

PHENIX Perspectives for the RHIC Energy Scan

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for the **PHENIX** Collaboration

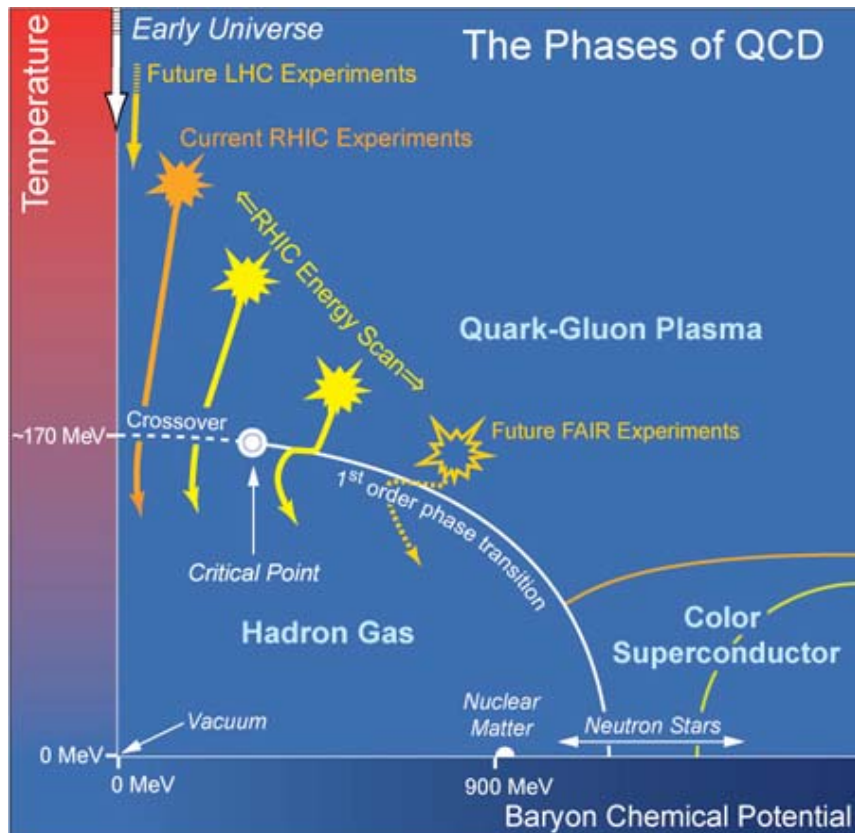
Symposium on "The Physics of Dense Baryonic Matter"

GSI, Darmstadt, March 9-10, 2009

- Introduction
- Search for the QCD Critical Point
- Search for the Onset of sQGP Production
- Summary and Outlook

QCD phase diagram

- goal of high energy heavy-ion physics
 - identify phases of matter and their properties
 - locate transitions and their properties



- vanishing μ_B
 - sQGP at top RHIC energy
 - evolution to hadron gas through a continuous rapid crossover transition
 - larger μ_B
 - possibility of a 1st order phase transition
 - critical point?
 - phase coexistence line?
- energy scan at RHIC

Questions for an energy scan

• search for the critical point

• where should we look?

→ guidance from lattice QCD

– critical point in reach at

- » FAIR
- » SPS
- » RHIC

• what (T, μ_B) for given \sqrt{s} ?

→ constraints from experiment

• what to measure?

• evolution of the sQGP

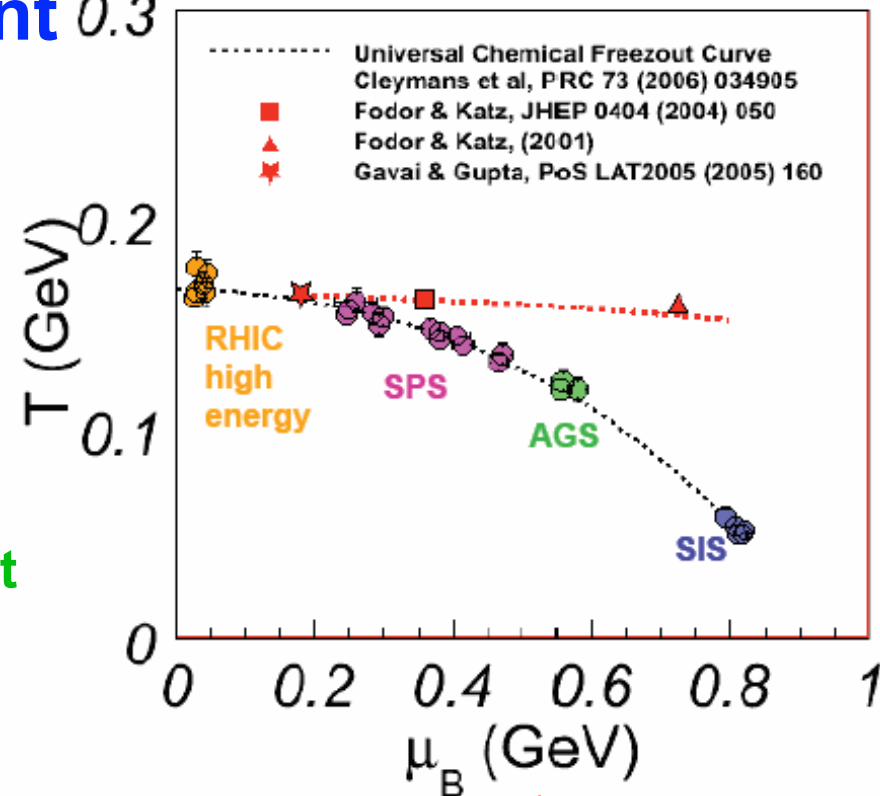
• how do the medium properties evolve with \sqrt{s} ?

• where do individual sQGP signatures "turn off"?

• experimental boundary conditions

• performance of RHIC at lower \sqrt{s} ?

• what are the constraints in PHENIX?

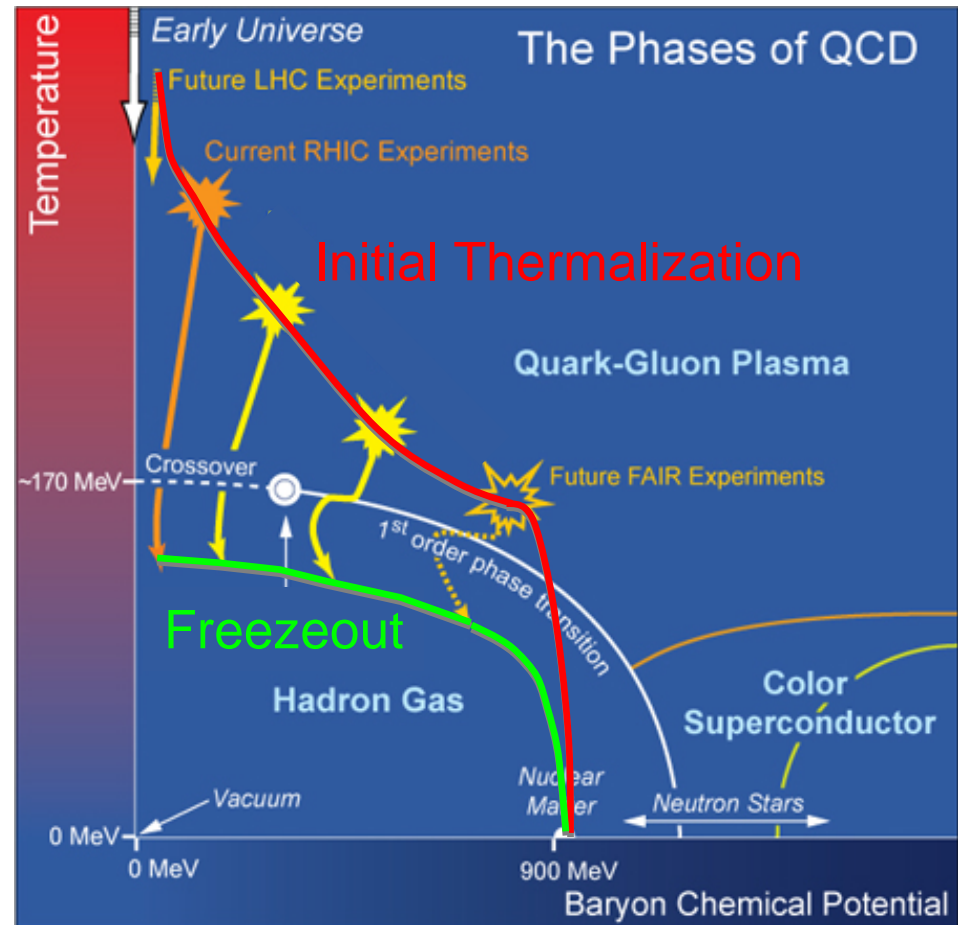
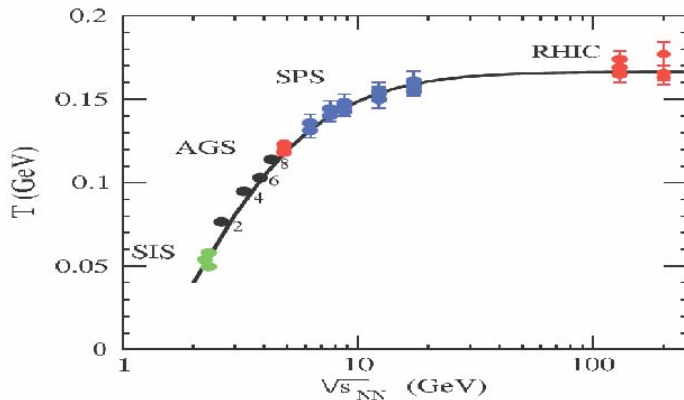
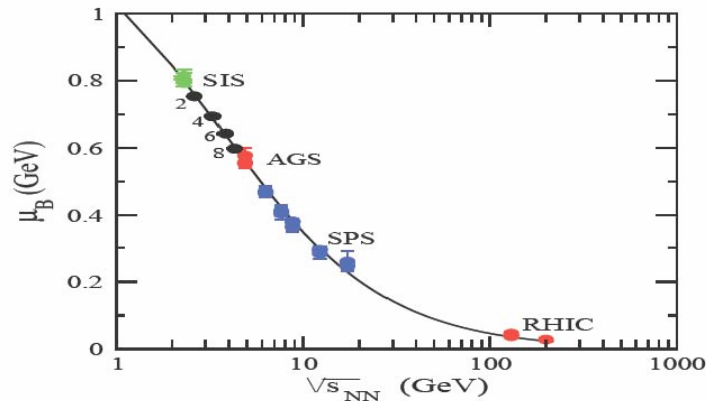


Where are we in (T, μ_B) ?

- important prerequisite

- initial thermalization in partonic world

– some idea of T_{initial} ?

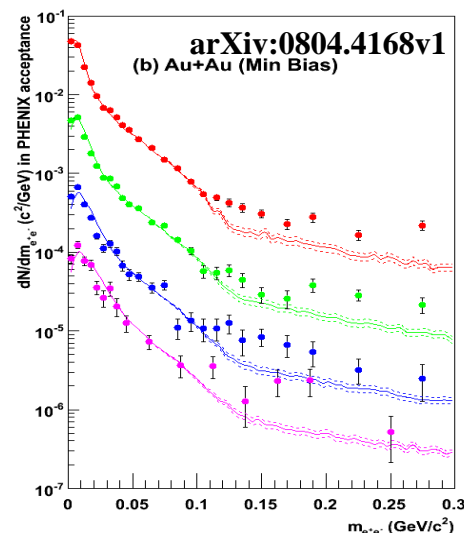
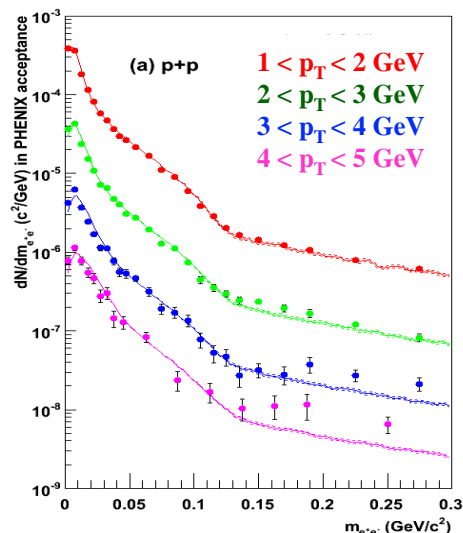
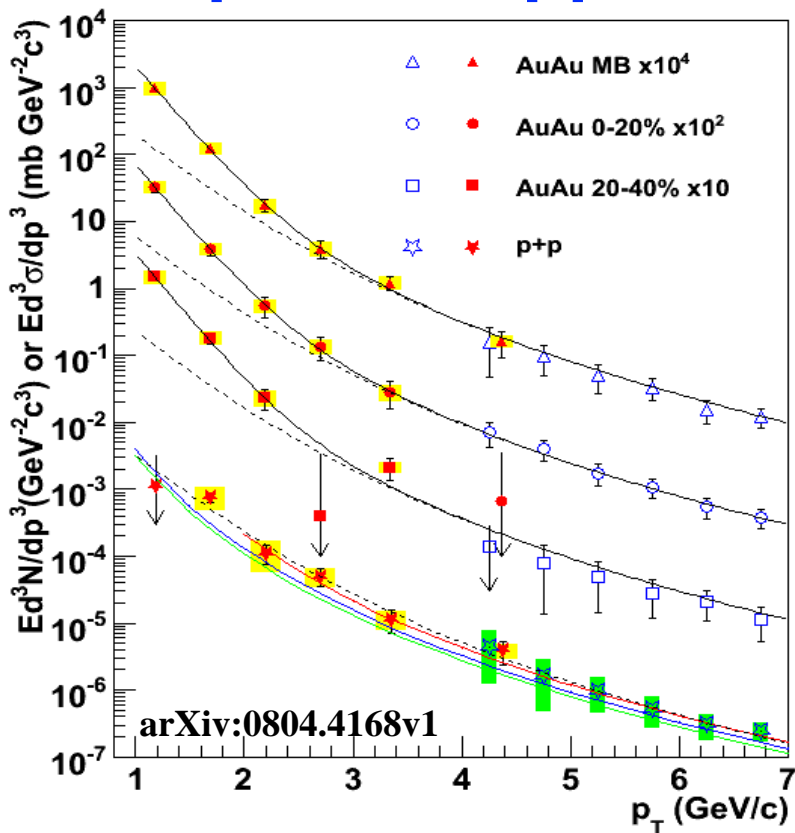


- evolution into hadronic world

– determine (T, μ_B) at freezeout from particle species ratios

Initial T from thermal photons

- enhanced emission of "soft" low-mass virtual photons in Au+Au compared to pp



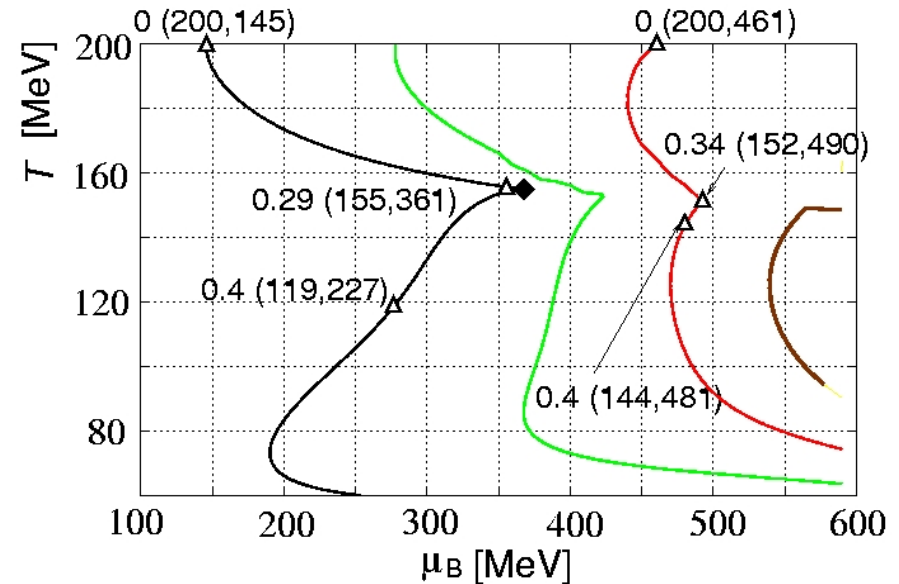
- consistent with hydrodynamic model calculation assuming $300 \text{ MeV} < T_{\text{initial}} < 600 \text{ MeV}$
- difficulties at low \sqrt{s}
 - signal/background
 - interaction rate at RHIC
- feasible at higher end of RHIC energy scan

Finding the critical point

Nonaka & Asakawa, PRC 71(2005)044904

● hydro prediction

- critical point "attracts" isentropic trajectories in the (T, μ_B) plane
- focusing causes a broadening of the signal region in (T, μ_B)



→ not necessary to exactly "hit" the critical point in an energy scan!

Stationary state variables

● properties

- divergence of stationary state variables at critical point

– compressibility

$$k_T \propto \left(\frac{T - T_C}{T_C} \right)^{-\gamma}$$

– heat capacity

$$C_V \propto \left(\frac{T - T_C}{T_C} \right)^{-\alpha}$$

- related to event-by-event fluctuations of observables

– multiplicity fluctuations $\frac{\sigma^2}{\mu^2} = k_B (T / V) k_T$

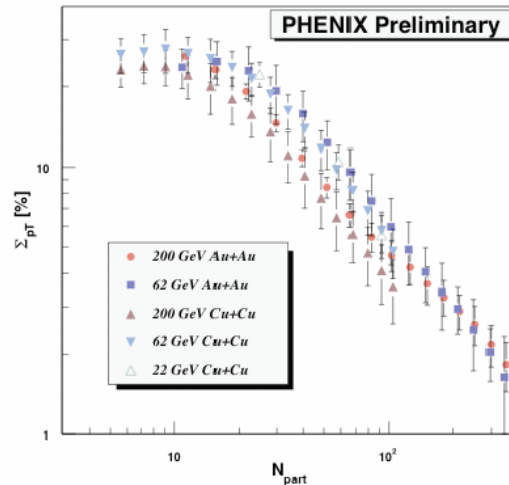
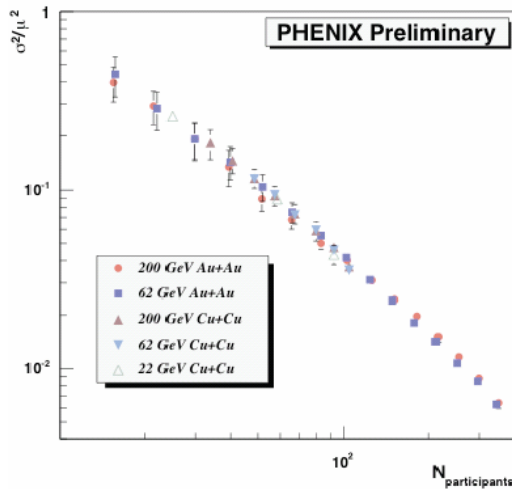
– $\langle p_T \rangle$ fluctuations $4 \sum_{pT} = \frac{1}{C_V}$

● strategy

- study fluctuations as function of μ_B (\sqrt{s})
- search for anomalies, i.e. large critical fluctuations

Fluctuations

- PHENIX measures fluctuations



- no compelling evidence for critical fluctuations yet

critical point search needs further observables

- limits and caveats

- fluctuations σ and correlation length ξ
(Stephanov, Rajagopal, Shuryak: PRD 60(1999)114028)

$$\sigma \propto \xi^2$$

- finite system size
- finite evolution time
- divergence of ξ (and σ) limited

- system slows down near critical point
→ fluctuations damped
(Berdnikov and Rajagopal: PRD 61(2000)105017)
- do critical fluctuations survive hadronization?

Antiproton-to-proton ratio

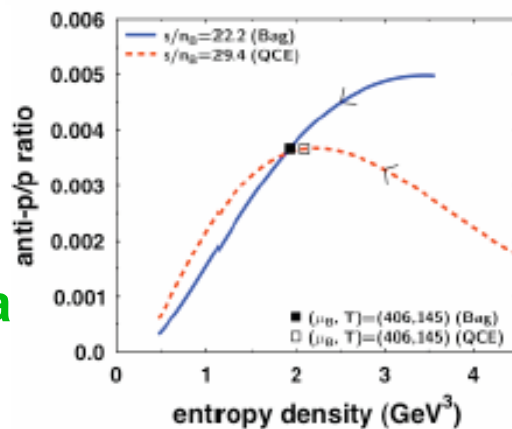
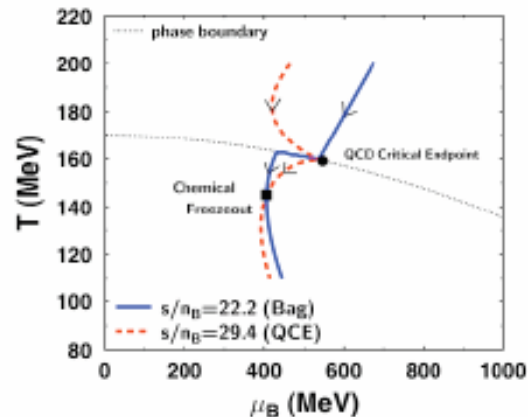
- back to hydro

- critical point deforms ("attracts") isentropic trajectories in the (T, μ_B) plane
- antiproton-to-proton ratio

$$\frac{\bar{p}}{p} \sim \exp(-2\mu_B/T)$$

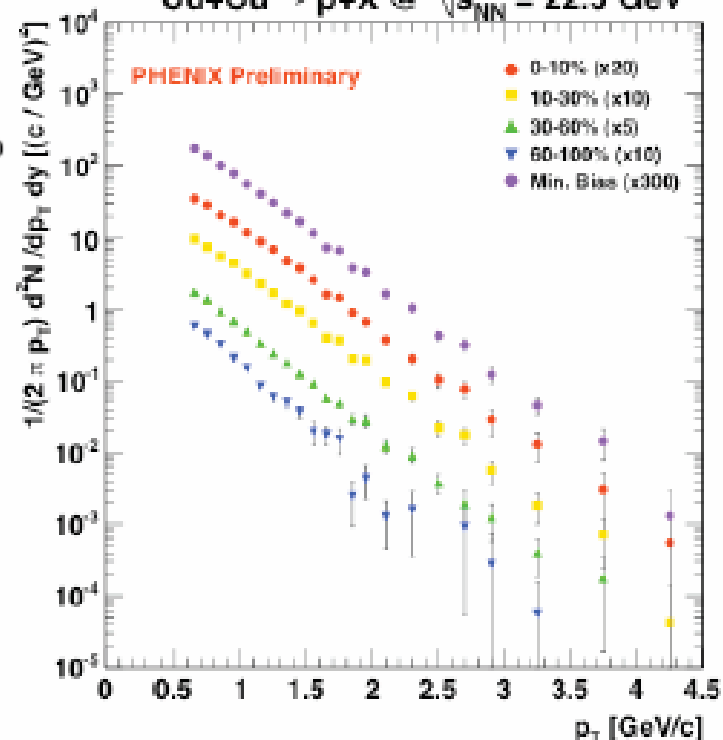
- prediction (Asakawa et al., arXiv:0803.2449)

- antiproton spectra are steeper than proton spectra at high p_T
- more robust than fluctuation observables



- PHENIX measures identified hadron spectra

Cu+Cu \rightarrow p+X @ $\sqrt{s_{NN}} = 22.5$ GeV



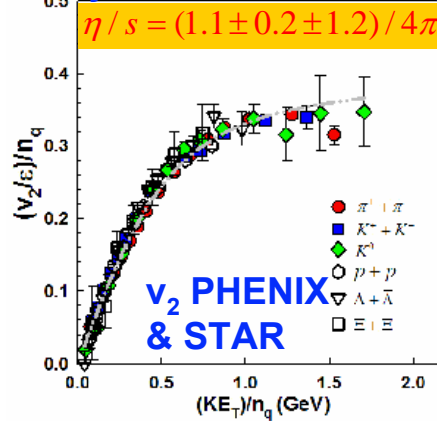
Dynamic variables

- again: correlation length ξ is important
- relation between diffusion constant D and ξ (Son & Stephanov)
$$D \sim \xi^{-1}$$
 - large ξ near critical point
→ small diffusion constant D
 - small shear viscosity to entropy density ratio η/s
- bulk viscosity is different $\bar{\eta} \sim \xi^{0.05-0.06}$
- again
 - limited system size
 - no extreme effects
- expectation close to the critical point
 - minimum in shear viscosity to entropy ratio η/s
 - bulk viscosity only somewhat sensitive

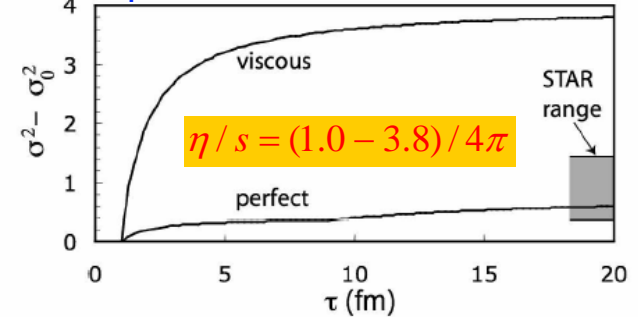
η/s measurements

- need observables that are sensitive to shear stress
- damping $\sim \eta/s$
- flow
- fluctuations
- heavy quark motion

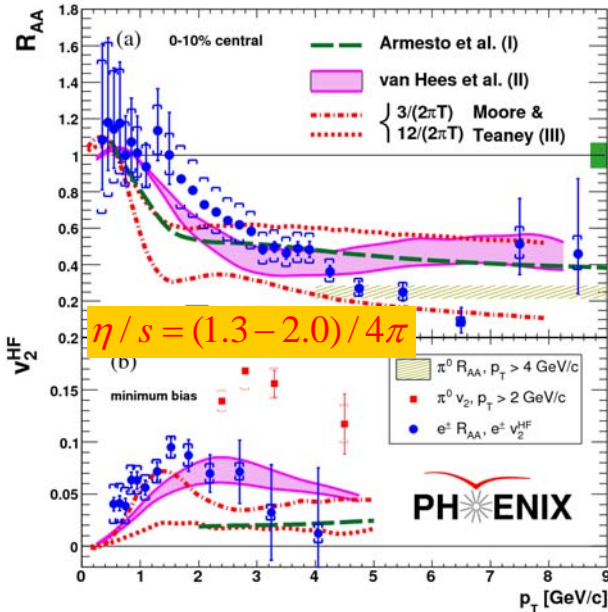
R. Lacey et al.: PRL 98:092301, 2007



S. Gavin and M. Abdel-Aziz:
PRL 97:162302, 2006
 p_T fluctuations STAR



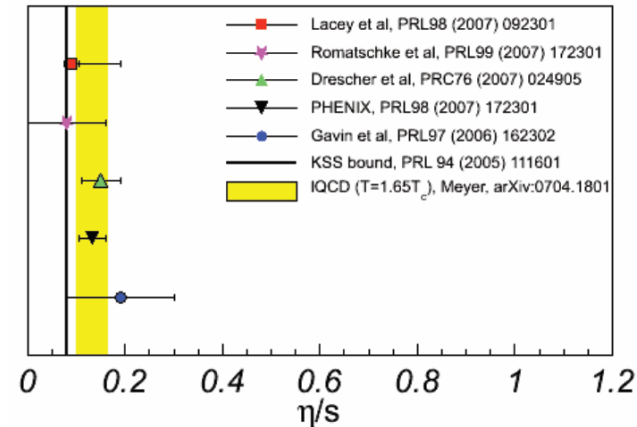
A. Adare et al.: PRL 98:092301, 2007



● top RHIC energy

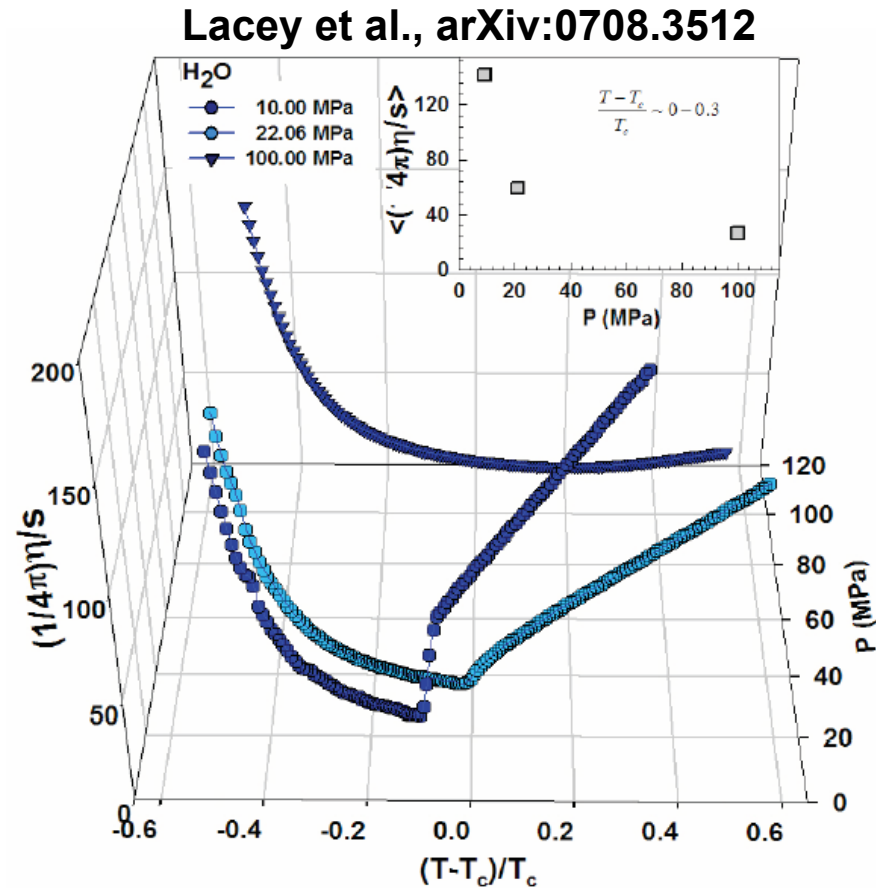
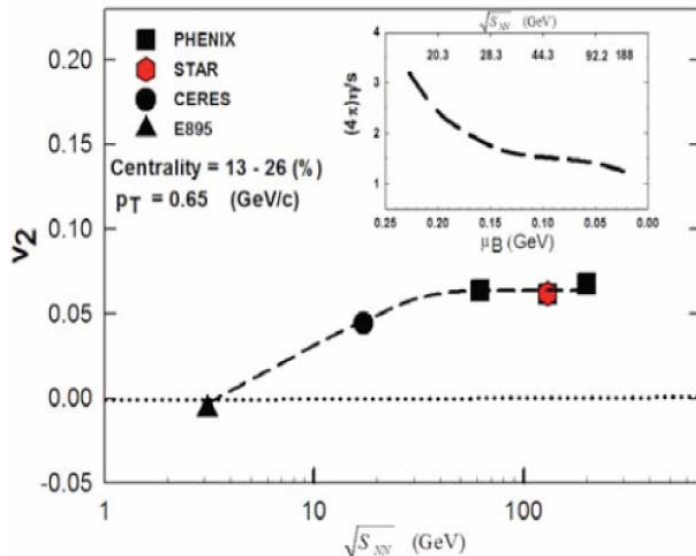
- η/s close to conjectured minimum $1/4\pi$

Shear viscosity to entropy density ratio (η/s) at RHIC



η/s near the critical point

- η/s goes through a minimum near the critical point
- estimate from Lacey et al. (based on v_2 systematics)
 - $T \sim 165\text{-}170$ MeV
 - $\mu_B \sim 120\text{-}150$ MeV



critical point search in the region $20 \text{ GeV} \leq \sqrt{s} \leq 62 \text{ GeV}$

Flow systematics

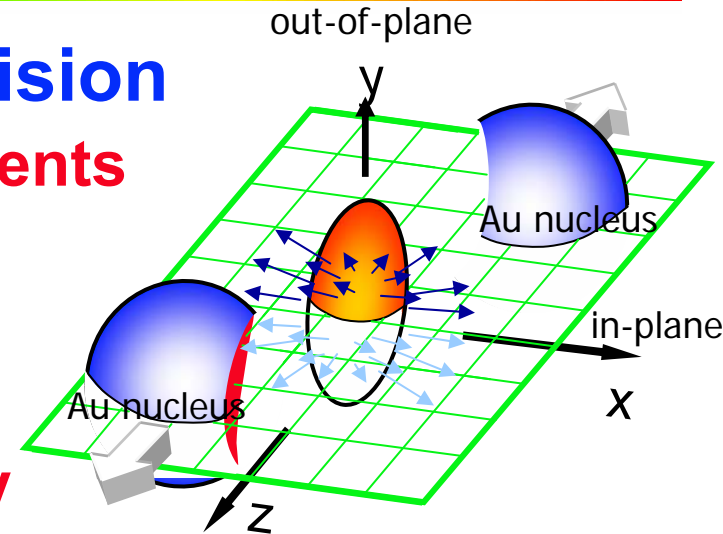
- initial state of non-central collision

- large asymmetric pressure gradients
→ hydrodynamic flow of partons

- control parameters: ε_0 , η , c_s

- translates into

- final state momentum anisotropy



$$E \frac{d^3 N}{d^3 p} = \frac{d^3 N}{p_T d\varphi dp_T dy} \sum_{n=0}^{\infty} 2v_n \cos(n(\varphi - \Psi_R)) \quad v_{2n} = \langle \cos(2n[\varphi - \Psi_R]) \rangle$$

- hydrodynamic flow exhibits scaling properties which can be validated (or invalidated), e.g.:

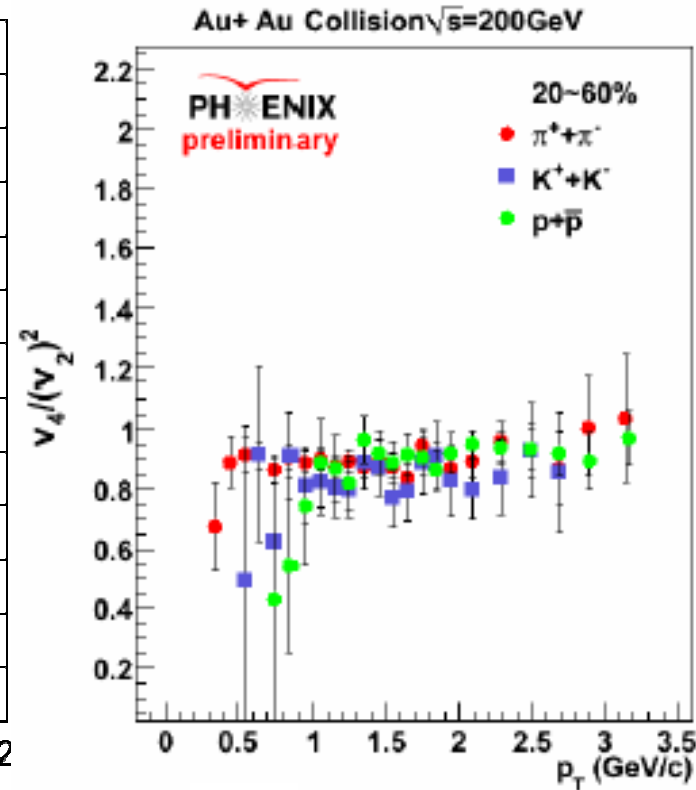
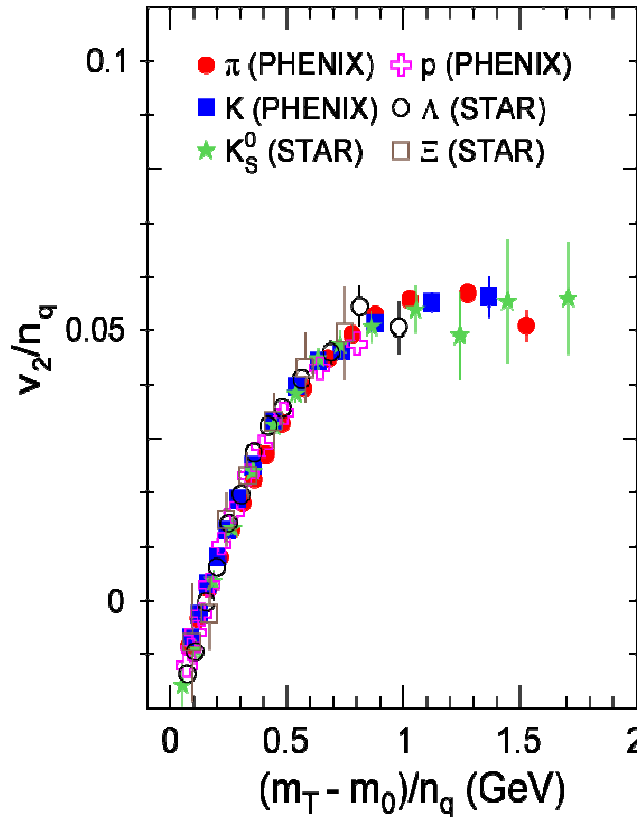
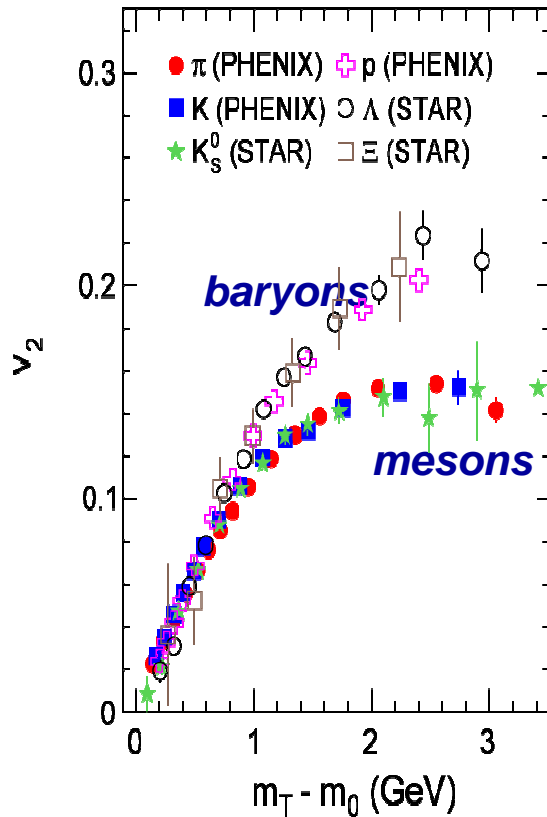
$$\frac{v_{4,M}(2p_T)}{v_{2,M}^2(2p_T)} \approx a \left(\frac{1}{4} + \frac{1}{2} \times \frac{v_{4,q}(p_T)}{v_{2,q}^2(p_T)} \right) \quad \frac{v_{4,q}(p_T)}{v_{2,q}^2(p_T)} \approx \frac{1}{2}$$

$$\rightarrow v_4/v_2^2 \approx 0.9$$

$$\frac{v_{4,B}(3p_T)}{v_{2,B}^2(3p_T)} \approx a \left(\frac{1}{3} + \frac{1}{3} \times \frac{v_{4,q}(p_T)}{v_{2,q}^2(p_T)} \right) \quad a \approx 1.8$$

Flow at RHIC

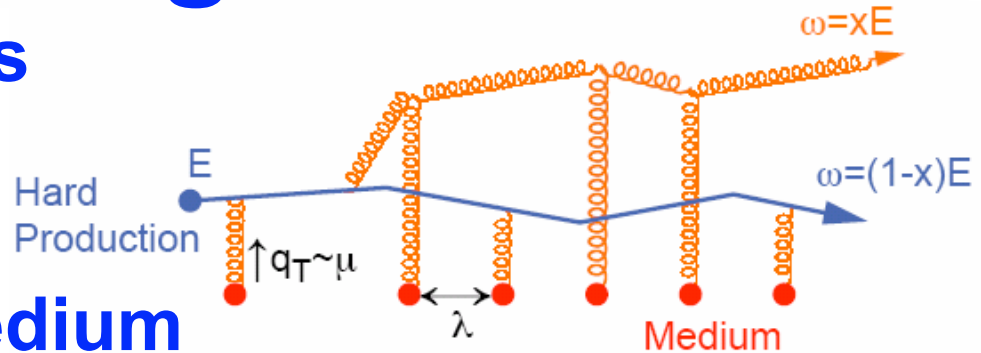
- flow shows KE_T and quark number scaling at top RHIC energy \rightarrow flow is dominantly pre-hadronic



- at what collision energy does scaling set in?

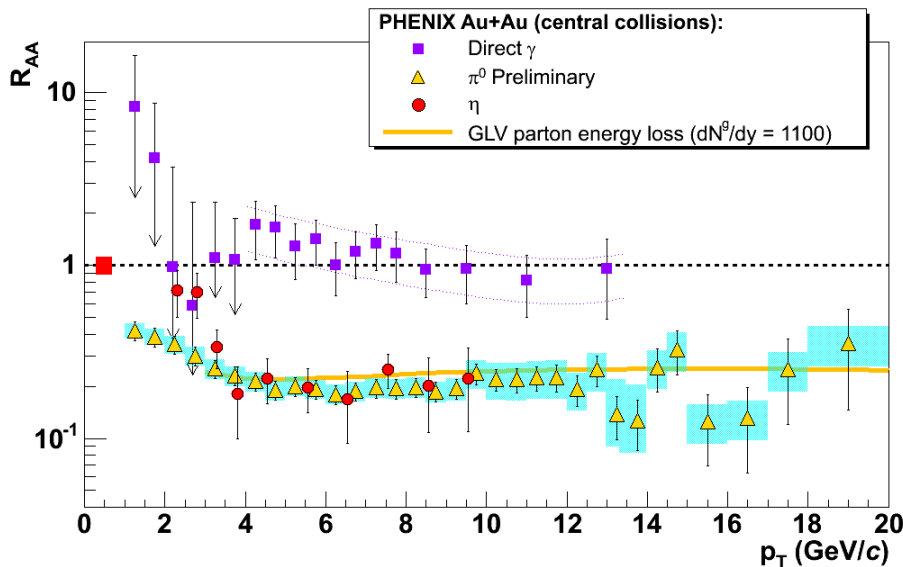
Jet quenching at RHIC

- energy loss of partons from hard scattering through re-scattering in the hot & dense medium

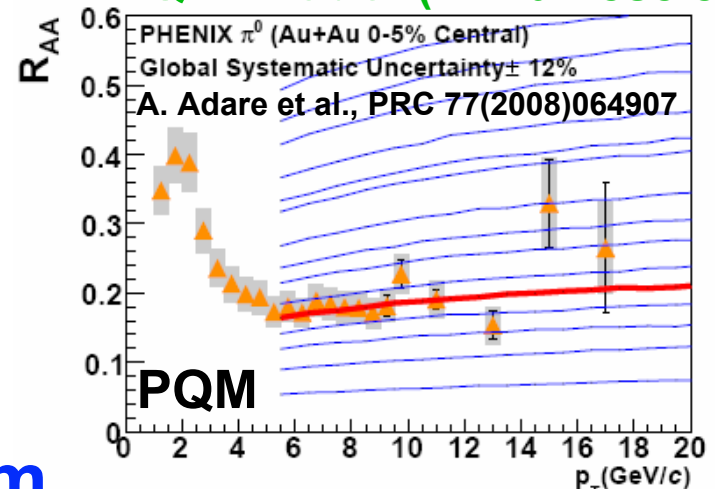


- nuclear modification factor $R_{AA} \ll 1$ at high p_T

- access medium properties through statistical analysis



- example: transport coefficient in PQM model (A. Dainese et al.)

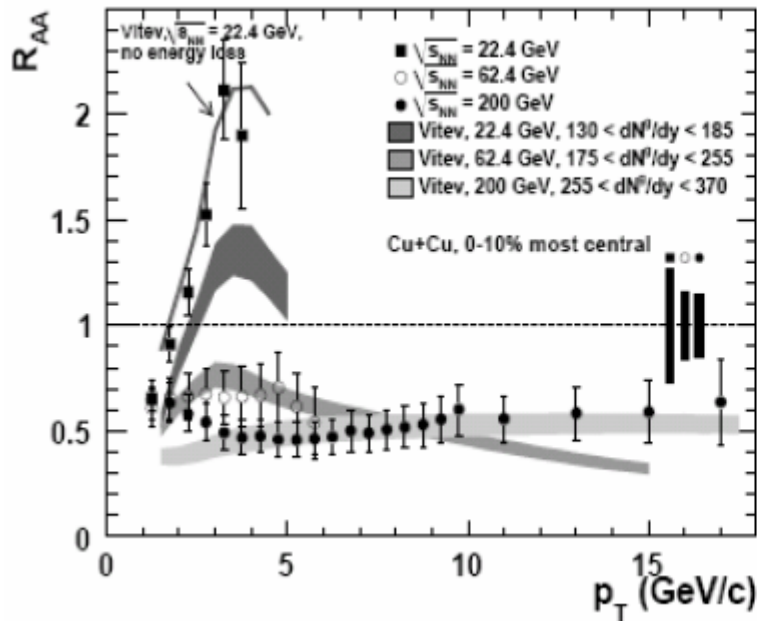


- huge opacity of the medium

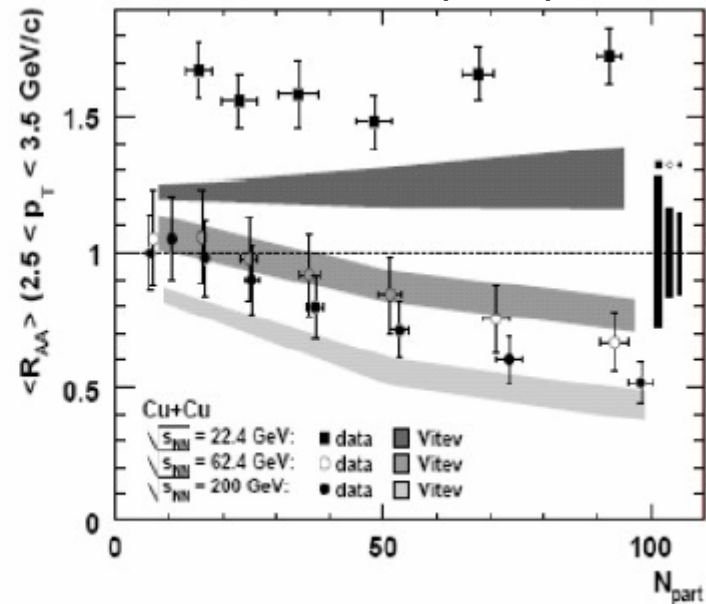
$$\rightarrow \langle \hat{q} \rangle = 13.2^{+2.1}_{-3.2} \text{ GeV}^2 / \text{fm}$$

Light quark opacity

- at what collision energy does the onset of light quark opacity occur?



A. Adare et al., PRL 101(2008)162301

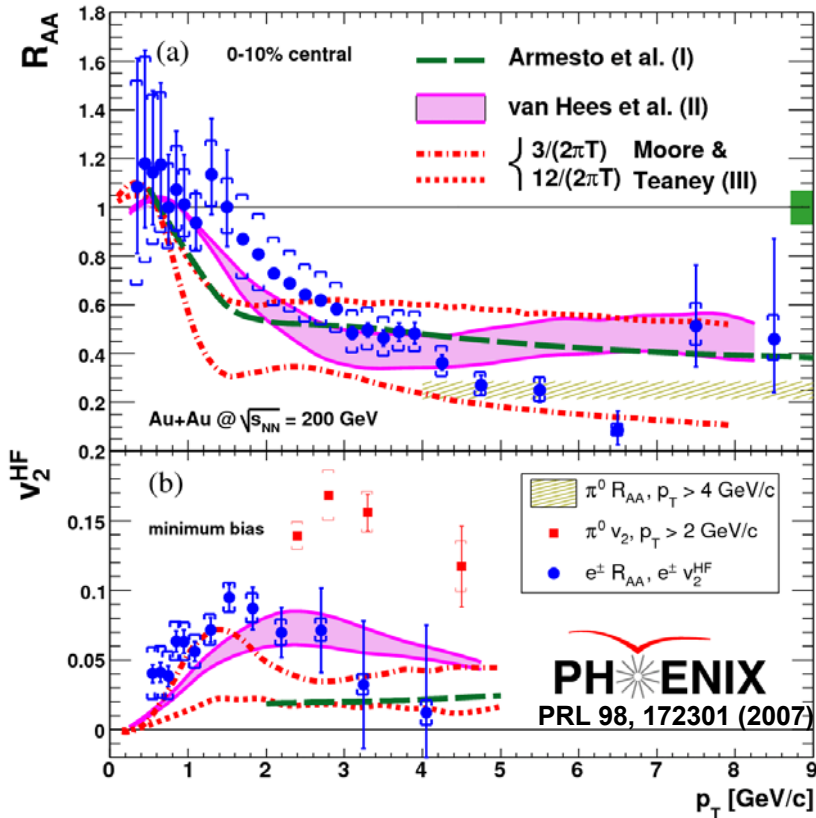


- PHENIX R_{AA} measurements in Cu+Cu collisions
 - onset for $22.4 \text{ GeV} \leq \sqrt{s_{NN}} \leq 62.4 \text{ GeV}$
- needs p+p and d+A samples in addition to A+A
- feasible only for SPS energies or higher

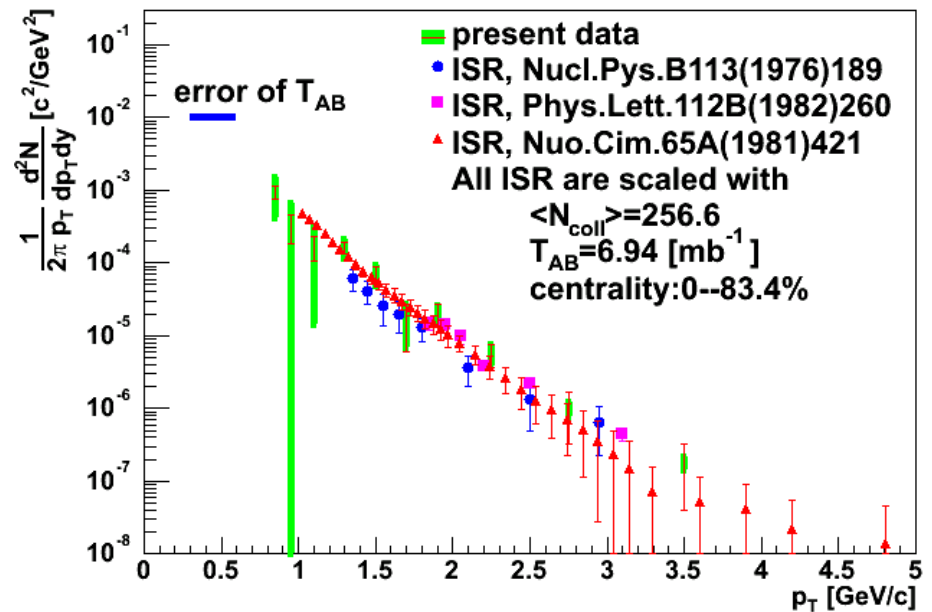
Heavy quark opacity

- where is the onset of heavy quark opacity?

R_{AA} for Au+Au @ $\sqrt{s_{NN}} = 200$ GeV



R_{AA} for Au+Au @ $\sqrt{s_{NN}} = 62.4$ GeV



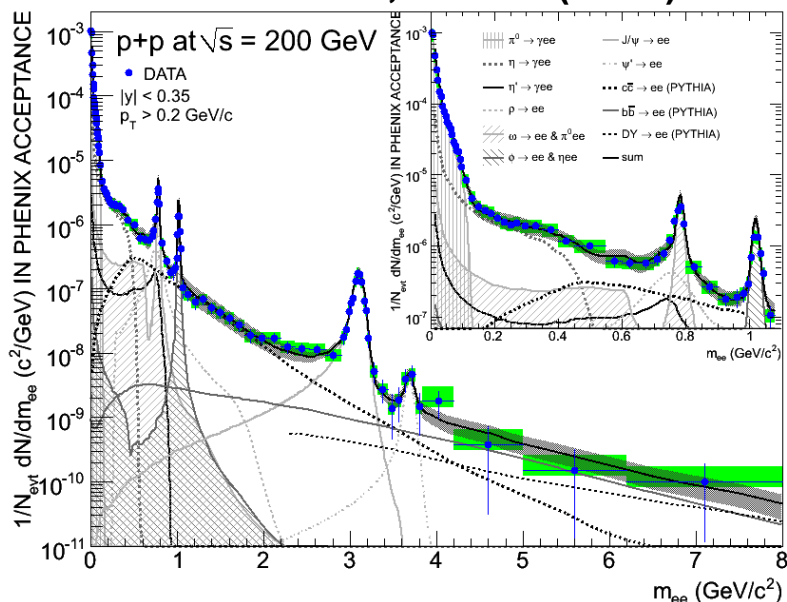
- R_{AA} consistent with unity
 - poor statistics
 - p+p reference missing

- interesting energies for heavy quark observables are above SPS energies, not below

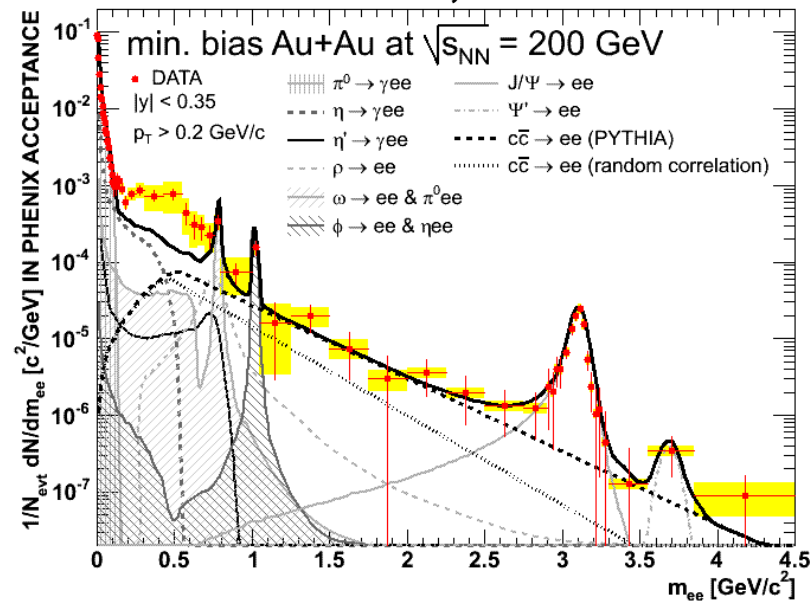
Low-mass dileptons at RHIC

- dielectrons from PHENIX in p+p and Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV

A. Adare et al., PLB 76(2009)313



S. Afanasiev et al., arXiv:0706.3034

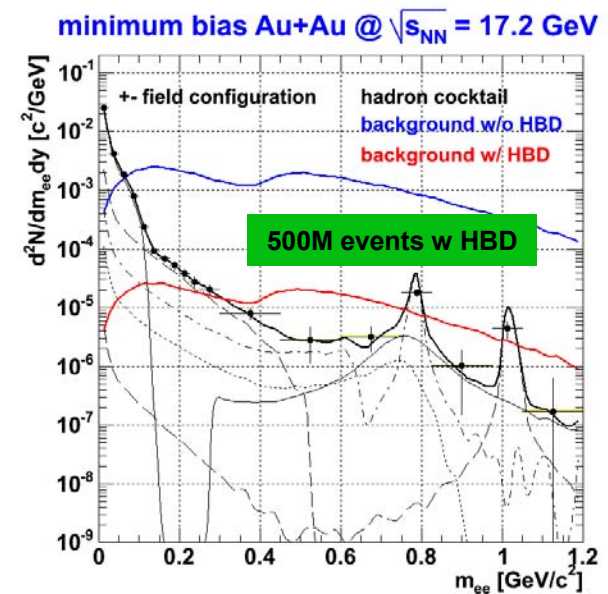
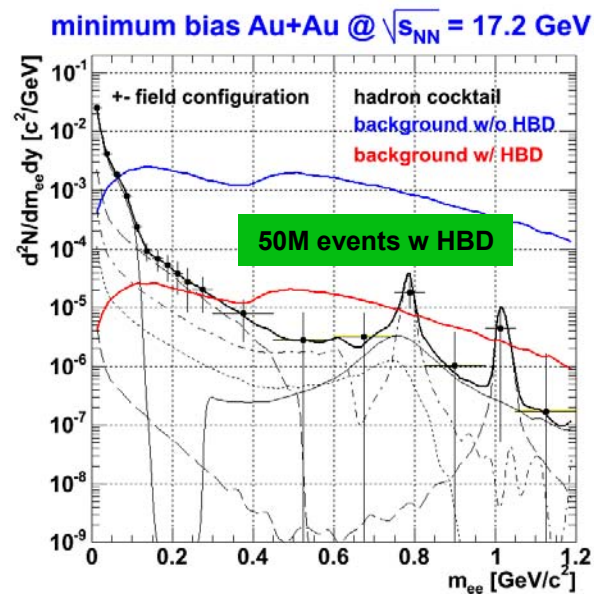
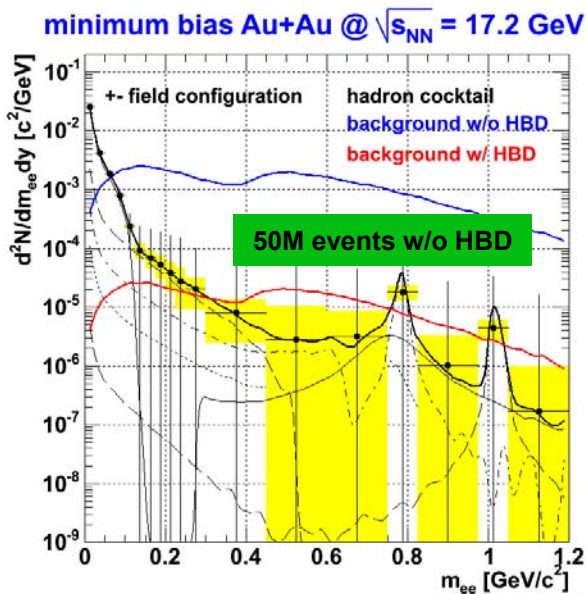


- agreement with expected e^+e^- sources in p+p
- enhancement observed in Au+Au collisions

→ can PHENIX measure e^+e^- in an energy scan?

e^+e^- at low RHIC energies

- dielectron cocktail calculation for Au+Au at $\sqrt{s} = 17.2$ GeV
 - assumptions
 - meson yields and phase space distributions as measured at SPS
 - no low-mass enhancement or any other medium effects
 - key ingredients
 - electron ID beyond PHENIX baseline is a must
 - Hadron Blind Detector (HBD)
 - increased luminosity (electron cooling) could have a huge impact



- e^+e^- measurements are possible with "CERES quality" (or better) at low RHIC energies!

RHIC boundary conditions

- life becomes difficult towards low energies

- key issues

- luminosity

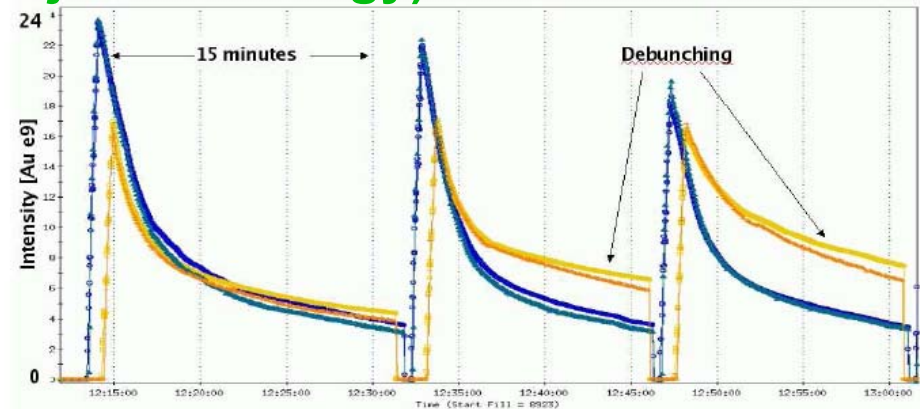
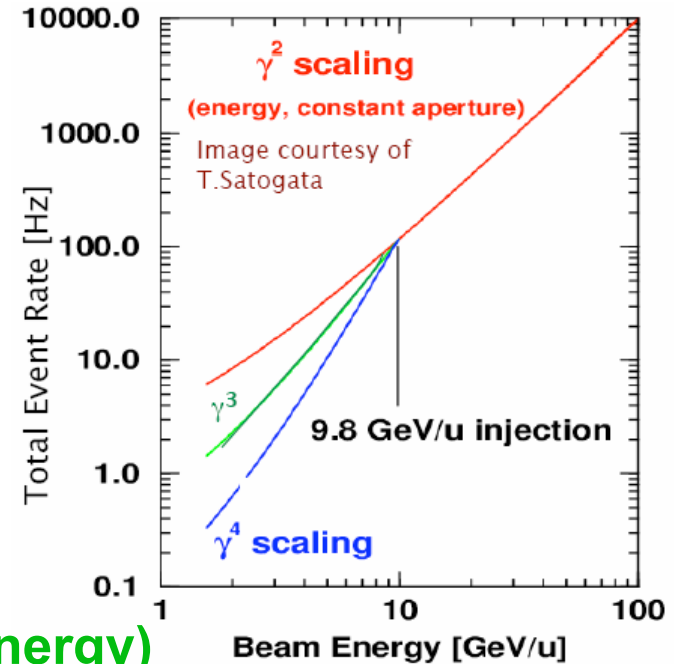
- limited by intra-beam scattering
 - below injection: γ^3 scaling
 - decent event rates above injection
 - difficult below injection energy
 - improvement: electron cooling

- lifetime

- only few minutes (below injection energy)
 - "continuous" injection?
 - improvement: electron cooling

- large "diamond" length

- spread of collision vertices along beam axis
 - improvement: electron cooling



Relevant PHENIX upgrades

- trigger and reaction plane measurement

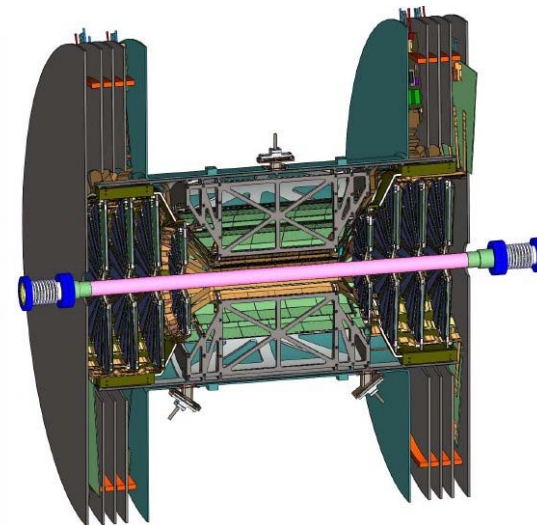
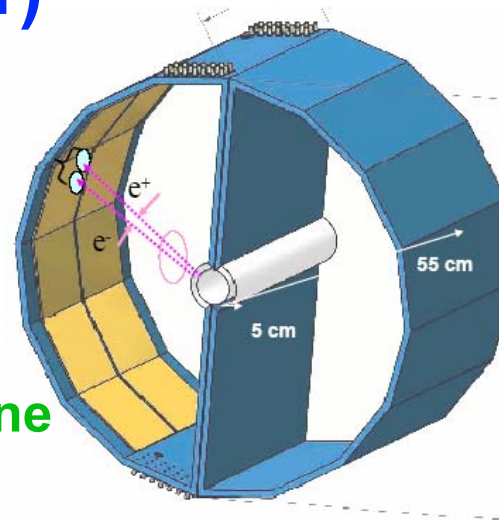
- reaction plane detector
 - already implemented
 - compatible with future upgrades

- electron identification

- Hadron Blind Detector (HBD)
 - commissioning in 2009 p+p run
 - Au+Au run at top energy: 2010

- the future (2010/2011)

- replace HBD with a barrel silicon vertex spectrometer (later: additional endcaps)
 - secondary vertices
 - trigger & reaction plane
 - limited electron ID



Summary & outlook

- PHENIX topics in a RHIC energy scan
 - above injection energy
 - strong program to
 - investigate onset of sQGP signatures
 - » hydrodynamic flow and scaling properties of flow parameters
 - » light/heavy quark opacity
 - » low-mass dielectron enhancement
 - » initial temperature
 - » (HBT & three/multi-particle correlations)
 - search for the QCD critical point
 - below injection energy
 - contribution to a search for the QCD critical point
 - no rare probe physics program unless
 - drastic improvement in RHIC performance
 - » luminosity
 - » lifetime
 - » length of collision diamond
- electron cooling could make a huge difference!