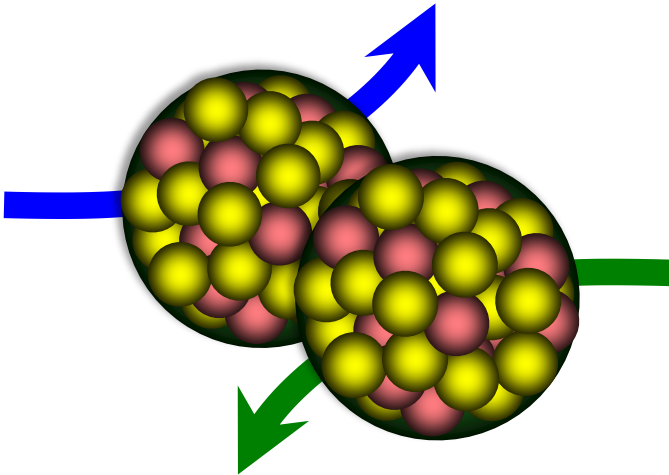
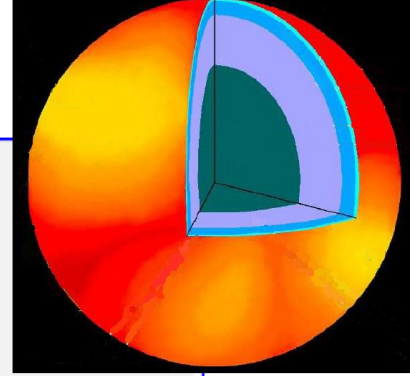


High-density nuclear matter probed by **nucleus-nucleus scattering**

- its implication to **neutron star mass** problem -



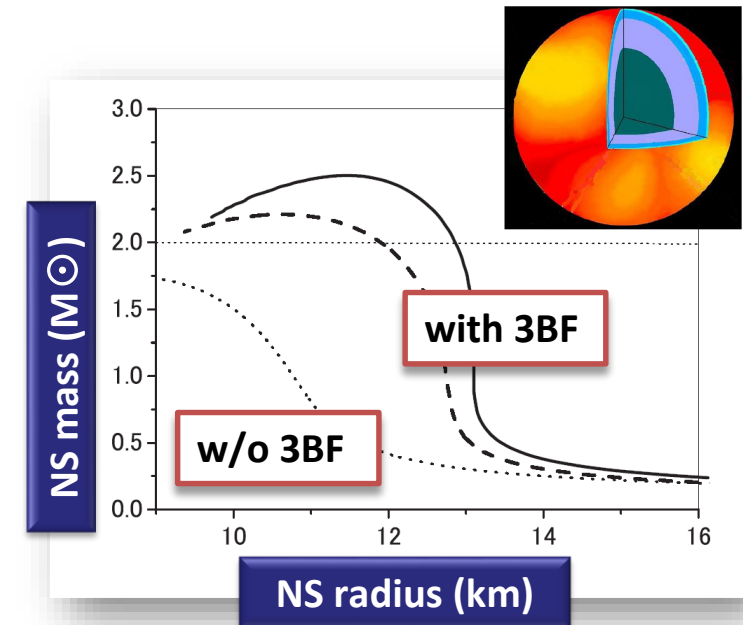
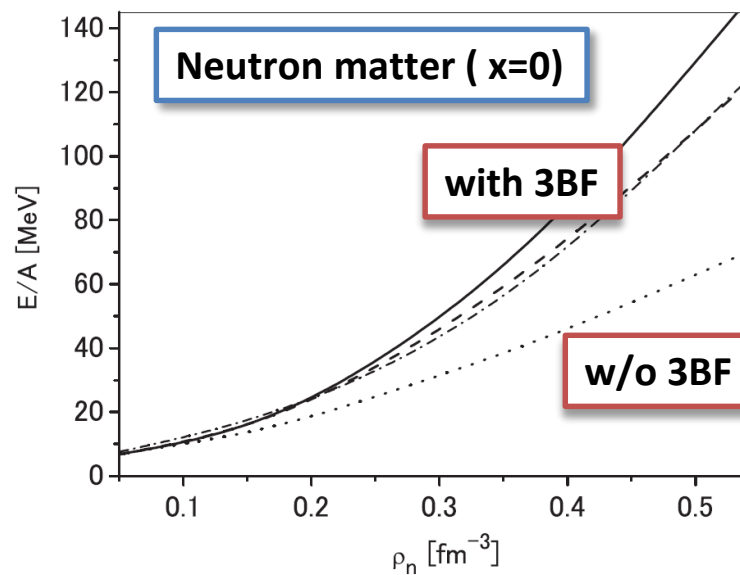
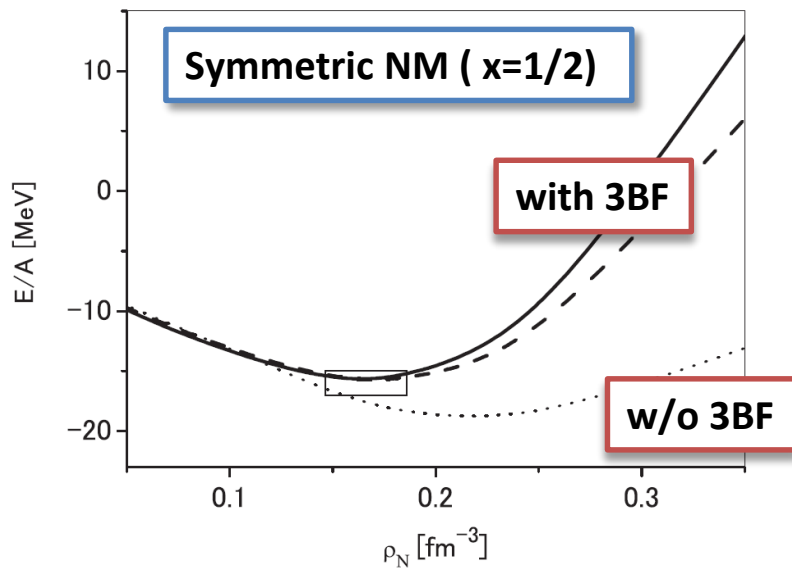
Y. Sakuragi (Osaka City University, Japan)

with **T. Furumoto, Y. Yamamoto,**

N. Yasutake, Th. A. Rijken

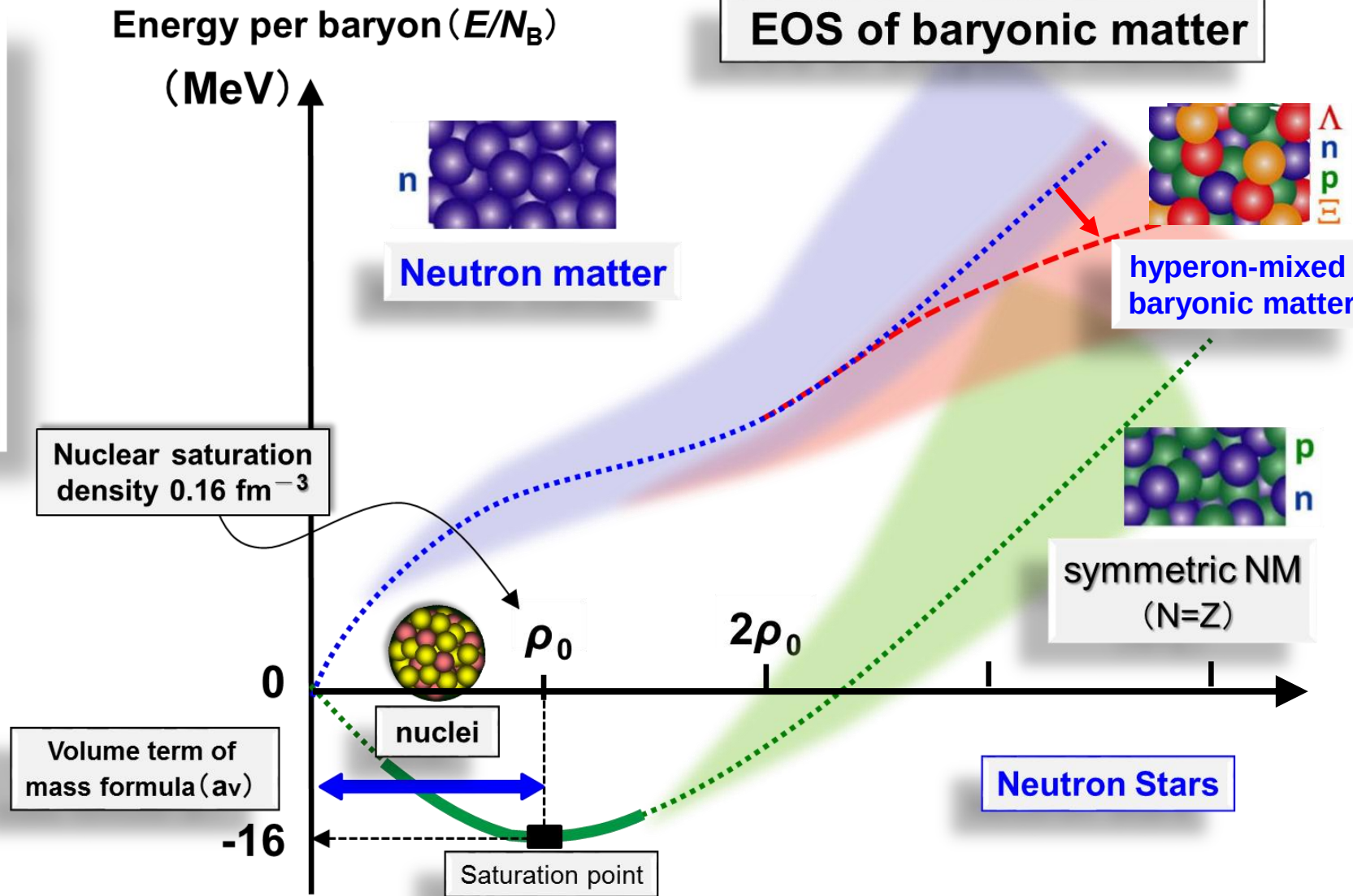
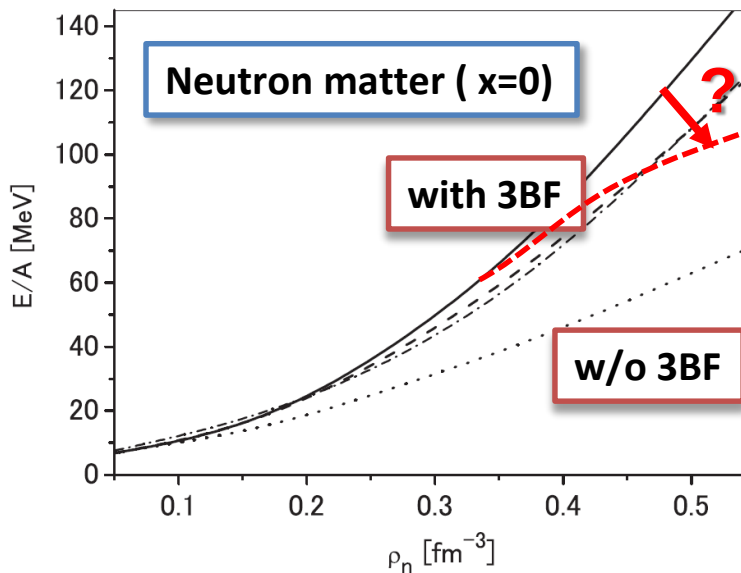
Background of the present study

- ✓ The new observation of the system J1614-2230 has brought a great impact on **the maximum-mass problem** of neutron stars. The large observed value of neutronstar mass, **$1.97 \times M_{\odot}$** , **sets a severe condition for the stiffness of the equation of state (EOS) of neutron-star matter**
- ✓ The **repulsive 3NF force** contribution in **high-density neutron matter** should play an essential role for **stiffening the EOS** of neutron-star matter, assuring the observed **maximum mass of neutron stars ($2M_{\odot}$)**



Background of the present study

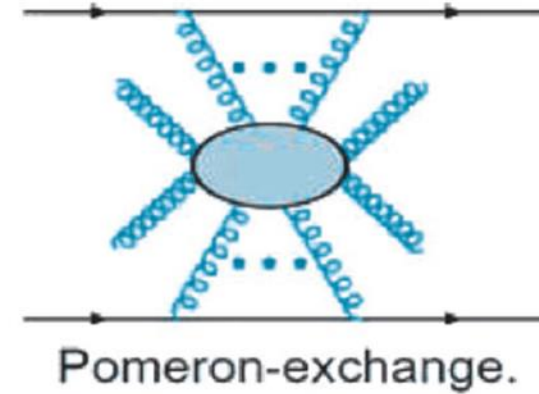
- ✓ However, **hyperon (Y) mixing** in neutron-star matter brings about a **remarkable softening** of the EOS at high- ρ region ($\rho > (2\sim 3) \times \rho_\odot$), which cancels the repulsive 3NF effect for the maximum mass.



Background of the present study

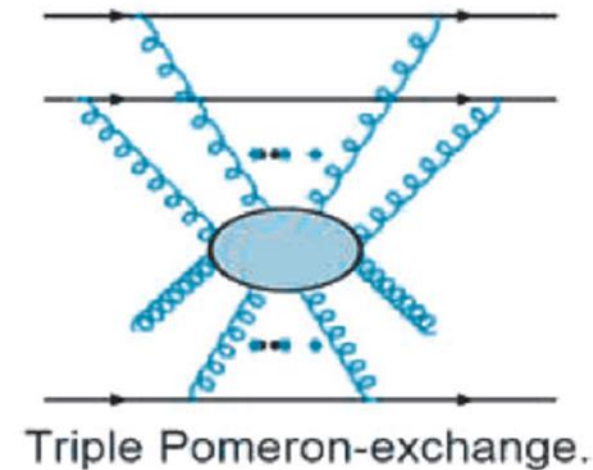
- ✓ In order to avoid this serious problem, we have proposed **a multi-Pomeron exchange potential (MPP)** that works universally not only in **NNN** states but also in **YNN**, **YYN**, and **YYY** states.
- ✓ Our basic concern is the existence of universal repulsions among three baryons (BBB), called here **3-baryon repulsion (BBB repulsion)**.
- ✓ The role of the repulsive BBB for stiffening the EOS can be recovered clearly by invoking this assumption.

What is pomeron?



Why pomeron?

SU3 scalar



Pomeron is a model for multi-gluon exchange

Our strategy

- ✓ The 3BF is very sensitive to **nucleus-nucleus (A-A) elastic scattering**
- ✓ We determine the strength of universal **3,4-baryon repulsion** by analyses of **A-A elastic scattering** with the **G-matrix folding model** (→explained soon)
- ✓ Then, construct the **EOS of neutron-star matter**,
- ✓ Examine if it **results in a stiff EOS**, enough to give the **observed neutron-star mass** ($M_{NS} \approx 2M_{\odot}$) with hyperon mixing.

Energy per baryon (E/N_B)
(MeV)

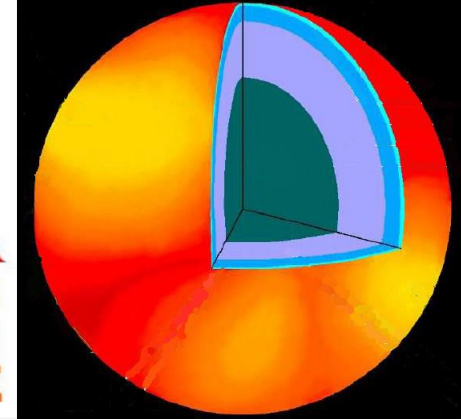
EOS of baryonic matter



Neutron matter

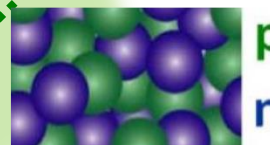


hyperon-mixed baryonic matter



Information on
 $\rho = \rho_0 \sim 2\rho_0$
is missing !

Nuclear saturation
density 0.16 fm^{-3}



symmetric NM
($N=Z$)

0

nuclei

ρ_0

$2\rho_0$

$5\rho_0$

Volume term of
mass formula (av)

-16

Saturation point

Neutron Stars

Baryon density (ρ_B)

Energy per baryon (E/N_B)
(MeV)

Probe **high-density nuclear matter** ($\rho = \rho_0 \sim 2\rho_0$)
with **nucleus-nucleus scattering**

Neutron matter

Information on
 $\rho = \rho_0 \sim 2\rho_0$
is **missing** !

hyperon-mixed baryonic matter

Nuclear saturation
density 0.16 fm^{-3}

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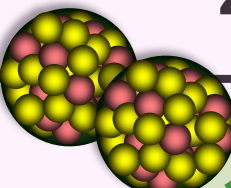
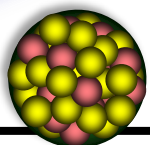
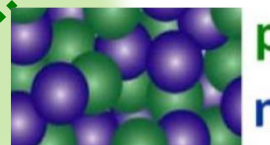
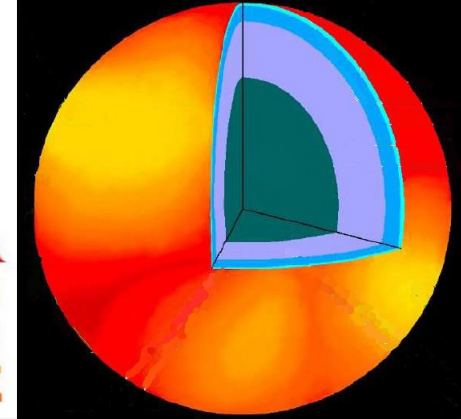
symmetric NM
($N=Z$)

$5\rho_0$

Baryon density (ρ_B)

Neutron Stars

Saturation point



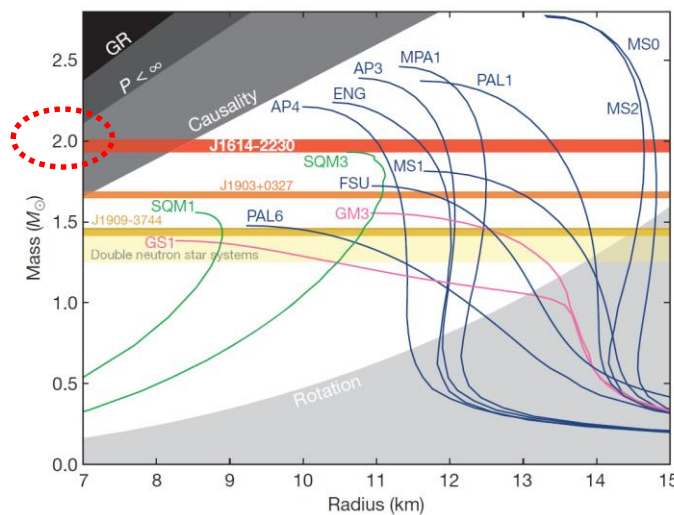
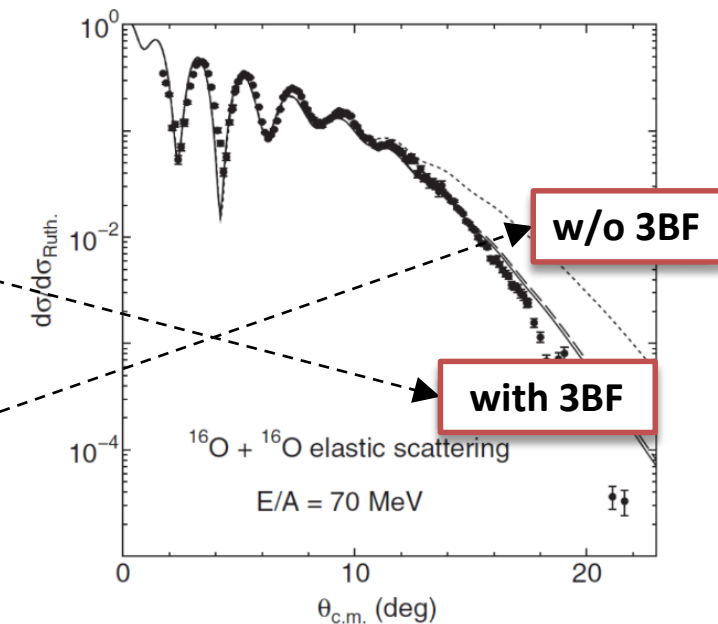
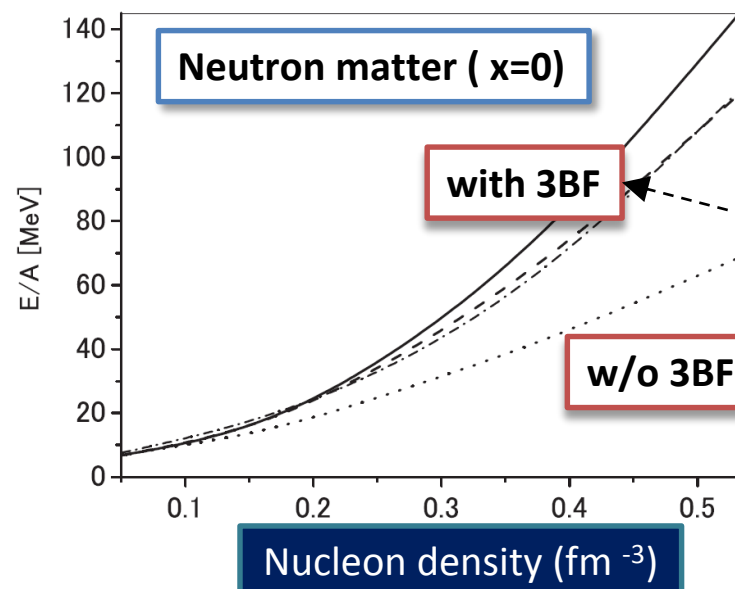
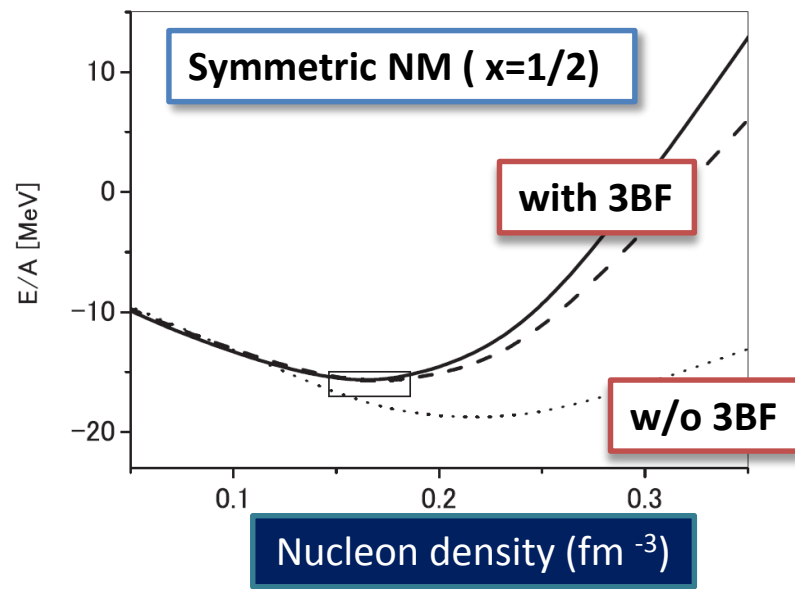
Purpose

High-density nuclear matter probed by **nucleus-nucleus scattering**
- *its implication to neutron star mass problem* -

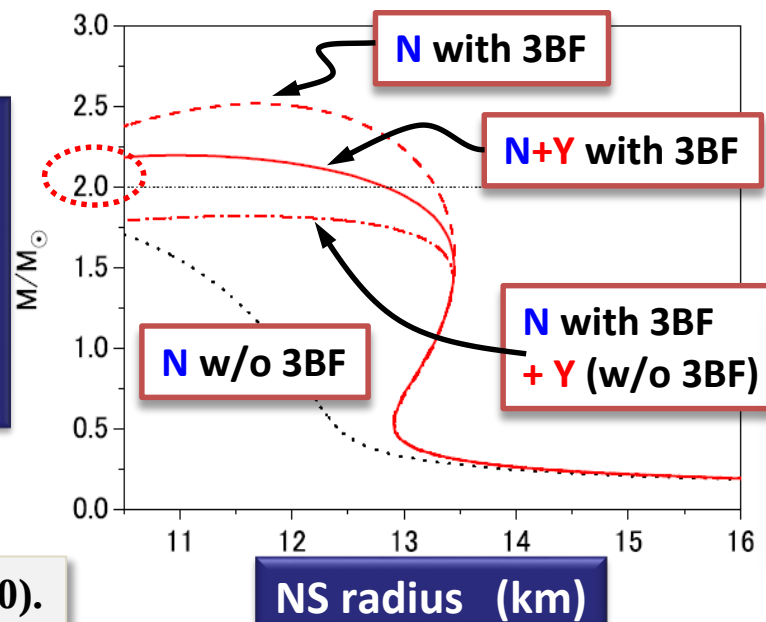
- ✓ To probe the property of **high-density nuclear matter** ($\rho = \rho_0 \sim 2\rho_0$) from the analyses of **nucleus-nucleus scattering** using **Microscopic interaction model** between finite nuclei on the basis of **Brueckner G-matrix** theory including the effects of **3-body & 4-body forces** in the **Multi-Pomeron exchange Potential** (MPP) model.



- ✓ **To fix the uncertain strength** of **3-body & 4-body forces** that are essential to determine the **stiffness of baryonic matter at high- ρ** and to see whether it is consistent with $M_{NS}^{(\max)} \cong 2 \times M_{\odot}$.



NS mass ($M_{\odot}$)



Complex G -matrix
> ESC+MPP model

P.B. Demorest et al., Nature 467, 1081 (2010).

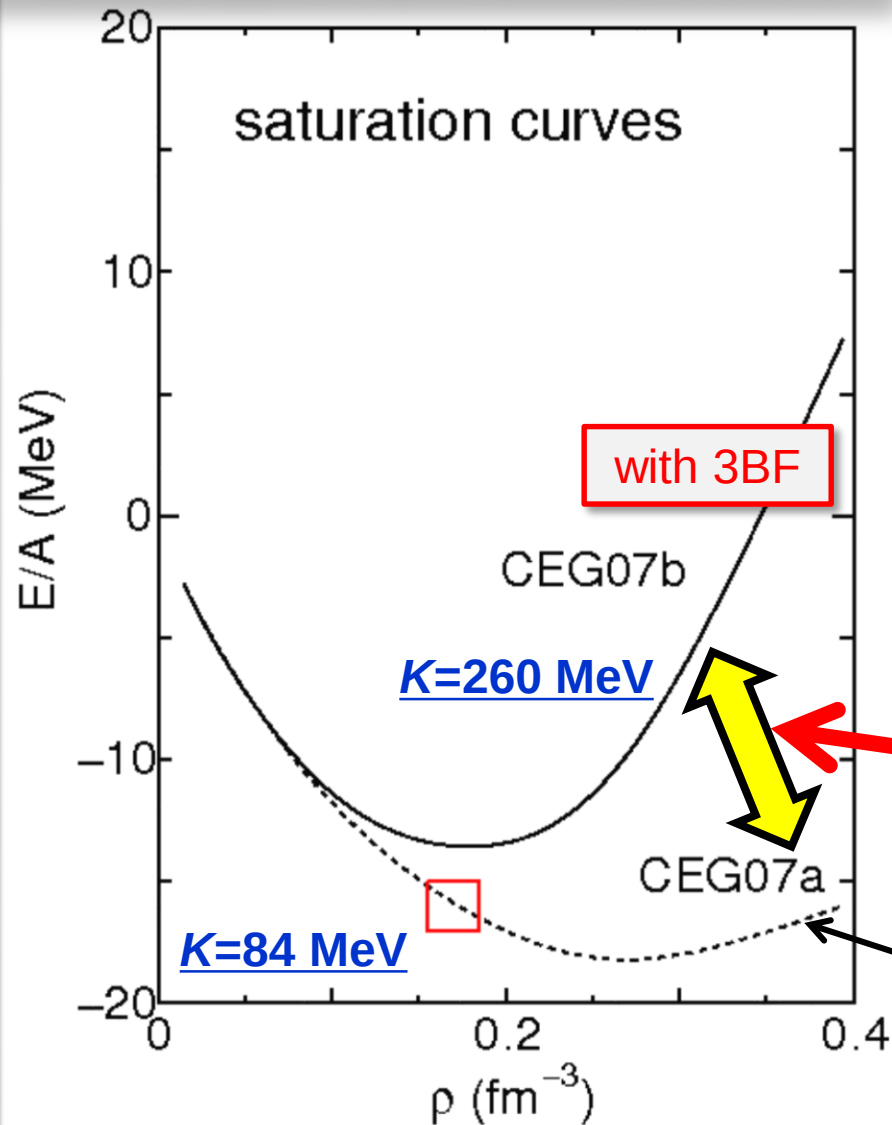
- Y. Yamamoto, T. Furumoto, N. Yasutake, Th. A. Rijken, PRC 88, 022802(R) (2013)
- Y. Yamamoto, T. Furumoto, N. Yasutake, Th. A. Rijken, PRC 90, 045805 (2014)

Brueckner G-matrix generated by **SU(3)-based BB interaction model (ESC08c)**:

● Th. A. Rijken et al., Prog. Theor. Phys. Suppl. **185**, 14 (2010)

- 1) **NN, NY, YY** interactions are treated **on the same footing: SU(3)**
- 2) Effects of the **3-body force (3BF)** and **4-body force (4BF)** are introduced based on the **Multi-Pomeron-Exchange Potential (MPP)** model, **consistently with the ESC08 modeling** for two-baryon system.
- 3) **Universal repulsion** among three baryons (**NNN, YNN, YYN, YYY**)
→ the strength of **the universal 3BF** can be determined through the **sensitivity** of **3BF (NNN)** effects on the **nucleus-nucleus scattering experiments**.

Saturation curves of Symmetric NM with G-matrix interaction (CEG07)



Brueckner Theory (G-matrix)

- **Effective interaction in infinite nuclear matter (NM)**
- Solve the **Bethe-Goldstone equation** in NM with a constant density ρ
- Summing-up of the **ladder diagram**
with **Medium effects** (Pauli effects, Binding effect etc.)

$$\left| \hat{G} \right| = \left| \begin{array}{c} \text{---} \\ \text{---} \end{array} \right| + \left| \begin{array}{c} \text{---} \\ \text{---} \end{array} \right| + \left| \begin{array}{c} \text{---} \\ \text{---} \end{array} \right| + \dots$$

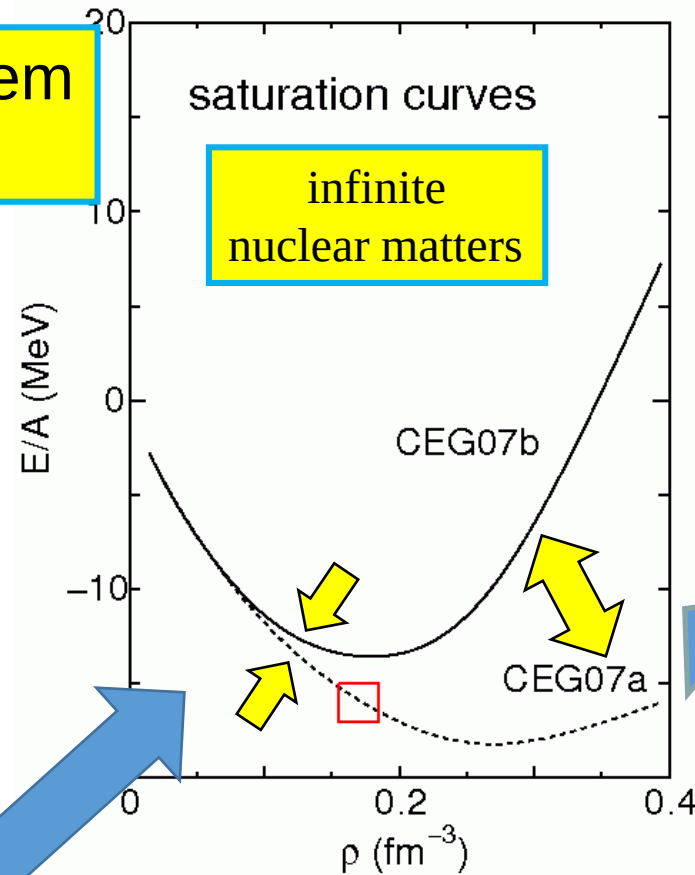
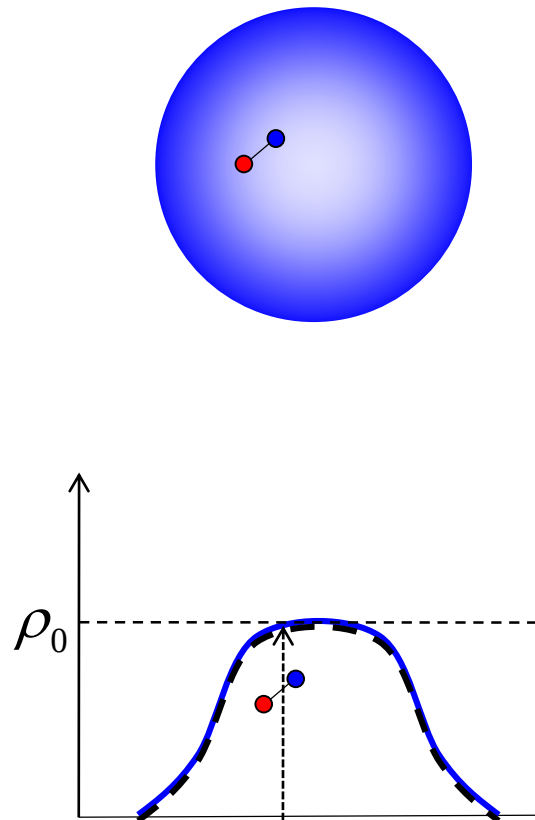
Importance of 3BF effect

w/o 3BF

T.Furumoto, Y. Sakuragi, Y. Yamamoto, PRC 79, 011601(R) (2009)

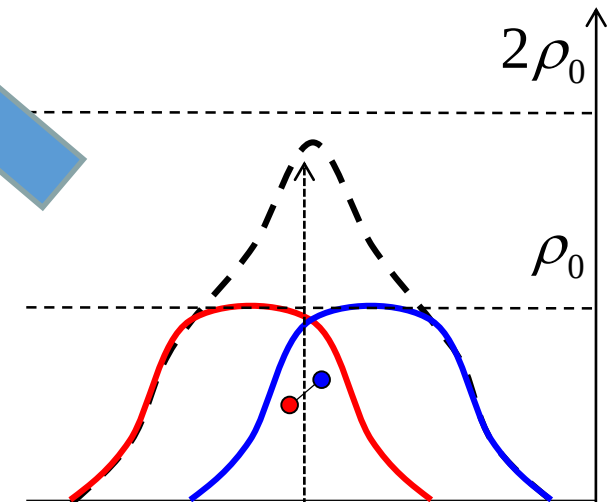
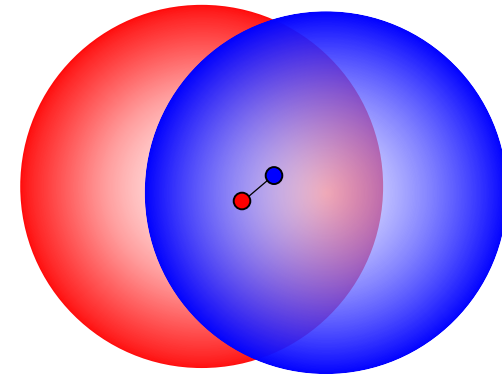
Effective interaction in infinite NM $\rightarrow$ applied to finite system in local-density approx.(LDA)

nucleon-nucleus (NA) system
(one nuclear matter)



low- $\rho \leftrightarrow$ high- ρ

nucleus-nucleus (A-A) system
(overlapped nuclear matters)



Energy per baryon (E/N_B)
(MeV)

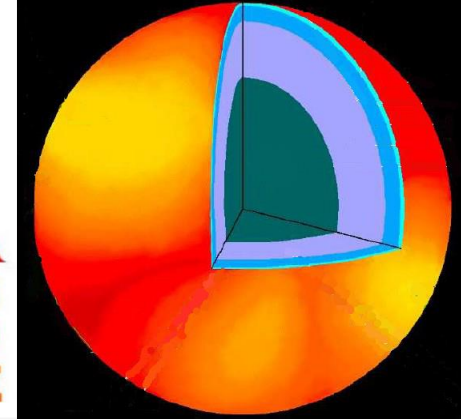
EOS of baryonic matter



Neutron matter



hyperon-mixed baryonic matter



Information on
 $\rho = \rho_0 \sim 2\rho_0$
is missing !

Nuclear saturation
density 0.16 fm^{-3}

0

Volume term of
mass formula (av)

-16

nuclei

ρ_0

$2\rho_0$

symmetric NM
($N=Z$)

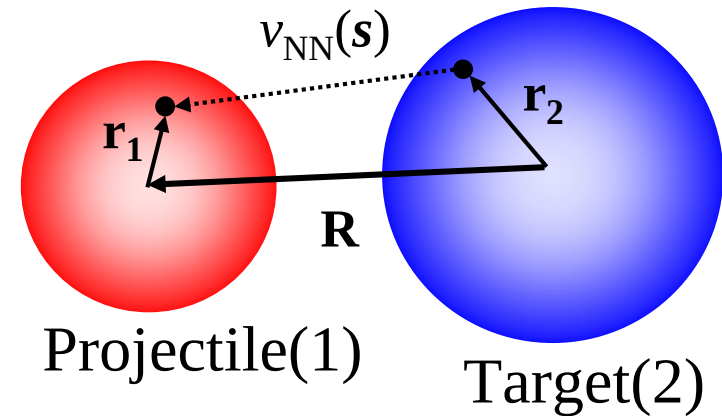
$5\rho_0$

Baryon density (ρ_B)

Neutron Stars

Saturation point

Double-Folding model Potential (DFM Potential)



$$\begin{aligned}
 U(\mathbf{R}) &= \int \rho_1(\mathbf{r}_1) \rho_2(\mathbf{r}_2) v_D(\mathbf{s}; \rho, E) d\mathbf{r}_1 d\mathbf{r}_2 \\
 &+ \int \rho_1(\mathbf{r}_1, \mathbf{r}_1 - \mathbf{s}) \rho_2(\mathbf{r}_2, \mathbf{r}_2 + \mathbf{s}) v_{EX}(\mathbf{s}; \rho, E) \exp\left[i \frac{\mathbf{K} \cdot \mathbf{s}}{M}\right] d\mathbf{r}_1 d\mathbf{r}_2 \\
 &= V_{DFM}(\mathbf{R}) + iW_{DFM}(\mathbf{R})
 \end{aligned}$$

Frozen-density approx. (FDA)

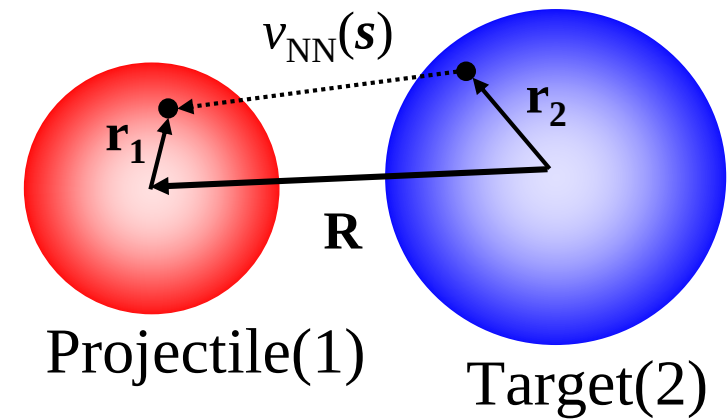
$$\rho = \rho_1 + \rho_2$$

Complex G-matrix interaction (CEG07/MP)

$$V_{D,EX} = v_{D,EX}^{(real)} + i v_{D,EX}^{(imag)}$$

- 1) T. Furumoto, Y. Sakuragi, Y. Yamamoto, *PRC* 78, 044610 (2008); *PRC* 80, 044614 (2009)
- 2) Y. Yamamoto, T. Furumoto, N. Yasutake, and Th. A. Rijken, *PRC* 88, 045805 (2014)

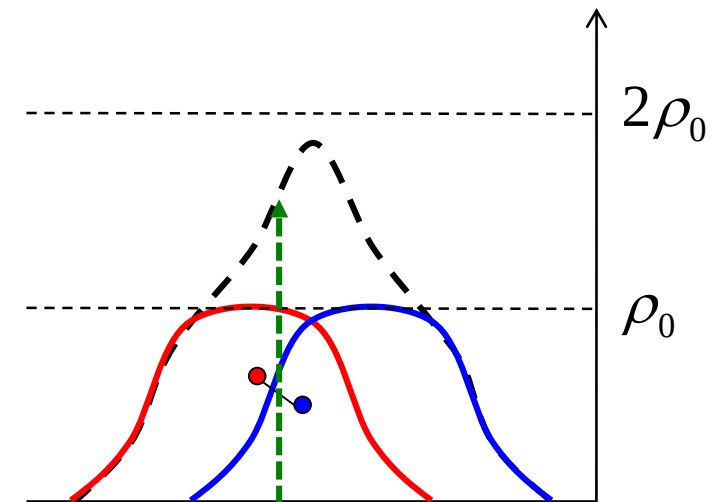
Double-Folding model Potential (DFM Potential)



$$\begin{aligned}
 U(\mathbf{R}) &= \int \rho_1(\mathbf{r}_1) \rho_2(\mathbf{r}_2) v_D(\mathbf{s}; \rho, E) d\mathbf{r}_1 d\mathbf{r}_2 \\
 &+ \int \rho_1(\mathbf{r}_1, \mathbf{r}_1 - \mathbf{s}) \rho_2(\mathbf{r}_2, \mathbf{r}_2 + \mathbf{s}) v_{EX}(\mathbf{s}; \rho, E) \exp\left[i \frac{\mathbf{K} \cdot \mathbf{s}}{M}\right] d\mathbf{r}_1 d\mathbf{r}_2 \\
 &= V_{DFM}(\mathbf{R}) + iW_{DFM}(\mathbf{R})
 \end{aligned}$$

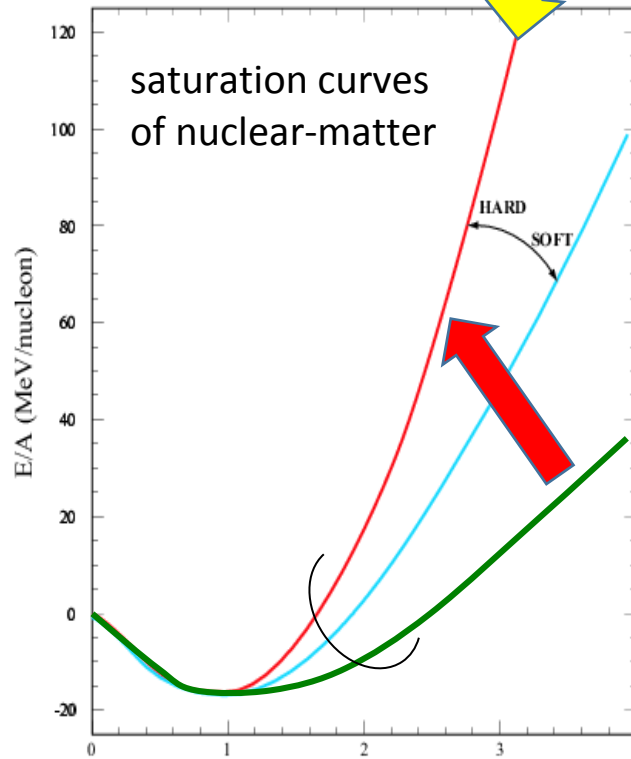
Frozen-density approx. (FDA)

$$\rho = \rho_1 + \rho_2$$

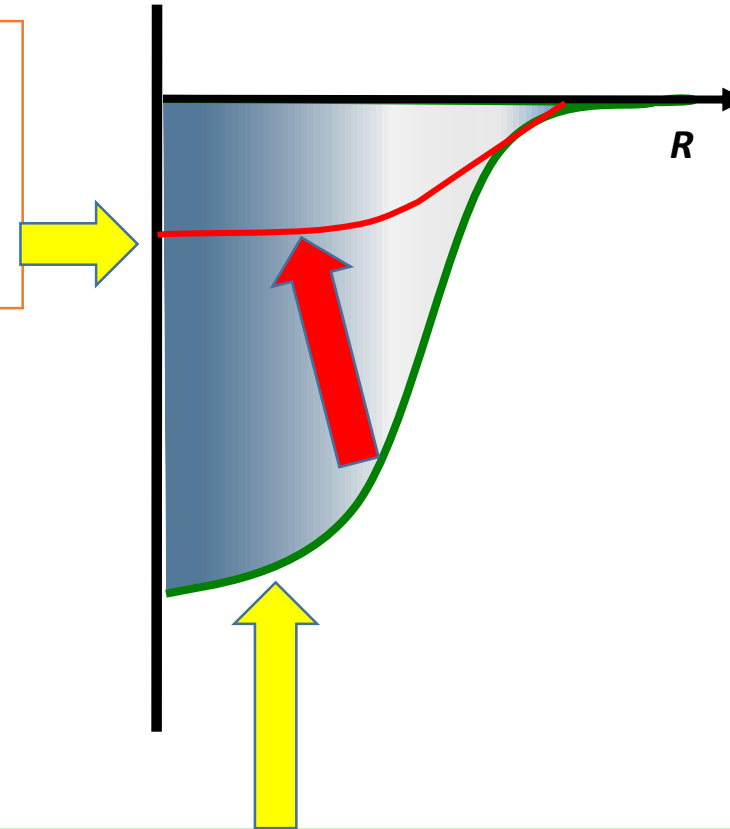


- If the nuclear matter is **hard**, the central depth of the potential become **shallow**.

B.E. per nucleon (E/A)



Nucleon density (ρ / ρ_0)



- If the nuclear matter is **soft**, the central depth of the potential may become **deep**.

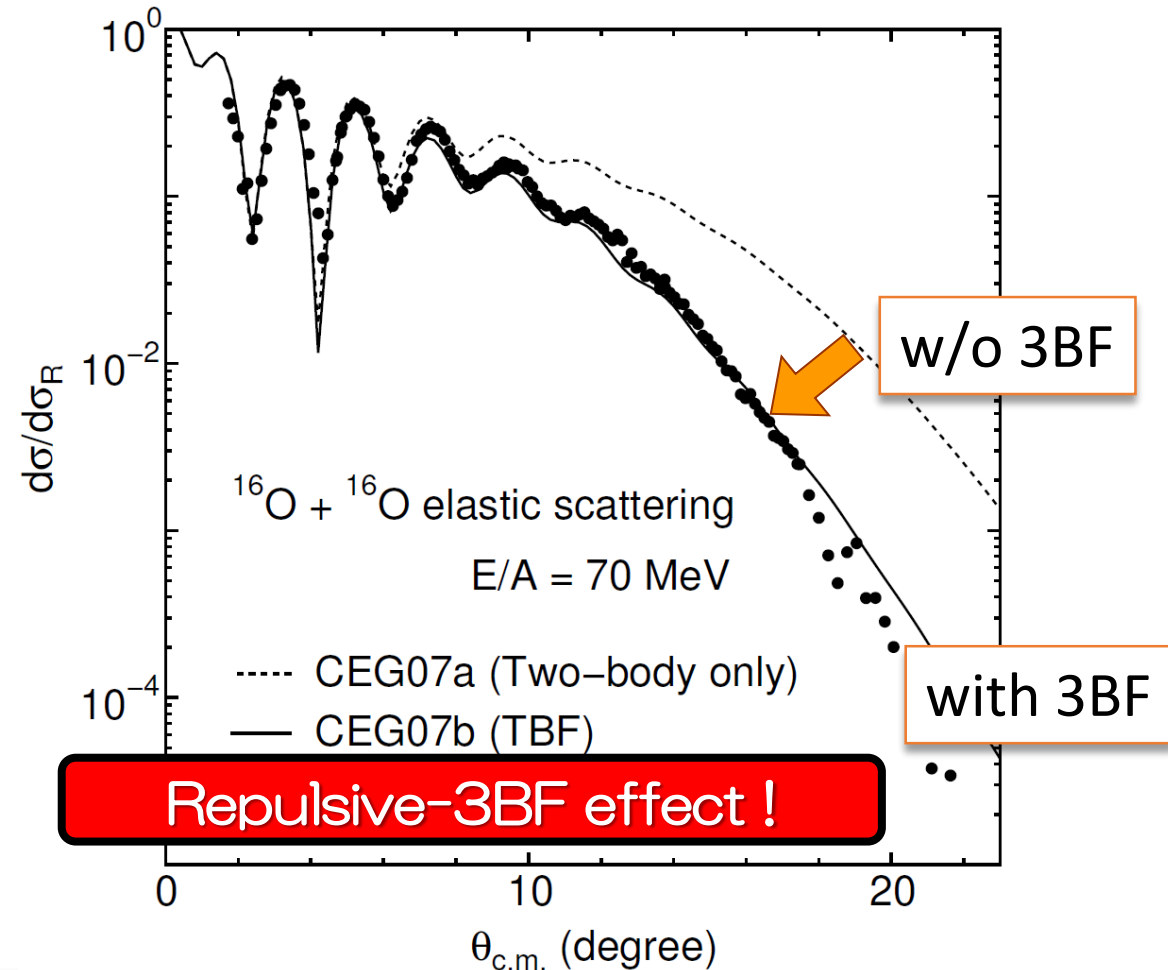
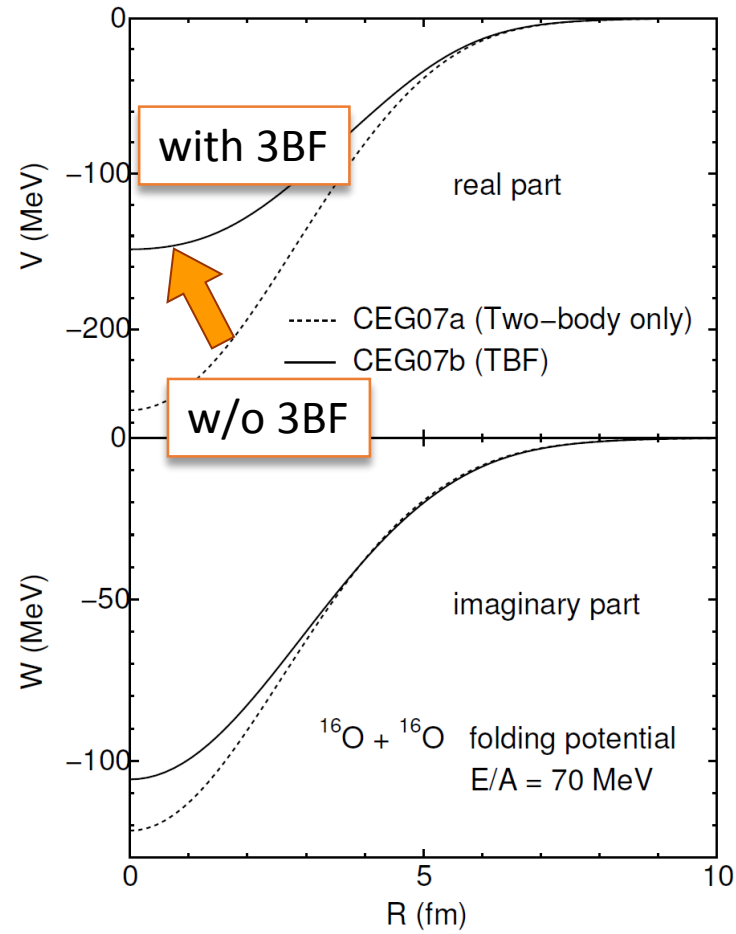
$^{16}\text{O}+^{16}\text{O}$: bench-mark system to test DFM

Frozen-density approx. (FDA)

$$\rho = \rho_1 + \rho_2$$

$^{16}\text{O}+^{16}\text{O}$ Elastic Scattering at $E/A = 70$ MeV : DFM with CEG07

⇒ decisive effect of Three-body force (mainly TBR) is clearly observed !



T.Furumoto, Y. Sakuragi, Y. Yamamoto, *PRC* 79, 011601(R) (2009)

Exp.: F. Nuoffer et al., *Nuovo Cimento A* 111, 971 (1998).

Multi-Pomeron Potential (MPP) :

• Y. Yamamoto, T. Furumoto, N. Yasutake, and Th. A. Rijken,
PRC 88, 022802(R) (2013); PRC 88, 045805 (2014)

✓ For N -tuple Pomeron vertex,
we take the Lagrangian :

$$\mathcal{L}_N = g_P^{(N)} \mathcal{M}^{4-N} \sigma_P^N(x) / N!$$

⇒ N -body local potential by
Pomeron exchange is written as →

$$\begin{aligned} V(\mathbf{x}'_1, \dots, \mathbf{x}'_N; \mathbf{x}_1, \dots, \mathbf{x}_N) &\equiv V(\mathbf{x}_1, \dots, \mathbf{x}_N) \prod_{i=1}^N \delta(\mathbf{x}'_i - \mathbf{x}_i), \\ V(\mathbf{x}_1, \dots, \mathbf{x}_N) &= g_P^{(N)} g_P^N \left\{ \int \frac{d^3 k_i}{(2\pi)^3} e^{-i\mathbf{k}_i \cdot \mathbf{x}_i} \right\} \\ &\times (2\pi)^3 \delta \left(\sum_{i=1}^N \mathbf{k}_i \right) \\ &\times \prod_{i=1}^N [\exp(-\mathbf{k}_i^2)] \mathcal{M}^{4-3N}, \end{aligned}$$

⇒ **Effective two-body potential** in
baryonic medium (ρ -dependent) →

$$\begin{aligned} V_{eff}^{(N)}(\mathbf{x}_1, \mathbf{x}_2) &= \rho_{NM}^{N-2} \int d^3 x_3 \dots \int d^3 x_N V(\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_N) \\ &= g_P^{(N)} g_P^N \frac{\rho_{NM}^{N-2}}{\mathcal{M}^{3N-4}} \frac{1}{\pi \sqrt{\pi}} \left(\frac{m_P}{\sqrt{2}} \right)^3 \\ &\times \exp \left(-\frac{1}{2} m_P^2 r_{12}^2 \right), \end{aligned} \tag{3}$$

Multi-Pomeron Potential (MPP) :

• Y. Yamamoto, T. Furumoto, N. Yasutake, and Th. A. Rijken,
PRC 88, 022802(R) (2013); PRC 88, 045805 (2014)

✓ **N=3 (effective 3BF)**

✓ **N=4 (effective 4BF)**

$$V_{\text{eff}}^{(3)}(r) = g_P^{(3)} (g_P)^3 \frac{\rho}{\mathcal{M}^5} F(r),$$

$$V_{\text{eff}}^{(4)}(r) = g_P^{(4)} (g_P)^4 \frac{\rho^2}{\mathcal{M}^8} F(r),$$

$$F(r) = \frac{1}{4\pi} \frac{4}{\sqrt{\pi}} \left(\frac{m_P}{\sqrt{2}} \right)^3 \exp \left(-\frac{1}{2} m_P^2 r^2 \right)$$

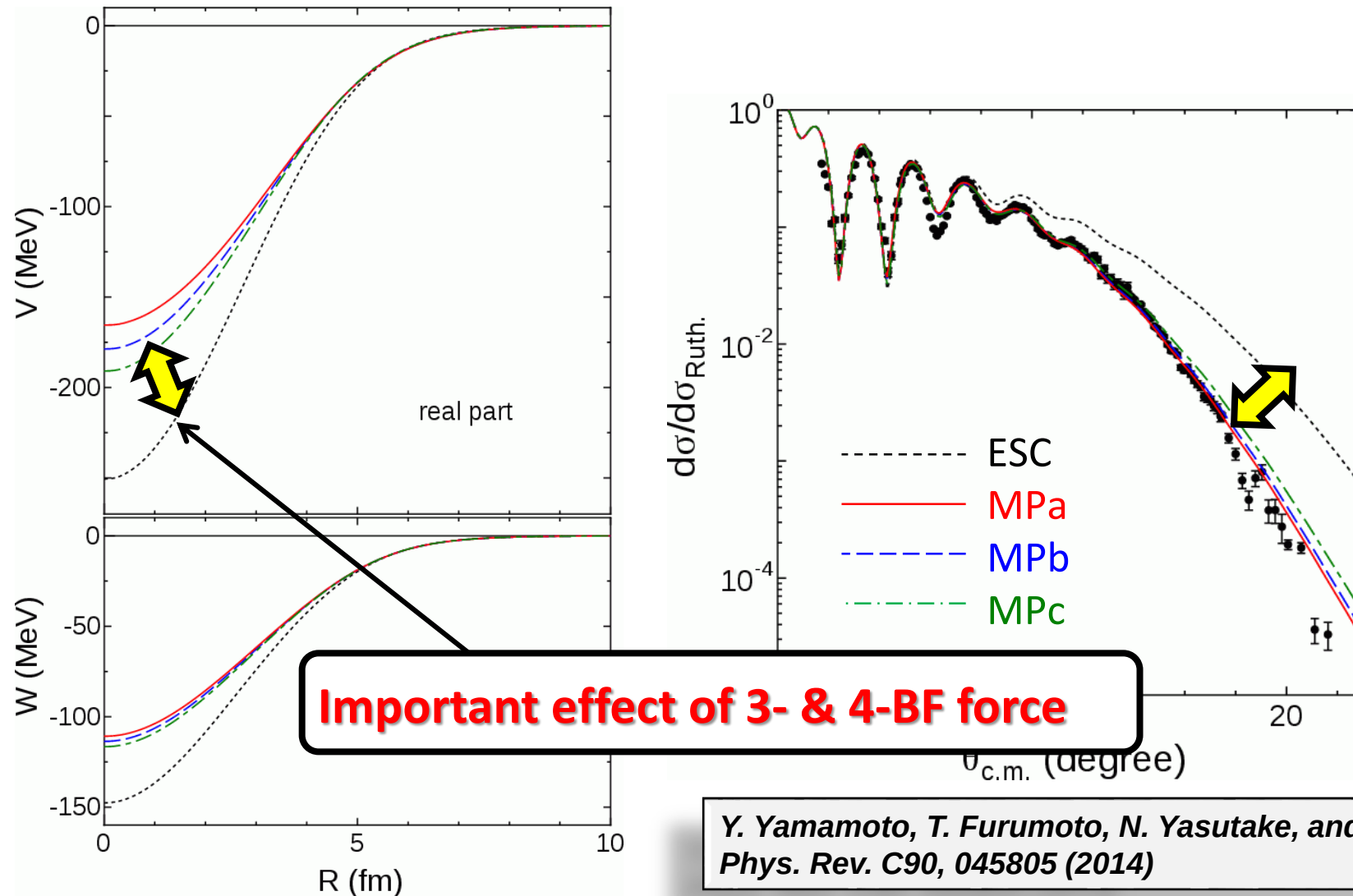
N=3 & N=4

⇒ **Effective two-body potential** in
baryonic medium (**ρ -dependent**) →

$$V_{\text{eff}}^{(N)}(\mathbf{x}_1, \mathbf{x}_2) = \rho_{NM}^{N-2} \int d^3x_3 \dots \int d^3x_N V(\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_N)$$

$$= g_P^{(N)} g_P^N \frac{\rho_{NM}^{N-2}}{\mathcal{M}^{3N-4}} \frac{1}{\pi \sqrt{\pi}} \left(\frac{m_P}{\sqrt{2}} \right)^3 \times \exp \left(-\frac{1}{2} m_P^2 r_{12}^2 \right), \quad (3)$$

$^{16}\text{O} + ^{16}\text{O}$ elastic scattering cross section at $E/A = 70$ MeV

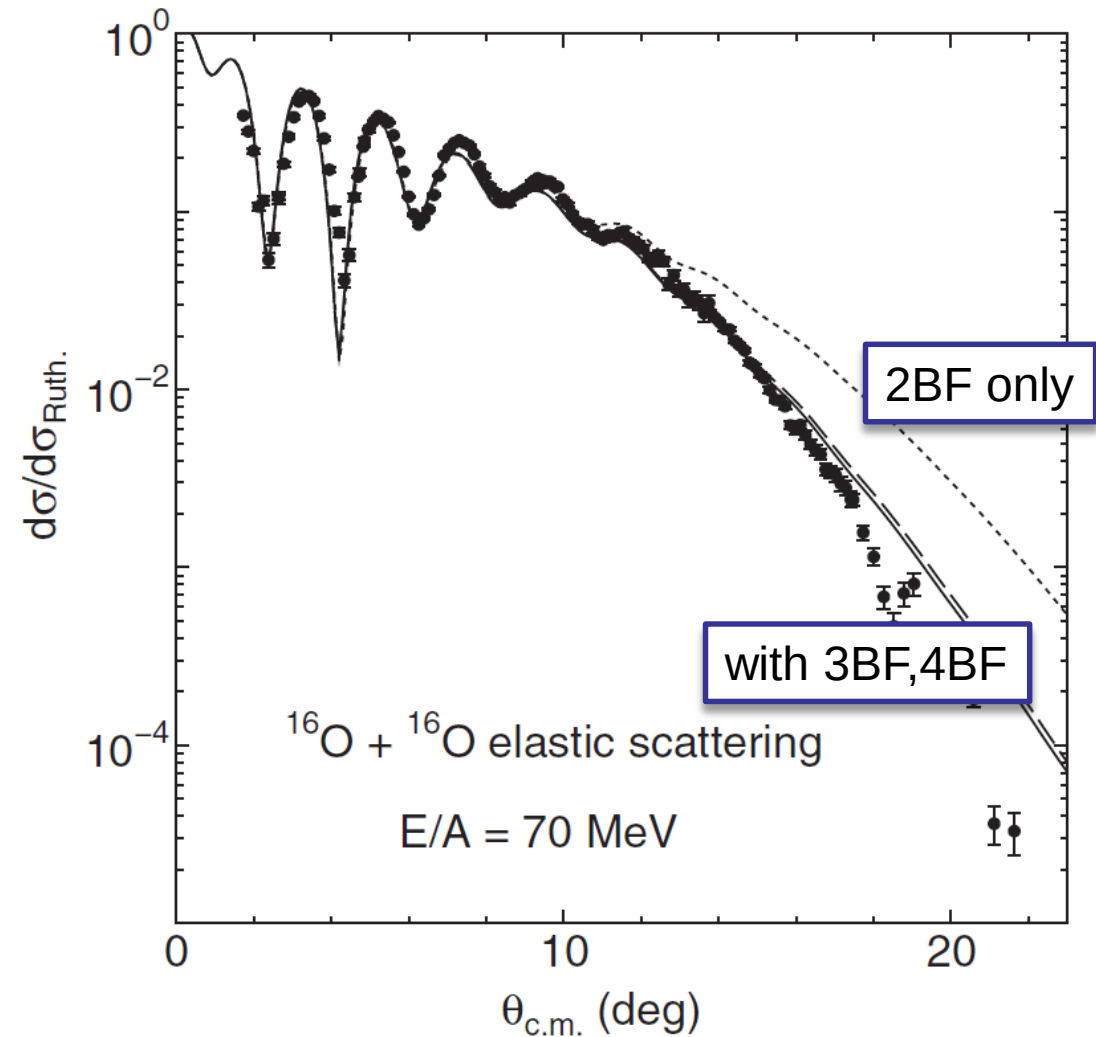
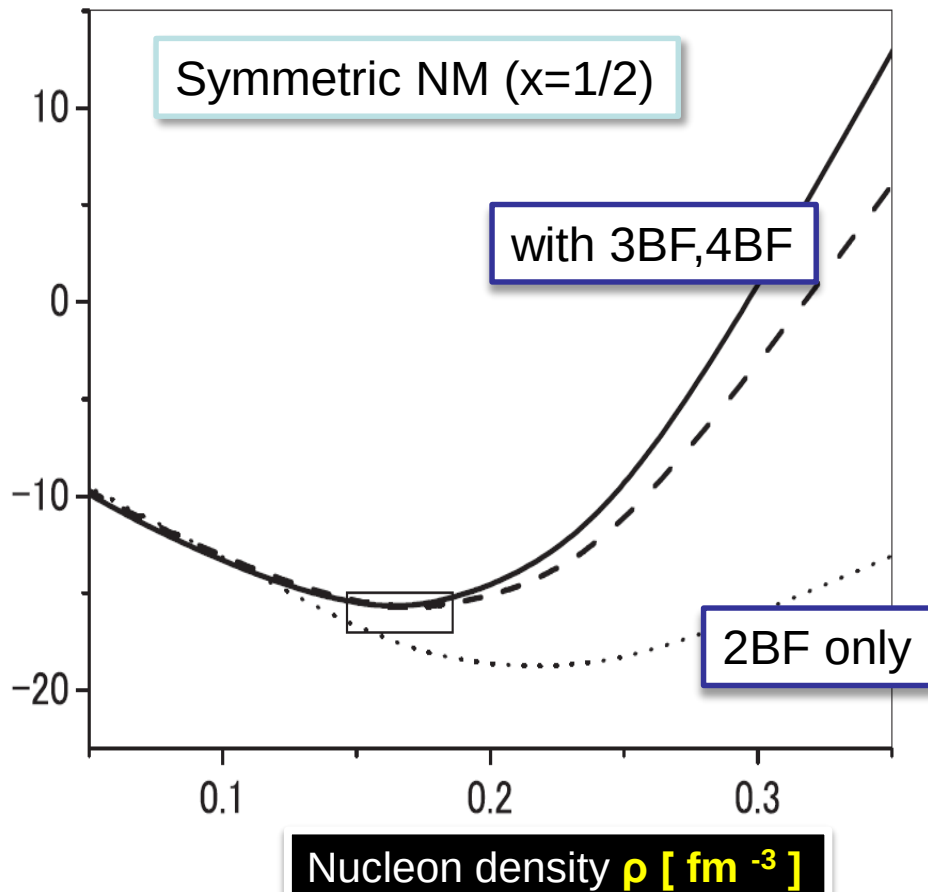


Effects of Three & Four-body force (with multi-Pomeron exchange model)

Y. Yamamoto, T. Furumoto, N. Yasutake, Th. A. Rijken, *PRC* 88, 022802(R) (2013)

MULTI-POMERON REPULSION AND THE NEUTRON-STAR MASS!

Binding energy per nucleon: E/A [MeV]



Parameters for MPP model for **3BF and 4BF** determined **by A-A scattering analysis** and the corresponding values (E_0 , K , E_{sym} , L) for saturation curves of symmetric nuclear matter at $\rho = \rho_0$

at $\rho = \rho_0$	MPa	MPb	MPc	ESC (NN only)	Exp. ¹⁾
$g_P^{(3)}$	2.34	2.94	2.34	0	-
$g_P^{(4)}$	30.0	0	0	0	-
E_0 (MeV)	-15.8	-15.8	-15.6	-	-16
K (MeV)	310	280	260	84	
E_{sym} (MeV)	33.1	33.1	32.7	-	32.5 ± 0.5 ¹⁾
L (MeV)	70	69	67	-	70 ± 15 ¹⁾

$$V_{\text{eff}}^{(3)}(r) = g_P^{(3)}(g_P)^3 \frac{\rho}{\mathcal{M}^5} F(r),$$

$$V_{\text{eff}}^{(4)}(r) = g_P^{(4)}(g_P)^4 \frac{\rho^2}{\mathcal{M}^8} F(r),$$

$$F(r) = \frac{1}{4\pi} \frac{4}{\sqrt{\pi}} \left(\frac{m_P}{\sqrt{2}} \right)^3 \exp \left(-\frac{1}{2} m_P^2 r^2 \right)$$

Y. Yamamoto, T. Furumoto, N. Yasutake,
Th. A. Rijken, *PRC*90, 045805 (2014)

1) P. Möller, et al., *PRL*108, 052501 (2012).

We have proposed the **complex G-matrix folding model**

- the complex G-matrix is derived from the **ESC08 NN interaction**
- includes the effect of **the multi-body repulsive force**
- consist of the **repulsive** and **attractive** parts

Y. Yamamoto, T. Furumoto, N. Yasutake, and Th. A. Rijken,
Phys.Rev.C 90, 045805 (2014)

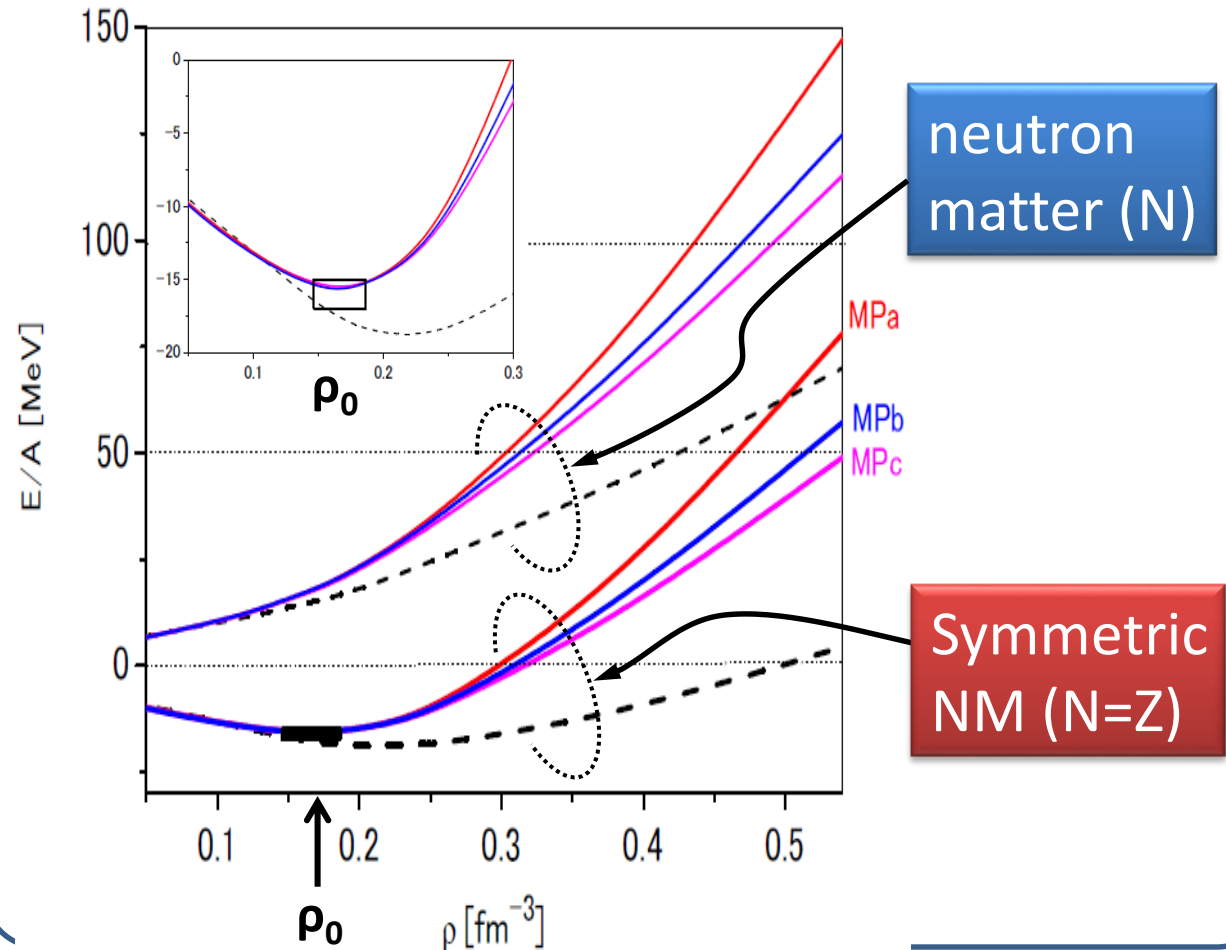
MP-model (MPa, MPb, MPc)
(multi-pomelon exchange model)

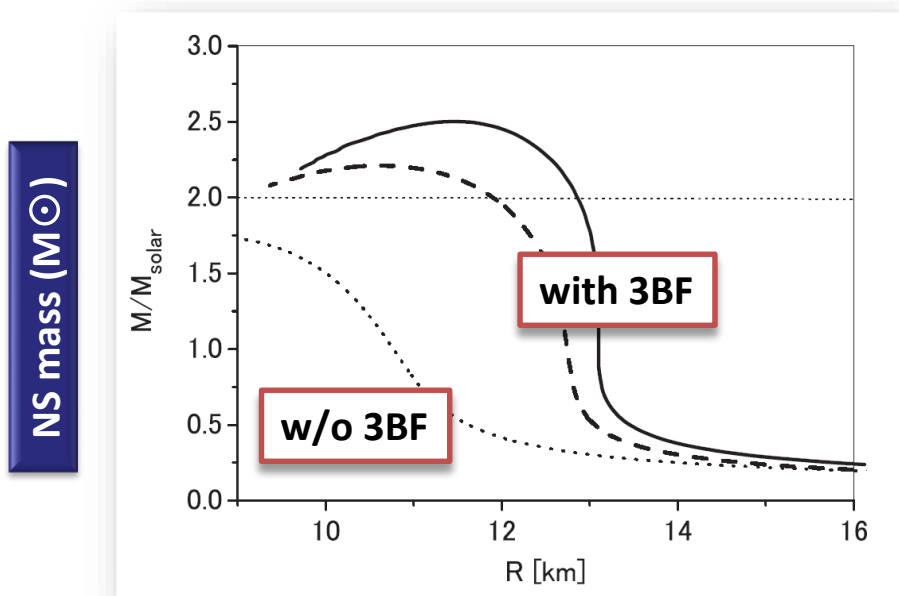
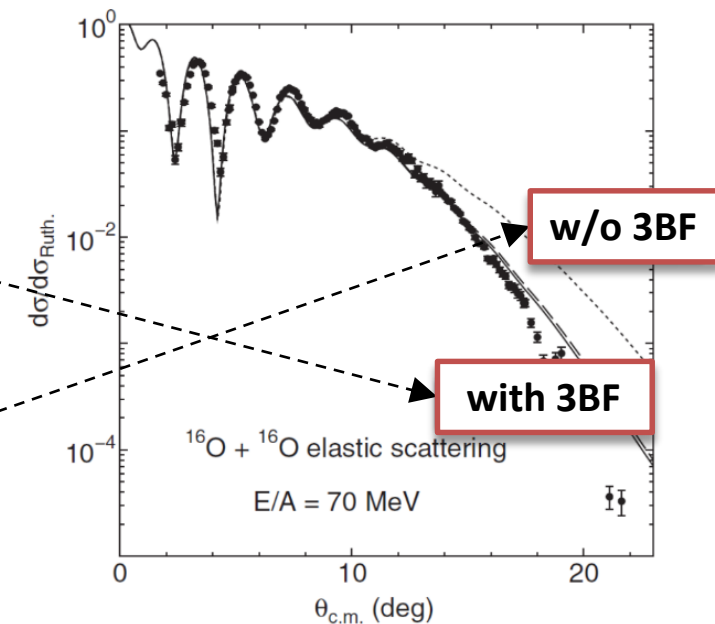
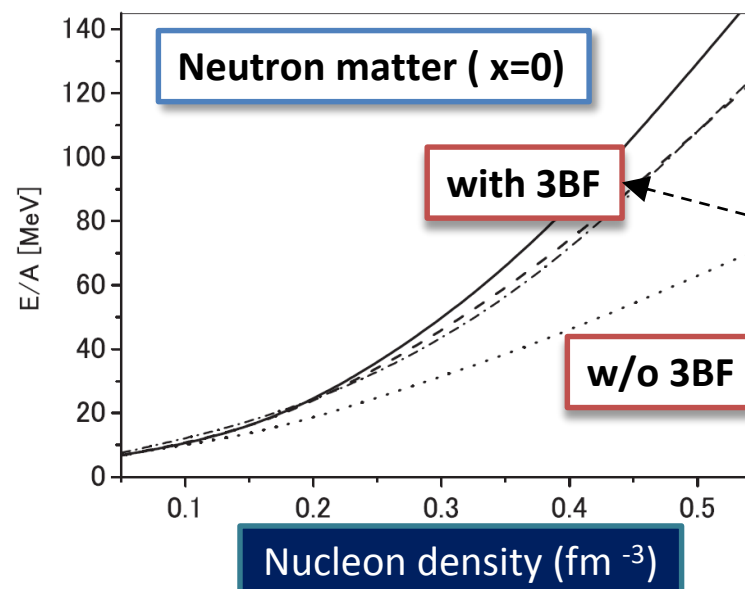
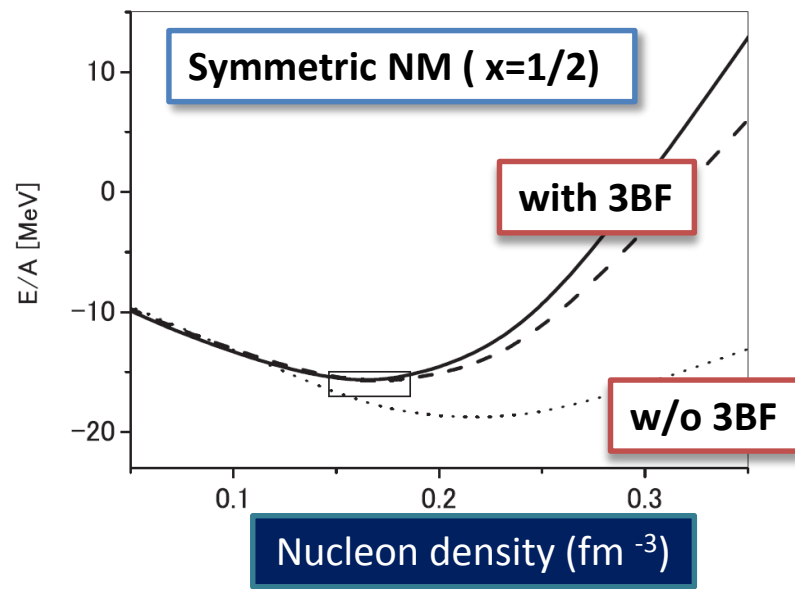
ESC : *two-body* only

MPa : with **three- & four-body** forces

MPb : with **three-body** forces

MPc : with **three-body** forces





Complex G -matrix
> ESC+MPP model

- Y. Yamamoto, T. Furumoto, N. Yasutake, Th. A. Rijken, *Phys.Rev.C* 88, 022802(R) (2013)
- Y. Yamamoto, T. Furumoto, N. Yasutake, Th. A. Rijken, *Phys.Rev.C* 90, 045805 (2014)

Hyperon mixing ratio as a function of baryon density

MPa (with 3&4BF)

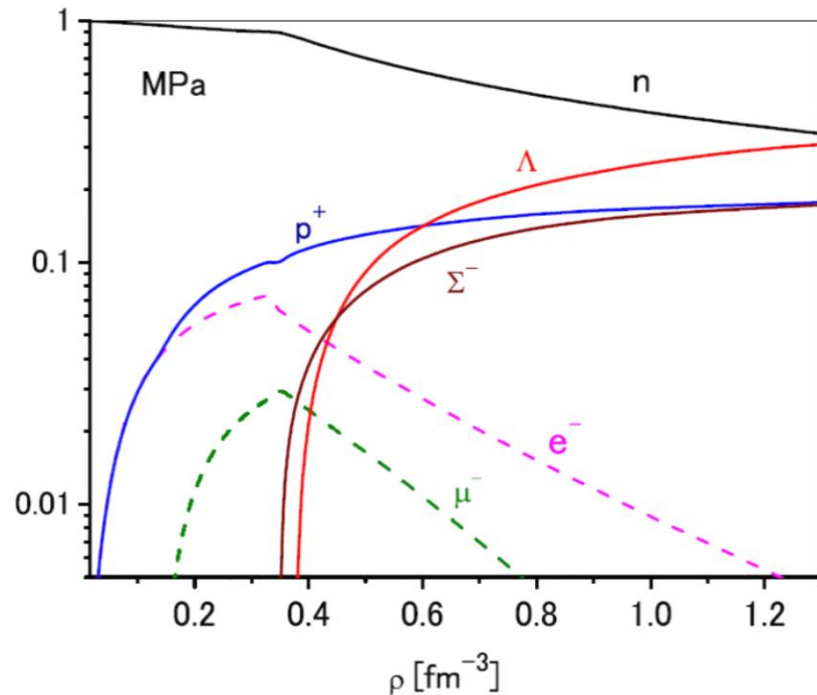


FIG. 8. (Color online) Composition of hyperonic neutron-star matter in the case of MPa.

ESC (2BF only)

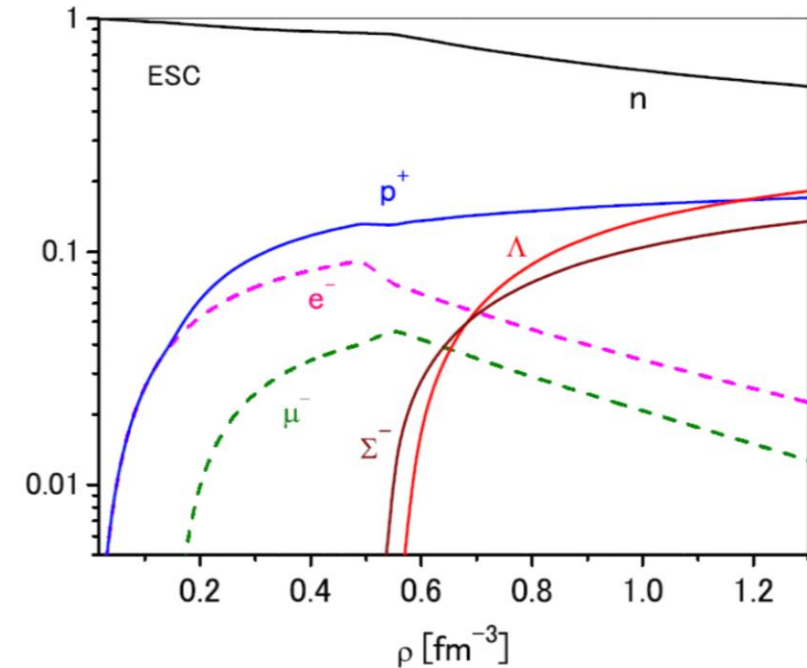
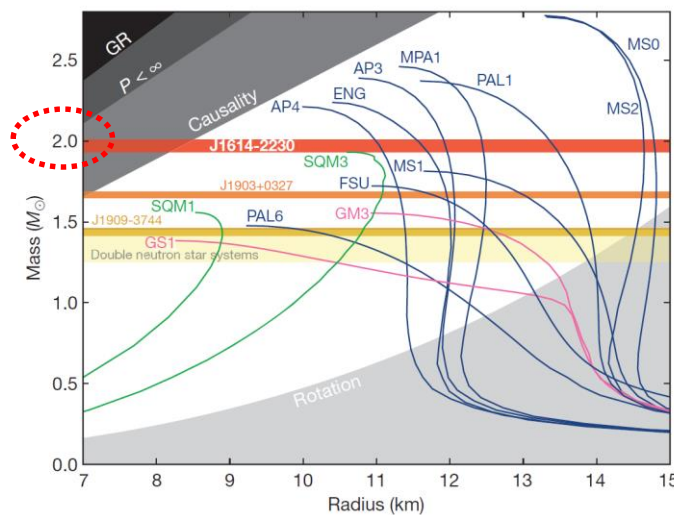
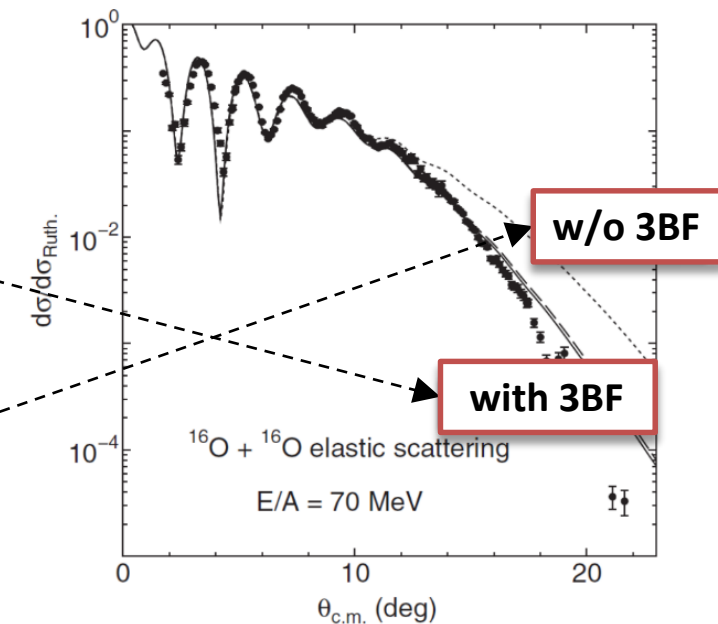
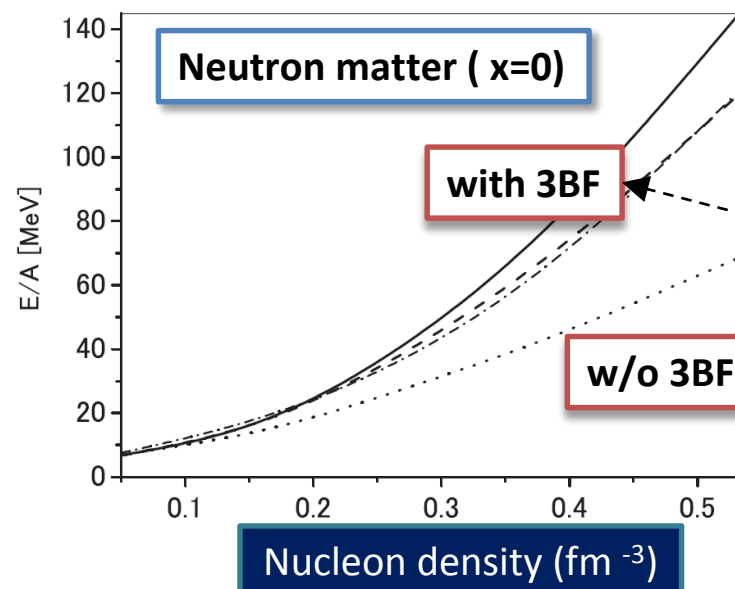
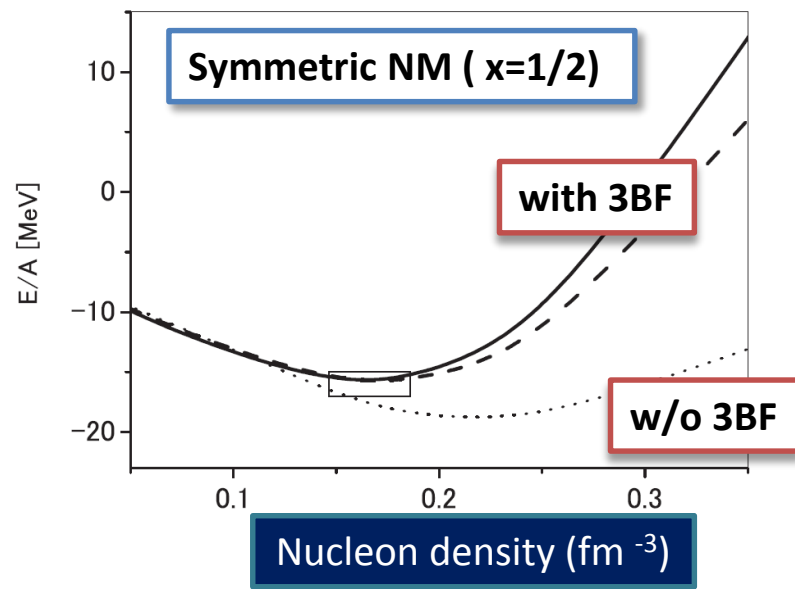
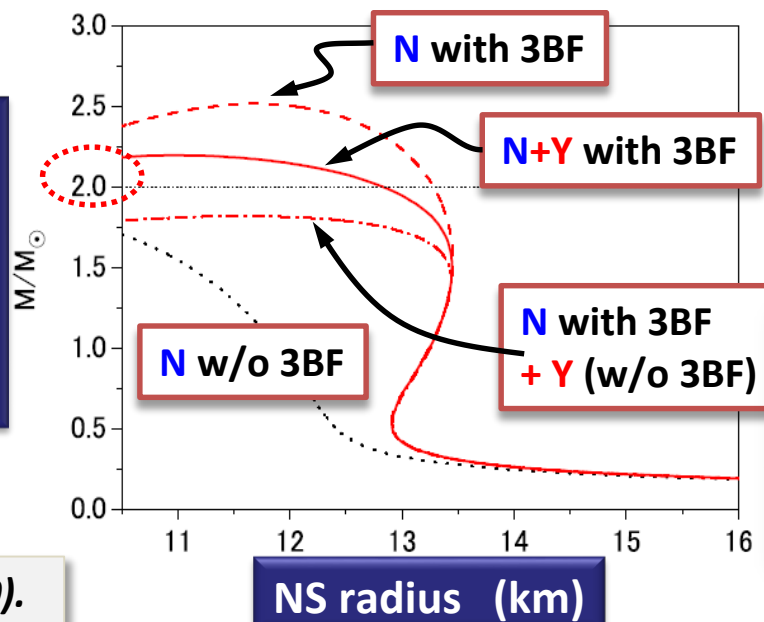


FIG. 9. (Color online) Composition of hyperonic neutron-star matter in the case of ESC.

Y. Yamamoto, T. Furumoto, N. Yasutake and Th. A. Rijken,
Phys. Rev. C 90, 045805 (2014)



NS mass ($M_{\odot}$)



Complex G -matrix
> ESC+MPP model

- Y. Yamamoto, T. Furumoto, N. Yasutake, Th. A. Rijken, *PRC* 88, 022802(R) (2013)
- Y. Yamamoto, T. Furumoto, N. Yasutake, Th. A. Rijken, *PRC* 90, 045805 (2014)

P.B. Demorest et al., *Nature* 467, 1081 (2010).

Conclusion: ← Our strategy

- ✓ The 3BF is very sensitive to **nucleus-nucleus (A-A) elastic scattering**.
- ✓ We determine the strength of universal **3,4-baryon repulsion** by the analyses of **A-A elastic scattering** with the **G-matrix folding model**.
- ✓ Then, construct the **EOS of neutron-star matter**,
- ✓ That leads to **a stiff EOS**, enough to give the **observed neutron-star mass** ($M_{NS} \approx 2M_{\odot}$) even if we take account of the **hyperon mixing**.