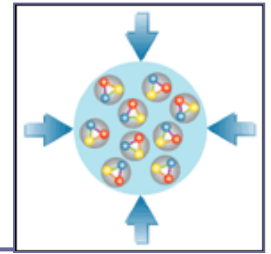




Symposium and 10th CBM Collaboration Meeting
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Forschungszentrum
Dresden Rossendorf



Charm dynamics from transport calculations

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for Advanced Studies



Introduction

FAIR energies are well suited to study **dense and hot nuclear matter** –

- a phase transition to QGP ,
- chiral symmetry restoration,
- in-medium effects

Way to study:

Experimental **energy scan** of different **observables** in order to find an **,anomalous‘** behavior in comparison with theory

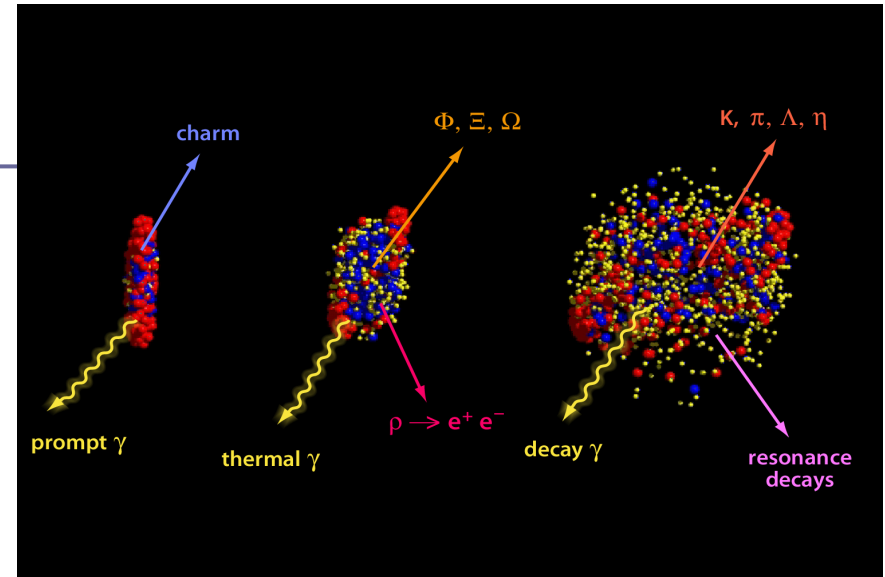
Observables for CBM:

- Excitation function of particle yields and ratios
- Transverse mass spectra
- Collective flow
- Dileptons
- Open and hidden charm
- Fluctuations and correlations
- ...

Microscopic transport models

Signals of the phase transition:

- Strangeness enhancement
- Multi-strange particle enhancement
- Charm suppression
- Collective flow (v_1, v_2)
- Thermal dileptons
- Jet quenching and angular correlations
- High p_T suppression of hadrons
- Nonstatistical event by event fluctuations and correlations
- ...



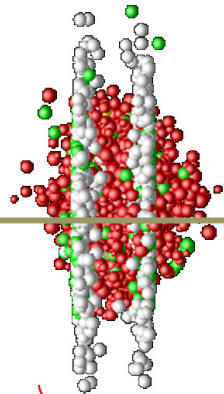
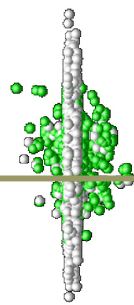
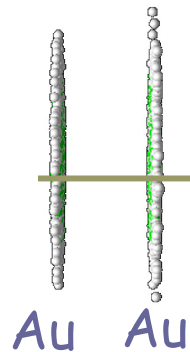
Experiment: measures final hadrons and leptons

How to learn about physics from data?

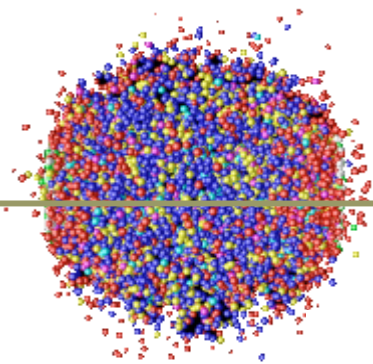
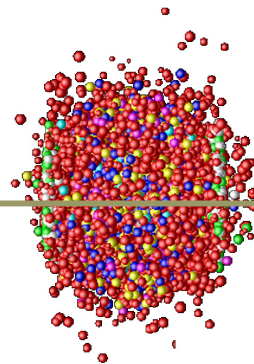
Compare with theory!

Models for heavy ion collisions

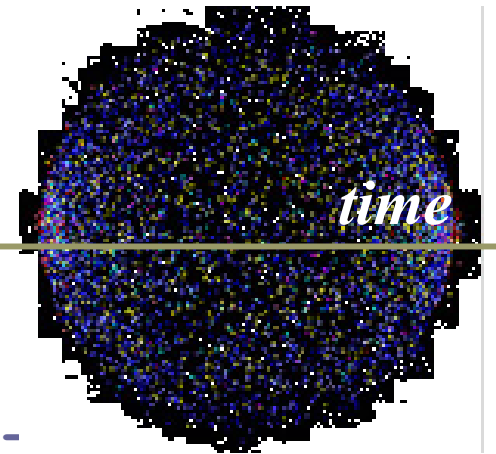
Initial State



Hadronization



Freeze-



Quark-Gluon-Plasma ?

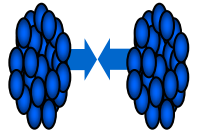
Thermal models

Hydro models (local equilibrium)

Transport models

Microscopical transport models provide the dynamical description of **nonequilibrium** effects in heavy-ion collisions

Basic concepts of Hadron-String Dynamics



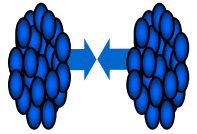
- for each particle species i ($i = N, R, Y, \pi, \rho, K, \dots$) the phase-space density f_i follows the **transport equations**

$$\left(\frac{\partial}{\partial t} + \left(\nabla_{\vec{p}} H \right) \nabla_{\vec{r}} - \left(\nabla_{\vec{r}} H \right) \nabla_{\vec{p}} \right) f_i(\vec{r}, \vec{p}, t) = I_{coll}(f_1, f_2, \dots, f_M)$$

with the **collision terms** I_{coll} describing:

- elastic and inelastic **hadronic reactions** $BB \leftrightarrow B'B'$, $BB \leftrightarrow B'B'm$, $mB \leftrightarrow m'B'$, $mB \leftrightarrow B'$
- formation and decay of baryonic and mesonic **resonances**
- **string** formation and decay (for inclusive production: $BB \rightarrow AX$, $mB \rightarrow AX$, $X = \text{many particles}$)
- Implementation of detailed balance on the level of $1 \leftrightarrow 2$ and $2 \leftrightarrow 2$ reactions (+ $2 \leftrightarrow n$ **multi-meson fusion reactions**)
- **Off-shell dynamics** for short living states

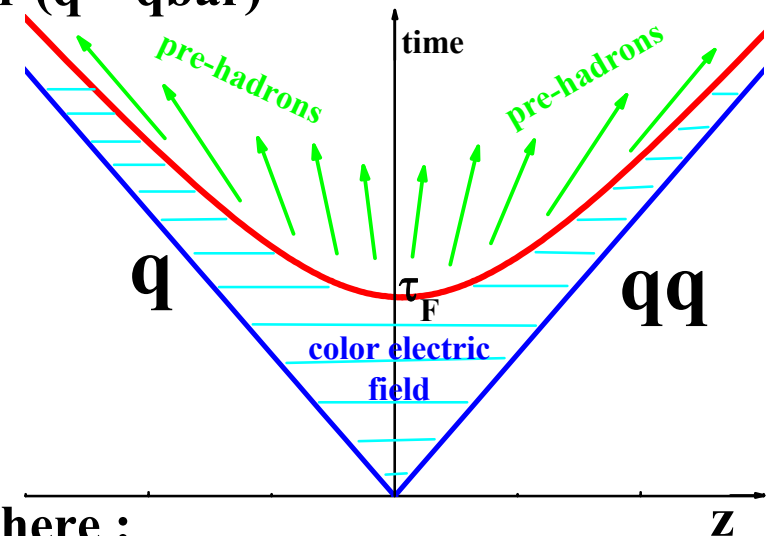
Degrees of freedom in HSD



- hadrons - baryons and mesons including excited states (resonances)
- strings – excited colour singlet states ($qq - q$) or ($q - q\bar{q}$)

Based on the LUND string model
& perturbative QCD via PYTHIA

- leading quarks ($q, q\bar{q}$) & diquarks ($q-q, q\bar{q}-q\bar{q}$)



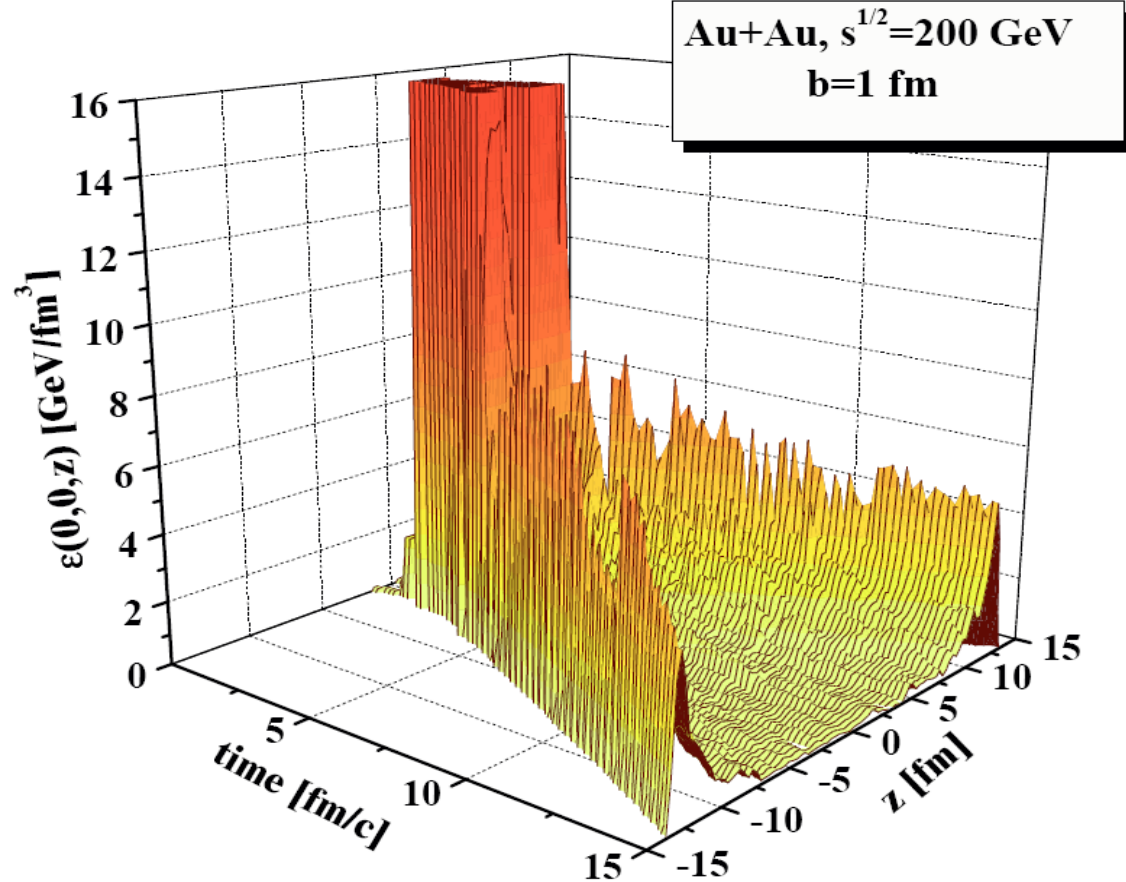
NOT included in the transport models presented here :

- no explicit **parton-parton interactions** (i.e. between quarks and gluons) outside strings!
- no **QCD EoS** for partonic phase

under construction:

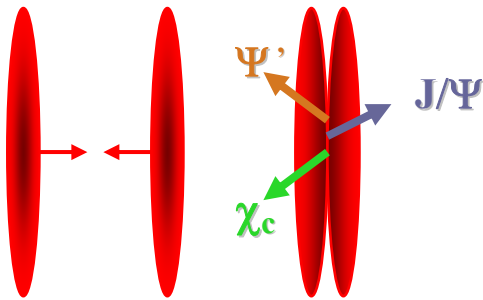
PHSD – Parton-Hadron-String-Dynamics W. Cassing arXiv:0704.1410

Time evolution of the energy density



HSD transport model allows to calculate the energy momentum tensor $T^{\mu\nu}(\mathbf{x})$ for all space-time points \mathbf{x} and thus the energy density $\epsilon(\mathbf{r},t)$ which is identified with $T^{00}(\mathbf{r},t)$

Local energy density ε vs Bjorken energy density ε_{Bj}



- transient time for central Au+Au at 200 GeV:
 $t_r \sim 2R_A/\gamma_{cm} \sim 0.13 \text{ fm}/c$
- $c\bar{c}$ formation time:
 $\tau_C \sim 1/M_T \sim 1/4\text{GeV} \sim 0.05 \text{ fm}/c < t_r$
- $c\bar{c}$ pairs are produced in the **initial hard NN** collisions
in time period t_r

Bjorken energy density:

$$\varepsilon_{Bj} = \frac{1}{A_{\perp} \tau} \frac{dE_T}{dy}$$

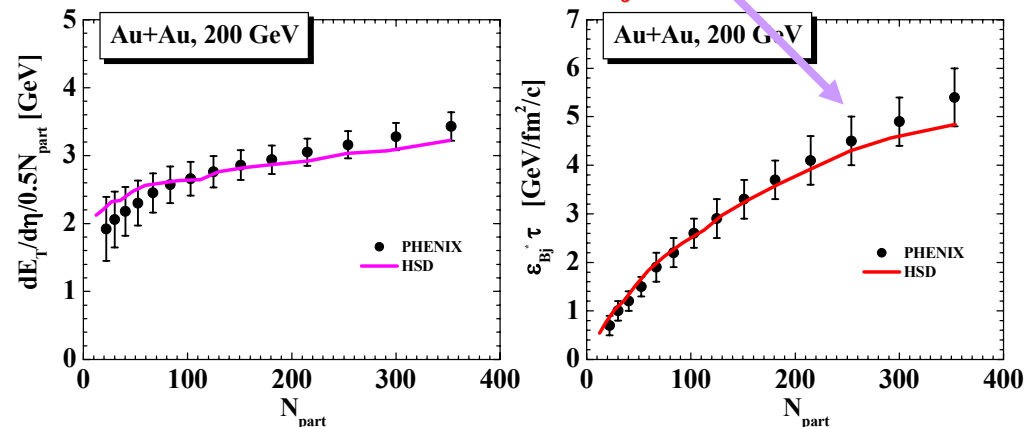
A_T is the nuclei transverse overlap area
 τ is the formation time of the medium

at RHIC $\varepsilon_{Bj} \tau \sim 5 \text{ GeV}/\text{fm}^2/c$

„Local“ energy density ε during
transient time t_r :

$$\varepsilon \sim 5 [\text{GeV}/\text{fm}^2/c] / [0.13 \text{ fm}/c] \\ \sim 30 \text{ GeV}/\text{fm}^3$$

accounting τ_C : $\varepsilon \sim 28 \text{ GeV}/\text{fm}^3$



- ✓ HSD reproduces PHENIX data for Bjorken energy density very well
- ✓ HSD results are consistent with simple estimates for the energy density

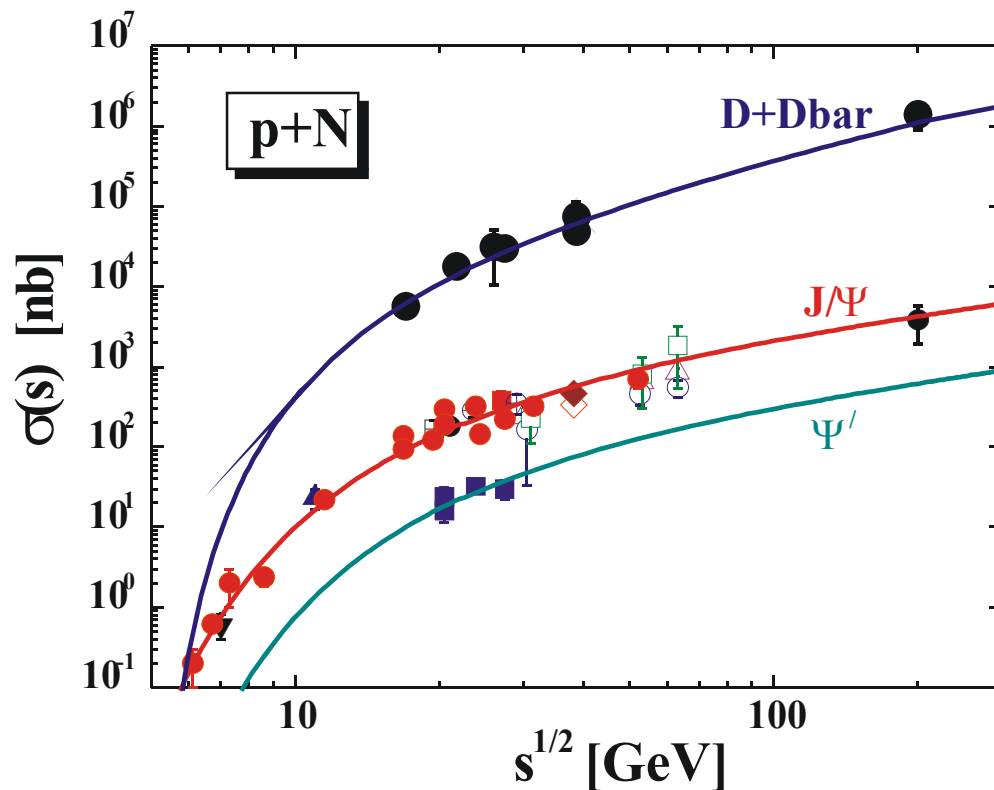
Charmonium production in pN

Hard probe $0A$ **binary scaling!**

$\sigma(J/\Psi)$ and $\sigma(\Psi')$:
 parametrization of the
 available
 experimental data

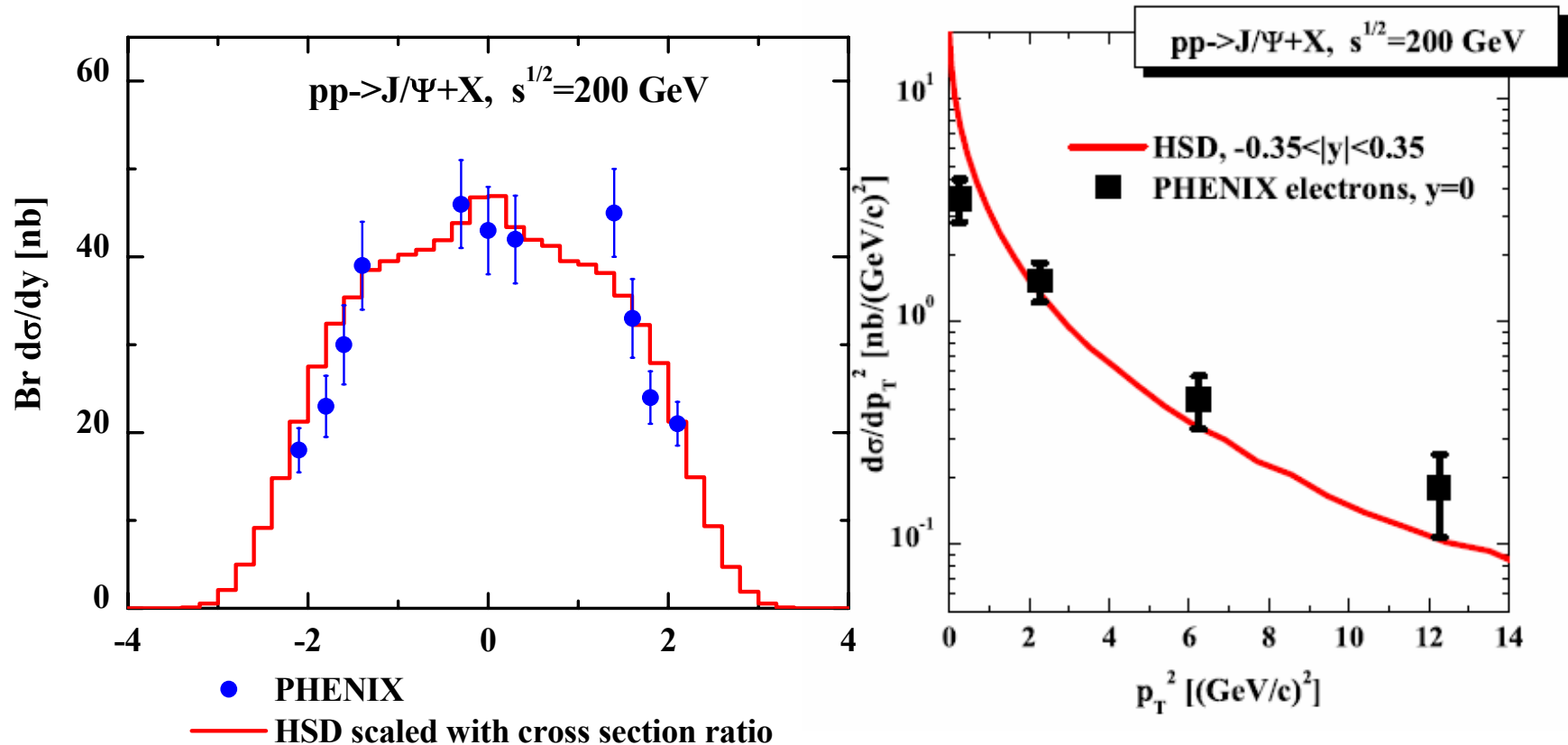
But data close to
 threshold
 are still needed !

FAIR at GSI



$$\sigma_{J/\Psi}^{\text{exp}} = \sigma_{J/\Psi} + B(\chi_c \rightarrow A J/\Psi) \sigma_{\chi_c} + B(\Psi' \rightarrow J/\Psi) \sigma_{\Psi'}$$

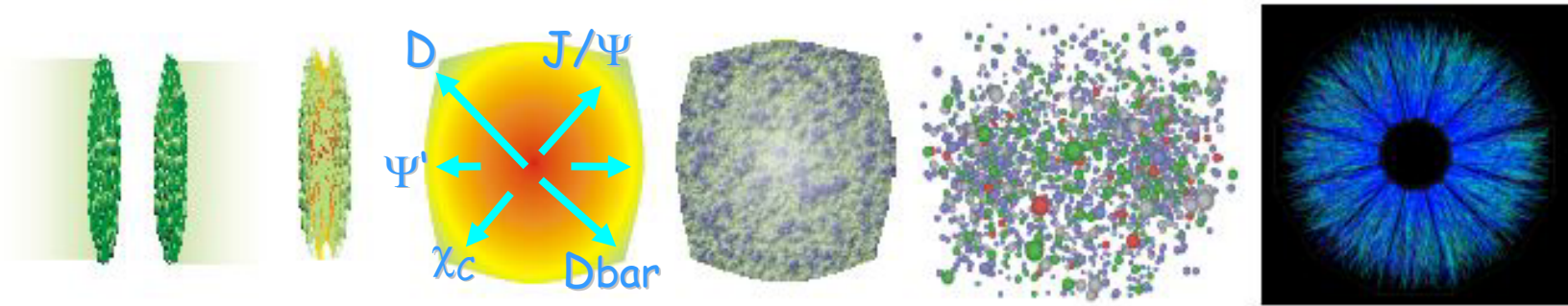
Charmonium production in pN



Differential cross section of charm production is successfully parametrized, too

Charmonium production *vs* absorption

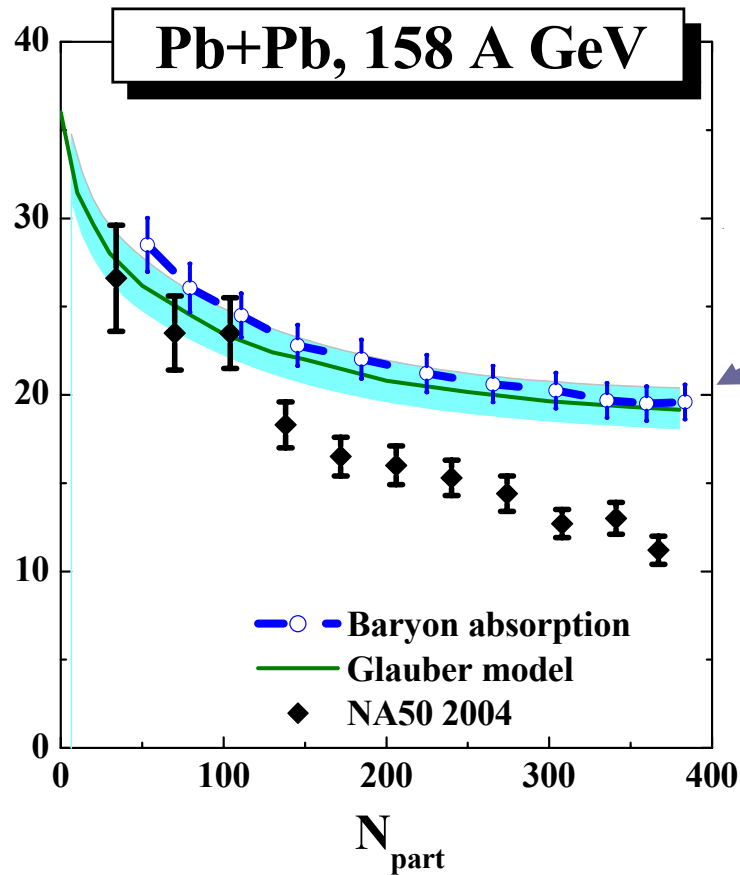
Charm sector reflects the dynamics **in the early phase** of heavy-ion collisions !



Charmonium is absorbed by

- Scattering on nucleons (**normal nuclear absorption**, as in pA)
- Interaction with secondary hadrons (**comovers**)
- Dissociation in the deconfined medium (**suppression in QGP**)

Anomalous J/Ψ suppression



J/Ψ ,normal' absorption by nucleons
(Glauber model)

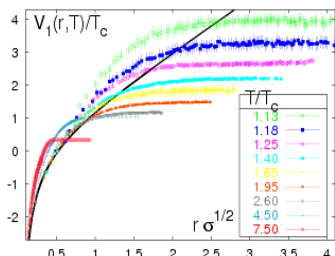
Experimental observation (NA38/50/60):
extra suppression in A+A collisions;
increasing with centrality

Scenarios for anomalous charmonium suppression

• QGP colour screening

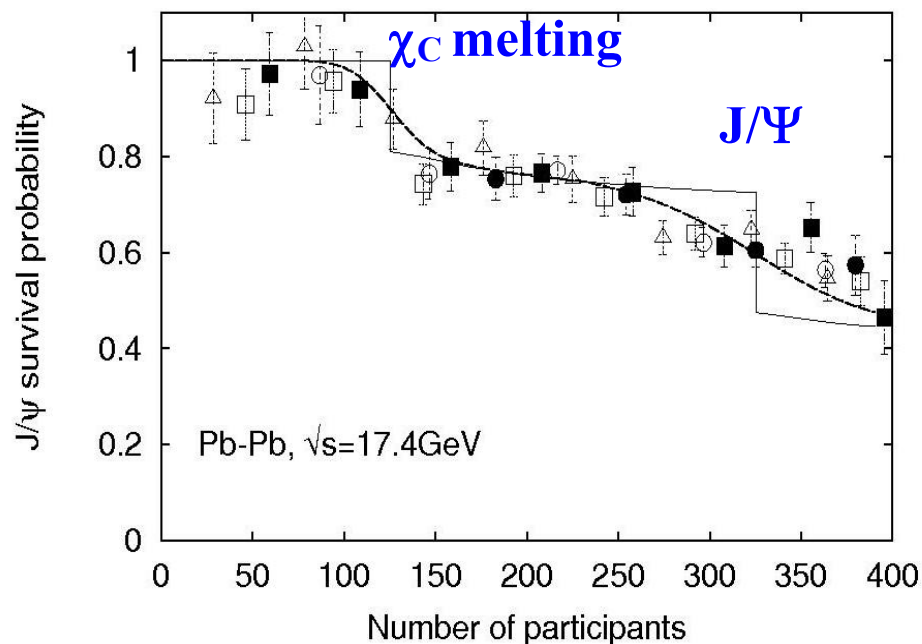
[Digal, Fortunato, Satz '03]

Quarkonium dissociation T :



state	$J/\psi(1S)$	$\chi_c(1P)$	$\psi'(2S)$
T_d/T_c	2.10	1.16	1.12

Dissociation energy density $\epsilon_d \sim 2(T_d/T_c)^4$



• Comover absorption

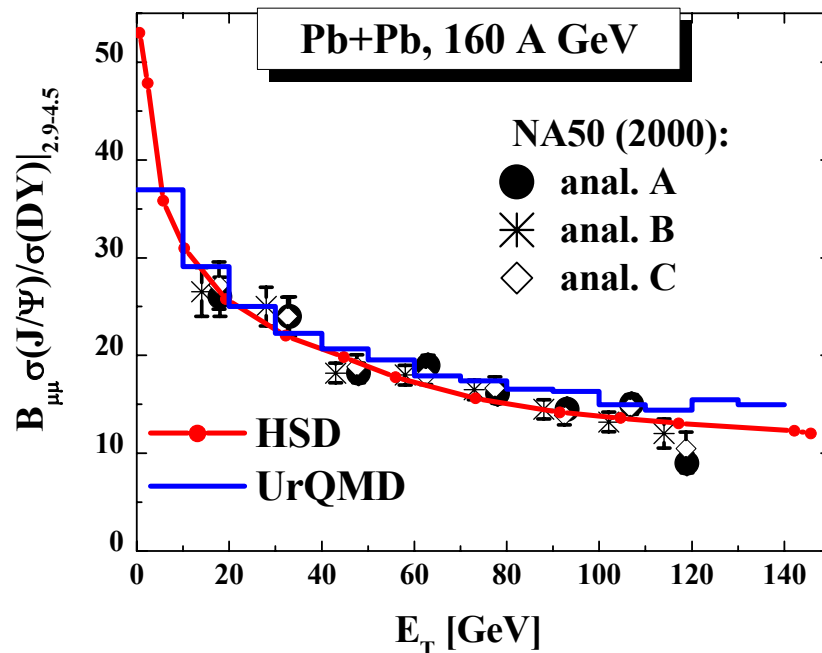
[Gavin & Vogt, Capella et al. '97]

absorption by low energy inelastic scattering with 'comoving' mesons ($m=\pi, \eta, \rho, \dots$)

$J/\Psi + m \rightarrow 0A \ D + Dbar$

$\Psi' + m \rightarrow 0A \ D + Dbar$

$\chi_c + m \rightarrow 0A \ D + Dbar$



Modelling the **comover** scenario in **HSD**

1. Charmonia **dissociation** cross sections with π, ρ, K and K^* mesons

$J/\Psi (\chi_c, \Psi') + \text{meson } (\pi, \rho, K, K^*) \leftrightarrow D + Dbar$

- **Phase-space model** for charmonium + meson dissociation:

$$\sigma_{1+2 \rightarrow 3+4}(s) = g_{\text{isospin}} 2^4 \frac{E_1 E_2 E_3 E_4}{s} |M_i|^2 \left(\frac{m_3 + m_4}{\sqrt{s}} \right)^6$$

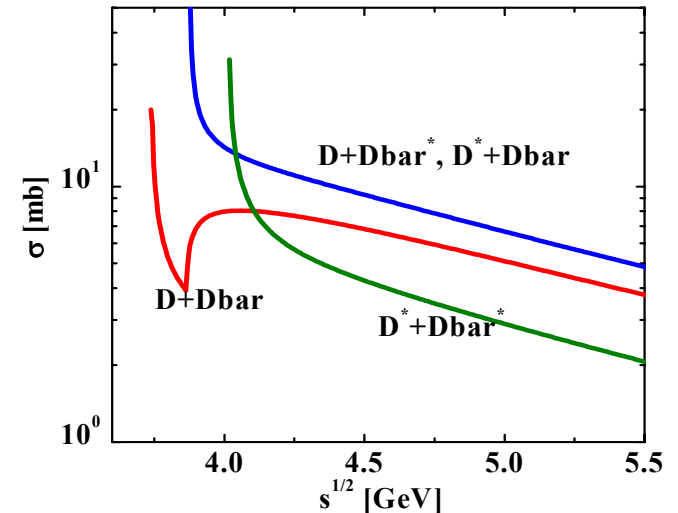
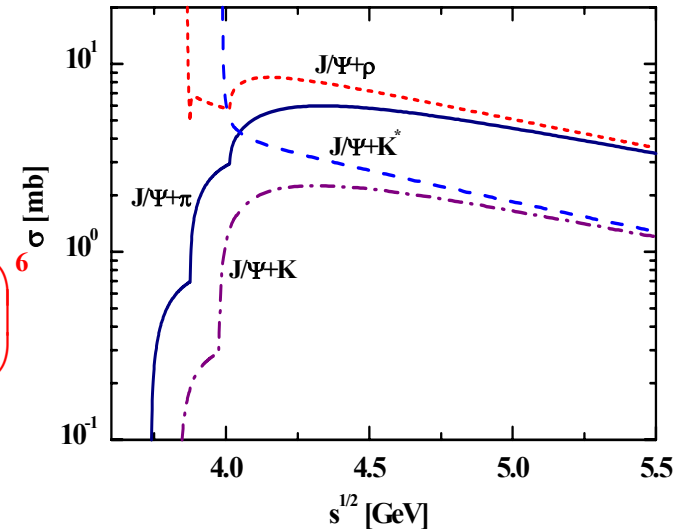
$$i = \chi_c, J/\Psi, \Psi'$$

$$|M_{J/\Psi}|^2 = |M_{\chi_c}|^2 = |M_{\Psi'}|^2 = |M_0|^2$$

constant matrix element

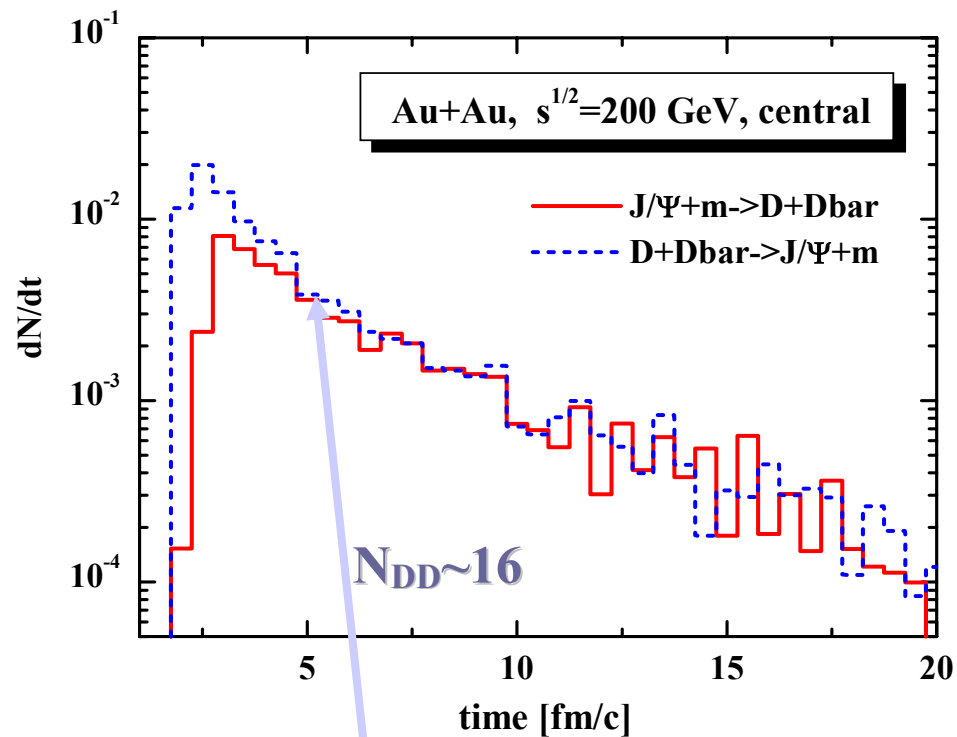
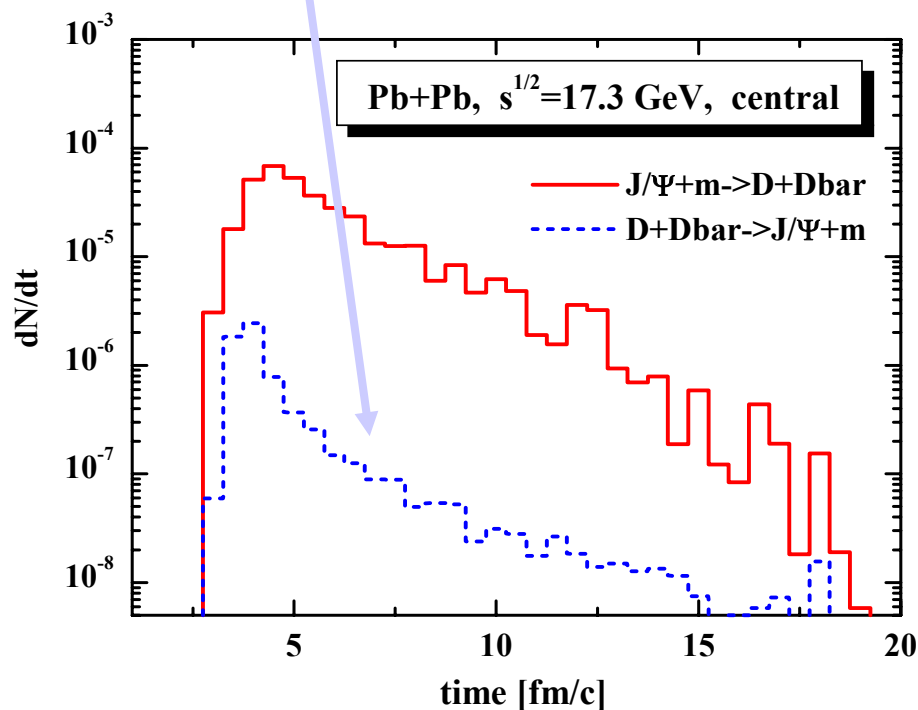
2. J/Ψ **recombination** cross sections by $D + Dbar$ annihilation:

$D + Dbar \rightarrow J/\Psi (\chi_c, \Psi') + \text{meson } (\pi, \rho, K, K^*)$
are determined by **detailed balance!**



Charmonium recombination by DDbar annihilation

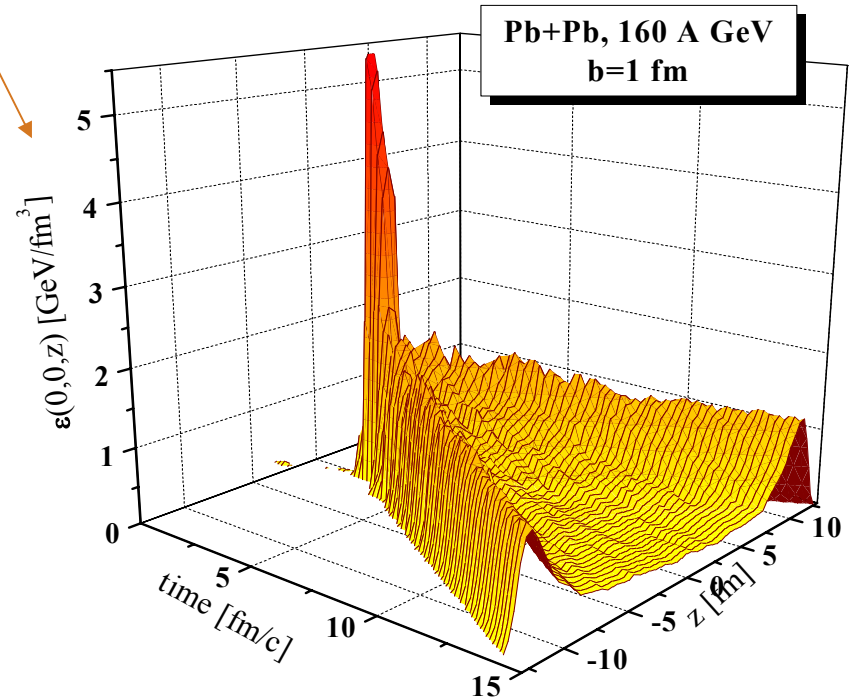
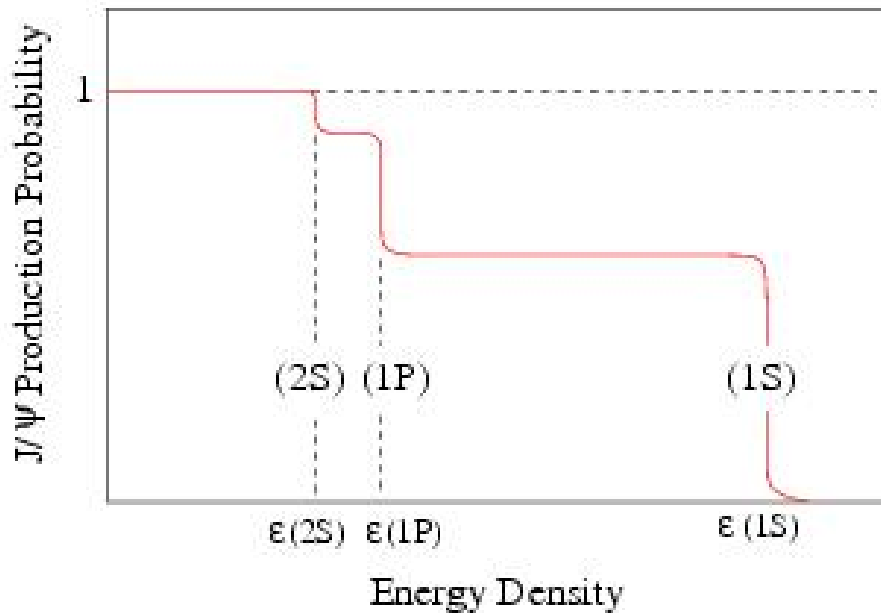
At **SPS** recreation of J/Ψ by D-Dbar annihilation is **negligible**



But at **RHIC** recreation of J/Ψ by D-Dbar annihilation is **strong!**

Modeling the QGP melting in HSD

Energy density $\varepsilon(x=0, y=0, z; t)$ from HSD



Threshold energy densities:

J/Ψ melting: $\varepsilon(J/\Psi) = 16 \text{ GeV/fm}^3$

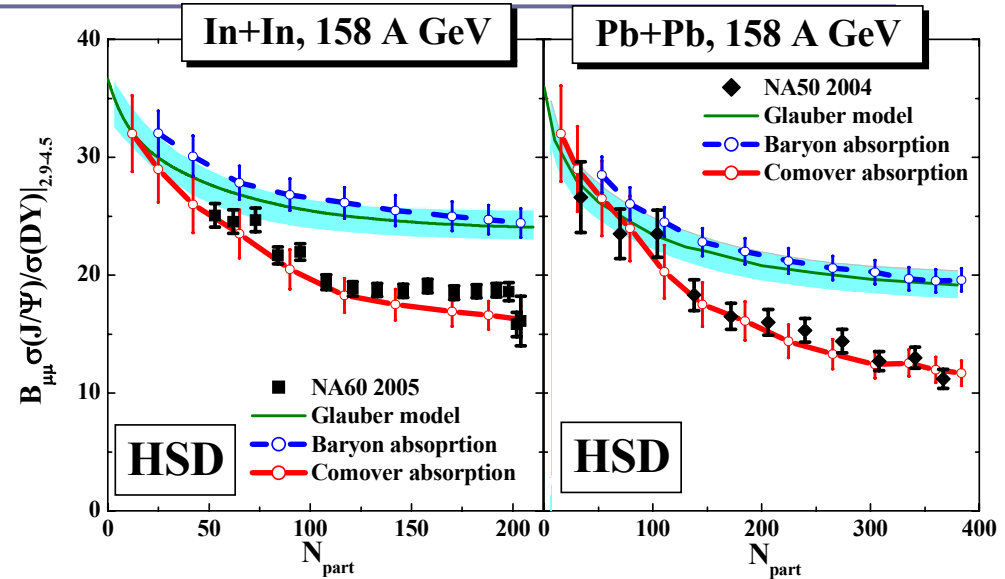
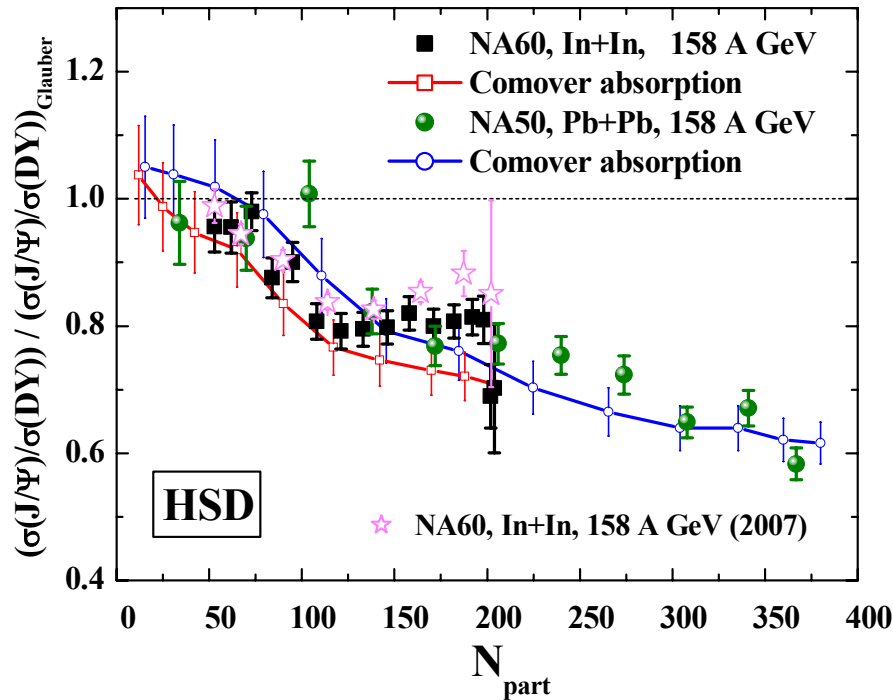
χ_c melting: $\varepsilon(\chi_c) = 2 \text{ GeV/fm}^3$

Ψ' melting: $\varepsilon(\Psi') = 2 \text{ GeV/fm}^3$

Comparison to data at SPS energy

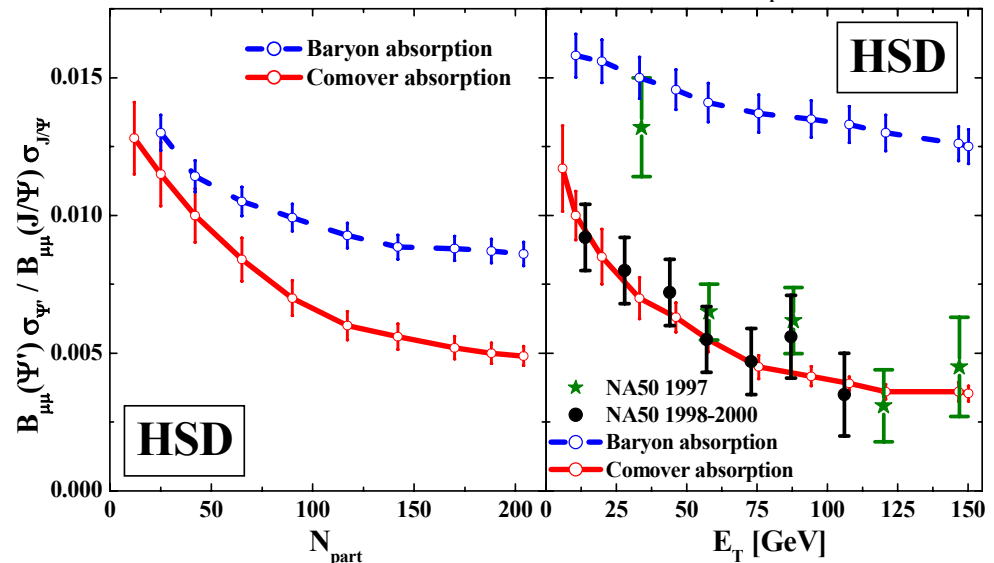
Pb+Pb and In+In @ 158 A GeV

comover absorption



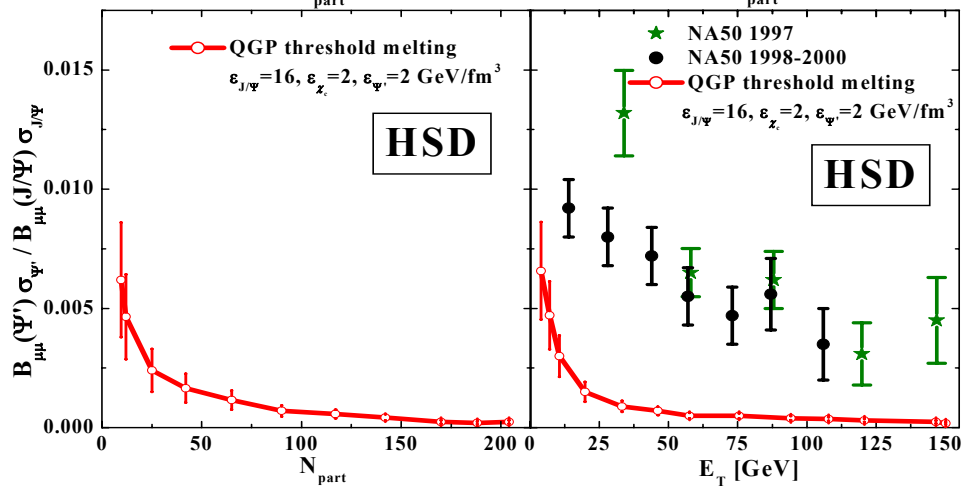
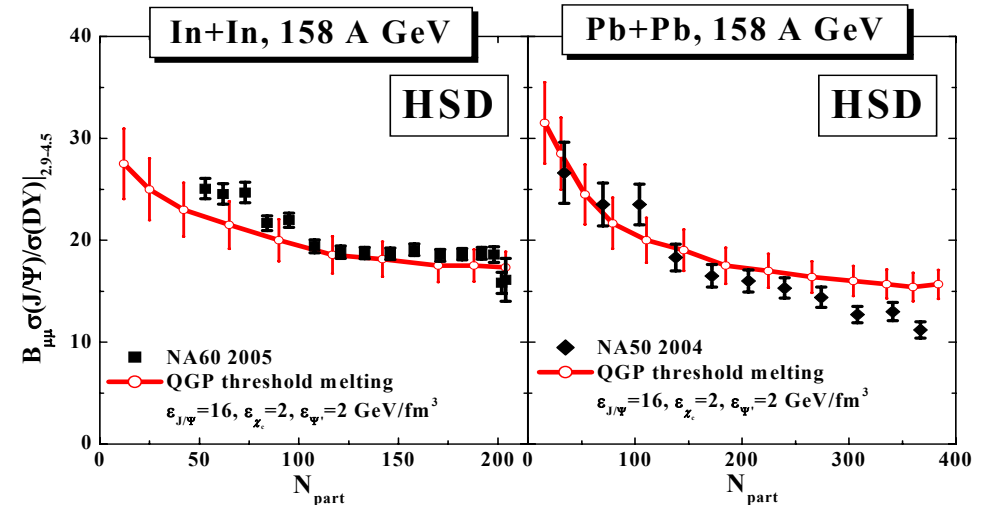
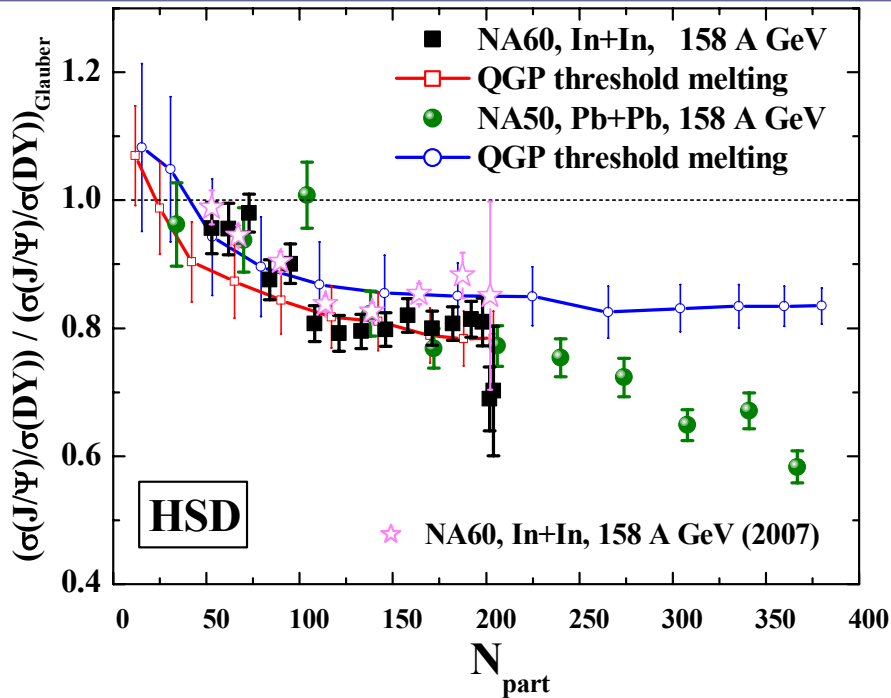
Pb+Pb and In+In @ 160 A GeV
consistent with the comover absorption
for the same parameter set!

[OL et al NPA786 (2007) 183]



Pb+Pb and In+In @ 158 A GeV

QGP threshold melting

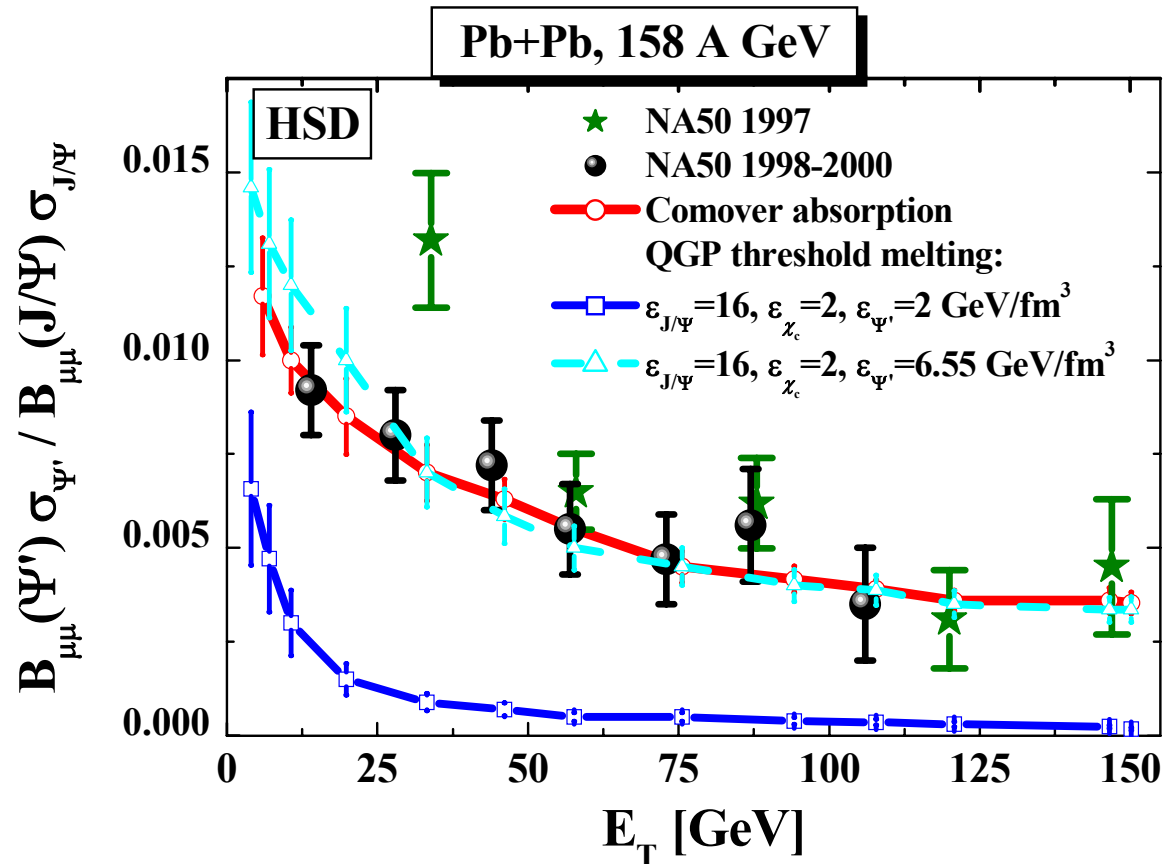


absorption too strong,
which contradict data

Ψ' data contradict threshold melting scenario with lQCD ε^d

$\varepsilon(J/\Psi) = 16 \text{ GeV/fm}^3$,
 $\varepsilon(\chi_c) = 2 \text{ GeV/fm}^3$,
 $\varepsilon(\Psi') = 6.55 \text{ GeV/fm}^3$

- Set 2: an increase of the melting energy density $\varepsilon(\Psi') = 6.55 \text{ GeV/fm}^3$ reduces the Ψ' suppression, but contradicts LQCD predictions for $T^d(\Psi') \sim 1.2 T_C!$**



Comparison to data at RHIC energy

Comover absorption + regeneration

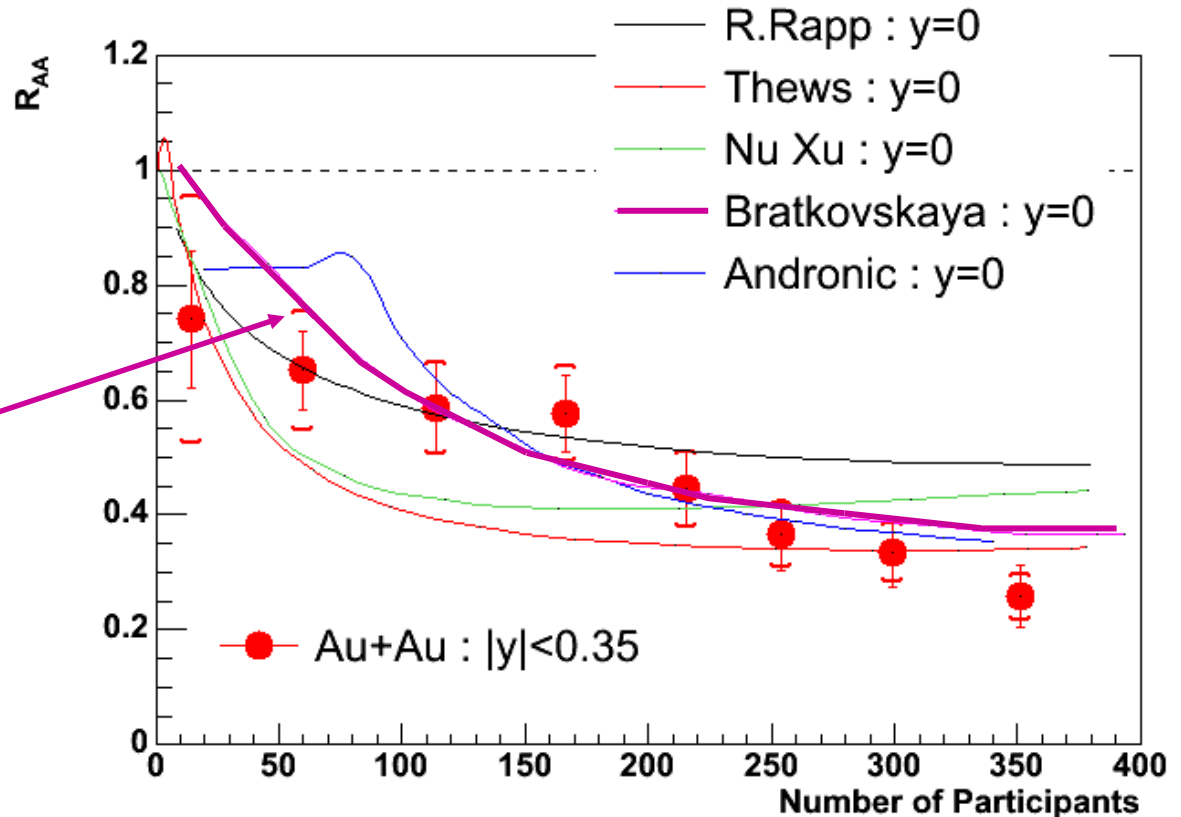
A successful prediction

Regeneration is essential!

HSD

NB:

obtained assuming
the existence of **comovers**
throughout the collision,
i.e. **at all energy densities.**

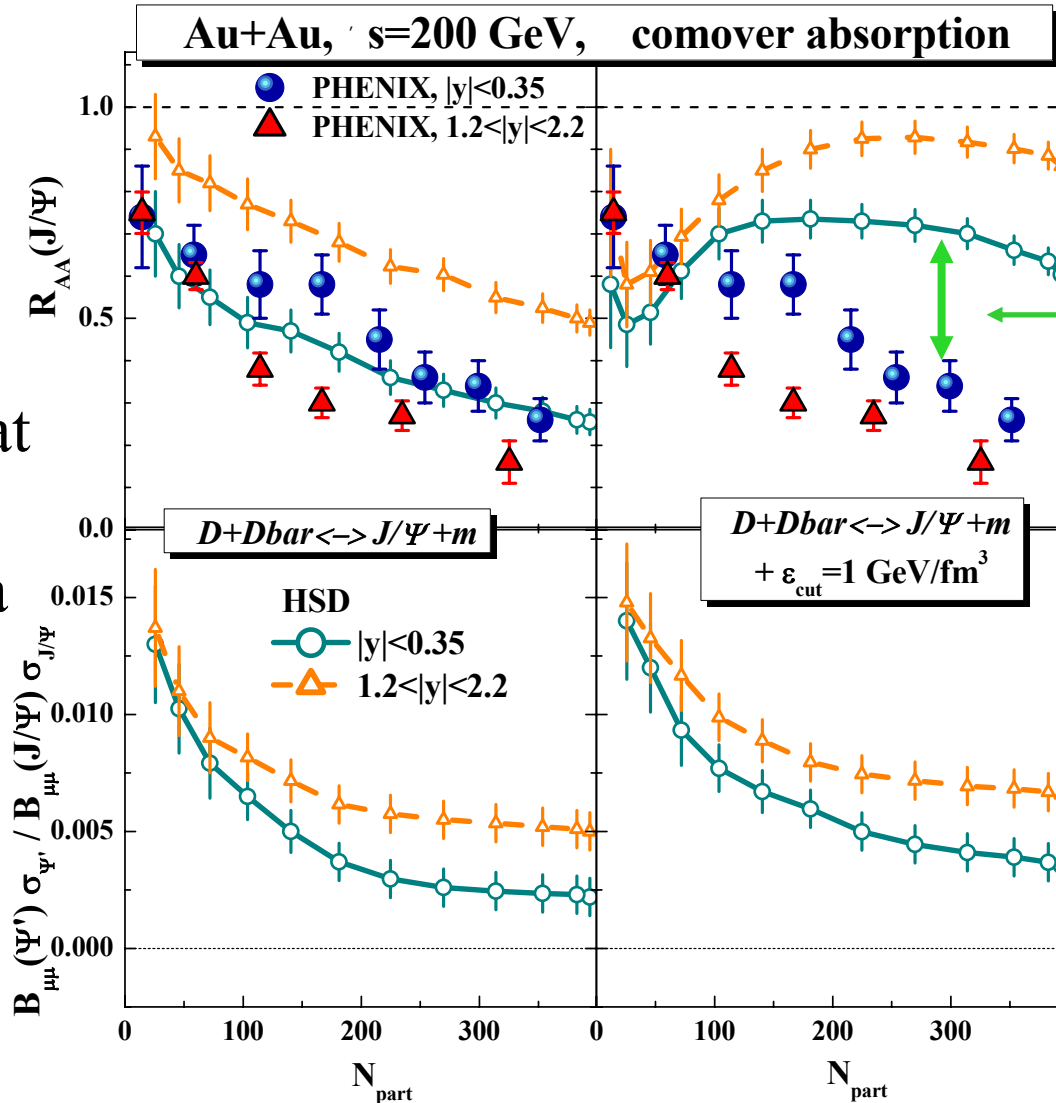


R. Rapp et al. PRL 92, 212301 (2004)
R. Thews et al, Eur. Phys. J C43, 97 (2005)
Yan, Zhuang, Xu, PRL97, 232301 (2006)
Bratkovskaya et al., PRC 69, 054903 (2004)
A. Andronic et al., NPA789, 334 (2007)

Au+Au @ $s^{1/2}=200$ GeV

Comover absorption + regeneration

In comover scenario, suppression at **mid-y** stronger than at **forward y**, unlike the data

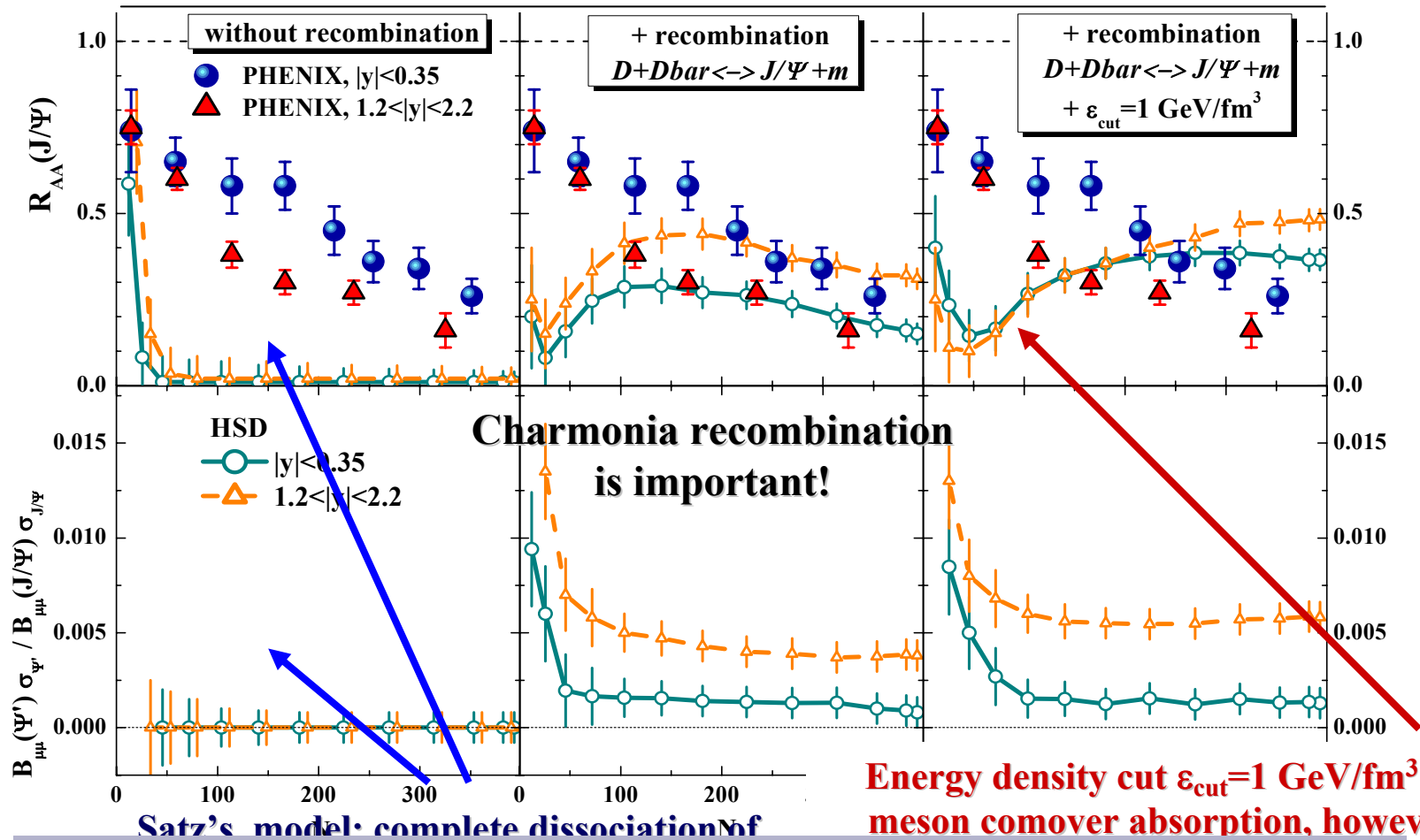


Space for parton phase effects

Energy density cut $\epsilon_{cut} = 1 \text{ GeV/fm}^3$ reduces the meson comover absorption

Au+Au @ $s^{1/2}=200$ GeV

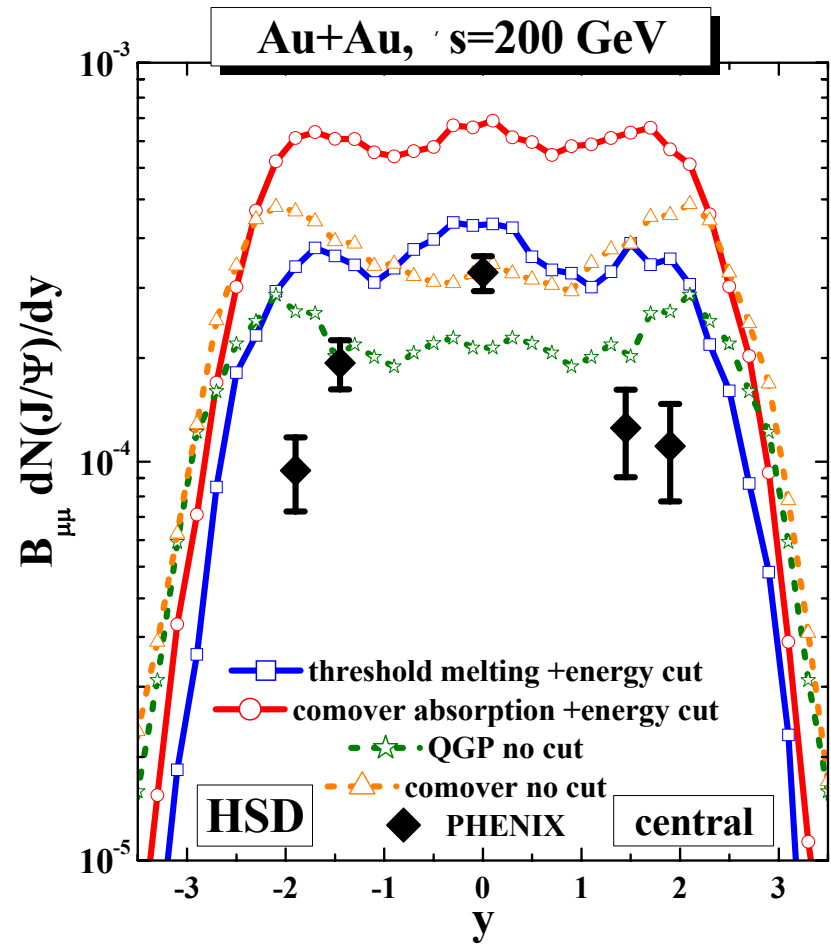
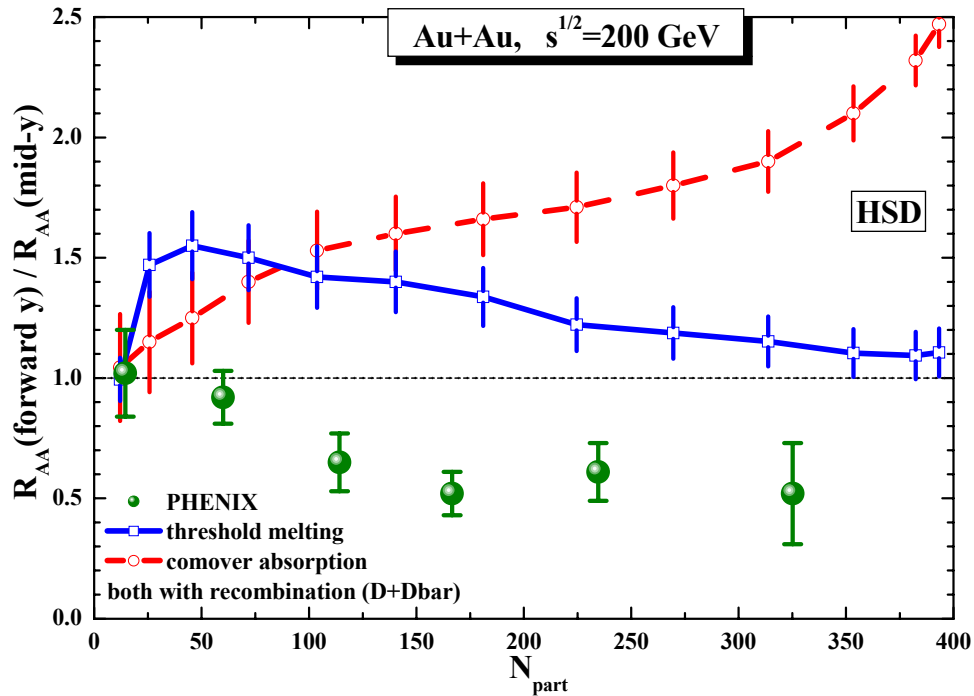
Threshold melting



QGP threshold melting scenario is ruled out by PHENIX data! ... high charmonia, ... especially for peripheral collisions!

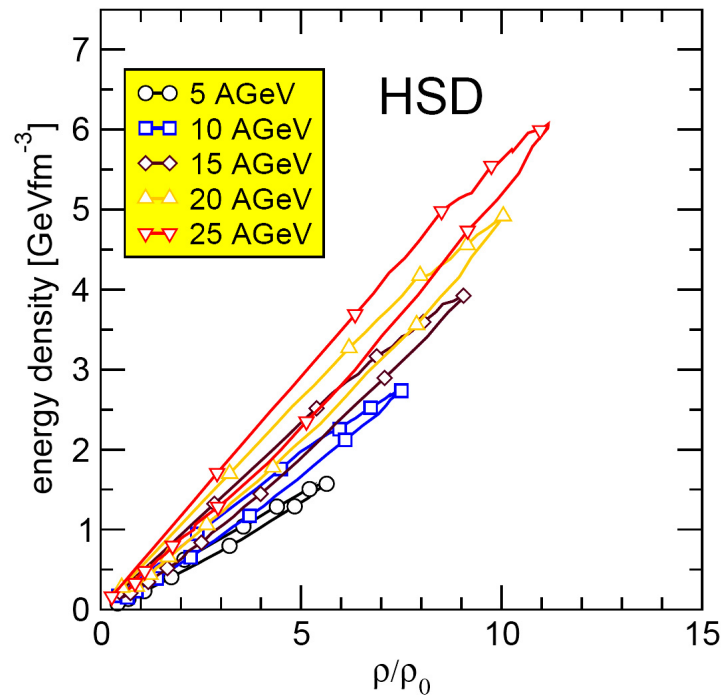
local energy densities!

Rapidity !



HSD predictions for FAIR energy

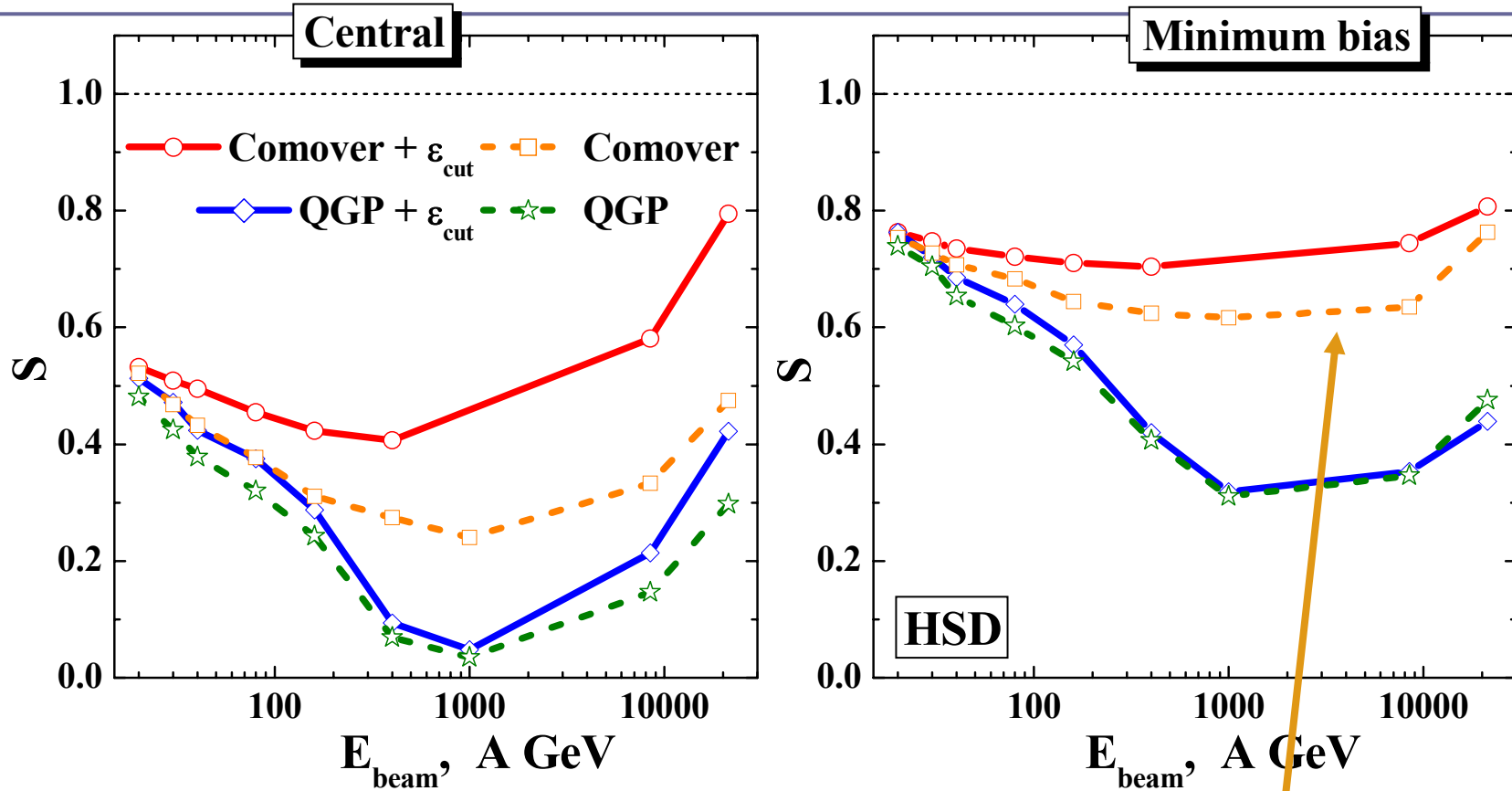
Energy density at FAIR



Huge energy density is reached ($\epsilon > \epsilon_{\text{crit}} = 1 \text{ GeV}/\text{fm}^3$) also at FAIR ($> 5 \text{ A GeV}$).

Additionally, high baryon density.

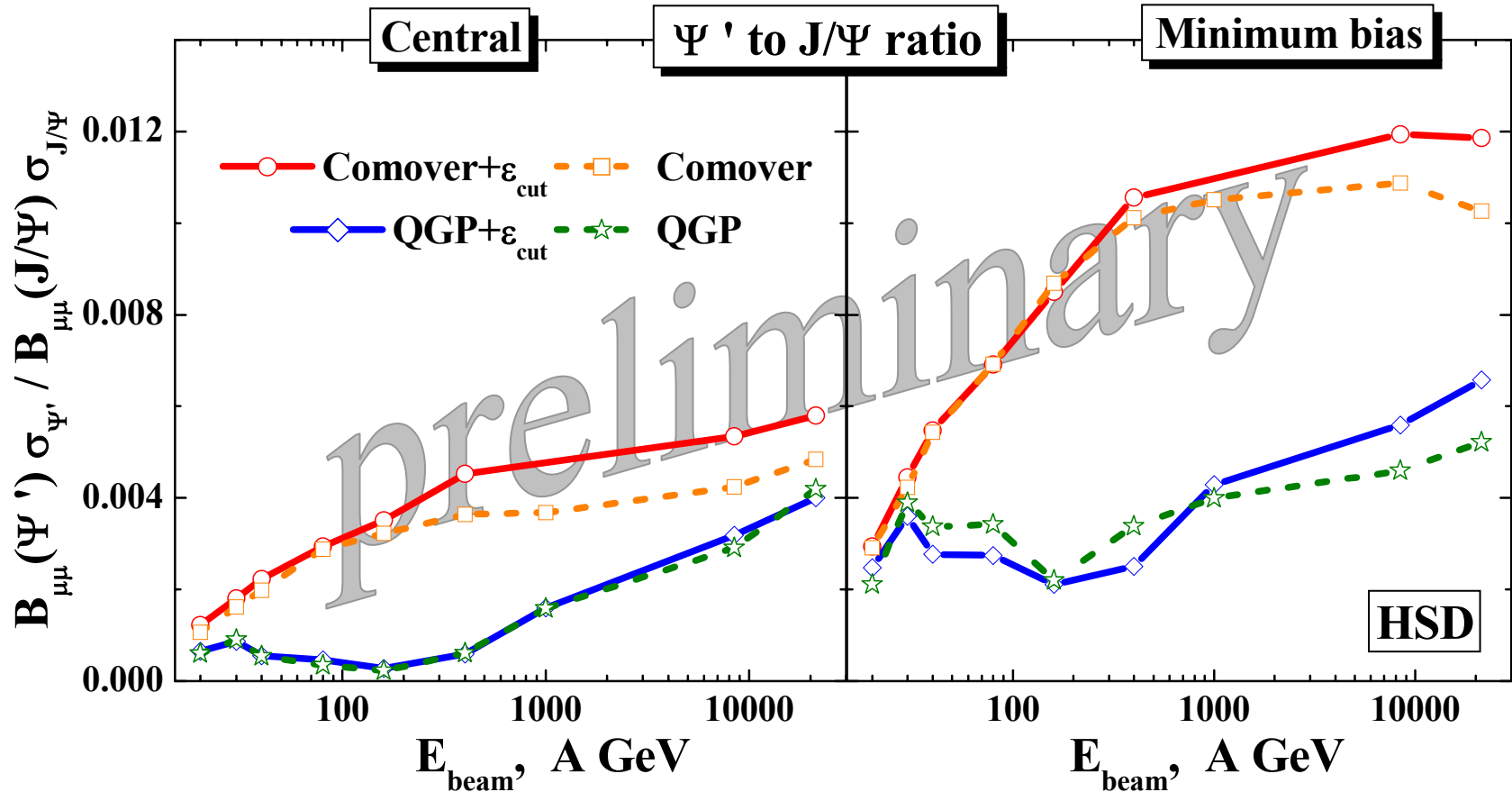
J/ψ excitation function



Comover reactions in the hadronic phase give almost a constant suppression; pre-hadronic reactions lead to a larger recreation of charmonia with E_{beam} .

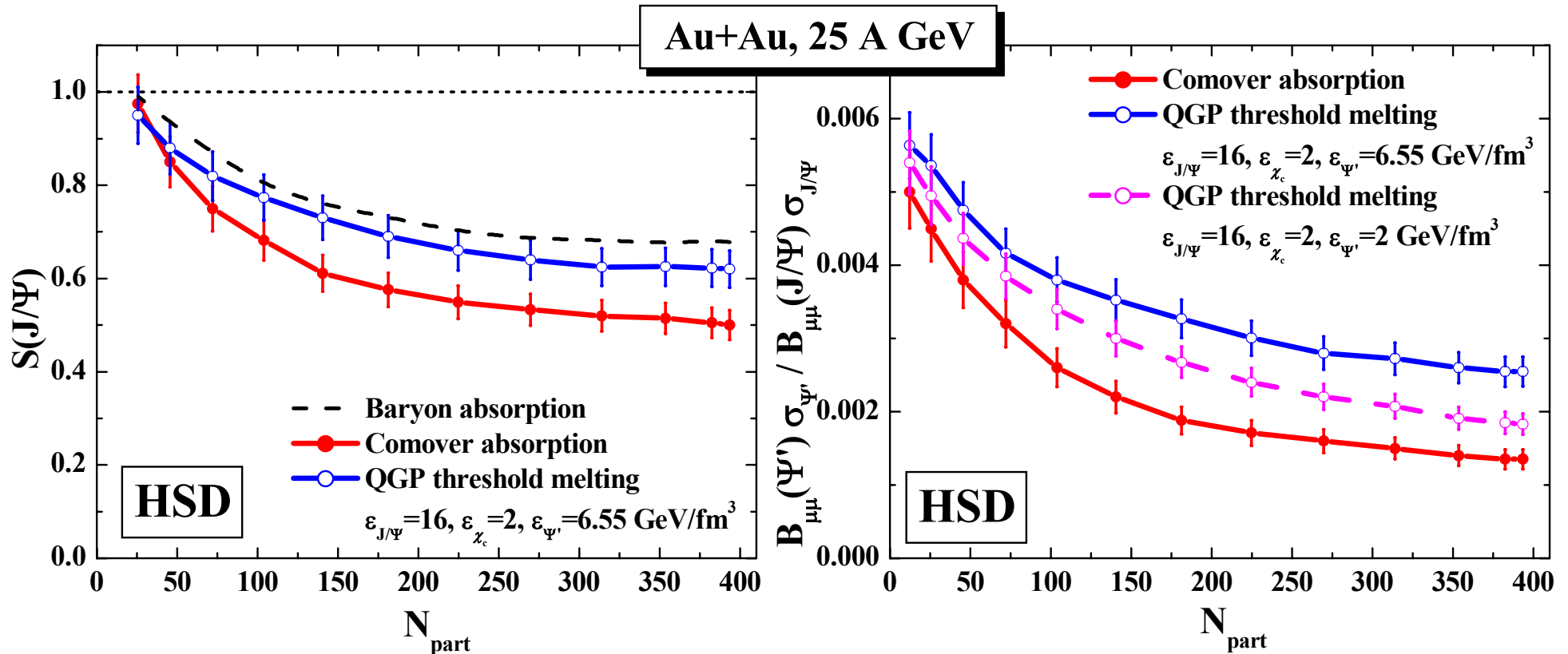
The J/ψ **melting scenario** with hadronic comover recreation shows a **maximum suppression** at $E_{\text{beam}} = 1 \text{ A TeV}$; **exp. data ?**

excitation function



Different scenarios can be distinguished at FAIR energies:
Comover scenario predicts a smooth excitation function whereas the
 ‘**threshold melting**’ shows a step in the excitation function

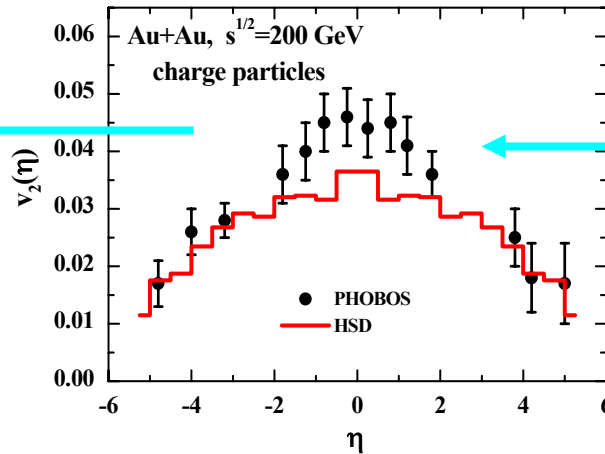
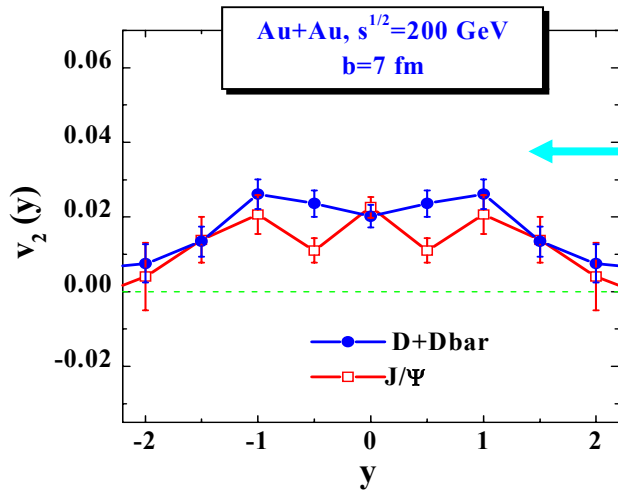
Predictions for J/Ψ and Ψ' suppression in Au+Au at **CBM**



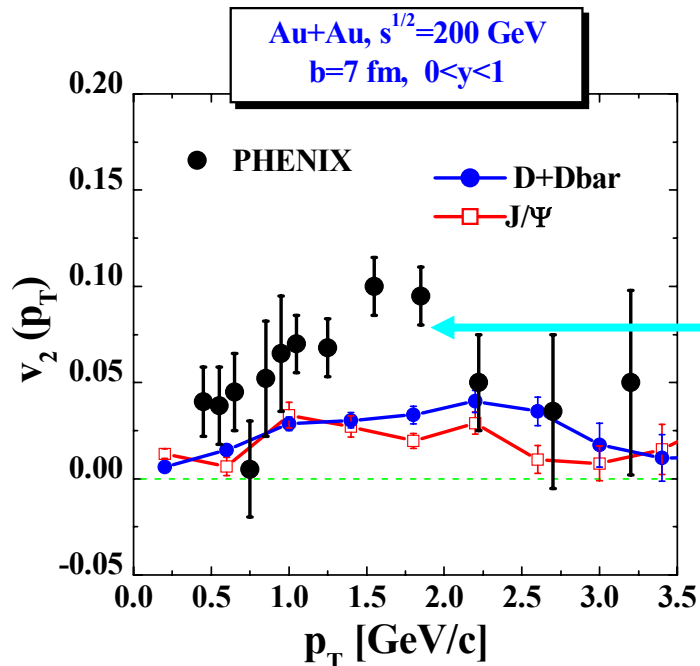
Possible mechanisms can be disentangled:

$\Psi'/(J/\Psi)$ is lower in the 'comover absorption' since the average comover density decreases only moderately with lower bombarding energy whereas the energy density falls rapidly

HSD: v_2 of D+Dbar and J/ Ψ from Au+Au versus p_T and y at RHIC



Collective flow from **hadronic interactions** is too low at midrapidity !



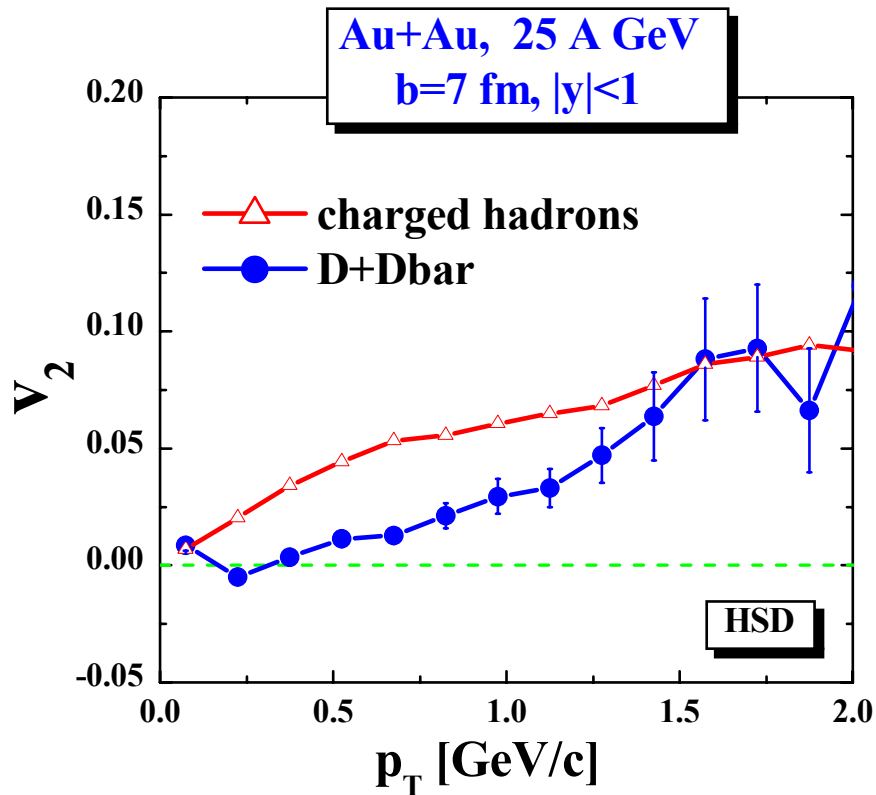
• **HSD:** D-mesons and J/ Ψ follow the charged particle flow \Rightarrow **small $v_2 < 3\%$**

• **Exp. data at RHIC** show **large** collective flow of D-mesons up to **$v_2 \sim 10\%$** !

\Rightarrow strong initial flow of **non-hadronic nature!**

HSD predictions for CBM

elliptic flow at 25 A GeV



•HSD: D-mesons and J/ Ψ
follow the charged particle flow
 \Rightarrow small v_2

Possible observation at CBM:
strong initial flow of D-mesons
and J/ Ψ due to partonic
interactions!

Challenge for CBM!

Summary

- J/Ψ probes early stages of fireball and HSD is the tool to model it.
- Comover absorption and threshold melting both reproduce J/Ψ survival in $Pb+Pb$ as well as in $In+In$ @ 158 A GeV, while Ψ' data are in conflict with the melting scenario.
- Comover absorption and colour screening fail to describe $Au+Au$ at $s^{1/2}=200$ GeV at mid- and forward rapidities simultaneously.
- Deconfined phase is clearly reached at RHIC, but a theory having the relevant/proper degrees of freedom in this regime is needed to study its properties (\rightarrow PHSD).

PHSD - transport description of the partonic and hadronic phases

E. Bratkovskaya, W. Cassing, H. Stöcker

Thank you!

Transport approach (HSD, UrQMD, ...)

- Non-equilibrium
 - > full evolution of the collision
- Universality
 - > large range of $s^{1/2}$ from one code
 - > predictions
 - > excitation functions
- High precision
 - > distinguish physical mechanisms
 - > possibility of verification by exp