

Overview of ALICE results on antinuclei and anti-hypernuclei

Maximiliano Puccio

Università & INFN Torino

for the ALICE Collaboration

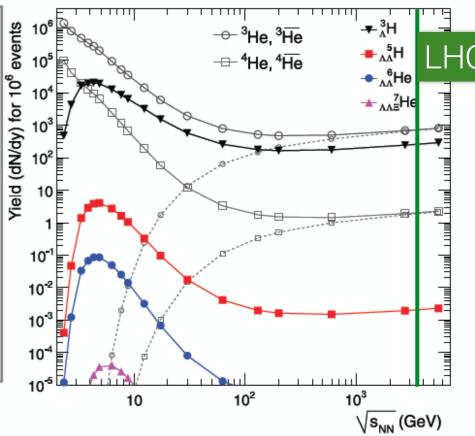
Nuclei production: theoretical approaches

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Thermal models

- Hadrons emitted from the interaction region in statistical equilibrium when the fireball reaches limiting temperature
- Abundances fixed at chemical freeze-out
- Freeze-out temperature T_{chem} is a key parameter
- Abundance of a species $\propto \exp(-m/T_{\text{chem}})$:
 - → For nuclei (large *m*) strong dependence on *T*_{chem}

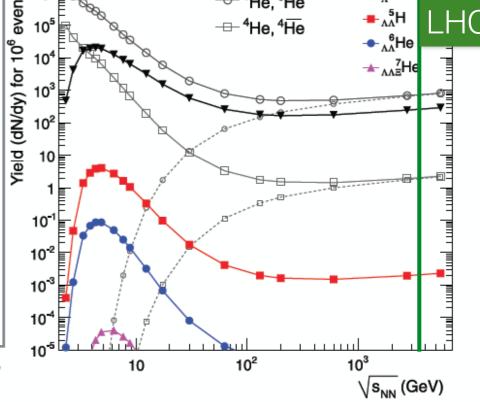
A. Andronic, P. Braun-Munzinger, J. Stachel and H. Stoecker, Phys. Lett. B607, 203 (2011), 1010.2995



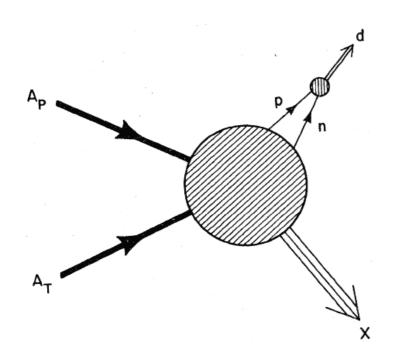
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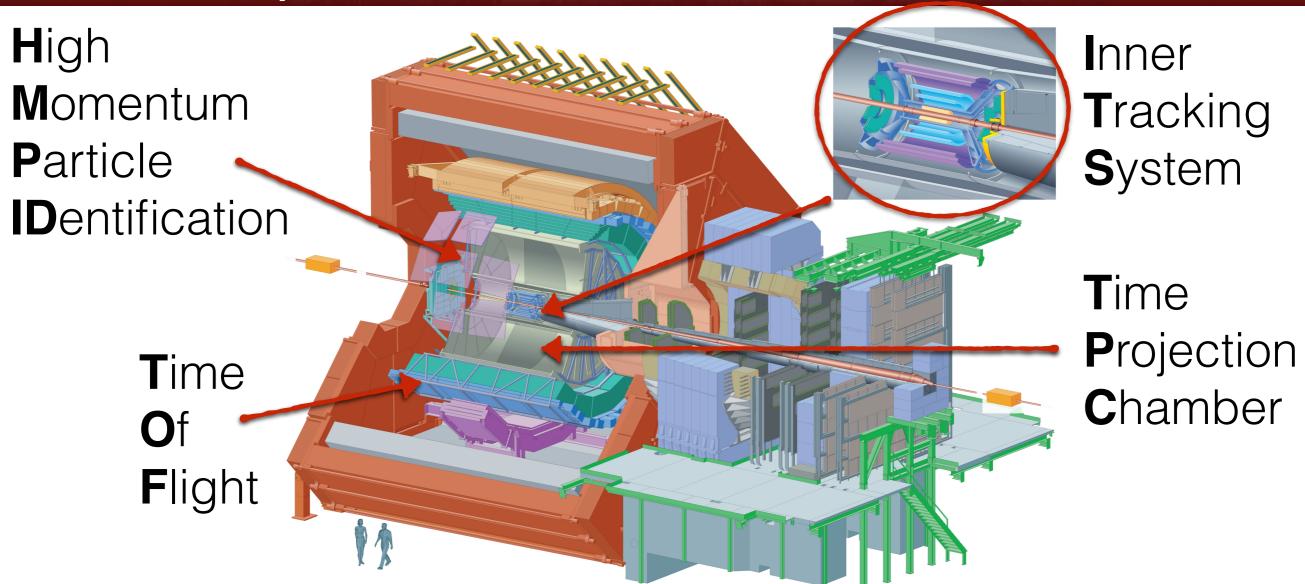


Coalescence models

- If (anti-)baryons are close in phase space after the kinetic freeze-out they can form a (anti-)nucleus
- (Anti-)nuclei produced at the chemical freeze-out might break and re-form during the time between the chemical freeze-out and the kinetic freeze-out.

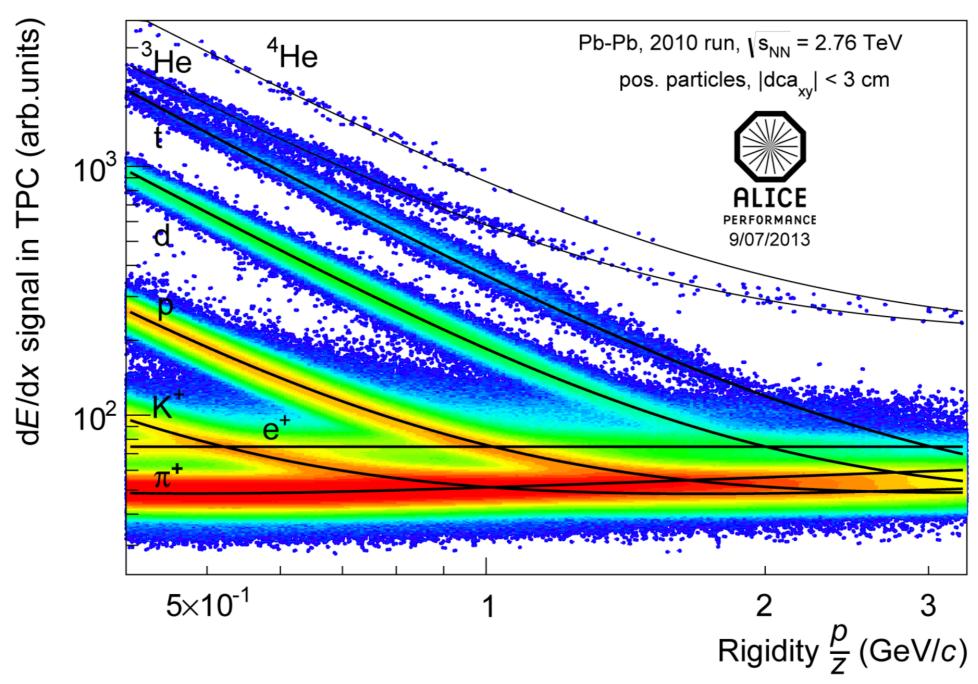
J. I. Kapusta, Phys.Rev. C21, 1301 (1980)

ALICE experiment



- General purpose heavy ion experiment
- Excellent particle identification (PID) capabilities and low material budget
- Most suited detector at the LHC to study the nuclei produced in the collisions

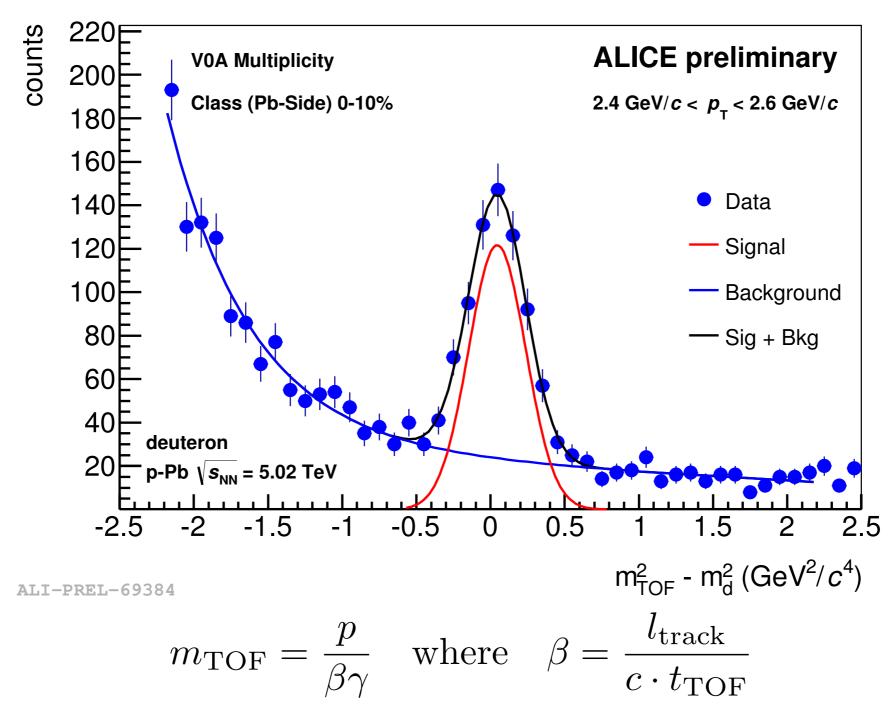
ALICE PID performance: TPC



Specific energy loss in the **T**ime **P**rojection **C**hamber volume provides an excellent PID for light nuclei.

 $\neg \sigma_{dE/dx} \sim 7\%$ (in Pb-Pb collisions)

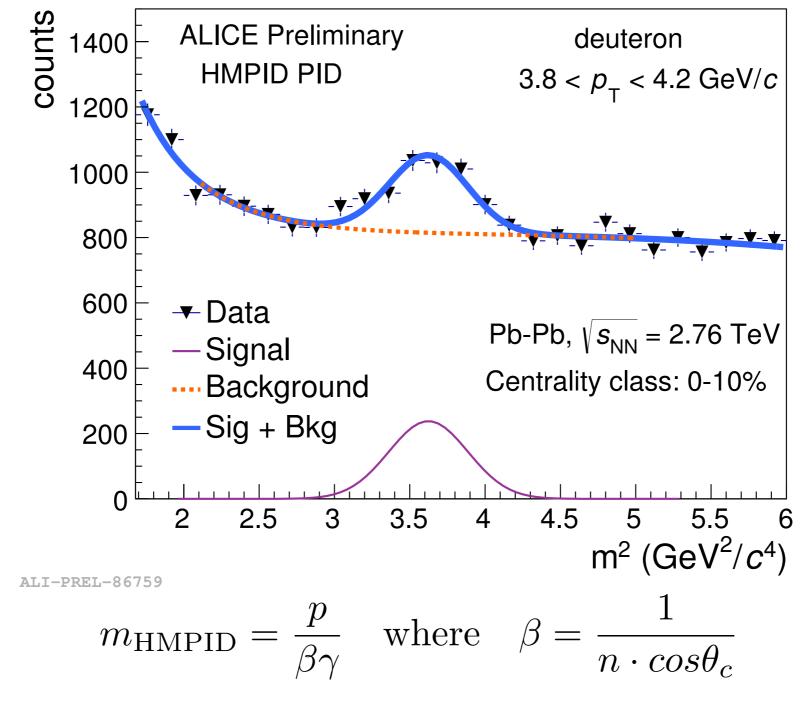
ALICE PID performance: TOF



Tracking information + time of flight measured by the **Time Of Flight** detector

- \rightarrow *m* of the particle
- → σ_{time-of-flight} ~ 85 ps (in Pb-Pb)

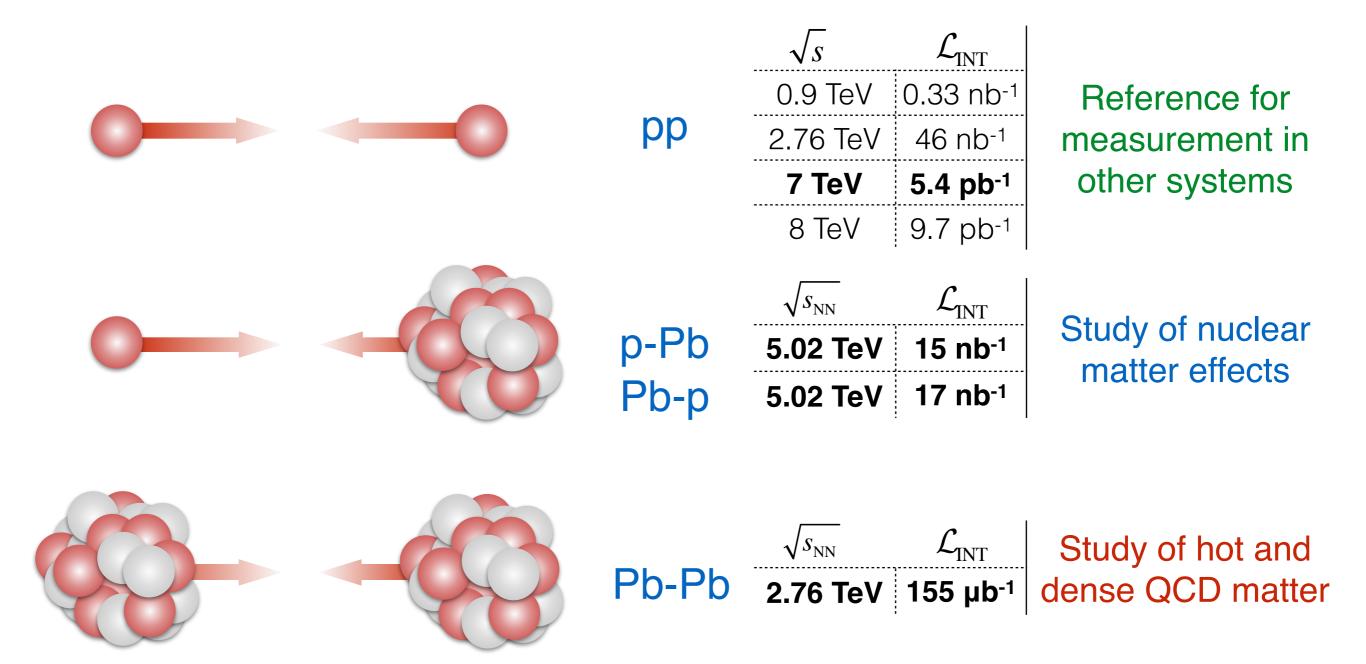
ALICE PID performance: HMPID



Tracking information + High Momentum Particle IDentification detector (Cherenkov light detector) signal

 \rightarrow m of the particle

Collision systems at the LHC



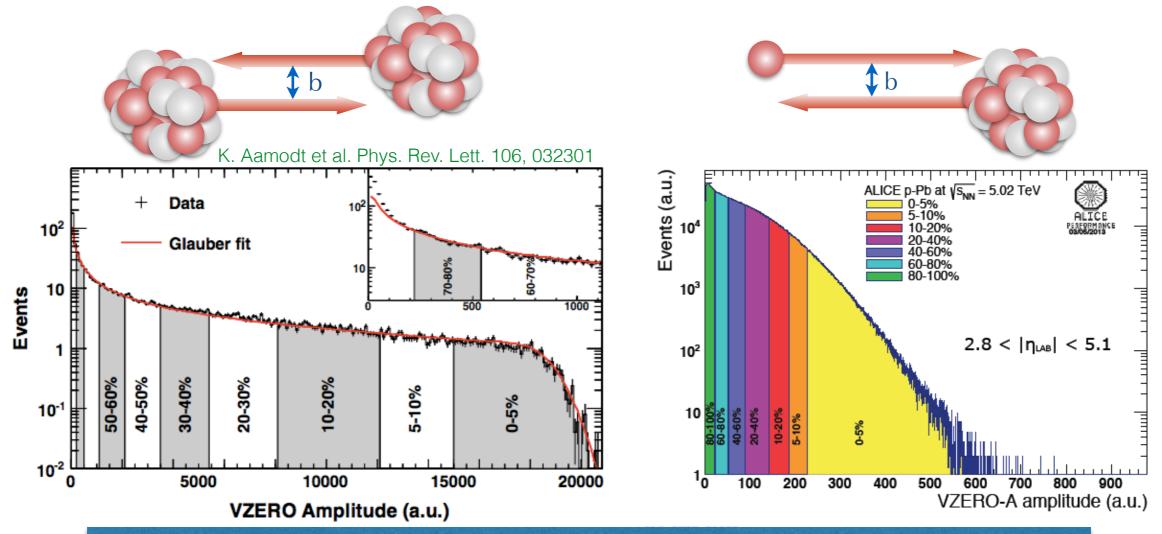
Three collision systems: unique opportunity to further study hadronisation and the strong interaction at extreme regimes of energy density

Centrality of a collision

The centrality of a collision is defined by the impact parameter *b*: *Most central collision* ⇔ *Smallest b*

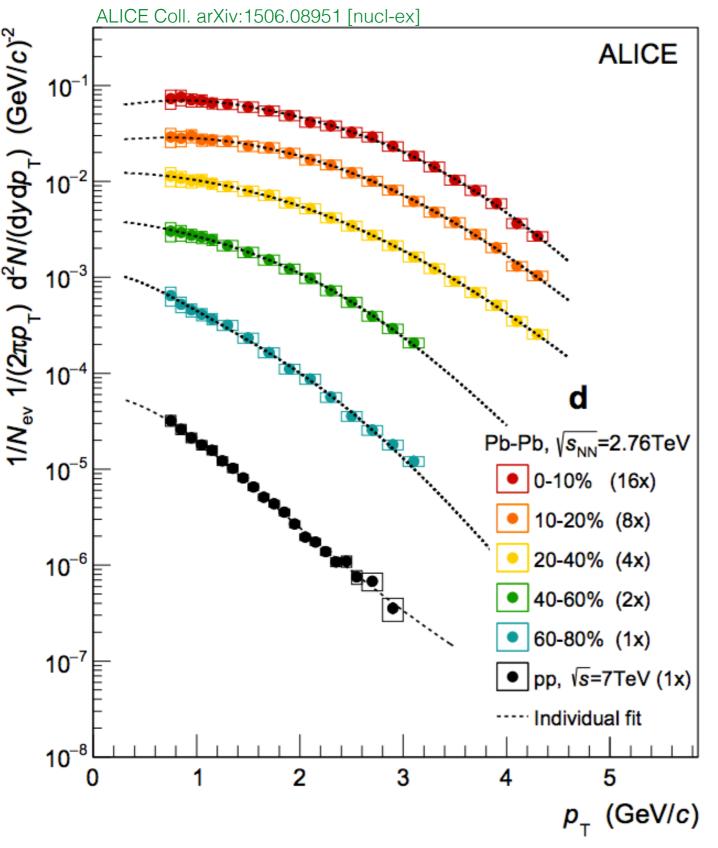
Experimentally it is possible to correlate the charged particle multiplicity to *b* by fitting data with the function shape predicted by the Glauber model.

The correlation between charged particle mult. and impact parameter in p-Pb is broader.



V-ZERO is a scintillator hodoscope used for centrality estimation and for the trigger.

Deuteron production in pp and Pb-Pb



pp

Invariant production spectrum is well fitted by the Levy-Tsallis function in pp

$$\frac{\mathrm{d}^2 N}{\mathrm{d}p_{\mathrm{T}}\mathrm{d}y} = p_{\mathrm{T}} \frac{\mathrm{d}N}{\mathrm{d}y} \frac{(n-1)(n-2)}{nC(nC+m_0(n-2))} \left(1 + \frac{m_{\mathrm{T}} - m_0}{nC}\right)^{-1}$$

where m_0 is the reference mass of the deuteron and n,C are fit parameters.

C. Tsallis. J.Statist.Phys. 52 479-487(1988)

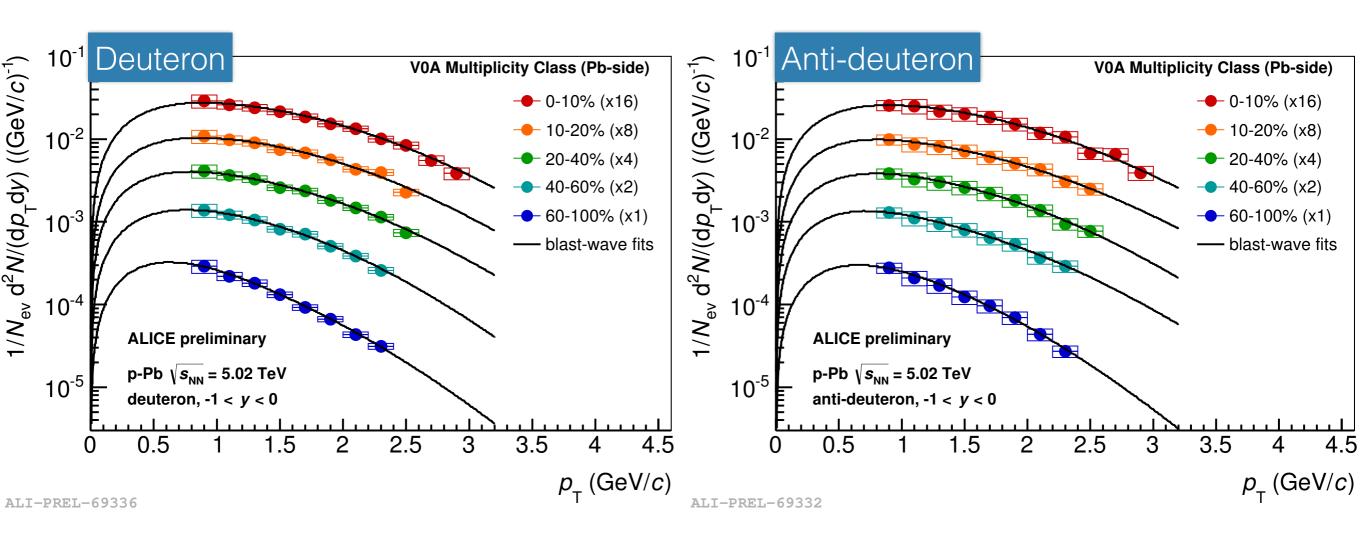
Pb-Pb

- The Blast-Wave (BW) function fits well the data.
- Characteristic hardening of the spectrum with increasing centrality.
- These fits are used for the extrapolation of the yield to the unmeasured region at low and high $p_{\rm T}$.

Blast wave model: E. Schnedermann et al., Phys. Rev. C48, 2462 (1993)

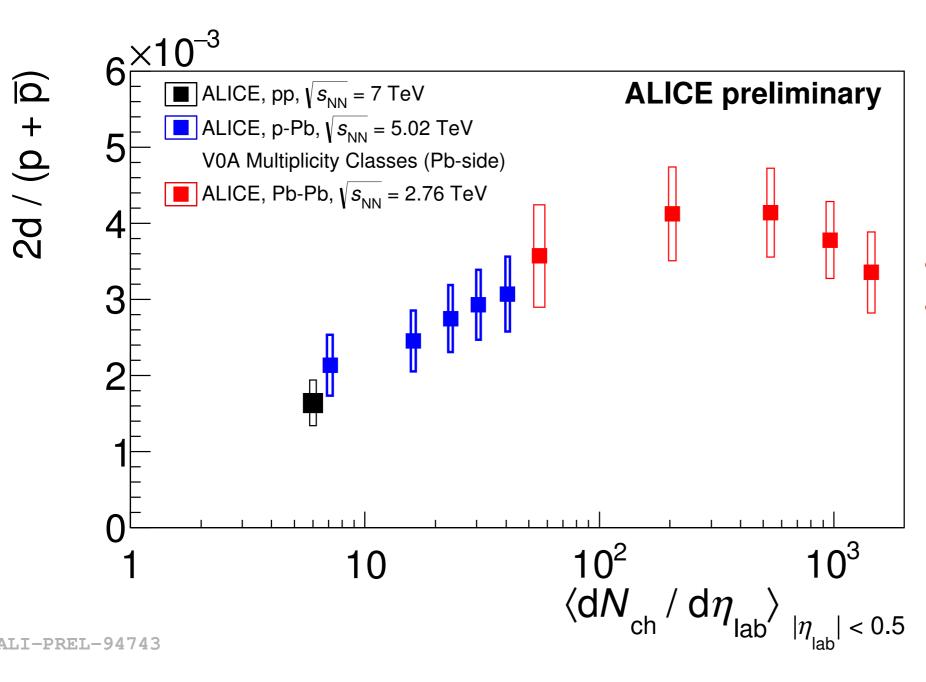
(Anti-)deuteron production in p-Pb

The Blast-Wave function fits well the data also in this case.



The spectra become harder with increasing multiplicity

d/pratio

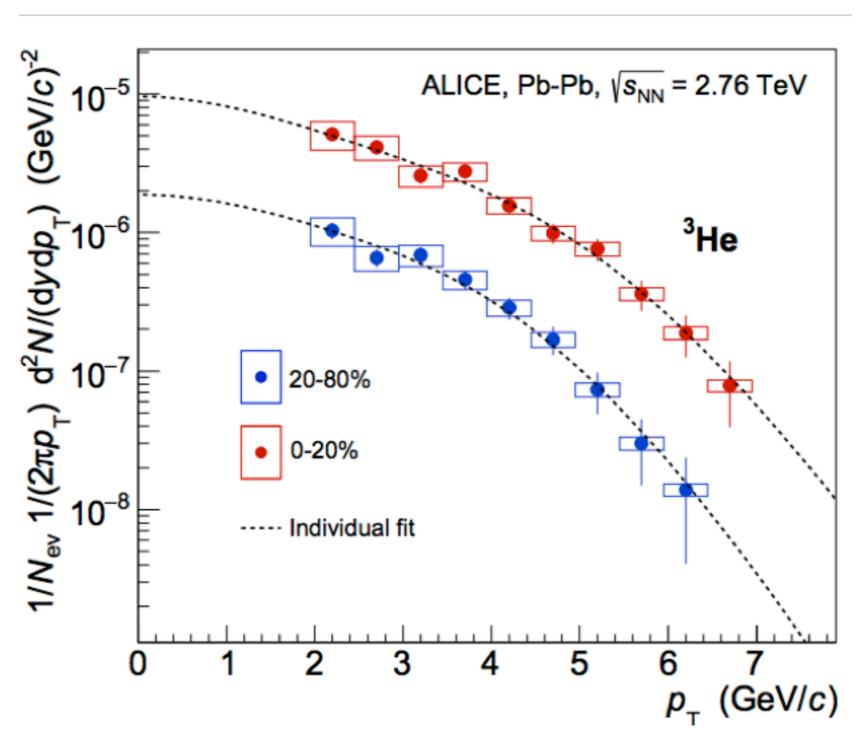


- Rise with multiplicity
- No further increase in Pb-Pb collisions within errors

The d/p ratio increases with the charged particle multiplicity: this is consistent with the coalescence picture

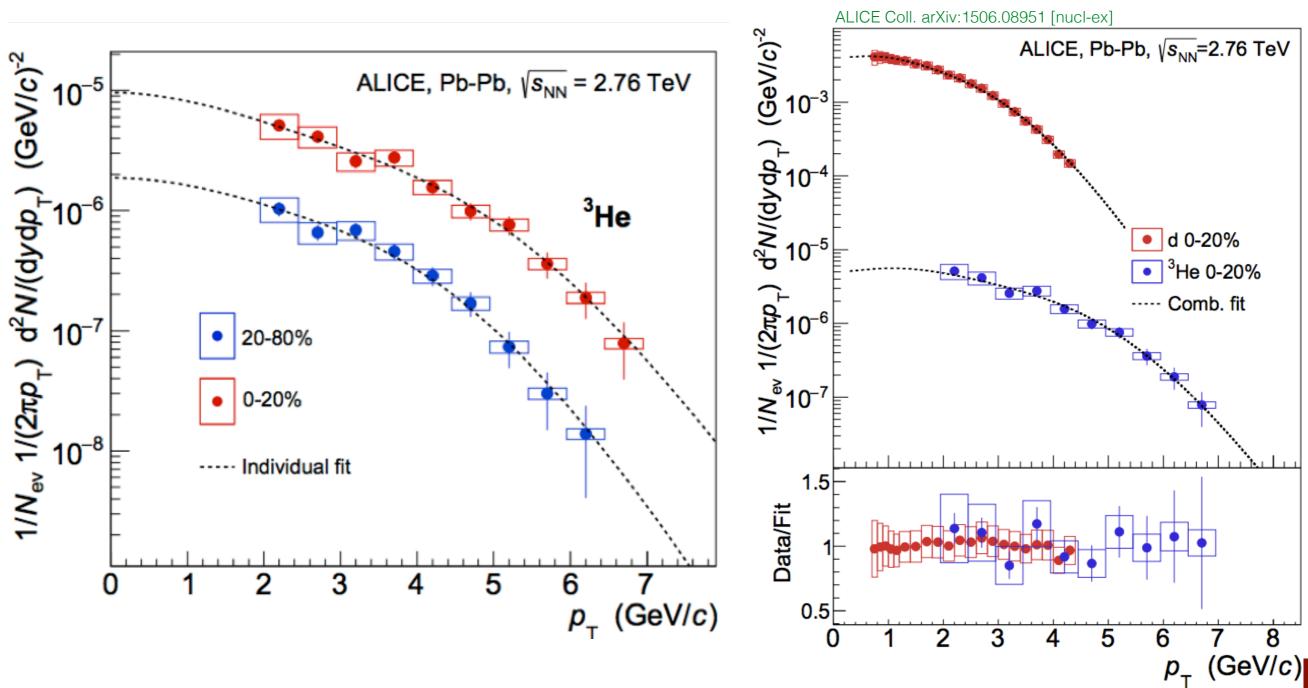
(Anti-)3He production in Pb-Pb

The individual Blast-Wave fits describe well the data.



(Anti-)3He production in Pb-Pb

It is possible to fit simultaneously deuteron and ${}^{3}\text{He}$ to extract common parameters. The simultaneous fit describes well the data and the common kinetic freeze-out temperature is $T_{kin}=83\pm22$ MeV



Combined BW fit

BW model fit:

- gives insight into the kinetic freeze-out conditions
- does not describe hard processes that contribute to particle production at high p_T

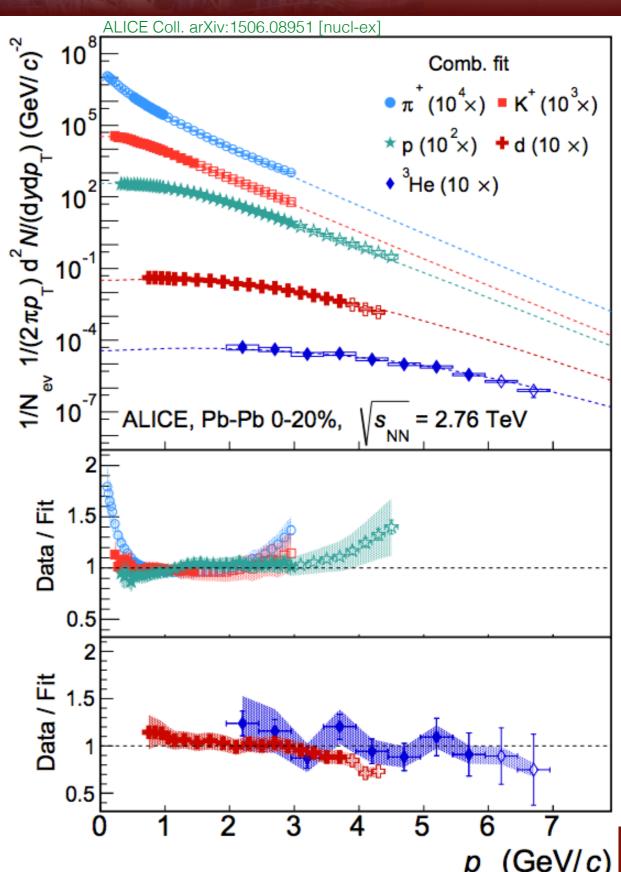
Fit parameters

$$<\beta> = 0.632 \pm 0.01$$

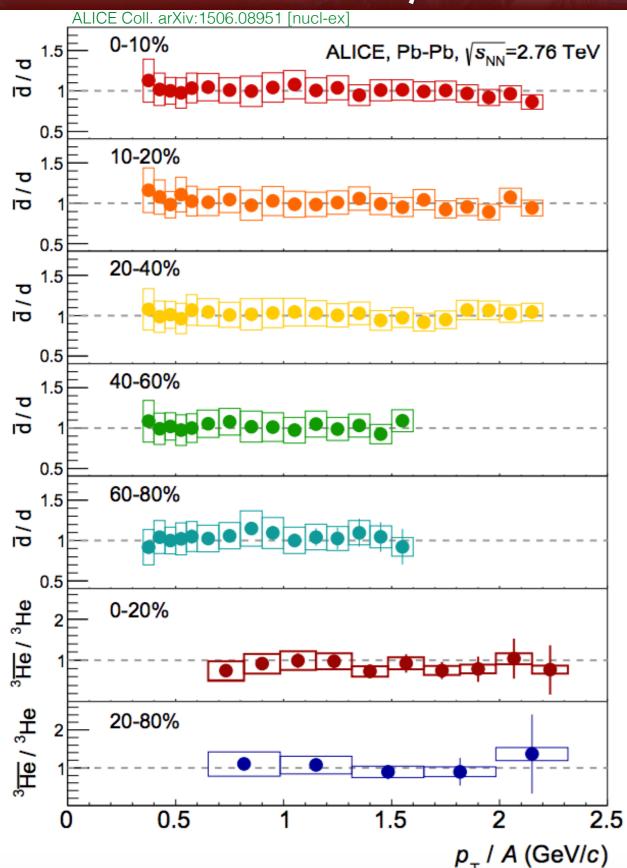
 $T_{kin} = 113 \pm 12 \text{ MeV}$
 $n = 0.72 \pm 0.03$

With respect to the fit performed without the nuclei the $<\beta>$ decreased while the T_{kin} increased but they are compatible within the uncertainties.

Solid symbols denote the spectra points used for the fit.



Ratio matter / anti-matter



The ratio nuclei / anti-nuclei is compatible with one



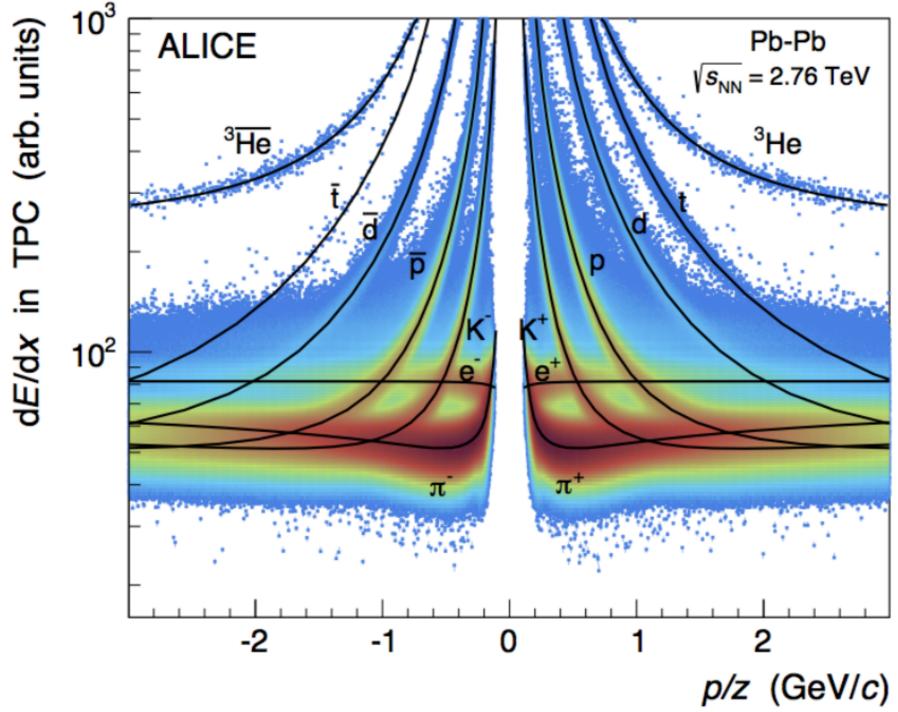
The same ratio is seen for other particle species measured at the LHC

A large fraction of the systematic uncertainties on the determination of the ratios is due to the limited knowledge of the cross sections of anti-nuclei interacting with the material of the detector.

Observation of the anti-triton

ALICE Coll. arXiv:1506.08951 [nucl-ex]

Only those tracks which pass the pre-selection done by applying a 3σ cut on the TPC dE/dx are considered as anti-triton candidates.

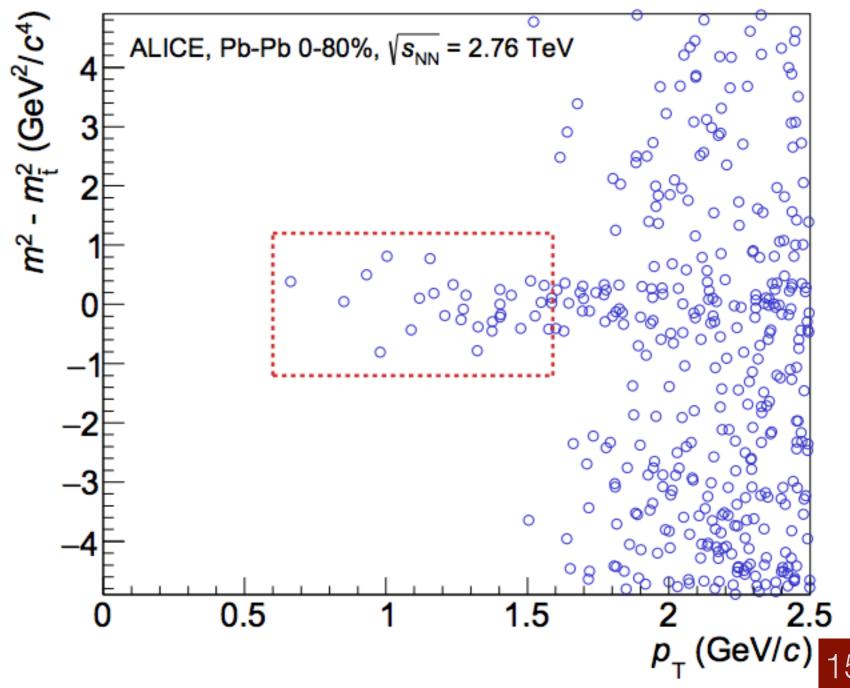


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- 31 anti-tritons can be identified unambiguously
- In this kinematic range, extrapolations place the expected anti-triton yield somewhere between 11 and 40 particles

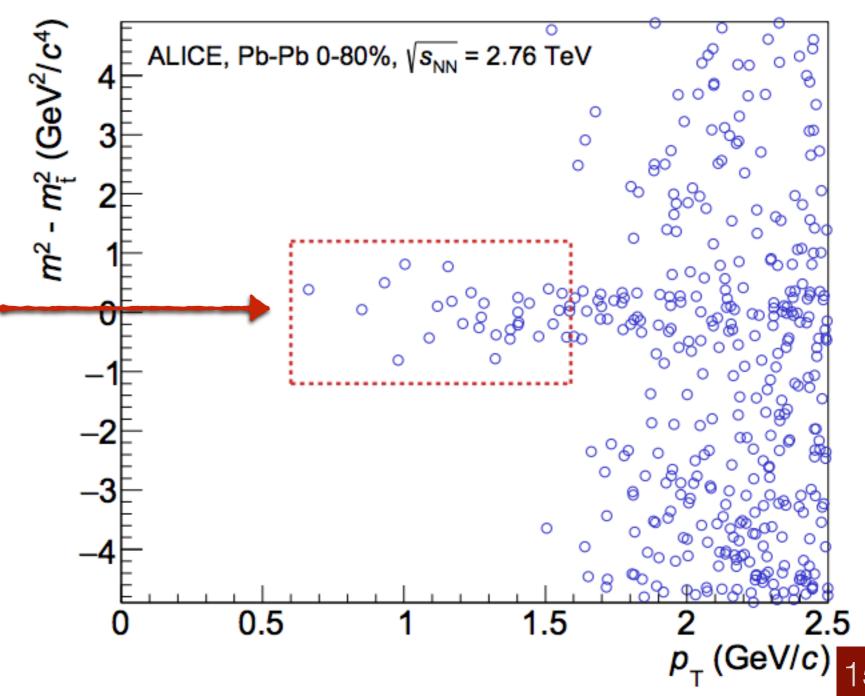


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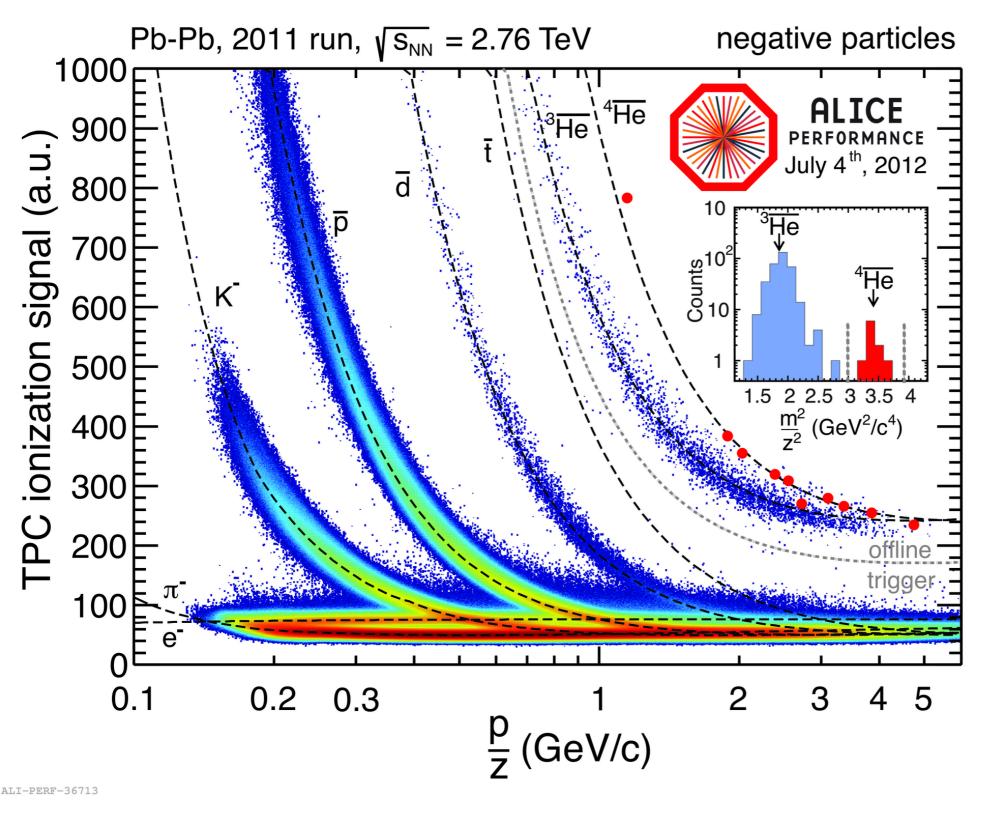
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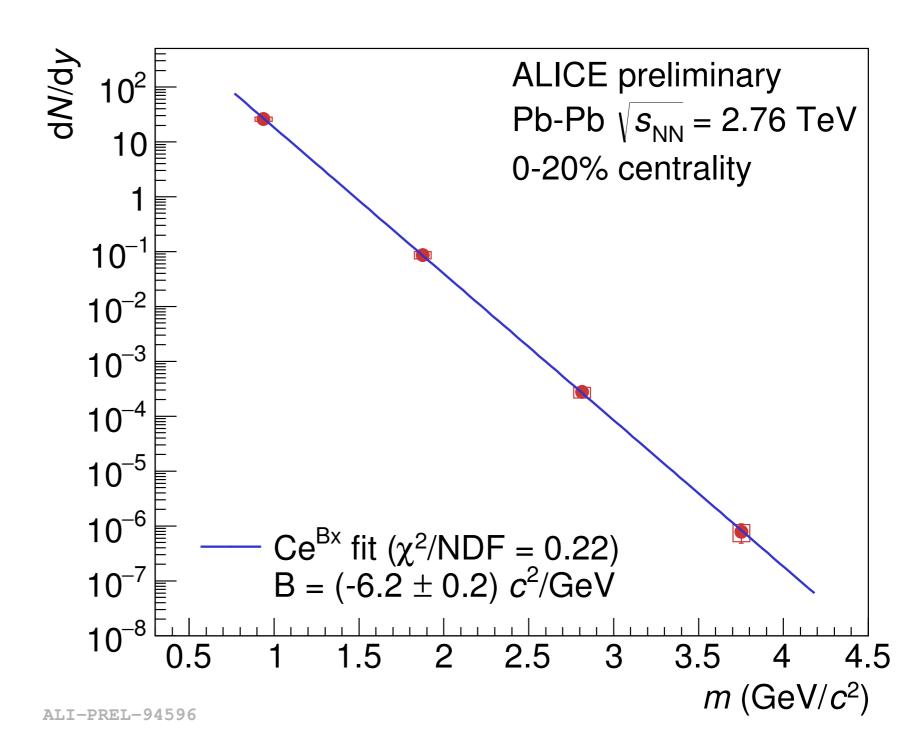
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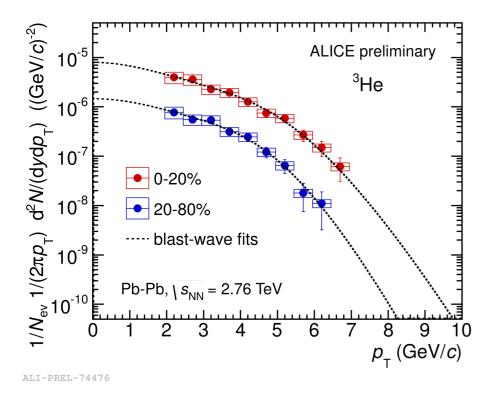


Observation of the anti-4He

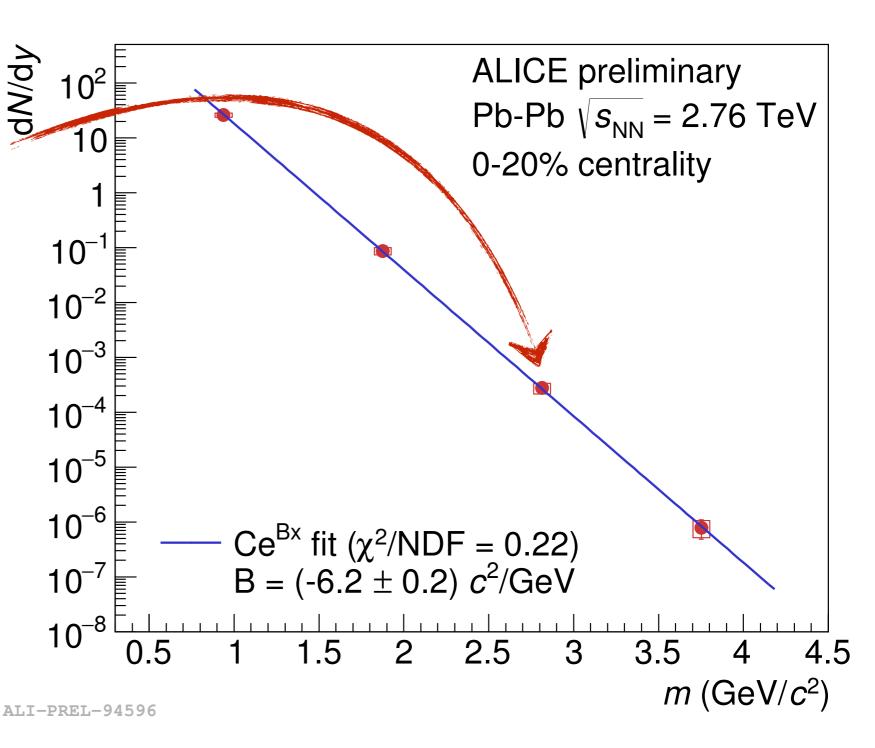


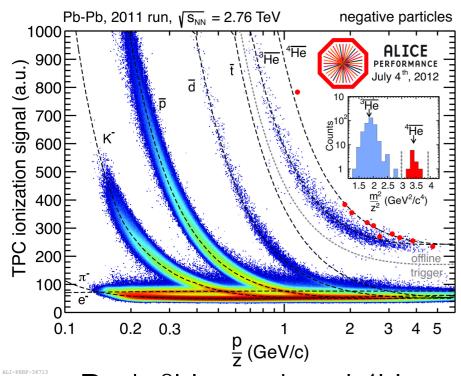
- TPC allows to separate particles with Z=2 from Z=1 particles over the full momentum range.
- Using also the TOF PID it is possible to measure about 10 anti-⁴He



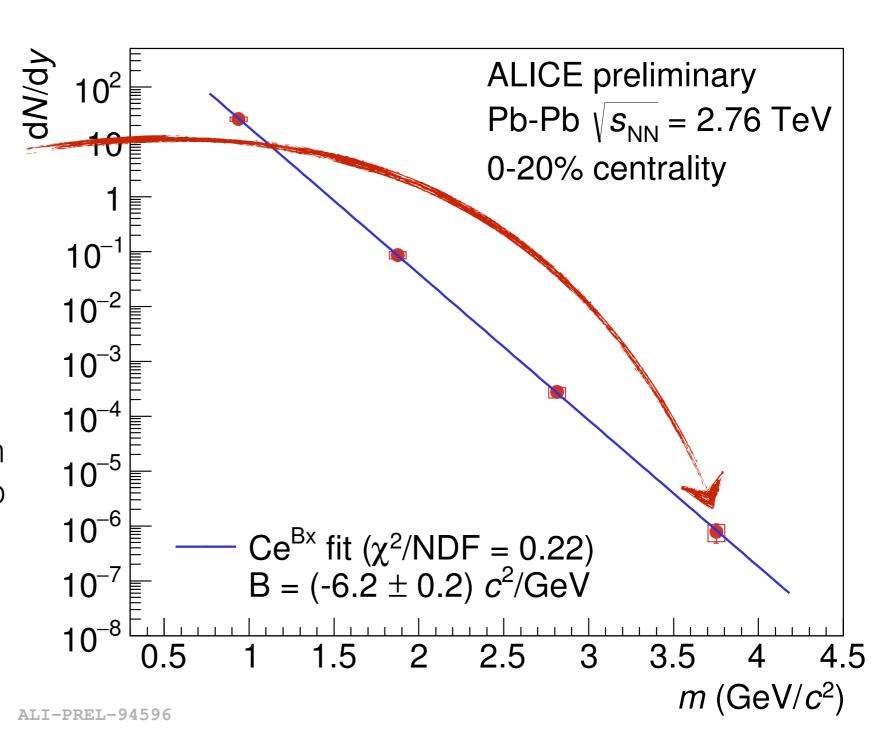


Both ³He and anti-⁴He integrated yields have been measured and compared to deuteron and proton integrated yields





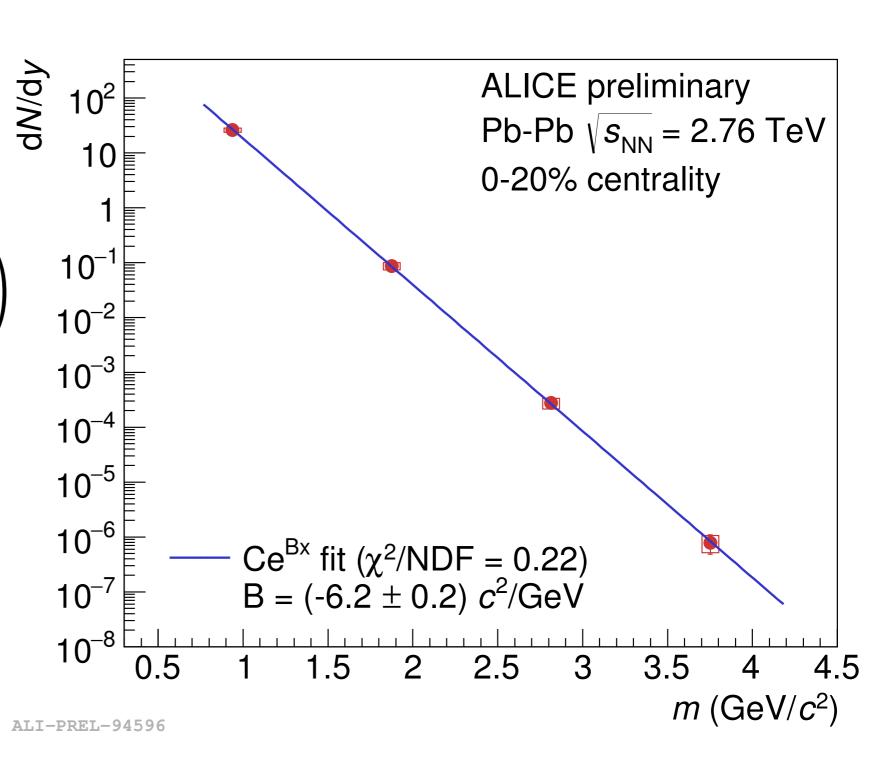
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Thermal model predicts:

$$\frac{\mathrm{d}N}{\mathrm{d}y} \propto exp\left(-\frac{m}{T_{\mathrm{chem}}}\right)$$

- Verified with nuclei!
- Each nucleon added gives a penalty of a factor ~300 in the integrated yield



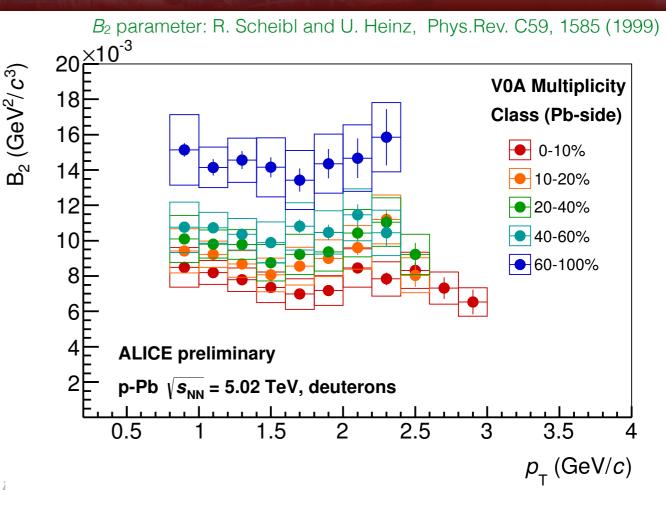
Coalescence parameter B2

The coalescence parameter, defined as:

$$B_2 = \frac{E_{\mathrm{d}} \frac{\mathrm{d}^3 N_{\mathrm{d}}}{\mathrm{d}p_{\mathrm{p}}^3}}{\left(E_{\mathrm{p}} \frac{\mathrm{d}^3 N_{\mathrm{p}}}{\mathrm{d}p_{\mathrm{p}}^3}\right)^2}$$

is predicted to be p_T independent by the simplest formulation of coalescence model.

This is observed in p-Pb collisions.



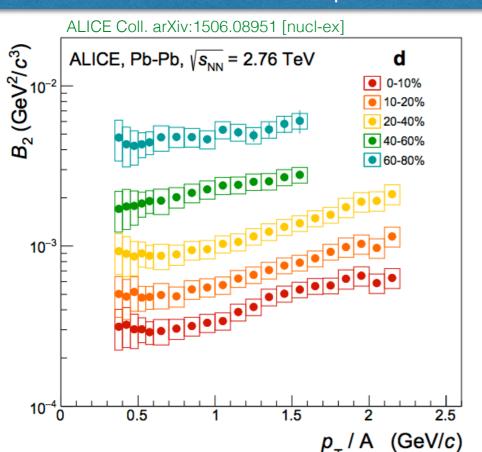
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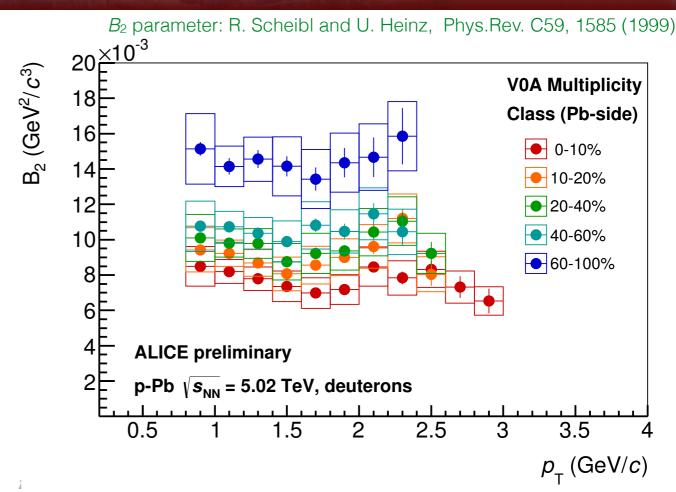
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The coalescence parameter gets smaller as the events are more central. This is due to the increasing size of the emitting source.

The B_2 depends on p_T in central events, which could be explained by looking at the Hanbury Brown and Twiss radii dependence of the B_2 .

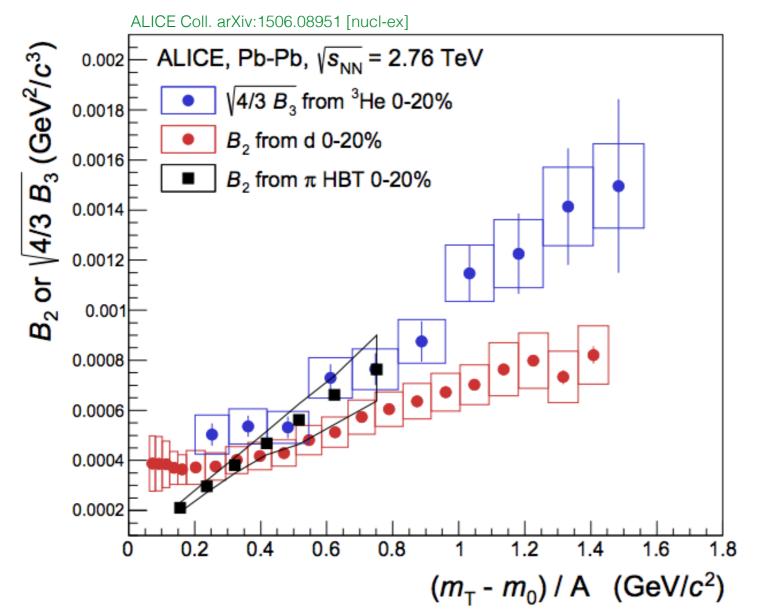
B₂ as a function of the HBT radii

The coalescence parameter can be expressed as a function of the HBT radii:

$$B_2 = \frac{3\pi^{3/2} \langle C_{\rm d} \rangle}{2m_{\rm T} R_{\perp}^2(m_{\rm T}) R_{\parallel}(m_{\rm T})}$$

R. Scheibl and U. Heinz, Phys.Rev. C59, 1585 (1999)

A rough agreement is found in terms of magnitude and the dependence on p_T .



The coalescence parameter for a nucleus *i* with A nucleons is defined as:

$$E_i \frac{\mathrm{d}^3 N_i}{\mathrm{d}p_i^3} = B_A \left(E_\mathrm{p} \frac{\mathrm{d}^3 N_\mathrm{p}}{\mathrm{d}p_\mathrm{p}^3} \right)^A$$

From this it is possible to derive a recursive formula. For instance for the ³He:

$$B_3 = B_2^2 \left(\frac{M_{^3He} \cdot m}{M_d^2} \right) \approx \frac{3}{4} B_2^2$$

(Anti-) Hyper-triton in Pb-Pb

Hyper-triton is the lightest hyper-nucleus. ALICE collaboration measured its production in the charged 2 body decay channel.

Mass = 2.991 GeV/c^2 Lifetime ~ 215 ps

Signal Extraction:

- Identify 3 He and π
- Evaluate (³He,π) invariant mass
- Apply topological cuts in order to:
 - identify secondary decay vertex
 - reduce combinatorial background

Decay modes

$$^{3}_{\Lambda}H \rightarrow^{3}He + \pi^{-}$$
 $^{3}_{\Lambda}H \rightarrow^{3}H + \pi^{0}$
 $^{3}_{\Lambda}H \rightarrow d + p + \pi^{-}$
 $^{3}_{\Lambda}H \rightarrow d + n + \pi^{0}$

+ anti-hypertriton counterpart

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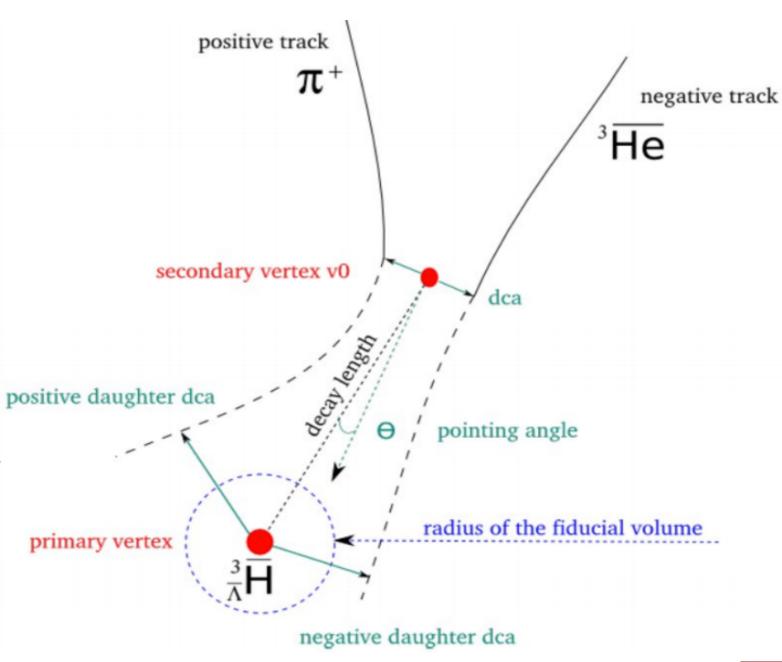
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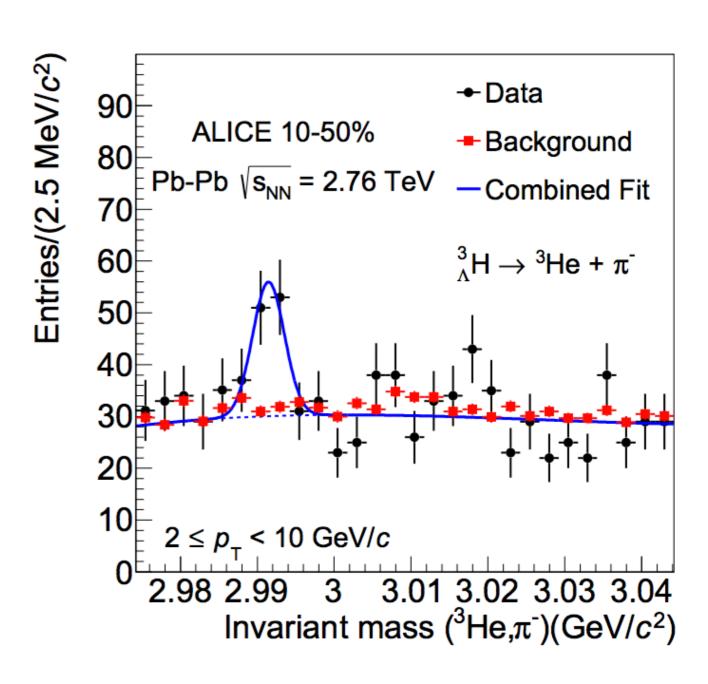
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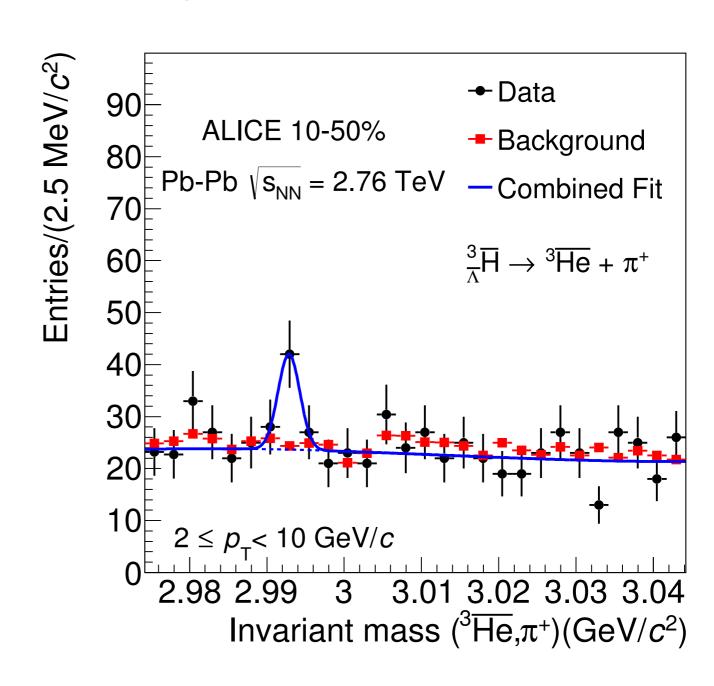
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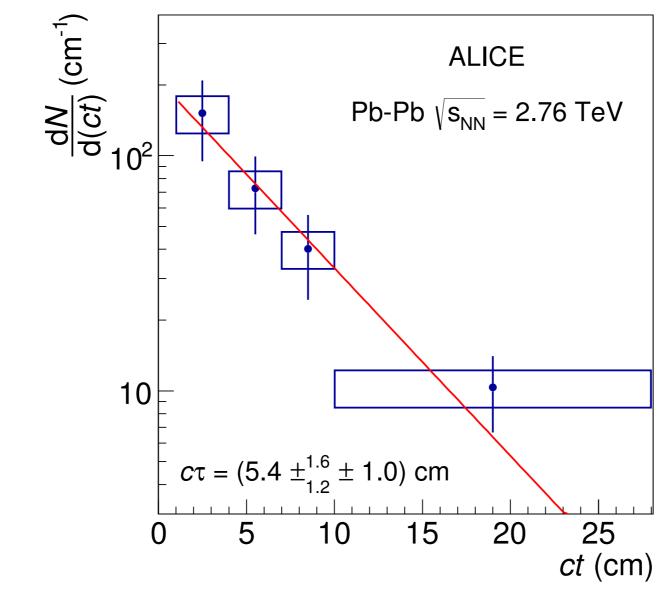
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Measurement of the hyper-triton lifetime

ALICE Coll. arXiv:1506.08453 [nucl-ex]

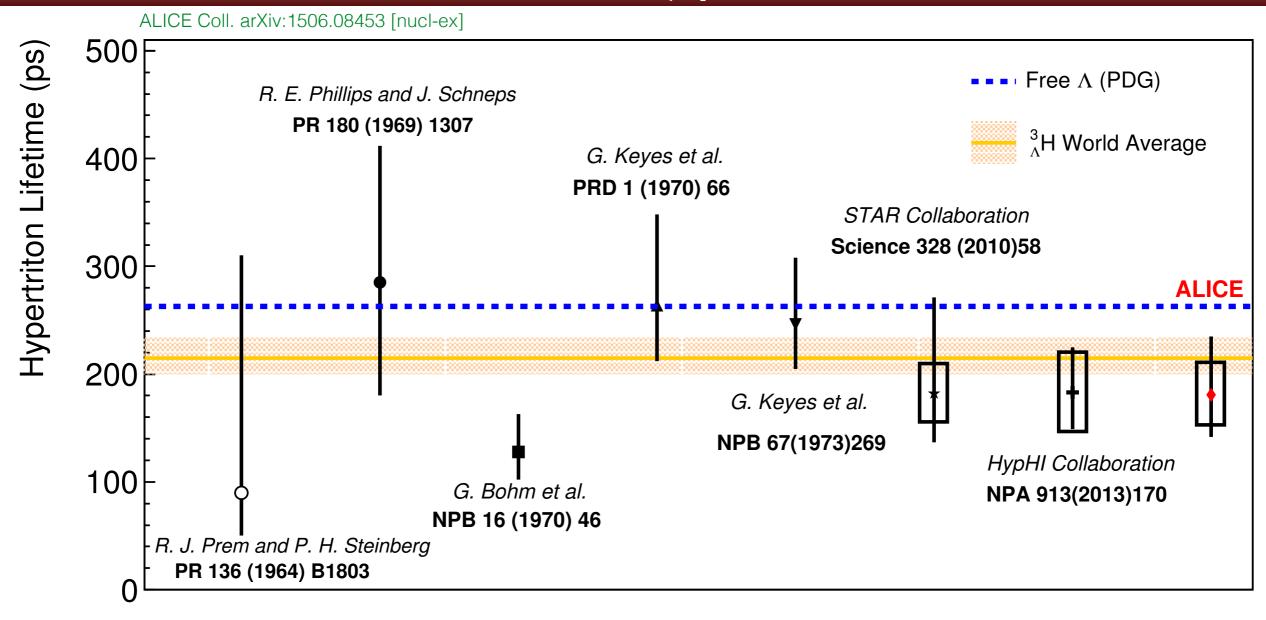


From the exponential fit to the differential yield in different ct bins it is possible to extract the lifetime of the hyper-triton

ALICE
$$\tau(181^{+54}_{-39}({\rm stat.}) \pm 33({\rm syst})) \, {\rm ps}$$

World average
$$\tau = 215^{+18}_{-16} \, \mathrm{ps}$$

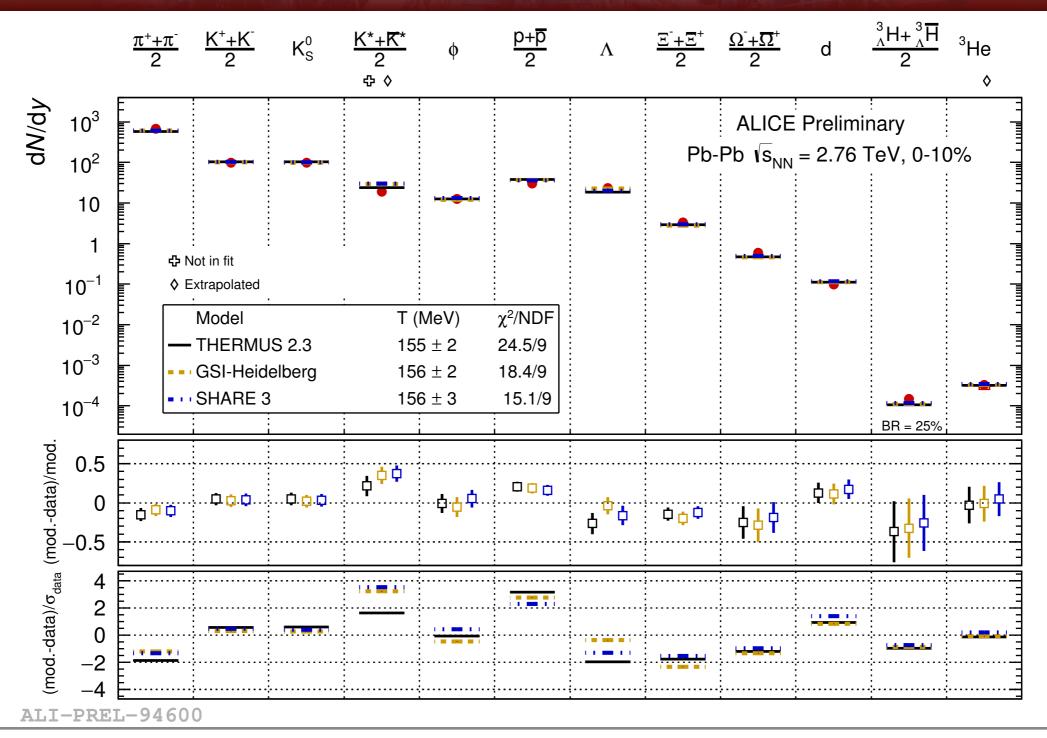
Measurement of the hyper-triton lifetime



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 $\tau=215^{+18}_{-16}\,{\rm ps}$

Thermal model



Grand canonical thermal fit for 0-10% central Pb-Pb collisions, with different thermal models. All models fit the abundances with T \approx 156 MeV. The predictions of all of them miss the yield of protons and K*.

Search for $\Lambda\Lambda$ and Λn

Thermal models describe with high accuracy the abundance of different particle species. They can also predict the abundances of exotic weakly decaying dibaryon states, such as:

- $\Lambda\Lambda$ (H-dibaryon):
 - predicted by Jaffe in bag model calculations

R. L. Jaffe, PRL 38, 195 (1977)

- lattice calculations suggest a lightly bound state or resonance close to Ξp threshold

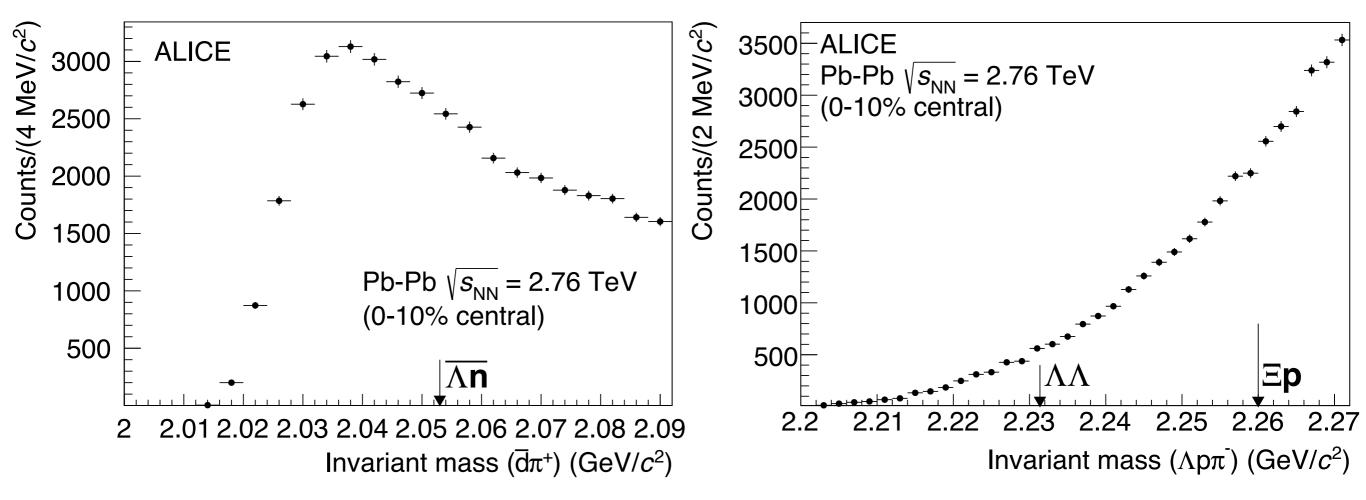
 T. Inoue et al., PRL 106, 162001 (2011) and S. R. Beane et al., PRL 106, 162002 (2011)
- Predicted yield: 211 2110 (in 14M events)
- An-bar bound state
 - Predicted yield: ~4000 (in 14M events)

Both $\Lambda\Lambda$ and Λ n-bar should be measurable with the statistics available in ALICE according to the thermal models predictions.

Invariant mass spectra for $\Lambda\Lambda$ and Λ n-bar

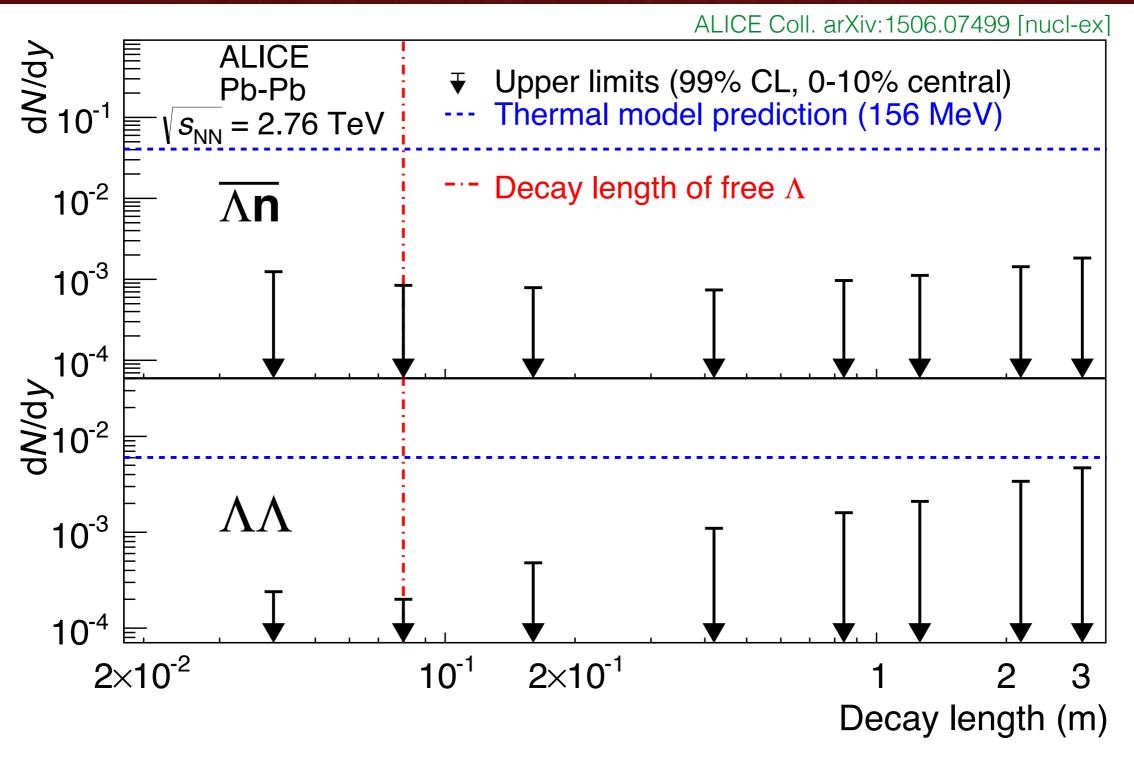
ALICE Coll. arXiv:1506.07499 [nucl-ex]

The arrows indicate the sum of the masses of the constituents.



No signal visible in the invariant mass spectra. From the nonobservation it is possible to compute the upper limits for the production rates.

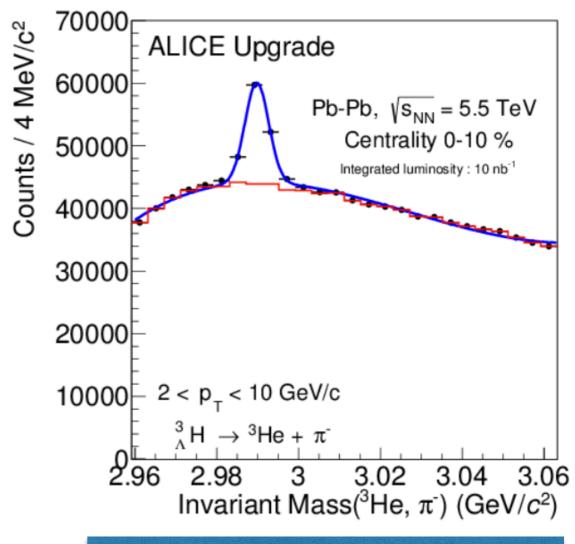
Upper limits



Upper limit of the dN/dy in the assumption of 64% branching ratio for the H-dibaryon and 54% for the Λ n-bar state.

ALICE Upgrade

- Run 3 & Run 4 of LHC will deliver much more statistics (50 kHz Pb-Pb collision rate)
- Upgraded ALICE apparatus will be able to cope with the new high luminosity environment
- ITS Upgrade: less material budget and more precise tracking for the identification of hyper-nuclei.
- All the physics which is now done for A = 2 and A = 3 (hyper-)nuclei will be done for A = 4.



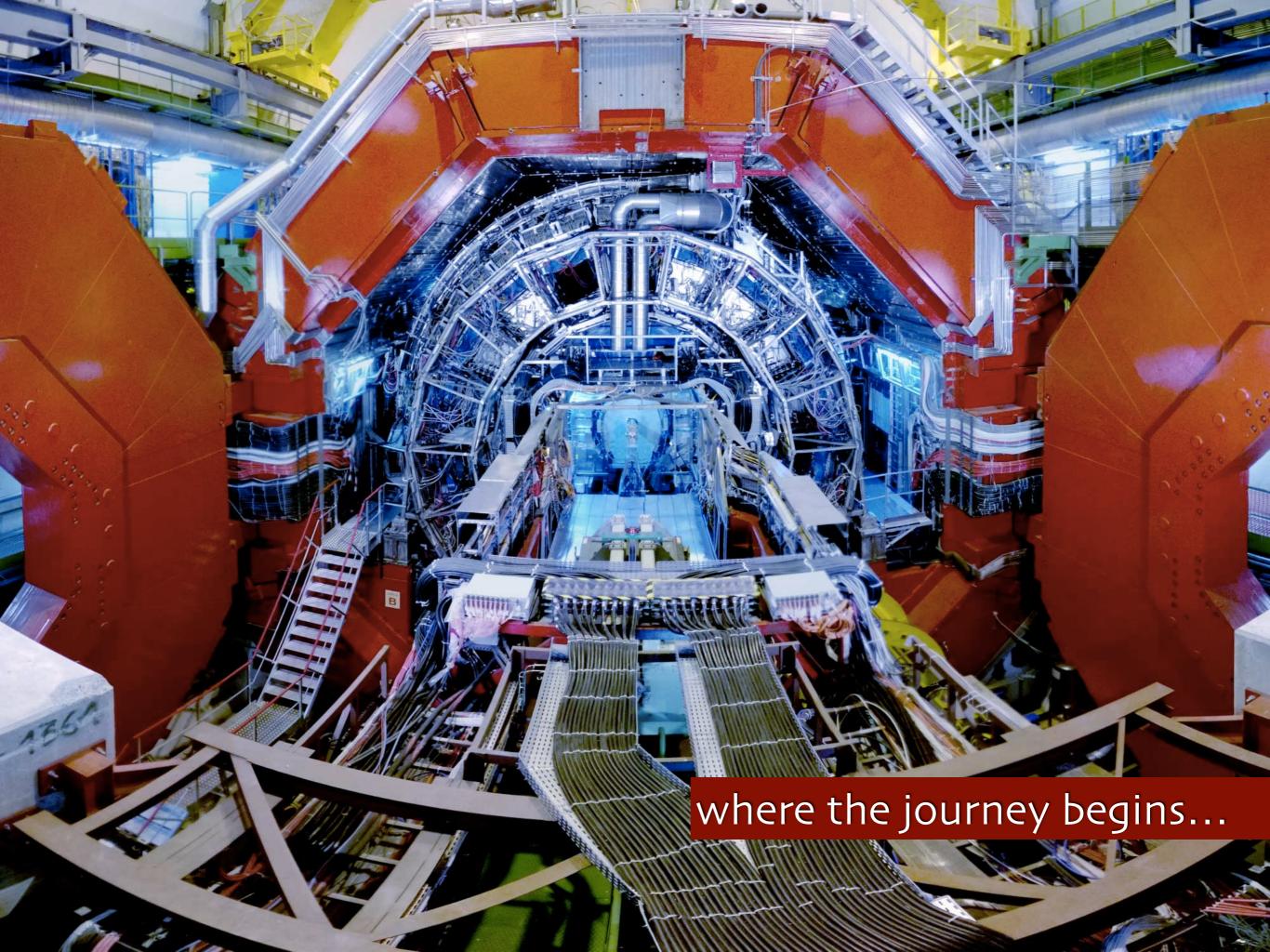
ALICE Coll. J. Phys. G 41 (2014)

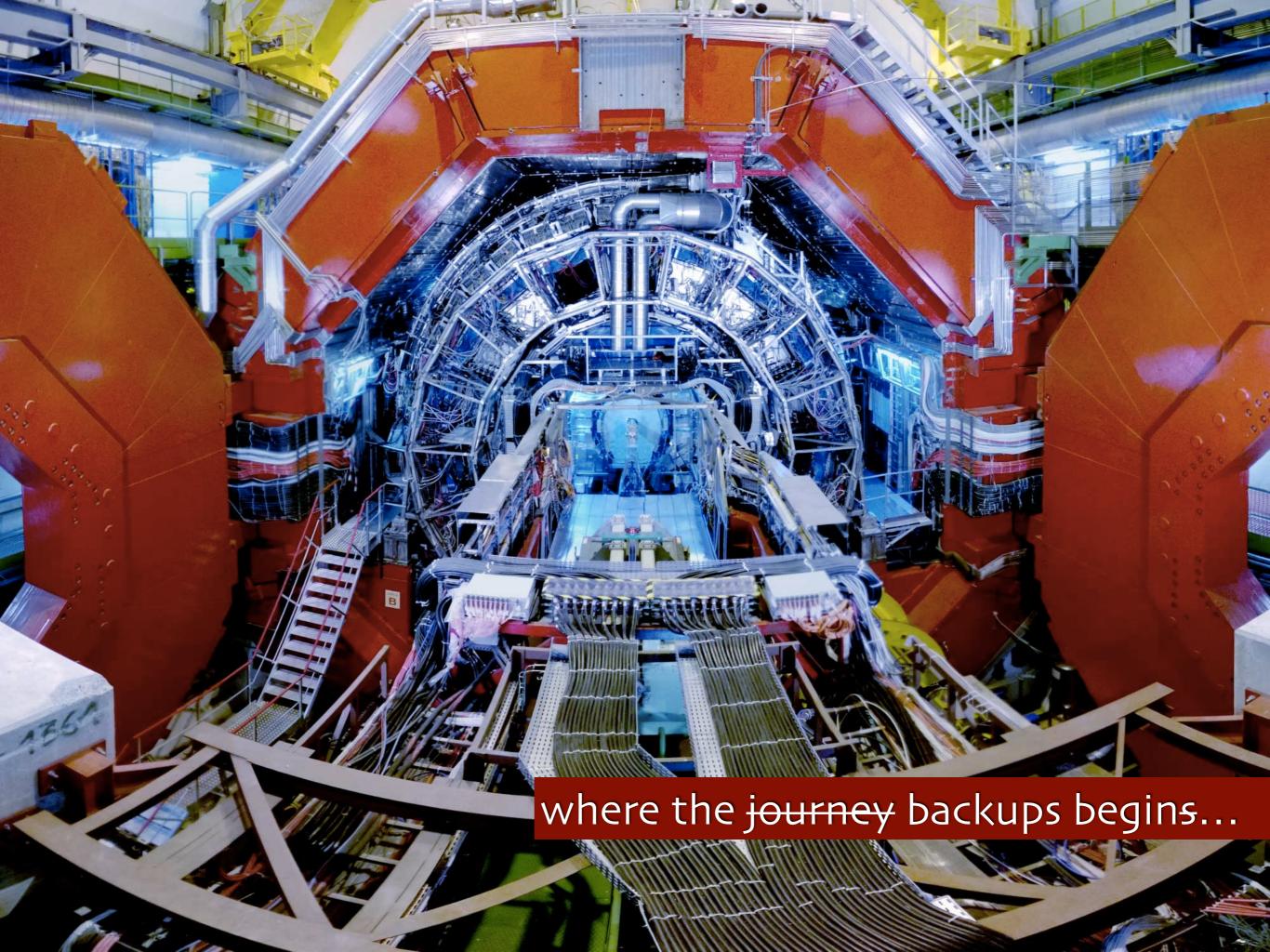
Numbers for $L_{int} = 10 \text{ nb}^{-1}$

State	$\mathrm{d}N/\mathrm{d}y$	B.R.	$\langle \mathrm{Acc} \times \epsilon \rangle$	Yield
$\overline{\ _{\Lambda}^{3}H}$	1×10^{-4}	25%	11 %	44000
$^{4}_{\Lambda}H$	2×10^{-7}	50%	7%	110
$\frac{4}{\Lambda}He$	2×10^{-7}	32%	8 %	130

Conclusions and perspectives

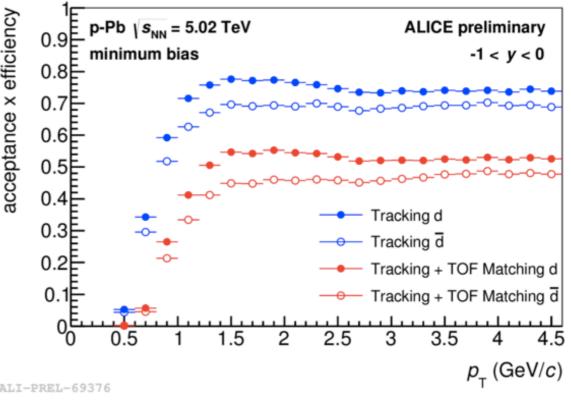
- The nuclei production (up to A=4) has been measured by ALICE experiment.
- Both coalescence and thermal model are successful in the description of particular aspects of the measurements:
 - Integrated yields are well described by thermal models
 - The trend of the coalescence parameter and of the d/p ratio could be described by the coalescence model
- Hyper-triton measurement shows the high performance of the experimental apparatus and will further benefit from the increase of statistics.
- The data from ALICE do not support the existence ΛΛ and Λn-bar.
- ALICE is looking forward to LHC Run 2 to extend the current measurements and put tighter constraints on theoretical models.

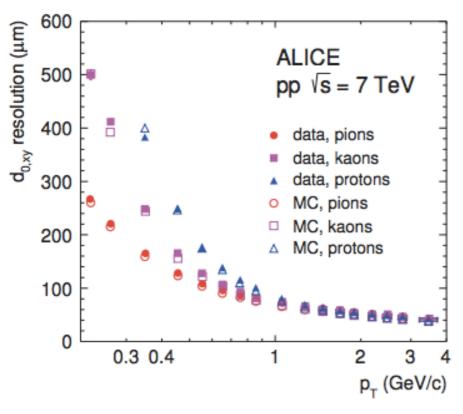


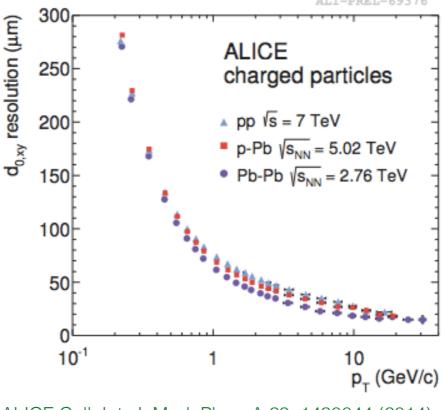


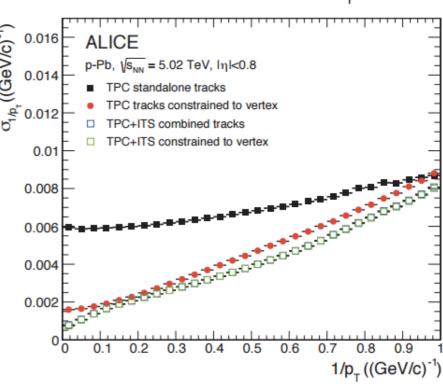
Tracking performance of ALICE

- Tracking efficiency x
 acceptance is fundamental to
 measure the production
 spectra of charged particle.
- Tracking efficiency depends on charged particle multiplicity.



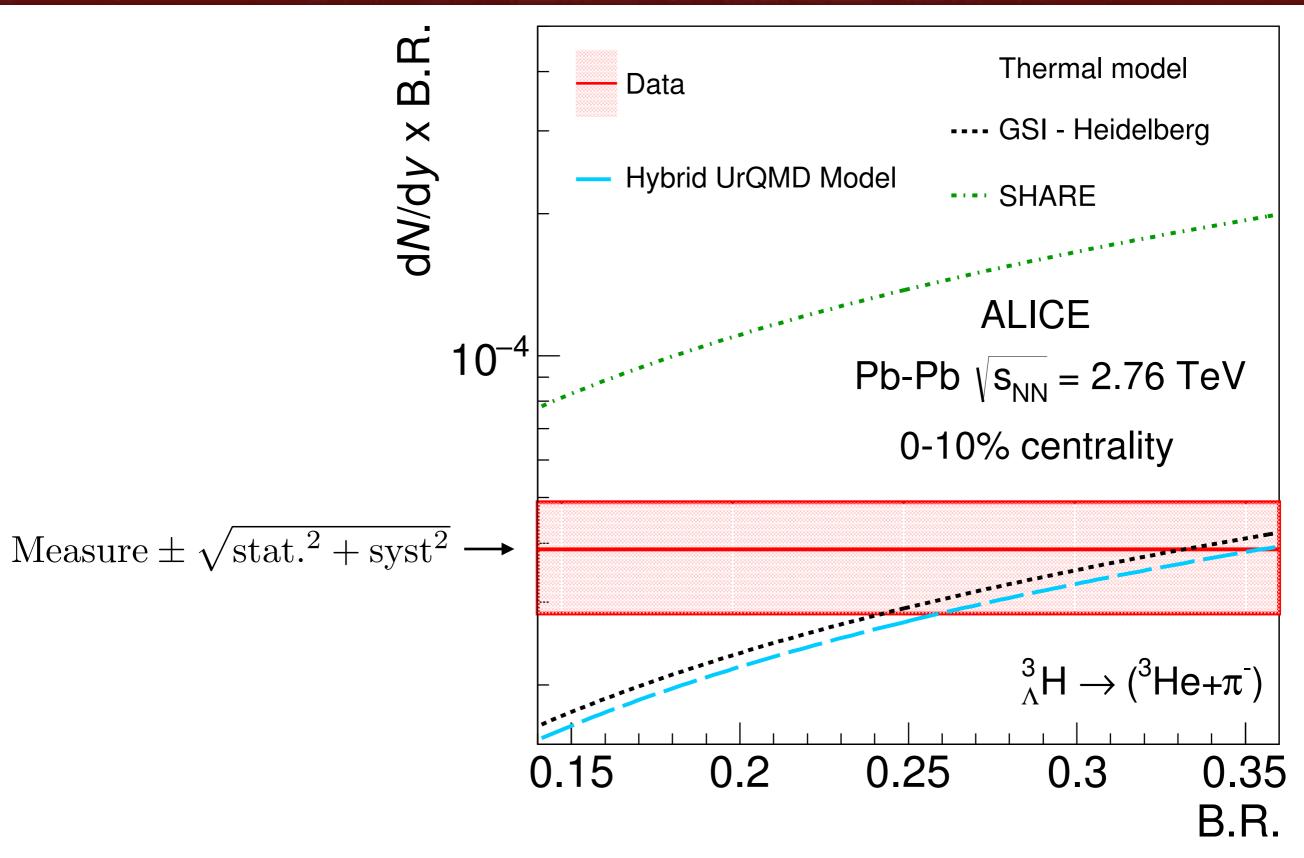






ALICE Coll. Int. J. Mod. Phys. A 29, 1430044 (2014)

Hyper-triton yield compared to thermal models



Upper limits to AA and An-bar production

