

D mesons in hot nuclear matter

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Outline

1. Motivation

2. Model: Self-consistent coupled-channel approach for DN scattering → $\Lambda_c(2593)$ and $\Sigma_c(2800)$ are dynamically generated

3. Open-charm mesons in hot dense matter

4. Conclusions & Outlook

Motivation I (experimental facts)

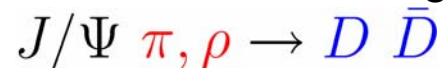
1. J/Ψ suppression: it has been advocated as signature of QGP due to color screening

T.Matsui and H. Satz, PLB 178 (1986) 416

NA50 Collaboration, M.Gonin et al., NPA 610 (1996) 404c

but there are alternative explanations in terms of conventional hadronic models:

→ “comover” scattering:



N. Armesto, A. Capella, PLB430 (1998) 23

B.Zhang, C.M.Ko, B.A.Li, Z.W.Lin, S.Pal, PRC65 (2002) 054909

W. Cassing, E.L. Bratkovskaya, S. Juchem, NPA674 (2000) 249

O. Lynnyk, E.L. Bratkovskaya, W. Cassing, H. Stöcker, NPA786 (2007) 183

→ statistical hadronization model:

A. Andronic, P. Braun-Munzinger, K. Redlich, J. Stachel.

PLB571 (2003) 36; NPA789 (2007) 334; nucl-th/0701079

2. Open-charm enhancement: $NN \rightarrow NN D^+ D^-$ $E_{\text{th}} = 2.8 \text{ GeV}$

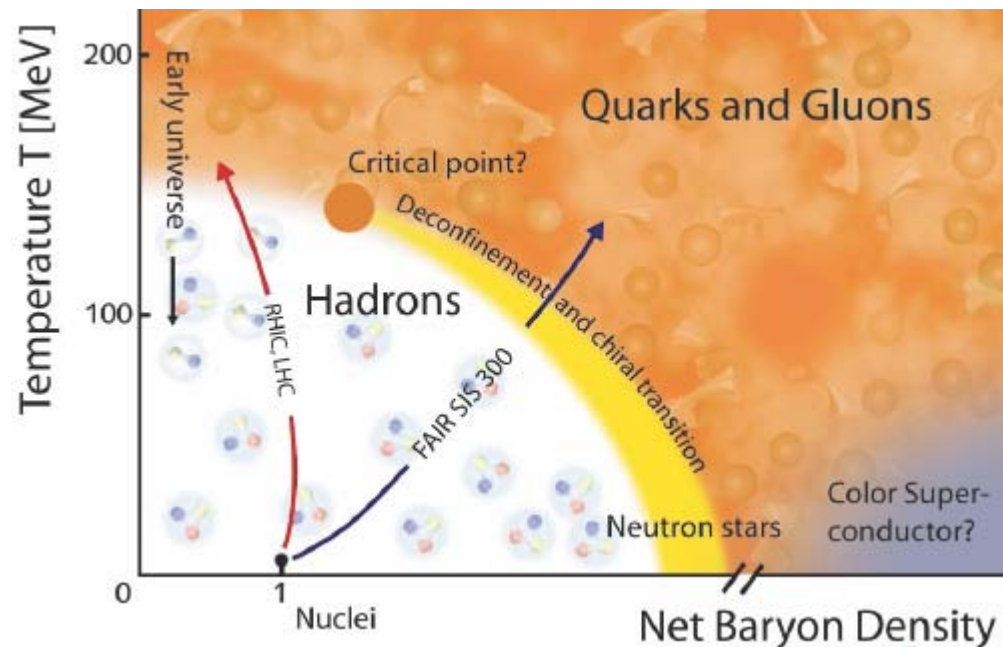
NA50 Collaboration, M.C.Abreu et al., EPJ C14 (2000) 443

it would facilitate J/Ψ suppression but ... recent debate because of different interpretation using the improved data of the NA60 collaboration

NA60 Collaboration, E.Scomparin, talk @ QM2005

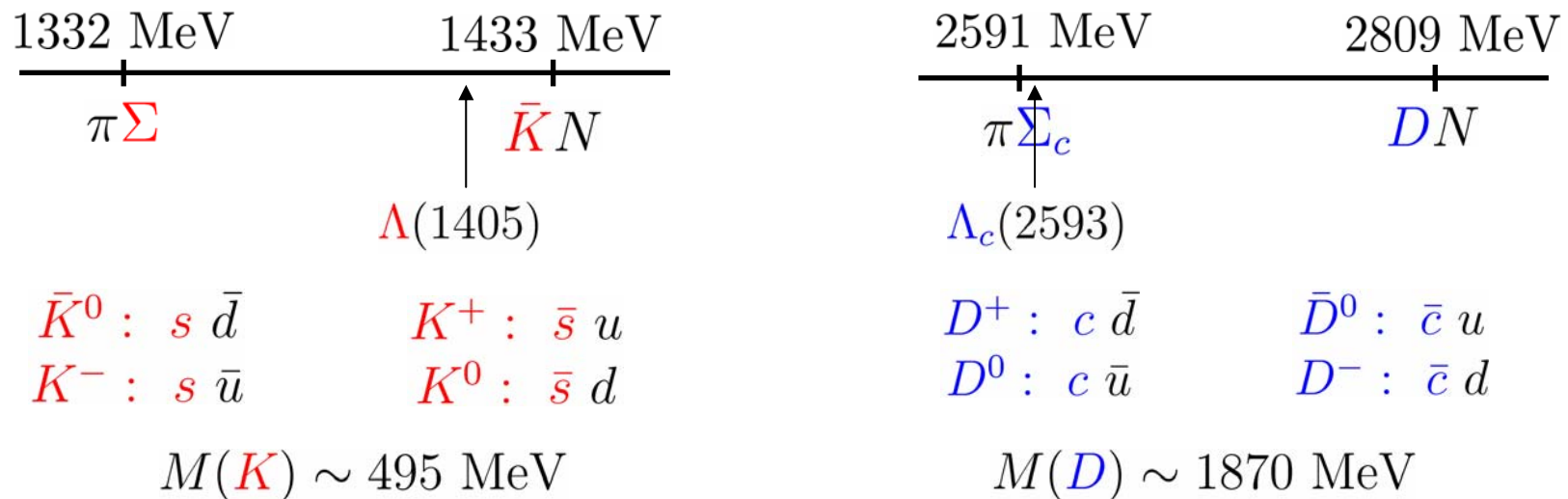
To understand these issues it is imperative to have a realistic picture of the properties of **charmonia** and **open charm mesons** in a **hot nuclear environment**.

This is especially important for the conditions of the **CBM@FAIR** heavy ion experiment, where both density ($\rho \sim 1-2 \text{ fm}^{-3}$) and temperature ($T \sim 100 \text{ MeV}$) are relevant parameters.



Motivation II (theoretical issues)

1. **DN interaction:** similar features as the thoroughly studied $K\bar{N}$ interaction. In the charm sector we also find a subthreshold $l=0$ resonance, the $\Lambda_c(2593)$ (udc), that bears a strong resemblance to the $\Lambda(1405)$ (uds).



May the $\Lambda_c(2593)$ be generated also dynamically, through a coupled-channel unitarized model?

2. Mean-field models: → predict the D^+, D^- mass shift

QCD sum rule: The in-medium mass shift is proportional to the mass of the charmed quark (m_c) and the light meson q - q bar condensate:

A. Hayashigashi, Phys. Let. B487, 96 (2000)

P. Morath, W. Weise, S.H. Lee, 17 Autumn school on QCD, → $\delta M_D \sim \delta M_{\bar{D}} \sim -50$ MeV
Lisbon 1999 (World Scientific, Singapore, 2001) 2001

Quark Meson Coupling approach: Hadron interactions mediated by the exchange of scalar-isoscalar (σ) and vector (ρ and ω) medium modified mesons among the light constituent quarks.

A.Sibirtsev, K.Tsushima, and A.W.Thomas,
Eur. Phys. J. A6, 351 (1999)

$$\delta M_{D^+} = U_s + U_v \sim -140 \text{ MeV}$$

$$\delta M_{D^-} = U_s - U_v \sim +25 \text{ MeV}$$

Nuclear Mean Field approach (NMFA): σ - ω model supplemented by **chiral terms** (of vector and scalar-isoscalar nature)

A.Mishra, E.L. Brakovskaya, J. Schaffner-Bielich,
S. Schramm, and H. Stoecker, PRC 70, 044904 (2004)

$$\delta M_{D^+} = -250 \text{ to } -70 \text{ MeV}$$

$$\delta M_{D^-} = -100 \text{ to } -20 \text{ MeV}$$

3. Spectral function models: → predict the full energy distribution of the D^+ meson strength in a self-consistent coupled channel approach

D self-energy with a SU(3) separable DN interaction with **u, d, c** content

L. Tolós, J. Schaffner-Bielich, and A. Mishra, *Phys. Rev.C* 70, 025203 (2004) (T=0 MeV)
L. Tolós, J. Schaffner-Bielich, and H. Stöcker, *Phys. Lett. B* 635, 85 (2006) (finite T)

D and Dbar self-energy with an improved SU(4) and chiral DN interaction

J.Hofmann and M.F.M.Lutz, *Nucl. Phys. A* 763, 90 (2005)
M.F.M.Lutz, and C.L.Korpa, *Phys. Lett. B* 633,43 (2006)

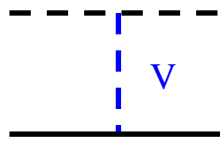
D self-energy using a revised SU(4) and chiral DN interaction, supplemented by an attractive scalar-isoscalar Σ_{DN}

T. Mizutani, A. Ramos, *Phys. Rev. C* 74, 065201 (2006)

HERE: we extend the model to **Dbar** mesons and implement **finite T** effects

Model:

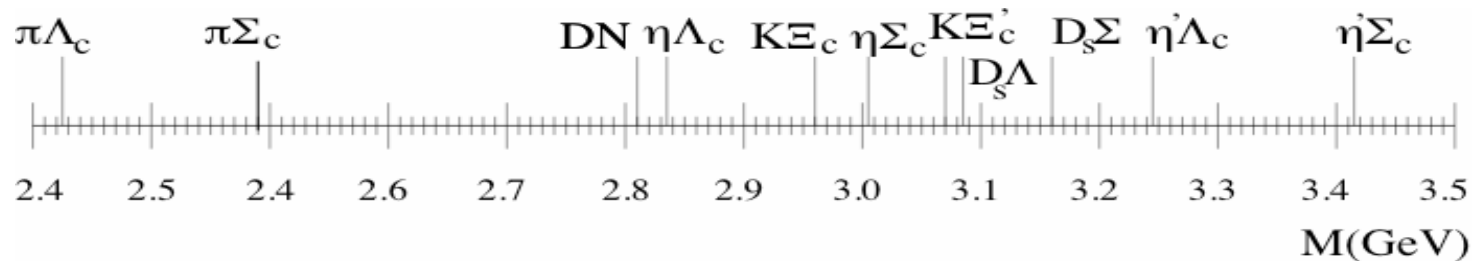
1. Transition potential V built from the meson-baryon Lagrangian at lowest order



$$V_{ij} = -\kappa C_{ij} \frac{1}{4f^2} (2\sqrt{s} - M_i - M_j) \left(\frac{M_i + E}{2M_i} \right)^{1/2} \left(\frac{M_j + E'}{2M_j} \right)^{1/2}$$

SU(4) symmetry broken by the use of physical masses.

$$\begin{aligned} \kappa &= 1 && \text{(non-charm exchange)} && DN \rightarrow DN, D_s Y \\ &= \left(\frac{m_\rho}{m_D} \right)^2 \sim 1/4 && \text{(charm exchange)} && DN \rightarrow \pi \Sigma_c, K \Xi_c \end{aligned}$$



\mathbf{V} is also supplemented by a scalar-isoscalar interaction (Σ_{DN} term)

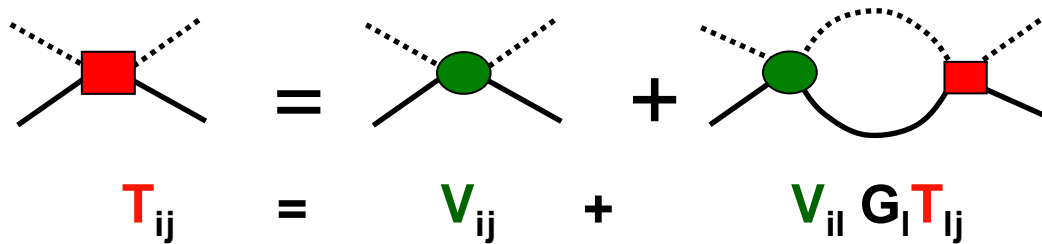
$$\mathcal{L}_\Sigma \equiv \frac{\Sigma_{DN}}{f_D^2} \bar{N} N \bar{D} D \longrightarrow V_\Sigma(\sqrt{s}) = -\frac{\Sigma_{DN}}{f_D^2} \left(\frac{M_N + E}{2M_N} \right) \sim -0.05 \text{ MeV}^{-1}$$

$$f_D \sim 200 \text{ MeV}$$

$$\Sigma_{DN} \sim 2000 \text{ MeV} \quad (\text{from QCDSR})$$

2. Unitarization: N/D method

→ equivalent to Bethe-Salpeter coupled-channel equations with on-shell amplitudes



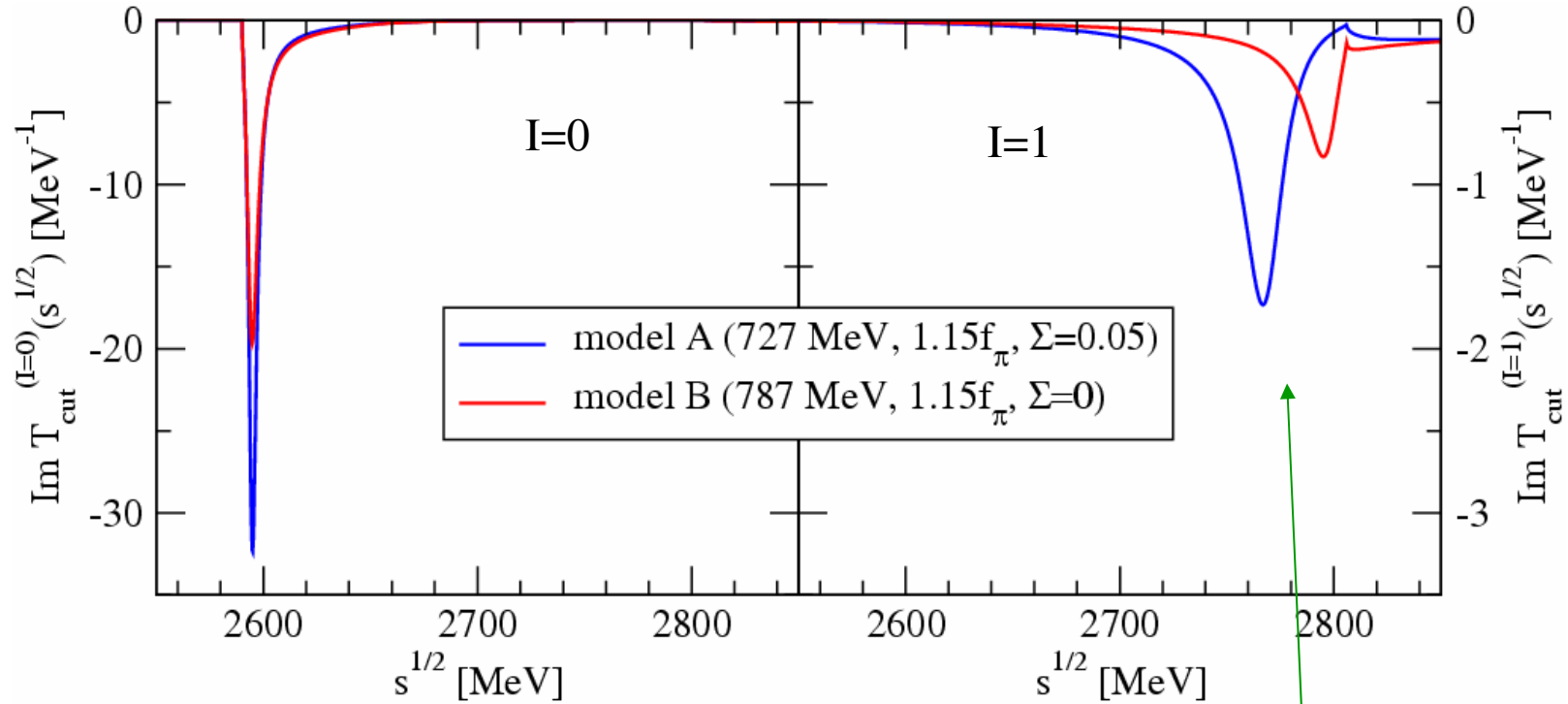
$$\mathbf{T}_{ij} = \mathbf{V}_{ij} + \mathbf{V}_{il} \mathbf{G}_l \mathbf{T}_{lj}$$

The loop function \mathbf{G} is regularized with a cut-off Λ [adjusted to reproduce $\Lambda_c(2593)$]

Model A: $\Lambda = 727 \text{ MeV}, f = 1.15 f_\pi, V_\Sigma = -0.05 \text{ MeV}^{-1}$

Model B: $\Lambda = 787 \text{ MeV}, f = 1.15 f_\pi, V_\Sigma = 0$

Free space DN amplitudes

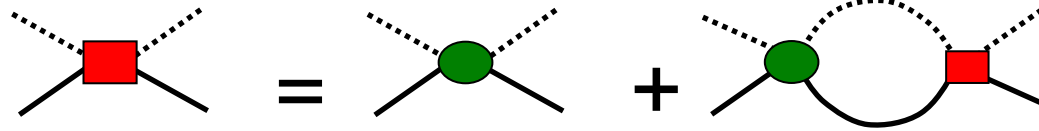


R. Mizuk et al. [Belle Collaboration]
 Phys.Rev.Lett.94, 122002(2005)
 $\Sigma_c(2800)$, $\Gamma \sim 60 \text{ MeV}$

The model generates the $I=0$ $\Lambda_c(2595)$ and another resonance in $I=1$ around the nominal $\Sigma_c(2800)$!

Self-consistent coupled channels procedure

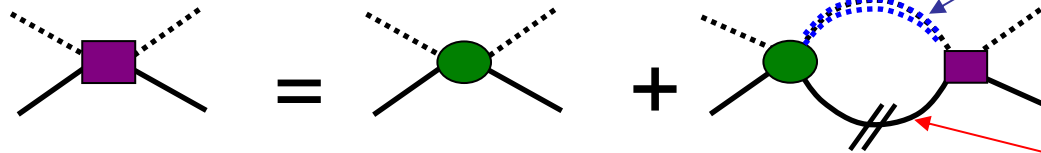
Free space



$$T_{ij} = V_{ij} + V_{il} G_l T_{lj}$$

meson dressing

Medium



$$T_{ij}(\rho) = V_{ij} + V_{il} G_l(\rho) T_{lj}(\rho)$$

Pauli blocking and baryon dressing

Dressed D meson:



Loop function (depends on ρ and T)

$$n(k) = \frac{1}{1 + e^{(\varepsilon(k) - \mu)/T}}$$

$$\tilde{G}_{DN}(P_0, \vec{P}, \rho) = \int_{|\vec{q}| < \Lambda} \frac{d^3q}{(2\pi)^3} \frac{M_N}{E_N(\vec{P} - \vec{q})} \times \left[\int_0^\infty d\omega S_D(\omega, \vec{q}) \frac{1 - n(\vec{P} - \vec{q})}{P_0 - \omega - E_N(\vec{P} - \vec{q}) + i\varepsilon} + \int_0^\infty d\omega S_D(\omega, \vec{q}) \frac{n(\vec{P} - \vec{q})}{P_0 + \omega - E_N(\vec{P} - \vec{q}) - i\varepsilon} \right]$$

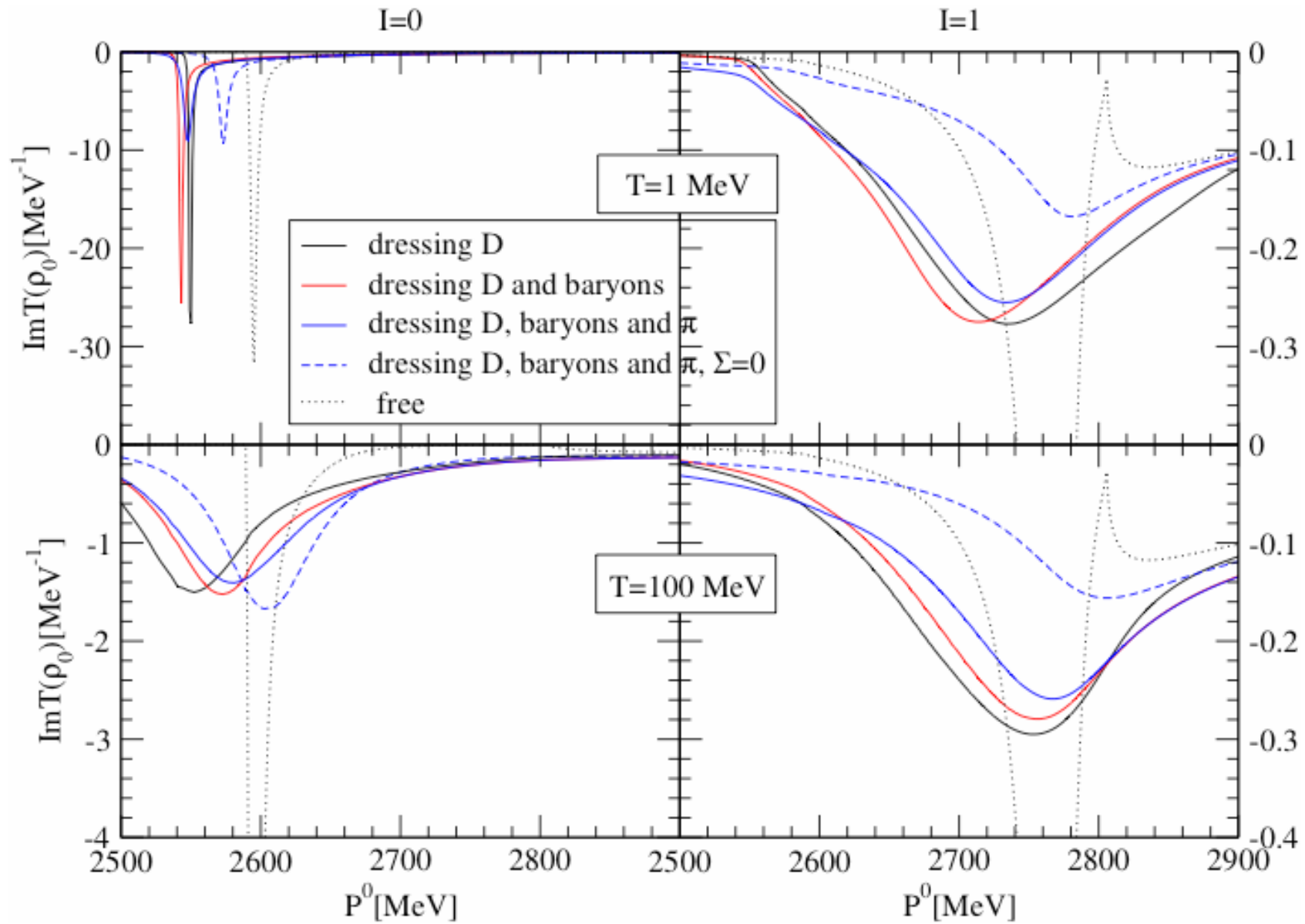
spectral density

$$S_D(q_0, \vec{q}) = -\frac{1}{\pi} \text{Im} D_D(q_0, \vec{q}) = -\frac{1}{\pi} \frac{\text{Im} \Pi_D(q_0, \vec{q})}{|q_0^2 - \vec{q}^2 - m_D^2 - \Pi_D(q_0, \vec{q})|^2}$$

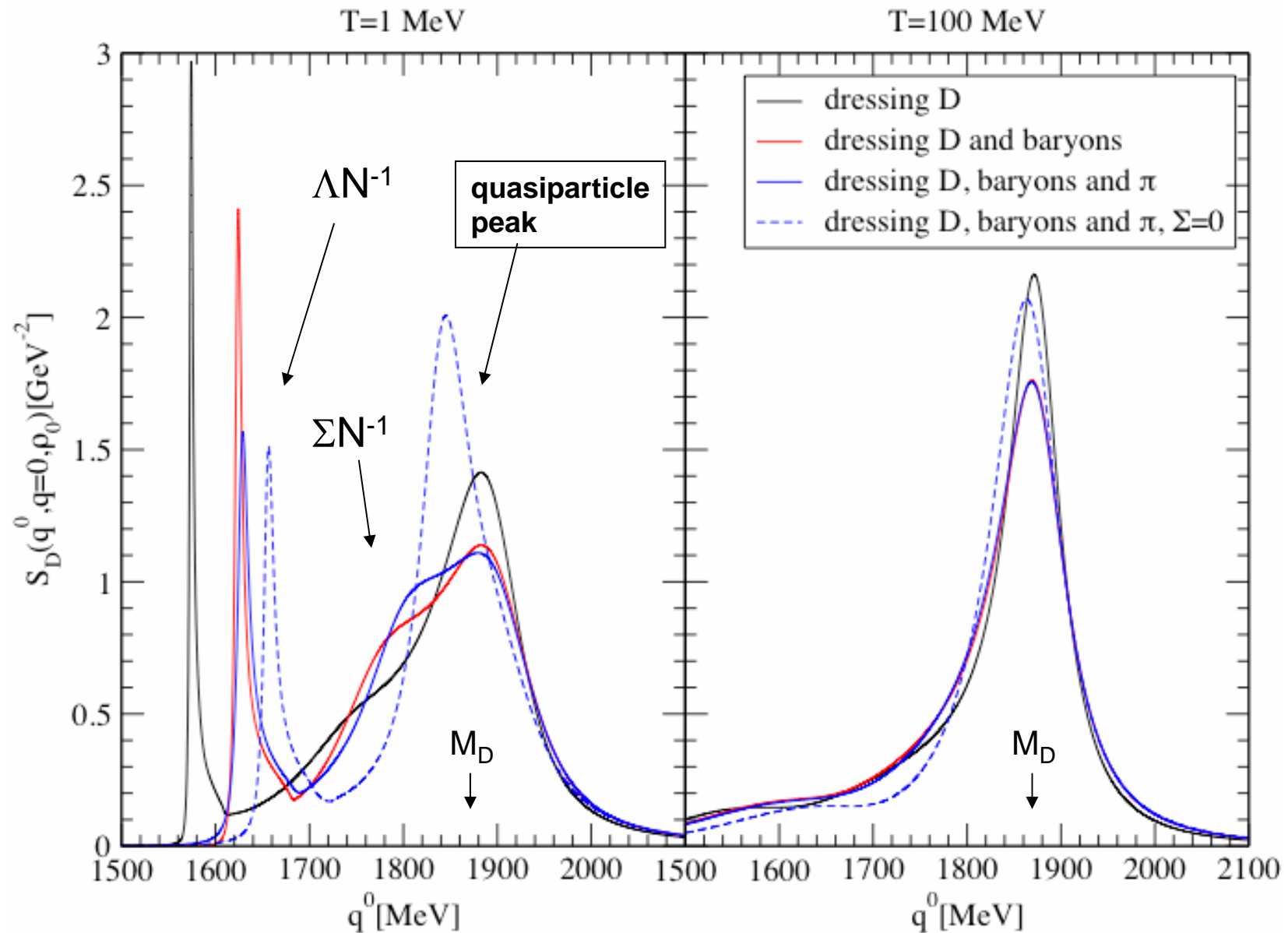
D-meson self-energy

$$\Pi_D(q_0, \vec{q}) = \int \frac{d^3p}{(2\pi)^3} n(\vec{p}) [\tilde{T}_{DN}^{(I=0)}(P_0, \vec{P}) + 3\tilde{T}_{DN}^{(I=1)}(P_0, \vec{P})]$$

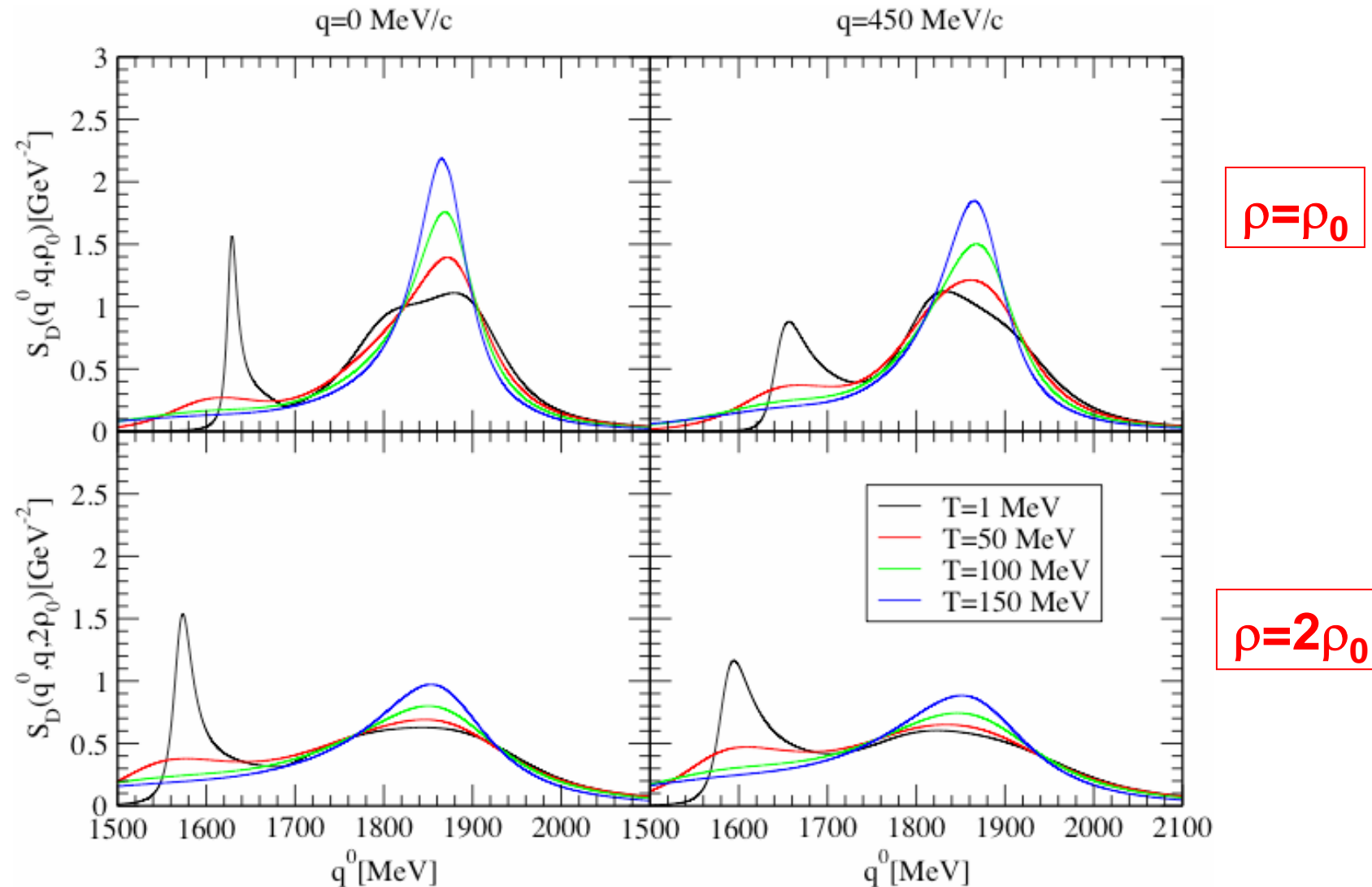
DN amplitudes in hot dense matter ($\rho=\rho_0$)



D meson spectral in hot and dense matter ($\rho=\rho_0$)



Evolution with density and temperature



Similar trend to previous finite temperature results:

LT, J. Schaffner-Bielich and H. Stoecker PLB 635 (2006) 85

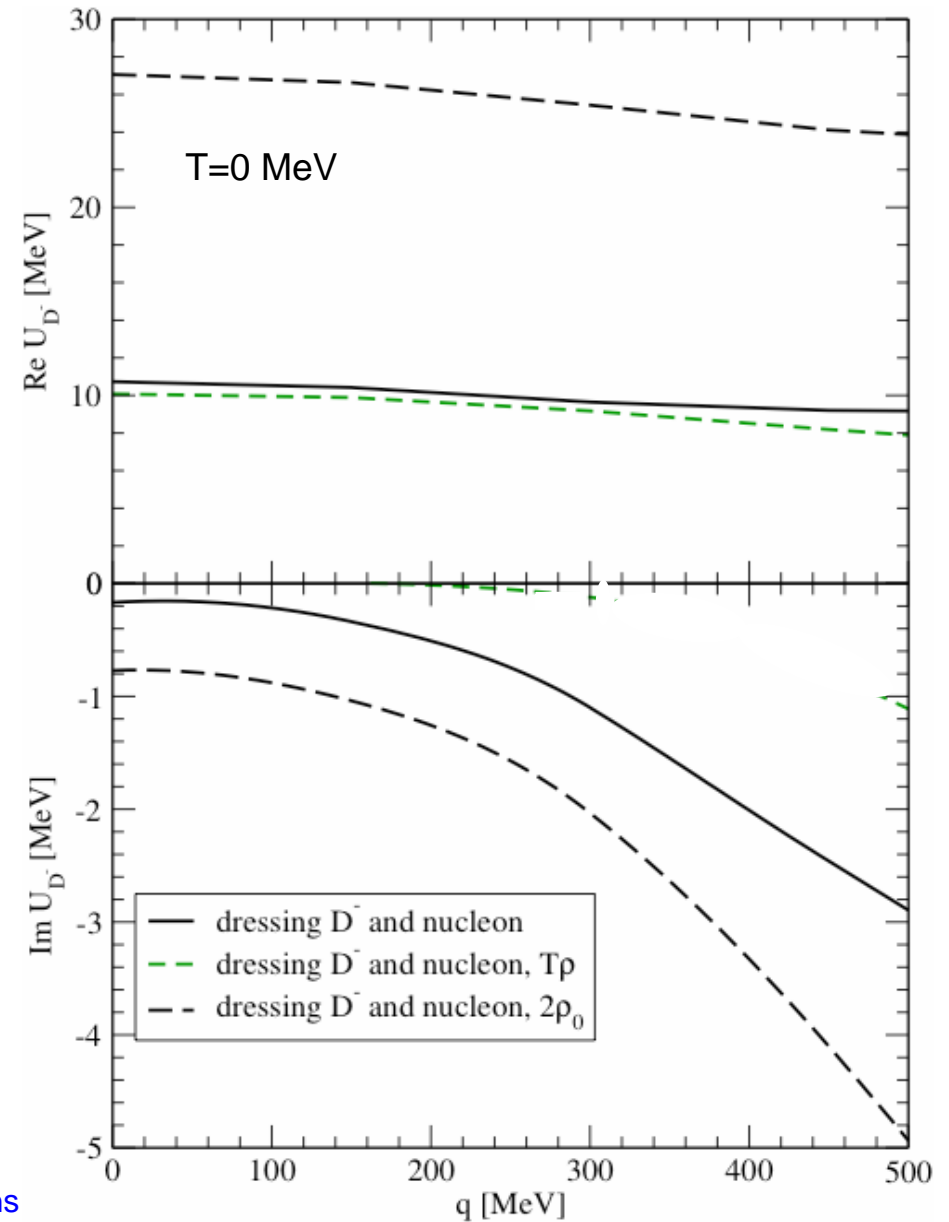
$\bar{D}N$ interaction

$$a_{\bar{D}N} = -\frac{1}{4\pi} \frac{M_{\bar{D}N}}{\sqrt{s}} T_{\bar{D}N \rightarrow \bar{D}N}$$

Table 1: $\bar{D}N$ scattering lengths (fm)

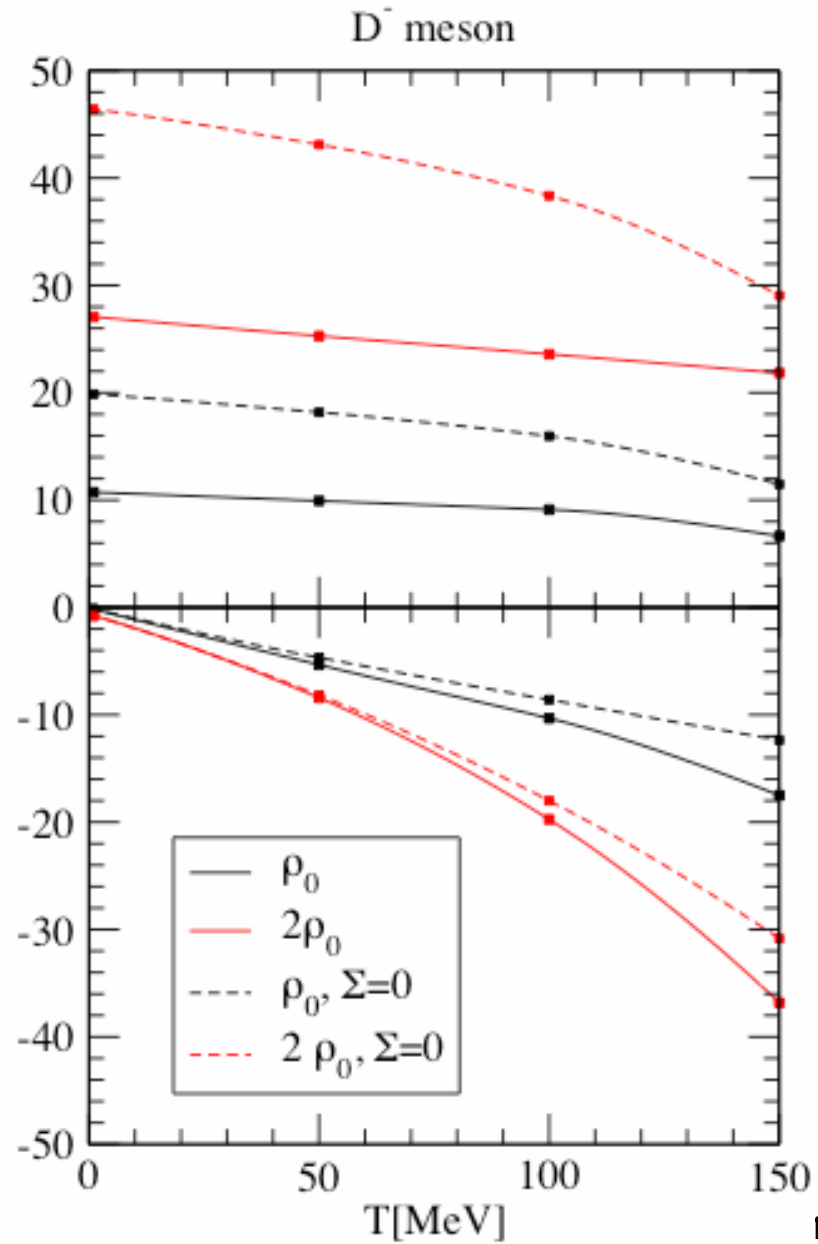
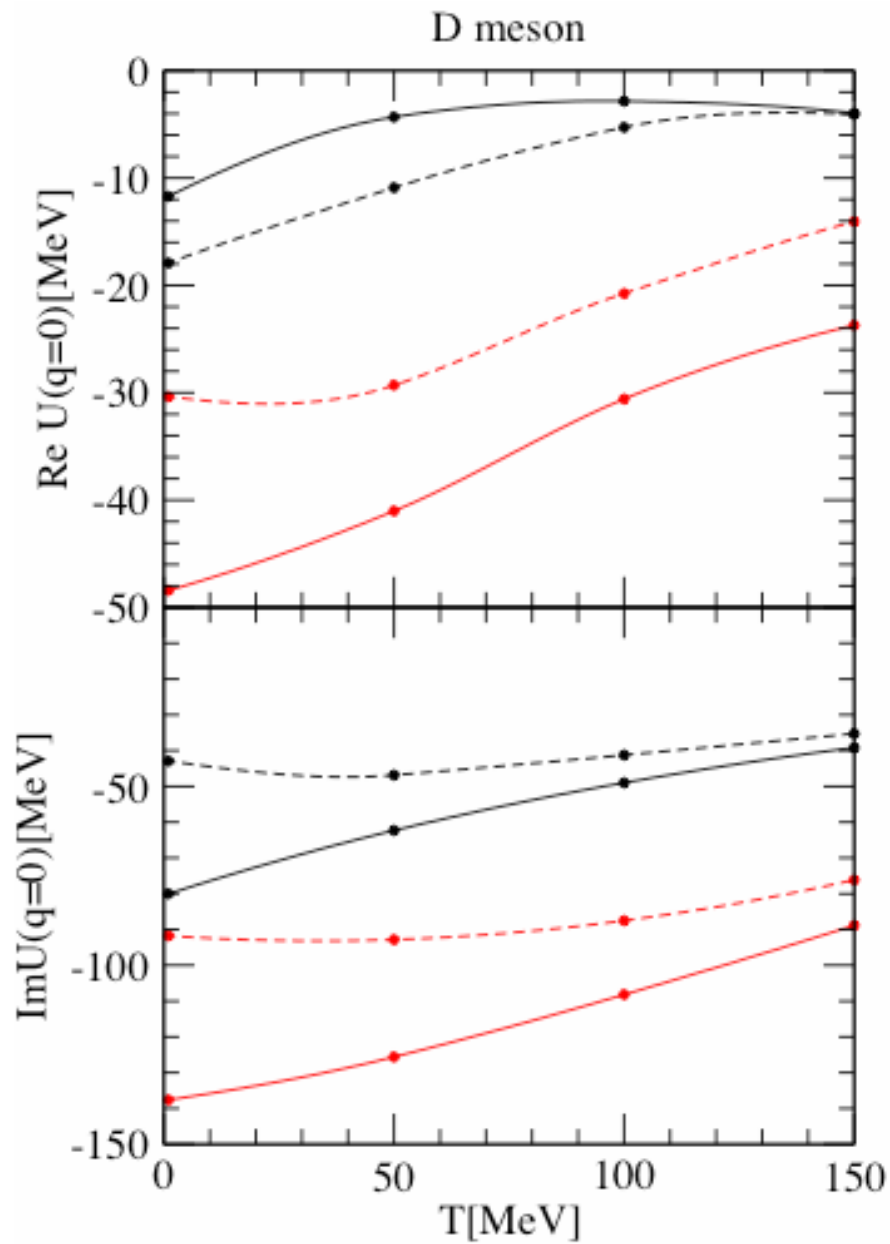
	Model A	Model B
$I = 0$	0.607	0
(Born approx.)	0.262	0
$I = 1$	-0.264	-0.289
(Born approx.)	-0.614	-0.876

\bar{D} optical potential



D-mesons

D and \bar{D} meson potentials



Conclusions & Outlook

We have performed a self-consistent coupled-channel calculation of the D and \bar{D} self-energies in symmetric nuclear matter at finite temperature taking, as bare interaction, the $SU(4)$ TW contribution supplemented by a Σ_{DN} term

✓ In hot dense matter, $\Lambda_c(2593)$ and $\Sigma_c(2800)$ stay close to their free position but develop a remarkable width

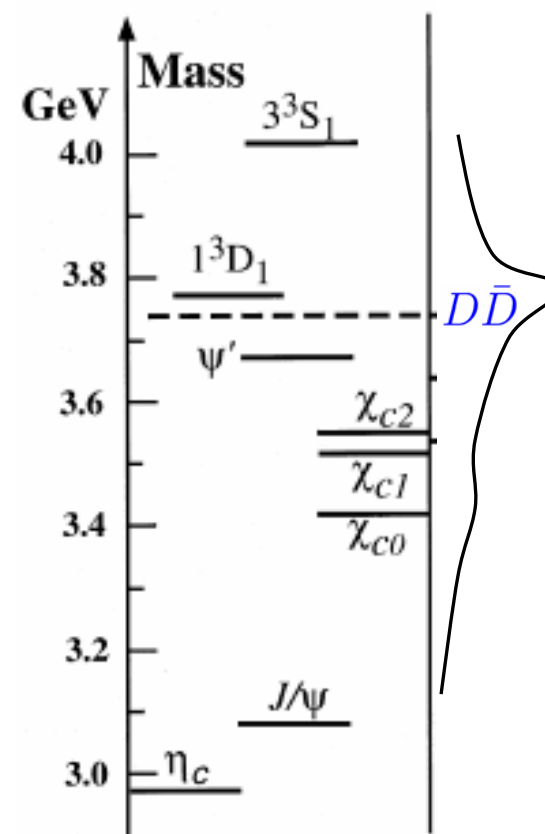
✓ The D meson spectral density shows a single pronounced peak at finite temperature that melts with increasing density

✓ Up to ρ_0 the low-density theorem is a good approximation for the DN , where the repulsive $I=1$ component dominates

✓ Due to the distinct resonant structure of the interaction, temperature induces a stronger change in the properties of the D than the \bar{D} meson

Open questions?

- J/Ψ suppression
- Open-charm enhancement
- D-mesic nuclei



Some answers expected from CBM@FAIR !