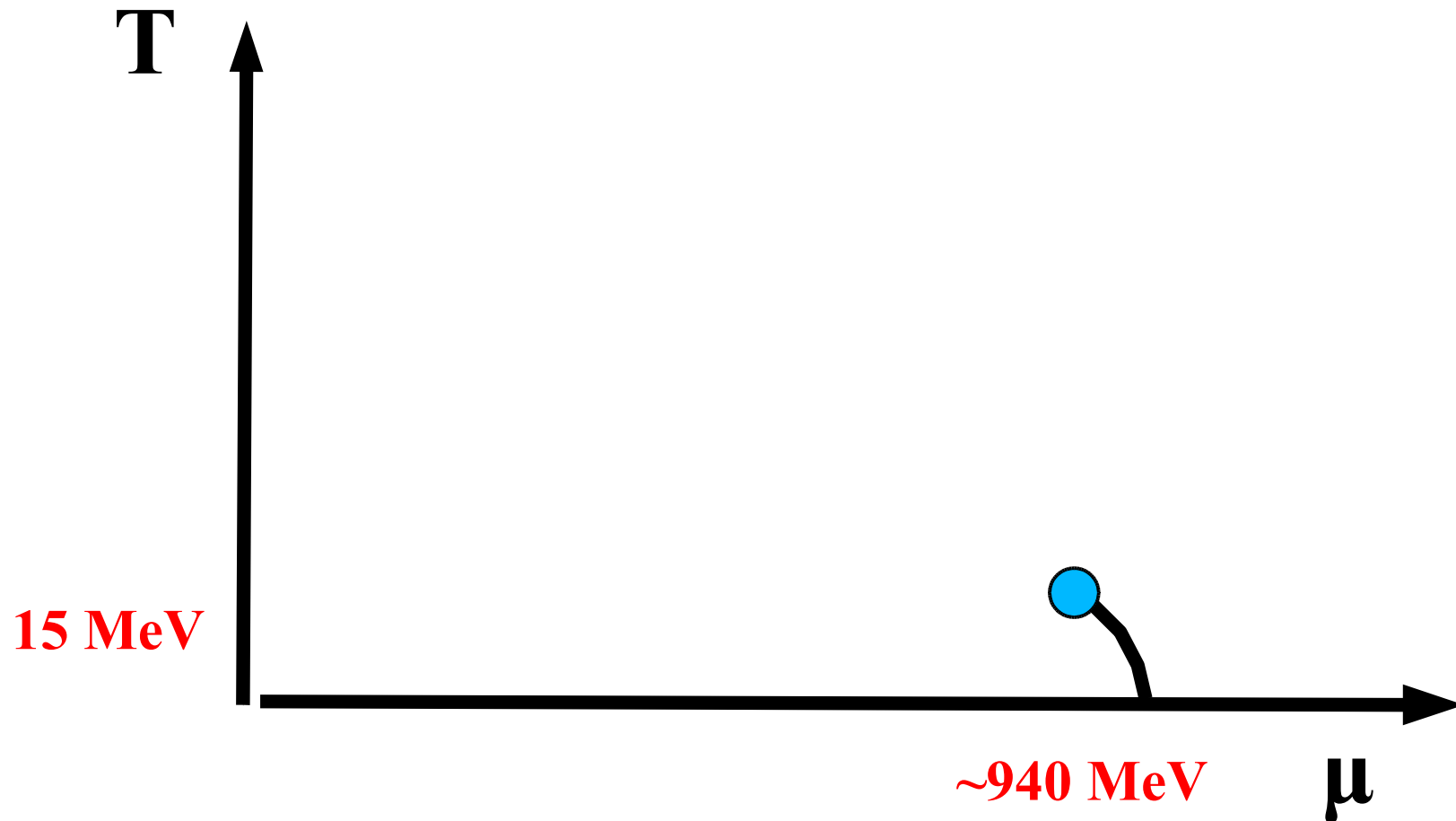


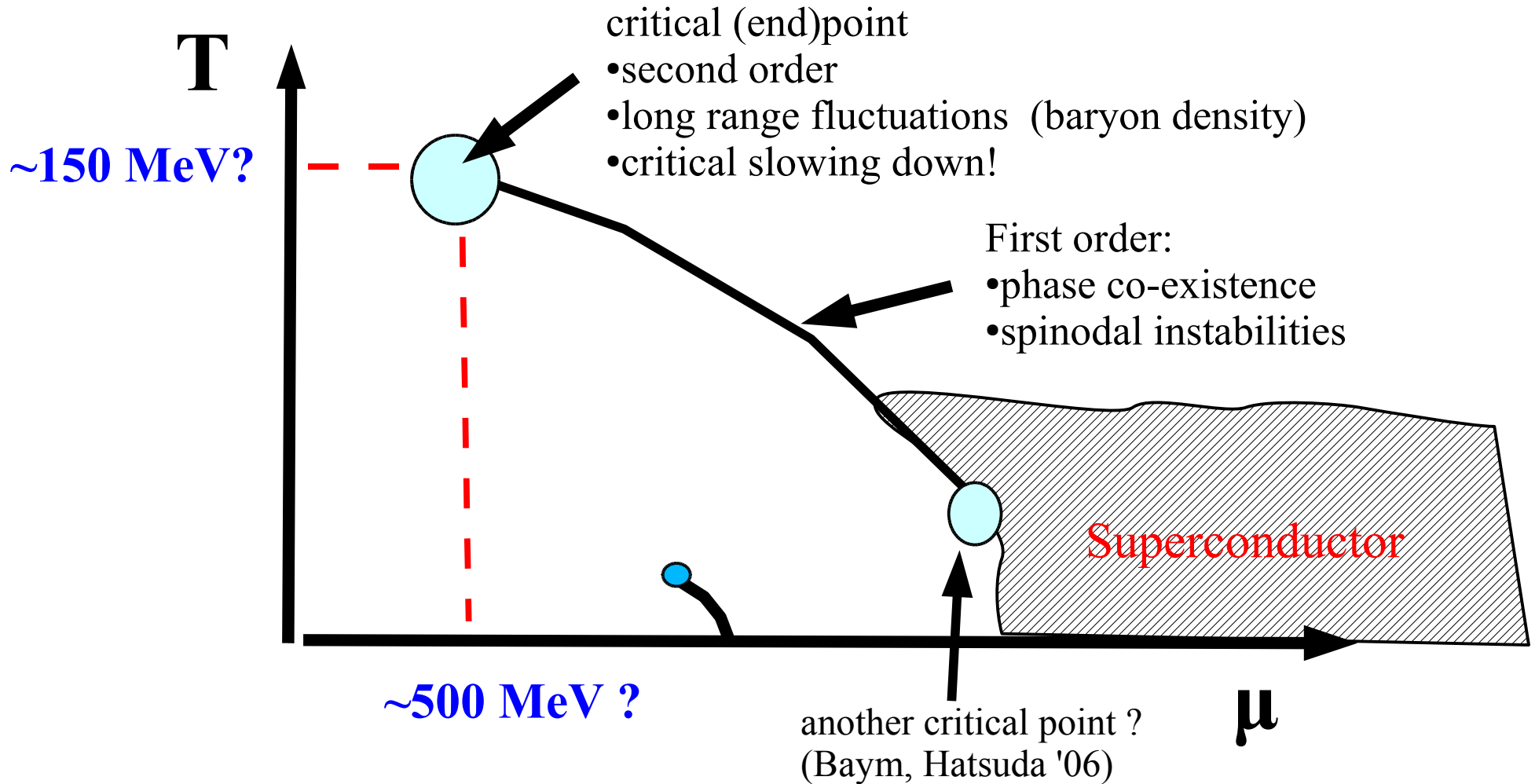
Phase Transition Phenomena in Nuclear Matter

V. Koch, LBNL, Berkeley

The QCD Phase Diagram from experiment

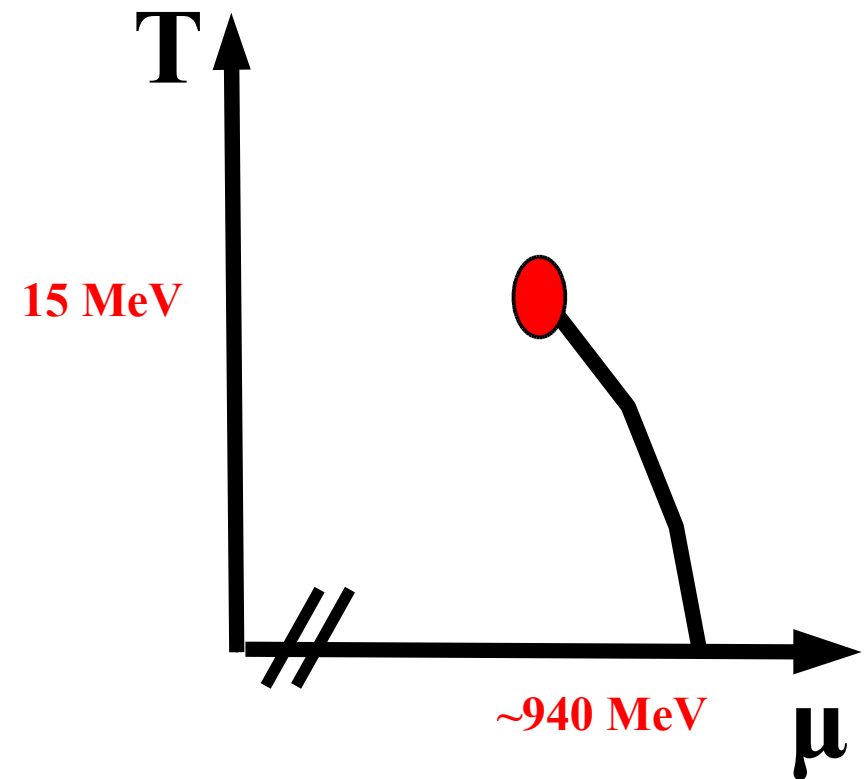
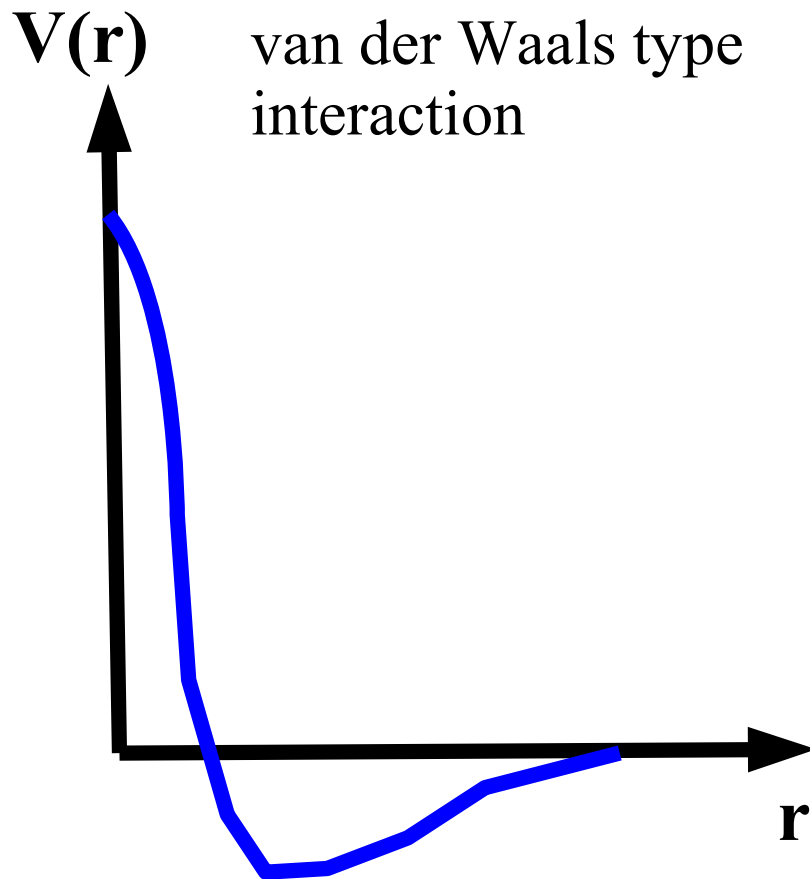


The QCD Phase Diagram (from a theorist's perspective)

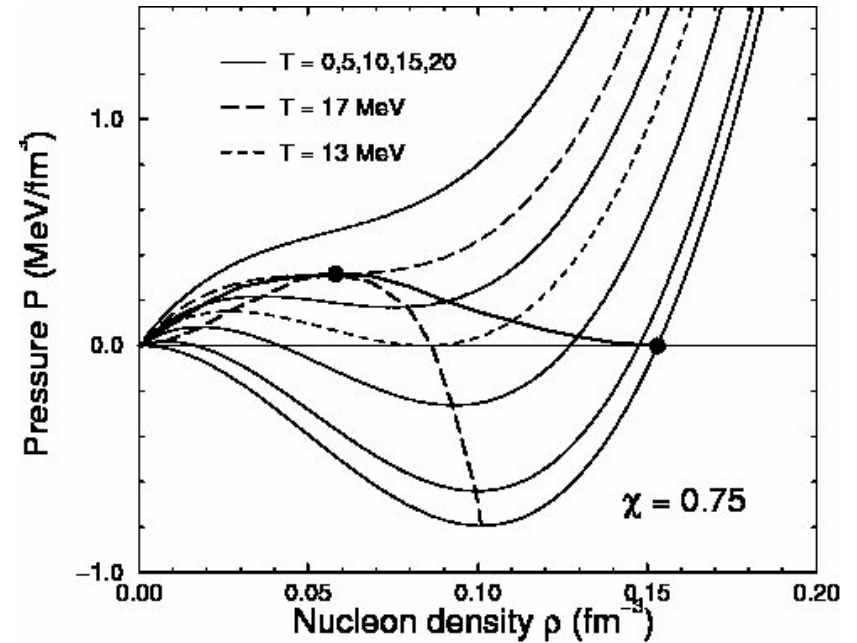
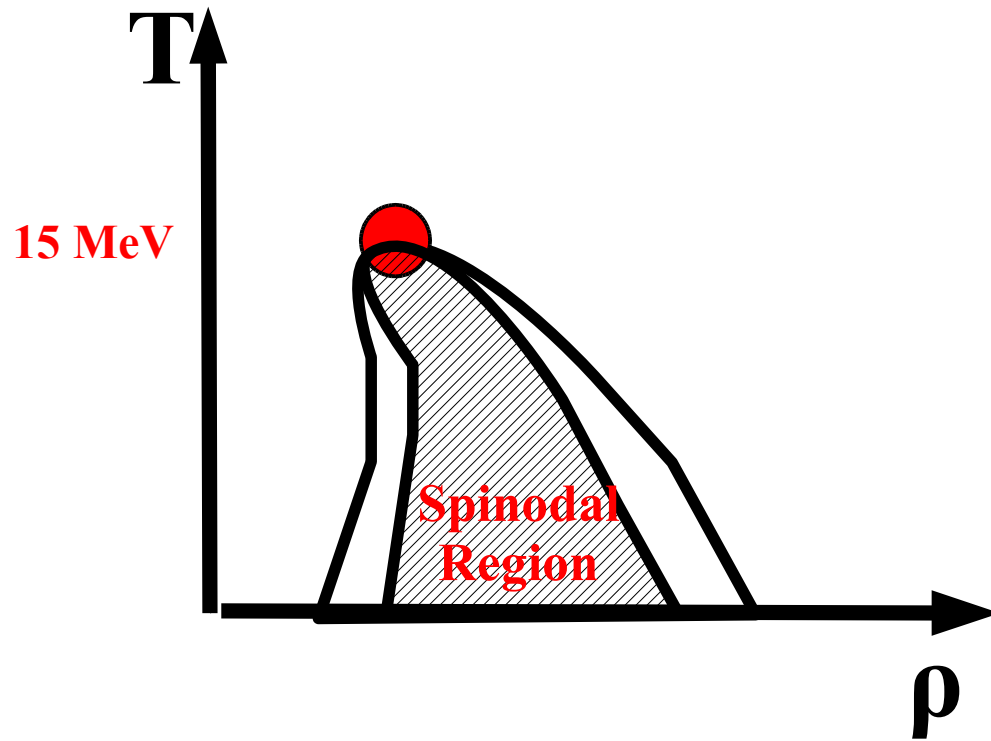


N.B.: Critical point of water: $T_c = 647.096$ K, $p_c = 22.064$ MPa, $\rho_c = 322$ kg/m³

The Nuclear Liquid Gas Phase Transition

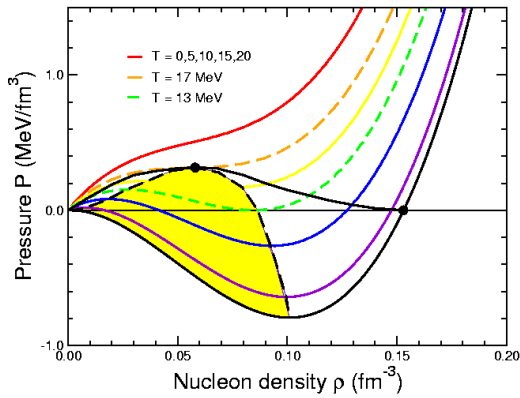


Nuclear Liquid-Gas Transition

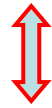


Spinodal Multifragmentation

Nuclear EoS:

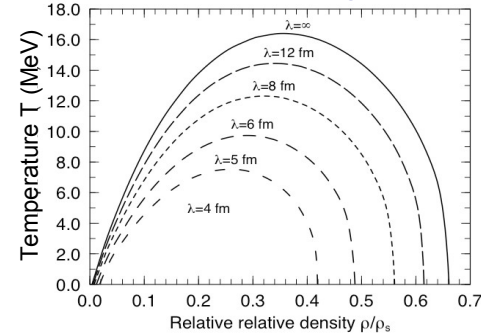


1st order phase transtion

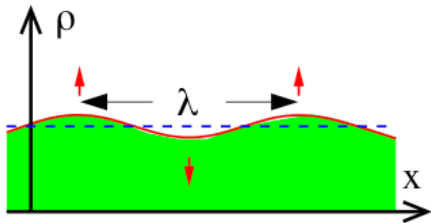
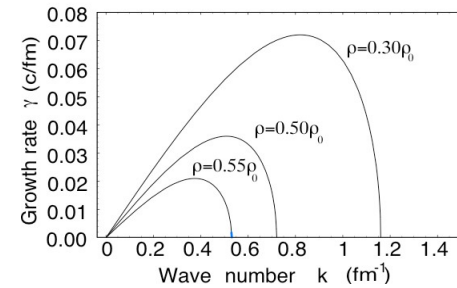


Spinodal instability

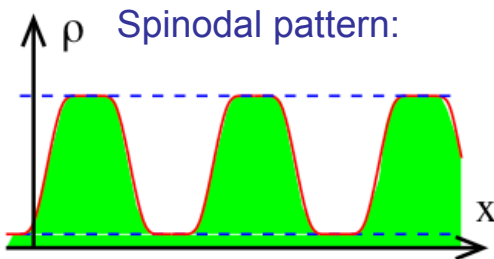
Spinodal region:



Growth rates:



Density undulations may be amplified



Ph Chomaz, M Colonna, J Randrup
Nuclear Spinodal Fragmentation
Physics Reports 389 (2004) 263



Fragments
≈ equal!

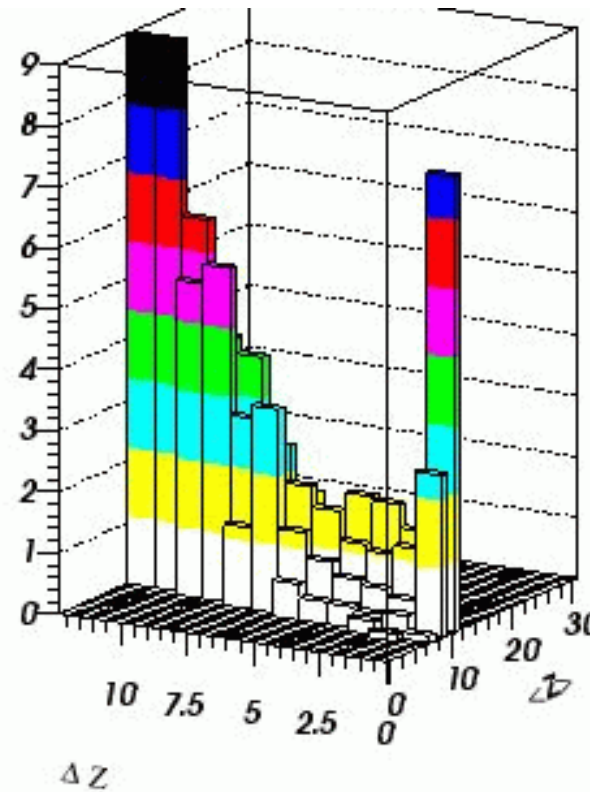
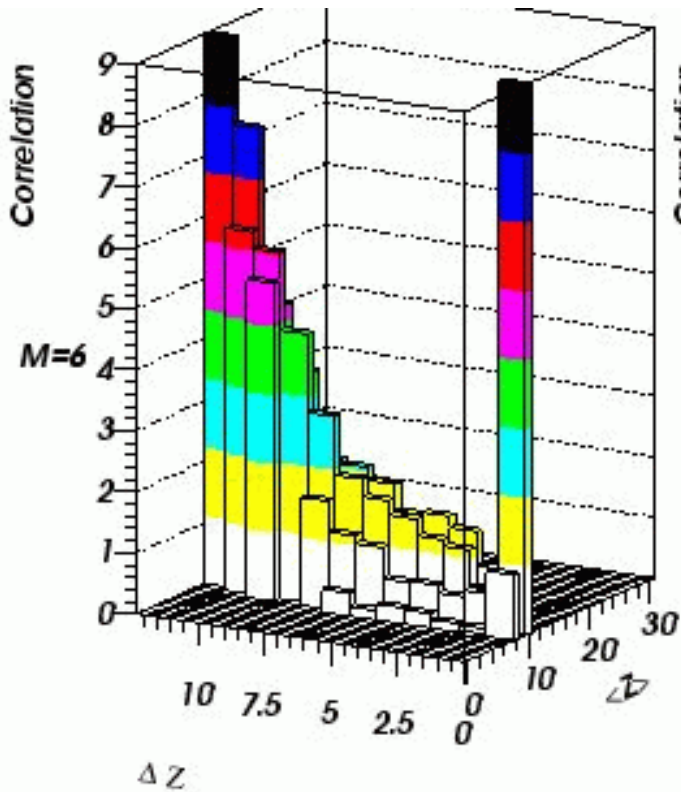


Highly non-statistical => Good candidate signature

Spinodal decomposition in nuclear multifragmentation

32 MeV/A Xe + Sn (b=0)
(select events with 6 IMFs)

Bin wrt $\left\{ \begin{array}{l} \langle Z \rangle : \text{average IMF charge} \\ \Delta Z : \text{dispersion in IMF charge} \end{array} \right.$



Experiment (*INDRA @ GANIL*)
Borderie *et al*, PRL 86 (2001) 3252

Theory (*Boltzmann-Langevin*)
Chomaz, Colonna, Randrup, ...

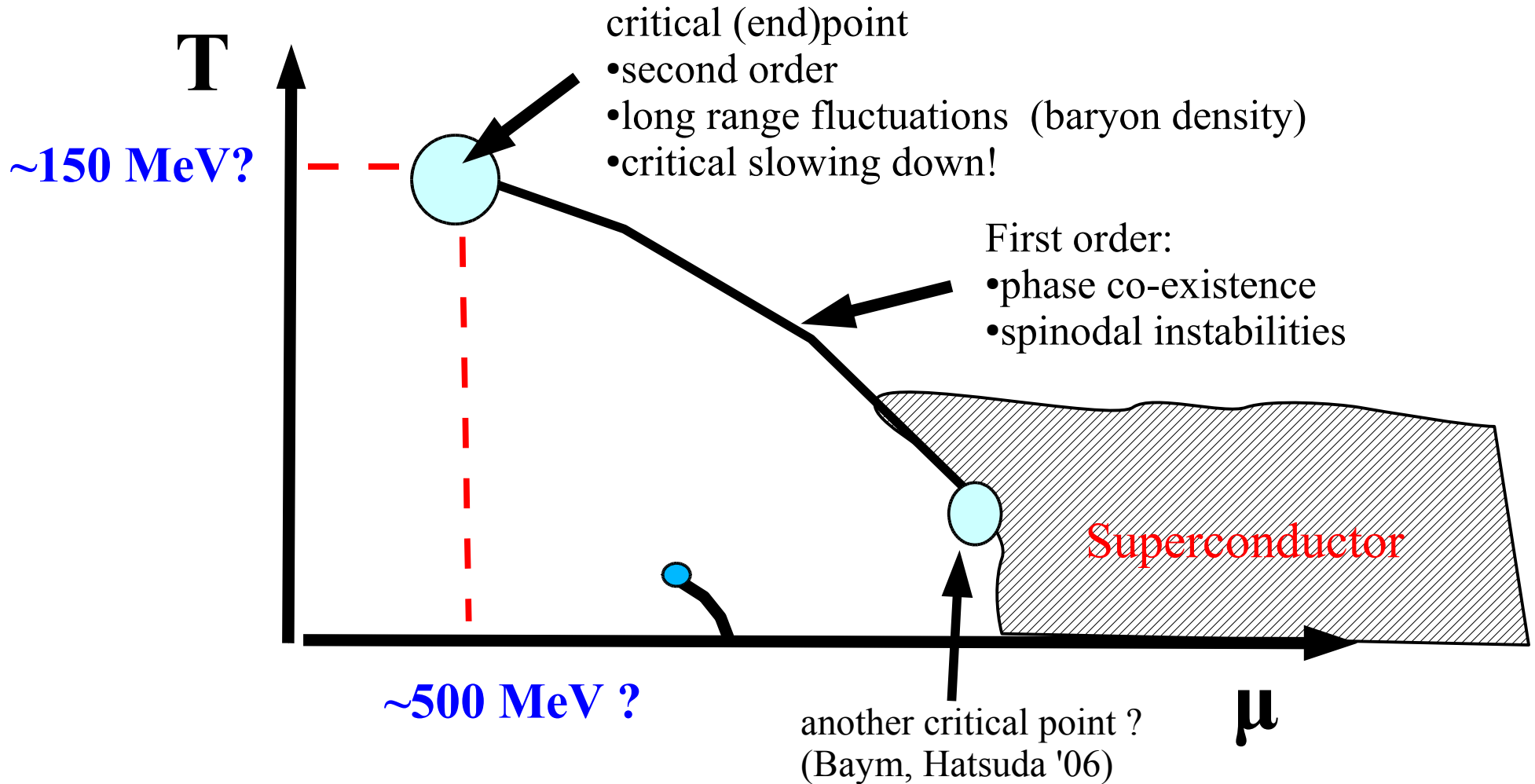
J. Randrup

Summary Nuclear Liquid Gas

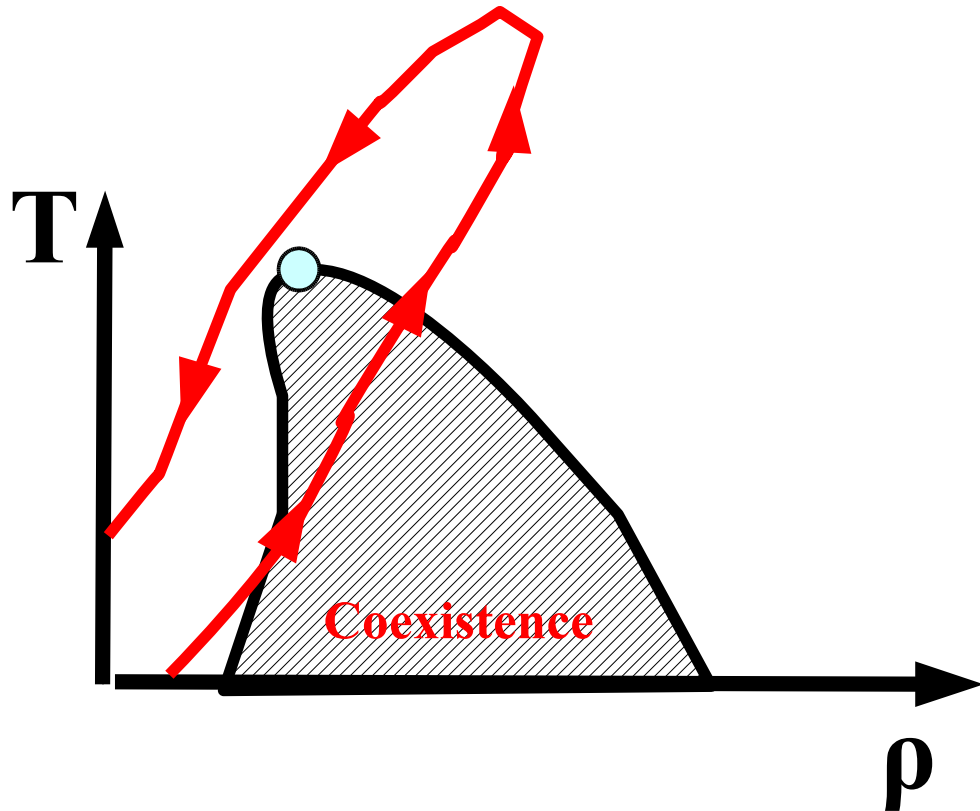
- Conceptually very straightforward
 - Force of van der Waals type
- Signs for co-existence has been found
 - Spinodal
 - Systematics of fragment distribution follows Fisher model
 - Extrapolate to critical point
- Phases are rather well defined
- >20 years of work !

The QCD Phase Diagram

(from theory)



Critical point vs co-existence



- Difficult to “hit” a point!
- Lesson learned from nuclear Liquid gas:
 - Establish co-existence and extrapolate to CP
 - Carefully chose energy such that system stalls in co-existence region

Fluctuations and Correlations in thermal system

e.g. Lattice QCD

$$Z = \text{Tr}[\exp(-\beta(H - \mu_Q Q - \mu_B B - \mu_S S))]$$

Mean :

$$\langle \alpha \rangle = T \frac{\partial}{\partial \mu_\alpha} \log(Z) = - \frac{\partial}{\partial \mu_\alpha} F$$

Variance:

$$\langle (\delta \alpha)^2 \rangle = T^2 \frac{\partial^2}{\partial \mu_\alpha^2} \log(Z) = -T \frac{\partial^2}{\partial \mu_\alpha^2} F$$

$\alpha, \beta = Q, B, S$

Co-Variance:

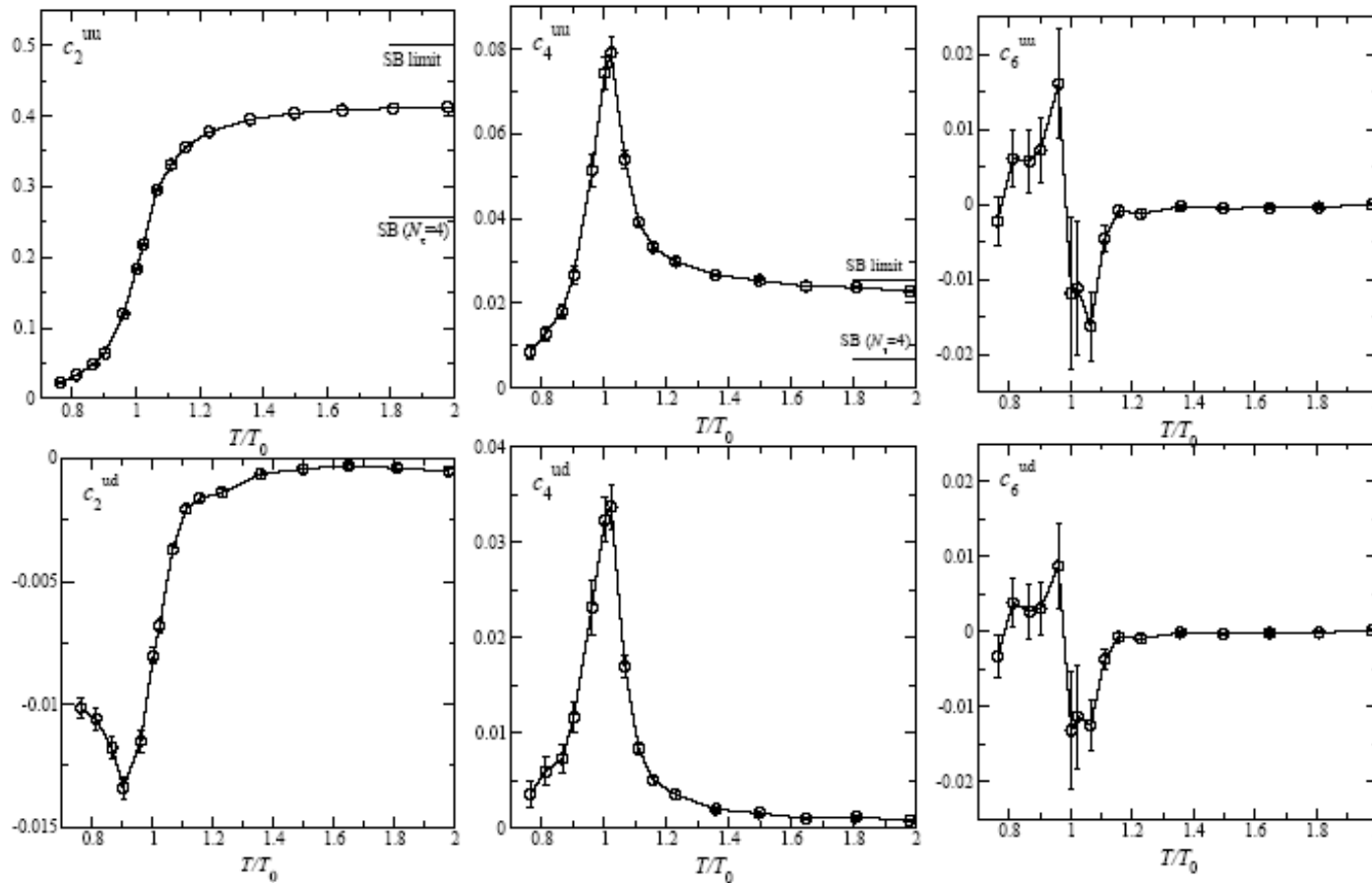
$$\langle (\delta \alpha)(\delta \beta) \rangle = T^2 \frac{\partial^2}{\partial \mu_\alpha \partial \mu_\beta} \log(Z) = -T \frac{\partial^2}{\partial \mu_\alpha \partial \mu_\beta} F$$

Susceptibility:

$$\chi_{\alpha\beta} = -\frac{1}{V} \frac{\partial^2}{\partial \mu_\alpha \partial \mu_\beta} F = -\frac{1}{V} \frac{\partial}{\partial \mu_\alpha} \langle \beta \rangle$$

Lattice-QCD susceptibilities

$$\frac{\chi(T, \mu_q)}{T^2} = 2c_2 + 12c_4 \left(\frac{\mu_q}{T}\right)^2 + 30c_6 \left(\frac{\mu_q}{T}\right)^4 + \dots$$



Rule of thumb:

$$c_n \sim \langle X^n \rangle$$

$$X = B, Q, S, \dots$$

Alton et al, PRD 66 074507 (2002)

Susceptibilities and Phasetransitions

$$Z = \text{Tr}[\exp(-\beta(H - \mu N))]$$

Susceptibility: $\chi \sim \frac{1}{V} \frac{\partial^r}{\partial \mu^r} \log(Z) = \frac{1}{V} (\langle N^r \rangle - \langle N \rangle^r)$

Poisson: $\chi \sim \frac{\langle N \rangle}{V}$ independent of volume  $\langle (\delta N)^2 \rangle = N \sim V$

In general: $\chi \sim \frac{1}{V} \int d^r x d^r y \langle \rho(x) \rho(y) \rangle_{connected} = \int d^r r \langle \rho(r) \rho(\cdot) \rangle_{connected} \sim \xi^r$

$$\langle \rho(r) \rho(\cdot) \rangle_{connected} \sim \frac{e^{(-r/\xi)}}{r} \quad \xi = \text{correlation length}$$

Cross-over: $\xi = \text{const} \rightarrow \chi = \text{const} \rightarrow \langle (\delta N)^2 \rangle \sim V$

Second Order: $\xi \sim V^{(1/3)} \rightarrow \chi \sim V^{(2/3)} \rightarrow \langle (\delta N)^2 \rangle \sim V^{(5/3)}$

First Order: $\langle \rho(r) \rho(0) \rangle = \text{const} \rightarrow \chi \sim V \rightarrow \langle (\delta N)^2 \rangle \sim V^2$

Susceptibilities and Observables

Susceptibility:
$$\chi \sim \frac{1}{V} \frac{\partial^2}{\partial \mu^2} \log(Z) = \frac{1}{V} (\langle N^2 \rangle - \langle N \rangle^2)$$

Fluctuations of some sort!

Cross-over:
$$\xi = \text{const} \rightarrow \chi = \text{const} \rightarrow \langle (\delta N)^2 \rangle \sim V$$

Second Order:
$$\xi \sim V^{(1/3)} \rightarrow \chi \sim V^{(2/3)} \rightarrow \langle (\delta N)^2 \rangle \sim V^{(5/3)}$$

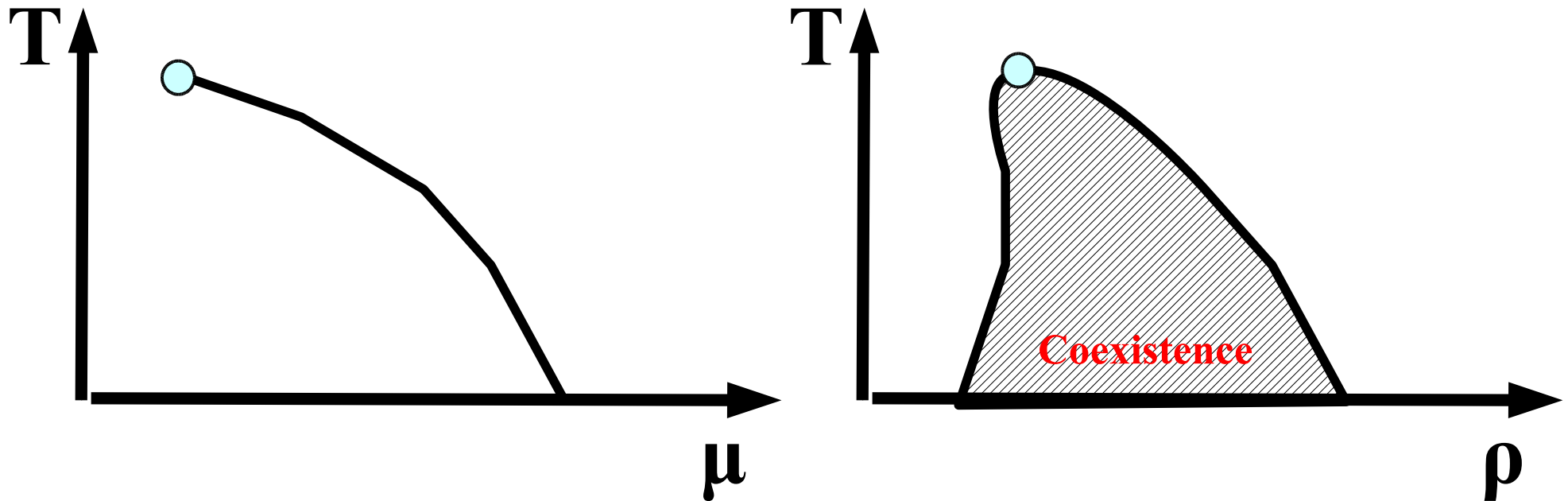
First Order:
$$\langle \rho(r) \rho(0) \rangle = \text{const} \rightarrow \chi \sim V \rightarrow \langle (\delta N)^2 \rangle \sim V^2$$

Since fluctuations diverge at phase transition **any** sort will do!

System size dependence!

Note of caution: Co-variances also diverge; trigger?

Order Parameter

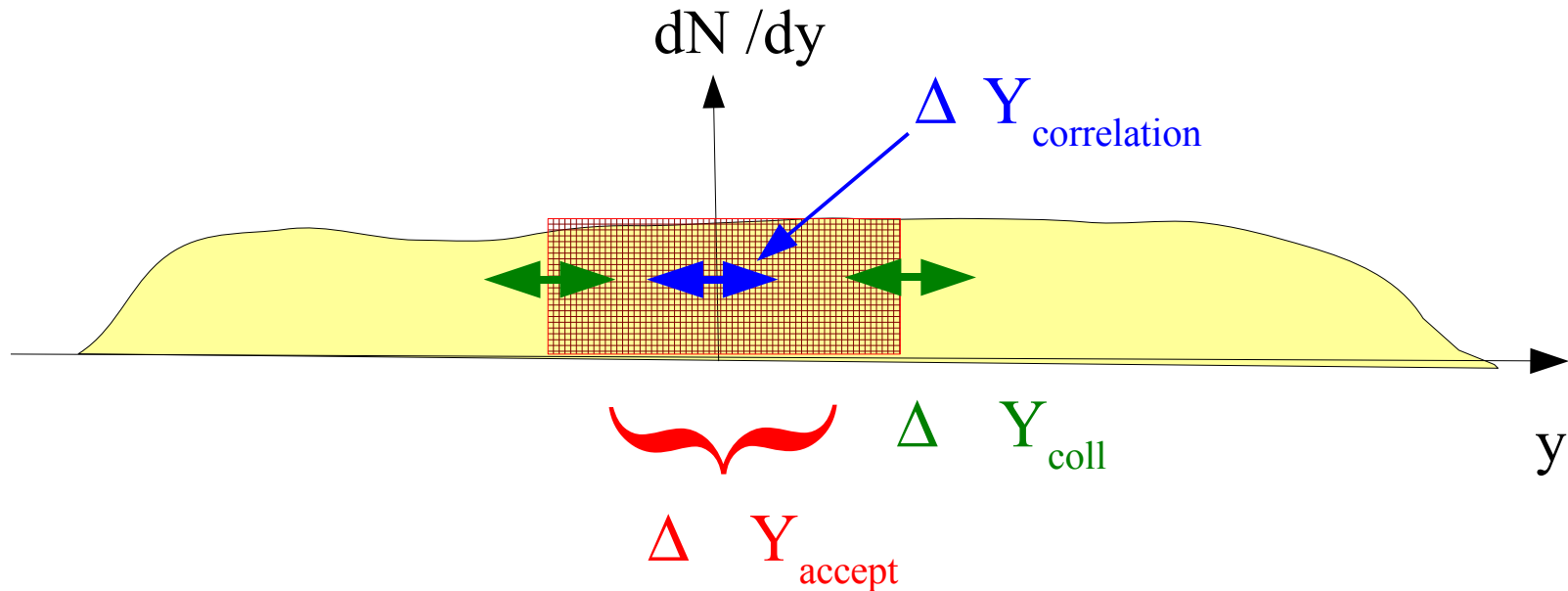


Baryon density is a good order parameter
density fluctuations are a good observable (theoretically...)



Baryon Number fluctuations also good in principle, but
global baryon number conservation is an issue at low energies

“Charge” fluctuations

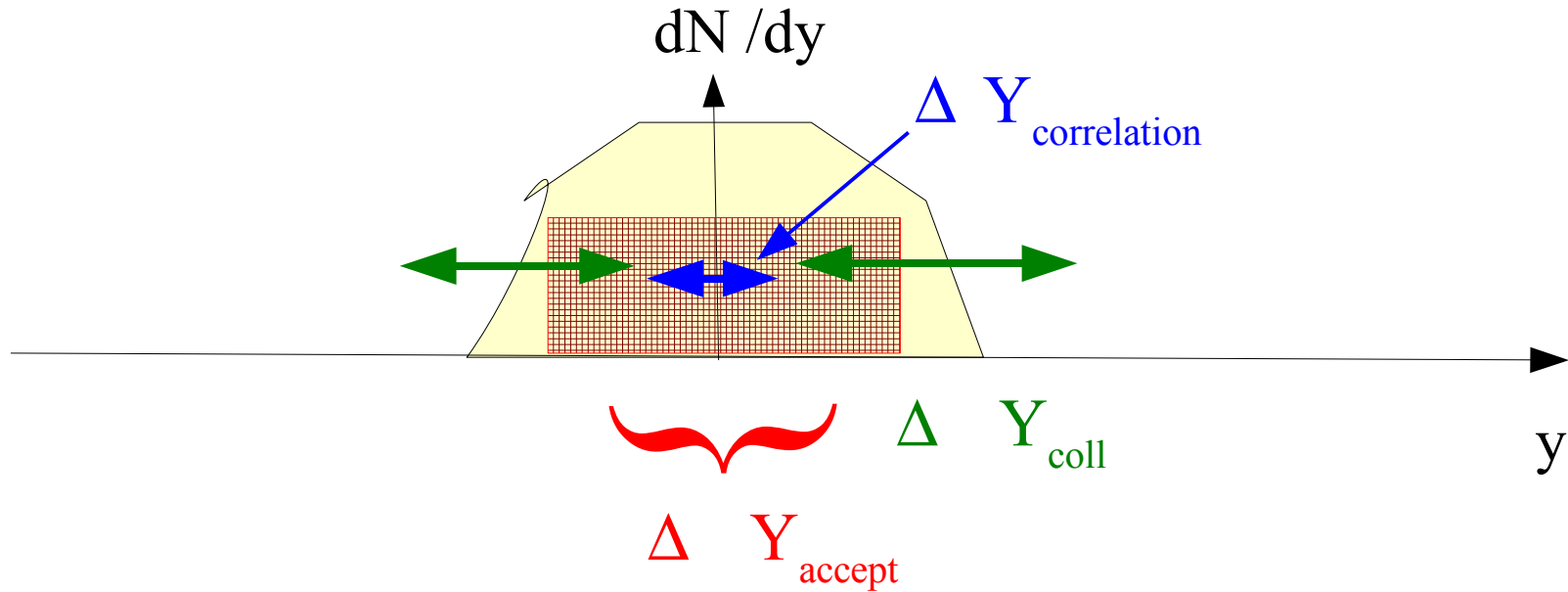


Condition for “charge” fluctuations:

1) $\Delta Y_{\text{correlation}} \ll \Delta Y_{\text{accept}}$ (catch the physics)

3) $\Delta Y_{\text{total}} \gg \Delta Y_{\text{accept}} \gg \Delta Y_{\text{coll}}$ (keep the physics)

“Charge” fluctuations at SPS and below

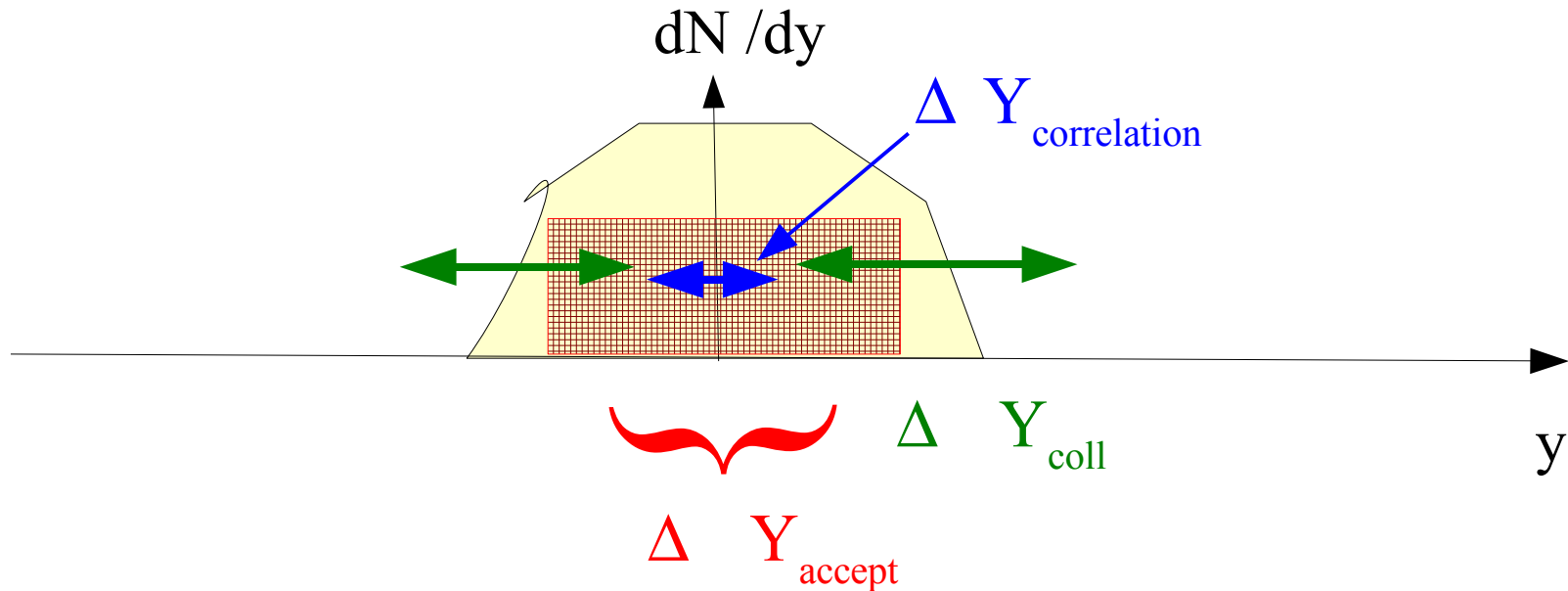


Condition for “charge” fluctuations:

1) $\Delta Y_{\text{correlation}} \ll \Delta Y_{\text{accept}}$ (catch the physics)

3) $\Delta Y_{\text{total}} \gg \Delta Y_{\text{accept}} \gg \Delta Y_{\text{coll}}$ (keep the physics)

“Charge” fluctuations at SPS and below

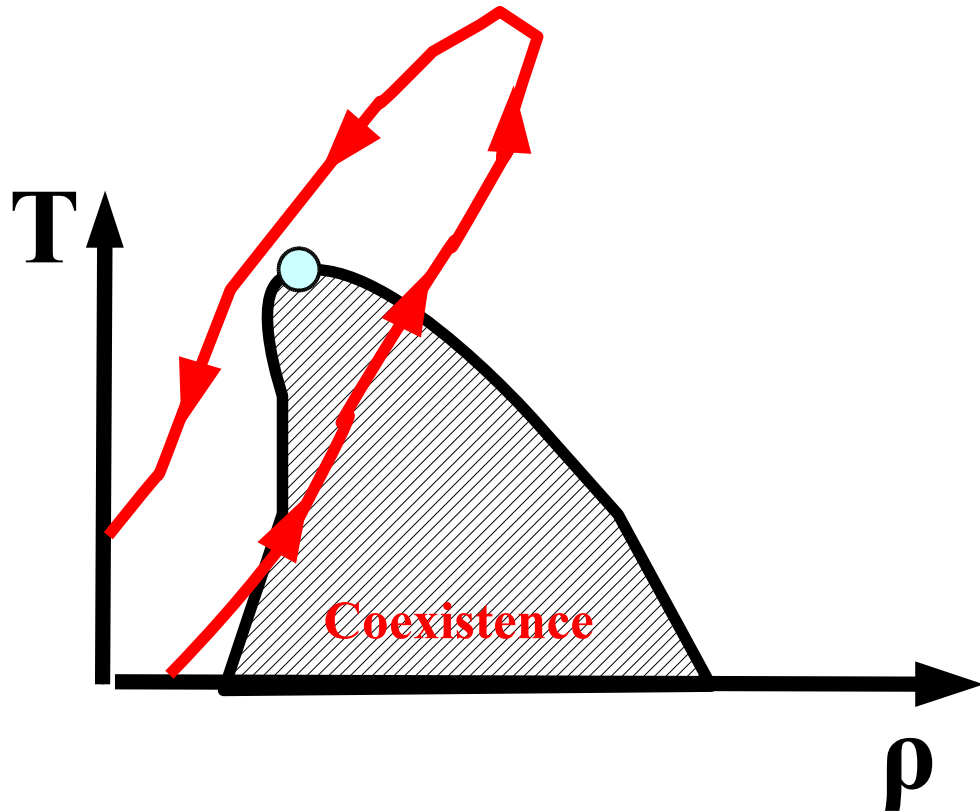


Condition for “charge” fluctuations:

1) $\Delta Y_{\text{correlation}} \ll \Delta Y_{\text{accept}}$ (catch the physics)

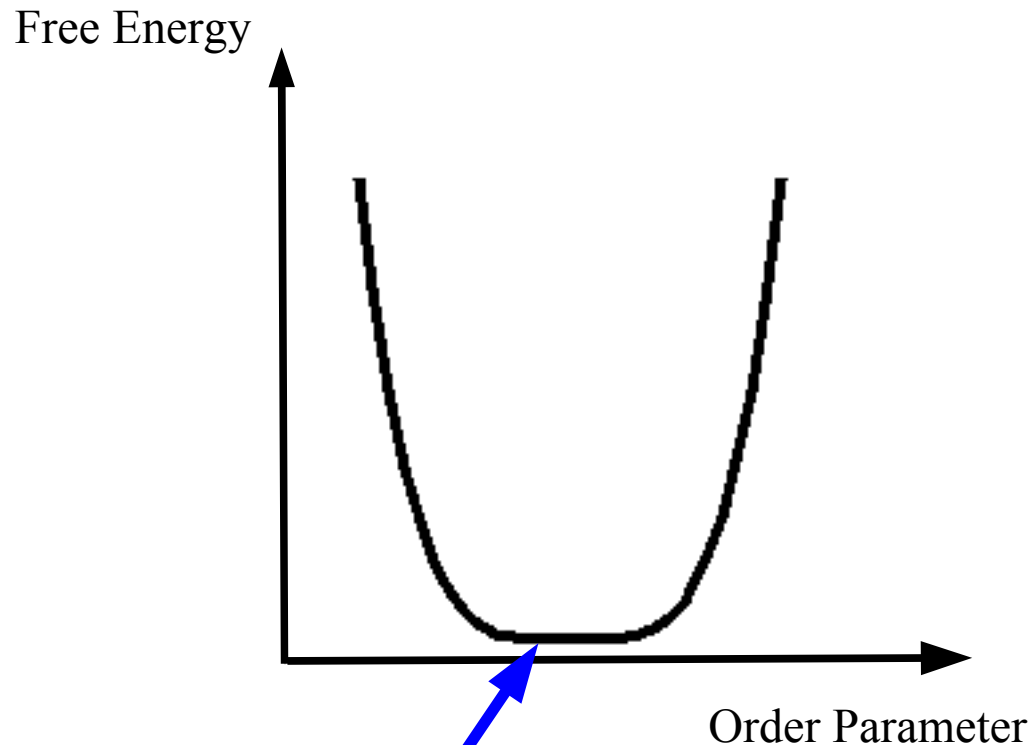
3) $\Delta Y_{\text{total}} \gg \Delta Y_{\text{accept}} \gg \Delta Y_{\text{coll}}$ (keep the physics)

Critical point vs co-existence



- Difficult to “hit” a point!
- Lesson learned from nuclear Liquid gas:
 - Establish co-existence and extrapolate to CP
 - Carefully chose energy such that system stalls in co-existence region

Second order

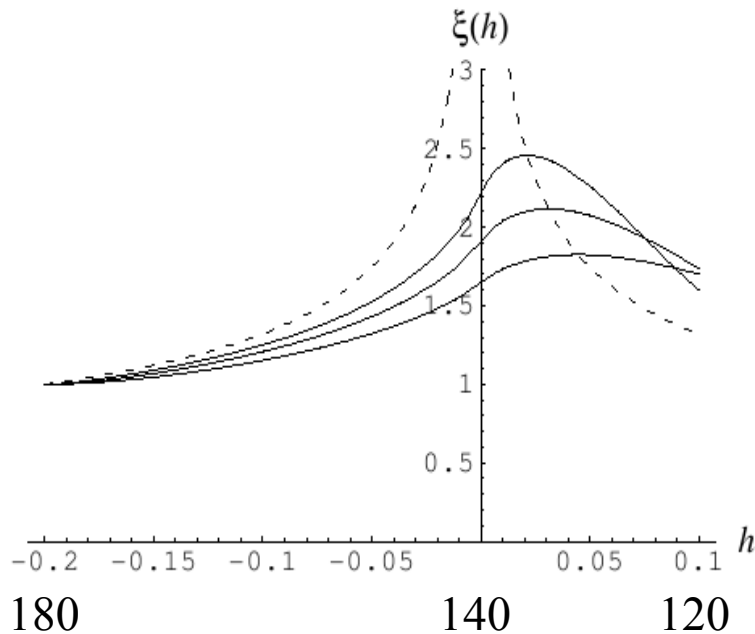


No curvature, "Mass" $=0$

- Fluctuation of order parameter at all scales
- Diverging susceptibilities
 $\sim 1/(\text{"Mass"})^2$
- Diverging correlation length
 $\sim 1/(\text{"Mass"})$
- Universality
- Critical slowing down !

Second order

correlation length $\sim 1/m_\sigma$

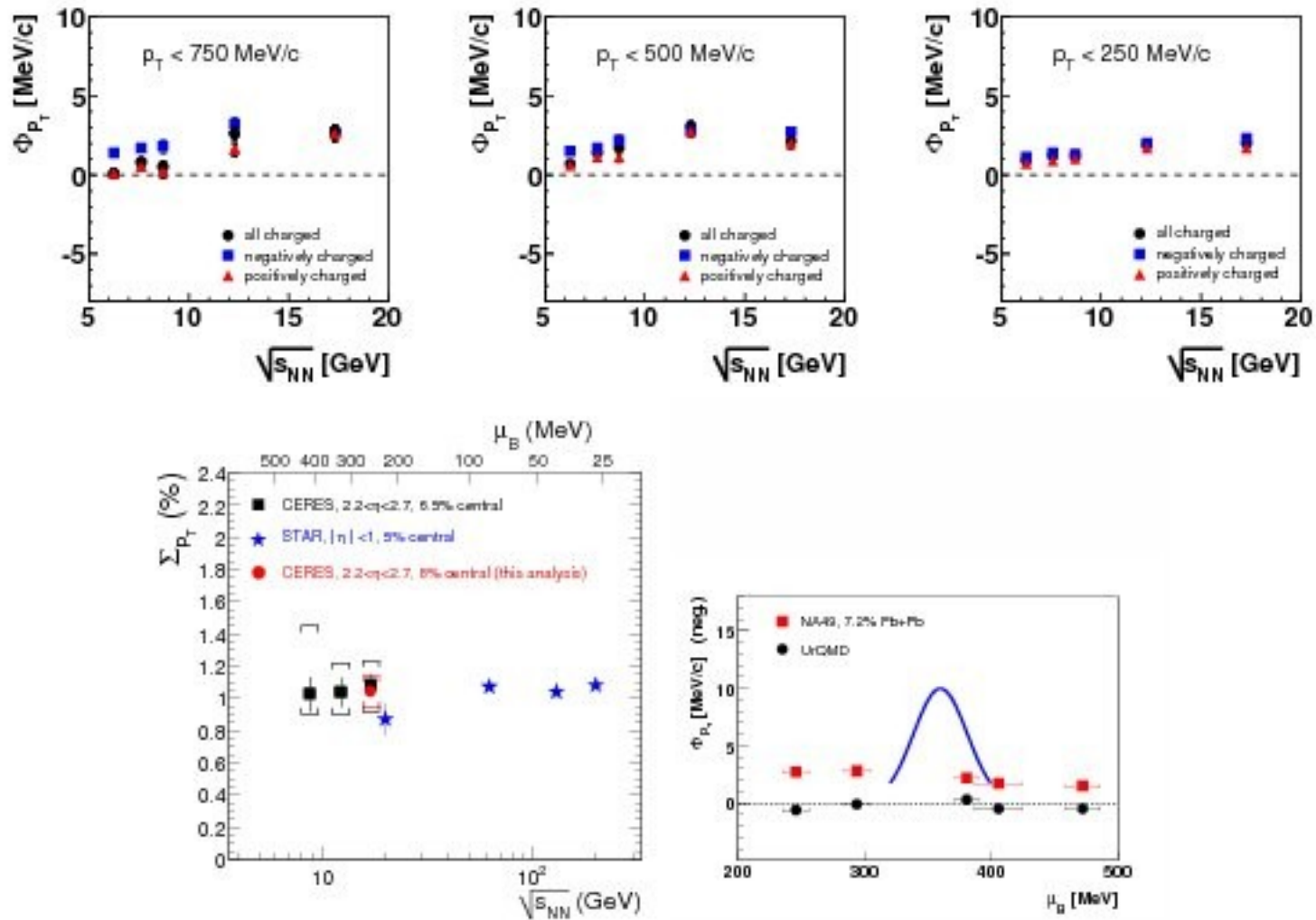


Bernikov, Rajagopal, hep-ph/9912274

- Critical slowing down
- limited sensitivity on model parameters
- Max. correlation length 2-3 fm
- Translates in **3-5%** effect in p_t -fluctuations

Expect:
Maximum in excitation function
of p_t -fluctuations at low p_t

What does experiment say?



Higher cumulants?

Stephanov
arXiv:0809.3450

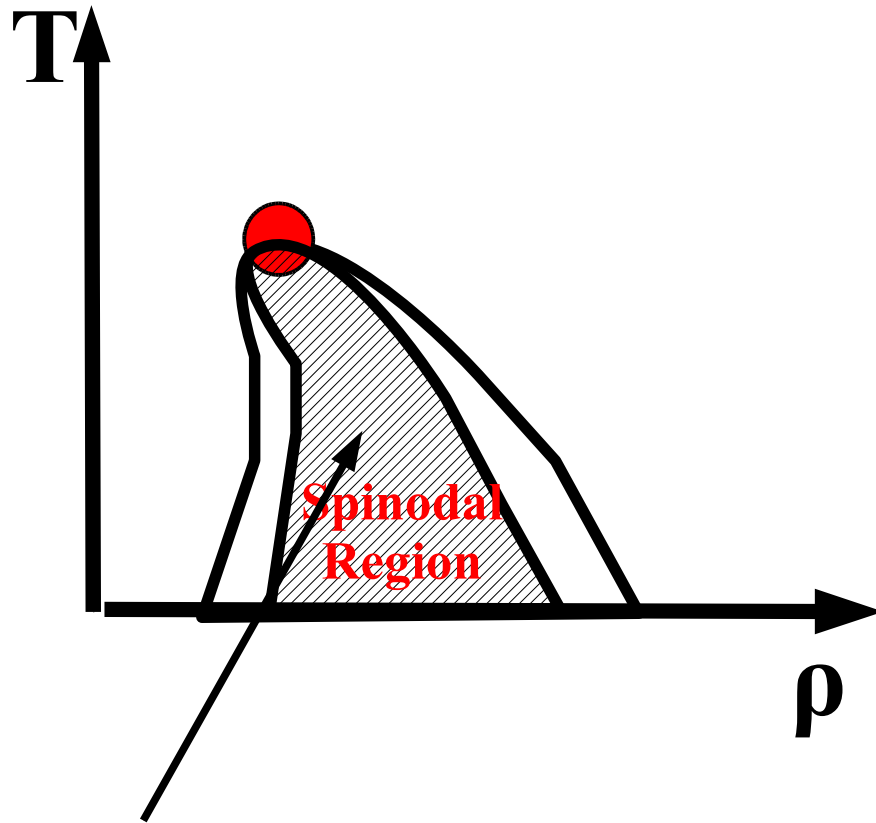
$$\omega_2 = \frac{\langle\langle \delta N \rangle^2\rangle}{\langle N \rangle} \sim \xi^2$$
$$\omega_4 = \frac{\langle\langle \delta N \rangle^4\rangle}{\langle N \rangle} \sim \xi^7$$

Higher cumulants diverge
with higher power:

5% in second order translates
20% in fourth order

Question: How does critical slowing down affect higher cumulants

Co-existence region



System should spent long time
in spinodal region

Relativistic fluid dynamics

Energy-momentum tensor:

$$T^{\mu\nu} = (\varepsilon + p)u^\mu u^\nu - pg^{\mu\nu}$$

$$p(\varepsilon)$$

Equation of motion:

$$\partial_\mu T^{\mu\nu} = 0$$

$$u^\mu = (\gamma, \gamma\mathbf{v})$$

Small disturbance in a uniform stationary fluid

$$\varepsilon(x, t) = \varepsilon_0 + \delta\varepsilon(x, t), \quad \delta\varepsilon \ll \varepsilon_0$$

First order in $\delta\varepsilon$:

$$\partial_t \delta\varepsilon(x, t) \approx (\varepsilon_0 + p_0) \partial_x v_x(x, t)$$

$$p_0 \equiv p(\varepsilon_0)$$

$$(\varepsilon_0 + p_0) \partial_t v_x(x, t) \approx \partial_x p(x, t) \approx \frac{\partial p_0}{\partial \varepsilon_0} \partial_x \delta\varepsilon(x, t)$$

Sound waves!

$$\partial_t^2 \delta\varepsilon(x, t) = \frac{\partial p_0}{\partial \varepsilon_0} \partial_x^2 \delta\varepsilon(x, t)$$

$$v_s^2 = \frac{\partial p}{\partial \varepsilon}$$

Growth rates γ_k

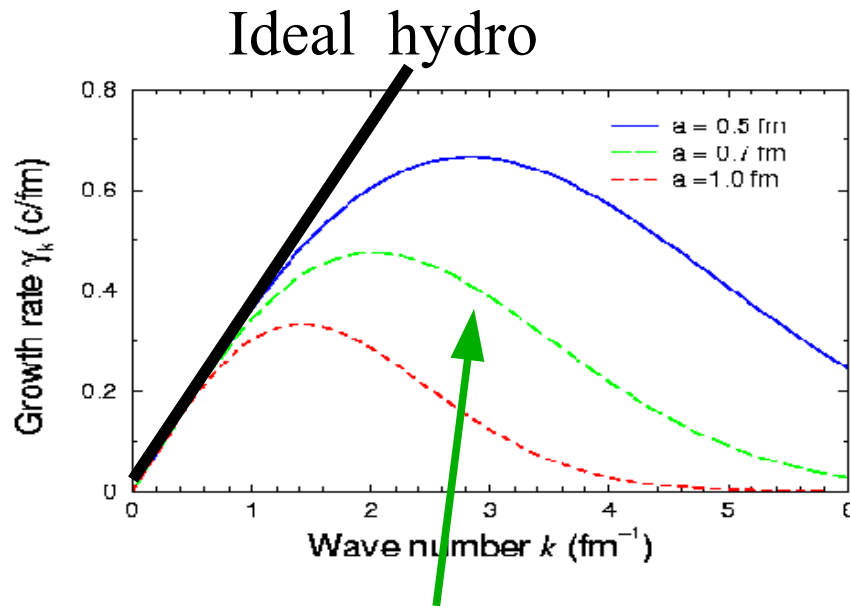
Small disturbance: $\varepsilon(x, t) = \varepsilon_0 + \delta\varepsilon(x, t)$, $\delta\varepsilon \ll \varepsilon_0$

Evolution: $\partial_t^2 \delta\varepsilon(x, t) = \frac{\partial p_0}{\partial \varepsilon_0} \partial_x^2 \delta\varepsilon(x, t) \Rightarrow \delta\varepsilon_k(x, t) \sim e^{ikx - i\omega_k t}$

Dispersion relation: $\omega_k^2 = \frac{\partial p_0}{\partial \varepsilon_0} k^2 = -\gamma_k^2 k^2 \Rightarrow \gamma_k = |v_s| k$

Local average: $\rho(r) = \langle \rho(\varepsilon(r)) \rangle$ $\omega_k^2 = \frac{\partial p_0}{\partial \varepsilon_0} g_k k^2$, $g_k = e^{-a^2 k^2 / 2}$

$\gamma \sim k$ OK for small k
 But what about $k \rightarrow \infty$?
 Ideal hydro has no scale!



a : smearing range suppresses large k



γ_k has a maximum



Spinodal pattern may develop

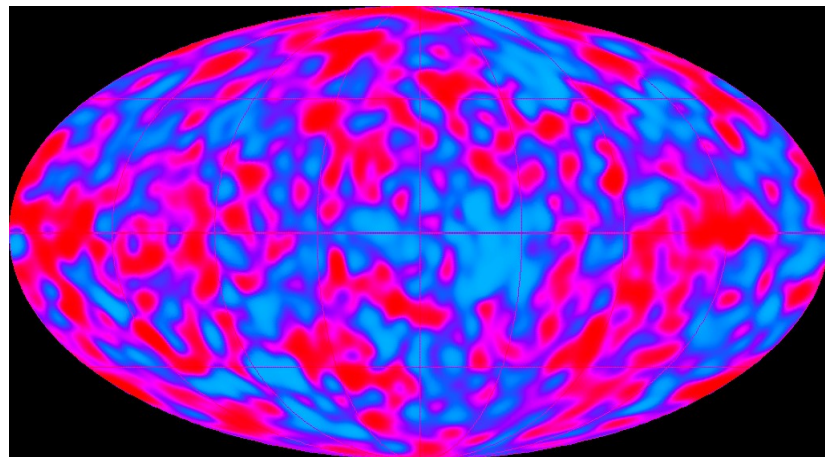
- *if* there is enough time!

Need a length scale!!
 Interface tension from lattice?!

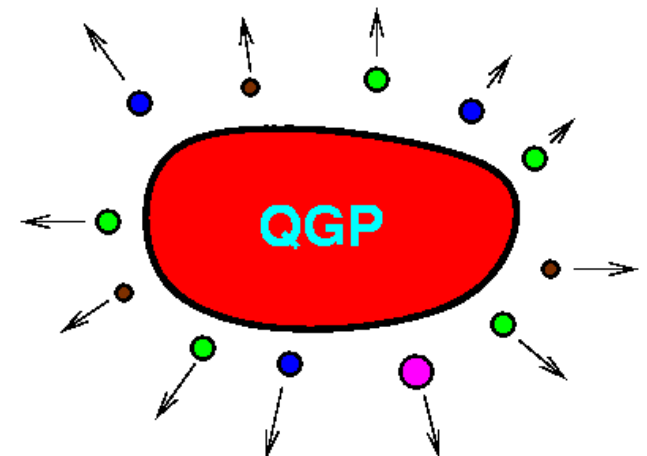
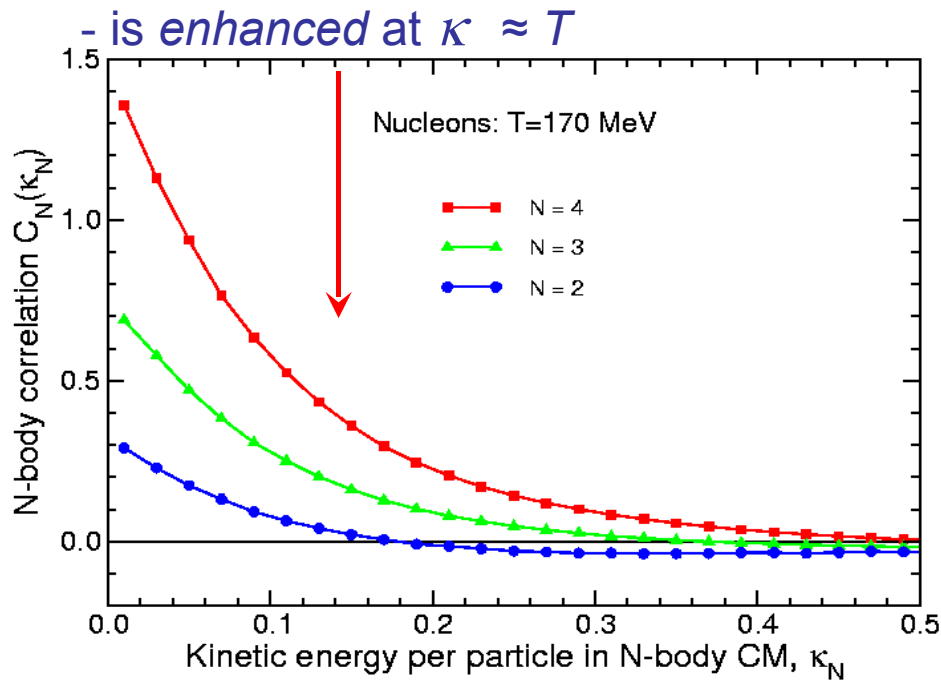
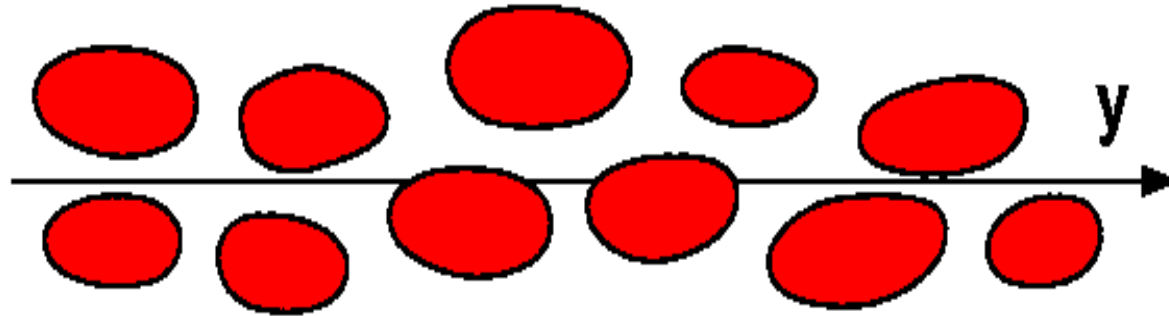
J. Randrup

How to detect clumping?

- No obvious candidates for clumps contrary to nuclear liquid gas
 - Kinematic correlations
 - Flavor correlations
- Fluctuations due to clumping

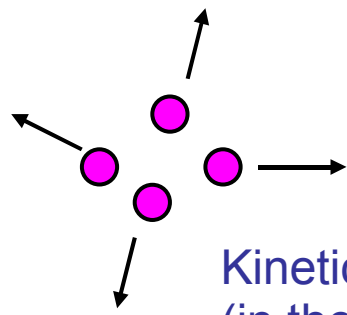


N-particle correlations



[J. Randrup, J. Heavy Ion Physics 22 (2005) 69]

Kinematic clumping =>



Invariant-mass correlations

Total four-momentum:

$$P\{\mathbf{p}_n\} = \sum_n (E_n, \mathbf{p}_n)$$

Kinetic energy per particle
(in the N -body CM frame):

$$\kappa_N\{\mathbf{p}_n\} = \frac{1}{N} \left[[P\{\mathbf{p}_n\} \cdot P\{\mathbf{p}_n\}]^{\frac{1}{2}} - \sum_n m_n \right]$$

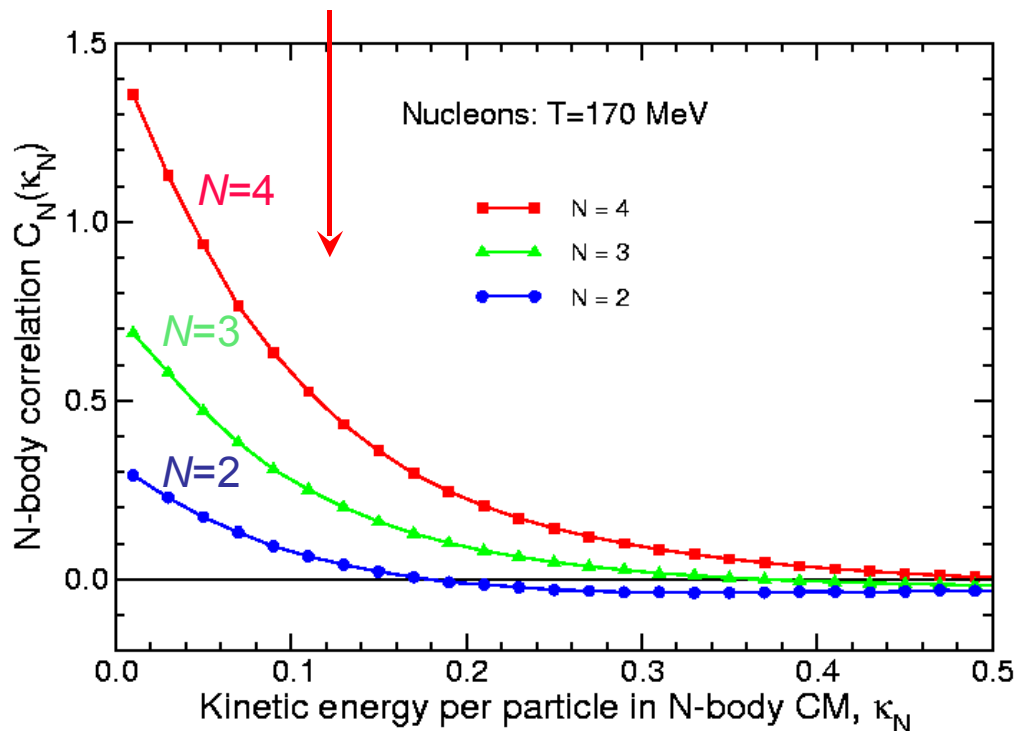
Distribution of κ :

$$P_N(\kappa) \equiv \langle \delta(\kappa - \kappa_N\{\mathbf{p}_n\}) \rangle$$

Correlation function:

$$C_N(\kappa) \equiv P_N(\kappa) / P_N^0(\kappa) - 1$$

- is enhanced at $\kappa \approx T$



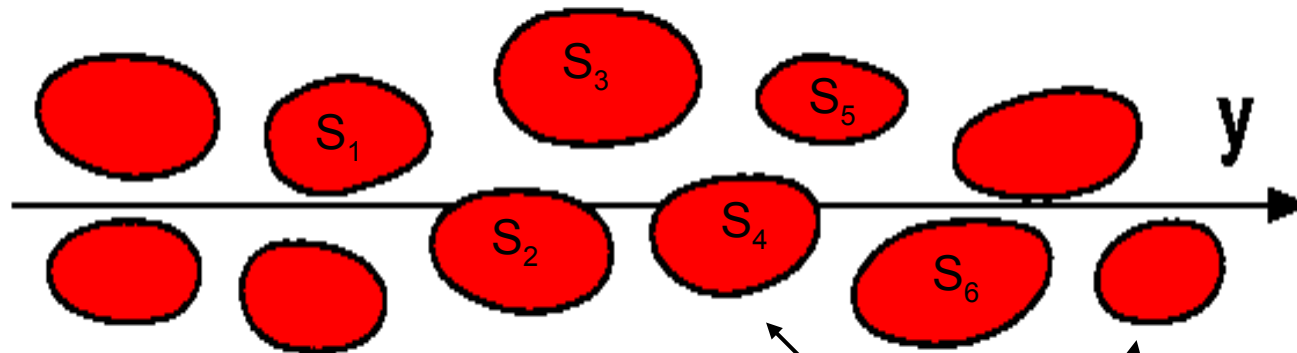
Same event / Mixed events

Higher-order correlations stand out more clearly!

(but require larger samples)

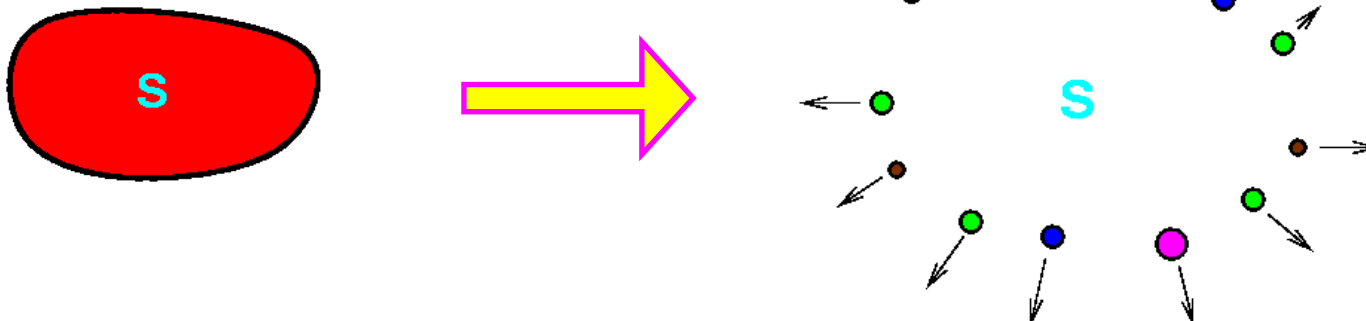
Strangeness correlations

The expanding system decomposes into plasma blobs which each contain a certain amount of strangeness:

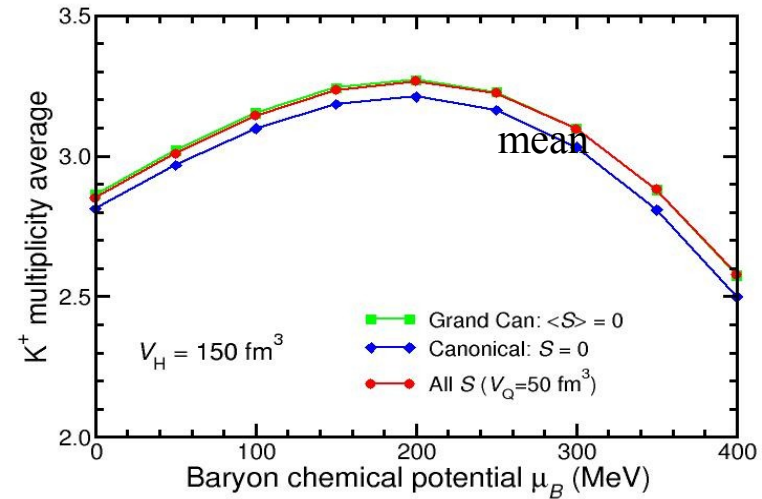
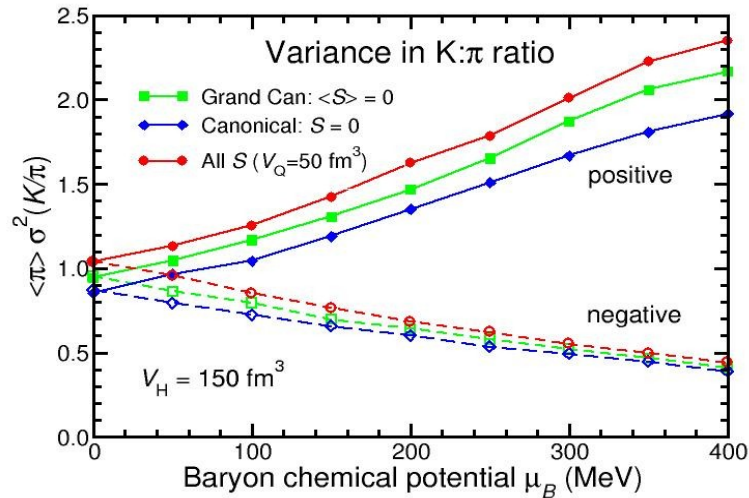


The hadronization of each isolated blob conserves strangeness:

$$S_n \neq 0$$



Some numbers



Variance: enhanced by $\sim 10\%$

$$V_{\text{QGP}} = 50 \text{ fm}^3$$

$$V_{\text{hadron}} = 150 \text{ fm}^3$$

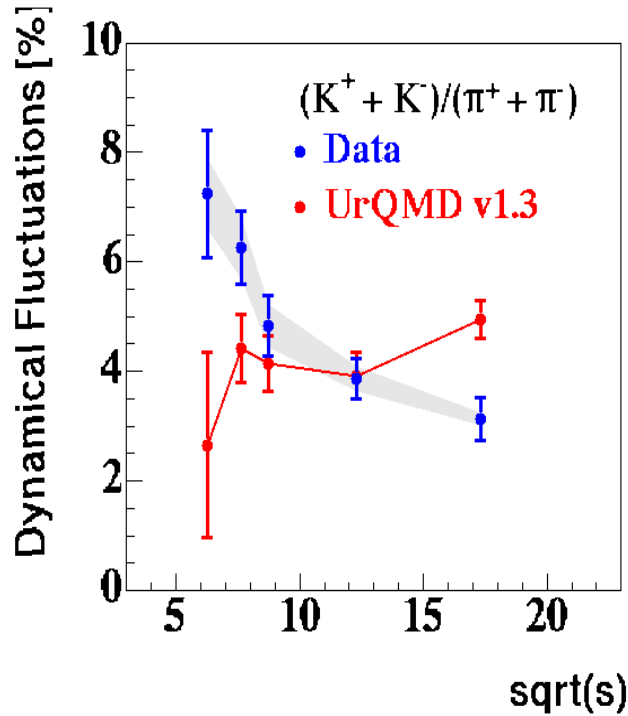
$$T = 170 \text{ MeV}$$

Generally: variance is more enhanced than mean

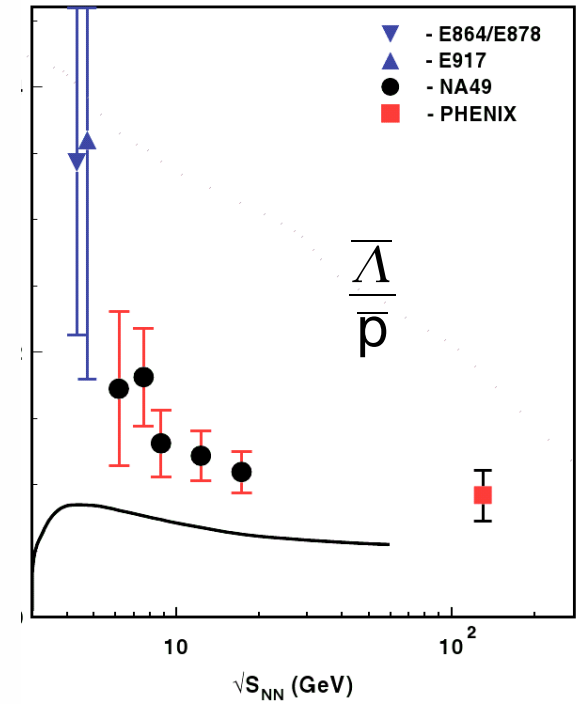
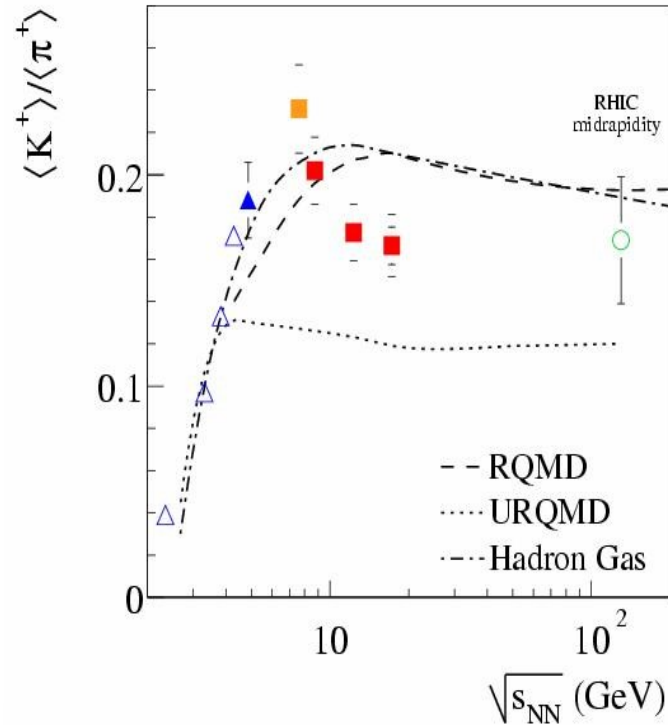
Strange things...

NA49, PRC73, 044910 (2006)

C. Roland, QM04, NA49

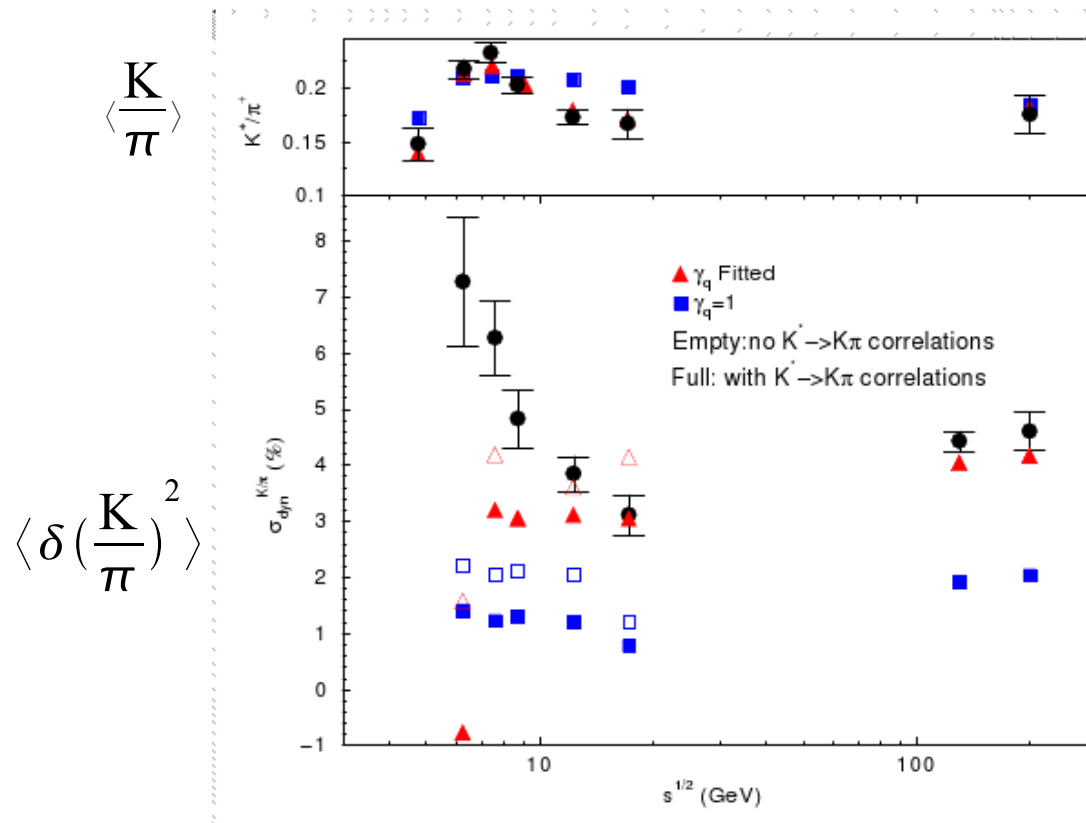


NA49

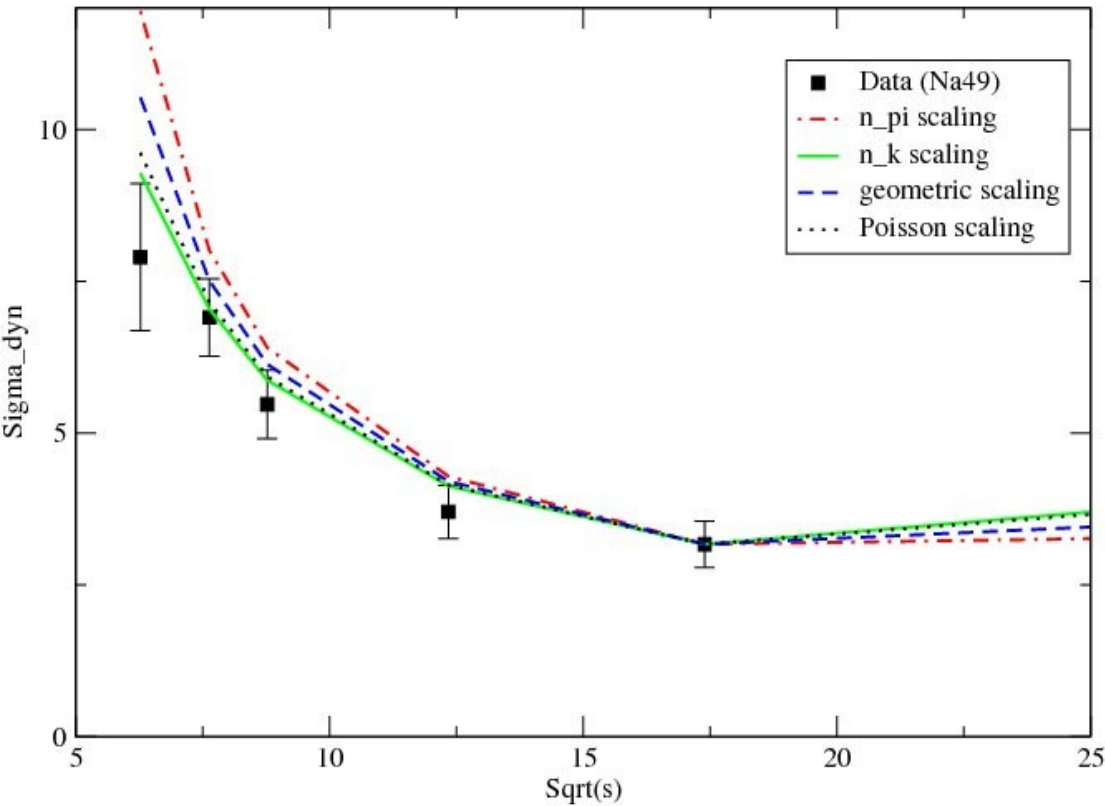


Hadron gas predictions

G. Torrieri, QM2006



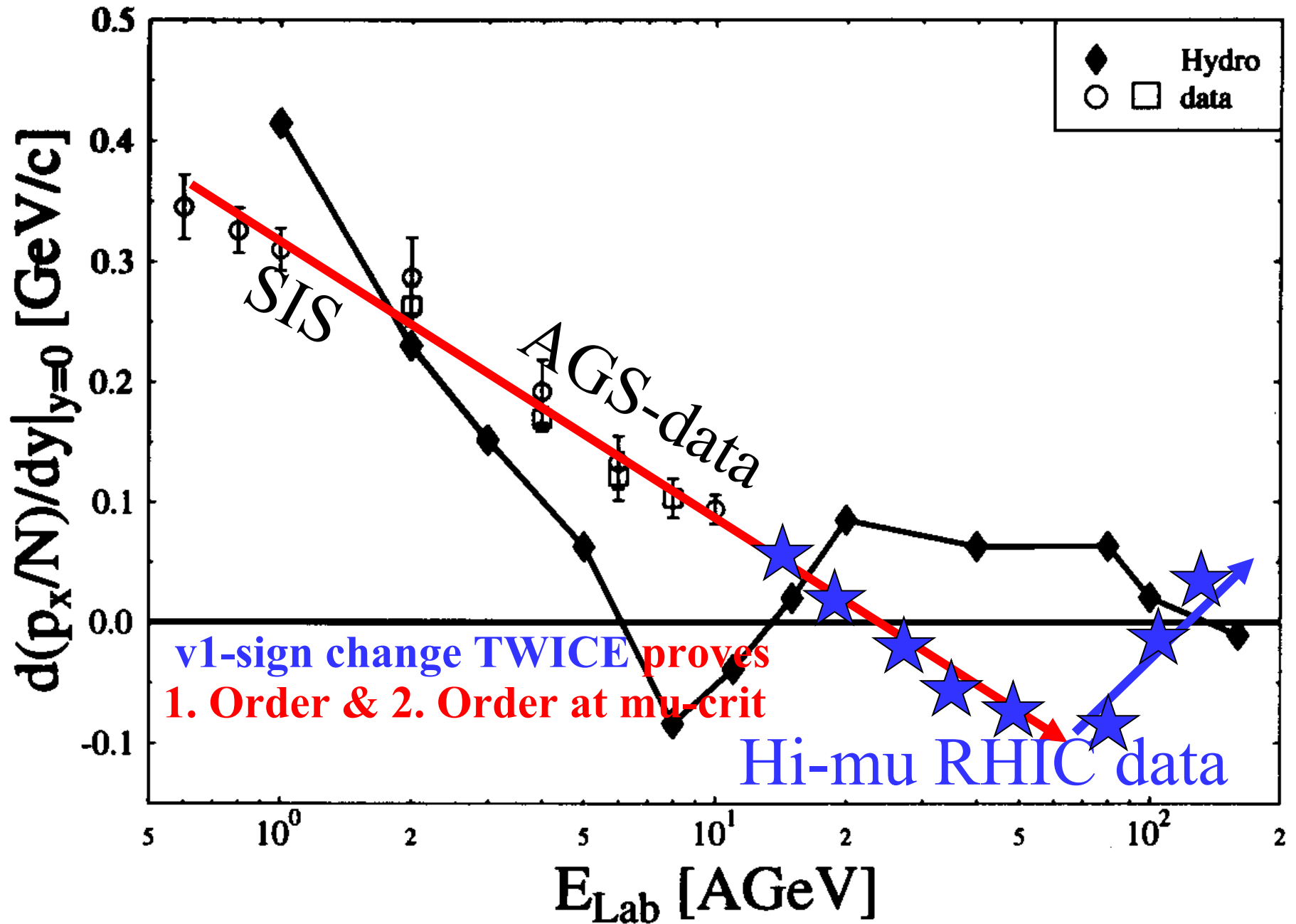
Some trivial effects...



$$\begin{aligned} \sigma_{\text{dyn}}^2 &= \frac{\langle (\delta K)^2 - K \rangle}{\langle K \rangle^2} + \frac{\langle (\delta \pi)^2 - \pi \rangle}{\langle \pi \rangle^2} - 2 \frac{\langle \delta K \delta \pi \rangle}{\langle K \rangle \langle \pi \rangle} \\ &= \frac{(\omega_K - 1)}{\langle K \rangle} + \frac{(\omega_\pi - 1)}{\langle \pi \rangle} - 2 \frac{(\omega_{K\pi} - 1)}{\sqrt{\langle K \rangle \langle \pi \rangle}} \\ &\sim 1 / (\text{accepted Multiplicity}) \end{aligned}$$

Other (“indirect”) observables

- Flow measurements (EOS, viscosity?)
- Lepton pairs? Only in conjunction with something else, such as baryon number fluctuations
 - Correlate baryon number with lepton yield in order to get after density fluctuations



v1-sign change TWICE proves
1. Order & 2. Order at mu-crit

Hi-mu RHIC data

Critical Point and viscosities

CP is in universality class of liquid gas (Son, Stephanov)

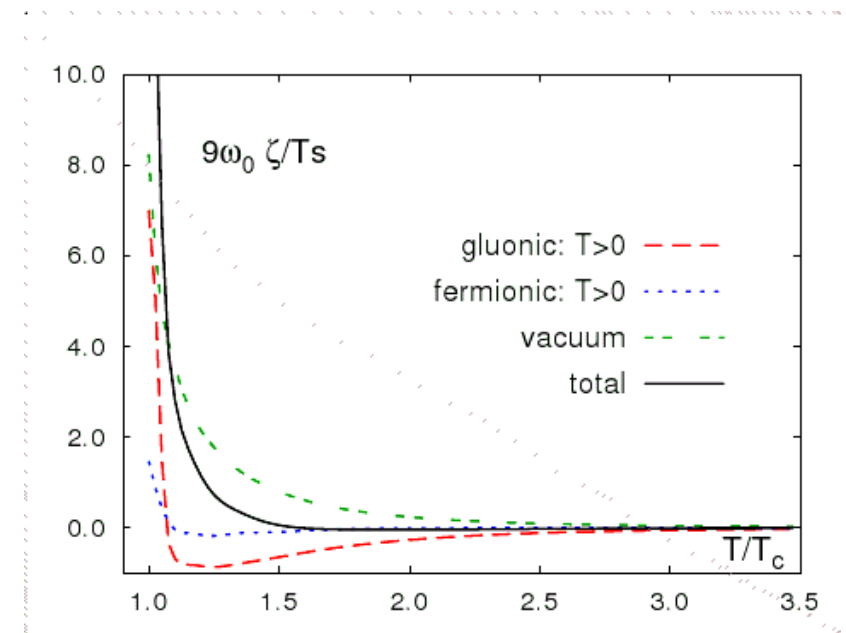
Hohenberg - Halperin Model H (Rev. Mod. Phys 49 (1977)):

$$\eta \sim \xi^{0.065}, \quad \xi = \text{Correlation Length}$$

Shear viscosity **diverges** at CP

Bulk viscosity also **diverges**:
(Kharzeev, Turchin, Karsch arXiv:0711.0914)

Note: even large increase without PT
due to vacuum contribution



QCD critical point

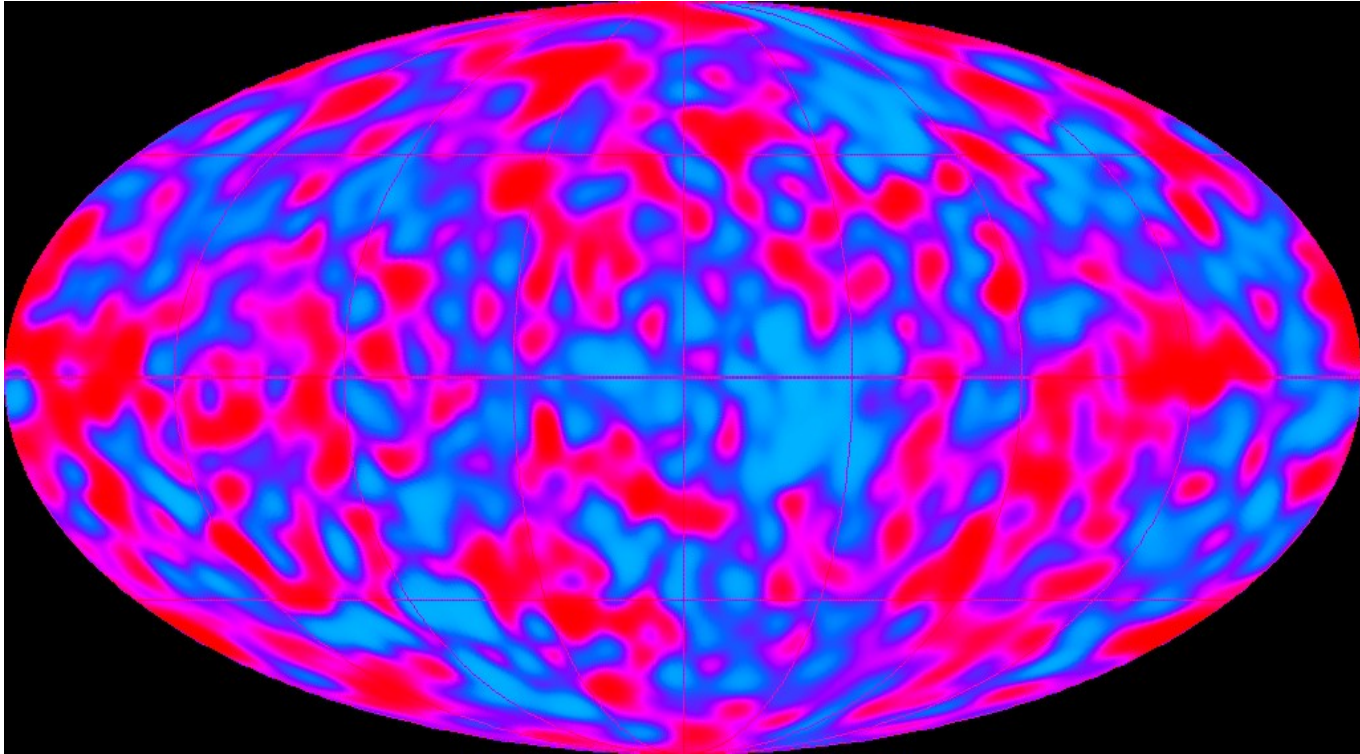
- Order parameter: baryon density or scalar density
 - Actually it is a superposition
- Both scalar (chiral) and quark number susceptibilities diverge
- **Screening** (“space like”) masses vanish (“omega”, “sigma”)
 - not accessible by (time-like) dileptons
- Is it related to chiral transition at $m_q = 0$?
- The transition is in same universality class as liquid gas! (Son, Stephanov)
 - Fluctuations are driven by density fluctuations; chiral field is just tagging
- CP “just” the end of of 1st Order transition
 - Spinodal instabilities

Observables for CP and co-existence

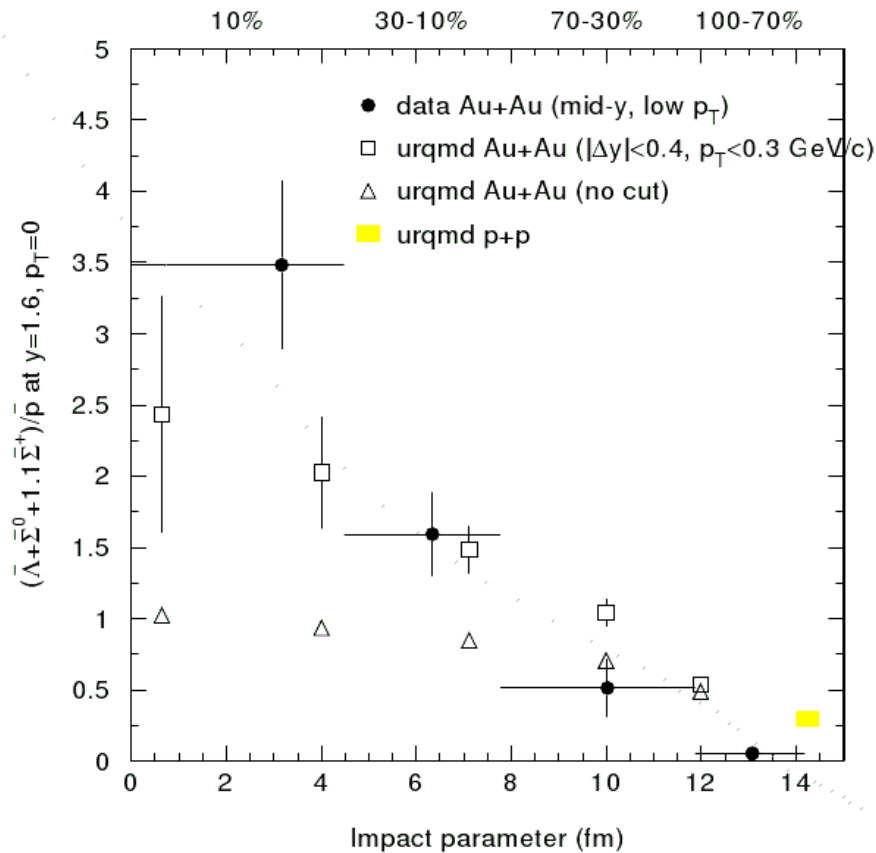
- Fluctuations (probably not of conserved charges)
- Correlations (spiondal blobs)
- Energy scan
- System size dependence (finite volume scaling)
 - centrality may not do
- Be prepared to measure everything
 - not clear (yet?) which observable couples strongest to baryon density
 - Would like to see finite volume scaling in more than one observable
- So far NOTHING seen

Summary

- Sign of phase co-existence CAN be seen in these type of experiments (Liquid Gas)
- Situation for QCD PT rather unsatisfactory
 - No firm theoretical guidance (Not even qualitative!)
 - Not clear how the phases present themselves (What are the “droplet”?)
 - So far no evidence for or against PT of whatever kind
-



UrQMD and Lambda-bar / p-bar



F. Wang
nucl-ex/0010002

Strong enhancement
mostly an effect of
acceptance cut !?

