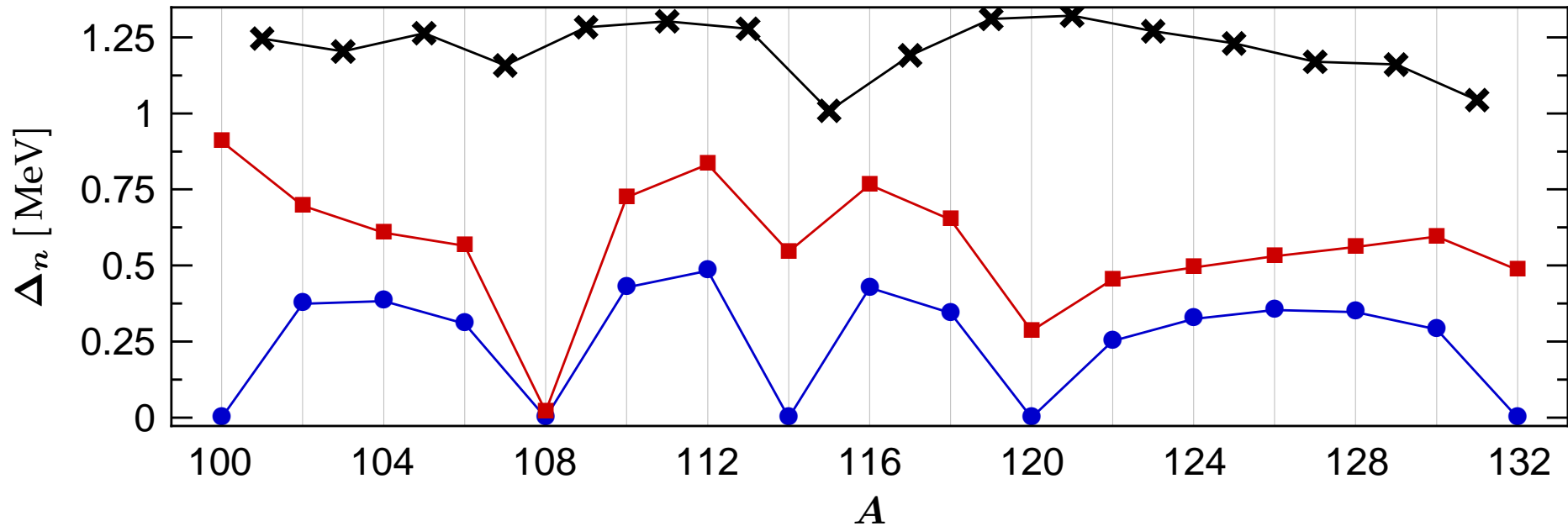


- consistent inclusion of all two-body terms (**crucial** for particle-number projection)
- projection includes additional **pairing correlations**
- sizable effects at **(major) shell closures**

PNP: Density-Dependent Interaction



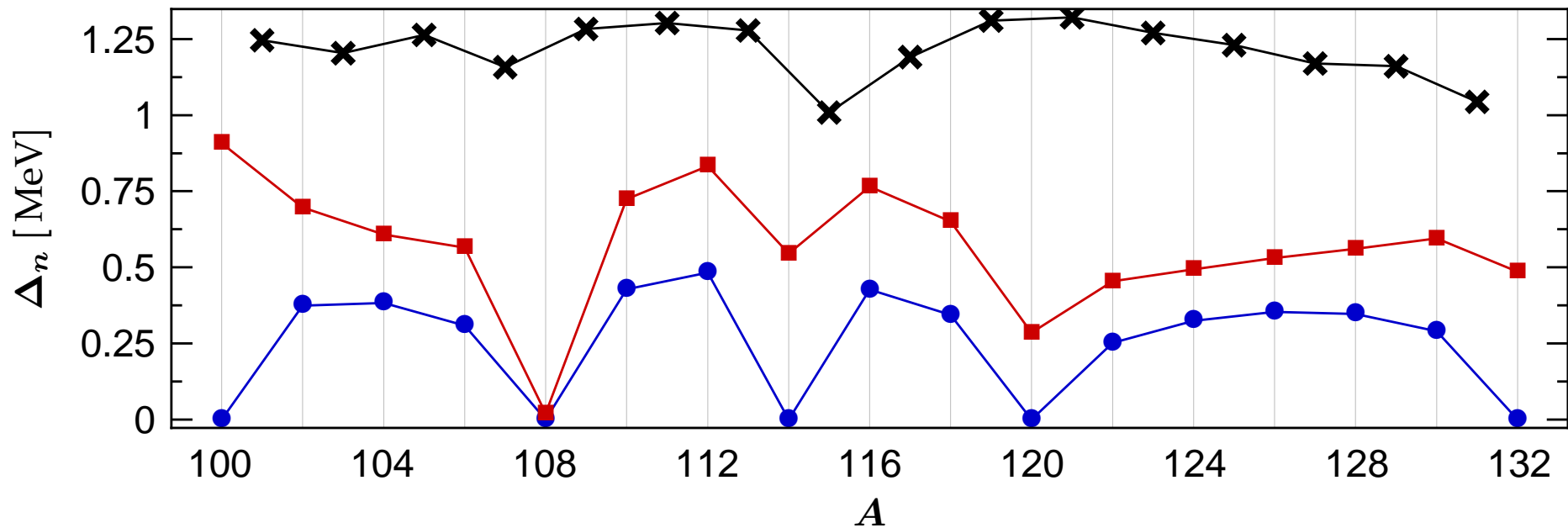
PNP: Density-Dependent Interaction



$\bar{\alpha} = 0.04 \text{ fm}^4, C_{3N} = 1.2 \text{ GeV fm}^6$ ● HFB ■ VAP

- **linear density-dependence:** isolated poles, check by projecting from neighbouring nuclei

PNP: Density-Dependent Interaction



$$\bar{\alpha} = 0.04 \text{ fm}^4, C_{3N} = 1.2 \text{ GeV fm}^6 \quad \bullet \text{ HFB} \quad \blacksquare \text{ VAP}$$

- **linear density-dependence:** isolated poles, check by projecting from neighbouring nuclei

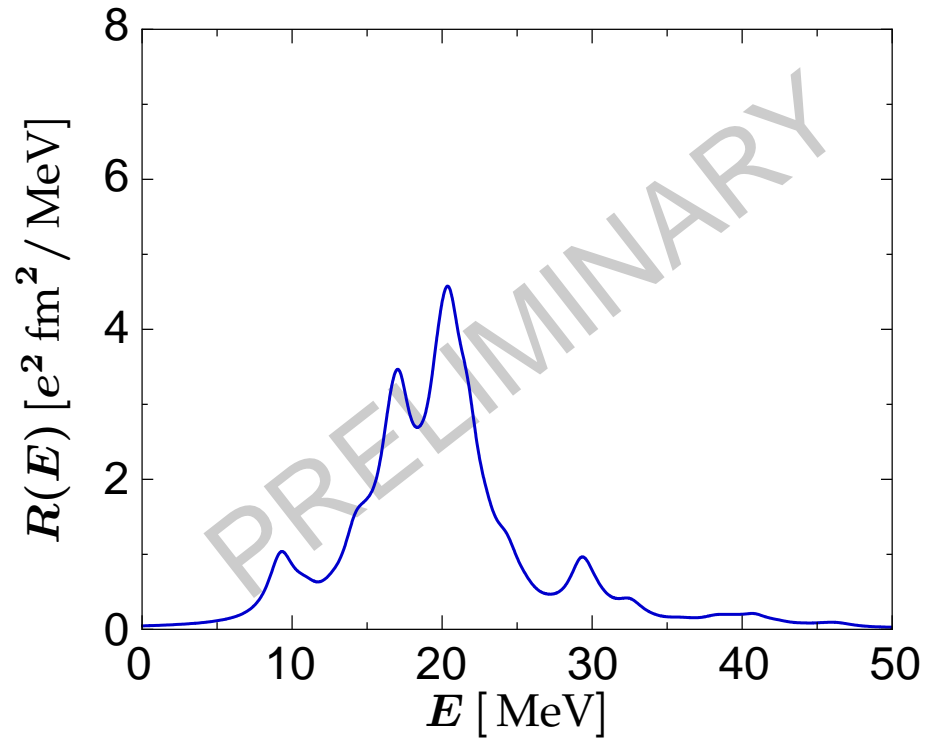
☞ implement explicit correction for isolated spurious poles

Duguet, Lacroix et al.

Pairing in the UCOM Framework

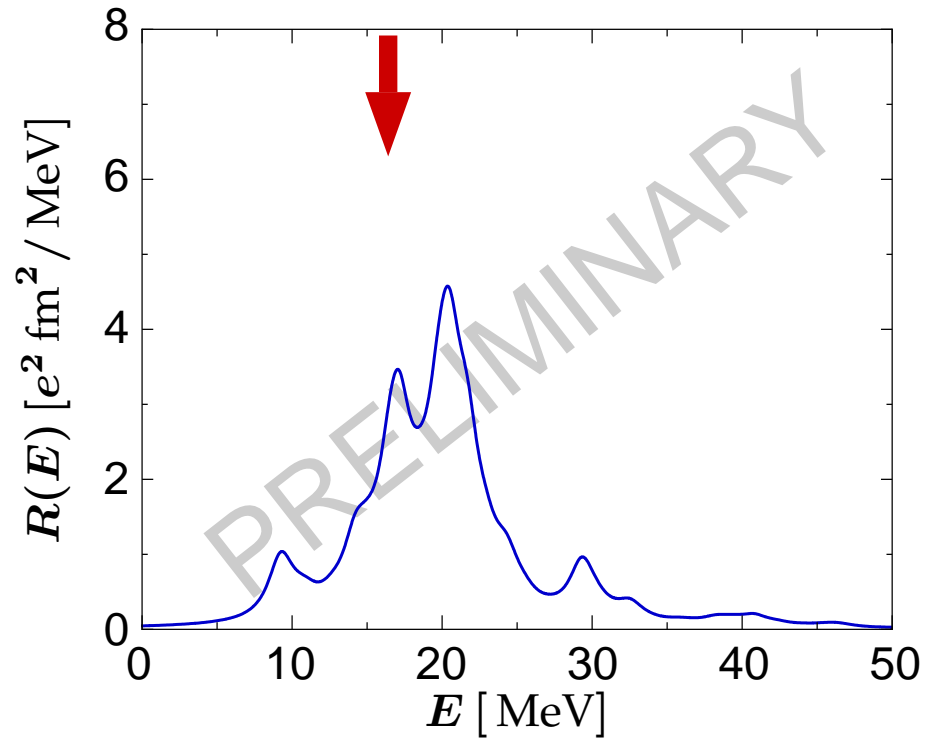
Quasiparticle RPA

QRPA: Dipole Response of ^{130}Sn



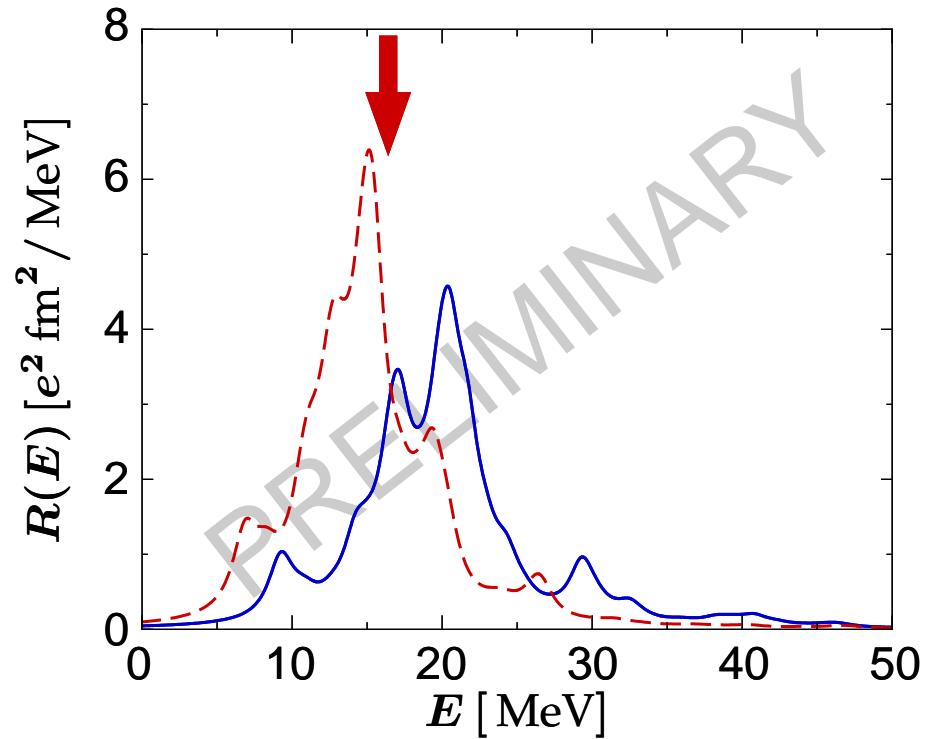
$(\bar{\alpha} [\text{fm}^4], C_{3N} [\text{GeV fm}^6])$ — (0.04, 0.0)

QRPA: Dipole Response of ^{130}Sn



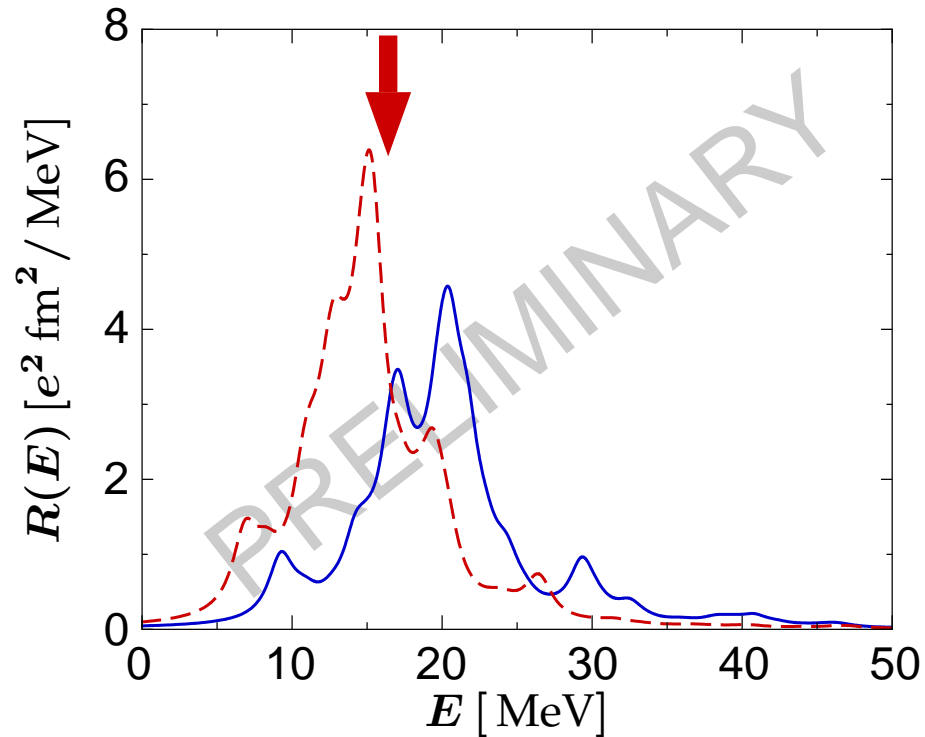
$(\bar{\alpha}[\text{fm}^4], C_{3N}[\text{GeV fm}^6])$ — (0.04, 0.0)

QRPA: Dipole Response of ^{130}Sn



$(\bar{\alpha}[\text{fm}^4], C_{3N}[\text{GeV fm}^6])$ — (0.04, 0.0) - - - (0.04, 1.2)

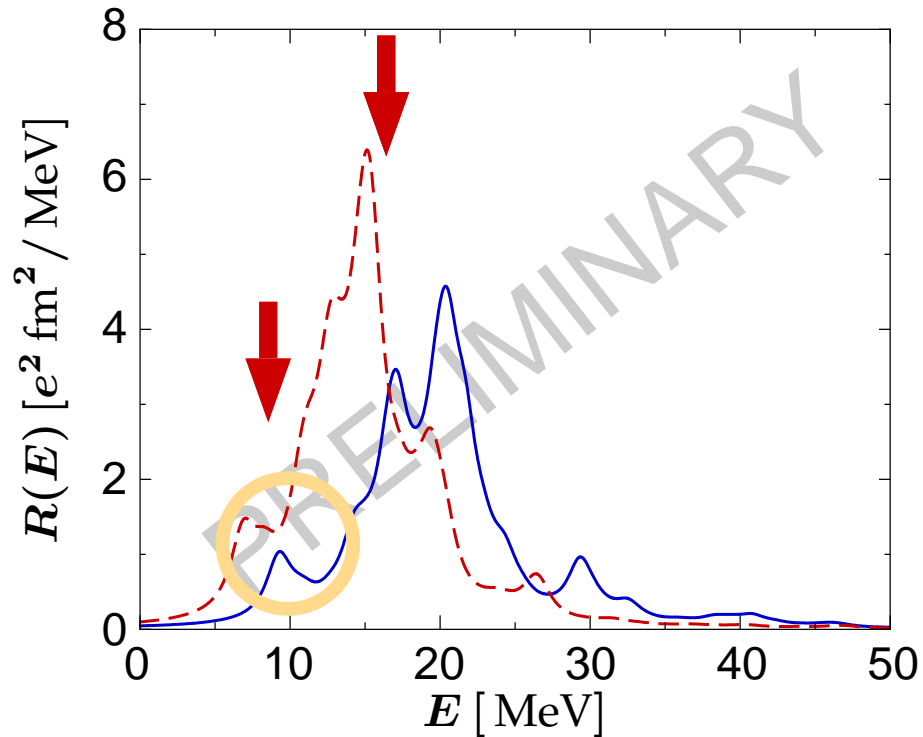
QRPA: Dipole Response of ^{130}Sn



E [MeV]	TRK [%]	N_{neut} [%]
12.86	~ 9	20.2
15.03	~ 14	52.2
15.43	~ 27	24.7

$(\bar{\alpha} [\text{fm}^4], C_{3N} [\text{GeV fm}^6])$
 — (0.04, 0.0)
 - - - (0.04, 1.2)

QRPA: Dipole Response of ^{130}Sn

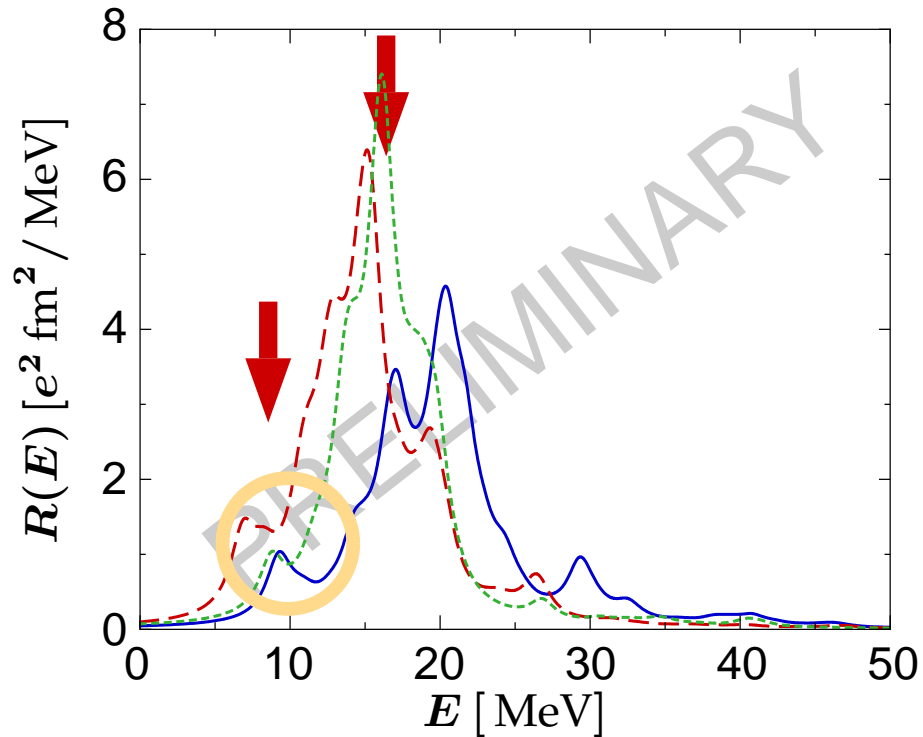


E [MeV]	TRK [%]	N_{neut} [%]
12.86	~ 9	20.2
15.03	~ 14	52.2
15.43	~ 27	24.7
6.86	~ 5	94.7
$\nu 2s_{1/2} \rightarrow \nu 2p_{3/2}$		36.0
$\nu 1d_{3/2} \rightarrow \nu 2p_{1/2}$		17.6
$\nu 2s_{1/2} \rightarrow \nu 2p_{1/2}$		10.6

$(\bar{\alpha} [\text{fm}^4], C_{3N} [\text{GeV fm}^6])$
— (0.04, 0.0)
- - - (0.04, 1.2)

Pygmy Dipole Resonance:
 significantly enhanced collectivity in QRPA

QRPA: Dipole Response of ^{130}Sn



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$(\bar{\alpha} [\text{fm}^4], C_{3N} [\text{GeV fm}^6])$ — (0.04, 0.0) - - (0.04, 1.2) ··· (0.05, 1.2)

Pygmy Dipole Resonance:
significantly enhanced collectivity in QRPA

Conclusions

Conclusions

Status

- **fully consistent** framework for HF(B), PNP, like-particle & charge-exchange (Q)RPA
- **inclusion of $3N$ forces**: density-dependent interaction for HFB, QRPA, ...

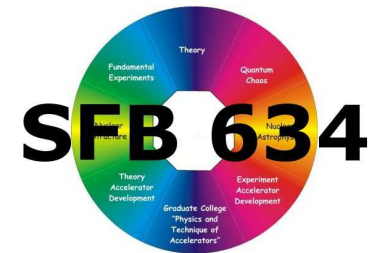
Outlook & Challenges

- density-dependent interactions in projection methods, GCM, ... (multi-reference scenarios)
- **dressed/renormalized single-particle energies** — e.g. self-consistent coupling to surface modes (HFB+QRPA)
- odd nuclei

Epilogue...

My Collaborators

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Gesellschaft für Schwerionenforschung (GSI)



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