Density Fluctuations as a Signal for First Order Phase Transitions

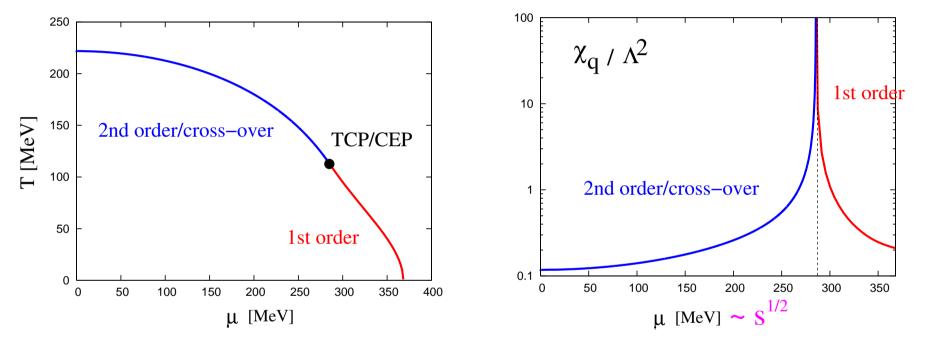
Chihiro Sasaki (GSI)

based on C. Sasaki, B. Friman and K. Redlich, hep-ph/0702254

QCD phase structure and fluctuations

- conserved charge fluctuations : critical behaviors
- net quark number/electric charge susceptibilities : search for the CEP

NJL model calculation : C. S., B. Friman, K. Redlich, Phys. Rev. D, 2007



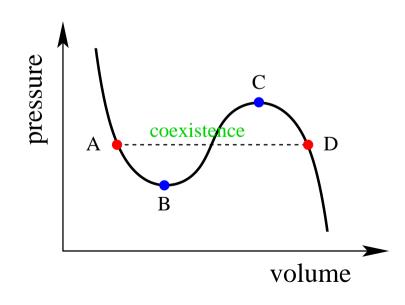
 $\chi_q \rightarrow \infty$ at CEP while $\chi_q \sim$ finite for the 1st order transition in equilibrium

• heavy-ion collisions : deviation from equilibrium, instabilities

The nature of first order phase transition

• change of thermodynamic potential from broken to symmetric phase

stability of a system



 $\partial P/\partial V < 0$: stable $\partial P/\partial V > 0$: unstable $\partial P/\partial V = 0$: spinodals

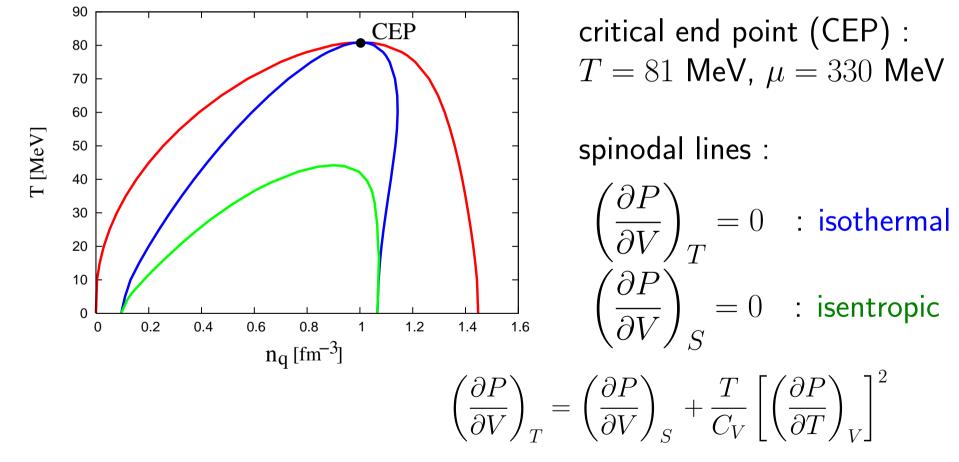
A-B : supercooling (symmetric phase)B-C : non-equilibrium stateC-D : superheating (broken phase)

Phase diagram in the Nambu–Jona-Lasinio model

• NJL model with two flavors

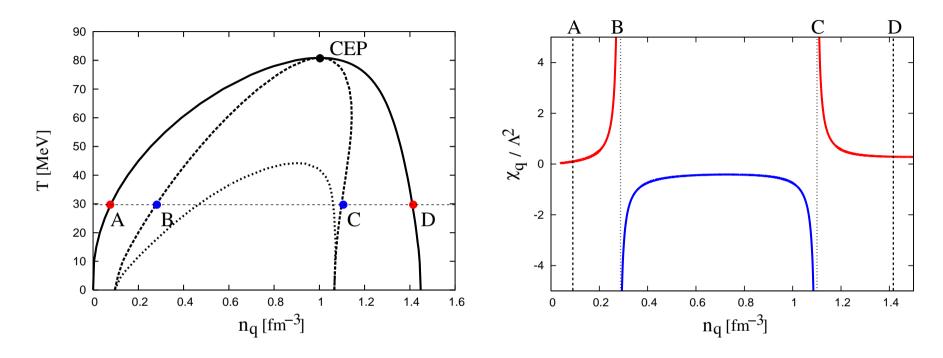
$$\mathcal{L} = \bar{\psi}(i\partial \!\!\!/ - m)\psi + \bar{\psi}\mu\gamma_0\psi + G_S \Big[(\bar{\psi}\psi)^2 + (\bar{\psi}i\vec{\tau}\gamma_5\psi)^2 \Big]$$
$$m = 5.6 \,\mathrm{MeV} \,, \, G_S \Lambda^2 = 2.44 \,, \, \Lambda = 587.9 \,\mathrm{MeV}$$

• phase diagram



Quark number susceptibility

• deviation from equilibrium, large fluctuations induced by instabilities



- at 1st order transition point (A, D) : χ_q is finite
- at isothermal spinodal point (B, C) : χ_q diverges and changes its sign $\frac{\partial P}{\partial V} < 0$ for stable/meta-stable state $\Rightarrow \frac{\partial P}{\partial V} > 0$ for unstable state
- in unstable region (B-C) : χ_q is finite and negative

• divergence of χ_q :

$$\left(\frac{\partial P}{\partial V}\right)_T = -\frac{n_q^2}{V}\frac{1}{\chi_q} = 0 \quad \text{at any spinodal points}$$

 $\Rightarrow~\chi_q$ diverges along the isothermal spinodal lines

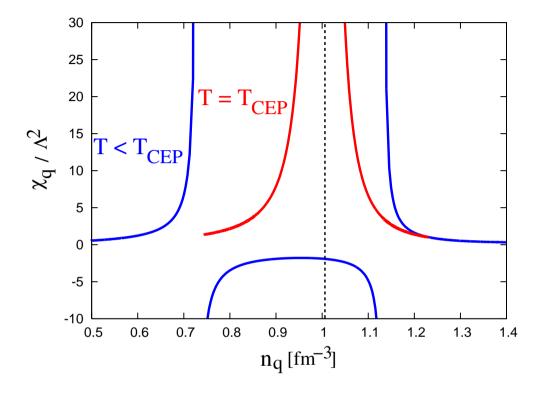
• electric charge susceptibility also diverges:

$$\chi_Q = \frac{1}{36}\chi_q + \frac{1}{4}\chi_I + \frac{1}{6}\frac{\partial^2 P}{\partial \mu_q \partial \mu_I}$$

 spinodal decomposition for the chiral/deconfinement phase transition in heavy-ion collisions :

instabilities, enhancement of baryon and strangeness fluctuations

Heiselberg, Pethick and Ravenhall (1988); Bower and Gavin (2001); Chomaz, Colonna and Randrup (2004); Koch, Majumder and Randrup (2005) • toward the critical end point



- two positive branches are approaching
- instability region shrinks
- strength of divergence remains the same because of the same critical exponents :

$$\chi_q \sim (\mu - \mu_c)^{-\gamma},$$

$$\gamma = 1/2 \ (m_q = 0)$$

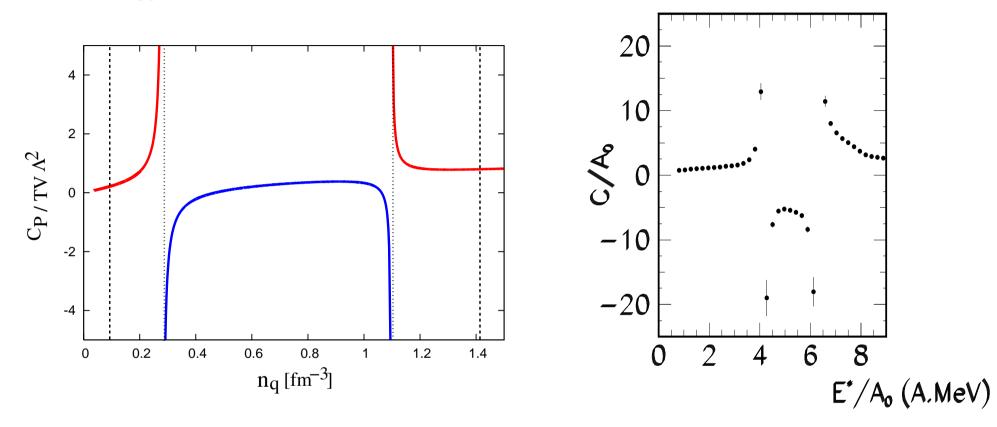
$$\gamma = 2/3 \ (m_q \neq 0)$$

• specific heat for constant pressure

$$C_P = T\left(\frac{\partial S}{\partial T}\right)_P = TV\left[\frac{\chi_{TT} - \frac{2s}{n_q}\chi_{\mu T} + \left(\frac{s}{n_q}\right)^2\chi_q\right]$$

Experimental evidence

• low-energy nuclear collisions



M. D'Agostino et al., Phys. Lett. B 473, 219 (2000)

negative heat capacity : anomalously large fluctuations \Rightarrow an evidence of the 1st order liquid-gas phase transition

Summary and conclusions

• The quark number susceptibility diverges if spinodal phase separation occurs.

cf. finite in the equilibrium transition

- a signal not only for the CEP but also for the 1st order phase transitions
 ⇒ large fluctuations will be seen in a wider range of the phase diagram.
- \bullet CBM energy \sim the 1st order phase transition large fluctuations of baryon number, electric charge, proton number will be expected.

