Thermodynamics of Dense Baryonic Matter in EFTs

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Outline:

- 1. Introduction: baryons near chiral symmetry restoration
- 2. Thermodynamics of hadronic matter in a parity doublet model
- 3. "Confinement" in models, anomaly matching

Baryons near chiral symmetry restoration?

- standard (naive): $D\chi SB$ generates masses $m_N \stackrel{\sigma \to 0}{\to} 0$
- parity doublet (mirror): $D\chi SB$ generates mass difference $m_{N_+} \stackrel{\sigma \to 0}{\to} m_{N_-} = m_0 \neq 0$ [Detar-Kunihiro (89)]
- emergence of a scale in QCD: trace anomaly $\Theta^{\mu}_{\mu} = \frac{\beta}{2g}G^2 + m(1+\gamma)\bar{q}q$ $\Rightarrow M_B \propto \langle B|G^2|B \rangle \& \langle G^2 \rangle^{\text{lattice}}_{T_{\chi}} \neq 0$
- naive vs. mirror: not yet discriminated
 - axial couplings: $g_A^{++} = g_A^{--}$ (naive) $g_A^{++} = -g_A^{--}$ (mirror) cf. other chiral invariant operators allowed [Jaffe-Pirjol-Scardicchio (06)] cf. lattice QCD: $g_A^{--} = 0.2 \pm 0.3$ [Takahashi-Kunihiro (07)] AdS/QCD: $g_A^{++} = 0.73$, $g_A^{--} = 0.38$ [Hashimoto-Sakai-Sugimoto (08)] - which state is the true chiral partner of N(940)? if N(1535) then $m_0 = 270$ MeV (from $\Gamma^{(\exp)}(N^* \to N\pi) = 70$ MeV) \Leftrightarrow cannot reproduce $\Gamma^{(\exp)}(N^* \to N\eta) \sim 80$ MeV a speculative candidate closer to N? and/or large OZI-violation?

Dense nuclear matter in chiral models

- nuclear matter: known properties
 - -binding energy: $E/A(\rho_0) m_N = -16 \text{ MeV}$
 - -saturation density: $\rho_0 = 0.16 \text{ fm}^{-3}$
 - incompressibility: $K = 9\rho_0^2 \partial^2 (E/A)/\partial \rho^2|_{\rho=\rho_0} = 200\text{-}400 \text{ MeV}$
- parity doublet model vs. other models
 - LSM: no stable ground state corr. to nuclear matter [Kerman-Miller (74)]
 - nucleonic NJL: possible if 4F vector and 8F scalar-vector int. incld. [Koch-Biro-Kunz-Mosel (87), Buballa (96), Mishustin-Satarov-Greiner (03)]

- PDM: possible if $m_0 \sim 800$ MeV [Zschiesche-Tolos-Schaffner-Bielich-Pisarski (07)]



• cold nuclear matter in parity doublet model [Zschiesche et al. (07)]

-2 nucleon fields

$$\psi_{1L} : (1/2, 0) \quad \psi_{1R} : (0, 1/2) \\ \psi_{2L} : (0, 1/2) \quad \psi_{2R} : (1/2, 0)$$

- Lagrangian

$$\mathcal{L} = \bar{\psi}_{1} i \partial \!\!\!/ \psi_{1} + \bar{\psi}_{2} i \partial \!\!/ \psi_{2} + m_{0} \left(\bar{\psi}_{2} \gamma_{5} \psi_{1} - \bar{\psi}_{1} \gamma_{5} \psi_{2} \right) + a \bar{\psi}_{1} (\sigma + i \gamma_{5} \vec{\tau} \cdot \vec{\pi}) \psi_{1} + b \bar{\psi}_{2} (\sigma - i \gamma_{5} \vec{\tau} \cdot \vec{\pi}) \psi_{2} - g_{\omega} \bar{\psi}_{1} \psi \psi_{1} - g_{\omega} \bar{\psi}_{2} \psi \psi_{2} + \mathcal{L}_{M} , \mathcal{L}_{M} = \frac{1}{2} \partial_{\mu} \sigma \partial^{\mu} \sigma + \frac{1}{2} \partial_{\mu} \pi \partial^{\mu} \pi + \frac{1}{2} \bar{\mu}^{2} (\sigma^{2} + \vec{\pi}^{2}) - \frac{\lambda}{4} (\sigma^{2} + \vec{\pi}^{2})^{2} + \epsilon \sigma - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m_{\omega}^{2} \omega_{\mu} \omega^{\mu} + (g_{4})^{4} (\omega_{\mu} \omega^{\mu})^{2}$$

-masses:
$$m_{\pm} = \frac{1}{2} \left[\sqrt{(a+b)^2 \sigma^2 + 4m_0^2} \mp (a-b) \sigma \right]$$

- thermodynamics of the model [CS-Mishustin (2010)] liquid-gas transition and chiral crossover/transition • phase diagram in PDM: $m_{N_{-}} = 1.5 \text{ GeV}$



• pion decay constant at an intermediate temperature



• phase diagram in PDM: $m_{N_{-}} = 1.5 \text{ GeV}$



• phase diagram in PDM: $m_{N_{-}} = 1.2 \text{ GeV}$



• in-medium quark condensate and the low-energy theorem present MF model



up to the leading order in ρ :

$$R(\rho) = \frac{\langle \bar{q}q \rangle_{\rho}}{\langle \bar{q}q \rangle_{\text{vac}}} = 1 - \frac{\Sigma_N}{m_\pi^2 f_\pi^2} \rho \,, \quad \Sigma_N = 45 \pm 8 \,\text{MeV}$$

beyond linear-density-approximation:

 $R^{\rm (PDM)}(\rho \sim 0.2\,{\rm fm}^{-3}) \sim 0.5\,, \quad R^{\rm (ChPT)}(\rho \sim 0.2\,{\rm fm}^{-3}) \sim 0.7$

 \Rightarrow importance of two-pion exchange correlations with $\Delta(1232)$

Exotic phase? role of tetra-quark states

• 2 phases with broken symmetry: distinguished by n_B

[Harada-CS-Takemoto (09)]



- symmetry breaking: $SU(N_f)_L \times SU(N_f)_R \rightarrow SU(N_f)_V \times Z_{N_f} \rightarrow SU(N_f)_V$

- order parameters: 2-quark state $\sigma \sim \bar{q}q$ and 4-quark state $\chi \sim (\bar{q}q)^2 + \bar{q}\bar{q}$ -qq
- -3 phases from a Ginzburg-Landau potential ($V = A\sigma^2 + B\chi^2 + \cdots$)

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- -3 phases from a Ginzburg-Landau potential I-II: χ_B max. ($\sigma \rightarrow 0$)

II-III: χ_B no much change (no Yukawa term $\bar{N}N\chi$ in phase II) - χ_B max. along Z_2 restoration line : baryons more activated

What is constraint on phases in quantum field theories?

• phase diagram from PNJL models: 3 regions



anomaly matching between UV and IR theories



- chirally restored phase: no NG boson thus no WZW term

- triangle diagrams with baryons: matched for $N_f = 2$ but not for $N_f = 3$ triangle graph \propto tr $[T^a{Q,Q}]$ [Shifman (89)] $N_f = 2$: (quark) 3Nc/9 = 1 (hadron) = 1 $N_f = 3$: (quark) 3Nc/9 = 1 (hadron) = 0

anomalies can match only if the system is deconfined.

Summary and prospects

dense nuclear matter and its modeling

- saturation properties \Rightarrow parity doublet model
- meson-baryon "transition": a trace of LG transition
- $-\,SU(N_f)_L \times SU(N_f)_R \times Z_{N_f}$ in dense matter \Rightarrow a model for 2- and 4-quark states
- enhancement of χ_B associated with Z_{N_f} symmetry restoration baryons are more activated in this *broken* phase.

anomaly matching in matter

- any gapless mode other than pions?
- origin of hadron masses
 - trace anomaly and hadron mass generation?
 - -how to discriminate two scenarios (naive vs. mirror)?