





## Neutrino-matter interactions for core-collapse supernovae – Precision at subleading order

NAVI Physics Days 2016 – January 18– GSI Andreas Lohs (Univ. Basel)







TECHNISCHE UNIVERSITÄT DARMSTADT Core collapse supernovae release huge amount of energy.



Neutrino spectra and interactions with matter are maior determinants of nucleosynthesis conditions.

#### **Neutrino-Interactions: Two Regimes**

#### Interior of the neutron star:

Neutrino spectra formation

$$p + e^- \rightleftharpoons \nu_e + n$$

$$n + e^+ \rightleftharpoons \bar{\nu}_e + p$$





Spectrum determines composition

$$\nu_e + n \rightarrow p + e^{-}$$

$$\bar{\nu}_e + p \to n + e^+$$

## **Uncertainties in Neutrino Physics**

#### What is the correct Equation of state?

Which reactions are relevant?

- Not obvious for  $\overline{v_e}$  and  $v_x$
- Answer may vary for different SNe

#### How to compute neutrino interactions?

-inelasticity, relativity, medium effects, weak magnetism ...

 $\mu_n$ 

 $\mu_{p}$ 

 $\mu_{\rho}$ 

#### Nucleosynthesis in Neutrino Driven Wind





Long term simulations indicate proton rich late NDW
Ye≤0.5 possible during first seconds
No full r-process, but weak r-process possible?

# A more detailed picture requires (among other things) inclusion of subleading order effects in neutrino rates

#### Mean Free Path for Neutrino Absorption

**Elastic Approximation** 

- Lowest order expression for nonrelativistic nucleons
- Analytic formula for  $\lambda(E_{
  u})$
- Can be corrected to include recoil, weak magnetism, ...

Nucleons as quasi-free fermions – Hartree response

- Relativistic kinematics, "full" matrix element, no correlations
- Mostly 2-D numerical integrals to obtain  $\lambda(E_{
  u})$

Structure function from RPA / Linear response theory

- Fully consistent with RMF-EOS, correlations (can be) included
- Requires 3-D numerical integrals to obtain  $\lambda(E_{
  u})$

#### **Elastic Approximation for Neutrino Absorption**

Mean-free path for (quasi-) free particles:

$$\lambda(E_{\nu})^{-1} \sim \int d^3 p_e \left[1 - f_e(E_e)\right] \int d^3 p_n \int d^3 p_p \frac{\left< |M|^2 \right>}{16E_{\nu}E_n E_e E_p} f_n(E_n) \left[1 - f_p(E_p)\right] \delta^4$$

Assume non-relativistic nucleons and elastic collision:

$$E_{n,p} \simeq m_{n,p} \Rightarrow \frac{\left\langle |M|^2 \right\rangle}{16E_{\nu}E_nE_eE_p} \simeq G_A^2 \left(3-x\right) + G_V^2 \left(1+x\right)$$
$$E_n - E_p \simeq m_n - m_p + U_n - U_p$$

#### Mean-free path reduces to

$$\lambda(E_{\nu})^{-1} \sim \left(3G_{A}^{2} + G_{V}^{2}\right) \left(E_{\nu} + \Delta m + \Delta U\right)^{2} \left[1 - f_{e}(E_{\nu} + \Delta m + \Delta U)\right] \frac{n_{n} - n_{p}}{1 - \exp\left[\left(\eta_{p} - \eta_{n}\right)/T\right]}$$

#### **Recoil and Weak Magnetism Corrections**

[Horowitz, PRD 65 (2002) 043001] pointed out:

- "Elastic Approximation" is more simplified than necessary
- Kinematics/Recoil can be treated relativistically

$$E_n = m_n \Rightarrow E_e = \frac{E_\nu}{1 + \frac{E_\nu}{m_n} \left(1 - x\right)}$$

- To include in phase space factor and matrix element
- Gives rise to analytic correction factor for cross-section

$$R = \left\{ G_V^2 \left( 1 + 4e + \frac{16}{3}e^2 \right) + 3G_A^2 \left( 1 + \frac{4}{3}e \right)^2 \pm 4G_A \left( G_V + F_2 \right) e \left( 1 + \frac{4}{3}e \right) + \frac{8}{3}G_V F_2 e^2 + \frac{1}{3}F_2^2 e^2 \left( 5 + 2e \right) \right\} / \left[ \left( 1 + 2e \right)^3 \left( G_V^2 + 3G_A^2 \right) \right]$$

#### **Recoil and Weak Magnetism Corrections**

[Horowitz, PRD 65 (2002) 043001] pointed out:

- "Elastic Approximation" is more simplified than necessary
- Kinematics/Recoil can be treated relativistically

$$E_n = m_n \Rightarrow E_e = \frac{E_\nu}{1 + \frac{E_\nu}{m_n} \left(1 - x\right)}$$

- To include in phase space factor and matrix element
- Gives rise to analytic correction factor for cross-section

Correction factors are ratios of vacuum cross sections with exact and approximated matrix element and phase space, assuming the target nucleon at rest.

#### **Correction Factor for Cross-Section**



#### **Correction Factor for Cross-Section**



- In NDW, equilibrium shifted towards neutrino absorption
- In decoupling region, neutrino sphere moves inwards for antineutrinos, resulting in higher average energy

#### **Improvement: Consider Mass and Potential Differences**

- Masses and strong interaction potentials of nucleons differ
- At large densities effective masses decrease

$$E_e = \frac{E_{\nu} + \frac{M_*^2 - m_p^{*2}}{2M_*}}{1 + \frac{E_{\nu}}{M_*} (1 - x)} \qquad M_* = m_n^* + U_n - U_p$$

- Analytic correction factor can still be derived the same way
- In the matrix element, additional terms can be included
- For neutrino scattering, only difference is exchange of rest mass with effective mass

#### **Improved Correction Factor**

$$\begin{split} R &= \left\{ G_V^2 \left[ 1 + 4e_* + \frac{16}{3} e_*^2 + \frac{4}{3} e_* \xi + \left( 1 + \frac{2}{3} e_* \right) (\xi - q_*) \right] \right. \\ &+ G_A^2 \left[ 3 + 8e_* + \frac{16}{3} e_*^2 - \frac{4}{3} e_* \xi - \left( 1 + \frac{2}{3} e_* \right) (\xi + q_*) \right] \\ &\pm G_A \left[ G_V + F_2 \frac{M_*}{m_N} \left( 1 - \frac{\xi}{2} \right) \right] \left[ 4e_* + \frac{16}{3} e_*^2 + q_* \left( 2 + \frac{4}{3} e_* \right) \right] \\ &+ G_V F_2 \frac{M_*}{m_N} \left[ \left( 1 + \frac{q_*}{e_*} - \frac{\xi}{2} \right) \frac{8}{3} e_*^2 + \xi q_* \left( 1 + 2e_* + \frac{4}{3} e_*^2 \right) \right] \\ &+ F_2^2 \frac{M_*^2}{m_N^2} \left[ \frac{5}{3} e_*^2 + \frac{2}{3} e_*^3 + \left( \frac{1}{2} + e_* \right) \tilde{A} + \left( \frac{1}{2} + \frac{1}{3} e_* \right) \tilde{B} + \frac{2}{3} e_* \tilde{C} \right] \right\} \\ &+ \left[ \left( 1 + 2e \right)^3 \left( G_V^2 + 3G_A^2 \right) \right] \\ \xi &= \frac{\Delta m^* + \Delta U}{M_*}, \quad q = \frac{m_n^{*2} - m_p^{*2}}{2M_*^2}, \quad q_* = \frac{M_*^2 - m_p^{*2}}{2M_*^2} \end{split}$$

 $M_*$ 

ξ

#### Improved Correction Factor at Low Densities



#### Improved Correction Factor at High Densities



#### Improved Correction Factor at High Densities



#### Neutron decay at high density



#### Elastic Approximation and Corrections for Neutron Decay

Elastic approximation for neutron decay similar to absorption

$$\lambda(E_{\nu})^{-1} \sim \left(3G_{A}^{2} + G_{V}^{2}\right) \left(\Delta m + \Delta U - E_{\nu}\right)^{2} f_{e}(\Delta m + \Delta U - E_{\nu}) \frac{n_{p} - n_{n}}{1 - \exp\left[\left(\eta_{n} - \eta_{p}\right)/T\right]}$$

Kinematic relation for inverse neutron decay, assuming proton at rest

$$E_e = \frac{-E_{\bar{\nu}} + \frac{M_*^2 - M_f^2}{2M_*}}{1 + \frac{E_{\bar{\nu}}}{M_*} \left(1 - x\right)}$$

Physical meaning only for different nucleon masses and/or potentials

Elastic Approximation and Corrections for Neutron Decay

Result: Correction factor for neutron decay is EXACTLY the same as for antineutrino absorption

- Neutron decay and antineutrino absorption refer to different neutrino energies
- Absorption:  $m_e + \Delta m_* + \Delta U < E_{\bar{\nu}} < \infty$
- Decay:  $0 < E_{\bar{\nu}} < -m_e + \Delta m_* + \Delta U$
- For decay, corrections increase or decrease rate, depending on antineutrino energy

#### Improved Correction Factor at High Densities



## Limit of Approximations

- Elastic opacities with weak magnetism corrections should be very good for zero temperature vacuum
- How good are the approximations for finite T and  $\rho$ ?
  - at conditions of NDW?
  - at conditions of neutrinosphere?
- Compare approximations with "exact" opacity
- How to compute "exact" opacity?

Computing "exact" neutrino opacities in CCSN

Hartree approximation for nucleon response:

- nucleon-nucleon interaction described by RMF-potentials and effective masses
- nucleons are quasi-free particles with modified energy

$$E_{n,p} = \sqrt{\mathbf{p}^2 + m_{n,p}^{*2} + U_{n,p}}$$

relativistic kinematics, "full" matrix element, weak magnetism included

$$\lambda(E_{\nu})^{-1} \sim \int d^3 p_e \left[1 - f_e(E_e)\right] \int d^3 p_n \int d^3 p_p \frac{\left< |M|^2 \right>}{16E_{\nu}E_n E_e E_p} f_n(E_n) \left[1 - f_p(E_p)\right] \delta^4$$

• No correlations, but always better than elastic approximation













## Limit of Approximations

- Elastic opacities with weak magnetism corrections are indeed very good for zero temperature vacuum
- For temperatures of several MeV, approximation underestimates opacities (~10%)
- At neutrinosphere, additional significant deviations for neutrino energies of several 10 MeV
- Approximation "fails" at the level of weak magnetism corrections for higher densities/temperatures
  - What is the reason for the failure? (target at rest; inelasticity; relativity)
  - Is there a "cure"?

## **Alternativ Approximations and Corrections**

#### Reminder:

Correction factors are ratios of vacuum cross sections with exact and approximated matrix element and phase space, assuming the target nucleon at rest.

Idea:

- Study correction factors isolated from elastic approximation
- Choose different/lesser approximations of matrix element

#### Example 1:

- Calculate "exact" opacity, but with coupling constant  $F_2=0$
- Correct by corresponding ratio of vacuum cross sections







#### **Alternativ Approximations and Corrections**

Example 1:

- Calculate "exact" opacity, but with coupling constant  $F_2=0$
- Correct by corresponding ratio of vacuum cross sections
- Correction factors precise also at higher densities and temperatures

Example 2:

- Inelastic, nonrelativistic opacity as base
- Most simple matrix element ~  $(G_V^2 + 3G_A^2)$
- Similar approaches currently applied in some simulations









## Limit of Approximations for Neutrino Opacities

- For densities up to NDW-conditions and temperatures below several MeV, exact neutrino opacities can be reproduced by elastic approximation + correction factors.
- For higher temperatures or for neutrinosphere densities, the approximation "fails" at the level of the correction.
- For inelastic opacities, "good" corrections can be found also at higher densities and temperatures.
- Calculation of inelastic opacities equally demanding as "exact" opacities
  - For precision at 10% level, "exact" opacity generally favourable over elastic approximation
  - When interested in correlations, inelastic but approximated opacity + corrections can be suitable

Correction Factors for neutrino-nucleon interactions:

- Can be extended to include strong interaction potentials and effective masses, and to describe neutron decay
- With elastic appoximation, validity of corrections restricted to to "lower" densities and temperatures
- For inelastic opacities, precise corrections exist also at neutrinospheres (decoupling region)

 "Exact" opacity preferable when interested in precise neutrino spectra

 Outlook: Include momentum dependence of coupling constants into "exact" opacity